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FINITE ELEMENT FORMULATION AND SOLUTION OF CONTACT-  
IMPACT PROBLEMS IN CONTINUUM MECHANICS

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Prepared for:

Naval Civil Engineering Laboratory  
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May 1974

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Naval Construction Battalion Center  
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**FINITE ELEMENT FORMULATION AND SOLUTION  
OF CONTACT-IMPACT PROBLEMS IN CONTINUUM  
MECHANICS**

May 1974

An Investigation Conducted by  
**STRUCTURAL ENGINEERING LABORATORY**  
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Berkeley, California

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## Introduction

In this report we consider the general problem of contact and impact between two bodies. The report is divided into three basic parts. These parts describe: (I) The general theory of contact-impact problems, (II) A numerical scheme for the analysis of contact-impact problems, and (III) The description of computer program FEAP 74 for the solution of contact-impact problems. In an appendix we include the program subroutines and general input description for FEAP 74.

In Sections 1 to 6, Part I, we deal with spatial aspects of the theory and in Section 7, Part I, we deal with temporal aspects. This splitting of the theory is motivated by the way we intend to numerically solve the equations, i.e., the finite element method spatially and a finite difference method temporally.

Part II considers a numerical implementation of the theory given in Part I. Section 9 deals with spatial notions of the numerical problem and Section 10 the temporal. The solution scheme for the resulting algebraic problem is discussed in Sections 11 and 12.

The computer program FEAP 74 was modified to incorporate the numerical contact-impact model. The program modifications and capabilities together with two numerical examples are contained in Part III.

Finally, in the appendix we give listings for the contact subroutines together with the data input instructions.

PART I  
VARIATIONAL FORMULATION OF CONTACT-IMPACT  
PROBLEMS IN CONTINUUM MECHANICS

1. Preliminaries

Our conventions on indices are as follows:

Superscripts indicate to which body an entity pertains. Summation is to take place only when explicitly indicated.

Latin subscripts range over 1,2,3, while Greek subscripts range over 1,2. The summation convention is assumed to hold for both.

A body  $B$  is a nice connected region of  $\mathbb{R}^3$  with a piecewise smooth boundary  $\partial B$ . A contact\* problem is a boundary value problem, or an initial-boundary value problem, in which two bodies,  $B^1$  and  $B^2$ , interact according to the principles of mechanics. Thus the primary kinematic axiom of a contact problem is that configurations  $G^1$  and  $G^2$ , of  $B^1$  and  $B^2$ , respectively, do not penetrate each other, i.e.,

$$\begin{aligned} (G^1)^\circ \cap G^2 &= \emptyset, \\ G^1 \cap (G^2)^\circ &= \emptyset, \end{aligned} \tag{1}$$

where  $(G^\alpha)^\circ$  denotes the interior of  $G^\alpha$ ,  $\alpha = 1,2$ .

On the other hand the unique condition which characterizes contact problems is that material points on the boundaries of  $B^1$  and  $B^2$  may coalesce during the motion of the bodies. Thus we say  $B^1$  and  $B^2$  are in contact if  $\partial G^1 \cap \partial G^2 \neq \emptyset$ , and we define the contact surface  $e$  by

---

\* It is usual for the term contact to have a static connotation while the term impact has a dynamic connotation. We shall use contact in the general sense to include static as well as dynamic phenomena.



$$\mathcal{E} = \partial B^1 \cap \partial B^2 . \quad (2)$$

If  $B^1$  and  $B^2$  are never in contact then  $\mathcal{E} = \emptyset$  for all configurations  $B^1$  and  $B^2$ , and in this case an initial-boundary value problem for  $B^1$  and  $B^2$  reduces to one in which  $B^1$  and  $B^2$  may be treated separately. Thus a non-trivial contact problem is one in which  $\mathcal{E} \neq \emptyset$  for at least one instant during the motion of  $B^1$  and  $B^2$ . The picture (Fig. 1) illustrates these notions.

Equation (1) implies that  $\mathcal{E}$  is a material surface with respect to both bodies, i.e., one which is not crossed by material particles. From this we may deduce the interface conditions on  $\mathcal{E}$ .

Let  $\underline{x}$  be a persistent point of  $\mathcal{E}$  (one at which joining or releasing of the bodies is not instantaneously occurring) and  $\underline{v}$  be the velocity of  $\underline{x}$  ( $\underline{v} = \dot{\underline{x}}$ ). Note that only the normal part of  $\underline{v}$  is independent of the parametrization of  $\mathcal{E}$ . Let  $\underline{v}^1$  and  $\underline{v}^2$  be the velocities of the material particles located at the points  $\underline{x}^1$  and  $\underline{x}^2$ , contained in  $\partial B^1$  and  $\partial B^2$ , respectively, such that  $\underline{x} = \underline{x}^1 = \underline{x}^2$  at the present instant. Then since  $\mathcal{E}$  is material and  $\underline{x}$  is persistent

$$\underline{v} \cdot \underline{n} = \underline{v}^1 \cdot \underline{n} = \underline{v}^2 \cdot \underline{n} , \quad (3)$$

where  $\underline{n}$  is a unit normal vector to  $\mathcal{E}$  at  $\underline{x}$ . From this it follows that a necessary condition for momentum to be balanced at  $\underline{x}$  is that

$$(\underline{t}^1 + \underline{t}^2) \cdot \underline{n} = \underline{0} , \quad (4)$$

where  $\underline{t}^m$  is the Cauchy traction vector with respect to  $\partial B^m$ .

In addition we assume that no tensile tractions can occur on  $\mathcal{E}$ ,

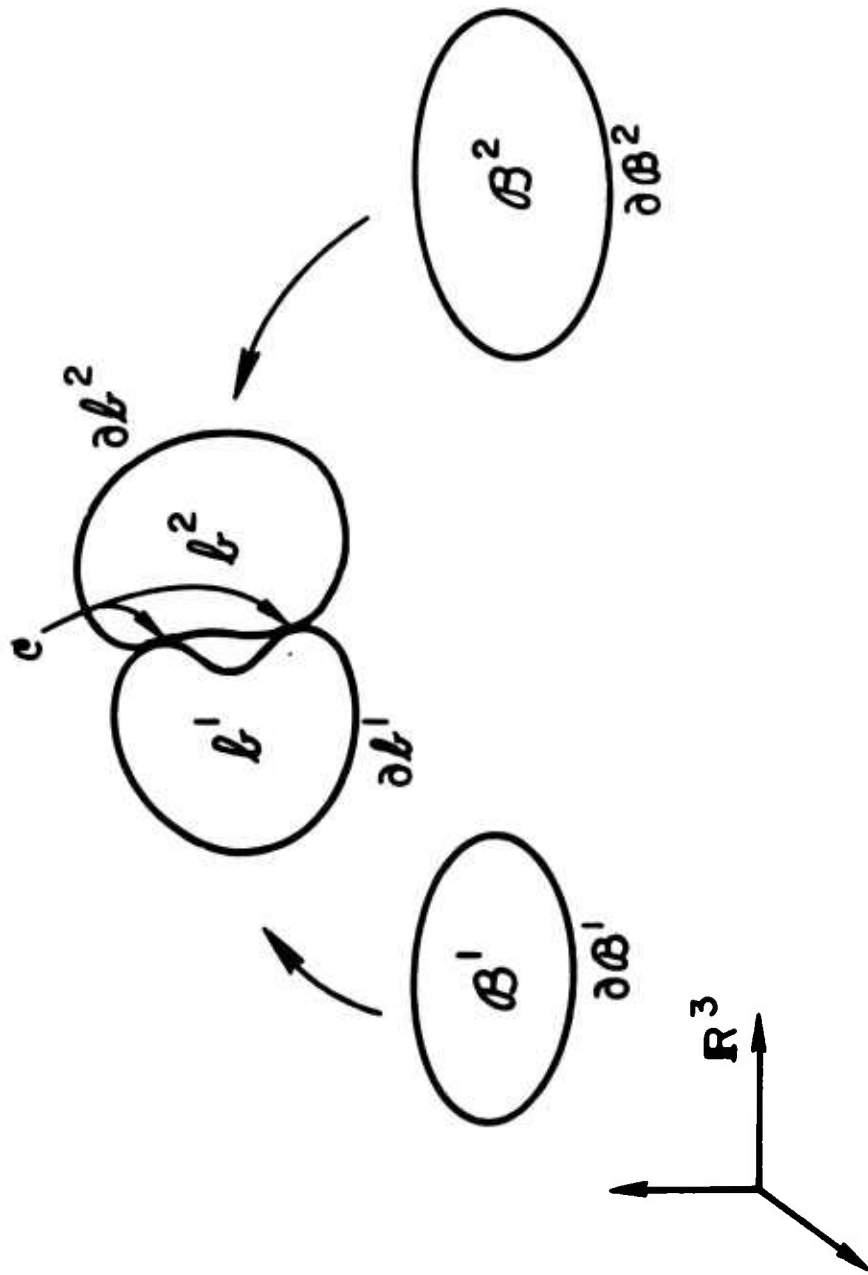


Figure 1

$$\underline{t}^{\alpha} \cdot \underline{n}^{\alpha} \leq 0, \quad (5)$$

where  $\underline{n}^{\alpha}$  is the outward unit normal vector to  $\partial \mathcal{B}^{\alpha}$ . This condition excludes the possibility of the two bodies being glued together. Conditions (1-5) characterize our notion of a contact problem.

Note that thus far we have said nothing about the tangential parts of  $\underline{v}^{\alpha}$  and  $\underline{t}^{\alpha}$ . These remaining conditions are determined by the frictional nature of the contact. We shall study two simple cases.

Case I: If we assume that points, once in contact, move with  $\mathcal{C}$  until released, we have that

$$\underline{v}^1 = \underline{v}^2, \quad (6)$$

and therefore

$$\underline{t}^1 + \underline{t}^2 = \underline{0}. \quad (7)$$

For this model we say that a no-slip, or perfect friction, condition is achieved on  $\mathcal{C}$ . Thus condition (5) and equations (6) and (7) are the interface conditions for this case.

Case II: We may create the interface conditions for a frictionless, sliding contact by asserting that the tangential part of each  $\underline{t}^{\alpha}$  is identically zero,

$$\underline{t}^{\alpha} - (\underline{t}^{\alpha} \cdot \underline{n}^{\alpha}) \underline{n}^{\alpha} = \underline{0}. \quad (8)$$

Eq. (8), along with (3-5), are the interface conditions for this case.

## 2. Variational Theorems

We will formulate a variational theorem for the contact problem of finite elastodynamics. We point out, however, that our treatment is entirely general and could be used in conjunction with any field theory, as the only unique feature of the formulation involves the handling of interface conditions. At the same time finite elastodynamics, though lending itself to a clean and simple variational statement, is a case of wide practical interest.

We shall first obtain a variational theorem for the usual initial-boundary value problem of finite elastodynamics by a trivial generalization of some work done by S. Nemat-Nasser [1].

For notational simplicity let  $\mathcal{C}$  denote  $\partial\mathcal{B}$ , and let  $d\mathcal{A}$  and  $d\mathcal{V}$  denote area and volume forms for  $\mathcal{B}$  and  $\mathcal{C}$ , respectively. Let  $\mathcal{A}_\tau \subset \mathcal{A}$  be that part of  $\mathcal{C}$  where surface tractions are prescribed, and denote by  $\bar{\mathbf{T}}$  the Piola - Kirchhoff traction vector representing these prescribed tractions. Call  $\rho_0$  the density of  $\mathcal{B}$  in the initial configuration,  $\underline{\mathbf{F}}$  the extrinsic body force vector and let  $\underline{\mathbf{x}} = \underline{\mathbf{x}}_0(\underline{\mathbf{X}})$  represent the position at time  $t$  of the material particle located at  $\underline{\mathbf{X}}$  in the initial configuration. For convenience we take  $\mathcal{B}$  to be the initial configuration. We denote by  $\partial\underline{\mathbf{x}}/\partial\underline{\mathbf{X}}$  the deformation gradients and by  $\Phi(\partial\underline{\mathbf{x}}/\partial\underline{\mathbf{X}})$  the strain energy density. Then if  $\underline{\mathbf{x}}$  satisfies the kinematic boundary conditions

$$\underline{\mathbf{x}} = \bar{\underline{\mathbf{x}}} \quad (9)$$

on  $\mathcal{A}_\alpha \subset \mathcal{A}$ , where

$$\begin{aligned} \mathcal{A}_\alpha \cup \mathcal{A}_\tau &= \mathcal{A}, \\ \mathcal{A}_\alpha \cap \mathcal{A}_\tau &= \emptyset, \end{aligned}$$

the functional  $\Pi$  defined by

$$\begin{aligned} \Pi(\underline{x}) = \int_0^t \left\{ \int_{\mathcal{B}} \left( \Phi(\partial \underline{x} / \partial \underline{X}) - \rho_0 \dot{\underline{x}} \cdot \dot{\underline{x}} / 2 \right. \right. \\ \left. \left. - \rho_0 \underline{F} \cdot \underline{x} \right) d\mathcal{B} - \int_{\mathcal{A}_r} \underline{x} \cdot \underline{T} d\mathcal{A} \right\} dt, \end{aligned} \quad (10)$$

is stationary, i.e., its first variation vanishes

$$\begin{aligned} 0 = \delta \Pi(\underline{x}, \delta \underline{x}) = \int_0^t \left\{ \int_{\mathcal{B}} \left( \rho_0 (\dot{\underline{x}} - \underline{F}) - \text{DIV } \underline{P} \right) \cdot \delta \underline{x} d\mathcal{B} \right. \\ \left. + \int_{\mathcal{A}_r} \left( \underline{T} - \underline{\bar{T}} \right) \cdot \delta \underline{x} d\mathcal{A} \right\} dt, \end{aligned} \quad (11)$$

subject to the constraint on variations  $\delta \underline{x}_t = \delta \underline{x}_0 = \underline{0}$ , (12)

if and only if the equations of motion and traction boundary conditions are satisfied

$$\rho_0 (\dot{\underline{x}} - \underline{F}) = \text{DIV } \underline{P}, \quad \text{in } \mathcal{B}, \quad (13)$$

$$\underline{T} = \underline{\bar{T}}, \quad \text{on } \mathcal{A}_r, \quad (14)$$

where  $\underline{P} = \partial \Phi / \partial (\partial \underline{x} / \partial \underline{X})$  is the first Piola - Kirchhoff stress tensor,  $\underline{T} = \underline{N} \cdot \underline{P}$  is the Piola - Kirchhoff traction vector, and  $\underline{N}$  is the outward unit normal vector to  $\mathcal{A}$ . The solution to the initial-boundary value problem must also satisfy the given initial conditions

$$\left. \begin{aligned} \underline{x} &= \underline{x}_0 \\ \underline{v} &= \underline{v}_0 \end{aligned} \right\} \quad \text{in } \mathcal{B}. \quad (15)$$

To interpret this variational theorem for two (non-interacting) bodies set

$$\begin{aligned} \mathcal{B} &= \mathcal{B}^1 \cup \mathcal{B}^2, \\ \mathcal{A} &= \mathcal{A}^1 \cup \mathcal{A}^2, \quad \text{etc.}, \end{aligned}$$

and write

$$\mathbb{I}(\underline{x}) = \mathbb{I}^1(\underline{x}^1) + \mathbb{I}^2(\underline{x}^2).$$

The next step is to add to  $\mathbb{I}$  terms manifesting the interface conditions on  $\mathcal{C}$  and to stipulate the constraints under which the vanishing of the first variation of the appended functional corresponds to a solution of the contact problem. To do this we must consider further the kinematics and geometry of  $\mathcal{C}$ .

Define two piecewise smooth, invertible maps  $\underline{x}^1, \underline{x}^2$  by the condition

$$(\underline{x}^\alpha)^{-1}: \mathcal{C} \longrightarrow \mathcal{C}^\alpha \subset \mathcal{Q}^\alpha, \quad (16)$$

where each  $\underline{x}^\alpha$  identifies points on the boundary of the initial configuration  $\mathcal{B}^\alpha$  which map into the contact surface  $\mathcal{C}$  at each instant of time. If  $\underline{x} \in \mathcal{C}$ , then  $\underline{x}^1 = (\underline{x}^1)^{-1}(\underline{x})$  and  $\underline{x}^2 = (\underline{x}^2)^{-1}(\underline{x})$  are the positions of particles in  $\mathcal{Q}^1$  and  $\mathcal{Q}^2$ , respectively, which have coalesced at  $\underline{x} \in \mathcal{C}$ . It is clear what the  $\underline{x}^\alpha$ 's really are, viz., if  $\underline{x}^\alpha = \underline{x}_t^\alpha(\underline{X}^\alpha)$ , for all  $\underline{X}^\alpha \in \mathcal{B}^\alpha$ , represents the motion of body  $\mathcal{B}^\alpha$  from the original configuration  $\mathcal{B}^\alpha$  to the present one  $\mathcal{B}^\alpha$ , then  $\underline{x}^\alpha$  is the restriction of  $\underline{x}^\alpha$  to  $\mathcal{C}^\alpha$ ,

$$\underline{x}^\alpha(\underline{X}^\alpha) = \underline{x}^\alpha(\underline{X}^\alpha), \quad (17)$$

for each  $\underline{X}^\alpha \in \mathcal{C}^\alpha$ ,  $\alpha = 1, 2$ . For the time being we consider the  $\underline{x}^\alpha$ 's as maps defined independently of the  $\underline{x}^\alpha$ 's and consider (17) a constraint on possible motions.

We are interested in to what extent the relation

$$\underline{x} = \underline{x}^1(\underline{X}^1) = \underline{x}^2(\underline{X}^2) , \quad (18)$$

is smooth in time and analogously under what circumstances the variations of the  $\underline{x}^a$  's are equal. In general the  $\underline{x}^a$  's will not even be continuous in time since contact surfaces can be instantaneously created or destroyed. If we eliminate such exceptional instants and consider only persistent points, the bodies still may slide with respect to each other, as depicted in Fig. 2. Thus tangential velocities are seen to be unequal in general. However, when  $\underline{x}$  is persistent, the impenetrability condition (1) forces the normal velocity components to be equal, and concomitantly the normal components of variations of the  $\underline{x}^a$  's are also equal

$$\delta \underline{x}^1 \cdot \underline{n} = \delta \underline{x}^2 \cdot \underline{n} \quad (19)$$

For sliding contact (Case II), Eq. (19) characterizes the constraint on variations of the  $\underline{x}^a$  's equivalent to the velocity constraint (3).

For no-slip contact (Case I),

$$\delta \underline{x}'^1 = \delta \underline{x}'^2 , \quad (20)$$

is easily seen to be the condition on variations equivalent to Eq. (6). We shall see that Eqs. (19) and (20) lead to the proper interface conditions in the variational theorems.

Introduce vector valued Lagrange multipliers  $\underline{\tau}^a$  , and add

$$\mathcal{X} = - \sum_{a=1}^2 \int_0^t \int_{\mathcal{C}^a} \underline{\tau}^a \cdot (\underline{x}^a - \underline{x}^a) d\mathcal{C}^a dt , \quad (21)$$

to the functional  $\Pi$  (Eq. 10). Note that when  $\mathcal{C} \neq \emptyset$  ,

$$\mathcal{C} = \mathcal{C}_x \cup \mathcal{C}_T \cup \mathcal{C} ,$$

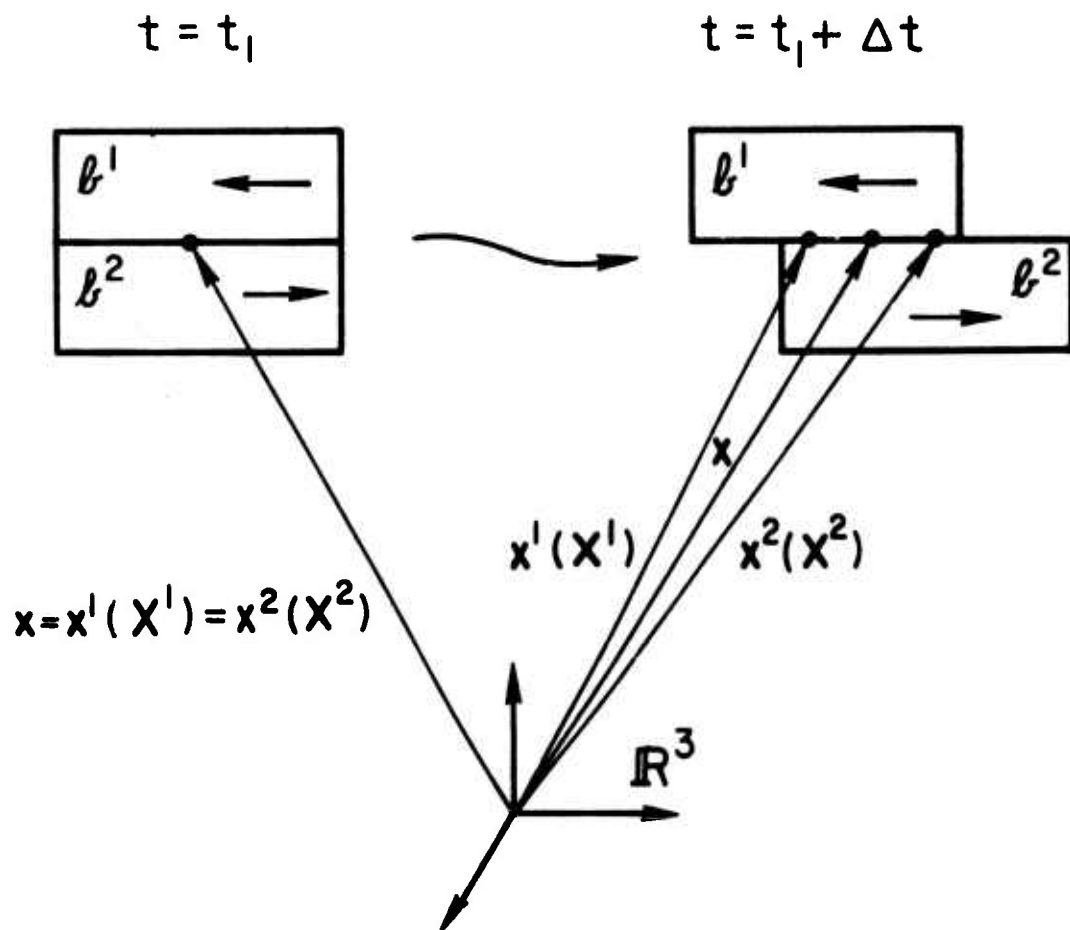


Figure 2



and assume for consistency's sake that

$$\begin{aligned} \mathcal{C}^{\alpha} &\subset \mathcal{A}_{\tau}^{\alpha}, \\ \underline{\mathbb{T}} &= \underline{\mathbb{0}} \quad \text{on} \quad \mathcal{C}^{\alpha}. \end{aligned} \tag{22}$$

This condition will preclude the ambiguous circumstance of non-zero tractions being specified on the contact area. Upon taking variations of

$\mathcal{J} = \mathbb{I} + \mathcal{X}$  we get Eqs. (13), (14) and,

$$\begin{aligned} 0 = & - \sum_{\alpha=1}^2 \int_0^t \int_{\mathcal{C}^{\alpha}} \left\{ \delta \underline{\mathbb{C}}^{\alpha} \cdot (\underline{x}^{\alpha} - \underline{\mathcal{L}}^{\alpha}) + \right. \\ & \left. + \delta \underline{x}^{\alpha} \cdot (\underline{\mathbb{C}}^{\alpha} - \underline{\mathbb{T}}^{\alpha}) - \delta \underline{\mathcal{L}}^{\alpha} \cdot \underline{\mathbb{C}}^{\alpha} \right\} d\mathcal{C}^{\alpha} dt + \\ & + \text{transversality condition.} \end{aligned} \tag{23}$$

The transversality condition is the classical terminology for variations associated with the domain  $\mathcal{C}^{\alpha}$ .

