



Queensland University of Technology
Brisbane Australia

This may be the author's version of a work that was submitted/accepted for publication in the following source:

[Li, Tong, Oloyede, Adekunle, & Gu, YuanTong](#)
(2013)

F-actin crosslinker: A key player for the mechanical stability of filopodial protrusion.

Journal of Applied Physics, 114(21), pp. 1-5.

This file was downloaded from: <https://eprints.qut.edu.au/220363/>

© Consult author(s) regarding copyright matters

This work is covered by copyright. Unless the document is being made available under a Creative Commons Licence, you must assume that re-use is limited to personal use and that permission from the copyright owner must be obtained for all other uses. If the document is available under a Creative Commons License (or other specified license) then refer to the Licence for details of permitted re-use. It is a condition of access that users recognise and abide by the legal requirements associated with these rights. If you believe that this work infringes copyright please provide details by email to qut.copyright@qut.edu.au

Notice: *Please note that this document may not be the Version of Record (i.e. published version) of the work. Author manuscript versions (as Submitted for peer review or as Accepted for publication after peer review) can be identified by an absence of publisher branding and/or typeset appearance. If there is any doubt, please refer to the published source.*

<https://doi.org/10.1063/1.4839715>

F-actin crosslinker: a key player for the mechanical stability of filopodial protrusion

Tong Li, Adekunle Oloyede and Y.T. Gu*

School of Chemistry, Physics and Mechanical Engineering, Queensland University of Technology,
Brisbane, Australia

Abstract: Filopodial protrusion initiates cell migration, which decides the fate of cells in biological environments. In order to understand the structural stability of ultra-slender filopodial protrusion, we have developed an explicit modeling strategy that can study both static and dynamic characteristics of microfilament bundles. Our study reveals that the stability of filopodial protrusions is dependent on the density of F-actin crosslinkers. This cross-linkage strategy is a requirement for the optimization of cell structures, resulting in the provision and maintenance of adequate bending stiffness and buckling resistance while mediating the vibration. This cross-linkage strategy explains the mechanical stability of filopodial protrusion and helps understand the mechanisms of mechanically induced cellular activities.

Keywords: F-actin crosslinker, Filopodial protrusion, Mechanical stability, Biomechanics

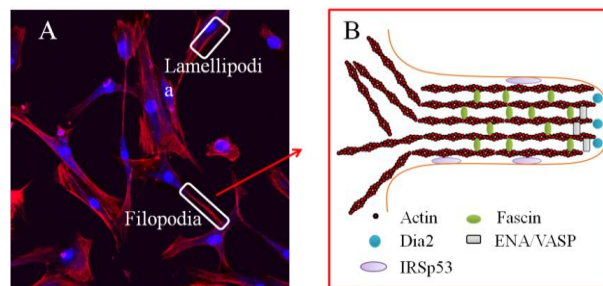
* Corresponding author. YuanTong Gu

Tel.: +61-7-31381009; fax: +61-7-31381469.

E-mail address: yuantong.gu@qut.edu.au (Y.T. Gu).

1 I. INTRODUCTION

2 Throughout our life, multicellular organism spatiotemporally coordinates various physiological processes,
3 such as embryonic morphogenesis and tissue formation¹, by elaborating the spreading and migration of
4 cells^{2,3}. Cell migration also initiates with leading edge protrusion, and finalizes with cell movements yielded
5 by intracellular cytoskeleton contraction from cell protrusions⁴. Fig 1 shows the *in vivo* morphology of
6 osteoblasts during cell migration. The flat, sheet like F-actin network in cells is lamellipodia, from where
7 cell protrusion initiates⁵. Various functional proteins comprehensively stimulate the lamellipodia to form
8 needle like, highly dynamic cell protrusions: filopodia⁶. The filopodia is made up of F-actin that are bundled
9 by actin binding proteins⁷ and acts as the mechanical unit of cell migration frontier. Abnormal cell
10 protrusions due to improper mechanical properties of filopodia can initiate unhealthy cell migration, which
11 leads to human diseases such as immune disorders and proliferation of tumor cells^{8,9}. Therefore,
12 understandings of the mechanical stability of filopodial protrusions are significant to cell pathology studies.



