

# **ARTICLE**

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# Enhanced electronic properties in mesoporous TiO<sub>2</sub> via lithium doping for high-efficiency perovskite solar cells

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Perovskite solar cells are one of the most promising photovoltaic technologies with their extraordinary progress in efficiency and the simple processes required to produce them. However, the frequent presence of a pronounced hysteresis in the current voltage characteristic of these devices arises concerns on the intrinsic stability of organo-metal halides, challenging the reliability of technology itself. Here, we show that n-doping of mesoporous TiO<sub>2</sub> is accomplished by facile post treatment of the films with lithium salts. We demonstrate that the Li-doped TiO<sub>2</sub> electrodes exhibit superior electronic properties, by reducing electronic trap states enabling faster electron transport. Perovskite solar cells prepared using the Li-doped films as scaffold to host the CH<sub>3</sub>NH<sub>3</sub>Pbl<sub>3</sub> light harvester produce substantially higher performances compared with undoped electrodes, improving the power conversion efficiency from 17 to over 19% with negligible hysteretic behaviour (lower than 0.3%).

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erovskite-based solar cells (PSCs) have made impressive strides in just a few years with maximum power conversion efficiencies (PCEs) jumping from 3.8% (ref. 1) in 2009 to 20.1% (ref. 2) in 2015. Even though further improvements are still expected<sup>3</sup>, such rapid progress is unprecedented for any photovoltaic (PV) material. For instance, silicon, GaAs, CIGS and CdTe required decades to fully realize their potential as solar cells<sup>4</sup>.

Perovskites comprise a large family of crystalline materials, where the most commonly used for solar cells have an ABX<sub>3</sub> chemical composition containing an organic cation A, such as methylammonium (MA) or formamidinium (FA)<sup>5,6</sup>, a divalent metal B, such as Pb or Sn<sup>7,8</sup>, and a halide X, such as Br or I. These organicinorganic perovskites can be processed by a large number of techniques ranging from spin coating<sup>5</sup>, dip coating<sup>9</sup>, two-step interdiffusion<sup>10</sup>, chemical vapour deposition<sup>11</sup>, spray pyrolysis<sup>12</sup>, atomic layer deposition<sup>13</sup>, ink-jet printing<sup>14</sup>, to thermal evaporation<sup>15,16</sup>. The PV performances have been attributed to their outstanding optoelectronic properties such as remarkably high absorption over the visible spectrum<sup>7</sup>, charge carrier diffusion lengths in the micrometre-range<sup>17–19</sup> implying a sharp optical band edge, and a tuneable band gap from 1.1 to 2.3 eV by interchanging the above cations<sup>2,20</sup>, metals<sup>21,22</sup> and/or halides<sup>23</sup>.

Recently, Jeon et al. achieved one of the highest certified PCEs of 17.9% by using the mixed halide and cation formulation, (FAPbI<sub>3</sub>)<sub>0.85</sub>(MAPbBr<sub>3</sub>)<sub>0.15</sub> (ref. 24). Their record solar cell architecture contains a mesoporous TiO2 layer, which is infiltrated by a liquid perovskite precursor solution forming the solid perovskite film after subsequent annealing. Mesoporous TiO<sub>2</sub> has been widely used for high-surface-area electrodes, in optoelectronic applications<sup>25</sup> and, in particular, in dye-sensitized solar cells (DSSCs)<sup>26</sup>, where they have been demonstrated to collect and transport electrons photoinjected from a surfaceadsorbed sensitizer. One strategy to improve DSSCs was to enhance the electron transport of the mesoporous TiO2 making use of substitutional dopants<sup>27–37</sup>. Also, lithium intercalation has been employed to lower the conduction band edge of TiO2 facilitating electron injection and transport in the mesoporous TiO<sub>2</sub> (refs 38-41). By analogy, it may be expected that the introduction of n-dopants would also enhance the performance of PSCs. However, so far very few studies have examined the doping effect on the electron transport within the mesoporous TiO2 scaffold employed as PCSs42. This may be attributed to the perovskites already being excellent charge transporting materials inherently suited for electron conduction in high-efficiency solar cells<sup>17</sup>. This makes it challenging to investigate the effects of enhanced charge transport in mesoporous TiO2 as this improvement would only have a discernible impact for embodiments when the perovskite is operating at its limit.

