Evidence for photometric activity cycles in 3203 Kepler stars*

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ABSTRACT

Context. In recent years it has been claimed that the length of stellar activity cycles is determined by the stellar rotation rate. It has been observed that the cycle period increases with rotation period along two distinct sequences, known as the active and inactive sequences. In this picture the Sun occupies a solitary position between the two sequences. Whether the Sun might undergo a transitional evolutionary stage is currently under debate.

Aims. Our goal is to measure cyclic variations of the stellar light curve amplitude and the rotation period using four years of *Kepler* data. Periodic changes in the light curve amplitude or the stellar rotation period are associated with an underlying activity cycle.

Methods. Using a recent sample of active stars we compute the rotation period and the variability amplitude for each individual *Kepler* quarter and search for periodic variations of both time series. To test for periodicity in each stellar time series we consider Lomb-Scargle periodograms and use a selection based on a false alarm probability (FAP).

Results. We detect amplitude periodicities in 3203 stars between $0.5 < P_{cyc} < 6$ yr covering rotation periods between $1 < P_{rot} < 40$ days. Given our sample size of 23 601 stars and our selection criteria that the FAP is less than 5%, this number is almost three times higher than that expected from pure noise. We do not detect periodicities in the rotation period beyond those expected from noise. Our measurements reveal that the cycle period shows a weak dependence on rotation rate, slightly increasing for longer rotation periods. We further show that the shape of the variability deviates from a pure sine curve, consistent with observations of the solar cycle. The cycle shape does not show a statistically significant dependence on effective temperature.

Conclusions. We detect activity cycles in more than 13% of our final sample with a FAP of 5% (calculated by randomly shuffling the measured 90-day variability measurements for each star). Our measurements do not support the existence of distinct sequences in the $P_{\text{rot}} - P_{\text{cyc}}$ plane, although there is some evidence for the inactive sequence for rotation periods between 5–25 days. Unfortunately, the total observing time is too short to draw sound conclusions on activity cycles with similar lengths to that of the solar cycle.

Key words. Sun: activity - stars: activity - starspots - stars: rotation - techniques: photometric

1. Introduction

The origin of the 11-year solar activity cycle is one of the most important questions in solar and stellar physics. Over the course of one cycle the Sun undergoes phases of strong and weak activity. During the active phase of the solar cycle, dark sunspots appear on the solar surface. These spots have lifetimes lasting from days to a few months (Petrovay & van Driel-Gesztelyi 1997), decaying into smaller magnetic concentrations which are bright (called faculae). Individual faculae have short lives, as the magnetic flux associated with them is buffeted by turbulent convective motions; however, since magnetic flux is conserved, the faculae reform and the lifetime of extended patches of faculae have much longer lifetimes. The magnetic cycle of the Sun is thus accompanied by brightness variations. This variability, which is an evolving combination of dark and bright features on the rotating surface of the Sun, is weak during times of minimum activity and strong during times of maximum activity. Although these activity phenomena act on different timescales, they can all be detected in the total solar irradiance (TSI) data (see, e.g., Domingo et al. 2009).

On stars other than the Sun, cyclic activity has also been observed through (long-term) brightness changes caused by epochs of increased occurrence of active regions on the surface or in the low stellar atmosphere. The magnetic field of the active regions transports energy into the chromosphere, which leads to increased chromospheric emission, most evident in the cores of the CaII H&K lines. The strength of this enhancement in the emission is usually described by the Mount Wilson S-index (Vaughan et al. 1978) or the related quantity $R'_{\rm HK}$ (Linsky et al. 1979). A long-term study of chromospheric activity of mainsequence stars in the solar neighborhood was initiated at the Mount Wilson observatory in the late 1960s. Wilson (1978) presented CaII H&K flux measurements with the first evidence of cyclic variations. Vaughan & Preston (1980) found that FGKM stars exhibit different levels of chromospheric activity with a deficiency of stars of intermediate activity. These authors explained the dearth of intermediate active stars in terms of an activity decline with age. In contrast, Noyes et al. (1984a) found a tight correlation between $R'_{\rm HK}$ activity and the Rossby number Ro = $P_{\rm rot}/\tau_{\rm c}$ (defined as rotation period $P_{\rm rot}$ divided by the convective overturn time τ_c), but could not detect evidence of the – nowadays called – Vaughan-Preston gap. Noyes et al. (1984b) detected cyclic variations of the chromospheric activity levels in 13 main-sequence stars providing the relation $P_{\rm cyc} \propto P_{\rm rot}^{1.25}$ between the cycle period and the rotation period. Baliunas et al. (1995)