The first summand of (23) gives us (17) which insures that the  $\underline{x}^{\alpha}$ 's map into  $\mathcal{C}$  properly. The second summand identifies  $\underline{\mathbb{C}}^{\alpha}$  as the Piola - Kirchhoff traction vector  $\underline{\mathbb{T}}^{\alpha}$  on  $\mathcal{C}^{\alpha}$ . Let us investigate the third summand.

Consider first Case I and define

$$\delta \underline{\mathcal{L}}^{\alpha} = \delta \underline{\mathcal{L}}^{\alpha}, \quad \alpha = 1, 2, \tag{24}$$

which makes sense because of Eq. (20). This condition is equivalent to insisting

$$\dot{\underline{\mathcal{L}}}^{\alpha} = \dot{\underline{\mathcal{L}}}^{\alpha}, \quad \alpha = 1, 2,$$

thus the first summand of (23) also implies (6) holds whenever we have a

persistent point. Let  $j^{\alpha}$  denote the Jacobian determinant associated with  $\chi^{\alpha}$ ,

$$de = j^{\alpha} dC^{\alpha}. \quad (25)$$

Notice then that since  $\underline{C}^{\alpha}$  is the Piola - Kirchhoff traction vector,  $(1/j^{\alpha}) \underline{C}^{\alpha}$  is the corresponding Cauchy traction vector. With these we have for the third summand,

$$0 = \sum_{\alpha=1}^2 \int_{C^{\alpha}} \delta \chi^{\alpha} \cdot \underline{C}^{\alpha} dC^{\alpha} = \int_{\mathcal{E}} \delta \chi^{\alpha} \cdot ((1/j^1) \underline{C}^1 + (1/j^2) \underline{C}^2) de, \quad (26)$$

which in words means the Cauchy traction vectors are in equilibrium. Thus the momentum balance, Eq. (7), is satisfied on  $\mathcal{E}$ .

In Case II we only have that (19) holds, so define

$$\delta \chi^{\alpha}(n) = \delta \chi^{\alpha} \cdot n, \quad \alpha = 1, 2. \quad (27)$$

This requirement also insures that,

$$\dot{\chi}^1 \cdot n = \dot{\chi}^2 \cdot n,$$

thus the first summand of (23) implies (3). For this case the third summand takes the form,

$$0 = \sum_{\alpha=1}^2 \int_{C^{\alpha}} \delta \chi^{\alpha} \cdot \underline{C}^{\alpha} dC^{\alpha} = \int_{\mathcal{E}} \delta \chi^{\alpha}(n) ((1/j^1) \underline{C}^1 \cdot n + (1/j^2) \underline{C}^2 \cdot n) de + \sum_{\alpha=1}^2 \int_{C^{\alpha}} (\delta \chi^{\alpha} - \delta \chi^{\alpha}(n) n) \cdot \underline{C}^{\alpha} dC^{\alpha}. \quad (28)$$

The integral over  $\mathcal{E}$  gives us Eq. (4). The significance of the second integral hinges on the observation that  $(\delta \chi^{\alpha} - \delta \chi^{\alpha}(n) n)$  is a tangent vector to  $\mathcal{E}$  for each  $\alpha$ . Thus the tangential part of each  $\underline{C}^{\alpha}$  is identically zero, which is equivalent to the shear free condition, Eq. (8),

which we require for Case II.

A standard calculation enables us to write the transversality condition as,

$$0 = \sum_{\alpha=1}^2 \int_{\partial C^{\alpha}} (\delta X^{\alpha} \cdot \underline{T}^{\alpha} \underline{C}^{\alpha} \cdot (\underline{x}^{\alpha} - \underline{x}^{\alpha})) \, d(\partial C^{\alpha}), \quad (29)$$

where the transversal  $\underline{T}^{\alpha}$  is a unit vector field tangent to  $C^{\alpha}$ , and perpendicular and pointing outward with respect to  $\partial C^{\alpha}$ , Fig. 3. Thus (29) implies that

$$\underline{C}^{\alpha} \cdot (\underline{x}^{\alpha} - \underline{x}^{\alpha}) = 0 \quad \text{on } \partial C^{\alpha}, \quad \alpha=1,2 \quad (30)$$

Assuming continuity of the integrands of (21) on the closure of  $C^{\alpha}$ , condition (30) is already implied by the first summand of (23). This assumption precludes  $\underline{C}^{\alpha}$  taking the form of a  $\delta$ -distribution on  $\partial C^{\alpha}$ .

Although this assumption is warranted here it may not be true when one employs certain approximate theories in mechanics. For instance consider the case where a Bernoulli-Euler beam is uniformly loaded and sits on a rigid parabolic surface (Fig. 4). At the contact points  $a, a'$ , concentrated reactions must exist to balance shear forces. This example is actually from a completely different class of contact problems in that contact is made along a part of the interior rather than the boundary. Such problems as the contact of plates and shells also fall into this class. We could summarize such situations by the description --  $m$ -dimensional contact of  $m$ -dimensional bodies, e.g., for the beam  $m=1$ , and for plates and shells  $m=2$ . The case under investigation in this paper ( $m=3$ ) is an example of the  $(m-1)$ -dimensional contact of  $m$ -dimensional bodies.

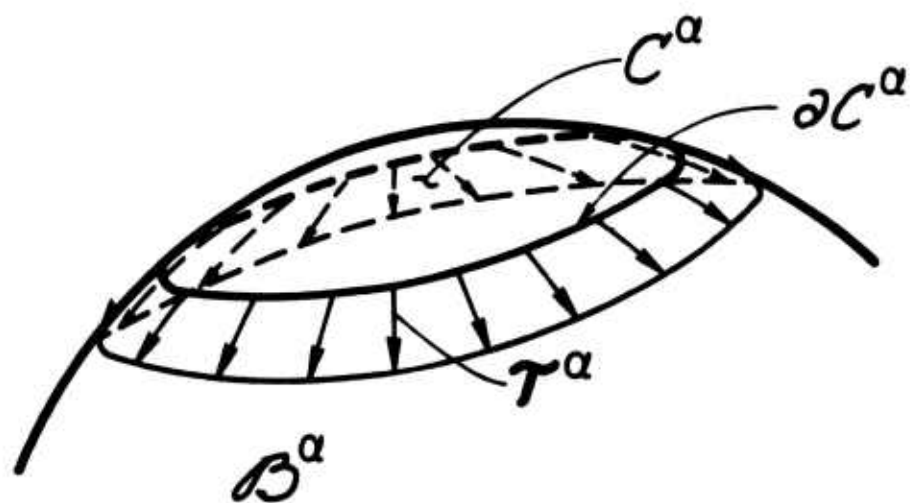


Figure 3

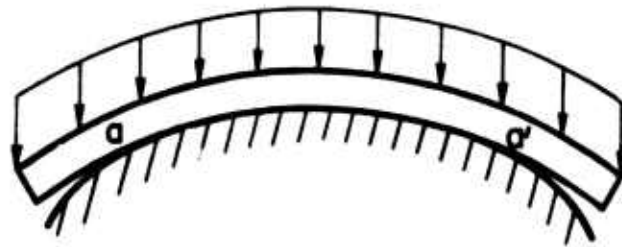


Figure 4

It is good to keep in mind cases such as that illustrated in Fig. 4 when considering specific boundary value problems.

A further point worth mentioning here is that the transversality condition will in general be an independent one in a numerical algorithm. For example, if the fields in the integrand of (21) are approximated by a family of trial functions, Eq. (23) only implies that some weighted integrals over the  $\zeta^\alpha$ 's vanish. The condition (29) requires that weighted integrals over the  $\delta\zeta^\alpha$ 's also vanish.

We now summarize our results in the following theorems:

Theorem I: Let (1), (2), (5), (9), (12), (15), and (20) hold. Then  $\underline{x}$  is a solution to the no-slip contact problem (Case I), that is, (6), (7), (13), (14) and (17) also hold if, and only if,  $\delta J = 0$  for arbitrary variations of  $\underline{x}^\alpha$ ,  $\underline{y}^\alpha$  and  $\zeta^\alpha$ ,  $\alpha = 1, 2$ .

Theorem II: Let (1), (2), (5), (9), (12), (15), and (19) hold. Then  $\underline{x}$  is a solution to the sliding contact problem (Case II), that is, (3), (4), (8), (13), (14) and (17) also hold if, and only if,  $\delta J = 0$  for arbitrary variations of  $\underline{x}^\alpha$ ,  $\underline{y}^\alpha$  and  $\zeta^\alpha$ ,  $\alpha = 1, 2$ .

### 3. Consideration of Theorems I and II as Computational Tools

Theorems I and II may be employed to generate numerical algorithms for the solution of contact problems. The basic idea is to represent  $\mathbf{x}^n$ ,  $\mathbf{z}^n$  and  $\mathbf{c}^n$  as the product of known functions on  $\mathbb{R}^3$  with unknown parameters depending on time. Then Theorems I and II provide us with a method for generating an approximate system of equations (e.g., by the classical Ritz-Galerkin technique) in terms of these unknown parameters, which then can be solved incrementally and/or iteratively, subject to the side conditions of the theorems. The constraints (1) and (5) will both take the form of inequalities in actual computations, thus the ideas of optimization theory will probably be useful in the actual construction of a numerical algorithm.

The finite element method is a powerful technique for obtaining a system of approximate equations, and it is of interest to find out how amenable are Theorems I and II to a finite element formulation. Unfortunately the term  $\chi$  would result in a terrible mess if the integrand was represented by typical finite element functions. This is because the boundaries of the  $\mathcal{C}^n$ 's are unknown and thus a parametric integration would bury the defining parameters of the  $\mathcal{C}^n$ 's in the arguments of Heaviside functions representing the supports of the elements. Note that a classical Ritz-Galerkin approximation would not be subject to this pitfall, since the associated trial functions could be chosen to be real analytic and thus easily integrated parametrically to a relatively simple form. However, such a formulation is restricted to a geometrically simpler class of problems. Thus it is desirable to seek a generalization that will lend itself cleanly to a finite element formulation.

#### 4. Variational Theorems Without Transversality Conditions

Let  $\tilde{C}^n$  be a fixed part of  $A_T^n$  such that

$$\tilde{C}^n \supset C^n, \quad (31)$$

and

$$\bar{T} = 0 \quad \text{on} \quad \tilde{C}^n \sim C^n. \quad (32)$$

Define a scalar valued function  $\eta^n$  on  $\tilde{C}^n$  such that

$$\eta^n(x^n) = 0 \quad \text{if} \quad x^n \in \tilde{C}^n \sim C^n. \quad (33)$$

Let  $\tilde{e} \supset e$ , and define the maps  $\tilde{x}^n$  by the condition

$$(\tilde{x}^n)^{-1} : \tilde{e} \rightarrow \tilde{C}^n,$$

where, as before,  $\tilde{x}^n$  represents  $x^n$  on  $C^n$ ; but on  $\tilde{C}^n \sim C^n$  we place no physical interpretation on  $\tilde{x}^n$ . Thus on  $\tilde{C}^n$  we will always have that,

$$\eta^n(\tilde{x}^n - x^n) = 0, \quad (34)$$

since  $\tilde{x}^n = x^n$  on  $C^n$  and  $\eta^n = 0$  on the relative complement  $\tilde{C}^n \sim C^n$ .

Introduce vector valued Lagrange multipliers  $\sigma^n$  and let  $\mathcal{L} = \mathbb{I} + \mathcal{M}$  where

$$\mathcal{M} = - \sum_{a=1}^k \int_0^1 \int_{\tilde{C}^n} \sigma^n \cdot \eta^n(\tilde{x}^n - x^n) \, d\tilde{C}^n dt. \quad (35)$$

We require that the variations of  $\tilde{x}^n$  satisfy the same conditions as before, but now for all  $\tilde{C}^n$ :

$$\left. \begin{array}{l} \text{Case I:} \quad \delta \tilde{x}^n \equiv \delta x^n \\ \text{Case II:} \quad \delta \tilde{x}^n(n) \equiv \delta x^n \cdot n \end{array} \right\} \quad \text{on} \quad \tilde{C}^n \quad (36)$$



where  $\underline{n}$  is a unit normal vector to  $\tilde{E}$ . Computing the first variation of  $\mathcal{L}$  we have the usual conditions emanating from II and

$$\begin{aligned}
 0 = & - \frac{\partial \mathcal{L}}{\partial \underline{x}} \int_0^t \int_{\tilde{E}} \left\{ \delta \underline{\sigma} \cdot (\underline{x} - \underline{x}') \right. \\
 & + \delta \underline{\eta} \cdot (\underline{\sigma} \cdot (\underline{x} - \underline{x}')) \\
 & + \delta \underline{x} \cdot (\underline{n} \cdot \underline{\sigma} - \underline{T}) \\
 & \left. - \delta \underline{x}' \cdot (\underline{n} \cdot \underline{\sigma}') \right\} d\tilde{E} dt .
 \end{aligned} \tag{37}$$

The first summand gives us (34) and we define

$$\mathcal{C} = \{ \underline{x} \in \tilde{E} : \underline{x}(\underline{x}) = \underline{x}'(\underline{x}) \} . \tag{38}$$

The third summand defines  $\underline{\eta} \cdot \underline{\sigma}$  as the Piola - Kirchhoff traction vector. Note that this insures that  $\underline{T} = \underline{0}$  on  $\tilde{E} \sim \mathcal{C}$  since  $\underline{\eta} = \underline{0}$  there. The fourth summand gives us the appropriate Cauchy traction condition across  $\mathcal{C}$  for each case of (36). The second summand is identically satisfied on  $\mathcal{C}$  since  $\underline{x} = \underline{x}'$ . On  $\tilde{E} \sim \mathcal{C}$  it tells us that  $\underline{\sigma}$  is orthogonal to  $\underline{x} - \underline{x}'$ , but this is of no physical interest.

Thus we can state the following theorems:

Theorem I': Let (1), (2), (5), (9), (12), (15), (31), (32) and (36)<sub>1</sub> hold. Then  $\underline{x}$  is a solution to the no-slip contact problem (Case I), that is, (6), (7), (13), (14) and (17) also hold where  $\mathcal{C}$  is defined by (38), if  $\delta \mathcal{L} = 0$  for arbitrary variations of  $\underline{x}$ ,  $\underline{x}'$ ,  $\underline{\eta}$  and  $\underline{\sigma}$ ,  $\alpha = 1, 2$ .

Theorem II': Let (1), (2), (5), (9), (12), (15), (31), (32) and (36)<sub>2</sub> hold. Then  $\underline{x}$  is a solution to the sliding contact problem (Case II), that is, (3), (4), (8), (13), (14) and (17) hold where  $\mathcal{C}$  is defined by (38), if  $\delta \mathcal{L} = 0$  for arbitrary variations of  $\underline{x}$ ,  $\underline{x}'$ ,  $\underline{\eta}$  and  $\underline{\sigma}$ ,  $\alpha = 1, 2$ .

The important feature of these theorems is that the regions  $\tilde{E}^+$  are fixed. Thus transversality conditions are absent, and the theorems may be applied to finite element formulations. In fact one would naturally take  $\tilde{E}^+$  to be a union of elements in  $\mathcal{Q}^+$ , large enough to contain  $\tilde{E}^+$  throughout the motion.

Thus far our considerations have been quite general and, in fact, more general than would be required for the solution of particular classes of contact problems. In the next section we illustrate the many simplifications which can be made in the application of the preceding theorems to a class of problems of wide practical interest.

### 5. Hertzian Contact Problems

We wish to characterize contact problems in which the contact surface is approximately planar and the bodies have undergone small deformations in the neighborhood of the contact surface.

Assume the following:

- (1)  $\underline{n} \approx n_i \underline{e}_i \approx \underline{e}_3$  on  $\mathcal{C}$ , where the  $n_i$  indicate components with respect to the standard basis  $\{\underline{e}_i\}_1^3$  for  $\mathbb{R}^3$ ,  
(see Fig. 5).

- (2)  $\underline{g}^\alpha \approx 1, \alpha = 1, 2$ , thus  $\underline{t}^\alpha \approx \underline{T}^\alpha$  on  $\mathcal{C}^\alpha$ .

Assumptions (1) and (2) together imply that,

$$\underline{t}_3^\alpha \approx \underline{t}^\alpha \cdot \underline{n} \approx \underline{T}^\alpha \cdot \underline{n} \approx T_3^\alpha,$$

and that,

$$(\underline{t}_1^\alpha, \underline{t}_2^\alpha, 0) \approx \underline{t}^\alpha - (\underline{t}^\alpha \cdot \underline{n})\underline{n} \approx \underline{T}^\alpha - (T_3^\alpha)\underline{n} \approx (T_1^\alpha, T_2^\alpha, 0).$$

- (3) Material points which eventually contact have, to the first order, the same initial coordinates  $\underline{z}_1$  and  $\underline{z}_2$ . Explicitly we manifest this idea by requiring that the  $\chi^\alpha$ 's satisfy

$$\chi^1(\underline{z}_1, \underline{z}_2, X_3^1(\underline{z}_1, \underline{z}_2)) = \chi^2(\underline{z}_1, \underline{z}_2, X_3^2(\underline{z}_1, \underline{z}_2)). \quad (39)$$

This is depicted in Fig. 6. Since  $X_3^\alpha$  are given functions which define the surfaces  $\mathcal{C}^\alpha$ , it follows from (39) that,

$$\delta \chi^1 = \delta \chi^2.$$

We term problems for which these assumptions hold Hertzian, since these assumptions are implicit in Hertz' classical theory [2] (see

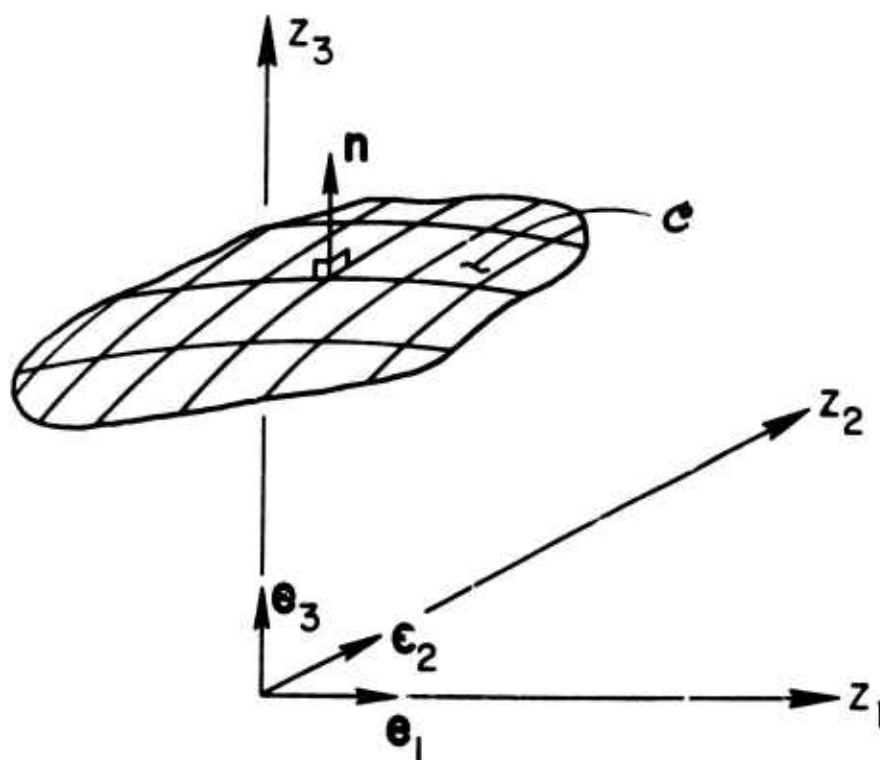


Figure 5

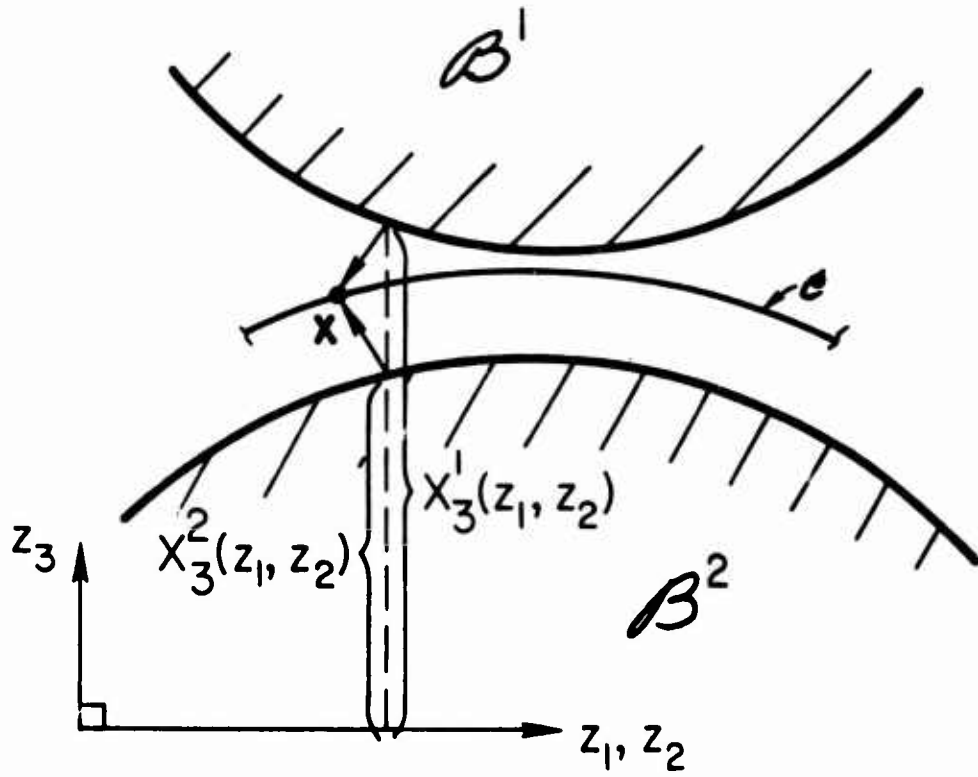


Figure 6

Goldsmith [3] for an excellent exposition of this work and also many applications of Hertz' theory to impact problems). It should be pointed out that the formulation we are about to give is still considerably more general than those to which Hertz' theory applies.

We now show how these assumptions allow us to make simplifications in the preceding theorems.

Theorems I and II:

Due to assumption (3) the term  $\mathcal{K}$  can be replaced by an integral over a region in the  $z_1, z_2$  -plane. This region, say  $c$ , is the projection of  $\mathcal{E}$  onto the  $z_1, z_2$  -plane, and due to assumption (2) it coincides, to the first order, with the projections of the  $\mathcal{E}^a$  's. Thus  $\mathcal{K}$  can be written

$$\mathcal{K} = - \sum_{a=1}^2 \int_0^t \int_c \underline{\tau}^a \cdot (\underline{x}^a - \underline{x}^a) \, dc \, dt . \quad (40)$$

Since, for Case I, we know that the momentum balance on  $\mathcal{E}$  requires that

$$\underline{\tau}^1 + \underline{\tau}^2 = \underline{0} ,$$

we may make use of this relation immediately. Thus define

$$\underline{\tau} = \underline{\tau}^1 = -\underline{\tau}^2 ,$$

and substitute into (40). Employing (39), the integrand simplifies to

$$\underline{\tau} \cdot (\underline{x}^2 - \underline{x}^1) . \quad (41)$$

The analog of (23) becomes

$$\begin{aligned}
0 = \int_0^t \int_C \left\{ \delta \tau \cdot (x^2 - x^1) + \right. \\
\left. + \delta x^1 \cdot (T^1 - \tau) + \delta x^2 \cdot (T^2 + \tau) \right\} dc dt \\
+ \text{transversality condition.}
\end{aligned} \tag{42}$$

Thus the same conclusions of Theorem I can be drawn. However, from a numerical standpoint things are considerably different. First of all, since the  $x^m$ 's are absent in this formulation, we do not get a uniquely defined  $\tau$ ;  $x^1$  and  $x^2$  will not in general be the same pointwise. If the graph of  $\tau$  is important it could be constructed by averaging  $x^1$  and  $x^2$ , which, if the solution is any good, should be reasonably close pointwise. On the other hand, the  $x^m$ 's being absent engenders a considerable saving in the number of equations to be solved and in their complexity.

