



Published in final edited form as:

Nat Geosci. 2017 ; 10: 646–651. doi:10.1038/ngeo3016.

Equatorial jet in the lower to middle cloud layer of Venus revealed by Akatsuki

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Author contributions

T.H. developed automated cloud tracking and error evaluation methods, corrected the bore-sight of IR2 nightside images, conducted tracking, and interpreted the results. S.M. and K.O. contributed cloud tracking programing. T.S., T.M.S., K.S., T.I., and M.N. conducted IR2 observations and contributed to the operation of Akatsuki and observation planning. K.O., T.K., H.K., and M.T. developed the bore-sight correction applied to dayside images, and they also developed geographical mapping. J.P. conducted manual tracking with IR2 data based on independent geographical mapping. T.S., J.P., T.K., S.S.L., M.T. and E.F.Y. helped scientific interpretation and the review of previous studies. S.W., M.Y., and A.Y. conducted UVI observations. E.F.Y. conducted IRTF observations and the tracking with them. K.M. contributed by optical and cloud-physical interpretation.

Competing financial interests

The authors declare no competing financial interests.

Supplementary animation GIF file

Unfiltered two-hourly 2.26- μm radiance observed with IR2 from 12 h UTC, July 11 to 2 h UTC, July 12, 2016. It is false-colored as shown at the bottom ($\text{W m}^{-2} \text{sr}^{-1} \mu\text{m}^{-1}$). The data are map-projected as if it were observed from an infinite distance at the longitude of 355° and the equator. The original images between 18 and 22 h include the dayside, but the white-out regions are not entirely saturated, since we set color-scales for good nightside representation. The other original images were obtained by removing the dayside from the field of view.

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Abstract

The Venusian atmosphere is in a state of superrotation where prevailing westward winds move much faster than the planet's rotation. Venus is covered with thick clouds that extend from about 45 to 70 km altitude, but thermal radiation emitted from the lower atmosphere and the surface on the planet's night-side escapes to space at narrow spectral windows of near-infrared. The radiation can be used to estimate winds by tracking the silhouettes of clouds in the lower and middle cloud regions below about 57 km in altitude. Estimates of wind speeds have ranged from 50 to 70 m/s at low- to mid-latitudes, either nearly constant across latitudes or with winds peaking at mid-latitudes. Here we report the detection of winds at low latitude exceeding 80 m/s using IR2 camera images from the Akatsuki orbiter taken during July and August 2016. The angular speed around the planetary rotation axis peaks near the equator, which we suggest is consistent with an equatorial jet, a feature that has not been observed previously in the Venusian atmosphere. The mechanism producing the jet remains unclear. Our observations reveal variability in the zonal flow in the lower and middle cloud region that may provide new challenges and clues to the dynamics of Venus's atmospheric superrotation.

Introduction

The planet Venus rotates westward with a very low angular speed corresponding to a period of 243 days, but its atmosphere rotates in the same direction at much higher angular speeds¹. This superrotation reaches its maximum near the cloud top located at around the altitude of 70 km, where the rotational periods are 3 to 5 days, several tens of times faster than the planetary rotation^{1,2}. Measurements by entry probes like Veneras, Pioneer Venus Multiprobe, and VEGA balloons revealed that zonal wind speeds below the cloud top decreases quasi-linearly with depth²⁻⁴. Despite the long history of studies of the superrotation⁵⁻¹², its mechanism continues to be a focus of research^{13,14}.

Two distinct atmospheric spectral windows between the wavelengths of 1.6 and 2.6 μm enable tracking clouds from ground-based observations¹⁵. It has been noted¹⁶ that the observed near-infrared (NIR) radiance has spatial variation owing to the spatial variation of optical thickness of condensational clouds below 57 km where static stability is low¹⁷, the region traditionally called the lower (47–50 km) and middle (50–57 km) cloud regions¹⁸. The NIR image tracking was later applied for spacecraft observations during the Galileo flyby¹⁹ and with the Venus Express orbiter^{20,21}. These and other previous studies consistently measured the zonal wind speeds in the lower to middle cloud layer in the range of 50 to 70 m/s over low to mid latitude, weakly peaking at mid-latitude¹⁵⁻²⁴. In contrast, Crisp et al²⁵ performed a two-month campaign of ground-based observations in 1990 covering Galileo's Venus flyby and found that features with horizontal scales greater than 2,000 km tend to move faster than smaller-scale features. These large-scale features moved roughly at a uniform angular velocity (solid-body rotation) and westward speeds around 80 m/s at the equator. They speculated that the apparent coexistence of distinct rotational speeds

may reflect motions at different altitudes; for example, large and small features may reflect motions in the middle and lower cloud regions, respectively. The speculation, however, has not been confirmed by later studies.