^{*} A table containing all cycle periods and time series is only available in electronic form at the CDS via anonymous ftp to cdsarc.u-strasbg.fr(130.79.128.5) or via http://cdsarc.u-strasbg.fr/viz-bin/qcat?J/A+A/603/A52

investigated stars of spectral type G0-K5 and found that these can be roughly separated by their mean S-index: fast rotators are assumed to be young objects with high S-indices, whereas older stars rotate more slowly and show lower values of the Sindex, again providing evidence for the Vaughan-Preston gap. Baliunas et al. (1996) tried to understand cyclic activity in terms of stellar dynamo theory, proving a relation between the ratio of the cycle and rotation period and the dynamo number D according $P_{\rm cyc}/P_{\rm rot} \propto D^{0.74}$. Brandenburg et al. (1998) and Saar & Brandenburg (1999) studied the relation between magnetic dynamo cycle period, rotation period, activity level, and stellar age. These authors were the first to claim the existence of the active (A) and inactive (I) sequences, two almost parallel branches separated by a factor of six in $P_{\rm cyc}$ following the relation $P_{\rm cyc}/P_{\rm rot} \propto {\rm Ro}^{0.5}$. Additionally, for fast rotators with $P_{\rm rot} < 3$ d these authors found evidence for a third branch with opposite slope $P_{\rm cyc}/P_{\rm rot} \propto {\rm Ro}^{-0.4}$. Böhm-Vitense (2007) confirmed that chromospherically active stars follow the A- and Isequences in the $P_{\rm rot} - P_{\rm cyc}$ plane. Slow rotators exhibit shorter cycle periods populating the I-sequence, whereas stars rotating faster than the Sun exhibit long cycle periods located on the Asequence. Some A-sequence stars exhibit secondary cycle periods located on the I-branch. The Sun is located in between the two branches. In contrast, do Nascimento et al. (2015) found that the Sun does not have a solitary position between the two branches, but lies on a newly proposed solar-analog sequence.

All the cycle period measurements above are based on cyclic CaII H&K emission, but long-term brightness variations also indicate cyclic activity (see, e.g., Baliunas & Vaughan 1985, which nicely links solar activity to photometric and chromospheric activity on stars other than the Sun). Oláh et al. (2000) and Oláh & Strassmeier (2002) investigated activity cycles from long-term V-band photometry. Messina & Guinan (2002) found evidence of cyclic activity in six solar-analog stars also using long-term V-band observations. In a subsequent paper, Messina & Guinan (2003) found that the mean rotation period also changes over the course of the activity cycle. Lockwood et al. (2007) provided parallel measurements of photometric and chromospheric variability, but unfortunately did not determine any cycle periods. Oláh et al. (2009) found a positive correlation between rotational and cycle periods for 20 active stars using time-frequency analysis. A similar approach was applied to *Kepler* light curves of very fast rotators ($P_{rot} \le 1 d$) by Vida et al. (2014). These authors found hints of cyclic activity on timescales of 300-900 days analyzing cyclic changes of the mean rotation period. Using CoRoT data Ferreira Lopes et al. (2015) detected cyclic variability in 16 FGK type stars. These authors found evidence of a possible third sequence, parallel to the previously identified active and inactive sequences, but shifted by a factor of 1/4 towards shorter cycle periods. Lehtinen et al. (2016) found evidence of activity cycles in 18 solar-type stars using decades-long photometric observations. The cycle periods follow the sequences defined by Saar & Brandenburg (1999), bridging chromospherically and photometrically derived activity cycles. In recent years, activity cycles have also been detected from asteroseismology through periodic shifts in the mode frequencies (García et al. 2010; Salabert et al. 2016).

The primary *Kepler* mission provides continuous, highprecision photometry of thousands of stars over four years of observation. Owing to the presence or absence of active regions on the visible hemisphere (e.g., spots, plage) the stellar flux gradually changes. Such long-term changes are observed in the TSI data, showing significant changes over the course of the 11-year cycle. In addition, the preferred latitude of active region occurrence changes. Owing to differential rotation equatorial spots have a shorter rotation period than higher latitude spots. Our main goal is to detect cyclic changes of the light curve amplitude and the stellar rotation period, which we interpret as an indication of a photometric activity cycle. The largest sample with accurately determined rotation periods is provided by McQuillan et al. (2014, hereafter McQ14). This huge sample of 34 030 active stars forms the basis of our activity cycle survey.

