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# Formation and growth of sub-3-nm aerosol particles in experimental chambers

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Atmospheric new particle formation (NPF), which is observed in many environments globally, is an important source of boundary-layer aerosol particles and cloud condensation nuclei, which affect both the climate and human health. To better understand the mechanisms behind NPF, chamber experiments can be used to simulate this phenomenon under well-controlled conditions. Recent advancements in instrumentation have made it possible to directly detect the first steps of NPF of molecular clusters (~1-2 nm in diameter) and to calculate quantities such as the formation and growth rates of these clusters. Whereas previous studies reported particle formation rates as the flux of particles across a specified particle diameter or calculated them from measurements of larger particle sizes, this protocol outlines methods to directly quantify particle dynamics for cluster sizes. Here, we describe the instrumentation and analysis methods needed to quantify particle dynamics during NPF of sub-3-nm aerosol particles in chamber experiments. The methods described in this protocol can be used to make results from different chamber experiments comparable. The experimental setup, collection and post-processing of the data, and thus completion of this protocol, take from months up to years, depending on the chamber facility, experimental plan and level of expertise. Use of this protocol requires engineering capabilities and expertise in data analysis.

### Introduction

New particle formation (NPF) is a major source of atmospheric aerosol particles<sup>1,2</sup>. It contributes substantially to global cloud condensation nuclei concentrations<sup>3,4</sup> and may also contribute to haze formation<sup>5</sup>. NPF involves the formation of sub-3-nm charged and neutral clusters from atmospheric vapors and their growth to stable aerosol particles<sup>6–8</sup>. This phenomenon is observed in a wide range of environments with varying levels of precursor vapors and different meteorological conditions<sup>9,10</sup>. The mechanisms leading to cluster formation and subsequent growth are currently under investigation, as are the particle-formation potentials of different biogenic and anthropogenic precursor vapors and their relative importance. However, studying these processes on the basis of field measurements is challenging because the contributions of different factors cannot be isolated and studied independently. There is also a lack of reliable atmospheric observations from many environments because of difficulties in operating all the necessary instrumentation under challenging field conditions. Consequently, scientists studying NPF have relied on laboratory experiments to validate their hypotheses, unveil hidden mechanisms and make new discoveries.

### Development of the protocol

In general, chambers have been operated to evaluate atmospheric gas-phase chemical mechanisms governing particle formation and to characterize secondary organic aerosols  $(SOAs)^{11-49}$ . However, older laboratory measurements were often hampered by unmeasured contaminant levels, atmospherically irrelevant high vapor concentrations and insufficient instrumentation to detect the forming clusters and their precursors<sup>50–53</sup>. In addition, all these studies relied on detection of particles >3 nm. More recent laboratory experiments have allowed for highly controlled experimental

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conditions (in terms of precursor concentrations and external conditions such as temperature and humidity) for studying the basic physical and chemical mechanisms behind particle formation and growth at a molecular level<sup>23,54–101</sup>. These experiments were also used for imitating different present-day<sup>34,102</sup> or preindustrial environments<sup>103–105</sup>, as well as for making projections into the future by screening changes induced by air pollution mitigation and global warming. Scientists have even simulated clouds on other planets, such as Mars<sup>106</sup>. In experimental reactors, it is usually possible to change one variable at a time to study its effects on particle properties and formation and growth mechanisms. Laboratory experiments can also be repeated under constant meteorological conditions, unlike atmospheric observations.

Collectively, chambers and flow tubes are referred to as environmental reactors. Flow tubes are tubular; the reactants are introduced at one end into the mixing area, and the products are sampled from the other end for analysis<sup>107</sup>. A chamber usually has a higher volume-to-surface ratio than do flow tubes, which minimizes the wall loss effect described in the following sections<sup>108</sup>. In addition, chambers have a vertical orientation, in contrast to flow tubes, which results in smaller particle losses<sup>109</sup>. Reactants are mixed actively (using a fan) or passively into the whole chamber volume before sample collection. Chambers also enable collection of higher sample volumes than do flow tubes. Chambers can be operated in either batch or continuous mode. In batch mode, the reactants are introduced into the chamber, left to react and then sampled for analysis. Chambers with flexible walls are more suitable for such experiments. In continuous mode, a constant flow of synthetic air (or a similar gaseous mixture) is introduced into the reactor in order to maintain a steady pressure to compensate for the flow rate of the sampling instruments. In continuous mode, concentrations of reactants eventually reach a steady state. The chambers operated in continuous mode are known as continuously mixed flow reactors (CMFRs) regardless of whether reactants are actively mixed or not<sup>110</sup>. Flow tubes are discussed further in the 'Chamber considerations and requirements' section, but this protocol focuses on the use of experimental chambers.

The chemical systems that have been studied in experimental flow reactors vary between simple binary acid-water systems and more complex systems involving several precursor vapors and oxidants. Sulfuric acid binary nucleation with water has been the focus of the vast majority of NPF experiments<sup>51,52,59-65,67,69,73,83,111</sup>. Other studies concentrated on the ternary nucleation of sulfuric acid-water with ammonia<sup>52,64,71,73</sup> or amines<sup>52,64,71-73,76,78,81</sup>, and aromatic vapors<sup>58</sup>. A small number of studies have focused on methane sulfonic acid (MSA) nucleation<sup>79,80,84</sup>. Similar chemical systems involving sulfuric acid have also been studied in chambers<sup>54,58,74,75,82,90,95,96,102</sup>. Recently, many chamber experiments have studied NPF involving organic compounds, such as monoterpenes, sesquiterpenes and isoprene<sup>23,70,75,85,87,89,94,95,97,98,102,112</sup>. Some chamber studies used emissions from real plants or trees<sup>86,91,113</sup>. A large number of other chamber and flow tube studies have concentrated on particle mass yields and formation of secondary organic aerosols, but they are out of the scope of this protocol.

The procedure described in this protocol has been used by Wagner et al.<sup>97</sup> and Lehtipalo et al.<sup>102</sup> for obtaining particle formation rates ( $J_{dp}$  values) and growth rates (GRs) from experiments in the CLOUD (cosmics leaving outdoor droplets) chamber (Table 1), and it builds upon the earlier NPF studies by the CLOUD collaboration<sup>54,74,77,96,98,103</sup>, as well as experience from earlier laboratory and field studies. Many methods have been proposed to quantify NPF, and these could potentially yield different results from the same dataset. Similar to the protocol for analyzing  $J_{dp}$  and GR from atmospheric data presented by Kulmala et al.<sup>114</sup>, we describe in this protocol a standard procedure for acquiring particle dynamics from chamber measurements.

In earlier chamber studies<sup>77,98,102</sup>, the error on the  $J_{dp}$  values was calculated on the basis of the propagation of error, by taking into account the statistical and systematic uncertainties and run-to-run repeatability in the chamber (assumed to be 30% for the CLOUD chamber). The systematic errors include errors in the concentration measurement, dilution and wall loss, whereas the statistical errors include uncertainty in dN/dt (time-derivative of the total particle concentration above a certain threshold) and coagulation sink. In this protocol, in addition to analyzing  $J_{dp}$  and GR values from chamber experiments, we present a recommended method for calculating the error in their measurement.

#### Overview of the procedure

In this protocol, we introduce a standard method for measuring, correcting and analyzing particle formation dynamics from chamber experiments in order to make results from different chamber experiments comparable. With the advancement of particle counters and mass spectrometers,

#### Table 1 | Abbreviations

D	Diffusion coefficient
dp	Mobility diameter of particle
d <sub>p,mean,i</sub>	Mean diameter of the size bin <i>i</i>
d <sub>p,ref</sub>	Mobility diameter of the reference particle
d <sub>u</sub>	Upper diameter of size bin
F	Experimentally determined correction factor for wall loss in chamber
GR	Particle growth rate
J	Formation rate
J <sub>ap</sub>	Apparent formation rate
J <sub>dp</sub>	Formation rate of particles with diameter $d_p$
$J^{\pm}_{dp}$	Formation rate of ions (charged particles) with diameter $d_p$
J <sub>n</sub>	Formation rate of neutral particles
J <sub>n,tot</sub>	Neutral fraction of the total formation rate
J <sub>rec</sub>	Formation rate of neutral particles by recombination of ions
J <sub>tot</sub>	Total particle formation rate
$K(d_{\rm p},d_{\rm p}')$	Coagulation coefficient between particles of $d_p$ and $d_p'$ sizes
k <sub>coag</sub>	Coagulation coefficient
k <sub>dil</sub>	Dilution coefficient specific for a specific chamber volume and total flow
$k_{\text{wall}}(d_{\text{p}},t)$	Wall-loss coefficient of particles with diameter $d_p$ at time t
N	Particle number concentration
N <sub>dp</sub>	Particle number concentration of particles with diameter $d_p$
N <sub>&gt;dp</sub>	Particle number concentration above a certain diameter
N <sup>±</sup> <sub><dp< sub=""></dp<></sub>	Charged particle number concentration of charged particles with smaller than $d_{ m p}$
N <sup>±</sup> <sub>dp - du</sub>	Charged particle number concentration in size bin of diameters $(d_p)$ and upper diameter $(d_u)$
N <sub>dpi</sub> – <sub>dpi+1</sub>	Particle number concentration in a size bin
n(d <sub>p</sub> ,t)	Number distribution of particles with diameter $d_{\rm p}$ at time t
N(t)	Number concentration of particles at time t
$Q(d_{\rm p},t)$	Source term for particle with diameter $d_p$ at time t
S(d <sub>p</sub> ,t)	Sink terms for particle with diameter $d_p$ at time $t$
S <sub>att</sub>	Production rate of ions by ion-neutral attachment
$S_{coag}$	Coagulation loss rate in the chamber
S <sub>dil</sub>	Dilution loss rate
Sgrowth	Loss rate of ions due to growth out of the size bin
S <sub>rec</sub>	Ion-ion recombination loss rate
S <sub>wall</sub>	Diffusional loss rate to the chamber walls
Т	Temperature
T <sub>ref</sub>	Reference temperature
t	Time
to	Time at the beginning of the experiment
t <sub>app,i</sub>	Time when concentration of size bin <i>i</i> starts to rise
t <sub>app50,i</sub>	Time when the concentration in size bin <i>i</i> reaches 50% of its maximum
t <sub>max</sub>	Time when the particle concentration reaches the maximum
t <sub>max,i</sub>	Time when the concentration in size bin <i>i</i> reaches the maximum
α	Ion-ion recombination coefficient
$\Delta t$	Time difference
χ	Ion-aerosol attachment coefficient

direct observation of clusters and freshly formed particles (<3 nm) is now possible. Here, we present a step-by-step procedure to calculate  $J_{dp}$  values and GRs from chamber experiment data. Following equipment setup and calibration, the procedure starts by determining particle GR, followed by calculating different loss corrections (dilution loss, wall loss and coagulation loss) needed for obtaining the final particle  $J_{dp}$ . The last step is to estimate the error on  $J_{dp}$  values and GRs. We present the required instruments and their operation for obtaining accurate data, explain the procedures for calculating  $J_{dp}$  and GR, and troubleshoot errors that might occur during chamber experiments or analysis.

