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FAST FINITE-DIFFERENCE SOLUTION OF
BIHARMONIC PROBLEMS

by

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1. Introduction

The advent of the high speed digital computer has been conducive to the development of new, and renewed interest in old, numerical techniques for biharmonic problems (see, e.g., references [1]-[4], [7]-[15]). Of these, the direct finite-difference method has been restrictive in that iteration for the resulting linear algebraic system is inherently slow [10]. Our objective here is to describe a finite-difference method for biharmonic problems which utilizes the fast iterative methods available for harmonic problems. The biharmonic then will be treated as a system of second order elliptic equations. However, as will be shown, it will be necessary, in addition, to incorporate a simple smoothing process to assure convergence [6]. After the method is described and illustrated, it will be shown how to adapt it easily to crack-type boundary value problems [2], [13].

2. A Prototype Problem

For simplicity, we will consider first a prototype biharmonic problem. Let S be a unit square and let R be the interior of S . Consider the boundary value problem on $R \cup S$ in which one seeks a function $u(x,y)$ which is continuous on $R \cup S$, which satisfies the biharmonic equation on R , and which satisfies prescribed function and normal derivative con-

ditions on S . More precisely, let S have vertices $A(0,0)$, $B(1,0)$, $C(1,1)$ and $D(0,1)$, as shown in Figure 2.1. Let f_i

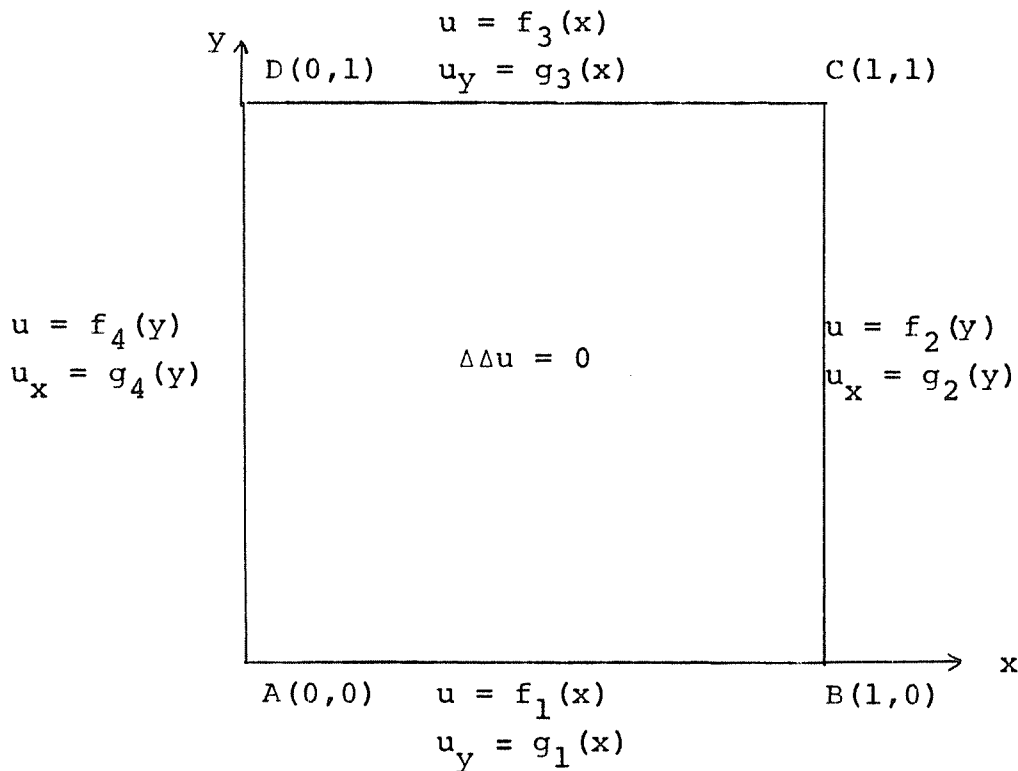


Figure 2.1

and g_i , $i = 1, 2, 3, 4$, be continuous functions of a single variable. Then our problem is to find a function $u(x, y)$ which is continuous on $R \cup S$, which is a solution on R of the biharmonic equation.

$$(2.1) \quad \Delta\Delta u = 0$$

and which satisfies the boundary conditions

$$\begin{aligned}
u(x,0) &= f_1(x), & u_y(x,0) &= g_1(x); & 0 \leq x \leq 1 \\
u(1,y) &= f_2(y), & u_x(1,y) &= g_2(y); & 0 \leq y \leq 1 \\
u(x,1) &= f_3(x), & u_y(x,1) &= g_3(x); & 0 \leq x \leq 1 \\
u(0,y) &= f_4(y), & u_x(0,y) &= g_4(y); & 0 \leq y \leq 1,
\end{aligned}$$

as shown in Figure 2.1.

