

Flow analysis in a channel with flexible vegetation using double-averaging method

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Abstract

The paper addresses the problem of the resistance due to vegetation in an open channel flow, characterized by partially and fully submerged vegetation formed by colonies of bushes. The flow is characterized by significant spatial variations of velocity between vertical profiles that make the traditional approach based on time averaging of turbulent fluctuations inconvenient. A more useful procedure, based on time and spatial averaging (Double-Averaging Method) is applied for the flow field analysis and characterization. The vertical distribution of mean velocity and turbulent stresses at different spatial locations has been measured with a 3D Acoustic Doppler Velocimeter (ADV) for two different vegetation densities where fully submerged real bushes (*salix pentandra*) have been used. Velocity measurements were completed together with the measurements of drag exerted on the flow by bushes at different flow depths. The analysis of velocity measurements allows depicting the fundamental characteristics of both the mean flow field and turbulence. The experimental data show that the contribution of form-induced stresses to the momentum balance cannot be neglected. The mean velocity profiles and the spatially averaged turbulent intensity profiles allow inferring that the vegetation density is a driving parameter for the development of a mixing layer at the canopy top in the case of submerged vegetation. Moreover, the net upward turbulent momentum flux, evaluated with the methodology proposed by Lu and Willmarth (1973), appears to be damped for increased vegetation density; this finding can rationally explain the reduction of the suspended sediment transport capacity typically observed in free surface flows over a vegetated bed.

Key words: vegetated channels, turbulence, Double-Averaging Method.

1. INTRODUCTION

The importance of vegetation as a factor of improvement of the environmental value of rivers is widely recognized and the use of living bushes in restoring regulated rivers and mountain streams is encouraged. Therefore, hydrophilous riparian plants, like willows, are not only naturally present on the banks and floodplains and sporadically also in riverbeds, but are often used as living material for banks preservation as well. The proper maintenance practices tend to preserve the plants elasticity characteristics (e.g., by limiting plant growth by means of periodical lopping), in order to allow vegetation to promptly bend and adapt to floods and to assume a more hydrodynamic shape. Therefore, the problem of hydrodynamic behaviour of flexible bushes fully submerged in the flow is quite common in hydraulics of both natural and restored streams. Willow height is of same order of magnitude as the flow depth, causing drag and energy losses. The resistance that willows create to the flow can be evaluated as the combined effect of the hydrodynamic drag of single plants (Petryk and Bosmanijan 1975), according to the equation:

$$F = C_D A_p \rho \frac{U^2}{2}, \quad (1)$$

where F is the drag force exerted on a single plant, A_p is the area of the plant projected in streamwise direction (**frontal area** in the following), ρ is the water density, C_D is the drag coefficient of a single plant, and U is the mean flow velocity. The plant drag per unit bed surface area and unit plant height can be evaluated with the following equation:

$$f_V = nF/h_p, \quad (2)$$

where n is the number of plants per unit bed surface area and h_p is the plant height. Therefore, the net bed shear stress τ'_b is given by the equation:

$$\tau'_b = \rho g R_h S_b - f_V h_p \approx \rho g h S_b - f_V h_p, \quad (3)$$

where g is the gravitational acceleration, R_h is the hydraulic radius, S_b is the bed slope, and h is the flow depth.

The drag can be included into rational formulations for the evaluation of the flow resistance (see e.g., Ming and Shen 1973, Petryk 1975, Righetti and Armanini 2002). The physical conditions of riverine vegetation, e.g., plant species, age and vegetation density, are extremely variable, which strongly affects the resistance coefficient and has to be properly taken into account in the hydraulic calculations. Moreover, the flow velocities are usually very low in vegetated areas or wetlands compared to rivers, so turbulent characteristics become important because of the sediment deposition and erosion in the vegetation zone (Lopez and Garcia 2001, Nepf and Vivoni 1999). The presence of plants strongly affects the structure of the turbulent flow through

and above the vegetation: the straining of motion due to stems and the wake turbulence they generate, significantly affect the turbulent energy dynamics (Nepf 1999). The effect of fully submerged vegetation on the flow field can be interpreted as a low relative submergence roughness, characterized by spatial variability of the length scales of the same order of magnitude as the flow depth. In this case, a proper approach derives from a double-averaging procedure on the classical Navier-Stokes equations, where the classical time-averaged Reynolds equations should be supplemented by spatial area averaging procedure in the plane parallel to the average bed. A similar procedure has been previously applied to terrestrial canopies (Wilson and Shaw 1977, see Finnigan 2000) and can be considered to be a reference framework also for vegetated rivers – even though differences exist in the turbulence structure of the unconfined flow field above canopies and confined free surface flows (Nepf and Vivoni 1999) – as well as to rough bed open channel flows with small relative submergence (Nikora *et al.* 2001).

Studies on terrestrial canopies (Raupach *et al.* 1996) and experimental observation in laboratory open channel flows (Ikeda and Kanazawa 1996, Carollo *et al.* 2002, Ghisalberti and Nepf 2002, Ghisalberti and Nepf 2004) demonstrate that the flow within and above vegetation shows a characteristic *S* shaped velocity profile, thus a mixing layer analogy seems to be more appropriate than a boundary layer analogy for its interpretation. These studies have been mainly based on laboratory experiments conducted with artificial models of vegetation, such as rigid cylinders or plastic strips, or with small scale flexible plants, such as macrophytes, but flows over riparian willows have not been studied yet. In the literature, studies of real vegetation generally focused on wetland's plants (e.g., sedges), marshes, or on rigid plants, usually schematised as cylinders, or groups of cylinders, in the laboratory tests. Indeed, as a matter of fact, the plants flexibility increases the complexity of both the vegetation resistance and the turbulent flow field characterization. Experiments concerning real flexible riverine plants, such as salix, are relatively rare and mainly refer to the drag evaluation for isolated plants (Oplatka 1998, Armanini *et al.* 2005) or on group of plants partially or totally submerged by the flow (Fathi-Maghadam and Kouwen 1997, Freeman *et al.* 2000, Jarvela 2002).

What comes out from the data in literature is the difficulty to extend the laboratory results to the field (as well as paucity of data concerning real flexible bushes). Indeed, living materials naturally have highly variable properties even within one single species; therefore, the extension of the laboratory results to field scale is subject to a high degree of uncertainty.

In this paper, measurements of flow velocity and hydrodynamic forces acting on real willows, fixed at the bottom of a laboratory channel, are presented. The aim of the present study is to improve knowledge of the interac-

tion between the flow and *real* flexible bushes, by analyzing the main features of the turbulent flow field in the frame of the Double-Averaging Method. Real bushes have been studied both in partial and in fully submerged conditions, at two different densities. The willows used in the experiments were *Salix Pentandra*, a kind of plants mostly used in restoration works along mountain streams. Drag force measurements have been made during all experiments, with detailed velocity measurements in two experiments with fully submerged plants. The force measurements allow to quantify the relative contribution to the flow resistance of the bed friction and plants drag, for the different plants densities. Moreover, for the two runs with plants fully submerged by the flow, the main features of the turbulent flow field is analysed in the frame of the Double-Averaging Method.

