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Emerging cool white light emission from Dy³⁺ doped single phase alkaline earth niobate phosphors for indoor lighting applications

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Single-phase cool white-light emitting BaNb₂O₆:Dy³⁺ phosphors have been synthesized *via* a conventional solid-state reaction method and characterized using X-ray diffraction (XRD), scanning electron microscopy (SEM) observations and spectrofluorophotometric measurements. XRD and Rietveld structural refinement studies confirm that all the samples exhibit pure orthorhombic structure [space group – $C222_1(20)$]. SEM observations reveal the dense particle packaging with irregular morphology in a micron range. The as-prepared phosphors exhibit blue (482 nm) and yellow (574 nm) emissions under 349, 364, 386 and 399 nm excitations corresponding to ${}^4F_{9/2} \rightarrow {}^6H_J$ (J = 15/2, 13/2) transitions of Dy³⁺ ions. The energy transfer mechanism between Dy³⁺ ions has been studied in detail and the luminescence decay lifetime for the ${}^4F_{9/2}$ level was found to be around 146.07 µs for the optimized phosphor composition. The calculated Commission Internationale de L'Eclairage (CIE) chromaticity coordinates for the optimized phosphor are (x = 0.322, y = 0.339), which are close to the National Television Standard Committee (NTSC) (x = 0.310, y = 0.316) coordinates. The values of CIE chromaticity coordinates and correlated color temperature (CCT) of 5907 K endorse cool white-light emission from the phosphor. The study reveals that BaNb₂O₆:Dy³⁺ phosphor could be a potential candidate for near ultra-violet (NUV) excited white-LED applications.

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1. Introduction

In recent years, rare earth ion doped inorganic luminescent materials have been extensively studied in the fields of materials science, physics, chemistry and life sciences due to their various potential applications in display devices (e.g., cathode ray tubes, vacuum fluorescent displays, and field emission displays), lighting gadgets (e.g., fluorescent tubes and white-light emitting diodes), solid-state lasers, biological labelling, X-ray, medical devices, ionization radiation and so on. Among them, white light emitting diodes (w-LEDs) have been considered to be the next generation illumination sources in the field of solid-state lighting instead of traditional incandescent and currently implemented fluorescent lamps due to their numerous advantages such as small size, high energy efficiency, energy-saving, robustness, high brightness, fast switching, longer life time (>100 000 h) and environ-

mental friendliness. 3-5 Currently, two approaches have been implemented to achieve white light through solid-state lighting (SSL). The first one is a phosphor-free SSL approach employing RGB-LEDs, which consist of red, green and blue monochromatic LEDs to obtain white-light. The main drawback of this approach is that every LED must be adjusted by individual power supply to balance the emission intensity of each color. However, the second approach involves phosphor integrated to the SSL device. The phosphor-converted (pc) LED uses an ultra-violet (UV)/near ultra-violet (NUV) light in combination with single/multiple phosphors that convert a part of the light emitted by the UV/NUV LED into white-light. 6,7 At present, most of the commercially available w-LEDs are based on the second approach because of the simplicity in operation. The combination of blue LED (InGaN) coated with yellowemitting (Y3Al5O12:Ce3+) phosphor is one of the widely used approaches currently to produce w-LEDs. However, this approach encounters some serious issues such as the halo effect of blue/yellow color separation, color dependence on chromaticity, and poor color-rendering index (<65) due to the lack of green and red-emitting phosphor components at long wavelength regions, which limits the LED applications further.8 On the other hand, UV/NUV LED coated with multi-

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color-emitting phosphors is an alternative approach to obtain white-light emission with an excellent color-rendering index (>90) as the luminous efficiency of NUV/UV chip pumped w-LEDs is higher than the blue chip.9 However, UV/NUV excitable multiple color-emitting phosphors have some disadvantages such as low luminous efficiency due to blending together of blue, green and red-emitting phosphors and different degradation schedules of each phosphor.^{9,10} Therefore, single phase phosphor has become a necessary pre-requisite for the fabrication of w-LEDs using UV/NUV LED chips to overcome the problems mentioned above. There are different methodologies to obtain white light emission from a single phase host lattice by (i) doping a single rare earth (RE) ion, (ii) doping of two or more RE ions, which are excited simultaneously, (iii) co-doping of different ions and controlling the emission via energy transfer processes, and (iv) controlling the concentration of the defect and reaction conditions of defect related luminescent materials.11

