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# Fatigue Crack Propagation in Microcapsule Toughened Epoxy\*

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## Abstract

The addition of liquid-filled urea-formaldehyde (UF) microcapsules to an epoxy matrix leads to significant reduction in fatigue crack growth rate and corresponding increase in fatigue life. Mode-I fatigue crack propagation is measured using a tapered double-cantilever beam (TDCB) specimen for a range of microcapsule concentrations and sizes: 0, 5, 10, and 20% by weight and 50, 180, and 460  $\mu\text{m}$  diameter. Cyclic crack growth in both the neat epoxy and epoxy filled with microcapsules obeys the Paris power law. Above a transition value of the applied stress intensity factor  $\Delta K_T$ , which corresponds to loading conditions where the size of the plastic zone approaches the size of the embedded microcapsules, the Paris law exponent decreases with increasing content of microcapsules, ranging from 9.7 for neat epoxy to approximately 4.5 for concentrations above 10 wt% microcapsules. Improved resistance to fatigue crack propagation, indicated by both the decreased crack growth rates and increased cyclic stress intensity for the onset of unstable fatigue-crack growth, is attributed to toughening mechanisms induced by the embedded microcapsules as well as crack shielding due to the release of fluid as the capsules are ruptured. In addition to increasing the inherent fatigue life of epoxy, embedded microcapsules filled with an appropriate healing agent provide a potential mechanism for self-healing of fatigue damage.

*Keywords: fatigue-crack propagation, microcapsule toughening, tapered double-cantilevered beam, brittle fracture of epoxy*

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## 1. Introduction

Highly crosslinked epoxy resins have low strain-to-failure and exhibit poor resistance to crack propagation. Fatigue loading is particularly problematic, causing small cracks to initiate and grow rapidly. These cracks often lead to catastrophic failure. An extensive body of work exists for the general area of fatigue of polymers [1–3], which focuses on understanding the mechanisms of fatigue and predicting the rates of fatigue-crack growth.

Fatigue crack propagation studies are performed with the cyclic-crack-tip stress state varying over a range defined by  $\Delta K_I \equiv (K_{\max} - K_{\min})$ . Dependence of the fatigue-crack-growth rate  $da/dN$  on the applied range of stress intensity factors  $\Delta K_I$  is generally described by the empirical Paris law equation [4]

$$\frac{da}{dN} = C_0 \Delta K_I^n, \quad (1)$$

where  $C_0$  and  $n$  are material constants that depend on the ratio of applied stress intensity  $R \equiv K_{\min}/K_{\max}$ , the loading frequency  $f$ , and the testing environment. The typical crack growth behavior described by Eq. (1) yields a linear log–log plot that is bounded by a threshold stress intensity range  $\Delta K_{th}$  below which a crack ceases to propagate, and the critical stress intensity  $K_{IC}$  above which crack growth is unstable.

Several researchers [5–7] have successfully measured fatigue-crack propagation in epoxy resins and obtained values of the Paris law exponent  $n$  on the order of 10. Incorporation of either a rubbery second phase [8–11] or solid particles [7,12–13] significantly improves the resistance to fatigue-crack propagation. Several of these studies [5,8–9,11,13–14] suggest that improvements in the resistance to fatigue crack propagation behavior are also associated with increased toughness in monotonic fracture [1–2,14–15].

Previously, we investigated the effect of embedded urea-formaldehyde (UF) microcapsules on the monotonic fracture properties of a self-healing epoxy [15]. In addition to providing an efficient mechanism for self-healing [16–18], the presence of liquid-filled microcapsules increased the virgin monotonic-fracture toughness of epoxy by up to 127% [15,17]. The increased toughening was correlated with a change in the fracture plane morphology from mirror-like to hackle markings with subsurface microcracking. The inherent fracture toughness as well as the healing efficiency both depended strongly on the size and concentration of microcapsules. In the current work, we extend this investigation to examine the influence of microcapsules on the fatigue crack propagation behavior of epoxy, with the effects of self-healing precluded. Consistent with the monotonic fracture studies, the addition of microcapsules to an epoxy matrix significantly increased the resistance to crack growth under dynamic loading conditions.

## 2. Experimental procedure

### 2.1. Materials and sample preparation

Urea-formaldehyde microcapsules containing dicyclopentadiene (DCPD) monomer were manufactured with average diameters of 50, 180, and 460  $\mu\text{m}$  using the emulsion *in situ* polymerization microencapsulation method outlined by Brown *et al.* [19]. Shell wall thickness was  $190 \pm 30$  nm for all batches. Tapered double-cantilever beam specimens

were cast from EPON<sup>®</sup> 828 epoxy resin (DGEBA) and 12 pph Ancamine<sup>®</sup> DETA (diethylenetriamine) curing agent with a prescribed concentration of microcapsules mixed into the resin. The epoxy mixture was degassed, poured into a closed silicone rubber mold and cured for 24 hours at room temperature, followed by 24 hours at 30° C. Relevant physical and mechanical properties of the microcapsules and neat epoxy are listed in Table 1. The tensile modulus and mode I critical stress intensity factor,  $K_{IC}$ , of the microcapsule toughened epoxy were measured as a function of capsule concentration by Brown *et al.* [15, 20] and Rzeszutko *et al.* [21] and summarized in Table 2.

