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Evaluation of the modified S -transform for time-frequency synchrony analysis and source localisation

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Abstract

This article considers the problem of phase synchrony and coherence analysis using a modified version of the S -transform, referred to here as the Modified S -transform (MST). This is a novel and important time-frequency approach to study the phase coupling between two or more different spatially recorded entities with non-stationary characteristics. The basic method includes a cross-spectral analysis to study the phase synchrony of non-stationary signals, and relies on some properties of the MST, such as phase preservation. We demonstrate the usefulness of the technique using simulated examples and real newborn EEG data. The results show the advantage of using the cross-MST in the study of the connectivity between different signals using the time-frequency coherence. The MST led to improvements in resolution of almost twofold over the standard S -Transform in the examples presented in the article.

Keywords: time-frequency coherence, cross time-frequency, modified S -transform, phase synchrony, array processing, EEG signal, time-frequency signal analysis, Quadratic Time-Frequency Distributions, instantaneous frequency

1 Introduction

1.1 Time-frequency methods

Non-stationary signals have statistical properties that vary with time and hence the traditional time averaged amplitude spectrum obtained using Fourier transform is inadequate to track changes in signal magnitude, frequency or phase. In analyzing non-stationary and multi-component signals, time-frequency-based techniques were shown to outperform classical techniques based on either time or frequency domains [1] (Chapter 1). The basic idea of time-frequency analysis is to understand and describe situations where the frequency content of a signal is changing in time. Although time-frequency analysis had its origin almost 50 years ago, significant advances have occurred in the past 20 years or so. In particular, the time-frequency representation has received considerable attention as a powerful high resolution and precision tool for analyzing a variety of bio-signals and systems such as speech, ECG, EEG, PCG, EMG, as well as signals arising from other fields [2]. A

time-frequency distribution (TFD) is used to analyze and process non-stationary signals in the joint time-frequency domain. Several TFDs exist in the literature [1]. Most of them are based on the Wigner-Ville distribution (WVD) [3], as all the other TFDs can be expressed as a smoothed version of the WVD. A popular candidate of this class is the spectrogram, which is the square modulus of the short time Fourier transform (STFT). The spectrogram is the WVD smoothed in time and frequency by the ambiguity function of the window used in the STFT [4]; and all the quadratic TFDs (QTFD) are 2D-smoothed versions of the WVD (and therefore of the spectrogram which is the square modulus of the STFT). The spectrogram has been widely used as an initial investigative tool [1] as it has the property of being a cross-terms free TFD; but it also suffers from the undesirable trade-off between time concentration and frequency concentration. To address the problem of cross-terms suppression, while keeping a high time-frequency resolution, other TFDs have been proposed. Among these, one can cite the smoothed pseudo WVD (SPWVD) [1], the Cone shaped kernel TFD (ZAMD) [5], Gaussian TFD [6] and the Modified B -distribution

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(MBD) [1] (Chapter 3) just to name a few. These methods form a class commonly referred to as QTFDs or Bilinear TFDs, as they represent and distribute the energy of a signal in the time-frequency domain. A different method, the continuous wavelet transform (CWT) has a wide area of application and can be seen as an extension of the spectrogram in a wide sense, except for representing signals in time-scale space instead of time-frequency space. As it is linear and based on amplitude decomposition, it is also a useful technique for filtering in the time-scale space. A more recently developed method called *S*-transform (ST) is conceptually a hybrid of the STFT and CWT. The ST uses a variable analyzing window length but preserves the phase information by using a Fourier kernel in the signal decomposition [7].

The above indicates that various methods of performing a time-frequency analysis exist. No method is superior to another, since they all have benefits for specific types of signals and applications to which they are adapted, depending on signal characteristics such as spectral contents.

The novel contributions in this article are mostly based on

- The design of the modified ST (a signal-dependent version of the standard ST with an improved time-frequency resolution).
- The design and application of the cross-modified ST to investigate the phase synchrony between signals in the time-frequency plane and source location.

1.2 Phase synchrony and source location

Phase synchrony analysis is a useful measure of linear dependence between two stochastic signals. This approach is based on the concept of phase synchronization of chaotic oscillators studied by Rosenblum et al. [8]. The phase synchrony (coefficient) takes on values between 0, for two signals at different frequencies, and 1, for signals that exhibit a constant difference in instantaneous phase (representing the situation where a signal and its time-shifted version are observed). So, phase synchrony refers to the interdependence between the instantaneous phases of two signals; the instantaneous phases may be strongly synchronized even when the amplitudes of the two signals are statistically independent [8]. Estimating this measure requires pre-filtering at a frequency of interest. The interesting properties of the phase synchrony coefficient are its independence of the signal amplitudes and that no assumptions about the nature of the signals are made. Another useful measure of linear dependence between two stochastic signals is the coherence function, which quantifies linear

correlations in frequency domain based measure between two wide sense stationary signals [9,10]. The magnitude squared coherence function “measures the degree to which one process can be represented as the output of a linear filter operating on the other process” [10] and varies from 0—for two statistically independent processes—to 1, when one process is the result of linear filtering performed on the other. The phase coherence is usually interpreted as a phase lead of one signal over the other; it finds many useful applications but the results based on coherence depend on several factors like stationarity of the signal, segment length, number of segments, etc. [11]. It is necessary to emphasize that although phase coherence and phase synchrony are quite similar and often are mixed up; they are two principally different measures. Phase coherence can be interpreted as phase shifts and amplitude changes over frequency between two correlated sequences, while phase synchrony indicates whether the phase shift is close to a constant over the specified time interval. This interpretation justifies narrow-band filtering in the case of phase synchrony. The concept of a phase shift (either lead or lag) between two signals is only applicable when both signals are at the same frequency. One way to accomplish phase synchrony is via the Hilbert transform (HT); which consists of transforming the original signal, using an auxiliary function, into a complex-valued signal called an analytical form, from which the instantaneous phase is easily obtained [12]. This method relies on the assumption that the signal is composed of a narrow band of frequencies. Hence, it requires the bandpass filtering of the signal around a frequency of interest and then applies the HT to obtain the instantaneous phase. Another approach to estimate the phase synchrony is by performing a time-varying complex energy spectrum using either the CWT with a complex Morlet wavelet [13] or the short-time Fourier transform (STFT) [14]. Being Gaussian and well localised in time and frequency domains, the Morlet wavelet has an optimal time and frequency resolution [13]. Basically, both CWT and STFT outperform the HT-based methods with the prior giving higher resolution phase synchrony estimates over time and frequency, especially at the low frequency range [8], as they take into account the non-stationarity of the signal. So, in contrast to HT, the wavelet-based function has compact support in such way that it is defined only for those frequencies close to the frequency of interest, so it is equivalent to band-pass filtering at this frequency, which makes the pre-filtering unnecessary. However, the phase information resulting from the wavelet transform is currently not fully understood, and is largely dependent on the specific wavelet, as it is essentially a time-scale method and not a time-frequency method; it has the useful property of being

linear in time for sinusoidal signals, and can thus be applied for phase synchrony. Although the wavelet and STFT-based phase synchrony approaches consider the nonstationarity issue, they suffer from a number of drawbacks. In the case of the wavelet transform, where we have a progressive time-frequency resolution the estimate won't be the same for the whole bandwidth of the signal, which means a non uniform time-frequency tiling on the analyzed signal and thus results in biased energy representations and corresponding phase estimates. In the case of STFT, there is a trade-off between time and frequency resolution due to the window function. For these reasons, there is a need for a higher time-frequency resolution phase distributions that can track dynamic changes in phase synchrony over the whole time-frequency plane. As cited above, and given a measurements made from an array of electrodes or sensors, phase synchronization methods are used to detect the phase difference between two signals from two different electrodes or sensors, this can be applied for source location. First, appropriate neighborhood of each channel or electrode is constructed. Then phase synchronization is measured between the channel and each neighbor. Only those channels have been identified as closest to the potential sources where strong phase synchronization is detected will be considered for source location.

The article is organized as follow: The Section 2 introduces the ST definition, its key properties and provides a brief comparison between ST, CWT, and QTFD will be also reviewed. The Section 3 presents the concept and formulation of generalized ST and shows the improved resolution using the modified ST on synthetic examples. The Section 4 discusses the Cross-MST for time-frequency coherence analysis on synthetic examples. The Section 5 demonstrates the application of the cross-MST for phase synchrony detection to simulated EEG seizure and real newborn EEG data.

2 Advantages and limitations of the ST

The advent of time-frequency analysis techniques using QTFDs, STFT, and wavelet transforms made the analysis of non-stationary signals more precise. The low resolution of the STFT and the absence of phase information in the CWT led to the development of the ST, which has the property of retaining the absolute phase information, while preserving a good time-frequency resolution for all frequencies. Even though the ST has better time-frequency resolution compared to STFT, the resolution is far from perfect and needs improvement. The CWT uses a basis function which dilates and contracts with frequency; it does not retain the absolute phase information; further, the visual analysis of the time-scale plots that are produced by the

CWT is intricate. The ST combines the good features of STFT and CWT and can be viewed either as a variable sliding window STFT or as a phase-corrected CWT [7], as detailed below.

