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### **Supplemental Material**

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# Engineering high-k/SiGe interface with ALD oxide for selective GeO<sub>x</sub> reduction

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#### Abstract

Suppression of electronic defects induced by GeO<sub>x</sub> at the high-k gate oxide/SiGe interface is critical for implementation of high mobility SiGe channels in CMOS technology. Theoretical and experimental studies have shown that a low defect density interface can be formed with an SiO<sub>x</sub>rich interlayer on SiGe. Experimental studies in literature indicates better interface formation with Al<sub>2</sub>O<sub>3</sub> in contrast to HfO<sub>2</sub> on SiGe however the mechanism behind this is not well understood. In this study, the mechanism of forming a low defect density interface between Al<sub>2</sub>O<sub>3</sub>/SiGe is investigated using atomic layer deposited (ALD) Al<sub>2</sub>O<sub>3</sub> insertion into or on top of ALD HfO<sub>2</sub> gate oxides. To elucidate the mechanism, correlations are made between the defect density determined by impedance measurements and the chemical and physical structure of the interface determined by high resolution scanning transmission electron microscopy and electron energy loss spectroscopy (STEM – EELS). Compositional analysis reveals an  $SiO_x$  rich interlayer for both Al<sub>2</sub>O<sub>3</sub>/SiGe and HfO<sub>2</sub>/SiGe interfaces with insertion of Al<sub>2</sub>O<sub>3</sub> into or on top of the HfO<sub>2</sub> oxide. The data is consistent with the Al<sub>2</sub>O<sub>3</sub> insertion inducing decomposition of the GeO<sub>x</sub> from the interface to form an electrically passive, SiO<sub>x</sub> rich interface on SiGe. This mechanism shows that nanolaminate gate oxide chemistry cannot be interpreted as resulting from a simple layer by layer ideal ALD process because the precursor or its reaction products can diffuse though the oxide during growth and react at the semiconductor interface. This result shows that in scaled CMOS, remote oxide ALD (oxide ALD on top of the gate oxide) can be used to suppress electronic defects at gate-oxide semiconductor interfaces by oxygen scavenging.

#### Introduction

SiGe alloys are employed as stressor layers in mainstream complementary metal–oxide– semiconductor (CMOS) transistors and are being investigated as p-type field effect transistor (FET) channels due to their high mobility<sup>1</sup> and ease of integration into CMOS<sup>2</sup>. Thermally stable HfO<sub>2</sub> gate oxides with high dielectric constants reduce CMOS device power consumption<sup>3-4</sup>. SiGe p-FETs with high-k gate dielectrics which have low defect interfaces can provide better electrostatic control of the channel and higher drive current for low gate bias voltage. Conversely, a high density of interface defects between the high k gate oxide and the SiGe channel degrades device performance metrics such as subthreshold slope and reduces the on/off current ratio<sup>5</sup>. The main challenge for implementing SiGe FETs is the binary atom termination (Si-Ge) of the surface which results in formation of SiGeO<sub>x</sub> mixed oxides and dangling bonds on both Si and Ge atoms<sup>6-8</sup>. GeO<sub>x</sub> and associated dangling bonds are the main sources of defects producing interface trapped charge (D<sub>it</sub>), while SiO<sub>x</sub> is a stable oxide that forms a nearly defect-free interface according to theoretical calculations<sup>9</sup>.

Previously, several techniques such as nitride and sulfur passivation on Si<sub>0.7</sub>Ge<sub>0.3</sub>(001) were studied with Al<sub>2</sub>O<sub>3</sub> gate oxides and reduction in the interface defect density via suppression of GeO<sub>x</sub> formation was reported <sup>10-11</sup>. However similar low defect density interfaces could not be established with HfO<sub>2</sub> gate oxide. This is because oxygen containing species such as excess H<sub>2</sub>O, OH, and/or O can diffuse through HfO<sub>2</sub> during atomic layer deposition (ALD), forming GeO<sub>x</sub> defects on the SiGe surface; in addition, the nature of HfO<sub>2</sub> allows diffusion of Ge and GeO<sub>x</sub> to the top surface of the oxide as a result of reaction with HfO<sub>2</sub> and GeO<sub>2</sub> decomposition<sup>11-14</sup> Recently, HfO<sub>2</sub>/SiGe interfaces formed with Al<sub>2</sub>O<sub>3</sub>-HfO<sub>2</sub> nanolaminate gate dielectric stacks were found to have a low interface state density, and it was hypothesized that the mechanism was

reduction of GeO<sub>x</sub> out-diffusion during ALD<sup>15</sup>. Theoretical DFT models of the amorphous  $HfO_2/Si_{0.5}Ge_{0.5}(001)$  interface have shown that low-defect interfaces may be formed even before hydrogen passivation with short anneals (<10 ps) when the interface is comprised solely of SiO (silicon monoxide)<sup>16</sup>. Experimental studies have shown a 10× D<sub>it</sub> reduction at the Al<sub>2</sub>O<sub>3</sub>/SiGe and  $HfO_2/SiGe$  interfaces via selective oxygen scavenging by using an oxygen-scavenging, metallic Al gate<sup>17</sup>; the selectivity is due to the difference in formation enthalpy of GeO<sub>x</sub> compare to SiO<sub>x</sub> facilitating transfer of oxygen from GeO<sub>x</sub> to Al.<sup>6, 17</sup> However, this process induced thicker gate oxides and reduced the maximum capacitance density (C<sub>max</sub>) resulting from Al oxidation on top of ALD grown gate oxide.

In the present study, formation of low interface defect density HfO<sub>2</sub>/SiGe gate stacks using inorganic and organometallic based ALD Al<sub>2</sub>O<sub>3</sub> insertion in bilayers and nanolaminates of Al<sub>2</sub>O<sub>3</sub> and HfO<sub>2</sub> was investigated with STEM-EELS analysis. It is found that insertion of the highly oxygen reactive trimethyl aluminum (TMA) ALD precursor for Al<sub>2</sub>O<sub>3</sub> in HfO<sub>2</sub> containing gate stacks reduces defects consistent with remote selective oxygen scavenging from the interface. This new selective oxygen scavenging technique is most effective when Al<sub>2</sub>O<sub>3</sub> layers are uniformly distributed across the HfO<sub>2</sub> in a nanolaminate (NL) structure, but it also is effective when the Al<sub>2</sub>O<sub>3</sub> ALD deposition occurs on top of the HfO<sub>2</sub>. In Al<sub>2</sub>O<sub>3</sub> ALD, during each TMA half cycle, the TMA is available to reduce additional species either by diffusion through HfO<sub>2</sub> or remotely at the growth surface. This is consistent with GeO<sub>x</sub> out diffusing to the top of HfO<sub>2</sub> gate oxides<sup>11, 13, 18</sup>, and the TMA remotely reducing the GeO<sub>x</sub> during each half cycles<sup>13</sup>. It is found that using the Al<sub>2</sub>O<sub>3</sub> ALD prior to HfO<sub>2</sub> deposition is not as effective as using the Al<sub>2</sub>O<sub>3</sub> ALD in the nanolaminate, consistent with the suggestion that gettering is most effective after deposition of sufficient oxide to act as an