TABLE 1 Properties of the constituent materials [15]

Properties	Epoxy	Urea-formaldehyde microcapsules
Density (kg/m <sup>3</sup> )	1160	~1000
Diameter (μm)	—	50±20 180±40 460±80
Wall thickness (nm)	—	190±30
$K_{IC}$ (MPa m <sup>1/2</sup> )	0.55±0.04	—
Young's modulus (GPa)	3.4±0.1	—
Ultimate stress (MPa)	39±4	—

TABLE 2 Mechanical properties of neat epoxy and epoxy with embedded microcapsules [15]

Microcapsule concentration (wt%)	Diameter, $d$ (μm)	Young's modulus, $E$ (GPa)	Critical stress intensity factor, $K_{IC}$ (MPa m <sup>1/2</sup> )	Critical strain energy release rate, $G_C$ (J/m <sup>2</sup> )
0	—	3.4±0.1	0.55±0.04	88±14
5	50±20	3.2±0.1	1.1±0.1	350±70
10	50±20	—	1.2±0.2	—
20	50±20	—	1.1±0.1	—
5	180±40	3.2±0.1	0.78±0.16	190±90
10	180±40	2.8±0.1	1.2±0.2	430±170
20	180±40	2.7±0.1	1.0±0.2	400±130
10	460±80	—	0.92±0.07	—
20	460±80	—	1.2±0.1	—

## 2.2 Mechanical testing

The fatigue-crack propagation behavior of the microcapsule-modified epoxy was investigated using the tapered double-cantilever beam (TDCB) specimen shown in Fig. 1. Side grooves ensured controlled crack growth along the centerline of the brittle specimen. The TDCB geometry, developed by Mostovoy *et al.* [22], provided a crack-length-independent relationship between applied stress intensity factor  $K_I$  and load  $P$ ,

$$K_I = \alpha P, \quad (2)$$

which only required knowledge of the coefficient  $\alpha$ . For the TDCB sample geometry in Fig. 1,  $\alpha = 11.2 \times 10^3 \text{ m}^{-3/2}$  was determined experimentally [17]. A constant range of Mode-I stress intensity factor  $\Delta K_I$  was achieved by applying a constant range of load  $\Delta P$ , independent of crack length. The constant- $K$  region of the TDCB specimen enables crack-growth-rate measurements over a range of cycles, rather than requiring use of the modified secant formulation commonly employed for changing  $\Delta K_I$  of a compact tension specimen [14]. Moreover, the constant- $K$  region is of great importance for observing the time-dependent effects of self-healing during growth of a fatigue-crack [20].

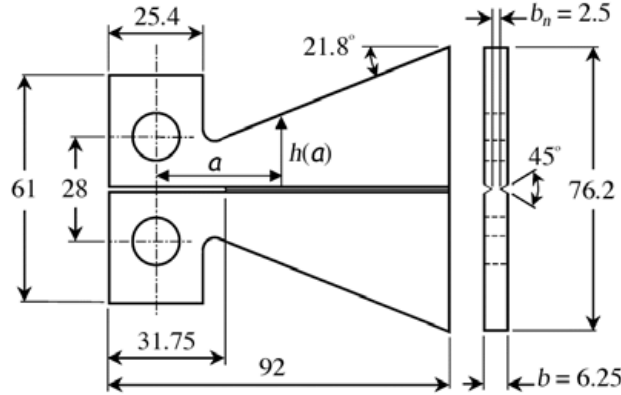


Figure 1. Tapered-double-cantilever-beam geometry [17]. All dimensions in mm.

Fatigue crack propagation studies were performed using an Instron DynoMight 8841 low-load frame with 250 N load-cell. Samples were precracked with a razor blade while ensuring the precrack tip was centered in the groove and then pin loaded. A triangular frequency of 5 Hz was applied with a load ratio ( $R = K_{\min}/K_{\max}$ ) of 0.1. Crack lengths were determined by optical measurements at finite times and by compliance-inferred measurements [23] acquired approximately every 256th cycle.