Alkaline earth niobates have emerged as novel materials with huge technological and scientific importance due to their excellent non-linear optical, photocatalytic, piezoelectric, ionic conductive and photorefractive properties for the applications involving acoustic transducers, delay lines in filters, optical modulators, beam deflectors etc. 12,13 Among all inorganic phosphors, oxide based phosphors are a preferred choice for display and SSL applications due to their exceptional chemical stability, inertness and moisture resistance. Moreover, rareearth doped alkaline earth niobate phosphors have attracted much attention and have been applied for light emitting diode (LED) and plasma display panel (PDP) applications. 14 The metaniobate ceramics, with the general formula M2+Nb2O6 (M²⁺ = divalent alkaline earth or transition metals) are subcomponents of the complex perovskite family, A(M_{1/3}Nb_{2/3})O₃ and they mostly exist in an isostructural form of orthorhombic structure with a columbite mineral group with the exception of the M = Sr, Ba and Pb analogues, which crystallize in different orthorhombic structures. 15-17 In this quest, barium metaniobate (BaNb₂O₆) has been selected as a host lattice due to its smaller band gap and exhibits higher charge generation under UV light irradiation compared to other binary niobates. Moreover, BaNb₂O₆ has been widely used as a high quality refractory material for photocatalytic and microwave dielectric applications in recent years. 18 The luminescent properties of various rare-earth ion doped niobate phosphors have been investigated and reported elsewhere. 14,19-21 However, to the best of our knowledge, luminescent properties of Dy³⁺ ion doped BaNb2O6 phosphors have never been reported in the literature. Hence, in the current study, the said phosphor has been chosen to be synthesized and investigated thoroughly for the first time. In addition, it is well-known that Dy3+ ions with 4f9 electronic configuration have complex energy levels and various possible transitions between f-f levels that are highly selective and exhibit sharp line spectra.²² It gives emission in blue and vellow bands and the intensity ratio of these two emission bands depends on the host crystal structure. 23,24

In the current work, a series of Dy3+ ion doped single-phase BaNb₂O₆ phosphors have been prepared by a solid-state reaction method to explore their possibility as potential phosphors for white-LEDs by investigating photoluminescence and colorimetric properties in detail.

Experimental procedure

 $Ba_{(1-x)}Nb_2O_6:xDy^{3+}$ (where x = 0.01, 0.1, 0.5, 1.0, 1.5, 2.0 and 2.5 mol%) were synthesized by the conventional solid-state reaction method. The precursor chemicals namely, BaCO₃ (Fisher Scientific, 99%), Nb₂O₅ (Fisher Scientific, 99.9%) and Dy₂O₃ (Sigma Aldrich, 99.9%) all of AR grade were taken as starting materials. A stoichiometric amount of precursor materials was thoroughly mixed using an agate mortar and pestle for an hour with acetone as a dispersing medium. The raw powder sample was kept in an alumina crucible and then heated in a programmable muffle furnace at 625 °C for an hour to remove CO2 and then sintered at 1200 °C at a heating rate of 6 °C min⁻¹ for 5 hour under an air atmosphere. Finally, the sample was naturally cooled to room temperature (~25 °C) in the furnace itself.

The structure of the prepared sample was determined through XRD. The crystalline phases were identified by using a X-ray diffractometer (Rigaku make, model-Mini flex-II), using nickel-filtered Cu K_{α} radiation ($\lambda = 1.54056$ Å) in the range of $20^{\circ} \le 2\theta \le 60^{\circ}$ and the accelerating voltage was maintained at 30.0 kV and the tube current at 15 mA. The FullProf suite programme was used to reveal the structural refinement. The morphological observations were carried out by SEM (Hitachi, Model-S-3700N). The photoluminescence excitation (PLE) and photoluminescence (PL) spectra were recorded using a Shimadzu spectrofluorophotometer (model: RF-5301PC) fitted with a Xenon flash lamp. The lifetime measurements were carried out using a time-resolved luminescence spectrometer (model-F900 Edinburgh), equipped with a time correlated single photon counting system and a microsecond xenon flash lamp as the source of excitation.

