

Fourier – transform spectroscopy of $^{13}\text{C}^{17}\text{O}$ and deperturbation analysis of the $\text{A}^1\Pi$ ($v = 0 - 3$) levels

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ABSTRACT

The high-resolution $\text{B}^1\Sigma^+ \rightarrow \text{A}^1\Pi$ (0, 0) and (0, 3) emission bands of the less-abundant $^{13}\text{C}^{17}\text{O}$ isotopologue have been investigated by Fourier-transform spectroscopy in the visible region using a Bruker IFS 125HR spectrometer at an accuracy 0.003 cm^{-1} . These spectra are combined with high-resolution photoabsorption measurements of the $^{13}\text{C}^{17}\text{O}$ $\text{B}^1\Sigma^+ \leftarrow \text{X}^1\Sigma^+$ (0, 0), $\text{B}^1\Sigma^+ \leftarrow \text{X}^1\Sigma^+$ (1, 0) and $\text{C}^1\Sigma^+ \leftarrow \text{X}^1\Sigma^+$ (0, 0) bands recorded with an accuracy of 0.01 cm^{-1} using the vacuum ultraviolet Fourier-transform spectrometer, installed on the DESIRS beamline at the SOLEIL synchrotron. In the studied $17,950 - 22,500\text{ cm}^{-1}$ and $86,800 - 92,100\text{ cm}^{-1}$ regions, 480 transitions have been measured. These new experimental data were combined with data from the $\text{C} \rightarrow \text{A}$ and $\text{B} \rightarrow \text{A}$ systems, previously analyzed in $^{13}\text{C}^{17}\text{O}$. The frequencies of 1003 transitions derived from 12 bands were used to analyze the perturbations between the $\text{A}^1\Pi$ ($v = 0 - 3$) levels and rovibrational levels of the $\text{d}^3\Delta_i$, $\text{e}^3\Sigma^-$, $\text{a}^3\Sigma^+$, $\text{I}^1\Sigma^-$ and $\text{D}^1\Delta$ states as well as to a preliminary investigation of weak irregularities that appear in the $\text{B}^1\Sigma^+$ ($v = 0$) level. Deperturbed molecular constants and term values of the $\text{A}^1\Pi$ state were obtained. The spin-orbit and L -uncoupling interaction parameters as well as isotopologue-independent spin-orbit and rotation-electronic perturbation parameters were derived.

Keywords: FT-spectroscopy; High resolution VUV-VIS spectra; Perturbation analysis; Carbon monoxide; $^{13}\text{C}^{17}\text{O}$ isotopologue.

1. Introduction

The spectroscopy of the carbon monoxide (CO) molecule remains of central interest to a variety of subfields in science, from the submm [1,2] to the X-ray ranges [3]. It is of relevance, inter alia, to many astronomical environments of the interstellar medium (ISM) [4–9], planetary [10] and exoplanetary atmospheres [11], cosmology [12], as well as in combustion processes [13], atmospheric physics [14], and plasma science [15,16]. Various CO isotopologues are often used as tracers to determine the content of molecular clouds and to map the distribution of matter [17–19]. An important description of modeling CO photochemistry was made by Visser *et al.* [20]. The large CO abundance gives a detectable signal for even the least-abundant isotopologues [21–27], including $^{13}\text{C}^{17}\text{O}$ [23,27]. The first observation of $^{13}\text{C}^{17}\text{O}$ in the ISM was made in 2001 by Bensch *et al.* [23] in the ρ -Ophiuchi molecular cloud with the ESO Submillimetre Telescope (SEST) at La Silla, Chile. Under laboratory conditions the $^{13}\text{C}^{17}\text{O}$ isotopologue was studied several times [14,28–31], recently observed in the vacuum ultraviolet (VUV) by Lemaire *et al.* [32], and investigated in the visible (VIS) by Hakalla *et al.* [33,34] and Hakalla [35].

The $A^1\Pi$ state of carbon monoxide is among the most extensively studied electronic states in molecular spectroscopy. While it is important in a variety of applications, the $A^1\Pi$ state is a benchmark for the study of multistate perturbations [36]. These perturbations are caused by rovibrational levels of the $e^3\Sigma^-$, $a^3\Sigma^+$, and $d^3\Delta_i$ triplet states as well as the $I^1\Sigma^-$ and $D^1\Delta$ singlet states. A systematic classification of the perturbations occurring in the $A^1\Pi$ state of the main $^{12}\text{C}^{16}\text{O}$ isotopologue was conducted by Krupenie [37] and Simmons *et al.* [38]. A considerable contribution to the identification and classification of the $A^1\Pi$ state perturbations has been made by Keĉpa and Rytel in a number of investigations over the years [39–41]. A comprehensive analysis and deperturbation calculations were conducted by Field *et al.* [42,43]. Later, Le Floch *et al.* [44] carried out a detailed study of perturbations appearing in the lowest $A^1\Pi$ ($v = 0$) vibrational level. Perturbations occurring in the $A^1\Pi$ ($v = 0 - 4$) levels were analyzed by Le Floch [45,46], who also derived very precise values of the terms for the $A^1\Pi$, $v = 0-8$ levels. Recently, the lowest vibrational levels, $v = 0 - 4$, of the $A^1\Pi$ state were deperturbed by Niu *et al.* [47,48] by means of accurate two-photon Doppler-free excitation using narrow band lasers with relative accuracy up to $\Delta\lambda/\lambda = 2 \times 10^{-8}$, as well as by a vacuum ultraviolet Fourier-transform spectrometer (VUV-FTS) installed on the DESIRS beamline at the SOLEIL synchrotron.

Deperturbation analyses of the $A^1\Pi$ state, in other CO isotopologues, were conducted by: Haridass *et al.* [49], Gavilan *et al.* [50], and Niu *et al.* [51] in $^{13}\text{C}^{16}\text{O}$; Haridass *et al.* [52,53] in $^{12}\text{C}^{18}\text{O}$ and $^{13}\text{C}^{18}\text{O}$; Beaty *et al.* [54] in $^{12}\text{C}^{18}\text{O}$; Hakalla *et al.* [55] in $^{12}\text{C}^{17}\text{O}$ by means of high-accuracy dispersive optical spectroscopy ($\Delta\lambda/\lambda$ up to about 1×10^{-7}) and VUV Fourier-transform measurements. For the less-abundant $^{13}\text{C}^{17}\text{O}$ isotopologue, the deperturbation analysis of the $A^1\Pi$ state has never been performed.

The $B^1\Sigma^+$ state is the first member of the $n\sigma$ Rydberg series of CO. The rotational perturbations occurring in the $B^1\Sigma^+$ ($v = 0$) level were studied by Amiot et al. [56] in the infrared $C^1\Sigma^+ - B^1\Sigma^+$ spectra of the $^{12}\text{C}^{16}\text{O}$ and $^{14}\text{C}^{16}\text{O}$ isotopologues. Observed level–shift irregularities were very weak: $0.013 - 0.017 \text{ cm}^{-1}$ for $J = 4$ and 6 in $^{12}\text{C}^{16}\text{O}$ as well as $0.003 - 0.005 \text{ cm}^{-1}$ for $J = 18$ and 20 in $^{14}\text{C}^{16}\text{O}$, while the measurement accuracy was about 0.005 cm^{-1} . Janjić et al. [57] found a ‘well marked’ perturbation in the $B^1\Sigma^+$ ($v = 1, J = 7 - 8$) of $^{13}\text{C}^{16}\text{O}$, without assigning the perturber state(s).

This work focuses on a comprehensive deperturbation analysis of the $A^1\Pi$ ($v = 0, 1, 2,$ and 3) rovibrational levels in the $^{13}\text{C}^{17}\text{O}$ isotopologue. In order to perform this analysis, information on the following bands is included: the $^{13}\text{C}^{17}\text{O}$ $B^1\Sigma^+ \rightarrow A^1\Pi$ ($0, 0$), ($0, 3$), $B^1\Sigma^+ \leftarrow X^1\Sigma^+$ ($0, 0$), ($1, 0$), and $C^1\Sigma^+ \leftarrow X^1\Sigma^+$ ($0, 0$) [58], as well as previously published studies of the $B^1\Sigma^+ \rightarrow A^1\Pi$ ($0, 1$), ($0, 2$), ($1, 0$), ($1, 1$) [33,34] and $C^1\Sigma^+ \rightarrow A^1\Pi$ ($0, 1$), ($0, 2$), and ($0, 3$) [35]. The $B^1\Sigma^+ \leftarrow X^1\Sigma^+$, and $C^1\Sigma^+ \leftarrow X^1\Sigma^+$ systems are included in the study in order to determine level energies of $A^1\Pi$ in relation to the $X^1\Sigma^+$ ($v = 0, J = 0$) ground state of $^{13}\text{C}^{17}\text{O}$, to search for an occurrence of possible perturbations in the $B^1\Sigma^+$ ($v = 0$) level and to introduce an independent set of improved constants for the $B^1\Sigma^+$ ($v = 0, 1$) and $C^1\Sigma^+$ ($v = 0$) levels. A comprehensive fit of the bands is performed to determine accurate deperturbed rotational constants and energy levels of the states under consideration.

2. Experiments

2.1. VIS–FTS data

The $B^1\Sigma^+ - A^1\Pi$ ($0, 0$) and ($0, 3$) bands were measured in emission from a water–cooled hollow–cathode discharge lamp with two anodes [59]. The lamp was initially filled with a mixture of helium and acetylene $^{13}\text{C}_2\text{D}_2$ (Cambridge Isotopes, 99.98% of ^{13}C) at a pressure of approximately 7 Torr. A dc electric current was passed through the mixture for about 200 h. During that time, a small quantity of ^{13}C carbon became deposited on the inside of the cathode. Next, the lamp was evacuated, and then oxygen containing the $^{17}\text{O}_2$ isotope (Sigma–Aldrich, $^{17}\text{O}_2$ 70%) was admitted. The anodes were operated at $2 \times 650 \text{ V}$ and $2 \times 50 \text{ mA}$ dc in the conditions of a static gas pressure of 2 Torr inside the lamp. During the discharge process, the $^{17}\text{O}_2$ molecules react with the ^{13}C –carbon atoms ejected from the cathode, thus forming the $^{13}\text{C}^{17}\text{O}$ molecules in the gas phase [34,35]. The temperature of the dc–plasma formed at the center of the cathode was about 650 K.

The high–resolution (HR) visible (VIS) emission spectra of the $^{13}\text{C}^{17}\text{O}$ $B^1\Sigma^+ \rightarrow A^1\Pi$ ($0, 0$) and ($0, 3$) bands were observed in the Materials Spectroscopy Laboratory, University of Rzeszów, with a 1.71–m Bruker IFS 125HR Fourier–transform spectrometer (FTS). The FTS was equipped with a visible quartz beamsplitter and a photomultiplier tube with housing and preamplifier, running in integration pulse mode. The optical bench of the FTS was evacuated to

a pressure lower than 0.002 hPa. The emission of the external CO source was focused onto the 1.0 mm diameter of the entrance aperture (field stop) of the FTS by a plano-convex quartz lens. In practice, the aperture is chosen to select as much as possible of the central region of the pattern of the interferogram. The permanently aligned interferometer has a 30° angle-of-incidence. The narrow angle-of-incidence makes more effective use of the beamsplitter by reducing beam polarization effects. The maximum optical path difference (OPD) of the FTS is 258 cm. The interferogram was obtained while recording the interferometric signal versus the OPD in a forward-backward acquisition mode. The inverse FT operation was then implemented to recover the spectrum from the interferogram. A 1.2 mW frequency-stabilized, single mode He-Ne laser controlled the position of the moving interferometer mirror (scanner), as well as provided a convenient internal wavenumber standard. The relative frequency stability of the laser beam was $\pm 2 \times 10^{-9}$ over 1 h.