Therefore, at first glance, it may be assumed that highly efficient PSCs could be achieved by depositing a perovskite film directly on a thin TiO2 compact layer, which would effectively work as an electron-selective contact<sup>15</sup>. However, flat junction PSCs prepared on a compact TiO<sub>2</sub> suffer from relatively poor charge collection efficiency under steady-state forward voltage bias<sup>43</sup>. Although Wojciechowski et al. have shown significant improvement of this architecture by modifying the TiO<sub>2</sub>-perovskite flat heterojunction with fullerene derivatives<sup>44</sup>, the steady-state PCE of TiO<sub>2</sub>-based flat PSCs is still substantially lower than those based on mesoporous TiO2. The planar PSC architecutres are also plagued by severe hysteresis in the J-V curve rendering it very difficult to determine their solar to electric PCEs. For this reason, none of the TiO2-based flat PSC architectures have been certified so far. Recently, Guillén et al. demonstrated that the charge collection in a PSC with mesoporous TiO<sub>2</sub> involves two separate electron transport

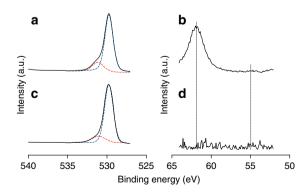
pathways: one through the perovskite, and one through the mesoporous  ${\rm TiO_2}$  (refs 45,46). Following this, Heo *et al.* reported that surface deposition by sintering of bis(trifluoromethane)sulfonimide lithium salt (Li-TFSI) onto mesoporous  ${\rm TiO_2}$  results in substantial improvements in the PCE and suppression of the hysteresis.

In this work, we show that mesoporous  ${\rm TiO_2}$  can be n-doped in a facile and effective way by a similar lithium salt surface treatment. We show that the Li-doping enables faster electron transport within the mesoporous  ${\rm TiO_2}$  electrodes and demonstrate that PSCs prepared on such electrodes achieve substantially higher performances compared with undoped electrodes, improving PCEs from 17 to 19.3% with negligible hysteresis behaviour (lower than 0.3%).

### **Results**

**X-ray photoelectron spectroscopy analysis**. We applied the lithium ion surface treatment of the meso-porous  ${\rm TiO_2}$  layer via spin coating of a LiTFSI solution (see the Methods section). After the deposition and solvent evaporation, the substrates were sintered at 450 °C for 30 min. The introduction of Li  $^+$  ions by the thermal diffusion modifies the surface of the particles. Together with the doping effect of the  ${\rm TiO_2}$  an overlayer of LiO $_2$  or LiOH could be formed  $^{47}$ . Interestingly, also the formation of spinel structures like  ${\rm Li_4Ti_5O_{12}}$  was observed with a synthetic procedure similar to the one employed here  $^{48}$ .

We used X-ray photoelectron spectroscopy (XPS) to study the elemental composition of the Li-treated and untreated TiO2 after sintering. No traces of sulphur or fluorine from the LiTFSI precursor for the treated sample were detected and the highresolution spectra for C 1s showed no difference between the Li-treated and untreated samples as seen in Supplementary Fig. 1. The O 1s spectra in Fig. 1a,c show that Li-treated TiO<sub>2</sub> (Fig. 1a) has a more pronounced shoulder at the higher energy of the main peak compared with the untreated TiO2 (Fig. 1c). The deconvolution of this signal reveals a second small peak at 531.2 eV for the untreated sample and a much more pronounced one for the Li-treated one that has been previously assigned to the oxygen interaction with the lithium<sup>47</sup>. The Ti 2p spectra can be found in Supplementary Fig. 2, and no difference was detected for the treated and untreated samples. On the right hand side of Fig. 1, the spectrum from 50 to 65 eV for the Li-treated TiO<sub>2</sub> (Fig. 1b) electrodes shows a peak corresponding to the Ti 3s and a weak, but distinct, signal for Li 1s. In the untreated TiO<sub>2</sub> (Fig. 1d),