Akatsuki observations and cloud tracking

Akatsuki was launched in 2010, and after the failure of the first orbit insertion in the same year, it was successfully maneuvered to start orbiting Venus on Dec 7, 2015^{26,27}. Its orbital period is about 11 days, and its low orbital inclination ($<10^\circ$) makes it suitable to observe low latitudes. Akatsuki is equipped with cameras designed to sense the Venusian atmosphere and clouds at multiple wavelengths. One of them, the Longwave Infrared Camera (LIR) detected a surprisingly large planetary-scale stationary gravity wave²⁸.

In this study, we mainly use data from the IR2 camera, the 2- μm infrared camera onboard Akatsuki, which images the planet between 1.74 to 2.32 μm ^{26,27,29} through four narrow band filters. IR2 started imaging in March 2016. As the orientation of Akatsuki's elongated elliptical orbit is nearly fixed in the inertial space, times favored for day-side or night-side observations alternate in a Venus year. April-June 2016 was the season for day-side observations, while March, July, and August, 2016 were favorable for night-side observations. In this study, we use data obtained when the satellite was relatively close to Venus in March, July, and August, 2016. For comparison, we also use data from the instrument UVI, the ultraviolet imager of Akatsuki, which images the cloud top with scattered sunlight^{26,27}.

In this study, we use a novel automated cloud tracking method^{30,31}, which is more precise and more reliable than previously used methods³². The template size used is $7.5^\circ \times 7.5^\circ$. To rule out the tracking of cloud features other than passive tracers (like atmospheric waves^{28,31}), the results were manually verified by human eyes. Error evaluation is made with a novel method³⁰. We also conduct independent manual tracking³³ to double-check. See Methods for more details.

Equatorial jet on July 11-12, 2016

Figure 1a shows the horizontal velocities (arrows) obtained from four 2.26 μm images taken at every two hours from 20:03 UTC, July 11, 2016. During the six-hours, the sub-spacecraft points moved from 4°W , 5°N to 13°W , 8°N , close to the longitude of sunset ($2\text{--}3^\circ\text{E}$). The winds are obtained between the local times 19:00 and 22:30 (see the abscissa of Fig. 2a and the tracer longitudes at end of cloud tracking shown in Fig. 1d). The wind field displayed in Fig. 1a is dominated by westward flow as reported in the previous studies with NIR observations^{15–22}. However, the trajectories associated with the results (“+” marks) indicate that rotational speed peaks near the equator, which is also apparent in the clouds' displacements in high-pass filtered images (Fig. 1b-d).

The longitudinal average of zonal wind (Fig. 2b) displays a clear wind-speed maximum of about 85 m/s slightly to the north of the equator, corresponding to a rotation period of about 5 days. It is not only the wind speed but also the angular speed around the rotational axis peaks near the equator, a feature previously unseen at any altitude. Accordingly, the angular

momentum is also maximized there. Manual tracking independently confirmed this result. The present equatorial zonal flow is much faster than previously reported^{15–22} except for the motion of large-scale features reported by Crisp et al¹⁶. Note that our results are obtained by tracking small-scale features by using $7.5^\circ \times 7.5^\circ$ templates.

We shall call the feature exhibiting a rotational speed maximum near the equator as the equatorial jet. Since it has not been observed previously, to assert its reality would require further examination. Specifically, one may suspect that the angular speed contrast is mainly in the vertical direction rather than in the meridional direction and the near equatorial winds we obtained reflect the flow at higher altitude than that at higher latitude. To examine it, we focus on small regions where NIR radiance is especially large as shown in Figure 3. The bright spots are called “holes”^{34,35}, since it indicates that the cloud optical thickness is much smaller than their surroundings. Figure 3a–c shows the time evolution of a narrow (~ 100 km) hole extending over 19°N – 27°N . The hole is zonally elongated with time, with the southern and northern ends advected westward by 81 and 75 m/s, respectively. The difference is likely due mainly to horizontal shear, since the hole is continuous, so it is rather unnatural to suppose that it is due to a shallow (say, 2–3 km) cloud clearing whose mean altitude changes gradually and significantly with latitude (say, from around the condensational cloud base at the northern end to around the middle cloud region at the southern end). Figure 3d–f shows a hole at 14°N advected westward by 86 m/s, which is close to the peak speed of the jet. Since it is isolated, one cannot rule out that it is associated with cloud clearing at altitude much higher than that of the other hole at 20°N (advected by 82 m/s) or the hole in Fig. 3a–c mentioned above. However, it is rather ad hoc to suppose so. Therefore, we can conclude that the jet feature is likely due mainly to horizontal wind differences. It is also noteworthy that the bright holes we mentioned above can possibly extend over a significant depth (say, > 5 km) in the lower and middle cloud regions, as can be inferred from modeling studies^{35,36}. In this case, the fact that the narrow hole persisted over more than 12 hours would either mean that the vertical shear is weak or that there is a dynamical/microphysical mechanism to maintain it against shear as in the convective clouds on the Earth. In the latter case, cloud movement would represent a mean velocity over the depth.