14 Fig. 1 In vivo cell migration and the illustration of filopodia structure. A. Osteoblasts morphology evaluation by laser Confocal
15 microscope (Nikon A1R Confocal system) after 24 hours culturing on 22mm×22mm glass cover. DAPI is employed for the
16 visualization of nucleus (in blue) and rhodamine-phalloidin is employed for the visualization of F-actin (in red). B: Molecular
17 structure of filopodia. Dia2 can nucleate the formation of new, unbranched F-actin. ENV/VASP crosslinks F-actin in the tip of
18 protruding filopodia. IRSp53 might sense the negative membrane curvature and initiates new filopodia. Fascin is the major F-actin
19 cross-linkage protein in filopodia.

20 The physiological environment of living cells includes various potential mechanical loadings, such as
21 extracellular interstitial fluids and intercellular cytoskeleton self-contraction. Fig 2 shows typical loading
22 models on living cell structures: bending model and buckling model. The focal adhesions of living cells are
23 mainly distributed at leading and trailing edges of cells¹⁰, which can be imaged as mechanical supports to
24 cell structures (e.g. filopodia and lamellipodia) under the condition of mechanical loadings. The diameter of
25 single F-actin is under 10nm¹¹, while the length of filopodia can be up to 40μm in various organisms¹². The
26 microscale F-actin length and nanoscale filament cross section result in large aspect ratio, making single F-
27 actin fragile to undergo the challenging mechanical conditions.

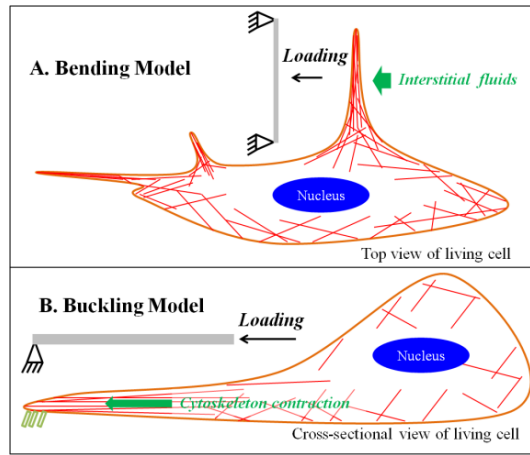


Fig. 2 Potential loading models of filopodia in biological environments. A: Bending model of filopodia. The filopodia can be simplified as thin beam. The tip of stress fiber is fixed because of focal adhesions while the other end belongs to the main body of cell. The interstitial fluids can apply transient or continuous transverse loadings to the filopodia. B. Buckling model of stress fibers (F-actin bundle), e.g. filopodia and pseudopodia. Contraction due to Myosin II acts as axial compression on the slender beam.

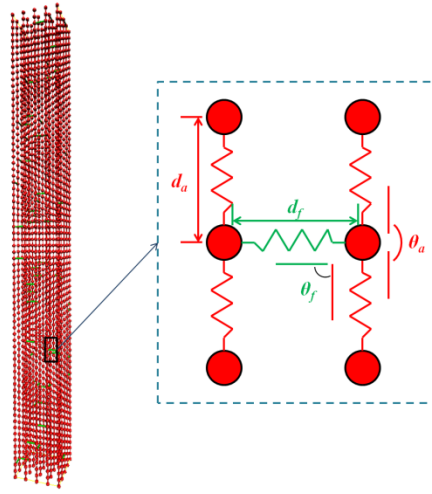
In living cells, single F-actin filaments are tightly bundled by crosslinkers to mediate the mechanical performances of filopodial protrusion^{13,14}, and the main cross-linkage protein in filopodia is fascin¹⁵. The mechanical properties of F-actin bundle have been explored by pure bending theory of slender beam¹⁶ and worm-like chain (WLC) characterization¹³. Based on experimental findings, a mechanics model for this semiflexible biological organism was proposed¹⁷ and numerical simulations successfully predicted the bending stiffness and buckling resistance with respect to the features of F-actin bundle structures^{18,19}. However, these continuum mechanics based models have difficulties in modeling the large, nonlinear deformation of semiflexible F-actin bundles which is comparable to the thickness of filopodia. Also, the dynamic response of filopodia protrusion after excitation is hard to be captured by adopting finite element method based modeling strategy. New modeling technique is needed to meet the requirements in capturing nonlinear deformations and dynamic response of semiflexible filopodia. Recently, multiscale modeling method based on experiments and molecular simulation shows great potential in the dynamics simulation of bio-inspired materials²⁰. A coarse-grained model is specifically developed for the mechanical deformation modeling of F-actin based on atomistic modeling of F-actin filaments²¹.