2. Data and method

2.1. Motivation from the Sun

As mentioned in the introduction, the brightness of the Sun (as measured by the TSI) varies due to the passage of dark sunspots and bright plage regions across the solar disk. The top panel of Fig. 1 shows the TSI time series from Shapiro et al. (2016) produced with the SATIRE-S model (Yeo et al. 2014). It is readily seen that at times of activity maxima (1979, 1989, 2000, 2014) the TSI varies rapidly, while there are only small variations during times of activity minima (1976, 1986, 1996 and 2008).

To quantify the time-dependent variability in a manner which we can easily extend to the *Kepler* data, we begin by rebinning the TSI data from a 1-min cadence to a 1-day cadence. We then break the data into N = 73 segments of 90-day length, and following Basri et al. (2011), we define the variability range $R_{\rm var}$ for that 90-day period as the difference between the 5th and 95th percentiles of the TSI measurements in the 90-day segment. This activity measure is similar to the magnetic index S_{ph} defined by Mathur et al. (2014). For each of the N = 73 segments we thus obtain a measurement of the variability $R_{var}(n)$, n = 1, ..., N. Using the midpoints of each 90-day time segment we create a time series $R_{var}(t)$, shown in the lower panel of Fig. 1. Computing the Lomb-Scargle periodogram of the time series $R_{var}(t)$ reveals a peak at $P_{\text{cyc, var}} = 10.98 \pm 0.26$ yr, which is compatible with other estimates of the solar cycle length. This shows that this method is capable of detecting the correct cycle period from the variability of a star's light curve.

2.2. Kepler data

Most *Kepler* data is released in segments of ~90 days (known as quarters) with exceptions for the quarters Q0 (~10 days), Q1 (~33 days), and Q17 (~32 days). There is an ongoing debate regarding how to remove instrumental effects from the time series. The data studied in this paper rely on the PDC-MAP pipeline using the following versions¹: Version 2.1 for Q0-Q4 & Q9-Q11; Version 3.0 for Q5-Q8 & Q12-Q14; and Version 5.0 for Q15-Q17. For each star we create a time sequence t(n), n = 0, ..., 17, consisting of the midpoints of the time segments covered by the quarters *Qn* when a star was observed.

The upper panel of Fig. 2 shows the light curve of the *Kepler* star KIC 2714077. Each light curve is normalized and centered around zero by dividing the flux by its median value and subtracting unity. This normalization is necessary because the *Kepler* spacecraft rolls between consecutive quarters, so the stars fall on different CCDs accompanied by offsets of the PDC-MAP flux between the quarters. Because the stellar variability we are interested in is that due to the rotation of spots and faculae across the disk, we further fit and subtract a second-order polynomial from the data in each quarter. This will substantially

¹ The version number can be found in the primary headers of the FITS files using the *FILEVER* keyword.

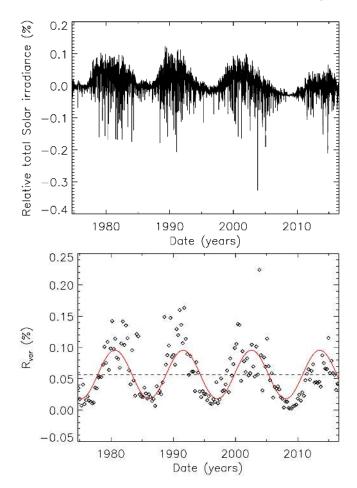


Fig. 1. Upper panel: measured total solar irradiance, binned in 1-day intervals, covering more than 40 yr of observations. The approximately 11-year solar cycle can be seen as the Sun's brightness changes from almost constant during activity minima to highly variable during activity maxima. *Lower panel*: variability range of the TSI data, R_{var} , as a function of time. We define R_{var} as the difference between the 5th and 95th percentiles of the 1-day binned intensity in each 90-day period in time. The red curve shows a sine fit to the R_{var} time series with a period $P_{cyc, var} = 10.98 \pm 0.26$ yr.

reduce the influence of possible longer term instrumental issues, while only slightly affecting modulation comparable to or faster than the rotation period of the star. We then use the R_{var} defined in the same way as for the Sun as a measure of the variability. The red triangles in the upper panel of Fig. 2 indicate the 5th and 95th percentiles of the relative flux in each quarter. The difference between the 5th and 95th percentiles equals R_{var} , which is shown as time series $R_{var}(t)$ in the middle panel of Fig. 2. The time series $R_{var}(t)$ shows strong periodicity with a period of $P_{cyc, var} = 3.51$ yr, shown as red sine fit to the data.

