

EQUIVARIANT HOPF BIFURCATION FOR NEUTRAL FUNCTIONAL DIFFERENTIAL EQUATIONS

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ABSTRACT. In this paper we employ equivariant Lyapunov-Schmidt procedure to give a clearer understanding of the one-to-one correspondence of the periodic solutions of a system of neutral functional differential equations with the zeros of the reduced bifurcation map, and then set up equivariant Hopf bifurcation theory. In the process we derive criteria for the existence and direction of branches of bifurcating periodic solutions in terms of the original system, avoiding the process of center manifold reduction.

1. INTRODUCTION

Suppose that $\tau \geq 0$ is a given constant (or possibly infinity), and \mathbb{R}^n is the n -dimensional Euclidean space with norm $|\cdot|$. Let $C([-\tau, 0], \mathbb{R}^n)$ denote the Banach space of continuous mapping from $[-\tau, 0]$ into \mathbb{R}^n equipped with the supremum norm $\|\phi\| = \sup_{-\tau \leq \theta \leq 0} |\phi(\theta)|$ for $\phi \in C([-\tau, 0], \mathbb{R}^n)$. If $a \geq 0$, and $x : [-\tau, a] \rightarrow \mathbb{R}^n$ is a continuous mapping, then for any $t \in [0, a]$, $x_t \in C([-\tau, 0], \mathbb{R}^n)$ is defined by $x_t(\theta) = x(t + \theta)$ for $\theta \in [-\tau, 0]$. In this paper, we consider the following parameterized system of neutral functional differential equations (NFDEs)

$$(1.1) \quad \frac{d}{dt}h(\alpha, x_t) = f(\alpha, x_t),$$

where $h, f : \mathbb{R} \times C([-\tau, 0], \mathbb{R}^n) \rightarrow \mathbb{R}^n$ are two continuously differentiable mappings satisfying that $f(\alpha, 0) = 0$ for all $\alpha \in \mathbb{R}$. We say that (1.1) is equivariant with respect to a group Γ if there exists a representation ϱ of Γ such that

$$(1.2) \quad h(\alpha, \varrho(\gamma)\phi) = \varrho(\gamma)h(\alpha, \phi), \quad f(\alpha, \varrho(\gamma)\phi) = \varrho(\gamma)f(\alpha, \phi)$$

for $(\alpha, \gamma, \phi) \in \mathbb{R} \times \Gamma \times C([-\tau, 0]; \mathbb{R}^n)$, where $\varrho(\gamma)\phi \in C([-\tau, 0]; \mathbb{R}^n)$ is given by $(\varrho(\gamma)\phi)(s) = \varrho(\gamma)\phi(s)$ for $s \in [-\tau, 0]$. Recall that a representation ϱ of a group Γ is a group homomorphism $\varrho : \Gamma \rightarrow GL(n, \mathbb{R})$. Condition (1.2) implies that system (1.1) is invariant under the transformation $(x, t) \rightarrow (\varrho(\gamma)x, t)$. Namely, $x(t)$ is a solution of (1.1) if and only if $\varrho(\gamma)x(t)$ is a solution. Throughout this paper we always assume that Γ is a compact Lie group and system (1.1) is Γ -equivariant.

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The phenomenon of Hopf bifurcation concerns the birth of a periodic solution from an equilibrium solution through a local oscillatory instability. The existence of such a bifurcation was found in the context of ordinary differential equations (ODEs) by Hopf in 1942 [11], after whom this type of bifurcation has been subsequently named. For delay differential equations (DDEs), the first results on Hopf bifurcation date back to work by Chafee [1] in 1971. However, according to Hale [9], the first proof of the Hopf bifurcation theorem for DDEs under analytically computable conditions was presented by Chow and Mallet-Paret in 1977 [2]. The first results on equivariant Hopf bifurcation of DDEs date back to the work by Wu [14]. In [12], Krawcewicz, Ma and Wu employ equivariant degree theory developed by Geba *et al* (cf. [4]) to study the existence, multiplicity and global continuations of symmetric periodic solutions for one parameter family of NFDEs with dihedral symmetry.

In the standard Hopf bifurcation theory ([9, 10, 13]), the central hypothesis is that at the critical value of the bifurcation parameter the infinitesimal generator has a complex conjugate pair of simple purely imaginary eigenvalues. The presence of symmetry may cause purely imaginary eigenvalues to arise with higher multiplicities which causes the bifurcation problem to become more complicated, see for instance [6]. The most common approach to study bifurcation problems in function differential equations involves the computation of (normal forms of) reduced bifurcation equations on center manifolds. For example, Faria and Magalhães [3], and Weedermann [13] compute the reduced system on the center manifold associated with the pair of purely imaginary solutions of the characteristic equation.

In this paper, we present a treatment of generic codimension-one Hopf bifurcation for equivariant NFDEs on the basis of equivariant Lyapunov-Schmidt reduction, following the spirit of the treatment of Golubitsky and Stewart [5, 6] in the case of equivariant ODEs. Moreover, our results generalize those obtained by Wu [14], Krawcewicz, Ma and Wu [12]. In the process we obtain explicit expressions in terms of the original system that determine the monotonicity of the period and Hopf bifurcation direction of branches of bifurcating symmetric periodic solutions. With these expressions at our disposal, the study of equivariant Hopf bifurcation in explicit examples can be performed without having to resort to lengthy computations associated to center manifold reduction.

