

# Forced Convection Heat Transfer Enhancement Using a Self-Oscillating Impinging Planar Jet

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*Impinging jets are widely used in the local enhancement of heat removed from internal passages of gas turbine blades. Arrays of stationary jets are usually impinged on surfaces of internal cooling passages. The current practice is to benefit from the high heat transfer coefficients existing in the vicinity of the jet impingement region on a target wall. The present study shows that a self-oscillating impinging-jet configuration is extremely beneficial in enhancing the heat removal performance of a conventional (stationary) impinging jet. In addition to a highly elevated stagnation line Nusselt number, the area coverage of the impingement zone is significantly enhanced because of the inherent sweeping motion of the oscillating coolant jet. When an oscillating jet ( $Re=14,000$ ) is impinged on a plate normal to the jet axis ( $x/d=24$  hole to plate distance), a typical enhancement of  $Nu$  number on the stagnation line is about 70 percent. The present paper explains detailed fluid dynamics structure of the self-oscillating jet by using a triple decomposition technique on a crossed hot wire signal. The current heat transfer enhancement levels achieved suggest that it may be possible to implement the present self-oscillating-impinging-jet concept in future gas turbine cooling systems, on rotating disks, glass tempering/quenching, electronic equipment cooling, aircraft de-icing, combustors and heat exchangers. [DOI: 10.1115/1.1471521]*

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## Introduction

The advancement of high performance thermal systems has stimulated interest in methods to improve heat transfer rates, especially on the coolant side of hot surfaces. Considerable efforts have been made to increase heat transfer rates by implementing passive enhancement methods that require no direct consumption of external power. The present study deals with a heat transfer enhancement method using planar cooling jets in a self-sustained oscillatory flow mode. Utilizing an impinging jet in oscillatory mode has great cooling potential for hot section components. This method is based on the enhancement of turbulent mixing processes on the coolant side of hot surfaces by generating additional turbulent mixing via periodic oscillations in the impingement region. The term "oscillation" is used for the periodic flapping motion of the jet in a direction normal to orifice/nozzle axis. An obvious advantage of using self-sustained oscillations is that the power for oscillation is drawn from the existing flow field. The oscillations can easily be generated by a simple orifice/nozzle configuration that has no moving components even in relatively small coolant passages. The flapping motion of the turbulent planar jet spreads the effective thermal transport features over a relatively large impingement area on the target plate. A typical enhancement of  $Nu$  number near the stagnation line can be as high as 70 percent over stationary jet values.

Stationary impinging jets are widely used to provide high local mass/momentum/heat transfer in a variety of applications including gas turbine cooling, paper drying, helicopter rotor blade de-icing system, glass manufacturing, food processing etc. Impingement jet systems are popular because of the relative ease of con-

trolling the specific area needed to be cooled/heated/dried and high local heat transfer rates on the target surface. Reviews of earlier heat transfer studies on stationary impinging jets are given in Martin [1] and Goldstein et al. [2]. Gardon and Cobonpue [3] described the rate of heat transfer to an impinging jet as proportional to the difference between the target plate temperature and the adiabatic wall temperature, which varied from point to point on the impingement surface. Gardon and Akfirat [4] noted that the level of turbulence in the jet had a significant effect on the rate of heat transfer between the target plate and the stationary jet. Heat transfer from a flat surface to an oblique impinging jet was studied by Goldstein and Franchett [5] and Foss [6]. Most studies have involved jets of the same temperature as the ambient fluid. However, when the temperature of the jet is different from that of the ambient air, entrainment of ambient air into the jet complicates the heat transfer problem, Goldstein et al. [7] and Goldstein and Seol [8]. Stationary impinging jet based cooling arrangements are currently extremely popular in gas turbine cooling engineering and process engineering.

Fluidic oscillatory nozzles as simple and effective mixing devices have been used by Viets [9]. An extension of the fluidic nozzle concept to supersonic flow conditions is described by Raman et al. [10]. These are detailed fluid mechanics studies geared towards using the unsteady motion to actively control viscous flow near surfaces. The current study uses the fluidic oscillating nozzle as a heat transfer augmentation method for general heat transfer applications. The self-oscillating impinging jet improves thermal transport via periodic flapping motion of the jet. The area coverage of the impingement zone is also greatly enlarged by the flapping motion. Details of the analytical flow model for the self-oscillating-impinging jet used in this study is given by Herr and Camci [11]. The analytical model based on a corrected similarity technique is validated via hot wire measurements using a crossed wire. A decomposition approach was used to identify the contribution due to the deterministic jet oscillations and

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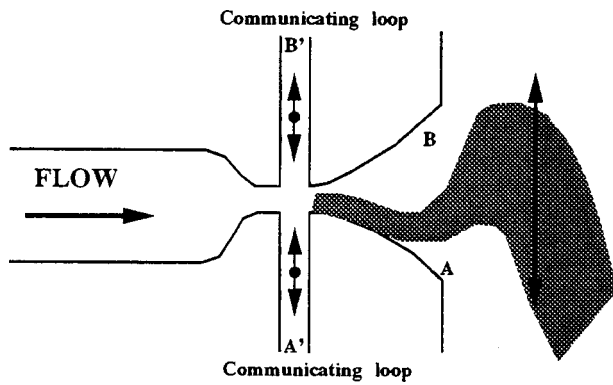


Fig. 1 A planar jet with self-sustained oscillations

stochastic turbulent fluctuations. The governing equations are derived for oscillating flows with fixed flapping frequency. At a given time, the oscillatory flow field can be approximately described by the solution of a stationary jet flow obtained from a nozzle/orifice attached to a relative reference frame that is oscillating at a prescribed frequency. The corrected similarity solution predicts the axial/lateral mean velocity measurements and oscillation kinetic energy data well, Herr [12].

Oscillation of the planar jet flow enhances transport process near the target plate because of the existence of additional inertia forces (oscillation stress) and additional thermal fluxes (oscillation heat flux). Experiments show that most of the kinetic energy in the oscillatory jet is contained in the deterministic unsteady part. The magnitude of the unsteady kinetic energy convection due to the flapping motion of the jet is about three times of that due to turbulence. The individual terms of the kinetic energy conservation equation sampled/evaluated from the measured crossed hot wire data show that the convection of kinetic energy in oscillatory flows is stronger than the convection due to turbulent fluctuations.

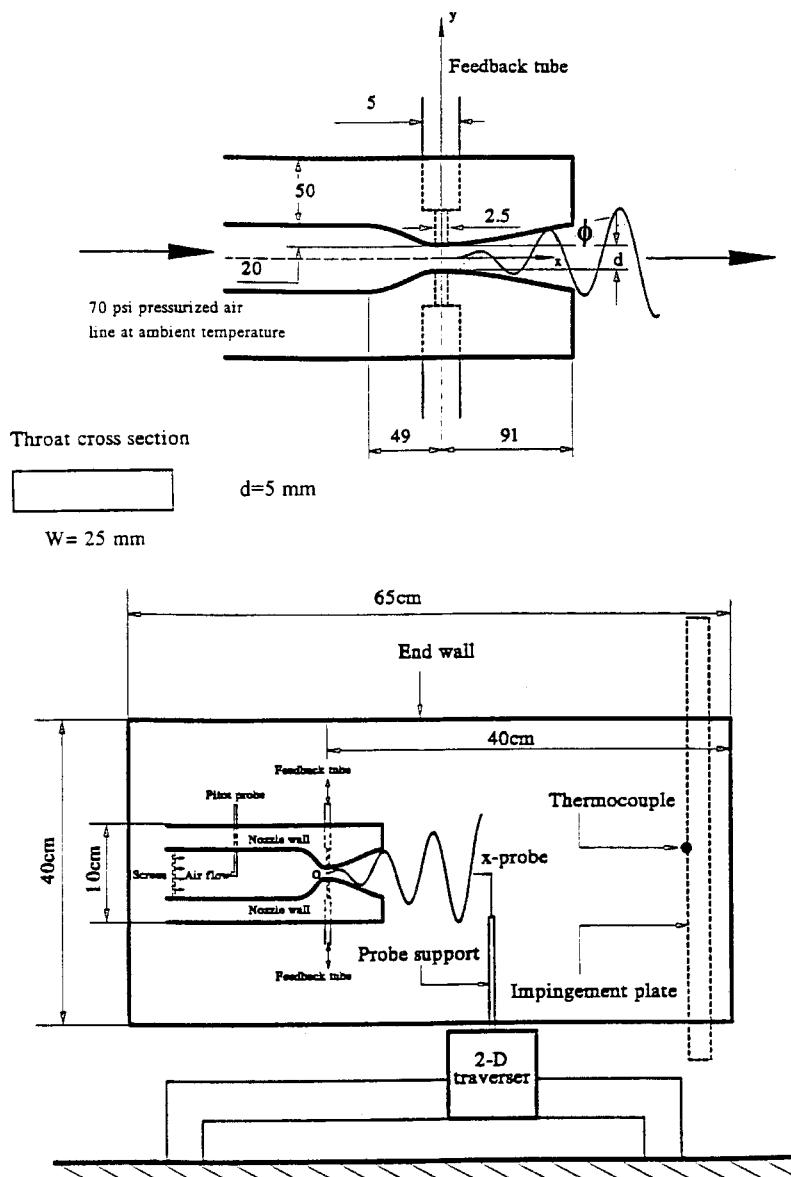


Fig. 2 Geometrical details of the nozzle and the experimental setup for the self-oscillating impinging jet experiments

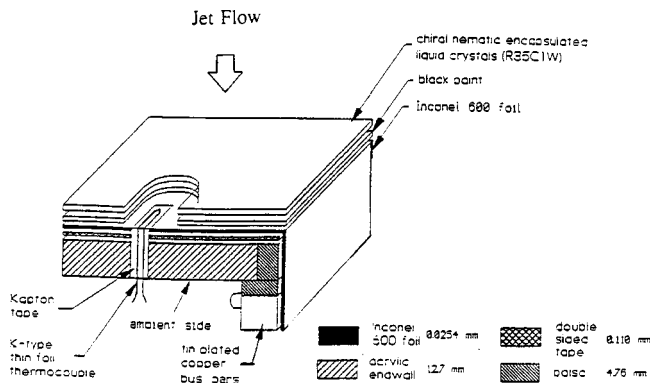


Fig. 3 Heat transfer surface construction

It is concluded that deterministic unsteadiness enhances transport process significantly, Herr and Camci [13]. Oscillation increases convection and diffusion significantly.

Enhancement of impingement heat transfer using a self-oscillating circular nozzle is described in Page et al. [14]. This acoustic excitation technique is based on the matching of the shear tone and organ pipe frequency of the pipe's finite length. The flow visualization studies revealed free transverse wave oscillations and enhanced vortex shedding. A collar extension approach results in enhancements in surface transport phenomena. An oscillation frequency of 5000 Hz from a self-oscillating nozzle resulted in measurable heat transfer coefficient enhancement. The nozzle-to-plate distance varied between diameters. Azevedo et al. [15] used a mechanical rotating ball valve to generate pulsations in axial velocity of a circular jet. The axial velocity in time varies between a peak and a minimum in each period created by valve rotations. This approach results in degradation in heat transfer coefficients on the impingement plate. The nozzle to plate distance varies between 2 to 10 diameters. The authors claim that the type of pulsations generated in this study adds energy primarily on the large scale, with little superimposed small scale turbulence. Liu and Sullivan [16] studied heat transfer from an acoustically excited circular impinging jet. The investigation is performed with a nozzle having an exit diameter of 12.7 mm and a nozzle to plate spacing less than two diameters. A loudspeaker attached to one side of a rectangular plenum chamber induces organ pipe resonance in the chamber producing plane-wave excitation at the jet exit. The perturbation velocity can be adjusted between 0.05 percent and 0.18 percent in a frequency range from 600 to 2500 Hz.

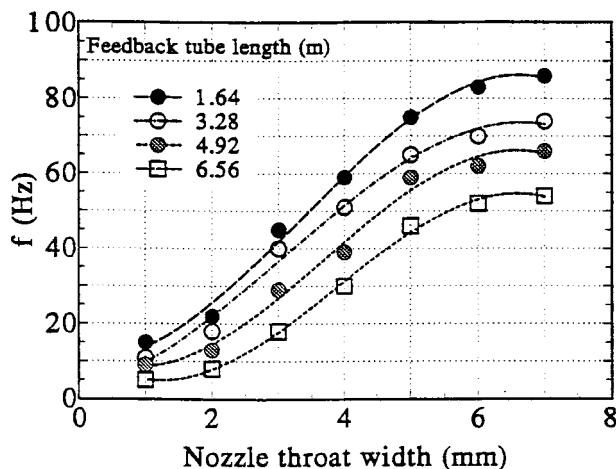


Fig. 4 Effect of nozzle throat width and communicating tube length on oscillation frequency of the jet, ( $\phi=20^\circ$ )

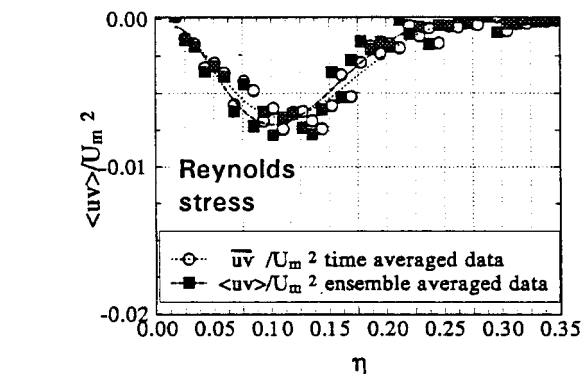
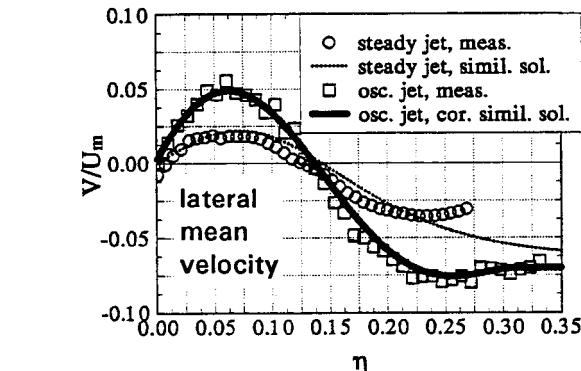
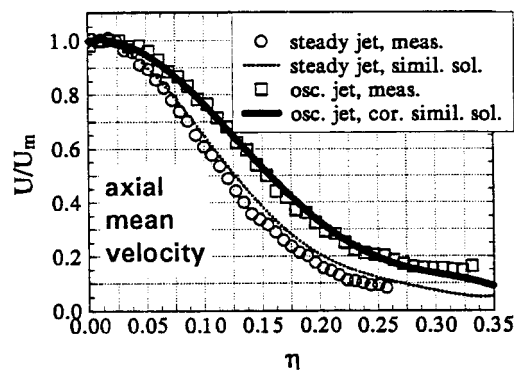


Fig. 5 Axial, lateral mean velocity profiles and Reynolds stresses in oscillating jet

A small nozzle-to-plate spacing (less than two diameters), the local heat transfer in the wall jet region can be significantly affected by exciting the impinging jet. The heat transfer near the stagnation point remains unchanged. The random vortical structures enhance the local heat transfer. The strong large-scale well-organized vortices formed after the stable pairing induce the unsteady separation of the wall boundary layer and hence, lead to the local heat transfer reduction. Mladin and Zumbrennen [17] use a rotating ball mechanism to generate a planar air jet with pulsations. The nozzle-to-plate separation distance varies between 0 to 10 nozzle widths with a pulse amplitude ranging from 0 to 50 percent of the mean flow velocity. Pulsation frequencies range from 0 to 80 Hz corresponding to Strouhal numbers below 0.106 based on nozzle width and jet discharge velocity. Heat transfer enhancements up to 12 percent near the nozzle mid-plane due to surface renewal effects and up to 80 percent at distances downstream due to increased turbulence levels are measured where both the pulse amplitude and Strouhal number are the highest. Pulses of small amplitude and low frequency induces a quasi-steady behavior at small separation distances, with no significant effect on the time averaged heat transfer near the nozzle mid plane. However, at

**Table 1** Estimates of the precision index and bias error values used for the uncertainty analysis

	Precision index	Bias Error	Uncertainty
$q_{gen}$	1.0 %	0.3 %	1.0 %
$q_{cond}$	5.5 %	5.9 %	8.1 %
$q_{rad}$	1.6 %	1.9 %	2.5 %
$(T_w - T_{char})$	2.7 %	2.6 %	3.4 %

larger separation distances, increased turbulence associated with the decay of the flow pulse enhanced the time averaged heat transfer away from the nozzle mid plane by up to 20 percent.

All of the studies reviewed in the previous paragraph focus on the nozzle to plate distances less than 10 jet diameters (or widths). Although Page et al.'s [14] study uses a self-sustained oscillation

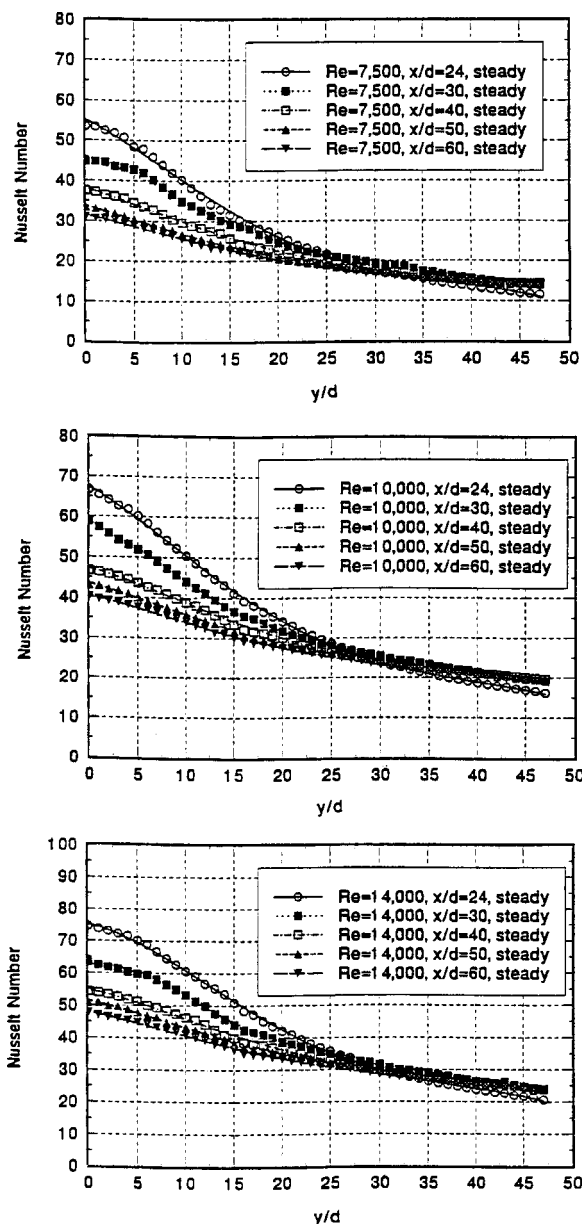
mode that does not require external power to maintain pulsations, Liu and Sullivan [16], Mladin and Zumbrennen [17] and Azevedo et al. [15] use loudspeaker and rotating ball valve arrangements that are powered externally. The current study is different from the studies summarized in this paragraph for two main reasons. First, the nozzle-to-impingement plate distance of the current study ( $24 < x/d < 60$ ) is geared towards engineering applications that can be found in food processing, drying, gas turbine cooling, glass tempering and aeronautical de-icing systems. The second main difference is in the character of unsteady jet motion imposed. The current study benefits from a self-oscillating impinging jet configuration that improves thermal transport via periodic flapping motion of the jet. The area coverage of the flapping fluid motion that is influencing the impingement plate zone is much greater than the coverage of all other studies reviewed in this section. Although the jets acoustically excited (or mechanically oscillated by a rotating ball valve) usually develop strong longitudinal oscillations they influence a limited impingement area on the target plate. The current self-oscillating planar jet in flapping mode has relatively larger area coverage on the target plate.