3. Colorimetric theory

The color of any object (self-luminous or reflecting) can be conveniently specified via Commission International de L' Eclairage (CIE) chromaticity coordinates marked on a chromaticity diagram.²⁵ These color coordinates are calculated from the PL emission spectra using $\bar{x}(\lambda)$, $\bar{y}(\lambda)$ and $\bar{z}(\lambda)$ color matching functions defined in the CIE 1931. For a specified power spectral density $P(\lambda)$, the degree of stimulation required to match the color of $P(\lambda)$ is given by three equations:

$$X = \int_{\lambda} \bar{x}(\lambda) P(\lambda) d\lambda$$

$$Y = \int_{\lambda} \bar{y}(\lambda) P(\lambda) d\lambda$$

$$Z = \int_{\lambda} \bar{z}(\lambda) P(\lambda) d\lambda$$

where X, Y, and Z are the tristimulus values. The tristimulus values provide stimulation (i.e. power) values for each of three primary (red, green, blue) colors to match the color of $P(\lambda)$.

The tristimulus values specifying the color are stored in the ratio of the primary colors and not in the specific amounts of each individual primary color, the number of dimensions used to match the color of $P(\lambda)$. This was done for the CIE 1931 XYZ space, and the resultant CIE 1931 x-y chromatic diagram is the most commonly used tool, to describe color, in spectroscopy today. From (X, Y, Z), the (x, y) chromaticity coordinates are calculated as:

$$x = \frac{X}{X + Y + Z}$$

$$y = \frac{Y}{X + Y + Z}$$

Further, the quality of white light is evaluated from the chromaticity coordinates using McCamy's relation, 26 by first evaluating the ratio between the inverse slope line and the chromaticity epicentre as:

$$n = \frac{(x - x_{\rm e})}{(y - y_{\rm e})}$$

where $x_e = 0.3320$ and $y_e = 0.1858$. Then the Correlated Color Temperature (CCT) was evaluated as:

$$CCT = -449n^3 + 3525n^2 - 6823.3n + 5520.33$$

The color purity or color saturation of a light source is the distance in the chromaticity diagram between the (x, y) colorcoordinate point of the test source and the coordinate of the equal-energy point divided by the distance between the equalenergy point and the dominant wavelength point. The color purity is thus given by:

Color purity =
$$\frac{\sqrt{(x - x_{ee})^2 + (y - y_{ee})^2}}{\sqrt{(x_d - x_{ee})^2 + (y_d - y_{ee})^2}}$$

where (x, y), (x_{ee}, y_{ee}) and (x_d, y_d) represent the chromaticity coordinates of the light source under test, equal-energy reference illuminant and dominant-wavelength respectively.27

4. Results and discussion

Structural and morphological analysis

The crystal structure of the BaNb₂O₆ compound was reported for the first time by Sirotinkin et al. in 1990, and has an orthorhombic [space group C2221 (20)] structure with cell parameters a = 7.880 Å, b = 12.215 Å, c = 10.292 Å and $Z = 8.^{28} \text{ The}$ phase purity of the as-prepared phosphors was characterized using XRD. Fig. 1 illustrates the XRD patterns of undoped and Dy3+ ion doped BaNb2O6 phosphors that are in agreement with the PDF-4 + (ICDD) standard card no. 04-012-8861. It was found that there is no change in peak positions among all

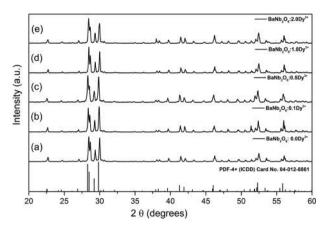


Fig. 1 Powder XRD patterns of $Ba_{(1-x)}Nb_2O_6:xDy^{3+}$ (x = 0.0, 0.1, 0.5, 1.0, 2.0 mol%) phosphors.

XRD patterns, which indicate pure phase formation for undoped and doped BaNb2O6 phosphors with respect to their corresponding (hkl) planes. This fact could be due to the large ionic radius of Ba²⁺ (1.42 Å) than Dy³⁺ (1.02 Å) and Dy³⁺ ions may occupy Ba2+ sites when they enter into the BaNb2O6 host lattice. Hence, no additional peaks were found up to 2 mol% doping of Dy₂O₃, which means that Dy³⁺ ions were successfully substituted for Ba²⁺ ions without changing the crystal structure of the host lattice.

