



Title	Existence and structure of infinitely degenerate zero-energy ground states of a Wess-Zumino type model in supersymmetric quantum mechanics
Author(s)	Ogurusu, O.
Citation	Hokkaido University Preprint Series in Mathematics, 179, 1-26
Issue Date	1993-01
DOI	10.14943/83323
Doc URL	http://hdl.handle.net/2115/68925
Type	bulletin (article)
File Information	pre179.pdf



[Instructions for use](#)

**EXISTENCE AND STRUCTURE OF
INFINITELY DEGENERATE ZERO-
ENERGY GROUND STATES OF A
WESS-ZUMINO TYPE MODEL IN
SUPERSYMMETRIC QUANTUM
MECHANICS**

Osamu Ogurisu

Series #179. January 1993

HOKKAIDO UNIVERSITY
PREPRINT SERIES IN MATHEMATICS

- # 151: G. Ishikawa, Developable of a curve and determinacy relative to osculation-type, 15 pages. 1992.
- # 152: H. Kubo, Global existence of solutions of semilinear wave equations with data of non compact support in odd space dimensions, 25 pages. 1992.
- # 153: Y. Watatani, Lattices of intermediate subfactors, 33 pages. 1992.
- # 154: T. Ozawa, On critical cases of Sobolev inequalities, 11 pages. 1992.
- # 155: M. Ohnuma, M. Sato, Singular degenerate parabolic equations with applications to geometric evolutions, 20 pages. 1992.
- # 156: S. Izumiya, Perestroikas of optical wave fronts and graphlike Legendrian unfoldings, 13 pages. 1992.
- # 157: A. Arai, Momentum operators with gauge potentials, local quantization of magnetic flux, and representation of canonical commutation relations, 11 pages. 1992.
- # 158: S. Izumiya, W.L. Marar, The Euler number of a topologically stable singular surface in a 3-manifold, 11 pages. 1992.
- # 159: T. Hibi, Cohen-Macaulay types of Cohen-Macaulay complexes, 26 pages. 1992.
- # 160: A. Arai, Properties of the Dirac-Weyl operator with a strongly singular gauge potential, 26 pages. 1992.
- # 161: A. Arai, Dirac operators in Boson-Fermion Fock spaces and supersymmetric quantum field theory, 30 pages. 1992.
- # 162: S. Albeverio, K. Iwata, T. Kolsrud, Random parallel transport on surfaces of finite type, and relations to homotopy, 8 pages. 1992.
- # 163: S. Albeverio, K. Iwata, T. Kolsrud, Moments of random fields over a family of elliptic curves, and modular forms, 9 pages. 1992.
- # 164: Y. Giga, M. Sato, Neumann problem for singular degenerate parabolic equations, 12 pages. 1992.
- # 165: J. Wierzbicki, Y. Watatani, Commuting squares and relative entropy for two subfactors, 18 pages. 1992.
- # 166: Y. Okabe, A new algorithm driven from the view-point of the fluctuation-dissipation theorem in the theory of KM_2O -Langevin equations, 13 pages. 1992.
- # 167: Y. Okabe, H. Mano and Y. Itoh, Random collision model for interacting populations of two species and its strong law of large numbers, 14 pages. 1992.
- # 168: A. Inoue, On the equations of stationary processes with divergent diffusion coefficients, 25 pages. 1992.
- # 169: T. Ozawa, Remarks on quadratic nonlinear Schrödinger equations, 19 pages. 1992.
- # 170: T. Fukui, Y. Giga, Motion of a graph by nonsmooth weighted curvature, 11 pages. 1992.
- # 171: J. Inoue, T. Nakazi, Finite dimensional solution sets of extremal problems in H^1 , 10 pages. 1992.
- # 172: S. Izumiya, A characterization of complete integrability for partial differential equations of first order, 6 pages. 1992.
- # 173: T. Suwa, Unfoldings of codimension one complex analytic foliation singularities, 49 pages. 1992.
- # 174: T. Ozawa, Wave propagation in even dimensional spaces, 15 pages. 1992.
- # 175: S. Izumiya, Systems of Clairaut type, 7 pages. 1992.
- # 176: A. Hoshiga, The initial value problems for quasi-linear wave equations in two space dimensions with small data, 25 pages. 1992.
- # 177: K. Sugano, On bicommutators of modules over H-separable extension rings III, 9 pages. 1993.
- # 178: T. Nakazi, Toeplitz operators and weighted norm inequalities, 17 pages. 1993.

EXISTENCE AND STRUCTURE OF INFINITELY DEGENERATE
ZERO-ENERGY GROUND STATES OF A WESS-ZUMINO TYPE
MODEL IN SUPERSYMMETRIC QUANTUM MECHANICS

OSAMU OGURISU

Department of Mathematics, Hokkaido University
Sapporo 060, Japan

1, October 1992

ABSTRACT

It is known that the $N=2$ Wess-Zumino supersymmetric quantum mechanical model with the superpotential $V(z) = \lambda e^{\alpha z}$ ($\lambda \in \mathbb{C} \setminus \{0\}$, $\alpha > 0$) has infinitely many bosonic zero-energy ground states and no fermionic zero-energy ground states [A. Arai, J. Math. Phys. **30** (1989), 1164]. In this paper, we extend these results to a more general model. The main results include the following: (1) identification of the spectra of the Hamiltonian H of the model; (2) non-Fredholmness of a supercharge of the model, which is a Dirac type operator; (3) existence of infinitely many bosonic zero-energy states of H ; (4) non-existence of fermionic zero-energy states of H .

PACS: 02.30.Jr, 02.30.Tb, 03.65.Db, 03.65.Ge

to appear in Journal of Mathematical Physics

I. INTRODUCTION

It is known that the $N=2$ Wess-Zumino (WZ) model in supersymmetric quantum mechanics,^{1,2} which describes the interaction between a complex bosonic degree $z \in \mathbb{C}$ and two fermionic degrees of freedom, has degenerate zero-energy ground states for some classes of superpotentials $V(z)$.³⁻⁶ The following (V1)–(V3) hold:

- (V1) If $V(z)$ is a polynomial of z , then the number of the bosonic zero-energy ground states n_b is equal to $\deg V - 1$.^{3,4};
- (V2) If $V(z) = \lambda z^{-p}$ ($\lambda \in \mathbb{C} \setminus \{0\}, p \in \mathbb{N}$), then $n_b = p - 1$.⁵;
- (V3) If $V(z) = \lambda e^{\alpha z}$ ($\lambda \in \mathbb{C} \setminus \{0\}, \alpha > 0$), then $n_b = \infty$.⁶

In all the cases, the number of the fermionic zero-energy states is equal to 0 (“vanishing theorem”).³⁻⁶

In the case where $V(z) = \lambda z^p$ ($\lambda \in \mathbb{C} \setminus \{0\}, p \in \mathbb{N}$), a special case of (V1), and in the case (V2), the mathematical structure of the degenerate ground states has been analyzed in detail.^{4,5} Both of these cases have some common structures:

- (i) There exists a strongly continuous one parameter unitary group acting as a symmetry group in the quantum system governed by the bosonic part H_+ of the Hamiltonian H of the model.
- (ii) The generator L of the symmetry group has purely discrete spectrum and infinitely many eigenspaces, each of which reduces H_+ .
- (iii) There exist some eigenspaces of L such that H_+ has a unique zero-energy state on each of them. (Hence the number of such eigenspaces is equal to the number of the bosonic zero-energy states.)
- (iv) H_+ has generalized zero-energy states on each of the eigenspaces of L .

However, in the case (V3), it has not been known whether the ground states have a structure similar to the cases (V1) and (V2). In this paper we present a new class of superpotentials containing the case (V3) such that each superpotential in the class defines a WZ type model with the following structures:

- (i') There exists a symmetry group of the bosonic Hamiltonian H_+ .
- (ii') The generator N_+ of the symmetry group has no point spectrum and every $p \in \mathbb{R}$ is a generalized eigenvalue of N_+ . Each generalized eigenspace \mathcal{H}_p of N_+ is regarded as

a Hilbert space and “reduces” H_+ .

- (iii') There exists an interval I in \mathbb{R} such that for every $p \in I$, H_+ has a unique zero-energy state on \mathcal{H}_p .
- (iv') There exist infinitely many zero-energy ground states of H_+ .
- (v') There exist infinitely many generalized zero-energy states of H_+ .

For another class of superpotentials containing the case (V3), we prove in a way similar to the case of H_+ the vanishing theorem and that there exist infinitely many generalized fermionic zero-energy states.

We also prove that the spectrum of H_{\pm} is equal to $[0, \infty)$, which implies that a supercharge of the model is not semi-Fredholm. This suggests the instability of $\ker H_-$. (See Remark 3.7)

The plan of this paper is as follows. In Sec. II, we introduce a class of superpotentials containing $\lambda e^{\alpha z}$ and define the WZ type model with a superpotential in the class. In Sec. III, we identify the spectra of the Hamiltonian of the model. In Sec. IV, we consider conditions for the superpotential under which there exist infinitely many zero-energy ground states of H_+ , and give a concrete class of superpotentials which satisfy these conditions. In Sec.V, we consider conditions for the superpotential under which there exists no fermionic zero-energy ground states. In Appendix we summarize a few known results in a theory of ordinary differential equation, which are used in the main text.