The present study focuses on the heat transfer aspects of self-oscillating impinging jets used in thermal system engineering. High-resolution heat transfer measurements are presented in a Re number range between 7500 and 14,000 ( $24 < x/d < 60$ ). This range is sufficiently wide for many current thermal engineering applications. It is possible to implement the present self-oscillating-impinging-jet concept in future gas turbine cooling systems, on rotating disks, in electronic equipment cooling, aircraft de-icing systems and heat exchanger systems.

## Experimental Setup and Procedures

**Oscillating Jet Assembly.** The operation principle of an oscillatory jet is based on the fact that a jet exiting into space between two sufficiently near walls is bi-stable, i.e., may attach to either wall. In addition, a small pressure gradient across the jet at the throat may cause the jet to detach one wall and attach to the opposite one. A simple fluidic nozzle representation of the oscillatory jet used in this study is shown in Fig. 1. The air flows through a nozzle shaped contraction, past two communication ports, into an expansion section. The output ports are attached to each other via a communication loop. Due to the proximity of the wall at the nozzle throat, the jet is bi-stable and must attach to one of the walls. Consider the jet to be attached to wall A. Due to large amount of entrainment into the jet, the pressure at the control port A' is relatively low while the pressure at the port B' is relatively high. Since the ports are attached to each other by a feedback loop, a compression wave travels from port B' to port A' tending to raise the pressure there and push the jet off the wall. Simultaneously, an expansion wave originates at port A' and travels to port B', tending to lower the pressure there and pull the jet onto the wall B. Thus, in a well designed nozzle, the jet will oscillate between walls A and B at a frequency determined by the length of the line connecting ports A' and B' and the line diameter. The frequency is also a weak function of the stagnation pressure entering the nozzle. The improved mixing characteristics of the oscillatory jet are seen in the increased spread angle and in the voids of stagnant gas that are entrained between the waves of the oscillation, Herr and Camci [11,13]

The flow and heat transfer tests for the current oscillating jet concept have been performed using precision Aluminum nozzles. The planar nozzle sections are precision-machined using fine, spline fitted coordinate points. The nozzle pieces are enclosed in 1.25 cm thick plexiglass sidewalls as shown in Fig. 2. A regulated laboratory compressed air supply is fed into the nozzle. The jet emerges from the nozzle having  $d = 5$  mm throat width (nozzle height is  $W = 25$  mm) at a typical mean velocity of 25 m/s ( $Re = 10^4$ ). To study the oscillation caused by the mechanism described above, the unsteadiness of other means must be minimum. With two communicating ports closed, the velocity at the throat



**Fig. 6** Stationary jet Nusselt number distributions on impingement plate, influence of Reynolds number and impingement plate location  $x/d$ , baseline data set

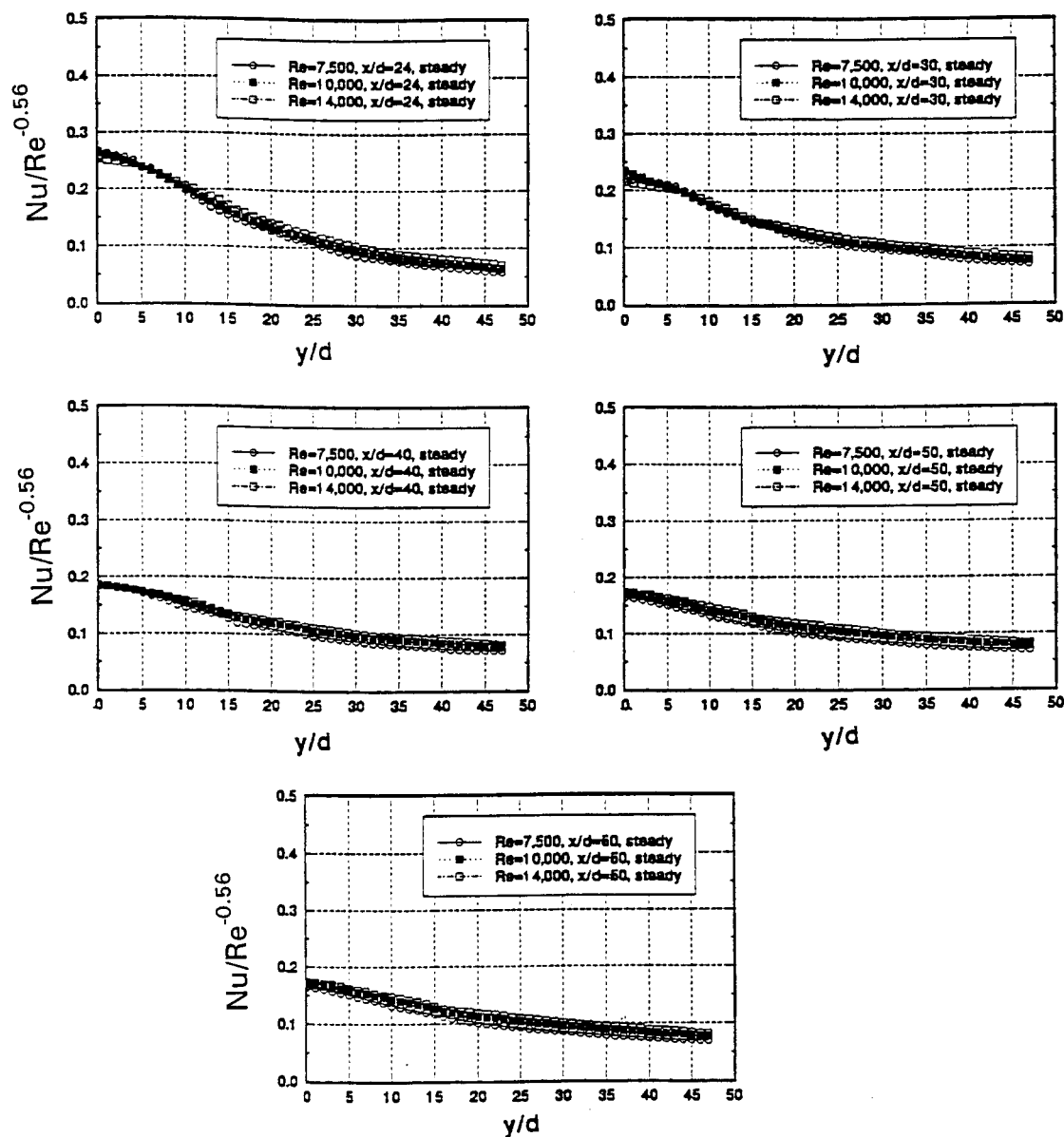


Fig. 7 Impingement plate heat transfer distribution from stationary jet,  $Nu/Re^{-0.56}$  distribution for all Reynolds numbers in function of  $y/d$

has been kept constant within 1 percent. For steady-flow situation, the jet is laminar, having a turbulence intensity of approximately 0.1 percent at the exit plane. The oscillation frequency of the jet is in the range of 10–100 Hz. This frequency is controlled by the stagnation pressure, nozzle throat width, and the length/diameter of the communicating loop. Other details of the experimental set-up can be found in Herr [12].

**Steady/Unsteady Flow Field Measurements.** Instantaneous velocity measurements have been taken with a TSI 1249A-T1.5 miniature 90 deg crossed hot wire probe and a TSI 1210-T1.5 general purpose single sensor probe using TSI series 1050 constant temperature anemometer. The single sensor probe has been used to measure the jet exit flow at the nozzle throat plane. The crossed wire probe has been calibrated in a velocity range from 0.5 to 40 m/s using a full velocity-yaw angle calibration scheme ( $\pm 45$  deg). The specific calibration scheme is the one defined by Lueptow et al. [18] and modified by Panchepakesan and Lumley [22]. This stretching coordinate transformation approach improves the measurement resolution at low velocities. Unsteady flow data

are sampled at a rate of 10 KHz for 5 seconds. A local averaging technique is used to detect the fronts of the periodical signals and 100 periods are ensemble averaged. Details of the current crossed hot-wire measurement method are presented in Herr [12].

**Heat Transfer Measurements on the Target Plate.** The target heat transfer surface, located downstream of the nozzle is shown in Fig. 2. The target plate to nozzle distance has been varied from 24 to 60 nozzle throat width “ $d$ ” during the experiments. The impingement plate is made up of 1.25 cm thick clear acrylic. The heat transfer surface is a composite system that includes a double-sided tape layer, Inconel heater foil, black backing paint, a liquid crystal thermo-indicator layer and flush mounted surface thermocouples. Details of the composite heat transfer surface are shown in Fig. 3. A low resistivity rectangular steel foil (Inconel-600) has been used as a constant heat flux surface on the target plate. The resistance of the steel foil has been accurately determined using a Keithley 580 micro-ohmmeter to an uncertainty of  $\pm 0.001 \Omega$ . A DC current reversal method removing unwanted offset voltages has been used. This technique lowers

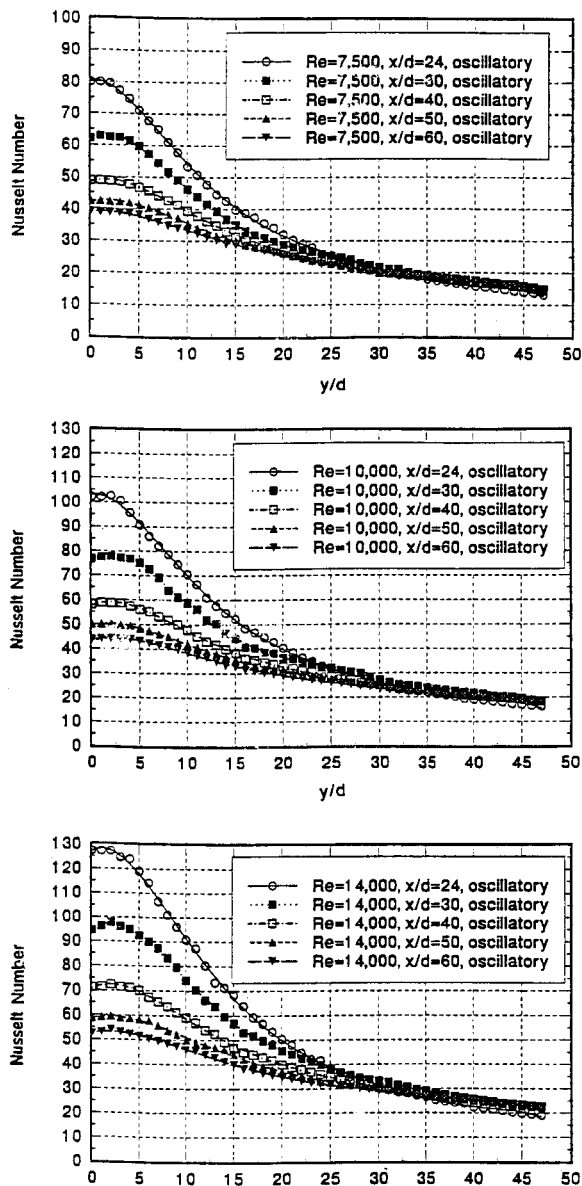


Fig. 8 Self-oscillating-jet Nusselt number distributions on impingement plate, influence of Reynolds number and impingement plate location  $x/d$

random noise in resistance measurements. Wall temperature measurement on the heat transfer has been made by using several flush-mounted thin-foil thermocouples of K type. Details of the heat transfer surface construction area given in Wiedner and Camci [19] and Herr [12]. Flush mounted thermocouples have been inlaid in several spanwise locations along the impingement plate and the junctions have been located at mid-point, Fig. 3, Herr [12]. The impingement plate shown in Fig. 2 has been traversed in the  $y$ -direction. The ambient side temperature of the heat transfer surface has also been monitored continuously for conduction heat loss calculations. The losses from the impingement side to the ambient side have been considered in the calculation of the convective heat transfer coefficient. A variable current DC power supply has been used to heat the constant heat flux surface. The jet and the ambient air in the laboratory are at the same temperature. The convective heat transfer coefficient is defined as,

$$h = q_{\text{conv}} / (T_w - T_{\text{char}}) \quad (1)$$

where  $q_{\text{conv}}$  is convective heat flux rate,  $T_w$  the wall temperature and  $T_{\text{char}}$  the characteristic temperature. The Nusselt number for this problem is based on the nozzle throat width  $d$ ,

$$\text{Nu} = hd/k. \quad (2)$$

In the present study, the reference temperature  $T_{\text{char}}$  is taken as the total temperature of the jet,  $T_o$ . The uniform wall heat flux value over the rectangular surface was fixed at a constant level for all tests. The heat flux generated by Joule heating of the steel foil is balanced by convective, conductive and radiative heat flux,

$$q_{\text{gen}} = q_{\text{conv}} + q_{\text{cond}} + q_{\text{rad}}. \quad (3)$$

The conduction heat loss from the impingement side to the back side of the plate is assumed to be one-dimensional,

$$q_{\text{cond}} = k_{\text{plate}} \partial T / \partial y = k_{\text{plate}} (T_w - T_{w,\text{back}}) / H \quad (4)$$

where  $T_w$  is the wall temperature on the heat flux surface,  $T_{w,\text{back}}$  is the temperature of the back side of the target plate, and  $H$  is the thickness of the target plate. During a typical run, conduction losses account for approximately 6–12 percent of the generated heat flux. Radiation heat transfer was estimated using the assumptions of black body emissivity and thermal equilibrium between the jet and sidewalls. The shape factors were approximated using an enclosure model and considering each surface as black. The radiation heat flux is approximately 3–9 percent of the local generated heat flux.

**Experimental Uncertainties.** The final experiment uncertainty estimates of the current study are described using the method given by Kline and McClintock [20]. The uncertainty levels presented are analogous to 95 percent coverage or 20:1 odds. A crossed hot wire sensor has been used to measure  $x$  and  $y$ -components of the instantaneous velocity field in which temperature variation in the flow field is not significant. The Inconel strip heater used in constant heat flux measurements on the target plate has been disconnected from the DC power supply during hot wire measurements. This approach assures an isothermal flow field that is essential for high quality hot-wire measurements. The instantaneous field has then been de-composed into mean, deterministic and random parts for further analysis. The sources of error in hot wire measurement include flow field temperature variations, thermal inertia of the sensors, probe misalignment, sampling time errors, constant temperature anemometer noise. The cut-off frequency used has been selected as a 4 kHz with an overheat ratio of 1.6. It has been noted from a power spectral density calculation that, energy content of the flow is extremely small for the frequencies higher than 600 Hz. A 0.5 deg alignment error is assumed in the uncertainty analysis. This error contributes 0.1 percent to the total error. The predicted error in instantaneous velocity measurements is about 1 percent. A detailed account of the error analysis on hot wire measurements is presented in Herr [12].

The heat transfer measurements on the constant heat flux surface are influenced from the errors made in the determination of wall temperature  $T_w$ , total temperature  $T_o$  of the jet (ambient), heat flux measurement, throat width measurement and absolute error in thermal conductivity  $k$  of air. Our experiments by varying the black paint and liquid crystal layer thickness showed that the wall temperature error  $\delta(T_w - T_{\text{char}}) / (T_w - T_{\text{char}})$  is less than 3.4 percent. This error is already included in the final uncertainty value of Nusselt number 6.3 percent approximately. Estimates of the precision index and bias error values for the main measured quantities in the uncertainty analysis are given in Table 1. Further details of the uncertainty analysis are presented in Herr [12].

## Experimental Results and Discussion

**Jet Half Width and Effective Mixing.** A useful indication of effective mixing in a jet is the variation of the jet half width in axial direction. The definition of the half width at any streamwise

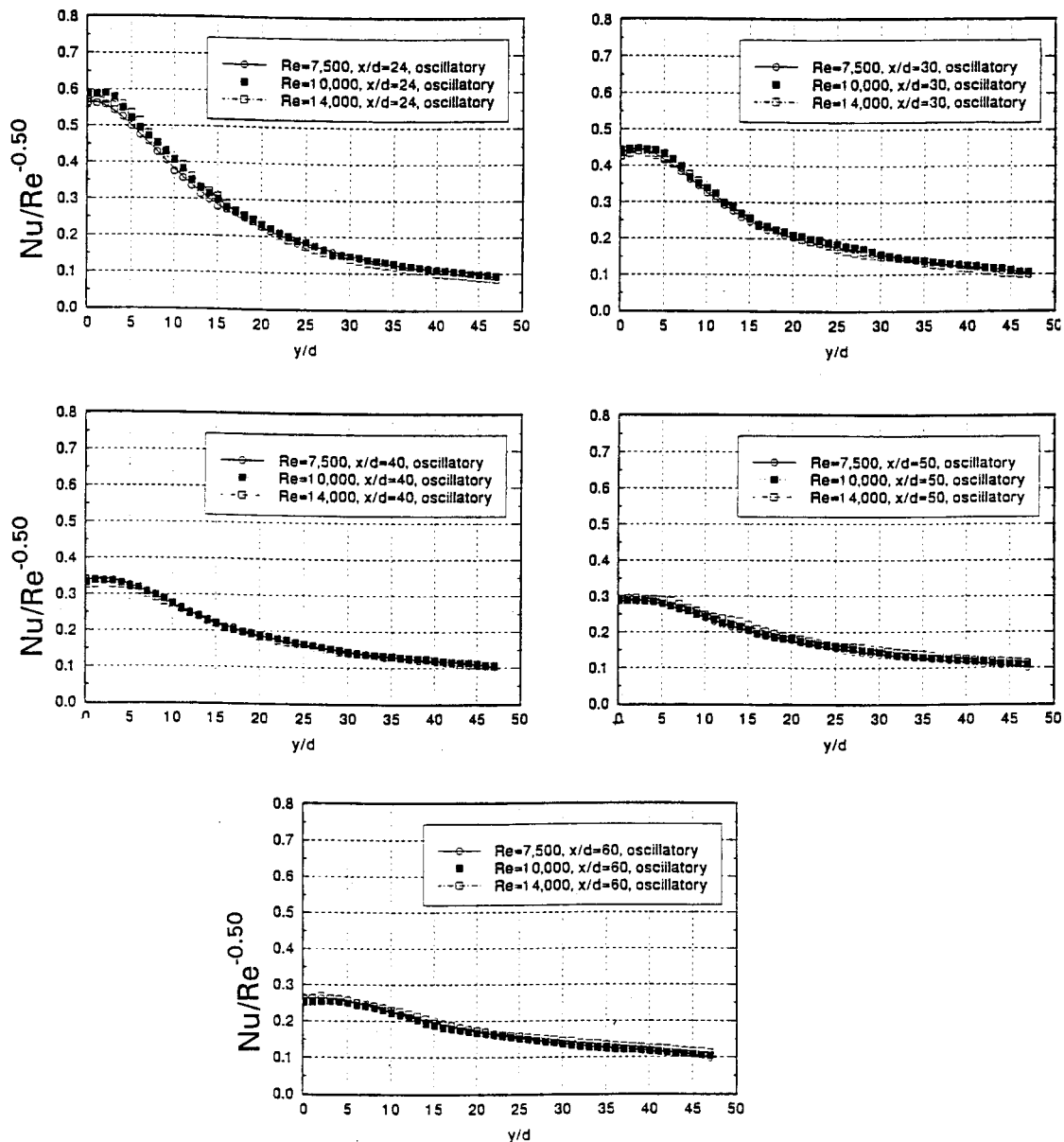


Fig. 9 Impingement plate heat transfer distribution from oscillating jet,  $Nu/Re^{-0.50}$  distribution for all Reynolds numbers in function of  $y/d$

position is the distance between the jet centerline and the point where the local velocity is equal to half of the centerline velocity. Current experiments clearly show that the introduction of the oscillatory motion dramatically increases the jet half width. At  $x/d = 60$ , ( $Re = 10,000$ ), the half width from a stationary jet is about 8 times the nozzle throat width  $d$ . However, when the same jet is put into a self-oscillatory motion, the half width increases to 13 times the throat width, Herr [12]. This observation is closely related to the fact that a cooling jet in a flapping motion covers a much wider impingement area when compared to a stationary jet. Enhanced turbulent transport with effective mixing is spread over a much larger impingement zone on the target surface via the use of an oscillating jet.