2. MAIN RESULTS

Let $D(\alpha), L(\alpha) : C([-\tau, 0]; \mathbb{R}^n) \rightarrow \mathbb{R}^n$ be the two linearized operators of $h(\alpha, \cdot)$ and $f(\alpha, \cdot)$, respectively. Furthermore, we assume that $D(\alpha)$ is atomic at zero (see [9] for more details). By the Riesz representation theorem, there exists an $n \times n$ matrix-valued function $\mu, \eta : [-\tau, 0] \rightarrow \mathbb{R}^{n^2}$ whose components each have bounded variation and such that for $\varphi \in C([-\tau, 0], \mathbb{R}^n)$,

$$D(\alpha)\varphi = \varphi(0) - \int_{-\tau}^0 d\mu(\alpha, \theta)\varphi(\theta), \quad L(\alpha)\varphi = \int_{-\tau}^0 d\eta(\alpha, \theta)\varphi(\theta),$$

where $\text{Var}_{[s, 0]}\mu(\alpha, \theta) \rightarrow 0$ as $s \rightarrow 0$ (see [9] for more details). For each fixed α , the linear system

$$(2.1) \quad \frac{d}{dt}D(\alpha)x_t = L(\alpha)x_t,$$

generates a strongly continuous semigroup of linear operators with infinitesimal generator \mathcal{A}_α . The spectrum of \mathcal{A}_α , denoted by $\sigma(\mathcal{A}_\alpha)$, is the point spectrum. Moreover, λ is an eigenvalue of \mathcal{A}_α , i.e., $\lambda \in \sigma(\mathcal{A}_\alpha)$, if and only if λ satisfies that $\det \Delta(\alpha, \lambda) = 0$, where the characteristic matrix $\Delta(\tau, \lambda)$ is given by

$$\Delta(\alpha, \lambda) = \lambda D(\alpha)(e^{\lambda(\cdot)} \text{Id}) - L(\alpha)(e^{\lambda(\cdot)} \text{Id}).$$

It is well-known that $\phi \in C([-\tau, 0]; \mathbb{C}^n)$ is an eigenvector of \mathcal{A}_α associated with the eigenvalue λ if and only if $\phi(\theta) = e^{\lambda\theta} b$ for $\theta \in [-\tau, 0]$ and some vector $b \in \mathbb{R}^n$ such that $\Delta(\alpha, \lambda)b = 0$. Let $E_{\alpha, \lambda}$ be the eigenspace of \mathcal{A}_α associated with the eigenvalues λ and $\bar{\lambda}$. Assume that \mathcal{A}_0 has a pair of purely imaginary eigenvalues $\pm i\beta_0$. The symmetry group Γ often causes purely imaginary eigenvalues to be multiple. So, we always assume that

(H1) \mathcal{A}_0 has a pair of purely imaginary eigenvalues $\pm i\beta_0$, each of multiplicity m , and all other eigenvalues of \mathcal{A}_0 are not the integer multiple of $i\beta_0$.

In studying the bifurcation problem we wish to consider how the eigenvalues of \mathcal{A}_α cross the imaginary axis at $\alpha = 0$ and to describe the structure of the associated eigenspace $E_{\alpha, \lambda}$. We consider the following nontrivial restrictions on the corresponding imaginary eigenspace of \mathcal{A}_0 .

(H2) The imaginary eigenspace $E_{0, i\beta_0}$ of \mathcal{A}_0 is Γ -simple.

Thus, we make use of the implicit function theorem and Lemma 1.5 in Page 265 of [6] and obtain the following results about the multiplicity of this eigenvalue and its associated eigenvectors of \mathcal{A}_α .

Theorem 2.1. *Under conditions (H1)–(H2), for sufficiently small α , infinitesimal generator \mathcal{A}_α has one pair of complex conjugate eigenvalues $\sigma(\alpha) \pm i\rho(\alpha)$, each of multiplicity m . Moreover, σ and ρ are smooth functions of α and satisfy that $\sigma(0) = 0$ and $\rho(0) = \beta_0$.*

In view of (H1), the purely imaginary eigenvalues of \mathcal{A}_0 has high multiplicity, so the standard Hopf bifurcation theorem can not be applied directly. So, we first develop the equivariant Lyapunov-Schmidt reduction for (1.1) to consider the existence of periodic solutions. Let $\omega_0 = 2\pi/\beta_0$, and \mathcal{C}_{ω_0} (respectively, $\mathcal{C}_{\omega_0}^1$) be the set of continuous (respectively, differentiable) n -dimensional ω_0 -periodic mappings. If we denote

$$\|x\|_0 = \max_{1 \leq i \leq n} \max_{t \in [0, \omega_0]} \{|x_i(t)|\}$$

for $x = (x_1, x_2, \dots, x_n)^T \in \mathcal{C}_{\omega_0}$, and $\|x\|_1 = \max\{\|x\|_0, \|\dot{x}\|_0\}$ for $x \in \mathcal{C}_{\omega_0}^1$, then \mathcal{C}_{ω_0} and $\mathcal{C}_{\omega_0}^1$ are Banach spaces when they are endowed with the norms $\|\cdot\|_0$ and $\|\cdot\|_1$, respectively. It is easy to see that \mathcal{C}_{ω_0} is a Banach representation of the group $\Gamma \times S^1$ with the action given by

$$(\gamma, \theta)u(t) = \varrho(\gamma)u(t + \theta), \quad \text{for } (\gamma, \theta) \in \Gamma \times S^1.$$

In view of complexity in analyzing NFDEs, we introduce two kinds of bilinear forms. One is the inner product $\langle \cdot, \cdot \rangle : \mathcal{C}_{\omega_0} \times \mathcal{C}_{\omega_0} \rightarrow \mathbb{R}$ defined by $\langle v, u \rangle = \frac{1}{\omega_0} \int_0^{\omega_0} \bar{v}^T(t)u(t)dt$, for $u, v \in \mathcal{C}_{\omega_0}$. The other is $(\cdot, \cdot) : C([-\tau, 0]; \mathbb{R}^n) \times C([0, \tau]; \mathbb{R}^n) \rightarrow \mathbb{R}$ defined by

$$(2.2) \quad (\psi, \varphi) = \bar{\psi}^T(0)\varphi(0) - \int_{-\tau}^0 \left[\frac{d}{ds} \int_0^s \bar{\psi}^T(\xi - s) d\mu(0, \theta)\varphi(\xi) d\xi \right]_{s=\theta} - \int_{-\tau}^0 \int_0^\theta \bar{\psi}^T(\xi - \theta) d\eta(0, \theta)\varphi(\xi) d\xi$$

for $\psi \in C([0, \tau]; \mathbb{R}^n)$ and $\varphi \in C([-\tau, 0]; \mathbb{R}^n)$. Let $\beta \in (-1, 1)$, $x(t) = u((1 + \beta)t)$. Then equation (1.1) can be rewritten as

$$(1 + \beta) \frac{d}{dt} h(\alpha, u_{t,\beta}) = f(\alpha, u_{t,\beta}),$$

where $u_{t,\beta}(\theta) = u(t + (1 + \beta)\theta)$ for $\theta \in [-\tau, 0]$. Define $F : \mathcal{C}_{\omega_0}^1 \times \mathbb{R}^2 \rightarrow \mathcal{C}_{\omega_0}$ by

$$(2.3) \quad F(u, \alpha, \beta) = -(1 + \beta) \frac{d}{dt} h(\alpha, u_{t,\beta}) + f(\alpha, u_{t,\beta}),$$

