

## Evidence for Dirac Fermions in a Honeycomb Lattice Based on Silicon

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Silicene, a sheet of silicon atoms in a honeycomb lattice, was proposed to be a new Dirac-type electron system similar to graphene. We performed scanning tunneling microscopy and spectroscopy studies on the atomic and electronic properties of silicene on Ag(111). An unexpected  $\sqrt{3} \times \sqrt{3}$  reconstruction was found, which is explained by an extra-buckling model. Pronounced quasiparticle interferences (QPI) patterns, originating from both the intervalley and intravalley scatter, were observed. From the QPI patterns we derived a linear energy-momentum dispersion and a large Fermi velocity, which prove the existence of Dirac fermions in silicene.

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Group IV (Si, Ge) analogs of graphite have been discussed for a long time even before the synthesis of isolated graphene [1], and recently there has been renewed interest in this topic due to the novel concepts and applications brought on by graphene. The silicon version of graphene in which Si atoms replace C atoms in a two-dimensional honeycomb lattice is named silicene [2–7]. Theoretical calculations show that silicene also has graphenelike electronic band structure, supporting charge carriers behaving as massless Dirac fermions [2,3,7]. Compared with graphene [8,9], silicene has a larger spin-orbit coupling strength, which may lead to a larger energy gap at the Dirac point and favor a detectable quantum spin Hall effect [7,10,11]. Currently a quantum spin Hall effect has been realized only in HgTe-CdTe quantum wells and further studies have been hindered by the challenging material preparation [12]. Easy preparation and compatibility with silicon-based nanotechnology make silicene particularly interesting for applications like quantum spin Hall effect devices.

Despite the rapidly increasing amount of theoretical works on silicene, there have been only a few experiments on silicene or a silicene nanoribbon [4–6]. The growth of a monolayer silicene sheet on Ag(111) had been claimed by Lalmi *et al.* [6]. A scanning tunneling microscopy (STM) study revealed a hexagonal honeycomb structure, which was distinct from known surface structures of bulk silicon, and resembles that of graphene [6]. However, the reported Si-Si distance of 17% shorter than that for the bulk and for the theoretical model implies unrealistic high compression of the silicene lattice, which remains to be confirmed and understood [6]. Apart from the preparation and structural studies, there are still no experiments on the electronic structure of monolayer silicene. Experimentally establishing a common understanding of the basic atomic and electronic properties of silicene is therefore highly desirable.

In this Letter we report a study on silicene by low temperature STM and scanning tunneling spectroscopy (STS). We observed an unexpected  $\sqrt{3} \times \sqrt{3}$  reconstruction on the silicene surface, in contrast to the  $1 \times 1$  structure reported previously [6]. Despite the structural difference, the electronic property measured by STS is consistent with theory very well. For example, quasiparticle interference (QPI) patterns suggesting intervalley and intravalley scattering of charge carriers were observed, and a linear energy-momentum dispersion relation and a large Fermi velocity were derived. Our results provide a solid basis for further studies on the electronic property and device applications of silicene.

Experiments were performed in a homemade low-temperature STM equipped with *in situ* sample preparation facilities. A clean Ag(111) surface was prepared by cycles of argon ion sputtering and annealing. Silicon was evaporated from a heated Si wafer with a deposition flux of about 0.05 monolayers per minute (ML/min), and the Ag(111) substrate maintained at  $\approx 500$  K. The differential conductance ( $dI/dV$ ) spectra were measured as the in-plane ac component in the tunneling current with a lock-in amplifier by superimposing an ac voltage of 10 mV and 676 Hz on the given dc bias of the substrate-tip gap. We performed STM measurements at both room temperature and 77 K and found no difference. All the data presented in this Letter were taken at 77 K.