II. MODELS WITH SUPERPOTENTIALS SIMILAR TO $\exp(\alpha z)$

In this section we define a class of functions on \mathbb{C} and discuss the N=2 WZ type model with a superpotential V in the class. We identify $z \in \mathbb{C}$ with $x + iy$ where $x, y \in \mathbb{R}$. We denote $\partial/\partial z$ and $\partial/\partial \bar{z}$ by ∂ and $\bar{\partial}$, respectively.

Let α be an arbitrary positive constant.

Definition 2.1. A function $V: \mathbb{C} \rightarrow \mathbb{C}$ is said to be in \mathcal{V}_0 if there exists $v \in C^\infty(\mathbb{R})$ such that

$$\partial V(z) = v(x)e^{i\alpha y}.$$

The Hilbert space of state vectors of the N=2 WZ type model with a superpotential $V \in \mathcal{V}_0$ is realized as

$$\mathcal{H} = L^2(\mathbb{C}; \mathbb{C}^4) = L^2(\mathbb{R}^2; \mathbb{C}^4) = \mathcal{H}_+ \oplus \mathcal{H}_-,$$

where

$$\mathcal{H}_+ = \left\{ \begin{pmatrix} f \\ g \end{pmatrix} \middle| f, g \in L^2(\mathbb{C}) \right\} = L^2(\mathbb{C}; \mathbb{C}^2)$$

and

$$\mathcal{H}_- = \left\{ \begin{pmatrix} f \\ g \end{pmatrix} \middle| f, g \in L^2(\mathbb{C}) \right\} = L^2(\mathbb{C}; \mathbb{C}^2)$$

are the subspaces of ‘‘bosonic’’ and ‘‘fermionic’’ states, respectively.

Let $V \in \mathcal{V}_0$. One of the supercharges of the model with the superpotential V is given by

$$Q = \begin{pmatrix} 0 & Q_- \\ Q_+ & 0 \end{pmatrix},$$

acting in \mathcal{H} , where Q_- and Q_+ are operators defined by

$$Q_- = \begin{pmatrix} i\partial V & -\partial \\ -\bar{\partial} & -i(\partial V)^* \end{pmatrix} \quad \text{in } \mathcal{H}_-$$

and

$$Q_+ = \begin{pmatrix} -i(\partial V)^* & \partial \\ \bar{\partial} & i\partial V \end{pmatrix} \quad \text{in } \mathcal{H}_+.$$

We put $D = C_0^\infty(\mathbb{R}^2; \mathbb{C}^4)$. In the same way as in the proof of Proposition 2.1 in Ref. 6, we can prove the following proposition, because the holomorphic property of $V(z)$ is not used in the proof there.

Proposition 2.2. *The operator Q is essentially self-adjoint on D . Further, every power of Q is essentially self-adjoint on D .*

We denote the closure of Q by the same symbol Q and define the non-negative self-adjoint operator H by

$$H = Q^2 = \begin{pmatrix} H_+ & 0 \\ 0 & H_- \end{pmatrix}$$

with

$$H_+ = Q_+^* Q_+, \quad H_- = Q_- Q_-^*.$$

The operator H is the Hamiltonian of the model, called the *supersymmetric Hamiltonian*.

We call H_+ (resp. H_-) the *bosonic* (resp. *fermionic*) *Hamiltonian*. We have

$$H_+ = H_0 + \begin{pmatrix} 0 & -i\partial^2 V \\ i(\partial^2 V)^* & 0 \end{pmatrix}$$

and

$$H_- = H_0 + \begin{pmatrix} 0 & -i(\bar{\partial}\partial V)^* \\ i\bar{\partial}\partial V & 0 \end{pmatrix}$$

on D , where

$$H_0 = (-\partial\bar{\partial} + |\partial V|^2) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}.$$

III. SPECTRA OF THE HAMILTONIANS

We denote by $\sigma(A)$, $\sigma_p(A)$ and $D(A)$ the spectrum, the point spectrum and the domain of an operator A , respectively.

In this section we identify the spectra of the Hamiltonian of the models with potentials in a subclass of \mathcal{V}_0 . Moreover, we show that the Hamiltonians H_{\pm} and the supercharge Q are not semi-Fredholm.

We first prove these facts in the case where $\partial V(z) = \lambda e^{\alpha z}$. In this case we can identify the point spectrum of H_- , too.

Theorem 3.1. *Let $\partial V(z) = \lambda e^{\alpha z}$ with $\lambda \in \mathbb{C} \setminus \{0\}$. Then*

$$\sigma(H_{\pm}) = [0, \infty), \quad \sigma_p(H_-) = \emptyset.$$

Proof. In the present case, we have

$$H_- = -\partial\bar{\partial} + |\partial V|^2 = \begin{pmatrix} K & 0 \\ 0 & K \end{pmatrix},$$

where

$$K = -\frac{1}{4} \frac{\partial^2}{\partial x^2} - \frac{1}{4} \frac{\partial^2}{\partial y^2} + |\lambda|^2 e^{2\alpha x}.$$

It follows that $\sigma(H_-) = \sigma(K)$, $\sigma_p(H_-) = \sigma_p(K)$. Let H_x and H_y be the self-adjoint operators in $L^2(\mathbb{R})$ defined by

$$H_x = -\frac{1}{4} \frac{d^2}{dx^2} + |\lambda|^2 e^{2\alpha x}, \quad H_y = -\frac{1}{4} \frac{d^2}{dy^2}.$$

We denote by I the identity on $L^2(\mathbb{R})$. Then

$$K = H_x \otimes I + I \otimes H_y$$

as an operator on $L^2(\mathbb{R}^2) = L^2(\mathbb{R}) \otimes L^2(\mathbb{R})$. It is well known that, for self-adjoint operators A_1 and A_2 ,

$$\begin{aligned} \sigma(A_1 \otimes I + I \otimes A_2) &= \overline{\{\mu = \mu_1 + \mu_2 \mid \mu_j \in \sigma(A_j), j = 1, 2\}}, \\ \sigma_p(A_1 \otimes I + I \otimes A_2) &= \overline{\{\mu = \mu_1 + \mu_2 \mid \mu_j \in \sigma_p(A_j), j = 1, 2\}} \end{aligned}$$

and that $\sigma(H_y) = [0, \infty)$ and $\sigma_p(H_y) = \emptyset$. Hence we see that $\sigma_p(K) = \emptyset$. Thus, $\sigma_p(H_-) = \emptyset$.

We next prove that $\sigma(H_x) \supset [0, \infty)$.

We use the following basic criterion: Let A be a self-adjoint operator. Then $\nu \in \sigma(A)$ if and only if there exist a sequence $\{\phi_n\}_{n \in \mathbb{N}} \subset D(A)$ such that $\|\phi_n\| = 1$ and $\|(A - \nu)\phi_n\| \rightarrow 0$ as $n \rightarrow \infty$. Let $\mu \geq 0$. We prove that $\mu^2 \in \sigma(H_x)$. Let $c = \sqrt{\pi}$ and for each $n \in \mathbb{N}$

$$\phi_n(x) = \frac{1}{\sqrt{nc}} e^{2i\mu x} e^{-x^2/2n^2}.$$

Then $\phi_n \in D(d^2/dx^2) \cap D(e^{2\alpha x})$ and $\|\phi_n\| = 1$. Let T_n be a shift operator on $L^2(\mathbb{R})$ defined by

$$T_n f(x) = f(x + n^3).$$

Then the following hold:

- (1) $\|T_n f\| = \|f\|$ for $f \in L^2(\mathbb{R})$,
- (2) $\|e^{2\alpha x} T_n f\| = e^{-2\alpha n^3} \|e^{2\alpha x} f\|$ for $f \in D(e^{2\alpha x})$,
- (3) $T_n \left(-\frac{1}{4} \frac{d^2}{dx^2} - \mu^2 \right) = \left(-\frac{1}{4} \frac{d^2}{dx^2} - \mu^2 \right) T_n$ on $D\left(\frac{d^2}{dx^2}\right)$.