**Figure 1 | X-ray photoelectron spectroscopy.** Doped and undoped mesoporous  $TiO_2$  layers for the O 1s peaks Li doped (**a**) and the undoped control (**c**), Ti 3s and Li 1s peaks slightly visible at 55 eV for Li-doped (**b**) and the signal for the undoped  $TiO_2$  here shown as reference, which reveals the absence of the peaks related to Ti 3s (dashed line at 61.5 eV) and Li 1s (**d**), dashed line at 54.9 eV.

no such signals, measured over a series of samples, could be detected in this energy region and it is here shown as noise. The Ti 3s peak at 61.5 eV reveals the presence of Ti³+ for Li-treated electrodes and thus indicates that the Li⁺ treatment induces a partial reduction of Ti⁴+ to Ti³+ within the TiO₂ lattice³9,49, which is not seen for the untreated electrode. Pathak *et al.* demonstrated that a small amount of species with valency +3 can passivate the electronic defects or trap states that originate from oxygen vacancies within the TiO₂ lattice³7. Accordingly, the Li⁺ doping mechanism is consistent with a passivation of electronic trap states resulting in improved charge transport properties and thus in better performing mesoporous TiO₂ electrodes⁵0.

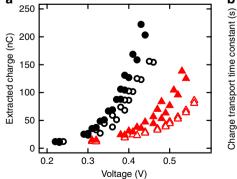
Charge extraction and electron transport analysis. To study the impact of the Li<sup>+</sup> doping on the electronic states and the charge transport within the TiO<sub>2</sub>, we prepared solid-state DSSCs using Li<sup>+</sup>-doped mesoporous TiO<sub>2</sub> as electron transporting layer. DSSCs were prepared according to the previously reported procedures<sup>51</sup>.

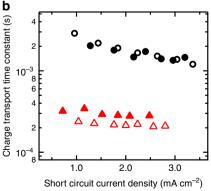
In the DSSC field, charge extraction is a well-established light-assisted technique, which qualitatively draws the density of state distribution below the TiO<sub>2</sub> conduction band<sup>52</sup>. In Fig. 2a, we report the charge extracted from the DSSCs at open circuit condition as a function of the open circuit voltage. At the same open circuit voltage, the devices employing Li:TiO<sub>2</sub> hold significantly less charges than the undoped TiO<sub>2</sub> (for example at 0.46 V, Li:TiO<sub>2</sub> and undoped TiO<sub>2</sub> holds 44.3 and 154.3 nC, respectively). This suggests that the Li<sup>+</sup> doping reduces the concentration of sub-bandgap states in the TiO<sub>2</sub> supporting our previous conjecture that a partial reduction of Ti<sup>4+</sup> to Ti<sup>3+</sup> passivates the trapping states associated with oxygen vacancies within the TiO<sub>2</sub> lattice.

From the same DSSCs, we extracted the charge transport time constant by intensity-modulated photocurrent spectroscopy<sup>45,53,54</sup>. The cell was biased at short circuit under light and the time constants were measured for different light intensities. In Fig. 2b, we compare the transport time constants for DSSCs prepared with Li<sup>+</sup>-treated and untreated TiO<sub>2</sub> electrodes. Li<sup>+</sup>-treated devices display up to over one order of magnitude faster charge transport than the untreated devices over the whole range of current densities. To explain this trend, we note that charge transport in DSSCs is controlled by the electron transport in the TiO<sub>2</sub> (ref. 41). In particular, it has been shown