**Frequency of Self-Sustained Jet Oscillations.** The oscillation frequency of the jet is determined by taking the fast Fourier transform of instantaneous hot wire data from the exit of the nozzle at,  $\eta = y/x = 0$ . Figure 4 shows the variation of oscillation frequency with respect to nozzle throat width  $d$  for constant stagnation pressure in the plenum chamber, ( $\phi = 20$  deg). Experiments

performed at  $x = 180$  mm also show that larger throat widths at a fixed communication line length induce higher oscillation frequencies. Similar trends have been observed at  $y/x = 0$ ,  $y/x = 0.05$ , and  $y/x = 0.3$ . Increasing  $d$  from 1 mm to 5 mm provides a five-fold frequency increase from 15 to 75 Hz. Increasing the length of the communication line slightly reduces the oscillation frequency. For the case of increased throat width, the amount of fluid entraining into the throat area increases. Thus the ambient air enters the feedback tube sooner than the case which has a smaller throat width. If the width is too large, the pressure gradient across the jet is not strong enough to move the jet and the jet would not bend to either wall, so that no oscillation will occur. Increasing plenum chamber stagnation pressure and shorter communication lines has a tendency to increase the frequency of oscillations, Herr [12].

**Axial and Lateral Mean Velocity Components.** Mean velocity components have been measured by time averaging the instantaneous hot wire signal at various  $\eta = y/x$  locations at  $x/d = 60$  for  $f = 75$  Hz. It should be noted that the data discussed from

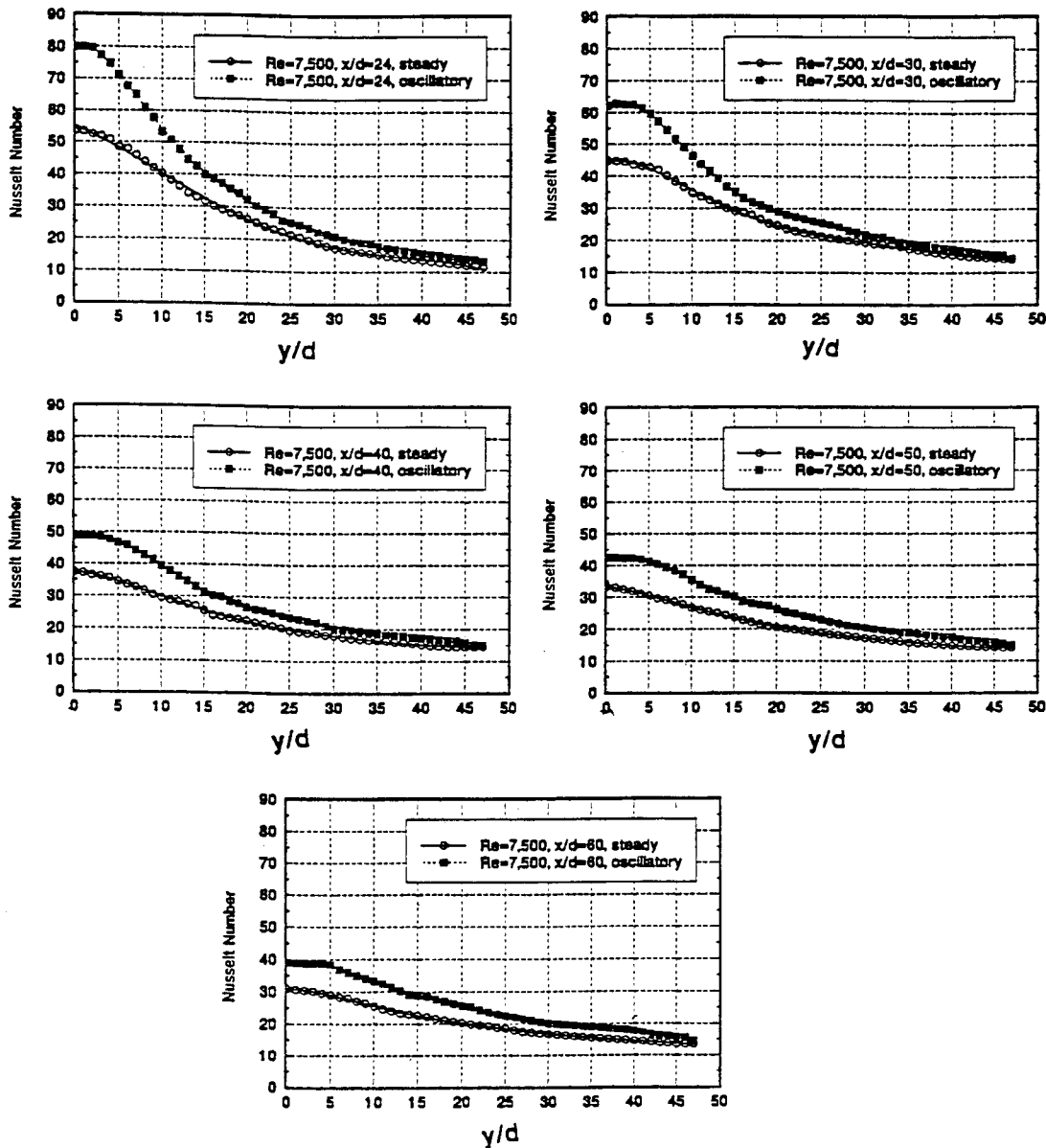


Fig. 10 Comparison of oscillating jet and stationary jet heat transfer distributions on the impingement plate,  $Re=7500$

this point on include only the results for the nozzle that has a throat width of  $d=5$  mm and a nozzle half angle of  $\phi=20$  deg. The measured axial mean velocity ( $U/U_m$ ) distribution of the oscillating jet is relatively fuller than that of the stationary jet as shown in Fig. 5. This is expected because oscillation increases the spreading rate of the jet. The lateral mean velocity ( $V/U_m$ ) distribution of the oscillating jet is enhanced when  $0 < y/x < 0.14$ , ( $\eta = x/d = 60$ ). The flapping motion of the jet issued from the fluidic nozzle provides increased transport of mean kinetic energy in the lateral direction. This feature is directly related to the observed enhancements in jet half width leading to increased transport in the lateral direction, Herr [12]. The thick solid lines in Fig. 5 represents the predictions from “the corrected similarity method” developed by Herr and Camci [11]. “The corrected similarity method” is an analytical extension of the stationary turbulent jet solution with special emphasis paid to deterministic oscillations of the jet. Figure 5 shows that “the corrected similarity method” based prediction of mean velocities is in very good agreement with hot wire measurements where  $0 < y/x < 0.275$ . The stationary

turbulent jet solution for  $U$  and  $V$  (based on a conventional similarity method) is represented by tiny dots in Fig. 5.

**Additional Stresses and Heat Fluxes Due to Deterministic Oscillations.** The instantaneous velocity and temperature can be decomposed into three parts, Hussain and Reynolds [21], a time averaged quantity ( $U_i, \Theta$ ), ensemble averaged deterministic oscillations ( $u_i^*, \theta^*$ ) and stochastic turbulent fluctuations ( $u_i, \theta$ ),

$$\begin{aligned} \tilde{u}_i &= U_i + u_i^* + u_i \\ \tilde{\theta} &= \Theta + \theta^* + \theta \end{aligned} \quad (5)$$

In an oscillatory flow, the continuity equation can be decomposed as,

$$\frac{\partial U_i}{\partial x_i} + \frac{\partial u_i}{\partial x_i} + \frac{\partial u_i}{\partial x_i} = 0 \quad (6)$$

Finally, the time averaged momentum equation for a flow with self-sustained oscillations is as follows,



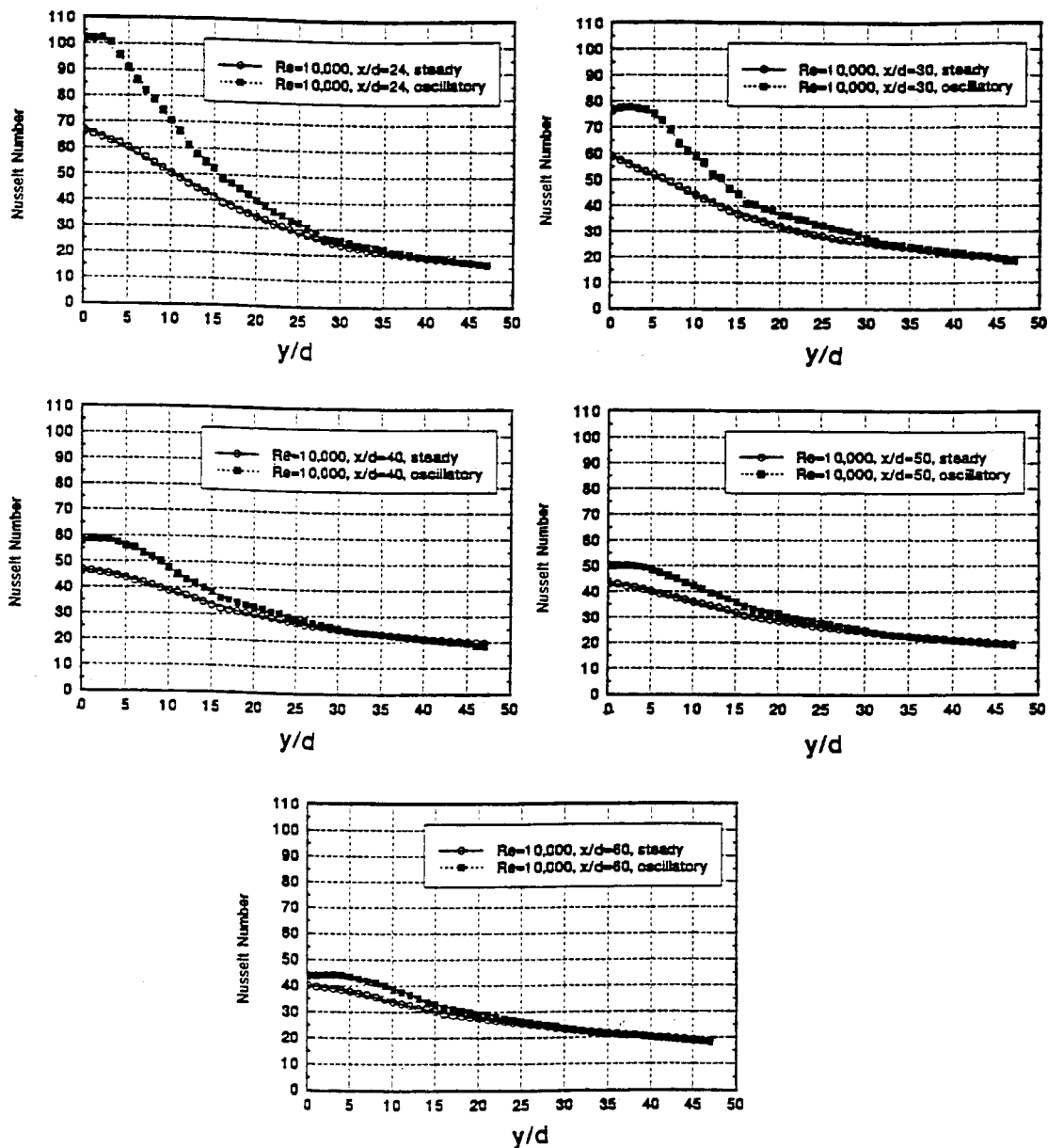


Fig. 11 Comparison of oscillating jet and stationary jet heat transfer distributions on the impingement plate,  $Re=10,000$

$$\frac{\partial(U_i U_j)}{\partial x_j} = -\frac{1}{\rho} \frac{\partial P}{\partial x_i} + \frac{\partial}{\partial x_j} \left( \nu \frac{\partial U_i}{\partial x_j} - \overline{u_i^* u_j^*} + \overline{u_i u_j} \right) \quad (7)$$

$$\frac{\partial(U_j \theta)}{\partial x_j} = \frac{\partial}{\partial x_j} \left( \gamma \frac{\partial \theta}{\partial x_j} - \overline{u_j^* \theta^*} + \overline{u_j \theta} \right) \quad (8)$$

Equation 7 clearly shows that in addition to shear stress due to molecular viscosity, there are two additional apparent stresses termed as  $-\overline{u_i^* u_j^*}$  and  $-\overline{u_i u_j}$ . The apparent stress terms are generated by the nonlinear inertia term (convective acceleration) that is the left hand side of the instantaneous momentum equation. Additional inertia force due to random oscillations of the turbulent field is usually termed as Reynolds stress  $-\overline{u_i u_j}$ . The term  $-\overline{u_j^* u_j^*}$  is due to the inertia force originating from deterministic variations of the oscillating jet flow. This extra stress term is named as oscillation stress  $-\overline{u_j^* u_j^*}$  throughout this study. Similar to the development leading to Eq. 7, one can obtain a time aver-

aged thermal energy equation, Eq. 8, Herr and Camci [13]. The three terms on the right hand side of the energy equation are conduction heat flux  $\gamma(\partial\theta/\partial x_j)$  due to molecular conductivity, oscillation heat flux  $-\overline{u_j^* \theta^*}$  due to flapping motion of the jet and turbulent heat flux  $-\overline{u_j \theta}$  due to random fluctuations of the flow field. Figure 5 also shows a comparison of the time averaged and ensemble averaged turbulent stress  $\langle uv \rangle / U_m^2$  in function of  $\eta$  at  $x/d=60$ ,  $f=80$  Hz. Reynolds stresses due to turbulent fluctuations reach a peak around  $\eta=0.10$  and diminish near the centerline and near the outer edge of the jet where  $\eta>0.30$ .

**Heat Transfer From Stationary Impinging Jet.** Stationary jet related Nusselt number distributions on the impingement plate are given in Fig. 6 as baseline distributions. Experiments at three different Reynolds numbers (7500, 10,000, and 14,000) are presented for various impingement plates to nozzle distance values ( $x/d=24,30,40,50$ , and 60). A strong influence of plate to nozzle distance on stagnation point heat transfer is shown in Fig. 6. Nus

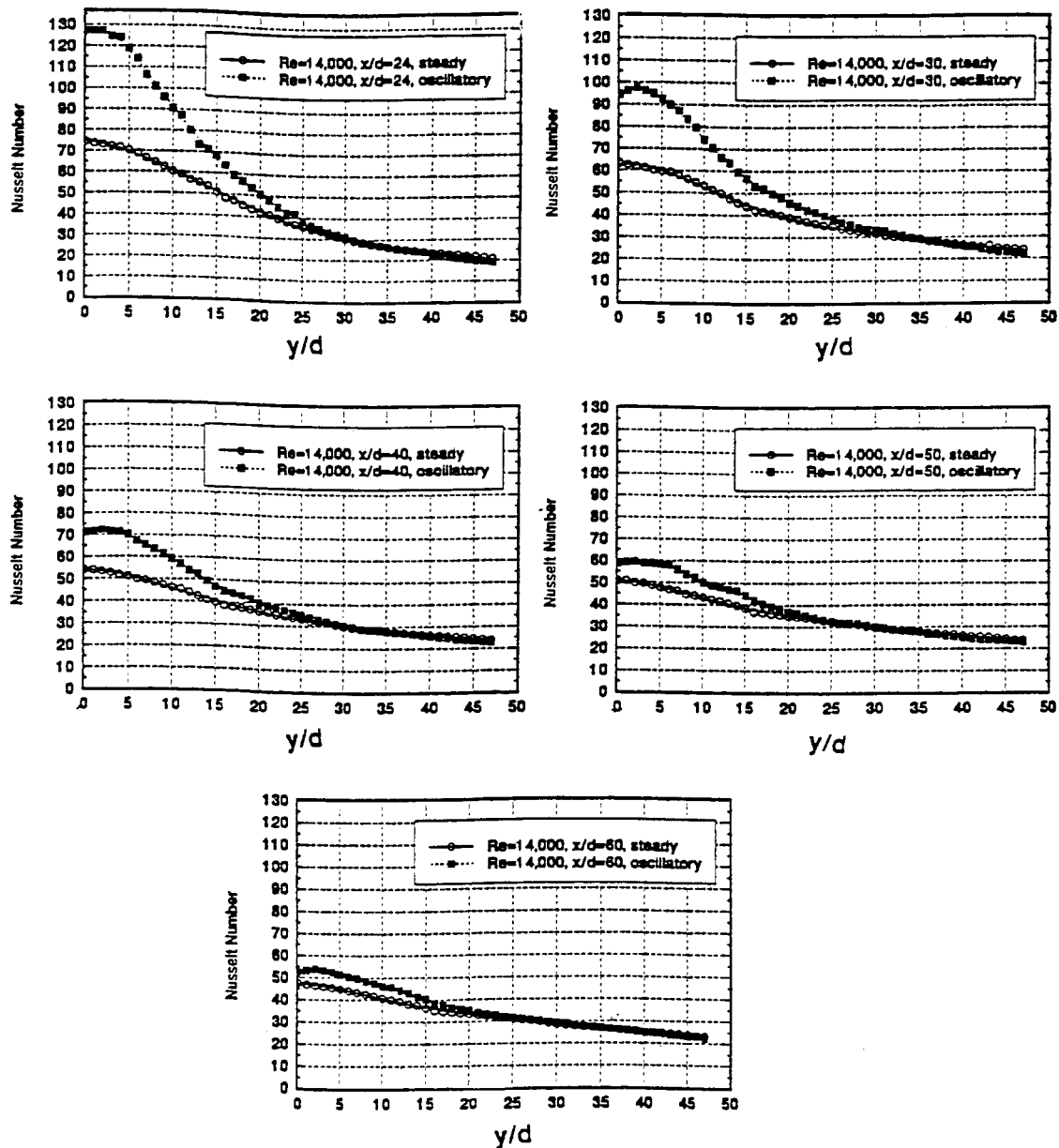


Fig. 12 Comparison of oscillating jet and stationary jet heat transfer distributions on the impingement plate,  $Re=14,000$

selt number at the stagnation point increases from 30 to 55 when  $x/d$  is reduced from 60 to 24 at  $Re=7500$ . This trend repeats itself for  $Re=10,000$  and  $14,000$ . The Nusselt number depends on  $x/d$  in the near impingement region  $0 < y/d < 20$ , while in the wall jet region, the heat transfer is almost independent of  $x/d$  for all  $Re$  numbers. At the minimum  $x/d$  value of 24, the jet is predicted to be fully turbulent for all Reynolds numbers. In other words, potential core does not exist at this  $x/d$ . The heat transfer information given in Fig. 6 provide the reference stationary heat transfer distributions when oscillating jet results are compared to stationary jet. In general the local heat transfer on the impingement plate is a function of  $x/d$ ,  $y/d$ ,  $Re$ , and  $Pr$ . This functional dependency can be represented by a power law in the form,

$$Nu = \sim Re^n Pr^m \quad (9)$$

The exact form of Eq. 9 varies from study to study depending on the heat transfer configuration, flow and thermal boundary conditions, local turbulence production and unsteadiness in the flow. The molecular Prandtl number can be taken constant for air under

present conditions. For the stationary impinging jet, the most representative value for  $n$  has been found to be  $n=0.56$ , when  $m$  is unity. At a fixed impingement plate to nozzle distance, all three heat transfer distributions for different  $Re$  values will collapse into a unique curve, if the results are plotted  $Nu/Re^{-0.56}$  versus  $y/d$ . Figure 7 clearly demonstrates this feature for 5 different  $x/d$  values ( $x/d=24,30,40,50,60$ ).