By varying the newly introduced small variable β , one keeps track not only of solutions of (1.1) with period ω_0 but also of solutions with nearby period. In fact, solutions to $F(u, \alpha, \beta) = 0$ correspond to $\frac{\omega_0}{1+\beta}$ -periodic solutions of (1.1). It follows that the Γ -equivariance of L and f that F is $\Gamma \times S^1$ -equivariant:

$$(\gamma, \theta)F(u, \alpha, \beta) = F((\gamma, \theta)u, \alpha, \beta),$$

for all $(\gamma, \theta) \in \Gamma \times S^1$. Define

$$\mathcal{L}u = -\frac{d}{dt}D(0)u_t + L(0)u_t.$$

The elements of $\text{Ker}\mathcal{L}$ correspond to solutions of the linear system $\mathcal{L}u = 0$ satisfying $u(t) = u(t + \omega_0)$. For convenience in computation we shall allow functions with range \mathbb{C}^n instead of \mathbb{R}^n . With respect to the inner product $\langle \cdot, \cdot \rangle : \mathcal{C}_{\omega_0} \times \mathcal{C}_{\omega_0} \rightarrow \mathbb{R}$, the adjoint operator of \mathcal{L} is

$$\mathcal{L}^*u = \frac{d}{dt} \left[u(t) - \int_{-\tau}^0 d\mu^T(0, \theta)u(t - \theta) \right] + \int_{-\tau}^0 d\eta^T(0, \theta)u(t - \theta).$$

It follows from (H1) that $\text{Ker}\mathcal{L} \cong E_{0,i\beta_0}$ and $\text{Ker}\mathcal{L}^* \cong E_{0,i\beta_0}^*$, both of which are $2m$ -dimensional. Here and in the sequel, \cdot^* denotes the adjoint map with respect to the inner product $\langle \cdot, \cdot \rangle$. Furthermore, we have

Lemma 2.2. *Spaces $\text{Ker}\mathcal{L}$, $\text{Range}\mathcal{L}$, and $\mathfrak{W} = (\text{Ker}\mathcal{L}^*)^\perp \cap \mathcal{C}_{\omega_0}^1$ are $\Gamma \times S^1$ -invariant subspaces of \mathcal{C}_{ω_0} . Moreover, $\mathcal{C}_{\omega_0} = \text{Ker}\mathcal{L} \oplus \text{Range}\mathcal{L}$ and $\mathcal{C}_{\omega_0}^1 = \text{Ker}\mathcal{L} \oplus \mathfrak{W}$.*

Let P and $I - P$ denote the projection operators defined by

$$P : \mathcal{C}_{\omega_0} \rightarrow \text{Range}\mathcal{L}, \quad I - P : \mathcal{C}_{\omega_0} \rightarrow \text{Ker}\mathcal{L}.$$

Obviously, P and $I - P$ are $\Gamma \times S^1$ -equivariant. Thus, $F(u, \alpha, \beta) = 0$ is equivalent to the following system:

$$(2.4) \quad PF(u, \alpha, \beta) = 0, \quad (I - P)F(u, \alpha, \beta) = 0.$$

According to the above direct sum decomposition, for each $u \in \mathcal{C}_{\omega_0}^1$, there is a unique decomposition such that $u = v + w$, where $v \in \text{Ker}\mathcal{L}$ and $w \in \mathfrak{W}$. Thus, the first equation of (2.4) can be rewritten as

$$G(v, w, \alpha, \beta) \equiv PF(v + w, \alpha, \beta) = 0.$$

Notice that $G(0, 0, 0, 0) = PF(0, 0, 0) = 0$ and $D_w G(0, 0, 0, 0) = P\mathcal{L} = \mathcal{L}$. When \mathcal{L} is restricted in \mathfrak{W} , it is invertible, and so $D_w G(0, 0, 0, 0)$ is invertible. Applying the implicit function theorem, we obtain a continuously differentiable $\Gamma \times S^1$ -equivariant map $W : \text{Ker}\mathcal{L} \times \mathbb{R}^2 \rightarrow \mathfrak{W}$ such that $W(0, 0, 0) = 0$ and

$$(2.5) \quad PF(v + W(v, \alpha, \beta), \alpha, \beta) \equiv 0.$$

Substituting $w = W(v, \alpha, \beta)$ into the second equation of (2.4), we have

$$(2.6) \quad \mathcal{B}(v, \alpha, \beta) \equiv (I - P)F(v + W(v, \alpha, \beta), \alpha, \beta) = 0.$$

Thus, we reduce our Hopf bifurcation problem to the problem of finding zeros of the map $\mathcal{B} : \text{Ker}\mathcal{L} \times \mathbb{R}^2 \rightarrow \text{Ker}\mathcal{L}$. We refer to \mathcal{B} as the bifurcation map of the system (1.1). It follows from the $\Gamma \times S^1$ -equivariance of F and W that the bifurcation map \mathcal{B} is also $\Gamma \times S^1$ -equivariant. Moreover, $\mathcal{B}(0, 0, 0) = 0$ and $\mathcal{B}_v(0, 0, 0) = 0$.