In this paper, crosslinkers are implemented to the coarse-grained model of F-actin bundle²¹ to investigate the mechanical performances of filopodial protrusion. This model aims to explicitly extract the complex mechanical behaviors of F-actin bundles considering influences from fascin-actin cross-linkage. The significance of F-actin crosslinkers in stabilizing the filopodial protrusion is evaluated based on static and dynamic characteristics of fascin binding F-actin bundles.

II. COMPUTATIONAL MODEL

In the modeling strategy, single filaments are simplified as particle strings that are connected by F-actin crosslinkers, as shown in Fig 3. The interaction between actin clusters on the same F-actin and the cross-

1 linkage between neighboring F-actin are assumed to be harmonic pair potential energy: $E = \frac{1}{2}k(r - r_0)^2$, where
 2 k is the energy scale parameter, r and r_0 correspondingly denote the actual distance and equilibrium distance
 3 between actin clusters. Similar characterizations for the angular potential energy functions are defined with
 4 respect to the angle between two neighboring connections. This explicit is capable to obtain dynamic
 5 responses of F-actin bundles after mechanical excitation. The stiffness of stress fibers that consists of actin
 6 filament has dependency on temperature²². However, we only adopt the filament stiffness parameters at
 7 human body temperature²¹ in this study. The randomly distributed weak interactions between G-actin
 8 monomers, e.g. H-bond and disulfide bond, are not considered since they are negligible comparing to
 9 mechanical loadings.



10
 11 Fig. 3 The schematic of granular simulation strategy for F-actin bundle modeling. Each virtual particle (red circle) consists of two
 12 actin monomers, and the equilibrium distance between neighboring particles is 5.53nm. The longitude stiffness of a 1 μ m long F-
 13 actin follows the experimental finding of 43pN/nm,²³ the equilibrium angle is 180 $^\circ$ between actin-actin bonds, and the angular
 14 stiffness is characterized to be 7630 kCal/mol rad².²¹ The tensile stiffness of 37.5nm long fascin is assumed as 1.5pN/nm, the
 15 equilibrium angle is 90 $^\circ$ for actin-fascin connection, and the angular stiffness is 500kCal/mol rad².

16 The cross section of bundle is assumed quadrate, consisting of 25 (5 \times 5) filaments. The length of this actin
 17 bundle is 2 μ m and the transverse distance between filaments is 37.5nm¹⁸, making the total bundle thickness
 18 150nm. The profile of F-actin bundle in this work follows the filopodia length prediction from literature¹⁹.
 19 The cross-linkage randomly occurs between two neighboring particles at the same longitudinal position
 20 from different filaments²⁴. Five scenarios of F-actin bundles with different crosslinker densities are selected
 21 to understand sensitivity of mechanical behaviors with respect to the crosslinker density. The crosslinker
 22 density ratio is calculated as $\alpha = n_f/n_a$, where, n_f and n_a respectively denote quantities of F-actin
 23 crosslinkers and actin monomers. Canonical ensemble (NVT) is a thermal bath, in which particle numbers,
 24 system volume and temperature are constant. Herein, NVT ensemble is employed with a 303K temperature
 25 to model the bending and buckling of F-actin bundle. Langevin dynamics algorithm²⁵ is employed to model
 26 the friction from implicit solvent. In the Langevin dynamics algorithm, two terms are added to the force
 27 calculation on each particle: viscous damping term due to solvent and a randomly bumping term due to

1 temperature. The combination of these two terms is $F_d = -\frac{m}{C_d}v + \sqrt{\frac{mk_B T}{dtC_d}}$. Where, m is the mass of particle, v is
2 the velocity, k_B is the Bozeman constant, T is the temperature, dt is the time step and C_d is the damping
3 factor with a time unit. In this simulation, C_d is chosen as $1ps$ to understand the sensitivity of dissipative
4 force to the density of F-actin crosslinkers.