Figure 2c shows the absolute vorticity computed from the longitudinal-mean zonal wind shown in Fig. 2b and the planetary vorticity; here, contribution of the meridional wind component, which is found small, is not included. Its absolute value is about $2 \times 10^{-5} \text{ s}^{-1}$ at 30°N and 20°S . If we naively compute its meridional gradient ($7 \times 10^{-12} \text{ m}^{-1} \text{ s}^{-1}$), it is roughly 30% of the “beta” value (meridional gradient of the vertical component of planetary vorticity) at the equator of the Earth ($2.3 \times 10^{-11} \text{ m}^{-1} \text{ s}^{-1}$).

Cloud top winds

Figure 4 shows zonal velocity obtained from dayside cloud-top images of reflected sunlight on July 11, 2016, corresponding to a region located about 80° eastwards (upstream of the mean zonal flow). Wind speeds from the UVI (365 nm, Fig. 4a) and IR2 ($2.02 \mu\text{m}$, Fig. 4b) are consistent with each other and exhibits a weak wind-speed minimum at around the

equator, as sometimes observed³⁷. The associated meridional gradient zonal absolute vorticity is roughly 1/10 of the Earth's beta value (not shown).

Temporal variability

Figure 5 shows examples of mean zonal flow obtained with 2.26- μm IR2 nightside images at different occasions. During March 25, 2016 (Fig. 5a), zonal wind speeds are smaller than on July 11. Since the rotational angular speed does not exhibit a peak, we do not call the low-latitude wind-speed maximum as the equatorial jet. The result is consistent with previous observations^{15–22} within uncertainties. Nightside IR2 observations before July, 2016 are rather limited. A preliminary analysis for another day, April 15 (not shown), suggests that the zonal wind speeds near the equator are roughly between those on March 25 and July 11.

During August 2016, IR2 nightside observations were more frequent. The zonal wind at low latitude changes with time, but the equatorial jet seems persistent over August. A clear example can be found in Fig. 5b, which shows the results for August 13 and 15. Although there is a time difference of 2.5 days, corresponding to a half of the rotational period, equatorial jet features are quite similar, indicating their persistency. In contrast, Fig. 5c shows that zonal wind profiles obtained for August 25 and 26 have profiles markedly different compared to the other dates, and the latitude where wind speed is maximized are changed from the southern to the northern hemisphere during the interval of 1.5 days. Future wind measurements of existing data to complete the time coverage of August are expected to provide a more detailed insight about the transitions among the different wind profiles exhibited here.

We earlier investigated the holes in the July 11–12 images to examine the altitude issue. Similar characteristics are observed in the above August cases (not shown); although significant holes are not observed at the jet core, deformation of holes on its flank suggests that the jet is predominantly a horizontal feature.

The appearance and disappearance of the equatorial jet suggest that it may be subject to a type of long-term episodic event, what might also explain why it has not been reported in previous works. Actually, previous works are quite limited in terms of temporal coverage: Galileo/NIMS observed the Venus nightside only for some hours during the flyby¹⁹, and the wind measurements for low latitude with Venus Express/VIRTIS are severely limited^{20,21}. Ground-based observations helped to extend the time coverage, but measurement errors due to the low spatial resolution of the images precluded the identification of clear wind features²². Additional cloud-tracking using ground-based NIR observations with NASA's Infrared Telescope Facility (IRTF) in September, 2007 also suggest an equatorial westward flow faster than 80 m/s.

