



CHORUS

This is the accepted manuscript made available via CHORUS. The article has been published as:

Four-phonon scattering significantly reduces intrinsic thermal conductivity of solids

Tianli Feng, Lucas Lindsay, and Xiulin Ruan

Phys. Rev. B **96**, 161201 — Published 27 October 2017

DOI: [10.1103/PhysRevB.96.161201](https://doi.org/10.1103/PhysRevB.96.161201)

Four-phonon scattering significantly reduces intrinsic thermal conductivity of solids

Tianli Feng,¹ Lucas Lindsay,² and Xiulin Ruan^{1,*}

¹*School of Mechanical Engineering and the Birck Nanotechnology Center,
Purdue University, West Lafayette, Indiana 47907-2088, USA*

²*Materials Science and Technology Division, Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831, USA*
(Dated: October 9, 2017)

For decades the three-phonon scattering process has been considered to govern thermal transport in solids, while the role of higher-order four-phonon scattering has been persistently unclear and so ignored. However, recent quantitative calculations of three-phonon scattering have often shown significant overestimation of thermal conductivity as compared to experimental values. In this work we show that four-phonon scattering is generally important in solids and can remedy such discrepancies. For silicon and diamond, the predicted thermal conductivity is reduced by 30% at 1,000 K after including four-phonon scattering, bringing prediction in excellent agreement with measurements. For the projected ultrahigh-thermal conductivity material, zincblende BAs, a competitor of diamond as a heat sink material, four-phonon scattering is found to be strikingly strong as three-phonon processes have an extremely limited phase space for scattering. The four-phonon scattering reduces the predicted thermal conductivity from 2,200 W/m-K to 1,400 W/m-K at room temperature. The reduction at 1,000 K is 60%. We also find that optical phonon scattering rates are largely affected, being important in the applications such as phonon bottleneck in equilibrating electronic excitations. Recognizing that four-phonon scattering is expensive to calculate, in the end we provide some guidelines on how to quickly assess the significance of four-phonon scattering, based on energy surface anharmonicity and the scattering phase space. Our work clears the decades-long fundamental question of the significance of higher-order scattering, and points out new ways to improve thermoelectrics, thermal barrier coatings, nuclear materials, and radiative heat transfer.

Phonons are quasi-particles that quantize lattice vibrations. They interact with each other through scattering events with the most significant scattering processes known as three-phonon processes. Recently, thermal conductivity (κ) calculations based on density functional theory have found incredible agreement with measured κ values for a variety of systems and made striking predictions of thermal processes despite considering only lowest order perturbative intrinsic scattering from three phonon interactions[1–4]. However, a persistent fundamental question for decades has been: what is the impact of four-phonon and higher-order scattering? Due to the lack of formalism and computational power, four-phonon scattering has been ignored in previous studies. However, this question has become particularly relevant, since first principles methods have overestimated the measured thermal conductivities of a number of materials [2, 3, 5–7]. For example, while some predictions give reasonable accuracy with measured data at low temperature, they overpredict significantly at higher temperature, diminishing the predictive power for applications such as thermal barrier coatings and high-temperature thermoelectrics [2]. Moreover, such deviations can become quite large even at room temperature (RT) for some technologically important materials such as *c*-BN [3, 5] and SnSe [6, 7]. As an example for thermal management applications, Lindsay *et al.* have predicted that the zincblende structure, boron arsenide (BAs), may have a thermal conductivity $\sim 2,200$ W/m-K at room temperature, higher than the known best heat conductor diamond [3]. This might open new opportunities for passive cooling and other thermal management applications. The BAs system has been synthesized, however the measured thermal conductivity has only reached ~ 200 W/m-K. [8, 9]. Extrinsic defects and grain boundaries are a possible cause for this discrepancy, or higher order intrinsic scattering may become significant in determining the upper limit of k in this material.

The prediction of four-phonon scattering has been pursued for a long time. Lindsay *et al.* examined the phase space for

four-phonon scattering processes [10]. Turney *et al.* studied the fourth and higher order anharmonicity of the interatomic potential of argon, by truncating the potential in molecular dynamics (MD) simulations [11]. Sapna and Singh estimated the four-phonon scattering rates in carbon nanotubes using an analytical model involving approximations such as the Callaway model, the Debye model, etc [12]. Despite those efforts, a