#### Alternative methods

Generally, the main parameter describing the NPF intensity is the formation rate,  $J_{dp}$ , that is, the rate at which particles are formed per unit volume per unit time at a given particle diameter  $d_p$  (ref. <sup>7</sup>). Preferably,  $d_{\rm p}$  is in the range of 1.5–2 nm, which is close to the size of the critical cluster, that is, the smallest stable particle. In this case, the formation rate can be called the nucleation rate. In many previous laboratory studies, the ability of a given system to produce new particles was characterized by the rate at which new particles appeared, dN/dt, termed the apparent formation rate ( $J_{ap}$ ). As such, this method may be internally consistent in a single experiment, but the comparability of the results to different reactors or field data is limited, owing to different loss processes affecting the measured concentration. No unified method has been introduced in the literature to correct dN/dt for particle losses due to dilution and scavenging onto existing particle surfaces or reactor walls to obtain  $J_{dp}$  for chamber measurements. In addition, to obtain the nucleation rate,  $J_{dp}$  has often been measured at the size of 3 nm, or even larger, and then theoretically extrapolated to 1.5 nm<sup>115,116</sup>. Although theoretically sound, this method relies on several assumptions, such as a constant particle GR<sup>117</sup> and negligible self-coagulation<sup>118</sup>. Other parameters that have been used for describing the NPF intensity in chamber or flow tubes include total number concentration of particles generated and apparent formation rate<sup>57,110</sup>.

The variety of parameters and methods to quantify NPF makes it difficult to compare obtained results, because different methods can yield different results, even for the same dataset. Kulmala et al.<sup>114</sup> presented a standard procedure to analyze  $J_{dp}$  values and GRs from atmospheric data. However, to our knowledge, a similar protocol for chamber observations has been lacking until now.

#### Advantages and limitations

The main advantages of this protocol compared to previous approaches are that  $J_{dp}$  values and GRs are measured and analyzed directly at the size of the forming clusters and that the protocol is specific to chamber experiments.

The main limitation of this protocol is that, although the same basic principles apply for all experiments, the precise calculations need to be modified depending on the specific characteristics of the chamber in question and the instrumentation used. In addition, experimental data from chambers might not always be directly comparable to atmospheric observations because of the different aerosol and chemical reaction dynamics, missing vapor and aerosol constituents, or poor representation of atmospheric processes such as oxidation.

### Applications

The methods for analyzing NPF from atmospheric and chamber data differ because of differences in the dynamics and the spatial and temporal scales of NPF in the chamber as compared with the atmosphere<sup>118</sup>. In the atmosphere, particle formation often occurs over a large geographical area and particle concentrations can change because of processes other than nucleation, such as primary particle emissions and horizontal or vertical transport of particles<sup>119–122</sup>. In addition, the particle sinks are different. The main sink for newly formed particles in the atmosphere is coagulation onto the existing particle population<sup>123</sup>, whereas in chamber experiments, particle losses to chamber walls are normally much more important than coagulation<sup>118</sup> because the existing particle population is usually absent or small. Furthermore, in the atmosphere, the intensity of particle formation changes during the course of the day because the concentrations and properties of precursor vapors are constantly changing depending on environmental conditions such as solar radiation intensity. In a chamber, the production rate of precursor vapors can usually be kept constant for hours, and several variables characterizing NPF can be averaged over this period.

#### **Experimental design**

To quantify the particle dynamics during a controlled NPF chamber experiment, a set of instruments is needed to measure the size distribution of the particles from the initial cluster formation size (~1 nm) to the maximum size they reach during the particle growth process. Whereas measurements of particle number concentration and size distribution at small sizes (preferably <3 nm) are needed for calculating the  $J_{dp}$  value and initial GR, a measurement of the size distribution extending to larger sizes is required for the determination of coagulation sink and growth to larger sizes (Table 2). To quantify the gas phase concentrations of the vapors participating in NPF and their precursors, as

#### Table 2 | Recommendations for the measurement of relevant trace gas concentrations

Quantity to be measured	Importance
Total particle number concentration above ~1.5 nm	Required for determining $J_{dp}$
Particle size distribution covering the size range of ~1-1,000 nm (upper diameter depends on the experimental plan)	Required for determining the particle GRs, and coagulation and condensation sinks
lon size distribution	Required for studying of role of ions in particle dynamics
Trace gas concentrations	Required to quantify the precursor concentrations and to monitor chamber cleanliness
Concentration and composition of vapors and clusters directly participating in NPF	Required to understand the mechanism and intensity of the NPF process
Accurate measurement of temperature and RH	Required for determining the exact characteristics of the experiment

well as to ensure the cleanliness of the chamber, a measurement of relevant trace gas concentrations is required (Table 2).

#### **Requirements for particle measurements**

Two categories of particles should be measured: (i) particles that result from the NPF and growth process (e.g., precursors, oxidation products) and (ii) background particles that should be monitored to ensure the cleanliness of the chamber (to ensure that NPF is a consequence of the precursors and the processes associated with them).

The concentration measurement accuracy and cutoff diameter of the condensation particle counters (CPCs) should be verified.  $\pm 10\%$  accuracy or better is recommended for the concentration measurement, and users should adhere to a  $\pm 0.2$ -nm accuracy for the CPC cutoff diameter to ensure an accurate derivation of growth and nucleation rates. To minimize errors, we recommend that the cutoff diameter of the CPC be the same as the diameter at which the  $J_{dp}$  value is determined. The verification should be done by using a particle composition and concentration similar to those produced in the chamber experiments. See 'Instrument calibration' in the 'Equipment setup' section in the Materials for further discussion of accurate measurements of sub-3-nm particles.

Multiple size distribution measurement instruments are required to cover the whole size range of 1–1,000 nm (upper size diameter depends on experimental design). An overlap in the size range between instruments is important to ensure comparability. Laboratory calibrations, side-by-side comparisons and combined size distributions are essential to obtaining an agreement between the instruments at overlapping size ranges. Multi-instrument inversion routines are beneficial, if available.

Time resolution should be as high as possible, especially during periods of rapid particle formation and growth. For instance, in the case that the GR in the chamber is 60 nm hr<sup>-1</sup>, an instrument with 1-min time resolution is needed to capture the concentration at each  $d_p$ . The time resolution should be optimized considering the particle concentration, instrument sensitivity and statistics<sup>124</sup>. Table 3 summarizes the commercially available particle number concentration instruments and their respective time resolutions.

Total particle concentrations can vary from very low (<10 cm<sup>-3</sup>) to very high (>10<sup>6</sup> cm<sup>-3</sup>). For an accurate determination of  $J_{dp}$  and GR, the instruments need to measure the total particle concentrations and size distributions accurately over a wide concentration range. At low concentrations (<1,000 cm<sup>-3</sup>), the performances of both particle size magnifiers (PSMs) and condensation particle counter batteries (CPCbs) for particles <5 nm are better than those of electric mobility spectrometers (EMSs), which suffer from high particle losses and low charging probabilities. On the other hand, at high concentrations (>10<sup>6</sup> cm<sup>-3</sup>), the performance of EMSs is usually not affected, whereas the accuracy of CPCs in measuring total particle concentrations is affected by a coincidence in the optics<sup>125,126</sup>. Depending on the CPC type, a coincidence correction can be applied for concentrations up to ~10<sup>5</sup>-10<sup>6</sup> cm<sup>-3</sup>; above that, the CPC signal typically becomes saturated.

Particle losses, especially in the sub-10-nm size range, are high. This should be considered when designing the sampling lines and selecting instruments. See 'Equipment setup' in the Materials section for further discussion of particle loss prevention measures and corrections.

### NATURE PROTOCOLS

#### Table 3 | Instruments for measuring sub-3-nm particle formation and growth

Abbv.	Instrument	Size range	Type of particle measurement	Time resolution	Details	Reference
PSM	Particle size magnifier	>1 nm/1-3 nm	Total concentration/ size distribution	1 s/2 min (scanning)	DEG & butanol/ water	Vanhanen et al. <sup>173</sup>
DEG CPC <sup>a</sup>	Diethylene glycol-condensation particle counter	>1 nm	Total concentration	1 s	DEG & butanol/ water	Wimmer et al. <sup>176</sup> , Jiang et al. <sup>180</sup>
CPC <sup>a</sup>	Condensation particle counter	Typically >2.5 nm or >7 nm, depending on model	Total concentration	1 s	Butanol or water	Stolzenburg and McMurry <sup>174</sup> , Hering et al. <sup>201</sup>
CPCb	Condensation particle counter battery	1-10 nm, depending on CPCs	Total concentration/ size distribution	1 s	DEG, butanol and/or water	Kulmala et al. <sup>185</sup>
DMA train	Differential mobility analyzer-train	1.6-8 nm	Size distribution	10 s	DEG and butanol and/or water	Stolzenburg et al. <sup>179</sup>
SMPS <sup>a</sup>	Scanning mobility particle sizer	1.5-1,000 nm (smaller for each instrument type/model)	Size distribution	~1-5 min, depending on model and size range	DEG, butanol or water	Wang and Flagan <sup>177</sup>
NAIS	Neutral and air ion spectrometer	2.5-42 nm (total), 0.8-42 nm (charged)	Size distribution, ion size distribution	10 s	NA	Mirme and Mirme <sup>178</sup>
CIC	Cluster ion counter	<3 nm	lon size distribution	10 s	NA	
AIS	Air ion spectrometer	0.8-42 nm	lon size distribution	10 s	NA	Mirme et al. <sup>202</sup>

NA, not applicable. <sup>a</sup>Many different commercially available or in-house-built models/sub-types exist.

