Fabrication of Thick Periodically-poled Lithium Niobate Crystals by Standard Electric Field Poling and Direct Bonding

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We fabricated 1 mm-thick periodically-poled lithium niobate (PPLN) crystals by using a high-voltage amplifier for standard electric field poling combined with a voltage multiplier. Furthermore, two 1 mm-thick PPLNs were directly bonded to make a 2 mm-thick PPLN. The large aperture allowed broad angular tuning, and a broad spectral range of quasi-phase matched second-harmonic generation can be achieved in a single channel. High-power applications are also expected.

Keywords: Periodically-poled lithium niobate, Quasi-phase-matching, Poling OCIS codes : (190.2620) Harmonic generation and mixing; (160.3730) Lithium niobate

I. INTRODUCTION

The quasi-phase-matching (QPM) technique provides not only high efficiency, but also many degrees of freedom in designing nonlinear optical frequency conversion devices [1]. It has been successfully implemented in ferroelectric crystals such as periodically-poled lithium niobate (PPLN) crystals, providing efficient optical second-harmonic generation (SHG) and other quadratic optical parametric processes, due to the large nonlinear optical tensor element d_{33} [2-3]. Although congruent lithium niobate (CLN) crystals are produced with reliable quality [4], the high coercive field allows poling of thin crystals only (< 1 mm). As a result, tight focusing is always needed, limiting their high power or intra-cavity applications. Ferroelectric crystals with low coercive fields have been developed as a solution to this problem, by reducing the microscopic defects which are the origins of the majority of the coercive fields. Thick crystals (1~5 mm) of MgO-doped and/or stoichiometric lithium niobate and lithium tantalate have been utilized in high-power QPM applications [5-6].

However, some problems still remain with the newlydeveloped ferroelectric crystals: (1) poling is not easy, sometimes producing low-quality QPM devices, (2) the reversed domains can relax due to the low coercive field, and (3) the uniformity of the domain width throughout the thickness is not as good as PPLNs made of CLN. Missey et al. reported diffusion bonding of three 1 mm-thick PPLNs, producing a large aperture PPLN for optical parametric oscillator [7]. In their work, the CLN stacks were pressed and heated at high temperature during bonding. Directbonding of various materials has been investigated for fabricating optical devices such as waveguides [8]. A buried waveguide was made by direct-hetero-bonding of lithium niobate and lithium tantalate crystals [9], and efficient SHG has been obtained with a similarly fabricated buried PPLN waveguide [10]. In this work, we utilized a similar directbonding method to make a large-aperture PPLN device for the demonstration of "slanted QPM" SHG. We poled 1 mm-thick CLN crystals by a modified electric field poling method, and by directly bonding two of them, we fabricated a 2 mm-thick PPLN. The quality of the bonded PPLN was good enough to perform a slanted QPM SHG experiment by tuning the wave propagation direction between the x-axis and the z-axis within the aperture of the thick PPLN. [Usually QPM is designed for non-critical propagation

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(wave vector perpendicular to the z- axis)].

II. ELECTRIC FIELD POLING OF 1 mm-CLN

We used the standard electric field poling technique [11] for poling 1 mm-thick CLN. However, the coercive field of CLN is ~ 22 kV/mm, while the maximum capacity of our high-voltage power amplifier is only ± 20 kV (TREK, 20/20B). We added a voltage-multiplier circuit in order to generate successive pulses of -5 kV and +17 kV [12], providing 22 kV to a z-cut 1 mm-CLN substrate (Crystal Technology), as described in Fig. 1. The QPM period was 18.6 μ m, the channel width was 3.0 mm (defined by a photomask), and the channel length was 15 mm along the x-axis.

III. DIRECT BONDING OF TWO 1 mm-PPLN CRYSTALS

We directly bonded the two 1 mm-PPLNs, making a 2 mm-thick PPLN. First, the 1 mm-PPLNs were cleaned and chemically treated in a H_2O_2 : NH₄OH: H_2O (1:1:6)

solution in order to form a hydrophilic OH-layer on the surfaces for bonding (z-faces) [13]. The two PPLNs were aligned (+ domains on top of +domains, etc.), and gently pressed in DI-water, followed by an annealing process at 700 °C for an hour. After annealing, the bonding was so strong that the original PPLNs could not be separated. Photographs of the bonded sample are shown in Fig. 2.

We verified excellent alignment between the ferroelectric domains in the top 1 mm-PPLN and the ones in the bottom PPLN throughout the bonded region as demonstrated in Fig. 2(a). As a result, we could fabricate a PPLN with thickness 2 mm, channel width 3 mm, and length 15 mm, with a only small unbounded region near the edge (Fig. 2(b)).

In order to check the optical loss of the bonding layer quantitatively, we measured the transmission spectrum, which was compared with those of single 1 mm-thick PPLN, and two unbounded 1 mm-thick PPLNs. As shown in Fig. 3, we could conclude that the transmission of the bonded sample (2 mm) is nearly the same as that of single PPLN in the non-absorbing spectral range between 450 and 1000 nm, indicating that the losses are mostly due to the Fresnel reflections at the air-crystal interfaces. Because the loss at the bonding interface is negligible, we can confirm that the optical quality of the bonded PPLN is excellent.