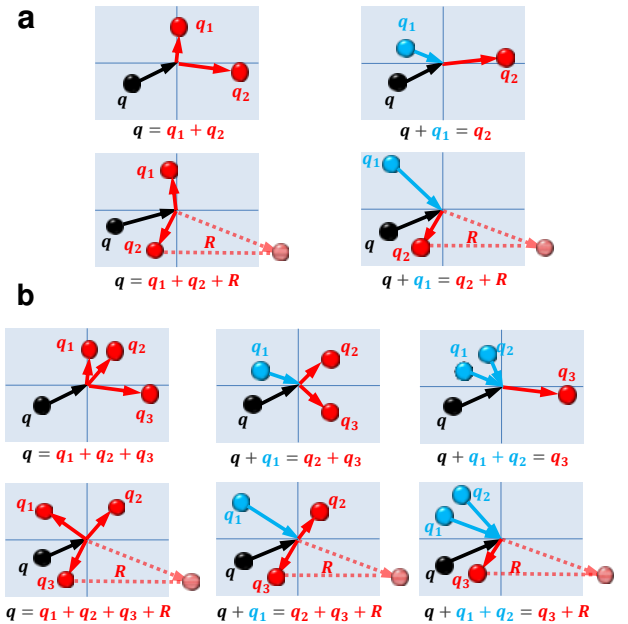


FIG. 1: **Three- and four-phonon scattering diagrams.** **a**, Three-phonon splitting and combination processes. **b**, Four-phonon splitting, redistribution, and combination processes. The shaded rectangles represent first Brillouin zone (BZ). The phonon momentum is $\hbar\mathbf{q}$ with \mathbf{q} standing for wave vector. The processes with momentum conserved are Normal processes. The others with momentum non-conserved are Umklapp processes, in which the resulting phonons are folded back by reciprocal lattice vectors \mathbf{R} .

direct and rigorous calculation of four-phonon scattering rates was not available until the recent work [13]. However, the anharmonic force constants therein were obtained from classical force fields though, thus the results are only qualitative and cannot be validated against experiments. In this work, for the first time we have calculated four-phonon scattering rates fully from first principles and examined its impact on thermal conductivity and optical phonon lifetimes of BAs, Si and diamond. We demonstrate that the four-phonon scattering resistance can reduce their predicted thermal conductivities and optical phonon lifetimes significantly, manifesting the general impact of four-phonon scattering in solids.

Intrinsic phonon scattering is caused by lattice anharmonicity [14]. From perturbation theory, the lowest-order anharmonic couplings involve three phonons: a single phonon may split into two phonons, or two phonons may combine to create a new phonon, as shown in Fig. 1 a. In addition to these scatterings, the second-order anharmonicity gives rise to four-phonon processes given in Fig. 1 b. The three- and four-phonon scattering rates for each phonon mode are calculated by summing up the probabilities of all the possible scattering events involving this mode. Each transition probability is calculated by Fermi's golden rule from density functional theory [13]. The computational cost is very high but can be mitigated as described in our earlier work [13]. See Sections Methods and Computational Challenge in the Supplemental Information [15].

We have obtained the scattering rates for each of the 24,576 (870 irreducible) phonon modes in the first BZ discretized by a $16 \times 16 \times 16$ \mathbf{q} mesh from 10 to 1,300 K for BAs, Si and diamond. The results as a function of phonon frequency at 300 K and 1,000 K are shown in Fig. 2 a-f. The insets show the low-frequency behavior. Contradictory to the prevailing general notion that four-phonon scattering is negligible, we find that for many frequencies the four-phonon scattering is comparable or even much stronger than the lower-order three-phonon scattering, as highlighted by the blue ovals.

Most surprisingly, the four-phonon scattering in BAs is quite strong, as shown in Fig. 2 a. This behaves against the general notion of perturbation theory, in which the strength of four-phonon scattering is driven by the magnitude of higher-order terms of the Hamiltonian, which are small in BAs. The origin of the strong four-phonon scattering in BAs is illustrated in Fig. 3 a. As shown in Ref. [3], the number of three-phonon processes is partly restricted by a large acoustic-optical (a-o) energy gap, which prevents two acoustic phonons from combining into an optical phonon as well as the reverse process since the summation of the energies of two acoustic phonons cannot reach that of the optical phonon. This coupled with closely bunched acoustic branches contributes to weak three-phonon scattering and the predicted ultra-high thermal conductivity in BAs [3]. Further evidence of this can be found near 21 THz in Fig. 2 a, in which the three-phonon scattering rates have a deep valley. These phonons are the optical modes near the Γ point and have high energy and small momentum. They can hardly find two other phonon modes that satisfy energy conservation and momentum conservation simultaneously. Such a large a-o gap, however, does not forbid four-phonon processes between acoustic and optical phonons. For example, at around 21 THz, the possible number of three-phonon scattering configurations is smaller than 20 while the number of four-phonon processes