Concluding remarks

This study revealed that the zonal flow in the lower and middle cloud regions has greater variability than was previously thought. The mechanism to create the equatorial jet is yet to be elucidated. The angular speed maximum indicates that it will not be created by a simple horizontal eddy diffusivity like the one supposed in theoretical studies^{7,8,38}, which smoothes

angular velocity. Since the cloud-top level zonal wind is faster, downward momentum transport could create the equatorial jet. However, it is not straightforward to occur, since the upper cloud layer has stable stratification³⁹. Vertical or horizontal momentum transport by atmospheric waves is a candidate. Interestingly, a numerical study¹¹ showed that atmospheric thermal tides created weak equatorial jets in the middle to upper cloud layer. Similar and more pronounced equatorial jets were found recently in a Venus general circulation model⁴⁰.

The mechanism of the superrotation in the Venusian atmosphere can be separated into two basic unknowns: one, the long-term torque balance between the whole atmosphere and the solid part of Venus, and, the other, the angular-momentum redistribution within the atmosphere to create a superrotation much faster than planetary rotation. The present discovery of the equatorial jet is expected to provide valuable hints and challenges on the latter aspect.

Methods

Automated cloud tracking

We used the version 0.2.1 IR2 nightside data, the version 0.2.1.1 IR2 dayside data, and the version 0.2.2.1 UVI data. The dayside image data are mapped onto a longitude-latitude grid with a resolution of 0.125° after correcting the bore-sight direction from the original navigation data by limb fitting⁴¹, and a correction of the directional dependence of sunlight reflection is applied⁴¹. The nightside data are mapped onto the same grid, but the bore-sight direction from the original navigation data is corrected by maximizing the radial (inward, with respect to the sub-spacecraft point) component of the gradient of radiance. Before conducting cloud tracking, all the gridded data are applied with a two-dimensional band-pass filter with Gaussian functions with the sigma values 0.25° for low-pass and 3° for high-pass with latitude and $0.25^\circ/\cos\phi$ for low-pass and $3^\circ/\cos\phi$ for high-pass with longitude, where ϕ is latitude³¹ (the $\cos\phi$ factor is to compensate the length change in the circles of latitude, which is not applied in Fig. 1).

Cloud tracking was conducted with an enhanced cross-correlation method, where multiple cross-correlation surfaces are combined to estimate the horizontal velocity at a specific location³⁰, and a novel error correction/elimination method is applied³¹. Specific settings are as in our latest paper³¹ except for the template size and the number of images used along with their time interval, which depend on quality considerations and data availability.

The size of the template images (sub-regions used to compute cross-correlation) is $7.5^\circ \times 7.5^\circ$ in this study; we checked the consistency among the results with smaller template sizes. Derived cloud-motion vectors are defined at the center of the template region at the initial time (see Fig. 1a); the wind derivation is conducted every 3° (Fig. 1a). Note that we use the spatial superposition of cross-correlation surfaces associated with nearby template regions^{30,31}. As a result, a wind vector represents motion over an area greater than the template size, say 12° . No smoothing is applied to derived vectors.

The obtained velocity vectors are screened to retain only the ones with the peak correlation-coefficient values greater than 0.6 and the measures of precision, termed ϵ in our previous works^{30,31} (computed from the sharpness of cross-correlation surfaces; see below for a further explanation) smaller than 15 m/s for nighttime tracking (10 m/s for daytime tracking). Also, a post-process to eliminate inconsistent results is applied³¹. The results are verified by human-eyes in a way similar to conduct manual tracking.

The actual images used are as follows: four two-hourly images for IR2 nighttime tracking starting at 20 h UTC, July 11, 0 h UTC, August 13, 11 h UTC, August 15, 4 h UTC, August 25, or 16 h UTC, August 26. The results presented in this study are based on the images obtained with the 2.26- μm filter, but we also conducted cloud tracking with the other two filters for nighttime (1.74 and 2.32 μm) and confirmed that the results are consistent with each other. As for the nighttime tracking for March 25, we only used two 2.26- μm images at 7 and 11 h UTC, since these two are the only available 2.26- μm images on the day. We also have another image at 7 h UTC with the 1.74- μm filter, so we conducted cloud tracking with this image and the 2.26- μm image at 11 h UTC for comparison. The two tracking results are remarkably consistent with each other. Daytime cloud tracking was conducted using three two-hourly images obtained at 7, 9, 11 h UTC, July 2016 for each of UVI 365-nm and IR2 2.02- μm images.