3. The Numerical Method

The method to be developed for prototype problem (2.1)-(2.2) can be described in general as follows. First set

$$(3.1) \quad \Delta u = -\omega$$

on R , so that from (2.1)

$$(3.2) \quad \Delta \omega = 0.$$

Consideration of (2.1) will be replaced by consideration of the system (3.1)-(3.2). Physically, one can interpret u and ω , from, say, the fluid dynamics point of view, as stream and vorticity functions, respectively. Next, replace R and S by sets of grid points R_h and S_h in the usual way [5]. We will construct on R_h a sequence of discrete functions

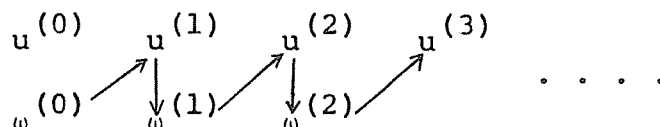
$$(3.3) \quad u^{(0)}, u^{(1)}, u^{(2)}, u^{(3)}, \dots$$

and on $R_h \cup S_h$ a sequence of discrete functions

$$(3.4) \quad \omega^{(0)}, \omega^{(1)}, \omega^{(2)}, \omega^{(3)}, \dots$$

which, with the aid of an elementary smoothing procedure, will converge to yield the numerical solution. The functions $u^{(k)}$

of (3.3) will result from a discretization of (3.1) while the functions $\omega^{(k)}$ of (3.4) will result from a discretization of (3.2). The double sequences (3.3) and (3.4) will be generated in an alternating fashion as follows:



Let us then now give a precise formulation of the algorithm:

Step 1.

For n a positive integer, let $h = \frac{1}{n}$ be the grid size, and construct and number in the usual way the interior grid points R_h and the boundary grid points S_h .

Step 2.

$$\begin{array}{l}
 \text{Set} \quad u^{(0)} = 0 \quad , \quad \text{on } R_h \\
 \quad \quad \omega^{(0)} = 0 \quad , \quad \text{on } R_h \cup S_h .
 \end{array}$$

Step 3.

To produce the second iterate $u^{(1)}$ of (3.3) on R_h proceed as follows. At each point of R_h write down the difference analogue

$$\begin{aligned}
 (3.5) \quad & -4u(x,y) + u(x+h,y) + u(x,y+h) + u(x-h,y) + u(x,y-h) \\
 & = -h^2 \omega^{(0)}(x,y)
 \end{aligned}$$

of (3.1) and, whenever possible, insert the given boundary values f_i from (2.2). Solve the resulting linear algebraic system by

SOR with over-relaxation factor r_1 and denote the solution by $\bar{u}^{(1)}$. Finally, define $u^{(1)}$ on R_h by the smoothing formula

$$(3.6) \quad u^{(1)} = \rho u^{(0)} + (1-\rho)\bar{u}^{(1)}, \quad 0 \leq \rho \leq 1.$$

Step 4.

To produce the second iterate $\bar{\omega}^{(1)}$ of (3.4) on $R_h \cup S_h$ proceed as follows. At each point of S_h of the form $(ih, 0)$, $i = 0, 1, 2, \dots, n$, set ([5], pp. 127-129)

$$(3.7) \quad \begin{aligned} \bar{\omega}^{(1)}(ih, 0) &= \frac{4}{h^2} f_1(ih) - \frac{1}{h^2} f_1(ih+h) - \frac{2}{h^2} u^{(1)}(ih, h) \\ &\quad - \frac{1}{h^2} f_1(ih-h) + \frac{2}{h} g_1(ih); \end{aligned}$$