The analogous case for Theorem II is constructed simply by setting

$$\tau_1 = \tau_2 = 0, \quad \tau \stackrel{\text{def.}}{=} \tau_3.$$

Then the integrand of  $\mathcal{K}$  becomes

$$\tau (x_3^2 - x_3^1) \tag{43}$$

and (42) reduces to

$$\begin{aligned}
0 = \int_0^t \int_C \left\{ \delta \tau (x_3^2 - x_3^1) + \right. \\
+ \delta x_3^1 (T_3^1 - \tau) + \delta x_3^2 (T_3^2 + \tau) \\
+ \delta x_m^1 T_m^1 + \delta x_m^2 T_m^2 \left. \right\} dc dt \\
+ \text{transversality condition.}
\end{aligned} \tag{44}$$

Hence the conclusions of Theorem II hold.

Thus in the case of Hertzian contact we can add the simplifications manifested in (41) and (43) to the conditions of Theorems I and II, respectively, and still garner the same conclusions.

Theorems I' and II':

For these cases  $\mathcal{M}$  can be written as an integral over  $\tilde{\mathcal{C}}$ , the projection of  $\tilde{\mathcal{E}}$ :

$$\mathcal{M} = - \int_{t_0}^t \int_{\tilde{\mathcal{C}}} \underline{\sigma}^{\alpha} \cdot \underline{n}^{\alpha} (\underline{x}^{\alpha} - \underline{x}^{\alpha'}) d\tilde{\mathcal{C}} dt.$$

Due to the present geometric situation, it is appropriate to take

$$\underline{n}^1 = \underline{n}^2,$$

and thus define

$$\underline{n} = \underline{n}^{\alpha}, \quad \alpha = 1, 2.$$

Analogous to the considerations for Theorems I and II, the momentum balance across  $\tilde{\mathcal{E}}$  motivates the simplification

$$\underline{\sigma}^{\alpha} \stackrel{\text{def.}}{=} \underline{\sigma}^{\alpha} = -\underline{\sigma}^{\beta}.$$

With these and (39), the integrand of  $\mathcal{M}$  can be written

$$\underline{\sigma} \cdot \underline{n} (\underline{x}^2 - \underline{x}^1).$$

A further simplification can be made by setting\*

$$\sigma_3 = -\underline{n}.$$

This eliminates one unknown function and, as we shall see, has the effect of satisfying (5) naturally. Thus the integrand of  $\mathcal{M}$  becomes

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\* This is a standard ploy of optimization theory, see p. 82, [4].



$$\sigma_a \eta (x_a^2 - x_a^1) - (\eta)^2 (x_3^2 - x_3^1) , \quad (45)$$

and the analog of (23) is

$$\begin{aligned} 0 = \int_0^t \int_{\bar{c}} \{ & \delta \eta (\sigma_a (x_a^2 - x_a^1) - 2\eta (x_3^2 - x_3^1)) \\ & + \delta \sigma_a (\eta (x_a^2 - x_a^1)) \\ & + \delta x_a^1 (T_a^1 - \eta \sigma_a) + \delta x_a^2 (T_a^2 + \eta \sigma_a) \\ & + \delta x_3^1 (T_3^1 + (\eta)^2) + \delta x_3^2 (T_3^2 - (\eta)^2) \} d\bar{c} dt . \end{aligned} \quad (46)$$

Summand two tells us that either  $\eta = 0$  or  $x_a^1 = x_a^2$ , on  $\bar{c}$ .

Suppose  $\eta \neq 0$ , then  $x_a^1 = x_a^2, a=1,2$ . Summand one then gives us that  $x_3^1 = x_3^2$  on  $\bar{c}$ . Thus we have

$$\eta (x^2 - x^1) = 0 , \quad \text{on } \bar{c} ,$$

as required, and  $\mathcal{E}$  is defined as the subset of  $\bar{c}$  where  $x^1 = x^2$ .

The last four summands give the momentum balance conditions, as usual, and, in addition, the last two summands imply that the normal tractions are compressive (since  $(\eta)^2 \geq 0$ ). Thus we have the conclusions of Theorem I' and condition (5).

The analogous set up for Theorem II' is accomplished by setting  $\sigma_a = 0$  in (45) yielding

$$- (\eta)^2 (x_3^2 - x_3^1) \quad (47)$$

for the integrand of  $\eta \eta$ . With this Eq. (46) becomes

$$\begin{aligned}
0 = \int_0^t \int_{\bar{c}} \{ & -2 \delta \tau (\tau (x_3^2 - x_3^1)) + \\
& + \delta x_a^1 T_a^1 + \delta x_a^2 T_a^2 \\
& + \delta x_3^1 (T_3^1 + (\tau)^2) + \delta x_3^2 (T_3^2 - (\tau)^2) \} d\bar{c} dt
\end{aligned}$$

In this case we achieve the conclusion of Theorem II' and condition (5).

Thus to Theorems I' and II' we can delete condition (5), add the simplifications manifested in (45) and (47), and achieve the conclusions of Theorems I' and II', respectively, plus condition (5).

## 6. Contact Problems for One, Two and Three-dimensional Bodies

The previous work needs only trivial modification to be made applicable to contact problems involving bodies of different dimensions. There are many cases of considerable interest which fall into this category. For example, models consisting of a shell and a plate, or a solid and a plate, are useful for the study of head impact. The modifications necessary are essentially interpretative. An example illustrates this assertion.

Consider the frictionless Hertzian contact of a three-dimensional solid and a two-dimensional plate. Let  $B^1$  represent the solid and  $B^2$  the plate. In evaluating  $\Pi$ , the  $B^1$  part is as before while the  $B^2$  part would manifest the particular plate theory used. The contact term  $\mathcal{K}$  (or  $\mathcal{K}$ ) would be exactly as before. However note that  $c$  (or  $\bar{c}$ ) is, in this case, also identifiable with part of the two-dimensional "volume" of the plate, rather than its boundary. Taking variations, everything is as before except that the term  $\tau \delta x_3^1$  (or  $-(\tau)^2 \delta x_3^2$ ) contributes to the transverse momentum equation of the plate, rather than to its boundary conditions. The interpretation of  $\tau$  (or  $-(\tau)^2$ ) is thus two-fold, i.e., it is the normal component of the traction vector with respect to  $B^1$ , as before, and it is also the equivalent normal "body force" with respect to  $B^2$ , manifested by the interaction with  $B^1$  (Fig. 7).

This interpretation is general, namely, for one and two-dimensional bodies the contact force is an equivalent "body force" which contributes to the momentum equations, rather than the boundary conditions. With this interpretation in mind, the construction of variational theorems, analogous to the ones constructed in Sections 2, 4 and 5, for the class of one, two and three-dimensional contact problems, is just a formal deductive exercise involving only appropriate definitions for  $\Pi$ .

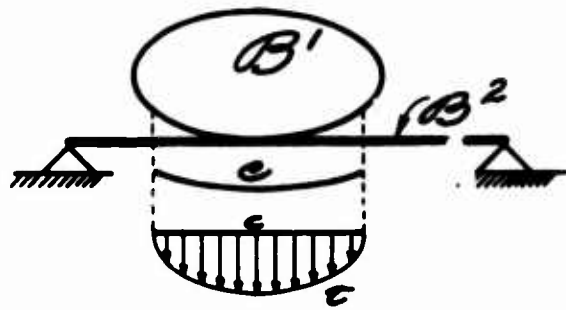


Figure 7

## 7. Impact

The previous sections deal with spatial aspects of contact problems. In this section we investigate temporal considerations, i.e., those phenomena which are unique to dynamic contact or impact. To manifest the problem encountered in such situations consider the following hypothetical situation. Assume that we are in the process of numerically solving some impact problem and suppose that it is discovered as we monitor the motion of the bodies that they impact somewhere in the time interval  $(t_1, t_2)$ . At time  $t_1$  we know the states of both bodies and we know that somewhere between  $t_1$  and  $t_2$  they have coalesced over a portion of their boundaries. Assume for the moment we know the geometry of the contact surface  $e$ . The question which arises then is what is the state of  $e$  at time  $t_2$ , i.e., what are the velocity and traction vectors on  $e$ ? It is necessary to know this information to carry forth the step forward time integration. The question though seems improperly posed without specifying considerable data about the nature of the impact. To get a handle on things, we will initially formulate a simple one-dimensional problem involving the impact of two elastic rods. Although this problem is trivial, it provides considerable insight into the general nature of impact of continuum bodies. Since we are interested in the state of  $e$  (in this case a point) immediately after impact, whether the rods are finite or semi-infinite is immaterial.

Assume that the pre-impact states of the two bodies are given by the following data:

$$B^{\alpha} \quad v_{-}^{\alpha} \quad , \quad (\partial x / \partial X)_{-}^{\alpha} \quad , \quad P_{-}^{\alpha} \quad ; \quad \alpha = 1, 2 \quad . \quad (48)$$

At impact the rods coalesce at  $e$ , and for some finite time interval thereafter (at least)  $x \in e$  is persistent. At the moment of impact shock waves begin to propagate in each body. The space-time picture is depicted in Fig. 8. As discussed in section 1, since  $e$  is material and  $x$  is persistent, we have

$$\underline{v} \stackrel{\text{def}}{=} \underline{v}_+^1 = \underline{v}_+^2, \quad P \stackrel{\text{def}}{=} T_+^1 = -T_+^2, \quad (49)^*$$

for the post-impact state ( $t_+$ ). In addition to (49), the well known shock conditions must hold across the wave fronts:

$$\begin{aligned} [\underline{v}^{\alpha}] + U^{\alpha} [(\partial x / \partial X)^{\alpha}] &= 0, \\ \rho_0^{\alpha} U^{\alpha} [\underline{v}^{\alpha}] &= [P^{\alpha}], \end{aligned} \quad (50)$$

where  $U^{\alpha}$  is the material velocity of the shock in  $\mathcal{B}^{\alpha}$ , and  $[ \ ]$  is the wave-front jump operator which assigns to a function the difference in its values behind and in front of the wave, i.e.,  $[f(X,t)] = f(X^-,t) - f(X^+,t)$  where  $X$  is a material point denoting the location of the wave-front. As can be deduced from Fig. 8, the states into which the shocks initially propagate are the pre-impact states given by (48), and the state at  $e$ , immediately after the shocks pass, is given by the post-impact state (49). These observations in conjunction with (50) yield,

$$\begin{aligned} \underline{v}_-^{\alpha} - \underline{v} + U^{\alpha} \{ (\partial x / \partial X)_-^{\alpha} - (\partial x / \partial X)_+^{\alpha} \} &= 0, \\ \rho_0^{\alpha} U^{\alpha} (\underline{v}_-^{\alpha} - \underline{v}) + P_-^{\alpha} - P &= 0. \end{aligned} \quad (51)^{**}$$

\*For convenience we choose the initial state to be the pre-impact state, thus we need not distinguish between Cauchy and Piola tractions.

\*\*A consistency condition for these equations is that  $\underline{v}_-^1 - \underline{v}_-^2 > 0$ . Otherwise the impact would not occur.

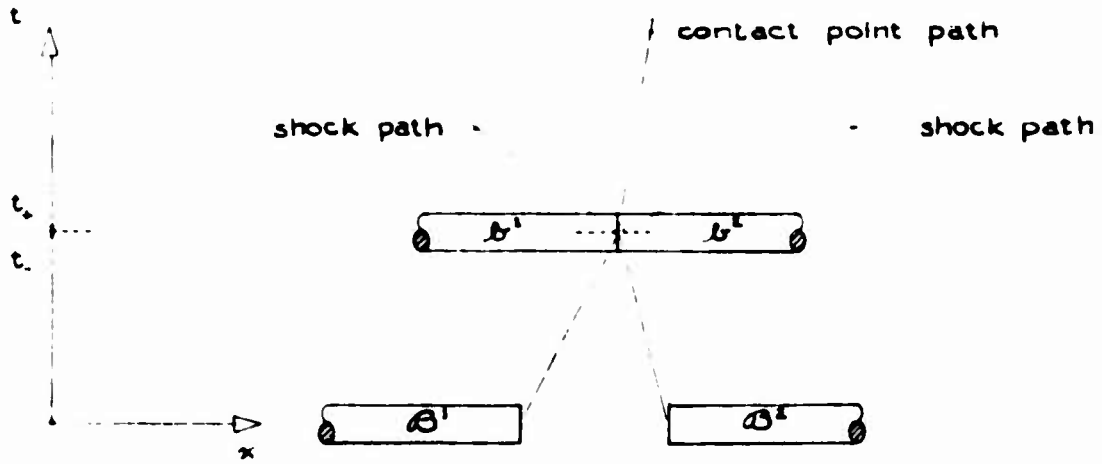


Figure 8

The four Eqs. (51) and constitutive equations relating  $P^\alpha$  to  $(\partial x / \partial X)^\alpha$  yield a formally deterministic system of six equations in the six unknowns  $v$ ,  $P$ ,  $U^\alpha$ ,  $(\partial x / \partial X)^\alpha$ . Thus we see that the desired quantities  $v$  and  $P$  depend on the pre-impact states and material properties of both  $B^\alpha$ . The precise form of this relationship depends upon the constitutive equations of the bodies. As a simple example, assume we have linear constitutive equations  $P^\alpha = E^\alpha \{ (\partial x / \partial X)^\alpha - 1 \}$ ,  $E^\alpha$  constant, and let the pre-impact state be given by

$$\begin{aligned} v_-^\alpha &= V^\alpha, \\ (\partial x / \partial X)_-^\alpha &= 1, \\ P_-^\alpha &= 0. \end{aligned} \quad (52)$$

These conditions, when inserted in Eqs. (51), lead to:

$$\begin{aligned} v &= \frac{\rho_0^2 U^2 V^2 - \rho_0^1 U^1 V^1}{\rho_0^2 U^2 - \rho_0^1 U^1}, \\ P &= \frac{V^2 - V^1}{\left( \frac{U^2}{E^2} - \frac{U^1}{E^1} \right)}, \\ (U^\alpha)^2 &= E^\alpha / \rho_0^\alpha. \end{aligned} \quad (53)$$

Note that the denominators in Eqs. (53)<sub>1,2</sub> present no problems since  $\rho_0^\alpha > 0$ ,  $E^\alpha > 0$  and  $+1 = \text{sgn } U^2 = -\text{sgn } U^1$ .

This result is also appropriate whenever the intensity of the impact is small enough such that the non-linear constitutive equation can be replaced by its linear approximation about the pre-impact state. In this case  $E^\alpha$  is a tangent modulus evaluated at the pre-impact strain

$\{ (\partial x / \partial X)_-^\alpha - 1 \} = 0$ . To further simplify, consider the case when both rods have identical properties (i.e.,  $\rho_0 = \rho_0^\alpha$ ,  $E = E^\alpha$ ,  $\alpha = 1, 2$ ). Then



$$\begin{aligned}
 v &= \frac{V^1 + V^2}{2} , \\
 P &= \rho U (V^2 - V^1) / 2 , \\
 (U)^2 &= E / \rho .
 \end{aligned}
 \tag{54}$$

In Eqs. (54)  $U$  is positive, and since consistency requires  $V^1 - V^2 > 0$ ,  $P$  is compressive.

Thus for the one-dimensional case at least the problem of computing the post-impact state is easily achieved. The solution of (51) for the fully non-linear case can be automated as part of a numerical algorithm. Although this problem is trivial, it serves to indicate that the post-impact problem, the solution of which is essential in a numerical algorithm, is one of wave propagation.

In the analysis of higher dimensional bodies the solution of the post-impact problem becomes greatly complicated due to the geometric variety of impact conditions. However, considerable simplifications can be taken advantage of if one keeps in mind the nature of the discrete problem. For instance, if a certain portion of the boundaries of two bodies have coalesced in  $e$ , each interior point of  $e$ , at which the tangent plane is well defined, may be treated, to the first order, as a point on the mating surface of two impacting half-spaces. As long as time steps are kept small enough, the local behavior is well represented. The post-impact problem for the general case, analogous to (51), can be automated as part of the numerical algorithm, and for many simple cases can be solved explicitly.

With these notions in mind, let us return to the case of main interest in this report, namely three-dimensional continuum bodies. We shall consider only the case of a frictionless contact surface (Case II), and leave the solution of the post-impact problem for the no-slip case (Case I), which is more difficult, for future work. With the proper interpretations, the one-dimensional rod formulation (Eqs. (48-54)) suffices to completely characterize this case. This is so because no tangential motions or stresses may be communicated across a frictionless surface, and thus we need only consider the configuration of normal incidence. In this case the requisite constitutive functionals in (51) would be those relating  $P^\alpha$ , the normal Piola stress, to the normal component of strain, holding all other components of strain fixed at the pre-impact values. For example, in the linear isotropic case,  $E^\alpha$  (Young's modulus) in Eqs. (53,54) would be replaced by  $\lambda^\alpha + 2\mu^\alpha$  ( $\lambda^\alpha, \mu^\alpha$  are the Lamé and shear moduli, respectively) and the propagation velocity would be that of dilatational waves.

## PART II

A NUMERICAL SCHEME FOR ANALYSIS OF  
CONTACT-IMPACT PROBLEMS8. Numerical Solution of Contact-Impact Problems

In performing numerical computations based on the above described variational formulation for contact-impact problems we have employed three distinct levels of approximation: (1) a spatial discretization of the bodies and contact surfaces, (2) a temporal discretization to determine the response of the discretized bodies, and (3) a numerical solution for the resulting system of nonlinear algebraic equations.

In the following sections we shall restrict our attention to the Hertzian contact problem described in Section 5. Significant numerical difficulties are encountered in the solution of impact problems; to complicate the problem further by introducing the additional steps necessary to determine the contact surface maps for the full kinematically nonlinear case is left for a future study. While this is a simple impact problem in terms of determining the contact surface and the full power of the preceding theory is neither necessary nor exploited in its solution, many of the features of the general problem are employed here.

### 9. Spatial Discretization of the Bodies and Contact Surface

The bodies  $\mathcal{B}^1$  and  $\mathcal{B}^2$  are discretized using standard finite element methods, (e.g., see [5]). In order to facilitate the computation of a discrete Hertzian contact surface the nodes of  $\mathcal{B}^1$  are arranged so that they align with the nodes of  $\mathcal{B}^2$ . This is consistent with the notions of condition 3 of Section 5 and ensures that during determination of the approximation to the contact surface contiguous nodes of the two bodies will meet. Thus, the simulation of the contact surface is trivial. The development of a numerical model for Hertzian contact problems is based upon the form of Theorem II' which uses (47) for the integrand of  $\mathcal{M}$ . For numerical computations we introduce the displacement vector  $\underline{u}$  such that

$$\underline{x} = \underline{X} + \underline{u} \quad (55)$$

For a compatible finite element displacement field the integrand of  $\mathcal{M}$  can be approximated by taking  $\mathcal{M}^2(\underline{x}, t)$  as the product of  $\mathcal{E}^2(t)$  and  $\delta(\underline{x} - \underline{x}_i)$  (i.e., Dirac delta functions in space). This corresponds to taking  $\mathcal{M}^2$  as "concentrated nodal loads" which are the generalized forces of the contact pressure. With this discretization we can describe pseudo contact elements between each pair of candidate contact nodes. Let these nodes be denoted as  $(i)$  and the generalized force as  $(\mathcal{E}_i)^2$ ; then

$$\mathcal{M} = \int_0^t \sum_i (\mathcal{E}_i(t))^2 (u_{3i}^2(t) - u_{3i}^1(t) + X_{3i}^2 - X_{3i}^1) dt \quad (56)$$

where  $\{i\}$  are the set of candidate contact nodes which span  $\tilde{\mathcal{C}}$ ;  $u_{3i}^m$  are the nodal displacements in the  $z_3$  direction and  $X_{3i}^m$  are the nodal coordinates of the candidate contact nodes.\*

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\*We assume here that 3 is the direction nominally normal to the contact surfaces, e.g., see Fig. 5.

Use of the finite element method in Theorem II' with  $\mathcal{M}$  given by (56) produces a set of nonlinear second order ordinary differential equations which together with the impenetrability conditions define the discretized contact impact problem. These equations take the form:

$$\underline{\underline{M}} \ddot{\underline{u}} + \underline{\underline{K}}(\underline{u}) = \underline{\underline{R}} \quad , \quad (57)$$

where  $\underline{\underline{M}}$  is the usual finite element mass matrix,  $\underline{\underline{K}}$  represents the elastic stiffness forces together with the contact terms,  $\underline{\underline{R}}$  is the set of generalized forces resulting from boundary tractions and  $\underline{u}$  is the set of time dependent nodal displacements (which also include the  $(\epsilon_i)^2$ ). For inelastic materials Theorem II' can be extended by treating the first variation as a Galerkin method (principle of virtual work) and replacing the elastic constitutive model by more general theories, e.g., viscoelastic, elastoplastic, viscoplastic, etc. In this case

$$\underline{\underline{K}}(\underline{u}) \rightarrow \underline{\underline{K}}(\underline{u}, \dot{\underline{u}}) \quad (58)$$

in (57).

## 10. Temporal Discretization

A temporal discretization of the second order ordinary differential equations which result from a finite element spatial discretization of the contact-impact problem is accomplished herein by using the Newmark family of methods [6]. The Newmark family of methods is a one-step integration method with two free parameters which can be used to control stability and numerical damping. The method is essentially a difference method in time. The behavior of the method for linear elasto-dynamics problems is discussed in [6,7]. The algorithm is given by

$$\underline{u}_{n+1} = \underline{u}_n + \Delta t \underline{\dot{u}}_n + \left(\frac{1}{2} - \beta\right) \Delta t^2 \underline{\ddot{u}}_n + \beta \Delta t^2 \underline{\ddot{u}}_{n+1}, \quad (59)$$

$$\text{and } \underline{\dot{u}}_{n+1} = \underline{\dot{u}}_n + (1-\gamma) \Delta t \underline{\ddot{u}}_n + \gamma \Delta t \underline{\ddot{u}}_{n+1},$$

where  $\underline{u}_n = \underline{u}(t_n)$ ,  $\Delta t = t_{n+1} - t_n$ , and  $\beta, \gamma$  are the two parameters. For linear problems  $\gamma = .5 + \delta = .5$  produces no artificial viscosity and  $\beta \geq \frac{1}{4} (1 + \delta)^2$  produces unconditional stability (i.e., the method is stiffly stable). Such generalization is not possible for nonlinear problems and during solution it may be necessary to monitor the solution for any signs of instability. In (59)  $\beta = 0$  produces an explicit method for  $\underline{u}_{n+1}$  and if  $\underline{M}$  is diagonal (lumped mass) with  $\underline{K}$  and  $\underline{R}$  independent of  $\underline{\dot{u}}$  the solution can be advanced without solving a large set of simultaneous equations; for all other cases the method is implicit and equations must be solved. Solution of (59)<sub>1</sub> for  $\underline{\ddot{u}}_{n+1}$  in terms of the solution at  $t_n$  and  $\underline{u}_{n+1}$  gives

$$\underline{\ddot{u}}_{n+1} = \frac{1}{\beta \Delta t^2} (\underline{u}_{n+1} - \underline{u}_n) - \frac{1}{\beta \Delta t} \underline{\dot{u}}_n - \left(\frac{1-2\beta}{2\beta}\right) \underline{\ddot{u}}_n \quad (60)$$

which can also be used in  $(59)_2$  to express the velocity in terms of the solution at  $t_n$  and  $\underline{u}_{n+1}$ . Since in this process we divide by  $\beta$  and  $\Delta t$  it is no longer possible to consider zero  $\beta$  or zero time steps.