is about $10^7 \sim 10^8$. For these four-phonon processes, 94% are $\mathbf{q} + \mathbf{q}_1 \rightarrow \mathbf{q}_2 + \mathbf{q}_3$ and 6% are $\mathbf{q} \rightarrow \mathbf{q}_1 + \mathbf{q}_2 + \mathbf{q}_3$. Therefore, the optical phonon relaxation time is brutally over-predicted by the three-phonon picture. When the four-phonon term is included, the relaxation time is reduced from 10^4 ps to 40 ps at room temperature, see Supplemental Fig. S1 [15]. Since these optical phonons near the Γ point are critical for interactions with electrons and photons such as in laser heating [16] and for infrared optical properties [17], the long lifetime predicted from three-phonon scattering is misleading, while four-phonon scattering is critical and should be included.

As temperature increases to 1,000 K, the four-phonon scattering of BAs becomes much more important especially for the phonons with higher frequencies as highlighted by the blue ovals in Fig. 2 b. The variation with temperature and frequencies are shown in Fig. 3 b and c, from which we determine the scaling law of four-phonon scattering $\tau_4^{-1} \sim T^2 \omega^4$. Compared to three-phonon scattering, four-phonon scattering is more important at higher temperatures and for higher-energy phonons as the phonon population increases with temperature and the phase space increases with phonon energy (frequency). Due to their simplicity, scaling laws are of great importance in thermal nanoengineering as well. For example, the power law of three-phonon scattering $\tau_3^{-1} \sim T\omega^2$ and phonon-defect scattering $\tau_d^{-1} \sim \omega^4$ has been widely used in understanding the experimental thermal conductivity in advanced thermoelectric materials [18, 19].

In the other two materials without phonon bandgaps, Si and diamond, four-phonon scattering is not as strong as in BAs but certainly not negligible. At 300 K, τ_4^{-1} is well below τ_3^{-1} for most of the acoustic phonons. This obeys the general notion in perturbation theory since the anharmonicity (the higher-order terms of the Hamiltonian) is small for Si and diamond. Nevertheless, the optical modes marked by the blue ovals still have large four-phonon rates, and this may explain why calculated infrared optical linewidths considering only three-phonon scattering are narrower than experiments [17, 20]. As T increases to 1,000 K, four-phonon rates of the low-frequency phonons remain insubstantial, however, higher-energy longitudinal acoustic (LA) modes and all the optical modes exhibit large τ_4^{-1} , comparable to τ_3^{-1} . The large τ_4^{-1} of the heat-carrying LA phonons will have a substantial effect on the thermal conductivity of these materials. In all the materials we note that the low-frequency limits for both three- and four-phonon scatterings obey the law $\lim_{\omega \rightarrow 0} \tau^{-1} = 0$ resulting from translational invariance, indicating the accuracy of our four-phonon calculations. The same as that in BAs, the power law fittings for diamond and silicon give $\tau_4^{-1} \sim \omega^4$, which are shown in Supplemental Fig. S2 [15].

We have also calculated κ beyond the relaxation time approximation (RTA) by exactly solving the phonon Boltzmann transport equation (BTE) using an iterative scheme mixing three-phonon interactions [1, 3]. In this work, due to the high computational cost, the four-phonon scattering rates are computed at the RTA level only and inserted into the iterative scheme that determines the nonequilibrium phonon distributions from mixing of the three-phonon processes. This is similar to employing phonon-isotope and phonon-boundary scattering terms in the full BTE solution [1, 21]. We will show that such an approximation is likely valid as the four-phonon