2. EXPERIMENTAL INSTALLATION AND MEASUREMENT TECHNIQUE

Experiments were conducted at the Experimental Centre for Hydraulic Models of Consorzio Venezia Nuova, in Voltabarozzo (Padova-Italy). The laboratory flume used in the tests is 150 m long, 2 m wide, 2 m deep and the bed and the walls are of concrete, except for a length of 12 m where one of the banks is made of glass. Through this glass windows, the plants bending has been registered by a digital photo camera and analyzed with image analysis techniques. The final section of the flume ends with a rectangular adjustable weir. The water flows through a closed circuit by means of a group of four pumps, up to a maximum discharge of $1.57 \text{ m}^3\text{s}^{-1}$. The channel is equipped with a double-track carriage which moves all along the flume. The tracks rest at the top of the concrete lateral walls. The carriage was used for the ADV velocity measurements, as explained in the following. The discharge is measured by means of an electromagnetic flow meter. The slope of the channel bed S_b is 0.18%. The downstream part of the flume was uniformly covered with salix bushes for a length of 45 m. The bushes had heights h_p ranging between 0.6 and 0.7 m. The bushes were placed at the bed in staggered configuration, with two different densities, named *sparse* and *dense* configuration in the following.

The geometrical characteristics of the plants positioning is reported in Figs. 1 and 2. In the *sparse* configuration, 1.1 bushes per square meter were placed, while the *dense* configuration is characterised by the presence of 2.5 bushes per square meter. Figure 1 also shows the position of the force transducer used for plant drag measurements.

Uniform inlet conditions were achieved using a perforated brick wall as a flow straightener to eliminate swirl. Just upstream the 45 m long reach with bushes, two rows of rigid wood sticks, 60 cm high, were placed, in order to minimize the bending of the first rows of salix invested by the flow.

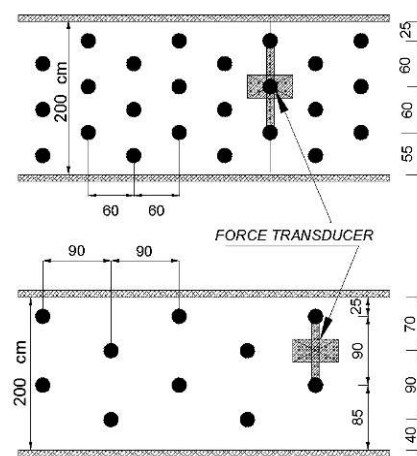


Fig. 1. Sketch of the plan view of the open channel with the two vegetation densities considered.

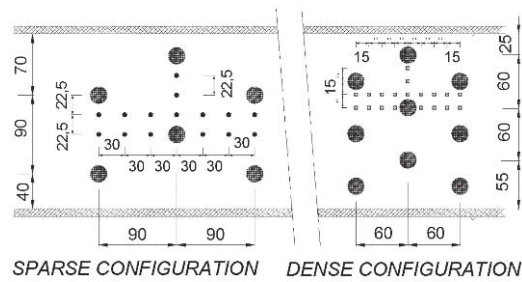


Fig. 2. Location of the verticals along which the ADV measurements were performed in the *dense* and *sparse* configurations.

For each discharge, the water level was measured all along the vegetated area and for every flow condition the formation of three reaches was observed. A transition profile was observed to develop in the first 10-15 m of the vegetated area, then a constant 20-27 m long flow depth region followed and finally a 5-10 m long, mild accelerating M3 profile developed. The length of each reach depends on the discharge values and on the deflected plants height-flow depth ratio h'_p/h . Therefore, for each discharge the water level at the downstream end of the flume was checked and adjusted by means of a sluice gate, in order to minimize the development of the M3 profile. The measuring zone was located 33 m downstream from the beginning of the vegetated area. Therefore, the flow conditions were not completely uniform all along the vegetated area, but it can reasonably be assumed that the uniform (or very close to uniform) flow conditions were achieved in the measuring zone, adjusting the sluice gate.

The selected plants were similar in height and also in shape, diameter of the main trunk, and principal and second order branches. Moreover, for the plants placed on the force transducer, the variation of the plants frontal area A with the distance z to the bed has been measured by means of image analysis techniques. Therefore, this area $A(z)$ represents the frontal area of the plant that would be submerged by a flow having depth z . It has to be remarked that, like in Armanini *et al.* (2005), only the trunk and the main branches have been considered for the frontal area definition. The normalized plant frontal area, that is, the ratio of $A(z)$ to $A_p = A(z = h_p)$, versus the relative height z/h_p , is plotted in Fig. 3 and compared with analogous measurements performed on willows sampled by plant cutting along the banks of some restored mountain streams and ranging between 1.5 and 3.5 m in height (Nobile 2007). It can be noted that the similitude between the laboratory and field willows is quite good.

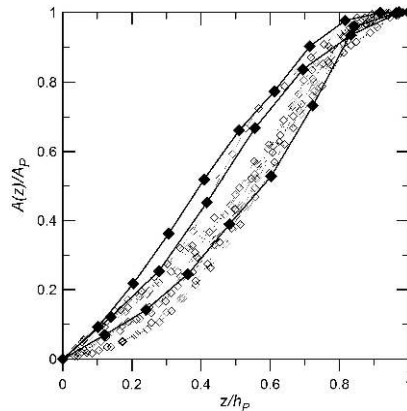


Fig. 3. Cumulative frontal area as a function of plants height. $-\diamond-$ real salix, 3 m tall, sampled along the banks of mountain streams in Trentino Region (Nobile 2007), $-\bullet-$ willows sampled in the laboratory channel.

A detailed analysis of the distribution of plant biomass with height and of plant density – defined as the number of plants per unit surface of the bed – has been performed for naturalized torrents. This analysis had the aim of reproducing the same vegetation density and plant biomass distribution at laboratory scale as sampled in the field (Nobile 2007). The results of the analysis will be discussed in the next section. The *sparse* vegetation density was chosen to be similar to that measured on the banks of some mountain streams. For each configuration a series of three runs was performed, for which the discharge was ranging between about 0.4 and $1.5 \text{ m}^3\text{s}^{-1}$, as reported in Table I, where also the ratio of deflected to undeflected plants height, h'_p/h_p , is reported. During each run, also the force exerted by the flow on a sample of bushes (three bushes for the *dense* configuration and

Table 1

Main features of the runs performed

<i>Sparse</i> configuration ($h_p = 0.70$ m)					<i>Dense</i> configuration ($h_p = 0.70$ m)				
Q [m ³ /s]	Q/A [m/s]	H [m]	h'_p [m]	h'_p/h	Q [m ³ /s]	Q/A [m/s]	H [m]	h'_p [m]	h'_p/h
0.556	0.524	0.53	0.69	0.99	0.422	0.402	0.525	0.69	0.99
0.867	0.647	0.67	0.67	0.96	0.711	0.531	0.67	0.67	0.96
1.533 *	0.824 *	0.93 *	0.45 *	0.65 *	1.311 *	0.728 *	0.90 *	0.58 *	0.83 *