In addition to measuring R_{var} , for each star we also measure the stellar rotation period in each quarter using the Lomb-Scargle periodogram resulting in a time series $P_{rot}(t)$, which is shown for the same stars in the lower panel of Fig. 2. The quarterly measured rotation periods scatter around the mean rotation period $P_{rot} = 13.88$ d. In contrast to R_{var} the time series $P_{rot}(t)$ does not show strong periodicity. The Lomb-Scargle periodogram finds a peak at a period of less than one year, which is visualized by the red sine fit to the data. As we show in the following, strong periodicity of the variability range and random scatter of the rotation period time series is quite common for many stars in our sample.

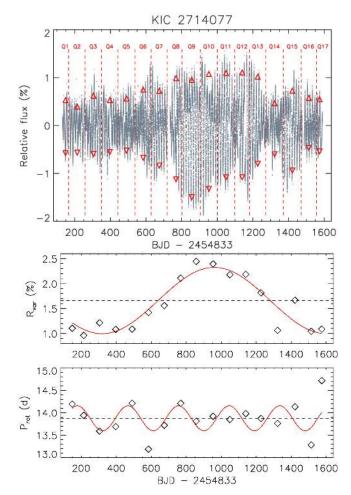


Fig. 2. Upper panel: Kepler light curve of the star KIC 2714077. Vertical red lines indicate the quarters Q1-Q17. Red triangles at the bottom and the top show the 5th and 95th percentile of the intensity, respectively. For each quarter we measure the intensity difference between the upper and the lower triangle, which equals R_{var} . Middle panel: time series of the variability range $R_{var}(t)$ of the same star, defined as the difference between the 5th and 95th percentiles in each quarter. Lower panel: time series of the rotation periods $P_{rot}(t)$ of the same star. The red curve in the middle and lower panels show the best sine fit to the $R_{var}(t)$ and $P_{rot}(t)$ time series, respectively.

Owing to the limited quarter length we set upper limits on the detectable rotation periods: 10 days for Q0, 30 days for Q1 and Q17, and 45 days for the other quarters. All period measurements deviating by more than 30% from the median period of all quarters were discarded. The R_{var} measurements of these quarters were also discarded because instrumental artifacts might dominate the photometry. To achieve meaningful variability and rotation period time series we require measurements in at least N = 12 quarters. These conditions shrink the initial McQ14 sample to 23 601 stars. We suspect that there might be more candidates with detectable activity cycle that have been missed, due to our constraints.

2.3. Analysis of the time series

For each star we have three time series, t(n), $R_{var}(n)$, and $P_{rot}(n)$ of length N, which is the number of quarters where a rotation period could be measured. The Lomb-Scargle periodogram is a straightforward way to look for periodicities in unevenly sampled data. We compute the Lomb-Scargle periodogram of the

KIC 10355856 KIC 10904491 0.16 2.2 (a) (b) 2.0 0.14 1.8 R_{vor} (%) 8 1.6 0.12 R. 1.4 1.2 0.10 1.0 0.08 0.8 800 1000 1200 1400 1600 600 800 1000 1200 1400 1600 200 400 600 200 400 BJD - 2454833 BJD - 2454833 KIC 4166395 KIC 3446150 0.07 (c) (d) ♦ 0 1.4 0.06 R_{vor} (%) 1.2 8 0.05 L, S 1.0 0.04 0.8 0.03 0.6 600 800 1000 1200 1400 1600 200 400 600 800 1000 1200 1400 1600 200 400 BJD 2454833 BJD 2454833

Fig. 3. Time series $R_{var}(t)$ of the quarterly measured variability ranges of four *Kepler* stars. The time series in panels **a**)–**c**) pass our selection criterion that randomly permuting the time series R_{var} of this star yields a higher peak in the Lomb-Scargle periodogram in less than 5% of all tested cases (FAP < 5%), whereas panel **d**) shows a time series with (FAP > 5%). The red curve shows the best sine fit to the data, and the dashed black line indicates the mean variability of the star. The $R_{var}(t)$ time series shown in panels **a**)–**d**) exhibit false alarm probabilities FAP = 1.0%, FAP = 0.5%, FAP = 0.3%, and FAP = 17.9%, respectively.

time series R_{var} and P_{rot} (using the time series *t*) to test for periodicity. We call these periods associated with the highest peaks in the Lomb-Scargle periodograms $P_{\text{cyc, var}}$ and $P_{\text{cyc, rot}}$, respectively.