5 The size of time step is $0.1ps$ in bending/buckling simulations and $0.2ps$ in the vibration simulation. The
6 granular simulations are performed with Lammmps²⁶ on HP Z600 workstation (Intel Xeon X5570) and all the
7 atomistic visualizations are finished on VMD²⁷.

8III. RESULTS AND DISCUSSION

9 The mechanics models of transverse loading conditions and corresponding characterization results are
10 detailed in Fig 4. In the bending test simulation, a rigid cylinder is designed to indent the F-actin bundle to
11 model the transverse bending of F-actin bundle. The velocity of indentation is $150nm$ per microsecond,
12 which is a slow loading rate in molecular simulations to allow the relaxation of stress-wave. The evaluation
13 of bending stiffness κ_B is based on Euler-Bernoulli beam theory under double clamped boundary condition:

14 $\kappa_B = \frac{Fl^3}{192\delta}$, where, F is the reaction due to rigid cylinder indentation, l is the bundle length and δ is the

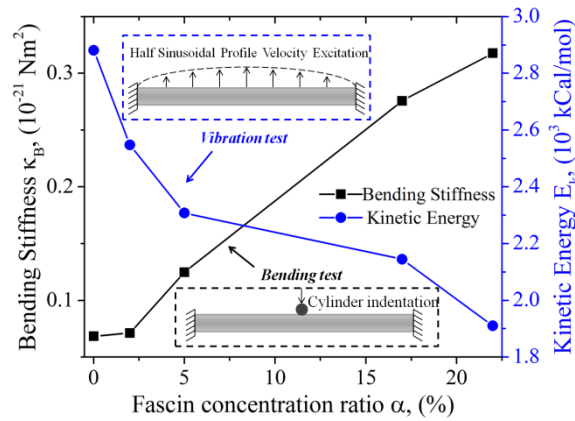
15 deflection of F-actin bundle. It should be noted that the distribution of crosslinkers on the F-actin bundle
16 also has effects on the mechanical behaviors of the whole bundle structure²⁸. With the same crosslinker
17 density, F-actin bundle with uniformly distributed crosslinkers presents higher mechanical stiffness in
18 bending simulation. Based on the nature of filopodial protrusion, the cases with uniformly distributed
19 crosslinkers are adopted for the mechanical characterization in this paper.

20 The bending stiffness of F-actin bundles(black solid square in Fig 4) increases with the density of
21 crosslinkers, which is consistent with previous experimental findings¹³. Refer to the interlaminar shear stress
22 in continuum mechanics model, the deformation of actin crosslinkers between different F-actin filaments
23 can store potential energy that is caused by external transverse loadings and increase the bending stiffness of
24 F-actin bundles. Direct indentation tests of *in vivo* F-actin bundles have provided the evaluation of the
25 bending stiffness, which is at the scale of $\kappa_B = 10^{-22}$ to $10^{-21} Nm^2$ ¹⁶, while WLC characterization of F-actin
26 bundles leads to the stiffness evaluation at scale of $\kappa_B = 10^{-23} Nm^2$ ^{13,14}. The bending stiffness prediction in this
27 paper, as shown in Fig 4, agrees with these experimental findings and is close to the characterization by
28 direct indentation¹⁶, as our simulation is also based on beam bending model.