2.1 Definition and key relationships

Given a time series $x(t)$, the local spectrum at time $t = \tau$ can be determined by multiplying $x(t)$ with a Gaussian window located at $t = \tau$ and taking the Fourier transform of the product. This defines the ST as follows: [7]

$$S(t, f, \sigma) = \int_{-\infty}^{+\infty} x(\tau)g(t - \tau, \sigma)e^{-j2\pi f\tau} d\tau \quad (1)$$

The original ST uses $g(t - \tau, \sigma)$ as the scaled Gaussian window whose midpoint is $\tau = t$. At any time t and frequency f , the ST can be seen as a set of localised Fourier coefficients, obtained by considering only the portion of the primary function lying within a few cycles on either side of $\tau = t$. The scaled contraction of $g(t - \tau, \sigma)$ causes the relevant range of τ to become more localised around t as f increases.

$$g(t - \tau, \sigma) = \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{(t - \tau)^2}{2\sigma^2}} \quad (2)$$

If we select $\sigma = \frac{1}{|f|}$, Equation (2) becomes

$$g(t - \tau, 1/|f|) = \frac{|f|}{\sqrt{2\pi}} e^{-\frac{(t - \tau)^2 f^2}{2}} \quad (3)$$

and its Fourier transform with respect to τ (which is gaussian as well) becomes

$$G(v, f) = e^{-\frac{2\pi^2 v^2}{f^2}} \quad (4)$$

By choosing $\sigma = \frac{1}{|f|}$ makes $g(t - \tau, \sigma)$ with two control parameters rather than three and add some constraints to the frequency. Hence, the ST can be defined as an adaptative STFT or as a CWT with a specific mother wavelet-like term multiplied by a phase factor.

$$S(t, f) = e^{-j2\pi ft} WT(t, d) \quad (5)$$

where the scale factor d is inversely proportional to the frequency f and

$$WT(t, d) = \int_{-\infty}^{+\infty} x(\tau)g(t - \tau, d)d\tau \quad (6)$$

is the CWT of a function $x(t)$ with a mother wavelet-like term $g(t - \tau, f)$. Using Equations (3) and (5), this term can be defined as

$$g(t - \tau, f) = \frac{|f|}{\sqrt{2\pi}} e^{-\frac{f^2(t-\tau)^2}{2}} e^{-j2\pi f(t-\tau)} \quad (7)$$

The ST separates the mother wavelet-like into two parts, the slowly varying envelope (the Gaussian function) which localises in time the amplitude modulated (AM) component, and the oscillatory exponential kernel $e^{-j2\pi ft}$ which selects the frequency being localised, the frequency modulated (FM) component. Hence, this separate AM-FM decomposition is similar in concept to the empirical mode decomposition (EMD), which is a local and fully data-driven technique aimed at decomposing non-stationary multicomponent signals in “intrinsic” AM-FM contributions [15]. The ST analyzing window is actually not a mother wavelet in a strict sense. It is the time localizing Gaussian parameter that is translated while the oscillatory exponential kernel remains stationary. By not translating the oscillatory exponential kernel in Equation (7), the ST localises the real and the imaginary components of the spectrum independently, therefore localizing the phase spectrum as well as the amplitude spectrum. This is referred to as absolutely referenced phase information [7]. Unlike the CWT, the ST produces a time-frequency representation instead of a time-scale representation.

2.2 The ST and signal spectrum

Integrating the ST over time results in the Fourier transform. This direct relation to the Fourier transform and the linearity property of the ST make both the inversion to time domain and the filtering process in the time-frequency domain possible. The ST is thus a representation of local spectrum; The relationship between the ST and the signal spectrum $X(f)$ can be expressed as follows:

$$\int_{-\infty}^{+\infty} S(t, f) dt = X(f) \quad (8)$$

where $X(f)$ is the Fourier transform of $x(t)$. Hence, $x(t)$ can be retrieved from ST by:

$$x(t) = \int_{-\infty}^{+\infty} \left\{ \int_{-\infty}^{+\infty} S(t, f) dt \right\} e^{j2\pi ft} df \quad (9)$$

This relationship shows that the concept of the ST is different from both the wavelet transform and QTFDs. The concept of recovering back the signal from the ST time-frequency plane is due to the fact that the phase is referenced at the origin, which means that the phase information given by the ST refers to the argument of the sinusoid at zero time as it is the case for Fourier transform. In contrast, The phase spectrum of the signal

can be extracted from its wavelet transform CWT, as a convolution of the signal with a complex wavelet, where the phase of the CWT is relative to the center (in time) of the analyzing wavelet (the mother wavelet). Thus as the wavelet translates, the reference point of the phase translates giving a locally referenced phase which is distinguishable from the phase properties of the ST and the signal can't be recovered from the time-scale plane. On the other hand, the fact that QTFDs are bilinear and real valued energy distributions describing the energy of the signal over time and frequency, simultaneously, they cannot be used for estimating the phase of an individual signal and the phase synchrony between two signals. Moreover, the inverse process also is not straightforward, as it requires some initial conditions.

The relation between the ST and the signal spectrum $X(f)$ is given in [7] and can be written as

$$S(t, f) = \int_{-\infty}^{+\infty} X(v + f) e^{-\frac{2\pi^2 v^2}{f^2}} e^{j2\pi vt} dv, \quad f \neq 0 \quad (10)$$

The equivalent discrete version of Equation (10) can be used to compute the discrete ST by taking advantage of the computational efficiency of the fast Fourier transform (FFT) and the convolution theorem.

2.3 Comparison of ST, CWT, and QTFDs

The similarities and differences between the ST, CWT, and QTFDs are reviewed below.

(1) *ST and Morlet wavelet*: Although the differences between the ST and Morlet wavelet are suitable, the only difference between ST and Morlet wavelet decomposition is that the ST time-frequency function is scaled by the carrier frequency f . Additionally, the amplitude modulation term f smoothes the ST spectrum, but may cause computational artifacts in the Morlet spectrum at low frequencies. Furthermore, the phase correction term in ST enable the ST to preserve absolute referenced phase information of a signal. Also, ST approach has advantages over the Morlet wavelet approach in terms of easy interpretation and fast computation.

(2) *Sampling frequency*: The discrete Fourier transform (DFT) has a very well defined sampling of the frequencies, in order to be both complete and orthonormal. The discrete ST has the identical sampling of the frequency space and retains the sampling of the signal in the time domain, Similar to a STFT this is redundant. On the other hand, CWT normally employs an octave scaling for frequencies, which results in an oversampled representation at the low frequencies and an under sampled representation at the higher frequencies. QTFDs have a sampling that is similar to the DFT.

(3) *Invertibility and recoverable signal parameters*: As the ST output is complex, we can read the amplitude $A(t) = \text{abs}(S(t, f))$, the frequency f , and the phase information $\phi(t) = \text{atan}\left(\frac{\text{Im}(S(t, f))}{\text{Re}(S(t, f))}\right)$ for each time step from the ST. This allows us to extract the corresponding signal by reproducing it in the time domain as follows: $\text{signal}(t) = A(t)\cos(2\pi f(t)t + \phi(t))$. The signal parameters were recoverable due to the combination of absolutely referenced phase information and frequency invariant amplitude of the ST, and such direct extraction cannot be done with CWT-based method. Moreover, the direct measurement of phase information makes the ST a potential candidate to estimate the phase synchrony between two signals.

For QTFDs, a signal $x(t)$ can be recovered exactly, apart from a complex scaling factor [1, p. 61].

(4) *ST phase*: The ST retains the absolute phase information, where as the phase information is lost in the CWT. The absolutely referenced (at time $t = 0$) phase of the ST leads to a generalization of the instantaneous frequency to broadband signals and can be used as a local peak-finding algorithm [16,17], for example. In contrast, in the wavelet approach; the phase is relative to the center (in time) of the analyzing wavelet, and as the wavelet translates, the reference point of the phase translates, hence the relative phase becomes meaningless. With the ST, the sinusoidal component of the basis function remains stationary, while the Gaussian envelope translates in time. Thus, the reference point for the phase remains stationary and the phase has the same meaning as in the Fourier domain [18]. For QTFDs, the phase information is linked to the time delay or instantaneous frequency [12]. The phase information can be extracted using the cross-WVD or related methods [19].

(5) *ST amplitude*: The time domain localizing window (the Gaussian function) in Equation (7) is normalized by the factor $\frac{|f|}{2\pi}$ to the unit area. In contrast to the CWT, this makes the amplitude response of the ST invariant to the frequency, which means that for a sinusoid with an amplitude A ($x(t) = A \cos(2\pi ft)$), the ST returns an amplitude A regardless of the frequency f in a similar concept as the amplitude of the Fourier transform. On the other hand, the amplitude of the CWT is large for the lower frequencies and diminishes at the higher frequency components. This lower amplitude estimation at the high frequencies is mainly due to the normalization of the CWT [18]. For QTFDs, it is the signal energy that is distributed in t and f , not the signal amplitude. Hence, the interpretation is different. The spectrogram is the square modulus of the STFT; and as the QTFDs are a 2D convolution of the spectrogram, the spectrum

$X(f)$ can also be recovered subject to the initial complex constant $X(f = 0)$.