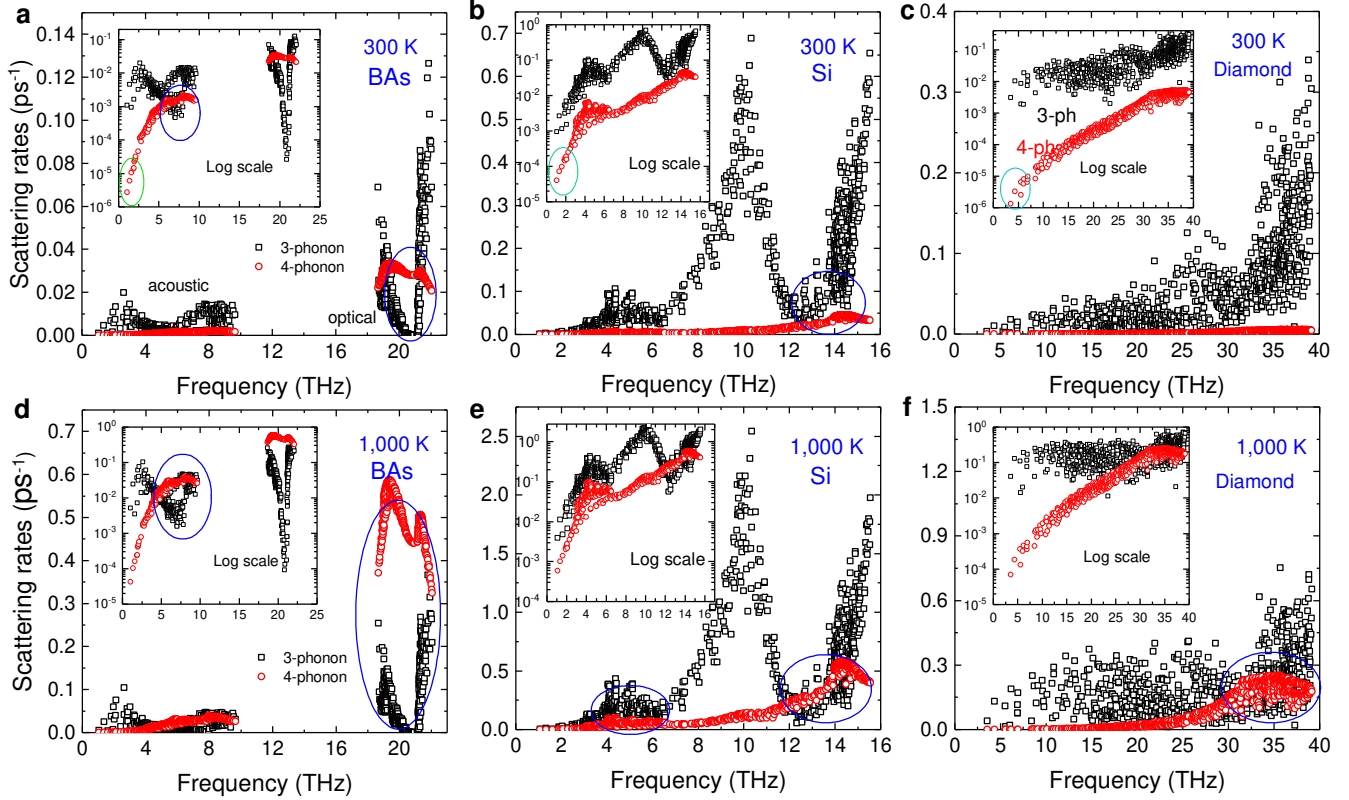


FIG. 2: **First-principles three-phonon (black dots) and four-phonon (red dots) scattering rates of BAs, Si and diamond at 300 K and 1,000 K.** The insets are in log-linear scales to give a better view of the low frequency regions. Blue ovals indicate the regions where four-phonon plays an important role. Green ovals in the insets indicate that our four-phonon results reproduce well the universal law $\lim_{\omega \rightarrow 0} \tau^{-1} = 0$, which is a critical mark of the calculation accuracy.