* denotes the two runs for which also the velocity measurements were performed

two bushes for the *sparse*) was measured by means of the submerged force transducer. The aim of these measurements was twofold: the characterization of the plants drag evolution at increasing water discharge and thus for increasing bending, and also evaluation of the relative importance of plant drag to bed friction resistance. The force transducer was made up of two steel plates (30×15 cm) bound together by four aluminium foils at which eight strain gauges were stuck in order to form four Wheatstone bridges. The willows were made perfectly integral to the superior plate of the force transducer. The Wheatstone bridges allowed to measure the three components of strength and the bending moment produced by the force acting on the willows along the main flow direction. The circuits were calibrated and carefully waterproofed to allow the use of this device in completely submerged conditions. The force transducer was placed at the bottom of the flume within the measuring zone (see Fig. 1). For further details about the flume and the force transducer, the reader can refer to Armanini *et al.* (2005).

As reported in Table 1, in four of the six runs the plants were partially submerged by the flow. However, for two runs, one for the *dense* configuration and one for the *sparse*, the bending of the plants makes them to be fully submerged by the flow, and the height of the bowed vegetation h'_p was substantially less than the flow depth h . For these two runs, 15 velocity profiles for the *sparse* configuration and 19 velocity profiles for the *dense* configuration were measured. The exact locations of the measured profiles are shown in Fig. 2. A 3-D acoustic Doppler velocimeter (Sontek ADV) was used to measure the streamwise (u), transverse (v), and vertical (w) velocity components in 20 points along each of the verticals. Therefore, the data at 300 measurement points for the *sparse* configuration and at 380 measurement points for the *dense* configuration were collected. The analysis of these results in the DAM framework deserves particular attention and will be discussed in the next sections.

3. ANALYTICAL METHOD

Fully submerged willows act on the flow field as a low relative submergence roughness, characterized by spatial variability of the length scales of the same order of magnitude as the flow depth. In this case, a proper approach derives from a double-averaging procedure of the traditional Navier-Stokes equations, where the time-averaged Reynolds equations should be supplemented by spatial area averaging procedure in the plane parallel to the average bed (see, e.g., Wilson and Shaw 1977, Finnigan 2000, Nikora *et al.* 2001, Righetti and Armanini 2002). Consider the case of a 2D, steady, uniform, spatially averaged flow over a fully submerged vegetation, where x is the longitudinal direction, y is the transversal direction and z is the vertical coordinate. Applying the double-averaging technique (in time and in the space, that is in the x, y plane) to the Navier-Stokes equations, and neglecting viscous terms, two sets of differential equations are obtained:

$$gS_b - \left\langle \frac{\partial \tilde{p}}{\partial x} \right\rangle - \frac{1}{\phi} \frac{\partial \phi \langle u'w' \rangle}{\partial z} - \frac{1}{\phi} \frac{\partial \phi \langle \tilde{u}\tilde{w} \rangle}{\partial z} = 0, \quad (4)$$

$$g + \frac{1}{\rho} \frac{\partial \langle \bar{p} \rangle}{\partial z} + \frac{1}{\rho} \left\langle \frac{\partial \tilde{p}}{\partial z} \right\rangle + \frac{1}{\phi} \frac{\partial \phi \langle w'^2 \rangle}{\partial z} + \frac{1}{\phi} \frac{\partial \phi \langle \tilde{w}^2 \rangle}{\partial z} = 0, \quad (5)$$

where u , v , w are the instantaneous velocity components along the three coordinate axes x , y , z , respectively, p is the pressure. The prime (X') represents the point temporal fluctuation of X , the brackets ($\langle X \rangle$) and the linear overbar (\bar{X}) refer, respectively, to the spatial averaging along (x, y) plane and to time-averaging, S_b is the bed slope, g is the gravitational acceleration, Φ is the combined spatial and temporal porosity of vegetation (if present) in the (x, y) layer where the space averaging is performed, that is, the ratio between the area occupied by fluid and the total area of the averaging region in the x, y plane at level z : its values range between 0 and 1 in the vegetated region. Wavy overbars denote the difference between time averaged and double-averaged values ($\tilde{X} = \bar{X} - \langle \bar{X} \rangle$) and represents the form-induced stresses due to spatial averaging of the Navier-Stokes equations. Equation (4) represents the momentum balance in the longitudinal direction, while eq. (5) the balance in the vertical one. The mean velocity and turbulence measurements, till the third-order turbulence moments, carried out during the experiments were analyzed in the frame of the double-averaged momentum equations; the main features will be presented in the next section.

4. RESULTS AND DISCUSSION

Porosity of vegetation and plants bending

An analysis of the distribution of plants biomass with height and of plants density – defined as the number of plants per unit surface of the bed – allows

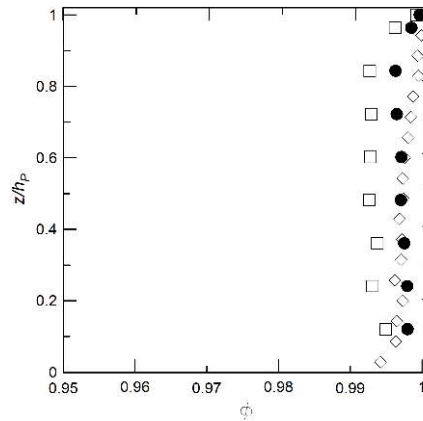


Fig. 4. Porosity of vegetation. \diamond measurements on real salix, sampled along the banks of mountain streams in Trentino Region (Italy); \bullet willows sampled in the laboratory channel, *sparse* configuration; \square willows sampled in the laboratory channel, *dense* configuration.

to assess the influence of vegetation porosity function on the balance equations. Figure 4 shows the porosity of vegetation $\Phi(z)$ evaluated on the banks of naturalized torrents (Nobile 2007) and as reproduced in the laboratory experiment for the *sparse* and *dense* configurations. In all the cases, porosity has been evaluated following the definition inferred in eqs. (4) and (5), that is, evaluating the area occupied by the trunks and main branches of plants over an averaging region containing a number of plants, so that $\Phi(z)$ is the porosity of the area used for spatial averaging and includes in its definition also the number of plants per unit surface of the averaging region. Moreover, it is worth of note that the vegetation porosity tends to assume values almost constant and very close to unity all along the plants height; therefore, despite the fact that the density of plants and so the porosity of vegetation could surely affect the flow resistance and the turbulence structure, the values of the parameter Φ can be assumed almost constant and substantially equal to one in eqs. (4) and (5), thus leading to an undoubted simplification in the balance equations.