The growth of Si on Ag(111) is very sensitive to the Si coverage and the substrate temperature. We have explored a wide range of Si coverage and substrate temperature, and several metastable phases were found at a low temperature range [13,14]. On the other hand, at a sufficiently high substrate temperature and Si coverage, silicon routinely forms one-atom thick, closely packed islands which we identify as silicene sheets. A typical STM topographic image of the thus obtained monolayer islands is shown in Fig. 1(a). The line profile across the edge of a silicene

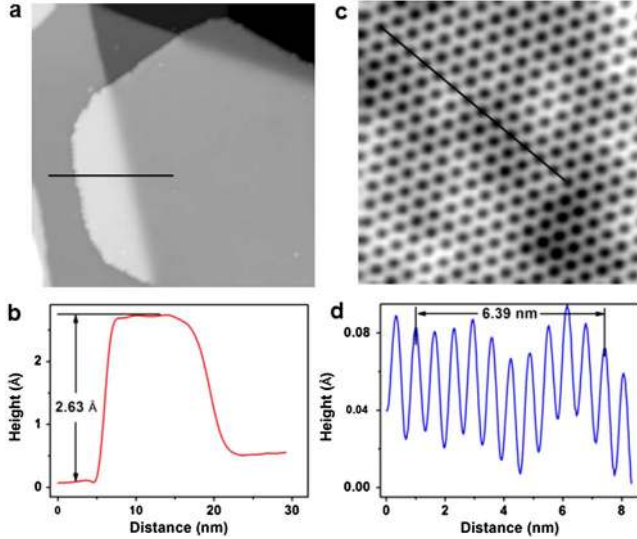


FIG. 1 (color online). The STM image of a large area ( $65 \text{ nm} \times 65 \text{ nm}$ ) consisting of a sheet of silicene on Ag(111) crossing two substrate steps. (b) The line profile as indicated by the black line in (a) shows that the island is one atom thick. (c) The high-resolution STM image ( $10 \text{ nm} \times 10 \text{ nm}$ ) of the silicene surface taken at tip bias 1.0 V. The honeycomb structure is clearly observed. (d) The line profile as indicated by the black line in (c) showing both the lateral and vertical corrugation of the structure observed by STM.

island indicates that it has a height of  $2.63 \text{ \AA}$ , corresponding to one atom thickness. Such islands can run across steps of the Ag(111) substrate without losing the continuity of the atomic lattice, similar to that of a graphene sheet falling on a stepped surface. As expected, the high resolution STM image in Fig. 1(c) shows a honeycomb structure reflecting the threefold symmetry of this film, which is not changed in STM images with different bias voltage. However, surprisingly, after a very careful calibration we derived a periodic constant of  $0.64 \pm 0.01 \text{ nm}$  [shown in line profile in Fig. 1(d)], which is approximately  $\sqrt{3}a$  ( $a$  is the lattice constant of silicene,  $\approx 0.38 \text{ nm}$ , proposed in several theoretical works [3,7]). Such a large difference is not a possible result of strain-induced lattice expansion of compression. It is most likely that we observed a  $(\sqrt{3} \times \sqrt{3})R30^\circ$  superstructure, which frequently shows up in threefold symmetric systems.

To account for this observation, we have carried out first-principles calculations using the projector augmented wave pseudopotential method and Perdew-Burke-Ernzerhof exchange-correlation potential [15] implemented in the VASP package [16]. The convergence criteria for energy and force were set to  $10^{-5} \text{ eV}$  and  $0.001 \text{ eV/\AA}$ , respectively. We performed an extensive search for different geometric configurations with  $(\sqrt{3} \times \sqrt{3})R30^\circ$  periodicity, and concluded that a  $(\sqrt{3} \times \sqrt{3})R30^\circ$  superstructure cannot be stabilized in a freestanding, fully relaxed silicene model. We also considered the incorporation of Ag atoms

in  $\sqrt{3} \times \sqrt{3}$  sites in the silicene lattice, but it was proven to be energetically unfavorable. The period constant  $0.64 \text{ nm}$  also cannot result from the commensuration between the lattices of silicene and Ag(111). Actually the theoretical modeling only gives  $4 \times 4$  or  $\sqrt{7} \times \sqrt{7}$  superstructure (with respect to the  $1 \times 1$  Ag substrate). We also exclude the possibility that the honeycomb  $(\sqrt{3} \times \sqrt{3})R30^\circ$  structure is simply due to intervalley quasiparticle interference which was known in graphene to create a  $(\sqrt{3} \times \sqrt{3})R30^\circ$  pattern [17]. In fact, we observed such interference patterns (to be shown in the following text), but they can be distinguished from the  $(\sqrt{3} \times \sqrt{3})R30^\circ$  superstructure observed here, indicating that they have different origins.