Error estimation of the automated cloud tracking

We estimated the uncertainty of the mean zonal winds shown in Figs. 2a, 3, and 4 as follows:

$$e = \left(\left(\epsilon_u^2 \right) / N + b^2 \right)^{1/2}$$

The error bars in the figures show $\pm e$ of the estimated mean wind. Here, ϵ_u represents the zonal-direction contribution to the parameter ϵ mentioned above (specifically, ϵ_u is the greatest among ϵ_{1u} and ϵ_{2u} defined in our previous work³⁰, but it is very close to ϵ in most cases). These ϵ values are derived from the lower confidence bound of cross-correlation at the significance level of 90 %³⁰. To evaluate the uncertainty of zonal mean, ϵ_u at each wind-vector grid point is first binned over four grids (a bin consists of four consecutive grid points) by considering the oversampling as stated above. Their square values are averaged along longitude, which is expressed as $\langle \epsilon_u^2 \rangle$. It is then divided by the number of bins, N , at each latitudinal grid, which represents the effective degrees of freedom. The term b represents possible uncertainty due to incorrect positioning. We have confirmed that it is negligible for daytime images (manuscript in preparation), so it is simply set to zero. We estimated the uncertainty for the IR2 nighttime images as described in the next paragraph. The resultant b values are generally much smaller than the $\left[\langle \epsilon_u^2 \rangle / N \right]^{1/2}$ values, but it is not necessarily negligible.

The planetary extent in the IR2 nighttime images is sometimes obscure, so careful treatment is needed to evaluate the accuracy of the pointing correction. The task was undertaken by

comparing cloud tracking results conducted with two images at the time interval of two hours (for example, from two-hourly images between 8 and 16 h of a day, four wind estimates were obtained by using the image pairs taken at 8&10, 10&12, 12&14, and 14&16 h). We then computed the unbiased standard deviation of the areal mean winds equatorward of 30°N and 30°S within a day. The results over multiple days consistently suggested that the standard error associated with the combined six-hourly tracking conducted in this study is roughly equal to or smaller than 1 m/s (it is 3 m/s if two-hourly tracking is made). The error arises not only from the insufficient positioning correction but also from other factors. However, we simply attributed it to the former for safety. For consistency with the definition of ϵ , the standard error was expanded by a factor of 1.7 to comply with the 90-percentile. Therefore, we set $b = 1.7$ m/s for the six-hourly tracking (July and August, 2016) and $b = 2.5$ m/s for the four-hourly tracking (March 25, 2016).

Manual cloud tracking

A manual cloud tracking was conducted following an earlier study³³ for the nightside IR2 observation in July 11—12, 2016. The tracking was done independently, and we did not share the pre-processes including bore-sight correction.

Cloud tracking with the 2007 IRTF data

The cloud tracking was made with two methods; one is the automated tracking using the computer program described on a Website (<http://www.skycoyote.com/FITSTFlow/>), and the other is manual tracking by a human operator.

Data availability

The data that support the findings of this study are available from the Akatsuki science data archive site <http://darts.isas.jaxa.jp/planet/project/akatsuki/>. The data will also be available from NASA's Planetary Data System (<https://pds.nasa.gov/>): Atmospheres Node.

Code availability

The computer program for the automated cloud tracking is currently not available, but will be made available in the future at <http://darts.isas.jaxa.jp/planet/project/akatsuki/>.

Supplementary Material

Refer to Web version on PubMed Central for supplementary material.

Acknowledgments

We sincerely thank the numerous people who contributed to create and operate Akatsuki spacecraft. We also thank the two anonymous reviewers for providing valuable comments that helped improve this paper. This study is supported by the following grants: JSPS KAKENHI 15K17767, 16H02225, and 16H02231, 16K17816; NASA Grant NNX16AC79G; JAXA's International Top Young Fellowship (ITYF). All figures and the supplementary movie were created using the GFD-Dennou Club graphic library, DCL.