(6) *Cross-ST analysis*: As the ST is complex and its phase characteristic is referenced at the origin, it can be employed in a cross spectrum analysis in a local manner to estimate the phase synchrony between two signals. Consider two signals measured by two receivers (transducers) separated by a known distance. Let a sinusoidal wave propagate through the medium of view of both receivers. What we will get is the same signal in both sensors with an additive noise (assuming a controlled environment conditions), but also with a time shift between the two signals. Since the ST is a linear operation on the signal and localizes spectral components in time, the cross correlation of specific events on two spatially separated STs will give the phase difference, hence the time delay between the two signals can be measured [18]. This concept is also found in cross-TFDs using quadratic methods, where the signal and its complex conjugate are used to calculate the cross-spectrum [19].

(7) *Co-ST and quadrature-ST*: As the local phase information can be extracted from the ST, we can use the cross-ST function to analyze the in-phase and the out-of-phase components in time-frequency space. This is a very useful characteristic for cross-spectral and phase synchrony. This property is exploited later for an array of EEG signals to estimate the phase synchrony between signals recorded from different electrodes and assess the source location of the brain activities. The cross-ST phase can be used to analyze this synchrony or asynchrony, thus providing a significant discriminating feature for potential abnormal brain activity. As with classical co-spectrum analysis, the real part of the cross-ST function gives the in-phase components of the local spectra. The imaginary part of the cross-ST function gives the in-quadrature components. This property is discussed in detail in Section 3.

The above paragraphs indicate that, unlike the QTFDs, the ST is linear with progressive resolution (see Section 3.3.2 for more details). Unlike the CWT and STFT, the ST has a variable analysis window and an absolute phase reference. These useful properties of the ST make it a natural candidate to study the phase synchrony between two different spatially recorded signals.

2.4 Limitations of the ST

The ST was defined with two unnecessary restrictions on the window function detailed below.

- Firstly, only a Gaussian window $g(t, \sigma)$ is considered.
- Secondly, frequency dependence of the analyzing window of the ST has been through horizontal and vertical dilations of the Gaussian window.

- Also, that window has no parameters to allow its width in time or frequency to be adjusted, as shown in Equation (2).

The resolution, on the ST, of the onset times of events can be improved by using a narrower window, for example, using a better controlled parameters of the Gaussian window. However, when a window is narrowed in the time domain, it inevitably widens in the frequency domain, with consequent loss of resolution in the frequency direction on the ST compromising the identification of the whole event. One way of addressing this problem is to use extra parameters controlling the scale and the shape of the analyzing window rather than just only making the frequency inversely proportional to the standard deviation σ .

3 Extension of the ST for resolution improvement

3.1 Concept and definition

To remedy the limitations discussed in the previous section, a variant of the original ST, introduced by Mansinha et al. [20], replaces f with $\frac{f}{\gamma}$ in Equation (7). By doing this, one standard deviation of the Gaussian window contains γ wavelengths of the Fourier sinusoid at all frequencies. This allows a better control of the time and frequency resolution of the ST on the time-frequency plane. More recently, McFadden et al. [21] presented a generalized ST which includes windows which are asymmetrical. They applied an asymmetrical window to analyze gearbox vibration signals. Later on, Pinnegar and Mansinha [22] have extended the generalization idea in [21] of the ST to include windows which have complicated scaling properties, including frequency-dependent shape and applied this approach to a broadband earthquake seismogram contaminated by noise. This generalized form retains the property of invertibility of the original S -transform [7] and its relation to the Fourier spectrum. The generalized ST is parameterized by a width control adjustment coefficient γ ; it is given by

$$S(t, f, \gamma) = \int_{-\infty}^{+\infty} x(\tau)g(t - \tau, f, \gamma)e^{-j2\pi f\tau} d\tau \quad (11)$$

where $g(t - \tau, f, \gamma)$ is the generalized window function of the ST and γ denotes the set of parameters that determine the shape and the properties of the window function, and the parameter t controls the position of the generalized window on the time axis. The Gaussian window is given by

$$g(t - \tau, f, \gamma) = \frac{|f|}{\gamma\sqrt{2\pi}} e^{-\frac{f^2(t-\tau)^2}{2\gamma^2}} \quad (12)$$

Note that, as in Equation (10), the generalized ST can also be obtained from the signal spectrum as shown below.

$$S(t, f, \gamma) = \int_{-\infty}^{+\infty} X(v + f)G(v, f, \gamma)e^{j2\pi vt} dv \quad (13)$$

3.2 Key properties

The generalized window $g(t - \tau, f, \gamma)$ satisfies the following normalization criteria,

$$\int_{-\infty}^{+\infty} g(t - \tau, f, \gamma) dt = 1 \quad (14)$$

This ensures that the time averaging of the ST spectrum $S(t, f, \gamma)$ yields the spectrum $X(f)$, as shown in (15) below

$$\int_{-\infty}^{+\infty} S(t, f, \gamma) dt = \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} x(\tau)g(t - \tau, f, \gamma)e^{-j2\pi f\tau} d\tau dt \quad (15)$$

$$= \int_{-\infty}^{+\infty} x(\tau)e^{-j2\pi f\tau} \int_{-\infty}^{+\infty} g(t - \tau, f, \gamma) dt d\tau \quad (16)$$

$$= \int_{-\infty}^{+\infty} x(\tau)e^{-j2\pi f\tau} d\tau = X(f) \quad (17)$$

where $X(f)$ is the Fourier transform of the signal $x(t)$.

3.3 Modified ST

3.3.1 Formulation of the modified ST

Based on this generalized form, we can examine the Gaussian window at a different scaling rule [23,24]. This can be achieved by getting the scaling parameter γ defined above to vary linearly with the frequency to get better progressive control of the window width as follows

$$\gamma(f) = mf + k \quad (18)$$

where m is the slope and k is the intercept for a linear change in frequency. Hence, the resolution in time and in frequency will be tuned depending on both the parameters m and k . These parameters will be able to control the Gaussian window length and variance to be appropriate for localizing the low and high frequencies. Figure 1 shows the Gaussian windows with a different value of γ obtained by varying the two parameters m and k . Using this new formulation, the ST becomes

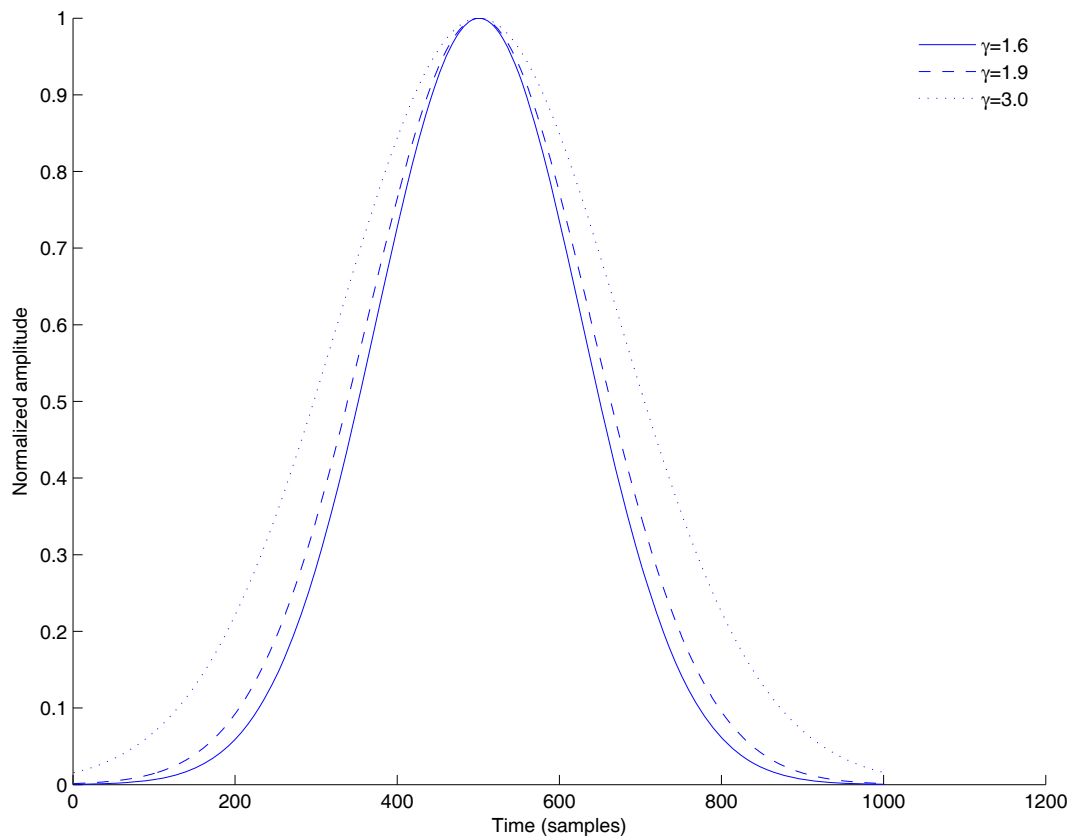


Figure 1 The Gaussian window varying with the parameter γ for frequency $f = 10\text{Hz}$.

$$\text{MST}(t, f, m, k) = \int_{-\infty}^{+\infty} x(\tau)g(t - \tau, f, m, k)e^{-j2\pi f\tau} d\tau \quad (19)$$

where $g(t - \tau, f, m, k)$ denotes the window function of the modified ST (MST), given as follows

$$g(t - \tau, f, m, k) = \frac{|f|}{(mf + k)\sqrt{2\pi}} e^{\frac{-f^2(t - \tau)^2}{2(mf + k)^2}} \quad (20)$$

Using Equations (19) and (20), the modified ST is given by:

$$\text{MST}(t, f, m, k) = \int_{-\infty}^{+\infty} x(\tau) \frac{|f|}{(mf + k)\sqrt{2\pi}} e^{\frac{-f^2(t - \tau)^2}{2(mf + k)^2}} e^{-j2\pi f\tau} d\tau \quad (21)$$

The MST also satisfies the normalization condition for ST window defined in Equation (14) and hence is

invertible.

$$\int_{-\infty}^{+\infty} \frac{|f|}{(mf + k)\sqrt{2\pi}} e^{\frac{-f^2(t - \tau)^2}{2(mf + k)^2}} d\tau = 1 \quad (22)$$

3.3.2 Improved resolution of the MST

In a similar concept to [20], The parameter $\frac{f}{\gamma}$ represents the number of cycles (periods) of a frequency that can be contained within one standard deviation σ of the Gaussian window given by Equation (2). Hence, we have a progressive improved resolution in this case. When too small, the Gaussian window retains very few cycles of the sinusoid and the frequency resolution degrades at low frequencies. In contrast, if it is too large, the window retains more cycles within it and in consequence, the time resolution degrades at high frequencies. This trade-off between time and frequency resolution is

governed by the “Heisenberg” uncertainty principle in a similar way as in the FT, ST or the CWT.