Finding periodic solutions to (1.1) rests on prescribing in advance the symmetry of the solution we seek. This can often be used to select a subspace on which the eigenvalues are simple. In addition, we should take temporal phase-shift symmetries in terms of the circle group S^1 into account as well as spatial symmetries. Here, we place emphasis on two-dimensional fixed-point subspaces and assume that

(H3) $\dim \text{Fix}(\Sigma, E_{0, i\beta}) = 2$ for some subgroup Σ of $\Gamma \times S^1$.

(H4) $\sigma'(0) \neq 0$.

Assumption (H4) is the transversality condition analogous to those of the standard Hopf bifurcation theorem. Now, we can present our main results about equivariant Hopf bifurcation.

Theorem 2.3. *Under conditions (H1)–(H4), in every neighborhood of $(x = 0, \alpha = 0)$ system (1.1) has a bifurcation of periodic solutions whose spatio-temporal symmetry can be completely characterized by Σ .*

Proof. We consider the restriction mapping $\tilde{\mathcal{B}} : \text{Fix}(\Sigma, \text{Ker}\mathcal{L}) \times \mathbb{R}^2 \rightarrow \text{Ker}\mathcal{L}$ of $\mathcal{B} : \text{Ker}\mathcal{L} \times \mathbb{R}^2 \rightarrow \text{Ker}\mathcal{L}$ on $\text{Fix}(\Sigma, \text{Ker}\mathcal{L}) \times \mathbb{R}^2$, i.e.,

$$\tilde{\mathcal{B}}(v, \alpha, \beta) = (I - P)F(v + W(v, \alpha, \beta), \alpha, \beta)$$

for $v \in \text{Fix}(\Sigma, \text{Ker}\mathcal{L})$, $\alpha \in \mathbb{R}$, and $\beta \in \mathbb{R}$. Clearly, $\tilde{\mathcal{B}}$ is also $\Gamma \times S^1$ -equivariant, and satisfies that

$$(2.7) \quad \tilde{\mathcal{B}}(0, 0, 0) = 0, \quad \tilde{\mathcal{B}}_v(0, 0, 0) = 0.$$

Moreover, it is easy to see that $\text{Range}\tilde{\mathcal{B}} \subseteq \text{Fix}(\Sigma, \text{Ker}\mathcal{L})$. Namely, $\tilde{\mathcal{B}}$ maps $\text{Fix}(\Sigma, \text{Ker}\mathcal{L}) \times \mathbb{R}^2$ to $\text{Fix}(\Sigma, \text{Ker}\mathcal{L})$. Therefore, we only need to consider the existence of nontrivial zeroes of $\tilde{\mathcal{B}}$.

Without loss of generality, assume that $\text{Fix}(\Sigma, \text{Ker}\mathcal{L}) = \text{span}\{q, \bar{q}\}$, where $q(\theta) = Ae^{i\beta_0\theta}$ and $A \in \mathbb{C}^n$ satisfies $\Delta(0, i\beta)A = 0$. Obviously, $\text{Fix}(\Sigma, \text{Ker}\mathcal{L})^* = \text{Fix}(\Sigma, \text{Ker}\mathcal{L}^*)$.

Thus, there exists $B \in \mathbb{C}^n$ such that $\overline{B}^T A = 1$ and $\text{Fix}(\Sigma, \text{Ker}\mathcal{L}^*) = \text{span}\{p, \bar{p}\}$, where $p(\theta) = Be^{i\beta_0\theta}$. As stated in Theorem 2.1, for sufficiently small α , the infinitesimal generator \mathcal{A}_α has one pair of complex conjugate eigenvalues $\lambda(\alpha) = \sigma(\alpha) \pm i\rho(\alpha)$, each of multiplicity m . Moreover, there exists a C^1 -continuous function $\tilde{A}(\alpha)$ such that $\tilde{A}(0) = A$ and $\Delta(\alpha, \lambda(\alpha))\tilde{A}(\alpha) \equiv 0$ for all sufficiently small α , we differentiate it with respect to α and obtain

$$[\Delta_\alpha(\alpha, \lambda(\alpha)) + \lambda'(\alpha)\Delta_\lambda(\alpha, \lambda(\alpha))]A(\alpha) + \Delta(\alpha, \lambda(\alpha))A'(\alpha) = 0.$$

In particular, we have

$$[\Delta_\alpha(0, i\beta_0) + \lambda'(0)\Delta_\lambda(0, i\beta_0)]A + \Delta(0, i\beta_0)A'(0) = 0.$$

This, together with the fact that $\overline{B}^T \Delta(0, i\beta_0) = 0$, implies that

$$(2.8) \quad \overline{B}^T \Delta_\alpha(0, i\beta_0)A + \lambda'(0)\overline{B}^T \Delta_\lambda(0, i\beta_0)A = 0.$$

For each $\phi \in \text{Fix}(\Sigma, \text{Ker}\mathcal{L})$, $\phi = zq + \bar{z}\bar{q}$, where $z = \langle p, \phi \rangle$. Let

$$g(z, \alpha, \beta) := \langle p, \tilde{\mathcal{B}}(zq + \bar{z}\bar{q}, \alpha, \beta) \rangle.$$

Thus, we only need to consider the existence of nontrivial solutions to $g(z, \alpha, \beta) = 0$. It follows from (2.7) that

$$(2.9) \quad g_z(0, 0, 0) = 0, \quad g_{\bar{z}}(0, 0, 0) = 0.$$

It is easy to see that $g(z, \alpha, \beta)$ is S^1 -equivariant. Using a similar arguments to that in [6], we can find two functions $\mathfrak{R}, \mathfrak{S} : \mathbb{R}^3 \rightarrow \mathbb{R}$ such that

$$(2.10) \quad g(z, \alpha, \beta) = \mathfrak{R}(|z|^2, \alpha, \beta)z + \mathfrak{S}(|z|^2, \alpha, \beta)iz.$$