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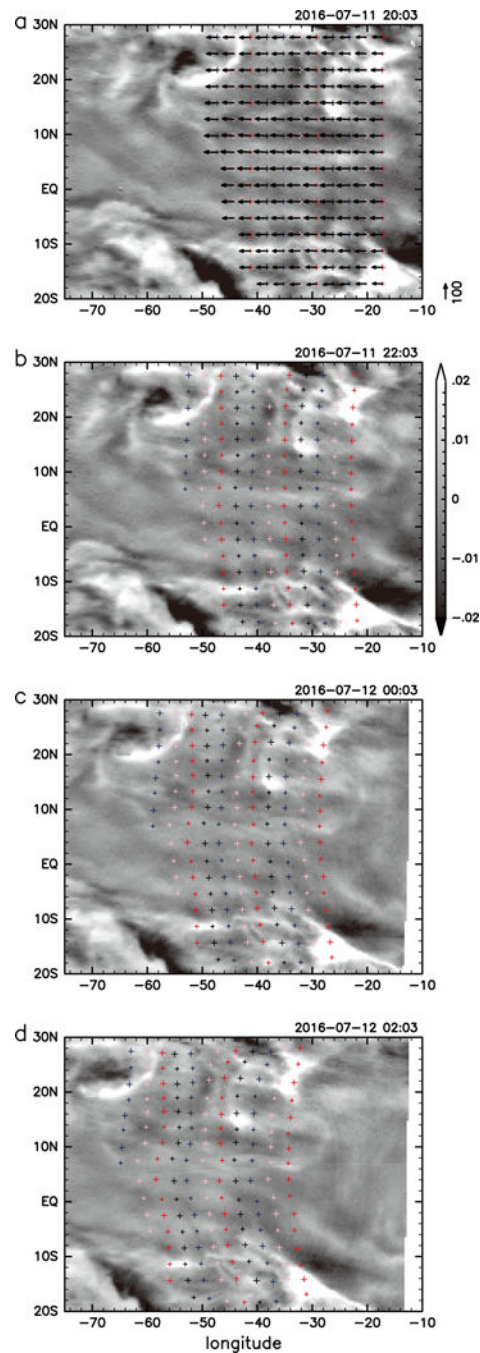


Figure 1. High-pass filtered radiance at 2.26 μm , horizontal velocities and associated trajectories for July 11-12, 2016

Grey-scale shading (**a-d**) represents the four two-hourly nightside IR2 radiance from 20:03 UTC, July 11 used for cloud tracking ($\text{W m}^{-2} \text{sr}^{-1} \mu\text{m}^{-1}$). It is high-pass filtered by the two-dimensional Gaussian filter whose half-width at half-maximum is 4° in both longitude and latitude, which is similar to the filtering for cloud tracking (see Methods). Arrows (**a**) show the horizontal velocities (scale in m/s shown on the right-hand side). Colored “+” symbols at

the initial time (**a**) show the centers of the template regions for cloud tracking, while those at later times (**b-d**) show their positions advected linearly by the velocities.

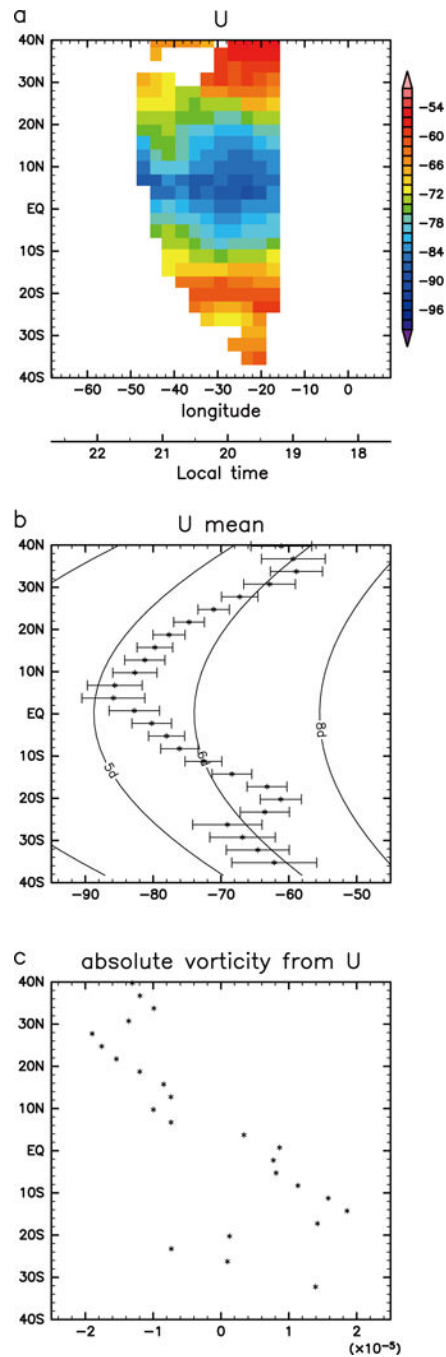


Figure 2. Zonal wind and vorticity in the lower and middle cloud regions in July 11-12, 2016
 Zonal (east-west) wind is defined as positive if it is eastward. **a:** Longitude (as well as local time)—latitude plot of zonal wind (m/s) in Fig. 1a. **b:** Longitudinal average of the zonal wind in **a**. The error bars show the uncertainty limit derived from the sharpness of cross-correlation peaks and estimated pointing accuracy (see Methods). Curves show the corresponding rotation periods: 4, 5, 6, and 8 Earth days from left to right. **c:** absolute vorticity (s^{-1}) derived only from the mean zonal wind in **b** by using the central differentiation and the planetary vorticity.