**Heat Transfer Enhancement From Self-Oscillating Impinging Jet.** The stationary jet flow can be converted into an oscillating jet flow by opening the communication lines at points  $A'$  and  $B'$  which are sealed for stationary jet experiments. The communication line diameter at  $A'$  and  $B'$  is 6 mm. The half angle for the nozzle exit section is kept constant at  $\theta=20$  deg. The flow conditions inside the nozzle are identical within quoted uncertainty values for both stationary and oscillating jets, for comparison purposes.

Figure 8 shows Nusselt number distributions over the impingement plate in function of  $y/d$  at different  $Re$  and impingement

plate to nozzle distance ( $x/d$ ) values. When compared to the stationary jet values as presented in Fig. 6, the oscillatory jet heat transfer results are significantly enhanced as shown in Fig. 8.

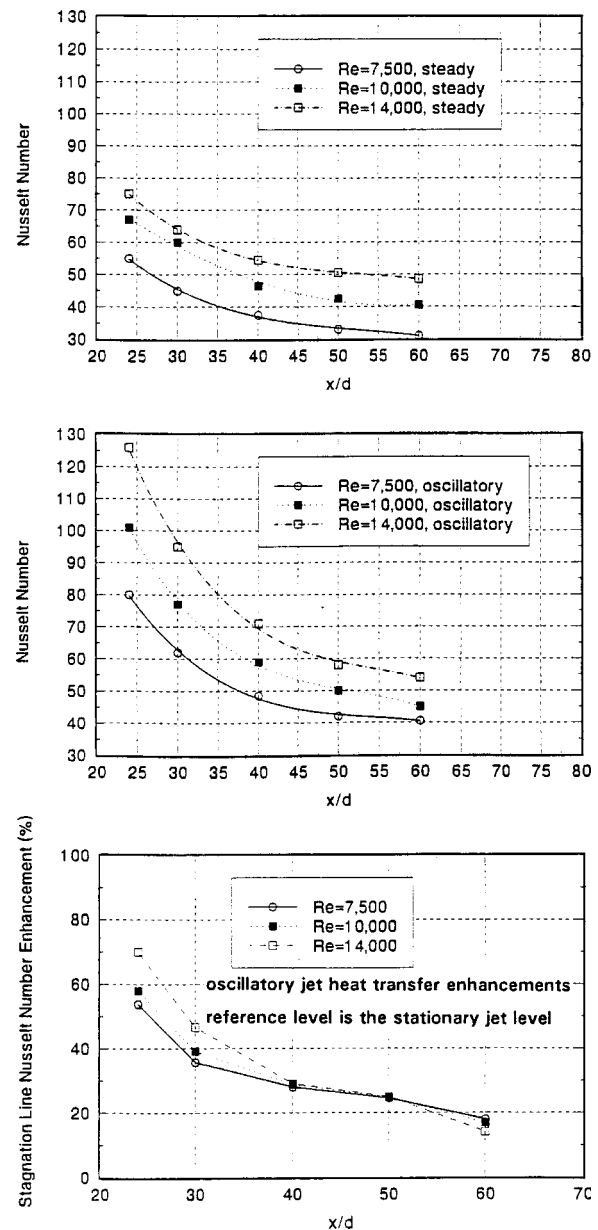
Another observation is the relatively flat distributions of Nusselt number around the stagnation line of the impingement plate compared to the stationary jet distributions. The well known high heat transfer effectiveness of the stagnation region is now spread over a relatively larger area round the stagnation point via organized flapping motion of the jet, Fig. 8 ( $x/d=40,50,60$ ).

When shorter plate to nozzle distances ( $x/d=24,30$ ) are examined, one can observe that the maximum heat transfer does not occur at the exact center line  $y/d=0$ . Off-center peaks are visible in Fig. 9 for  $x/d=24$  and 30. The stagnation line Nusselt numbers are slightly lower than the peak values. The heat transfer enhancement around the centerline point is spread over a much wider area by double peaks presented in Fig. 8, especially at shorter plate to nozzle distances ( $x/d=24$  and 30). This observation is consistent with the fact that the oscillating flow now sweeps through a larger interaction area on the impingement plate with enhanced thermal transport features due to imposed deterministic unsteady motion.

Current experiments suggest that the impingement plate positions where  $x/d < 40$  are the most effective heat transfer enhancement cases when the oscillating jet is impinged on a flat surface. Removing the Reynolds number dependency from the oscillating jet heat transfer results can be achieved by introducing the same type of functional dependency as presented in Eq. 9. For the oscillating jet, the most representative value for  $n$  has been found to be  $n=0.50$  when  $m$  is taken as 1. At a fixed impingement plate to nozzle distance, all three heat transfer distributions for three different Re values again collapse into a unique curve ( $Nu/Re^{-0.50}$  in function of  $y/d$ ). Figure 9 presents this feature for 5 different  $x/d$  values ( $x/d=24,30,40,50,60$ ).

Figures 10, 11, and 12 show significant heat transfer improvement when a stationary jet is converted into an oscillating jet at fixed Reynolds number and impingement plate to nozzle distance. For  $Re=7500$  and  $x/d=24$ , the stagnation line local Nusselt number for oscillatory impinging jet is 30 percent higher than that of the stationary jet. When Reynolds number is increased to 14,000 the enhancement from oscillations is about 71 percent higher compared to stationary impinging jet, at  $x/d=24$ . The heat transfer enhancement almost diminishes in the outer region of the impingement plate where  $y/d > 25$  for all Re numbers and impingement plate to nozzle distance values. In general, the enhancements become more pronounced as the Reynolds number is gradually increased from 7500 to 14,000.

**Comparison of Heat Transfer Enhancements at the Stagnation Line.** Local Nusselt numbers at  $y/d=0$  for both the stationary jet and the oscillating jet are given in Figs. 13(a) and 13(b). For the stationary jet, at  $x/d=30$ , the stagnation point Nusselt number is about 20 percent less than the case for  $x/d=24$  regardless of the Re number of the jet. For larger spacing, where  $x/d > 40$  (40,50,60), dependence of Nusselt number on  $x/d$  becomes weaker as shown in Fig. 13(a). At large nozzle to impingement plate distances, flow near the stagnation point is influenced from ambient air entrainment at a more significant rate compared to short  $x/d$  cases. All of the current heat transfer measurements show that Nusselt number levels at the stagnation point are drastically increased, when the stationary jet is switched into its oscillating mode, Fig. 13(b). The enhancement rate has a slight Re number dependency. However,  $x/d$  dependency at a given Re number is much stronger. When  $x/d < 40$ , oscillation induced enhancements are all greater than 30 percent in reference to the stationary jet Nusselt number values. Figure 13(c) summarizes the relative heat transfer enhancement values in reference to the stationary jet for all Re numbers and  $x/d$  values. The minimum enhancement from oscillations has been measured as approximately 18 percent at  $x/d=60$ . However, if one reduces the plate to nozzle distance, significant Nusselt number enhancement values are encountered. For example at  $x/d=40$  the enhancement



**Fig. 13 Stagnation line heat transfer at various nozzle to impingement plate distances, without and with jet oscillations**

from oscillations is about 30 percent. After this point, the enhancements become stronger for smaller plate to nozzle distances. At  $x/d=24$ , there is about 70 percent heat transfer augmentation on the stagnation line for  $Re=14,000$ .

The oscillation driven enhancements discussed in this paragraph may be explained by using time averaged momentum and thermal energy equations as discussed in previous paragraphs. When unsteady deterministic fluid motion governs the heat transfer process, the additional inertia forces (oscillation stresses) due to the existence of the periodic (deterministic) flow play a significant role in the conservation of linear momentum, Eq. 7. The thermal energy transfer is also influenced because of the existence of the oscillations. For the case of heat transfer, oscillation heat flux terms become as significant as conventional turbulent heat flux terms in the energy conservation statement, Eq. 8, Herr [12]. When  $x/d$  is small, the oscillation heat fluxes and stresses dominate the enhancement process, since the large amount of ambient air entrainment does not affect the thermal energy transfer process. When  $x/d$  is large, the ambient air entrainment overweighs

the oscillation heat fluxes. Thus the enhancement caused by jet oscillations become almost Re number independent.

## Conclusions

It is possible to convert a stationary impinging cooling jet into a self-oscillating-impinging jet by adding two communication ports at the throat section. The unsteady travel of compression and expansion waves in the communication loop can easily put a stationary planar jet into a flapping mode for more effective thermal transport enhancement.

Oscillation frequencies ranging from 20 Hz to 100 Hz are possible by adjusting the proper length and diameter of the communication line and the stagnation pressure in the nozzle plenum chamber.

A triple-decomposition technique applied to the instantaneous velocity field and temperature field clearly indicates the contributions of additional fluid stresses (oscillation stress) and thermal energy transfer due to oscillatory motion (oscillation heat flux). The flapping motion of the jet clearly improves thermal transport near the impingement plate via the imposed deterministic/periodic motion.

Oscillation of the jet increases convection and diffusion significantly, however the kinetic energy production is only slightly modified. Previous studies performed by the authors show that most of the kinetic energy is contained in the deterministic part.

For the case of oscillating-impinging-jet, the area containing the enhanced heat transfer zone on the impingement plate is enlarged effectively when compared to the stationary jet.

Oscillating jet heat transfer results are compared against stationary jet results at three different Reynolds number levels (7500, 10,000, 14,000) at various nozzle to plate distances ( $x/d = 24,30,40,50,60$ ). Significant heat transfer coefficient enhancements ranging from 20 percent to 70 percent over the stationary jet values exist because of the oscillation motion of the impinging jet.

When the jet is put into the oscillating mode, the Reynolds number dependency of oscillating jet heat transfer distribution can be removed by plotting  $Nu/Re^{-0.50}$  against  $y/d$  for a prescribed  $x/d$  value. The data show that at a fixed nozzle to plate distance  $x/d$ , all three heat transfer distributions for individual Re number values collapse into a unique curve.

The current heat transfer enhancement achieved suggest that it is possible to implement the present self-oscillating-impinging-jet concept in future turbine blade internal cooling systems, on rotating disks, in electronic cooling, aircraft/helicopter de-icing systems and heat exchanger systems.

## Acknowledgments

The authors would like to acknowledge the assistantship provided to Frank Herr during the execution of this project at the Turbomachinery Heat Transfer Laboratory by the Department of Aerospace Engineering, the Pennsylvania State University. Thanks are also due to Mr. George Sayers for the precision machining of Aluminum nozzles used in this study.

## Nomenclature

$d$	= nozzle throat width, see Fig. 2 (linear dimension)
$f$	= oscillation frequency of the jet
$h$	= convective heat transfer coefficient
$H$	= thickness of the plexiglass impingement plate
$i, j, k$	= summation indices
$k$	= thermal conductivity
$m, n$	= exponents in Nu number correlation
Nu	= Nusselt number, $Nu = hd/k$
$p$	= instantaneous pressure
$P$	= time averaged pressure
Pr	= molecular Prandtl number
$q$	= wall heat flux rate

$q$	= turbulent kinetic energy using $u, v, w$
$q^*$	= oscillation kinetic energy using $u^*, v^*, w^*$
Re	= Reynolds number, $Re = Ud/v$ , based on bulk velocity of the jet at the throat section
$t$	= time
$T$	= temperature
$U, V, W$	= mean velocity components
$u, v, w$	= velocity fluctuations due to stochastic fluctuations (turbulence)
$u^*, v^*, w^*$	= deterministic (periodic) velocity component
$-\langle uv \rangle$	= ensemble averaged Reynolds stress, $-\langle u_i u_j \rangle$
$-\overline{uv}$	= time averaged Reynolds stress, $-\overline{u_i u_j}$
$x, y, z$	= Cartesian coordinates
$x/d$	= impingement plate to nozzle distance (normalized)
$x$	= axial direction for the jet
$y$	= lateral direction for the jet

## Greek letters

$\phi$	= nozzle half angle at the divergent exit section
$\Theta$	= time averaged local temperature in the flow field
$\theta$	= temperature fluctuations due to stochastic fluctuations (turbulence)
$\theta^*$	= deterministic (periodic) temperature
$\gamma$	= thermal diffusivity
$\eta$	= similarity variable, $\eta = y/x$
$\mu$	= absolute viscosity
$\nu$	= kinematic viscosity $\nu = \mu/\rho$
$\rho$	= density

## Subscripts

char	= characteristic
cond	= conduction
conv	= convection
gen	= generated
$m$	= centerline quantity (at $y=0$ )
$o$	= total condition
rad	= radiated
$t$	= turbulent quantity
$w$	= wall
$w, \text{back}$	= back side of the wall

## Superscripts

*	= deterministic (periodic) quantity
$m, n$	= correlation constants
$\sim$	= instantaneous quantity

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# Forced Convection Heat Transfer Enhancement Using a Self-oscillating Impinging Planar Jet

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Impinging jets are widely used in the local enhancement of heat removed from internal passages of gas turbine blades. Arrays of stationary jets are usually impinged on surfaces of internal cooling passages. The current practice is to benefit from the high heat transfer coefficients existing in the vicinity of the jet impingement region on a target wall. The present study shows that a self-oscillating impinging-jet configuration is extremely beneficial in enhancing the heat removal performance of a conventional (stationary) impinging jet. In addition to a highly elevated stagnation line Nusselt number, the area coverage of the impingement zone is significantly enhanced because of the inherent sweeping motion of the oscillating coolant jet. When an oscillating jet ( $Re=14,000$ ) is impinged on a plate normal to the jet axis ( $x/d=24$  hole to plate distance), a typical enhancement of Nu number on the stagnation line is about 70 %. The present paper explains detailed fluid dynamics structure of the self-oscillating jet by using a triple decomposition technique on a crossed hot wire signal. The current heat transfer enhancement levels achieved suggest that it may be possible to implement the present self-oscillating-impinging-jet concept in future gas turbine cooling systems, on rotating disks, glass tempering/quenching, electronic equipment cooling, aircraft de-icing, combustors and heat exchangers.

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## NOMENCLATURE

$d$	=	nozzle throat width , see Figure 2 (linear dimension)
$f$	=	oscillation frequency of the jet
$h$	=	convective heat transfer coefficient
$H$	=	thickness of the plexiglass impingement plate
$i,j,k$	=	summation indices
$k$	=	thermal conductivity
$m,n$	=	exponents in $Nu$ number correlation
$Nu$	=	Nusselt number, $Nu=hd/k$
$p$	=	instantaneous pressure
$P$	=	time averaged pressure
$Pr$	=	molecular Prandtl number
$q$	=	wall heat flux rate
$q$	=	turbulent kinetic energy using $u,v,w$
$q^*$	=	oscillation kinetic energy using $u^*,v^*,w^*$
$Re$	=	Reynolds number, $Re=Ud/\nu$ , based on bulk velocity of the jet at the throat section
$t$	=	time
$T$	=	temperature
$U,V,W$	=	mean velocity components
$u,v,w$	=	velocity fluctuations due to stochastic fluctuations (turbulence)
$u^*,v^*,w^*$	=	deterministic (periodic) velocity component
$-\langle uv \rangle$	=	ensemble averaged Reynolds stress, $-\langle u_i u_j \rangle$
$-\overline{uv}$	=	time averaged Reynolds stress , $-\overline{u_i u_j}$
$x,y,z$	=	Cartesian coordinates
$x/d$	=	impingement plate to nozzle distance (normalized)
$x$	=	axial direction for the jet
$y$	=	lateral direction for the jet

### Greek letters

$\phi$	=	nozzle half angle at the divergent exit section
$\Theta$	=	time averaged local temperature in the flow field
$\theta$	=	temperature fluctuations due to stochastic fluctuations (turbulence)
$\theta^*$	=	deterministic (periodic) temperature
$\gamma$	=	thermal diffusivity
$\eta$	=	similarity variable, $\eta=y/x$
$\mu$	=	absolute viscosity
$\nu$	=	kinematic viscosity $\nu=\mu/\rho$

$\rho$  = density

### subscripts

*char* = characteristic  
*cond* = conduction  
*conv* = convection  
*gen* = generated  
*m* = centerline quantity (at  $y=0$ )  
*o* = total condition  
*rad* = radiated  
*t* = turbulent quantity  
*w* = wall  
*w,back* = back side of the wall

### superscripts

\* = deterministic (periodic) quantity  
*m,n* = correlation constants  
 $\sim$  = instantaneous quantity

## INTRODUCTION

The advancement of high performance thermal systems has stimulated interest in methods to improve heat transfer rates, especially on the coolant side of hot surfaces. Considerable efforts have been made to increase heat transfer rates by implementing passive enhancement methods that require no direct consumption of external power. The present study deals with a heat transfer enhancement method using planar cooling jets in a self-sustained oscillatory flow mode. Utilizing an impinging jet in oscillatory mode has great cooling potential for hot section components. This method is based on the enhancement of turbulent mixing processes on the coolant side of hot surfaces by generating additional turbulent mixing via periodic oscillations in the impingement region. The term “oscillation” is used for the periodic flapping motion of the jet in a direction normal to orifice/nozzle axis. An obvious advantage of using self-sustained oscillations is that the



power for oscillation is drawn from the existing flow field. The oscillations can easily be generated by a simple orifice/nozzle configuration that has no moving components even in relatively small coolant passages. The flapping motion of the turbulent planar jet spreads the effective thermal transport features over a relatively large impingement area on the target plate. A typical enhancement of Nu number near the stagnation line can be as high as 70 % over stationary jet values.