The value of m and k need to be selected properly to get better resolution depending on the properties of the signal under consideration. Selected values of k equals to $\frac{1}{N}$ and m equals to four times the variance of the signal were found empirically to be appropriate; where N is the length of the signal in samples. The method is therefore data dependent or adaptive. As an example for illustration, let's consider a three components signal composed of: a low frequency (100Hz), medium frequency (200Hz) and a high frequency burst (400Hz) as shown in Figure 2a. Figure 2b,c represent the ST and the MST, respectively, of the signal in Figure 2a.

Figure 2b,c show clearly the difference in time-frequency resolution when the ST and the MST are applied. Using the MST leads to an improvement of almost twofold in resolution over the ST.

3.3.3 Simulated results, discussion and interpretation

In the previous example, the values of m and k were chosen depending on the length and the variance of

the signal to get a better resolution (see Appendix 1 for more details). This indicates that as is the case with power spectral estimation, the approach is not restricted to a Gaussian window and any relevant window or apodising function may be employed to improve resolution. The method is not restricted to a Gaussian window and can be generalized in a similar concept to the CWT for which different windows can be used. Hence, the ST can be seen as a special case of the CWT using a Morlet-type wavelet, with a phase and amplitude correction. This generalized form of ST is more flexible, as it uses two added parameters for varying the frequency in a linear way, resulting in adjustable windows for improved time or frequency resolution. Similar ideas led to the development of transforms based on groups generated by translations, modulations, and dilation of a mother wavelet [25]. All of these approaches give a more versatile choice of transform suitable to particular cases and applications. However, the particular characteristic of absolute phase, make the MST more attractive than wavelet

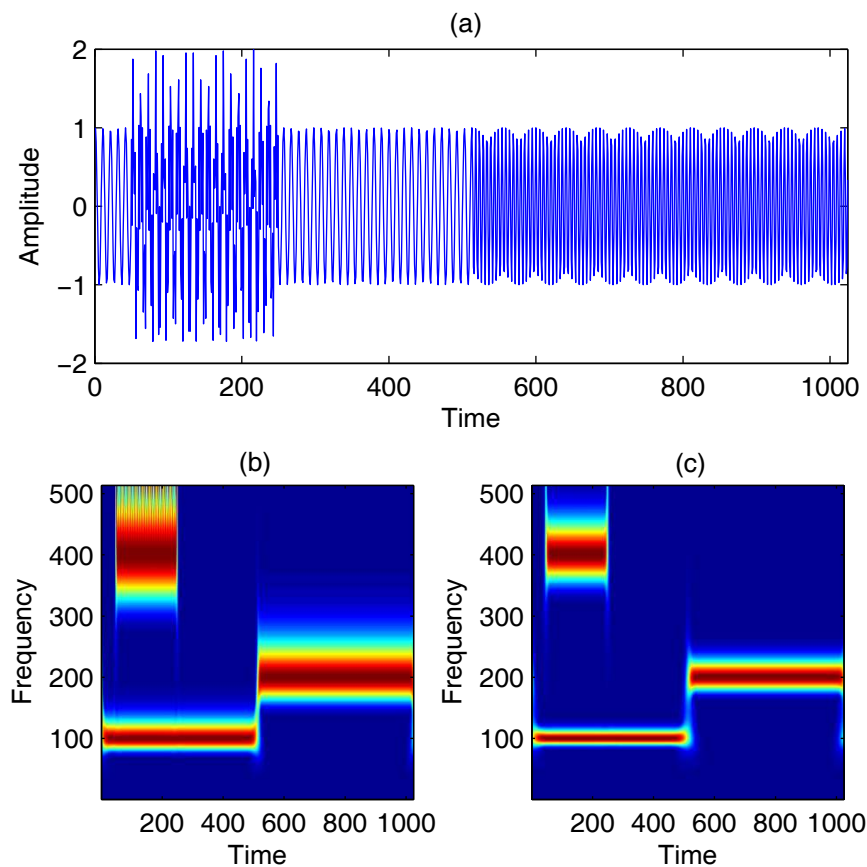


Figure 2 ST and MST comparison. (a) Time plot of a three-component signal. (b) The ST of the signal in (a). (c) The MST of the signal in (a).

when phase need to be estimated (e.g., in coherence analysis and cross-spectral analysis). It can simultaneously estimate the local amplitude spectrum and the local phase spectrum, whereas a wavelet approach is only capable of probing the local amplitude/power spectrum. The MST fully estimates the amplitude of the signal, in contrast to the CWT which attenuates high frequencies. The MST can be further improved by dealing with issues such as fast algorithms, redundancy of representations, Hilbert space properties, resolution of the identity, among others [7].

4 Cross-MST and phase synchrony

4.1 Motivation and illustration

As mentioned above, the MST phase preservation characteristic is an advantage that allows a cross spectrum analysis in a local manner. Let's consider two signals measured by two sensors separated by a known distance. What these two sensors see is the same signal, but with a time shift in the arrivals as well as noise. Since the MST localizes spectral components in time, the cross correlation of specific events on two spatially separated MSTs gives the phase difference and hence the phase synchrony can be estimated. The amplitude of the cross-MST indicates coincident or partly overlapped signals. The phase of the cross-MST at local maxima indicates the phase difference between them. The cross-MST of two signal $x(t)$ and $y(t)$ is defined as follows [18]

$$\text{crossMST}(t, f) = \text{MST}_x(t, f) \cdot \text{MST}_y(t, f)^* \quad (23)$$

where $()^*$ denotes the complex conjugate.

The phase of the cross-MST is given by

$$\arg(\text{crossMST}) = \phi_x(t, f) - \phi_y(t, f) \quad (24)$$

Three key properties of the MST that make the cross-MST useful are detailed below.

(1) *The MST translation property in time and frequency*: As mentioned in [18], the ST (or MST) time translation property is similar to what Fourier transform states that is: if

$$x(t) \Leftrightarrow \text{MST}(t, f) \quad (25)$$

As MST preserve the same translation property as Fourier transform, if we translate $x(t)$ by an amount r we obtain

$$x(t - r) \Leftrightarrow \text{MST}(t - r, f) e^{-j2\pi fr} \quad (26)$$

The above equation shows that the phase difference between two signals is equivalent to the cross-power spectrum. By Taking the inverse MST transform of the representation in the frequency domain, the

displacement between two signals can be easily obtained (phase correlation technique).

(2) *The MST phase shift*: According to [18], if we consider a sinusoidal signal $x(t)$

$$x(t) = e^{j2\pi f_0 t} \quad (27)$$

with a spectrum

$$X(v) = \frac{1}{2} \delta(v - f_0) \quad (28)$$

Using Equation (13), The MST of $x(t)$, is given by

$$\text{MST}_{x(t)}(t, f) = \int \delta(v + f - f_0) e^{\frac{-2\pi^2 v^2}{f^2}} e^{j2\pi vt} dv \quad (29)$$

which results in:

$$\text{MST}_{x(t)}(t, f) = \frac{1}{2} e^{\frac{-(f - f_0)^2}{\left(\frac{f}{2\pi}\right)^2}} e^{j2\pi(f - f_0)t} \quad (30)$$

(3) *Phase in the cross-MST*: if we introduce a constant phase shift into $x(t)$, we get

$$y(t) = x(t) e^{j\phi} = e^{j2\pi f_0 t + j\phi} \quad (31)$$

with a resulting spectrum

$$Y(v) = \frac{1}{2} \delta(v - f_0) e^{j\phi} \quad (32)$$

this will also introduce a phase shift in the modified ST

$$\text{MST}_{y(t)}(t, f) = \frac{1}{2} e^{\frac{-(mf + k - f_0)^2}{\left(\frac{mf+k}{2\pi}\right)^2}} e^{j2\pi(mf+k-f_0)t} e^{j\phi} \quad (33)$$

The three above properties indicate that we can perform a cross spectral analysis between $x(t)$ and $y(t)$ by multiplying the MST of $x(t)$ with the complex conjugate of the modified ST of $y(t)$ as follows

$$\text{crossMST}(t, f) = \frac{1}{2} e^{\frac{-(mf + k - f_0)^2}{\left(\frac{mf+k}{2\pi}\right)^2}} e^{j2\pi(mf+k-f_0)t} \frac{1}{2} e^{\frac{(mf+k-f_0)^2}{\left(\frac{mf+k}{2\pi}\right)^2}} e^{-j2\pi(mf+k-f_0)t} e^{-j\phi} \quad (34)$$

Which simplifies to:

$$\text{crossMST}(t, f) = \frac{1}{4} e^{-j\phi} \quad (35)$$

hence, the phase of cross-MST is clearly given by

$$\text{Phase}\{\text{MST}_x\text{MST}_y^*\} = -\phi \quad (36)$$

The above derivation shows that we can use the cross spectral property to detect a time lag between the two signal as a function of t and f . A similar attempt was made in [26,27] by using the cross-WVD for representing non-stationary processes.