 $H_2O$  barrier to reduce additional formation of GeO<sub>x</sub> during water-based ALD. This hypothesis of the gate oxide acting as a  $H_2O$  barrier and TMA being a GeO<sub>x</sub> reducing agent was supported by experiments with ALD of purely  $Al_2O_3$  gates, because  $Al_2O_3$  is a better diffusion barrier than  $HfO_2^{13}$ . Oxygen vacancies in  $HfO_2$  is well documented and prevention of Ge out diffusion with  $Al_2O_3$  incorporation into gate oxide was shown<sup>11, 13</sup>. Once the  $Al_2O_3$  gate oxide reaches a critical thickness, an additional 25% increase in oxide thickness results in a 4× reduction in  $D_{it}$  and nearly complete elimination of the SiGeO<sub>x</sub> interfacial layer. For  $HfO_2$ , the selective scavenging process benefits from the difference in formation enthalpy of SiO<sub>x</sub> in comparison to GeO<sub>x</sub> and reduces the interface trapped charge density by forming Si-rich SiO<sub>x</sub> at the interface, consistent with the predictions of the DFT models<sup>16</sup>. This is also consistent with the known ability of TMA to reduce low enthalpy of formation oxides on substrates at the start of ALD, a process known as ALD cleanup<sup>19-20</sup>.

#### Methods

Metal oxide semiconductor capacitor (MOSCAP) devices were fabricated with HfO<sub>2</sub> and Al<sub>2</sub>O<sub>3</sub> oxides onto Si<sub>0.7</sub>Ge<sub>0.3</sub>(100) epitaxially grown on p-type Si (100) (Applied Materials). Samples were degreased by sonication in methanol for one minute followed by 10 second rinses in acetone, IPA, and deionized H<sub>2</sub>O. Native oxides were removed by cycling 2.5 times through 1 min 2% HF solution and 1 min DI H<sub>2</sub>O, ending with 2% HF. The samples were dried in N<sub>2</sub>, passivated in an ammonium sulfide solution (25% (NH4)<sub>2</sub>S) for 15 minutes, rinsed with water for 30 seconds, and dried with N<sub>2</sub>. Subsequently, samples were transferred to the atomic layer deposition (ALD) chamber and oxide structures were grown by using a Beneq TFS200 cross flow reactor at 275°C using tetrakis dimethyl amino hafnium (TDMAH), HfCl<sub>4</sub> and TMA metal precursors and water as an oxidant. Each HfO<sub>2</sub> cycle consisted of 1s of a TDMAH or HfCl<sub>4</sub> pulse and a 500ms of H<sub>2</sub>O

pulse, each Al<sub>2</sub>O<sub>3</sub> cycle consisted of a 1s TMA pulse and a 500ms H<sub>2</sub>O pulse. Ar was used as carrier gas for all processes, and 6s purges were employed between each pulse. A set of samples were fabricated in bilayer, tri-layer and nanolaminate (NL) structures formed by Al<sub>2</sub>O<sub>3</sub> and HfO<sub>2</sub> in gate stack for MOSCAPs as shown in Fig. S1 and the insets of Fig. 1. Nickel gates (50nm thick, 150um diameter) were deposited with a shadow mask onto the oxide surface using a Denton 502A thermal evaporator in vacuum  $< 2 \times 10^{-6}$ . Al back contacts were deposited by sputtering after native oxide removal with Ar plasma at 100W, 5mTorr in Denton Discovery 635 sputtering system. Samples were annealed using an optimized recipe for 30 minutes total (10 min at 300°C, 10 min at 330°C, and 10 min at 350°C) in forming gas (5% H<sub>2</sub>, 95%N<sub>2</sub>) in a Ulvac MILA-3000 Minilamp annealer at 3 slpm at 1atm. MOSCAP characterization at 300 K was performed with Keysight B1500. Multi-frequency capacitance-voltage (C-V) and conductance-voltage (G-V) measurements were obtained from 2 kHz to 1MHz with 30mV AC signal superimposed on DC gate bias varied from 2V to -2V. For the structural and compositional analysis, TEM specimens (<50 nm) were prepared from MOSCAP devices using a FEI-Scios Ga focus ion beam (FIB). STEM/TEM analysis on high-k/SiGe gate stacks were carried out with FEI Titan 80-300, FEI Metrios and JEOL ARM 200CF. STEM High-angle annular dark field (HAADF) and bright field (BF) along with TEM images were obtained. Compositional analysis was performed with electron energy loss spectroscopy (EELS) both at 80keV and 200keV with JEOL ARM 200CF equipped with Gatan Quantum EELS spectrometer.

#### **RESULTS AND DISCUSSION**

Electrical characterization of the MOSCAP devices with various gate oxide structures after forming gas annealing (FGA) are shown in Fig. 1. The C-V analysis for 45 cycles (~4.5nm) of HfO<sub>2</sub> control sample with a  $C_{max}$  of 2.5 uF/cm<sup>2</sup> and corresponding G-V characterization are

presented in Fig.1a and 1f; a peak defect density of ( $D_{it}$ ) 3.74×10<sup>12</sup> eV<sup>-1</sup>cm<sup>-2</sup> (see Figure 2) is obtained. The full interface state  $model^{21-22}$  is used to calculate the energy distribution of the  $D_{it}$  by fitting multi frequency C-V and G-V curves for each bias point from inversion (2V) to accumulation (-1V). The result peak D<sub>it</sub> values from the full interface state model are compared with the D<sub>it</sub> values from conductance method<sup>23</sup> and shown to be in agreement as documented in table 1. Fig. 1b and 1c show the C-V for 5 cycle Al<sub>2</sub>O<sub>3</sub> insertion below and above 45 cycles of HfO<sub>2</sub>. As expected, Al<sub>2</sub>O<sub>3</sub> insertion decreases C<sub>max</sub> in both cases due to an increase in total oxide thickness and the lower dielectric constant of Al<sub>2</sub>O<sub>3</sub> in comparison to HfO<sub>2</sub>. However, the magnitude of the dispersive, depletion capacitance D<sub>it</sub> feature also decreases as shown in Fig. 1b and 1c. Al<sub>2</sub>O<sub>3</sub> insertion below or above HfO<sub>2</sub> has a nearly identical effect on interface defect density corresponding to peak  $D_{it}$  values of  $3.30 \times 10^{12}$  $eV^{-1}cm^{-2}$  and  $3.15 \times 10^{12} eV^{-1}cm^{-2}$ . Note full  $D_{it}$  distributions as a function of the thickness are shown below confirming the trends. Because interface trap response involves defects at the semiconductor oxide interface, a change in D<sub>it</sub> resulting from the addition of 5 ALD cycles of  $Al_2O_3$  on top of a 4.5 nm thick HfO<sub>2</sub> layer is unexpected. Moreover, as shown in Fig 1d,  $Al_2O_3$ insertion both above and below HfO2 in an Al2O3/HfO2/Al2O3 tri-layer structure further decreases the interface trapped charge density by 10% to  $2.53 \times 10^{12}$  eV<sup>-1</sup>cm<sup>-2</sup>. Furthermore, as shown in Fig 1e, when Al<sub>2</sub>O<sub>3</sub> layers dispersed across the HfO<sub>2</sub> in the nanolaminate structure (NL), the D<sub>it</sub> decreases further down to  $2.22 \times 10^{12} \text{ eV}^{-1} \text{ cm}^{-2}$ , 12% lower than the tri-layer with a small increase in C<sub>max</sub>. The C<sub>max</sub> reduction due to an increase in total oxide thickness is expected, but D<sub>it</sub> reduction with more Al<sub>2</sub>O<sub>3</sub> incorporation remote from the interface is not.