Let $f_n = T_n \phi_n$. Then $\|f_n\| = 1$ and $f_n \in D(d^2/dx^2) \cap D(e^{2\alpha x})$. We have

$$\|(H_x - \mu^2)f_n\| \leq \left\| \left(-\frac{1}{4} \frac{d^2}{dx^2} - \mu^2 \right) \phi_n \right\| + |\lambda|^2 e^{-2\alpha n^3} \|e^{2\alpha x} \phi_n\|. \quad (3.1)$$

We have

$$\left(-\frac{1}{4} \frac{d^2}{dx^2} - \mu^2 \right) \phi_n = \frac{1}{4n^2} \left(-\frac{x^2}{n^2} + 4i\mu x + 1 \right) \phi_n.$$

Hence

$$\left\| \left(-\frac{1}{4} \frac{d^2}{dx^2} - \mu^2 \right) \phi_n \right\| \leq \frac{1}{4n^2} \left\{ \frac{1}{n^2} \|x^2 \phi_n\| + 4\mu \|x \phi_n\| + \|\phi_n\| \right\}.$$

We can easily verify that

$$\|x^2 \phi_n\| = n^2 \sqrt{\frac{1}{c} \int_{\mathbb{R}} y^4 e^{-y^2} dy} \quad \text{and} \quad \|x \phi_n\| = n \sqrt{\frac{1}{c} \int_{\mathbb{R}} y^2 e^{-y^2} dy}$$

for all $n \in \mathbb{N}$. Therefore, the first term on the right hand side of (3.1) converges to 0 as $n \rightarrow \infty$. We have $\|e^{2\alpha x} \phi_n\| = \exp(2\alpha^2 n^2)$. Thus, since α is positive, the second term on the right hand side of (3.1) converges to 0 as $n \rightarrow \infty$. Hence

$$\|(H_x - \mu^2)f_n\| \rightarrow 0 \quad \text{as} \quad n \rightarrow \infty.$$

Hence $\mu^2 \in \sigma(H_x)$. Therefore we obtain that $[0, \infty) \subset \sigma(H_x)$. Since H_x is non-negative self-adjoint, it follows that $\sigma(H_x) = [0, \infty)$, which implies $\sigma(K) = [0, \infty)$. Hence $\sigma(H_-) = [0, \infty)$.

In general, for any densely defined closed linear operator A from a Hilbert space to another one, $\sigma(A^*A) \setminus \{0\} = \sigma(AA^*) \setminus \{0\}$. (See Ref.7) Hence $\sigma(H_+) \setminus \{0\} = \sigma(H_-) \setminus \{0\}$. It follows that $\sigma(H_+) = [0, \infty)$. \square

Corollary 3.2. *Let $\partial V(z) = \lambda e^{\alpha z}$ with $\lambda \in \mathbb{C} \setminus \{0\}$. Then H_{\pm} is not semi-Fredholm.*

Proof. It is known that a densely defined closed linear operator A is semi-Fredholm if and only if (1) $\inf \sigma(A^*A) \setminus \{0\} > 0$ and (2) at least one of $\dim \ker A^*A$ and $\dim \ker AA^*$ is finite. By Theorem 3.1, it follows that $\sigma(H_{\pm}^*H_{\pm}) = \sigma(H_{\pm}^2) = [0, \infty)$, which implies that H_{\pm} is not semi-Fredholm. \square

We introduce a subclass \mathcal{V}_1 of \mathcal{V}_0 .

Definition 3.3. We say that V is in \mathcal{V}_1 if $V \in \mathcal{V}_0$ and there exist constants c and $\beta > 0$ such that

$$|v^{(j)}(x)| \leq ce^{\beta x}, \quad x \in \mathbb{R}, \quad j = 0, 1$$

where $v^{(j)}(x) = d^j v(x)/dx^j$.

For this subclass we obtain the following results.

Theorem 3.4. *Suppose $V \in \mathcal{V}_1$. Then*

$$\sigma(H_{\pm}) = [0, \infty).$$

Proof. Let

$$H_0 = -\frac{1}{4} \frac{d^2}{dx^2} \otimes I - I \otimes \frac{1}{4} \frac{d^2}{dy^2} + |v(x)|^2 \otimes I \quad \text{in } L^2(\mathbb{R}^2).$$

We can write

$$H_- = \begin{pmatrix} H_0 & 0 \\ 0 & H_0 \end{pmatrix} + \frac{i}{2} \begin{pmatrix} 0 & -\overline{(v' - \alpha v)} \otimes e^{-i\alpha y} \\ (v' - \alpha v) \otimes e^{i\alpha y} & 0 \end{pmatrix}.$$

Let $\mu \geq 0$. We prove that $2\mu^2 \in \sigma(H_-)$. Let ϕ_n and f_n be the functions in the proof of Theorem 3.1. Then,

$$\left\| \left(-\frac{1}{4} \frac{d^2}{dx^2} - \mu^2 \right) \phi_n \right\| \rightarrow 0, \quad \left\| \left(-\frac{1}{4} \frac{d^2}{dx^2} - \mu^2 \right) f_n \right\| \rightarrow 0 \quad \text{as } n \rightarrow \infty.$$

Moreover,

$$\begin{aligned} \| |v|^2 f_n \| &\leq c^2 \| e^{2\beta x} f_n \| \rightarrow 0 \quad \text{as } n \rightarrow \infty, \\ \| (v' - \alpha v) f_n \| \quad \text{and} \quad \| \overline{(v' - \alpha v)} f_n \| &\leq (c + \alpha c) \| e^{\beta x} f_n \| \rightarrow 0 \quad \text{as } n \rightarrow \infty. \end{aligned}$$

Let

$$\Phi_n = \frac{1}{\sqrt{2}} \begin{pmatrix} f_n \otimes \phi_n \\ f_n \otimes \phi_n \end{pmatrix}.$$

Then $\|\Phi_n\| = 1$. We have

$$\begin{aligned} \|(H_- - 2\mu^2)\Phi_n\|^2 &= \left\| \left(H_0 - 2\mu^2 - \frac{i}{2} \overline{(v' - \alpha v)} \otimes e^{-i\alpha y} \right) f_n \otimes \phi_n \right\|^2 \\ &\quad + \left\| \left(H_0 - 2\mu^2 + \frac{i}{2} (v' - \alpha v) \otimes e^{i\alpha y} \right) f_n \otimes \phi_n \right\|^2. \end{aligned} \quad (3.2)$$

We have

$$\begin{aligned} &\left\| \left(H_0 - 2\mu^2 - \frac{i}{2} \overline{(v' + \alpha v)} \otimes e^{-i\alpha y} \right) f_n \otimes \phi_n \right\| \\ &\leq \left\| \left(-\frac{1}{4} \frac{d^2}{dx^2} - \mu^2 \right) f_n \right\| + \| |v|^2 f_n \| + \left\| \left(-\frac{1}{4} \frac{d^2}{dy^2} - \mu^2 \right) \phi_n \right\| + \frac{1}{2} \| \overline{(v' - \alpha v)} f_n \| \\ &\rightarrow 0 \quad \text{as } n \rightarrow \infty. \end{aligned}$$

In the same way, the second term on the right hand side of (3.2) converges to 0 as $n \rightarrow \infty$.

Hence

$$\|(H_- - 2\mu^2)\Phi_n\| \rightarrow 0 \quad \text{as } n \rightarrow \infty.$$

Therefore, for every $\mu \geq 0$, $2\mu^2 \in \sigma(H_-)$. Since H_- is non-negative self-adjoint, $\sigma(H_-) = [0, \infty)$. It follows in the same way as in the proof of Theorem 3.1 that $\sigma(H_+) = [0, \infty)$. \square

Corollary 3.5. *Suppose $V \in \mathcal{V}_1$. Then H_{\pm} is not semi-Fredholm.*

Proof. We can prove this corollary in the same way as in the proof of Corollary 3.2. \square

Corollary 3.6. *The operator Q_+ is not semi-Fredholm.*

Proof. We have $\sigma(Q_+^* Q_+) = \sigma(H_+) = [0, \infty)$. Hence $\inf \sigma(Q_+^* Q_+) \setminus \{0\} = 0$. It implies that Q_+ is not semi-Fredholm. \square

Remark 3.7. We have $\text{index} Q_+ = \dim \ker H_+ - \dim \ker H_-$. By Corollary 3.6, $\text{index} Q_+$ may be unstable under perturbations. This suggests that $\ker H_-$ may be also unstable under perturbations.

IV. ZERO-ENERGY GROUND STATES OF THE BOSONIC HAMILTONIAN

In this section we discuss zero-energy ground states of the bosonic Hamiltonian H_+ and the structure of them, and give a concrete class of superpotentials with which the models have infinitely many bosonic zero-energy ground states.

A. Symmetry Group and Direct Integral Decomposition.

Let $\partial/\partial y$ be the generalized derivative and put

$$N_+ = i \frac{\partial}{\partial y} + \frac{\alpha}{2} \sigma_3$$

with $D(N_+) = D(\partial/\partial y) \subset \mathcal{H}_+$, where $\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$. Then N_+ is self-adjoint. Hence, by Stone's theorem it generates the strongly continuous one parameter unitary group $\{e^{isN_+} \mid s \in \mathbb{R}\}$ on \mathcal{H}_+ .