### 11. Solution of the Nonlinear Algebraic Problem

Use of the Newmark method in (57) (including (58)) yields the set of nonlinear algebraic equations:

$$\frac{1}{\beta \Delta t^2} \underline{M} \underline{u}_{n+1} + \underline{K}(\underline{u}_{n+1}, \underline{u}_n, \dot{\underline{u}}_n, \ddot{\underline{u}}_n) = \underline{R}_{n+1} + \underline{M} \underline{A}_n, \quad (61)$$

where

$$\underline{A}_n = \frac{1}{\beta \Delta t^2} \underline{u}_n + \frac{1}{\beta \Delta t} \dot{\underline{u}}_n + \left( \frac{1-2\beta}{2\beta} \right) \ddot{\underline{u}}_n.$$

A Newton-Raphson iterative solution to this set of equations can formally be constructed, giving:

$$\left( \frac{1}{\beta \Delta t^2} \underline{M} + \partial_u \underline{K} - \partial_u \underline{R} \right) \Delta \underline{u}^{(i)} = \underline{R} - \underline{K}(\underline{u}_{n+1}^{(i)}, \underline{u}_n, \dot{\underline{u}}_n, \ddot{\underline{u}}_n) - \underline{M} \ddot{\underline{u}}_{n+1}^{(i)}, \quad (62)$$

where  $\partial_u \underline{R}$  is the effect of loads varying with the deformation and

$$(\partial_u \underline{K})_{ij} = \partial K_i / \partial u_j, \quad (63)$$

is the tangent stiffness matrix. The coefficient to  $\Delta \underline{u}^{(i)}$  is generally called the Jacobian matrix of the Newton-Raphson iteration. The solution is advanced by taking

$$\underline{u}_{n+1}^{(i+1)} = \underline{u}_{n+1}^{(i)} + \Delta \underline{u}^{(i)}, \quad (64)$$

and iterating until a norm of the solution satisfies

$$\|\Delta \underline{u}^{(i)}\| \leq \epsilon \|\underline{u}_{n+1}^{(i)}\|, \quad (65)$$



where  $\epsilon$  is some small positive error tolerance. In the work reported here the norm  $\| \cdot \|$  is taken as the Euclidian norm

$$\| \underline{x} \| = \left( \sum_i x_i^2 \right)^{1/2}, \quad (66)$$

and the load vector  $\underline{R}$  is assumed to be independent of  $\underline{u}$ . For stable elastic materials the resulting tangent stiffness is then symmetric and positive definite, consequently, standard direct solution methods normally employed in the solution of linear finite element problems can be used. For inelastic materials or deformation dependent loads the tangent stiffness may be asymmetric. In these cases some special methods may be necessary to effect a solution.

## 12. Discretized Impact Conditions

In the previous numerical development  $\tilde{c}$  has been defined by discrete points which correspond to nodes along the boundaries of  $B^1$  and  $B^2$ . When, during the course of advancing the solution in time, any one of these points violates the impenetrability condition a re-solution must be obtained in which the  $(\epsilon_i)^2$  are now non-zero and the  $u_3^*$  satisfy the impenetrability condition. Some control and monitoring are required to effect this in a computer program. In addition to satisfying these conditions, the impact relations denoted in Section 7 must be invoked. In the present study these conditions are applied to the solution at the end of a time step in which points first go into contact. Accordingly we compute from (50)\*

$$\dot{u}_+ = \frac{e_2^2 U^2 \dot{u}_-^2 - e_1^1 U^1 \dot{u}_-^1}{e_2^2 U^2 - e_1^1 U^1}, \quad (67)$$

and assign this value to the appropriate node of  $B^1$  and  $B^2$ .

To determine the solution vector  $\underline{u}$  at  $t_{n+1}$ , we have solved the set of equations (61). As described above the shock conditions are then used to determine the value of the velocity at time  $t_{n+1}$  for all points which have come into contact during the time interval. In order to get a consistent solution at these points we must modify the accelerations and contact force to reflect the shock conditions. This is accomplished by re-solving the equilibrium conditions of  $B^1$  and  $B^2$  at point  $i$ . The expanded forms of the appropriate equations are:

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\*The  $( )_-$  denotes a value which is computed before impact, whereas  $( )_+$  denotes the value after impact.

$$M^1 \ddot{u}_-^1 + K^1(\underline{u}) + (\mathcal{E}_i)_-^2 = R^1, \quad (68)$$

and

$$M^2 \ddot{u}_-^2 + K^2(\underline{u}) + (\mathcal{E}_i)_-^2 = R^2.$$

For nodes which have come into contact we must enforce the condition on acceleration

$$\ddot{u}_+^1 = \ddot{u}_+^2 = \ddot{u}_+, \quad (69)$$

and compute the contact force  $(\mathcal{E}_i)_+^2$ . The solution for these is obtained from

$$M^1 \ddot{u}_+ + K^1(\underline{u}) + (\mathcal{E}_i)_+^2 = R^1,$$

and

$$M^2 \ddot{u}_+ + K^2(\underline{u}) + (\mathcal{E}_i)_+^2 = R^2.$$

These are two equations in two unknowns which can be solved for the  $\ddot{u}_+$  and  $(\mathcal{E}_i)_+^2$ . If  $K^{\alpha}(\underline{u})$  is independent of velocity the stiffness forces and  $R^{\alpha}$  will remain unchanged during the impact, hence we can solve the simpler problem

$$M^1 \ddot{u}_+ + (\mathcal{E}_i)_+^2 = M^1 \ddot{u}_-^1 + (\mathcal{E}_i)_-^2$$

$$M^2 \ddot{u}_+ - (\mathcal{E}_i)_-^2 = M^2 \ddot{u}_-^2 - (\mathcal{E}_i)_-^2$$

whose solution is

$$\ddot{u}_+ = \frac{M^1 \ddot{u}_-^1 + M^2 \ddot{u}_-^2}{M^1 + M^2},$$

and

$$2(\mathcal{E}_i)_+^2 = 2(\mathcal{E}_i)_-^2 + M^1 (\ddot{u}_-^1 - \ddot{u}_+) - M^2 (\ddot{u}_-^2 - \ddot{u}_+).$$

This completes the numerical specification of the solution at  $t_{n+1}$ ; this solution process is now repeated for each of the succeeding time steps.

At this point it is important to compare the solution procedure for impact of a continuum discretized by a finite element method with the solution procedure for a physically discrete body, i.e., a body composed of mass points joined by massless elastic springs. Both problems may be described by algebraic equations of the form of (57). The impenetrability condition is also identical. The impact conditions, however, are different. For the discretized continuum the procedure is described above. The study of the impact of mass points is considered in elementary mechanics books, e.g. [8]. The impact of two mass points is described by impulsive motion such that at  $t_-$  the velocities of the two mass points are  $V_-^1$  and  $V_-^2$ ; after impact at time  $t_+$ , the two points have velocities  $V_+^1$  and  $V_+^2$ . The two points will not in general stay in contact (i.e.,  $V_+^1 \neq V_+^2$ ) but will rebound. The conditions used to compute the  $V_+^1$  and  $V_+^2$  are:

Balance of Momentum\*

$$M^1 \{V^1\} + M^2 \{V^2\} = 0, \quad (71)$$

and use of an equation involving the "coefficient of restitution",  $e$ :

$$\frac{V_+^2 - V_+^1}{V_-^1 - V_-^2} = e. \quad (72)$$

For  $e=1$  energy is conserved whereas for  $e=0$  the points "stick" and energy is dissipated. We must comment in passing that (72) is the energy

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\*  $\{f(t)\} = f(t_+) - f(t_-)$ .

equation in disguise. To see this we can write the jump conditions for energy as

$$\frac{1}{2} M^1 \{(\dot{v}^1)^2\} + \frac{1}{2} M^2 \{(\dot{v}^2)^2\} = \{\dot{v}\} \quad . \quad (73)$$

The term  $\{\dot{v}\}$  can exist only if other energies are dissipated during the jump. We rewrite (73) by using

$$\frac{1}{2} \{(\dot{v}^1)^2\} = [\dot{v}^1] \langle \dot{v}^1 \rangle \quad ,$$

where

$$\langle \dot{v}^1 \rangle = \frac{1}{2} (\dot{v}_+^1 + \dot{v}_-^1) \quad . \quad (74)$$

Use of the momentum equation (71) then gives, after dividing by  $M^1 \{\dot{v}^1\}$

$$\langle \dot{v}^1 \rangle - \langle \dot{v}^2 \rangle = \frac{\{\dot{v}\}}{M^1 \{\dot{v}^1\}} \quad ,$$

or after recollecting terms and dividing by  $(\dot{v}_-^1 - \dot{v}_-^2)$  we obtain:

$$\frac{\dot{v}_+^2 - \dot{v}_+^1}{\dot{v}_-^1 - \dot{v}_-^2} = 1 - \frac{\{\dot{v}\}}{M^1 \{\dot{v}^1\} (\dot{v}_-^1 - \dot{v}_-^2)} \quad . \quad (75)$$

The significance of the coefficient of restitution then is associated with the right hand side of (75).

It is clear from the above developments that the numerical simulation of the discretized continuum and the physically discrete system involve two distinct methods for treating the impact conditions. It is imperative then to associate the correct method for the problem at hand. In the

present study we are interested in the impact of continua, and in this case we shall employ the discrete shock condition to effect the solution. This a priori assumes that the response we are computing involves a time scale associated with wave propagation problems. Consequently, we cannot expect the computation procedure for advancing the solution in time to be accurate if we take time steps greatly in excess of transit times through each body. In this context it may be important to consider an "explicit" time integration procedure in future work. The stability restrictions may be too severe to make this feasible.

## PART III

FEAP 74 - A COMPUTER PROGRAM FOR  
SOLUTION OF CONTACT-IMPACT PROBLEMS13. Development of a Contact-Impact Model for FEAP

In order to incorporate an ability to compute solutions to contact-impact problems using a finite element method as described above it is necessary to have available a computer program which can solve the nonlinear equations of motion given by (61). The computer program FEAP is a general program to solve finite element problems. The program has a capability of solving both quasistatic and dynamic problems and can incorporate several types of elements simultaneously. The nonlinear capabilities required for the solution of contact-impact problems have been incorporated into FEAP and currently includes the user options (see Input Instructions):

- (1) Selection of quasistatic or dynamic option: The dynamic option will integrate the equations of motion using the one-step Newmark method to advance the solution in time. Quasistatic analysis is accommodated by any one-step algorithm. The algorithm employed is incorporated into each element routine and thus is defined by the developer of each element. Impact problems require description of the contact surface and wave speeds.
- (2) Selection of the nonlinear method to advance the solution: Options include:
  - (a) No iterations in each time step. Unbalanced forces at each time are added to the next time step.
  - (b) Iterations in each time step to achieve a balance of force within each time step. In this option the user can select to reform

the Jacobian matrix for each iteration or only at the first iteration in each time step.

In the impact problems solved to date it has been necessary to use the general form of the Newton-Raphson algorithm. This includes a complete forming and factoring of the Jacobian matrix for each iteration of each time step in the analysis. If the method described herein is to become computationally effective improvements in the computer program are paramount. Undoubtedly the most important aspect in reducing computer times is to introduce a substructuring system so that the highly nonlinear equations in the vicinity of the contact surface can be isolated from the remainder of the bodies. This will normally involve only a small number of equations in the total system of (62). The solution of a large finite element problem will generally concentrate the computer solution time in the forming and factoring of the tangent stiffness matrix. The fewer times that it is necessary to perform this costly step the more efficient the solution algorithm. Substructuring can be used then to restrict the part of the equations which must be formed and factored often, and thus greatly reduce the computer costs in analyzing impact problems.

The version of FEAP which can currently be used to analyze contact-impact problems includes, in addition to the nonlinear Newton-Raphson iterative algorithm, a new special contact-impact element and a new subroutine to describe impact surfaces and the discrete shock conditions described in Section 12. These are described in the following sections.



#### 14. Contact Element for Hertzian Contact

The contact-impact element which has been developed is called ELMT05 and can be used along any coordinate direction. As developed it cannot be used along normals which are in non-coordinate directions. The development of the contact element assumes that within the framework of linear elasticity theory a node on  $\omega^1$  will impact on a node of  $\omega^2$ . In using this contact element we shall assume that the contact surface on  $\omega^2$  is located at larger coordinate values than the contact surface of  $\omega^1$ . The contact element is described by three nodes. Node 1 is associated with  $\omega^1$ , Node 3 is associated with  $\omega^2$ , and Node 2 is used as storage for the contact force  $(\mathcal{E})^2$ . The user can select the direction of contact motion by specifying the degree of freedom of the nodal unknowns to which the contact is to be measured; this must agree with the physical direction of the element (see Fig. 9). The degree of freedom for the contact element is specified during the MATERIAL data input and consists of a single card in I5 format. The element nodes are described along with all other elements according to Section 4 of the Input Instructions. The Node order as shown in Fig. 9 must be observed.

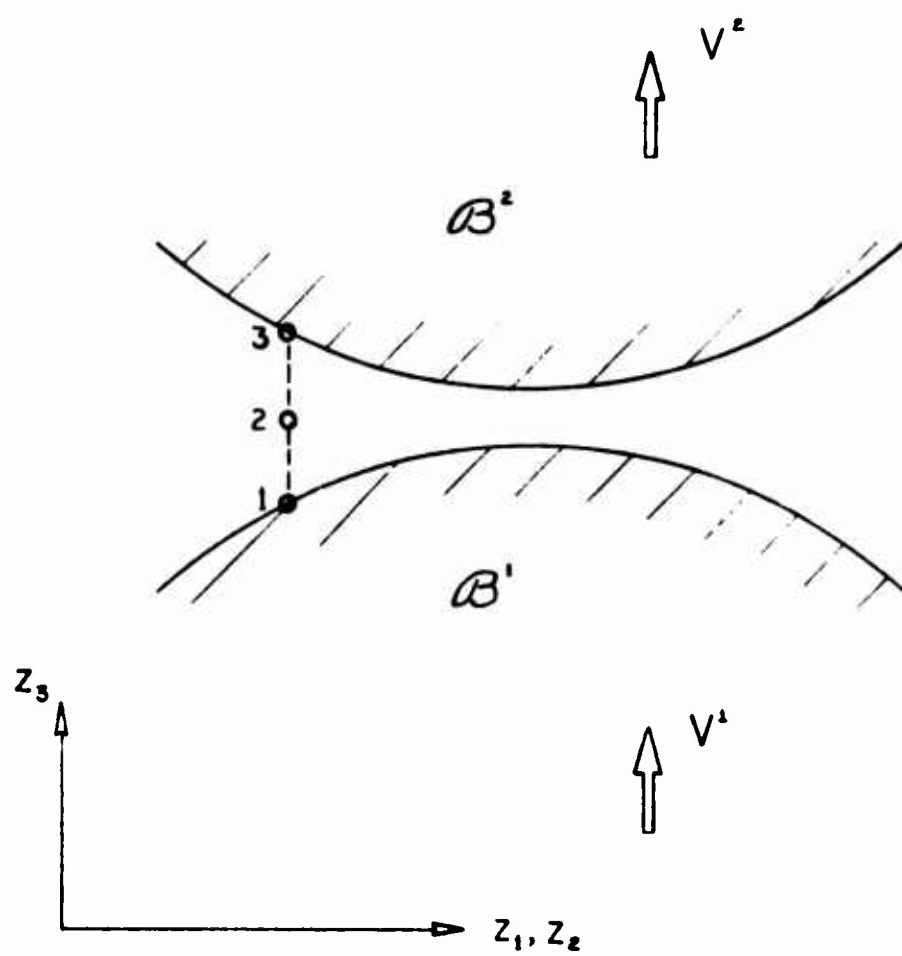


Figure 9

### 15. Impact Surface Description

The definition of the impact surface includes a list of all elements on the contact surface together with the degree of freedom describing the direction of contact motion (as described above). In addition, the product of mass density and wave velocity (always a positive number) for each body is input. This assumes, currently, that (1) each contact surface belongs to a linear material, and (2) the same material exists along all of the contact surface. This data need be prescribed only for impact problems, quasistatic contact problems do not require this data since no velocity or acceleration computations are performed in this class of problems. Data to be input for the impact surface is given in Table I.

Table I - Impact Surface Data

CARD 1) (6X,A6)		
COL. 7 to 12		Must contain CONTACT
CARD 2) (2F10.0)		
COL. 1 to 10	$\rho U$	of body 1
COL. 11 to 20	$\rho U$	of body 2
CARD 3) (I5)		
COL. 1 to 5		NLIST, number of elements on contact surface
CARD 4) (2I5)		
Repeat NLIST times		
COL. 1 to 5		Contact element number
COL. 6 to 10		Degree of freedom of this contact element

## 16. Example Problems

Two example problems are included to illustrate the characteristics of the methodology and the associated computer program described above for Hertzian contact problems. The first problem is a quasistatic contact problem which is used to demonstrate the ability of the computer program to compute an evolving contact surface. The second problem will demonstrate the ability of the program to properly model the temporal response of an impact problem.

To model a problem in which a contact surface will change under different load levels we consider two beams with an initial parabolic curvature. A symmetric configuration is analyzed and the resulting finite element model is shown in Fig. 10. Each element is nominally one unit by one unit. The gap at the load end is initially 0.5 units. The material properties used are  $E = 500$  and  $\nu = 0$ . The load  $P$  is applied as shown and allowed to increase linearly in time. The problem then is to determine the contact surface at various load levels. In order to eliminate a singularity in the system of equations it was necessary to permanently attach the two nodes at the symmetry axis of the contact surface. All other nodes along the boundaries between the two bodies are assumed to be possible contact points and contact elements are assigned between vertical nodal pairs. The load was varied from 0.2 to 0.8 in increments of 0.1 and the computed contact surface and forces were computed. These are given in Table II. The deformed shape at a load of 0.4 is also shown in Fig. 10 as a dotted form. The attached node at the center has influenced the solution at loads above 0.3 since the contact pressure there is tensile (negative). The force is small and should not greatly affect the actual contact region computed. As the load increases the

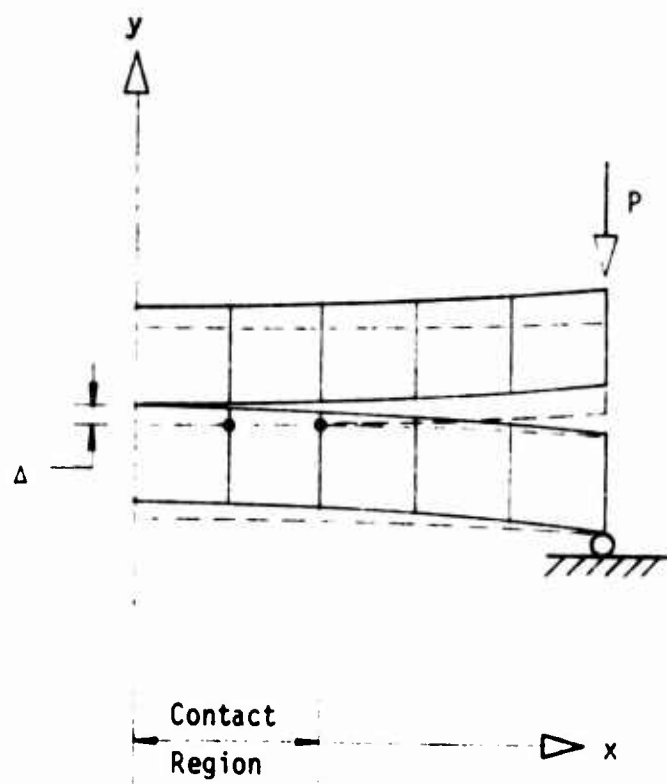


Figure 10

contact surface moves toward the load. This is conceptually correct since if the beams were modeled according to Euler-Bernoulli theory the contact force would be a point load which gradually moves from the center to the outer edge according to the relationship (using the above values for sizes and material properties)

$$x = 5 \left( 1 - \frac{1}{6P} \right) .$$

This relation predicts that the contact point will be non-zero only after P exceeds 1/6. The finite element model is in qualitative agreement with this beam theory, but since shear deformations are included the finite element solution gives a distributed load on the contact surface. It is interesting to also note that the contact force over the center of the beams is zero, just as in the beam theory.

Table II - Contact Forces

LOAD	X-COORDINATE						BEAM THEORY-X
	0	1	2	3	4	5	
0.2	0.2	-	-	-	-	-	0.83
0.3	.07	.23	-	-	-	-	2.22
0.4	-.01	.07	.34	-	-	-	2.92
0.5	-0.00	.02	.23	.25	-	-	3.33
0.6	-.01	-	.09	.51	.01	-	3.61
0.7	-.01	-	.07	.45	.19	-	3.81
0.8	-.01	-	.05	.38	.38	-	3.96

This problem demonstrates that the computer program can model the evolution of a contact surface. Of particular importance is to note that as the load increases the program can both attach and detach a contact point. This is an essential requirement for the analysis of

impact problems as is shown in the next problem.

As a simple example we consider the impact against a rigid wall of a finite, linear elastic rod traveling at constant velocity. The rod has a modulus of elasticity  $E$  of 100, and a mass density  $\rho$  of 0.1. The arrival velocity is taken to be 0.1 (units may be assigned in any convenient system). The rod is taken to be 10 units long and is divided into 10 elements plus one contact element as shown in Fig. 11. At time zero the rod is just arriving at the wall. The exact solution predicts a contact duration of 0.2 time units. This corresponds to the time required for a wave to travel from the contact point to the left end and back to the contact point at which time the rod will part from the wall. The problem was analyzed using FEAP with time steps of 0.01 unit (transit time across an element) and the rod remains in contact until time 0.20 units and has rebounded at time 0.21. Thus the program can predict accurately the contact duration of the rod. The finite element solution obtained is compared with the exact solution in Fig. 12. The agreement of stresses and contact force is good. The largest discrepancy exists in defining the shock front, which is "smeared" by the finite element method and ordinary differential equation solution method used here. This is the same type of solutions which are commonly obtained with numerical solutions of this type even without impact. Solutions such as the impact shocks generated are probably one of the most difficult responses to accurately calculate by a finite element method.