To better document the effects of bottom vs top  $Al_2O_3$ , samples were growth with 10 cycles of  $Al_2O_3$  inserted either below the  $HfO_2$  or above the  $HfO_2$ . This caused very large decreases in  $D_{it}$  as shown in Fig. 2. Again, the deposition of  $Al_2O_3$ , on top of the  $HfO_2$  induced a 57% decrease

in the peak  $D_{it}$  as well as a 54% decrease in the integrated  $D_{it}$  while the Al<sub>2</sub>O<sub>3</sub> inserted below the HfO<sub>2</sub> only induced a 40% decrease in the peak  $D_{it}$  as well as a 40% decrease in the integrated  $D_{it}$ . The integrated  $D_{it}$  are obtained by summation of the defect densities across the band gap in Fig.4; this is equivalent to integration of the areas under the  $D_{it}$  curves; these integrated  $D_{it}$  are a measure of the total defects induced in the band gaps. urves provides total defects induced in band gap.

For the 5 cycle  $Al_2O_3$  insertion, although absolute defect density quantification is challenging, small decreases in the magnitude of the interface trap-related depletion capacitance feature while maintaining an essentially constant C<sub>max</sub> suggests a lower defect density at the interface oxide/semiconductor interface for the nanolaminate structure compared to tri-layer sample or the bilayer samples. For each processing condition, 5-10 devices were studied. The 2-3 devices with the most consistent C-V were chosen for further analysis. Therefore, although absolute D<sub>it</sub> calculations are accurate only within 30%; even 10% changes in D<sub>it</sub> decay with Al<sub>2</sub>O<sub>3</sub> insertion are reliable as confirmed via fabrication multiple sample sets as shown in Fig S8. The D<sub>it</sub> standard errors of the mean (in /cm<sup>2</sup>-eV and percentage) for samples were calculated as shown in table 1. The typical standard error is 3.9%, and changes in D<sub>it</sub> of 10% are significant in this comparison. The  $D_{it}$  standard error analysis shown documents that changes in  $D_{it}\ of\ 10\%$  can reliably be determined. Error bars are not given for each data point since the absolute D<sub>it</sub> is less precisely determined in part because the uncertainty in Cox is about 10% which can translate into 10-20% uncertainly in absolute D<sub>it</sub><sup>24</sup>. Additional details for standard error analysis are shown in Table 1. As shown in Table 1, there is no significant change for V<sub>fb</sub> as a function of processing conditions. In addition, there is very little dependence of V<sub>fb</sub> on frequency; however, when there is frequency dependence, the highest frequency C-V is employed to access V<sub>fb</sub> since it is the least affected by traps.

To investigate the mechanism of D<sub>it</sub> reduction due to insertion of Al<sub>2</sub>O<sub>3</sub> ALD layers, MOSCAP devices with only Al<sub>2</sub>O<sub>3</sub> gate oxides of varying thickness were prepared, as shown in

Fig. 3. As expected, C-V and G-V measurements from Al<sub>2</sub>O<sub>3</sub> devices after FGA as a function of deposition cycles or oxide thickness show a decrease in  $C_{max}$  from 1.3uF/cm<sup>2</sup> to 1.0 uF/cm<sup>2</sup> by increasing the oxide thickness. The depletion capacitance D<sub>it</sub> feature also decreases in amplitude. Once the ALD-grown oxide thickness increases above a critical value, there is a super-linear decrease in D<sub>it</sub>. For an increase of 25% in the number ALD cycles (40 to 50 cycles), the D<sub>it</sub> decreases by  $> 4 \times$  from 2.93  $\times 10^{12}$  to 0.67  $\times 10^{12}$  eV<sup>-1</sup>cm<sup>-2</sup>. The lower D<sub>it</sub> from increasing the number of ALD cycle from 40 to 50 is not the result of a sudden increase in oxide thickness due to an ALD onset delay. The C<sub>max</sub> for 50 cycles is just 25% lower than for 40 cycles consistent with nearly linear growth. Because the D<sub>it</sub> originates from defects at the interface with the semiconductor, D<sub>it</sub> reduction by growing additional Al<sub>2</sub>O<sub>3</sub> layers is consistent with chemical modification of the interface by exposure to the TMA-based ALD environment. The C-V curves obtained from MOS capacitors containing Al<sub>2</sub>O<sub>3</sub> before FGA shown in Fig. S4 exhibit a similar trend. It is noted that for Si interfaces, Al<sub>2</sub>O<sub>3</sub> ALD is reported to reduce defects by a hydrogen passivation mechanism. However, for SiGe, the most important dangling bonds are on Ge which are not readily passivated by hydrogen as shown by the lowest D<sub>it</sub> devices being ones which eliminate interfacial GeO<sub>x</sub> using a gettering gate. This is also consistent with previous reports of extremely low Dit for HfO2/SiGe by deposition of Al gettering metal on top of HfO<sub>2</sub> in which the FGA was performed after Al deposition<sup>17</sup>.

The suppression of the interface defects with  $Al_2O_3$  insertion into gate oxide influences is not limited to a single energy but is distributed across the band gap as shown in the  $D_{it}$  energy distributions extracted from measured C-V and conductance-voltage (G-V) data using the full interface state model<sup>21</sup> as shown in Fig. 4. Note that, interface defect distribution across the band gap for all the devices in this work is summarized in fig. 4 side by side for better comparison. In addition, the insertion of 1, 3, 5 and 10 monolayers of  $Al_2O_3$  below HfO<sub>2</sub> (Fig. S2) and on top of HfO<sub>2</sub> (Fig. S3) was studied via C-V and G-V and indicate similar trends. Trapped charge energy distributions of the corresponding capacitors shown in Fig. 4a and b are consistent with Fig. 1 and show a decrease of  $D_{it}$  with incorporation of  $Al_2O_3$  ALD layers but below and above HfO<sub>2</sub> layers. Inset  $D_{it}$  values in Fig. 4 c and d denote the integrated defect density across the band gap and are in good agreement with the behavior observed for the peak  $D_{it}$ . As shown in Fig 2 and 4, the shapes of the  $D_{it}$  distributions for HfO<sub>2</sub>/SiGe and  $Al_2O_3$ /SiGe are similar (not identical), but the  $D_{it}$  values vary more than 2x; this is consistent with the source of  $D_{it}$  being similar for both oxides (for example GeO<sub>x</sub>) but present in greater concentration for HfO<sub>2</sub>/SiGe than Al<sub>2</sub>O<sub>3</sub>/SiGe.