The average crystallite size (D) and strain (ε) of the samples were calculated using the most reliable Williamson-Hall (W-

H) equation^{29,30}
$$\left[\beta\cos\theta = \left(\frac{K\lambda}{D}\right) + 4\varepsilon\sin\theta\right]$$
, where $K = \text{shape}$

factor (0.94), D is the average crystallite size, λ is the wavelength of CuK α radiation, θ is Bragg's diffraction angle of the planes and β is the corrected full width at half maximum (FWHM). The average crystallite size of $Ba_{(1-x)}Nb_2O_6:xDy^{3+}$ (x = 0.0, 0.1, 0.5, 1.0 and 2.0 mol%) samples was found to be in the range 44 and 55 nm. The strain present in the lattice was calculated and the values for x = 0.0, 0.1, 0.5, 1.0 and 2.0 mol% Dy³⁺ ion doped BaNb₂O₆ sintered at 1200 °C were found to be 0.081, 0.114, 0.113, 0.061 and 0.091, respectively. Further the average crystallite size calculated by Debye-Scherer's formula³¹ $[D = K\lambda/\beta \cos \theta]$ was found to be in the range 40 and 89 nm and is in good agreement with the calculations performed using the W-H method.³² The Rietveld refinement of undoped BaNb₂O₆ was carried out using FullProf software shown in Fig. 2 and the resulting parameters are summarized in Table 1. The results indicate a good agreement between the observed and calculated diffraction patterns of orthorhombic phase [space group C2221 (20)] without any anonymous peak.³³ The unit cell parameters were determined and are found to be a = 7.8707 Å, b = 12.2096 Å, c = 10.2881 Å and cell volume (\vec{V}) = 988.680 Å³, which are close to those reported by Sirotinkin et al.28 The refinement finally converges to goodness-of fit parameters (χ^2) = 4.2%, $R_{\rm wp}$ = 31.0% and $R_{\rm p}$ = 25.6%. Fig. 3 shows the orthorhombic structure of a BaNb₂O₆

lattice projected onto the *b*–*c* plane.

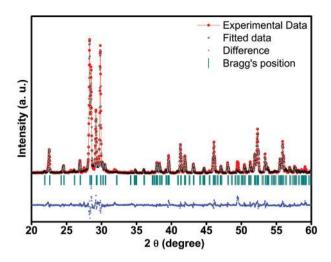


Fig. 2 Experimental, calculated and difference in X-ray diffraction patterns of BaNb₂O₆ powder after Rietveld refinement.

Table 1 Calculated crystallographic data of BaNb₂O₆ by the Rietveld refinement method

Formula	$\mathrm{BaNb_2O_6}$	
Radiation	Си Ка	
2θ (°)	20-60	
Symmetry	Orthorhombic	
Space group	$C222_{1}(20)$	
a (Å)	7.8707	
b (Å)	12.2096	
c (Å)	10.2881	
α (°)	90	
β (\circ)	90	
γ (°)	90	
Z	8	
$R_{\rm p}$	25.6	
$R_{\rm p}$ $R_{\rm wp}$ X^2	31.0	
X^{2P}	4.2	
$V(\mathring{A}^3)$	988.680	
Density (mg m ⁻³)	5.62	

Fig. 4(a) and (b) represent SEM micrographs of 0.5 mol% Dy3+ ion doped BaNb2O6 phosphor. The micrograph reveals an inhomogeneous and uneven dense morphology in the micrometer range. The typical crystalline particle size is in the range of 4-6 micrometers in dimension. It is obvious that, micrometer sized crystalline powder would be more suitable to produce an efficient white light for solid-state lighting $applications. ^{34} \\$

4.2 Photoluminescence studies

Fig. 5 illustrates the PLE spectrum of the $Ba_{(1-x)}Nb_2O_6:xDy^{3+}$ (x = 0.5 mol%) phosphor by monitoring the emission wavelength at 574 nm. The PLE spectrum consists of seven sharp peaks due to intra 4f-4f transitions located at 326, 349, 364, 386, 427, 454 and 474 nm, which are attributed from the ground state ${}^{6}H_{15/2}$ to the different excited states (${}^{6}P_{3/2}, {}^{4}M_{17/2}$),

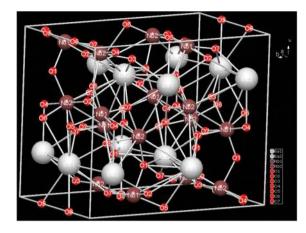


Fig. 3 Crystal structure of BaNb₂O₆ in the b-c plane.