Stationary impinging jets are widely used to provide high local mass/momentum/heat transfer in a variety of applications including gas turbine cooling, paper drying, helicopter rotor blade de-icing system, glass manufacturing, food processing, etc. Impinging jet systems are popular because of the relative ease of controlling the specific area needed to be cooled/heated/dried and high local heat transfer rates on the target surface. Reviews of earlier heat transfer studies on stationary impinging jets are given in Martin (1977) and Goldstein et al. (1986). Gardon and Cobonpue (1962) described the rate of heat transfer to an impinging jet as proportional to the difference between the target plate temperature and the adiabatic wall temperature, which varied from point to point on the impingement surface. Gardon and Akfirat (1965) noted that the level of turbulence in the jet had a significant effect on the rate of heat transfer between the target plate and the stationary jet. Heat transfer from a flat surface to an oblique impinging jet was studied by Goldstein and Franchett (1988) and Foss (1979). Most studies have involved jets of the same temperature as the ambient fluid. However, when the temperature of the jet is different from that of the ambient air, entrainment of ambient air into the jet complicates the heat transfer problem, Goldstein, Sobolik and Seol (1990) and Goldstein and Seol (1991). Stationary impinging jet based cooling arrangements are currently extremely popular in gas turbine cooling engineering and process engineering.

Fluidic oscillatory nozzles as simple and effective mixing devices have been used by Viets (1975). An extension of the fluidic nozzle concept to supersonic flow conditions is described by Bremhorst, Hailye and Rice (1990) and Raman, Hailye and Rice (1993). These are detailed fluid mechanics studies geared towards using the unsteady motion to actively control viscous flow near surfaces. The current study uses the fluidic oscillating nozzle as a heat transfer augmentation method for general heat transfer applications. The self-oscillating impinging jet improves thermal transport via periodic flapping motion of the jet. The area coverage of the impingement zone is also greatly enlarged by the flapping motion. Details of an analytical flow model for the self-oscillating-

impinging jet used in this study is given by Herr and Camci (1994.a). The analytical model based on a corrected similarity technique is validated via hot wire measurements using a crossed wire. A decomposition approach was used to identify the contribution due to the deterministic jet oscillations and stochastic turbulent fluctuations. The governing equations are derived for oscillating flows with fixed flapping frequency. At a given time, the oscillatory flow field can be approximately described by the solution of a stationary jet flow obtained from a nozzle/orifice attached to a relative reference frame that is oscillating at a prescribed frequency. The corrected similarity solution predicts the axial/lateral mean velocity measurements and oscillation kinetic energy data well, Herr (1995).

Oscillation of the planar jet flow enhances transport process near the target plate because of the existence of additional inertia forces (oscillation stress) and additional thermal fluxes (oscillation heat flux). Experiments show that most of the kinetic energy in the oscillatory jet is contained in the deterministic unsteady part. The magnitude of the unsteady kinetic energy convection due to the flapping motion of the jet is about three times of that due to turbulence. The individual terms of the kinetic energy conservation equation sampled/evaluated from the measured crossed hot wire data show that the convection of kinetic energy in oscillatory flows is stronger than the convection due to turbulent fluctuations. It is concluded that deterministic unsteadiness enhances transport process significantly, Herr and Camci (1994.b). Oscillation increases convection and diffusion significantly.

Enhancement of impingement heat transfer using a self-oscillating circular nozzle is described in Page, Chinnock and Seyed-Yagoobi (1996). This acoustic excitation technique is based on the matching of the shear tone and organ pipe frequency of the pipe's finite length. The flow visualization studies revealed free transverse wave oscillations and enhanced vortex shedding. A collar extension approach results in enhancements in surface transport phenomena. An oscillation frequency of 5000 Hz from a self-oscillating nozzle resulted in measurable heat transfer coefficient enhancement. The nozzle-to-plate distance varied between 4 to 6 diameters. Azevedo, Webb and Queiroz (1994) used a mechanical rotating ball valve to generate pulsations in axial velocity of a circular jet. The axial velocity in time varies between a peak and a minimum in each period created by valve rotations. This approach results in degradation in heat transfer coefficients on the impingement plate. The nozzle to plate distance varies between 2 to 10 diameters. The authors

claim that the type of pulsations generated in this study adds energy primarily on the large scale, with little superimposed small scale turbulence. Liu and Sullivan (1996) studied heat transfer from an acoustically excited circular impinging jet. The investigation is performed with a nozzle having an exit diameter of 12.7 mm and a nozzle to plate spacing less than two diameters. A loudspeaker attached to one side of a rectangular plenum chamber induces organ pipe resonance in the chamber producing plane-wave excitation at the jet exit. The perturbation velocity can be adjusted between 0.05 % and 0.18 % in a frequency range from 600 to 2500 Hz. At small nozzle-to-plate spacing (less than two diameters), the local heat transfer in the wall jet region can be significantly affected by exciting the impinging jet. The heat transfer near the stagnation point remains unchanged. The random vortical structures enhance the local heat transfer. The strong large-scale well-organized vortices formed after the stable pairing induce the unsteady separation of the wall boundary layer and hence, lead to the local heat transfer reduction. Mladin and Zumbrennen (1997) use a rotating ball mechanism to generate a planar air jet with pulsations. The nozzle-to-plate separation distance varies between 0 to 10 nozzle widths with a pulse amplitude ranging from 0 to 50 % of the mean flow velocity. Pulsation frequencies range from 0 to 80 Hz corresponding to Strouhal numbers below 0.106 based on nozzle width and jet discharge velocity. Heat transfer enhancements up to 12 % near the nozzle mid-plane due to surface renewal effects and up to 80 % at distances downstream due to increased turbulence levels are measured where both the pulse amplitude and Strouhal number are the highest. Pulses of small amplitude and low frequency induces a quasi steady behavior at small separation distances, with no significant effect on the time averaged heat transfer near the nozzle mid plane. However, at larger separation distances, increased turbulence associated with the decay of the flow pulse enhanced the time averaged heat transfer away from the nozzle mid plane by up to 20 %.

All of the studies reviewed in the previous paragraph focus on the nozzle to plate distances less than 10 jet diameters (or widths). Although Page et al.'s (1996) study uses a self sustained oscillation mode that does not require external power to maintain pulsations, Liu and Sullivan (1996), Mladin and Zumbrennen (1997) and Azevedo et al. (1994) use loudspeaker and rotating ball valve arrangements that are powered externally. The current study is different from the studies summarized in this paragraph for two main reasons. First, the nozzle-to-impingement plate distance

of the current study ( $24 < x/d < 60$ ) is geared towards engineering applications that can be found in food processing, drying, gas turbine cooling, glass tempering and aeronautical de-icing systems. The second main difference is in the character of unsteady jet motion imposed. The current study benefits from a self-oscillating impinging jet configuration that improves thermal transport via periodic flapping motion of the jet. The area coverage of the flapping fluid motion that is influencing the impingement plate zone is much greater than the coverage of all other studies reviewed in this section. Although the jets acoustically excited (or mechanically oscillated by a rotating ball valve) usually develop strong longitudinal oscillations they influence a limited impingement area on the target plate. The current self-oscillating planar jet in flapping mode has relatively larger area coverage on the target plate.

The present study focuses on the heat transfer aspects of self-oscillating impinging jets used in thermal system engineering. High-resolution heat transfer measurements are presented in a Re number range between 7,500 and 14,000 ( $24 < x/d < 60$ ). This range is sufficiently wide for many current thermal engineering applications. It is possible to implement the present self-oscillating-impinging-jet concept in future gas turbine cooling systems, on rotating disks, in electronic equipment cooling, aircraft de-icing systems and heat exchanger systems.

## EXPERIMENTAL SETUP AND PROCEDURES

**Oscillating jet assembly:** The operation principle of an oscillatory jet is based on the fact that a jet exiting into space between two sufficiently near walls is bi-stable, i.e. may attach to either wall. In addition, a small pressure gradient across the jet at the throat may cause the jet to detach one wall and attach to the opposite one. A simple fluidic nozzle representation of the oscillatory jet used in this study is shown in Figure 1. The air flows through a nozzle shaped contraction, past two communication ports, into an expansion section. The output ports are attached to each other via a communication loop. Due to the proximity of the wall at the nozzle throat, the jet is bi-stable and must attach to one of the walls. Consider the jet to be attached to wall A. Due to large amount of entrainment into the jet, the pressure at the control port A' is relatively low while the pressure at the port B' is relatively high. Since the ports are attached to each other by a feedback loop, a compression wave travels from port B' to port A' tending to raise the pressure there and push the

jet off the wall. Simultaneously, an expansion wave originates at port A' and travels to port B', tending to lower the pressure there and pull the jet onto the wall B. Thus, in a well designed nozzle, the jet will oscillate between walls A and B at a frequency determined by the length of the line connecting ports A' and B' and the line diameter. The frequency is also a weak function of the stagnation pressure entering the nozzle. The improved mixing characteristics of the oscillatory jet are seen in the increased spread angle and in the voids of stagnant gas that are entrained between the waves of the oscillation, Herr and Camci (1994.a, 1994.b).

The flow and heat transfer tests for the current oscillating jet concept have been performed using precision Aluminum nozzles. The planar nozzle sections are precision-machined using fine, spline fitted coordinate points. The nozzle pieces are enclosed in 1.25 cm thick plexiglass sidewalls as shown in Figure 2. A regulated laboratory compressed air supply is fed into the nozzle. The jet emerges from the nozzle having  $d=5$  mm throat width (nozzle height is  $W=25$  mm) at a typical mean velocity of 25 m/s ( $Re=10^4$ ). To study the oscillation caused by the mechanism described above, the unsteadiness of other means must be minimum. With two communicating ports closed, the velocity at the throat has been kept constant within 1 %. For steady flow situation, the jet is laminar, having a turbulence intensity of approximately 0.1 % at the exit plane. The oscillation frequency of the jet is in the range of 10-100 Hz. This frequency is controlled by the stagnation pressure, nozzle throat width, and the length/diameter of the communicating loop. Other details of the experimental set-up can be found in Herr (1995)

**Steady/unsteady flow field measurements:** Instantaneous velocity measurements have been taken with a TSI 1249A-T1.5 miniature  $90^\circ$  crossed hot wire probe and a TSI 1210-T1.5 general purpose single sensor probe using TSI series 1050 constant temperature anemometer. The single sensor probe has been used to measure the jet exit flow at the nozzle throat plane. The crossed wire probe has been calibrated in a velocity range from 0.5 to 40 m/s using a full velocity-yaw angle calibration scheme ( $\pm 45^\circ$ ). The specific calibration scheme is the one defined by Lueptow et al. (1988) and modified by Pancheakesan and Lumley (1993). This stretching coordinate transformation approach improves the measurement resolution at low velocities. Unsteady flow data are sampled at a rate of 10 KHz for 5 seconds. A local averaging technique is used to detect the

fronts of the periodical signals and 100 periods are ensemble averaged. Details of the current crossed hot-wire measurement method are presented in Herr (1995).

**Heat transfer measurements on the target plate:** The target heat transfer surface, located downstream of the nozzle is shown in Figure 2. The target plate to nozzle distance has been varied from 24 to 60 nozzle throat width “d” during the experiments. The impingement plate is made up of 1.25 cm thick clear acrylic. The heat transfer surface is a composite system that includes a double-sided tape layer, Inconel heater foil, black backing paint, a liquid crystal thermo-indicator layer and flush mounted surface thermocouples. Details of the composite heat transfer surface are shown in Figure 3. A low resistivity rectangular steel foil (Inconel-600) has been used as a constant heat flux surface on the target plate. The resistance of the steel foil has been accurately determined using a Keithley 580 micro-ohmmeter to an uncertainty of  $\pm 0.001 \Omega$ . A DC current reversal method removing unwanted offset voltages has been used. This technique lowers random noise in resistance measurements. Wall temperature measurement on the heat transfer surface has been made by using several flush-mounted thin-foil thermocouples of K type. Details of the heat transfer surface construction are given in Wiedner and Camci (1996) and Herr (1995). Flush mounted thermocouples have been inlaid in several spanwise locations along the impingement plate and the junctions have been located at mid-point, Figure 3, Herr (1995). The impingement plate shown in Figure 2 has been traversed in the y direction. The ambient side temperature of the heat transfer surface has also been monitored continuously for conduction heat loss calculations. The losses from the impingement side to the ambient side have been considered in the calculation of the convective heat transfer coefficient. A variable current DC power supply has been used to heat the constant heat flux surface. The jet and the ambient air in the laboratory are at the same temperature. The convective heat transfer coefficient is defined as,

$$h = q_{conv}/(T_w - T_{char}) \quad (1)$$

where  $q_{conv}$  is convective heat flux rate,  $T_w$  the wall temperature and  $T_{char}$  the characteristic temperature. The Nusselt number for this problem is based on the nozzle throat width d,

$$Nu = hd/k . \quad (2)$$

In the present study, the reference temperature  $T_{char}$  is taken as the total temperature of the jet,  $T_o$  . The uniform wall heat flux value over the rectangular surface was fixed at a constant level for all tests. The heat flux generated by Joule heating of the steel foil is balanced by convective, conductive and radiative heat flux,

$$q_{gen} = q_{conv} + q_{cond} + q_{rad} . \quad (3)$$

The conduction heat loss from the impingement side to the back side of the plate is assumed to be one-dimensional,

$$q_{cond} = k_{plate} \partial T / \partial y = k_{plate} (T_w - T_{w,back}) / H \quad (4)$$

where  $T_w$  is the wall temperature on the heat flux surface,  $T_{w,back}$  is the temperature of the back side of the target plate and  $H$  is the thickness of the target plate. During a typical run, conduction losses account for approximately 6-12 % of the generated heat flux. Radiation heat transfer was estimated using the assumptions of black body emissivity and thermal equilibrium between the jet and sidewalls. The shape factors were approximated using an enclosure model and considering each surface as black. The radiation heat flux is approximately 3-9 % of the local generated heat flux.

**Experimental uncertainties:** The final experimental uncertainty estimates of the current study are described using the method given by Kline and McClintock (1953). The uncertainty levels presented are analogous to 95 % coverage or 20:1 odds. A crossed hot wire sensor has been used to measure  $x$  and  $y$  components of the instantaneous velocity field in which temperature variation in the flow field is not significant. The Inconel strip heater used in constant heat flux measurements on the target plate has been disconnected from the DC power supply during hot wire measurements. This approach assures an isothermal flow field that is essential for high quality hot-wire measurements. The instantaneous field has then been de-composed into mean, deterministic and random parts for further analysis. The sources of error in hot wire measurement include flow field temperature

variations, thermal inertia of the sensors, probe misalignment, sampling time errors, constant temperature anemometer noise. The cut-off frequency used has been selected as 4 KHz with an overheat ratio of 1.6. It has been noted from a power spectral density calculation that, energy content of the flow is extremely small for the frequencies higher than 600 Hz. A 0.5 ° alignment error is assumed in the uncertainty analysis. This error contributes 0.1 % to the total error. The predicted error in instantaneous velocity measurements is about 1 %. A detailed account of the error analysis on hot wire measurements is presented in Herr, (1995)

The heat transfer measurements on the constant heat flux surface are influenced from the errors made in the determination of wall temperature  $T_w$ , total temperature  $T_o$  of the jet (ambient), heat flux measurement, throat width measurement and absolute error in thermal conductivity  $k$  of air. Our experiments by varying the black paint and liquid crystal layer thickness showed that the wall temperature error  $\delta(T_w - T_{char}) / (T_w - T_{char})$  is less than 3.4 %. This error is already included in the final uncertainty value of Nusselt number 6.3 % approximately. Estimates of the precision index and bias error values for the main measured quantities in the uncertainty analysis are given in Table 1. Further details of the uncertainty analysis are presented in Herr (1995).

	Precision index	Bias Error	Uncertainty
$q_{gen}$	1.0 %	0.3 %	1.0 %
$q_{cond}$	5.5 %	5.9 %	8.1 %
$q_{rad}$	1.6 %	1.9 %	2.5 %
$(T_w - T_{char})$	2.7 %	2.6 %	3.4 %

Table 1

Estimates of the precision index and bias error values used for the uncertainty analysis



## EXPERIMENTAL RESULTS AND DISCUSSION

**Jet half width and effective mixing:** A useful indication of effective mixing in a jet is the variation of the jet half width in axial direction. The definition of the half width at any streamwise position is the distance between the jet centerline and the point where the local velocity is equal to half of the centerline velocity. Current experiments clearly show that the introduction of the oscillatory motion dramatically increases the jet half width. At  $x/d=60$ , ( $Re=10,000$ ), the half width from a stationary jet is about 8 times the nozzle throat width  $d$ . However, when the same jet is put into a self-oscillatory motion, the half width increases to 13 times the throat width, Herr (1995). This observation is closely related to the fact that a cooling jet in a flapping motion covers a much wider impingement area when compared to a stationary jet. Enhanced turbulent transport with effective mixing is spread over a much larger impingement zone on the target surface via the use of an oscillating jet.

**Frequency of self-sustained jet oscillations:** The oscillation frequency of the jet is determined by taking the fast Fourier transform of instantaneous hot wire data from the exit of the nozzle at,  $\eta=y/x=0$ . Figure 4 shows the variation of oscillation frequency with respect to nozzle throat width  $d$  for constant stagnation pressure in the plenum chamber, ( $\phi=20^\circ$ ). Experiments performed at  $x=180$  mm also show that larger throat widths at a fixed communication line length induce higher oscillation frequencies. Similar trends have been observed at  $y/x=0$ ,  $y/x=0.05$  and  $y/x=0.3$ . Increasing  $d$  from 1 mm to 5 mm provides a five-fold frequency increase from 15 to 75 Hz. Increasing the length of the communication line slightly reduces the oscillation frequency. For the case of increased throat width, the amount of fluid entraining into the throat area increases. Thus the ambient air enters the feedback tube sooner than the case which has a smaller throat width. If the width is too large, the pressure gradient across the jet is not strong enough to move the jet and the jet would not bend to either wall, so that no oscillation will occur. Increasing plenum chamber stagnation pressure and shorter communication lines has a tendency to increase the frequency of oscillations, Herr (1995).