Finally, in Fig. 5 the typical trunk bending of two laboratory plants belonging to the *dense* and *sparse* configurations are reported for different mean flow velocities. The data in Fig. 5 refers to the runs reported in Table 1, for which the corresponding bending has been estimated by image analysis techniques. The deflection of the plants is significant, inducing a significant reduction of the thickness of the vegetated layer during the flows. The plot shows that *sparse* configuration tends to exhibit lower flow depths and more deflected plants with respect to the *dense* configuration, for the same discharge. Therefore, it can be argued that for a given discharge the drag acted by each plant increases for decreasing plant density thus giving a more

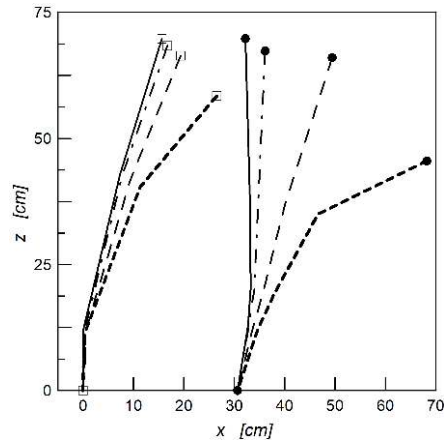


Fig. 5. Typical trunk bending for a willow belonging to the *dense* and *sparse* configurations, at increasing discharges. **Dense** configuration: —□— $Q/A = 0$, —□— $Q/A = 0.40 \text{ m}^3\text{s}^{-1}$, $h = 0.52 \text{ m}$, —□— $Q/A = 0.53 \text{ m}^3\text{s}^{-1}$, $h = 0.67 \text{ m}$, —□— $Q/A = 0.728 \text{ m}^3\text{s}^{-1}$, $h = 0.90 \text{ m}$. **Sparse** configuration: —●— $Q/A = 0$, —●— $Q/A = 0.52 \text{ m}^3\text{s}^{-1}$, $h = 0.53 \text{ m}$, —●— $Q/A = 0.65 \text{ m}^3\text{s}^{-1}$, $h = 0.67 \text{ m}$, —●— $Q/A = 0.824 \text{ m}^3\text{s}^{-1}$, $h = 0.93 \text{ m}$.

pronounced deflection, but the drag per unit bed surface decreases because the number of plants per unit bed surface is lower.

Drag resistance, mean motion and turbulent shear stress

In Fig. 6, the ratio between the plant drag per unit bed surface and the streamwise momentum component, $r_\tau = f\nu/\rho g S_b R_h$, is reported for the six runs as a function of the relative flow depth h/h_p . In Fig. 7 the measured drag coefficient as a function of the plant Reynolds number $Re_p = d_p U/\nu$ is reported, where d_p is the mean trunk diameter. It has to be pointed out that, in the case of willows fully submerged by the flow (i.e., one point corresponding to the higher Reynolds number for each vegetation density), the reference velocity for the evaluation of Re_p is assumed as evaluated by the application of the DAM to the velocity profiles measured in the vegetated layer. It can be observed that, for the considered vegetation densities, r_τ ranges between about 0.5 and 0.95, with the higher values for the higher vegetation density. Therefore, at least for the *sparse* configuration, the plants drag is of the same order of magnitude as the flow resistance due to shear at the bed and the latter is not negligible. This circumstance clashes with the assumption usually made in many simple models proposed in literature for the prediction and parameterization of a velocity profile in flows with submerged vegetation (see e.g., Kouwen *et al.* 1969, Christensen 1985, Raupach *et al.* 1996, Klopstra *et al.* 1997, Stephan and Gitknecht 2002, Righetti and Armanini 2002, Katul *et al.* 2002, Carollo *et al.* 2002), for which the bed shear can be considered negligible with respect to the plants drag.

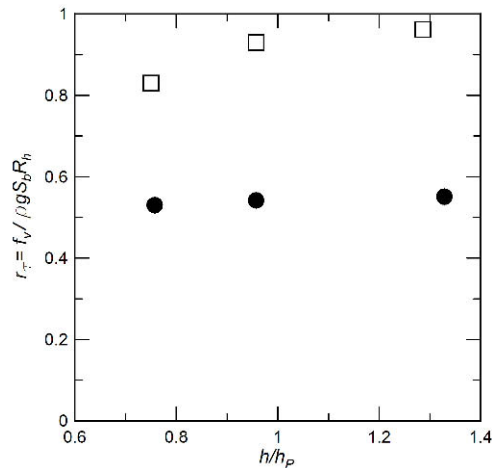


Fig. 6. Ratio between the plant drag per unit bed surface and streamwise momentum component for *dense* (□) and *sparse* (●) configurations.

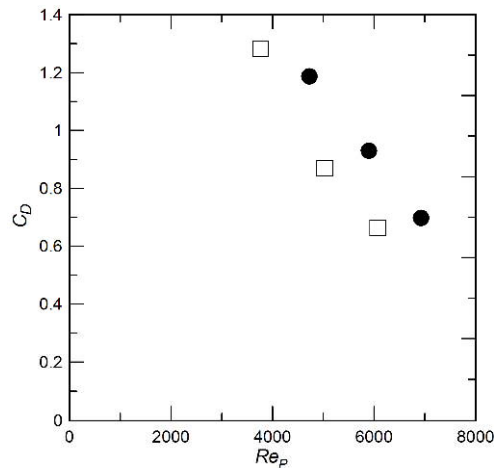


Fig. 7. Plant drag coefficient as a function of the plant Reynolds number $Re_p = d_p U / \nu$, for *dense* (□) and *sparse* (●) configurations.

Moreover, the drag coefficient reported in Fig. 7 shows a strong dependence on the flow velocity (that is the plants Reynolds number), due to the plant elasticity (see, Armanini *et. al.* 2005) and a slight dependence on the density of vegetation. The dependence on vegetation density can be ascribed also to the mutual interaction of plants, that is, sheltering effect and to the modifications in the turbulent structure of the flow field induced by surrounding plants. Indeed, this reduction on the drag coefficient has already been investigated for rigid cylinders and has been ascribed to the wake interaction between adjacent cylinders (see e.g., Lindner 1982, Nepf 1999). Cor-

rective procedures have been proposed for arrays of rigid cylinders, in order to correct the C_D values (Pasche and Rouvè 1985, Nepf 1999), the applicability of these procedures to the case of flexible plants with significant bending is an open question that goes beyond the aim of the present paper.

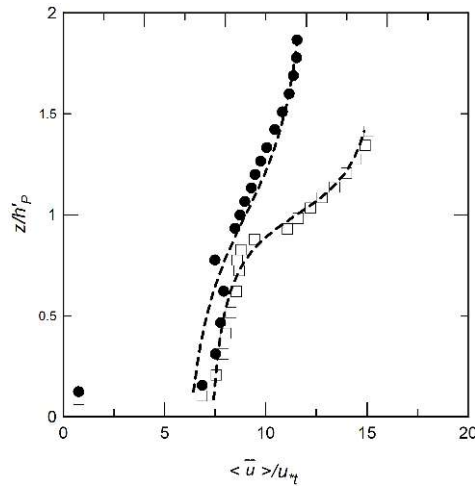


Fig. 8. Double-averaged dimensionless profiles. Free surfaces are located at $h/h'_p = 2.07$ and 1.55 for *dense* (\square) and *sparse* (\bullet) configurations. Dashed lines: inverse tangent formula, eq. (8) with suitable choice of the parameters a_i, b_i appearing in eq. (8).