Lemma 4.1. *Let $V \in \mathcal{V}_0$. Then, for all $t, s \in \mathbb{R}$,*

$$e^{itH_+} e^{isN_+} = e^{isN_+} e^{itH_+}.$$

Proof. We can prove this lemma in the same way as in the proof of Lemma 3.1 in Ref. 6. \square

Remark. In the case $V(z) = \lambda z^p = \lambda r^p e^{ip\theta}$ (polar coordinate; $z = re^{i\theta}$), the generator of a symmetry group of H_+ is given by $L = i\partial/\partial\theta + (p-2)\sigma_3/2$. In the case $V(z) = \lambda/z^p = \lambda r^p e^{-ip\theta}$, the operator $M = i\partial/\partial\theta - (p+2)\sigma_3/2$ is the generator of a symmetry group of H_+ . Both of the spectrum of L and of M are equal to $\{n - p/2 \mid n \in \mathbb{N}\}$. (See Refs.4 and 5.)

The generator N_+ has no point spectrum and the spectrum of N_+ is equal to \mathbb{R} . This is different from the case of L and M . However every $p \in \mathbb{R}$ is a generalized eigenvalue of N_+ and the generalized eigenspace with eigenvalue p is given by

$$\mathcal{H}_p = \left\{ \begin{pmatrix} \phi(x)e^{-i(p-\alpha/2)y} \\ \psi(x)e^{-i(p+\alpha/2)y} \end{pmatrix} \mid \phi, \psi \in L^2(\mathbb{R}) \right\} \cong \oplus^2 L^2(\mathbb{R}).$$

This suggests that the bosonic Hilbert space \mathcal{H}_+ may be written as the constant fiber direct integral over \mathbb{R} with fiber $\oplus^2 L^2(\mathbb{R})$ and H_+ may be decomposable with respect to the direct integral decomposition. In fact, this is true as we show in the following.

The operator $U: L^2(\mathbb{R}^2; \mathbb{C}^2) \rightarrow L^2(\mathbb{R}; \oplus^2 L^2(\mathbb{R}))$ given by

$$U \begin{pmatrix} f \\ g \end{pmatrix} (p) = \begin{pmatrix} (2\pi)^{-1/2} \int_{\mathbb{R}} e^{-ipy} f(y, x) dy \\ (2\pi)^{-1/2} \int_{\mathbb{R}} e^{-i(p-\alpha)y} g(y, x) dy \end{pmatrix}, \quad \begin{pmatrix} f \\ g \end{pmatrix} \in L^2(\mathbb{R}^2; \mathbb{C}^2),$$

is unitary. Then this operator gives a direct integral decomposition of $L^2(\mathbb{R}; \oplus^2 L^2(\mathbb{R}))$. We denote by f' the derivative of f with respect to x .

Lemma 4.2. *Let $V \in \mathcal{V}_0$ with $\partial V(z) = v(x)e^{i\alpha y}$. Let*

$$h(p) = \begin{pmatrix} -\frac{1}{4} \frac{d^2}{dx^2} + |v|^2 + \frac{1}{4} p^2 & -\frac{i}{2} (v' + \alpha v) \\ \frac{i}{2} \overline{(v' + \alpha v)} & -\frac{1}{4} \frac{d^2}{dx^2} + |v|^2 + \frac{1}{4} (p - \alpha)^2 \end{pmatrix}$$

acting in $\oplus^2 L^2(\mathbb{R})$. Then

$$UH_+U^* = \int_{\mathbb{R}}^{\oplus} h(p) dp. \quad (4.1)$$

Proof. Let $\mathcal{D} = C_0^\infty(\mathbb{R}; \mathbb{C}^2)$ and A_p be an operator given by

$$A_p = \begin{pmatrix} -i\bar{v} & \frac{1}{2} \left(\frac{d}{dx} + p - \alpha \right) \\ \frac{1}{2} \left(\frac{d}{dx} - p \right) & iv \end{pmatrix}$$

with $D(A_p) = \mathcal{D}$. The operator

$$T = \begin{pmatrix} 0 & A_p \\ A_p^* & 0 \end{pmatrix}$$

can be written

$$T = \frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix} \frac{d}{dx} + B$$

with $D(T) = \mathcal{D} \oplus \mathcal{D}$ and B being a C^∞ 4×4 matrix-valued function. It is obvious that T is symmetric on $D(T)$. By Theorem 2.2 in Ref.8,

$$T^2 = \begin{pmatrix} A_p A_p^* & 0 \\ 0 & A_p^* A_p \end{pmatrix}$$

is essentially self-adjoint on $\mathcal{D} \oplus \mathcal{D}$. Hence $A_p^* A_p$ is essentially self-adjoint on \mathcal{D} . Since $h(p) = A_p^* A_p$ on \mathcal{D} , $h(p)$ is essentially self-adjoint on \mathcal{D} . The operator R given by

$$R = \int_{\mathbb{R}}^{\oplus} \{h(p) + i\}^{-1} dp$$

is bounded and decomposable on $L^2(\mathbb{R}; \oplus^2 L^2(\mathbb{R}))$. Moreover we have $R = (UH_+U^* + i)^{-1}$. Hence UH_+U^* is decomposable and has the representation (4.1). \square

The following lemma will be needed.

Lemma 4.3. *Let $\langle M, \mu \rangle$ be a σ -finite measure space. Let \mathcal{H} and \mathcal{K} be separable Hilbert spaces and $C(\mathcal{H}, \mathcal{K})$ be the set of densely defined closed linear operators from \mathcal{H} to \mathcal{K} . Let $A(m)$ be a $C(\mathcal{H}, \mathcal{K})$ -valued, measurable function on M . Let $H(m) = A(m)^*A(m)$, $m \in M$, and*

$$H = \int_M^\oplus H(m) d\mu$$

which acts in $\int_M^\oplus \mathcal{H} d\mu$, the constant fiber direct integral of \mathcal{H} . Then

$$\ker H = \int_M^\oplus \ker A(m) d\mu.$$

Proof. In general, $\ker A^*A = \ker A$. By the general theory of constant fiber direct integral, $\ker H = \int_M^\oplus \ker H(m) d\mu$. Hence we obtain the desired result. \square

The identity $\ker A^*A = \ker A$ implies that $\ker(UH_+U^*) = \ker(UQ_+U^*)$.

Lemma 4.4. *Let $V \in \mathcal{V}_0$ with $\partial V(z) = v(x)e^{i\alpha y}$. Let $\Phi \in \ker(UQ_+U^*)$. Write $\Phi(p) = (\phi(p, \cdot), -i\psi(p, \cdot))$. Then the following equations hold a.e. (almost everywhere) p in the generalized sense:*

$$\overline{v(\cdot)}\phi(p, \cdot) + \frac{1}{2} \left(\frac{\partial}{\partial x} + p - \alpha \right) \psi(p, \cdot) = 0, \quad (4.2)$$

$$\frac{1}{2} \left(\frac{\partial}{\partial x} - p \right) \phi(p, \cdot) + v(\cdot)\psi(p, \cdot) = 0. \quad (4.3)$$

Proof. By direct computations. \square

We fix p to solve the equation $UQ_+U^*\Phi = 0$ on each of fiber spaces. We rewrite $(\phi(p, \cdot), \psi(p, \cdot))$ as $((\phi(\cdot), \psi(\cdot)))$.

Lemma 4.5. *Let $V \in \mathcal{V}_0$ with $\partial V(z) = v(x)e^{i\alpha y}$. Suppose $v'/v \in C^\infty(\mathbb{R})$.*

(i) Every solution (ϕ, ψ) of (4.2) and (4.3) is a C^∞ -function and satisfies the following ordinary differential equations:

$$-\phi'' + \left(\frac{v'}{v} + \alpha\right)\phi' + \left(4|v|^2 - p\frac{v'}{v} + p(p - \alpha)\right)\phi = 0, \quad (4.4)$$

$$-\psi'' + \left(\frac{\bar{v}'}{\bar{v}} + \alpha\right)\psi' + \left(4|v|^2 + (p - \alpha)\frac{\bar{v}'}{\bar{v}} + p(p - \alpha)\right)\psi = 0, \quad (4.5)$$

$$\frac{1}{2} \left(\frac{d}{dx} - p\right)\phi + v\psi = 0. \quad (4.6)$$

(ii) For every solution ϕ of (4.4), the function ψ satisfying (4.6) is a solution of (4.5).

Proof. Since v and v'/v are in $C^\infty(\mathbb{R})$, it follows from elliptic regularity that ϕ and ψ are in $C^\infty(\mathbb{R})$. Therefore we have the desired results. \square

B. Zero-energy states in each of fiber spaces.

We introduce a subclass of superpotentials in \mathcal{V}_0 .

Definition 4.6. We say that V is in \mathcal{V}_2 if there exist a constant $\lambda \in \mathbb{C} \setminus \{0\}$ and $v \in C^\infty(\mathbb{R})$ such that $v > 0$, $\partial V(z) = \lambda v(x)e^{i\alpha y}$ and

$$\lim_{x \rightarrow \infty} e^{\alpha x} v(x) > 0. \quad (4.7)$$

Let $V \in \mathcal{V}_2$. It is well known that the equations (4.4) and (4.5) can be transformed into the equations of the form

$$y'' - Gy = 0. \quad (4.8)$$

Let $h(x) = \sqrt{e^{\alpha x} v(x)}$ and define $f(x) = \phi(x)/h(x)$ and $g(x) = \psi(x)/h(x)$. Then we have from (4.4) and (4.5)

$$f'' - G_p f = 0, \quad (4.9)$$

$$g'' - G_{\alpha-p} g = 0, \quad (4.10)$$

where

$$G_p(x) = 4|\lambda|^2 v^2 + \frac{1}{4} \left(\frac{v'}{v} - 2p + \alpha\right)^2 + \frac{1}{2} \left(\frac{v'}{v}\right)^2 - \frac{1}{2} \frac{v''}{v}.$$

We next investigate properties of solutions of the equation (4.8).