at each point of S_h of the form $(1, ih)$, $i = 1, 2, \dots, n-1$, set

$$(3.8) \quad \begin{aligned} \bar{\omega}^{(1)}(1, ih) &= \frac{4}{h^2} f_2(ih) - \frac{1}{h^2} f_2(ih+h) - \frac{2}{h^2} u^{(1)}(1-h, ih) \\ &\quad - \frac{1}{h^2} f_2(ih-h) - \frac{2}{h} g_2(ih); \end{aligned}$$

at each point of S_h of the form $(ih, 1)$, $i = 0, 1, 2, \dots, n$, set

$$(3.9) \quad \begin{aligned} \bar{\omega}^{(1)}(ih, 1) &= \frac{4}{h^2} f_3(ih) - \frac{1}{h^2} f_3(ih+h) - \frac{1}{h^2} f_3(ih-h) \\ &\quad - \frac{2}{h^2} u(ih, 1-h) - \frac{2}{h} g_3(ih); \end{aligned}$$

and at each point of S_h of the form $(0, ih)$, $i = 1, 2, \dots, n-1$, set

$$(3.10) \quad \bar{w}(0, ih) = \frac{4}{h^2} f_4(ih) - \frac{2}{h^2} u^{(1)}(h, ih) - \frac{1}{h^2} f_4(ih+h) \\ - \frac{1}{h^2} f_4(ih-h) + \frac{2}{h} g_4(ih) .$$

Next, at each point of R_h write down the difference analogue

$$(3.11) \quad -4\omega(x, y) + \omega(x+h, y) + \omega(x, y+h) + \omega(x-h, y) + \omega(x, y-h) = 0$$

of (3.2), inserting wherever possible the boundary values obtained from (3.7)-(3.10). Solve the linear algebra system generated by (3.11) by SOR with over-relaxation factor r_2 and denote the solution by $\bar{w}^{(1)}$. Finally, on all of $R_h \cup S_h$, define $\omega^{(1)}$ by the smoothing formula

$$(3.12) \quad \omega^{(1)} = \mu \omega^{(0)} + (1 - \mu) \bar{w}^{(1)}, \quad 0 \leq \mu \leq 1.$$

Step 5.

Proceed next on R_h to determine $u^{(2)}$ from $\omega^{(1)}$ in the same fashion as $u^{(1)}$ was determined from $\omega^{(0)}$. Then on $R_h \cup S_h$ construct $\omega^{(2)}$ from $u^{(2)}$ just as $\omega^{(1)}$ was determined from $u^{(1)}$. In the indicated fashion construct sequences (3.3) and (3.4).

Step 6.

Given $\epsilon_1 > 0$ and $\epsilon_2 > 0$, terminate the iteration of Step 5 when both the following are valid:

$$(3.13) \quad |u^{(k)} - u^{(k+1)}| < \epsilon_1 \quad \text{uniformly on } R_h$$

$$(3.14) \quad |\omega^{(k)} - \omega^{(k+1)}| < \epsilon_2 \quad \text{uniformly on } R_h \cup S_h ,$$

and let $u^{(k+1)}$ be the approximate solution on R_h of boundary value problem (2.1)-(2.2).

4. Illustrative Example

Typical of the examples run on the UNIVAC 1108 is the following. For the boundary data $f_1 = x^3$, $f_2 = 2y + 1 - 3y^2$, $f_3 = x^3 + 2x - 3$, $f_4 = -3y^2$, $g_1 = 2x$, $g_2 = 2y + 3$, $g_3 = 2x - 6$, and $g_4 = 2y$, the method of Section 3 was executed with $h = 0.05$, $\rho = 0.2$, $u = 0.85$, $\epsilon_1 = 10^{-4}$, $\epsilon_2 = 10^{-3}$, $r_1 = 1.8$, $r_2 = 1.0$. Convergence criteria (3.13) and (3.14) were satisfied at $k = 93$ and the numerical solution agreed with the exact solution, $u = x^3 - 3y^2 + 2xy$, to three decimal places. The total running time was under 6 minutes. Deletion of smoothing, that is, setting $\rho = \mu = 0$, in this example resulted in divergence in three outer iterations.