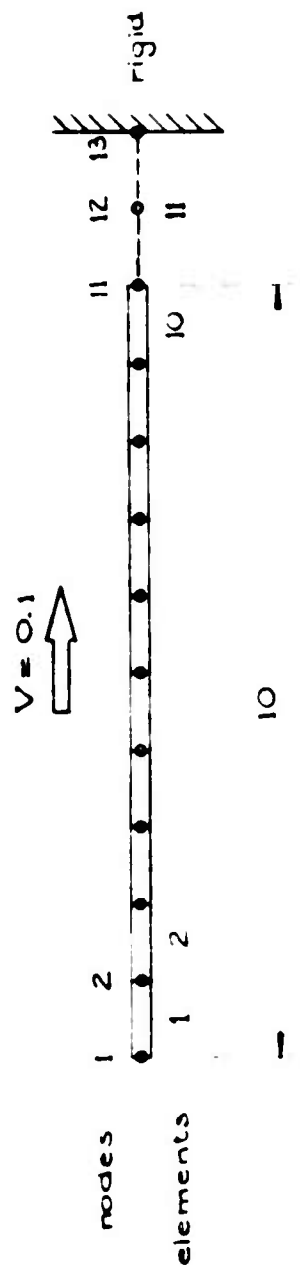


Figure 11



Comparison of Numerical Data with Exact Solution for Bar Impacting Rigid Wall

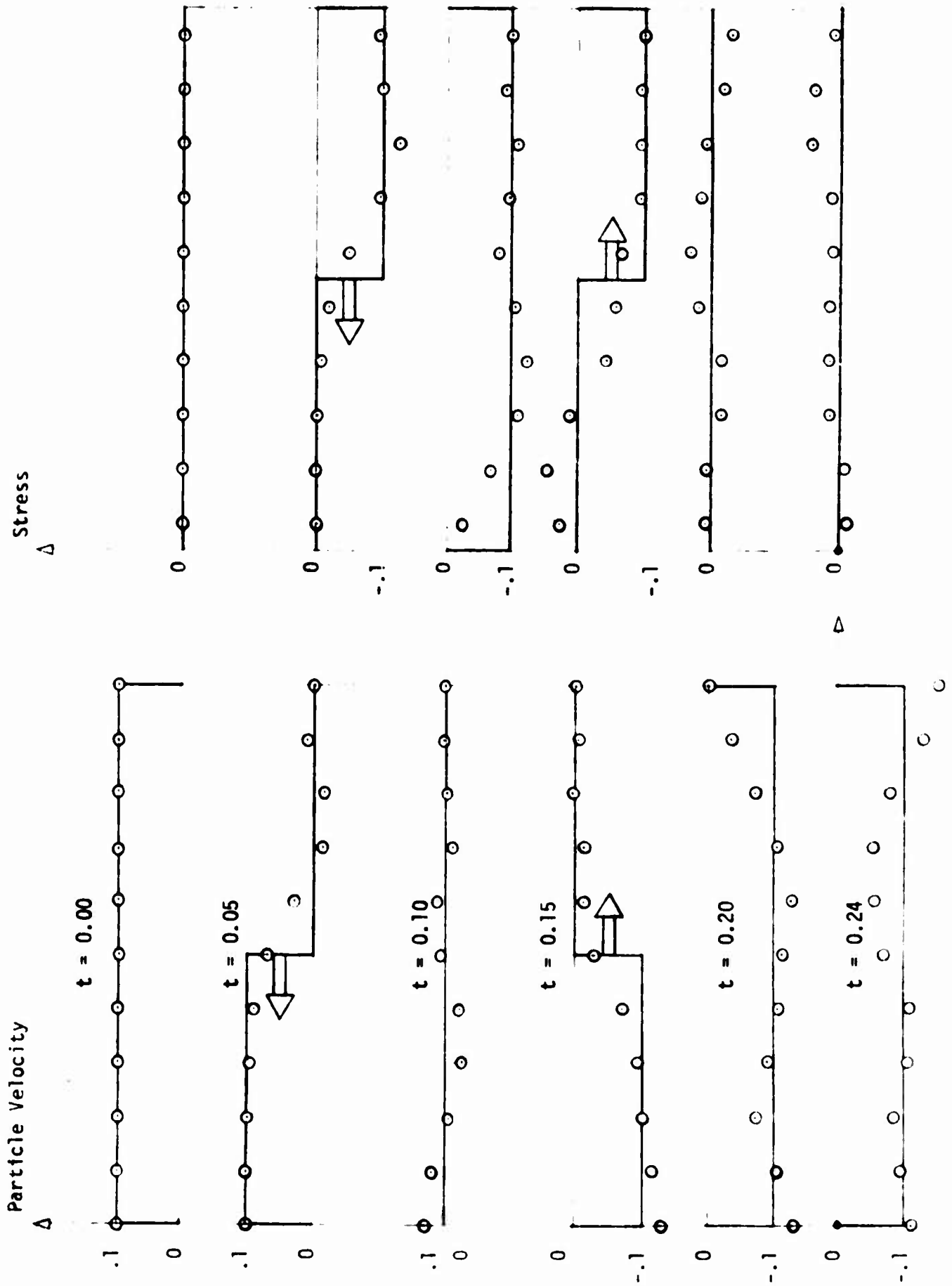


Figure 12

## 17. Closure and Recommendations for Future Work

In the preceding sections we have presented a theory for contact-impact problems together with the numerical development of a Hertzian contact-impact model. The computer program FEAP 74 has been modified to include the model and has successfully solved a contact problem and an impact problem. The work reported herein must be considered initiatory; the general theories and their numerical implementation have not been completed. The problem is of such a complicated nature and the literature existing prior to this study was so meager that we consider it fitting to document the work completed thus far.

We have attempted to qualify each stage of the development throughout the report, however, it may be fitting to reiterate future work which we consider to be essential for numerical models to be effective and efficient tools for predictive analyses.

- (1) The restriction of Hertzian type contact must be removed. This involves the non-trivial task of finding appropriate numerical methods to handle the  $\chi^4$  maps.
- (2) Improved methods for solving the set of nonlinear algebraic equations must be found. We have suggested two methods which should be considered: (a) Substructure the problem about the contact regime so that a more efficient forming and factoring of the tangent stiffness can be performed; and (b) Since the impact problem is a wave propagation problem an explicit time integration of the equations of motion should be explored. In complex situations the explicit integration method may have severe stability limitations which could make it unacceptable.

- (3) Methods of utilizing the shock conditions need to be explored further. We have noted some peculiar anomalies when the bodies separate. These appear to be caused by a shock like separation phenomena.
- (4) When the wave propagation property of the impact problem is ignored by taking time steps greatly in excess of the transit times in a body the computed response is meaningless. Under such situations the bodies rebound within a single time step. Currently the rebound velocity is much too large. When the shock conditions are used for a class of problems where the response desired is in the target instead of in the impactor, it may be expedient to take a large time step. Methods should be explored to accomplish this capability.

The above recommendations for future work should in no way minimize what has been accomplished by the present study. For the first time a contact-impact theory in the form of a variational problem has been presented in a general form. This formulation was motivated by the fact that numerical solutions would be obtained by a finite element method. In addition the necessary foundation for the numerical solution has been thought out and within this context a computer program has been developed for Hertzian contact-impact problems.

The implementations considered here have produced results which are hopeful signs for the eventual success of the more general impact problems.

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## APPENDICES

### A. Input Instructions for Contact/Impact Problems

In order to analyze contact/impact problems in FEAP, users must prepare the data for a time dependent analysis. This will include the following Data Type Identification Cards (see Section 1, Appendix B):

FEAP 74  
 MATERIAL  
 NODAL  
 ELEMENT  
 CONTACT (for impact problems only)  
 loadings  
 INITIAL CONDITIONS (if non-zero)

and

VISCOE (for quasi-static contact problems)

or

IMPLICIT (for impact problems)

In performing the necessary solution to (62) the full Newton-Raphson method must be employed; this is controlled by the data in col. 76-80 of the first card following the VISCOE or IMPLICIT card, and consists of a negative number (negative uses full Newton-Raphson iteration with the absolute value of the number giving the number of iterations to be performed before going to the next time step). As an example of the required input data Table A shows the input data used for the impact problem reported in Section 16.

14.19.14 OUTPUT  
FEAP74 \* \* ONE DIMENSIONAL BAR CONTACT PROBLEM

```

1 X
1 U
3 NODAL
11 0.
11 10.
13 10.
11 111111
11 ELEMENTS : 1
1 1
1 2
11 12 13
11 3 INITIAL CONDITIONS
DISPLACEMENT
VELOCITY
ACCELERATIONS : 1
11 0

```

```

2 MATERIALS
1 ELMO9 1. .01
2 ELMO5
1 CONTACT 1.0 1.0E+30
11 1

```

```

1 IMPLICIT
1.0 5 1 1 13 1 11 0 0 0.25 -3
.01 25 1 1 13 1 11 0 0 0.25 -3
.01 STOP 10 1 1 13 9 11 0 0 0.25 -3

```

#### B. Input Instructions for FEAP 74

The input instructions for the description of a finite element mesh, together with the initial and boundary conditions, is described by subroutine MANUAL listed on the following pages. The input of the contact surface for impact problems is described in Table I of this report.

The description of material properties for the contact element is described in Section 14. For material properties for other elements in FEAP special input instructions must be supplied.

SUBROUTINE MANUAL

C \*\*\*\*\* USER INSTRUCTIONS AND INPUT FORMATS FOR FEAP-74 \*\*\*\*\*

C FINITE ELEMENT ANALYSIS PROGRAM

C \*\*\*\*\* I N D E X \*\*\*\*\*

SECTION	WORD	C O N T E N T
1.0		DATA DIRECTORY WORDS
2.0	FEAP74	PROBLEM INITIATION AND CONTROL CARDS
2.1	REMARK	REMARK CARD
2.2	TITLE	TITLE CHANGE
2.3	STOP	EXECUTION TERMINATION
3.0	MATERI	MATERIAL CHARACTERIZATION
4.0	NODAL	SEQUENTIAL NODAL GENERATION
4.1	GENERA	NON-SEQUENTIAL NODAL GENERATION
4.2	BOUNDA	BOUNDARY CODE INPUT
4.3	POLAR	POLAR TO CARTESIAN CONVERSION
5.0	ELEMEN	ELEMENT GENERATION
5.1	BLOCK	BLOCK MESH GENERATION
6.0	VECTOR	USER VECTOR INPUT
6.1	INITIA	INITIAL CONDITION
7.0	FORCE	GENERALIZED NODAL FORCE
7.1	BLOADS	SURFACE LOADS
7.2	ELOADS	ELEMENT LOADS
7.3		PROPORTIONAL LOADS
8.0	SOLVE	INITIATION OF STATIC SOLUTION
	RESOLV	STATIC SOLUTION
8.1	EXPLIC	EXPLICIT DYNAMIC INTEGRATION
8.2	IMPLIC	IMPLICIT INTEGRATION
	VISCOE	IMPLICIT TIME INTEGRATION
8.3	MESH	MESH CHECK
	PLOT	PLOT MESH
8.4	FOURIE	FOURIER SERIES HARMONICS
9.0		OUTPUT CONTROL

FEAP74 IS A GENERAL (F)INITE (E)LEMENT (A)NALYSIS (P)ROGRAM WHICH FURNISHES TO THE USER MESH INPUT/OUTPUT, ELEMENT ASSEMBLY AND SOLUTION OF EQUATIONS (LINEAR, IMPLICIT AND EXPLICIT TIME DEPENDENT, NONLINEAR), PRESCRIBED GENERALIZED NODAL FORCES, PRESCRIBED NODAL AND ELEMENT DATA, AND OUTPUT OF THE GENERALIZED DISPLACEMENTS AND FORCES. ELEMENT MATRICES FOR TWO AND THREE DIMENSIONAL LINEAR ELASTICITY, SHELLS, PLATES, AND FIELD (LAPLACE EQUATION) PROBLEMS ARE AVAILABLE. ALTERNATIVELY USERS MAY SUPPLY THEIR OWN ELEMENT LIBRARY BY PROVIDING A SUBROUTINE CALLED ELEMTH. WHERE NH IS A TWO DIGIT NUMBER (01-10). IDENTIFYING THE ELEMENT SUBROUTINE. EACH ELEMENT SUBROUTINE HAS AT LEAST FOUR BASIC FUNCTIONS WHICH ARE DELINEATED BY A SWITCHING PARAMETER. ISW. IN THE SUBROUTINE.

MAN 1C  
MAN 2C  
MAN 3C  
MAN 4C  
MAN 5C  
MAN 6C  
MAN 7C  
MAN 8C  
MAN 9C  
MAN 10C  
MAN 11C  
MAN 12C  
MAN 13C  
MAN 14C  
MAN 15C  
MAN 16C  
MAN 17C  
MAN 18C  
MAN 19C  
MAN 20C  
MAN 21C  
MAN 22C  
MAN 23C  
MAN 24C  
MAN 25C  
MAN 26C  
MAN 27C  
MAN 28C  
MAN 29C  
MAN 30C  
MAN 31C  
MAN 32C  
MAN 33C  
MAN 34C  
MAN 34C  
MAN 34C  
MAN 35C  
MAN 35C  
MAN 36C  
MAN 37C  
MAN 38C  
MAN 39C  
MAN 40C  
MAN 41C  
MAN 42C  
MAN 43C  
MAN 44C  
MAN 45C  
MAN 46C  
MAN 47C  
MAN 48C  
MAN 49C  
MAN 50C



ELMTN(N,MA,NDIM,NDF,NEL,NEL1,NSTF,NSIZV,NVEC,NCT,IM,D,XYZ,  
IX,F,FORCE,ESTIF,U,VECT,ISW)

N IS ELEMENT NUMBER.  
MA IS THE MATERIAL NUMBER.  
NDIM IS SPATIAL DIMENSION, 1-2 OR 3.  
NDF IS NUMBER OF DEGREES OF FREEDOM PER NODE.  
NEL IS THE NUMBER OF EXTERNAL NODES PER ELEMENT  
NEL1 IS DIMENSION OF ELEMENT PROPERTY ARRAY.  
NSTF IS THE SIZE OF THE ELEMENT STIFFNESS.  
NSIZV IS THE SIZE OF UTILITY VECTORS.  
NVEC IS THE NUMBER OF UTILITY VECTORS.  
NCT IS A PRINTER LINE COUNTER.  
DM IS A PARAMETER FOR MATERIAL IDENTIFICATION.  
D(L,I) IS MATERIAL PROPERTY MATRIX (63 CELLS).  
XYZ(NDIM,I) ARE NODAL COORDINATES.  
IX(NEL,I) ARE ELEMENT PROPERTIES, NODES, ETC.  
F(NDF,I) ARE NODAL GENERALIZED FORCES.  
FORCE(NSTF,2) IS ELEMENT FORCE VECTOR TO BE  
COMPUTED. COLUMN 2 IS LUMPED MASS  
ESTIF(NSTF,NSTF) IS ELEMENT MATRIX TO BE  
COMPUTED.  
VECT(NSIZV,I) ARE PRESCRIBED NODAL OR ELEMENT  
QUANTITIES, TEMPERATURES ETC.  
U(NDF,I) IS SOLUTION VECTOR.  
ISW IS SWITCHING PARAMETER.

ISW=1. \*\* MATERIAL CHARACTERIZATION\*\*  
ISW=2. \*\* CHECK ELEMENT FOR POSITIVE AREA \*  
ISW=3. \*\* ELEMENT STIFFNESS AND  
LOAD COMPUTATION\*\*  
ISW=4. \*\* ELEMENT STRESSES AND PRINTOUT\*\*  
ISW=5. \*\* ELEMENT LOAD COMPUTATION ONLY\*\*  
ISW=6. \*\* NONLINEAR GENERALIZED FORCES\*\*  
OTHER ISW MAY BE USED FOR SPECIAL PURPOSES.

USERS CAN GENERATE SURFACE LOADINGS BY PROVIDING SLDNN  
SUBROUTINES (WHERE NH IS A TWO DIGIT NUMBER BETWEEN 01 AND 05)  
THAT SPECIFY THE LOAD ROUTINE. THE SUBROUTINE IS ACCESSED BY  
THE CALL TO

SLDNN(NDIM,NDF,NDR,NPRES,IPRES,PP,XYZ,FS)

WHERE IN ADDITION TO QUANTITIES DEFINED ABOVE  
FOR ELMTN.

NDF IS THE DIMENSION OF LOADED SURFACE

IPRES IS NUMBER OF LOADED NODES (MAX 8)

PP(8) ARE NODE NUMBERS OF LOADED NODES.

PRINT ARE LOAD VALUES AT CORRESPONDING IPRES  
NODE.

FOR(8) ARE THE COMPUTED GENERALIZED (NODAL)  
FORCES FOR EACH DEGREE OF FREEDOM AT EACH  
IPRES NODE.

MAN 51C  
MAN 52C  
MAN 53C  
MAN 54C  
MAN 55C  
MAN 56C  
MAN 57C  
MAN 58C  
MAN 59C  
MAN 60C  
MAN 61C  
MAN 62C  
MAN 63C  
MAN 64C  
MAN 65C  
MAN 66C  
MAN 67C  
MAN 68C  
MAN 69C  
MAN 70C  
MAN 71C  
MAN 72C  
MAN 73C  
MAN 74C  
MAN 75C  
MAN 76C  
MAN 77C  
MAN 78C  
MAN 79C  
MAN 80C  
MAN 80C  
MAN 81C  
MAN 82C  
MAN 83C  
MAN 84C  
MAN 85C  
MAN 86C  
MAN 87C  
MAN 88C  
MAN 89C  
MAN 90C  
MAN 91C  
MAN 92C  
MAN 93C  
MAN 94C  
MAN 95C  
MAN 96C  
MAN 97C  
MAN 98C  
MAN 99C  
MAN 99C  
MAN 99C  
MAN 99C  
MAN 99C

SEE SECTION 7.1 FOR DATA INPUT DETAILS.

INTEGRATION TABLE IS ACCESSED BY THE CALL

CALL INTEGL(LIM,NCI,NDIM,LINT,STWJ)

STWJ(4,M) INTEGRATION POINTS AND WEIGHTS.

\*\*NOTE\*\* M MUST BE SET EXPLICITLY AND BE LARGER THAN OR EQUAL TO LINT.

LINT - RETURNS WITH NUMBER INTEGRATION POINTS.

NCI = 0 RETURNS GAUSS POINTS AND WEIGHTS IN

STWJ.

LIM = 1 TO 5 IS NUMBER OF GAUSS POINTS DIRECTION.

NCI = 1 RETURNS A SPECIAL 3-D GAUSS FORMULA.

SET LIM = 1 FOR 6 PT. CUBIC ACCURACY

SET LIM = 2 FOR 14 PT. QUINTIC ACCURACY.

NCI = 2 RETURNS TRIANGULAR INTEGRATION FORMULA

SET LIM = 1 FOR 1 PT. LINEAR ACCURACY.

SET LIM = 2 FOR 3 PT. QUADRATIC ACCURACY.

SET LIM = 3 FOR 7 PT. QUARTIC ACCURACY.

1.) DATA TYPE IDENTIFICATION CARDS (15,1X,12A6).

EACH DATA SEGMENT IS PRECEDED BY A CARD WHICH IDENTIFIES THE TYPE OF DATA AND LIMITS ON THE AMOUNT OF DATA WHICH IMMEDIATELY FOLLOWS THE CARD. EXCEPT AS NOTED THE DATA SEGMENTS MAY APPEAR IN ANY ORDER. THE IDENTITY CARDS MAY ALSO AID THE USER IN INTERPRETTING THE INPUT DATA CARDS. AS SUPPLIED THERE ARE TWENTY-FIVE DIFFERENT DATA IDENTIFICATION CARDS. THESE ARE

COL 7 TO 12 IDENTITY(RESTRICTIONS)

FEAP74 START OF EACH PROBLEM (MUST PRECEDE ALL OTHER DATA).

TITLE CHANGE OUTPUT PAGE HEADINGS

REMARK COMMENTS ON OUTPUT

MATEPI MATERIAL CHARACTERIZATION.

NODAL NODAL CARDS

POLAR POLAR CONVERSION. (PRECEDE BY NODAL, GENERA. OR BLOCK)

ELEMEN ELEMENT CONNECTION CARDS.

GENERA GENERATE NODES IN A LINEAR PATH BY ANY INCREMENT

BLOCK GENERATE ALL MESH DATA (BOTH NODAL AND ELEMENT)

BOUNDARY FOR 4, 2, OR 3 DIMENSIONAL REGION WHOSE BOUNDARY

MAY BE DEFINED BY 4(S) OR 8(20) COLLOCATED POINTS

BOUNDARY CODE PRESCRIPTION (PRECEDE BY NODAL OR

GENERATE OF BLOCK)

VECTOP PRECEDED NODAL OR ELEMENT DATA (PRECEDE BY

NODAL, OF GENERA AND ELEMEN. OF BLOCK)

FORCE NODAL GENERALIZED FORCES (PRECEDE BY NODAL OR

GENERATE OF BLOCK)

- MAN104C
- MAN105C
- MAN106C
- MAN107C
- MAN108C
- MAN109C
- MAN110C
- MAN111C
- MAN112C
- MAN113C
- MAN114C
- MAN115C
- MAN116C
- MAN117C
- MAN118C
- MAN119C
- MAN120C
- MAN121C
- MAN122C
- MAN123C
- MAN124C
- MAN125C
- MAN126C
- MAN127C
- MAN128C
- MAN129C
- MAN130C
- MAN131C
- MAN132C
- MAN133C
- MAN134C
- MAN135C
- MAN136C
- MAN137C
- MAN138C
- MAN139C
- MAN140C
- MAN141C
- MAN142C
- MAN143C
- MAN144C
- MAN145C
- MAN146C
- MAN147C
- MAN148C
- MAN149C
- MAN150C
- MAN151C
- MAN152C
- MAN153C
- MAN154C
- MAN155C
- MAN156C
- MAN157C

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 C MAN204C  
 C MAN205C  
 C MAN206C  
 C MAN207C  
 C MAN208C  
 C MAN209C  
 C MAN210C  
 C MAN211C

C SURFACE LOADINGS (SAME AS FORCE).  
 C ELEMENT LOADINGS (SAME AS FORCE).  
 C CHECK CONSISTENCY OF MESH ONLY (SAME AS SOLVE)  
 C PLOT MESH (SAME AS SOLVE)  
 C INITIAL CONDITION PRESCRIPTION FOR DYNAMIC  
 C ANALYSIS (PRECEDE BY NODAL, GENERA OR BLOCK)  
 C COMPLETE FORMULATION AND SOLUTION FROM ELEMENTS  
 C (PRECEDE BY MATERI, NODAL OR GENERA, AND ELEMEN  
 C OR PRECEDE BY MATERI AND BLOCK)  
 C USE PREVIOUS PROBLEM DESCRIPTION WITH NEW LOAD  
 C ONLY (PRECEDE BY SOLVE AND NEW LOADING CARDS).  
 C DYNAMIC SOLUTION BY EXPLICIT INTEGRATION. (SAME  
 C AS SOLVE)  
 C IMPLICIT INTEGRATION OF DYNAMIC PROBLEMS  
 C (PRECEDE BY SAME DATA AS FOR SOLVE)  
 C QUASI-STATIC LINEAR VISCOELASTIC INTEGRATION  
 C (PRECEDE BY SAME DATA AS FOR SOLVE)  
 C FOURIER COMPOSITION (SAME AS SOLVE)  
 C ACCUMULATE FOURIER SOLUTION (AFTER FOURIE)  
 C NORMAL EXIT (MUST FOLLOW ALL DATA)

C \*\*\*NOTE\*\*\* EACH IDENTIFIER IS PUNCHED STARTING IN COL 7 (LEFT  
 C JUSTIFIED).

C EXCESS CARDS MAY EXIST BETWEEN EACH SECTION OF DATA, HOWEVER,  
 C THE DATA TO BE USED MUST IMMEDIATELY FOLLOW THE TYPE CARD AND  
 C MUST BE IN PROPER ORDER. NO PARTICULAR ORDER OF THE TYPE  
 C CARDS IS NECESSARY EXCEPT THAT THE FEAP74 CARD MUST ALWAYS BE  
 C THE FIRST CARD IN EACH SET OF DATA, AND RESTRICTIONS MUST BE  
 C OBSERVED.

2.) PROBLEM INITIATION AND CONTROL CARDS

C CARD 1. (6X,12A6)  
 C COL 7 TO 12 MUST CONTAIN WORD FEAP74  
 C COL 13 TO 78 OUTPUT PAGE HEADER  
 C CARD 2. (15,1X,3A6)  
 C COL 1 TO 5 NDIM - SPATIAL DIMENSION OF PROBLEM (1 TO 3)  
 C COL 7 TO 12 NAMES TO BE PRINTED AS OUTPUT HEADERS TO  
 C COL 13 TO 18 COORDINATES - IF BLANK SET TO 1.2.3 AS NEEDED.  
 C COL 19 TO 24  
 C CARD 3. (15,1X,6A6)  
 C COL 1 TO 5 NDF - NUMBER OF UNKNOWN PER NODE (1 TO 6)  
 C COL 7 TO 12 NAMES TO BE PRINTED AS OUTPUT HEADERS OF THE  
 C COL 13 TO 18 GENERALIZED DISPLACEMENTS AND FORCES - IF  
 C BLANK SET TO 1.2.3.4.5.5 AS NECESSARY.  
 C COL 19 TO 24

C CARD 4. (615,5F10.0)

C COL 1 TO 5 NEN - MAXIMUM NUMBER OF NOIES CONNECTED TO ANY  
 C ELEMENT (1 TO 20).

C COL 6 TO 10 NEXTRA - INCREASES ELEMENT MATRIX SIZE FROM  
 C NDF\*NEN TO NDF\*NEN + NEXTRA

C COL 11 TO 15 IREC - COMPUTE GENERALIZED FORCE CHECK IF  
 C NONZERO (FOR TIME INVARIANT ANALYSIS ONLY)

C COL 16 TO 20 MBAN - MAXIMUM EXPECTED BANDWIDTH, DEFAULT IS  
 C SET TO 100. USED AS AN ERROR CHECK TO PREVENT  
 C RUNNING WITH AN OBVIOUS ERROR.