The spectrum in the 11,000 – 25,000 cm^{-1} wavenumber region was recorded at an instrumental resolution of 0.018 cm^{-1} by coadding 128 scans in about 2.5 h of integration. Higher resolution was not necessary, since the full width at half maximum (FWHM) Doppler width of the CO lines is about 0.06 cm^{-1} for the medium-strength, and single lines. The resolving power of the instrument was better than 10^6 . The recorded spectrum had a signal-to-noise ratio (SNR) of 70:1 and 100:1 for the $^{13}\text{C}^{17}\text{O B}^1\Sigma^+ \rightarrow \text{A}^1\Pi(0, 0)$ and the (0, 3) band, respectively.

The spectrum was calibrated using the He-Ne 633 nm line (0.002 cm^{-1} FWHM, 3000:1 SNR) of the internal frequency-stabilized laser. The line positions were measured by fitting Voigt lineshape functions to the experimental lines in a Levenberg-Marquardt procedure [60,61] included in the Bruker OPUS™ software package [62]. The estimated absolute 1σ calibration uncertainty was 0.003 cm^{-1} . Finally, we expect the absolute accuracy of the line frequency measurements to be about 0.003 cm^{-1} for single, medium-strength lines and about 0.01–0.02 cm^{-1} for weak/blended ones.

While determining values of the molecular wavenumbers, blending due to spectral lines associated with other isotopologues, $^{13}\text{C}^{16}\text{O B} \rightarrow \text{A}(0, 0)$ [51] and $\text{B} \rightarrow \text{A}(0, 3)$ [63], $^{12}\text{C}^{17}\text{O B} \rightarrow \text{A}(0, 0)$ (theoretically evaluated) and $\text{B} \rightarrow \text{A}(0, 3)$ [55], as well as $^{12}\text{C}^{16}\text{O B} \rightarrow \text{A}(0, 0)$, and (0, 3) [64] bands was taken into account. In the studied regions 17,950 – 18,450 cm^{-1} and 22,150 – 22,500 cm^{-1} , in total 315 molecular emission lines of $^{13}\text{C}^{17}\text{O}$ were analyzed, 189 of which were interpreted as belonging to the (0, 0) and (0, 3) bands of the $\text{B}^1\Sigma^+ \rightarrow \text{A}^1\Pi$ system, and 122 were identified as lines deriving from the $\text{B}^1\Sigma^+ \rightarrow \text{d}^3\Delta_i$, $\text{B}^1\Sigma^+ \rightarrow \text{e}^3\Sigma^-$, $\text{B}^1\Sigma^+ \rightarrow \text{a}^3\Sigma^+$, $\text{B}^1\Sigma^+ \rightarrow \text{I}^1\Sigma^-$, $\text{B}^1\Sigma^+ \rightarrow \text{D}^1\Delta$, $\text{C}^1\Sigma^+ \rightarrow \text{e}^3\Sigma^-$, $\text{C}^1\Sigma^+ \rightarrow \text{a}^3\Sigma^+$, and $\text{C}^1\Sigma^+ \rightarrow \text{D}^1\Delta$. These 122 “extra-lines” refer to the spectral emission lines terminating on perturber states and gaining intensity from mixing with the $\text{A}^1\Pi$ state. Transition frequencies for the $^{13}\text{C}^{17}\text{O B}^1\Sigma^+ \rightarrow \text{A}^1\Pi(0, 0)$ and (0, 3) bands are listed in Table 1, and wavenumbers connected with perturber states are listed in Table 2. During the deperturbation analysis, we extended and corrected the assignment of some heavily perturbed or weak $^{13}\text{C}^{17}\text{O}$ lines, previously analyzed [33–35]. They are collected in Table 3. An

overview of the observed $^{13}\text{C}^{17}\text{O}$ $\text{B}^1\Sigma^+ (v = 0) \rightarrow \text{A}^1\Pi (v = 0, 3)$ spectra, together with extra-lines, rotational assignments and simulated spectra are shown in Figs. 1 and 2. The ratios of the molecular gas compositions used to obtain the spectra were $^{13}\text{C}^{17}\text{O} : ^{13}\text{C}^{16}\text{O} : ^{12}\text{C}^{17}\text{O} : ^{12}\text{C}^{16}\text{O} = 1 : 0.76 : 0.34 : 0.25$.

2.2. VUV-FTS data

VUV spectra are recorded with the Fourier-transform spectrometer installed on the DESIRS beamline at the SOLEIL synchrotron. This spectrometer is described in great detail in [65–68]. $^{13}\text{C}^{17}\text{O}$ $\text{A}^1\Pi \leftarrow \text{X}^1\Sigma^+$ photoabsorption bands have been recently observed for the first time by Lemaire et al. [32] in the course of a $^{13}\text{C}^{18}\text{O}$ study. $^{13}\text{C}^{17}\text{O}$ bands have been also observed in their $^{13}\text{C}^{16}\text{O}$ sample. The occurrence of these $^{13}\text{C}^{17}\text{O}$ bands is due to $< 0.1\%$ isotopic contamination of $^{13}\text{C}^{17}\text{O}$ in samples of $^{13}\text{C}^{16}\text{O}$ and $^{13}\text{C}^{18}\text{O}$, and are accompanied by much stronger absorption from the main isotopologues. Full details including transition frequencies, term values and molecular constants of the multi-isotopologue measurements of $^{13}\text{C}^{17}\text{O}$ $\text{B}^1\Sigma^+ \leftarrow \text{X}^1\Sigma^+ (0, 0)$, $\text{B}^1\Sigma^+ \leftarrow \text{X}^1\Sigma^+ (1, 0)$ and $\text{C}^1\Sigma^+ \leftarrow \text{X}^1\Sigma^+ (0, 0)$ and other bands will be published in a separate paper [58]. Figures 3, 4, and 5 show the corresponding $^{13}\text{C}^{17}\text{O}$ spectra obtained after subtracting from a $^{13}\text{C}^{16}\text{O}$ raw spectrum a model including $^{12}\text{C}^{16}\text{O}$, $^{13}\text{C}^{16}\text{O}$, and $^{13}\text{C}^{18}\text{O}$. At the high pressures used to record spectra with nominally $^{13}\text{C}^{16}\text{O}$ or $^{13}\text{C}^{18}\text{O}$ gas species, lines of the dominant gas are strongly saturated, and lines from the trace isotopologues, particularly $^{13}\text{C}^{17}\text{O}$, can be identified. Based on the relative strengths of the absorption features, the ratio of the isotopologues in the $^{13}\text{C}^{16}\text{O}$ gas sample, sorted by concentration, was found to be $^{13}\text{C}^{16}\text{O} : ^{13}\text{C}^{18}\text{O} : ^{13}\text{C}^{17}\text{O} : ^{12}\text{C}^{16}\text{O} = 1 : 0.041 : 0.073 : 0.0045$. Relevant to this study are the J -range of $^{13}\text{C}^{17}\text{O}$ rotational energy levels and their accuracy deduced from these spectra. We note that the $^{13}\text{C}^{17}\text{O}$ results extracted from $^{13}\text{C}^{16}\text{O}$ and $^{13}\text{C}^{18}\text{O}$ transition frequencies are less accurately determined close to the band origin (i.e. at low J' values) due to the saturation effects inducing blending: $J' < 2$ for B(0), $J' < 7$ for C(0). This effect becomes much weaker or absent for the high J' of B(0) and C(0), and all J' values of B(1) as the lines of the different isotopologues become more separated while v' and J' increase and the saturation of the dominant species disappears. Thus the results for low J' can be reliably extrapolated from the observed levels (up to $J = 23$ for B(0), 21 for B(1) and 39 for C(0)). At J' higher than the above mentioned values, the absolute accuracy on the $^{13}\text{C}^{17}\text{O}$ transition frequencies is similar to that of $^{13}\text{C}^{16}\text{O}$ and $^{13}\text{C}^{18}\text{O}$. The total uncertainty of the VUV-FTS-measured B(0), B(1), and C(0) level energies is 0.04 cm^{-1} on average and always lower than 0.1 cm^{-1} . These uncertainties are a combination of statistical fitting errors (different for every line, estimated to fall between 0.001 and 0.09 cm^{-1}) and calibration error of the wavenumber scale (different for each band, estimated to be 0.01 cm^{-1}). The wavenumber calibration was made with reference to various CO absorption lines measured in the highly-accurate laser-based work of Drabbels et al. [69], Ubachs et al. [30], and Cacciani et al. [70].

3. Results

3.1. Perturbations in CO

The CO molecular orbitals are conventionally labeled in energy and symmetry order: 1σ ($\sigma 1s$), 2σ ($\sigma^* 1s$), 3σ ($\sigma 2s$), 4σ ($\sigma^* 2s$), 1π ($\pi 2p$), 5σ ($\sigma 2p$), 2π ($\pi^* 2p$), 6σ ($3s\sigma$), 7σ ($3p\sigma$), where $(1 - 2)\sigma$ are core orbitals, $(3 - 5)\sigma$ and $(1 - 2)\pi$ are valence orbitals, and $(6 - 7)\sigma$ are Rydberg orbitals. The principal configurations of the electronic states considered in this paper are:

$$\begin{aligned} X^1\Sigma^+ &: && 1\sigma^2 2\sigma^2 3\sigma^2 4\sigma^2 1\pi^4 5\sigma^2, \\ A^1\Pi \text{ and } a^3\Pi &: && 1\sigma^2 2\sigma^2 3\sigma^2 4\sigma^2 1\pi^4 5\sigma^1 2\pi^1 \\ a^3\Sigma^+, e^3\Sigma^-, d^3\Delta_i, I^1\Sigma^-, D^1\Delta &: && 1\sigma^2 2\sigma^2 3\sigma^2 4\sigma^2 1\pi^3 5\sigma^2 2\pi^1 \\ B^1\Sigma^+ \text{ or } C^1\Sigma^+ &: && 1\sigma^2 2\sigma^2 3\sigma^2 4\sigma^2 1\pi^4 5\sigma^1 (6\sigma^1 \text{ or } 7\sigma^1) \end{aligned}$$

The spin-orbit and L -uncoupling perturbations between the $A^1\Pi$ state and the $a^3\Sigma^+$, $e^3\Sigma^-$, $d^3\Delta_i$, and $I^1\Sigma^-$, $D^1\Delta$ states, respectively, involve a $1\pi \rightarrow 5\sigma$ orbital promotion. The possible L -uncoupling perturbation between the $A^1\Pi$ state and the $B^1\Sigma^+$ state involves a $2\pi \rightarrow 6\sigma$ orbital promotion. The magnitude of this $2\pi \rightarrow 6\sigma$ perturbation matrix element is likely to be weaker than that for the L -uncoupling $1\pi \rightarrow 5\sigma$ perturbations because the 1π and 5σ orbitals are mostly of O-atom localized $2p$ character while the 2π orbital resembles a $d\delta$ Rydberg orbital and the 6σ orbital is an $s\sigma$ Rydberg orbital. The possible perturbation between the $e^3\Sigma^-$ and $B^1\Sigma^+$ states involves a $(2\pi, 5\sigma) \rightarrow (1\pi, 6\sigma)$ double-promotion. Since the spin-orbit operator is a one-electron operator, which has a $\Delta(s-o) = 0$ or 1 change in spin-orbital ($s-o$) selection rule, the $e \sim B$ spin-orbit perturbation matrix element is zero in first-order. A second-order interaction, mediated by a nearby vibrational level of the $a^3\Pi$ or $A^1\Pi$ state, could result in a very weak $e \sim a/A \sim B$ indirect perturbation. The possibility of a perturbation of the $B^1\Sigma^+$ state by *direct* spin-orbit interaction with the $a^3\Sigma^+$ state can be ruled out by the $\Delta\Omega = 0$ and $\pm \leftrightarrow \pm$ parity selection rule for spin-orbit perturbations. The $a^3\Sigma^+$ state contains only an 0^- component and the $B^1\Sigma^+$ state consists exclusively of an 0^+ component. It is in principle possible to have an *indirect* interaction between the a' and B states “mediated by” the $\Omega = 1$ component of an intermediate $a^3\Pi$ or $A^1\Pi$ state:

$$\langle {}^3\Sigma_1^+ | \mathbf{H}^{SO} | {}^3\Pi_1 \rangle \langle {}^3\Pi_1 | \mathbf{H}^{rot}(S - uncoupling) | {}^3\Pi_0 \rangle \langle {}^3\Pi_0 | \mathbf{H}^{SO} | {}^1\Sigma_0^+ \rangle, \quad (1)$$

$$\text{or} \quad \langle {}^3\Sigma_1^+ | \mathbf{H}^{SO} | {}^1\Pi_1 \rangle \langle {}^1\Pi_1 | \mathbf{H}^{rot}(L - uncoupling) | {}^1\Sigma_0^+ \rangle. \quad (2)$$

These two perturbation paths can be ruled out by the small computed values of the vibrational overlap integrals between near-degenerate vibrational levels of the B and A or a electronic states, $\langle v_B | v_A \rangle$ and $\langle v_B | v_a \rangle$. The possibility of a direct spin-orbit interaction between the $B^1\Sigma^+$ state and a near-degenerate vibrational level of the $a^3\Pi_{0+}$ state can be ruled out by the small value of the computed $\langle v_B | v_a \rangle$ vibrational overlap integral.