5. Extension of the Numerical Method

The method of Section 3 generalizes easily to rectangular type regions in any number of dimensions [5]. Let us then proceed in a different direction and show how to extend the method to a difficult problem of applied interest in which S need not be a simple closed curve. The problem is that of determining an Airy stress function u when one has reentrant boundaries [2], and it can be formulated in general as follows.

Consider the rectangular type plate $O A B C D E F A$, shown in Figure 5.1, where O is the geometric center of the figure

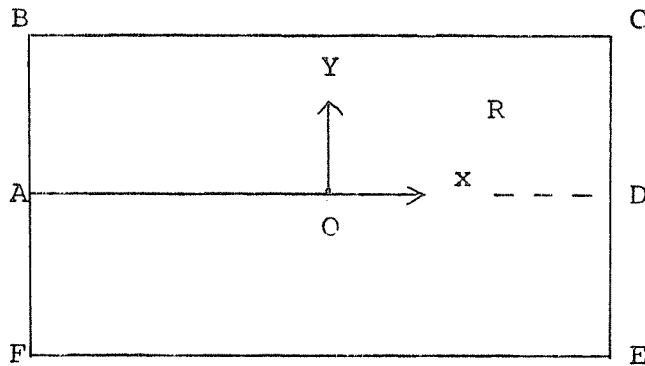


Figure 5.1

and where the boundary sections are parallel to the axes. For convenience, let S represent the connected segments $O A B C D E F A$, and let R be the finite open region bounded by S . Then if u and its normal derivatives are prescribed in S , one must find $u(x,y)$ which is continuous on $R \cup S$, satisfies the prescribed boundary conditions on S , and satisfies (2.1) on R .

The fundamental difficulty with the above problem is that the point O is a singular point in the sense that higher-order derivatives of u become unbounded in any neighborhood of O . The standard finite-difference methods still work, but the ultrasimplictic convergence proofs based on Taylor expansions are no longer valid. Also, one may have to exert special effort to obtain a high accuracy in a neighborhood of O . In [2], the form of such an effort was based in special series expansions developed by Williams [14]. Here, we shall replace five-point

difference equations by nine point approximations at selected points near O . Let us then consider a particular problem and describe the method precisely.

As in [2], consider a problem which is symmetric about AD , so that attention need be restricted only to $O A B C D O$. Set $AB = 0.7$, $BC = 0.8$, $OA = 0.4$ and fix the boundary conditions as

$$(5.1) \quad u = 0 \quad , \quad u_y = 0 \quad ; \quad \text{on } OA$$

$$(5.2) \quad u = 0 \quad , \quad u_x = 0 \quad ; \quad \text{on } AB$$

$$(5.3) \quad u = 5000x^2 \quad , \quad u_y = 0 \quad ; \quad \text{on } BC$$

$$(5.4) \quad u = 3200 \quad , \quad u_x = 8000 \quad ; \quad \text{on } CD .$$

For the boundary value problem defined by (2.1) and (5.1)-(5.4), the method of Section 3 was modified only in the following ways. At grid points $(h,0)$, (h,h) , $(0,h)$ and $(-h,h)$ about O , the nine point difference analogue of the Laplace operator [1] was used. At grid points between O and D which were different from $(h,0)$, the five point analogue was used. Symmetry of the solution about AD was incorporated into all formulas. For the parameter choices $h = 0.1$, $\rho = 0.2$, $\mu = 0.85$, $r_1 = 1.8$, $r_2 = 1.0$, $\epsilon_1 = 0.003$, and $\epsilon_2 = 2.0$, the iteration converged at $k = 251$ with a running time of two minutes. For $h = 0.05$, $\rho = 0.2$, $\mu = 0.85$, $r_1 = 1.8$, $r_2 = 1.0$, $\epsilon_1 = .001$, and $\epsilon_2 = 2.0$, the iteration converged at $k = 248$ with a running time of 5 minutes. For $h = \frac{1}{80}$, $\rho = 0.2$, $\mu = 0.85$, $r_1 = 1.8$, $r_2 = 1.0$,

$\epsilon_1 = .01$, $\epsilon_2 = 20$, and with the previous results with linearly interpolated values used for $u^{(0)}$ and $\omega^{(0)}$, the iteration converged at $k = 196$ with a running time of 15 minutes. Shown in the Table are typical results from the case $h = \frac{1}{80}$.