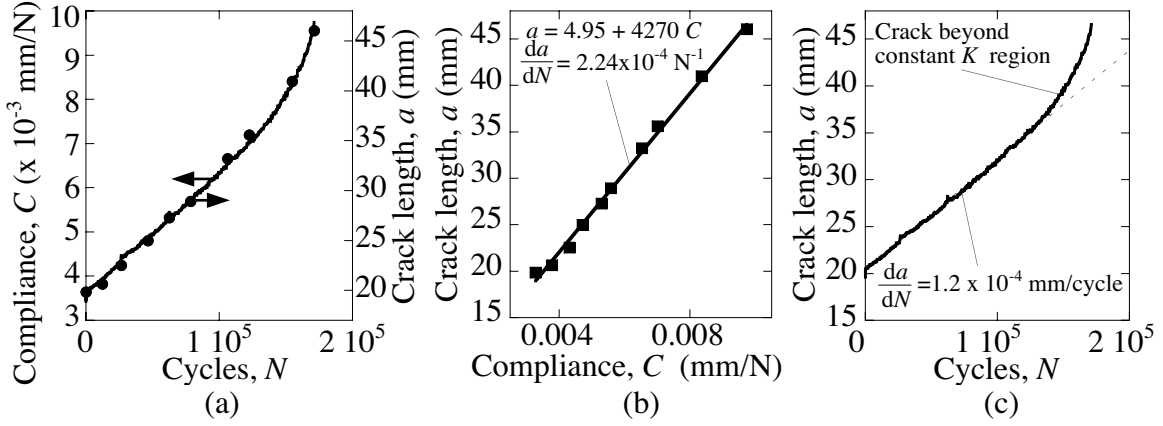
The optically measured crack-tip position and specimen compliance are plotted against number of cycles in Fig. 2a. The linear relationship between optically measured crack length and specimen compliance (Fig. 2b) is used to calculate the crack-tip position at all times during the experiment (Fig. 2c). Crack-growth data are generated under constant  $\Delta P$  (*i.e.* constant  $\Delta K_I$ ), with a complete set of loading conditions measured on a single specimen by incrementally increasing  $\Delta P$ . Crack-growth rates are obtained from the number of cycles  $N$  required to grow a crack a distance  $\Delta a$  of approximately 1 mm for a given range of Mode-I stress intensity factor  $\Delta K_I$ . Statistically equivalent values of crack-growth rate are obtain from using either the optically measured crack length prior to,  $a_{i=0}^{\text{opt}}$ , and following,  $a_{i=S}^{\text{opt}}$ , the application of  $N$  cycles at a give  $\Delta K_I$ ,

$$\left. \frac{da}{dN} \right|_{\Delta K} = \frac{\Delta a}{\Delta N} = \frac{a_{i=S}^{\text{opt}} - a_{i=0}^{\text{opt}}}{N_{i=S}}, \quad (3)$$

or from using a linear fit of the compliance-inferred measurements  $a_i^{\text{comp}}$ ,

$$\frac{da}{dN}_{\Delta K} = \frac{S \sum_{i=1}^S N_i a_i^{\text{comp}} - \sum_{i=1}^S N_i \sum_{i=1}^S a_i^{\text{comp}}}{S \sum_{i=1}^S N_i^2 - \left( \sum_{i=1}^S N_i \right)^2}, \quad (4)$$

where  $S$  is the total number of cycle samples acquired.



*Figure 2.* Plots illustrating method to calculate the continuous crack-tip position from compliance data (—) and finite optical measurements (●) (a). (b) Linear fit of optically measured crack vs. specimen compliance. The squares (■) represent compliance values corresponding to the optical data and the line (—) represents the linear best fit. (c) Crack length calculated from measured compliance using the relationship obtained from (b) plotted vs. number of cycles. The specimen is neat epoxy; the test parameters are  $f = 5$  Hz,  $R = 0.1$ , and  $\Delta K_I = 0.472 \text{ MPa m}^{1/2}$ .

Fracture surface morphologies of the fatigue samples were examined with an environmental scanning electron microscope (Philips XL30 ESEM-FEG). After failure, specimens were mounted and sputtered with gold/palladium. Micrographs were obtained using 10kV secondary electrons in high vacuum mode.

### 3. Results

The effect of embedded microcapsule concentration on fatigue crack growth is shown in Fig. 3 for  $180 \mu\text{m}$  diameter microcapsules. The relationship between the crack growth rate  $da/dN$  of epoxy and the applied range of Mode-I stress intensity factor  $\Delta K_I$  clearly follows the Paris power law (Eq. 1). The measured Paris exponent is  $n = 9.7$  for crack propagation in neat EPON<sup>®</sup> 828–Ancamine<sup>®</sup> DETA (no microcapsules). At applied stress intensity factors greater than  $\Delta K \sim 0.35$  to  $0.4 \text{ MPa m}^{1/2}$ , a distinct transition in crack growth is observed for the microcapsule filled epoxy. Above a transition value  $\Delta K_T$ , epoxy with microcapsules exhibits a higher resistance to fatigue crack growth than neat epoxy, accompanied by a reduction of the Paris law exponent  $n$ . At applied load levels below the transition point, the microcapsules have little influence on the crack growth rate.