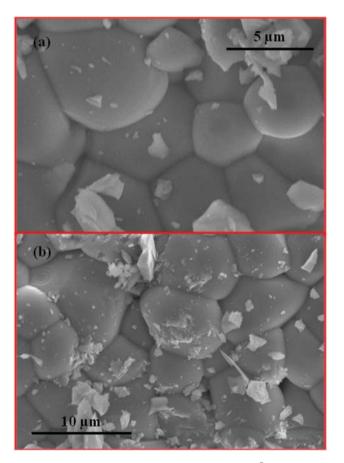


Fig. 4 (a and b) SEM micrographs of $Ba_{(1-x)}Nb_2O_6:xDy^{3+}$ (x = 0.5 mol%) phosphor sintered at 1200 °C.

 $^{6}P_{7/2}$, $(^{4}I_{11/2}, ^{6}P_{5/2})$, $(^{4}F_{7/2}, ^{4}I_{13/2})$, $^{4}G_{11/2}$, $^{4}I_{15/2}$ and $^{4}F_{9/2}$, respectively.35,36

Fig. 6(a) shows the PL emission spectra of $Ba_{(1-x)}Nb_2O_6$: xDy^{3+} (x = 0.01, 0.1, 0.5, 1.0, 1.5, 2.0, 2.5 mol%) phosphors with different doping concentrations of Dy3+ ions at 386 nm

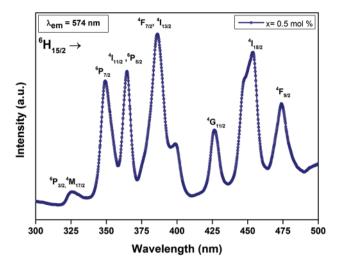


Fig. 5 Photoluminescence excitation spectra (λ_{em} $Ba_{(1-x)}Nb_2O_6:xDy^{3+}$ (x = 0.5 mol%) phosphor.

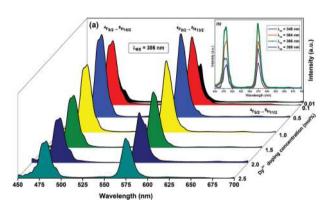


Fig. 6 (a) Emission spectra (λ_{ex} = 386 nm) of BaNb₂O₆:Dy³⁺ phosphor at different mol% concentrations of Dy3+ ions. (b) Emission spectra of $Ba_{(1-x)}Nb_2O_6:xDy^{3+}$ (x = 0.5 mol%) phosphor at different excitations (λ_{ex} = 349, 364, 386, 399 nm) (in the inset).

excitation wavelength. Emission spectra exhibit two intense peaks at 482 and 574 nm and a very weak peak at 664 nm corresponding to the ${}^4F_{9/2} \rightarrow {}^6H_{15/2}, \, {}^4F_{9/2} \rightarrow {}^6H_{13/2},$ and ${}^4F_{9/2} \rightarrow$ $^{6}\mathrm{H}_{11/2}$ transitions, respectively. 35 The $^{4}\mathrm{F}_{9/2} \rightarrow ^{6}\mathrm{H}_{15/2}$ transition belongs to the magnetic dipole allowed and its intensity does not depend on the crystal field of the host.³⁷ On the other hand, the intensity of the hypersensitive transition $(\Delta L = 2; \Delta J = 2)$ $^{4}\text{F}_{9/2} \rightarrow ^{6}\text{H}_{13/2}$ belongs to a forced electric dipole transition, which is allowed in the case where Dy³⁺ ions are located at the local sites with non-inversion center symmetry.³⁸⁻⁴⁰ In Dy³⁺ doped BaNb₂O₆ phosphor, the intensity of yellow emission (${}^4F_{9/2} \rightarrow {}^6H_{13/2}$) is stronger than blue $({}^4F_{9/2} \rightarrow {}^6H_{15/2})$ that confirms the location of the active ions (Dy3+) in low symmetry environment without the inversion centre in the host. As the radius of Dy3+ ions is less than Ba²⁺, Dy³⁺ ions can easily enter into Ba²⁺ sites having low symmetry. This is in good agreement with the results obtained

from XRD analysis.41 Moreover, the emission spectra of the sample Ba_(1-x)Nb₂O₆:xDy³⁺ were measured at 349, 364, 386 and 399 nm excitations as shown in the inset as Fig. 6(b). A similar profile of emission lines has been observed with different intensities for each excitation. The emission intensity at 386 nm excitation is considered as optimum as its intensity is higher than the emission intensity observed for other excitation wavelengths. This may be due to relatively higher absorption at that wavelength.