TABLE

x	y	u	x	y	u	x	y	u
0.1	0	481	0.2	0.2	1500	0.3	0.4	2417
0.2	0	1333	0.3	0.2	2351	-0.3	0.5	39
0.3	0	2299	-0.3	0.3	11	-0.2	0.5	162
-0.3	0.1	-4	-0.2	0.3	66	-0.1	0.5	383
-0.2	0.1	-5	-0.1	0.3	207	0	0.5	718
-0.1	0.1	11	0	0.3	488	0.1	0.5	1176
0	0.1	124	0.1	0.3	958	0.2	0.5	1755
0.1	0.1	590	0.2	0.3	1612	0.3	0.5	2436
0.2	0.1	1387	0.3	0.3	2389	-0.3	0.6	48
0.3	0.1	2315	-0.3	0.4	26	-0.2	0.6	190
-0.3	0.2	-1	-0.2	0.4	119	-0.1	0.6	432
-0.2	0.2	18	-0.1	0.4	307	0	0.6	777
-0.1	0.2	96	0	0.4	622	0.1	0.6	1230
0	0.2	316	0.1	0.4	1087	0.2	0.6	1788
0.1	0.2	785	0.2	0.4	1697	0.3	0.6	2447

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APPENDIX - RE-ENTRANT BOUNDARY FORTRAN PROGRAM

```

COMMON N,NPLUS1,M,MPLUS1
DIMENSION PSI(70,70),OMA(70,70),SVPSI(70,70),SVOMA(70,70),
1 SVOUT(70,70),VSS(70,70)
RW=1
F1=.2
C1=.85
REWIND 10
READ 3( ),N,M,H
300  FORMAT(2I2,F8.2)
MPLUS1=M+1
MMESH=M-1
NPLUS1=N+1
NMESH=N-1
MID=N/2+1
MID1=MID+1
MIDM=MID-1
H2=H*H
ISTOP=0
EPS=.2
EPSS=50.
C  INITIALIZE VECTORS
NZ=0
MIN=0
MP=19
R=0
104  CONTINUE
JM=0
KJ=0
PRINT 2323,C1,F1
2323  FORMAT(1H1,2F8.3)
ISTART =1
ISTART=0
233  FORMAT(7F10.6)
IF(ISTART .EQ. 1) GO TO 911
NE=0
NN=7
DO 330 JP=1,NPLUS1,7
LB=JP
NE=NE+NN
IF(NE .GT. NPLUS1) GO TO 331
GO TO 332
331  NE=NPLUS1
332  DO 330 J=1,MPLUS1
L=MPLUS1-J+1
READ(10,233)(PSI(I,L),I=LB,NE)
330  CONTINUE
NE=0
NN=7
DO 240 JP=1,NPLUS1,7
LB=JP
NE=NE+NN
IF(NE .GT. NPLUS1) GO TO 241
GO TO 242
241  NE=NPLUS1
242  DO 240 J=1,MPLUS1
L=MPLUS1-J+1

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      READ(10,233)(OMA(I,L),I=LB,NE)
240  CONTINUE
      PRINT 894,(OMA(I,3),I=31,39)
      PRINT 894,(OMA(I,2),I=31,39)
894  FORMAT(1X,9F11.2)
      PRINT 91
911  CONTINUE
      DO 2014 I=MID,N
2014  OMA(I,1)=OMA(I,3)
      CALL PRNTLT(PSI)
      CALL PRNTLT(OMA)
      REWIND 10
      DO 1001 I=2,N
      X=(I-1)*H
      X2=X*X
1001  PSI(I,MPLUS1)=10000.*X2/2.
      DO 1002 J=1,MPLUS1
1002  PSI(NPLUS1,J)=10000.*.8*.8/2.
      NM1=0
      NM2=0
      NM=0
      C2=1-C1
      F2=1-F1
C     BEGIN LOOP FOR OUTER ITERATIONS
23   DO 40 I=1,NPLUS1
      DO 40 J=1,MPLUS1
      VSS(I,J)=PSI(I,J)
40   SVOUT(I,J)=OMA(I,J)
      NM=NM+1
      NCOUNT=0
C     BEGIN INNER ITERATION FOR STREAM FUNCTION
11   DO 2 I=2,N
      DO 2 J=2,M
      IF(I.LE.MID.AND.J.EQ.2) GO TO 2
      SVPSI(I,J)=PSI(I,J)
      IF(I.EQ.MIDI.AND.J.EQ.3) GO TO 1428
      IF(I.EQ.MIDM.AND.J.EQ.3) GO TO 1428
      IF(I.EQ.MIDI.AND.J.EQ.2) GO TO 1428
      IF(I.EQ.MID.AND.J.EQ.3) GO TO 1428