Finally, we found a clue by comparing the obtained periodicity,  $0.64 \text{ nm}$ , with the theoretically proposed values, all around  $0.67 \text{ nm}$ , which equal to  $\sqrt{3}a$  ( $a$  is lattice constant of silicene  $0.38 \text{ nm}$ ) [3,7,18]. The experimental value is about 4% smaller than the theoretical ones, indicating that the substrate may exert some influence to result in the contraction of the silicene lattice. Although the detailed mechanism for such contraction still remains to be understood, for a phenomenological model we can fix the periodicity constant in a series of contracted silicene, and fully relax the inner coordinates of Si atoms. Interestingly, we found that the  $(\sqrt{3} \times \sqrt{3})R30^\circ$  reconstruction can indeed be stabilized in a contracted lattice. The Si atoms in previous theoretical low-buckled silicene [Fig. 2(a)] have two different heights, which we named *AB* configurations. To understand the experimentally observed  $(\sqrt{3} \times \sqrt{3})$

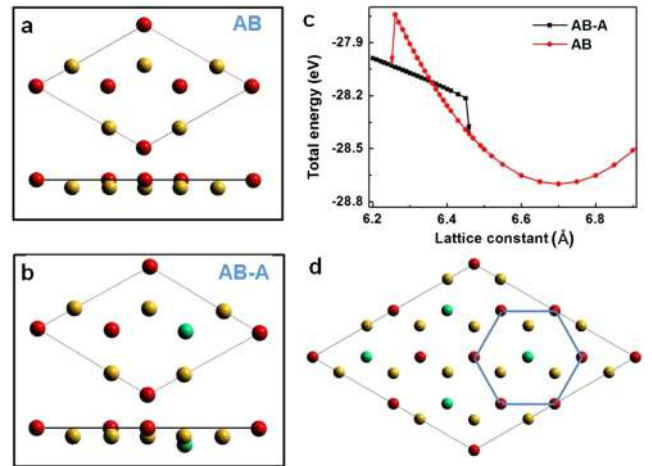


FIG. 2 (color online). (a),(b) The top view and side view of the lattice geometry of the low-buckled silicene structure (*AB*) and  $\sqrt{3} \times \sqrt{3}$  superstructure (*AB $\bar{A}$* ), respectively. Note that in the  $\sqrt{3} \times \sqrt{3}$  superstructure a Si atom with planar coordinate  $(2/3, 2/3)$  is pulled downward. The red (black), yellow (light gray), and green (dark gray) balls represent the *A*, *B*, and  $\bar{A}$  Si atoms, respectively. (c) The structural phase transition diagram of silicene depending on the lattice constants of  $\sqrt{3} \times \sqrt{3}$  superstructure. (d) A larger schematic model illuminating the honeycomb structure of  $\sqrt{3} \times \sqrt{3}$  reconstructed silicene.