C COL 21 TO 25 IBUF - BUFFER SIZE FOR STORAGE OF HISTORY  
 C EFFECTS IN TIME DEPENDENT ANALYSIS. DEFAULT IS  
 C IBUF = ISZDT\*20

C COL 26 TO 30 NC1 - USER INTEGER CONSTANT

C COL 31 TO 40 CON1 - USER DEFINED CONSTANT

C COL 41 TO 50 CON2 - USER DEFINED CONSTANT

C COL 51 TO 60 CON3(1) - USER DEFINED CONSTANT

C COL 61 TO 70 CON3(2) - USER DEFINED CONSTANT

C COL 71 TO 80 CON3(3) - USER DEFINED CONSTANT

C 2.1) REMARK \* USER COMMENTS ON OUTPUT. (6X,12A6)

C SUBSEQUENT CARDS

C COL 7 TO 12 MUST CONTAIN REMARK

C COL 13 TO 78 STATEMENTS TO BE OUTPUT , USE AS MANY REMARK  
 C CARDS AS DESIRED. INSERT BEFORE ANY TYPE CARD.

C 2.2) TITLE CHANGE ON OUTPUT (6X,12A6)

C COL 7 TO 12 MUST CONTAIN TITLE

C COL 13 TO 78 NEW TITLE DESCRIPTOR

C 2.3) EXECUTION TERMINATION (6X,44)

C COL 7 TO 10 MUST CONTAIN STOP. INSERT AFTER LAST PROBLEM.

C 3.) MATERIAL CHARACTERIZATION (15,1X,12A6)

C COL 1 TO 5 NUMMAT - NUMBER OF DIFFERENT MATERIAL CHARACT-  
 C ERIZATIONS TO FOLLOW.

C COL 7 TO 12 MUST CONTAIN WORD MATEP1

C THE FOLLOWING CARDS ARE SUPPLIED FOR EACH MATERIAL TO BE CHARAC  
 C TERIZED (MUST BE EXACTLY NUMMAT SETS OF CARDS)

C CASE 1. ELEMENT SELECTOR CARD (15,1X,45,11A6)

C COL 1 TO 5 MATERIAL NUMBER (1 TO NUMMAT)

C COL 7 TO 11 ELEMENT NUMBER (1 TO 10) IS NUMBER OF ELEMENT CLASS (0)

C COL 13 TO 78 TO 100 TO INDICATE THE CHARACTERIZATION BELONGS.  
 C ELEMENTS INFORMATION TO BE OUTPUT.

MAN212C  
 MAN213C  
 MAN214C  
 MAN215C  
 MAN216C  
 MAN217C  
 MAN218C  
 MAN219C  
 MAN220C  
 MAN221C  
 MAN222C  
 MAN223C  
 MAN224C  
 MAN225C  
 MAN226C  
 MAN227C  
 MAN228C  
 MAN229C  
 MAN230C  
 MAN231C  
 MAN232C  
 MAN233C  
 MAN234C  
 MAN235C  
 MAN236C  
 MAN237C  
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 MAN263C  
 MAN264C  
 MAN265C

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 MAN271C  
 MAN272C  
 MAN273C  
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 MAN277C  
 MAN278C  
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 MAN301C  
 MAN302C  
 MAN303C  
 MAN304C  
 MAN305C  
 MAN306C  
 MAN307C  
 MAN308C  
 MAN309C  
 MAN310C  
 MAN311C  
 MAN312C  
 MAN313C  
 MAN314C  
 MAN315C  
 MAN316C  
 MAN317C

CARD 2, 1, ETC. \*\* USER DEFINED FOR EACH ELEMENT TYPE PROVIDED.

EXCESS BLANK CARDS ARE PERMISSIBLE BETWEEN EACH MATERIAL SET.

4.) NODAL CARDS (15.14.45)

COL 1 TO 5 NUNIP - NUMBER OF NODAL POINTS  
 COL 7 TO 12 MUST CONTAIN NODAL

SUBSEQUENT CARDS LAST NODAL CARD MUST NOT BE GENERATED.  
 (15.115.3F10.0)

COL 1 TO 5 NODE NUMBER  
 COL 15 1 IF 1 DISPLACEMENT IS SPECIFIED  
 COL 16 1 IF 2 DISPLACEMENT IS SPECIFIED  
 COL 17 1 IF 3 DISPLACEMENT IS SPECIFIED  
 COL 18 1 IF 4 DISPLACEMENT IS SPECIFIED  
 COL 19 1 IF 5 DISPLACEMENT IS SPECIFIED  
 COL 20 1 IF 6 DISPLACEMENT IS SPECIFIED  
 COL 21 TO 30 1 COORDINATE VALUE  
 COL 31 TO 40 2 COORDINATE VALUE \* AS REQUIRED  
 COL 41 TO 50 3 COORDINATE VALUE

NODAL CARDS MUST BE IN ORDER. MISSING NODES ARE INTERPOLATED LINEARLY FROM INPUT NODES. IF SUCCEEDING CARDS HAVE IDENTICAL BOUNDARY CODES. THIS BOUNDARY CODE WILL BE ASSIGNED TO THE INTERVENING NODES. IN ALL OTHER CASES THE BOUNDARY CODE IS SET TO ZERO \*TERMINATE ON NODE NUNIP OR A BLANK CARD\*

4.1) NON SEQUENTIAL NODAL GENERATOR OPTION. (15.14.46)

COL 1 TO 5 NUMBER OF NODAL POINTS  
 COL 7 TO 12 MUST CONTAIN GENERA

SUBSEQUENT CARDS (215.110.3F10.0)

COL 1 TO 5 NODE-NUM\*\*\*  
 COL 6 TO 10 NODE-NUMBER-INCREMENT WHICH WILL BE SUCCESSIVELY ADDED TO NODE-NUMBER UNTIL SUM IS GREATER THAN NODE-NUMBER ON FOLLOWING CARD (ALGEBRAIC).  
 COL 15 TO 20 BOUNDARY CODE. SAME AS INPUT FOR NODAL.  
 IF SUCCEEDING CARDS HAVE IDENTICAL BOUNDARY CODES, THIS BOUNDARY CODE WILL BE ASSIGNED TO THE INTERVENING NODES. IN ALL OTHER CASES THE BOUNDARY CODE IS SET TO ZERO.

COL 21 TO 30 1 COORDINATE VALUE \*  
 COL 31 TO 40 2 COORDINATE VALUE \* AS REQUIRED \*  
 COL 41 TO 50 3 COORDINATE VALUE \*

\*TERMINATE WITH BLANK CARD \*

4.2) BOUNDARY CODE PATCHING OPTION. (15.14.47)

COL 7 TO 12 MUST CONTAIN BOUNDARY

SUBSEQUENT CARDS. (815)

COL 1 TO 5 N. NODE NUMBER TO HAVE REDEFINED BOUNDARY CODE.  
 COL 6 TO 10 NX. GENERATOR INCREMENT TO BE ADDED ALGEBRAICALLY TO  
 CALLY TO N. UNTIL SUM EXCEEDS (MAX OR MIN) THE  
 N OF THE FOLLOWING CARD.

COL 11 TO 15 IBC(I), (I=1,2,..,NDF) CODE FOR SPECIFYING FORCE  
 COL 16 TO 20 OR DISPLACEMENT BOUNDARY CONDITIONS.

COL ...  
 IBC(I) .EQ. 0. FORCE SPECIFIED.  
 IBC(I) .GT. 0. DISPLACEMENT SPECIFIED. NO  
 INTERVENING GENERATION.  
 IBC(I) .LT. 0. DISPLACEMENT SPECIFIED.  
 GENERATE BETWEEN MISSING NODES IN ALGEBRAIC  
 INCREMENTS OF NX.

\* TERMINATE WITH A BLANK CARD. \*

4.3) POLAR OR CYLINDRICAL COORDINATE CONVERSION TO CARTESIAN  
 COORDINATES (6X.A6)

COL 7 TO 12 MUST CONTAIN POLAR (LEFT JUSTIFIED)

CARD 1. (315.5%.2F10.0)

COL 1 TO 5 N1. FIRST NODE TO BE CONVERTED  
 COL 6 TO 10 N2. LAST NODE TO BE CONVERTED  
 COL 11 TO 15 N3. INCREMENT ADDED (ALGEBRAICALLY), N1 TO N2  
 COL 21 TO 30 X0. ORIGIN OF POLAR X-COORDINATE  
 COL 31 TO 40 Y0. ORIGIN OF POLAR Y-COORDINATE

5.) ELEMENT CARDS (15.1X.A6)

COL 1 TO 5 NUMEL - NUMBER OF ELEMENTS  
 COL 7 TO 12 MUST CONTAIN ELEMENT

SUBSEQUENT CARDS (415.2013/2014)

CARD 1.

COL 1 TO 5 ELEMENT NUMBER  
 COL 6 TO 10 MATERIAL NUMBER  
 COL 11 TO 15 NUMBER OF SUBSEQUENT ELEMENTS USING SAME  
 STIFFNESS MATRIX \* SAVES RECOMPUTATION OF  
 SIMILAR MATRICES. ELEMENT MUST ALSO HAVE  
 SAME ELEMENT FORCE VECTOR \* IF THESE ARE  
 IN THE STIFFNESS SUBROUTINE \*  
 PRINT ELEMENT MATRIX IF NONZERO.

COL 16 TO 20 INCR1 ELEMENT INCREMENT ARRAY OR MORE 1.  
 COL 21 TO 25 INCR2 \* IF NOT INPUT IS SET AUTOMATICALLY  
 COL 26 TO 30 INCR3

FOR IDENTIFY ELEMENTS \* SEE REPORT

MAN318C  
 MAN319C  
 MAN320C  
 MAN321C  
 MAN322C  
 MAN323C  
 MAN324C  
 MAN325C  
 MAN326C  
 MAN327C  
 MAN328C  
 MAN329C  
 MAN330C  
 MAN331C  
 MAN332C  
 MAN333C  
 MAN334C  
 MAN335C  
 MAN336C  
 MAN337C  
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 MAN387C  
 MAN388C  
 MAN389C  
 MAN390C  
 MAN391C  
 MAN392C  
 MAN393C  
 MAN394C  
 MAN395C  
 MAN396C  
 MAN397C  
 MAN398C  
 MAN399C  
 MAN400C  
 MAN401C  
 MAN402C  
 MAN403C  
 MAN404C  
 MAN405C  
 MAN406C  
 MAN407C  
 MAN408C  
 MAN409C  
 MAN410C  
 MAN411C  
 MAN412C  
 MAN413C  
 MAN414C  
 MAN415C  
 MAN416C  
 MAN417C  
 MAN418C  
 MAN419C  
 MAN420C  
 MAN421C  
 MAN422C  
 MAN423C  
 MAN424C  
 MAN425C  
 MAN426C

CARD 2.  
 COL 1 TO 4 NODE 1  
 COL 5 TO 8 NODE 2  
 COL 9 TO 12 NODE 3  
 COL 13 TO 16 NODE 4  
 CONTINUE IN 14 FORMAT TO A MAXIMUM  
 ELEMENT CARDS MUST BE IN ORDER. MISSING ELEMENTS ARE GENERATED BY INCREMENTING NODES.  
 LAST ELEMENT CARD MUST NOT BE GENERATED.  
 \* TERMINATE ON ELEMENT NUMBER OF A BLANK CARD \*

5.1) BLOCK GENERATOR. GENERATES ALL MESH DATA. (6X,A6)  
 COL 1 TO 5 NUMBER OF NODAL POINTS TO BE GENERATED. MUST CONTAIN BLOCK  
 COL 7 TO 12

SUBSEQUENT CARDS (10I5\*6(4X,16)/(10X,3F10.0))  
 CARD 1.  
 COL 1 TO 5 NN, NUMBER OF POINTS REQUIRED TO DEFINE BOUNDARY OF REGION. FOR 2-DIM., NN=4 OR 8. FOR 3-DIM., NN= 8 OR 20.  
 COL 6 TO 10 NUMBER OF ELEMENTS IN R-DIRECTION.  
 COL 11 TO 15 NUMBER OF ELEMENTS IN S-DIRECTION.  
 COL 16 TO 20 NUMBER OF ELEMENTS IN T-DIRECTION.  
 COL 21 TO 25 INITIAL NODE NUMBER, DEFAULT = 1.  
 COL 26 TO 30 INITIAL ELEMENT NUMBER, DEFAULT = 1.  
 COL 31 TO 35 MATERIAL NUMBER OVER REGION. DEFAULT = 1  
 COL 36 TO 40 BOUNDARY CODE SKIP. A NON-ZERO ENTRY WILL OMIT SETTING ALL INTERIOR BOUNDARY CODES TO ZERO.  
 COL 41 TO 45 IREUSE - REUSE ELEMENT STIFFNESS OPTION - USES EACH ELEMENT STIFFNESS IREUSE TIMES BEFORE GENERATING A NEW ELEMENT STIFFNESS MATRIX. WILL ELEMENT STIFFNESS-PRINT. A NON-ZERO ENTRY WILL CAUSE PRINT-OUT OF FIRST ELEMENT.  
 COL 46 TO 50 IMESH - IF NONZERO PRINT COMPUTED ELEMENTS.  
 COL 51 TO 55 IELM - IF NONZERO PRINT COMPUTED ELEMENTS.  
 COL 56 TO 60

CARD 2. (BOUNDARY CODE AS DEFINED IN NODAL CARD.)  
 COL 5 TO 10 BOUNDARY CODE OVER FACE -R.  
 COL 11 TO 15 BOUNDARY CODE OVER FACE +R.  
 COL 16 TO 20 BOUNDARY CODE OVER FACE -S.  
 COL 21 TO 25 BOUNDARY CODE OVER FACE +S.  
 COL 26 TO 30 BOUNDARY CODE OVER FACE -T.  
 COL 31 TO 35 BOUNDARY CODE OVER FACE +T.  
 COL 36 TO 40

CARD 3. REPEAT IN TIMES.

BOUNDARY-DEFINING POINT.

COL 21 TO 30 2-COORDINATE OF BOUNDARY-DEFINING-POINT.  
 COL 31 TO 40 3-COORDINATE OF BOUNDARY-DEFINING-POINT.

NOTES. 1. BLOCK GENERATES ONLY 4 PT. QUADRIATERALS OR 8 PT. BRICKS.  
 2. INPUT OF CARDS 3.) FOLLOW ORDER RULES FOR ELEMENT INPUT (SEE SECTION 5.1).  
 3. R-S-T APE LOCAL COORDINATES.  
 I.E. (-1 .LE. R.S.T .LE. 1). WHERE R IS DIRECTED FROM NODE 1 TO 2. S IS IN PLANE OF FIRST THREE NODES AND T IS NORMAL TO R-S PLANE.

6.) VECTOR CARDS, I.E. USER DEFINED INPUT (15.1X.A6)

COL 1 TO 5 NVEC. NUMBER OF DIFFERENT VECTORS (7 MAX)  
 COL 7 TO 12 MUST CONTAIN VECTOR

SUBSEQUENT CARDS

CARD 1. (215)

COL 1 TO 5 NS12V. VECTOR LENGTH.COMMON TO ALL NVEC VECTORS  
 COL 6 TO 10 IPICK. CODED PARAMETER.

IPICK = 0. VECTORS ASSOCIATED WITH NODES  
 IPICK = 1. VECTORS ASSOCIATED WITH DEG.FREEDOM  
 IPICK = 2. VECTORS ASSOCIATED WITH ELEMENTS

CARD 2. (6X. 2A6) REPEAT NVEC TIMES

COL 7 TO 18 DESCRIPTIVE TITLE FOR VECTOR

CARD 3. (215.7F10.0)

COL 1 TO 5 POSITION NUMBER OF VECTOR ELEMENT. 1 TO NS12V  
 COL 6 TO 10 GENERATOR INCREMENT  
 COL 11 TO 20 VECTOR ELEMENT VALUE OF VECTOR 1  
 COL 21 TO 30 VECTOR ELEMENT VALUE OF VECTOR 2  
 COL ..... AS PREPARED FOR NVEC VECTORS

LINEAR INTERPOLATION IS PERFORMED ON ALL VECTORS BETWEEN NON-CONSECUTIVE POSITION NUMBERS SPECIFIED IN COL 1 TO 5 IF INCREMENT IS NONZERO.  
 IF DESCRIPTIVE TITLES OF ALL VECTORS ARE BLANK CARDS. PRINTING OF THE VECTOR VALUES IS SUPPRESSED.

\* TERMINATE ON BLANK CARD \*

5.1. INITIAL CONDITIONS FOR TIME DEPENDENT ANALYSIS.

COL 1 TO 5 NVEC. NUMBER OF INITIAL CONDITION VECTORS  
 COL 7 TO 12 MUST CONTAIN VECTOR

SUBSEQUENT CARDS

- MAN427C
- MAN428C
- MAN429C
- MAN430C
- MAN431C
- MAN432C
- MAN433C
- MAN433C
- MAN433C
- MAN433C
- MAN433C
- MAN435C
- MAN436C
- MAN437C
- MAN438C
- MAN439C
- MAN440C
- MAN441C
- MAN442C
- MAN443C
- MAN444C
- MAN445C
- MAN446C
- MAN447C
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- MAN451C
- MAN452C
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- MAN454C
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- MAN456C
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- MAN459C
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- MAN461C
- MAN462C
- MAN463C
- MAN464C
- MAN465C
- MAN466C
- MAN467C
- MAN468C
- MAN469C
- MAN470C
- MAN471C
- MAN472C
- MAN473C
- MAN474C
- MAN475C
- MAN476C
- MAN477C



MAIN478C  
 MAIN479C  
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 MAIN598C  
 MAIN599C  
 MAIN600C

C 64 (6X.2A6) REPEAT NICK TIMES  
 C COL 7 TO 19 DESCRIPTIVE TITLE FOR INITIAL CONDITIONS  
 C CAPD 2. (215.7F10.0)  
 C COL 1 TO 5 POSITION NUMBER. AS IN VECTOR CARDS FOR IPICK=1  
 C COL 6 TO 10 GENERATOR INCREMENT  
 C COL 11 TO 20 INITIAL CONDITION 1  
 C COL 21 TO 30 INITIAL CONDITION 2  
 C COL ..... AS REQUIRED FOR NICK INITIAL CONDITIONS  
 C INTERPOLATION BETWEEN INPUT VALUES AS DESCRIBED IN VECTOR INPUT  
 C \*\*\* NOTE \*\*\* IF MISSING THE INITIAL CONDITIONS ARE SET ZERO  
 C  
 C 7.) FORCE CARDS (15.1X.A6)  
 C COL 1 TO 5 LAST NODE TO WHICH A FORCE IS TO BE SPECIFIED  
 C COL 7 TO 12 MUST CONTAIN FORCE  
 C SUBSEQUENT CARDS (15.5X.6F10.0)  
 C THE FOLLOWING VALUES ARE EACH INTERPRETTED AS FORCES IF THE  
 C CORRESPONDING BOUNDARY CODE IS A 0 \*ZERO\* AND AS A DISPLACEMENT  
 C IF THE CORRESPONDING BOUNDARY CODE IS 1 \*ONE\*.  
 C COL 1 TO 5 NODE TO WHICH FORCE OR DISPLACEMENT IS APPLIED  
 C COL 11 TO 20 VALUE OF 1 FORCE DISPLACEMENT  
 C COL 21 TO 30 VALUE OF 2 FORCE DISPLACEMENT \* AS \*  
 C COL 31 TO 40 VALUE OF 3 FORCE DISPLACEMENT \* REQUIRED \*  
 C COL 41 TO 50 VALUE OF 4 FORCE DISPLACEMENT  
 C COL 51 TO 60 VALUE OF 5 FORCE DISPLACEMENT  
 C COL 61 TO 70 VALUE OF 6 FORCE DISPLACEMENT  
 C  
 C 7.1) SURFACE LOAD CARDS (15.1X.A6)  
 C COL 1 TO 5 NUMBER OF LOADED FACE CARDS  
 C COL 7 TO 12 MUST CONTAIN BLOADS  
 C CAPD 1. (15.1X.A5.14.815.813)  
 C COL 1 TO 5 DIMENSION OF LOADING SURFACE. (1 OF 2).  
 C COL 7 TO 11 SURFACE ALPHA-NUMERIC NAME OF SURFACE LOADING  
 C SUBROUTINE THIS IS BETWEEN 1 AND 5.  
 C COL 13 TO 15 NET NUMBER OF ADDITIONAL ELEMENT LOAD  
 C SURFACES TO BE GENERATED FROM CURRENT MODEL.  
 C COL 16 TO 20 REFERENCE NODE NUMBERS DEFINING LOADING SURFACE  
 C COL 21 TO 25 OF CURRENT ELEMENT.  
 C COL 26 TO 30 IDENTIFY FROM 2 TO 8 AS REQUIRED  
 C COL 31 TO 35 INITIAL INCREMENT VALUE NOTED IN DEFSURF TO  
 C COL 36 TO 40 INITIAL INCREMENT VALUE OF SURFACE LOADS

MAN532C  
 MAN533C  
 MAN534C  
 MAN535C  
 MAN536C  
 MAN537C  
 MAN538C  
 MAN539C  
 MAN540C  
 MAN541C  
 MAN542C  
 MAN543C  
 MAN544C  
 MAN545C  
 MAN546C  
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 MAN580C  
 MAN581C  
 MAN582C  
 MAN583C  
 MAN584C  
 MAN585C

COL ... (IDENTIFY FROM 2 TO 8 AS REQUIRED)  
 CAPD 2. (BF10.0)  
 COL 1 TO 90 LOAD AT NODES GIVEN ON PREVIOUS CARD \*  
 MUST CORRESPOND IN SEQUENCE TO THE NODE NUMBERS

7.2) ELEMENT LOAD CARDS (15.1X.A6 )  
 COL 1 TO 5 NLD. NUMBER OF ELEMENT LOAD CARDS.  
 COL 7 TO 11 MUST CONTAIN ELOADS  
 SUBSEQUENT CARDS (15.1X.A5.14.15.6F10.0)  
 COL 1 TO 5 IEL. INITIAL ELEMENT OF A GENERATED SEQUENCE.  
 COL 7 TO 11 ELMTN). ALPHA-NUMERIC NAME OF ELEMENT  
 SUBROUTINE WHERE ELEMENT LOADS ARE COMPUTED.  
 USED AS CHECK TO INSURE IEL. ETC. ARE PROPER  
 ELEMENTS.  
 COL 12 TO 15 INC. INCREMENT NUMBER IN A GENERATED SEQUENCE.  
 (DEFAULT = 1).  
 COL 16 TO 20 JEL. TERMINAL ELEMENT NUMBER IN A GENERATED  
 SEQUENCE. IF JEL = 0, ONLY IEL IS COUNTED.  
 COL 21 TO 80 PR-USER DEFINED VALUES FOR DETERMINING BODY  
 LOADS IN THE ISU=5 PORTION OF ELMTN).

NOTE. USER MUST PROVIDE COMPUTATION OF LOADS IN ELMTN.  
 PR IS TRANSFERRED TO SUBROUTINE ELMTN IN THE U VECTOR.  
 WHEN ISU =5. ONLY.

7.3) PROPORTIONAL LOADS FOR TIME DEPENDENT ANALYSIS  
 TRANSFER TO THIS OPTION OCCURS ONLY FOR TIME ANALYSES.

ONE CARD FOR EACH PROPORTIONAL LOAD REQUIRED  
 COL 1 TO 5 PROPORTIONAL LOAD TYPE. 1,2 OR 3  
 COL 6 TO 10 K. TABLE CONSTANT  
 COL 11 TO 20 TIME. SMALLEST TIME LOADING IS VALID  
 COL 21 TO 30 TIME. LARGEST TIME LOADING IS VALID  
 COL 31 TO 40 A0  
 COL 41 TO 50 A1  
 COL 51 TO 60 A2  
 COL 61 TO 70 A3  
 COL 71 TO 80 A4  
 LOAD TYPE 1. T = TIME

PROF = A0 + A1\*T + A2\*T\*T + A3\*T\*T\*T + A4\*T\*T\*T\*T  
 LOAD TYPE 2.