      PSI(I,J)=(-.8*PSI(I,J))+.45*(PSI(I,J-1)+PSI(I,J+1)+PSI(I-1,J)+
1PSI(I+1,J)+H2*OMA(I,J))
      GO TO 2
1428 PSI(I,J)=(-.9*PSI(I,J))+.09*(4*PSI(I+1,J)+4*PSI(I,J+1)+4*PSI(I-1,J)
1)+4*PSI(I,J-1)+PSI(I+1,J+1)+PSI(I-1,J+1)+PSI(I-1,J-1)+PSI(I+1,J-1)
2+H2*OMA(I,J))
2   CONTINUE
      NM1=NM1+1
      DO 1003 I=MID,N
1003  PSI(I,1)=PSI(I,3)
C     TEST STREAM FUNCTION FOR CONVERGENCE
      DO 5 I=2,N
      DO 5 J=2,M
      DIFF=ABS(SVPSI(I,J)-PSI(I,J))
      IF(DIFF.GT.EPSS) GO TO 6
5   CONTINUE

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C      RECALCULATE STREAM FUNCTION USING WEIGHTING
      DO 114 I=2,N
      DO 114 J=2,M
114    PSI(I,J)=F1*VSS(I,J)+F2*PSI(I,J)
      DO 1114 I=MID,N
1114   PSI(I,1)=F1*VSS(I,1)+F2*PSI(I,1)
      GO TO 200
6      NCOUNT=NCOUNT+1
      IF(NCOUNT .GT. 100) GO TO 8
      GO TO 11
C      TEST STREAM FUNCTION FOR DIVERGENCE
8      IF(DIFF .GT. 10) GO TO 28
      PRINT 73
93     FORMAT(1H1,11H PSI VALUES)
      CALL PRNLT(PSI)
10     FORMAT(10F11.6)
      NCOUNT=0
      GO TO 11
28     PRINT 81
81     FORMAT(13H PSI DIVERGED)
      CALL PRNLT(PSI)
      CALL PRNLT(OMA)
      GO TO 699
C      BEGIN INNER ITERATION FOR VORTICITY
200    NCOUNT=0
30     HCONST=C2*(-2./H2)
      DO 12 I=1,MIDM
12     OMA(I,2)=C1*OMA(I,2)+HCONST*PSI(I,3)
      OMA(MID,2)=C1*OMA(MID,2)+HCONST*(PSI(MID,3)+.5*PSI(MID1,2))
      DO 13 J=2,MPLUS1
13     OMA(I,J)=C1*OMA(I,J)+HCONST*PSI(2,J)
      DO 1012 I=2,N
1012   OMA(I,MPLUS1)=C1*OMA(I,MPLUS1)+ -C2*(-(4./H2)*PSI(I,MPLUS1)+(
11./H2)*PSI(I+1,MPLUS1)+(1./H2)*PSI(I-1,MPLUS1)+(2./H2)*PSI(I,M))
      DO 1013 J=2,M
1013   OMA(NPLUS1,J)=C1*OMA(NPLUS1,J)+ -C2*(-(4./H2)*PSI(NPLUS1,J)
14./H2)*PSI(NPLUS1,J+1)+(2./H2)*PSI(N,J)+(1./H2)*PSI(NPLUS1,J-1)+
15./H2)*PSI(NPLUS1,J-1)+2(2./H)*10000.*.8)
90     CONTINUE
      NM2=NM2+1
      DO 14 I=2,N
      DO 14 J=2,M
      IF(I .LE. MID .AND. J .EQ. 2) GO TO 14
      A1=PSI(I+1,J)-PSI(I-1,J)
      B1=PSI(I,J+1)-PSI(I,J-1)
      A=ABS(A1)
      B=ABS(B1)
      W0=4+(A+B)*(R/2)
      IF(A1 .GE. 0) GO TO 15
      GO TO 16
15     W2=1+(R/2)*A
      W4=1
      GO TO 17
16     W2=1
      W4=1+A*(R/2)
17     IF(B1 .GE. 0) GO TO 18

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18      GO TO 19
        W1=1
        W3=1+B*(R/2)
        GO TO 19
19      W1=1+B*(R/2)
        W3=1
20      IF(ISTOP .EQ. 1) GO TO 425
        SVOMA(I,J)=OMA(I,J)
        GO TO 427
425     WM=ABS(SVOMA(I,J)-OMA(I,J))
        IF(WM .LT. WMIN) GO TO 426
        GO TO 427