29 Transverse impact is another typical loading on filopodia in cell surviving environments (e.g. from sudden
30 interstitial flow). The capability of filopodia to mediate the violent vibration after transient excitation is
31 significant in stabilizing healthy cell protrusion modes under impacts. We have employed the proposed
32 computational model to capture dynamic responses of F-actin bundles after transient excitation. In the

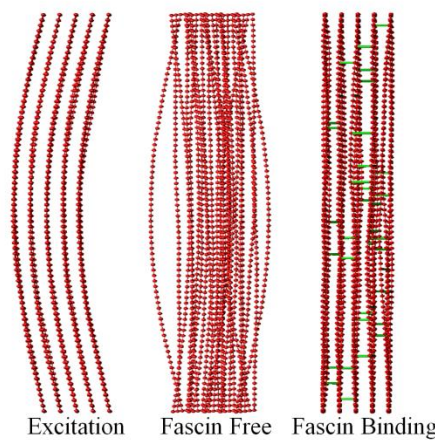
1 modeling, the F-actin bundle is fixed with the same boundary conditions as aforementioned modeling of
 2 transverse indentation. A transient velocity with half sinusoidal profile is applied on the whole bundle, and
 3 the velocity amplitude is 10^{-2} nm/ps .

4 According to the dynamic response extracted from numerical simulation, the kinetic energy of F-actin
 5 bundle decreases while the density of crosslinker increases, indicating that actin crosslinkers can mediate the
 6 vibration of ultra-slender filopodia protrusions, which is positive to enhance the stability of cell migration.



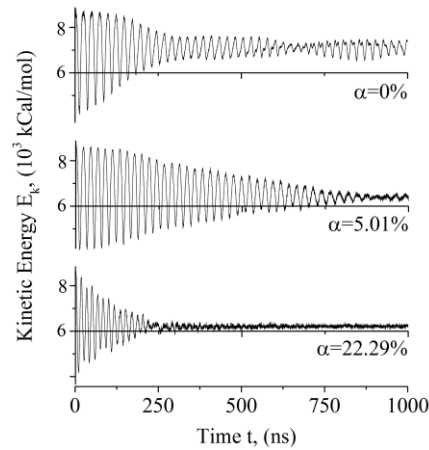
7
 8 Fig. 4 The calculation result from direct bending test and vibration test of F-actin bundles. The mechanics model of both
 9 simulation cases are double clamped slender beam. Different from indentation tests, half-sinusoidal profile velocity excitation is
 10 applied to the simulation system in vibration tests for $1 \mu\text{s}$, and the averaging kinetic energy in the last $0.1 \mu\text{s}$ (in which the kinetic
 11 energy is steady) is captured to evaluate the dynamic stability of the bundle system.

12 Fig 5 illustrates the simulation results of free vibration modes of F-actin bundles in different states without
 13 and with crosslinker binding. With the assistance of F-actin crosslinkers, F-actin filament is crosslinked to
 14 its neighboring filaments, providing mechanical supports to the filopodia protrusion and mediates the
 15 dynamic response of a whole F-actin bundle system. This characteristic improves the resistance of filopodial
 16 protrusion to undergo sudden impacts in complex biological environments that includes random transverse
 17 loadings.



18
 19 Fig. 5 Free vibration modes of F-actin bundles in different crosslinker binding states. The first illustration is the initial excitation
 20 applied on the bundle. The second and last illustrations correspondingly represent free vibration modes of F-actin without and
 21 with inner cross-linkages.

1 As a typical biological material, the semiflexible filopodia protrusion also presents highly overdamped
 2 property after transient mechanical loadings. By using the aforementioned Langevin dynamics algorithm,
 3 the dissipative characteristics of F-actin bundles can be estimated to understand the sensitivity of filopodia
 4 dissipation to F-actin crosslinker density. As shown in Fig 6, the kinetic energy dose not converge after $1\mu s$
 5 simulation in the scenario of zero F-actin crosslinker while both lower (5.01%) and higher (22.29%)
 6 crosslinker densities lead to the convergence of kinetic energy. Moreover, the kinetic energy dissipates faster
 7 with higher crosslinker density, indicating that crosslinker density contributes to the dissipative force on
 8 filopodia protrusion.



9

10 Fig. 6 Dissipative kinetic energy profiles of F-actin bundles with different crosslinker densities. The dissipation time for kinetic
 11 energy and the residue kinetic energy are both inversely proportional to the density of crosslinkers.