It follows from $g_z(0, 0, 0) = 0$ that $\mathfrak{R}(0, 0, 0) = 0$ and $\mathfrak{S}(0, 0, 0) = 0$. Let $z = re^{i\theta}$. Then solving g is equivalent to either solve $r = 0$ or $\mathfrak{R}(r^2, \alpha, \beta) = 0$ and $\mathfrak{S}(r^2, \alpha, \beta) = 0$. In view of the implicitly defined function $W(v, \alpha, \beta)$, which vanishes through first order in $v = zq + \bar{z}\bar{q}$, we have

$$F(v + W(v, \alpha, \beta), \alpha, \beta) = -(1 + \beta) \frac{d}{dt} D(\alpha) v_{t, \beta} + L(\alpha) v_{t, \beta} + O(|z|^2).$$

Therefore,

$$\begin{aligned} F_\alpha(v + W(v, 0, 0), 0, 0) &= \Omega v_t + O(|z|^2) \\ F_\beta(v + W(v, 0, 0), 0, 0) &= \Xi v_t + O(|z|^2). \end{aligned}$$

where operators $\Omega = -\frac{d}{dt} D'(0) + L'(0)$, $\Xi v_t = \frac{\partial}{\partial \beta} [-(1 + \beta) \frac{d}{dt} D(\alpha) v_{t, \beta} + L(\alpha) v_{t, \beta}]_{\beta=0}$, $D'(0) = \frac{d}{d\alpha} D(\alpha)|_{\alpha=0}$, and $L'(0) = \frac{d}{d\alpha} L(\alpha)|_{\alpha=0}$. Notice that

$$\begin{aligned} \langle p, \Omega q \rangle &= \langle B e^{i\beta_0 t}, i\beta_0 \int_{-\tau}^0 d\mu_\alpha(0, \theta) A e^{i\beta_0(t+\theta)} + \int_{-\tau}^0 d\eta_\alpha(0, \theta) A e^{i\beta_0(t+\theta)} \rangle \\ &= \bar{B}^T \int_{-\tau}^0 [i\beta_0 d\mu_\alpha(0, \theta) + d\eta_\alpha(0, \theta)] A e^{i\beta_0 \theta} \\ &= -\bar{B}^T \Delta_\alpha(0, i\beta_0) A, \\ \langle p, \Xi q \rangle &= -\langle B e^{i\beta_0 t}, i\beta_0 e^{i\beta_0 t} - i\beta_0 \int_{-\tau}^0 d\mu(0, \theta) A [1 + i\beta_0 \theta] e^{i\beta_0(t+\theta)} \rangle \\ &\quad + i\beta_0 \bar{B}^T \int_{-\tau}^0 d\eta(0, \theta) A \theta e^{i\beta_0 \theta} \\ &= -i\beta_0 \bar{B}^T A + i\beta_0 \bar{B}^T \int_{-\tau}^0 d\mu(0, \theta) A [1 + i\beta_0 \theta] e^{i\beta_0 \theta} \\ &\quad + i\beta_0 \bar{B}^T \int_{-\tau}^0 d\eta(0, \theta) A \theta e^{i\beta_0 \theta} \\ &= -i\beta_0(p, q), \end{aligned}$$

and $(p, q) = \bar{B}^T \Delta_\lambda(0, i\beta_0) A$. It follows from (2.8) that $\langle p, \Omega q \rangle = (p, q) \lambda'(0)$. Similarly, we have $\langle p, \Omega \bar{q} \rangle = \langle p, \Xi \bar{q} \rangle = 0$. Therefore,

$$\begin{aligned} g_\alpha(z, 0, 0) &= \langle p, F_\alpha(v, 0, 0) \rangle = z(p, q) \lambda'(0) + O(|z|^2), \\ g_\beta(z, 0, 0) &= \langle p, F_\beta(v, 0, 0) \rangle = -i\beta_0 z(p, q) + O(|z|^2). \end{aligned}$$

Then

$$\begin{aligned} \mathfrak{R}_\alpha(0, 0, 0) &= \text{Re}\{\lambda'(0)(p, q)\}, & \mathfrak{S}_\alpha(0, 0, 0) &= \text{Im}\{\lambda'(0)(p, q)\}, \\ \mathfrak{R}_\beta(0, 0, 0) &= \beta_0 \text{Im}\{(p, q)\}, & \mathfrak{S}_\beta(0, 0, 0) &= -\beta_0 \text{Re}\{(p, q)\}. \end{aligned}$$

So the Jacobi determinant of functions \mathfrak{R} and \mathfrak{S} with respect to α and β is

$$\det \begin{bmatrix} \mathfrak{R}_\alpha(0, 0, 0) & \mathfrak{R}_\beta(0, 0, 0) \\ \mathfrak{S}_\alpha(0, 0, 0) & \mathfrak{S}_\beta(0, 0, 0) \end{bmatrix} = -\beta_0 \text{Re}\{\lambda'(0)\} |(p, q)|^2.$$

Thus, under condition (H4), the above Jacobi determinant is nonzero. The implicit function theorem implies that there exists a unique function $\alpha = \alpha(r^2)$ and $\beta = \beta(r^2)$ satisfying $\alpha(0) = 0$ and $\beta(0) = 0$ such that

$$(2.11) \quad \mathfrak{R}(r^2, \alpha(r^2), \beta(r^2)) \equiv 0, \quad \mathfrak{S}(r^2, \alpha(r^2), \beta(r^2)) \equiv 0$$

for all sufficient small r . Therefore, $g(z, \alpha(|z|^2), \beta(|z|^2)) \equiv 0$ for z sufficiently near 0. Therefore, system (1.1) has a bifurcation of periodic solutions whose spatio-temporal symmetry can be completely characterized by Σ . This completes the proof of Theorem 2.3. \square

Remark 2.4. Theorem 2.3 implies that a Hopf bifurcation for (1.1) occurs at $\alpha = 0$. Namely, in every neighborhood of $(x = 0, \alpha = 0)$ there is a branch of Σ -symmetric periodic solutions $x(t, \alpha)$ with $x(t, \alpha) \rightarrow 0$ as $\alpha \rightarrow 0$. The period ω_α of $x(t, \alpha)$ satisfies that $\omega_\alpha \rightarrow \omega_0$ as $\alpha \rightarrow 0$. Moreover, Γ -equivariance implies that there are $(\Gamma \times S^1)/\Sigma$ different periodic solutions, which have isotopy subgroups conjugate to Σ in $\Gamma \times S^1$.