Lemma 4.7. Suppose that $G \in C(\mathbb{R})$ and there exist constants $c_-, c_+ \in \mathbb{R}$, positive constants δ, γ, σ and $\rho > 1/2$ such that

- (i) $\gamma^2 \leq G(x) \leq \delta^2$ on $(-\infty, c_-]$ and $\lim_{x \rightarrow -\infty} h(x) e^{(2\delta-\gamma)|x|} |x|^\rho = 0$,
- (ii) $\sigma^2 \leq G(x)$ on $[c_+, \infty)$ and $\lim_{x \rightarrow \infty} h(x) e^{-\sigma x} x^\rho = 0$.

Then there exists a unique solution y of (4.8), up to constant multiples, such that $hy \in L^2(\mathbb{R})$.

Proof. By Lemma A.2 (see Appendix), there exists one and only one solution y_+ of (4.8) on $(-\infty, c_-]$ such that

$$y_+(c_-) = 1 \quad \text{and} \quad y_+(x) \rightarrow 0 \quad \text{as} \quad x \rightarrow -\infty.$$

Moreover, by Lemma A.4 the following estimate holds:

$$e^{-\delta|x-c_-|} \leq y_+(x) \leq e^{-\gamma|x-c_-|} \quad \text{for} \quad x \leq c_-. \quad (4.11)$$

Let

$$z(x) = y_+(x) \int_{c_-}^x y_+(t)^{-2} dt \quad \text{for} \quad x \leq c_-.$$

Then z is a solution of (4.8) and linearly independent of y_+ . (4.11) gives the following estimate:

$$|z(x)| \leq y_+(x) \left| \int_{c_-}^x y_+(t)^{-2} dt \right| \leq \text{const.} e^{(2\delta-\gamma)|x|}.$$

Since all the solutions y of (4.8) are linear combinations of y_+ and z , it follows that $hy \in L^2(-\infty, c_-)$.

By Lemma A.4 there exists a unique subdominant solution y_+ of (4.8), up to constant multiples, on $[c_+, \infty)$ such that

$$0 < y_+(x) < \text{const.} e^{-\sigma x} \quad \text{for} \quad x \geq c_+.$$

Obviously we have $hy_+ \in L^2(c_+, \infty)$. On the other hand, by Lemma A.3 any solution z of (4.8) linearly independent of y_+ tends to infinity as $x \rightarrow \infty$. Since $\lim_{x \rightarrow \infty} h(x) > 0$ by (4.7), it follows that $hz \notin L^2(c_+, \infty)$. Consequently, there exists a unique solution y of (4.8) on $[c_+, \infty)$ such that $hy \in L^2(c_+, \infty)$. Since every solution of (4.8) is continuous on $[c_-, c_+]$, it follows that there exists a unique solution y of (4.8) such that $hy \in L^2(\mathbb{R})$. \square

We are now ready to solve the equations (4.4)–(4.6).

Lemma 4.8. Let $V \in \mathcal{V}_2$. Suppose that the functions G_p and $G_{\alpha-p}$ satisfy the hypothesis in Lemma 4.7 with $G = G_p$, $G = G_{\alpha-p}$. Suppose that $\overline{\lim}_{x \rightarrow \infty} e^{\alpha x} v^{-1}(x) < \infty$ and $\overline{\lim}_{x \rightarrow \infty} |v'(x)v^{-2}(x)| < \infty$. Then there exists a unique solution of (4.4)–(4.6) in $\oplus^2 L^2(\mathbb{R})$, up to constant multiples.

Proof. By Lemma 4.7 we can construct a unique solution $\phi = hf$ of (4.4) in $L^2(\mathbb{R})$. We define ψ by (4.6) with this ϕ . Then by part (ii) of Lemma 4.5, ψ is a solution of (4.5). We can prove in the same way as in the proof of Lemma 4.7 that ψ is in $L^2(-\infty, c_-)$. Since $\phi = hf$ and $h = \sqrt{e^{\alpha x} v}$,

$$-2\lambda\psi = \frac{1}{v} \left(\frac{d}{dx} - p \right) \phi = \frac{1}{2} \left[\frac{\alpha - 2p}{v} + \frac{v'}{v^2} \right] hf + \sqrt{\frac{e^{\alpha x}}{v}} f' \quad (4.12)$$

holds. By the construction of ϕ , the function f is a subdominant solution of (4.9). Hence, by Lemma A.2, $f' \in L^2(c_+, \infty)$ and by Lemma A.4, $f \in L^2(c_+, \infty)$. Hence it follows that all the terms on the right hand side of (4.12) are in $L^2(c_+, \infty)$. Consequently, $\psi \in L^2(\mathbb{R})$. By Lemma 4.7, this ψ is a unique solution of (4.5) in $L^2(\mathbb{R})$. Hence the pair (ϕ, ψ) is a unique solution of (4.4)–(4.6) in $\oplus^2 L^2(\mathbb{R})$. \square

C. Zero-energy states in the whole space.

In the last subsection we got a zero-eigenvalue solution of $UH_+U^*\Phi = 0$ restricted to each fiber space. By bundling them, we shall construct solutions of $UH_+U^*\Phi = 0$. But it is not obvious whether the bundled function is Lebesgue measurable on \mathbb{R} or not. In this subsection we first examine this measurability and then prove that $\dim \ker H_+ = \infty$. The following is the main theorem in this section.

Theorem 4.9. Let $V \in \mathcal{V}_2$. Suppose that there exists an interval I of \mathbb{R} with non-zero Lebesgue measure, satisfying the following properties:

- (i) there exist $c_-, c_+, \delta, \gamma, \sigma$ and ρ such that for all $p \in I$, $G = G_p$ satisfies the hypothesis in Lemma 4.7;
- (ii) there exist $c_-, c_+, \delta, \gamma, \sigma$ and ρ such that for all $p \in I$, $G = G_{\alpha-p}$ satisfies the hypothesis in Lemma 4.7;
- (iii) $\overline{\lim}_{x \rightarrow \infty} e^{\alpha x} v^{-1}(x) < \infty$ and $\overline{\lim}_{x \rightarrow \infty} |v'(x)v^{-2}(x)| < \infty$.

Then

$$\dim \ker H_+ = \infty.$$

The proof of this theorem is based on the following proposition.

Proposition 4.10. *Suppose that hypothesis (i) for G_p in Theorem 4.9 is satisfied. Then there exists a solution y_p of*

$$y'' - G_p y = 0 \quad \text{on } \mathbf{I} \times \mathbb{R} \quad (4.13)$$

such that $h(\cdot)y_p(\cdot) \in L^2(\mathbb{R})$ for every $p \in \mathbf{I}$ and $hy_p: \mathbf{I} \rightarrow L^2(\mathbb{R})$, $p \mapsto h(\cdot)y_p(\cdot)$ is measurable. Moreover if the measure of \mathbf{I} is finite, then $hy_p \in L^2(\mathbf{I}; L^2(\mathbb{R}))$.

Remark. The function $hy_p: \mathbf{I} \rightarrow L^2(\mathbb{R})$ is measurable if and only if for all $g \in L^2(\mathbb{R})$ the mapping $\mathbf{I} \ni p \mapsto \langle hy_p, g \rangle_{L^2(\mathbb{R})} \in \mathbb{C}$ is measurable. (See Ref. 9, vol. I.)

Proof. Without loss of generality we can assume \mathbf{I} is a bounded closed interval.

Step 1. We prove the continuity of the functions y_p and y'_p on $\mathbf{I} \times \mathbb{R}$.

We assume $y_p(c_+) = 1$ for every $p \in \mathbf{I}$. For each $p \in \mathbf{I}$, let $y_{p,1}$ and $y_{p,0}$ be the solutions of (4.13) with $y_{p,1}(c_+) = 1$, $y'_{p,1}(c_+) = 0$, $y_{p,0}(c_+) = 0$ and $y'_{p,0}(c_+) = 1$. We put

$$\alpha(p) = - \lim_{x \rightarrow \infty} \frac{y_{p,1}(x)}{y_{p,0}(x)}.$$

From the proof of Lemma 4.7, y_p must be the normalized subdominant solution of (4.13) on (c_+, ∞) . By Lemma A.2, we can write

$$y_p(x) = y_{p,1}(x) + \alpha(p) y_{p,0}(x).$$

The function $G_p(x)$ is continuous in $x \in \mathbb{R}$ and we have for all $p, q \in \mathbb{R}$

$$G_p(x) - G_q(x) = (q - p) \left(\frac{v'(x)}{v(x)} + \alpha - p - q \right).$$

Since v'/v is continuous and \mathbf{I} is bounded, G_p satisfies the Lipschitz condition with respect to p on every compact interval of \mathbb{R} . Hence, by a general theory of ordinary differential equation, the solutions $y_{p,1}$, $y'_{p,1}$, $y_{p,0}$ and $y'_{p,0}$ are continuous in p at each $x \in \mathbb{R}$. (See e.g. Ref. 10.)