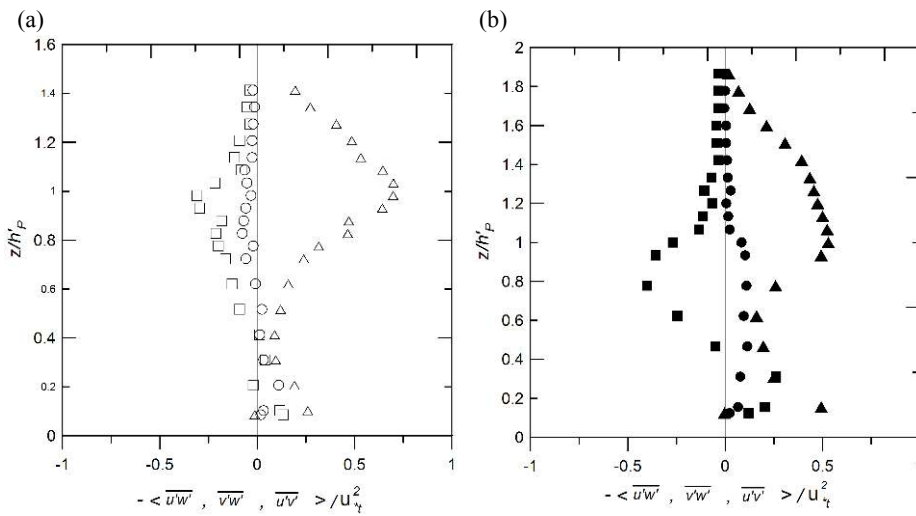


Fig. 9. Spatially averaged turbulent Reynolds profiles for *dense* and *sparse* configurations. (a) *Dense* configuration: $\Delta \langle \bar{u}w' \rangle / u_y^2$, $\square \langle \bar{u}v' \rangle / u_y^2$, $\circ \langle \bar{v}w' \rangle / u_y^2$.
 (b) *Sparse* configuration: $\blacktriangle \langle \bar{u}w' \rangle / u_y^2$, $\blacksquare \langle \bar{u}v' \rangle / u_y^2$, $\bullet \langle \bar{v}w' \rangle / u_y^2$.

The spatially averaged dimensionless mean velocity profiles for the *sparse* and *dense* configurations are reported in Fig. 8; Figure 9 shows the spatially averaged turbulence shear stresses profiles, while Fig. 10 shows the corresponding form-induced components. In the following the height of bended vegetation h'_p and the shear velocity at the top of vegetation, $u_{*t} = gS_b(h - h'_p)$, are chosen as length and velocity scales in order to make the plots dimensionless. The ratio between the flow depth and the vegetation height ranges between 1.55 and 2.07, corresponding to low relative submergence. Following the definition of Nikora *et al.* (2001), the studied flows can be classified as a flow type II. The double-averaged velocity distributions are qualitatively similar for both flows: with an almost constant sublayer close to the free surface, a quasi linear distribution in the central region of the non-vegetated layer, a strong gradient close to the vegetation crests, the gradient being stronger for *dense* configuration, and an almost uniform distribution in the vegetated layer with a decreasing mean value for increasing vegetation density. Finally, closer to the bed the velocity tends to decrease due to the no-slip condition at the bed.

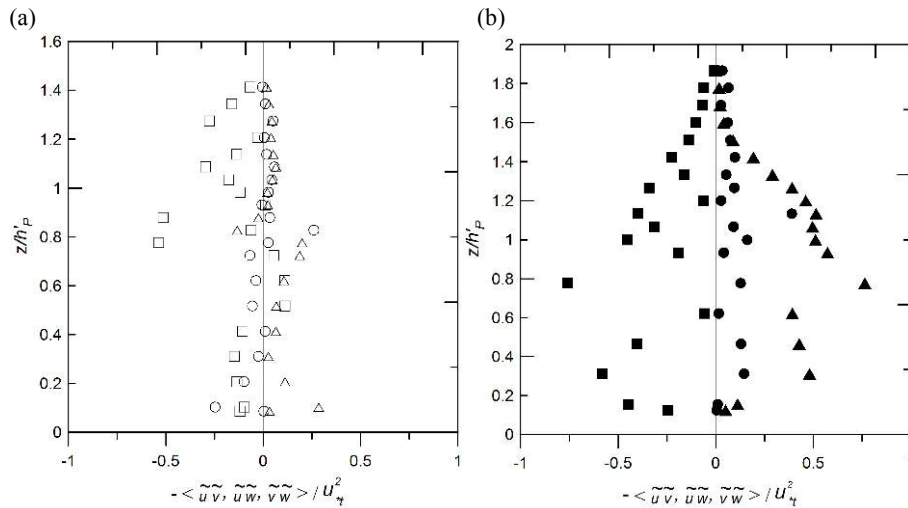


Fig. 10. Form-induced turbulent Reynolds profiles for *dense* and *sparse* configurations. (a) *Dense* configuration: $\Delta \langle \tilde{u}\tilde{w} \rangle / u_{*t}^2$, $\square \langle \tilde{u}\tilde{v} \rangle / u_{*t}^2$, $\circ \langle \tilde{v}\tilde{w} \rangle / u_{*t}^2$.
 (b) *Sparse* configuration: $\blacktriangle \langle \tilde{u}\tilde{w} \rangle / u_{*t}^2$, $\blacksquare \langle \tilde{u}\tilde{v} \rangle / u_{*t}^2$, $\bullet \langle \tilde{v}\tilde{w} \rangle / u_{*t}^2$.

The vegetated layer corresponds to the interfacial sublayer defined by Nikora *et al.* (2001) for low submergence flows. If $f_p = \langle \partial \tilde{p} / \partial x \rangle$ and $\phi(\langle u'w' \rangle + \langle \tilde{u}\tilde{w} \rangle)$ in eq. (4) can be considered approximately constant along z in the interfacial sublayer (wake turbulence dynamics), then a constant velocity profile can be derived (Nikora *et al.* 2004). Indeed, this seems to be

just our case, because for both the runs the porosity ϕ appears to be constant all along the vertical and also the spatially averaged and the form-induced shear stresses profiles (Figs. 9 and 10) seem to be almost constant.