In Al<sub>2</sub>O<sub>3</sub> ALD, during each TMA half cycle, the TMA is dosed in excess; therefore, after the surface hydroxyl groups are eliminated, the TMA is available to reduce additional species. However, during each H<sub>2</sub>O half cycle, the H<sub>2</sub>O is present in excess at the end of the pulse, and some H<sub>2</sub>O or H<sub>2</sub>O-derived species may diffuse through thin Al<sub>2</sub>O<sub>3</sub> gate stack. Also, although the  $Al_2O_3$  is better diffusion barrier in comparison to the HfO<sub>2</sub>, it has been reported previously that  $GeO_x$  can diffuse through thin  $Al_2O_3$  gate stacks<sup>10</sup>. The  $Al_2O_3$  – SiGe interface is known to include  $GeO_x$  species which are source of interface defects<sup>6</sup>. It is hypothesized that above a critical thickness, the  $Al_2O_3$  becomes a good barrier to  $H_2O_3$ ; therefore, excess TMA exposure during ALD can scavenge oxygen from GeO<sub>x</sub> species and convert it into Ge which might be redeposited on SiGe without formation of new GeO<sub>x</sub> from the water pulses.<sup>6</sup> This mechanism is consistent with decreasing D<sub>it</sub> by insertion of Al<sub>2</sub>O<sub>3</sub> into HfO<sub>2</sub>-containing gate dielectric stacks or deposition of  $Al_2O_3$  on top of HfO<sub>2</sub> gate dielectrics. However, as noted earlier,  $Al_2O_3$  is a more effective diffusion barrier in comparison to HfO<sub>2</sub> for oxygen containing species; consequently, the D<sub>it</sub> suppression with HfO<sub>2</sub> dielectrics is not expected to be as great as with Al<sub>2</sub>O<sub>3</sub> dielectrics of similar thickness<sup>13</sup>. This suggests that the number of  $Al_2O_3$  layers incorporated into oxide is mainly responsible for the interface defect density reduction observed in these experiments. Similar to  $6 \times (9 \text{HfO}_2 + 1 \text{Al}_2 \text{O}_3)$  NL gate dielectric structure, gate oxide without oxidant during Al<sub>2</sub>O<sub>3</sub> ALD is performed by only TMA dosing and D<sub>it</sub> reduction along with C<sub>max</sub> observed indicating Al<sub>2</sub>O<sub>3</sub> layer formation by oxygen scavenging." It is noted that for HfO<sub>2</sub> based gate stacks with incorporated  $Al_2O_3$  layers the TMA may diffuse to the interface to directly reduce the GeO<sub>x</sub> to Ge since HfO<sub>2</sub> is a poor diffusion barrier; however, the mechanism likely coexists with the remote scavenging mechanism.

Fig. 5 shows STEM high-angle annular dark-field (HAADF), bright field (BF) and TEM images of the HfO<sub>2</sub>-only, HfO<sub>2</sub> - Al<sub>2</sub>O<sub>3</sub> bilayers and nanolaminate gate stacks. White and black arrows in HAADF and BF images indicate the estimated oxide interfacial layer (IL) thickness. The SiGe-HfO<sub>2</sub> interface in Fig. 5a shows a 0.8 nm IL. Insertion of Al<sub>2</sub>O<sub>3</sub> between SiGe and HfO<sub>2</sub> increases the apparent IL thickness (Fig. 5b) which is expected because the interface now consists of both Al<sub>2</sub>O<sub>3</sub> and SiGeO<sub>x</sub>. EELS or another spatially-resolved composition profiling method is needed to differentiate between these layers because of the similar atomic mass contrast of these oxides in HAADF-mode imaging. Al<sub>2</sub>O<sub>3</sub> grown on HfO<sub>2</sub> and the Al<sub>2</sub>O<sub>3</sub> - HfO<sub>2</sub> nano-laminate have thickness similar that shown in Fig 5.

Elemental profiles across the gate stacks were investigated with STEM-EELS analysis as shown in Fig. 6. A multiple linear least square (MLLS) fitting procedure<sup>25</sup> is used to resolve spectroscopic feature overlay issues, especially for Al, Hf and Si. The red dash line intercepts the half max of oxygen peak and is employed to indicate the SiGe surface. Black and green arrow denote Si and Ge composition of the SiGe surface. It should be noted that the determination of interface location in STEM-EELS is subject to interpretation especially due to surface roughness. However even with the surface roughness, the IL should be within few angstroms of the dashed lines, and this metric was confirmed by determining the location of the Si/SiGe interface since these layers are epitaxially grown. Any uncertainty in the exact location of the interface will not impact the trends in the elemental composition across the oxide which indicating suppression of Ge out diffusion with Al<sub>2</sub>O<sub>3</sub> incorporation. Fig. 6a and 6b show 200keV EELS analysis of the same structures fabricated with hafnium tetrachloride at 300°C and tetradimethylamido hafnium (IV) ALD precursors at 275°C; no significant differences were observed as seen in Fig. 6 a and b consistent with interface defects obtained with multifrequency C-V analysis. Both the HfCl<sub>4</sub> and TDMAH gate stacks have nearly identical elemental distribution across the device and have a Ge tail extending about 2 nm into the HfO<sub>2</sub> layer in contrast to the capacitors in Fig. 6c-f which show a diminished or zero Ge tail. This may be attributed to the Al<sub>2</sub>O<sub>3</sub> layer impeding Ge out-diffusion<sup>11,</sup> <sup>13</sup>. It is surprising that Al<sub>2</sub>O<sub>3</sub> reduces Ge out diffusion even when Al<sub>2</sub>O<sub>3</sub> is deposited onto HfO<sub>2</sub> (Fig 6c - 6f). This may be result of TMA reaction products (either AlO<sub>x</sub> or AlC<sub>x</sub>H<sub>y</sub>) diffusion into HfO<sub>2</sub> reaching IL.

Further documentation that even remote Al<sub>2</sub>O<sub>3</sub> ALD can control the GeO<sub>x</sub> in the interlayer is observed in the composition of the interlayers. For all the gate stacks with Al<sub>2</sub>O<sub>3</sub>, the amount of Ge and the ratios of Si to Ge at the interface (between full max and 1/2 height of the oxygen peak) are greater than the control HfO<sub>2</sub>/SiGe (Fig 6a); exact comparison of the interfacial Si/Ge ratio between the samples with Al<sub>2</sub>O<sub>3</sub> (Fig 6b-f) is challenging since the ratio can vary with slight adjustment of the nominal interface position. All the samples have an interlayer above the SiGe as shown by the gap between the Hf and O edges; however, the Si signal falls off most steeply for Si on the nanolaminate samples (Fig S5-7) consistent with these samples having the most abrupt interface.

STEM-EELS at both 80 keV and 200keV is employed to investigate the AlO<sub>x</sub> distribution in the ALD HfO<sub>2</sub> layer, while controlling beam induced Al damage due to high energy electrons<sup>26</sup>. For the HfO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub>/SiGe structures (Figs 6c and 6d), the 200 keV spectrum shows complete diffusion of the Al while the 80 keV spectrum shows only a small retention of Al at the interfaces likely due to the limited signal to noise of the 80 keV spectra; the data is consistent AlO<sub>x</sub> diffusing during HfO<sub>2</sub> deposition. For the Al<sub>2</sub>O<sub>3</sub>/HfO<sub>2</sub>/SiGe structures (Figs 6e and 6f), the 200 keV spectrum shows two Al peaks, one above and one below HfO<sub>2</sub> while the 80 keV spectrum shows a just small retention of AlO<sub>x</sub> above the HfO<sub>2</sub> layer, consistent with the limited signal to noise of the 80 keV spectra. Overall both are consistent with diffusion of TMA or its reaction products into the  $HfO_2$  when the  $Al_2O_3$  ALD is performed after  $HfO_2$  ALD. The nanolaminate gate stack (Fig.1e) composition studied with 80 keV EELS shows dispersed Al (Fig. S5).