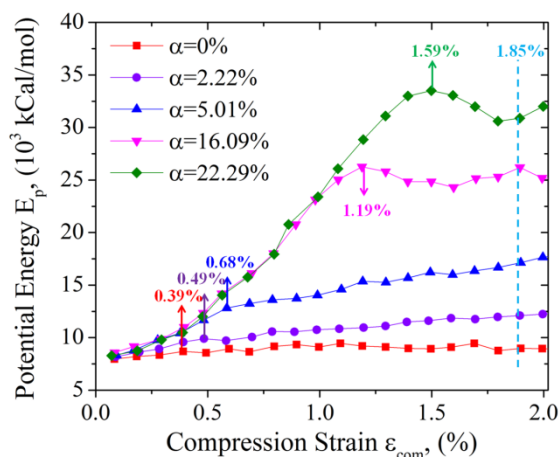
12 The filopodial protrusion also undergoes axial forces due to the intracellular cytoskeleton contraction from
 13 Myosin-II²⁹ or extracellular mechanical loadings. Another set of computational scenarios are set up to
 14 investigate the dynamic buckling behaviors of F-actin bundle in filopodia with different F-actin crosslinkers
 15 density. In the simulation models, one end of the F-actin bundle is fixed and a constant velocity is applied at
 16 the other end to model axial compression. The potential energy profile of F-actin bundle is summarized to
 17 characterize the buckling resistance of F-actin bundle model with different F-actin cross-linkage conditions,
 18 as shown in Fig 7.

19 The critical buckling strain of a continuum slender beam with quadrate cross section can be derived from

20 Euler's equation as: $\varepsilon_{cr} = \frac{\pi^2 a^2}{3l^2}$, where, a is the length of quadrant edge and l is the length of slender bundle.

21 Adopting 150nm edge length and $2\mu m$ bundle length, theoretical solution of the buckling strain of a
 22 nonporous F-actin bundle is 1.85%. The approximate buckling strain for porous F-actin bundle in the
 23 modeling can be directly extracted from the transitions on potential energy profile, where the increase of
 24 potential energy slows down or even stops. After this critical loading strain, the F-actin bundle is incapable
 25 to fully carry further axial compression. The buckling strain of porous F-actin bundles in modeling increases
 26 with the density of crosslinkers, and approaches the theoretical solutions of a continuum beam. Take the
 27 scenario of $\alpha = 22.29\%$ for example, by adopting the design strategy of cross-linkage, porous F-actin bundle

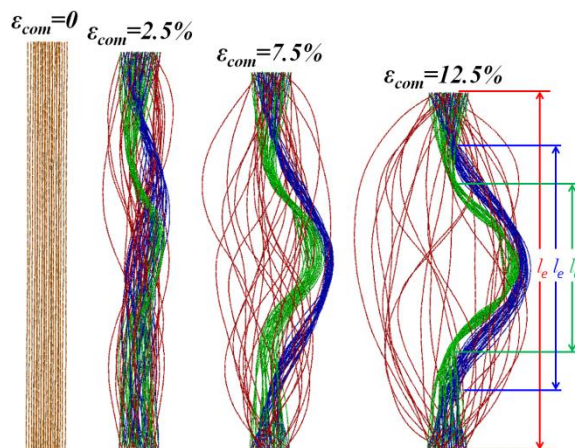
1 can approximately reduce 90% of the mass comparing to a nonporous bundle of F-actin filaments. However,
 2 the critical buckling strain only decreases 14% (from 1.85% to 1.59%).



3

4 Fig. 7 The characterization of F-actin bundles buckling with different crosslinker density. Buckling strain is marked with different
 5 colors in corresponding to the computational scenario and the light blue dotted line (1.85%) denotes the theoretical solution of
 6 buckling strain of gapless F-actin bundle with the same profile.