superstructure, we considered a configuration with extra buckling of the Si atom with planar coordinate  $(2/3, 2/3)$  downward in one  $(\sqrt{3} \times \sqrt{3})$  silicene unit cell, which is named the  $AB\bar{A}$  configuration, as shown in Fig. 2(b). We carried on a comprehensive study of the structure and stability of a series of contracted silicene structures from the  $AB$  and  $AB\bar{A}$  configuration with and without small perturbation as the initial structures, respectively. Figure 2(c) shows the structure phase transition diagram of the  $(\sqrt{3} \times \sqrt{3})$  superstructure and the low-buckled  $AB$  phase with minimum energy and stability, through structural optimization by keeping the periodicity constant in the contracted region. When the lattice constant is smaller than  $6.25 \text{ \AA}$ , both the  $AB\bar{A}$  configuration and the  $AB$  configuration as the initial structures return to the  $(\sqrt{3} \times \sqrt{3})$  superstructure ( $AB\bar{A}$  configuration). In contrast, when the lattice constant is larger than  $6.45 \text{ \AA}$ , an initial  $AB\bar{A}$  configuration will spontaneously transfer to the low-buckled  $AB$  configuration. The  $\sqrt{3} \times \sqrt{3}$  superstructure ( $AB\bar{A}$  configuration) is robust and stable when the lattice constant is smaller than  $6.35 \text{ \AA}$ . Taking into account the scale error, our experimental silicene sample may be in the  $AB\bar{A}$  phase with the  $(\sqrt{3} \times \sqrt{3})$  superstructure.

Silicene has been predicted to have Dirac-type electron structure around the Fermi energy [3,7], similar to that of graphene. This however remains to be experimentally confirmed. On the other hand, measuring the characteristic electronic structure of our silicene film can also provide additional proof of the basic atomic structure of our film. For graphene, the Dirac cones are located at high symmetric  $K$  points in the 2D Brillouin zone. The constant energy contours in reciprocal space cut through the electron or hole conical sheets near  $E_F$  resulting in small circles centered at the  $K$  points [shown in Fig. 3(d)]. Free carriers can be scattered within the small circles (intravalley scattering) or between circles (intervalley scattering) [19–21], resulting in QPI patterns in real space, which have been observed by STS. Similar QPI patterns are expected to exist for silicene based on the same analysis.

We have observed short-wavelength interference patterns in STM images of silicene, consistent with an intervalley scattering. From Fig. 3(d) one can see that the wave vectors of intervalley scattering,  $q_2$ , are close in length to the wave vectors  $K_{1(2)}$ , which have a length of  $1/\sqrt{3}$  in reciprocal space, corresponding to a  $(\sqrt{3} \times \sqrt{3})R30^\circ$  periodicity in real space. Such a  $(\sqrt{3} \times \sqrt{3})R30^\circ$  interference pattern had been observed in areas near step edges and defects which serve as scattering centers, as illustrated in Figs. 3(a)–3(c). The red circle in (b) is drawn near the step edge, and in (c) it is along the close-packing direction away from the point defect. Within the red circles the surface is imaged as close-packed protrusions. In contrast, within the white circles the surface is imaged as a honeycomb structure, as commonly observed in large area silicene islands. Therefore, we can find a clear phase shift between the

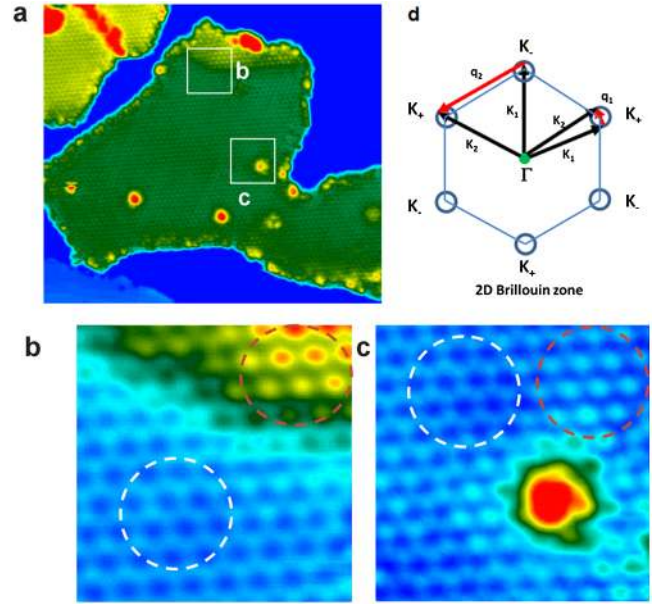


FIG. 3 (color online). (a) The STM image ( $50 \text{ nm} \times 45 \text{ nm}$ ) of a silicene island taken at tip bias  $-1.1 \text{ V}$ . The white squares label two typical defect sites, namely the step edge and point defect, where scattering patterns with hexagonal close packing structures were observed, as enlarged in (b) and (c). (d) Schematic of the 2D Brillouin zone [blue (thin gray) lines], constant energy contours [blue (gray) rings at  $K$  points], and the two types of scattering vectors:  $q_1$  [intravalley, short red (dark gray) arrow] and  $q_2$  [intervalley, long red (dark gray) arrow].