Claim. *Let $c \in (0, \infty]$. We consider a family of C^1 -functions such that*

$$\{f_\alpha: (0, c) \rightarrow \mathbb{R} \mid f'_\alpha \geq 0, \exists \lim_{x \rightarrow c} f_\alpha(x)\}_{\alpha \in A}.$$

If there exists a function $F: (0, c) \rightarrow \mathbb{R}$ such that $F(x) \geq f'_\alpha(x)$ on $(0, c)$ for every $\alpha \in A$ and $\int_0^c F(t)dt < \infty$, then the convergence, $\lim_{x \rightarrow c} f_\alpha(x) = f_\alpha(c)$, is uniform in $\alpha \in A$.

Proof. For every $\alpha \in A$, we have

$$0 \leq f_\alpha(c) - f_\alpha(t) = \int_t^c f'_\alpha(s)ds \leq \int_t^c F(s)ds.$$

As $t \rightarrow c$, the right hand side of this inequality converges to 0 uniformly in α . Hence the desired result follows. \square

Since $y_{p,0} > 0$ by Lemma A.1, the function $y_{p,1}/y_{p,0}$ is differentiable on (c_+, ∞) . We have

$$\frac{d}{dx} \left(\frac{y_{p,1}}{y_{p,0}} \right) = \frac{W(y_{p,1}, y_{p,0})}{(y_{p,0})^2} = \left(\frac{1}{y_{p,0}} \right)^2 > 0,$$

where W is Wronskian. For the solution Y of $y'' - \sigma^2 y = 0$ with $Y(c_+) = 0$ and $Y'(c_+) = 1$, Lemma A.4 implies that $0 < Y(x) \leq y_{p,0}$ on (c_+, ∞) for every $p \in \mathbf{I}$. Hence

$$\left(\frac{1}{Y(x)} \right)^2 \geq \frac{d}{dx} \left(\frac{y_{p,1}}{y_{p,0}} \right).$$

Since $Y(x) = \sinh(\sigma(x - c_+))/\sigma$, we have

$$\int_{c_++1}^{\infty} \left(\frac{1}{Y(x)} \right)^2 dx < \infty.$$

By the above claim, $-y_{p,1}(x)/y_{p,0}(x)$ converges to $\alpha(p)$ as $x \rightarrow \infty$ uniformly in $p \in \mathbf{I}$. Since $y_{p,1}(x)/y_{p,0}(x)$ is continuous in p at each $x \in \mathbb{R}$, $\alpha(p)$ is continuous. Consequently, y_p and y'_p are continuous in p and x . This fact will be used also in the proof of Theorem 4.9.

Step 2. We prove the measurability of hy_p . To do this, we construct a function F on \mathbb{R} such that $hF \in L^2(\mathbb{R})$ and $|y_p| \leq |F|$ on \mathbb{R} for every $p \in \mathbf{I}$. Then, it follows from Step 1 and Lebesgue dominate convergence theorem that for all $g \in L^2(\mathbb{R})$, $\langle hy_p, g \rangle$ is continuous in p .

On the interval (c_+, ∞) , y_p is the normalized subdominant solution of (4.13) for each p . Hence we define F on (c_+, ∞) to be the normalized subdominant solution of $y'' - \sigma^2 y = 0$. Then, by Lemma A.4, we have

$$0 \leq y_p(x) \leq F(x) \leq e^{-\sigma(x-c_+)} \quad \text{on } [c_+, \infty).$$

It is clear that $hF \in L^2(c_+, \infty)$.

Let $M = \sup\{|y_p(x)|; p \in \mathbf{I}, c_- \leq x \leq c_+\}$ and $N = \sup\{|y_p(c_-)|; p \in \mathbf{I}\}$. Since y_p and y'_p are continuous in p and x and \mathbf{I} is bounded, M and N are finite.

We put $F(x) = M$ on (c_-, c_+) . Then $hF \in L^2(c_-, c_+)$ and $|y_p| \leq F$.

For each $p \in \mathbf{I}$, there exists the normalized subdominant solution z_p of (4.13) on $(-\infty, c_-)$. The solution w_p of (4.13) given by

$$w_p(x) = z_p(x) \int_{c_-}^x z_p(t)^{-2} dt$$

is linearly independent of z_p . Then $w_p(c_-) = 0$ and $w'_p(c_-) = z_p(c_-)^{-1} = 1$. Hence we have

$$y_p(x) = y_p(c_-)z_p(x) + \{y'_p(c_-) - y_p(c_-)z'_p(c_-)\}w_p(x).$$

Lemma A.4 implies

$$e^{\delta(x-c_-)} \leq z_p(x) \leq e^{\gamma(x-c_-)}.$$

We obtain

$$\begin{aligned} |w_p(x)| &\leq e^{\gamma(x-c_-)} \left| \int_{c_-}^x e^{-2\delta(t-c_-)} dt \right| \\ &= \frac{1}{2\delta} e^{(2\delta-\gamma)c_-} e^{\gamma x} (e^{-2\delta x} - e^{-2\delta c_-}). \end{aligned}$$

These estimates do not depend on p . In the same way as in Step 1, we see that $z'_p(c_-)$ is continuous in p . Hence, the quantity $L = \sup\{|z'_p(c_-)|; p \in \mathbf{I}\}$ is finite. Consequently, we have for every $p \in \mathbf{I}$

$$|y_p(x)| \leq M e^{\gamma(x-c_-)} + (N + ML) \frac{1}{2\delta} e^{(2\delta-\gamma)c_-} e^{\gamma x} (e^{-2\delta x} - e^{-2\delta c_-}).$$

Hence there exists a constant C not depending on p such that

$$|y_p(x)| \leq C e^{(\gamma-2\delta)x} \quad \text{on } (-\infty, c_-].$$

We put $F(x) = C e^{(\gamma-2\delta)x}$ on $(-\infty, c_-]$. By the hypothesis of this proposition, $hF \in L^2(-\infty, c_-)$.

In this way we obtain a function F on \mathbb{R} such that $hF \in L^2(\mathbb{R})$ and $|y_p| \leq |F|$. Hence, the function $\mathbf{I} \ni p \mapsto h y_p \in L^2(\mathbb{R})$ is measurable. Since \mathbf{I} is bounded, $hF \in L^2(\mathbf{I}; L^2(\mathbb{R}))$. Hence $h y_p \in L^2(\mathbf{I}; L^2(\mathbb{R}))$. \square

Remark 4.11. It is easy to see that it is sufficient to assume, in Step 1, that $y_p(c_+)$ is continuous in p . Then we have $y_p(x) = y_p(c_+)\{y_{p,1}(x) + \alpha(p)y_{p,0}(x)\}$.

Proof of Theorem 4.9. By Lemma 4.8, for each $p \in \mathbf{I}$, there exists a unique solution $(\phi(p, \cdot), \psi(p, \cdot))$ of (4.4)–(4.6), up to constant multiples. We normalize the solution with $\phi(p, c_+) = 1$. By Step 1 in the proof of Proposition 4.10, $\phi(p, x)$ and $\phi'(p, x)$ are continuous in $\mathbf{I} \times \mathbb{R}$. Hence by (4.6), $\psi(p, c_+)$ is continuous in $p \in \mathbf{I}$, too. By Proposition 4.10 together with Remark 4.11, the function $\mathbf{I} \ni p \mapsto (\phi, \psi) \in \oplus^2 L^2(\mathbb{R})$ is measurable.