#### Requirements for measurements of gaseous species

Two categories of gaseous species should be measured: (i) compounds that participate in the NPF and growth process (e.g., precursors, oxidation products) and (ii) compounds that are monitored to ensure the cleanliness of the chamber (to ensure that NPF is a consequence of the precursors and the processes associated with them).

The concentrations of relevant trace gases that may be introduced into the chamber, such as SO<sub>2</sub>, NO<sub>x</sub> and O<sub>3</sub>, must be measured. SO<sub>2</sub> usually serves as a precursor for H<sub>2</sub>SO<sub>4</sub>, NO<sub>x</sub> affects the formation of oxidized organic vapors, and O<sub>3</sub> is the main oxidant of many unsaturated organic species and is also used to produce OH radicals through photolysis (with UV light of ~248-nm wavelength) or by reacting with alkenes (e.g., tetramethylethylene). In most chamber experiments, concentrations of these trace gases are well above 1 p.p.b. and can be accurately measured by trace-gas monitors, which typically have a detection limit of ~0.5 p.p.b. In experiments that aim to simulate pristine environments, very low concentrations of SO<sub>2</sub> and NO<sub>x</sub> are required and more advanced instruments are needed. Cavity-attenuated phase shift (CAPS)<sup>127</sup> can be used to measure NO<sub>2</sub> with a detection limit of 0.1 p.p.b. (3- $\sigma$ , 10-s time interval). The CLD 780 TR (chemiluminescence detector)<sup>128</sup> is able to measure NO with the lowest detection limit of 3 p.p.t. SO<sub>2</sub> can be also measured with a chemical ionization mass spectrometer (CIMS) using CO<sub>3</sub><sup>-</sup> as the primary ion<sup>98</sup>, which has a detection limit of 15 p.p.t.

Volatile organic compounds (VOCs) should be measured, and the most commonly used instrument for this purpose is a proton-transfer-reaction time-of-flight mass spectrometer (PTR-TOF)<sup>129-131</sup>. The detection limit of the PTR-TOF is a few p.p.t., much lower than the required VOC concentration in most chamber studies. Some VOCs oxidize to form HOMs (highly oxygenated molecules) and other low-volatility organic vapors, which are important precursors for NPF, whereas some are present in the chamber as contaminants and might not influence NPF.

 $H_2SO_4$ , which is generally regarded as the most important NPF precursor in the atmosphere, can be measured by a chemical ionization atmospheric-pressure-interface time-of-flight mass spectrometer (CI-APi-TOF) using NO<sub>3</sub><sup>--</sup> as the primary ion<sup>132</sup>. Calibrations need to be performed before and after the experiments<sup>133</sup>. HOMs can also be measured with the nitrate-ion-based CI-APi-TOF<sup>23</sup>. Because many different HOMs have very similar masses, the mass resolution of the instrument is critical. The detection limit of the CI-APi-TOF is ~1 p.p.q. We suggest a minimum mass resolution of 4,000 mass divided by mass difference ( $m/\Delta m$ ), but a higher mass resolution is always preferable.

### Table 4 | Chambers used for studying sub-3-nm particle formation and growth

		used in NPF analysis	
ess steel 26 m <sup>3</sup>	PSM (<2 nm), CPC (~3 nm), NAIS, SMPS	J <sub>1.7</sub> , GR	Kirkby et al. <sup>54</sup> , Almeida et al. <sup>74</sup> , Riccobono et al. <sup>77</sup> , Tröstl et al. <sup>85</sup> , Kirkby et al. <sup>98</sup> , Duplissy et al. <sup>96</sup> , Lehtipalo et al. <sup>102</sup> , Kürten et al. <sup>203</sup>
ilicate glass 1.45 m <sup>3</sup>	PSM, CPC (TSI 3025A, 3 nm), SMPS (15-600 nm)	J <sub>1</sub>	Dal Maso et al. <sup>86</sup>
1 4.2 m <sup>3</sup>	NAIS	J <sub>3,ap</sub>	Boulon et al. <sup>87</sup> , Wang et al. <sup>88</sup>
n 10 m <sup>3</sup>	DMA-TRAIN	GR	Pichelstorfer et al. <sup>89</sup>
n 27 m <sup>3</sup>	CPC battery, SMPS	J <sub>1.5</sub> <sup>a</sup>	Paulsen et al. <sup>93</sup> , Metzger et al. <sup>94</sup> , Riccobono et al. <sup>95</sup>
1	ilicate glass 1.45 m <sup>3</sup> n 4.2 m <sup>3</sup> n 10 m <sup>3</sup>	NAIS, SMPS NAIS, SMPS NAIS, SMPS NAIS, SMPS NAIS, SMPS NAIS, SMPS (15-600 nm) N 4.2 m <sup>3</sup> NAIS N 10 m <sup>3</sup> DMA-TRAIN	NAIS, SMPSilicate glass1.45 m³PSM, CPC (TSI 3025A, 3 nm), $J_1$ SMPS (15-600 nm)n4.2 m³NAISn10 m³DMA-TRAIN

Although it remains a challenge to perform direct calibrations for HOMs, it is reasonable to assume that the detection efficiency of HOMs is the same as that of  $H_2SO_4^{23}$ . In addition, because the mass-to-charge ratio (*m*/*z*) of HOMs varies widely over the mass spectrum, the mass-dependent transmission bias of the mass spectrometer needs to be corrected<sup>134</sup>.

Low-volatility or semi-volatile organic compounds, which probably do not directly nucleate but are likely to contribute to subsequent particle growth, should preferably also be monitored. These compounds can be measured with a high-resolution time-of-flight chemical ionization mass spectrometer (HR-tof CIMS) using, for example, iodide<sup>135</sup> or acetate<sup>136</sup> as reagent ions, or with a recently developed PTR<sup>137,138</sup>. As with HOM measurements, we suggest a minimum mass resolution of 4,000  $m/\Delta m$  be applied when using these instruments.

NH<sub>3</sub> and amines can be measured with an HR-tof CIMS using ethanol<sup>139</sup>, hydronium<sup>140</sup>, or nitrate<sup>141</sup> as reagent ions. Because these vapors are tend to stick to surfaces, a heated sample inlet is recommended. Masses of amine compounds are usually close to the masses of amide compounds, and a mass resolution of 4,000  $m/\Delta m$  is needed to separate them. Besides CIMS, ion chromatography is also a technique that can possibly be used to measure NH<sub>3</sub> and dimethylamine<sup>142</sup>. Both measurement techniques show a detection limit of a few p.p.t.v. for NH<sub>3</sub> and amines.

The chemical composition of charged embryonic clusters can be measured with an atmosphericpressure-interface time-of-flight mass spectrometer<sup>143</sup> (APi-TOF). This instrument is similar to the CI-APi-TOF, but does not use active chemical ionization (i.e., it measures naturally charged ion clusters). With a properly calibrated and tuned transmission efficiency, this instrument can measure the composition of charged clusters up to 3,000 Thomsons. This instrument can be operated in either a positive or negative mode. In the positive mode,  $NH_4^+$  is the most typical charge carrier<sup>144</sup>, whereas in negative mode,  $NO_3^-$  and  $HSO_4^-$  are the main charge carriers<sup>145,146</sup>.

In addition to the two main types of gases that are involved in the chemical reactions and NPF mechanism inside the chamber, a dilution tracer can be used to estimate the dilution lifetime of reactants and products within the chamber. The gases used for this purpose, such as argon or  $CO_2^{147-155}$  (although SF<sub>6</sub> was used in several early experiments, it has been subject to a worldwide ban since January 2006), should not be a byproduct of the reaction and should not stick to the wall. The dilution tracer is needed only when the chamber in use is not well mixed. If the chamber is well mixed and both total volume and flow are known, the dilution lifetime can be calculated using the inverse of Eq. 4.

#### Chamber considerations and requirements

Chamber characteristics vary widely, depending on the specific applications and building specifications (Table 4). The main varying features are size, irradiance, wall material, and temperature and

pressure control. Many of the currently available chambers are made of Teflon, Pyrex, quartz, aluminum or stainless steel<sup>46,88</sup>. Although Teflon chambers are expandable and non-sticky, they are not electrically conductive, causing efficient ion scavenging and preventing studies of atmospheric nucleation mechanisms, in which the particle charge might play a crucial role. Teflon has been found to be a sink for organic vapors and particles<sup>48,147,156</sup>. Yet this material can be used outdoors under 'real' natural radiation conditions or indoors under more controlled temperatures for studying particle growth and chemical mechanisms related to particle formation. Compared with Teflon chambers, stainless-steel reactors are more robust (i.e., leakproof) at various pressures and electrically conductive, so they are more suitable for studying ion processes. Similarly, glass chambers are robust and can be easily cleaned and operated at low-pressure conditions. However, the irradiance in such chambers cannot be controlled. Glass and Pyrex are typical materials for flow tubes<sup>110</sup>. Chamber walls are required to be as smooth as possible because rough surfaces have been found to enhance particle deposition<sup>109</sup>.