7 Explicit solutions of the buckling behaviors of F-actin bundle with different cross-linkage conditions are
 8 provided in Fig 8 to illustrate the molecular mechanisms of strengthened buckling resistance. F-actin
 9 bundles without cross-linkage (red bundles) are indeed individual actin filaments whose aspect ratio is too
 10 large to resist axial compression. With the help of crosslinkers, single filaments cooperate as a combined
 11 bundle to resist extracellular or intracellular mechanical challenges from its surviving environment. The
 12 effective lengths l_e (refer to Fig 8) of buckling models decreases with actin crosslinkers density, improving
 13 the buckling resistance. The post-buckling behaviors of crosslinked F-actin bundles also demonstrate that,
 14 with the assistance from actin crosslinkers (blue and green bundles), F-actin bundles can still partly carry the
 15 deformation energy caused by external loadings. Similar findings have been reported for carbon nano
 16 materials³⁰. This is another characteristic of ultra-slender filopodia protrusions to undergo complex
 17 mechanical conditions in biological environments.



18

19 Fig. 8 Buckling modes of F-actin bundles with respect to cross-linkage status. The red filaments represent non-bundled individual
 20 F-actin filaments, and the blue and green filaments respectively represent F-actin bundles with 5.01% and 22.29% crosslinker

1 density ratios. The effective length of the F-actin bundle (l_e) decreases with the density of crosslinkers, which explains the
2 increase of buckling resistance by crosslinkers.

3 Mechanics is one of the most primitive signaling systems for multicellular system³¹. The mechanical
4 stability is filopodial protrusion is critical to living cells as it is involved in various dynamic cellular
5 activities, such as spreading³, migration³² and adhesion³³. Take breast cancer cells for example, the high
6 density of fascin can regulate the invasion of cancer cells³⁴ whose mechanical stiffness is usually higher than
7 normal cells. The modeling results in this paper can theoretically prove that the mechanical performance of
8 filopodial protrusion has dependency on the quantity of crosslinker protein. By adjusting the cross-linkage
9 strategy between F-actin filaments, living cells can sensitively mediate the cytoskeleton mechanical
10 performance, which is significant for the pathological research and physics therapies of cell diseases.

11 IV. CONCLUSION

12 In the present study, the significant role F-actin crosslinkers (fascin) plays in enhancing the mechanical
13 stability of ultra-slender filopodia protrusion is investigated by developing an explicit granular simulation
14 strategy. The modeling of transverse and axial deformation demonstrates that the bending stiffness and
15 buckling resistance of ultra-slender filopodial protrusion are strengthened by the cross-linkage between
16 single F-actin filaments. The dynamics modeling of crosslinked F-actin bundle also proves that crosslinker
17 protein functions to mediate the vibration of filopodial protrusion after transient excitation. In summary, this
18 cross-linkage design of F-actin bundle in filopodial protrusion can stabilize the mechanical behaviors of
19 cellular activities in complex physiological environment.

20 ACKNOWLEDGEMENT

21 Support from the ARC Future Fellowship grant (FT100100172) is gratefully acknowledged. The authors
22 thank Dr. Ling Liu and Miss J.Q. Li for helpful discussion.