4.2 Illustration on simulated signals

Let's consider the signal $x(t)$ given in Figure 2 with 1,000 samples length and sampled at 1 kHz; it includes three lower (100Hz), medium (200Hz), and high frequency (400Hz) components; consider also another signal $y(t)$ with similar three components except that the higher component is translated and partly overlaps in time with the higher component of $x(t)$ (see Figure 2a). The higher components are out of phase (the $x(t)$ component is a cosine function, the $y(t)$ component is a sine function). The lower and medium components of the

two signals are in phase. Figure 3 shows the two signals $x(t)$ and $y(t)$ and their related MSTs. Figure 4 shows the cross-MST of the signals $x(t)$ and $y(t)$ corresponding to the two MSTs in Figure 3c,d. The figure indicates that, for the higher component, only the overlapped part in time is detected, as expected.

In a similar way to co-spectrum analysis, the real part of the cross-MST gives the in-phase components of the local spectra and the imaginary part of the cross-MST gives the in-quadrature components. Figures 5a,b show the Co-MST and the Quadrature-MST amplitudes corresponding to the two signals $x(t)$ and $y(t)$, respectively. In the co-MST (Figure 5a), only the features that are in phase are present—those being the low and the medium frequency components. The higher frequency (which is out of phase) does not appear. In the quadrature-MST plot, only the out-of-phase components appear, which is the higher frequency component. As the example above was only for illustration, where the used signals have

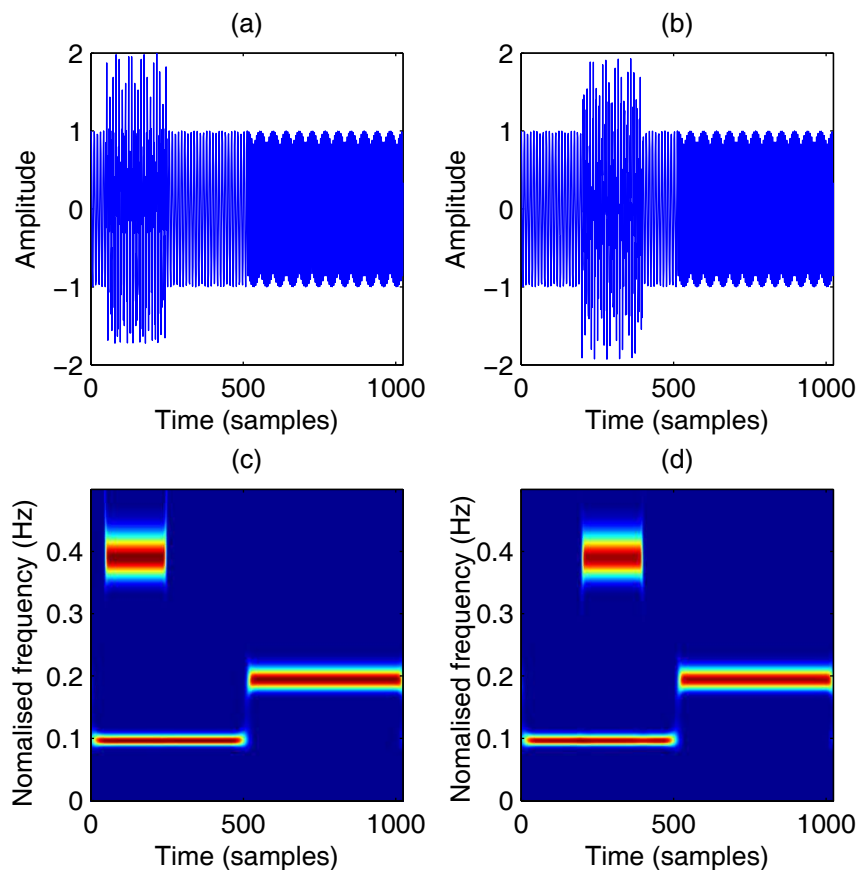


Figure 3 MST for the two signals $x(t)$ and $y(t)$. (a) The signal $x(t)$ with three-components (100, 200, and 400 Hz). (b) The signal $y(t)$ with three-components (100, 200, and 400 Hz). (c) The MST of the signal $x(t)$ given in (a). (d) The MST of the signal $y(t)$ given in (b).

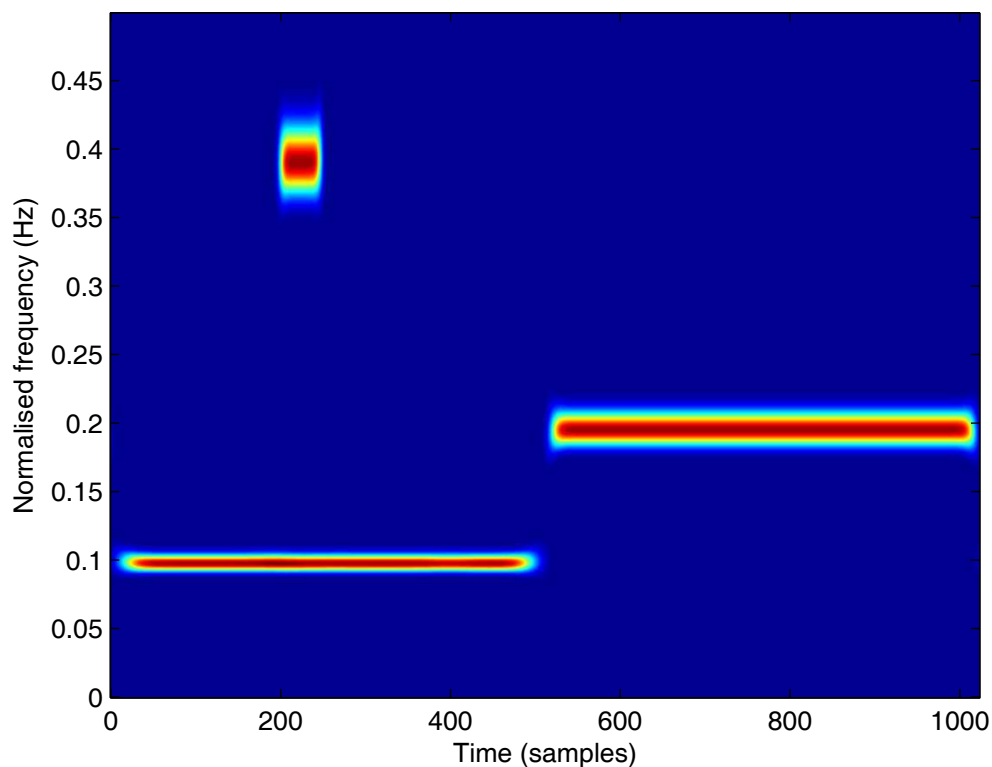


Figure 4 Cross-MST of the two signals $x(t)$ and $y(t)$ MSTs given in Figure 3c,d, respectively.

straight and linear IF components and there is no noise; these results show that it is possible to study the coherence with sufficient high resolution using the cross-MST, whether for amplitude or phase coupling between two different spatially recorded components or signals.

5 Application to EEG seizure prediction

5.1 Phase synchrony information

The nonlinear interdependencies between fluctuations of brain physiological activities EEG-based recordings were intensively studied in pairs, and the synchrony flow differences were compared [28]. It was shown in [29] that a neurological disorder in the brain can be detected by the random occurrence of its clinical manifestations, i.e., the seizure. There is also evidence that the nonseizure (interictal) to seizure (ictal) transition is not an abrupt phenomenon [29]. Hence, to provide a valuable insight into such mechanism, identification of early changes in EEG signals [28] can be used for prediction, prevention, and control of upcoming seizures. Different approaches have been employed for seizure prediction in the past. Some limited techniques based on visual inspection of the EEG signals or on linear methods have failed to detect specific and

sustained changes preceding seizures [28]. Other methods reported some changes based on the spectral contents [29], complexity [30], or spatio-temporal patterns of spikes [31]. Other recent approaches applied to newborn EEG-based seizure detection have been proposed such as time-frequency signal processing [32,33], and nonlinear signal processing [34,35].