None of the reactor types are flawless; each suffers from unwanted features that may vary depending on the chamber size and material. For instance, common problems for all chambers are the wall losses of particles and gaseous species, inhomogeneity, difficulty in attaining vapor and particle concentrations similar to the atmosphere, and possible contamination issues. Most of these problems can be largely avoided by using large-volume chambers with a small surface-to-volume ratio. Thorough characterizations of chamber reactors have already been presented in the literature<sup>20,46,88,110,151,157-159</sup>, but the critical characteristics of chambers for studying atmospheric NPF are as follows:

- Minimal contamination. Materials of the chamber walls and sampling ports must be cleanable (e.g., by water, heating and/or ozone treatment) in order to remove all chemical compounds that may affect the particle formation process. Note that certain compounds, such as ammonia, stick to surfaces and may be re-released when thermodynamic conditions change. See 'Chamber setup and cleaning' in the 'Equipment setup' section of the Materials section for a discussion of and recommendations for chamber cleanliness. In addition, the gas injection system (including mass flow controllers, valves and so on) must not introduce any additional contaminants. The concentration levels of potential particle precursor vapors (e.g., ammonia, amines, sulfuric acid, organics, iodic acid) need to be monitored with instruments that have low enough detection limits (parts per trillion volume level) to detect possible contaminants.
- *Homogeneous mixing.* There should be no gradients or hotspots of precursor concentrations or vapor supersaturations in the chamber. Characterization of homogeneity can be done via sampling from multiple ports complemented with flow simulations. A suitable configuration of fans (location, fan speeds) is needed to mix the air, depending on the chamber geometry. There should be a balance between mixing and particles losses. If a laminar flow is not available, a counterflow fan should be used for mixing<sup>20,46,88,160,161</sup>.
- Equal charge distribution. All materials should be electrically conductive to avoid ion losses. In addition, a high-voltage clearing field can be used to filter away ions to study purely neutral processes.
- *Large chamber volume*. The chamber volume needs to be large, so that the particle residence time is long enough, because the chamber air is constantly diluted to compensate for the air consumed by the measurement instruments. Typical time scales of the particle formation and growth processes are from minutes to hours when using atmospheric concentration levels. Although a spherical configuration would have the highest volume-to-surface ratio (i.e., maximum volume for minimum surface), such a configuration is very difficult to operate, clean and illuminate and is also associated with inefficient mixing<sup>162</sup>. A cylindrical configuration is recommended instead. The choice of a chamber volume depends on the aim of the experiment, especially the required maximum size of the particles, as well as on precursor concentrations and thus the dilution lifetime.
- *Chamber flow rate.* The total volumetric flow rate required by all the sampling instruments should be considered when designing the chamber flow system and volume. Efficient sampling of sub-5-nm particles requires high sample transport flows to the instruments (see 'Equipment setup' in the Materials section), which increases the required total flow rate from the chamber. To maintain chamber purity, no backflows from the instruments should be allowed, so there should always be a small overpressure inside the chamber. Particle and vapor measurement instruments should be placed as close to the chamber as possible in order to minimize losses in inlet lines.
- *Chamber stability.* The chamber temperature, irradiance (and other external conditions), and gas flow rates must remain stable, so that experiments are reproducible. Particle formation rates ( $J_{dp}$  values) should not vary by >5% for at least 12 min (three consecutive full PSM scans, depending on instrument time resolution).

#### Flow tube design

The design of flow tubes can vary, but they usually consist of four sections: an inlet system, a mixing unit, a nucleation unit and an outlet where sampling takes place. The inlet system design depends on the precursor gas. In  $H_2SO_4$  nucleation experiments,  $H_2SO_4$  can be point-produced from a liquid solution and then injected into the flow tube or it can be generated in situ from  $SO_2$  gas. The point-production from liquid solution is achieved either by atomizing a liquid solution followed by vaporization using a furnace<sup>51,65,67,69</sup> or by flowing carrier gas over a temperature-controlled reservoir<sup>52,58,68,73,78,111</sup>. Alternatively, in situ generation of  $H_2SO_4$  involves production of OH radicals from photo-dissociation of  $H_2O$  vapor<sup>61,63,71,72,83</sup> or ozone photolysis by UV light<sup>59,60,62,64,67</sup>. The injection of other gases is usually achieved by using the same techniques, by applying several dilutions from concentrated gas bottles or by using permeation tubes.

In the mixing unit, the precursor gas (or gas mixture) is mixed with the carrier flow. The method by which the gas flow is introduced, together with the profile of the mixing, determines the distance needed for laminar flow to develop downstream of the mixing unit. For this purpose, simple plugtype inlets, perforated Teflon manifolds, showerhead inlets, spoke inlets, diffusers and transition cones are commonly used.

After mixing, nucleation should take place in the nucleation unit, which is kept at an approximately constant temperature using cooling jackets. The nucleation unit is typically made of glass, although stainless-steel and Teflon-lined stainless-steel units can also be used (Table 5). In studies in which photolysis is needed, the nucleation unit can be irradiated with UV light<sup>59,60,62,64</sup>; otherwise, gases such as ozone can be irradiated before entering the nucleation unit<sup>61,63,71,72,83</sup>.

Finally, particle-measuring instruments are connected to the outlet of the flow tube. The choice of the particle-measuring instrument can have substantial effects on the measured particle number concentration and therefore on calculated  $J_{dp}$  values. In nucleation experiments, we are interested in measuring sub-3-nm particles, so the instrument must be chosen accordingly. Nevertheless, the measured particle number concentration at the end of the flow tube also depends on the nucleation and growth processes taking place across the length of the flow tube. For these reasons, the measured formation rate at the end of the flow tube might not be equal to the actual formation rate, so we refer to the 'apparent formation rate' ( $J_{ap}$ ) when using flow tubes.

The calculation of formation rates in flow tubes  $(I_{ap} = \frac{N}{\Delta t})$  is similar to that for formation rates in chambers, yet it does not include the particle loss mechanisms (wall loss and coagulation loss), which are usually considered to be negligible in short reaction times. In the calculation, N is the number concentration of particles measured at the outlet of the flow tube and  $\Delta t$  is the characteristic time frame during which nucleation occurs<sup>110</sup>. If nucleation is homogeneous across the nucleation unit of the flow tube, then  $\Delta t$  is equal to the residence time. Consequently,  $\Delta t$  is simply calculated from the flow rate and volume of the nucleation unit. If nucleation is not homogeneous across the nucleation unit, a proper characterization (through experiments or flow simulations) of the flow tube must be performed to determine the real nucleation zone and the real nucleation time ( $\Delta t$ ). Under these circumstances,  $\Delta t$  can range from 10% to 60% of the residence time<sup>51,52,63,65,72</sup>.

Multiple factors can affect the nucleation homogeneity inside the nucleation unit, for example:

- Buoyancy and thermally driven convection are likely to take place at the entrance of the nucleation unit, and the flow regime might not be fully laminar.
- Possible temperature gradients between the mixing unit and the nucleation unit might induce undesirable particle production at the entrance of the nucleation unit.
- Across the flow tube, the gas mixture composition is likely to change because precursor gases have reacted away or have been lost to the walls, rendering nucleation negligible beyond a certain point.
- Point production of gases from liquid samples is known to produce a non-uniform gas profile in the flow reactor as compared with in situ production<sup>63,67,71</sup>.

Error estimates for the apparent nucleation rate must include errors associated with particlecounting instruments and inaccuracies in the determination of the nucleation time. Additional errors can be caused by neglecting wall losses and coagulation losses.

Wall losses in flow tubes have a more severe effect on precursor gases than on particles, so proper corrections must be applied to determine the exact precursor gas concentration at which nucleation takes place. Particle wall losses become important in the case of turbulent flows, high residence times or large surface-to-volume ratios of the flow tube (e.g., due to small inner diameters). Ideally, one can assume that the particle wall loss rate is of first order and determine it experimentally by introducing particles of a certain size into the flow tube and measuring their concentrations at different positions

PR		Т	$\frown$	C	$\frown$	
ΓК	U		U	C	U	L

Institute/collaboration	Nucleation unit material	Unit dimensions; length	Particle measurement	Calculated quantities	Reference
		(cm) × i.d. (cm)	instrument (cutoff)	used in NPF analysis	
Finnish Meteorological Institute	Stainless steel	200 × 6	UCPC TSI 3025A (2.18 nm)	J <sub>an</sub>	Brus et al. <sup>65</sup>
1			UCPC TSI 3025A (2.8 nm) & PSM	J ap	Brus et al. <sup>69</sup>
			PHA-UCPC (<2 nm) & PSM	J <sub>ap</sub>	Sipilä et al. <sup>67</sup>
	Stainless steel with inner Teflon coating	255 × 10.95	CNC TSI 3020	Jap	Viisanen et al. <sup>51</sup>
Kent State University	Pyrex	82 × 5.08 or 80 × 2.54	SMPS TSI 3936N76/UCPC TSI 3786	J <sub>ap</sub>	Benson et al. <sup>61</sup> Young et al. <sup>63</sup>
		85 × 12.8	UCPC TSI 3786 (2.3 nm)	Jap	Benson et al. <sup>71</sup> , Yu et al. <sup>72</sup>
Nanjing University of Information Teflon Science and Technology	Teflon	90 × 5.08	CPC TSI 3776 & PSM	J <sub>1,7</sub> a	Yu et al. <sup>83</sup>
Leibniz Institute for Tropospheric Research (TROPOS)	Teflon	449 × 8	UCPC TSI 3025 (3 nm)	J <sub>ap</sub>	Berndt et al. <sup>59</sup> , Berndt et al. <sup>60</sup> , Berndt et al. <sup>62</sup>
			UCPC TSI 3025 (3 nm), PHA- UCPC (1.5 nm) & PSM	J <sub>ap</sub>	Berndt et al. <sup>64</sup> , Sipilä et al. <sup>67</sup>
Texas A&M University	Pyrex	60 × 2.45	UCPC TSI 3025A (3 nm)	N>3	Zhang et al. <sup>58</sup> , Wang et al. <sup>68</sup>
University of Alabama in Huntsville (TANGENT - FT1)	Pyrex	80 × 4.85	PSM TSI SMPS (3080) & CPC (3776) (>3 nm)	Job	Benson et al. <sup>61</sup> , Young et al. <sup>63</sup> , Benson et al. <sup>21</sup> , Tiszenkel et al. <sup>204</sup> , Benson et al. <sup>205</sup> , Erupe et al. <sup>206</sup>
University of Delaware (TANGENT - FT2)	Quartz	152 × 20	SMPS (TSI 3938, 3788)	GR <sub>&gt;2</sub>	Tiszenkel et al. <sup>204</sup> , Krasnomowitz et al. <sup>207</sup> , Stangl et al. <sup>208</sup>
Augsburg College University of Minnesota	Glass	105 × 5	UFCNC UCPC	J <sub>ap</sub>	Ball et al. <sup>52</sup> Zollner et al. <sup>73</sup> , Glasoe et al. <sup>78</sup> Jen et al. <sup>76</sup> , Jen et al. <sup>81</sup>
University of California Irvine	Borosilicate glass	110 × 7.6	TSI SMPS (2.5 nm) <sup>c</sup> TSI SMPS (2.5 nm) & PSM	N>2.5 Jan	Chen et al. <sup>79</sup> , Chen et al. <sup>80</sup> Chen and Finlavson-Pitts <sup>84</sup>
Caltech	FEP Teflon	2 × 28	CPC (TSI 3025: 3 nm; TSI 3010: 6 nm), SEMS (25-700 nm)		Jimenez et al. <sup>57</sup>

#### Box 1 | Recommendations for flow tube experiments

- Residence time inside the flow tube must be high enough to allow for particle growth to diameters larger than the cutoff of the particle-measurement instruments, yet low enough that wall losses will not become important.
  Increasing the inner diameter of the flow tube helps reduce wall losses.
- Experiments should be started with the lowest concentrations and highest RH, and ended with the highest concentrations and lowest RH.
- We recommend keeping a continuous flow of nitrogen through the flow reactor when it is not used in experiments to minimize the exposure of the reactor to room air.
- In the case of point production of  $H_2SO_4$ , care must be taken to minimize  $H_2SO_4$  decomposition to  $SO_3$  in the liquid reservoir or downstream of it.
- In the case of point production of  $H_2SO_4$ , Teflon filters or glass frits should be used in order to remove any liquid residue or particulate impurities.

inside the tube<sup>163</sup>. Introducing a laminar sheath flow to prevent the sample flow from contacting the wall may be an option to reduce wall losses<sup>164</sup>, yet a precise control is needed in such complex designs.