In the upper part of the vegetated layer, the constant velocity profile seems not to be valid anymore, and an inflection point can be observed close to the top of vegetation. This kind of profile is quite similar to the hyperbolic tangent profile suggested on the basis of mixing layer analogy for canopy flows by Raupach *et al.* (1996), for low submergence open channel flow by Katul *et al.* (2002) and for flows in vegetated rivers by Ghisalberti and Nepf (2002, 2004). Therefore, the experimental data have been compared with the following theoretical profile suggested by Ghisalberti and Nepf (2002):

$$\langle \bar{u} \rangle = U_0 + \frac{\Delta U}{2} \tanh\left(\frac{z - z_i}{2\theta}\right), \quad (6)$$

where $U_0 = (U_1 + U_2)/2$, U_1 and U_2 are the lowest and highest values of flow velocity, in the present case the mean velocity values in the vegetated layer and in the almost constant sublayer at the free surface, respectively, $\Delta U = U_2 - U_1$, z_i is the distance from the bed where $\langle \bar{u} \rangle = U_0$ and θ is the momentum thickness defined as (Rogers and Moser 1994, Ghisalberti and Nepf 2002):

$$\theta = \int_{-\infty}^{+\infty} \left(\frac{1}{4} - \frac{\langle \bar{u} \rangle - U_0}{\Delta U} \right) dz. \quad (7)$$

The momentum thickness has been determined by solving the integral of eq. (7) as a summation.

Figure 11 shows the velocity profiles made dimensionless with the mixing layer parameters. Concerning the upper layer, it can be noted that the experimental data seem to be fairly well described by the hyperbolic tangent model. However, when the lower layer is considered, the theoretical profiles tend to deviate from the experimental data, especially for the data concerning the *sparse* configuration.

In the same way, an inverse tangent profile (eq. 8), as proposed by Carollo *et al.* (2002), with suitable choice of a_i and b_i coefficients appearing in eq. (8), seems to successfully approximate the experimental data in the upper non-vegetated layer. Again, in the vegetated layer the theoretical profiles slightly deviate from the data, the higher differences occurring for the *sparse* configuration and close to the bed (Fig. 8).

$$\frac{\langle \bar{u} \rangle}{u_{*t}} = b_0 + b_1 \arctan\left(\frac{(z/h'_p) - a_1}{a_2}\right). \quad (8)$$

In order to better understand the features of the dynamical interaction between the willows and the flow, the analysis of the turbulent second and

third order velocity moments have been performed; the main results are reported in the following.

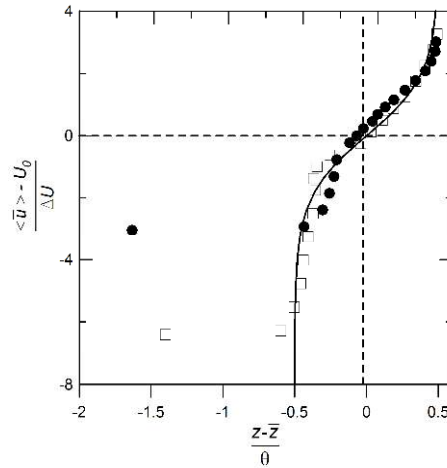


Fig. 11. Spatially averaged velocity dimensionless profiles with mixing layer variables. \square *Dense* configuration, \bullet *sparse* configuration; Solid line: hyperbolic tangent (eq. 6).

The spatially averaged Reynolds stress profile $-\langle u'w' \rangle / u_{*t}$ reported in Fig. 9 tends to assume an almost triangular shape in the upper non vegetated layer for both the configurations, with a maximum close to the top of the vegetated layer, and the maximum value for the *denser* vegetation. Further down into the vegetated layer, the profile tends to decrease up to a minimum in the middle of the layer and then slightly increases toward the bed.

The form-induced contributions to the Reynolds stresses are of the same order of magnitude as the spatially averaged ones for *sparse* configuration (Fig. 10a, b). For the *dense* configuration they are smaller than the spatially averaged ones but still non negligible. Moreover, while in the *sparse* configuration the form-induced shear stress component $-\langle \tilde{u}\tilde{w} \rangle$ is positive all along the flow depth, in the *dense* configuration it becomes negative (upward form-induced momentum flux) in a small region close to the top of the vegetation. This circumstance has already been reported in literature for a vegetated layer simulated by rigid cylinders (Poggi *et al.* 2004) and explained in (Pokrajac *et al.* 2008) using experiments with transverse bars. The authors ascribed this upward form-induced momentum flux to the vortex shedding phenomena that take place at the top of the roughness elements. Nevertheless, the changing sign for the *denser* configuration is apparently contradictory; the measurements of Poggi *et al.* (2004) and of Pokrajac *et al.* (2008) show an increase of the upward form-induced momentum flux for decreasing roughness element density, while the contrary is shown in Fig. 10.

The explanation lies in the effect of plants flexibility. The streamlined shape assumed by the plants in the *sparse* configuration tends to reduce the flow separation and vortex shedding in the region close to the top of vegetation. On the other hand, in the *dense* configuration willows are less bent by the flow; thus, the vortex shedding is enhanced at the crests, giving as a result a net upward form-induced momentum flux in that region. Whether the different hydrodynamic behaviour between the *sparse* and *dense* configurations could also explain the different attitude to obey the mixing layer analogy is not clear. An analysis of the turbulence intensities – which will be developed in the next section – could be helpful to this purpose.

Turbulent intensities

Figure 12 shows the spatially averaged turbulence intensity profiles and Fig. 13 the corresponding dispersive components. It can be noted that in the lower region of the vegetated layer the transversal component of the spatial averaged turbulent intensity is higher than the streamwise one while they are almost equal in the upper region of the vegetated layer. This effect can be ascribed to the Karman wake generation behind trunks and branches, which tends to enhance the transversal fluctuations, especially for the dense configuration. The comparison between the spatially averaged and the form-induced turbulent intensity profiles clearly shows that the contribution of the latter to the turbulent energy budget cannot be neglected. The values of the form-induced transversal and longitudinal components are higher in the vegetated layer than in the upper layer. The registered higher values in the

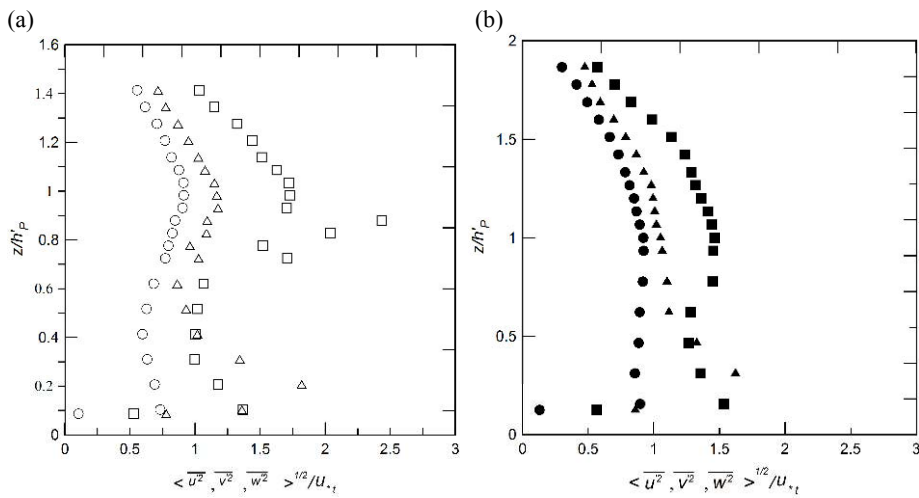


Fig. 12. Spatially averaged turbulence intensity components profiles for *dense* and *sparse* configurations.