It is always possible to evaluate a matrix element using basis functions expressed in any Hund's coupling case. For many reasons, we find case (a) to be most convenient for valence states of almost all diatomic molecules. The spin-orbit operator is diagonal in Ω in Hund's case (a). The spin-orbit matrix element between pure $\Omega = 0$ and 1 states is rigorously zero. This is true for all spin and all Λ basis states. All terms in the field free Hamiltonian are diagonal in parity. The eigenstates of a $^3\Sigma^+$ state for each value of J (J is rigorously conserved at zero electric field; nonzero nuclear spins can weakly break the J quantum number, but this effect is negligible for $^{13}\text{C}^{17}\text{O}$) are linear combinations of $\Omega = 0$ and 1 for f -symmetry and pure $\Omega = 1$ for e -symmetry. The $\Omega = 0$ basis state of $^3\Sigma^+$ has f symmetry but the $^1\Sigma^+$ state consists exclusively of e -symmetry $\Omega = 0$ basis states. Hence, there can be no perturbation of $^1\Sigma^+$ by a $^3\Sigma^+$ state. Contrary to the opinion of Kovács [71] (Fig. 4.55 p. 267), the F_2 component of a $^3\Sigma^+$ state is a pure $\Omega = 1$ state. It cannot perturb $^1\Sigma^+$ via spin-orbit because that is a pure $\Omega = 0$ state. In order for such a perturbation to occur, there needs to be a third electronic state, either $^3\Pi$ or $^1\Pi$, that can perturb via a combination of L -uncoupling and spin-orbit mechanisms. Also, rotational interactions between rovibronic levels of the $^1\Sigma^+$ state and levels of $^3\Delta$, $^1\Delta$, and $^1\Sigma^-$ states are forbidden by the $\Delta S = 0$, $\Delta\Omega = 0, \pm 1$, and $\pm \leftrightarrow \pm$ parity selection rules for rotational (L.J "L-uncoupling" and L.S "spin-electronic") interaction terms [36,72].

3.2. Irregularities of the $^{13}\text{C}^{17}\text{O}$ $A^1\Pi$ and $B^1\Sigma^+$ states

Predicted maxima of perturbations and identification of responsible perturbers result from the rovibronic term-crossing diagrams of the $^{13}\text{C}^{17}\text{O}$ $A^1\Pi$ ($v = 0 - 3$), $d^3\Delta$ ($v = 4 - 9$), $e^3\Sigma^-$ ($v = 1 - 6$), $a^3\Sigma^+$ ($v = 9 - 14$), $I^1\Sigma^-$ ($v = 0 - 5$), and $D^1\Delta$ ($v = 0 - 4$) levels, which were presented in our previous works [33–35]. A term-crossing diagram of the $^{13}\text{C}^{17}\text{O}$ $B^1\Sigma^+$ ($v = 0$) level containing irregularities and close lying $A^1\Pi$ ($v = 20 - 22$), $e^3\Sigma^-$ ($v = 29 - 33$), $a^3\Sigma^+$ ($v = 39 - 41$), and $a^3\Pi$ ($v = 31 - 32$) levels is shown in Fig. 6. The diagram was plotted on the basis of $^{13}\text{C}^{17}\text{O}$ equilibrium molecular constants given by Hakalla et al. [34] for the $A^1\Pi$ state, isotopically recalculated equilibrium molecular constants given by Field [42] for the $d^3\Delta$, $e^3\Sigma^-$, $a^3\Sigma^+$, and $I^1\Sigma^-$ states, by Havenith et al. [73] for the $a^3\Pi$ state, as well as by Kittrell et al. [74] for the $D^1\Delta$ state. The T_e values were taken from Refs. [74–76] and the $G(v = 0)$ value of the $^{13}\text{C}^{17}\text{O}$ $X^1\Sigma^+$ ground state comes from Coxon et al. [77]. Note that in some cases we will deal with observable effects of interactions that originate from some perturbers that do not intersect with the $A^1\Pi$ term curves.

Thanks to high accuracy VIS and VUV FTS measurements carried out in the $^{13}\text{C}^{17}\text{O}$ isotopologue, weak ($0.05 - 0.75 \text{ cm}^{-1}$) irregularities appearing in the $B^1\Sigma^+$ ($v = 0$) level at $J = 0, 7, 16, 21, 26,$ and $34 - 36$ were noticed. However, identification of the states responsible for these possible perturbations is complicated due to indirect interaction problems discussed in the previous subsection. We were not able to determine the dependence between the effective perturbation parameters and isotopologue-independent \mathbf{a} and \mathbf{b} parameters [55] for the $B^1\Sigma^+ \sim$

perturbers' interactions, because relevant single-orbital perturbation matrix element values were too small to extract perturbation parameters. In addition, the studied region is close to the (${}^3P_0 + {}^3P_2$) first dissociation limit of the CO molecule ($90,679 \text{ cm}^{-1}$ [78]), where many singlet and triplet states converge, forming a region with very densely spaced vibrational levels, which complicates the identification of a perturber. Moreover, the vibrational overlap integrals involving the $B^1\Sigma^+$ state and the $A^1\Pi$ and/or $a^3\Pi$ states as mediating ones are very small, at about $1 \times 10^{-3} - 1 \times 10^{-5}$. Therefore, an identification of the $B^1\Sigma^+$ perturbers was not possible. Thus, the transition frequencies which involved the unidentified perturbers were appropriately weighted in the final deperturbation fit, so that they would have a negligible influence on the quality of the results.

3.3. Deperturbation analysis

We model each observed ${}^{13}\text{C}^{17}\text{O}$ $B^1\Sigma^+(v=0, 1) \rightarrow A^1\Pi(v=0-3)$ [33,34], $C^1\Sigma^+(v=0) \rightarrow A^1\Pi(v=1-3)$ [35], $B^1\Sigma^+(v=0, 1) \leftarrow X^1\Sigma^+(v=0)$, and $C^1\Sigma^+(v=0) \leftarrow X^1\Sigma^+(v=0)$ band and interacting levels based on a local deperturbation analysis, similar as was done by Hakalla et al. [55]. In total, 1003 transitions from 12 bands of ${}^{13}\text{C}^{17}\text{O}$ were used in the procedure of disentangling $A^1\Pi$ from the $d^3\Delta$, $e^3\Sigma^-$, $a^3\Sigma^+$, $I^1\Sigma^-$, and $D^1\Delta$ states using the Pgopher software [79]. This procedure results in 63 deperturbed molecular parameters fitted for the less-abundant ${}^{13}\text{C}^{17}\text{O}$ isotopologue. The computed level positions, line wavenumbers and intensities are the result of a matrix diagonalization taking into account all interacting levels. The assignment of $A^1\Pi(v)$ and perturber levels, the selection of which interactions and parameters could be determined from the spectra, and their values, were iteratively optimized. The Pgopher program [79] is based on an effective Hamiltonian with matrix elements similar to Field [42], Bergeman et al. [80], and Le Floch et al. [44], as we presented in Ref. [55], retaining their symbols for the various molecular constants. The Hamiltonian matrix used in the present deperturbation analysis contains not only the perturber levels which cross the $A^1\Pi(v)$ level, but also those which interact with it without intersections. The wavenumbers of intense, moderately strong and isolated spectral lines were assigned relative weights of 1.0 during the fitting. However, the frequencies of the weak and/or blended lines are less accurate, so they were individually weighted between 0.5 – 0.1, depending on their weakness and/or extent of blending as well as failure to recognize the causes of irregularities in B ($v=0$).

The homogeneous interactions of the short-lived $A^1\Pi$ singlet levels with the long-lived $e^3\Sigma^-$, $a^3\Sigma^+$, and $d^3\Delta_i$ triplet levels result from the spin-orbit term coupling, represented by J -independent matrix elements. In turn, the interactions of $A^1\Pi$ with $I^1\Sigma^-$ and $D^1\Delta$ are caused by L -uncoupling and, therefore, produce heterogeneous interactions with J -dependent matrix elements [36]. The interaction parameters are denoted by $\eta_{v,v'}$ and $\xi_{v,v'}$ for the spin-orbit and L -uncoupling perturbations, respectively.

For some of the levels under consideration, there are insufficient term-value data to determine molecular constants independently of the deperturbation calculation. In such cases, we adopted isotopically recalculated molecular constants of these $d^3\Delta$, $e^3\Sigma^-$, $a^3\Sigma^+$, $I^1\Sigma^-$, and $D^1\Delta$ states, using Dunham's relationship within the Born–Oppenheimer approximation [81] from the data given in Refs. [42,74]. These values were held fixed during the fitting procedures. In the calculations we included all possible perturber levels, listed in Table 4, which have a non-negligible influence on the $^{13}\text{C}^{17}\text{O}$ $A^1\Pi$ ($v = 0-3$), even if some of them do not exhibit crossings, as it is possible to see in Fig. 7. The non-negligible influence means a shifting of rotational level with the value exceeding the mean residuals value.

The unweighted residuals of the fits are dominated by the uncertainties of the weakest B–A (1, 0) and (1, 1) bands obtained in Ref. [34]. The mean weighted residuals value of the deperturbation calculation based on the data of VIS–VUV FTS, is 0.007 cm^{-1} , so the fitting model acceptably reproduces the experimental data set.

In some cases, it was impossible to calculate the $\eta_{v,v'}$ and $\xi_{v,v'}$ interaction parameters, because their fitting was statistically unjustified. This occurs when we deal with an insufficient quantity of experimental transitions attributed to perturber states in the neighborhood of the avoided crossings or when interacting states are too distant. The solution to this problem was a semi-empirical method of estimation of the missing interaction parameters provided in Refs. [36,42–44]. It shows that the perturbation matrix element (α , β) is the product of a vibrational factor and an electronic perturbation parameter (\mathbf{a} , \mathbf{b}). The isotopologue-independent parameters \mathbf{a} and \mathbf{b} are characteristic of the relevant electronic configuration within the single-configuration approximation and they are strongly dependent on the internuclear separation of the molecule [82]. An approach to calculate of $\eta_{v,v'}$ and $\xi_{v,v'}$ for the $A^1\Pi(v)$ state and its interacting levels is described in our previous work [55]. The isotopologue-independent purely electronic perturbation parameters $\mathbf{a}_{A-d,e,a'}$ and $\mathbf{b}_{A-I,D}$ were taken from Le Floch [44]. The $\langle v_A | v_{d,e,a'} \rangle$ vibrational overlap integrals, as well as the $\langle v_A | \mathbf{B}(\mathbf{R}) | v_{I,D} \rangle$ rotational operator integrals of $^{13}\text{C}^{17}\text{O}$, were calculated on the basis of RKR parameters for the $^{13}\text{C}^{17}\text{O}$ A, d, e, a', I, and D states. The RKR values were obtained from isotopically recalculated equilibrium constants of Field [42], Field et al. [36,43], Le Floch et al. [44], Kittrell et al. [74] and Hakalla et al. [34] using the computer program 'FRACONB' of Jung [83] and Jakubek [84]. Finally, the missing perturbation parameters under consideration were estimated and held fixed during the fitting.