Coagulation losses become important and cannot be neglected if (i) all particles form immediately  $(\sim t_0)$  and the residence time inside the flow tube is sufficient for particles to interact, (ii) the time scale of the coagulation process is comparable to that of the nucleation process, or (iii) the particle concentration is high. An upper limit of coagulation can be determined from the measured particle size distribution and by estimating coagulation rate constants as a function of particle size. Additional recommendations for operating flow tubes are presented in Box 1.

#### Method for determining particle GR

The particle GR is defined as the rate of change of the diameter,  $d_p$ , representing the growing particle population (see also Kulmala et al.<sup>114</sup>):

$$GR = \frac{dd_p}{dt}$$
(1)

Particle GRs can be determined by following the time evolution of the particle number size distribution during a particle formation event. This can be done by using different methods, including the log-normal distribution function method (which is not covered in this protocol because it is often unsuitable for chamber experiments, being that there are no distinct particle modes)<sup>114</sup>, the maximum concentration method (Step 2A; Lehtinen and Kulmala<sup>165</sup>), the appearance time method (Step 2B; Lehtipalo et al.<sup>166</sup>), and different general dynamics equation (GDE)-based methods (Step 2C; Pichelstorfer et al.<sup>89</sup>, Kuang et al.<sup>167</sup>).

For GDE-based methods, the time evolution of the aerosol number distribution  $n(d_{p},t)$  is described by the so-called GDE, which in its continuous form can be written as

$$\frac{\partial n(d_{\rm p},t)}{\partial t} = \frac{1}{2} \int_{0}^{d_{\rm p}} K\left(\sqrt[3]{d_{\rm p}^{3} - d_{\rm p}^{\prime 3}}, d_{\rm p}^{\prime}\right) n\left(\sqrt[3]{d_{\rm p}^{3} - d_{\rm p}^{\prime 3}}, t\right) n(d_{\rm p}^{\prime}, t) dd_{\rm p}^{\prime} - n(d_{\rm p}, t) \int_{0}^{\infty} K(d_{\rm p}, d_{\rm p}^{\prime}) n(d_{\rm p}^{\prime}, t) dd_{\rm p}^{\prime} - n(d_{\rm p}, t) \int_{0}^{\infty} K(d_{\rm p}, d_{\rm p}^{\prime}) n(d_{\rm p}^{\prime}, t) dd_{\rm p}^{\prime} + Q(d_{\rm p}, t) - S(d_{\rm p}, t) dd_{\rm p}^{\prime} + Q(d_{\rm p}, t) dd_{\rm p}^{\prime} + Q(d_{\rm$$

Here  $K(d_p, d'_p)$  is the coagulation coefficient between particles of diameters  $d_p$  and  $d'_p$ , and  $Q(d_p,t)$  and  $S(d_p,t)$  are the source and sink terms, respectively, for particles with diameter  $d_p$ . In a typical chamber experiment, the only source of particles is nucleation and the sink term arises from wall deposition.

In our typical problem setup, the time evolution of the number distribution function  $n(d_p,t)$  is measured and the coagulation coefficients  $K(d_p, d'_p)$  are sufficiently well predicted by theory. Lehtinen et al.<sup>168</sup> applied simple least squares based optimization to solve the unknown  $GR(d_p)$  and  $Q(d_p,t)$  for atmospheric field data measured in Hyytiälä, Finland, assuming size-independent growth and neglecting deposition. This method was later improved (with more processes and fewer assumptions) by Verheggen et al.<sup>169</sup> and Kuang et al.<sup>167</sup>. None of these methods, however, is suitable for rigorous estimation of errors in GR (or Q).

Different GR methods have been compared using measurement and simulation data<sup>89,170–172</sup>, and they have been found to agree reasonably well in most conditions. GR methods can be applied to data measured with different particle-sizing instruments, which enables determination of GR for different size ranges or comparison of GR values for same-size particles from different instruments. GRs

can usually be determined more accurately from chamber experiments than from atmospheric measurements because there is less fluctuation in the data, as well as more accurate particle size distribution measurements. Estimation of uncertainties in GRs is explained in Steps 8 and 9.

### Materials

### Reagents

#### For operation of instruments (depending on setup)

- 1-Butanol, reagent grade (VWR 20808.325) **! CAUTION** Butanol is flammable and corrosive. It causes skin irritation and serious eye damage and may cause respiratory irritation. While using butanol, wear protective gloves and eye protection, and work in a well-ventilated area. Keep it away from heat and ignition sources.
- Diethylene glycol, reagent grade (DEG; VWR 8.03131.5000) **! CAUTION** Diethylene glycol is harmful if swallowed and may be combustible at high temperatures. Wear protective gloves and eye protection while handling.
- Nitric acid (Fisher Chemical N/2300/PB17) **! CAUTION** Nitric acid is flammable, highly corrosive and toxic. It causes severe skin burns and eye damage. Wear protective gloves, protective clothing, eye protection, and face protection while using nitric acid. Use it only in a well-ventilated area away from heat, hot surfaces, sparks, open flames and other ignition sources.

### For conducting the experiment

- Synthetic air (or atmospherically relevant mixture of O2 and N2)
- Water (Milli-Q ultrapure<sup>66</sup>)
- Ozone
- Suitable precursor gases, depending on which chemical system will be studied (e.g., SO<sub>2</sub>, H<sub>2</sub>SO<sub>4</sub>, NO, NH<sub>3</sub>, amines, volatile organic compounds, iodine)

### Equipment

**CRITICAL** A comparison of instrument types is shown in Table 3.

- Condensation particle counters (CPCs) together with a particle size magnifier (PSM), commonly known as a condensation nuclei counter<sup>173–176</sup>. A particle counter with a cutoff size ~1.5 nm is required for a total particle concentration measurement. It is preferable to have multiple CPCs at different cutoff sizes or a PSM operated in a scanning mode to measure size distributions between ~1 and 3 nm
- Electric mobility spectrometers (EMSs). Multiple EMSs are required for measuring particle and ion size distributions between ~1 and 1,000 nm. These include some varieties of scanning mobility particle sizer (SMPS)/differential mobility particle sizer (DMPS), neutral cluster and air ion spectrometer (NAIS), air ion spectrometer (AIS), cluster ion counter (CIC) and differential mobility analyzer-train (DMA-train)<sup>177-180</sup>
- Mass spectrometers (MSs). Multiple mass spectrometers are required for quantifying the precursors of NPF and for determining the composition of growing clusters. Suitable instruments include APi-TOF, CI-APiTOF and other chemical ionization MSs, and PTR-TOF instruments<sup>132,137,143</sup>
- Gas-measuring instruments. Trace gas monitors are required for measuring the relevant trace gases (e.g., O<sub>3</sub>, NO, NO<sub>2</sub>, SO<sub>2</sub>, NH<sub>3</sub>)
- Instrumentation to accurately monitor chamber conditions, such as temperature, pressure, relative humidity (RH), dew point, flow rates and (solar) radiation

### Equipment setup

### Chamber setup and cleaning

Cleanliness of the chamber is a key requirement for studying atmospheric NPF, as even a few p.p.t.v. of trace gases (e.g., ammonia) can have a tremendous effect on nucleation rates<sup>54</sup>. We recommend taking the following precautions:

- All gas supplies should be as clean as possible. One way to provide clean gases is to have them in liquid form. During evaporation, some of the contaminant gases will be trapped in the liquid.
- All pipelines going from the gas supply to the chamber should be made of stainless steel. Avoid connectors and mass flow controllers that contain plastic material and polytetrafluoroethylene (Teflon).
- A dedicated line for each trace gas is preferable to avoid memory effects from previous gases.
- Electropolishing of the chamber's inner surfaces is recommended to enable better cleaning (in stainless-steel chambers).

#### Box 2 | Recommendations for ensuring accuracy of the total concentration measurement

There are several ways to ensure the accuracy of the total concentration measurement, which is critical in accurately determining  $J_{dp}$ . These are as follows:

- 1 Calibrate the CPC cutoff size using particles of the same composition and a concentration similar to that of the chamber.
- 2 Compare total particle concentration measurements to integrated particle concentrations from size distribution measurements.
- 3 Measure the total concentration at multiple cutoff sizes (using scanning PSM or CPCb) and size distributions using an EMS (e.g., NAIS or nanoSMPS, which are usually composition independent) and compare their development over time (appearance times); see Fig. 1.
- 4 Use only straight laminar flow particle transport lines, and experimentally characterize the size-dependent particle penetration.
- 5 Correct the transport losses for the measured size-resolved particle size distribution.
- Clean water vapor can be made using a Permapore-type interface between the clean air carrier and clean recirculating liquid water, to ensure no contamination from the water supply. Note that a few parts per trillion of impurities in liquid water results in a few parts per quadrillion in the gas phase.
- Cleaning of the chamber walls before and between experiments can be done by using synthetic clean air of high humidity at high temperatures. High ozone concentrations can be used to remove contaminants from chamber wall surfaces with UV light on. The exact cleaning procedure depends on the chamber material and the chemical compounds used in the experiments.
- The chamber should be leakproof and have a slight overpressure to avoid any vapor and particle contamination from room air.
- Run a background or zero measurement before initiating any NPF process inside the chamber; measurement of precursor vapors and background particles should be done before precursor injection to ensure a contamination-free environment.

### Instrument calibration

For accurate determination of  $J_{dp}$ , the concentration, sizing and cutoff diameter calibration of the instruments measuring the total particle concentration and size distribution are critical. The concentration response of particle counters should be verified against a reference instrument in the concentration range that is expected during the measurements. The cutoff diameter of a CPC must be verified, preferably with test particles that are similar to the particles produced during the chamber experiments, because the cutoff diameter is affected by the particle chemical composition and charging state<sup>181-185</sup>. The magnitude of the effect depends on the working liquid of the instrument, as well as on the particle composition.