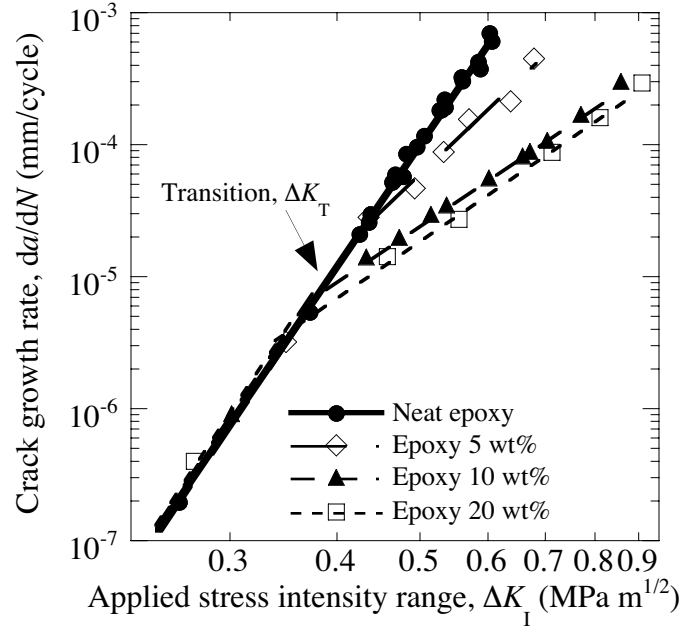


Figure 3. Influence of microcapsule concentration on the fatigue crack growth behavior for 180  $\mu\text{m}$  diameter microcapsules.

Measured Paris law parameters for epoxy with embedded 50, 180, and 460  $\mu\text{m}$  diameter microcapsules are summarized in Table 3. The Paris law exponent,  $n$  is plotted as a function of microcapsule concentration in Fig. 4. Above  $\Delta K_T$ , the value of  $n$  decreases significantly with increasing microcapsule concentration, independent of capsule diameter. For concentrations greater than 10 wt% microcapsules,  $n$  has a steady-state value of approximately 4.5. Others have observed similar behavior for rubber-modified epoxies [14, 24].

TABLE 3 Constants of the Paris power law,  $\Delta K_I^{\text{max}}$ ,  $\Delta K_T$ , and  $r_y$

Microcapsule concentration (wt%)	Diameter ( $\mu\text{m}$ )	$C_0$ (above $\Delta K_T$ )	$n$	$\Delta K_I^{\text{max}}$ ( $\text{MPa m}^{1/2}$ )	$\Delta K_T$ ( $\text{MPa m}^{1/2}$ )
0	—	$8.2 \times 10^{-2}$	9.7	0.60	—
10	$50 \pm 20$	$1.5 \times 10^{-3}$	4.9	0.82	0.41
20	$50 \pm 20$	$1.6 \times 10^{-3}$	4.6	0.81	0.44
5	$180 \pm 40$	$4.2 \times 10^{-3}$	6.1	0.64	0.46
10	$180 \pm 40$	$5.4 \times 10^{-4}$	4.4	0.82	0.40
20	$180 \pm 40$	$3.8 \times 10^{-4}$	4.3	0.80	0.36
10	$460 \pm 80$	$7.8 \times 10^{-4}$	4.4	0.64	0.40
20	$460 \pm 80$	$8.6 \times 10^{-4}$	4.7	0.82	0.39

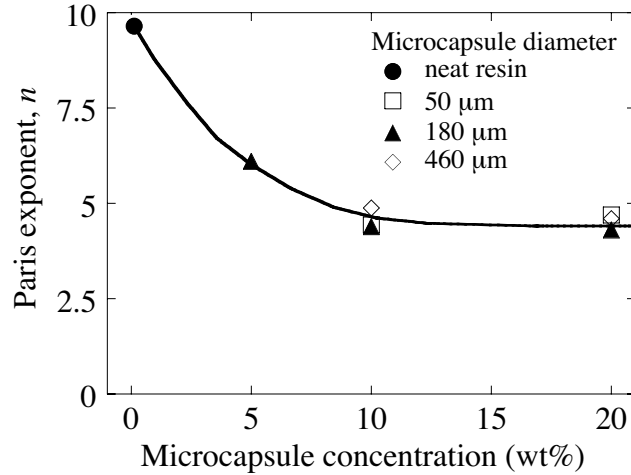


Figure 4. Influence of microcapsule concentration on Paris law exponent,  $n$ .

The effect of the microcapsules on life extension is shown more clearly in Fig. 5 by comparison of crack length as a function of loading cycles for neat epoxy and epoxy with 20 wt% microcapsules above  $\Delta K_T$ . The addition of microcapsules significantly increases fatigue life; for the loading condition of  $\Delta K_1 = 0.586 \text{ MPa m}^{1/2}$  the fatigue life increases from  $86 \times 10^3$  cycles for neat epoxy to  $239 \times 10^3$  for the microcapsule filled system.