Neuroscientists claim that when a focal seizure is generated, synchronized brain activity is initially observed only in a small area of the brain; and, from this focus, the activity spreads spatially out to other brain areas in the temporal lobe over time [29]. Involving the seizure focus spread into its surrounding, a hyper-synchronous state happens and neighboring areas lose their synchronization with the other cortical regions around them [36]. In consequence, the seizure focus becomes isolated from the rest of the brain dynamics, making the considered population of neurons inactive. Methods to track these changes using coherence and synchronization are being explored using EEG, for the purpose of predicting an impending seizure [36]. Phase synchronization was also shown to be a sensitive indicator of coupling between signals and as an important factor in the genesis of epileptic phenomenon [36]. Many studies suggest

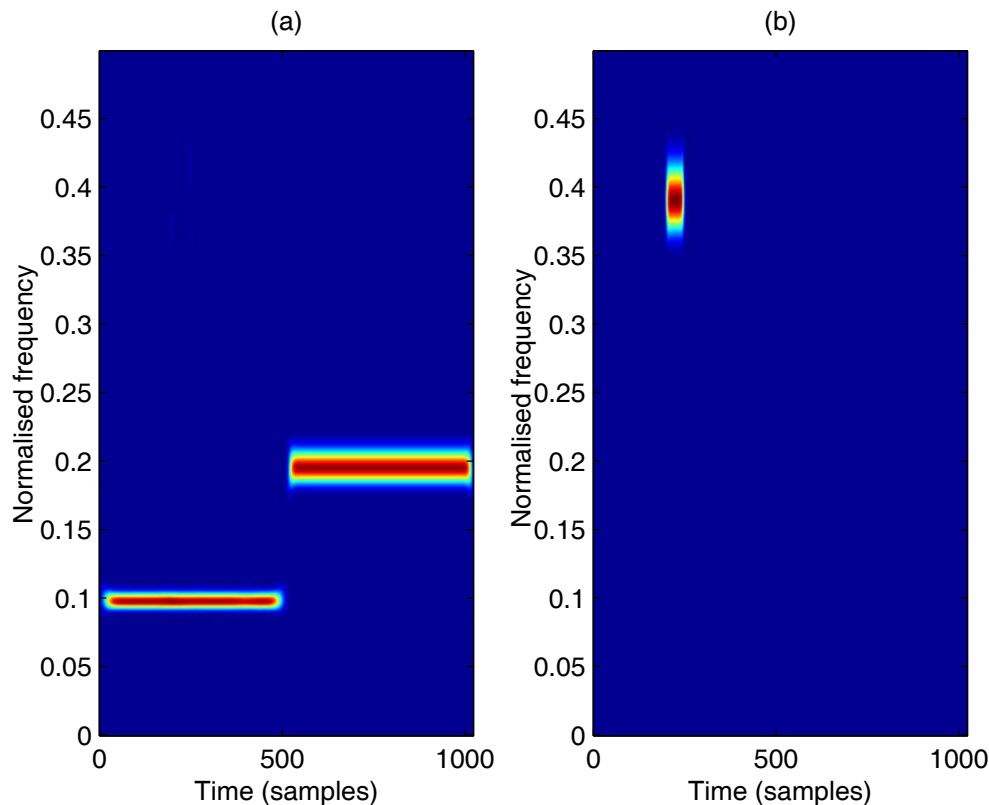


Figure 5 co-MST and quadrature-MST. (a) The extracted in-phase components of $x(t)$ and $y(t)$. **(b)** The extracted out-of-phase components of $x(t)$ and $y(t)$.

an underlying correlation between neuronal synchronization and seizure development and onset [36,37]. Thus, the concept of phase synchrony helped to measure the synchrony evolution while the amplitude of the signals remained uncorrelated. In contrast with cross-correlation that measures linear relationships only, a phase coupling approach is able to show the presence of non-linear coupling [38]. Moreover, phase synchrony can be used for non-periodic and for chaotic signals such as EEG [8,39]. Some authors used HT to measure phase synchrony between two signals [37,38,40]. However, the use of the HT assumes the signals have a narrow frequency band [1, p. 14], and it is not straightforward to extend the same analysis for broadband data. This assumption is usually seen to be ignored in the context of biomedical signals like the EEG [41]. As EEG signals are sometimes broadband (1-100 Hz) [41], the HT may not be able to correctly estimate the instantaneous phase of broadband signals. This raises the concern that the broadband phase synchronization analysis may lead to a misinterpretation of the results [41]. Therefore, when the signal is broadband it is necessary to pre-filter it in the frequency band of interest before applying the

HT, in order to get a better estimate of the phase [42]. Some authors used wavelet-based approach to analyze phase synchrony between pairs of EEG signals [13,43-45]. These time-varying measures of phase synchrony using wavelet or HTs are similar in their results. Both give a sharper phase synchrony estimates over time and frequency, especially at the low frequency range [46]. Although the wavelet and Hilbert based phase synchrony estimates address the issue of non-stationarity of EEG signals, they suffer from limited time-scale resolution due to the limited number of available scales in the case of wavelets and the narrowband assumption in the case of HT. Moreover, the role of phase in wavelet analysis is not as well understood as it is for the Fourier transform, especially for orthonormal wavelet representations. Current complex wavelets such as Daubechies, dual tree, and Shannon wavelets do not have a direct relationship to the signal spectrum [18]. For these reasons, this article proposes to use the MST and the cross-MST to analyze the phase synchrony with the required resolution and with meaning, as the phase in the ST is meaningful and has the same concept as in the Fourier transform approach.

5.2 Coherence analysis using cross-MST for simulated EEG seizure

Due to its non-stationarity characteristic, the newborn EEG seizure was modelled in the time-frequency domain [47-49]. using time-frequency characteristics previously identified in the newborn EEG seizure. Three models of newborn EEG seizure simulation were previously proposed. One model is based on some physiological parameters of the brain and utilises a stationary sawtooth waveform [50]. This approach is extended in [47] to incorporate a single linear frequency modulation (LFM) signal. Another piecewise LFM model was defined for seizure based on non-stationary inputs to a nonlinear model [51]. We propose here a method of newborn EEG seizure analysis using the piecewise LFM pattern outlined in [47-50] to simulate the nonstationary behavior of the newborn EEG seizures.

The simulated signal is composed of several piecewise LFM components with different LFM slope parameters. Figure 6a represents three LFM piecewise signal with added Gaussian noise simulating the EEG seizure and the corresponding MST analysis. Figure 6b represents another three LFM piecewise signal with its MST, similar to the one in Figure 5a except that the last component is out of phase compared to the two other components which are in phase in the two signals (More details are given in Appendix 2).

Key parts of the commented Matlab[®] code written to implement this method are given in Appendix 2.

Figure 7a represents the real part of the cross-MST of the two MSTs in Figure 6. We can see clearly that only the in-phase components are detected. However, in

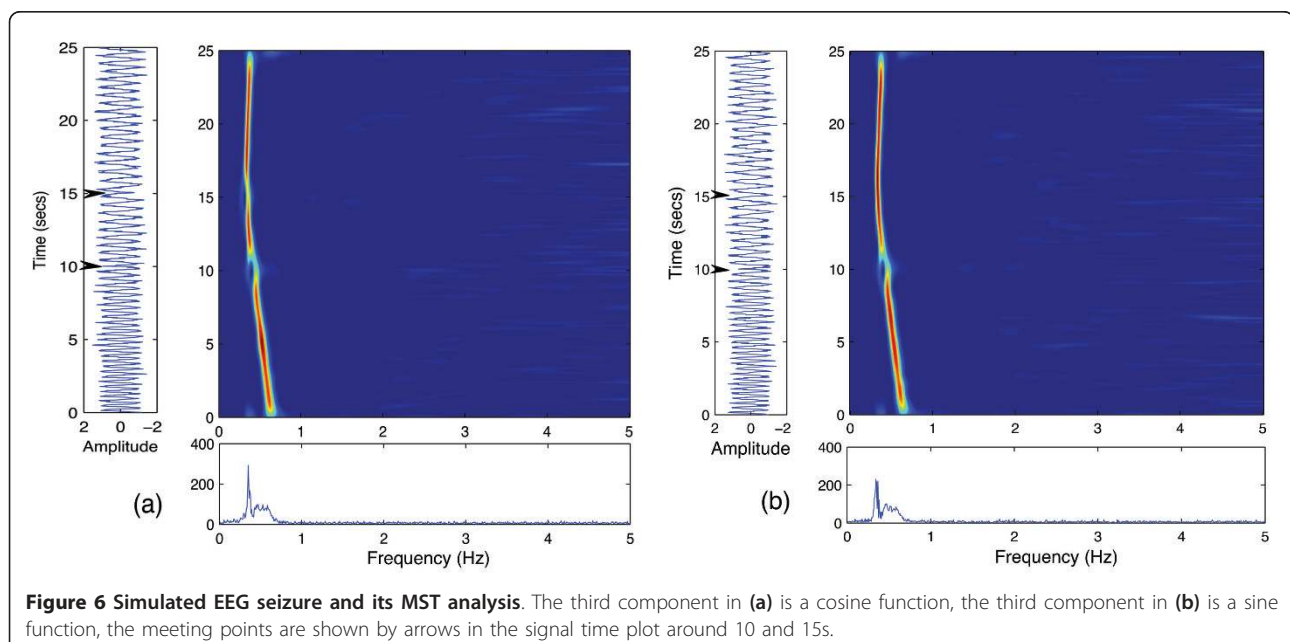
Figure 7b, only the out-of-phase component is detected (the third component). Key parts of the commented Matlab[®] code written to implement this method are given in Appendix 1. Cross-MST of the Simulated EEG seizure showing the in-phase components (a) and the out-of-phase component (b).