In what follows, we consider the bifurcation direction. Assuming sufficient smoothness of h and f , we write

$$\begin{aligned} h(0, \varphi) &= D(0)\varphi + \frac{1}{2}\mathcal{B}(\varphi, \varphi) + \frac{1}{6}\mathcal{C}(\varphi, \varphi, \varphi) + o(\|\varphi\|^3) \\ f(0, \varphi) &= L(0)\varphi + \frac{1}{2}\mathcal{E}(\varphi, \varphi) + \frac{1}{6}\mathcal{F}(\varphi, \varphi, \varphi) + o(\|\varphi\|^3), \end{aligned}$$

In view of (2.5), we have $PF(zq + \bar{z}\bar{q} + W(zq + \bar{z}\bar{q}, \alpha, \beta), \alpha, \beta) \equiv 0$. Write $W(zq + \bar{z}\bar{q}, 0, 0)$ and $g(z, 0, 0)$ as

$$W(zq + \bar{z}\bar{q}, 0, 0) = \sum_{s+l \geq 2} \frac{1}{s!l!} W_{sl} z^s \bar{z}^l \quad g(z, 0, 0) = \sum_{s+l \geq 2} \frac{1}{s!l!} g_{sl} z^s \bar{z}^l.$$

It follows from (2.10) that $g_{21} = \Re_1(0, 0, 0) + i\Im_1(0, 0, 0)$, where $\Re_1(u, \alpha, \beta) = \Re_u(u, \alpha, \beta)$ and $\Im_1(u, \alpha, \beta) = \Im_u(u, \alpha, \beta)$. Therefore, $\Re_1(0, 0, 0) = \Re\{g_{21}\}$ and $\Im_1(0, 0, 0) = \Im\{g_{21}\}$. From (2.11), we can calculate the derivatives of $\alpha(r^2)$ and $\beta(r^2)$ and evaluate at $r = 0$:

$$\alpha'(0) = -\frac{\Re\{\overline{(p,q)}g_{21}\}}{\Re\{\lambda'(0)\} |(p,q)|^2}, \quad \beta'(0) = -\frac{\Im\{\lambda'(0)(p,q)\overline{g_{21}}\}}{\Re\{\lambda'(0)\} |(p,q)|^2}.$$

The bifurcation direction is determined by $\text{sign}\alpha'(0)$, and the monotonicity of period of bifurcating closed invariant curve depends on $\text{sign}\beta'(0)$. Using a similar argument as that in [5], we have

$$\begin{aligned} g_{21} &= \langle p, \mathcal{F}(q, q, \bar{q}) - \frac{d}{dt}\mathcal{C}(q, q, \bar{q}) \rangle \\ &\quad + 2 \langle p, \mathcal{E}(q, W_{11}) - \frac{d}{dt}\mathcal{B}(q, W_{11}) \rangle \\ &\quad + \langle p, \mathcal{E}(\bar{q}, W_{20}) - \frac{d}{dt}\mathcal{B}(\bar{q}, W_{20}) \rangle. \end{aligned}$$

We still need to compute W_{11} and W_{20} . In fact, it follows that

$$\begin{aligned} W_{20} &= -\mathcal{L}^{-1}P \left\{ -\frac{d}{dt}\mathcal{B}(q, q) + \mathcal{E}(q, q) \right\}, \\ W_{11} &= -\mathcal{L}^{-1}P \left\{ -\frac{d}{dt}\mathcal{B}(q, \bar{q}) + \mathcal{E}(q, \bar{q}) \right\}, \end{aligned}$$

In order to evaluate function W_{20} , we must solve the following differential equations

$$(2.12) \quad \frac{d}{dt}D(0)W_{20} - L(0)W_{20} = P \left\{ -\frac{d}{dt}\mathcal{B}(q, q) + \mathcal{E}(q, q) \right\}.$$

Note that $\mathcal{B}(q, q) = \mathcal{B}(Ae^{i\beta_0(\cdot)}, Ae^{i\beta_0(\cdot)})e^{2i\beta_0 t}$ and $\mathcal{E}(q, q) = \mathcal{E}(Ae^{i\beta_0(\cdot)}, Ae^{i\beta_0(\cdot)})e^{2i\beta_0 t}$. So, $g_{20} = \langle p, -\frac{d}{dt}\mathcal{B}(q, q) + \mathcal{E}(q, q) \rangle = 0$. Namely, $-\frac{d}{dt}\mathcal{B}(q, q) + \mathcal{E}(q, q) \in \text{Range}\mathcal{L}$. Hence, the projection P on $-\frac{d}{dt}\mathcal{B}(q, q) + \mathcal{E}(q, q)$ acts as the identity, and (2.12) is an inhomogeneous difference equations with constant coefficients. Thus, there is a

particular solution of (2.12) of the form $W_{20}^*(t) = D_2 e^{2i\beta_0 t}$. Substituting W_{20}^* into (2.12) and comparing the coefficients, we obtain

$$(2.13) \quad D_2 = \Delta^{-1}(0, 2i\beta_0) \left\{ \mathcal{E}(Ae^{i\beta_0(\cdot)}, Ae^{i\beta_0(\cdot)}) - 2i\beta_0 \mathcal{B}(Ae^{i\beta_0(\cdot)}, Ae^{i\beta_0(\cdot)}) \right\}$$