- (a) *Dense* configuration: \square $\langle \overline{u^2} \rangle^{1/2} / u_{*t}$, Δ $\langle \overline{v^2} \rangle^{1/2} / u_{*t}$, \circ $\langle \overline{w^2} \rangle^{1/2} / u_{*t}$.
 (b) *Sparse* configuration: \blacksquare $\langle \overline{u^2} \rangle^{1/2} / u_{*t}$, \blacktriangle $\langle \overline{v^2} \rangle^{1/2} / u_{*t}$, \bullet $\langle \overline{w^2} \rangle^{1/2} / u_{*t}$.

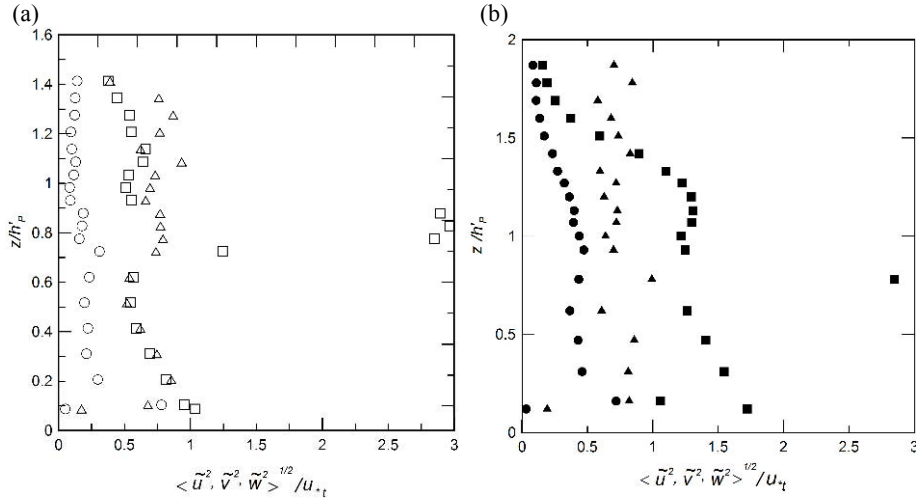


Fig. 13. Form-induced turbulent intensity components profiles for *dense* and *sparse* configurations.

(a) *Dense* configuration: $\square \langle \tilde{u}^2 \rangle^{1/2} / u_{*t}$, $\Delta \langle \tilde{v}^2 \rangle^{1/2} / u_{*t}$, $\circ \langle \tilde{w}^2 \rangle^{1/2} / u_{*t}$.

(b) *Sparse* configuration: $\blacksquare \langle \tilde{u}^2 \rangle^{1/2} / u_{*t}$, $\blacktriangle \langle \tilde{v}^2 \rangle^{1/2} / u_{*t}$, $\bullet \langle \tilde{w}^2 \rangle^{1/2} / u_{*t}$.

vegetated layer are a result of deviations induced by the presence of trunks and branches on the local (time averaged) motion with respect to the mean (spatially and temporally averaged) motion.

Higher order turbulence moments

Figure 14 shows the vertical profiles of the spatially averaged longitudinal third order moment $\overline{\langle u'^3 \rangle} / u_{*t}^3$ and the corresponding form-induced contribution. Figure 15 shows the vertical profile of the third order moment $\overline{\langle u'^2 w' \rangle} / u_{*t}^3$, which is related to the transport of turbulent energy (Raupach and Shaw 1982), and the corresponding form-induced contribution.

Both plots show that above the vegetation layer the form-induced contribution is significantly lower than the spatially averaged. On the contrary, in the vegetated layer the form-induced third order moments tend to increase due to the trunks and branches effect just mentioned in the previous paragraph.

The joint analysis of the spatially averaged third order moments deserve particular attention. Indeed, flow regions dominated by ejection turbulent events are associated with $\overline{\langle u'^3 \rangle} < 0$ together with $\overline{\langle u'^2 w' \rangle} > 0$ values. Flow regions dominated by sweep turbulent events are associated with $\overline{\langle u'^3 \rangle} > 0$ together with $\overline{\langle u'^2 w' \rangle} < 0$ values instead. Therefore, looking at Figs. 14 and 15 it can be deduced that:

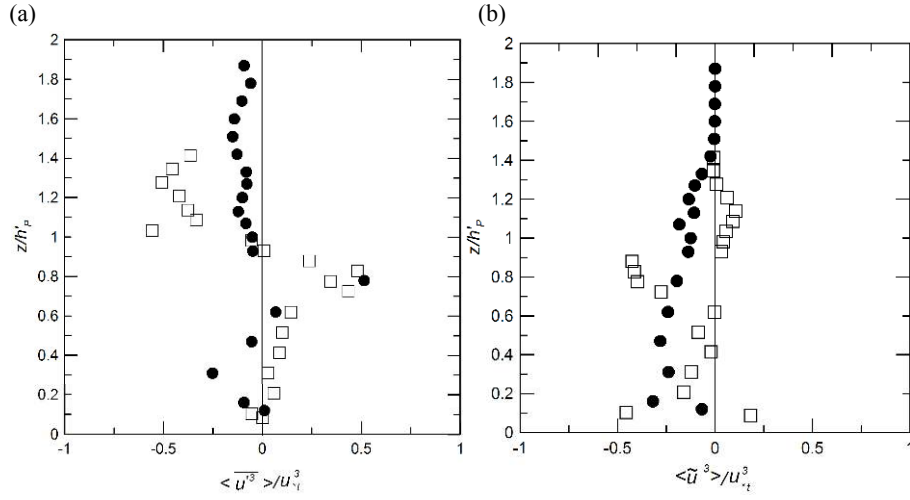


Fig. 14. Streamwise velocity third-order moment profiles for *dense* (\square) and *sparse* (\bullet) configurations: (a) spatially averaged contribution $\langle \bar{u}^3 \rangle / u_{s_t}^3$; (b) form-induced contribution $\langle \tilde{u}^3 \rangle / u_{s_t}^3$.

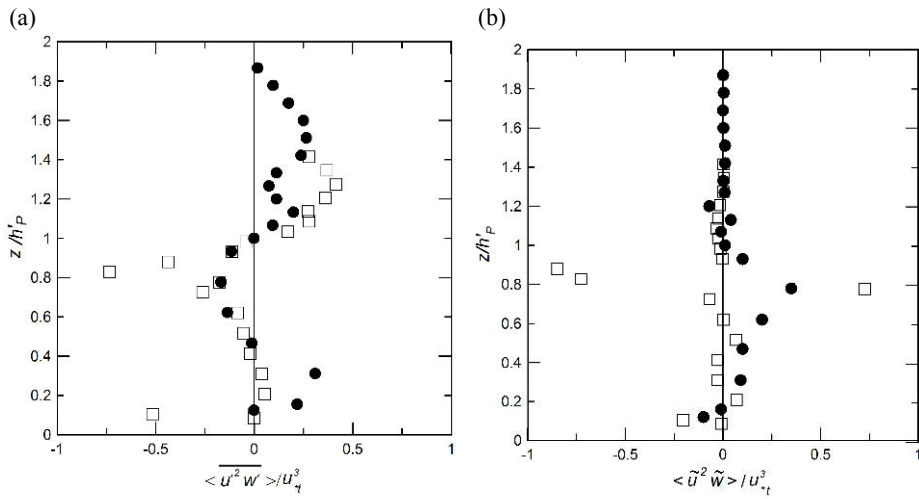


Fig. 15. Mixed third order velocity moment profiles for *dense* (\square) and *sparse* (\bullet) configurations: (a) spatially averaged contribution $\langle \bar{u}^2 \bar{w} \rangle / u_{s_t}^3$; (b) form-induced contribution $\langle \tilde{u}^2 \tilde{w} \rangle / u_{s_t}^3$.