that the electronic transport in the TiO<sub>2</sub> is limited by the temporary localization of electrons within sub-bandgap states, which can be passivated with different doping mechanisms<sup>53,55</sup>. Our results indicate that Li<sup>+</sup> doping also reduces the concentration of sub-bandgap states (Fig. 2a) and improves the electronic transport within mesoporous TiO<sub>2</sub> electrodes (Fig. 2b). As we observe the presence of Li<sup>+</sup> ions within the TiO<sub>2</sub> lattice (Fig. 1), we can regard this method as an effective way of n-doping via a facile post treatment of the mesoporous TiO<sub>2</sub> films. Even though such doping strategies may be effective for improving DSSC performance, they may not necessarily benefit PSCs, as the latter already accomplish fast charge transport within the perovskite layer<sup>17</sup>. To further elucidate the effect of Li doping on the electron transport in mesoscopic perovskite solar cells, we performed light intensity-modulated photocurrent spectroscopy also on PSC $^{56-58}$ .

In Fig. 3a,b, we report the intensity modulated photocurrent spectroscopy (IMPS) spectra at different light intensities for PSCs prepared with mesoporous TiO<sub>2</sub> films with and without the Li<sup>+</sup>doping. The Nyquist plot of the untreated sample shows two distinct semicircles related to two different transport processes. In the treated sample, on the contrary, the slow time component (semicircle on the right) merges with the faster component (semicircle on the left) confirming that the TiO<sub>2</sub> transport time is faster than that one observed in the control. These processes are resolved in frequency in Fig. 3b, where the imaginary part of IMPS is plotted against the frequency. The electron transport time for a given pathway equals the inverse of the frequency of the corresponding peak<sup>53</sup>. In the case of the untreated  $TiO_2$  films, we observe two distinct peaks at 10<sup>3</sup> and 10<sup>4</sup>-10<sup>5</sup> Hz. Guillén et al. correlated these peaks with two separate electron transport pathways, running in parallel, involving the mesoporous TiO2 and the perovskite, respectively<sup>45</sup>. The peak at low frequency  $(10^3 \text{ Hz})$  is directly linked to the electron transport in TiO<sub>2</sub> and it shifts to higher frequency (faster transport) as the photocurrent density increases, as revealed also by the trends in Fig. 2b. By contrast, the high-frequency peak (104-105 Hz) is related to the charge transport through the perovskite and does not shift significantly with increasing the photocurrent density<sup>45</sup>. Therefore, at higher current density, the TiO2 and the perovskite peaks overlap at around 10<sup>4</sup>-10<sup>5</sup> Hz. Interestingly, for the Li + -doped PSCs, the TiO<sub>2</sub> peak at low frequency does not show up even for extremely low current density, such as 0.07 mA cm<sup>-2</sup>. As Li<sup>+</sup>-doping of TiO<sub>2</sub> enables faster electron injection and transport in DSSCs, we infer from the lack of a low-





**Figure 2 | Charge extraction and electron transport time constant.** (a) Charge extracted at open circuit as function of the voltage for no treated samples (black) and Li-doped samples (red) and (b) charge transport lifetime as function of the short circuit current density for solid-state dye-sensitized solar cells prepared without and with Li-doped mesoporous TiO<sub>2</sub> (black and red markers, respectively). Two samples per condition are shown: open triangle, sample 1 Li-doped; closed triangle, sample 2 Li-doped; open circle, sample 1 no treated; closed circle, sample 2 no treated. For the charge extraction measurement (a), each sample was measured twice.

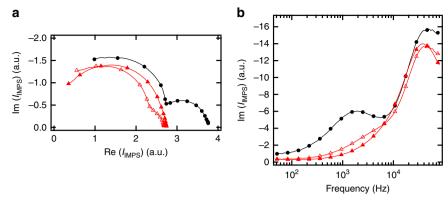
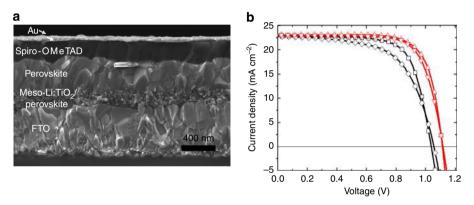


Figure 3 | IMPS spectra. Nyquist plot (a) for PSCs prepared with no treated TiO<sub>2</sub> electrode (in black at 0.15 mA cm<sup>-2</sup>) and with Li-doped TiO<sub>2</sub> electrode at 0.07 mA cm<sup>-2</sup> (open triangle) and 0.26 mA cm<sup>-2</sup> (closed triangle). (b) Imaginary component of the same IMPS shown in a versus the frequency.