During NPF, the size-resolved concentration of precursor vapors, molecular clusters and nanoparticles usually decreases rapidly with an increasing particle diameter. In the size range of 1–3 nm, this gradient is so high that a slightly erroneous CPC cutoff diameter causes a non-negligible error in the measured total particle concentration<sup>186</sup>, which further accumulates into errors in  $J_{dp}$  calculations. Furthermore, changes in external conditions, such as pressure or sample temperature and RH, can affect the cutoff size and instrument performance in general<sup>187</sup>. Therefore, we recommend that the instrument be calibrated before the start of the experiment under the same environmental conditions as used for the experiment itself, using the same chemical composition of particles to be characterized (Box 2 and Fig. 1).

The sampling lines should be as short as possible. For straight lines with tubular and laminar flow, the particle losses can be estimated using the Gormley–Kennedy equation<sup>188</sup>. At the sampling line outlet, a core sampling system is needed to reduce the sampling losses and to maximize the signal at the particle counters<sup>189</sup>. Using transport flows to reduce the time for diffusional particle losses during transport from the chamber to the instrument may decrease particles losses in the sampling lines. The transport lines should be designed carefully to maximize the particle transport efficiency and signal in the detectors<sup>190</sup>. Any bends, elbows, valves, or splits cause distortion of the parabolic flow profile and will enhance particle losses as compared to laminar tubular flow. If these cannot be avoided, the particle penetration in the sampling line must be experimentally characterized to correct the size-resolved particle concentrations.

Evaporation of particles during sampling may be an issue, depending on the particle composition, especially when measuring very small clusters. Changes in the carrier gas temperature or composition (e.g., due to dilution, drying or bringing a cold sample to room temperature) should be avoided by using thermally insulated sampling lines and chamber air as carrier gas.

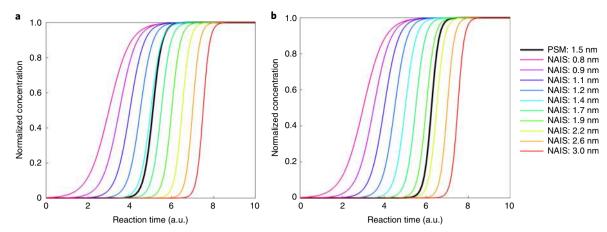


Fig. 1 | Simulated NPF experiment and instrument response. a,b, Comparison of the ion appearance times measured with the PSM at a nominal cutoff diameter of 1.5 nm (using the difference of total and neutral concentration from PSM measured without and with ion trap, respectively), and concentrations of ions measured with different size bins of the NAIS. The concentrations are normalized by the maximum concentration reached in the experiment (steady-state value). The PSM cutoff is determined correctly in **a**, whereas the cutoff is shifted (e.g., due to different particle composition compared to calibration) in **b**, and a correction is needed.

### Procedure

### Pre-injection measurement of chamber background

1 Start the NPF process by injecting the nucleating precursors and initiating their oxidation by injecting oxidant and/or illuminating the chamber with the required wavelength of radiation (Fig. 2).

▲ **CRITICAL STEP** Measurement of precursors, particles and other parameters should precede the initiation of the NPF process to ensure chamber cleanliness (see 'Equipment setup' in the Materials section).

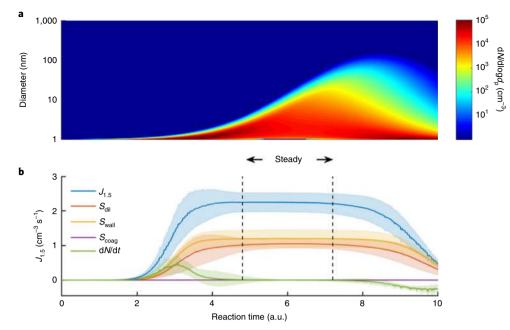
### ? TROUBLESHOOTING

### Determining the GR

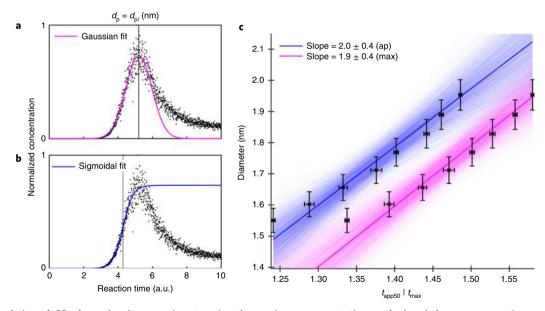
- 2 Particle GRs can be determined by either the maximum concentration method (option A), the appearance time method (option B) or the GDE method (option C). See the 'Experimental design' section for guidance on choosing one of these methods.
  - (A) Maximum concentration method
    - (i) Determine the times,  $t_{\max,i}$ , when the concentration in each size bin *i* reaches the maximum. To obtain an accurate  $t_{\max,i}$  value, fit the concentration time series with a Gaussian function. An example of applying this method to chamber experiment data is shown in Fig. 3.
    - (ii) Plot the mean diameters of the size bins,  $d_{p,mean,i}$ , as a function of the maximum times,  $t_{max,i}$ .
    - (iii) Apply a linear fit to the size range at which the GR is determined.
    - (iv) Obtain GR as a slope of the linear fit (Fig. 3c). The maximum concentration method can also be used to estimate a size-dependent GR. It must be noted, however, that the (t<sub>max,i</sub>, d<sub>p,mean,i</sub>) pairs would then correspond to different times, which means that a 'pure' size dependence cannot be determined.
      ? TROUBLESHOOTING

### (B) Appearance time method

(i) Determine the times,  $t_{app50,i}$ , when the concentration in each size bin *i* reaches 50% of the maximum concentration<sup>86,166,171</sup>. Unlike option A, this can be done by fitting a Sigmoid function to the concentration time series. Alternatively, one can determine the time,  $t_{app,i}$ , when the concentration in each size bin starts to rise. This can be done, for example, by determining the time when the concentration reaches 5% of the maximum concentration. It is also possible to determine  $t_{app50,i}$  and  $t_{app,i}$  from the total concentration measured with a CPC<sup>95</sup>, instead of using the concentration in a certain size bin. Lehtipalo et al.<sup>166</sup> compared different methods to determine appearance times and concluded that the most robust method is to either determine  $t_{app50,i}$  from size bin data or  $t_{app,i}$  from total



**Fig. 2 | Anticipated results from an NPF experiment performed in a chamber. a**, Simulated time evolution of particle size distribution during the experiment. **b**, Particle formation rate ( $J_{1,5}$ ) and its different components. The shaded areas correspond to  $\pm 1\sigma$  uncertainty, obtained from the Monte Carlo simulation of 10,000 runs. The time between the dashed lines shows the time of the stable formation rate of particles (steady state), for which the average  $J_{dp}$  value should be calculated. The magnitude of the components and time scales varies depending on the chamber specifications, experimental plan (e.g., gas concentrations) and  $J_{dp}$  value and GR (affecting the particle size distribution).



**Fig. 3** | **Calculation of GRs from chamber experiments using the maximum concentration method and the appearance time method. a**, The concentration in a size bin is normalized by dividing it by the maximum concentration reached during the experiment and then fitting using a Gaussian fit. The same process is repeated for all the size bins for which a GR is calculated. c, The times corresponding to maximum concentration are then plotted as diameter versus time  $(t_{max})$ , as shown in magenta. The x-axis uncertainty is the ±16 fit uncertainty from the Monte Carlo simulation of 10,000 runs, and the y-axis uncertainty is the estimated instrumental sizing uncertainty. GR is obtained as the slope of the linear fit to the  $d_p$  versus  $t_{max}$  data; GR = 1.9 ± 0.4 nm/h. The GR uncertainty is ±10 from the Monte Carlo simulation. b, the concentration in a size bin is normalized by dividing it by the maximum concentration reached during the experiment and then fitted using a sigmoidal fit. The same is repeated for all the size bins for which a GR is calculated. c, The midpoints of the fits are then plotted as diameter versus time  $(t_{app50})$ , as shown in blue. GR is obtained as the slope of the linear fit to the  $d_p$  versus  $t_{app50}$  data; GR = 2.0 ± 0.4 nm/h. Note that the maximum concentration method gives the GR at a later time step during the experiment, so particle size distribution and gas concentrations in the chamber might have changed. app50, 50% appearance time; max, maximum.

concentration data. An example of  $t_{app50,i}$  determined from the size bin data is shown in Fig. 3.

- (ii) Plot the mean diameters of the size bins, d<sub>p,mean,i</sub>, as a function of the appearance time t<sub>app50,i</sub> or t<sub>app,i</sub>.
- (iii) Apply a linear fit to the size range at which the GR is determined.
- (iv) Obtain GR as a slope of the linear fit (Fig. 3). Note that the GR value might change with size, especially during the beginning of the growth process<sup>85</sup>, in which case using a linear fit is a good assumption only in a narrow size range. It is also possible to fit a higher-order polynomial to the data points and obtain GR as a derivative of the curve.
   ? TROUBLESHOOTING
- (C) General dynamics equation method
  - (i) Calculate the optimal match between the measured data and the solution to the GDE (see equation in the 'Experimental design' section). To solve the GDE, which is a partial-differential-integral equation, the continuous function  $n(d_{pt}t)$  is typically approximated by a histogram, that is, by dividing the continuous-size spectrum into finite intervals (or bins), resulting in a set of ordinary differential equations for the bin concentrations N(t). These ordinary differential equations can then be solved by standard numerical time integration routines.
  - (ii) Find the growth rate  $GR(d_p)$  and source rate  $Q(d_p,t)$  corresponding to the optimal match. This can be done using different approaches. Our suggestion is to use the method by Pichelstorfer et al.<sup>89</sup>, in which the GDE is fitted to the measured size distribution step by step. Starting from a measured size distribution at any time step, all the other aerosol dynamical processes are simulated first, on the basis of measured conditions and theory; then  $GR(d_p)$  is estimated by moving the distribution in size space in an optimal way to match the measured distribution at the next time step. In this way, by marching step by step in time, the dependence of GR is on both time and size can be estimated.