The branching ratio (β) is a critical parameter to calculate the relative intensities of emission lines originating from the ⁴F_{9/2} excited state. The branching ratios for yellow and blue transitions originating from 4F9/2 were calculated by taking the integral under respective emission bands. For all Dy³⁺ ion concentrations, the sum of the branching ratios for the corresponding emission bands ${}^4F_{9/2} \rightarrow {}^6H_{15/2}$ and ${}^4F_{9/2} \rightarrow {}^6H_{13/2}$ was found to be unity $(\beta_{482} + \beta_{574})$ suggesting that both the transitions have wide possibility of attaining stimulated emission with higher efficiency. 42 Moreover, it could be noticed that the emission band ${}^4F_{9/2} \rightarrow {}^6H_{15/2}$ is broadened and also observed that this transition splits into a maximum number of $J + \frac{1}{2}$ Stark components in the blue emission region (450-500 nm), where *I* is the total angular momentum of electrons. ^{43,44}

As shown in Fig. 6(a), the emission intensity increases initially with an increase in concentration of Dv3+ ions and reaches to a maximum at x = 0.5 mol% and then gradually decreases beyond 0.5 mol% due to the concentration guenching phenomenon. As the doping concentration of Dy³⁺ ions increases, the distance between luminescent centres decreases that increases the possibility of non-radiative energy transfer. 45 The concentration quenching phenomenon resulted mainly by the non-radiative energy transfer among Dy3+ ions. In the present system, the energy transfer mechanism from one Dy3+ ion to another depends on the critical distance between Dy3+ and Dy3+ ions. Hence, it is necessary to calculate the critical distance (R_c) between the adjacent Dy³⁺ ions. According to Blasse^{45–47} the critical distance could be expressed as:

$$R_{\rm c} pprox 2 \left[rac{3V}{4\pi X_{
m c} N}
ight]^{rac{1}{3}}$$

where V is the volume of unit cell, X_c is the critical/optimised concentration (mole) of the activator ions and N is the number of cations per unit cell. By analyzing the experimental data, the values are found to be $V = 988.68 \text{ Å}^3$, N = 8 and $X_c = 0.005$. The calculated critical energy transfer distance is 36 Å for the current system. Van Uitert⁴⁸ has pointed out that energy transfer is generally associated with exchange interaction, radiation re-absorption, or multipolar interactions. The exchange interaction is usually accountable for the energy transfer for the forbidden transition and the critical distance of about 5 Å.47 Since, the distance between adjacent Dy3+ ions is larger than 5 Å, as a result the exchange interaction becomes ineffective and multipolar interaction will become important in this case. According to Dexter's theory, when the doping amount of the activator is large enough, the luminescence intensity I and the

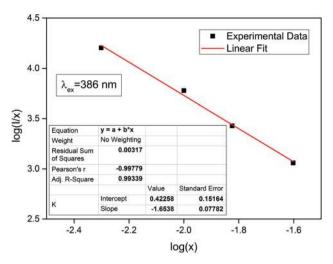


Fig. 7 Relationship of log(I/x) with log(x) in $BaNb_2O_6:Dy^{3+}$ phosphor under 386 nm excitation.

mole fraction of activator ions x could be related as follows:49,50

$$\log(I/x) = -\frac{Q}{3}\log x + A$$

where A is the constant and Q represents interaction type between rare-earth ions. If Q = 6, 8 and 10, the interactions may be corresponding to the electric dipole-dipole (d-d), dipole-quadrupole (d-q) and quadrupole-quadrupole (q-q) interactions, respectively.⁵¹ Depending on the emission spectra of Ba_(1-x)Nb₂O₆:xDy³⁺ excited at 386 nm, the correlation between log(I/x) and log(x) is shown in Fig. 7. The calculated value of Q is 4.95, which is close to 6 discloses that the concentration quenching mechanism in the BaNb₂O₆:Dy³⁺ phosphor occurs due to electric dipole-dipole (d-d) interactions.45,52