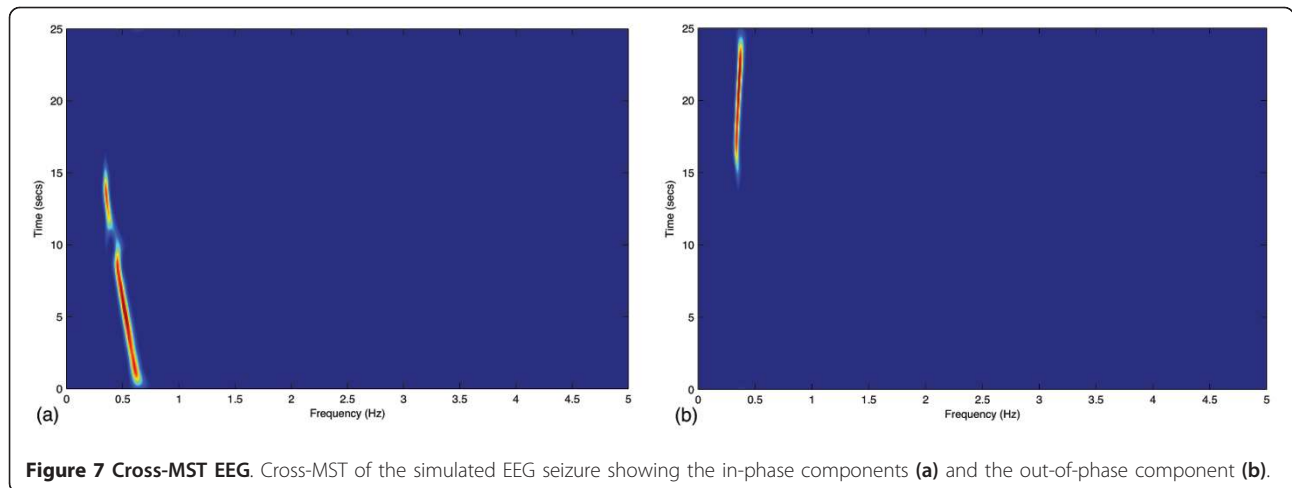
5.3 Coherence Analysis using cross-MST for newborn EEG data

5.3.1 Background and analysis

The data was acquired using a 20-channel Medelec Profile system (Medelec, Oxford Instruments, Old Woking, UK) at 256Hz sampling rate and marked by a pediatric neurologist from the Royal Children's Hospital, Brisbane, Australia. The experiment was conducted as part of a long term collaboration between Prof. Boashash, Director of the SPRC and Prof. Paul Colditz, Director of the Perinatal Research Centre at the University of Queensland. Measurements were made from a number of sick babies and this work was previously reported in [52]. Five monopolar channels (O1, O2, P3, P4, Cz) out of the 14 recorded signals according to the 10-20 standard [53] modified for newborn subjects were selected from newborn EEG dataset. The data were analysed to investigate the coherence and phase coupling interhemispheric and intrahemispheric interactions during an EEG seizure period. Figure 8a shows the 10-20 standard map of the 20-channel EEG recorder.

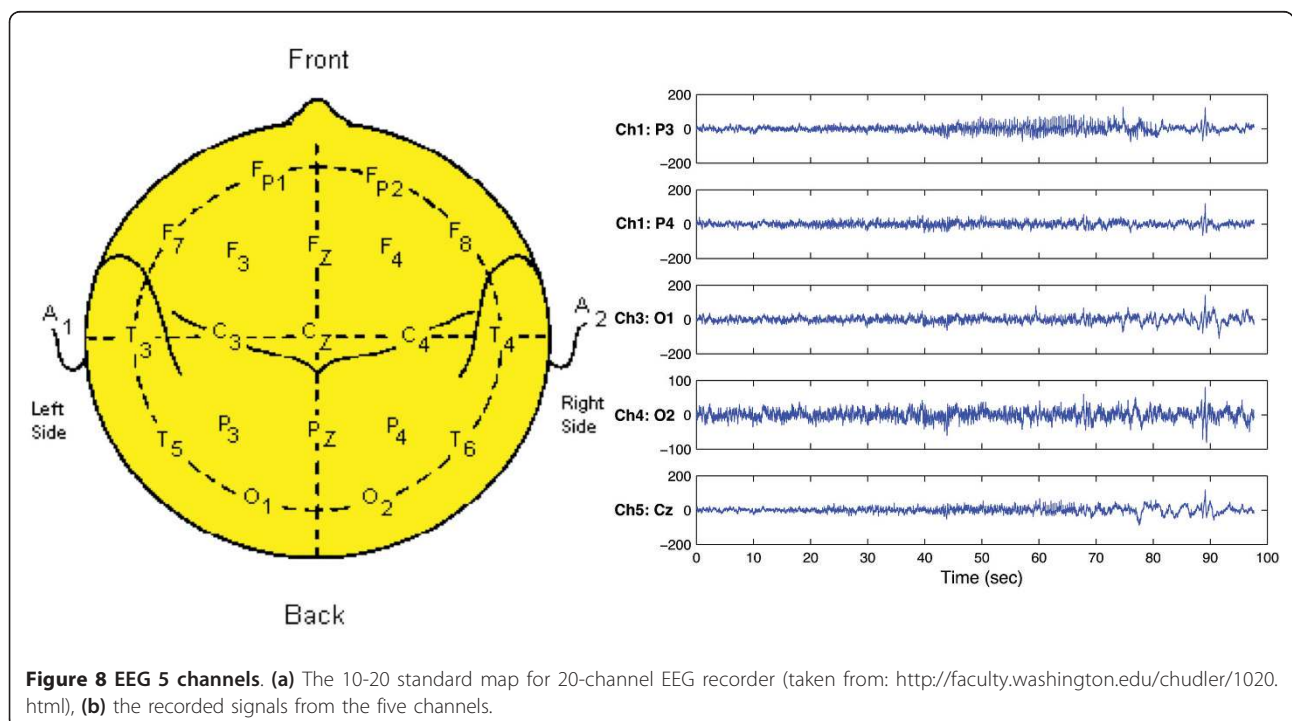
To study the phase coupling between the five channels; we apply the MST to the five signals recorded from the five channels corresponding to P3, P4, O1, O2, and Cz, respectively. The location of the electrodes for





the five channels is shown in Figure 8a. Then, we apply the cross-MST function to analyze the in-phase and out-of-phase components in time-frequency space between the five channels. As mentioned previously, the real part of the cross-MST function gives the in-phase components of the local spectra; and the imaginary part of the cross-MST function gives the in-quadrature components. Figure 8 show the in-phase components for the connectivity between channels. From the all 25 possible connectivities shown, Figure 9 indicates that a strong phase synchrony exist in the last few seconds of the recording. This is shown by a peak in the real cross-MST images around 90s. This is in agreement with the

recorded signals shown in the time domain (see Figure 8b) as at the corresponding time (90 s) a peak with high amplitude is appearing in the signals. This may explain that at this time (around 90s) the pulses coming from different group of neurone are in phase and when they are added together, the sum results in high peak amplitude. This interpretation will not be possible if we just look at the signals amplitude from Figure 8b. The time signals recorded amplitude does not give a phase information to justify the synchrony; however, the cross-MST shows clearly this phase synchrony by a high peak amplitude in the real part present in the connectivity images (Figure 9).



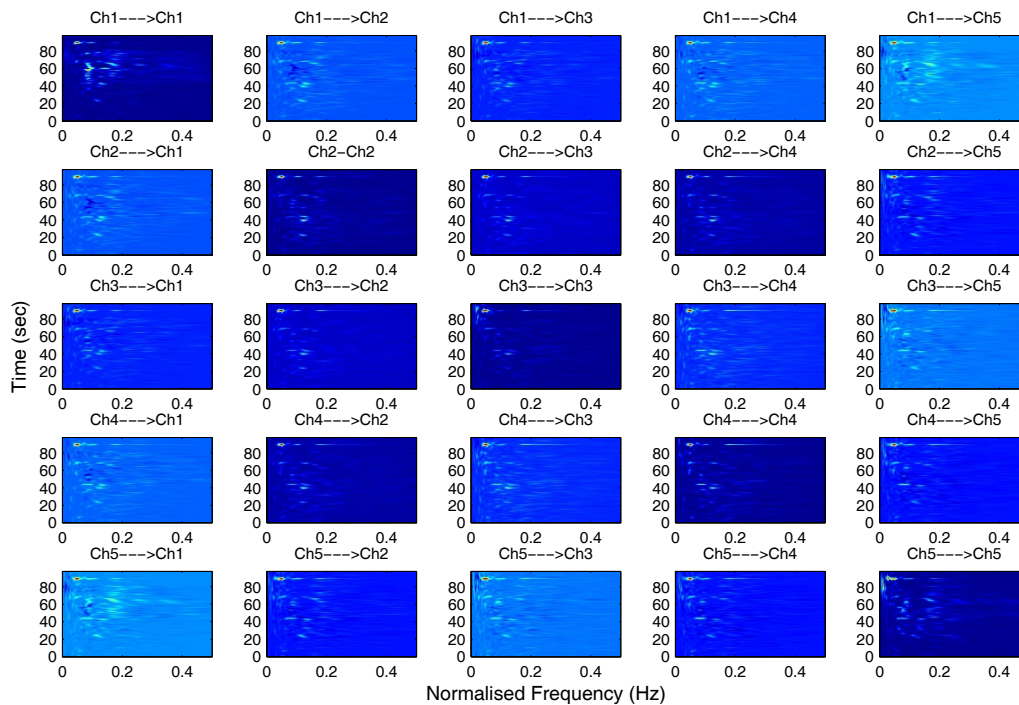


Figure 9 The phase synchrony between the five channels using the real part of the cross-MST.