### ? TROUBLESHOOTING

### Determination of particle losses

3 Determine dilution losses. If the chamber is operated in continuous mode, synthetic clean air should be continuously flowing into the chamber, and the instruments should be continuously sampling from the chamber. This leads to an artificially lower particle concentration in the chamber that is due to dilution, which needs to be corrected for when calculating  $J_{dp}$  values. The calculation for the loss rate of particles due to dilution is as follows:

$$S_{\rm dil} = N_{>dp} \cdot k_{\rm dil} [\rm cm^{-3} \, s^{-1}]$$
 (3)

with 
$$k_{\rm dil}[s^{-1}] = \frac{\rm Flow_{synthetic\,air}}{V_{\rm chamber}}$$
 (4)

 $N_{>dp}$  is the total particle concentration above the size at which you want to calculate the  $J_{dp}$  value,  $k_{dil}$  is the dilution rate, Flow<sub>synthetic air</sub> is the flow rate of clean air, and  $V_{chamber}$  is the volume of the chamber.

4 Determine wall losses. Diffusional losses of particles to the chamber walls ( $S_{wall}$ ) can be determined empirically by observing the decay of the concentration of a specific compound having a known diameter (e.g., decay of sulfuric acid monomer concentration after its photochemical production has been stopped by turning off the UV lights). The obtained loss rate coefficient is inversely proportional to the mobility diameter in a size range <100 nm, where diffusional losses are the most critical<sup>191</sup>, and can therefore be scaled and applied to correct for the losses of different-sized particles when calculating  $J_{dp}$  values. See also Schwantes et al.<sup>110</sup> and references therein. The calculation for the wall loss rate is as follows:

$$S_{\text{wall}}(T) = \sum_{i} N_{d\text{pi}-d\text{pi}+1} \cdot k_{\text{wall}}(d_{\text{p}}, T) [\text{cm}^{-3}\text{s}^{-1}]$$
(5)

 $N(d_{\rm p})$  is the number concentration of particles with a mobility diameter  $d_{\rm p}$  from the size distributions, and  $k_{\rm wall}$  is an experimentally determined factor based on mixing, chamber conditions and dark decay of the reference species in the absence of particles.

In previous chamber experiments<sup>97,102</sup>, the value of  $k_{\text{wall}}$  was determined from the theoretical temperature dependence of the diffusion coefficient,  $D \sim (T/T_{\text{ref}})^{1.75 \ 192}$  and the wall loss dependence

on the diffusion coefficient,  $k_{wall} \sim (D)^{0.5}$ ; see also McMurry and Grosjean<sup>108</sup> and McMurry and Rader<sup>159</sup>.

The calculation for this is as follows:

$$k_{\text{wall}}(d_{\text{p}}, T) = F \cdot \left(\frac{T}{T_{\text{ref}}}\right)^{0.875} \cdot \left(\frac{d_{\text{p,ref}}}{d_{\text{p}}}\right) [\text{s}^{-1}].$$
(6)

Here, *F* is an experimentally determined factor based on mixing, chamber conditions and dark decay of the reference species in the absence of particles,  $d_{p,ref}$  is the mobility diameter of the reference species,  $T_{ref}$  is the reference temperature at which the experimental loss rate was determined, and *T* is the studied chamber temperature. In this protocol, we recommend that this equation be used only when the mean diameter of the particle number size distribution is smaller than ~100 nm<sup>159</sup>.

5 Determine coagulation sink ( $S_{coag}$ ), which describes the loss rate of particles due to coagulation onto a pre-existing particle population. Calculate the coagulation sink for the particle size at which you want to calculate the  $J_{dp}$  value, using the measured particle number size distribution. The calculation for determining coagulation sink is as follows<sup>193</sup>:

$$S_{\text{coag}}(d_{\text{p}}) = \int k_{\text{coag}}(d_{\text{p}}, d'_{\text{p}}) n(d'_{\text{p}}) dd'_{\text{p}} \cong$$

$$\sum_{d'_{\text{p}}=d_{\text{p}}}^{d'_{\text{p}}=\max} k_{\text{coag}}(d_{\text{p}}, d'_{\text{p}}) N_{d'_{\text{p}}}[\text{cm}^{-3}\text{s}^{-1}]$$
(7)

 $k_{\text{coag}}(d_{p}, d_{p}')$  is the Brownian coagulation coefficient for particles sizes  $d_{p}$  and  $d_{p}'$ . It is usually calculated by using the Fuchs interpolation between continuum and free-molecule regimes. For chamber experiments, the coagulation sink is often negligible at the start of the experiment, but increases as the particles grow and more particles are formed in the chamber (Fig. 2b).

### Determining $J_{dp}$ value

6 Determine the total particle formation rate,  $J_{dp}$ , defined as the net flux of new particles into the measurable size range, that is, across the lower detection limit ( $d_p$ ) of the particle counter used. By integrating the GDE from the instrument detection limit up to infinity, we obtain the following balanced equation for calculating the total number concentration *N*:

$$\frac{dN}{dt} = J_{\rm dp} - S_{\rm dil} - S_{\rm wall} - S_{\rm coag} \tag{8}$$

By rearranging the terms in this equation, we can solve for  $J_{dp}^{114}$ , as shown in the equation below.

$$J_{\rm dp} = \frac{dN}{dt} + S_{\rm dil} + S_{\rm wall} + S_{\rm coag} \left[ \rm cm^{-3} s^{-1} \right]$$
(9)

dN/dt is the time derivative of the total particle concentration above a certain threshold (preferably close to 1.5 nm) and  $S_{dil}$ ,  $S_{wall}$ , and  $S_{coag}$  are the loss rate of particles, described in Steps 3–5.

For chamber experiments, the formation rate can be calculated from changes in the total particle number concentration measured with a PSM or some other CPC, because nucleation is the only source of particles. The  $d_p$  value at which the formation rate is determined depends on the cutoff size of the instrument, which is assumed to be a step function. For atmospheric data, a certain size bin close to nucleation size is used instead of the total particle concentration<sup>114</sup> in order to eliminate the effect of other particle sources or particle transport.

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7 (Optional) Determine the ion nucleation rate using an ion trap in front of a particle-measuring instrument. The ion nucleation rate  $(J^{\pm})$ , which is the formation rate of naturally charged ions only, can be determined using two particle-measuring instruments (e.g., CPC or PSM) in parallel, one of which is equipped with an ion trap to remove all ions<sup>97</sup>. The total  $J_{dp}$  value is then calculated from the instrument without the ion trap (see Steps 3–6). The neutral fraction of the total particle formation rate ( $J_{n,tot}$ ), which is different from the neutral formation rate ( $J_n$ ) introduced later in Step 13, is calculated using the instrument with the ion trap (see Steps 1–6). Note that in this case the neutral particle formation rate includes the particles that are formed by recombination of ions

and are therefore detected as neutral particles ( $J_{n,tot} = J_n + J_{rec}$ ). The ion formation rate (sum of both polarities) is the difference between  $J_{tot}$  and  $J_{n,tot}$ 

$$J^{\pm} = J_{\text{tot}} - J_{\text{n,tot}} \tag{10}$$

### Error estimation

- 8 Determine the error in the GR. When using the appearance time and maximum concentration methods, there are two sources of uncertainty: the fits used for determining the appearance times or maximum concentration times, and the particle diameter that is obtained from the instrument. If one of them is clearly greater than the other, apply weighted least square fit using the variable representing the smaller error as an explanatory variable (option A). If the two variables have errors of similar magnitude, a fitting method allowing for error in both variables (e.g., total least squares or geometric mean regression) should be used (option B).
  - (A) Weighted least square fit using the variable representing a smaller error as an explanatory variable
    - (i) Plot size-classified particle concentrations as a function of time and retrieve the appearance time for a given particle size at a time when the concentration is 50% of the maximum concentration<sup>166</sup>.
    - (ii) After retrieving the size-dependent appearance times, apply a linear regression model, for example, a weighted least square fit, to appearance time versus particle diameter data (for other options, see Mikkonen et al.<sup>194</sup>), and calculate both the GR and error estimate directly based on the fit and fit uncertainty.

### (B) Total least squares or geometric mean regression

**CRITICAL STEP** The error in GR is determined using Monte Carlo simulation (or some other numerical method<sup>195,196</sup>).

- (i) Assume normally distributed uncertainty that includes random and systematic errors in the measured particle concentration and particle diameter on the basis of instrument performance and reproduce 10,000 datasets for which the values are randomly picked from the estimated distribution around each data point.
- (ii) Calculate the GR for all reproduced datasets.
- (iii) Obtain the GR as the median value with uncertainty as  $\pm 1$  s.d. For the example calculations, see the Anticipated results section and Fig. 3.
- 9 Report the value of GR with one significant figure if the error is >20% and with two significant figures if the error is <20%.
- 10 Determine the error in the formation rate using the Monte Carlo method. Reproduce the formation rate 10,000 times at the plateau value, from which formation rate is normally determined. First, as the detected particle number concentration above a given cutoff diameter depends on the cutoff diameter ('Equipment setup' section), estimate the relation between the cutoff diameter and detected particle concentration. Thereafter, assume independent uncertainties for the cutoff, N,  $k_{dij}$ ,  $k_{\text{wall}}$ , and  $k_{\text{coag}}$ , which are normally distributed and include both random and systematic errors. Typically, in chamber experiments, the uncertainty in  $k_{\rm dil}$  can be estimated from the dilution flow rate, the uncertainty in  $k_{\text{wall}}$  can be estimated from a decay experiment to which the decay rate can be fitted, and the uncertainty in  $k_{\text{coag}}$  is estimated assuming 10% error in the size distribution. The Monte Carlo run should be constructed so that the first cutoff diameter is selected from the cutoff distribution, which determines N, and for that N the uncertainty is normally distributed and selected randomly. If the size distribution is obtained from two or more instruments, the uncertainties in cutoff and N should be estimated for each instrument. Then run a Monte Carlo simulation and obtain  $J_{dp}$  as the median value with uncertainty as ±1 s.d. For the example calculations, see the Anticipated results and Fig. 2b.
- 11 After determining the error in the formation rate, report the value of the formation rate with one significant figure if the error is >20% and two significant figures if the error is <20%.
- 12 (Optional) Determine the importance of charge in NPF by comparing the ion formation rate (determined from measured ion size distributions<sup>114</sup>) to the total  $J_{dp}$  value. When calculating the formation rate of charged particles, additional terms need to be added to Eq. 9 in order to account for the loss of ions due to ion-ion recombination ( $S_{rec}$ ) and the production of ions by the charging of neutral particles ( $S_{att}$ )<sup>197</sup>. As calculating recombination and charging between all sizes is rather complicated, we recommend calculating charged formation rates from ion size distribution in a size

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bin between diameters  $d_p$  and the upper diameter  $d_u$ . For this reason, a term describing the growth of ions out of the studied size bin ( $S_{growth}$ ) also needs to be included when calculating the ion formation rate. The equation for calculating the ion formation rate for positive (superscript +) and negative ions (superscript -) is as follows:

$$J_{\rm dp}^{\pm} = \frac{dN_{\rm dp-du}^{2}}{dt} + S_{\rm dil} + S_{\rm wall} + S_{\rm growth} + S_{\rm coag} + S_{\rm rec} - S_{att}[\rm cm^{-3} s^{-1}]$$
(11)

 $\frac{dN_{dp-du}^{z}}{dt}$  is the time derivative of the ion concentration in a certain size bin, measured with a NAIS or some other ion instrument. The terms describing the loss of ions due to dilution ( $S_{dil}$ ), deposition on chamber walls ( $S_{wall}$ ) and coagulation ( $S_{coag}$ ) are calculated as described in Steps 3–5, but instead of calculating them for all the particles larger than a certain threshold size, they are calculated for ions in a size bin between  $d_p$  and  $d_u$ .