atomic corrugation and the QPI pattern, although both of them exhibit  $(\sqrt{3} \times \sqrt{3})R30^\circ$  periodicity. The QPI pattern originating from the scattering center extends for a length of only a few nm. Such a characteristic decaying length cannot make us assign the  $(\sqrt{3} \times \sqrt{3})R30^\circ$  structure observed on larger islands simply as a QPI pattern. The observation of QPI patterns, consistent with analysis based on the theoretical band structure of silicene, is additional proof that the underlying atomic structure of our film is graphenelike, with only some buckling that does not change the basic electronic structure of the film.

To investigate the electronic structure of silicene in more detail, we performed STS measurements ( $dI/dV$  curves and maps) on the film. A typical  $dI/dV$  curve taken at  $77 \text{ K}$  is shown in Fig. 4(a). Beside the pronounced peak at  $0.9 \text{ V}$ , there is a small dip located at about  $0.5 \text{ V}$ , which is attributed to the position of the Dirac point (DP) of silicene. The dip in the  $dI/dV$  curve corresponding to the DP is not very obvious compared with that of graphene [22,23], which is due to the pronounced electronic density of states of the underlying Ag(111) substrate superimposed in the  $dI/dV$  spectra. Figures 4(c)–4(e) are  $dI/dV$  maps of a silicene island (second layer of silicene) on the single layer of silicene, reflecting the distribution of local density of states in real space. Wavelike QPI patterns near the boundary of the island were clearly observed. The wavelength changes as a function of tip bias voltage. As the bias