Because we are simply taking the highest peak in the Lomb-Scargle periodogram, $P_{\rm cyc, var}$ and $P_{\rm cyc, rot}$ will be defined for every star. To define a subset of the stars where the inferred periodicity is more likely to be real, we introduce a false alarm probability (FAP) threshold. We calculate the FAP by considering 1000 random permutations of $R_{\rm var}$ and $P_{\rm rot}$. For each permutation we have three time series t, $R_{\rm var}^{\rm permutation}$, and $P_{\rm rot}^{\rm permutation}$. We then create Lomb-Scargle periodograms for these new sequences and again concentrate on the highest peaks. We define the FAP for $R_{\rm var}$ as the number of cases where a higher peak appears in the Lomb-Scargle periodogram of $R_{\rm var}^{\rm permutation}$ than the peak found for $R_{\rm var}$ divided by the number of permutations considered.

We define the periodic sample by selecting those stars for which the FAP was less then 5%. The samples are defined independently for $P_{\text{cyc, var}}$ and $P_{\text{cyc, rot}}$. The time series $R_{\text{var}}(t)$ of three representative stars from the periodic sample are shown in panels (a)–(c) in Fig. 3. For comparison panel (d) shows a star with FAP > 5%. We discuss the issue of false positives in the results section.

2.4. Method sensitivity

To test the sensitivity of the method to different cycle periods, we created a set of 1000 synthetic time series $R_{\text{var}}^{\text{syn}}(t)$. We randomly selected N data points, $12 \le N \le 18$, from the time series t(n). The time series $R_{\text{var}}^{\text{syn}}(t)$ are pure sine functions with input periods $P_{\text{cyc, in}}$ uniformly distributed between 0.5 yr < $P_{\text{cyc, in}} < 10$ yr.

The amplitudes are uniformly distributed random numbers, comparable to the sine fit amplitudes of the periodic sample. We add noise to the time series $R_{var}^{syn}(t)$, which we define as the standard deviation of the difference between $R_{var}(t)$ and the best sine fit from the periodic sample. The noise distribution is Gaussian. Because the quantity R_{var} describes a relative flux difference, the sine amplitude and the noise are linearly correlated. For simplicity, the phase and the offset of the sine curves are set to zero. We analyze the set of synthetic light curves using the Lomb-Scargle periodogram, and compare the returned cycle periods $P_{cyc, out}$ to the input periods $P_{cyc, in}$ in Fig. 4. This Monte Carlo test reveals that we can trust cycle periods of up to six years.

The uncertainties of the derived cycle periods have been estimated in the same way. For different cycle period bins between 0.5–10 years we created a set of 1000 synthetic light curves for each bin, and applied the above analysis. In Fig. 4 the red crosses show the median value of $P_{\text{cyc,out}}$ and the error bars show the median absolute difference between $P_{\text{cyc, in}}$ and $P_{\text{cyc, out}}$. We omit error bars for $P_{\text{cyc, in}} > 7$ yr because the uncertainties become huge. The associated periods and uncertainties up to $P_{\text{cyc, in}} < 7$ yr are given in Table 1.

3. Results

3.1. Cycle periods

The cycle period distribution of $P_{\text{cyc, var}}$ for all stars satisfying FAP < 5% is shown in black in the left panel of Fig. 5, and the distribution of $P_{\text{cyc, rot}}$ for all stars satisfying FAP < 5% is shown in black in the right panel of Fig. 5. Assuming the time series $R_{\text{var}}(t)$ and $P_{\text{rot}}(t)$ do not contain any periodic signals, 1180 stars