**Axial and lateral mean velocity components:** Mean velocity components have been measured by time averaging the instantaneous hot wire signal at various  $\eta=y/x$  locations at  $x/d=60$  for  $f=75$  Hz. It should be noted that the data discussed from this point on include only the results for the

nozzle that has a throat width of  $d=5\text{mm}$  and a nozzle half angle of  $\phi=20^\circ$ . The measured axial mean velocity ( $U/U_m$ ) distribution of the oscillatory jet is relatively fuller than that of the stationary jet as shown in Figure 5. This is expected because oscillation increases the spreading rate of the jet. The lateral mean velocity ( $V/U_m$ ) distribution of the oscillating jet is enhanced when  $0 < y/x < 0.14$ , ( $\eta=x/d=60$ ). The flapping motion of the jet issued from the fluidic nozzle provides increased transport of mean kinetic energy in the lateral direction. This feature is directly related to the observed enhancements in jet half width leading to increased transport in the lateral direction, Herr (1995). The thick solid lines in Figure 5 represents the predictions from "*the corrected similarity method*" developed by Herr and Camci (1994.a). "*The corrected similarity method*" is an analytical extension of the stationary turbulent jet solution with special emphasis paid to deterministic oscillations of the jet. Figure 5 shows that "*the corrected similarity method*" based prediction of mean velocities is in very good agreement with hot wire measurements where  $0 < y/x < 0.275$ . The stationary turbulent jet solution for U and V (based on a conventional similarity method) is represented by tiny dots in Figure 5.

**Additional stresses and heat fluxes due to deterministic oscillations:** The instantaneous velocity and temperature can be decomposed into three parts, Hussain and Reynolds (1970), a time averaged quantity ( $U_i, \Theta$ ), ensemble averaged deterministic oscillations ( $u_i^*, \theta^*$ ) and stochastic turbulent fluctuations ( $u_i, \theta$ ),

$$\begin{aligned} \tilde{u}_i &= U_i + u_i^* + u_i \\ \tilde{\theta} &= \Theta + \theta^* + \theta \end{aligned} \quad (5)$$

In an oscillatory flow, the continuity equation can be decomposed as,

$$\frac{\partial U_i}{\partial x_i} + \frac{\partial u_i^*}{\partial x_i} + \frac{\partial u_i}{\partial x_i} = 0 \quad (6)$$

Finally, the time averaged momentum equation for a flow with self-sustained oscillations is as follows,

$$\frac{\partial(U_i U_j)}{\partial x_j} = -\frac{1}{\rho} \frac{\partial P}{\partial x_i} + \frac{\partial}{\partial x_j} \left( \nu \frac{\partial U_i}{\partial x_j} - \overline{u_i^* u_j^*} + \overline{u_i u_j} \right) \quad (7)$$

$$\frac{\partial(U_j \theta)}{\partial x_j} = \frac{\partial}{\partial x_j} \left( \gamma \frac{\partial \theta}{\partial x_j} - \overline{u_j^* \theta^*} + \overline{u_j \theta} \right) \quad (8)$$

Equation 7 clearly shows that in addition to shear stress due to molecular viscosity, there are two additional apparent stresses termed as  $-u_i^* u_j^*$  and  $-u_i u_j$ . The apparent stress terms are generated by the non-linear inertia term (convective acceleration) that is the left hand side of the instantaneous momentum equation. Additional inertia force due to random oscillations of the turbulent field is usually termed as Reynolds stress  $-u_i u_j$ . The term  $-u_i^* u_j^*$  is due to the inertia force originating from deterministic variations of the oscillating jet flow. This extra stress term is named as oscillation stress  $-u_i^* u_j^*$  throughout this study. Similar to the developments leading to equation 7, one can obtain a time averaged thermal energy equation, equation 8, Herr and Camci (1994.b). The three terms on the right hand side of the energy equation are conduction heat flux  $\gamma(\partial\theta/\partial x_j)$  due to molecular conductivity, oscillation heat flux  $-u_j^* \theta^*$  due to flapping motion of the jet and turbulent heat flux  $-u_j \theta$  due to random fluctuations of the flow field. Figure 5 also shows a comparison of the time averaged and ensemble averaged turbulent stress  $\langle uv \rangle / U_m^2$  in function of  $\eta$  at  $x/d=60$ ,  $f=80$  Hz. Reynolds stresses due to turbulent fluctuations reach a peak around  $\eta=0.10$  and diminish near the centerline and near the outer edge of the jet where  $\eta>0.30$ .

**Heat transfer from stationary impinging jet:** Stationary jet related Nusselt number distributions on the impingement plate are given in Figure 6 as baseline distributions. Experiments at three different Reynolds numbers (7,500, 10,000 and 14,000) are presented for various impingement

plates to nozzle distance values ( $x/d=24,30,40,50$  and  $60$ ). A strong influence of plate to nozzle distance on stagnation point heat transfer is shown in Figure 6. Nusselt number at the stagnation point increases from 30 to 55 when  $x/d$  is reduced from 60 to 24 at  $Re=7,500$ . This trend repeats itself for  $Re=10,000$  and  $14,000$ . The Nusselt number depends on  $x/d$  in the near impingement region  $0 < y/d < 20$ , while in the wall jet region, the heat transfer is almost independent of  $x/d$  for all  $Re$  numbers. At the minimum  $x/d$  value of 24, the jet is predicted to be fully turbulent for all Reynolds numbers. In other words, potential core does not exist at this  $x/d$ . The heat transfer information given in Figure 6 provide the reference stationary heat transfer distributions when oscillating jet results are compared to stationary jet. In general the local heat transfer on the impingement plate is a function of  $x/d$ ,  $y/d$ ,  $Re$  and  $Pr$ . This functional dependency can be represented by a power law in the form,

$$Nu = \sim Re^n Pr^m . \quad (9)$$

The exact form of equation 9 varies from study to study depending on the heat transfer configuration, flow and thermal boundary conditions, local turbulence production and unsteadiness in the flow. The molecular Prandtl number can be taken constant for air under present conditions. For the stationary impinging jet, the most representative value for  $n$  has been found to be  $n=0.56$ , when  $m$  is unity. At a fixed impingement plate to nozzle distance, all three heat transfer distributions for different  $Re$  values will collapse into a unique curve, if the results are plotted  $Nu/Re^{-0.56}$  versus  $y/d$ . Figure 7 clearly demonstrates this feature for 5 different  $x/d$  values, ( $x/d=24,30,40,50,60$ ).

**Heat transfer enhancement from self-oscillating impinging jet:** The stationary jet flow can be converted into an oscillating jet flow by opening the communication lines at points A' and B' which are sealed for stationary jet experiments. The communication line diameter at A' and B' is 6 mm. The half angle for the nozzle exit section is kept constant at  $\phi=20^\circ$ . The flow conditions inside the nozzle are identical within quoted uncertainty values for both stationary and oscillating jets, for comparison purposes.

Figure 8 shows Nusselt number distributions over the impingement plate in function of  $y/d$  at different Re and impingement plate to nozzle distance ( $x/d$ ) values. When compared to the stationary jet values as presented in Figure 6, the oscillatory jet heat transfer results are significantly enhanced as shown in Figure 8.

Another observation is the relatively flat distributions of Nusselt number around the stagnation line of the impingement plate compared to the stationary jet distributions. The well known high heat transfer effectiveness of the stagnation region is now spread over a relatively larger area around the stagnation point via organized flapping motion of the jet, Figure 8, ( $x/d=40,50,60$ ).

When shorter plate to nozzle distances ( $x/d=24,30$ ) are examined, one can observe that the maximum heat transfer does not occur at the exact center line  $y/d=0$ . Off-center peaks are visible in Figure 9 for  $x/d=24$  and 30. The stagnation line Nusselt numbers are slightly lower than the peak values. The heat transfer enhancement around the centerline point is spread over a much wider area by double peaks presented in Figure 8, especially at shorter plate to nozzle distances ( $x/d=24$  and 30). This observation is consistent with the fact that the oscillating flow now sweeps through a larger interaction area on the impingement plate with enhanced thermal transport features due to imposed deterministic unsteady motion.

Current experiments suggest that the impingement plate positions where  $x/d < 40$  are the most effective heat transfer enhancement cases when the oscillating jet is impinged on a flat surface. Removing the Reynolds number dependency from the oscillating jet heat transfer results can be achieved by introducing the same type of functional dependency as presented in equation 9. For the oscillating jet, the most representative value for  $n$  has been found to be  $n=0.50$  when  $m$  is taken as 1. At a fixed impingement plate to nozzle distance, all three heat transfer distributions for three different Re values again collapse into a unique curve ( $Nu/Re^{-0.50}$  in function of  $y/d$ ). Figure 9 presents this feature for 5 different  $x/d$  values, ( $x/d=24,30,40,50,60$ ).

Figures 10, 11, and 12 shows significant heat transfer improvement when a stationary jet is converted into an oscillating jet at fixed Reynolds number and impingement plate to nozzle distance. For  $Re=7,500$  and  $x/d=24$ , the stagnation line local Nusselt number for oscillatory impinging jet is 30 % higher than that of the stationary jet. When Reynolds number is increased to 14,000 the enhancement from oscillations is about 71 % higher compared to stationary impinging jet, at  $x/d=24$ .

The heat transfer enhancement almost diminishes in the outer region of the impingement plate where  $y/d > 25$  for all Re numbers and impingement plate to nozzle distance values. In general, the enhancements become more pronounced as the Reynolds number is gradually increased from 7,500 to 14,000.

**Comparison of heat transfer enhancements at the stagnation line:** Local Nusselt numbers at  $y/d=0$  for both the stationary jet and the oscillating jet are given in Figures 13.a and 13.b. For the stationary jet, at  $x/d=30$ , the stagnation point Nusselt number is about 20 % less than the case for  $x/d=24$  regardless of the Re number of the jet. For larger spacing, where  $x/d > 40$  (40,50,60), dependence of Nusselt number on  $x/d$  becomes weaker as shown in Figure 13.a. At large nozzle to impingement plate distances, flow near the stagnation point is influenced from ambient air entrainment at a more significant rate compared to short  $x/d$  cases. All of the current heat transfer measurements show that Nusselt number levels at the stagnation point are drastically increased, when the stationary jet is switched into its oscillating mode, Figure 13.b. The enhancement rate has a slight Re number dependency. However,  $x/d$  dependency at a given Re number is much stronger. When  $x/d < 40$ , oscillation induced enhancements are all greater than 30 % in reference to the stationary jet Nusselt number values. Figure 13.c summarizes the relative heat transfer enhancement values in reference to the stationary jet for all Re numbers and  $x/d$  values. The minimum enhancement from oscillations has been measured as approximately 18 % at  $x/d=60$ . However, if one reduces the plate to nozzle distance, significant Nusselt number enhancement values are encountered. For example at  $x/d=40$  the enhancement from oscillations is about 30 percent. After this point, the enhancements become stronger for smaller plate to nozzle distances. At  $x/d=24$ , there is about 70 percent heat transfer augmentation on the stagnation line for  $Re=14,000$ .

The oscillation driven enhancements discussed in this paragraph may be explained by using time averaged momentum and thermal energy equations as discussed in previous paragraphs. When unsteady deterministic fluid motion governs the heat transfer process, the additional inertia forces (oscillation stresses) due to the existence of the periodic (deterministic) flow play a significant role in the conservation of linear momentum, equation 7. The thermal energy transfer is also influenced because of the existence of oscillations. For the case of heat transfer, oscillation heat flux terms

become as significant as conventional turbulent heat flux terms in the energy conservation statement, equation 8, Herr (1995). When  $x/d$  is small, the oscillation heat fluxes and stresses dominate the enhancement process, since the large amount of ambient air entrainment does not affect the thermal energy transfer process. When  $x/d$  is large, the ambient air entrainment outweighs the oscillation heat fluxes. Thus the enhancement caused by jet oscillations become almost Re number independent.

## CONCLUSIONS

It is possible to convert a stationary impinging cooling jet into a self-oscillating-impinging jet by adding two communication ports at the throat section. The unsteady travel of compression and expansion waves in the communication loop can easily put a stationary planar jet into a flapping mode for more effective thermal transport enhancement.

Oscillation frequencies ranging from 20 Hz to 100 Hz are possible by adjusting the proper length and diameter of the communication line and the stagnation pressure in the nozzle plenum chamber.

A triple-decomposition technique applied to the instantaneous velocity field and temperature field clearly indicates the contributions of additional fluid stresses (oscillation stress) and thermal energy transfer due to oscillatory motion (oscillation heat flux). The flapping motion of the jet clearly improves thermal transport near the impingement plate via the imposed deterministic/periodic motion.

Oscillation of the jet increases convection and diffusion significantly, however the kinetic energy production is only slightly modified. Previous studies performed by the authors show that most of the kinetic energy is contained in the deterministic part.

For the case of oscillating-impinging-jet, the area containing the enhanced heat transfer zone on the impingement plate is enlarged effectively when compared to the stationary jet.

Oscillating jet heat transfer results are compared against stationary jet results at three different Reynolds number levels (7,500, 10,000, 14,000) at various nozzle to plate distances ( $x/d=24,30,40,50,60$ ). Significant heat transfer coefficient enhancements ranging from 20 % to 70 % over the stationary jet values exist because of the oscillation motion of the impinging jet.

When the jet is put into the oscillating mode, the Reynolds number dependency of oscillating jet heat transfer distribution can be removed by plotting  $Nu/Re^{-0.50}$  against  $y/d$  for a prescribed  $x/d$  value. The data show that at a fixed nozzle to plate distance  $x/d$ , all three heat transfer distributions for individual Re number values collapse into a unique curve.

The current heat transfer enhancement achieved suggest that it is possible to implement the present self-oscillating-impinging-jet concept in future turbine blade internal cooling systems, on rotating disks, in electronic equipment cooling, aircraft/helicopter de-icing systems and heat exchanger systems.

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Acknowledgments

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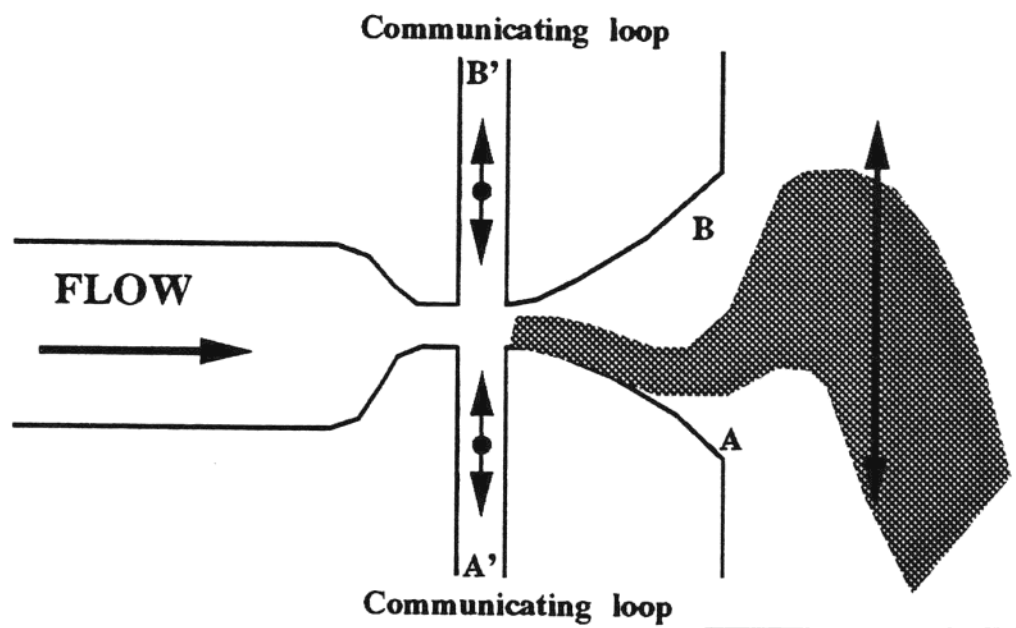