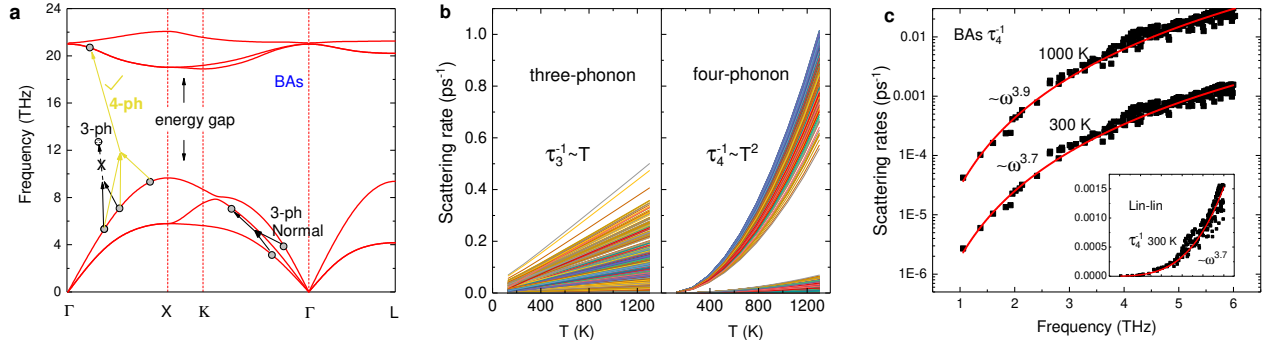


FIG. 3: **Impacts of phonon bandgap, temperature and frequency on three- and four-phonon scattering.** **a**, Three- and four-phonon processes in BAs. The large energy gap between acoustic and optical phonons forbids three-phonon processes, however, it is not as restrictive for four-phonon processes. The three-phonon process is dominated by intraband Normal scattering, while 4-phonon process is dominated by inner and interband Umklapp scattering. **b**, Temperature dependencies of three-phonon and four-phonon scatterings of BAs. The curves cover all the mode in the BZ, with each curve corresponding to a different mode. **c**, Power law fitting $\tau_4^{-1} = A\omega^n$ of the acoustic phonons in BAs. The log-linear scale gives a clear view of the low frequency behavior, while the inserted linear-linear scale is for the view of the high frequency behavior.

scattering is dominated by Umklapp processes. We also include phonon-isotope scattering [22] in these κ calculations of naturally occurring materials.

The thermal conductivity κ of the naturally occurring materials are shown in Fig. 4 b. For diamond and Si the three-phonon predictions agree well with measured data at low temperature (<600 K for Si, <900 K for diamond), however, significant deviations from experiment occur at high temperatures. For example, at 1,000 K three-phonon scattering alone

over-predicts κ of diamond and silicon by 31% and 26% as compared to experimental values, respectively. After including four-phonon scattering, we find that the predicted thermal conductivity agrees well with measurements throughout the entire temperature range. Such reduction is beneficial for thermoelectrics and thermal barrier coatings where low thermal conductivities are desired. As for naturally occurring BAs, even at room temperature the thermal conductivity is reduced significantly, from 2241 W/m-K to 1417 W/m-K, after in-

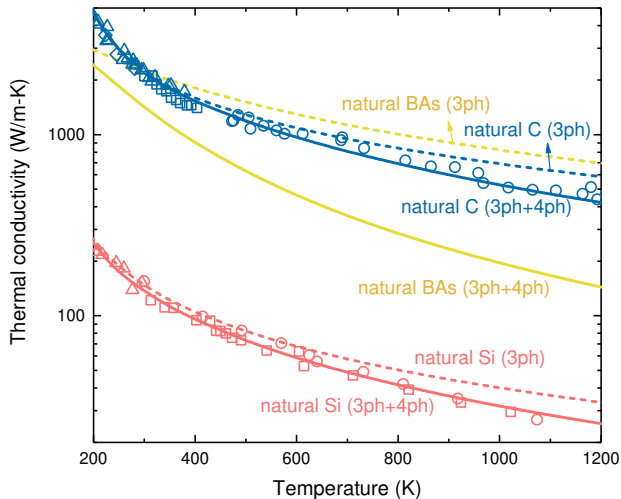


FIG. 4: **Thermal conductivities of naturally occurring BAs, Si and diamond.** Dashed lines give calculated κ_3 , while solid lines give κ_{3+4} . Note that these calculations employed the iterative solution of the BTE, beyond the RTA. Symbols represent measured data: red triangles [23], red squares [24], red circles [25], blue triangles [26], blue squares [27], blue circles [28], and blue-diamonds [29]. No experimental data are available for single-crystal BAs.

cluding four-phonon scattering. As temperature increases to 1,000 K, the reduction grows to over 60%. Nonetheless, the room temperature intrinsic thermal conductivity of BAs is still among the highest of known materials, and much higher than common metals. Moreover, the well-known temperature scaling $\kappa \sim 1/T$ at high temperature by only considering three-phonon scattering can now be modified to $\kappa \sim 1/(AT + BT^2)$ after adding four-phonon scattering.

Here we discuss important physical details regarding four-phonon scattering processes. More discussions on computational details can be found in Supplemental Information, i.e., the impact of broadening factors, \mathbf{q} -meshes, interatomic cutoff radius, translational invariance, force constant truncation, and the role of Normal and Umklapp processes [15].

Phase space. For a particular mode, the scattering phase space is the possible ways to find three- or four-phonon processes that conserve energy and momentum (or differs by \mathbf{R} for Umklapp processes). Four-phonon scattering can provide significant thermal resistance despite being higher order in perturbation theory with much lower scattering probability per scattering configuration than for three-phonon processes. The origin is found to be the huge phase space allowed for four-phonon scattering. The number of the possible scattering configurations of four-phonon process using a $16 \times 16 \times 16$ \mathbf{q} -mesh for the materials examined here are on the order of $10^7 - 10^8$, which is 4-5 orders more than the number of three-phonon configurations for the same grid density.