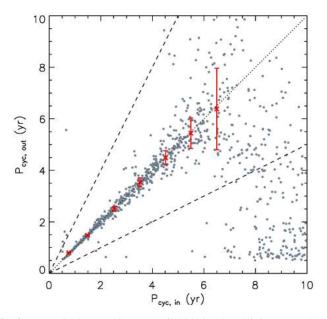


Fig. 4. Monte Carlo test using a set of 1000 simulated light curves showing the input periods $P_{\text{cyc,in}}$ and the returned periods $P_{\text{cyc,out}}$. The black lines show different $P_{\text{cyc,out}}$: $P_{\text{cyc,in}}$ period ratios: the dotted line shows the 1:1 period ratio, and the upper and lower dashed line show the 2:1 and 1:2 period ratio, respectively. The cycle period uncertainties (shown as red error bars) were estimated from ten different sets of 1000 simulated light curves for each input period bin between 0.5–10 yr. The red crosses show the median value of $P_{\text{cyc,out}}$ and the error bars show the median absolute difference between $P_{\text{cyc,in}}$ and $P_{\text{cyc,out}}$. We omit error bars for $P_{\text{cyc,in}} > 7$ yr because the uncertainties become huge.

Table 1. Intervals of the input periods $P_{\text{cyc, in}}$, median values of the output cycle periods $P_{\text{cyc, out}}$, and the associated absolute and relative uncertainties of $P_{\text{cyc, out}}$ as shown in Fig. 4.

<u></u>			
$P_{\rm cyc,in}$	$\langle P_{\rm cyc,out} \rangle$	$\langle \Delta P_{\rm cyc,out} \rangle$	$\langle \Delta P_{\rm cyc,out} / P_{\rm cyc,out} \rangle$
(yr)	(yr)	(yr)	(%)
[0.5, 1.0]	0.80	0.01	1.2
[1.0, 2.0]	1.48	0.03	2.0
[2.0, 3.0]	2.50	0.09	3.6
[3.0, 4.0]	3.54	0.17	4.9
[4.0, 5.0]	4.49	0.27	6.1
[5.0, 6.0]	5.44	0.59	11.2
[6.0, 7.0]	6.38	1.58	26.2

would be obtained after applying a limit of FAP < 5% to a total sample of 23 601 stars. For the cycle periods $P_{\text{cyc, var}}$ of the variability amplitude $R_{\text{var}}(t)$ we detect 3203 cycles. This number is significantly higher than the expected number of false-positive detections.

To test whether the different number of detections might be an artifact of the method, we performed a Monte Carlo simulation where we replaced the time series $R_{var}(t)$ of each star with a random permutation of that star's time series. We then applied exactly the same analysis as was applied to the original time series (see Sect. 2.3).

This randomization process is repeated 500 times for the whole sample. Comparing the peak height of the initially randomized time series to the peak heights of the remaining ones yields the false alarm probability FAP_{MC} for each star. We repeated this process 500 times, and found that on average 1122 stars satisfied our selection criterion. This is consistent

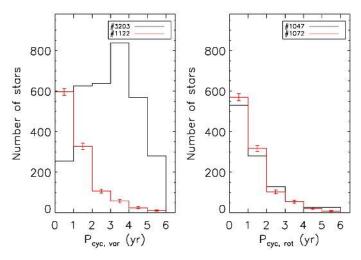


Fig. 5. *Left panel*: distribution of cycle periods $P_{\text{cyc, var}}$ derived from all time series R_{var} which satisfied our selection criterion FAP < 5%. The black distribution shows all cases where random permutations of time series R_{var} yielded a higher peak in the Lomb-Scargle periodogram than the original time series in less than 5% of all cases. The red curve shows the results from a Monte Carlo experiment where the analysis was performed on random permutations of each R_{var} time series for the whole sample. Error bars are based on the Monte Carlo simulations and show one standard deviation. *Right panel*: distribution of $P_{\text{cyc, rot}}$, the cycle period determined from variations in the determined rotation periods P_{rot} , satisfying the selection criterion of FAP < 5%. Again the black histogram shows the cycle period distribution of the 5% strongest periodicity data, and the red curve shows the period distribution obtained from the Monte Carlo experiments.

with the expectation from the condition $\text{FAP}_{\text{MC}} < 5\%$ applied to 23 601 stars.

The averaged cycle period distribution from the Monte Carlo experiment is shown in red in Fig. 5, with error bars showing the standard deviation of different realizations. It is obvious that the distributions of $P_{\text{cyc, var}}$ of the original and the randomized data are very different. The distribution of the real data has a maximum of around three years, whereas the randomized sample mostly contains short cycles between one and two years.

In the right panel the distribution of $P_{\text{cyc, rot}}$ reveals a completely different picture. The distributions of the real data and the randomized sample are largely consistent, and the total number of detections is almost identical. This shows that we cannot distinguish whether a cyclic change of the rotation period might indicate an activity cycle or if this change occurs just by chance. We conclude that we cannot use the quantity $P_{\text{cyc, rot}}$ to search for activity cycles. Thus, we concentrate on $P_{\text{cyc, var}}$ in the following.