The radiative emission process explained is that the radiation excites the Dy3+ ions to the higher excited levels and then quickly relaxes to the ⁴F_{9/2} level by non-radiative and radiative transfers from the ⁴F_{9/2} excited level as shown in the schematic energy level diagram in Fig. 8. The upward and downward arrows indicated in this figure represent excitation and emission, respectively. The possible non-radiative channels explained in Fig. 8 could be: (i) the possible resonant energy transfer (RET): $({}^4F_{9/2} + {}^6H_{15/2} \rightarrow {}^6H_{15/2} + {}^4F_{9/2})$ by considering the energy match rule, and (ii) cross relaxation channels (CRC1, CRC2 and CRC3) among Dy3+ ions are responsible for de-population of ⁴F_{9/2} energy levels by non-radiative such as $({}^{4}F_{9/2} + {}^{6}H_{15/2} \rightarrow {}^{6}F_{11/2}, {}^{6}H_{9/2} + {}^{6}F_{5/2}), ({}^{4}F_{9/2} + {}^{6}H_{15/2} \rightarrow$ $^6F_{9/2}, ^6H_{7/2}+^6F_{5/2})$ and $(^4F_{9/2}+^6H_{15/2}\to ^6F_{1/2}+^6F_{11/2}, ^6H_{9/2})$ for as-prepared phosphors. 41,45,53

Further, the luminescence intensity ratio of yellow to blue (Y/B) is essential for white-light emission. The calculated (Y/B)ratio for the optimized excitation wavelength of 386 nm was found to be close to unity for all doping concentrations. There has been slight variation in the value of the ratio near to unity

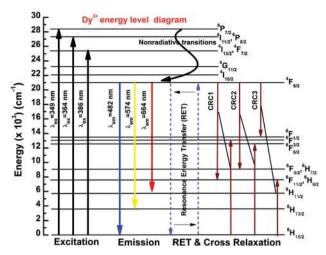


Fig. 8 Partial energy level diagram illustrating excitation, emission and energy transfer mechanisms of Dy3+ ions in BaNb2O6 phosphors.

for other three excitation wavelengths namely, 349, 364 and 399 nm, which confirms excellent stability of the color coordinates against different excitations and concentrations. The intensity ratio being almost constant was attributed to the local environment around Dy3+ ions and is invariant with the varying concentrations of Dy3+ ions.39

4.3 CIE chromaticity coordinates

Fig. 9 shows the CIE chromaticity coordinates for the optimised sample calculated from the emission spectra measured under different excitations. The CIE chromaticity coordinates for optimized phosphor were found to be (0.312, 0.343), (0.319, 0.364), (0.322, 0.339) and (0.312, 0.342) for corresponding excitations at 349, 364, 386, 399 nm, respectively and are indicated in Fig. 9. Excellent white-light chromaticity coordinates (0.322, 0.339) were observed for 386 nm excitation, which are very close to the standard equal energy white-light point (0.333, 0.333). The CIE chromaticity coordinates under

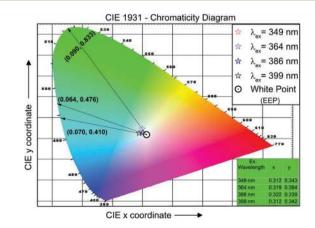


Fig. 9 CIE chromaticity diagram for Dy3+ ion doped BaNb₂O₆ phosphor.

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Table 2 Y/B ratio, CIE chromaticity coordinates and CCT for BaNb₂O₆: Dy3+ phosphors at various doping concentrations

$X (Dy^{3+} concentration)$			()
in mol%)	Y/B ratio	(x, y)	CCT (K)
0.01	1.11	(0.329, 0.339)	5689
0.10	1.06	(0.322, 0.338)	6000
0.50	1.01	(0.322, 0.339)	5907
1.00	1.05	(0.314, 0.341)	6373
1.50	1.06	(0.320, 0.340)	6056
2.00	1.11	(0.331, 0.332)	5623
2.50	1.07	(0.312, 0.319)	6652

optimized excitation (λ_{ex} = 386 nm) for different Dy³⁺ ions concentration are given in Table 2. It is interesting to note that the optimized $Ba_{(1-x)}Nb_2O_6:xDy^{3+}$ (x = 0.5 mol%) phosphor sample exhibited superior white luminescence coordinates compared to different Dy3+ doped phosphor hosts such as CaMoO4:Dy3+ (0.28, 0.27), ⁴¹ $Sr_3Gd(PO_4)_3:Dy^{3+}$ (0.25, 0.29), ⁴⁵ $YPO_4:Dy^{3+}$ (0.36, 0.36)0.42)22 and found to be extremely close to commercial pc-LED (Blue LED + YAG:Ce³⁺) and National Television System Committee (NTSC) white-light emission having (0.32, 0.32) and (0.310, 0.316), respectively.