We can decompose the interval \mathbf{I} into infinitely many subintervals $\{\mathbf{I}_n\}_{n \in \mathbb{N}}$ such that the measure of \mathbf{I}_n is non-zero. For each $n \in \mathbb{N}$ we define a $\oplus^2 L^2(\mathbb{R})$ -valued function Φ_n on \mathbb{R} by

$$\Phi_n(p) = \begin{cases} \begin{pmatrix} \phi(p, \cdot) \\ -i\psi(p, \cdot) \end{pmatrix}, & \text{if } p \in \mathbf{I}_n, \\ 0, & \text{otherwise.} \end{cases}$$

Then, by Proposition 4.10, Φ_n is in $L^2(\mathbb{R}; \oplus^2 L^2(\mathbb{R}))$ and by the construction of Φ_n , we have $UH_+U^*\Phi_n = 0$. It is clear that for $n \neq m$ Φ_n is orthogonal to Φ_m . Consequently, it follows that $\dim \ker UH_+U^* = \infty$. \square

Remark 4.12. From the discussions in this section, we can see that there exist infinitely many generalized zero-energy states of H_+ . For the generalized eigenvector

$$\begin{pmatrix} f(x)e^{-i(p-\alpha/2)y} \\ -ig(x)e^{-i(p+\alpha/2)y} \end{pmatrix}$$

of N_+ with eigenvalue p , we have

$$Q_+\Psi = \begin{pmatrix} -ie^{-i(p+\alpha/2)y} \left(\bar{\lambda}v(x)f(x) + \frac{1}{2} \left(\frac{\partial}{\partial x} - p - \frac{\alpha}{2} \right) g(x) \right) \\ e^{-i(p-\alpha/2)y} \left(\frac{1}{2} \left(\frac{\partial}{\partial x} + p - \frac{\alpha}{2} \right) f(x) + \lambda v(x)g(x) \right) \end{pmatrix}.$$

Hence for every solution (f, g) of

$$\bar{\lambda}vf + \frac{1}{2} \left(\frac{\partial}{\partial x} - p - \frac{\alpha}{2} \right) g = 0 \quad \text{and} \quad \frac{1}{2} \left(\frac{\partial}{\partial x} + p - \frac{\alpha}{2} \right) f + \lambda vg = 0,$$

the vector $\Psi = (f(x)e^{-i(p-\alpha/2)y}, -ig(x)e^{-i(p+\alpha/2)y})$ satisfies the equation $H_+\Psi = Q_+^*Q_+\Psi = 0$. But $\Psi \notin L^2(\mathbb{R}^2)$. It means that Ψ is a generalized zero-energy eigenvector of H_+ . Since every $p \in \mathbb{R}$ is a generalized eigenvalue of N_+ , H_+ has infinitely many generalized zero-energy states. This kind of phenomenon occurs also in the $N=2$ WZ model with $V(z) = \lambda z^p$ and $V(z) = \lambda/z^p$.^{4,5}

D. A Concrete Class of Superpotentials.

In this subsection we give a concrete class of superpotentials, which satisfy the hypotheses of Theorems 4.9. In other words, we give a class of the N=2 WZ type models which have infinitely many bosonic zero-energy ground states.

Definition 4.13. We say that $V: \mathbb{C} \rightarrow \mathbb{C}$ is in \mathcal{V}_3 if there exist constants $\lambda \in \mathbb{C} \setminus \{0\}$, $\beta \in \mathbb{R}$ with $\beta \geq \alpha$ and $w \in C^\infty(\mathbb{R})$ such that:

- (1) $\partial V(z) = \lambda(1 + w(x))e^{\beta x} e^{i\alpha y}$,
- (2) $1 + w > 0$ on \mathbb{R} ,
- (3) $w(x) \rightarrow 0$ as $x \rightarrow \infty$,
- (4) $w(x), w'(x)$ and $w''(x) \rightarrow 0$ as $x \rightarrow -\infty$,
- (5) $w'/(1 + w)$ and $w''/(1 + w)$ are bounded on \mathbb{R} .

It is clear that $\mathcal{V}_3 \subset \mathcal{V}_1 \cap \mathcal{V}_2$.

Theorem 4.14. Let $V \in \mathcal{V}_3$. Then

$$\dim \ker H_+ = \infty.$$

Proof. We must prove that every function in \mathcal{V}_3 satisfies the hypotheses of Theorems 4.9. Since $V \in \mathcal{V}_3$, we can write $\partial V(z) = \lambda(1 + w(x))e^{\beta x} e^{i\alpha y}$. Let $\mathbf{I} = (0, \alpha)$. We have

$$\begin{aligned} G_p(x) = & 4|\lambda|^2(1 + w)^2 e^{2\beta x} + \frac{1}{4} \left(\alpha + \beta - 2p + \frac{w'}{1 + w} \right)^2 \\ & + \frac{1}{2} \left(\frac{w'}{1 + w} + \beta \right)^2 - \frac{1}{2} \left(\frac{w''}{1 + w} + 2\beta \frac{w'}{1 + w} + \beta^2 \right). \end{aligned}$$

By (4.13.5), $G_p(x)$ tends to infinity as $x \rightarrow \infty$ uniformly in $p \in \mathbf{I}$. Hence there exist constants c_+ and $\sigma > (\alpha + \beta)/2$ such that for every $p \in \mathbf{I}$, $G_p \geq \sigma^2$ on $[c_+, \infty)$ and for all $\rho \geq 0$, by (4.13.3),

$$h(x)e^{-\sigma x} x^\rho = \sqrt{1 + w(x)} e^{(\alpha + \beta)x/2} e^{-\sigma x} x^\rho \rightarrow 0 \quad \text{as } x \rightarrow \infty.$$

By (4.13.4), $G_p(x) \rightarrow (\alpha + \beta - 2p)^2/4$ as $x \rightarrow -\infty$ uniformly in $p \in \mathbf{I}$. Thus, for each $\theta \in (0, 1)$ there exists a constant c_- such that for every $p \in \mathbf{I}$,

$$\theta^2 \left(\frac{\alpha + \beta - 2p}{2} \right)^2 \leq G_p(x) \leq \frac{1}{\theta^2} \left(\frac{\alpha + \beta - 2p}{2} \right)^2, \quad x \in (-\infty, c_-].$$

We put $\gamma = \theta(\alpha + \beta - 2p)/2$ and $\delta = \theta^{-1}(\alpha + \beta - 2p)/2$. We have

$$\frac{1}{\theta} - \frac{\theta}{2} > \frac{1}{2} \quad \text{and} \quad \frac{1}{\theta} - \frac{\theta}{2} \rightarrow \frac{1}{2} \quad \text{as} \quad \theta \rightarrow 1.$$

Hence for every $\varepsilon > 0$, there exists θ in $(0, 1)$ such that

$$\frac{1}{\theta} - \frac{\theta}{2} = \frac{1 + \varepsilon}{2}.$$

Thus,

$$2\delta - \gamma - \frac{\alpha + \beta}{2} = \frac{\varepsilon}{2}(\alpha + \beta) - (1 + \varepsilon)p.$$

Since $p > 0$, it follows that for all sufficiently small ε , $2\delta - \gamma - (\alpha + \beta)/2 < 0$, which implies that for all $\rho \geq 0$

$$h(x)e^{(2\delta - \gamma)|x|}|x|^\rho = \sqrt{1 + w(x)}e^{(\alpha + \beta)x/2}e^{(2\delta - \gamma)|x|}|x|^\rho \rightarrow 0 \quad \text{as} \quad x \rightarrow -\infty.$$

Hence hypothesis (i) of Theorem 4.9 is satisfied. Similarly we can show that hypothesis (ii) of Theorem 4.9 is satisfied. Hypothesis (iii) of Theorem 4.9 can easily be verified. We have thus proved Theorem 4.14. \square

It is obvious that $V(z) = \lambda e^{\alpha z}$ with $\lambda \in \mathbb{C} \setminus \{0\}$ is in \mathcal{V}_3 . Hence we obtain the following result, which has been proved in Ref. 6.

Corollary 4.15. *Let $V(z) = \lambda e^{\alpha z}$ with $\lambda \in \mathbb{C} \setminus \{0\}$. Then*

$$\dim \ker H_+ = \infty.$$

Example 4.16. Let $\mathcal{S}(\mathbb{R})$ be the set of rapidly decreasing C^∞ -functions on \mathbb{R} . Let

$$SP = \left\{ V: \mathbb{C} \rightarrow \mathbb{C} \left| \begin{array}{l} \exists \lambda \in \mathbb{C} \setminus \{0\}, \exists \beta \geq \alpha, \exists s \in \mathcal{S}(\mathbb{R}) \\ \text{s.t. } \inf_{x \in \mathbb{R}} \{1 + s(x)\} > 0, \\ \partial V(z) = \lambda(1 + s(x))e^{\beta x} e^{i\alpha y} \end{array} \right. \right\}.$$

Then it is clear that every function in SP is in \mathcal{V}_3 .

V. THE VANISHING THEOREM

In the last section, we formulated a condition under which $\dim \ker H_+ = \infty$. In this section we consider a condition under which $\dim \ker H_- = 0$. By Remark 3.7, we can not expect that the kernel of H_- is stable. Hence a class of V for which $\ker H_- = \emptyset$ may be very small.

We put

$$F_p(x) = \frac{v''}{v} - \left(\frac{v'}{v}\right)^2 + 2(p - \alpha) \left(\frac{v'}{v} - p\right).$$

Definition 5.1. We say that V is in \mathcal{V}_4 if there exist a constant $\lambda \in \mathbb{C} \setminus \{0\}$ and $v \in C^\infty(\mathbb{R})$ such that $\partial V(z) = \lambda v(x)e^{i\alpha y}$, v'/v and v''/v are bounded on \mathbb{R} and $v'/v \in C^\infty(\mathbb{R})$.