426     WMIN=WM
427     CONTINUE
        IF(ISTOP .EQ. 1) GO TO 305
        IF(I .EQ. MID .AND. J .EQ. 3) GO TO 1429
        IF(I .EQ. MID1 .AND. J .EQ. 2) GO TO 1429
        IF(I .EQ. MIDM .AND. J .EQ. 3) GO TO 1429
        IF(I .EQ. MID1 .AND. J .EQ. 3) GO TO 1429
        OMA(I,J)=((W1/W0)*OMA(I+1,J)+(W2/W0)*OMA(I,J+1)+(W3/W0)*OMA(I-1,J)
1+((W4/W0)*OMA(I,J-1)))*RW+(1-RW)*OMA(I,J)
        GO TO 14
1429    OMA(I,J)=(1-RW)*OMA(I,J)+(RW/20.)*(4*OMA(I+1,J)+4*OMA(I,J+1)
1+4*OMA(I-1,J)+4*OMA(I,J-1)+OMA(I+1,J+1)+OMA(I-1,J+1)+OMA(I-1,J-1)
2+OMA(I+1,J-1))
        GO TO 14
305     DIFF=((W1/W0)*OMA(I+1,J)+(W2/W0)*OMA(I,J+1)+(W3/W0)*OMA(I-1,J)
1+((W4/W0)*OMA(I,J-1))-OMA(I,J))
        DIF=ABS(DIFF)
        IF(DIF .GT. EPS1) GO TO 282
        GO TO 421
282     EPS1=DIF
421     IF(DIF .LT. EPS2) GO TO 420
        GO TO 14
420     EPS2=DIF
14      CONTINUE
        DO 1014 I=MID,N
1014    OMA(I,1)=OMA(I,3)
        IF(ISTOP .EQ. 1) GO TO 700
C       TEST VORTICITY FOR CONVERGENCE
        DO 21 I=2,N
        DO 21 J=2,M
        IF(I .EQ. MID .AND. J .EQ. 2) GO TO 21
        DIFF=ABS(SVOMA(I,J)-OMA(I,J))
        IF(DIFF .GE. EPS) GO TO 22
21      CONTINUE
C       RECALCULATE VORTICITY USING WEIGHTING
        DO 414 I=2,N
        DO 414 J=2,M
414     OMA(I,J)=C1*SVOUT(I,J)+C2*OMA(I,J)
        DO 1414 I=MID,N
1414    OMA(I,1)=C1*SVOUT(I,1)+C2*OMA(I,1)
        PRINT 1016,NM,NM1,NM2
1016    FORMAT(1X,3I6)
        JM=JM+1

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IF(JM .EQ. 30) GO TO 89
GO TO 59
89   JM=0
MIN=MIN+20
PRINT 79,NM
79   FORMAT(1H1,I3,17H OUTER ITERATIONS)
PRINT 91
CALL PRNTLT(Psi)
PRINT 92
CALL PRNTLT(OMA)
REWIND 10
NZ=0
C   TEST OUTER ITERATIONS FOR CONVERGENCE
59   CONTINUE
DO 45 I=2,N
DO 45 J=2,M
DIFF=ABS(SVOUT(I,J)-OMA(I,J))
IF(DIFF .GT. EPS) GO TO 7
45   CONTINUE
PRINT 99,NM
99   FORMAT(1H1,22H PROBLEM CONVERGED IN ,I4)
PRINT 91
91   FORMAT(1X,11H PSI VALUES)
CALL PRNTLT(Psi)
PRINT 92
92   FORMAT(1H1,14H OMEGA VALUES)
CALL PRNTLT(OMA)
EPS1=0
RMAX=0
EPS2=1
WMIN=1
RMIN=1
ISTOP=1
DO 181 II=2,N
DO 181 JJ=2,M
IF(I .LE. MID .AND. J .EQ. 2) GO TO 181
RES=ABS(Psi(II,JJ)-VSS(II,JJ))
IF(RES .GT. RMAX) GO TO 301
GO TO 302
301  RMAX=RES
302  CONTINUE
IF(RES .LT. RMIN) GO TO 422
GO TO 320
422  RMIN=RES
320  CONTINUE
A=4*Psi(II,JJ)+Psi(II+1,JJ)+Psi(II,JJ+1)+Psi(II-1,JJ)+Psi(II,JJ-
11)
B=-H*H*OMA(II,JJ)
D=ABS(A-B)
IF(D .GT. EPS1) GO TO 182
GO TO 183
182  EPS1=D
183  IF(D .LT. EPS2) GO TO 184
GO TO 181
184  EPS2=D
181  CONTINUE