Estimating the significance of four-phonon scattering. As discussed above, four-phonon scattering becomes increasingly important at higher temperatures in all materials; however, after careful examination of the three materials in this work, a natural question is: How important is four-phonon scattering generally in other solids? Like three-phonon scattering, four-phonon scattering is governed by two factors: anharmonicity of the energy surface and the scattering phase space. We

note that construction of both three- and four-phonon scattering rates is nontrivial, each a complex interplay of matrix elements, eigenvectors and frequencies combining various phonon modes. However, based on the phonon scattering formalism [13], we attempt to estimate the relative importance of the fourth order anharmonicity. Yue *et al.* have done some estimations by evaluating the fourth-order terms in Hamiltonians [30]. Roughly speaking, the n th order phonon scattering rate is proportional to $\tau_n^{-1} \sim |V_n|^2 \cdot f^{n-2}/\omega^{n+1}$, where V_n is the n th order transition matrix element $\sim \Phi_n/m^{n/2}$. Φ_n is the n th order force constant, m is the average atomic mass, and f is the Bose-Einstein distribution approximated as $\sim k_B T/\hbar\omega$, high temperature behavior. Approximating the frequency as $\omega \sim \sqrt{\Phi_2/m}$ we get $\tau_n^{-1} \sim |\Phi_n|^2 \cdot \sqrt{m}/|\Phi_2|^{n-1/2}$. Thus, the relative importance of the n th order anharmonicity to the third order anharmonicity is crudely evaluated as $\tau_n^{-1}/\tau_3^{-1} \sim |\Phi_n/\Phi_3|^2/|\Phi_2|^{n-3}$. Note that this captures the increasing importance of higher order scattering with increasing temperature and only requires single calculations of a second, third and higher order derivative of the potential energy.

Table I gives the relative strength of the fourth order anharmonicity ($|\Phi_4/\Phi_3|^2/|\Phi_2|$) for the materials studied in this work and a low thermal conductivity zinc blende material, CuCl (~ 1 W/m-K at room temperature) [31, 32]. Note that the temperature dependence is omitted. This simplistic formalism is intended to give a rough estimation for when higher order scattering may be important in one material over another. As shown in the table, the fourth order anharmonicity is predicted to be significantly more important in CuCl than in the high conductivity materials considered here. Our previous work based on interatomic potentials directly demonstrated a positive correlation between four-phonon scattering and anharmonicity [13]. In soft materials, such as those of interest for thermoelectric and thermal barrier coating applications [4], atoms can deviate significantly from equilibrium to sample higher anharmonicity, fourth-order terms in their Hamiltonians can be quite large [30]. Note that $|\Phi_4/\Phi_3|^2/|\Phi_2|$ is higher in Si than in diamond; and Si also has fourth order scattering that is relatively more important.

This simple formalism, however, does not account for the phase space for three- and four-phonon scattering processes. Thus, even though the fourth order anharmonicity is predicted to be relatively less important in BAs as demonstrated in Table I, the BAs conductivity is still significantly reduced by the four-phonon scattering resistance (Fig. 4). In materials where three-phonon scattering is weak due to dispersion features (e.g., phonon bandgap, acoustic bunching [3]) that reduce the phase space, we expect four-phonon scattering to also be important, for instance in BSb [3] and BeSe [33].

In summary, we have rigorously calculated four-phonon scattering rates and thermal resistance from fully first principles methods. Due to the large number of possible scattering configurations, four-phonon processes play an important role in determining intrinsic phonon transport. We find that four-phonon scattering is surprisingly strong in zincblende BAs, and reduces its thermal conductivity substantially from $\sim 2,200$ to $\sim 1,400$ W/m-K even at room temperature when compared to previous calculations. The room-temperature optical phonon lifetime is reduced from 10^4 to 40 ps. These impacts increase substantially with increasing temperature. Such strong

TABLE I: Anharmonicity ratio $|\Phi_4/\Phi_3|^2/|\Phi_2|$ to approximate the relative importance of four-phonon scattering in Diamond, Si, BAs and CuCl. Units of the representative n th order force constant are $eV/\text{\AA}^n$. The representative force constant values are taken at $\Phi_{0,0,x,x}$, $\Phi_{0,0,0,x,y,z}$, and $\Phi_{0,0,0,0,x,x,x,x}$. The index 0 indicates the origin primitive cell, see the supplemental information for more details [15]. The third order $\Phi_{0,0,0,0,x,x,x} = 0$ due to symmetries cannot be taken as the anharmonicity.