**Figure 4 | Morphology and** *J-V* **curves of the solar cells.** (a) SEM cross-sectional image of the device. (b) Current density-voltage curves of the solar cells with and without the Li<sup>+</sup> doping (red and black curve, respectively) collected under AM1.5 simulated sun light. Devices were masked with a black metal aperture of 0.16 cm<sup>2</sup> to define the active area. The curves were recorded scanning at 0.01 V s<sup>-1</sup> from forward bias to short circuit condition and *vice versa* with no device preconditioning such as light soaking or holding at forward voltage bias. Legend: black diamond, control forward scan; black square, control reverse scan; red down-pointing triangle, Li + -doped forward scan; red up-pointing triangle, Li + -doped reverse scan.

frequency TiO<sub>2</sub> peak for the Li<sup>+</sup>-treated PSCs (Fig. 3) that the electron transport through the TiO<sub>2</sub> occurs at a similar rate as through the perovskite. This result also suggests that in PSCs employing doped TiO<sub>2</sub> further improvements may originate from applying hole transporting materials with a higher hole mobility than the commonly used spiro-MeOTAD.

PV characterization. In Fig. 4a, we show the stack architecture used in this study composed of fluorinated tin oxide (FTO)/ compact-TiO<sub>2</sub>/mesoporous-TiO<sub>2</sub>/perovskite/Spiro MeOTAD/ gold. An X-ray diffraction pattern of the perovskite employed in solar cell fabrication is shown in Supplementary Fig. 3, where the peroskvite structure is detected. Figure 4b shows a typical current density-voltage (J-V) characteristic that we measured for PSCs with and without Li+-doped mesoporous TiO2. In Table 1, we summarize the device performance parameters, as extracted from the J-V curves in Fig. 4b, and the light intensity measured during each J-V scan. We note that the difference between the backward and the forward scans is significantly larger in the control than in the Li+-doped device. The short circuit currents are quite similar, whereas the Li+-doped device showed significantly larger open circuit voltage (about 60 mV) and 0.11 units higher fill factor. The current mismatch between the measured current density and the current density calculated by the incident photonto-electron conversion efficiency (IPCE) over the Solar AM1.5 G spectrum was less than 2% (Supplementary Figs 4 and 5). The

	Scan direction	$J_{\rm sc}$ (mA cm $^{-2}$ )	V <sub>oc</sub> (V)	FF	PCE (%)	Light intensity (mW cm <sup>- 2</sup> )
Li <sup>+</sup> doped	Backward Forward	23.0 23.1	1.114 1.118	0., .	19.3 19.0	98.1
Control	Backward Forward	22.7 22.7	1.038 1.056	0.72 0.61	17.1 14.7	99.4

FF, fill factor; PCE, power conversion efficiency. Short circuit photocurrent ( $J_{sc}$ ), open circuit voltage ( $V_{oc}$ ), FF extracted from the data in

overall power conversion efficiency of devices employing the Li<sup>+</sup>-doped TiO<sub>2</sub> electrodes was systematically higher than the devices employing the undoped scaffold (Supplementary Fig. 6). This result is consistent with the fact that the Li<sup>+</sup>-doping decreases the number of deep traps, which act as recombination centres and induces faster charge transport within the TiO<sub>2</sub>, improving the open circuit voltage and fill factor, respectively.