In addition, W_{20}^* is orthogonal to p , so it belongs to $\text{Range } \mathcal{L}$. Thus $W_{20}(0, 0, 0)$ is equal to W_{20}^* with D_2 determined by (2.13). Similarly, we have

$$g_{02} = g_{11} = 0, \quad W_{02} = \bar{D}_2 e^{-2i\beta_0 t}, \quad W_{11} = D_0,$$

where $D_0 = \Delta^{-1}(0, 0) \mathcal{E}(Ae^{i\beta_0(\cdot)}, \bar{A}e^{-i\beta_0(\cdot)})$. Therefore,

$$\begin{aligned} g_{21} &= \bar{B}^T \mathcal{F}(Ae^{i\beta_0(\cdot)}, Ae^{i\beta_0(\cdot)}, \bar{A}e^{-i\beta_0(\cdot)}) - i\beta_0 \bar{B}^T \mathcal{C}(Ae^{i\beta_0(\cdot)}, Ae^{i\beta_0(\cdot)}, \bar{A}e^{-i\beta_0(\cdot)}) \\ &\quad + 2\bar{B}^T \mathcal{E}(Ae^{i\beta_0(\cdot)}, D_0) - 2i\beta_0 \bar{B}^T \mathcal{B}(Ae^{i\beta_0(\cdot)}, D_0) \\ &\quad + \bar{B}^T \mathcal{E}(\bar{A}e^{-i\beta_0(\cdot)}, D_2) - i\beta_0 \bar{B}^T \mathcal{B}(\bar{A}e^{-i\beta_0(\cdot)}, D_2). \end{aligned}$$

We summarize the above discussion as follows.

Theorem 2.5. *In addition to conditions (H1)–(H4), assume that $L(\alpha)$ and $f(\alpha, \cdot)$ are sufficiently smooth. Then there exists a branch of Σ -symmetric periodic solutions, parameterized by α , bifurcating from the trivial solution $x = 0$ of (1.1). Moreover,*

- (i) $\text{Re}\{\lambda'(0)\} \text{Re}\{\overline{(p, q)}g_{21}\}$ determines the direction of the bifurcation: the bifurcation is supercritical (respectively, subcritical), i.e. the bifurcating periodic solutions exist for $\alpha > 0$ (respectively, < 0), if $\text{Re}\{\lambda'(0)\} \text{Re}\{\overline{(p, q)}g_{21}\} < 0$ (respectively, > 0). and
- (ii) $\text{Re}\{\lambda'(0)\} \text{Im}\{\lambda'(0)(p, q)\overline{g_{21}}\}$ determines the period of the bifurcating periodic solutions along the branch: the period is greater than (respectively, smaller than) ω_0 if it is positive (respectively, negative).

Example 2.6 (Application of Theorems 2.3 and 2.5 to ODEs). Consider the following system

$$(2.14) \quad \begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = \begin{bmatrix} \alpha & \beta_0 \\ -\beta_0 & \alpha \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} + \begin{bmatrix} f^1(x, y) \\ f^2(x, y) \end{bmatrix}$$

where $\beta_0 > 0$, f^j is three times differentiable and satisfying $f_x^j(0, 0) = f_y^j(0, 0) = 0$, $j = 1, 2$. It is easy to see that if $\alpha = 0$, then at the origin $(0, 0)$, the linearized system of (2.14) has a solution basis $\{e_1, e_2\}$, where $e_1(t) = (\cos(\beta_0 t), -\sin(\beta_0 t))^T$ and $e_2(t) = (\sin(\beta_0 t), \cos(\beta_0 t))^T$. If we choose $q = 2p = (1, i)^T e^{i\beta_0(\cdot)}$ and denote $\text{Re}\{g_{21}\}$ by Υ . Then, $(p, q) = 1$ and

$$\begin{aligned} \Upsilon &= f_{xxx}^1 + f_{xyy}^1 + f_{xxy}^2 + f_{yyy}^2 \\ &\quad + \frac{1}{\beta_0} [f_{xy}^1 (f_{xx}^1 + f_{yy}^1) - f_{xy}^2 (f_{xx}^2 + f_{yy}^2) - f_{xx}^1 f_{xx}^2 + f_{yy}^1 f_{yy}^2]. \end{aligned}$$

It follows from Theorem 2.5 that there exists a branch of periodic solutions of (2.14) bifurcating from the trivial solution $x = 0$, and that Υ determines the directions of the Hopf bifurcation: the Hopf bifurcation is supercritical (subcritical), i.e. the bifurcating periodic solutions exist for $\alpha > 0$ (respectively, < 0) if $\Upsilon < 0$ (respectively, > 0). It should be noticed that the above Υ is exactly the one derived by Hassard and Wan [10] by using center manifold reduction and normal form approach (for more details, see [7]).

3. APPLICATION TO A RING NETWORK

To illustrate the results presented above, we consider a ring network consisting of n identical elements with time delayed nearest-neighbour coupling [12]:

$$(3.1) \quad [u_j(t) - cu_j(t-1)]' = g(u_{j+1}(t-1)) + g(u_{j-1}(t-1)) - 3g(u_j(t-1)),$$

where $i \pmod n$, $g \in C^3(\mathbb{R}; \mathbb{R})$ with $g(0) = g''(0) = 0$ and $g'(0) = b > 0$, and $c \in [0, 1)$ is the bifurcation parameter. Define the action of the dihedral group \mathbb{D}_n on \mathbb{R}^n by

$$(3.2) \quad (\rho \cdot u)_j = u_{j+1} \quad \text{and} \quad (\kappa \cdot u)_j = u_{2-j}$$

for all $j \pmod n$ and $x \in \mathbb{R}^n$. It is easy to see that system (3.1) is \mathbb{D}_n -equivariant (see [8, 12]). Let $\mathcal{A}(c)$ be the infinitesimal generator of linear operator generated by the linearization of (3.1) about the trivial solution $u = 0$. It can be shown [12] that $\lambda \in \mathbb{C}$ is an eigenvalue of $\mathcal{A}(c)$ if and only if $\prod_{j=0}^{n-1} p_j(\lambda, c) = 0$, where $p_j(\lambda, c) = \lambda + (\vartheta_j - c\lambda)e^{-\lambda}$ and $\vartheta_j = b + 4b \sin^2(2j\pi/n) > 0$.