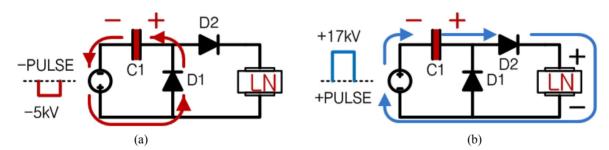


FIG. 1. Circuit diagram for voltage-multiplier (circle symbol: TREK amplifier, LN: 1 mm-CLN substrate). Diagrams (a) and (b) show the processes of generating a high-voltage pulse (22 kV); Charges of negative pulse (-5 kV) are stored at capacitor C1 (a), and then positive pulse (+17 kV) of the amplifier is added serially to C1. A total of 22 kV is applied to LCN. Arrows indicate the directions of electric current.

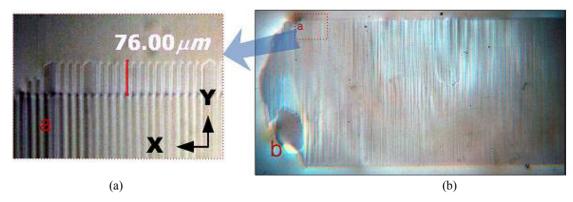


FIG. 2. Micrographs of the directly bonded PPLN (top view), showing good alignment of ferroelectric domains between top and bottom PPLNs (a), and larger region (lower magnification) of bonded sample (b) (channel width : 3.0 mm). Small unbonded area is also observed as indicated with 'b' at a corner.

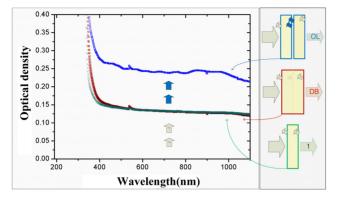


FIG. 3. Transmission spectra (optical density vs. wavelength) for bonded PPLN (red, 2 mm-thick), single 1 mm-thick PPLN (green), and two 1 mm-thick PPLNs (blue, unbounded). Each arrow in graph indicates the loss due to Fresnel reflection at an air-crystal interface.

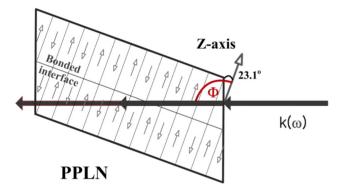


FIG. 4. Configuration of SHG experiment. Both fundamental and harmonic are extraordinary waves (polarized in the plane of paper). Φ : angle between the optic axis and internal wave vector (=113.1°).

IV. SHG EXPERIMENT

For slanted QPM SHG experiment, we cut and polished the bonded PPLN as shown in Fig. 4. The angle between the optic (z-) axis and the internal wave vector could be changed by rotating the sample in the xz-plane. As a light source, we used a β -barium borate optical parametric oscillator (OPO) pumped by the third harmonic of a Q-switched Nd:YAG laser. The idler wavelength from the OPO was scanned between 1580~1640 nm in order to obtain the SHG-tuning curve at a fixed angle of incidence (normal incidence in this work). The input intensity was attenuated and loosely focused to the sample.

The transmitted second harmonic intensity was measured as a function of fundamental wavelength at a fixed angle of incidence. The result is a typical sinc-square function centered at 806 nm, as shown in Fig. 5. The peak has been confirmed as the first order QPM wavelength by using the Sellmeier's formula of CLN for the extraordinary wave index $n_e(\Phi)$ [14]. The first order QPM wavelength is

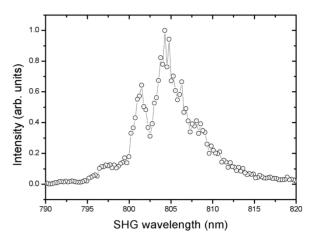


FIG. 5. Slanted QPM SHG intensity versus second-harmonic wavelength; $\Phi = 113.1^{\circ}$ (normal incidence).

767 nm for conventional PPLN (propagation along x-axis). A large shift of 39 nm (767 nm \rightarrow 806 nm) has been achieved by slanted QPM. We also note that the results are accurately predictable, due to the well-established Sellmeier formula [14] and reproducible material properties of CLN.

V. CONCLUSIONS

We fabricated a 2 mm-thick periodically-poled lithium niobate (PPLN) of high optical quality by direct bonding of two 1 mm-thick PPLNs made of congruent lithium niobate. As an application, we demonstrated "slanted" firstorder QPM SHG by properly choosing the wave propagation direction within the aperture of the thick PPLN. The large aperture allowed broad angular tuning, and a broad spectral range of quasi-phase matched second-harmonic generation can be achieved in a single channel. One can extend the tunable QPM range by stacking more PPLNs, fabricating thicker PPLN crystals. High-power applications such as optical parametric oscillators are also expected.

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