The equation for calculating the loss of ions due to growth out of the studied size bin is as follows:

$$S_{\text{growth}} = \frac{N}{(d_{\text{u}} - d_{\text{p}})} \times \text{GR}$$
(12)

where GR is the growth rate of ions out of the size bin, which can be determined from the ion size distribution (see Step 1).

The equation for calculating the loss rate of ions due to ion-ion recombination is as follows:

$$S_{\rm rec} = \alpha N_{dp-du}^{\pm} N_{$$

 $\alpha$  is the ion-ion recombination coefficient, for which the constant value of  $1.6 \times 10^{-6}$  cm<sup>3</sup> s<sup>-1</sup> is usually assumed<sup>198</sup>. Note that the recombination coefficient depends on both the size and the chemical composition of ions, as well as on environmental conditions such as temperature and RH<sup>199</sup>.

The equation for calculating the production rate of ions due to charging of neutral particles is as follows:

$$S_{\text{att}=\chi} N_{dp-du} N_{$$

 $\chi$  is the ion–aerosol attachment coefficient which, similar to the recombination coefficient, depends on the particle size and environmental conditions.  $\chi$  is usually assumed to equal  $0.01 \times 10^{-6}$  cm<sup>3</sup> s<sup>-1</sup> (ref. <sup>200</sup>).

- 13 (Optional) If a high-voltage field is available inside the chamber to remove all ions, calculate the contribution of charged species to the total particle number concentration. When the high-voltage field is off, the total  $J_{dp}$  value consists of both neutral and charged particles ( $J_{tot} = J_n + J^{\pm} + J_{rec}$ ). The neutral formation rate ( $J_n$ ) can be determined from the time when the high-voltage field is on and all particles are produced by neutral processes. The formation rate resulting from charged particles ( $J^{\pm} + J_{rec}$ ) involves the naturally charged ions and those that have lost their charge due to recombination. These two contributions cannot be distinguished using this method. The importance of ion processes can thus be assessed by comparing the total formation rates in these two cases. This method was introduced by Kirkby et al.<sup>54</sup> for the CLOUD chamber.
- 14 (Optional) If a high-voltage field is available and you are using an ion trap, calculate the formation rate of neutral particles due to recombination,  $J_{rec}$ , using the following equation:

$$J_{\rm rec} = J_{\rm tot} - J_{\rm n} - J^{\pm} \tag{15}$$

The total particle formation rate  $(J_{tot})$  consists of both neutral and charged parts and is calculated with the high-voltage field turned off, whereas the neutral formation rate  $(J_n)$  can be determined from the time when the high-voltage field is on and all particles are produced by neutral processes (see Steps 6 and 7).  $J^{\pm}$  is the ion formation rate, which is the difference between  $J_{tot}$  and  $J_{n,tot}$  acquired by using an ion trap in front of the particle measurement instrument (Step 7; Eq. 10).

### Troubleshooting

Troubleshooting advice can be found in Table 6.

Table 6	Troubleshooting table	
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Step	Problem	Solution
1	No NPF is observed	Check that precursor gas injection works and the gas concentration in the chamber is what you expect (when a new gas is introduced into the chamber, there could be a delay in the increase of the gas phase concentration due to saturation of walls and inlet lines)
		Check the performance of the particle-counting instruments. If everything works, the precursor vapor concentration might be too low for NPF
	Unexpected particle formation/particle burst is detected	Check that chamber mixing works (avoid introducing reactants through a shared inlet) and that there are no fluctuations in any of the flows or chamber conditions
		Check for contamination from the gas injection system, chamber wall or backflow from the sampling instruments
		Check that the precursor concentration is what you expect
		Check that there are no leaks in your instruments
1,6	Particle formation stops unexpectedly	Check that there is no interruption or drop in the gas injection and that the measured gas concentrations are stable
		Check that the measurement instruments are working correctly
		The sink from growing particles might be too high to completely suppress particle formation and growth
		A drop in the $J_{dp}$ value can sometimes be seen at the start of an experiment when the gas and particle concentrations first peak, and then a steady-state value is reached when production and losses stabilize
1,2(A,B,C),6	The size distribution plot does not look like a 'banana' or multiple 'bananas' are detected	Owing to different dynamics (see Introduction), size distribution during particle formation might look different in the chamber than in the atmosphere. Particle formation usually continues as long as the precursor concentration is high enough and the particle sink is low enough, so there is a threshold at which the first-formed particles grow, and continuous production and growth of particles occurs thereafter. Multiple bananas can be observed if there is fluctuation in conditions or the sink becomes lower during the experiment
2(A,B,C)	Particles form, but they do not grow past a few nanometers	There could be an insufficient amount of precursor vapors capable of growing the particles, or the growth is very slow in comparison to the loss rates and dilution lifetime of the particles in the chamber
		Calculate the expected GR on the basis of vapor concentrations
		Check that the instruments measuring growth (e.g., SMPSs) are able to detect the growing particles (losses in mobility spectrometers are often very high for the smallest particles and, if the particle concentration in the chamber is low, the growing particles might not be detected at all)
2(A,B)	Particle GR determined using maximum concentration or appearance time method appears to be negative	It is possible that the slope of the linear fit to $d_{p,mean,i}$ versus $t_{max,i}/t_{app50,i}/t_{app,i}$ data is negative. This can be caused by very fast particle growth, in which case the time difference between different sizes is too small to be detected. Try a different method or a wider size range. Negative GR can also be caused by errors in the measured particle size distribution
2,6	Results (e.g., J or GR at a certain gas concentration) do not match the values reported in previous studies (within error estimates)	Check the performance of the instruments (both precursor gas and particle measurements) and their calibration. If <i>J</i> or GR is lower than expected, there could be losses that are not accounted for (see 'Equipment setup' in the Materials section for inlet line losses) or the gas concentration may have been overestimated
		Check for contamination or some other unaccounted for precursor vapor; the particle losses could be overestimated or the gas concentration could be underestimated
6	J is not increasing with increased precursor vapor concentration	Check that the particle-counting instruments are detecting particles correctly (e.g., maximum detectable concentration is not exceeded, there is enough working fluid).
		If no technical problems are detected, J might be saturated with respect to this variable or not affected by it

### Timing

Step 1, pre-injection measurement of chamber background: 30 min; injection of precursors: 5 min Step 2, determination of the GR: 12 min to several hours, depending on the amount of precursors, the GR, the aimed-for maximum diameter and the chamber dilution lifetime. Typical GRs at atmospheric conditions and concentrations vary from <1 nm/h to several tens of nanometers per hour. Steps 3–5, determination of particle losses: 12 min to several hours, depending on the concentration of precursors that results in the formation of particles in the size range of interest Step 6, determination of formation rate: 12 min to several hours; see timing for Steps 1–5 Step 7, determination of ion nucleation rate: 12 min to several hours; see timing for Steps 1–5 Steps 8 and 9, determination of error on GR: 12 min to several hours; see timing for Step 2 Steps 10 and 11, determination of error on formation rate: 12 min to several hours; see timing for Step 2 Steps 1–5

Steps 12-14, determination of ion nucleation rate (alternative methods): 12 min to several hours; see timing for Steps 1-5

### Anticipated results

An example of an NPF experiment is shown in Fig. 2a. A surface plot exhibiting a banana-shaped NPF event represents the anticipated time evolution of particle size distribution in an experiment in which particles constantly form until the gas injection or oxidation is stopped and the gases and particles are left to decay. To determine  $J_{dp}$ , the experiment should be stopped only when a steady-state value is reached and kept for a long enough time to allow for averaging over a suitable period (e.g., 12 min, depending on the instrument's time resolutions); see Fig. 3b. For measuring GR, the experiment should be continued until the particle size distribution has reached the sizes of interest. If the experiment is stopped too early, particles will not reach large enough sizes to allow for the calculation of growth. An example of the calculation of particle formation is shown in Fig. 2b. The formation rate usually increases rapidly in the beginning of the experiment when gas concentrations in the chamber increase as the experiment is started (usually by turning on UV lights or by starting gas injection). To obtain a single formation rate value per experiment (representing a specific set of conditions), the formation rate can be averaged over the period with steady chamber conditions (steady state), demonstrated by 'steady' in Fig. 2. In our example, the anticipated GR =  $2.0 \pm 0.4$  nm h<sup>-1</sup>, and the anticipated  $J_{1.5} = 2.0 + 0.4$  cm<sup>-3</sup> s<sup>-1</sup>, as shown in Figs. 2b and 3c, respectively.

#### **Reporting Summary**

Further information on research design is available in the Nature Research Reporting Summary linked to this article.

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#### Author contributions

L.D., K. Lehtipalo, J. Kontkanen, T.N., K. Lehtinen, V.-M.K. and M.K. contributed to the development of the technique for calculating  $J_{dp}$  and GR. R.B., L.A., J.D., T.P., C.Y., B.C. and J. Kangasluoma contributed to development of the technique for calibrating and minimizing losses during particle measurement. All authors contributed to the writing of this protocol and to the scientific discussions related to it.

#### **Competing interests**

The authors declare no competing interests.

### Additional information

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