As expected from the extensive coverage of the $^{13}\text{C}^{17}\text{O}$ transition energies in the 12 bands, we fitted accurate deperturbed molecular constants for the $A^1\Pi$ ($v = 0, 1, 2$, and 3) levels and their perturbers, where the definitions of the constants are the same as in Ref. [47,48,55]. The results are collected in Table 4. We also obtained the molecular constants for the $B^1\Sigma^+$ ($v = 0, 1$) and $C^1\Sigma^+$ ($v = 0$) levels, results of which are given in Table 5. The constants of the $B^1\Sigma^+$ ($v = 0, 1$) and $C^1\Sigma^+$ ($v = 0$) levels are compared with analogous values derived in previous studies [33–35]. The

correlation matrix was checked carefully and it was found that it shows satisfactorily low correlations between fitted parameters of the model.

All perturbation parameters that are the result of the global deperturbation analysis of the $^{13}\text{C}^{17}\text{O}$ isotopologue are listed in Table 6. The electronic part of each perturbation parameter was obtained by dividing the spin–orbit parameter $\eta_{v,v'}$ by the vibrational overlap $\langle v|v'\rangle$, or by dividing the rotation–electronic parameter $\xi_{v,v'}$ by the rotational operator integral $\langle v|\mathbf{B}(r)|v'\rangle$. Table 6 also includes r –centroids, which are indicators of the internuclear distance of maximum vibrational overlap. It may be noted that for perturbations between any pair of electronic states, the r –centroids are approximately constant. In this table, for comparison, we also included the results of all the tests to determine the isotopologue–independent electronic perturbation parameters \mathbf{a} and \mathbf{b} , which have been carried out so far, that is, by Field et al. [43], Le Floch [44], Haridass et al. [52], and Hakalla et al. [55].

3.4. Level energies

The rotational terms of the $^{13}\text{C}^{17}\text{O}$ $\text{A}^1\Pi$ ($v = 0, 1, 2$ and 3) vibrational levels referenced to the $\text{X}^1\Sigma^+$ ($v = 0, J = 0$) ground state were calculated as differences of the $\text{B}^1\Sigma^+$ ($v = 0, 1$), $\text{C}^1\Sigma^+$ ($v = 0$) term values obtained from a VUV–FTS experiment [58] and $\text{B} \rightarrow \text{A}$ ($0 - v''$) (this work and [33,34]), $\text{C} \rightarrow \text{A}$ ($0 - v''$) [35] transition frequencies. The $^{13}\text{C}^{17}\text{O}$ ($\text{B}^1\Sigma^+, \text{C}^1\Sigma^+$) \rightarrow ($\text{d}^3\Delta_i, \text{e}^3\Sigma^-, \text{a}^3\Sigma^+, \text{I}^1\Sigma^-, \text{D}^1\Delta$) extra–lines, listed in Table 2, were used to determine level energies of the D, I, e, a', and d states in $^{13}\text{C}^{17}\text{O}$. The higher J –terms of the $^{13}\text{C}^{17}\text{O}$ $\text{A}^1\Pi$ ($v = 0 - 3$) levels were calculated using deperturbed T_v constants from Table 4 and relative level energies of the $\text{A}^1\Pi$ state obtained from $\text{B} \rightarrow \text{A}$ ($0 - v''$) [33,34], $\text{C} \rightarrow \text{A}$ ($0 - v''$) [35] bands by means of the linear least–squares method in the version given by Curl and Dane [85] and Watson [86]. The final term values were merged by using the weighted average method. The values for the $\text{A}^1\Pi$ ($v = 0, 1, 2$ and 3) terms are collected in Table 7 and for the $\text{d}(v), \text{e}(v), \text{a}'(v), \text{D}(v),$ and $\text{I}(v)$ levels in Table 8. The results for $\text{B}^1\Sigma^+$ ($v = 0, 1$), $\text{C}^1\Sigma^+$ ($v = 0$), and other levels in the 115–101 nm range will appear in a forthcoming publication [58].

The reduced term values $T(J) - BJ(J+1) + DJ^2(J+1)^2$ of the perturbed $^{13}\text{C}^{17}\text{O}$ $\text{A}^1\Pi$ ($v = 0 - 3$) and $\text{B}^1\Sigma^+$ ($v = 0$) levels, in relation to the $v = 0, J = 0$ level of the $\text{X}^1\Sigma^+$ ground state, together with the hypothetical unperturbed and crossing perturber levels, are presented in Fig. 7. These reduced terms of perturbers, which could not be derived from our current analysis, were calculated on the basis of isotopically recalculated equilibrium molecular constants given by Field [42] for $\text{d}^3\Delta_i, \text{e}^3\Sigma^-, \text{a}^3\Sigma^+$, and $\text{I}^1\Sigma^-$ states, by Kittrell et al. [74] for $\text{D}^1\Delta$ state as well as by Havenith et al. [73] for the $\text{a}^3\Pi$ state. T_e values were taken from Refs. [74–76] and the G ($v = 0$) value of $\text{X}^1\Sigma^+$ from Coxon et al. [77].

It should be noted that one could make a distinction between accuracy and precision. The first referring to the absolute wavenumber scale of importance for deriving term energies with

respect to X ($v = 0, J = 0$) state (not better than the 0.04 cm^{-1}) and the other refers to the fitting with respect to the relative scale (at about 0.003 cm^{-1}). The assessment of some of most of the perturbations relies (only) on this.

4. Discussion

The high accuracy VIS–VUV FTS measurements carried out on the $^{13}\text{C}^{17}\text{O}$ isotopologue reveal very weak irregularities appearing in the $\text{B}^1\Sigma^+$ ($v = 0$) level and contribute to a successful global deperturbation fit. As a result accurate deperturbed molecular constants, term values, spin–orbit and L –uncoupling interaction parameters as well as isotopologue–independent spin–orbit and rotation–electronic perturbation parameters of the $^{13}\text{C}^{17}\text{O}$ $\text{A}^1\Pi$ ($v = 0, 1, 2$ and 3) state could be derived. The analysis allowed also for verification and improvement of the perturbations in the $^{13}\text{C}^{17}\text{O}$ $\text{A}^1\Pi$ ($v = 0 - 3$) levels reported in Ref. [33–35]. Comparing Table 4 with the data given in the mentioned publications, we can see that the $\text{I}^1\Sigma^-$ ($v = 0$), $\text{d}^3\Delta_i$ ($v = 5$), and $\text{e}^3\Sigma^-$ ($v = 5$) rovibronic levels have a weak impact on the $v = 0, v = 1$, and $v = 3$ levels of the $\text{A}^1\Pi$ state, respectively. The reduced terms presented in Fig. 7 indicate that the rotational progressions of these perturbers do not cross with the $\text{A}^1\Pi$ ($v = 0, 1$, and 3) levels, and as a result they cause global energy shifts of the $\text{A}^1\Pi$ (v) levels. This phenomenon is precisely the same as a vibrational perturbation shifting an entire band systematically. Also, all J 's of maximum irregularities of the $\text{A}^1\Pi$ ($v = 0, 1$, and 3) levels, which have been observed so far [33–35], were verified.

We tried to find the alleged perturbers responsible for irregularities appearing in the $\text{B}^1\Sigma^+$ ($v = 0$) level. As shown in Fig. 7a, a possible perturber of B ($v = 0, J = 7$) level could be the $\text{e}^3\Sigma^-$ ($v = 29$) state, but this would require a second–order interaction beyond the single–configuration picture of diabatic electronic states, or else a mediating interaction with a close lying $\text{a}^3\Pi$ or $\text{A}^1\Pi$ state, i.e. $\text{a}(31)$ or $\text{A}(19)$ or $\text{A}(20)$. However, *i*) the $\text{a}(31)$, $\text{A}(19)$, and $\text{A}(20)$ levels lie too far from $\text{e}(29)$, *ii*) the $^{13}\text{C}^{17}\text{O}$ $\langle v_{\text{B}(0)} | \mathbf{B}(\mathbf{R}) | v_{\text{a}(31)} \rangle$, $\langle v_{\text{B}(0)} | \mathbf{B}(\mathbf{R}) | v_{\text{A}(19)} \rangle$, and $\langle v_{\text{B}(0)} | \mathbf{B}(\mathbf{R}) | v_{\text{A}(20)} \rangle$ rotational operator integrals equal $-9.3 \times 10^{-5} \text{ cm}^{-1}$, $-5.8 \times 10^{-3} \text{ cm}^{-1}$, and $4.2 \times 10^{-3} \text{ cm}^{-1}$, respectively, so they are too small to cause the expected effect. Similarly for B ($v = 0, J = 21$), where the $\text{A}(20)$ state is too distant to cause the $\text{B}(0) \sim \text{A}(20) \sim \text{e}(30)$ indirect perturbation and the $\langle v_{\text{B}(0)} | \mathbf{B}(\mathbf{R}) | v_{\text{A}(20)} \rangle$ vibrational overlap integral is too small ($-9.3 \times 10^{-5} \text{ cm}^{-1}$). The $J = 34$ level of B(0) could be affected by an indirect $\text{B}(0) \sim \text{a}(32) \sim \text{e}(32)$ perturbation, but $\langle v_{\text{B}(0)} | \mathbf{B}(\mathbf{R}) | v_{\text{a}(32)} \rangle = 7.7 \times 10^{-5} \text{ cm}^{-1}$ is far too small. Finally, the $J = 35$ level of B(0) could be directly perturbed via an L –uncoupling interaction with $\text{A}(21)$, but the vibrational overlap integral, $\langle v_{\text{B}(0)} | \mathbf{B}(\mathbf{R}) | v_{\text{A}(20)} \rangle = -3.1 \times 10^{-3} \text{ cm}^{-1}$, is also not big enough for this purpose. We note that the term values of the states suspected of perturbing the $\text{B}^1\Sigma^+$ ($v = 0$) level were calculated on the basis of long extrapolations to high rovibrational levels using equilibrium molecular constants. This results in large uncertainties in the level positions

determinations. For all these reasons we could not unequivocally find a cause for the alleged perturbations in the $B^1\Sigma^+$ ($v = 0$) level. Therefore, the molecular constants obtained for this level, presented in Table 5, should be considered as effective ones. The only other possibility is that another hitherto unknown Rydberg state is involved, of $^1\Pi$, $^3\Pi$, or $^3\Sigma^-$ character. We hope, that our future work [87] concerning a multi-isotopologue analysis of the CO molecule may shed more light on this problem.

The q_v constant of the $A^1\Pi$ ($v = 0$) level, describing the Λ -doublet splitting and presented in Table 4, has a large and erratic value as compared to other q_v values of this state. The problems with this constant for the $A^1\Pi$ state are well-known [44,47,48]. The reason for the unexpected behavior could be the strong dependence of the q_v constant on local perturbations [44] as well as negligible Λ -doubling effects excluding the crossing regions of the $A^1\Pi$ state with its perturbers.

We should note that the perturbation parameters that control the interactions between the $d^3\Delta_i$, $e^3\Sigma^-$, $a^3\Sigma^+$ states and the $A^1\Pi$ state are the cumulative result of the interactions of all three Ω -components of triplet states. It is obvious, however, that each of the F_1 , F_2 , and F_3 components interacts with appropriate components of the perturbed state to a varying degree. It is clearly seen in Fig. 7b for the $A^1\Pi$ ($v = 0$) \sim $d^3\Delta_i$ ($v = 4$) interaction, where the second and third Ω -component of the $d^3\Delta_i$ ($v = 4$) state causes more severe deviations (about 4.5 cm^{-1}) of the $A^1\Pi$ ($v = 0$) level than the first one (about 0.5 cm^{-1}). One of the causes could be the fact that strong spin-orbit interaction between the states in the Hund's case (a) requires $\Delta\Omega = 0$. Thus, the third component ($\Omega = 1$) of the $d^3\Delta_i$ state should have the strongest influence on the $A^1\Pi$ ($\Omega = 1$) state and that is what we have observed.