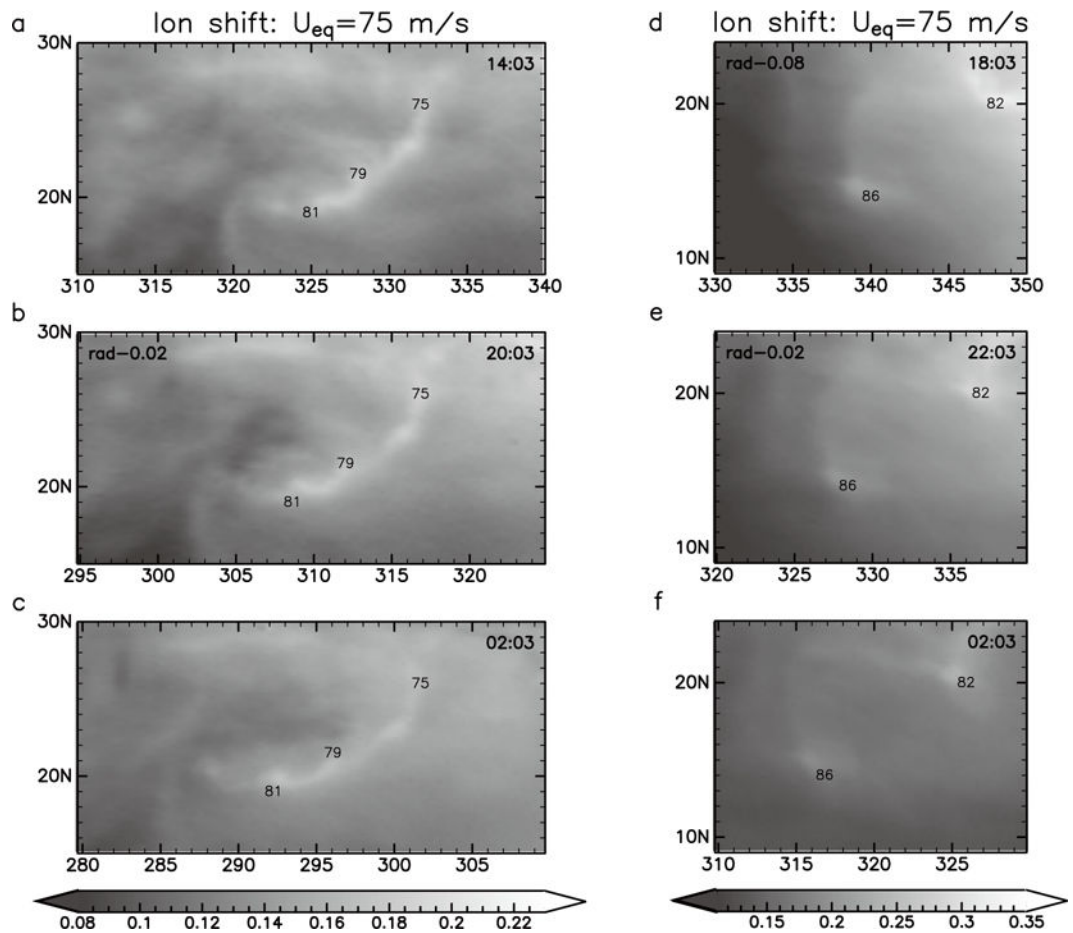


Figure 3. Close up to show the movement of radiance “holes” (bright spots) on July 11-12, 2016. Similar to Fig. 1, but gray-scale shading shows unfiltered radiance ($\text{W m}^{-2} \text{sr}^{-1} \mu\text{m}^{-1}$) over longer time-ranges (observational time is indicated in each panel). The numeric texts on bright spots (holes) such as “75” show the speeds in m/s at which they are moved westward across the panels **a,b,c** or **d,e,f** to indicate the speed the holes’ movement. In order to crudely adjust the offset owing to the leak from dayside brightness, constant values are subtracted in some panels (0.02 in **b** and **e**, 0.08 in **d**).