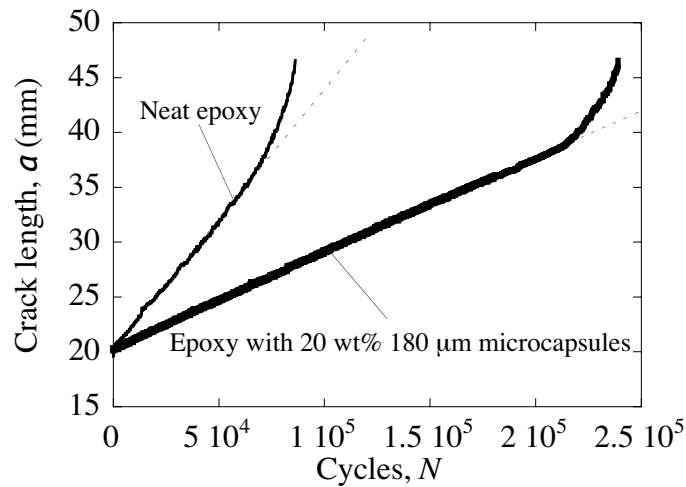
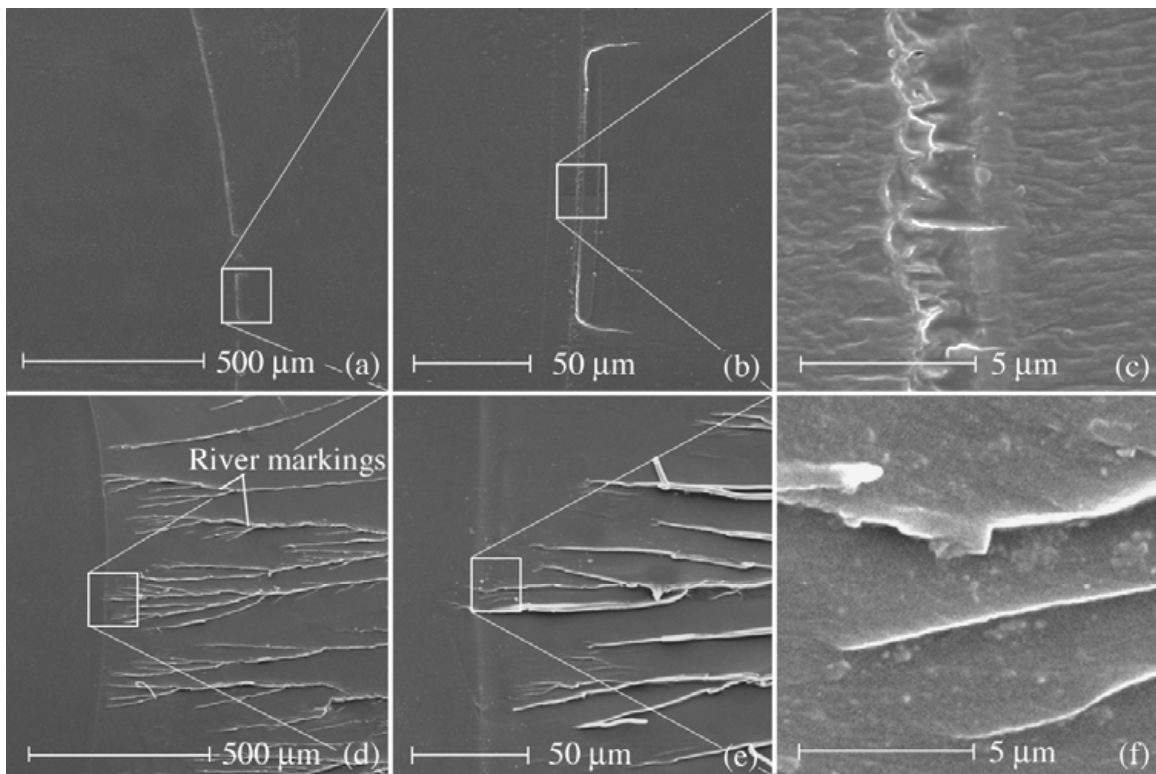