The calculated CCT values for Dy3+ ion doped BaNb2O6 phosphors were found to vary between 5689 and 6373 K and falls in the cool-white region. The calculated CCT was 5907 K for the optimized concentration (x = 0.5 mol%) of BaNb₂O₆: Dy³⁺ phosphor, which represents cool-white emission and is very close to the "ideal white" region of the chromaticity diagram. The higher value of CCT indicates better visual acuity and greater brightness perception as compared to lower values.⁵⁴ The CCT values lie in the cool white-light region signifying the possibility of the phosphor for application in w-LEDs for outdoor illumination.

Using the chromaticity coordinates given in Fig. 9 for the optimized concentration of as-prepared phosphor, the color purity was calculated and found to be around 7.89×10^{-2} , 5.39×10^{-2} , 5.39×10^{-2} and 7.89×10^{-2} corresponding to 349, 364, 386, and 399 nm excitation wavelengths, respectively. The low value of the color purity indicates the purity for white-light emission. ^{27,53} The above-mentioned results indicate that the as-prepared phosphor can be considered as a potential candidate for fabrication of w-LEDs based on NUV chips as the excitation source.

4.4 Luminescence decay curve analysis

The room temperature luminescence decay curve has been plotted for $Ba_{(1-x)}Nb_2O_6:xDy^{3+}$ (x = 0.5 mol%) phosphor and is shown in Fig. 10. It represents the decay curve measured for ${}^{4}\mathrm{F}_{9/2} \rightarrow {}^{6}\mathrm{H}_{13/2}$ emission for phosphor when excited under 386 nm wavelength. To understand the behaviour of luminescent decay, the decay curve was fitted with different equations and the best fit was observed for the bi-exponential equation:55,56

$$I = A_1 e^{-t/\tau_1} + A_2 e^{-t/\tau_2}$$

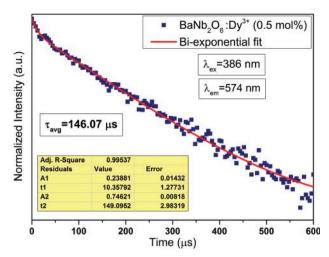


Fig. 10 Luminescence decay curve ($\lambda_{em} = 574$ nm) of $Ba_{(1-x)}Nb_2O_6$: xDy^{3+} (x = 0.5 mol%) phosphor under 386 nm excitation.

where I is the luminescence intensity; t is the time; τ_1 and τ_2 are the decay times for the exponential component and A_1 and A_2 are the fitting parameter constants, respectively. Thus, the average lifetime in the case of bi-exponential fitting can be determined by using the equation:⁴¹

$$au_{\text{avg}} = rac{A_1 { au_1}^2 + A_2 { au_2}^2}{A_1 { au_1} + A_2 { au_2}}$$

The fluorescent lifetime au_{avg} for the $^4F_{9/2}$ level for the optimized phosphor sample was found to be ~146.07 μs. Generally, the PL decay curves can be influenced by energy transfer between Dy³⁺ ions. If there is no interaction between the rareearth ions, the decay curves are usually fitted to a single exponential function. The bi-exponential fitting behaviour shows the possibility of interaction between Dy³⁺ ions in the BaNb₂O₆ lattice as discussed in the previous sections.

Conclusions

Single-phase BaNb₂O₆:Dy³⁺ white-light emitting phosphors were successfully synthesized using the solid-state reaction method. The crystallinity and pure phase of the as-prepared phosphors were examined by XRD and Rietveld refinement studies. All the prepared samples exhibited single-phase with an orthorhombic structure. The excitation spectra indicate that the phosphors could be effectively excited by NUV LED chips having an excitation wavelength of 386 nm. In order to determine the optimized doping concentration of BaNb₂O₆:Dy³⁺ phosphors, concentration dependent luminescence measurements were carried out. 0.5 mol% of Dy3+ has been found to be the optimum doping concentration under different excitation wavelengths. The blue (482 nm) and yellow (574 nm) emission bands corresponding to ${}^{4}F_{9/2} \rightarrow {}^{6}H_{I}$ (J = 15/2, 13/2) transitions and the value of Y/B ratio close to unity have been successfully

achieved. The combination of these emission bands emits white-light and the CIE chromaticity coordinates for the optimized phosphor are (x = 0.322, y = 0.339) with a CCT value of 5907 K, which are close to the standard white-lamp colorimetric point in the cool white region. All the above-mentioned results indicate that the Dy³⁺ ion doped BaNb₂O₆ phosphor could be used as a practical potential luminescent material for NUV based w-LED applications.

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