In Table 5 one can notice minor inconsistencies in the B_v and D_v rotational constants of the $B^1\Sigma^+$ ($v = 0, 1$) and $C^1\Sigma^+$ ($v = 0$) levels in relation to our previous results [33–35]. These minor discrepancies may result from: *i*) the first attempt of analysis of the weak perturbations occurring at $^{13}\text{C}^{17}\text{O}$ $B^1\Sigma^+$ ($v = 0$) level, *ii*) the differences in calculation methods used to obtain the constants, as we described in Ref. [55], and/or *iii*) band measurements that include higher- J levels, having a significant impact on the rotational distortion D_v constants, always strongly correlated with both B_v and T_v constants.

Table 6 shows the perturbation parameters obtained in this work. The isotopic independence of the a spin-orbit and b rotation-electronic perturbation parameters for each state that interacts with the $A^1\Pi$ state is evident. The parameters also remain satisfactorily consistent with analogous values determined in the deperturbation calculations of $^{12}\text{C}^{16}\text{O}$ [43,44], $^{12}\text{C}^{18}\text{O}$ [52], and $^{12}\text{C}^{17}\text{O}$ [55].

5. Conclusion

Two high precision experimental methods, photoemission in a discharge with Fourier–transform spectroscopy in the visible region as well as synchrotron absorption FTS in the vacuum ultraviolet region, were used to obtain first high–resolution spectra of the $B^1\Sigma^+ \rightarrow A^1\Pi(0, 0)$, $(0, 3)$, and the $B^1\Sigma^+ \leftarrow X^1\Sigma^+(0, 0)$, $(1, 0)$, $C^1\Sigma^+ \leftarrow X^1\Sigma^+(0, 0)$ bands, respectively. Combining the new information and our previous results [33–35], in total 1003 rotational lines in 12 bands of $^{13}\text{C}^{17}\text{O}$ ($B^1\Sigma^+ \rightarrow A^1\Pi$, $C^1\Sigma^+ \rightarrow A^1\Pi$, $B^1\Sigma^+ \leftarrow X^1\Sigma^+$, $C^1\Sigma^+ \leftarrow X^1\Sigma^+$) and in 18 bands of extra–lines, and literature data on CO isotopologues, allowed for a global deperturbation analysis of this complex of interacting levels and the fitting of model energy levels to all experimental data within their uncertainties. 63 independent adjustable parameters were required for this fit. The model describes the $A^1\Pi(v = 0 - 3)$, $B^1\Sigma^+(v = 0, 1)$, $C^1\Sigma^+(v = 0)$, $d^3\Delta(v = 4, 7, 8)$, $e^3\Sigma^-(v = 1, 2, 4)$, $a^3\Sigma^+(v = 10, 12, 13)$, $I^1\Sigma^-(v = 2, 3, 5)$, and $D^1\Delta(v = 1, 4)$ levels, as well as the spin–orbit and L –uncoupling parameters responsible for the interactions of the $A^1\Pi(v = 0 - 3)$ levels with their perturbers. Isotopologue–independent perturbation parameters of the $A^1\Pi$ interactions were also derived. With these combined data, 335 rotational terms of $A^1\Pi(v = 0 - 3)$ levels and their perturbers were determined up to $J = 36$.

The CO A–state has been a case study for diatomic molecular deperturbation studies. The mass–dependent shifts of energy levels mean each isotopologue exhibits different perturbation effects, but each arises from the same underlying rotational and spin–orbit interactions. The present study of $^{13}\text{C}^{17}\text{O}$ provides an additional perspective on the non–Born–Oppenheimer and relativistic effects common to all isotopologues.

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Table 1

$B^1\Sigma^+ \rightarrow A^1\Pi(0, 0)$ and $(0, 3)$ transitions identified in the VIS–FT spectra of $^{13}\text{C}^{17}\text{O}$.^{a, b, c}

J''	$B^1\Sigma^+ \rightarrow A^1\Pi(0, 0)$			$B^1\Sigma^+ \rightarrow A^1\Pi(0, 3)$		
	$P_e(J'')$	$Q_f(J'')$	$R_e(J'')$	$P_e(J'')$	$Q_f(J'')$	$R_e(J'')$
1	22159.76 ^{wb}	22163.32 ^b	22170.60 ^{wb}	17960.852	17964.481	17971.728
2	22157.590	22164.72 ^b	22175.717	17958.758	17966.010	17976.886
3	22156.280	22166.84 ^b	22181.662	17957.436	17968.310	17982.811
4	22155.944	22169.721	22188.573	17956.87 ^b	17971.374	17989.498
5	22156.963	22173.42 ^b	22196.853	17957.074	17975.20 ^b	17996.969
6	22147.991	22178.013	22195.16 ^b	17958.043	17979.817	18005.218
7	22150.998	22183.728	22205.38 ^b	17959.802	17985.19 ^b	18014.18 ^b
8	22153.749	22190.78 ^b	22215.342	17962.339	17991.322	18023.931
9	22156.698	22182.374	22225.560	17965.595	17998.229	18034.457
10	22160.24 ^b	22190.78 ^b	22236.344	17969.647	18005.912	18045.75 ^b
11	22164.62 ^b	22199.11 ^b	22247.964	17974.476	18014.378	18057.816
12	22170.255	22207.661	22260.829	17980.095	18023.654	18070.668
13	22165.997	22216.589	22263.797	17986.490	18033.79 ^b	18084.293
14	22173.42 ^b	22226.000	22278.472	17993.710	18045.329	18098.740
15	22180.607	22235.959	22292.854	18001.828	18054.479	18114.079
16	22187.95 ^b	22246.504	22307.416	18011.209	18067.406	18130.677
17	22195.722	22257.670	22322.404	18016.149	18080.649	18142.829
18	22204.059	22269.485	22337.954	18028.288	18094.594	18162.178
19	22213.056	22282.006	22354.144	18039.67 ^b	18109.531	18180.770
20	22223.03 ^b	22295.519	22371.317	18051.579	18122.03 ^b	18199.867
21	22231.94 ^b	22307.998	22387.43 ^b	18064.156	18139.766	18219.646
22	22243.48 ^b	22323.108	22406.144	18077.480	18156.885	18240.151
23	22255.776	22338.982	22425.622	18091.557	18174.610	18261.402
24	22270.567	22357.333	22447.60 ^b	18106.390	18193.058	18283.405
25	22274.727	22365.085	22458.93 ^b	18121.97 ^b	18212.258	18306.173
26	22290.412	22384.336	22481.76 ^b	18138.361	18232.214	18329.702
27	22306.03 ^w	22403.538		18155.503	18252.94 ^b	18354.001
28		22425.06 ^{wb}		18173.39 ^b	18274.436	18379.057
29				18192.108	18296.699	18404.887
30				18211.56 ^b	18319.73 ^b	18431.483
31				18231.824	18343.537	18458.85 ^b
32				18252.84 ^b	18368.127	18486.97 ^b
33				18274.657	18393.488	18515.90 ^{wb}
34				18297.291	18419.678	18545.63 ^{wb}
35				18320.53 ^w	18446.46 ^b	18575.95 ^{wb}
36					17964.48 ^b	18607.19 ^{wb}

^a All values in cm^{-1} .

^b Lines marked with ‘w’ and/or ‘b’ were weak and/or blended, respectively.

^c The estimated absolute calibration 1σ uncertainty was 0.003 cm^{-1} . The expected absolute accuracy of the line frequency measurements was 0.003 cm^{-1} for the medium–strength and single lines as well as $0.01\text{--}0.02 \text{ cm}^{-1}$ for the weak/blended lines.

Table 2

($C^1\Sigma^+$, $B^1\Sigma^+$) \rightarrow ($e^3\Sigma^-$, $d^3\Delta_i$, $a^3\Sigma^+$, $I^1\Sigma^-$, and $D^1\Delta$) transitions arising from perturbations of the $A^1\Pi$ ($v = 0 - 3$) levels in $^{13}C^{17}O$.^{a, b}

System	Band	J''	Branch / Wavenumber (cm^{-1})						
$B^1\Sigma^+ \rightarrow e^3\Sigma^-$	(1, 1) ^c	4	$^qP_{11ee}$	24151.28 ^{wb}					
		5		24154.85 ^{wb}	$^qQ_{12ef}$	24161.09 ^{wb}			
		6		24170.326		24167.77 ^b	$^sR_{11ee}$	24194.23 ^{wb}	
		7		24175.73 ^b		24175.250		24216.87 ^b	
		8	$^oP_{13ee}$	24140.18 ^{wb}		24183.193		24229.44 ^b	
		9		24145.16 ^{wb}		24208.545			
		10		24151.262		24218.995			
		11		24158.27 ^b		24231.37 ^w			
		12		24165.84 ^b			$^qR_{13ee}$	24255.25 ^b	
		13		24185.171				24281.720	
		14		24194.66 ^b				24298.35 ^{wb}	
		15		24206.33 ^b					
		(0, 1)	2	$^qP_{11ee}$	22135.25 ^w	$^qQ_{12ef}$	22136.38 ^b		
			3		22137.46 ^b		22140.06 ^b		
			4		22140.608		22144.87 ^b	$^sR_{11ee}$	22173.23 ^{wb}
	5			22144.361		22150.82 ^b		22184.24 ^{wb}	
	6					22157.78 ^b			
	7					22165.603			
	8		$^oP_{13ee}$	22130.53 ^{wb}		22173.907			
	9			22135.89 ^{wb}		22199.69 ^b			
	10			22142.39 ^{wb}		22210.596	$^qR_{13ee}$	22218.49 ^b	
	11			22149.871		22223.456		22233.203	
	12			22157.98 ^b		22238.06 ^b		22248.541	
	13			22177.86 ^b		22254.20 ^{wb}		22275.66 ^b	
	14			22187.95 ^b				22292.99 ^w	
	15							22312.52 ^{wb}	
	(0, 4)		27			$^qQ_{12ef}$	19590.088		
		28				19637.20 ^b			
	$B^1\Sigma^+ \rightarrow d^3\Delta$	(1, 4)	20		$^iQ_{11ef}$	24286.66 ^{wb}			
			21						
22					$^qQ_{12ef}$	24281.56 ^{wb}			
23			$^pP_{12ee}$	24228.09 ^{wb}					
24							$^qR_{13ee}$	24427.23 ^{wb}	
(0, 4)		21			$^iQ_{11ef}$	22316.56 ^{wb}			
		22							
		23					$^qR_{13ee}$	22398.85 ^b	
		24			$^qQ_{12ef}$	22341.319			
		25				22383.13 ^{wb}			
(0, 7)		31			$^iQ_{11ef}$	19668.59 ^b			
		32				19716.43 ^b			
(0, 8)		33	$^qP_{11ee}$	19645.78 ^b					
		10	$^qP_{11ee}$	17991.92 ^b					
		...							
16		$^oP_{13ee}$	17991.88 ^b						
$B^1\Sigma^+ \rightarrow a^3\Sigma^+$	(0, 10)	2			$^pQ_{13ef}$	20734.67 ^w			
		3				20736.98 ^w			
		4	$^pP_{12ee}$	20738.249		20740.61 ^{wb}			
	(0, 13)	14			$^iQ_{11ef}$	18036.857			
		15				18064.289			
		16	$^pP_{12ee}$	17997.984			$^iR_{12ee}$	18117.449	
		17		18024.949				18151.630	
		18		18046.91 ^b					
		19			$^pQ_{13ef}$	18096.55 ^w			
		20				18128.001			
$C^1\Sigma^+ \rightarrow a^3\Sigma^+$	(0, 10)	1			$^pQ_{13ef}$	25735.38 ^b			
		2				25736.48 ^b			
		3				25738.77 ^b			
		4				25742.34 ^b			
	(0, 13)	13			$^iQ_{11ef}$	23015.31 ^b			
		14				23037.92 ^b			
		15				23065.26 ^b	$^iR_{12ee}$	23090.90 ^b	
		16	$^pP_{12ee}$	22998.92 ^{wb}		23091.11 ^b		23118.17 ^b	
		17		23025.82 ^{wb}				23152.26 ^b	
		18							