Figure 6 shows the dependence of $P_{\text{cyc, var}}$ on the stellar rotation period P_{rot} for all 3203 stars of the periodic sample. The number of stars in each bin on the abscissa is given in black above the color bins.

We turn now to the issue of the cycle period's dependence on rotation period. Our observations only allow activity periods of less than six years to be detected. From Fig. 6 it is obvious that most cycle periods have a length between 2–4 years (compare Fig. 5), slightly clumping at rotation periods between 10– 20 days. Up to rotation periods of 25 days the cycle period slightly increases with rotation period. For stars rotating even more slowly the cycle period behavior is hard to determine from the data. Using unit weights for all data points a linear fit yields $P_{\rm cyc, var} = 0.04 P_{\rm rot} + 2.40$ yr, i.e., a weak dependence of the cycle

A&A 603, A52 (2017)

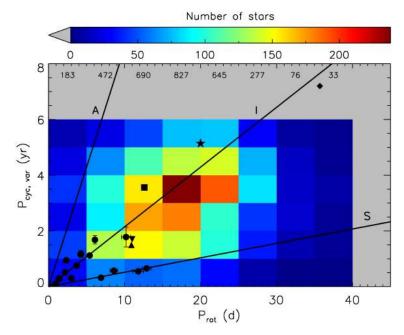


Fig. 6. Distribution of periodic stars with FAP < 5% in the $P_{\text{rot}} - P_{\text{cyc, var}}$ plane; the gray area indicates the limits of our detection reliability. The different black symbols represent cycle period measurements from other authors: circles: Ferreira Lopes et al. (2015), upside down triangle: Egeland et al. (2015), normal triangle: Salabert et al. (2016), square: Moutou et al. (2016), star: Flores et al. (2016), and diamond: Boro Saikia et al. (2016). The solid black lines show the active (A), inactive (I), and short-cycle (S) sequences as defined by Ferreira Lopes et al. (2015). The numbers above the color bins denote the total number of stars in each column.

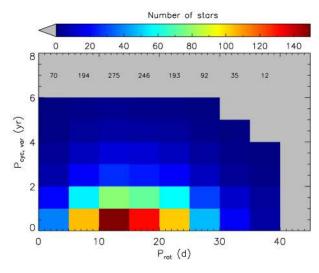


Fig. 7. Distribution of the cycle periods in the $P_{\text{rot}} - P_{\text{cyc,var}}$ plane expected from noise (determined using Monte Carlo experiments).

on rotation period. This result is confirmed by the small Pearson correlation coefficient r = 0.20 between the two quantities.

Some previous authors have reported a correlation between the rotation period of a star and the period of its activity cycle, with periods preferentially lying on active and inactive branches (shown in Fig. 6), and with some stars having detectable periodicities of their activity cycles lying on both branches. Clearly our detection scheme cannot find multiple activity periodicities, and we find that most of our detections are concentrated along the inactive branch; however, the distributions are broad.

For comparison photometric cycle measurements of 16 Co-RoT stars (Ferreira Lopes et al. 2015) are shown as black circles in Fig. 6. The upside down triangle, normal triangle, square, star, and diamond show measurements from Egeland et al. (2015),

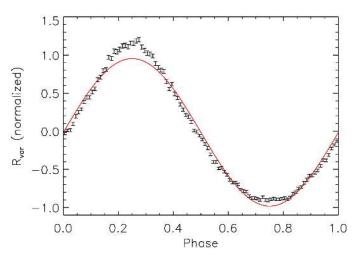
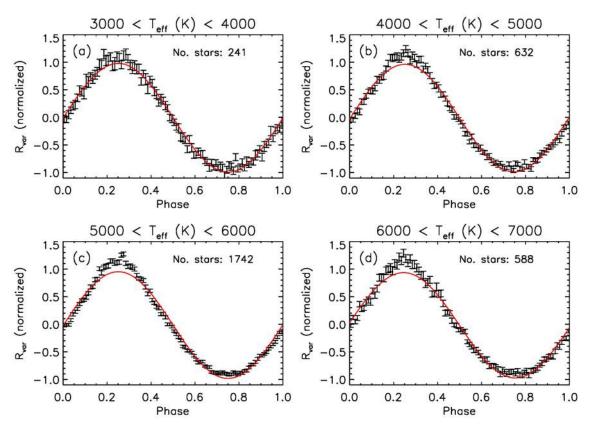


Fig. 8. Phase-folded $R_{var}(t)$ curves of all 3203 stars with strong periodicity. Error bars show the standard error of the $R_{var}(t)$ distribution in each phase bin.