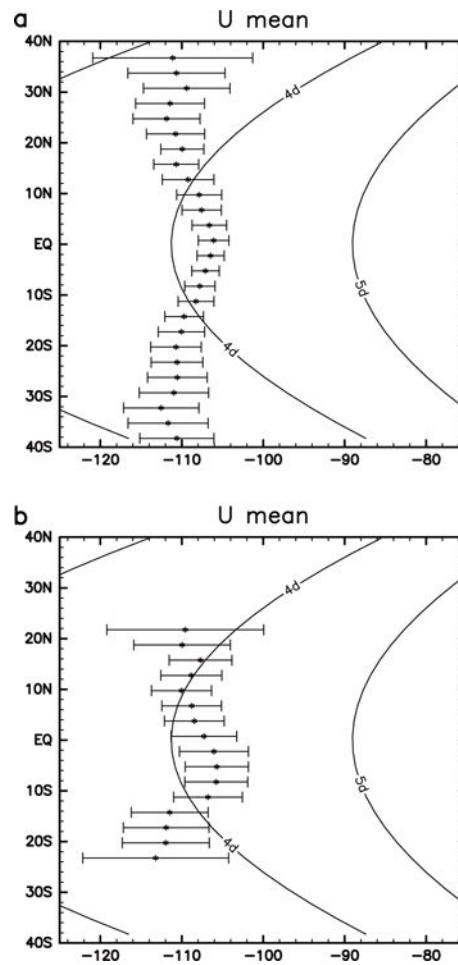


Figure 4. Zonal winds at the cloud top derived from UVI (365 nm) and IR2 (2.02 μm) data for July 11, 2016
a: Derived from three images obtained by UVI at 7, 9, and 11 h UTC, July 11, 2016 with the 365-nm filter; **b:** derived from the three images obtained by IR2 at 7, 9, and 11 UTC, July 11, 2016 with the 2.02-μm filter. Curves show the corresponding rotation periods: 3, 4, and 5 Earth days from left to right. See the caption of Fig.2b for error bars.

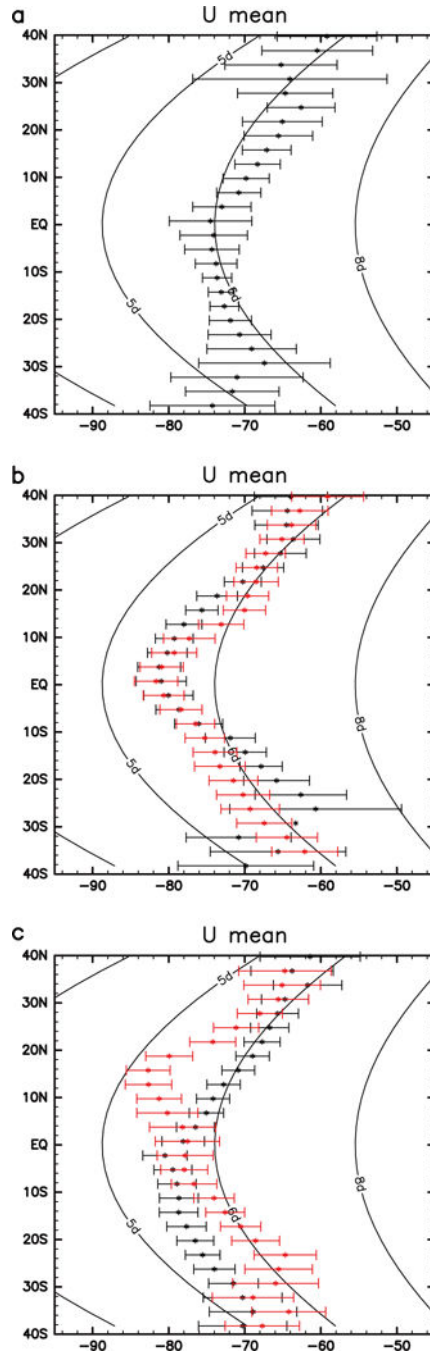


Figure 5. Zonal winds in the lower cloud at various times

The plot is as in Fig. 2b, but for the results of cloud tracking using images obtained on (a) March 25, (b) August 13 (black) and 15 (red), and (c) August 25 (black) and 26 (red). The local time ranges covered are approximately 0 to 4 h (March 25), 21:30 to 1:30 (August 13), 18:30 to 23 h (August 15), 18:30 to 23 h (August 25), and 19:30 to 24 h (August 26); the coverages for August 25 and 26 are in terms of the equator to southern hemisphere, and it is narrower in the northern hemisphere.