		19						
		20		${}^pQ_{13ef}$	23128.36 ^b			
		21						
		22			23187.95 ^b			
$B^1\Sigma^+ \rightarrow I^1\Sigma^-$	(0, 2)	35		${}^qQ_{11ef}$	21148.89 ^b			
		36			21199.64 ^b			
	(0, 3)	7		${}^qQ_{11ef}$	19347.351			
		8			19358.594			
(0, 5)	40		${}^qQ_{11ef}$	18556.25 ^{wb}				
	41			18615.13 ^{wb}				
$C^1\Sigma^+ \rightarrow I^1\Sigma^-$	(0, 3)	7		${}^qQ_{11ef}$	24348.96 ^b			
		8			24360.149			
$B^1\Sigma^+ \rightarrow D^1\Delta$	(0, 1)	25		${}^qQ_{11ef}$	20903.11 ^b			
		26			20937.80 ^w			
		27			20973.72 ^{wb}		${}^rR_{11ee}$	21074.75 ^{wb}
		28	${}^pP_{11ee}$	20910.76 ^b		21011.86 ^{wb}		21116.47 ^w
	(0, 4)	31		${}^qQ_{11ef}$	18271.42 ^{wb}			
		32			18317.20 ^{wb}			
		33						
		34	${}^pP_{11ee}$	18290.62 ^w		${}^rR_{11ee}$	18538.961	
$C^1\Sigma^+ \rightarrow D^1\Delta$	(0, 1)	27	${}^pP_{11ee}$	25875.66 ^b		${}^rR_{11ee}$	26073.80 ^b	
		28		25910.06 ^b			26115.35 ^b	

^a All values in cm^{-1} .

^b Lines marked with ‘w’ and/or ‘b’ were weak and/or blended, respectively. The estimated absolute calibration 1σ uncertainty was 0.003 cm^{-1} . The expected absolute accuracy of the line frequency measurements was 0.003 cm^{-1} for the medium–strength and single lines as well as $0.01\text{--}0.02 \text{ cm}^{-1}$ for the weak/blended lines. The left–superscripts *o*, *p*, *r*, *s*, or *q* denote the change of the total angular momentum excluding spin of the perturber states [88].

^c The wavenumbers of the ${}^{13}\text{C}^{17}\text{O } B^1\Sigma^+ \rightarrow e^3\Sigma^- (1, 1)$ band were already published by us in Ref. [34]. In this table, all the lines of the band under consideration are presented which were used in the current deperturbation fit. Some of the line assignments, whose values are given in italics, were corrected in relation to Ref. [34]. It is worth noticing a relevant change in the interpretation of the $J'' = 13$ lines, that is, lines $24185.1712 \text{ cm}^{-1}$ and $24281.7195 \text{ cm}^{-1}$ thought in Ref. [34] to belong to the ${}^{13}\text{C}^{17}\text{O } B^1\Sigma^+ \rightarrow A^1\Pi (1, 0)$ band, however, turned out to be its extra–lines, that is ${}^oP_{13ee}(13)$ and ${}^qR_{13ee}(13)$ of the $B^1\Sigma^+ \rightarrow e^3\Sigma^- (1, 1)$ transition. Thus, lines $24173.3064 \text{ cm}^{-1}$ and $24269.8541 \text{ cm}^{-1}$ [34] turned out to be the $P_{11ee}(13)$ and $R_{11ee}(13)$ lines belonging to the $B^1\Sigma^+ \rightarrow A^1\Pi (1, 0)$ band.

Table 3.

Extended and corrected assignment ^a of the heavily perturbed and/or extremely weak lines of the ¹³C¹⁷O
 $B^1\Sigma^+ \rightarrow A^1\Pi$ and $C^1\Sigma^+ \rightarrow A^1\Pi$ systems.^b

J''	$B^1\Sigma^+ \rightarrow A^1\Pi (0, 1)$			$B^1\Sigma^+ \rightarrow A^1\Pi (1, 0)$			$B^1\Sigma^+ \rightarrow A^1\Pi (1, 1)$		$C^1\Sigma^+ \rightarrow A^1\Pi (0, 1)$	
	$P_{11ec}(J'')$	$Q_{11ef}(J'')$	$R_{11ec}(J'')$	$P_{11ec}(J'')$	$Q_{11ef}(J'')$	$R_{11ec}(J'')$	$P_{11ec}(J'')$	$Q_{11ef}(J'')$	$P_{11ec}(J'')$	$R_{11ec}(J'')$
1										
2				24168.50 ^{wb}		24186.40 ^w				
3										
4										
5						24206.85 ^b				
6	20720.011		20767.15 ^b					22751.599	25721.72 ^b	25768.759
7					24193.36 ^b					
8	20723.16 ^b									
...										
13				24173.31 ^b		24269.85 ^b				
...										
17				24200.32 ^b		24325.35 ^b				
18							22783.20 ^b			
19							22791.741			
20	20798.86b		20947.157						25799.30 ^w	25947.32 ^{wb}
21				24233.02 ^b	24308.05^b		22810.72^w			
22				24243.58^{wb}			22821.17^w			
23				24254.86^b						
24				24268.54^b						
25	20859.569		21043.76 ^b	24271.57^b						
26	20873.795		21065.15 ^b							
27		20986.328								
28	20903.69 ^w	21004.76 ^{wb}	21109.365							
29	20920.255	21024.973	21133.09 ^b							
30	20937.300	21045.543								
31	20954.98 ^b	21066.813								
32	20973.34^b	21088.782								
33	20992.366	21111.478								
34		21134.94 ^b								
35		21159.96^b								
<hr/>										
J''	$B^1\Sigma^+ \rightarrow A^1\Pi (0, 2)$			$C^1\Sigma^+ \rightarrow A^1\Pi (0, 2)$			$C^1\Sigma^+ \rightarrow A^1\Pi (0, 3)$			
	$P_{11ec}(J'')$	$Q_{11ef}(J'')$	$R_{11ec}(J'')$	$P_{11ec}(J'')$	$Q_{11ef}(J'')$	$R_{11ec}(J'')$	$P_{11ec}(J'')$	$Q_{11ef}(J'')$	$R_{11ec}(J'')$	
1					24332.63 ^{wb}		22962.67 ^{wb}		22973.53 ^{wb}	
2					24334.06 ^w		22960.57 ^w		22978.66 ^{wb}	
3					24336.22 ^w		22959.23 ^b		22984.57 ^{wb}	
4					24339.09 ^w		22958.64 ^b		22991.216	
5					24342.678					
6										
7						24352.08 ^b				
8										
9		19362.60 ^b								
...										
18								23095.216		
19								23109.99 ^b		
20								23122.3558		
21								23139.9881		
22				24434.49 ^{wb}				23156.9625		
23	19447.82 ^b		19617.68 ^b	24447.911		24617.46^{wb}		23174.5270		
24			19640.782	24463.49 ^{wb}		24640.19^{wb}		23192.8315		
25	19473.191		19657.374	24472.82 ^{wb}				23211.89^b		
26			19680.883	24489.11 ^{wb}	24585.37 ^b			23231.68^b		
27	19505.94 ^b		19704.435		24607.04^b			23252.18^b		
28	19522.954	19619.564	19728.597		24618.66^b			23273.52^b		
29	19540.750	19642.995	19753.529							
30	19559.624	19665.54 ^b	19779.526							
31	19581.584	19688.515	19808.61^w							
32		19711.44 ^b								
33		19736.081								

^a An extension (in bold) and improvement over the previously analyzed [33–35] data. All values in cm⁻¹.

^b Lines marked with ‘w’ and/or ‘b’ were weak and/or blended, respectively. The estimated absolute calibration 1 σ uncertainty was 0.003 cm⁻¹. The expected absolute accuracy of the line frequency measurements was 0.003 cm⁻¹ for the medium–strength and single lines as well as 0.01–0.02 cm⁻¹ for the weak/blended lines.

Table 4

 Deperturbed rovibrational constants of the $A^1\Pi$ ($v = 0, 1, 2, 3$) levels as well as their perturbors in $^{13}\text{C}^{17}\text{O}$.^a

Constant / Level	$A^1\Pi(v=0)$	$A^1\Pi(v=1)$	$A^1\Pi(v=2)$	$A^1\Pi(v=3)$
T_v	64757.7221 (10)	66188.7843 (34)	67587.3349 (44)	68953.3740 (19)
B_v	1.493 035 (18)	1.472 586 (16)	1.451 969 (14)	1.431 118 (11)
$q_v \times 10^5$	1.77 (59)	-1.63 (41)	-1.25 ^c	-1.35 (32)
$D_v \times 10^6$	6.523 (28)	6.399 (17)	6.480 (21)	6.558 (14)
$H_v \times 10^{11}$	-1.10 ^b	-1.10 ^b	-1.10 ^b	-1.10 ^b
Constant / Level	$d^3\Delta_i(v=4)$	$d^3\Delta_i(v=5)$	$d^3\Delta_i(v=7)$	$d^3\Delta_i(v=8)$
T_v	64967.897 (47)	66009.64 ^h	68028.906 (84)	69011.413 (11)
B_v	1.150 629 (78)	1.135 07 ^f	1.104 643 (87)	1.089 99 ^f
A_v	-15.947 (16)	-16.13 ^f	-16.42 ^f	-16.57 ^f
λ_v	0.772 ^d	0.865 ^d	1.052 ^d	1.145 ^d
$\gamma_v \times 10^3$	-7.76 ^e	-7.76 ^e	-7.76 ^e	-7.76 ^e
$D_v \times 10^6$	5.61 ^f	5.59 ^f	5.57 ^f	5.56 ^f
$H_v \times 10^{13}$	-6.45 ^g	-6.45 ^g	-6.45 ^g	-6.45 ^g
$A_{Dv} \times 10^5$	-4.72 ^e	-4.72 ^e	-4.72 ^e	-4.72 ^e
η	-21.598 (12)	-15.98 ⁱ	9.57 (23)	-0.560 (59)
Constant / Level	$e^3\Sigma^-(v=1)$	$e^3\Sigma^-(v=3)$	$e^3\Sigma^-(v=4)$	$e^3\Sigma^-(v=5)$
T_v	64781.49523 (86)	66842.0375 ^j	67843.866 (65)	68829.515 ^j
B_v	1.170 759 (17)	1.139 058 ^f	1.123 537 (65)	1.108 336 ^f
λ_v	0.5139 (35)	0.528 ^d	0.626 (32)	0.561 ^d
$D_v \times 10^6$	5.90 ^f	5.83 ^f	5.79 ^f	5.77 ^f
$H_v \times 10^{12}$	-1.61 ^g	-1.61 ^g	-1.61 ^g	-1.61 ^g
η	14.5737 (28)	4.748 ⁱ	13.101 (12)	8.322 ⁱ
Constant / Level	$a^{13}\Sigma^+(v=9)$	$a^{13}\Sigma^+(v=10)$	$a^{13}\Sigma^+(v=12)$	$a^{13}\Sigma^+(v=13)$
T_v	65182.01 ^j	66180.0621 (83)	68123.35 ^j	69068.9897 (45)
B_v	1.104 62 ^f	1.089 35 (47)	1.061027 ^f	1.046 459 (21)
λ_v	-1.115 ^d	-1.0955 (78)	-1.096 ^d	-1.111 8 (67)
$\gamma_v \times 10^3$	-5.98 ^e	-5.98 ^e	-5.98 ^e	-5.98 ^e
$D_v \times 10^6$	5.43 ^f	5.43 ^f	5.41 ^f	5.41 ^f
$H_v \times 10^{13}$	-3.22 ^g	-3.22 ^g	-3.22 ^g	-3.22 ^g
η	-2.60 ⁱ	-5.157 (15)	6.02 (22)	7.2554 (59)
Constant / Level	$I^1\Sigma^-(v=0)$	$I^1\Sigma^-(v=2)$	$I^1\Sigma^-(v=3)$	$I^1\Sigma^-(v=5)$
T_v	64565.40 ^k	66613.073 (18)	67608.1757 (98)	69541.194 (31)
B_v	1.173 59 ^f	1.141 89 ^f	1.126 04 (18)	1.093 99 ^f
$D_v \times 10^6$	5.93 ^g	5.97 ^g	5.98 ^g	6.02 ^g
$H_v \times 10^{12}$	2.42 ^g	2.42 ^g	2.42 ^g	2.42 ^g
$\xi \times 10^2$	-3.92 ⁱ	-7.492 (32)	-5.264 (61)	2.52 (16)
Constant / Level	$D^1\Delta(v=1)$	$D^1\Delta(v=4)$		
T_v	66447.8209 (50)	69355.14 (11)		
B_v	1.145 99 ^l	1.098 394 (97)		
$D_v \times 10^6$	6.08 ^l	6.03 ^l		
$H_v \times 10^{13}$	-2.42 ^g	-2.42 ^g		
$\xi \times 10^2$	-6.295 (38)	-2.21 (11)		

^a All values in cm^{-1} . T_v denotes ($v = 0, J = 0$) term value of the excited state in relation to the $X^1\Sigma^+$ ($v = 0, J = 0$) ground state; η_{uv} —spin-orbit interaction parameter; $\zeta_{v,v} - L$ —uncoupling interaction parameter. The values without given uncertainties were fixed during the final fit.