23 REFERENCE

- 24 ¹ A. J. Ridley, M. A. Schwartz, K. Burridge, R. A. Firtel, M. H. Ginsberg, G. Borisy, J. T. Parsons, and A. R. Horwitz,
25 *Science* **302**, 1704 (2003).
- 26 ² F. Chamaroux, S. Fache, F. Bruckert, and B. Fourcade, *Physical Review Letters* **94**, 158102 (2005).
- 27 ³ Y. Li, G.-K. Xu, B. Li, and X.-Q. Feng, *Applied Physics Letters* **96**, 043703 (2010).
- 28 ⁴ A. Ponti, M. Machacek, S. L. Gupton, C. M. Waterman-Storer, and G. Danuser, *Science* **305**, 1782 (2004).
- 29 ⁵ J. V. Small, T. Stradal, E. Vignal, and K. Rottner, *Trends in Cell Biology* **12**, 112 (2002).
- 30 ⁶ K. Burridge and K. Wennerberg, *Cell* **116**, 167 (2004).
- 31 ⁷ K. Lee, J. L. Gallop, K. Rambani, and M. W. Kirschner, *Science* **329**, 1341 (2010).
- 32 ⁸ D. A. Lauffenburger and A. F. Horwitz, *Cell* **84**, 359 (1996).
- 33 ⁹ R. P. Stevenson, D. Veltman, and L. M. Machesky, *Journal of Cell Science* **125**, 1073 (2012).
- 34 ¹⁰ M. A. Wozniak, K. Modzelewska, L. Kwong, and P. J. Keely, *Biochimica et Biophysica Acta (BBA) - Molecular Cell*
35 *Research* **1692**, 103 (2004).
- 36 ¹¹ S. Sharma, E. E. Grintsevich, M. L. Phillips, E. Reisler, and J. K. Gimzewski, *Nano letters* (2011).
- 37 ¹² P. K. Mattila and P. Lappalainen, *Nature Reviews. Molecular Cell Biology* **9**, 446 (2008).
- 38 ¹³ M. M. A. E. Claessens, M. Bathe, E. Frey, and A. R. Bausch, *Nature Materials* **5**, 748 (2006).
- 39 ¹⁴ A. Zidovska and E. Sackmann, *Biophysical Journal* **100**, 1428 (2011).
- 40 ¹⁵ D. A. Fletcher and R. D. Mullins, *Nature* **463**, 485 (2010).
- 41 ¹⁶ J. H. Shin, L. Mahadevan, P. T. So, and P. Matsudaira, *Journal of Molecular Biology* **337**, 255 (2004).

1 17 O. Lieleg, M. M. A. E. Claessens, C. Heussinger, E. Frey, and A. R. Bausch, *Physical Review Letters* **99**, 088102 (2007).
2 18 M. Bathe, C. Heussinger, M. M. A. E. Claessens, A. R. Bausch, and E. Frey, *Biophysical Journal* **94**, 2955 (2008).
3 19 A. Mogilner and B. Rubinstein, *Biophysical Journal* **89**, 782 (2005).
4 20 L. Liu, L. Zhang, and J. Lua, *Applied Physics Letters* **101**, 161907 (2012).
5 21 T. Li, Y. T. Gu, X.-Q. Feng, P. K. D. V. Yarlagadda, and A. Oloyede, *Journal of Applied Physics* **113**, 194701 (2013).
6 22 K. Kroy and E. Frey, *Physical Review Letters* **77**, 306 (1996).
7 23 H. Kojima, A. Ishijima, and T. Yanagida, *Proceedings of the National Academy of Sciences* **91**, 12962 (1994).
8 24 O. Pelletier, E. Pokidysheva, L. S. Hirst, N. Bouxsein, Y. Li, and C. R. Safinya, *Physical Review Letters* **91**, 148102
9 (2003).
10 25 G. E. Uhlenbeck and L. S. Ornstein, *Physical Review* **36**, 823 (1930).
11 26 S. Plimpton, Sandia National Laboratories (2007).
12 27 W. Humphrey, A. Dalke, and K. Schulten, *Journal of Molecular Graphics* **14**, 33 (1996).
13 28 See supplementary material at [URL will be inserted by AIP] for modeling results with different crosslinker distributions.
14 29 C. A. Wilson, M. A. Tsuchida, G. M. Allen, E. L. Barnhart, K. T. Applegate, P. T. Yam, L. Ji, K. Keren, G. Danuser, and
15 J. A. Theriot, *Nature* **465**, 373 (2010).
16 30 L. Liu, G. Cao, and X. Chen, *Journal of Nanomaterials* **2008**, 12 (2008).
17 31 E. Jonietz, *Nature* **491**, S56 (2012).
18 32 M. L. Gardel, I. C. Schneider, Y. Aratyn-Schaus, and C. M. Waterman, *Annual review of cell and developmental biology*
19 **26**, 315 (2010).
20 33 G.-K. Xu, X.-Q. Feng, H.-P. Zhao, and B. Li, *Physical Review E* **80**, 011921 (2009).
21 34 M. Al-Alwan, S. Olabi, H. Ghebeh, E. Barhoush, A. Tulbah, T. Al-Tweigeri, D. Ajarim, and C. Adra, *PLoS ONE* **6**,
22 e27339 (2011).