This unexpected inter diffusion of Al<sub>2</sub>O<sub>3</sub> is also shown with EELS raw data in 3D semilog graph in Fig. 7 (Ni/Al<sub>2</sub>O<sub>3</sub>/HfO<sub>2</sub> /SiGe) and Fig S6 (Ni/HfO<sub>2</sub> / Al<sub>2</sub>O<sub>3</sub>/Ni). The oxide compositional profile can be seen from the electron energy loss peaks starting after element specific edges; for example, the Si K edge is at 1839 eV (orange arrow) and the Ge L edges at 1217 eV (pink arrow). The blue arrow indicates the SiGe/HfO<sub>2</sub> interface region. It is seen that the Ge peak decays earlier than the Si peak as a function of distance from the SiGe surface consistent with a SiO<sub>x</sub> rich region at the interface. The black arrow indicates energy loss due to Al (K edge 1560 eV) across the oxide. The peak intensity is significantly lower in comparison to other elements; however, it can be seen in several regions in the HfO<sub>2</sub> and reaches a maximum close to the  $HfO_2/Ni$  interface since it is deposited on  $HfO_2$  as shown in 3D data and the region of interest (RIO) in Fig 7 a-b. However, the Al peak can be also seen close to SiGe/HfO<sub>2</sub> interface, almost 4 nm away where it was deposited. The existence of Al signal in this region is consistent with AlO<sub>x</sub> diffusion through HfO<sub>2</sub>. A similar raw data analysis performed for Ni/HfO<sub>2</sub> /Al<sub>2</sub>O<sub>3</sub> /SiGe bilayer device also show interdiffusion (Fig S6). However, comparison of EELS analysis for both bilayer structure indicates enhanced Al<sub>2</sub>O<sub>3</sub>-HfO<sub>2</sub> interdiffusion when Al<sub>2</sub>O<sub>3</sub> is deposited prior to HfO<sub>2</sub> ALD.

Fig. 8 shows simultaneously acquired HAADF and BF-STEM images of Ni/Al<sub>2</sub>O<sub>3</sub>. /SiGe/Si (a-b) and Al/HfO<sub>2</sub>/SiGe/Si (c-d) MOSCAP devices along with corresponding C-V graphs (e-f). In contrast to the Ni/HfO<sub>2</sub>/SiGe/Si gate stack in Fig. 5, Al/HfO<sub>2</sub>/SiGe/Si has a gettering gate which is known to remove oxygen from the interfaces and reduces IL thickness; this is supported by EELS analysis (Fig. S6) indicating similar O and Hf peak decay profiles at the SiGe interface unlike with Ni gates which show offsets between O and Hf peak at the interface (Fig. 6a-f)<sup>17</sup>.

The gate stacks in Fig. 8e and f have two very dissimilar oxides and show very low depletion capacitance feature resulting from charging/discharging of interface traps. Similar integrated D<sub>it</sub> values are displayed in the inset. However, these two very different gate oxide structures deposited on SiGe show similar, almost abrupt interfaces with the underlying SiGe (Fig S6 and Fig 8g). For Al/HfO<sub>2</sub>/SiGe capacitors, it was previously shown that an Al metal gate can scavenge oxygen from the interfacial layer and reduce the interface defect density while also thinning the IL<sup>17</sup>. For the case of the Ni/Al<sub>2</sub>O<sub>3</sub>/SiGe device, it is most likely that introduction Al<sub>2</sub>O<sub>3</sub> ALD in the gate stack fabrication process provides an effect similar to that of an Al gate and scavenges oxygen from oxide/SiGe interface. The mechanisms are similar because TMA is a highly reactive precursor with oxygen which can interact with nearly all oxygen-containing molecules to form Al<sub>2</sub>O<sub>3</sub>. Insertion of each additional Al<sub>2</sub>O<sub>3</sub> layer using TMA precursor can scavenge excess oxygen from the gate oxide or the high-k/SiGe interface. Similar to 6×(9HfO<sub>2</sub>+1Al<sub>2</sub>O<sub>3</sub>) NL gate dielectric structure, gate oxide without oxidant during Al<sub>2</sub>O<sub>3</sub> ALD is performed by only TMA dosing and D<sub>it</sub> reduction along with C<sub>max</sub> observed indicating Al<sub>2</sub>O<sub>3</sub> layer formation by oxygen scavenging. It is hypothesized that TMA exposure scavenges weakly bound oxygen from the interface either by diffusing into the interface as TMA or TMA reaction products (for example monomethyl aluminum) or it decomposes the GeO<sub>x</sub> remotely, producing suboxide species that diffuse readily through even thin Al<sub>2</sub>O<sub>3</sub> (remote gettering)<sup>10</sup>. As TMA interacts and scavenges oxygen from the interface, it is likely that GeO<sub>x</sub> dissociates and donates oxygen to TMA due to the lower Gibbs free energy of formation of GeO<sub>x</sub> in comparison to SiO<sub>x</sub>.<sup>27</sup> Therefore, TMA can selectively scavenge

oxygen from the interface layer and reduce the interface trap density, while also thinning the IL, which has important benefits for gate stack dimensional scaling.

STEM-EELS compositional analysis for the Ni/Al<sub>2</sub>O<sub>3</sub>/SiGe/Si device shown in Fig. 8g supports this picture and indicates that a Si-rich interface forms, as shown by the intersection of the red dashed line marking the SiGe surface and half maximum count-rate of oxygen. The black and green arrows denote the Si and Ge composition at the oxide/SiGe interface, respectively. It is clearly seen that GeO<sub>x</sub> composition is diminished significantly at the SiGe surface consistent with less diffusion of diffusion of GeO<sub>x</sub> into the gate oxide and less diffusion of H<sub>2</sub>O through the Al<sub>2</sub>O<sub>3</sub> once the Al<sub>2</sub>O<sub>3</sub> reaches a critical thickness.

#### Conclusion

Novel gate oxide structures were investigated which suppress electronic defects at high-k/SiGe interfaces by employing an oxygen scavenging ALD precursor, TMA. The approach utilizes the difference in the heat of formation of SiO<sub>x</sub> and GeO<sub>x</sub>, achieving lower interface trap densities at the high-k / SiGe interface with just a modest reduction of  $C_{max}$ . Although metallic AI remains more effective at oxygen scavenging, it induces a much larger  $C_{max}$  reduction, demonstrating the benefit of TMA remote oxygen scavenging. The data is consistent with insertion of Al<sub>2</sub>O<sub>3</sub> into the HfO<sub>2</sub> gate oxide, removing a GeO<sub>x</sub> component of the interlayer between the channel and the deposited gate dielectric, and suggests that effective oxygen scavenging can be achieved with TMA during ALD. This TMA based oxygen scavenging technique is most effective when the Al<sub>2</sub>O<sub>3</sub> layers are uniformly distributed across the HfO<sub>2</sub> in a nanolaminate structure, but it also is effective when the Al<sub>2</sub>O<sub>3</sub> ALD deposition occurs on top of the HfO<sub>2</sub>. To achieve the effect, during each TMA half cycle of Al<sub>2</sub>O<sub>3</sub> ALD, TMA is dosed in excess to provide sufficient TMA for

reduction of additional chemical species after the surface hydroxyl groups are eliminated. In the present study, for the first time, the mechanism for low defect interface formation with  $Al_2O_3$  in contrast to  $HfO_2$  is explained with experimental data. In contrast to previous reports in which the good diffusion barrier properties of  $Al_2O_3$  are considered to be the source of  $D_{it}$  improvement, the mechanism for interface defect reduction with  $Al_2O_3$  in the present work is shown to be selective oxygen scavenging by diffusion of TMA reaction products in  $HfO_2$ . Furthermore, direct and remote oxygen scavenging using ALD  $Al_2O_3$  and etching back with atomic layer etching (ALE), it is likely that low  $D_{it}$  can be obtained in high aspect ratio structure such as on fins which is not possible with Al deposition for gettering gates.