^b Isotopically recalculated on the basis of values given by Le Floch.[45].

^c Isotopically recalculated on the basis of the individual constants q of $A^1\Pi$ ($v = 0-4$) by Niu et al.[47,48].

^d Isotopically recalculated on the basis of the spin-spin C constants of Field [42] taking into account the equation $\lambda = -(\frac{3}{2})C$ (see Table 3.4 in Ref. [36]).

^e Calculated on the basis of data by Field [42] using the 29979,2458 $\text{MHz}/\text{cm}^{-1}$ conversion [89] and isotopically recalculated.

^f Isotopically recalculated on the basis of values given by Field [42].

^g Isotopically recalculated on the basis of values given by Le Floch [44].

^h Evaluated by the means of the isotopically recalculated vibrational equilibrium constants of $d^3\Delta_i$ by Field [42], and $X^1\Sigma^+$ by Le Floch [89]. T_e of $d^3\Delta_i$ was taken from Huber and Herzberg [76].

ⁱ Calculated on the basis of the isotopologue-independent, electronic perturbation parameters a and b given by Field [43,90] and Le Floch [44] as well as $\langle v_A | v_{d,e, \text{ or } a} \rangle$ and $\langle v_A | B(r) | v_{r, \text{ or } D} \rangle$ vibrational integrals in $^{13}\text{C}^{17}\text{O}$, using a method given by Hakalla et al. [55] (eqns. 1–8 ibidem).

^j Obtained by means of the isotopically recalculated vibrational equilibrium constants of $e^3\Sigma^-$ or $a^{13}\Sigma^+$ by Field [42] and $X^1\Sigma^+$ by Le Floch [89]. T_e of $e^3\Sigma^-$ or $a^{13}\Sigma^+$ was taken from Tilford et al. [75].

^k Calculated on the basis of isotopically recalculated vibrational equilibrium constants of $I^1\Sigma^-$ by Field [42] and $X^1\Sigma^+$ by Le Floch [89]. T_e of $I^1\Sigma^-$ was taken from Herzberg et al. [91].

^l Isotopically recalculated from Kittrell et al. [74].

Table 5Parameters for the $B^1\Sigma^+$ and $C^1\Sigma^+$ states of $^{13}\text{C}^{17}\text{O}$.^a

Level / Constant	$B^1\Sigma^+(v=0)$	$B^1\Sigma^+(v=1)$	$C^1\Sigma^+(v=0)$
T_v	86917.1236 (14) ^b	88928.0904 (18) ^b	91918.9416 (17) ^b
B_v	1.813 234 9 (93) 1.813 194 1 (58) ^c	1.789 919 (21) 1.790 227 (23) ^d	1.809 582 1 (82) 1.809 711 3 (93) ^e
$D_v \times 10^6$	5.706 (13) 5.562 0 (46) ^c	6.106 (46) 6.233 (47) ^d	5.350 5 (71) 5.532 (12) ^e

^a All values in cm^{-1} . Numbers in parentheses denote one standard deviation in units of the last significant digit.^b T_v denotes position of the level with respect to the $X^1\Sigma^+(v=0, J=0)$ ground state of $^{13}\text{C}^{17}\text{O}$. Note that the T_v uncertainties refer to the fitting with respect to a relative scale (fitting), while the absolute accuracy is not better than the 0.04 cm^{-1} .^c After Hakalla et al. [33].^d After Hakalla et al. [34].^e After Hakalla et al. [35].

Table 6

Spin–orbit and rotation–electronic perturbation parameters resulting from the deperturbation analysis of the $A^1\Pi$ ($v = 0, 1, 2, 3$) levels in $^{13}\text{C}^{17}\text{O}$.^a

Interaction	r -centroid ^b (Å)	$\langle v v' \rangle$ ^b (unitless)	$\eta_{v,v'}$ ^c (cm ⁻¹)	$\eta_{v,v'}/\langle v v' \rangle$ (cm ⁻¹)	a ^d (cm ⁻¹)	\bar{a} (cm ⁻¹)
$A^1\Pi$ ($v = 0$) $\sim d^3\Delta_i$ ($v = 4$)	1.2288	0.3728	-21.598 (12)	-57.929 (32)	94.60 (52)	94.60 (52) ^e
$A^1\Pi$ ($v = 2$) $\sim d^3\Delta_i$ ($v = 7$)	1.2311	-0.1666	9.57 (23)	-57.4 (14)	93.8 (22)	95.59 (27) ^f
$A^1\Pi$ ($v = 3$) $\sim d^3\Delta_i$ ($v = 8$)	1.1324	0.0107	-0.560 (59)	-52.2 (56)	85.3 (90)	96.029 (59) ^g 95.8 ^h 95.31 (22) ^h
$A^1\Pi$ ($v = 0$) $\sim e^3\Sigma^-$ ($v = 1$)	1.3463	-0.3382	14.5737 (28)	-43.0888 (82)	99.51 (19)	99.56 (19) ^e
$A^1\Pi$ ($v = 2$) $\sim e^3\Sigma^-$ ($v = 4$)	1.2888	-0.2998	13.101 (12)	-43.700 (41)	100.92 (95)	98.90 (33) ^f 99.045 (36) ^g 99.0 ^h 98.910 (58) ^h 98(4) ⁱ
$A^1\Pi$ ($v = 1$) $\sim a'^3\Sigma^+$ ($v = 10$)	1.1401	-0.1409	-5.157 (15)	36.59 (11)	84.50 (25)	84.02 (66) ^e
$A^1\Pi$ ($v = 2$) $\sim a'^3\Sigma^+$ ($v = 12$)	1.1323	0.1601	6.02 (22)	37.6 (14)	86.88 (32)	83.62 (12) ^f
$A^1\Pi$ ($v = 3$) $\sim a'^3\Sigma^+$ ($v = 13$)	1.1385	0.1995	7.2554 (59)	36.364 (30)	83.98 (68)	82.0 (25) ^g 83.4 ^h 88.6(3) ⁱ
		$\langle v \mathbf{B} v' \rangle$ ^b (cm ⁻¹)	$\xi_{v,v'} \times 10^{2c}$ (cm ⁻¹)	$\xi_{v,v'}/\langle v \mathbf{B} v' \rangle$ (unitless)	b ^d (unitless)	\bar{b} (unitless)
$A^1\Pi$ ($v = 1$) $\sim I^1\Sigma^-$ ($v = 2$)	1.3041	0.4526	-7.492 (32)	-0.1655 (71)	0.2341 (10)	0.2348 (95) ^e
$A^1\Pi$ ($v = 2$) $\sim I^1\Sigma^-$ ($v = 3$)	1.3178	0.3103	-5.264 (61)	-0.1696 (20)	0.2399 (28)	0.2274 (46) ^f
$A^1\Pi$ ($v = 3$) $\sim I^1\Sigma^-$ ($v = 5$)	1.3120	-0.1729	2.52 (16)	-0.1458 (93)	0.206 (13)	0.2269 (23) ^g 0.228 ^h 0.2273 (38) ^h 0.212(16) ⁱ
$A^1\Pi$ ($v = 1$) $\sim D^1\Delta$ ($v = 1$)	1.3290	-0.5531	-6.294 (38)	0.1138 (69)	0.1138 (69)	0.1139 (68) ^e
$A^1\Pi$ ($v = 3$) $\sim D^1\Delta$ ($v = 4$)	1.3141	-0.1699	-2.21 (11)	0.130 (62)	0.130 (62)	0.1103 (14) ^f 0.0969 (44) ^g 0.137(15) ⁱ

^a Numbers in parentheses denote one standard deviation in units of the last significant digit.

^b The r -centroids and vibrational integrals were calculated on the basis of RKR of the $^{13}\text{C}^{17}\text{O}$ A, d, e, a', I, and D states obtained from isotopically recalculated equilibrium constants of Field [42], Le Floch et al. [44], Field et al. [36,43], and Kittrell et al. [74]. In order to perform calculations, the program 'FRACONB' of Jung [83] and Jakubek [84] was used.

^c The spin–orbit and L -uncoupling interaction parameters $\eta_{v,v'}$ and $\xi_{v,v'}$ were taken from the global deperturbation fit of $^{13}\text{C}^{17}\text{O}$ conducted in this work.

^d The isotopologue independent spin–orbit and rotation–electronic perturbation parameters a and b were calculated in this work using a method highlighted in Ref. [55] (eqns. 1–8 ibidem).

^e The weighted average values of the electronic perturbation parameters.

^f After Hakalla et al. [55] on the basis of $^{12}\text{C}^{17}\text{O}$ data [92–94].

^g Weighted average of the values obtained by Le Floch et al. [44] (in their Table IV) on the basis of $^{12}\text{C}^{16}\text{O}$ data (uncertainties of a and b were taken from previously unpublished data).

^h On the basis of $^{12}\text{C}^{16}\text{O}$ data, after Field et al. [43].

ⁱ After Haridass et al. [52] on the basis of $^{12}\text{C}^{18}\text{O}$ data.