– For the *dense* configuration, the upper non vegetated layer is mainly affected by ejection events. The most part of the vegetated layer is affected by sweep events. This is a typical characteristic of a mixing layer dynamics, as also reported by Ghisalberti and Nepf (2002): the roller type vortices that

develop are responsible for strong sweep events in the region just below the top of vegetation and for ejection events in the upper region.

– For the *sparse* configuration the scenario is more complicated. The upper region of the non vegetated layer is mainly dominated by ejection events down to about $z / h'_p \approx 1.3$. In the vegetated layer the sweeps-dominated region is limited only to the vegetation crests proximity, while in the lower half of the vegetated layer the ejection events tend to prevail.

These considerations concur to confirm the validity of the mixing layer analogy for the *dense* configuration while for the *sparse* configuration this similitude is not fully confirmed by the experiments.

In any case, the flow region above the plants is clearly dominated by ejection turbulent events, presumably due to the dynamics of coherent eddies which take place at the top of the plants (see, e.g., Raupach *et al.* 1996). The features of these coherent structures in the *dense* and *sparse* configurations can presumably differ, as suggested by the third order moments analysis, but after all they are responsible for the vertical transport of momentum and substances, such as nutrients, oxygen and sediment, between the vegetation layer and the outer flow (Ikeda and Kanazawa 1996, Ghisalberti and Nepf 2005). Focusing our attention on the transport by suspension of fine sediments, it is indeed well known that ejections are responsible for the maintaining into suspension of particles heavier than fluid, supplying the required net turbulent upward momentum flux (see, e.g., Sumer and Deigaard 1981, Kaftori *et al.* 1995, Lu and Willmart 1973, Righetti and Romano 2004). The impact of plants on the suspended sediment transport capacity can be therefore investigated generalizing the approach proposed by Lu and Willmart (1973) and analyzing the spatially averaged net upward turbulent momentum flux, evaluated by means of eq. (9):

$$\langle \text{Net Upward Turbulent Momentum Flux} \rangle = \left\langle (w^+)^2 T^+ - (w^-)^2 T^- \right\rangle, \quad (9)$$

where w^+ and w^- are the vertical upward and downward instantaneous vertical velocities in the measurement point, T^+ and T^- are the time intervals for which the velocity signal was upward and downward directed. The results are shown in Fig. 16. It can be noted that all along the upper non-vegetated layer the momentum flux is positive, but the maximum values tend to be reduced as the density of plants increases. On the other hand, the vegetated layer is mainly dominated by negative momentum flux which is associated to sweep events. The negative momentum flux is stronger for the dense configuration. Therefore, it can be argued that as the vegetation density increases, the suspended sediment transport capacity of the upper non vegetated flow tends to be damped. At the same time the vegetated layer tends to entrap the suspended sediments transported by the flow, and the entrapment efficiency increases for increasing vegetation density.

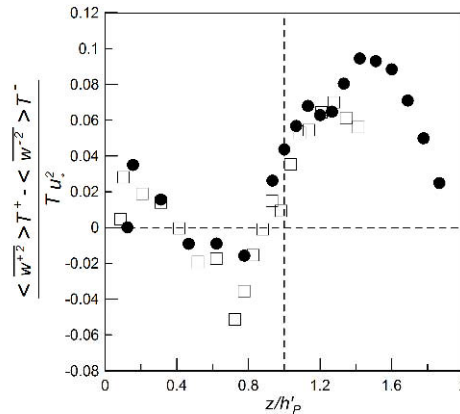


Fig. 16. Net Turbulent Upward Momentum Flux profiles for *dense* (□) and *sparse* (●) configurations.

5. CONCLUSIONS

An experimental study has been conducted using natural, flexible living willows to investigate the vegetation resistance for unsubmerged and submerged conditions at uniform flow condition. Real bushes, about 0.7 m high, have been used as vegetation, and the plants were uniformly distributed at the bed of a laboratory flume at two vegetation densities, named *sparse* and *dense* configurations. The *sparse* configuration reproduced the plants density and morphological features of natural willows (*salix pentandra*) as measured along the banks of some restored mountain streams in Trentino Region (Italy). The drag exerted to the flow by the bushes has been measured by means of a force transducer, and the plants bending has been evaluated by image analysis technique. For two experiments with fully submerged bowed plants, velocity measurements were performed and analyzed by means of the Double-Averaging Method. The following conclusion can be drawn by the analysis of the drag and velocity data:

- The plants drag per unit bed surface is of the same order of magnitude as the shear resistance at the bed, and the latter is not negligible. For a given discharge the drag per unit bed surface tends to increase with vegetation density, while the drag exerted by a single willow decreases.
- The form-induced Reynolds stresses and turbulence intensity contributions are not negligible with respect to the corresponding spatially averaged contributions.
- The average velocity, the vegetation density and the plant flexibility seem to be driving parameters for the development of a mixing layer analogy for the upper region of the flow field. This aspect has been confirmed both by the double-averaged velocity profiles and by the analysis of the double-

averaged third order turbulent moments $\langle \overline{u'^3} \rangle$, $\langle \overline{u'^2 w'} \rangle$. The mixing layer similitude seems to be applicable only to a limited portion of the vegetated layer, located toward the plants crests. The thickness of this region decreases as the plants density decreases.

□ For the lower portion of the vegetated layer, the analysis of the mean velocity profiles and of the spatial averaged and form-induced shear stresses, $[\langle \overline{u'w'} \rangle + \langle \overline{\tilde{u}\tilde{w}} \rangle]$ confirms the applicability of an almost constant velocity profile.

□ The analysis of the spatially averaged Net Upward Turbulent Momentum Flux shows that the upper non vegetated region is dominated by ejection events, while in the vegetated layer the sweep events tend to prevail. Once more the vegetation density influences also the turbulent momentum flux dynamics. As the vegetation density increases, the Net Upward Turbulent Momentum Flux and then the related suspended sediment transport capacity in the upper layer decreases. In the vegetated layer, the negative momentum flux increases for increasing vegetation density and so increases the attitude of the vegetated layer to act as a sink for the sediments in suspension transported by the flow.

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