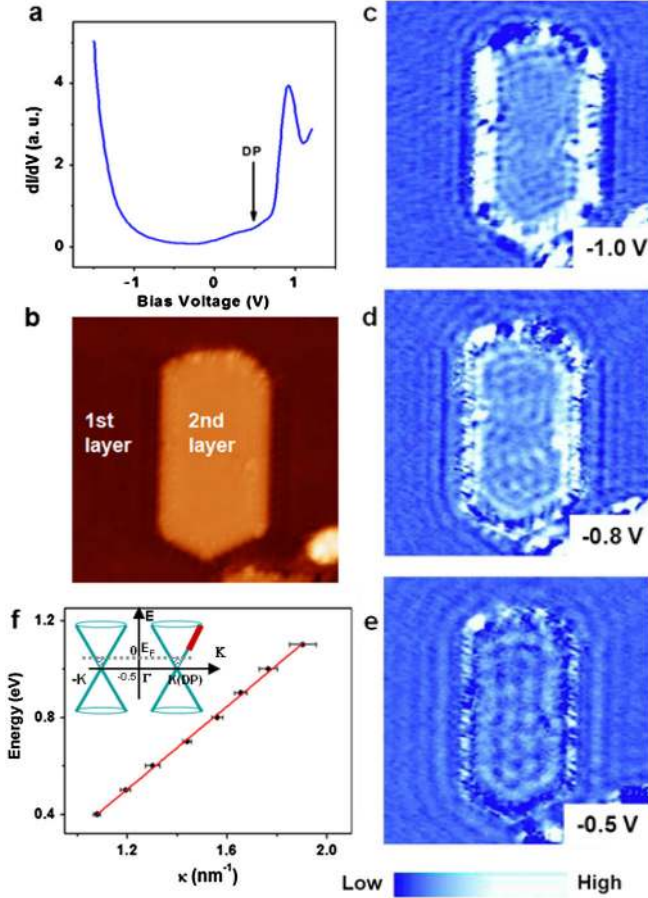


FIG. 4 (color online). (a)  $dI/dV$  curves taken at 77 K. The position of the DP is labeled. (b) The STM image ( $40 \text{ nm} \times 40 \text{ nm}$ ) of the 1 ML silicene surface containing an island of second layer taken at tip bias  $-1.0 \text{ V}$ . (c), (d), and (e)  $dI/dV$  maps of the same area as (b) taken at tip bias  $-1.0$ ,  $-0.8$ , and  $-0.5 \text{ V}$ , respectively. (f) Energy dispersion as a function of  $\kappa$  for silicene determined from the wavelength of QPI patterns. The inset shows a schematic drawing of the overall band structure, with the relative location of the DP,  $E_F$ , and our data points [red (thick gray) line].

increases from  $-0.4 \text{ V}$  to  $-1.1 \text{ V}$ , the wavelength decreases correspondingly from  $2.8 \text{ nm}$  to  $1.6 \text{ nm}$ .

The 2D constant-energy contours in reciprocal space [Fig. 3(d)] are used to understand the QPI patterns. The wave vector  $q_1$  of intravalley scattering connects points within a single constant-energy circle and determines the observed long wavelength interference pattern. In order to deduce the quasiparticle energy-momentum dispersion relation, we drew the  $E(\kappa)$  curve in Fig. 4(f), where  $\kappa$  is the radius of constant-energy circle at  $K$  point with  $2\kappa = q_1$ . The values of  $q_1$  are determined by measuring the wavelength of QPI patterns in  $dI/dV$  maps. We found  $\kappa$  varied linearly with energy, with Fermi velocity  $V_F = (1.2 \pm 0.1) \times 10^6 \text{ m/s}$ . The  $\kappa = 0$  energy intercept gives the Dirac energy,  $E_F - E_D = 0.52 \pm 0.02 \text{ eV}$ , consistent with the position of the DP in the  $dI/dV$  spectra [Fig. 4(a)].

The linear  $E-\kappa$  dispersion proves the existence of the Dirac cone in electronic band structures of silicene. The surprising large Fermi velocity, comparable with that of graphene [19,24], suggests the prospective applications comparable with those that have been proposed or realized in graphene.

In summary, we have obtained silicene film on Ag(111) and observed an unexpected  $\sqrt{3} \times \sqrt{3}$  buckled superstructure. QPI patterns resulting from both intervalley and intravalley scattering were found, and the study of the linear energy-momentum dispersion and the measured Fermi velocity, as large as  $10^6 \text{ m/s}$ , proves that quasiparticles in silicene behave as massless Dirac fermions. Based on these results, many further interesting studies may proceed, including gap opening in silicene, which can realize the quantum spin Hall effect [7], and high temperature electron-phonon superconductivity in hydrogenated silicene [25,26].

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*Note added.*—Recently, new results for silicene sheets on Ag(111) were published [27,28]. These works reported a  $4 \times 4$  phase and a  $\sqrt{13} \times \sqrt{13}$  phase, consistent with the phases *b* and *c* in our Supplemental Material, respectively [13]. The preparation of a silicene sheet on a  $\text{ZrB}_2$  substrate was also reported [29].

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- [1] K. Takeda and K. Shiraishi, *Phys. Rev. B* **50**, 14916 (1994).
- [2] G. G. Guzmán-Verri and L. C. L. Voon, *Phys. Rev. B* **76**, 075131 (2007).
- [3] S. Cahangirov, M. Topsakal, E. Akturk, H. Sahin, and S. Ciraci, *Phys. Rev. Lett.* **102**, 236804 (2009).
- [4] B. Aufray, A. Kara, S. Vizzini, H. Oughaddou, C. Léandri, B. Ealet, and G. L. Lay, *Appl. Phys. Lett.* **96**, 183102 (2010).
- [5] P. D. Padova *et al.*, *Appl. Phys. Lett.* **96**, 261905 (2010).
- [6] B. Lalmi, H. Oughaddou, H. Enriquez, A. Kara, S. Vizzini, B. Ealet, and B. Aufray, *Appl. Phys. Lett.* **97**, 223109 (2010).
- [7] C. C. Liu, W. X. Feng, and Y. G. Yao, *Phys. Rev. Lett.* **107**, 076802 (2011).
- [8] A. K. Geim and K. S. Novoselov, *Nature Mater.* **6**, 183 (2007).
- [9] A. H. Castro Neto, F. Guinea, N. M. R. Peres, K. S. Novoselov, and A. K. Geim, *Rev. Mod. Phys.* **81**, 109 (2009).
- [10] C. L. Kane and E. J. Mele, *Phys. Rev. Lett.* **95**, 226801 (2005).