Figure 1 Nozzle issuing a planar jet with self sustained oscillations

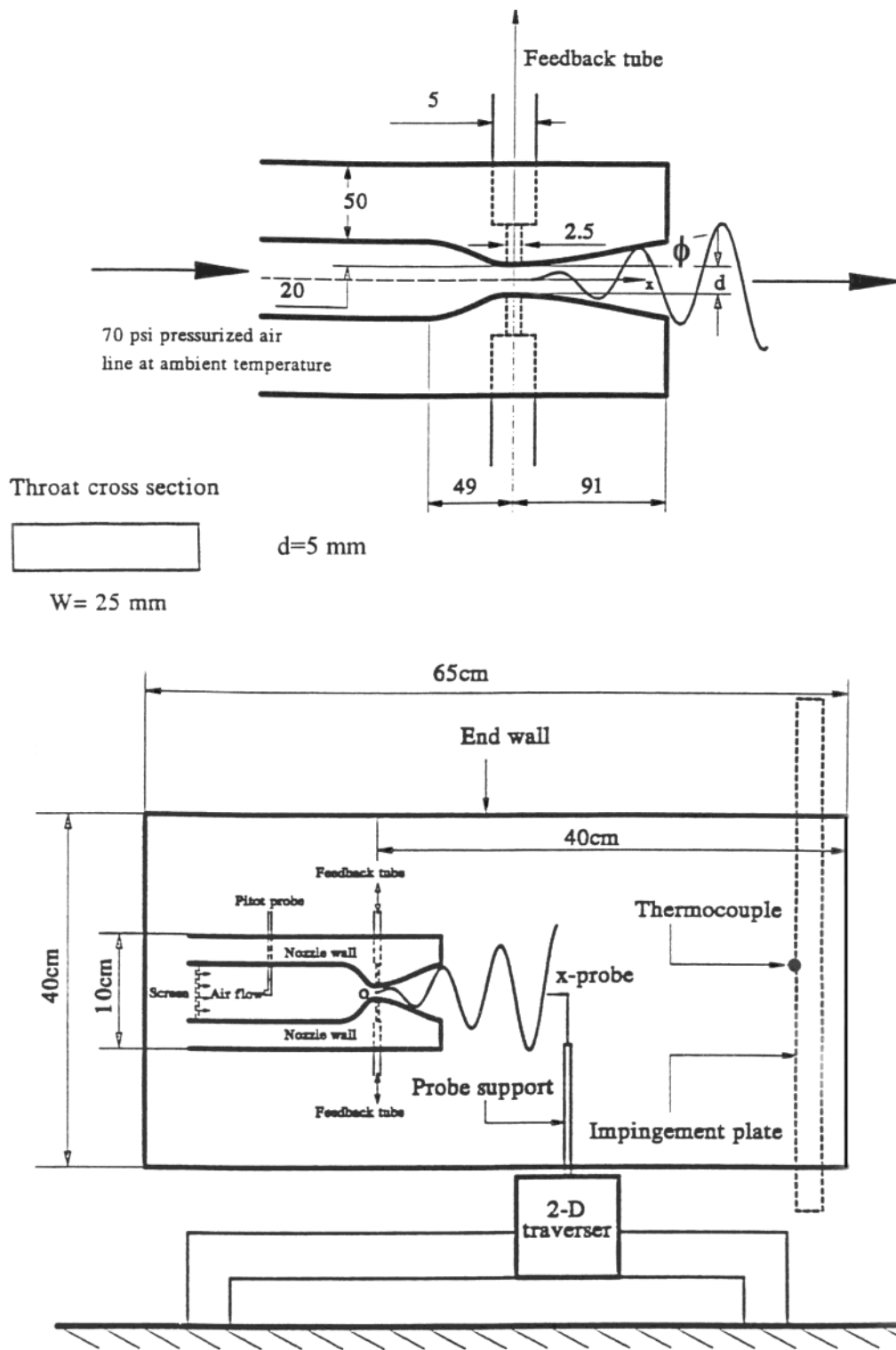


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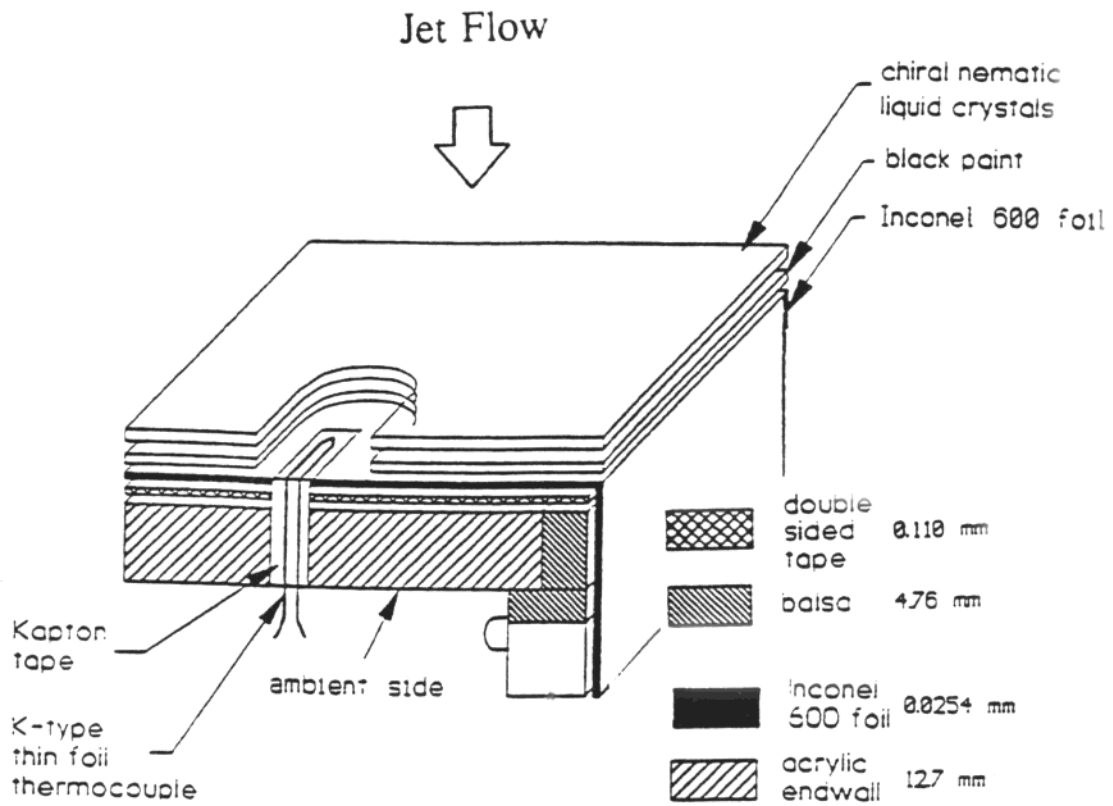


Figure 3 Heat transfer surface construction

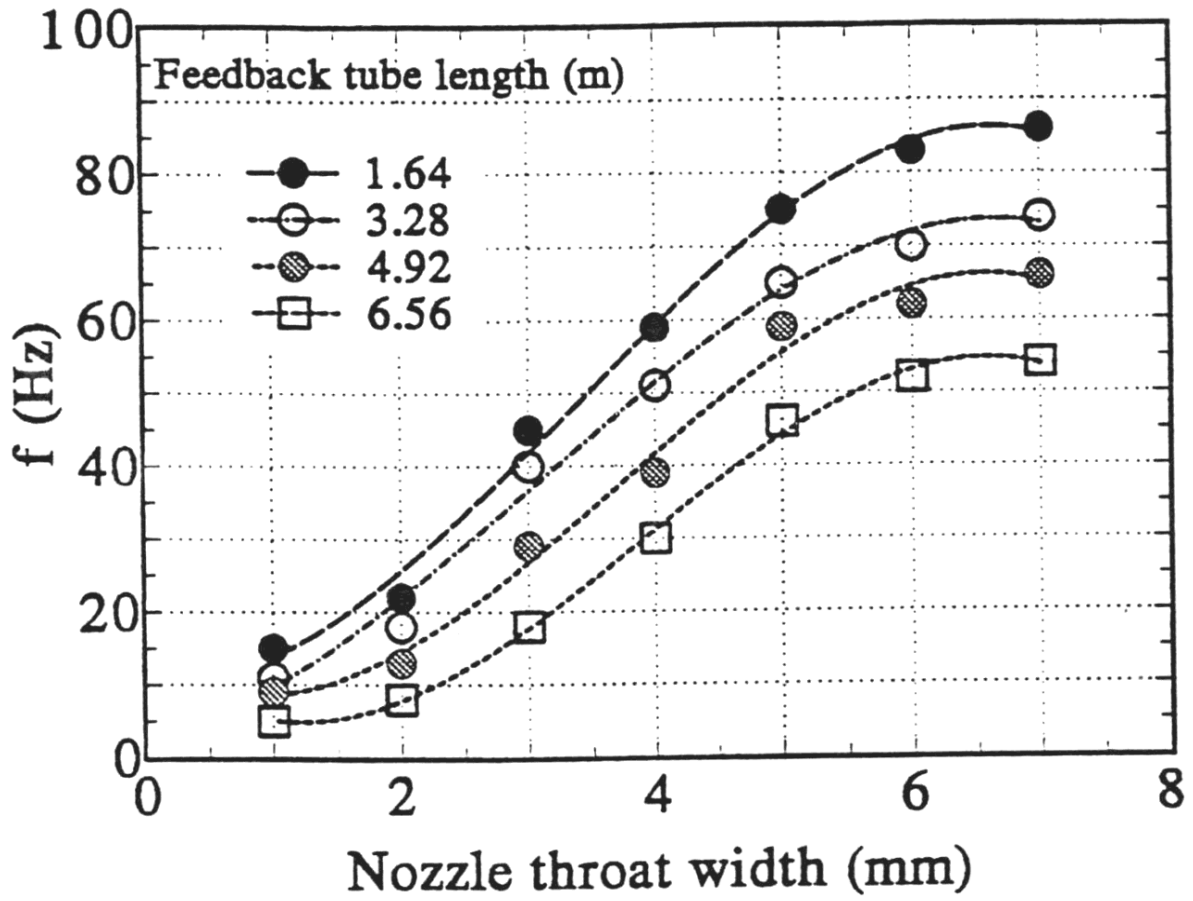


Figure 4 Effect of nozzle throat width and communicating tube length on oscillation frequency of the jet,  $\phi=20^\circ$

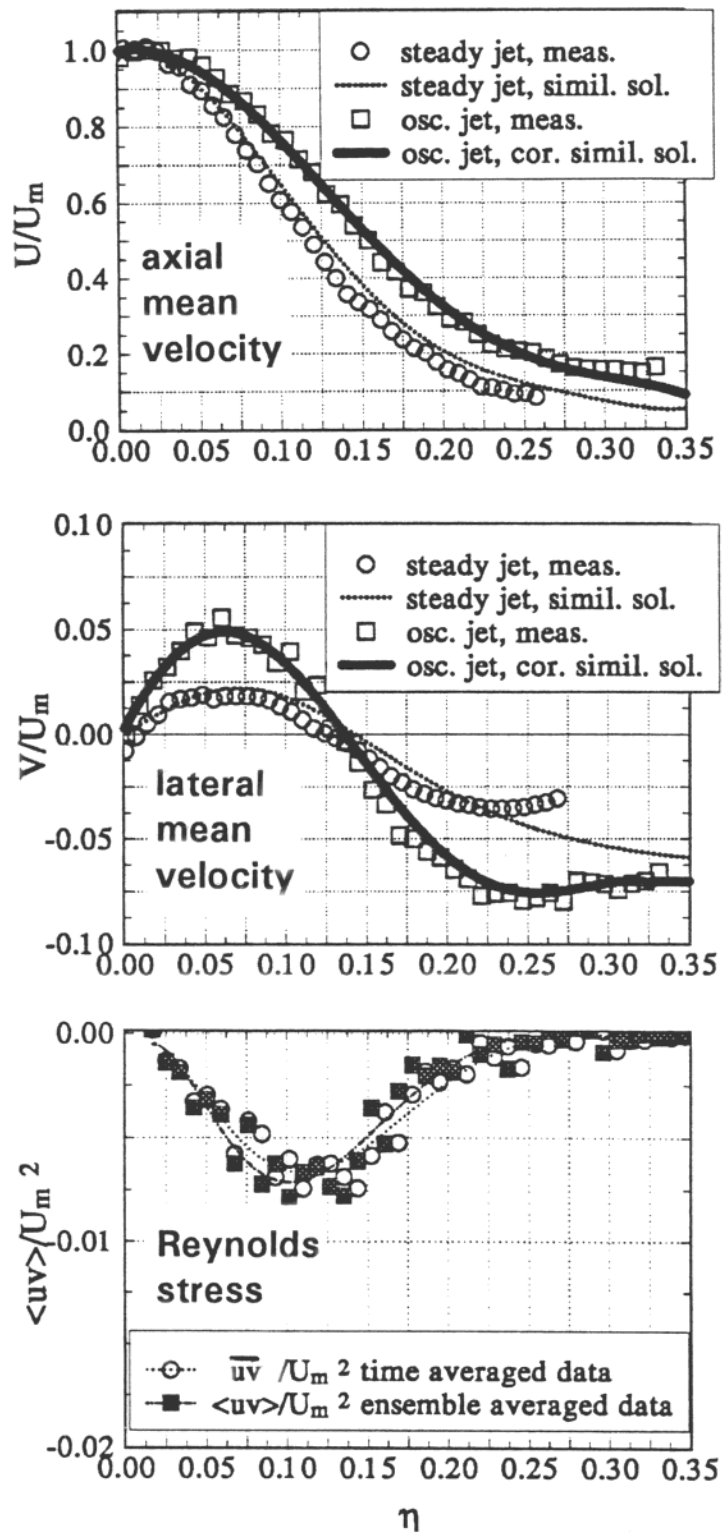


Figure 5 Axial, lateral mean velocity profiles and Reynolds stresses in the oscillating jet



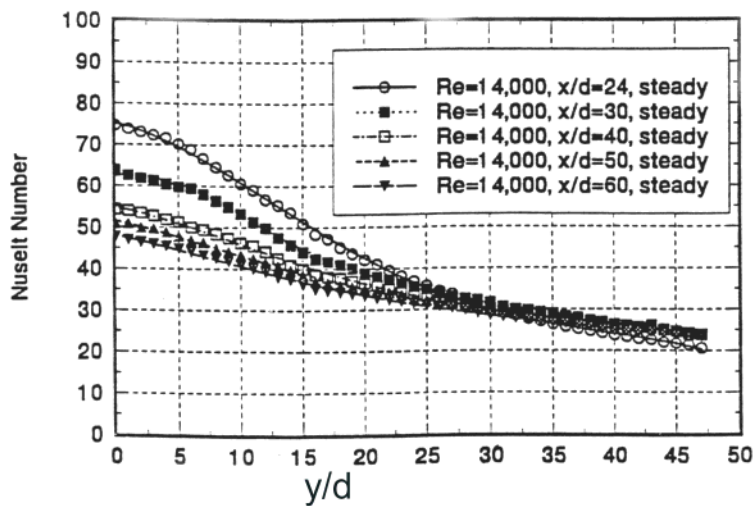
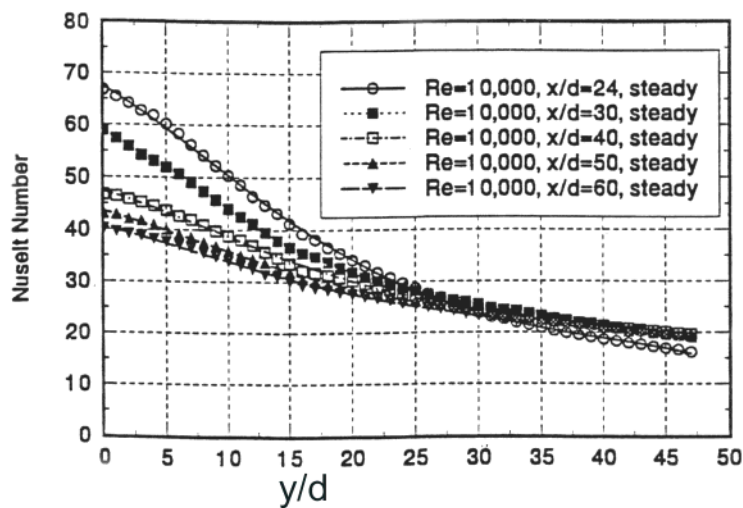
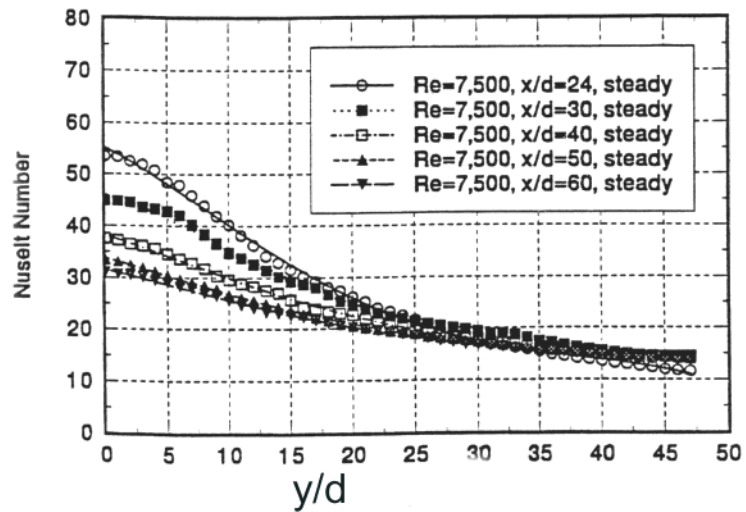


Figure 6 Stationary jet Nusselt number distributions on impingement plate, influence of Reynolds number and impingement plate location  $x/d$ , baseline data set

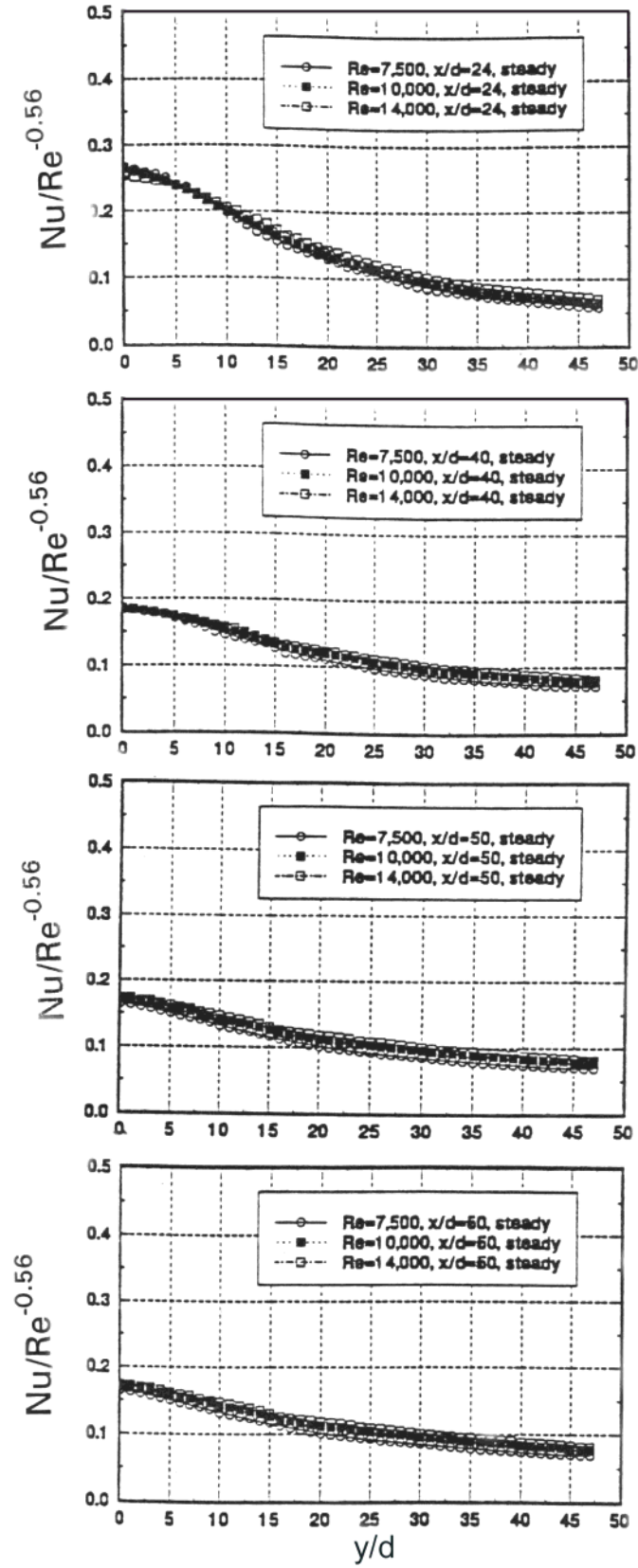


Figure 7

Impingement plate heat transfer distribution from stationary jet,  $Nu/Re^{-0.56}$  distribution for all Reynolds numbers in function of  $y/d$