MAN550C  
 MAN557C  
 MAN558C  
 MAN559C  
 MAN560C  
 MAN561C  
 MAN562C  
 MAN563C  
 MAN564C  
 MAN565C  
 MAN566C  
 MAN567C  
 MAN568C  
 MAN569C  
 MAN570C  
 MAN571C  
 MAN572C  
 MAN573C  
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 MAN616C  
 MAN617C  
 MAN618C  
 MAN619C  
 MAN620C  
 MAN621C  
 MAN622C  
 MAN623C  
 MAN624C  
 MAN625C  
 MAN626C  
 MAN627C  
 MAN628C  
 MAN629C  
 MAN630C  
 MAN631C  
 MAN632C  
 MAN633C  
 MAN634C  
 MAN635C  
 MAN636C  
 MAN637C  
 MAN638C  
 MAN639C  
 MAN640C

PROP = A8\*(CIN(H1+I1))\*\*K + A2\*(CLOS(A3+T))\*\*K + M5

LOAD TYPE 3.

PROF = USER DEFINED FUNCTION FROM SUBROUTINE  
 EXPRLD(CPOP,T,AY) WHERE H IS AN ARRAY WITH  
 INFORMATION FOR COLUMNS 6-80 OF DATA CARDS.

\*\*NOTE\*\* PROPORTIONAL LOADS CAN BE ACCUMULATED FROM DIFFERENT  
 TYPES AT THE SAME TIME.

8.) INITIATION OF TIME INDEPENDENT SOLUTION (15.1X.A6)

COL 1 TO 5 IOUT. OUTPUT CONTROL CODE.

IOUT .EQ. 0. ALL STRESSES AND DISP. PRINTED  
 IOUT .NE. 0. SELECTED PRINTOUT. MORE DATA INPUT  
 SEE SECTION 9 FOR DATA PREPARATION.

COL 7 TO 12 MUST CONTAIN SOLVE \*INDICATES ALL DATA INPUT\*  
 COMPLETE FORMULATION AND SOLUTION OF EQUATIONS.  
 COL 7 TO 12 MUST CONTAIN RESOLV TO OBTAIN SUBSEQUENT  
 SOLUTIONS WHERE BOUNDARY CODES DO NOT CHANGE  
 AND ALL PRESCRIBED DISPLACEMENTS ARE ZERO.

8.1) INITIATION OF DYNAMIC SOLUTION BY EXPLICIT INTEGRATION.

COL 1 TO 5 IPR. OUTPUT CONTROL FOR DISPLACEMENT AND  
 STRESS PRINTOUT. SEE SECT. 9 FOR DATA INPUT.  
 IOUT = 1 - MIN(I,IPR)  
 IF IOUT .NE. 0. THE SPATIAL CONTROL DATA  
 COMES AT THE END OF THE DYNAMIC SEGMENT.  
 COL 7 TO 12 MUST CONTAIN EXPLIC

SUBSEQUENT CARDS (215.2F10.0.215)

COL 1 TO 5 NUMBER OF TIME STEPS  
 COL 6 TO 10 PRINT INTERVAL  
 COL 11 TO 20 TIME INCREMENT  
 COL 21 TO 30 NEWMARK DELTA-DAMPING TERM (GAMMA - .5)  
 COL 31 TO 35 NUMBER OF TIME EVOLUTION ELEMENT VARIABLE PLOTS  
 COL 36 TO 40 NPROP. NUMBER OF PROPORTIONAL LOADS TO BE INPUT  
 COL 41 TO 45 NFORC. LAST NODE ON WHICH A FORCE IS CHANGED  
 DURING EACH TIME STEP.  
 COL 46 TO 50 KKK. STABILITY CHECK OVERRIDE \*\* CAUTION USE  
 ONLY WHEN A BETTER ESTIMATE OF THE STABLE TIME  
 STEP IS AVAILABLE THAN CAN BE PERFORMED BY CODE  
 KKK ZERO. USES INTEGRAL STABILITY CHECK.  
 KKK NONZERO. DISREGARDS STABILITY CHECK.  
 KKKL. 0 FOR LINEAR  
 1 FOR NON LINEAR  
 COL 51 TO 55 LPE. 0 PRINTS LOADS  
 1 COMPRESS PRINT  
 COL 56 TO 60

SUBSEQUENT CARDS (315) ONE FOR EACH STRESS PLOT.

COL 1 TO 5 ELEMENT NUMBER CONTAINING STRESS TO BE PLOTTED.  
 COL 6 TO 10 LOCAL COORDINATE POINT CODE, 1 TO 9, AS  
 PATTERNED AFTER, COL 11 TO COL 19, IN SECT. 9.  
 COL 11 TO 15 PLOT COMPONENT CODE, 1 TO 6 FOR SIGMA(I,J),  
 I.E., SIGMA(1,1)=1, SIGMA(1,2)=2, SIGMA(1,3)=3,  
 SIGMA(2,2) = 4, SIGMA(2,3) = 5, SIGMA(3,3) = 6.

IF(NPROP,NE.0) READ PROPORTIONAL LOAD CARDS, SEE SECT. 7.3

IF(NFORC,NE.0) READ FORCE CARDS AT EACH TIME STEP. IF OUTPUT IS  
 LIMITED BY TOUT NONZERO, THE FIRST FORCE CARD SET PRECEDES  
 OUTPUT CARDS AND THE REMAINDER FOLLOW THE OUTPUT CARDS NO BLANK  
 CARDS MAY BE USED BETWEEN SETS OF CARDS OTHER THAN THE USUAL  
 BLANK TERMINATOR CARD FOR FORCE INPUT CARDS.

IF(IOUT,NE.0) DATA FOR SPATIAL PRINTOUT CONTROL. SEE SECT.9.

SPECIAL COMMENTS FOR DYNAMIC OPTION

- (1) ONLY COLUMNS 1 TO 66 ARE AVAILABLE FOR PAGE HEADING.
- (2) MAXIMUM ADVANTAGE OF ELEMENT REUSE OPTION SHOULD BE TAKEN.
- (3) INITIAL CONDITIONS FOR DISPLACEMENT AND VELOCITY VECTORS,  
 AS WELL AS STORAGE FOR ACCELERATION VECTOR, MAY BE MADE  
 THROUGH INPUT OF AN INITIAL CONDITION CARD SET, WITHOUT  
 SPECIFIED INITIAL CONDITIONS THEY ARE AUTOMATICALLY SET ZERO.
- (4) SPATIAL LOADING IS INPUT THROUGH FORCE OR BOUNDARY  
 PRESSURE CARDS.
- (5) EXTREME CAUTION ON ORDER OF DATA CARDS MUST BE OBSERVED. NO  
 EXTRA CARDS ARE PERMITTED AND STRICT COUNTS ARE OBSERVED  
 EXCEPT FOR THE NUMBER OF FORCE CARDS USED IN EACH TIME STEP.

8.2) INITIATION OF IMPLICIT TIME INTEGRATIONS (15,IX:A6)

COL 1 TO 5 NSEQ, NUMBER OF TIME SEQUENCES  
 COL 7 TO 12 MUST CONTAIN VISCOE FOR LINEAR VISCOELASTIC  
 QUASI-STATIC PROBLEMS (ONE INITIAL CONDITION  
 ONLY MUST BE USED)  
 COL 7 TO 12 MUST CONTAIN IMPLICIT FOR DYNAMIC IMPLICIT  
 INTEGRATIONS (THREE INITIAL CONDITIONS ARE  
 REQUIRED, MORE CAN BE SPECIFIED WITHOUT ERROR)

SUBSEQUENT CARDS. ONE SET FOR EACH TIME SEQUENCE

CARD 1. F10.0.815.2F10.0.215)

COL 1 TO 10 DT, TIME INCREMENT (NONZERO FOR IMPLICIT)  
 COL 11 TO 15 NTS, NUMBER OF TIME STEPS IN SEQUENCE  
 COL 16 TO 20 INT, PRINT INTERVAL (DEFAULT 1)  
 COL 21 TO 25 NHI, FIRST NODE PRINTED  
 COL 26 TO 30 NHE, LAST NODE PRINTED  
 COL 31 TO 35 NEI, FIRST ELEMENT STRESS TO BE PRINTED

MAN635C  
 MAN636C  
 MAN637C  
 MAN638C  
 MAN639C  
 MAN640C  
 MAN641C  
 MAN642C  
 MAN643C  
 MAN644C  
 MAN645C  
 MAN646C  
 MAN647C  
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 MAN690C  
 MAN691C  
 MAN692C  
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 MAN694C  
 MAN695C  
 MAN696C  
 MAN697C  
 MAN698C



COL 11 TO 20 1 - FORCE/DISPL. MULTIPLIER  
 COL 21 TO 30 2 - FORCE/DISPL. MULTIPLIER  
 COL 31 TO 40 3 - FORCE/DISPL. MULTIPLIER  
 COL 41 TO 50 4 - FORCE/USER MULTIPLIER  
 COL 51 TO 60 5 - FORCE/USER MULTIPLIER  
 COL 61 TO 70 6 - FORCE/USER MULTIPLIER  
 COL 71 TO 80 USER CONSTANT.

9.) OUTPUT CONTROL FOR LIMITED PRINTS

DISPLACEMENT OUTPUT CONTROL. IF IOUT.NE. 0.

CARD 1. (15)

COL 1 TO 5 NUMDIS - NUMBER OF DISPLACEMENT PRINT CARDS  
 SUBSEQUENT CARDS (215) SKIP IF NUMDIS = 0

COL 1 TO 5 NODAL NUMBER TO BE OUTPUT.  
 COL 6 TO 10 HIGHER NODE NUMBER OF A GENERATED SEQUENCE,  
 IF ZERO JUST FIRST NODE IS COUNTED.  
 COL 11 TO 15 INCREMENT TO GENERATOR. DEFAULT = 1  
 \*\*\* REPEAT UNTIL NUMDIS CARDS HAVE BEEN READ

STRESS OUTPUT CONTROL. IF IOUT.NE. 0.

CARD 1. (15.5X.911)

COL 1 TO 5 NUMSTR - NUMBER OF STRESS OUTPUT CARDS  
 COL 11 TO 19 NSIG(9) - PRINT PATTERN WITHIN AN ELEMENT.  
 LOCAL POINTS OF EACH ELEMENT CAN BE  
 SUPPRESSED BY NON-ZERO ENTRIES AS FOLLOWS.  
 E.G.

COL 11 SUPPRESS PRINT AT LOCAL POINT 1. ( 0. 0. 0 )  
 COL 12 SUPPRESS PRINT AT LOCAL POINT 2. (-1. 0. 0 )  
 COL 13 SUPPRESS PRINT AT LOCAL POINT 3. ( 1. 0. 0 )  
 COL 14 SUPPRESS PRINT AT LOCAL POINT 4. ( 0.-1. 0 )  
 COL 15 SUPPRESS PRINT AT LOCAL POINT 5. ( 0. 1. 0 )  
 COL 16 SUPPRESS PRINT AT LOCAL POINT 6. ( 0. 0.-1 )  
 COL 17 SUPPRESS PRINT AT LOCAL POINT 7. ( 0. 0. 1 )  
 COL 18 SUPPRESS PRINT AT LOCAL POINT 8  
 COL 19 SUPPRESS PRINT AT LOCAL POINT 9

SUBSEQUENT CARDS (215) SKIP IF NUMSTR = 0

COL 1 TO 5 ELEMENT NUMBER TO BE PRINTED.  
 COL 6 TO 10 HIGHER ELEMENT NUMBER OF A GENERATED SEQUENCE.  
 IF ZERO ONLY FIRST ELEMENT IS COUNTED.  
 COL 11 TO 15 INCREMENT TO GENERATOR. DEFAULT = 1  
 \*\*\* REPEAT UNTIL NUMSTR CARDS HAVE BEEN READ

MAN736E  
 MAN736C  
 MAN736C  
 MAN736C  
 MAN736C  
 MAN736C  
 MAN736C  
 MAN736C  
 MAN737C  
 MAN738C  
 MAN739C  
 MAN740C  
 MAN741C  
 MAN742C  
 MAN743C  
 MAN744C  
 MAN745C  
 MAN746C  
 MAN747C  
 MAN748C  
 MAN749C  
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 MAN761C  
 MAN762C  
 MAN763C  
 MAN764C  
 MAN765C  
 MAN766C  
 MAN767C  
 MAN768C  
 MAN769C  
 MAN770C  
 MAN771C  
 MAN772C  
 MAN773C  
 MAN774C  
 MAN775C  
 MAN776C  
 MAN777C  
 MAN778C  
 MAN779C  
 MAN780C

C. Listing of the Contact/Impact Subroutines Added to FEAP 74

The listings for the subroutines which are added to FEAP 74 for the contact/impact theory described herein are given below.

```

SUBROUTINE CONTACT(ISW,IX,NEL1,NDF,NUMPF,DU,U,UD,UDD)
C*****CONTACT ***** 04/29/74 *****
LOGICAL CFLAG,FLAG
DIMENSION IX(NEL1,NUMPF),DU(NDF,NUMPF),UD(NDF,NUMPF),
C UDD(NDF,NUMPF)
DIMENSION RMP(10)
COMMON/CONTACT/FLAG,CFLAG,L1ST,ICLIST(10),ICDEG(10),PU1,PU2,
C RM(2,10)
COMMON/LABELS HEAD(12),O,IPG,HEAD(3),UHED(6),XH,FX,WH,NSTR,FLAG(7)
COMMON/TAPES/ITPS,ITP6
DATA CFLAG,FLAG, FALSE...FALSE.
GO TO (1,2,3),ISW
READ(ITP5,1000) L1ST,PU1,PU2
PF = PU1+PU2
IF (PP.EQ.0.0) PF = 1.
READ(ITP5,1001) ICLIST(L),ICDEG(L),L=1,L1ST
WRITE(ITP6,2000) O,HEAD,IPS,PU1,PU2,(ICLIST(L),L=1,L1ST)
IPG = IPG + 1
FLAG = .TRUE.
RETURN
2 DO 200 L = 1,L1ST
N = ICLIST(L)
IDEG = ICDEG(L)
I = IX(1,N)
PH(1,L) = DU(IDEG,I)
I = IX(3,N)
PH(3,L) = DU(IDEG,I)
WRITE(ITP6,2002)
DO 220 L = 1,L1ST
N = ICLIST(L)
RMP(L) = RM(1,L) + RM(2,L)
IF (PH(1,L).EQ.0.0) PH(2,L) = PH(L) * L.E+30
WRITE (ITP6,2003) N,PH(1,L),L=1,L1,L2,PMFIL)
DO 210 N = 1,NUMPF
DC 210 I = 1,NDF
DU(I,N) = 0.0
RETURN
3 DO 300 L = 1,L1ST
N=ICLIST(L)
IDEG = ICDEG(L)
K = IX(2,N)
UD(IDEG,K) = 0.0
UDD(IDEG,K) = 0.0
IF (UIDEG,K) = 0.0,300,310
310 I = I+1
I = I-24
CU = PU1+UDD(IDEG,1)+PU2+UDD(IDEG,1)+PP
WRITE (ITP6,2004) L,PH(1,L),L=1,L1,PHFIL)
DO 310 L = 1,L1ST
N=ICLIST(L)
L2=IX(2,N)
PH(2,L) = RM(1,L) + RM(2,L)
PH(3,L) = RM(3,L)
END

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451      UDD(IDEG,J) = CU
J454      300      CONTINUE
000457      RETURN
000457      1000   FORMAT(15.5X,2F10.0)
000457      1001   FORMAT(2I5)
000457      2000   FORMAT(A1,12A6,30X,4HPAGE,14//5X,25HCONTACT SHEET DESCRIPTION//
C          10X,13HBODY 1 RO*U =,E13.5//
C          10X,13HBODY 2 RO*U =,E13.5//
X          11X,7HELEMENT,1X,9HDIRECTION (15,110))
000457      2001   FORMAT(110,6E13.4)
000457      2002   FORMAT(10A8  ELEMENT,13H BODY 1 MASS,13H BODY 2 MASS)
000457      END

```

ELM 10

SUBROUTINE ELMT05(N,MA,ND,IM,NDF,NEL,NEL1,NSTF,NV,IZV,NVEC,FACT,DM,D, X XYZ,IX,F,FORCE,ESTIF,U,VECT,ISW)

C\*\*\*\*\*ELMT05 \*\*\*\*\* 04/29 74 \*\*\*\*\*\*\*\*\*\* \*\*\*\*\*

```

000027 LOGICAL NPF
000027 DIMENSION U(NDF,1),IX(NEL,1),PLOT(8),THED(8)
000027 COMMON LOCALS PUL(6,20),UL(6,20),DUDL(6,20)
000027 COMMON PRTPLT NSIG(8),NPLT(9,2),NT,NSTEP,NURPLT,NEDATA(20,3),NPF
000027 COMMON SHAPE NARC,SHAPE(4,20),SG(3,3),S(3,3),R(3,20)
000027 COMMON TAPES ITP6,NPS
000027 COMMON TIME,TIME,INT,CTR,NA,NSCH,CO,CI,CC,C3,C4,C5,C6,FACT
000027 EQUIVALENCE (PLOT,TAU),(PLOT,ETA)
000027 DATA THED,SHATHU,SHETHA,SHU1,SHU2,SHU3,SHU4,SHU5,SHU6,SHU7,SHU8,SHU9,SHU10,SHU11,SHU12,SHU13,SHU14,SHU15,SHU16,SHU17,SHU18,SHU19,SHU20
000027 GO TO (1,2,3,4,5,6,7,8,9,10,11,12,13,14,15,16,17,18,19,20)
000041 READ(ITP6,1000) IDEG,NPS
000051 IDEG = MAX(IDEG,1)
000055 WRITE(ITP6,2000) IDEG,NPS
000066 TOL = 1.E-06
000073 I = NDF + IDEG
000074 J = I + NDF
000075 RI = 0.5
000076 RJ = 0.5
000077 CV = 0.
000077 NV = NPS
000100 NA = NPS + NPS
000101 RETURN
000102 DM = X(IDEG,3) - X(IDEG,1)
000106 DM = DM + 1.E-25
000110 RETURN
000110 TAU = UL(IDEG,2)
000112 D0 = X(IDEG,3) - X(IDEG,1) + (UL(IDEG,3) - UL(IDEG,1))
000120 ETA = 0.
000121 IF(DD.LT.TOL) ETA = 1.
000125 IF(TAU.LT.0.0) ETA = 0.0
000130 IF(ISW.GT.7) GO TO 4
000134 ESTIF(IDEG,1) = ETA
000142 ESTIF(1,J) = -ETA
000151 ESTIF(J,1) = -ETA
000155 ESTIF(1,1) = 1.0 - ETA
000163 RETURN
000163 IF(ISW.EQ.4) GO TO 6
000166 FORCE(1,1) = ETA*D0
000174 IF(ETA.EQ.0.0) FORCE(1,1) = -TAU
000203 TAU = AMAX1(0.0,TAU)
000207 FORCE(IDEG,1) = -TAU
000215 FORCE(J,1) = TAU
000220 IF(.NOT.NPF) WRITE(ITP6,3000) N,DM,NA,TAU,ETA
000245 IF(NURPLT.LE.0) SE UPN
000246 PLOT(3) = J(1),IEG,1
000250 PLOT(4) = X(1),IEG,1
000262 PLOT(5) = U(1),IDEG,1
000263 PLOT(6) = U(1),IEG,1

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000256      PLOT(7) = UDL(IDEG,3)
000260      PLOT(8) = UDL(IDEG,3)
000262      DO 60 K = 1,NUMPLT
000263      KK = NPLT(K,2)
000265      IF (NPLT(K,1).GT.0) CALL FLDATA(NDIM,NPLT(K,1),THED(KK),N(1,2))
          C      PLOT(KK)
          C      CONTINUE
          C      RETURN
          C      FORMAT(2I5)
          C      FORMAT(5X,2IHP0INT CONTACT ELEMENT
          C      5X,27HC0NTACT DEGREE OF FREEDOM = 13.5X,5HP0S = 15)
          C      FORMAT(5X,7HELEMENT,15.5X,A5.5X,8HMATERIAL,13.5X,5HTAU =,E15.5,5X,
          C      5SHETHA =,F3.1)
          C      END
000306      60
000311      5
000312      1000
000312      2000
000312      2001
000312      END
000312

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SUBROUTINE ELMT09(N,IA,NDIM,NDF,NELM,NEL1,NSTF,NSIZV,NVEC,NCT,DM,
X D,XYZ,IX,F,FORCE,ESTIF,U,VECT,ISW)
LOGICAL NORPNT,NPR,NPL
DIMENSION ESTIF(NSIF,NSTF),D(3,21,1),IX(NEL1,1),U(NDF,1),
X V(3,20),NU(3),XY(3),FORCE(NSIF,2)
COMMON GAUSS,ELIM,SGAUSS(5,5),WGAUSS(5,5)
COMMON LABELS,HEAD(12),O,IPG,XHEAD(3),UHED(6),ZH,FA,UH,NSTP,FLAG(7)
COMMON LOCALS, IUL(6,20),UL(6,20),UPL(6,20),UDDL(6,20)
COMMON ORTALT,NSIG(9),NPLT(9,2),NT,NSTEP,NUMPLT,MEDOTM(20,3),NPP
COMMON SHAP, XJAC, SHAPE(4,20),SG(3,3),SK(3,3),Z(3,20)
COMMON TAPES, ITP5, ITPS
COMMON TIMHIS, TIME,DT,DTP,NA,ISEH,C0,C1,C2,C3,C4,C5,C6,NCT
DATA SH,EH,BL,SHSTRESS,SHSTRAIN,SH
GO TO (1,2,3,4,5,4), ISW
CONTINUE
1 C..... BAR STIFFNESS CHARACTERIZATION
READ(ITP5,1000) E,AA,RO
WRITE(ITP5,2100) E,AA,RO
D(1,1,NA) = E
D(1,2,NA) = AA
D(1,3,NA) = RO
C5 = 0.
RETURN
2 CONTINUE
RETURN
3 C..... COMPUTATION OF BAR STIFFNESS NO BENDING * ISOPARAMETRIC 1ST AND 2M
CONTINUE
C..... CONSTRUCT STIFFNESS AT INTEGRATION POINTS
EA = D(1,1,NA)*D(1,2,NA)
RA = D(1,3,NA)*D(1,2,NA)
DO 260 II = 1,NELM
55 = SGAUSS(II,NELM)
CALL LINEISS,NDIM,NELM)
UM=WGAUSS(II,NELM)
DVOL=EA*UMJ*XJAC*KJAC*KJAC)
C..... COMPUTE A LUMPED MASS MATRIX
IU = 0
PAH=PA*KJAC*UMJ
DO 217 I = 1,NELM
AA=RAH*SHAPE(2,I)
DO 215 FN = 1,NDIM
FORCE(IU+K*2) = FORCE(IU+K,2) + AA
215 IU = IU + NDF
217
DO 250 NY = 1,NDIM
DO 250 LY = 1,NELM
SGD=SG(1,1)*LLI*FVUL
II = II + 1
DO 240 I = 1,NELM
SHI=SHAPE(1,I)*SGD
L = L + 1
TEU = U(1,1)*L = 1
M = M + SHI*U(1,1)
L = L + 1
240 I = 1,NELM
250

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ELM 1C  
ELM 1C  
ELM 4C  
ELM 5C  
ELM 3C  
  
ELM 10C  
ELM 11C  
ELM 13C  
ELM 14C  
  
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ELM 47C  
ELM 44C  
ELM 45C  
ELM 46C

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000573      615      FORCE(KK+IU,1) = FORCE(IY+U,1) - AR*MOD(LETK,1) - BB*(S(FF),1)
000615      600      IU = IU + NDF
000621      400      CONTINUE
000624      5      RETURN
000635      1000     FORMAT(3F10.0)
000635      2000     FORMAT(14I:12H6,E13.5,17F:4HPAGE,13X:10H  ELEMENT,6:2PHMATER,10)
000625      2001     FORMAT(110,15.5%,45.5E12,4)
000625      2100     FORMAT(30HLINEAR ELASTIC MATERIAL, BAR ELEMENT//
000625      4000     X 5:1.7HE =,E15.5,5:1.6HAREH =,E15.5,5X:9HDENSITY =E15.5,1)
000625      4000     FORMAT(45.15,1F3E12,5.24X,110)
000625      4000     END
    
```