- [11] C. C. Liu, H. Jiang, and Y. G. Yao, *Phys. Rev. B* **84**, 195430 (2011).
- [12] B. A. Bernevig, T. L. Hughes, and S. C. Zhang, *Science* **314**, 1757 (2006).
- [13] See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevLett.109.056804> for experimentally observed different metastable phases formed by Si on Ag(111) at different growth conditions.
- [14] B. Feng, Z. Ding, S. Meng, Y. Yao, X. He, P. Cheng, L. Chen, and K. Wu, *Nano Lett.* **12**, 3507 (2012).
- [15] J. P. Perdew, K. Burke, and M. Ernzerhof, *Phys. Rev. Lett.* **77**, 3865 (1996).
- [16] G. Kresse and J. Furthmüller, *Phys. Rev. B* **54**, 11 169 (1996).
- [17] Y. Niimi, T. Matshui, H. Kambara, K. Tagami, M. Tsukada, and H. Fukuyama, *Phys. Rev. B* **73**, 085421 (2006).
- [18] Y. Ding and J. Ni, *Appl. Phys. Lett.* **95**, 083115 (2009).
- [19] G. M. Rutter, J. N. Crain, N. P. Guisinger, T. Li, P. N. First, and J. A. Stroscio, *Science* **317**, 219 (2007).
- [20] I. Brihuega, P. Mallet, C. Bena, S. Bose, C. Michaelis, L. Vitali, F. Varchon, L. Magaud, K. Kern, and J. Y. Veuillen, *Phys. Rev. Lett.* **101**, 206802 (2008).
- [21] J. Xue, J. Sanchez-Yamagishi, K. Watanabe, T. Taniguchi, P. Jarillo-Herrero, and B. J. LeRoy, *Phys. Rev. Lett.* **108**, 016801 (2012).
- [22] Y. B. Zhang, V. W. Brar, F. Wang, C. Girit, Y. Yayon, M. Panlasigui, A. Zettl, and M. F. Crommie, *Nature Phys.* **4**, 627 (2008).
- [23] L. Zhao *et al.*, *Science* **333**, 999 (2011).
- [24] C. Berger *et al.*, *Science* **312**, 1191 (2006).
- [25] A. Kara, C. Leandri, M. E. Davila, P. De Padova, B. Ealet, H. Oughaddou, B. Aufray, and G. Le Lay, *J. Supercond. Novel Magnetism* **22**, 259 (2009).
- [26] G. Savini, A. C. Ferrari, and F. Giustino, *Phys. Rev. Lett.* **105**, 037002 (2010).
- [27] P. Vogt, P. D. Padova, C. Quaresima, J. Avila, E. Frantzeskakis, M. Asensio, A. Resta, B. Ealet, and G. L. Lay, *Phys. Rev. Lett.* **108**, 155501 (2012).
- [28] C. L. Lin, R. Arafune, K. Kawahara, N. Tsukahara, E. Minamitani, Y. Kim, N. Takagi, and M. Kawai, *Appl. Phys. Express* **5**, 045802 (2012).
- [29] A. Florence, R. Friedlein, T. Ozaki, H. Kawai, Y. Wang, and Y. Yamada-Takamura, *Phys. Rev. Lett.* **108**, 245501 (2012).