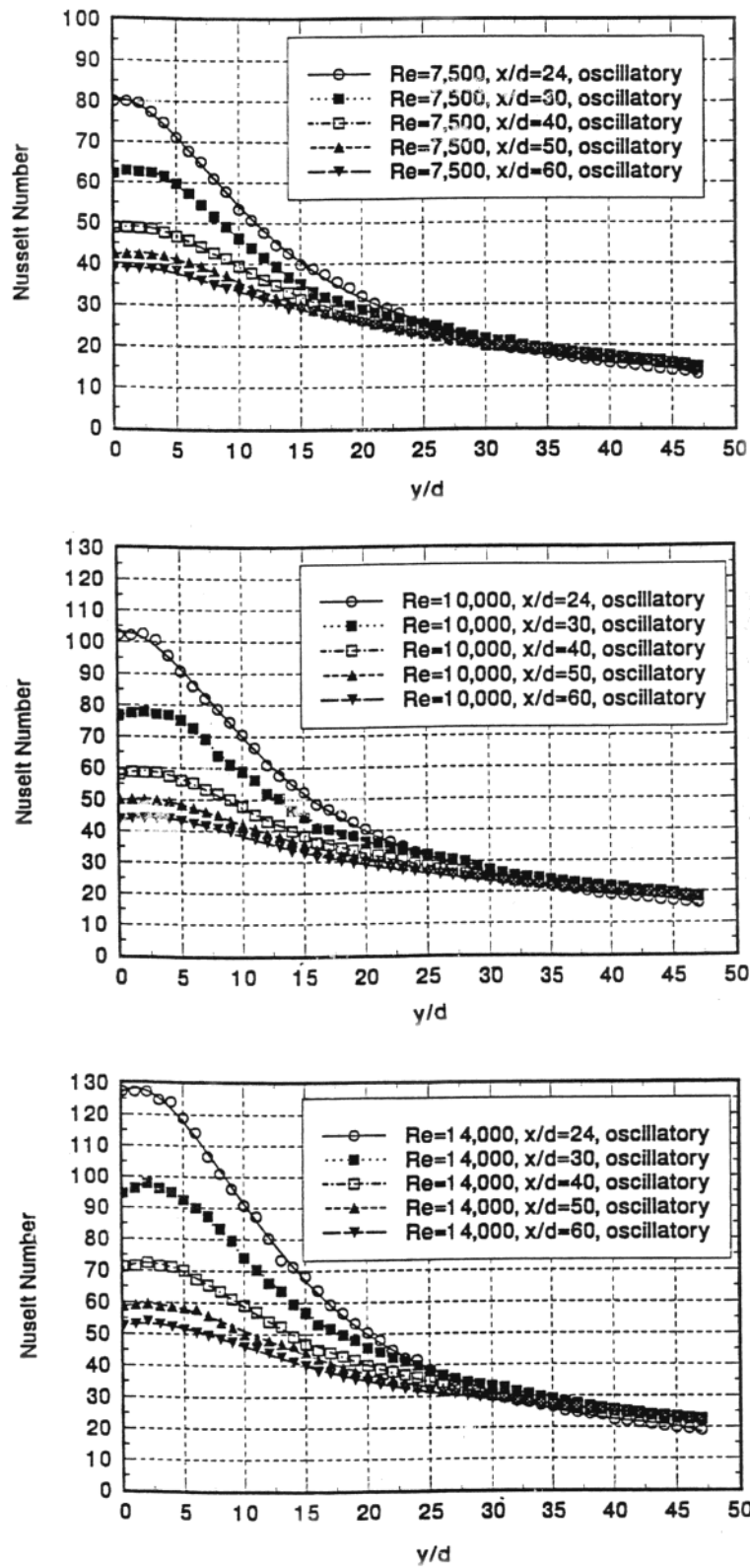


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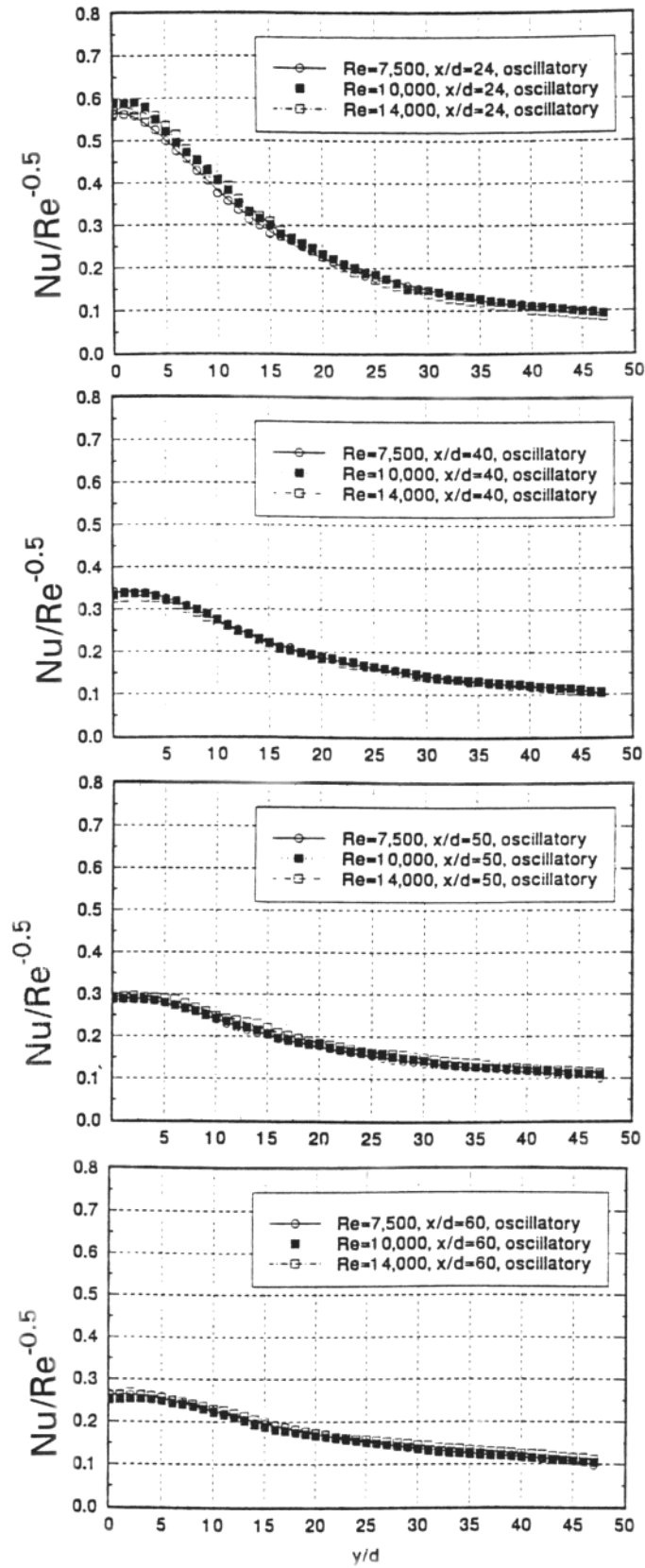


Figure 9 Impingement plate heat transfer distribution from oscillating jet,  $Nu/Re^{-0.50}$  distribution for all Reynolds numbers in function of  $y/d$

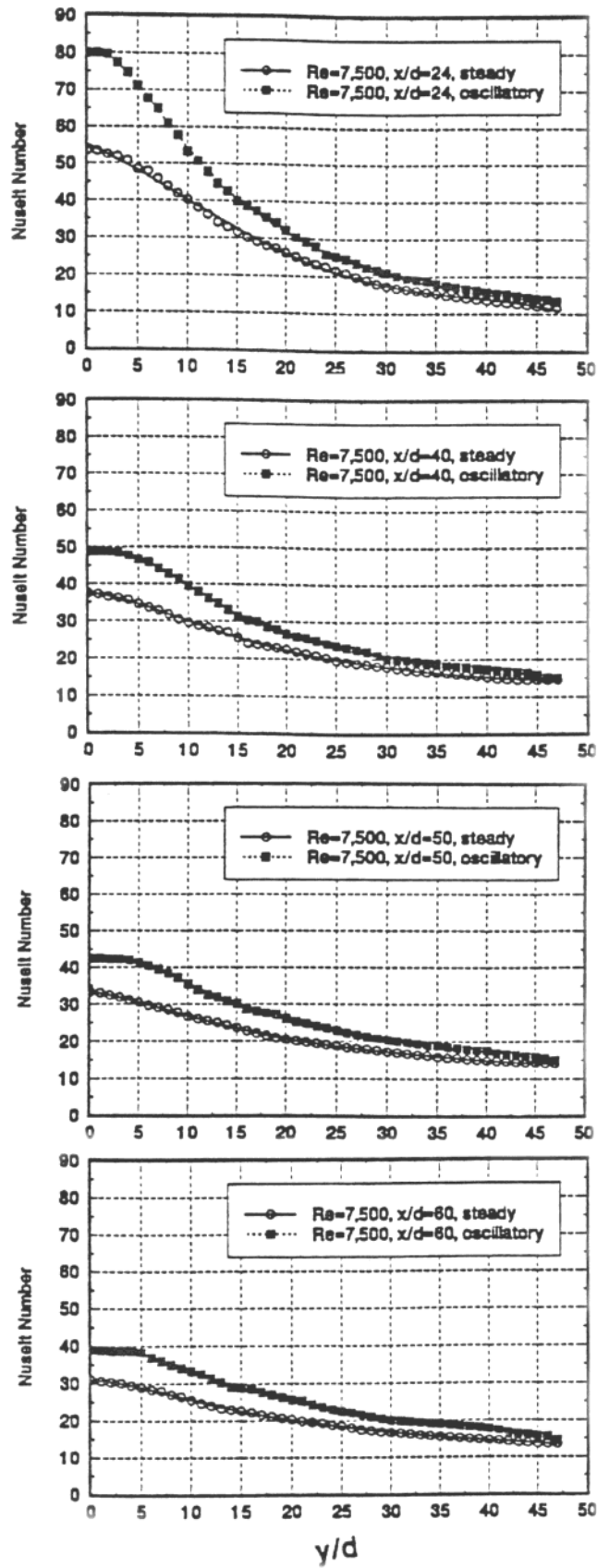


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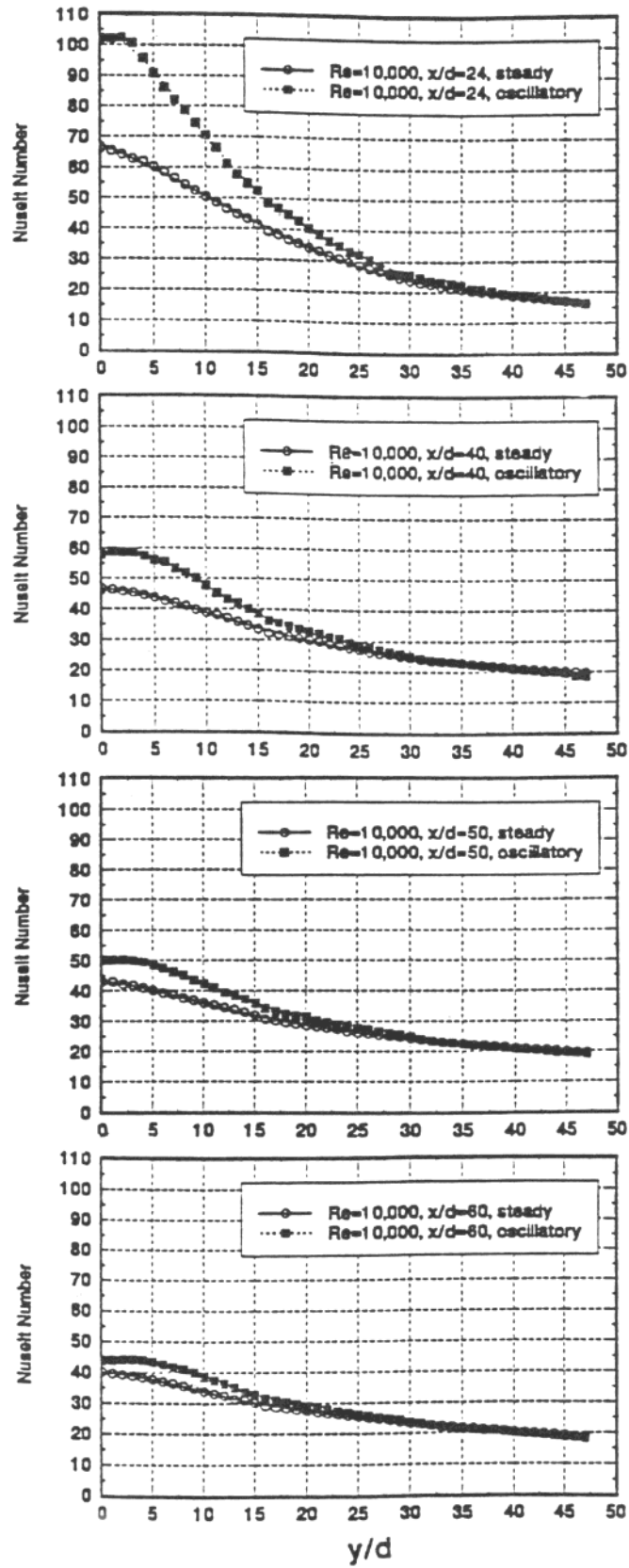


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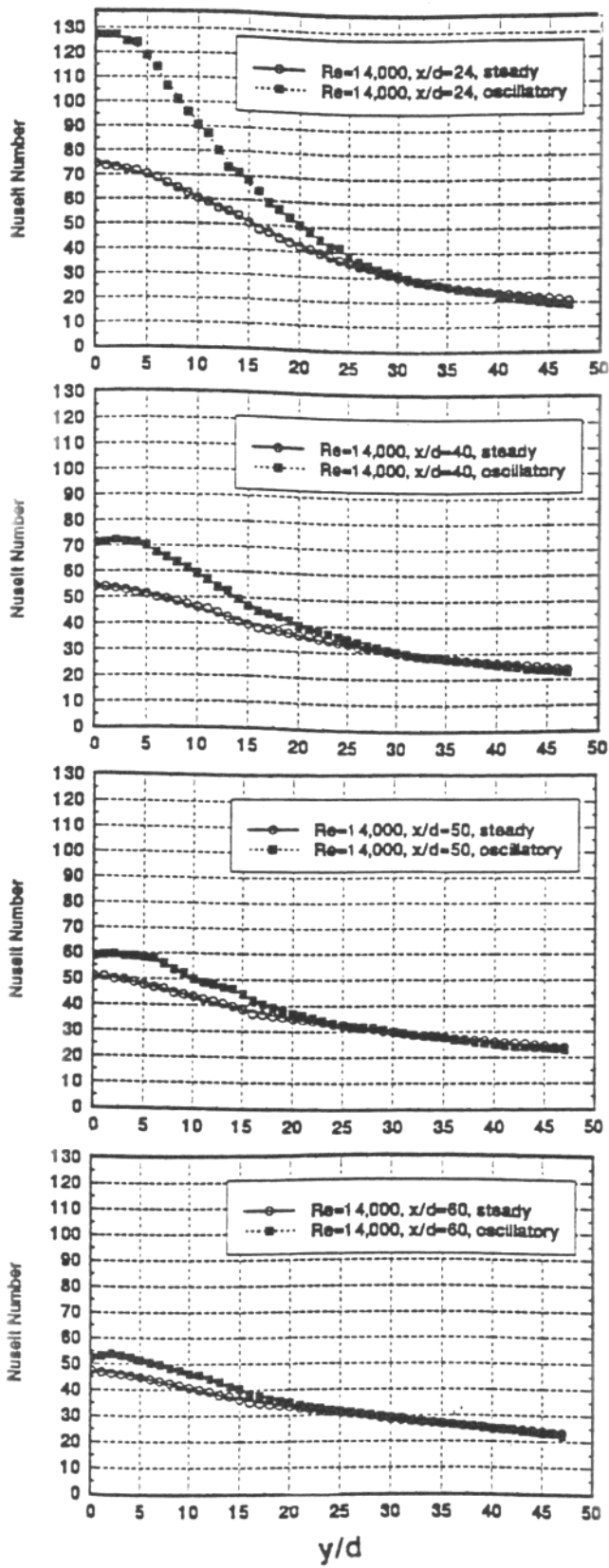


Figure 12 Comparison of oscillating jet and stationary jet heat transfer distributions on the impingement plate,  $Re = 14,000$

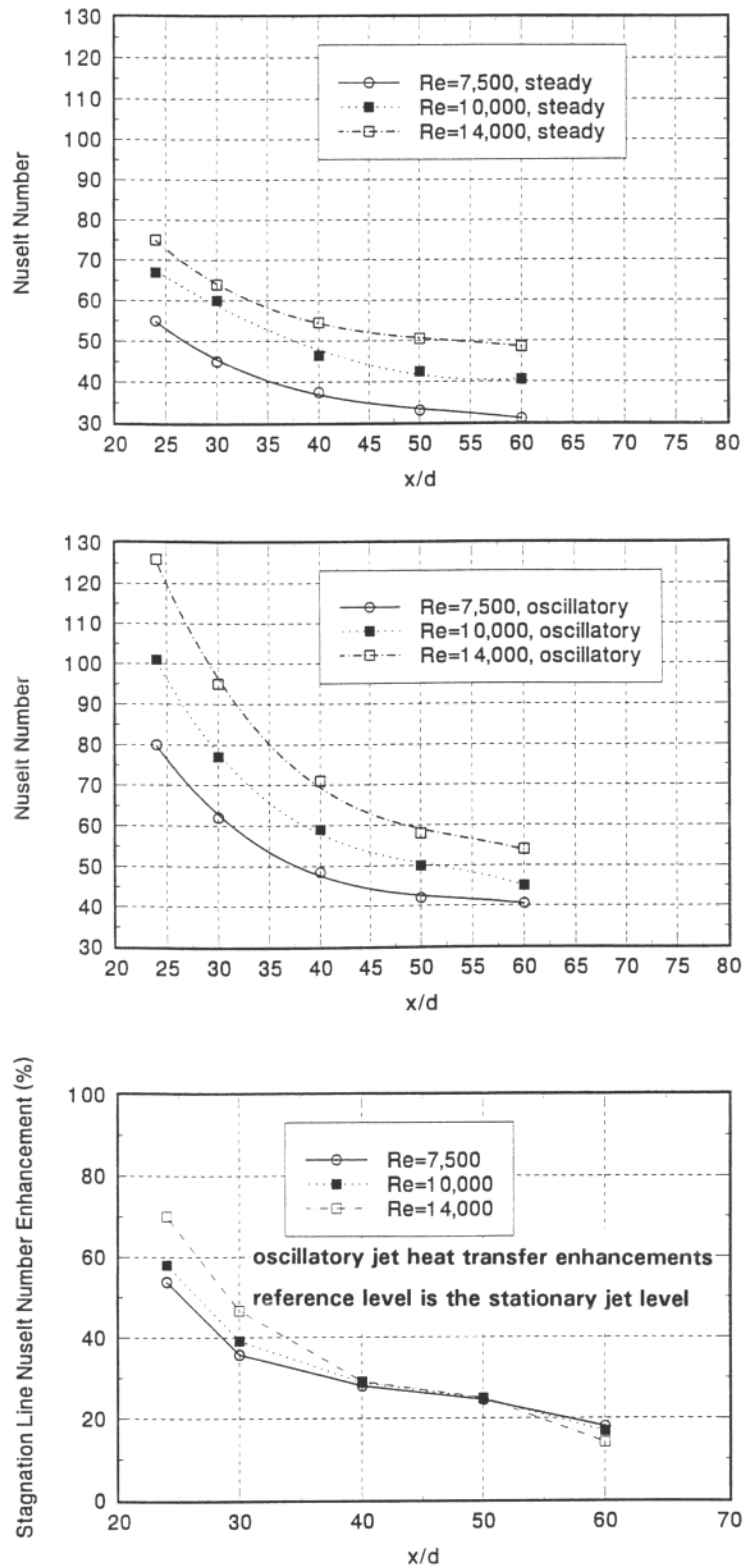


Figure 13 Stagnation line heat transfer at various nozzle to impingement plate distances, without and with jet oscillations