ELM 810  
 ELM 830  
 ELM 860  
 ELM 970  
 ELM 980  
 ELM 990

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000006 SUBROUTINE LINE(INDIM,NEL)
000007 C*****
000008 C..... SHAPE FUNCTION ROUTINE FOR 15 MEMBER QUADRATIC ELEMENTS
000009 C..... COMMON SHAP = JAC, SHAPE(1,20), SG(1,3), S(1,3), S(1,2)
000010 C..... FORM SHAPE FUNCTIONS
000011 SHAPE(1,1) = 0.5
000012 SHAPE(2,1) = 0.5*(1.0-S)
000013 SHAPE(1,NEL) = 0.5
000014 SHAPE(2,NEL) = 0.5*(1.0+S)
000015 IF (NEL.EQ.2) GO TO 350
000016 SHAPE(1,2) = -2.0*S
000017 SHAPE(2,2) = 1.0-S+S
000018 V = NEL-1
000019 DO 100 I = 1,NEL,K
000020 SHAPE(1,I) = SHAPE(1,1) - 0.5*K*SHAPE(1,2)
000021 SHAPE(2,I) = SHAPE(2,1) - 0.5*K*SHAPE(2,2)
000022 C..... FORM JACOBIAN MATRIX AND DETERMINANT
000023 DO 360 I = 1,3
000024 SK(1,I) = 0.0
000025 DO 400 J = 1,NEL
000026 SK(1,I) = SK(1,I) + X(I,J)*SHAPE(1,J)
000027 XJAC = 0.0
000028 DO 500 I = 1,NDIM
000029 XJAC = XJAC + SK(1,I)*S(I,1)
000030 DO 500 J = 1,NDIM
000031 SG(I,J) = SK(1,I)*S(I,J)
000032 XJAC = SOPT(XJAC)
000033 RETURN
000034 END

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LIN 10  
 LIN 20  
 LIN 30  
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 LIN 60  
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 LIN 100  
 LIN 110  
 LIN 120  
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 LIN 270  
 LIN 280  
 LIN 290

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000017 SUBROUTINE HLNORM( ISW, MDEG, NH, U, DU, DF, F, PROP, NCT, NTT, IDEST, IBLK )
000018 C***** NLNORM ***** 12 14 73 *****
000019 DIMENSION U(1), DU, MDEG, DF(1), F(1), IDEST(1)
000020 COMMON NORMS UNORM, DNORM, ANORM, CS, CSP, DP, DNP, IFL, XF, AG, EFF
000021 GO TO (100,200,200,100), ISW
000022 .00 CS=1.0
000023 CSF=1.0
000024 DP=0.75
000025 IFL=0
000026 XFLAG=.FALSE.
000027 ERR=1.0E-3
000028 RETURN
000029 200 UNORM = 0.
000030 DNORM = 0.
000031 ANORM=0.
000032 DO 500 N = 1, MDEG
000033 DUN = DU(N)
000034 UN = U(N)
000035 IF (ISW.EQ.3) UN = UN + DF(N) + DUN
000036 ANOR,1=ANOR,1+UN*DUN
000037 UNORM=UNORM+UN*DUN
000038 DNORM=DNORM+DUN*DUN
000039 UNORM=SQRT(UNORM)
000040 DNORM=SQRT(DNORM)
000041 IF (NH.EQ.1) DNP=UNORM
000042 CS=1.0
000043 IF (UNORM*DNORM.NE.0) CS=ANORM/(UNORM*DNORM)
000044 IF (IFL.EQ.1) DP=0.75
000045 IF (IFL.EQ.1.AND.XFLAG) CSP=CS
000046 IFL=0
000047 IF (DNORM.LE.0.5*UNORM) GO TO 550
000048 IFL=1
000049 DP=0.5*UNORM/DNORM
000050 550 IF (DP*DNORM.GT.DNP) DP=DNP/DNORM
000051 IF (DF.EQ.0) DF = 1.0
000052 IF (CS*CSP.LT.0) DP=DP*2.
000053 IF (CS*CSP.GT.0) IP=1.25*IP
000054 DNP=DF*DNORM
000055 CSF=CS
000056 IF (ISW.EQ.2) RETURN
000057 DO 600 N = 1, MDEG
000058 DF(N) = DF(N) + DUN(N)
000059 IF (NCT.EQ.0) DF(N) = U(N)
000060 DUN(N) = DF(N)
000061 PERM=H
000062 END

```

NLN 10  
NLN 20  
NLN 30  
NLN 40  
NLN 50  
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NLN 80  
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NLN 400  
NLN 410  
NLN 420  
NLN 430  
NLN 440  
NLN 450





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000152 DTP = 0.
000153 TIME = 0.
000154 MB = MAXBAN + 1
000156 NEI = NEN + 1
C..... INITIALIZE INCREMENTAL FORCE VECTOR
6 DO 6 N = 1,NSICD
DF(N) = 0.0
DO 10 N = 1,NUMEL
10 IX(NEI,N) = 0
C..... PUT INITIAL DATA ON TAPE
ISZH = IBUF
ITRD = ITP13
ITUR = ITP14
REWIND ITRD
IF(IBLK.GT.0) WRITE(ITRD) ((U(I,J),I=1,NDF),J=1,KU)
11 H(N) = 0.
NTB = (NH+ISZH-1)/ISZH
WRITE(ITRD) H
IF(NTB.LE.1) GO TO 13
NT = NTB*2
DO 12 N = 1,NT
12 WRITE(ITRD) H
REWIND ITRD
IF(IBLK.GT.0) READ(ITRD) DM
NEP = NUMNP + NUMNP
C6 = 0.0
NSTEP = 0
NST = 0
NUMPLT=0
PROP = 0.
CFLAG = .FALSE.
DO 900 M = 1,NSEQ
READ(ITP5,1000) DT,NTS,INT,NNI,NNF,NEI,NEF,NPROP,NFORC,C0,DM,NT,NL
IF(NICD.GT.1) C6= DM
IF(NL.NE.0) NST= NL
NTT = IABS(NST)
DFLAG = .FALSE.
IF (NTT.NE.0) DFLAG = .TRUE.
NSTEP=NSTEP+NTS
IF(M.EQ.1) NUMPLT=NT
IF(INT.LE.0) INT = 1
WRITE(ITP6,2000) 0,HEAD,TIME,IPG,DT,NTS,INT,NNI,INF,NEI,NEF
IF(DFLAG) WRITE(ITP6,2032) NST
IPG = IPG + 1
IF(NICD.EQ.1) GO TO 1
IF(OT.EQ.0.0) GO TO 901
CALL UPDATE(3)
IF(NPROP.GT.0) PROP = PROFLO(TIME,NPROP)
IF(NFORC.GT.0) AND,NPROP,GT.0) WRITE(ITP5,-4031)
C..... READ STRESS PLOT INFORMATION
1 IF(NUMPLT.LE.0) OR(M.GT.1) GO TO 50C
REWIND 12
WRITE(ITP6,2005)
0536
0537
0538

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TSO 44C
TSO 45C
TSO 46C
TSO 47C
TSO 48C
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TSO 88C
TSO 89C
TSO 90C
TSO 91C
TSO 92C
TSO 93C
TSO 94C
TSO 95C
TSO 96C

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00555 DO 501 N=1,NUMPLT
00562 READ(ITP5,1006) (NEDATA(N ,I),I=1,2)
000600 WRITE(ITP6,2006) N, (NEDATA(N ,I),I=1,3)
000632 CALL PLZERO
000633 VFLAG = .FALSC.
000634 DO 500 NT = 1,NTS
000641 PPROP = PROP
000643 PROP = 1.0
000644 IF (DFLAG) CALL NLNORM(4,MDEG,NT,U,DF,DU,F,PPROP,NCT,NTT, IDEST,IBLK)
000667 IF (NPROP.GT.0) PROP = PROPLD(TIME+DT,0)
000701 IF (NFORC.GT.0) CALL RESET(-NFORC,NUMNP,NDF,F)
000713 NCT = 0
14 REWIND ITR
000714 IF (IBLK.EQ.0) GO TO 15
000716 IF (.NOT.VFLAG) GO TO 20
000720 DO 21 I = 1,NDEG
000723 DF(I) = 0.
GO TO 20
000730 DO 19 N = 1,NUMNP
000732 DO 18 K = 1,NDF
000733 J = IDEST(K,N)
000740 IF (J.EQ.0) GO TO 18
000741 IF (VFLAG) GO TO 17
000743 DO 16 I = 1,MAXBAN
000744 A(J,I) = 0.0
000756 DF(J) = F(K,N)*PROP
18 CONTINUE
19 CONTINUE
000771 NH = 1
000773 IF (NTB.GT.1) READ(ITRD) H
000774 IF (NCT.GT.0) GO TO 44
C..... OUTPUT THE SOLUTION VECTOR FOR THE CURRENT TIME
IF (MINI.EQ.0.OR.(NT-1)/INT)*INT.NE.NT-1) GO TO 40
MCT = 0
DO 30 N = 1,NUMNP
MCT = MCT + 1
IF (MCT.GT.0) GO TO 31
IF (NICD.EQ.1) WRITE(ITP6,2001) 0,HEAD,TIME,IPG,PPROP,M,NT,
X (XHED(I),XH,I=1,NDIM),(UHED(I),UH,I=1,NDF)
IF (NICD.NE.1) WRITE(ITP6,2001) 0,HEAD,TIME,IPG,PPROP,M,NT,
X (XHED(I),XH,I=1,NDIM),(UHED(I),UH,I=1,NDF),(UHED(I),UH,I=1,NDF)
X (UHED(I),UH,I=1,NDF)
IPG = IPG + 1
MCT = 50
31 IF (NICD.EQ.1) WRITE(ITP6,LABLON,(XYZ(I,N),I=1,NDIM),(U(I,N)
X ,I=1,NDF)
IF (NICD.NE.1) WRITE(ITP6,LABLON,(XYZ(I,N),I=1,NDIM),(U(I,N),I=1,NDF
X ),(U(I,NUMNP,N),I=1,NDF),(U(I,NEP+N),I=1,NDF)
CONTINUE
30 UPDATE(I) * * FOR DYNAMIC SOLUTIONS ONLY
C..... IF (NICD.GT.1) CALL UPDATE(-MDEG,NICD,U,(1,NUMNP+1),(U(1,NEP+1),
X DU,F,DF,IDEST,PROP)
IF (IBLK.GT.0) WRITE(ITMP) (U(I,J),I=1,NDF), J=1+K
MP = 0
001235 TS0136C
001237 TSC.37C
001240
001311 TS0139C
001420 TS0140C
001423 TS0141C
001457 TS0142C
001512 TS0143C

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001513 TEMP = 0.
001514 DO 400 N = 1,NUMEL
001515 NPR = .TRUE.
001516 DO 43 I = 1,9
001520 NPLT(I,1) = 0
001521 NPLT(I,2) = 0
001524 43 IF(NCT.GT.0) GO TO 46
001526 IF(N.GE.NCI.AND.N.LE.NEF) NPR = .FALSE.
001540 IF(NUMPLT.LE.0) GO TO 46
001542 DO 45 I = 1,NUMPLT
001543 IF(NEDATA(I,1).NE.N) GO TO 45
001546 J = NEDATA(I,2)
001547 NPLT(J,1) = I
001551 NPLT(J,2) = NEDATA(I,3)
001553 CONTINUE
001556 45 MA = MOD(IX(HEL1,N),100)
001567 46 IF(MR.LE.0) MRR = IX(HEL1,N)/1000
001600 IF(MR.LE.0) MR = MRR
001603 DO 60 I = 1,NSTF
001605 FORCE(I,1) = 0.
001612 FORCE(I,2) = 0.
001616 LD(I) = 0
001620 IF(MR.NE.MRR.OR.VFLAG) GO TO 60
001625 DO 50 J = 1,NSTF
001627 ESTIF(I,J) = 0.0
001640 CONTINUE
001643 L = 0
001644 DO 110 I = 1,NEN
001645 K = IX(I,N)
001652 DO 90 J = 1,NDIM
001654 X(J,I) = 0.
001661 IF(K.EQ.0) GO TO 120
001662 NEL = I
001663 DO 100 J = 1,NDIM
001665 X(J,I) = XYZ(J,K)
001701 DO 110 J = 1,NDF
001702 L = L + 1
001704 DUL(J,I) = DU(J,K)
001704 UL(J,I) = U(J,K)
001722 IF(NCT.GT.0) UL(J,I) = UL(J,I) + DUL(J,I)
001731 IF(NICD.EQ.1) GO TO 110
001734 UDL(J,I) = U(J,K+NEP)
001744 UDDL(J,I) = U(J,K+NEP)
001754 IF(NCT.GT.0) UDDL(J,I) = UDDL(J,I) + CS*DUL(J,I)
001765 LD(I) = IDEST(J,K)
002001 CM = TYPE(MA)
002004 IF(MI.NE.TEMP) MCT = 0
002007 TEMP = TM
002010 CALL TICCTO(TIME,5)
C..... COMPUTE ELEMENT STRESSES AND UPDATE FORCES
C..... CALL ELMBIN(M,NDIM,NDI,HEL,HEL1,NSTF,NEIZV,NECV,MCT,DM,D,X,Y,Z)
C..... X,IX,H,FORCE,ESTIF,UMVECT,5)
C..... CALL TICCTO(TIME,5)
C..... FORM STIFFNESS IF NEEDED FOR THE NEXT TIME STEP

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002052 IF (.NOT. VFLAG.AND. NR.EQ. IPR) OR. (VFLAG.AND. IX.NE1.N.NE.EQ.1)
XCALL ELM18(N,MA,NDIM,NDF,NEL,NEL1,NSTF,NSIZY,NVEC,MCT,DM,D,X,Y,Z,
X IX,H,FORCE,ESTIF,U,VECT,3)
IF (MOD(IX,NEL1.N),1000)/100.GT.0.AND..NOT.VFLAG)
X CALL PPTMAT(HEAD,IPG,N,NSTF,ESTIF,FORCE,LD,NSTF,0)
C..... MODIFY FOR THE DISPLACEMENT B.C.
X CALL MODIF(NDF,NEL,NEL1,NEL2,IBLK,NSTF,PROP,IX,ICOD,F,FORCE,ESTIF,
X H)
IF (VFLAG) GO TO 300
IF (IBLK.EQ.0) CALL COMPIN(A,DF,NEOB,ESTIF,FORCE,LD,NSTF)
IF (IBLK.GT.0) WRITE(7) ESTIF,FORCE,LD
IF (.NOT.FLAGC.OR.CFLAG) GO TO 400
L = 1
DO 130 I = 1,NEL
K = IX(I,N)
DO 130 J = 1,NDF
DU(J,K) = DU(J,K) + FORCE(L,2)
L = L + 1
GO TO 400
300 CONTINUE
C..... ADD THE FORCE TO THE SOLUTION FOR A RESOLVE
DO 310 K = 1,NSTF
J = LD(K)
IF (J.GT.0) DF(J) = DF(J) + FORCE(K,1)
CONTINUE
CALL TICTOC(TYME,3)
IF (FLAGC.AND..NOT.CFLAG) CALL CONTAC(2,IX,NEL1,NDF,NUMNP,DU)
CFLAG = .TRUE.
IF (NTB.GT.1) WRITE(ITUP,H)
IF (IBLK.GT.0) WRITE(ITUR) (IX(NE1.N),N=1,NUNEL)
IF (.NOT.VFLAG) CALL SOLVEQ(NUMNP,NUNEL,NDF,IDI,M8,MAXBAN,9,NSTF,
1 ISZA,NEOB,IBLK,A,DU,DF,IDEST,FORCE,ESTIF,LD,MAXB,NDEG)
IF (VFLAG) CALL RESVEQ(NUMNP,NDF,M8,MAXBAN,ISZA,NEOB,IBLK,A,DF,DF,
1 IDEST,IDEST,MAXB,IFLG)
CALL TICTOC(TYME,4)
C..... UPDATE THE SOLUTION
IF (NT.EQ.NTS.AND.M.EQ.NSEQ) GO TO 960
IF (NCT.GT.0) GO TO 410
DTP = DT
I = ITRD
ITRD = ITUP
ITUP = I
IF (IBLK.GT.0) BACKSPACE ITRD
IF (IBLK.GT.0) READ(ITRD) (IX(NE1.N),N=1,NUNEL)
REWIND ITRD
IF (IBLK.GT.0) READ(ITRD) (DU(I,J),I=1,NDF,J=1,KU)
IF (VFLAG) CALL RESORP(N,NVEC,NNT,J,DF,DU,F,PEPOP,NCT,NTT,IDEST,IBLK)
VFLAG = .TRUE.
IF (NCT.GT.0) VFLAG = .FALSE.
NCT = NCT + 1
IF (.NOT.IFLG) GO TO 960
WRITE(6,900) (J,DU(I,J),I=1,NDF,J=1,KU)
IF (NCT.GT.0) WRITE(6,900) (J,DU(I,J),I=1,NDF,J=1,KU)
GO TO 14
900 IF (J-1) 45,15,900,900,900

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TS0190C  
TS0199C  
TS0200C  
TS0201C  
TS0202C  
TS0203C  
TS0204C  
TS0205C  
TS0206C  
TS0207C  
TS0208C  
  
TS0209C  
TS0210C  
TS0211C  
TS0212C  
TS0213C  
TS0214C  
TS0215C  
TS0216C  
  
TS0217C  
TS0219C  
TS0220C  
TS0221C  
TS0222C  
TS0223C  
TS0224C  
TS0225C  
TS0226C  
TS0227C  
TS0228C  
TS0229C  
TS0230C  
TS0231C  
TS0232C  
TS0233C  
TS0234C  
TS0235C  
TS0236C  
TS0237C  
TS0238C  
TS0239C  
  
TS0047C

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C..... UPDATE (2)
002733 CALL UPDATE(2,MDEG,NICD,U(U(1,NUMP+1)),U(1,NEF+1),DU,F,DF,IDEST,
X PROP)
002751 IF (FLAG) CALL CONTRA(3,1,4,NEL,1,DEF,NUMP,DU,U,U(1,NUMP+1),
X U(1,NEF+1))
003015 MP = MP - 1
003020 TIME = TIME + DT
003024 CONTINUE
003027 IF (NUMPLT.GT.0) CALL PLOT60(NDIM,NEL,DM,U)
003040 CALL TIC10C (TIME,6)
003043 WRITE (ITP6,2033) TIME
003052 RETURN
003053 WRITE (ITP6,2030)
003057 ITPG = 0
003060 RETURN
C.....
003061 FORMAT(F10.0,8I5,2F10.0,2I5)
003061 FORMAT(3I5)
003061 FORMAT(6I,1206,E13.5,17X,4HPAGE,14//5X,23HTIME DEPENDENT SOLUTION
X//10X,17HTIME INCREMENT = 1PE15.4/
X 10X,17HNUMBER OF STEPS = 15/
X 10X,17HPRINT INTERVAL = 15/
X 10X,14HPRINT NODES ,15,3H TO,15/
X 10X,14HPRINT ELEMENTS,15,3H TO,15)
003061 FORMAT(A1,12A6,E13.5,17X,4HPAGE,14//5X,12HMODAL VALUES,5X,
X 27HPROPORTIONAL LOAD FACTOR = ,F8.3,5X, 8HSEQUENCE,14,5X,
X 9HTIME STEP,14//12H MODAL POINT,9(1X,2A6), (12X,9(1X,2A6))) )
003061 FORMAT(10X,17HITERATION LIMIT =,15,1X)
003061 FORMAT(4X,50H DESCRIPTION OF STRESS EVOLUTION PLOTS TO BE MADE, /
X 5X,10H PLOT NO. 5X,8H ELEMENT ,7X, 9H XYZ-CODE,6X,5H SIG-CODE/)
003061 FORMAT(4(7X,15,3X))
003061 FORMAT(5X,53H**FATAL ERROR 33** TIME STEP ZERO FOR DYNAMIC PROBLEM
X )
003061 FORMAT (2X,12HRELEASE: TIME//10X,25HINPUT PROPERTIES AND MESH
X F10.3,10X,25HCHECK AND PUT INPUT DATA,F10.3,10X,14HFORN STIFFNESS
X F10.3,10X,21HSELECTION OF ELEMENTS,F14.3//10X,15HOUTPUT STEPS,
X F10.3,10X,27HIMPLICIT SOLUTION, 9I5PL.,F8.3,10X,19HTOTAL TIME//11
X PLT,F16.3,10X,
X F10.3,10X,25HRELEASE: MESHING 7M, BOTH THE PROPORTIONAL LOADING AND FOR
X ARE BEING PRESENT ON EACH TIME STEP )
003061 FORMAT(5X,53H**FATAL ERROR 33** TIME STEP ZERO FOR DYNAMIC PROBLEM
X )
003061 END

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TS0245C  
TS0246C  
TS0247C  
  
TS0248C  
TS0249C  
TS0250C  
TS0251C  
TS0252C  
TS0253C  
TS0254C  
TS0255C  
TS0256C  
TS0257C  
TS0258C  
TS0259C  
TS0260C  
TS0261C  
TS0262C  
TS0263C  
TS0264C  
TS0265C  
TS0266C  
TS0267C  
  
TS0270C  
TS0271C  
TS0272C  
TS0273C  
TS0274C  
TS0275C  
TS0276C  
TS0277C  
TS0278C  
TS0279C  
TS0280C  
TS0281C  
TS0282C  
TS0283C

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000016 SUBROUTINE UPDATE (ISU,NDEG,NICD,U,UD,UDD,DU,F,DF,IDEST,PROP)
000016 C***** UPDATE ***** 12/14/73 *****
000016 DIMENSION U(1),DU(1),F(1),DF(1),IDEST(1),UD(1),UDD(1)
000016 COMMON TAPES,ITP5,ITP6
000016 COMMON TIMHIS,TIME,DT,DTP,NH,ISZH,C0,C1,C2,C3,C4,C5,C6,NCT
000016 GO TO (100,300,500),ISU
000024 UPDATE (1) * * PREUPDATE OF ACCELERATIONS FOR DYNAMIC SOLUTIONS.
000026 DO 200 N = 1,NDEG
000035 UDD(N) = C4*UDD(N) - C5*U(N)
000035 C.....
000037 UPDATE (2) * * UPDATE THE SOLUTION FOR GENL. DISPL. VEL. ACCL.
000043 DO 400 N = 1,NDFG
000053 DU(N) = DF(N)
000056 IF (IDEST(N).EQ.0) DU(N) = F(N)*PROP - U(N)
000056 TEMP = DU(N)
000060 IF (NICD.EQ.1) GO TO 400
000061 P = UDD(N)
000067 UD(N) = C1*UD(N) + C2*P + C3*TEMP
000073 UDD(N) = P + C6*TEMP
000100 U(N) = U(N) + TEMP
000100 RETURN
000100 C.....
000102 UPDATE(3) * * SET INTEGRATION CONSTANTS
000102 C6 = C6 + 0.5
000104 IF (C0.EQ.0.0) C0 = 0.25
000104 WRITE (ITP6,2002) C0,C6
000114 C5 = 1./C0/DT
000116 C4 = 1. - 0.5/C0
000121 C3 = C5*C5
000127 C2 = 1. - C6/C0/2.)*DT/C4
000134 C1 = 1. - C6/C0 + C2+C5
000137 C6 = C5*DT
000137 RETURN
000137 FORMAT(10X,30HNEWMARK INTEGRATION PARAMETERS/
000137 C10X17HBETA VALUE = F6.3/
000137 C10X17HGAFFA VALUE = F6.3)
000137 END
000137
UPD 1C
UPD 2C
UPD 3C
UPD 4C
UPD 5C
UPD 6C
UPD 7C
UPD 8C
UPD 9C
UPD 10C
UPD 11C
UPD 12C
UPD 13C
UPD 14C
UPD 15C
UPD 16C
UPD 17C
UPD 18C
UPD 19C
UPD 20C
UPD 21C
UPD 22C
UPD 23C
UPD 24C
UPD 25C
UPD 26C
UPD 27C
UPD 28C
UPD 29C
UPD 30C
UPD 31C
UPD 32C
UPD 33C
UPD 34C
UPD 35C

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