#### **ASSOCIATED CONTENT: Supplementary materials**

The supporting information documenting variation of interface defect density on samples as well as comparison of interface defect calculation methods used in the article and the EELS analysis for few devices are available free of charge on the ACS Publication website at DOI:

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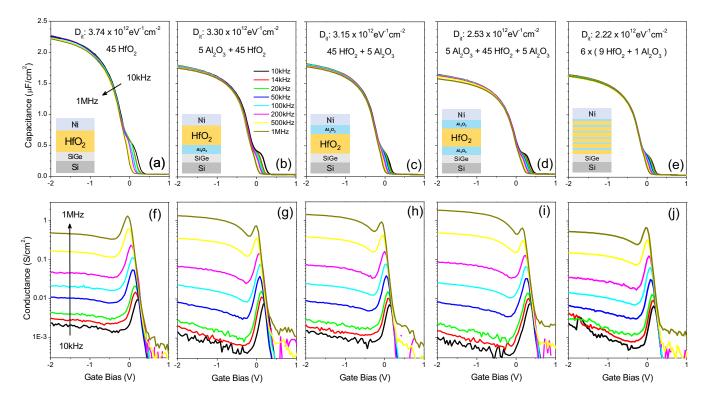


Fig. 1. C-V and G-V graphs of multilayer MOSCAP devices (a-j). Inset drawings indicate device structure for given graph. The inset  $D_{it}$  values indicates peak interface defect density value in bandgap obtained with full interface state model. The energy distribution of the defects density and the energy of the peak  $D_{it}$  values in band gap are shown in Fig.4. Control device with HfO<sub>2</sub> gate oxide shows highest interface defect density of  $3.74 \times 10^{12} \text{eV}^{-1} \text{cm}^{-2}$ . Al<sub>2</sub>O<sub>3</sub> insertion into HfO<sub>2</sub> reduces  $D_{it}$  gradually by going from HfO<sub>2</sub>-Al<sub>2</sub>O<sub>3</sub> bilayer into Al<sub>2</sub>O<sub>3</sub>-HfO<sub>2</sub>-Al<sub>2</sub>O<sub>3</sub> tri-layer and finally nanolaminate gate oxides to  $D_{it} = 2.22 \times 10^{12} \text{ eV}^{-1} \text{cm}^{-2}$ . Al<sub>2</sub>O<sub>3</sub> insertion reduces  $D_{it}$  regardless of its position shows

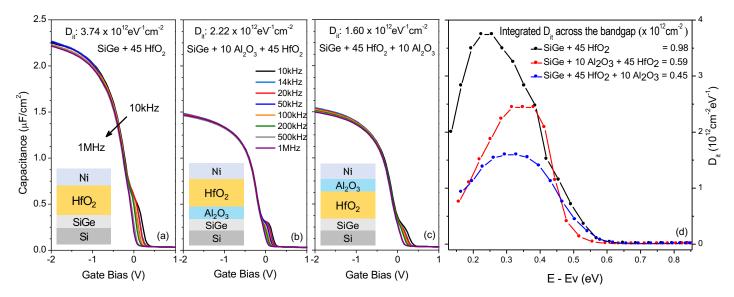


Fig. 2. C-V graphs and Interface defect density distributions across the band gap for MOSCAP devices calculated with the full interface state model. Inset drawings indicate device structure for given graph. The inset  $D_{it}$  values indicates peak interface defect density value in bandgap obtained with full interface state model (a) Control device with HfO<sub>2</sub> gate oxide shows highest interface defect density of  $D_{it}$  (peak) =  $3.74 \times 10^{12} \text{eV}^{-1} \text{cm}^{-2}$ . (b) Interface defects at SiGe oxide interface decrease by insertion of 10 cycles of Al<sub>2</sub>O<sub>3</sub> mono layers before HfO<sub>2</sub>  $D_{it}$  (peak) =  $2.22 \times 10^{12} \text{eV}^{-1} \text{cm}^{-2}$  and (c) after HfO<sub>2</sub>  $D_{it}$ (peak) =  $1.6 \times 10^{12} \text{eV}^{-1} \text{cm}^{-2}$  gate oxide. (d) Comparison of defect density at SiGe oxide interface across the band gap shows decrease in peak and integrated  $D_{it}$  (inset) by insertion of Al<sub>2</sub>O<sub>3</sub> layers.

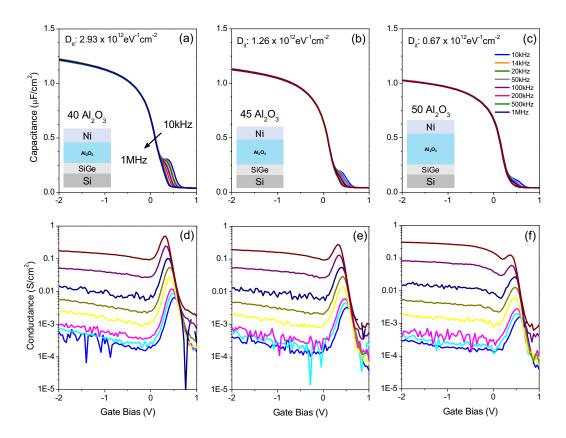


Fig. 3. C-V and G-V graphs of single oxide MOSCAP devices (a-f). Inset drawings indicate device structure for given graph. The inset  $D_{it}$  values indicates peak interface defect density value in bandgap obtained with full interface state model. Control device with 40 cycles of Al<sub>2</sub>O<sub>3</sub> gate oxide shows highest interface defect density of  $2.93 \times 10^{12} \text{ eV}^{-1} \text{cm}^{-2}$ . The small increase of Al<sub>2</sub>O<sub>3</sub> thickness reduces  $D_{it}$  by  $> 4 \times$ ; the lowest interface defect density of  $0.67 \times 10^{12} \text{ eV}^{-1} \text{cm}^{-2}$  is obtained with 50 cycles of Al<sub>2</sub>O<sub>3</sub>.