Figure 10 represents the imaginary part of the cross-STs which reveal the out-of-phase behavior between the five channels. Note that there is no activity in the diagonal images of Figure 10. This is because the diagonal represents the self-coherence between the five channels, hence there is no out-of-phase in this case; which is consistent as a signal can't be out-of-phase with itself. From the off-diagonal images, we can see that most of the channels exhibit an out of phase behavior mostly in the period 40-80 s corresponding probably to the important duration of the appearance of the seizure. Some out-of-phase components are also present before this 40-80 s period in some channels, which can be explained by the fact that the seizure is not appearing in the same time and with the same way in the spatially different parts of the brain from where the EEG signals are recorded. The end of seizure can be explained by the disappearance of the out-of-phase components and the appearance of the out-of-phase components.

5.3.2 Results, discussion and interpretation

The results on the real newborn EEG data show that the cross-ST can be used as an effective tool to study the phase coupling between channels. These results suggest that during the seizure period, all the components lose synchrony and are out of phase. This can be seen like a disorder in the neuronal activity during the seizure

interval. This observed irregularity in the time-frequency patterns during the seizure period (40-80 s) between two hemispheres in the posterior area suggest that different regions and not only one specific region may be causing the seizure. In contrast, the components (neural activity) became synchronized again after the seizure activity vanishes. This EEG seizure application demonstrates that the time-frequency coherence is an appropriate tool to study the phase coupling between different signals recorded from different spatially separated electrodes. This study is performed by the cross-MST which has a referenced phase at the origin, a property that allows the study of phase coupling and time-frequency coherence between the neuronal activities. Although the cross-MST represents a good candidate tool to analyze such coupling behavior, this methodology can be extended to other cross-time-frequency methods with high time-frequency resolution such as the QTfDs. In particular, the time-frequency coherence using the cross-WVD has a useful property for cross spectral analysis, and can be interpreted as a time and frequency dependent correlation coefficients [26].

6 Conclusions and perspectives

This article applies the concept of time-frequency coherence using the cross-MST which is a time-

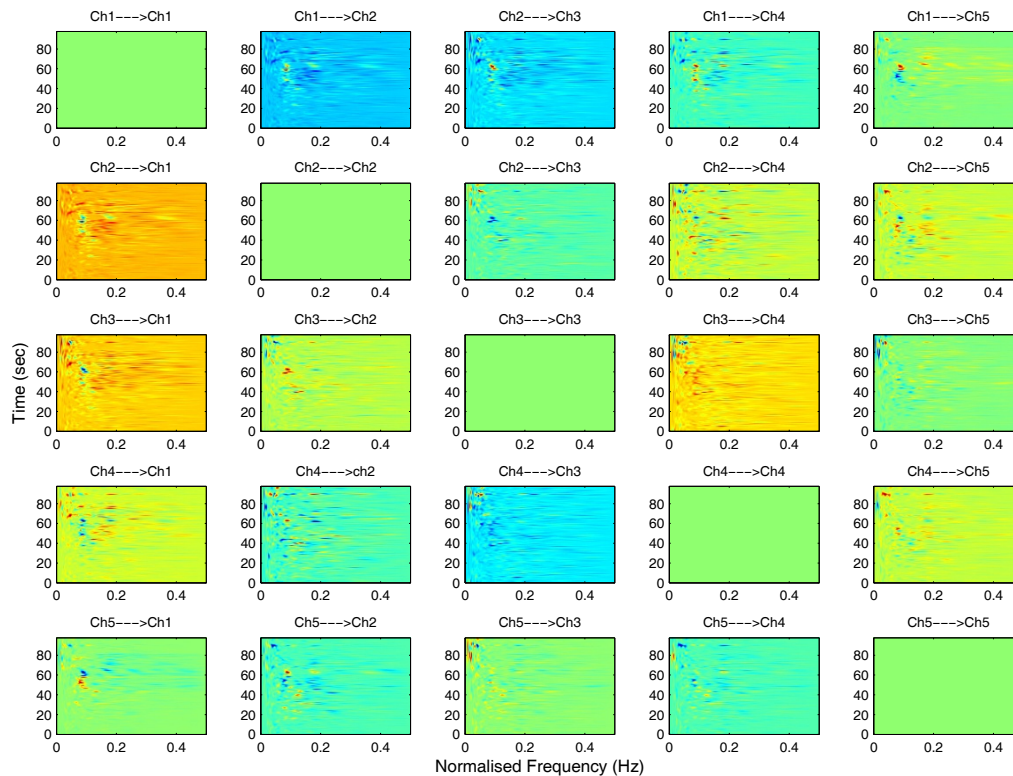


Figure 10 The out-of-phase (asynchrony) between the five channels using the imaginary part of the cross-MST.

varying cross-spectral analysis method obtained by extending the *S*-transform. The article demonstrates the ability of this time-frequency coherence using the cross-MST to study the functional neuronal phase synchrony between channels. The performance of the phase synchrony estimation using the cross-MST are evaluated both by simulated EEG seizure using a piecewise LFM signal to extract the in and out-of-phase components, and through the analysis of seizure EEG abnormalities in the newborn, where the dynamics of brain changes rapidly and the neurones lose their synchrony during seizure. Both the simulated and newborn results show cross-MST phase and synchrony estimation to be more robust to noise. The resolution obtained by the MST shows almost twofold improvement over the standard ST approach. Future work will concentrate on a comparison study between the proposed method and other existing time-frequency techniques such as Wigner-Ville and QTDFs based methods, for both simulated and real EEG data to develop more efficient time-varying phase and synchrony estimation for nonstationary signals, and such

results will appear elsewhere. The proposed method can be extended to other physiological problems where the phase coupling is relevant, such in cardio-respiratory relationship, for example, where the coupling can be affected by cardio-vascular diseases.

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Appendix

1 Computing the MST

```
% Compute the MST
[M,N]=size(sig); % get the
length of the signal
N2=fix(N/2); j = 1;
if N2*2==N; j = 0;end
```



```
f=[0:N2 -N2+1-j:-1]/N; %frequency
MST=zeros(N2+1,N); %allocate mem-
ory for positive frequencies of MST.
SIG=fft(sig,N); %compute the sig-
nal spectrum
g=(1/N)*f+4*var(sig); %parameter
gamma
for i = 2:N2
    SIGs=circshift(SIG,[0,-(i-1)]); %
circshift the spectrum SIG
    W=(g(i)/f(i))*2*pi*f; %
Scale Gaussian
    G=exp((-W.^2)/2); %W in
Fourier domain
    MST(i,:)=ifft(SIGs.*G); %Com-
pute the complex values of MST
end
imagesc(abs(MST)');
set(gca,'Ydir','Normal'); %default Y
axis in imagesc is inverted.
```

2 Generation of piece-wise LFM

```
%Piecewise LFM chirp generation for EEG
Seizure simulation
fs = 50; %Sampling frequency
t1=(-5:1/50:5); t2=(-2.5:1/50:2.5); t3=
(-5:1/50:5); % time for the first,
% the second and the
% third segment
t=[t1 t2 t3]; % total time
for the Piecewise LFM.
f00 = 2.626; f01 = 1.83; f10 = 1.751; %
start frequencies
a1=-0.7; a2=-0.6; a3 = 0.2; % chirp
rates or slope
ph1 = 2*pi*f00*t1+0.5*a1*t1.*t1; %
quadratic phase of piecewise LFM1
ph2 = 2*pi*f01*t2+0.5*a2*t2.*t2; %
quadratic phase of piecewise LFM2
ph3 = 2*pi*f10*t3+0.5*a3*t3.*t3; %
quadratic phase of piecewise LFM3
lfm1=exp(1i*ph1); lfm2=exp(1i*ph2);
lfm3=exp(1i*ph3); %generate piecewise
LFMs
lfm=[real(lfm1) real(lfm2) imag(lfm3)];
% put all in one signal LFM.
sig=lfm+0.2*randn(length(lfm),1)'; %
add gaussian noise
```

3 Computing the cross-MST

```
%The cross-MST, co-MST and quadrature-
MST
%Computing the Cross-MST of two signal x
(t) and y(t)
%having as modified ST: MST1 and MST2,
respectively.
```

```
% 1.Compute MST1 and MST2 using the
related code in appendix A.
% 2.Compute the cross-MST between MST1
and MST2
cross_MST = MST1.*conj(MST2); %Get the
cross-MST.
% 3. Get the out-of-phase components
imagesc(imag(cross_MST)); set
(gca,'Ydir','Normal')
% 4. Get the in-phase components
imagesc(real(cross_MST)); set
(gca,'Ydir','Normal')
```

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Competing interests

The authors declare that they have no competing interests.

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