Theorem 5.2. Let $V \in \mathcal{V}_4$. Suppose that for a.e. $p \in \mathbb{R}$, either $F_p \leq 0$ on \mathbb{R} or $F_{\alpha-p} \leq 0$ on \mathbb{R} holds. Then,

$$\dim \ker H_- = 0.$$

To prove Theorem 5.2, we prepare some lemmas. We first note that there exists a symmetry group of the fermionic Hamiltonian H_- . Let

$$N_- = i \frac{\partial}{\partial y} - \frac{\alpha}{2} \sigma_3 \quad \text{with} \quad D(N_-) = D\left(\frac{\partial}{\partial y}\right) \subset \mathcal{H}_-.$$

Then the following lemma holds (see Lemma 4.1).

Lemma 5.3. Let $V \in \mathcal{V}_0$. Then for all $s, t \in \mathbb{R}$,

$$e^{isH_-} e^{itN_-} = e^{itN_-} e^{isH_-}.$$

The operator $\tilde{U}: L^2(\mathbb{R}^2; \mathbb{C}^2) \rightarrow L^2(\mathbb{R}; \oplus^2 L^2(\mathbb{R}))$ given by

$$\tilde{U} \begin{pmatrix} f \\ g \end{pmatrix} (p) = \begin{pmatrix} (2\pi)^{-1/2} \int_{\mathbb{R}} e^{-i(p-\alpha)y} f(y, x) dy \\ (2\pi)^{-1/2} \int_{\mathbb{R}} e^{-ipy} g(y, x) dy \end{pmatrix}$$

is unitary. To identify the kernel of $H_- = Q_+ Q_+^*$, we solve $\tilde{U} Q_+^* \tilde{U}^* \Phi = 0$ on each of the fiber spaces of $L^2(\mathbb{R}; \oplus^2 L^2(\mathbb{R}))$ in a way similar to that in Sec. IV.

Let $V \in \mathcal{V}_4$. Let $\Phi \in \ker(\tilde{U}H - \tilde{U}^*)$ with $\Phi(p) = (\phi(p, \cdot), -i\psi(p, \cdot))$. In the same way as in the proof of Lemma 4.5, we can see that for a.e. p the following equations hold in the ordinary sense:

$$-\phi'' + \left(\frac{v'}{v} - \alpha\right)\phi' + \left(4|\lambda|^2v^2 - (p - \alpha)\frac{v'}{v} + p(p - \alpha)\right)\phi = 0, \quad (5.1)$$

$$-\psi'' + \left(\frac{v'}{v} - \alpha\right)\psi' + \left(4|\lambda|^2v^2 + p\frac{v'}{v} + p(p - \alpha)\right)\psi = 0, \quad (5.2)$$

$$\frac{1}{2}\left(\frac{d}{dx} - p + \alpha\right)\phi + \bar{\lambda}v\psi = 0 \quad \text{and} \quad \lambda v\phi + \frac{1}{2}\left(\frac{d}{dx} + p\right)\psi = 0. \quad (5.3)$$

We fix p in the following lemma.

Lemma 5.4. *Let $V \in \mathcal{V}_4$. Let ϕ (resp. ψ) be a solution of (5.1) (resp. (5.2)) in $D(d^2/dx^2) \cap D(v^2)$. Then the following hold:*

- (i) *If $F_p \leq 0$ on \mathbb{R} , then $\phi = 0$.*
- (ii) *If $F_{\alpha-p} \leq 0$ on \mathbb{R} , then $\psi = 0$.*

Remark. By (5.3), either $\phi = 0$ or $\psi = 0$ implies that both of them are equal to 0.

Proof. We prove part (i). Taking the $L^2(\mathbb{R})$ -inner product of (5.1) with ϕ yields that

$$\|\phi'\|^2 + \|(2\lambda v)\phi\|^2 + \left\langle \left(\frac{v'}{v} - \alpha\right)\phi', \phi \right\rangle - \left\langle (p - \alpha)\left(\frac{v'}{v} - p\right)\phi, \phi \right\rangle = 0. \quad (5.4)$$

Via integration by parts, we have

$$2\operatorname{Re} \left\langle \left(\frac{v'}{v} - \alpha\right)\phi', \phi \right\rangle = \left\langle \left\{ \left(\frac{v'}{v}\right)^2 - \frac{v''}{v} \right\} \phi, \phi \right\rangle.$$

(5.4) implies that

$$2\|\phi'\|^2 + 2\|(2\lambda v)\phi\|^2 = \left\langle \left\{ \frac{v''}{v} - \left(\frac{v'}{v}\right)^2 + 2(p - \alpha)\left(\frac{v'}{v} - p\right) \right\} \phi, \phi \right\rangle = \langle F_p \phi, \phi \rangle. \quad (5.5)$$

By the hypothesis of part (i) the right hand side of (5.5) is non-positive. But the left hand side of (5.5) is non-negative. Hence the both sides of (5.5) are equal to 0. Hence it follows that $\phi = 0$. In the same way we can prove part (ii). \square

We are now ready to prove Theorem 5.2.

Proof of Theorem 5.2. By the assumption of this theorem and Lemma 5.4, we immediately obtain the desired result. \square

A concrete example of superpotential satisfying the assumption in Theorem 5.2 is given by

$$\partial V(z) = \lambda e^{\beta x} e^{i\alpha y}$$

with constants $\lambda \in \mathbb{C} \setminus \{0\}$ and $\beta > \alpha/2$.

Remark 5.5. By a reason similar to that given in Remark 4.12, we can see that for H_- , there exist infinitely many generalized zero-energy states.

APPENDIX

In this appendix, we present some facts on the ordinary differential equation,

$$y'' - Gy = 0 \quad \text{on } [0, \infty). \quad (\text{A.1})$$

For the proof of them, we refer to, for example, Sec. 9.2 in Ref.10.

We suppose that the function G is positive, continuous and $xG(x) \notin L^1(0, \infty)$. Let y_0 and y_1 be the solutions of (A.1) with

$$y_1(0) = 1, \quad y_1'(0) = 0, \quad y_0(0) = 0 \quad \text{and} \quad y_0'(0) = 1.$$

Lemma A.1. $y_0(x) > 0$ and $y_1(x) > 0$ on $(0, \infty)$.

Lemma A.2. The limit $\alpha = -\lim_{x \rightarrow \infty} y_1(x)/y_0(x)$ exists and $y_+(x) = y_1(x) + \alpha y_0(x)$ is the unique solution of (A.1) such that

$$y_+(0) = 1 \quad \text{and} \quad y_+(x) \rightarrow 0 \quad \text{as} \quad x \rightarrow \infty.$$

Moreover, y_+ is monotone decreasing on $(0, \infty)$ and $y_+' \in L^2(0, \infty)$.

We call y_+ in Lemma A.2 the “normalized subdominant” solution.

Lemma A.3. $\lim_{x \rightarrow \infty} y_1(x) = \infty$ and $\lim_{x \rightarrow \infty} y_0(x) = \infty$. Hence every solution of (A.1) linearly independent of y_+ tends to infinity as $x \rightarrow \infty$.

The following comparison theorem holds.

Lemma A.4. Let H be a function such that $H \geq G$. Let Y_+ be the normalized subdominant solution of $y'' - Hy = 0$ and Y_0 be the solution of it with $Y_0(0) = 0$ and $Y_0'(0) = 1$. Then

$$0 < Y_+ \leq y_+ \quad \text{and} \quad 0 < y_0 \leq Y_0 \quad \text{on } (0, \infty).$$

In particular, if there exist γ and $\delta > 0$ such that $\delta^2 \leq G \leq \gamma^2$, then the following estimate holds:

$$e^{-\gamma x} \leq y_+(x) \leq e^{-\delta x} \quad \text{on } (0, \infty).$$

ACKNOWLEDGEMENT

The author would like to thank Professor Asao Arai for helpful discussions.

References.

1. E. Witten, *Dynamical breaking of supersymmetry*, Nucl. Phys. B 188 (1981), 513.
2. E. Witten, *Constraints on supersymmetry breaking*, Nucl. Phys. B 202 (1982), 253.
3. A. Jaffe, A. Lesniewski and M. Lewenstein, *Ground state structure in supersymmetric quantum mechanics*, Ann. Phys. 178 (1987), 313.
4. A. Arai, *On the degeneracy in the ground state of the $N=2$ Wess-Zumino supersymmetric quantum mechanics*, J. Math. Phys. 30 (1989), 2973.
5. A. Arai and O. Ogurisu, *Meromorphic $N=2$ Wess-Zumino supersymmetric quantum mechanics*, J. Math. Phys. 32 (1991), 2427.
6. A. Arai, *Existence of infinitely many zero-energy states in a model of supersymmetric quantum mechanics*, J. Math. Phys. 30 (1989), 1164.
7. P. Deift, *Applications of a commutation formula*, Duck Math. J. 45 (1978), 267.
8. Paul R. Chernoff, *Essentially self-adjointness of powers of generators of hyperbolic equations*, J. Funct. Anal. 12 (1973), 401.
9. M. Reed and B. Simon, *Methods of Modern Mathematical Physics Vol. I, II, IV*, Academic, New York, 1979.
10. E. Hille, *Lectures on Ordinary Differential Equations*, Addison-Wesley Publishing Company.