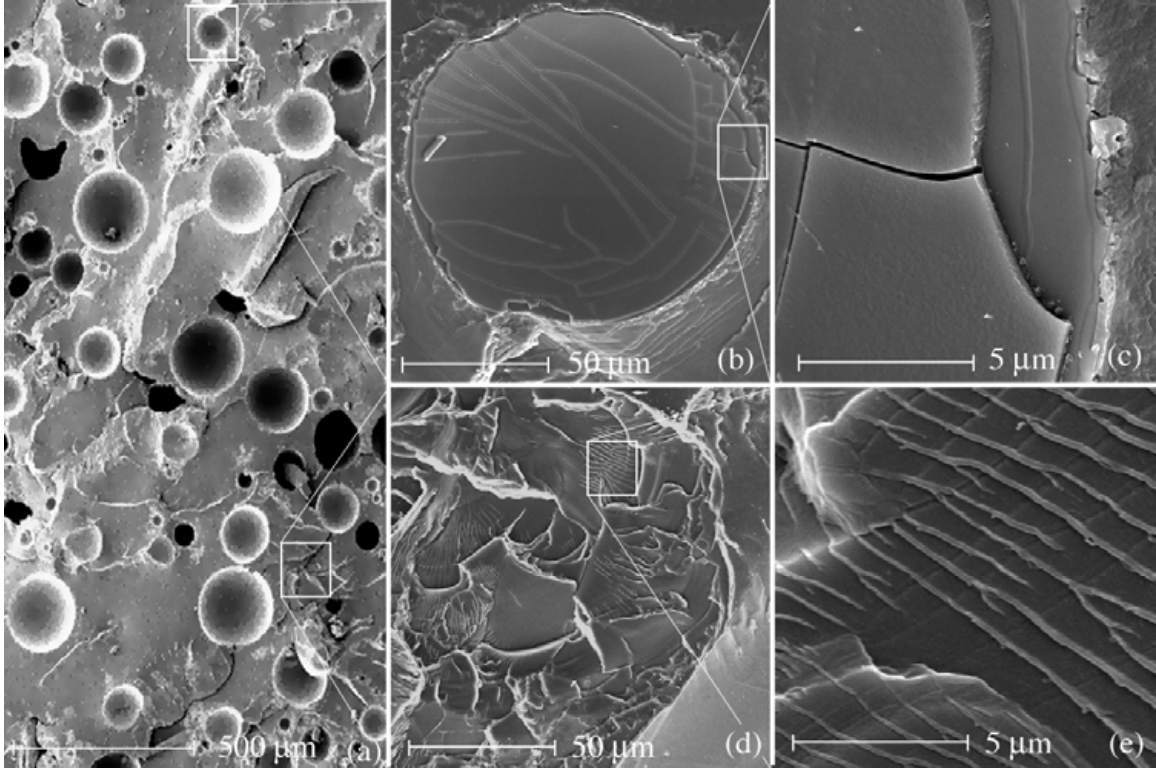
Figure 5. Fatigue crack extension vs. number of cycles for neat epoxy and epoxy with 20 wt% 180  $\mu\text{m}$  microcapsules. The test parameters are  $f = 5 \text{ Hz}$ ,  $R = 0.1$ , and  $\Delta K_1 = 0.586 \text{ MPa m}^{1/2}$ .



Examination of the fatigue-fracture surface for neat epoxy reveals a relatively featureless morphology (Fig. 6a–c). River markings [25] are observed on some fracture surfaces, but are only initiated at locations corresponding to the crack-tip position when  $\Delta K_I$  is incrementally increased (Fig. 6d–f). In contrast, the fatigue-fracture surface for epoxy with embedded microcapsules is characterized by substantial out-of-plane morphology at all length scales (Fig. 7). At the largest length scale, the microcapsules are ruptured at the fracture plane (similar to monotonic fracture), the fatigue-crack-growth path has significant deviation in and out of the plane, and the crack branches out of the plane at several locations (Fig. 7a). At smaller length scales the microcapsule walls have numerous secondary cracks (Fig. 7b,c) and the matrix has additional contortion and river markings at decreasing length scales (Fig. 7d,e).



*Figure 6.* Fatigue-fracture surfaces for neat epoxy. (a–c) The dominant surface morphology is featureless at different length scales and (d–f) has some local river markings. Note: The crack propagation is from left to right in all images.



*Figure 7.* Fatigue-fracture surfaces for epoxy with 20 wt% of 180  $\mu\text{m}$  microcapsules. (a) At large length scales the fracture plane has a contorted crack-growth path and ruptured microcapsules. At smaller length scales (b,c) the microcapsules exhibit cracking and (d,e) the matrix has complex morphology. Note: The crack propagation is from left to right in all images.

#### 4. Discussion

The addition of microcapsules significantly improves the fatigue response above a transition value of the applied stress intensity factor,  $\Delta K_T$ . Similar improvements reported in the literature for rubber modified epoxies have been explained by effects of plastic zone size [11,14] and increases of monotonic fracture toughness [5,8–9,11,13–14]. Azimi *et al.* [12,14] observed a transition point in the fatigue crack propagation behavior of DGEBA epoxies modified with CTBN and silicon rubber. Above a threshold value, the rubber-modified epoxy exhibited improved resistance to fatigue crack growth. Below the threshold, both neat epoxy and the rubber-modified epoxy had similar resistance. The transition was attributed to interactions between the rubber particles and the plastic zone present at the crack tip. Moreover, Azimi showed that the transition phenomenon was triggered when the size of the theoretical crack-tip plastic zone was of the order of the size of the filler.