material	atom	$ \Phi_2 $	$ \Phi_3 $	$ \Phi_4 $	$ \Phi_4/\Phi_3 ^2/ \Phi_2 $
Diamond	C	44.2	95.0	224	0.128
BAs	B	17.0	45.6	65.7	0.122
Si	Si	13.6	33.9	50.8	0.165
CuCl	Cu	1.93	13.1	48.0	6.956

four-phonon scattering in BAs originates from weak three-phonon scattering that arises from properties of its phonon dispersion coupled with fundamental conservation conditions. These conditions do not as strictly restrict the available scatterings for four-phonon process. With four-phonon scattering included, the predicted thermal conductivities of silicon and diamond reduces κ by $\sim 30\%$ at high temperature, and brings the prediction in agreement with measured data throughout the entire temperature range considered. Based on our findings of the general and significant impact of four-phonon processes on thermal transport and optic mode lifetimes, we expect future predictions of these properties will incorporate this important scattering mechanism, especially when considering engineering materials for energy transfer, conversion and storage applications.

The work on silicon is supported by the Defense Advanced Research Projects Agency (Award No. HR0011-15-2-0037), and the work on diamond and boron arsenide is supported by the National Science Foundation (Award No. 1150948). L.L. acknowledges support from the U. S. Department of Energy, Office of Science, Office of Basic Energy Sciences, Materials Sciences and Engineering Division.