For a given j , $p_j(\cdot, c)$ has a pair of purely imaginary zeros $\pm\beta_0$ if $c = \cos \beta_0$ and $\vartheta_j = \beta_0 \sin \beta_0$. This results in a family of bifurcation values $c_{j,k}$ in the interval $[0, 1)$, where $c_{j,k} = \cos \beta_{j,k}$ for $k \in \mathbb{N}$, and $\{\beta_{j,k}\}_{k=1}^\infty$ is a strictly increasing sequences of positive numbers satisfying $\vartheta_j = \beta_{j,k} \sin \beta_{j,k}$ for all $k \in \mathbb{N}$ and $\lim_{k \rightarrow \infty} \beta_{j,k} = \infty$. Moreover, if $\lambda(c)$ is a smooth curve of zeros of $p_j(\cdot, c)$ with $\lambda(c_{j,k}) = i\beta_{j,k}$, it is easy to see that

$$(3.3) \quad \lambda'(c_{j,k}) = D(c_{j,k}\beta_{j,k}^2 - i\vartheta_j\beta_{j,k}),$$

where $D = |(\cos \beta_{j,k} + i \sin \beta_{j,k})(1 + \beta_{j,k}) - c_{j,k}|^{-2}$. Therefore, for fixed j and k , $\mathcal{A}(c_{j,k})$ has a pair of purely imaginary eigenvalues $\pm i\beta_{j,k}$ with the associated eigenspace E_0 spanned by the eigenvectors $e^{i\beta_{j,k}(\cdot)}v_j$, $e^{i\beta_{j,k}(\cdot)}\bar{v}_j$, $e^{-i\beta_{j,k}(\cdot)}v_j$ and $e^{-i\beta_{j,k}(\cdot)}\bar{v}_j$, where $v_j = (1, e^{2\pi i/n}, \dots, e^{2(n-1)j\pi/n})$. Thus, assumptions (H1), (H2), and (H4) hold. If $j \neq 0$ and $j \neq n/2$, it furthermore follows from [8] that

$$(3.4) \quad \begin{aligned} \text{Fix}(\Sigma_\kappa^+) &= \text{span}\{w_1 \cos(\beta_{j,k}t), w_1 \sin(\beta_{j,k}t)\}, \\ \text{Fix}(\Sigma_\kappa^-) &= \text{span}\{w_2 \cos(\beta_{j,k}t), w_2 \sin(\beta_{j,k}t)\}, \\ \text{Fix}(\Sigma_\rho^+) &= \text{span}\{\text{Re}(\bar{v}_j e^{i\beta_{j,k}t}), \text{Im}(\bar{v}_j e^{i\beta_{j,k}t})\}, \\ \text{Fix}(\Sigma_\rho^-) &= \text{span}\{\text{Re}(v_j e^{i\beta_{j,k}t}), \text{Im}(v_j e^{i\beta_{j,k}t})\}. \end{aligned}$$

where $w_1 = \text{Re}(v_j)$, $w_2 = \text{Im}(v_j)$, and $\Sigma_\kappa^\pm = (\kappa, \pm 1)$ and $\Sigma_\rho^\pm = (\rho, e^{\pm 2ij\pi/(n\beta_{j,k})})$ are subgroups of $\mathbb{D}_n \times S^1$. Thus, all conditions of the equivariant Hopf bifurcation theorem (Theorem 2.3) are satisfied. Therefore, we apply Theorem 2.3 to system (3.1) and obtain the following results.

Theorem 3.1. (i) Near $c = c_{0,k}$ for each $k \in \mathbb{N}$, there exists a branch of synchronous periodic solutions of period ω near $(2\pi/\beta_{0,k})$ bifurcated from the zero solution of the system. (ii) Near $c = c_{j,k}$ for each $j \in \{1, 2, \dots, [(n-1)/2]\}$ and $k \in \mathbb{N}$, there exist $2(n+1)$ branches of asynchronous periodic solutions of period ω near $(2\pi/\beta_{j,k})$ bifurcated from the zero solution of the system and these are two phase-locked waves, n mirror-reflecting waves, and n standing waves.

In what follows, we start with the two phase-locked oscillations mentioned above, which are characterized by Σ_ρ^\pm . In view of (3.4), for the vectors A and B defined in Section 2, we choose $A = nB = \bar{v}_j$ or $A = nB = v_j$. We have

$$(3.5) \quad \text{Re}\{\overline{(p, q)g_{21}}\} = \vartheta_j^2 g'''(0) \quad \text{and} \quad \text{Im}\{\lambda'(0)(p, q)\overline{g_{21}}\} = -g'''(0)\vartheta_j\beta_{j,k}.$$

Similarly, we choose $A = \frac{n}{2}B = w_1$ for the the mirror-reflecting waves characterized by Σ_{κ}^+ , and $A = \frac{n}{2}B = w_2$ for the the mirror-reflecting waves characterized by Σ_{κ}^- . By a direct computation, we have

$$(3.6) \quad \operatorname{Re}\{\overline{(p, q)g_{21}}\} = \frac{m}{4}\vartheta_j^2 g'''(0) \text{ and } \operatorname{Im}\{\lambda'(0)(p, q)\overline{g_{21}}\} = -\frac{m}{4}g'''(0)\vartheta_j\beta_{j,k}.$$

where $m = 4$ if $4j = 0 \pmod{n}$, or $m = 3$ otherwise.

Finally, for the synchronous periodic solution mentioned in Theorem 3.1, we can show that (3.5) holds for $j = 0$. Thus, applying Theorem 2.5, we have the following results.

Theorem 3.2. *Near $c = c_{j,k}$ for each $j \in \{0, 1, \dots, n-1\}$ and $k \in \mathbb{N}$, system (3.1) undergoes a Hopf bifurcation, both the bifurcation direction and the period of bifurcating periodic solutions are determined by the sign of $g'''(0)$. More precisely, if $g'''(0) < 0$ (or > 0) then (i) the Hopf bifurcation is supercritical (respectively, subcritical) and all the bifurcating periodic solutions exist for $c > c_{j,k}$ (respectively, $< c_{j,k}$); (ii) the period of each branch of bifurcating periodic solutions is greater (respectively, smaller) than $(2\pi/\beta_{j,k})$.*

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