Applying Azimi's hypothesis [12,14] to microcapsule-toughened epoxy, the theoretical plastic zone size at a cyclically loaded crack tip is estimated based on Irwin's [26] formula for the size of a plastic zone,

$$r_y(\Delta K_T) = \frac{1}{2\pi} \left( \frac{\Delta K_T / (1-R)}{\sigma_{YS}} \right)^2, \quad (5)$$

where  $\sigma_{YS}$  is the yield stress of the epoxy with embedded microcapsules. In previous work, Irwin's formula accurately captured the plastic zone size in EPON<sup>®</sup> 828–Ancamine<sup>®</sup> DETA under monotonic fracture conditions [15]. The plastic zone size estimated by Eq. (5) for 5, 10, and 20 wt% of 180  $\mu\text{m}$  diameter microcapsules is c.a. 70  $\mu\text{m}$  (assuming  $\sigma_{YS} = 20$  MPa [15]). The calculated plastic zone size at the transition point for samples with 180  $\mu\text{m}$  is within an order of magnitude of the microcapsule diameter consistent with Azimi's findings for 2–3  $\mu\text{m}$  diameter rubber particles. When  $\Delta K_I$  is less than  $\Delta K_T$ , the plastic zone is smaller than the embedded microcapsules. In this regime the plastic zone is unaffected by the microcapsules—other than the presence of discontinuities along the crack front—resulting in little or no toughening and a fatigue response similar to that of the neat epoxy. When  $\Delta K_I$  is greater than  $\Delta K_T$ , the plastic zone is sufficiently large to encompass the majority of microcapsules in its path, activating toughening mechanisms similar to those observed for monotonic fracture.

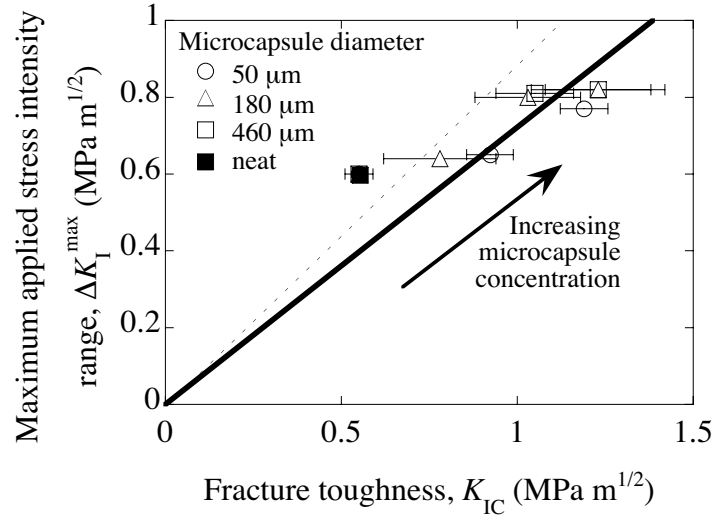
The onset of fatigue crack growth instability is generally considered to be equivalent to monotonic fracture. Hence, the maximum stress intensity value during the cycle when the fatigue crack propagation becomes unstable should correlate with the monotonic fracture toughness,

$$K_{\max}(\Delta K_I^{\max}) \equiv \frac{\Delta K_I^{\max}}{1-R} = K_{IC} \quad (6)$$

In Fig. 8,  $\Delta K_I^{\max}$  values from Table 3 are plotted against the corresponding  $K_{IC}$  values from Table 2 for the range of microcapsule sizes and concentrations tested. With the exception of neat epoxy, the data for microcapsule filled epoxy falls close to a line with a slope less than unity ( $K_{IC} > K_{\max}(\Delta K_I^{\max})$ ), indicating a trend of increasing fatigue crack growth resistance with increasing monotonic fracture toughness. A similar trend has been reported previously for rubber-modified epoxies [5,8–9,13], implying that the failure mechanisms for static fracture and fatigue failure in modified epoxies are comparable.

For neat epoxy the  $K_{\max}(\Delta K_I^{\max})$  value at fatigue crack instability exceeds the static fracture toughness values  $K_{IC}$ . This anomalous behavior has been reported previously for other filled polymer systems [5,11–12,14]. The inconsistency between  $K_{\max}(\Delta K_I^{\max})$  and  $K_{IC}$  is explained by differences in the crack-tip geometry [27] and to a lesser degree the loading rate [5,28] between monotonic and fatigue testing. Monotonic fracture toughness values are determined using precracks generated by a razor blade as prescribed in ASTM Standard D 5045. In contrast, the maximum stress intensity values  $K_{\max}$  are determined for a crack tip generated in fatigue by progressively increasing  $\Delta K_I$ . As described by Xiao *et al.* [27], yielding and damage at the crack tip under cyclic loading can cause apparent fracture toughness values for DGEBA epoxy with precracks generated by fatigue to exceed fracture toughness values measured in samples with razor blade generated precracks by as much as 31%. Hence, the high value of  $\Delta K_I^{\max}$  for neat epoxy is an artifact of the specimen loading history introducing progressive blunting at

the crack tip. This effect was only observed in the neat epoxy and is not present in the microcapsule modified systems.