Salabert et al. (2016), Moutou et al. (2016), Flores et al. (2016), and Boro Saikia et al. (2016), respectively. Egeland et al. (2015) have also detected a longer activity cycle of \sim 12 yr for the same star (HD 30495).

To check that the data are not due to a hidden selection bias, we prepared the same figure beginning from the Monte Carlo simulations discussed above. The result is shown in Fig. 7. As expected from pure noise, most cycle period detections lie in the range of 1–2 yr because the periodogram interprets the noise as short cycles. Furthermore, the sample shows a flat distribution, i.e., no dependence on rotation. The associated Pearson correlation coefficient equals r = 0.01, on average, showing that $P_{cyc, var}$ and P_{rot} are uncorrelated.



T. Reinhold et al.: Evidence for photometric activity cycles in 3203 Kepler stars

Fig. 9. Phase-folded $R_{var}(t)$ curves of all periodic stars split up into different temperature bins as indicated on the top of each panel.

3.2. Cycle shapes

So far we have focused on the cycle periods, but have not considered the shape of the variability. Figure 8 shows a phasefolded curve of the variability range R_{var} averaged over all stars of the periodic sample. All R_{var} time series are phase-folded with the measured cycle period and shifted so that they all lie in phase. The mean variability level is subtracted, and the curves are normalized by their amplitudes. The average over all curves is shown in black and the red curve shows a sine fit. The error bars show the standard error in each phase bin. It is clear that the shape of the variability shows some deviation from a pure sine curve. At maximum activity the observed variability curve is more spiky than a sine curve. At minimum activity the variability is rather flat. Applying the same analysis to the $R_{\rm var}$ time series of the TSI data (compare Fig. 1) we also find a flat shape at minimum activity, but a large scatter at maximum activity. To quantify the deviation of the data from the sine curve we compute the reduced chi-square and the associated *p*-value to be $\chi^2_{red} = 16.9$ and p = 0.002, indicating a high statistical significance for the deviation from a sinusoidal shape.

Figure 9 shows the same curve as Fig. 8 for different temperature bins. In each bin the stars have roughly the same cycle period distribution. The coolest stars are nicely represented by a sine curve. Neither the spiky top nor the flat bottom are visible. The shape of the curve gradually changes towards hotter stars. The top becomes more spiky and the bottom flattens out. To quantify the temperature dependence we calculate the difference between each of the six possible pairs of the average curves (a)–(d) in Fig. 9. We do not find a statistically significant difference between the pairs that would indicate a temperature dependence of the cycle period. The slightly different cycle shape in panel (a)

might be explained by the smaller number of stars in this temperature bin.

4. Summary

We detected periodicity of the variability amplitude R_{var} in 3203 stars of the McQ14 sample, which we interpreted as an indication of an underlying activity cycle. To account for random detections, the original sample was compared to a randomized sample. It was found that the number of activity cycle detections was three times higher than expected from pure noise. We did not detect cyclic changes in the rotation period beyond those expected from noise. The $P_{\rm rot} - P_{\rm cyc}$ plane revealed that the cycle period $P_{\rm cyc, var}$ shows weak dependence on the rotation period, slightly increasing with rotation period. In contrast to previous studies (Noves et al. 1984b; Baliunas et al. 1996; Saar & Brandenburg 1999; Böhm-Vitense 2007; Ferreira Lopes et al. 2015; Lehtinen et al. 2016) we did not find evidence of a tight functional dependence between the cycle and the rotation period. Our measurements show a large scatter around the inactive branch. Previous studies (Saar & Brandenburg 1999; Böhm-Vitense 2007) have shown stars with multiple periodicities on the A- and I-branches. These studies were based on chromospheric activity indices, which will have different sensitivities to spots and plages than the photospheric variability detected using Kepler data. Hence it is possible that our finding of a strong I-branch and little evidence of the A- and S-branches might indicate that the I-branch has a stronger signature in the photosphere than the A- or S-branches. In good agreement with the solar cycle the average shape of all activity cycles was found to deviate from a pure sine curve. We could not detect a statistically significant temperature dependence of the activity cycle shape.

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