```

```

      PMAX=EPS1
      EPS1=0
      PMIN=EPS2
      EPS2=1
      GO TO 90
C     TEST OUTER ITERATIONS FOR DIVERGENCE
7     IF(DIFF .GT. 9900) GO TO 199
      GO TO 23
22    NCOUNT=NCOUNT+1
      IF(NCOUNT .GT. 300) GO TO 24
      GO TO 90
C     TEST VORTICITY FOR DIVERGENCE
24    IF(DIFF .GT. 999) GO TO 29
      PRINT 94
94    FORMAT(1H1,14H OMAEGA VALUES)
      CALL PRNLT(OMA)
      PRINT 91
      CALL PRNLT(PSI)
32    FORMAT(10F11.6)
      NCOUNT=0
      GO TO 90
29    PRINT 82
82    FORMAT(13H OMA DIVERGED)
      CALL PRNLT(PSI)
      CALL PRNLT(OMA)
      GO TO 699
199   PRINT 189
189   FORMAT(26H OUTER ITERATIONS DIVERGED)
700   CONTINUE
      PRINT 500,RMAX,RMIN
      PRINT 501,EPS,WMIN
500   FORMAT(2X,8H PSIMAX=,E12.4,8H PSIMIN=,E12.4)
501   FORMAT(2X,8H OMAMAX=,E12.4,8H OMAMIN=,E12.4)
      PRINT 502,PMAX,PMIN
502   FORMAT(2X,8HRPSIMAX=,E12.4,8HRPSIMIN=,E12.4)
      PRINT 503,EPS2,EPS1
503   FORMAT(2X,8HROMAMAX=,E12.4,8HROMAMIN=,E12.4)
699   CONTINUE
      SUBROUTINE PRNLT(Z)
      COMMON N,NPLUS1,M,MPLUS1
      DIMENSION Z(70,70)
      IF(MIN .GT. MP) GO TO 5
      GO TO 6
5     NZ=NZ+2
      IF(NZ .GT. 6) GO TO 6
      NE=0
      NN=7
      DO 81 JP=1,NPLUS1,7
      LB=JP
      NE=NE+NN
      IF(NE .GT. NPLUS1) GO TO 10
      GO TO 12
10    NE=NPLUS1
12    DO 81 J=1,MPLUS1
      L=MPLUS1-J+1

```

```
81  WRITE(10,82)(Z(I,L),I=LB,NE)
82  FORMAT(7F10.1)
6   CONTINUE
    IF(N .GT. 11) GO TO 103
75  DO 61 J=1,MPLUS1
    L=MPLUS1-J+1
61  PRINT 52,(Z(I,L),I=1,NPLUS1)
    RETURN
103  NE=0
    NN=11
    DO 51 IP=1,NPLUS1,11
    IF( IP .GT. 1) GO TO 1
    GO TO 2
1   PRINT 3
3   FORMAT(///)
2   NB=IP
    NE=NE+NN
    IF(NE .GT. NPLUS1) GO TO 101
    GO TO 102
101  NE=NPLUS1
102  DO 51 J=1,MPLUS1
    L=MPLUS1-J+1
51  PRINT 52,(Z(I,L),I=NB,NE)
52  FORMAT(1X,11F10.3)
    RETURN
    END
```