### **Discussion**

In summary, we demonstrated a doping mechanism that allowed preparing mesoporous  ${\rm TiO_2}$  electrodes with superior electron

properties. The doping can be accomplished with a facile post treatment of the mesoporous  ${\rm TiO_2}$  making use of lithium salts to induce a partial reduction of  ${\rm Ti}^{4}{}^+$  to  ${\rm Ti}^{3}{}^+$  within the  ${\rm TiO_2}$  lattice and passivating electronic defect states acting as nonradiative recombination centres. We exploit the  ${\rm Li}^+$ -doped mesoporous  ${\rm TiO_2}$  electrodes to improve the maximum power conversion efficiency of perovskite solar cells from 17% to over 19%, which is comparable to the highest values reported in the literature.

# Methods

Substrate preparation and Li-doping TiO<sub>2</sub>. Nippon Sheet Glass of  $10\,\Omega\,\mathrm{sq}^{-1}$  was cleaned by sonication in 2% Hellmanex water solution for 30 min. After rinsing with deionized water and ethanol, the substrates were further cleaned with ultraviolet ozone treatment for 15 min. Then, 30-nm TiO<sub>2</sub> compact layer was deposited on FTO via spray pyrolysis at 450 °C from a precursor solution of titanium diisopropoxide bis(acetylacetonate) in anhydrous ethanol. After the spraying, the substrates were left at 450 °C for 45 min and left to cool down to room temperature. Then, mesoporous TiO<sub>2</sub> layer was deposited by spin coating for 20 s at 4,000 r.p.m. with a ramp of 2,000 r.p.m. s<sup>-1</sup>, using 30 nm particle paste (Dyesol 30 NR-D) diluted in ethanol to achieve 150- to 200-nm-thick layer. After the spin coating, the substrates were immediately dried at 100 °C for 10 min and then sintered again at 450 °C for 30 min under dry air flow.

Li-doping of mesoporous  $TiO_2$  was accomplished by spin coating a 0.1 M solution of Li-TFSI in acetonitrile. The solution was prepared freshly before the application in nitrogen atmosphere. 150 µl were poured on  $1.4 \times 2.4 \, \text{cm}^2$  substrate. After 5 s of loading time, the spinning programme started with an acceleration of  $1,000 \, \text{r.p.m. s}^{-1}$  to a final speed of  $3,000 \, \text{r.p.m.}$ , the substrate was left spinning for  $30 \, \text{s.}$  Both  $\text{Li}^+$ -doped and undoped electrodes were completed with a second calcination step at  $450 \, ^{\circ}\text{C}$  for  $30 \, \text{min.}$  After cooling down to  $150 \, ^{\circ}\text{C}$ , the substrates were immediately transferred in a nitrogen atmosphere glove box for the deposition of the perovskite films.

**DSSC preparation procedure.** Glass substrates for solid-state DSSCs were prepared following the same procedure used for PSC. 900 nm of mesoporous  ${\rm TiO_2}$  (Dyesol 30 nrd ethanol diluited ) were deposited by spin coating at 4,000 r.p.m. After sintering, the substrates were cooled down to 70 °C and immersed in 0.1 mM solution of dye (Y123) in 1:1 mixture of acetonitrile and tert-butyl alcohol for 30 min. After the dyed films were rinsed in abundant acetonitrile, the hole conductor was applied by spincoating at 2,000 r.p.m. for 20 s. The hole transport composition and the following steps to complete the DSSCs were identical to what used for PSCs.

**Perovskite precursor solution and film preparation.** The perovskite films were deposited from a precursor solution containing FAI (1 M), PbI<sub>2</sub> (1.1 M), MABr (0.2 M) and PbBr<sub>2</sub> (0.2 M) in anhydrous dimethylformamide/ dimethylsulphoxide (4:1 (v:v)) solution. The perovskite solution was spin coated in a two-step programme at 1,000 and 4,000 r.p.m. for 10 and 30 s, respectively. During the second step,  $100\,\mu$ l of clorobenzene was poured on the spinning substrate 15 s prior the end of the programme. The substrates were then annealed at  $100\,^{\circ}$ C for 1 h in nitrogen-filled glove box.