* Electronic address: ruan@purdue.edu

- [1] D. A. Broido, M. Malorny, G. Birner, N. Mingo, and D. A. Stewart, Applied Physics Letters **91**, 231922 (2007), ISSN 00036951, URL <http://aip.scitation.org/doi/10.1063/1.2822891>.
- [2] K. Esfarjani, G. Chen, and H. T. Stokes, Physical Review B **84**, 085204 (2011), ISSN 1098-0121, URL <http://link.aps.org/doi/10.1103/PhysRevB.84.085204>.
- [3] L. Lindsay, D. A. Broido, and T. L. Reinecke, Physical Review Letters **111**, 025901 (2013), ISSN 0031-9007, URL <http://link.aps.org/doi/10.1103/PhysRevLett.111.025901>.
- [4] A. Seko, A. Togo, H. Hayashi, K. Tsuda, L. Chaput, and I. Tanaka, Physical Review Letters **115**, 205901 (2015), ISSN 0031-9007, 1506.06439, URL <http://link.aps.org/doi/10.1103/PhysRevLett.115.205901>.
- [5] M. Novikov, T. Ositinskaya, O. Shulzhenko, O. Podoba, O. Sokolov, and I. Petrusha, Dopovidi Akademii Nauk Ukrain'skoi Rsr Seriya A-Fiziko-Matematichni Ta Technichni Nauki pp. 72–75 (1983).
- [6] L.-D. Zhao, S.-H. Lo, Y. Zhang, H. Sun, G. Tan, C. Uher, C. Wolverton, V. P. Dravid, and M. G. Kanatzidis, Nature **508**, 373 (2014).
- [7] R. Guo, X. Wang, Y. Kuang, and B. Huang, Physical Review B **92**, 115202 (2015).
- [8] H. Ma, C. Li, S. Tang, J. Yan, A. Alatas, L. Lindsay, B. C. Sales, and Z. Tian, Physical Review B **94**, 220303 (2016), ISSN 2469-9950, URL <http://link.aps.org/doi/10.1103/PhysRevB.94.220303>.
- [9] J. Kim, D. A. Evans, D. P. Sellan, O. M. Williams, E. Ou, A. H. Cowley, and L. Shi, Applied Physics Letters **108** (2016), ISSN 00036951, URL <http://dx.doi.org/10.1063/1.4950970>.
- [10] L. Lindsay and D. a. Broido, Journal of Physics: Condensed Matter **20**, 165209 (2008), ISSN 0953-8984, URL <http://stacks.iop.org/0953-8984/20/i=16/a=165209?key=crossref.fae35963b743f83e18ac9194e2b46aea>.
- [11] J. Turney, E. Landry, A. McGaughey, and C. Amon, Physical Review B **79**, 064301 (2009), ISSN 1098-0121, URL <http://link.aps.org/doi/10.1103/PhysRevB.79.064301>.
- [12] P. Sapna and T. J. Singh, Modern Physics Letters B **27**, 1350117 (2013), ISSN 0217-9849, URL <http://www.worldscientific.com/doi/abs/10.1142/S0217984913501170>.
- [13] T. Feng and X. Ruan, Physical Review B **93**, 045202 (2016), ISSN 2469-9950, URL <http://link.aps.org/doi/10.1103/PhysRevB.93.045202>.
- [14] A. Maradudin and A. Fein, Physical Review **128**, 2589 (1962), URL http://prola.aps.org/abstract/PR/v128/i6/p2589_1.
- [15] See Supplemental Material for: the methods and computational challenge, the relaxation time of the optical phonon at Γ point in BAs, the power law fittings for diamond and silicon, the impact of broadening factors, \mathbf{q} -meshes, interatomic cutoff radius, translational invariance, force constant truncation, and the role of Normal and Umklapp processes.
- [16] A. K. Vallabhaneni, D. Singh, H. Bao, J. Murthy, and X. Ruan, Phys. Rev. B **93**, 125432 (2016), URL <http://link.aps.org/doi/10.1103/PhysRevB.93.125432>.
- [17] H. Bao, B. Qiu, Y. Zhang, and X. Ruan, Journal of Quantitative Spectroscopy and Radiative Transfer **113**, 1683 (2012), ISSN 00224073, URL <http://linkinghub.elsevier.com/retrieve/pii/S0022407312002336>.
- [18] S. I. Kim, K. H. Lee, H. a. Mun, H. S. Kim, S. W. Hwang, J. W. Roh, D. J. Yang, W. H. Shin, X. S. Li, Y. H. Lee, et al., Science **348**, 109 (2015), ISSN 0036-8075, URL <http://science.sciencemag.org/content/348/6230/109>.
- [19] M. Hong, T. C. Chasapis, Z.-G. Chen, L. Yang, M. G. Kanatzidis, G. J. Snyder, and J. Zou, ACS Nano **10**, 4719 (2016), ISSN 1936-0851.
- [20] A. Debernardi, Phys. Rev. B **57**, 12847 (1998), URL <https://link.aps.org/doi/10.1103/PhysRevB.57.12847>.
- [21] L. Lindsay, D. Broido, and N. Mingo, Physical Review B **80**, 125407 (2009), ISSN 1098-0121, URL <http://link.aps.org/doi/10.1103/PhysRevB.80.125407>.
- [22] S.-i. Tamura, Physical Review B **27**, 858 (1983), URL <http://journals.aps.org/prb/abstract/10.1103/PhysRevB.27.858>.
- [23] T. Ruf, R. Henn, M. Asen-Palmer, E. Gmelin, M. Cardona, H.-J. Pohl, G. Devyatych, and P. Sennikov, Solid State Communications **115**, 243 (2000).
- [24] B. Abeles, D. Beers, G. Cody, and J. Dismukes, Physical review **125**, 44 (1962), URL <https://journals.aps.org/pr/abstract/10.1103/PhysRev.125.44>.
- [25] C. J. Glassbrenner and G. A. Slack, Physical Review **134**, A1058 (1964), URL http://prola.aps.org/abstract/PR/v134/i4A/pA1058_1.
- [26] L. Wei, P. K. Kuo, R. L. Thomas, T. R. Anthony, and W. F. Banholzer, Phys. Rev. Lett. **70**, 3764 (1993), URL <https://link.aps.org/doi/10.1103/PhysRevLett.70.3764>.
- [27] D. G. Onn, A. Witek, Y. Z. Qiu, T. R. Anthony, and W. F. Banholzer, Phys. Rev. Lett. **68**, 2806 (1992), URL <https://link.aps.org/doi/10.1103/PhysRevLett.68.2806>.
- [28] J. R. Olson, R. O. Pohl, J. W. Vandersande, A. Zoltan,

- T. R. Anthony, and W. F. Banholzer, Phys. Rev. B **47**, 14850(1993), URL <https://link.aps.org/doi/10.1103/PhysRevB.47.14850>.
- [29] R. Berman, P. R. W. Hudson, and M. Martinez, Journal of Physics C: Solid State Physics **8**, L430 (1975), URL <http://stacks.iop.org/0022-3719/8/i=21/a=003>.
- [30] S.-Y. Yue, X. Zhang, G. Qin, S. R. Phillpot, and M. Hu, Phys. Rev. B **95**, 195203 (2017), URL <https://link.aps.org/doi/10.1103/PhysRevB.95.195203>.
- [31] G. A. Slack and P. Andersson, Physical Review B **26**, 1873 (1982).
- [32] S. Mukhopadhyay, D. Bansal, O. Delaire, D. Perrodin, E. Bourret-Courchesne, D. J. Singh, and L. Lindsay, Physical Review B **96**, 100301 (2017).
- [33] L. Lindsay, D. A. Broido, and T. L. Reinecke, Physical Review B **88**, 144306 (2013).