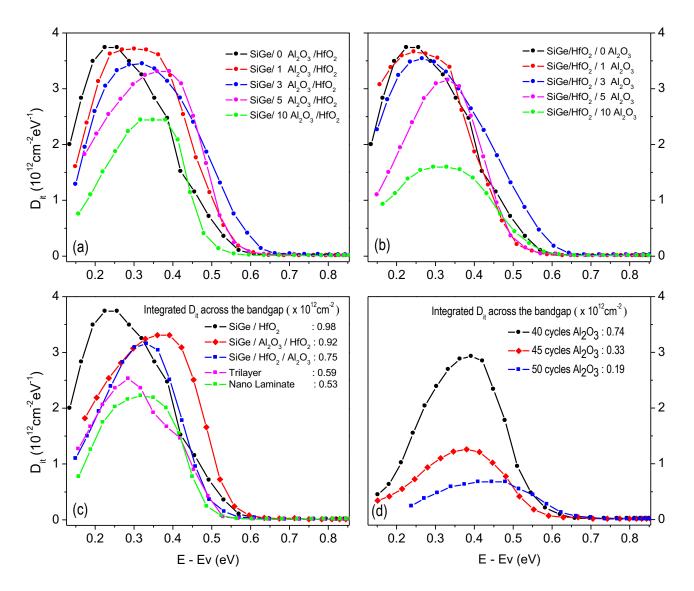


Figure 4. Interface defect density distributions across the band gap for MOSCAP devices calculated with the full interface state model. Interface defects at SiGe- oxide interface decrease by insertion of  $Al_2O_3$  layers before  $HfO_2$  (a) and after  $HfO_2$  (b) gate oxide. (c) Comparison of interface defects variation at SiGe- oxide interface by insertion of  $Al_2O_3$  layers into  $HfO_2$  gate oxide. (d) Interface defect density decreases by increase in  $Al_2O_3$  thickness. For 50 cycles of  $Al_2O_3$ , peak  $D_{it}$  reduces to  $6 \times 10^{11} \text{ eV}^{-1} \text{ cm}^{-2}$  and integrated defects across the bandgap is as low as  $0.19 \times 10^{11} \text{ eV}^{-1} \text{ cm}^{-2}$ .

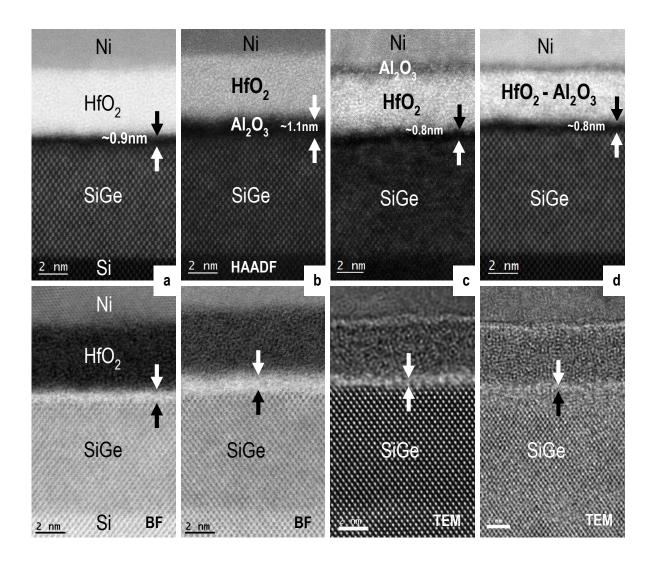
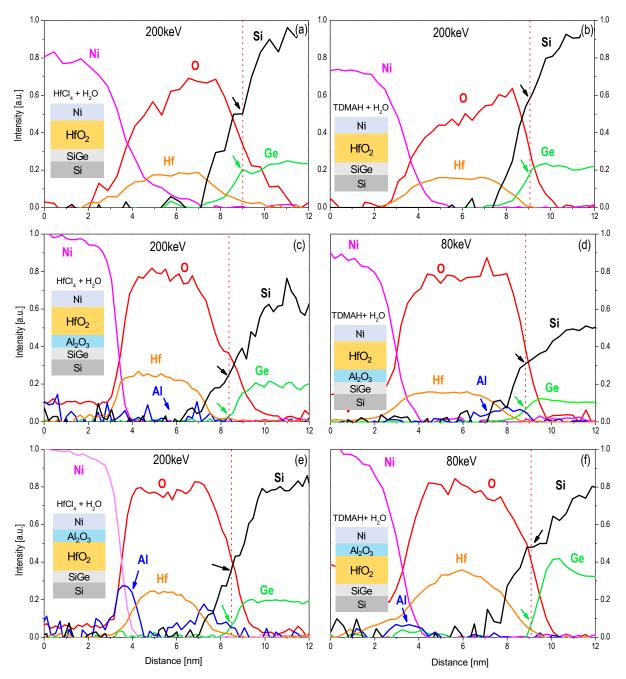
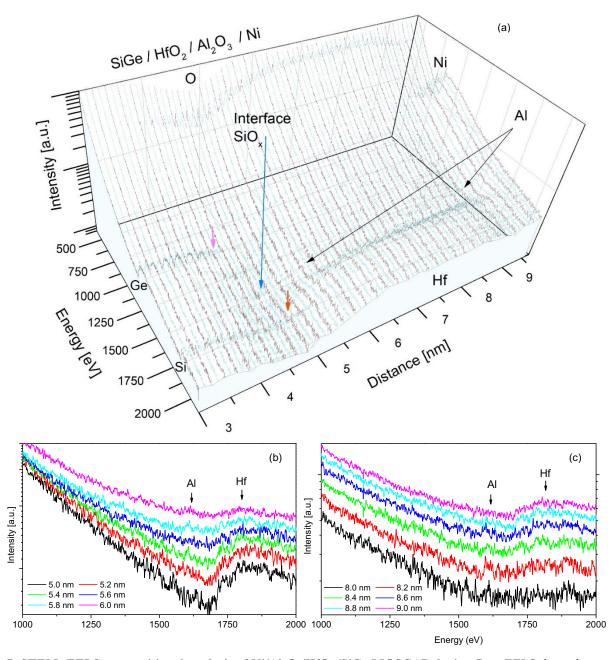


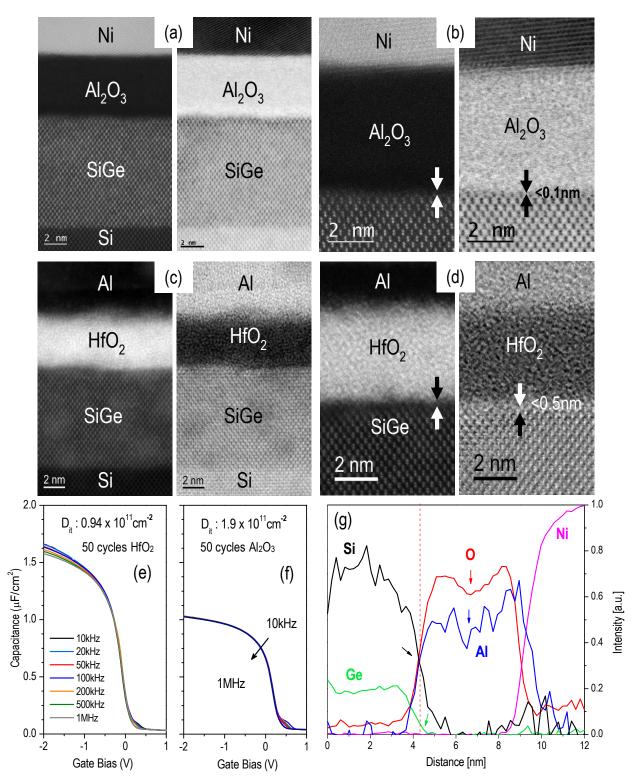
Figure 5. STEM HAADF (high-angle annular dark-field), BF (bright field) images of (a) control HfO<sub>2</sub>, (b)  $HfO_2/Al_2O_3/SiGe$  bilayer, (c)  $Al_2O_3/HfO_2/SiGe$  bilayer, (d) and  $Al_2O_3 - HfO_2$  Nanolaminate MOSCAPs. In these images, oxide structures and regions are defined according to z contrast. The interfacial layer between SiGe and oxide indicated with black and white arrows on corresponding STEM – HAADF, STEM - BF and TEM image. Note in (b) the interlayer consists of both SiGeO<sub>x</sub> and  $Al_2O_3$ , so it appears thicker than the control device in (a). In comparison to control device of  $HfO_2/SiGe$ , bilayer (c) and NL (d) shows thinner interface consistent with remote  $Al_2O_3$  insertion reducing IL.