*Figure 8.* Influence of  $K_{IC}$  on  $\Delta K_I^{\max}$ . The dashed line corresponds to the maximum applied stress intensity range  $\Delta K_I^{\max}$  for fatigue crack instability when  $K_{\max}(\Delta K_I^{\max}) \equiv \Delta K_I^{\max}/(1-R) = K_{IC}$ . The solid line is the best linear fit through the data points and the origin.

In addition to the toughening mechanisms induced by the embedded microcapsules, the flow of fluid released into the crack plane provides a crack tip shielding mechanism that can improve resistance to fatigue crack propagation. Several authors have reported that hydrodynamic pressure due to viscous flow within a fatigue crack reduces the effective range of Mode-I stress intensity and hence the fatigue crack growth rate for metal submerged in oil [29–32]. During cyclic loading, the crack volume changes significantly with time, requiring fluid flow into and out of the crack. When the crack contains a viscous fluid, the forces required to squeeze the fluid out of the crack during unloading and draw fluid into the crack during loading can be sufficient to shield the crack tip. Reduced crack growth rates in epoxy due to crack tip shielding from viscous fluid flow have been demonstrated for infiltration of both precatalyzed dicyclopentadiene and mineral oil [20].

## 5. Conclusions

Fatigue crack propagation was investigated in epoxy toughened with liquid-filled urea-formaldehyde (UF) microcapsules. The addition of microcapsules significantly decreased the fatigue crack growth rate and increased the fatigue life above a transition value of the stress intensity factor  $\Delta K_T$ . Below  $\Delta K_T$  the fatigue behavior was unaffected by the embedded microcapsules. The transition value between these two regimes corresponded to loading conditions where the size of the plastic zone approached the size of the embedded microcapsules. The fatigue-crack growth rate dependence on applied range of stress intensity  $\Delta K_I$  was accurately captured by the Paris power law in both neat epoxy and epoxy with embedded microcapsules. The Paris law exponent  $n$

was strongly dependent on the microcapsule concentration, varying from 9.7 for neat epoxy to approximately 4.5 above 10 wt% microcapsules, but was independent of microcapsule diameter. The onset of unstable fatigue-crack growth  $\Delta K_I^{\max}$  increased with monotonic fracture toughness, and was independent of microcapsule diameter. Improved resistance to fatigue crack propagation, was attributed to toughening mechanisms induced by the embedded microcapsules as evidenced by changes in the fatigue fracture plane morphology. In addition to increasing the inherent fatigue life of epoxy, embedded microcapsules filled with an appropriate healing agent offer a potential mechanism to further extend fatigue life through self-healing of fatigue damage.

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963	Dong, F., A. T. Hsui, and D. N. Riahi	A stability analysis and some numerical computations for thermal convection with a variable buoyancy factor— <i>Journal of Theoretical and Applied Mechanics</i> <b>2</b> , 19–46 (2002)	Jan. 2001
964	Phillips, W. R. C.	Langmuir circulations beneath growing or decaying surface waves— <i>Journal of Fluid Mechanics</i> (submitted)	Jan. 2001
965	Bdzil, J. B., D. S. Stewart, and T. L. Jackson	Program burn algorithms based on detonation shock dynamics— <i>Journal of Computational Physics</i> (submitted)	Jan. 2001
966	Bagchi, P., and S. Balachandar	Linearly varying ambient flow past a sphere at finite Reynolds number: Part 2—Equation of motion— <i>Journal of Fluid Mechanics</i> <b>481</b> , 105–148 (2003) (with change in title)	Feb. 2001
967	Cermelli, P., and E. Fried	The evolution equation for a disclination in a nematic fluid— <i>Proceedings of the Royal Society A</i> <b>458</b> , 1–20 (2002)	Apr. 2001
968	Riahi, D. N.	Effects of rotation on convection in a porous layer during alloy solidification—Chapter 12 in <i>Transport Phenomena in Porous Media</i> (D. B. Ingham and I. Pop, eds.), 316–340 (2002)	Apr. 2001
969	Damljanovic, V., and R. L. Weaver	Elastic waves in cylindrical waveguides of arbitrary cross section— <i>Journal of Sound and Vibration</i> (submitted)	May 2001
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1047	Bagchi, P., and S. Balachandar	Response of the wake of an isolated particle to isotropic turbulent cross-flow— <i>Journal of Fluid Mechanics</i> (submitted)	Mar. 2004
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