We note that the perovskite precursor solution for this paper was prepared with a different composition from what reported by Jeon *et al.*, who used an equimolar amount of FAI and PbI<sub>2</sub> to achieve a certified PCE of 17.9% with the mixed halide and cation formulation, (FAPbI<sub>3</sub>)<sub>0.85</sub>(MAPbBr<sub>3</sub>)<sub>0.15</sub> (ref. 24). Interestingly, we observed a systematic improvement in PCE moving away from the equimolar concentration for FAI and PbI<sub>2</sub> towards 10 mol% lower stoichiometric amount of FAI.

Hole transporting layer and top electrode. After the perovskite annealing, the substrates were cooled down for few minutes and a spirofluorene-linked methoxy triphenylamines (spiro-OMeTAD, from Merck) solution (70 mM in chlorobenzene) was spun at 4,000 r.p.m. for 20 s. The spiro-OMeTAD was doped with bis(trifluoromethylsulfonyl)imide lithium salt (Li-TFSI, from Aldrich), tris(2-(1H-pyrazol-1-yl)-4-*tert*-butylpyridine)- cobalt(III) tris(bis(trifluoromethylsulfonyl)imide) (FK209, from Dyenamo) and 4-*tert*-Butylpyridine (TBP, from Aldrich) <sup>40,41,59</sup>. The molar ratio of additives for spiro-OMeTAD was: 0.5, 0.03 and 3.3 for Li-TFSI, FK209 and TBP, respectively. Finally, 70–80 nm of gold top electrode was thermally evaporated under high vacuum.

**PV device testing.** The solar cells were measured using a 450-W xenon light source (Oriel). The spectral mismatch between AM1.5G and the simulated illumination was reduced by the use of a Schott K113 Tempax filter (Präzisions Glas & Optik GmbH). The light intensity was calibrated with a Si photodiode equipped with an IR-cutoff filter (KG3, Schott) and it was recorded during each measurement. Current-voltage characteristics of the cells were obtained by applying an external voltage bias while measuring the current response with a

digital source meter (Keithley 2400). The voltage scan rate was  $10\,\mathrm{mV}\,\mathrm{s}^{-1}$  and no device preconditioning was applied before starting the measurement, such as light soaking or forward voltage bias applied for long time. The starting voltage was determined as the potential at which the cells furnishes  $1\,\mathrm{mA}$  in forward bias, no equilibration time was used. The cells were masked with a black metal mask (0.16 cm²) to estimate the active area and reduce the influence of the scattered light. The devices were characterized 2 days after their preparation.

Charge extraction technique. Charge extraction measurement were performed with Autolab potentiostat PGSTAT30 driven by NOVA software. The procedure for the charge extraction comprised four steps. First, the cell was kept for 10 s in dark at short circuit. At this stage, the carriers eventually accumulated in the intrinsic capacitances of the device were discharged. Then, the potential was brought to open circuit and the light was switched on for an equilibration time of 10 s. The light was then switched off and the open circuit voltage decay was monitored for a defined decay time  $T_{\rm d}$ . In the last step, the cell was brought back to short circuit condition from  $V_{\rm TD}$  (voltage at time  $T_{\rm d}$ ) and the discharge current was measured. The integration over the time (starting from  $T_{\rm d}$ ) of this current gave the value of charge stored at the voltage  $V_{\rm TD}$ .

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# **Author contributions**

F.G. proposed the experiment. F.G., A.A., J.P.C.B., M.S. and T.M. prepared the PSC devices. J.P.C.B. performed the microscopy analysis and analysed the XPS data. F.G., A.A. and J.P.C.B. prepared the SSDSC devices, performed IMPS, charge extraction and analysed the data. All the authors contributed to the analysis of the data and the writing of the paper. M.G. supervised the experiment.

### **Additional information**

Supplementary Information accompanies this paper at http://www.nature.com/naturecommunications

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