**Figure 6. STEM- EELS compositional analysis of MOSCAP devices**. EELS experiment performed at 80keV and 200keV as indicated. The inset drawings illustrate corresponding gate stack structure along with the ALD chemistry above it. The compositions of the elements are averages area of ~6 x 0.2 nm parallel to sample surface. The red dashed line intercepts the half peak values of the O signals and indicate the SiGe - HfO<sub>2</sub> interface. Black and green arrows denote Si and Ge composition on the SiGe surfaces respectively. Blue arrow indicates the Al composition in the oxide. AlO<sub>x</sub>-HfO<sub>2</sub> interdiffusion is seen for bilayer samples regardless of the initial structure and confirmed with raw data analysis in Fig 6 and Fig. S5. This interdiffusion is prominent for EELS analysis at 80keV. In comparison to a-b, devices in c-f show lower Ge/Si ratio at the intersection with red dashed lines indicates Si rich interface formation with Al<sub>2</sub>O<sub>3</sub> incorporation into HfO<sub>2</sub>. Ni interdiffusion is seen in devices a and b and Al<sub>2</sub>O<sub>3</sub> insertion into HfO<sub>2</sub> impedes the Ni diffusion as seen in c-f.



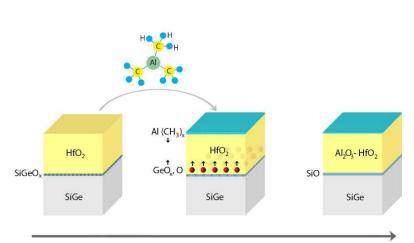
**Figure 7. STEM- EELS compositional analysis of Ni/Al<sub>2</sub>O<sub>3</sub>/HfO<sub>2</sub> /SiGe MOSCAP device**. Raw EELS data taken at 200 keV from sample in Fig. 6e is shown in a 3D semi-log graph (a) with the energy axis indicating the electron energy loss and corresponding intensity in arbitrary units. The axis labeled with distance indicates location of the electron beam on sample. The colored consecutive black and light blue lines indicate electron energy loss for the given location on samples and two colors chosen to enhance the image contrast. Each data line projects energy loss averaged from areas of  $5 \times 0.2$ nm parallel to the sample surface. The peaks appear on the graphs corresponds to Si K edge (1839 eV), Ge L edge 1217 eV, Hf M edge 1662 eV, O K edge 532 eV, Al K edge 1560 eV, Ni L edge 885 eV. The blue arrow indicates SiO<sub>x</sub> interface formation between SiGe and HfO<sub>2</sub>. Pink and red arrows indicate the Ge and Si compositions on SiGe surface. The Ge signal decays earlier than Si as it approaches the HfO<sub>2</sub> layer. Black arrows denote Al composition across the oxide. Al<sub>2</sub>O<sub>3</sub> insertion onto HfO<sub>2</sub> in bilayer structure forms intermixing by Al diffusion. To increase the visibility of Al peak and inter diffusion, semi-log 2D graph of raw EELS data with 1-2K energy loss range is presented in graph b (SiGe/HfO<sub>2</sub> interface region) and c (HfO<sub>2</sub>/Ni interface region) Two graphs prepared with offsets introduced between each curve to improve visibility.



**Figure 8.** Comparison of Ni/Al<sub>2</sub>O<sub>3</sub>/SiGe/Si and Al/HfO<sub>2</sub>/SiGe/Si MOSCAPs. Note the Al is a gettering gate metal. High resolution STEM HAADF and BF images at 80 keV of devices with (a-b) Ni/Al<sub>2</sub>O<sub>3</sub> and (c-d) Al/HfO<sub>2</sub> gate oxides; both have 50 ALD cycles. Nearly abrupt interfaces are observed in both bright field and dark field imaging for both devices. e-f) C-V graphs of Al/HfO<sub>2</sub> and Ni/Al<sub>2</sub>O<sub>3</sub> devices with very low depletion capacitance are shown, insets denote integrated D<sub>it</sub> values across the band gap extracted with full interface state model. EELS elemental composition of Ni/Al<sub>2</sub>O<sub>3</sub>/SiGe/Si is shown in g. The regions of Al<sub>2</sub>O<sub>3</sub> gate oxide defined with Z contrast in a-b are in good agreement with the EELS spectra in g. The red dashed line intercepts the half peak values of the O, Si and Ge signals and delineate the SiGe-HfO<sub>2</sub> interface. Black and green arrows denote Si and Ge composition on SiGe surface; a Si rich interfaces observed. The blue and red arrows indicates the electron beam damaged (80keV) region of Al<sub>2</sub>O<sub>3</sub>. A similar effect also observed for samples studied at 2000 keV. Ni gate metal overlaying with Al indicating Ni and Al intermix. Note at the SiGe interface, the Al and O profiles are nearly identical consistent with a near zero SiGeO<sub>x</sub> interlayer.

MOSCAD	C	N	D <sub>it</sub>	D <sub>it</sub>	Standard	Standard
MOSCAP devices in	C <sub>ox</sub>	V <sub>fb</sub>	C.M.	F.I.S.M	errors of	errors in
SiGe//Ni	µF/cm <sup>2</sup>	(V)	cm <sup>-2</sup> eV <sup>-1</sup>	cm <sup>-2</sup> eV <sup>-1</sup>	the mean	percentage
45 HfO <sub>2</sub>	2.5	-0.02	3.57	3.74	0.266	7.1%
5Al <sub>2</sub> O <sub>3</sub> /45HfO <sub>2</sub>	1.96	-0.03	2.43	3.30	0.074	2.2%
45 HfO <sub>2</sub> /5Al <sub>2</sub> O <sub>3</sub>	2.02	-0.01	2.35	3.15	0.10	3.1%
5Al <sub>2</sub> O <sub>3</sub> /45HfO <sub>2</sub> /5Al <sub>2</sub> O <sub>3</sub>	1.74	0.12	2.18	2.53	0.124	4.9%
6×(9HfO <sub>2</sub> +1Al <sub>2</sub> O <sub>3</sub> )	1.8	-0.05	2.02	2.22	0.05	2.2%

Table 1. Comparison of  $D_{it}$  values obtained with conductance and full interface state model.



Oxygen scavenging with TMA during Al<sub>2</sub>O<sub>3</sub> ALD process

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