Table 7Term values (in cm^{-1}) of the $A^1\Pi$, $v = 0, 1, 2,$ and 3 rovibronic levels in $^{13}\text{C}^{17}\text{O}$.^{a, b}

J	$A^1\Pi(v=0)$		$A^1\Pi(v=1)$		$A^1\Pi(v=2)$		$A^1\Pi(v=3)$	
	e	f	e	f	e	f	e	f
1	64757.411	64757.393	66193.082	66193.131	67589.927	67589.898	68956.287	68956.234
2	64763.153	64763.288	66198.866	66199.075	67595.717	67595.750	68961.984	68962.005
3	64771.710	64772.028	66207.608	66207.882	67604.430	67604.445	68970.568	68970.574
4	64782.931	64783.652	66219.274	66219.536	67616.056	67616.044	68982.021	68982.000
5	64796.398	64798.088	66233.891	66234.122	67630.551	67630.567	68996.291	68996.317
6	64823.529	64815.244	66251.503	66251.644	67648.002	67647.936	69013.468	69013.435
7	64842.257	64834.947	66272.019	66272.133	67668.297	67668.189	69033.461	69033.485
8	64864.915	64856.834	66295.498	66295.599	67691.497	67691.673	69056.332	69056.328
9	64890.945	64897.883	66321.932	66321.987	67717.607	67717.636	69082.050	69082.024
10	64920.018	64925.755	66351.310	66351.362	67746.609	67746.655	69110.612	69110.591
11	64951.892	64957.263	66383.646	66383.672	67778.512	67778.534	69142.030	69141.990
12	64986.119	64992.198	66418.900	66418.933	67813.302	67813.331	69176.281	69176.199
13	65033.861	65030.362	66457.106	66457.129	67850.988	67851.031	69213.377	69213.163
14	65073.495	65071.657	66498.232	66498.255	67891.552	67891.603	69253.243	69252.346
15	65117.055	65116.008	66542.310	66542.305	67935.026	67935.046	69295.846	69297.501
16	65164.021	65163.412	66589.286	66589.299	67981.338	67981.400	69340.764	69342.528
17	65214.190	65213.769	66639.201	66639.205	68030.530	68030.580	69393.770	69390.797
18	65267.370	65267.082	66692.032	66692.019	68082.621	68082.655	69443.164	69441.973
19	65323.512	65323.311	66747.769	66747.811	68137.538	68137.632	69496.907	69495.802
20	65382.323	65382.137	66806.512	66806.455	68195.348	68195.436	69553.776	69555.662
21	65445.721	65445.665	66868.065	66868.079	68255.954	68256.132	69613.521	69613.850
22	65510.169	65510.047	66932.577	66932.541	68319.351	68319.629	69676.166	69676.250
23	65577.379	65577.308	66999.967	66999.940	68385.299	68385.973	69741.611	69741.658
24	65645.723	65645.68 ^c	67070.277	67070.310	68452.609	68455.154	69809.913	69809.935
25	65728.37 ^c	65728.23 ^c	67143.502	67143.493	68529.803	68527.005	69881.031	69881.018
26	65802.91 ^c	65802.87 ^c	67219.579	67219.56 ^c	68603.790	68601.259	69954.97 ^c	69954.953
27	65881.17 ^c	65881.13 ^c	67298.41 ^c	67298.36 ^c	68681.31 ^c	68676.937	70031.70 ^c	70031.796
28		65960.65 ^c	67380.99 ^c	67380.91 ^c	68761.74 ^c	68766.16 ^c	70111.25 ^c	70111.26 ^c
29			67465.43 ^c	67465.38 ^c	68844.99 ^c	68847.36 ^c	70193.60 ^c	70193.61 ^c
30			67553.05 ^c	67552.99 ^c	68930.73 ^c	68932.99 ^c	70278.74 ^c	70278.75 ^c
31			67643.55 ^c	67643.46 ^c	69016.95 ^c	69021.77 ^c	70366.66 ^c	70366.68 ^c
32			67736.93 ^c	67736.80 ^c	69115.24 ^c	69114.15 ^c	70457.39 ^c	70457.39 ^c
33			67833.21 ^c	67832.96 ^c	69207.89 ^c	69208.37 ^c	70550.86 ^c	70550.87 ^c
34				67931.91 ^c		69313.44 ^c	70647.07 ^c	70647.08 ^c
35				68032.84 ^c		69403.36 ^c	70746.23 ^c	70746.23 ^c
36							70847.99 ^c	70848.02 ^c

^a The rotational terms were calculated with regard to the $X^1\Sigma^+$ ($v = 0, J = 0$) ground state from the combined data sets of two experiments: the VIS-FTS and VUV-FTS study of the $^{13}\text{C}^{17}\text{O}$ $B \rightarrow A, C \rightarrow A$ bands (this work and Ref. [33–35]), and $B \leftarrow X, C \leftarrow X$ transitions (this work and Ref. [58]), respectively. The final term values were obtained using the weighted average method.

^b Note that the the absolute level energies are only good to 0.04 cm^{-1} , but the internal consistency (precision) of this list is much better—for some levels even 0.005 cm^{-1} .

^c Higher absolute term values obtained by means of the $A^1\Pi(v)$ relative terms (calculated from $B \rightarrow A$ ($0-v''$) [33,34] and $C \rightarrow A$ ($0-v''$) [35] bands by means of the Curl and Dane [85] and Watson [86] least-squares method) and fitted T_v values of the $A^1\Pi(v)$ levels (from Table 4). The final values are obtained using the weighted average method.

Table 8

Term values of the $d^3\Delta_i$, $e^3\Sigma^-$, $a^3\Sigma^+$, $1^1\Sigma^-$, and $D^1\Delta$ perturber states of the $A^1\Pi$ ($v = 0, 1, 2, 3$) levels in $^{13}\text{C}^{17}\text{O}$.^{a, b}

J	$e^3\Sigma^-$			$a^3\Sigma^+$			$1^1\Sigma^-$	
	F_{1e}	F_{2f}	F_{3e}	F_{1f}	F_{2e}	F_{3f}	F_{1f}	
	$v = 1$			$v = 10$				
1						66187.17		
2	64785.47	64791.64				66193.33		
3	64790.56	64798.82				66201.89		
4	64798.28	64808.50		66200.63	66212.77			
5	64809.01	64820.69						
6	64811.44	64835.47						
7	64827.49	64853.06					$v = 3$	
8		64873.73	64888.13				67671.33	
9		64880.57	64911.76				67689.06	
10		64905.91	64937.86					
11		64932.91	64966.65					
12		64961.80	64998.42	$v = 13$				
13		64992.76	65021.99	69232.80				
14			65058.98	69260.81				
15			65097.40	69287.69	69319.89			
16				69319.68	69353.99			
17					69384.96			
18					69424.53			
19						69508.77		
20						69549.66		
21								
22						69645.26		
	$v = 4$							
27		68694.59						
28		68748.51						
...							$v = 2$	
35							68043.82	
36							68122.54	
	$d^3\Delta$						$D^1\Delta$	
	F_{1e}	F_{1f}	F_{2e}	F_{2f}	F_{3e}	F_{3f}	F_{1e}	F_{1f}
	$v = 4$						$v = 1$	
20		65392.11						
21		65437.10						
22				65550.68				
23			65604.16					
24				65661.70				
25				65710.21			67190.23	
26				65771.64			67249.44	
27							67310.97	67310.96
28							67373.90	67373.85
29								
30							$v = 4$	
31		69041.66					70438.837	
32		69109.11					70508.34	
33	69179.76							
34							70653.75	
	$v = 8$							
10	69088.34							
...								
16					69360.09			

^a All values in cm^{-1} .

^b The rotational terms were calculated with regard to the $X^1\Sigma^+$ ($v = 0, J = 0$) ground state from the combined data sets of two experiments: the VIS-FTS and VUV-FTS study of the $^{13}\text{C}^{17}\text{O}$ $B \rightarrow A, C \rightarrow A$ bands (this work and Ref. [33–35]), and $B \leftarrow X, C \leftarrow X$ transitions (this work and Ref. [58]) with their extra-lines, respectively. The final values of the terms were merged using the weighted average method.

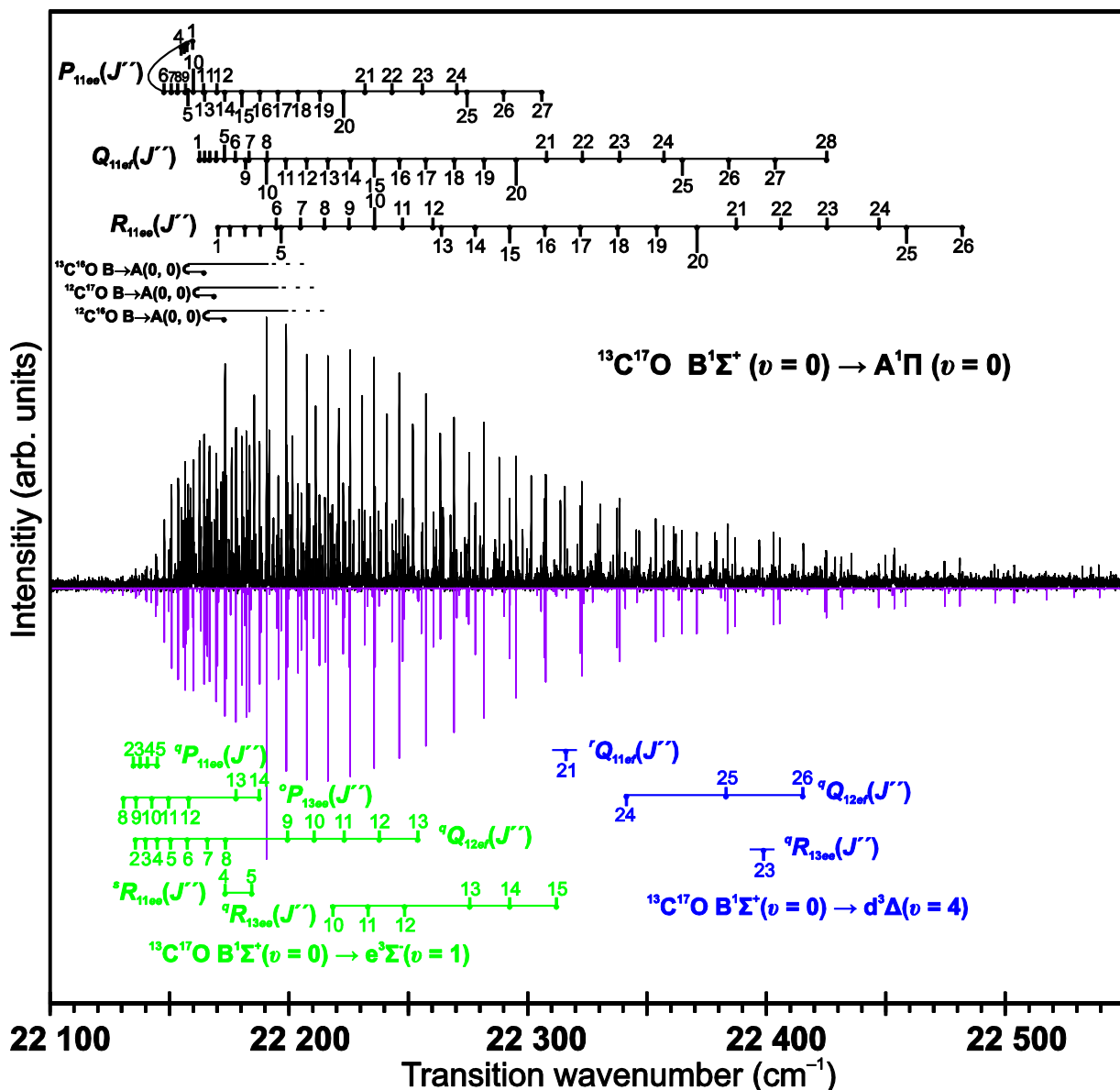


Fig. 1. High-resolution emission spectrum (upper trace) of the (0, 0) band of the Ångström ($B^1\Sigma^+ \rightarrow A^1\Pi$) system in the less-abundant $^{13}\text{C}^{17}\text{O}$ isotopologue recorded in the visible region with the Fourier-transform spectrometer at an instrumental resolution of 0.018 cm^{-1} . The extra-lines belong to the nominally forbidden $B^1\Sigma^+ \rightarrow e^3\Sigma^-(0, 1)$, and $B^1\Sigma^+ \rightarrow d^3\Delta_i(0, 4)$ intercombination systems, that are a result of the $A^1\Pi(v=0) \sim e^3\Sigma^-(v=1)$, and $A^1\Pi(v=0) \sim d^3\Delta_i(v=4)$ interactions, respectively. The lower trace presents a simulated spectrum [79] of the $^{13}\text{C}^{17}\text{O}$ B – A (0, 0) band. The ratio of the molecular gas composition used to obtain the spectrum was $^{13}\text{C}^{17}\text{O} : ^{13}\text{C}^{16}\text{O} : ^{12}\text{C}^{17}\text{O} : ^{12}\text{C}^{16}\text{O} = 1 : 0.76 : 0.34 : 0.25$. The estimated absolute calibration 1σ uncertainty was 0.003 cm^{-1} . The absolute accuracy of the line frequency measurements was about 0.003 cm^{-1} for the single, medium-strength lines and about 0.02 cm^{-1} for the weakest and/or most blended lines.

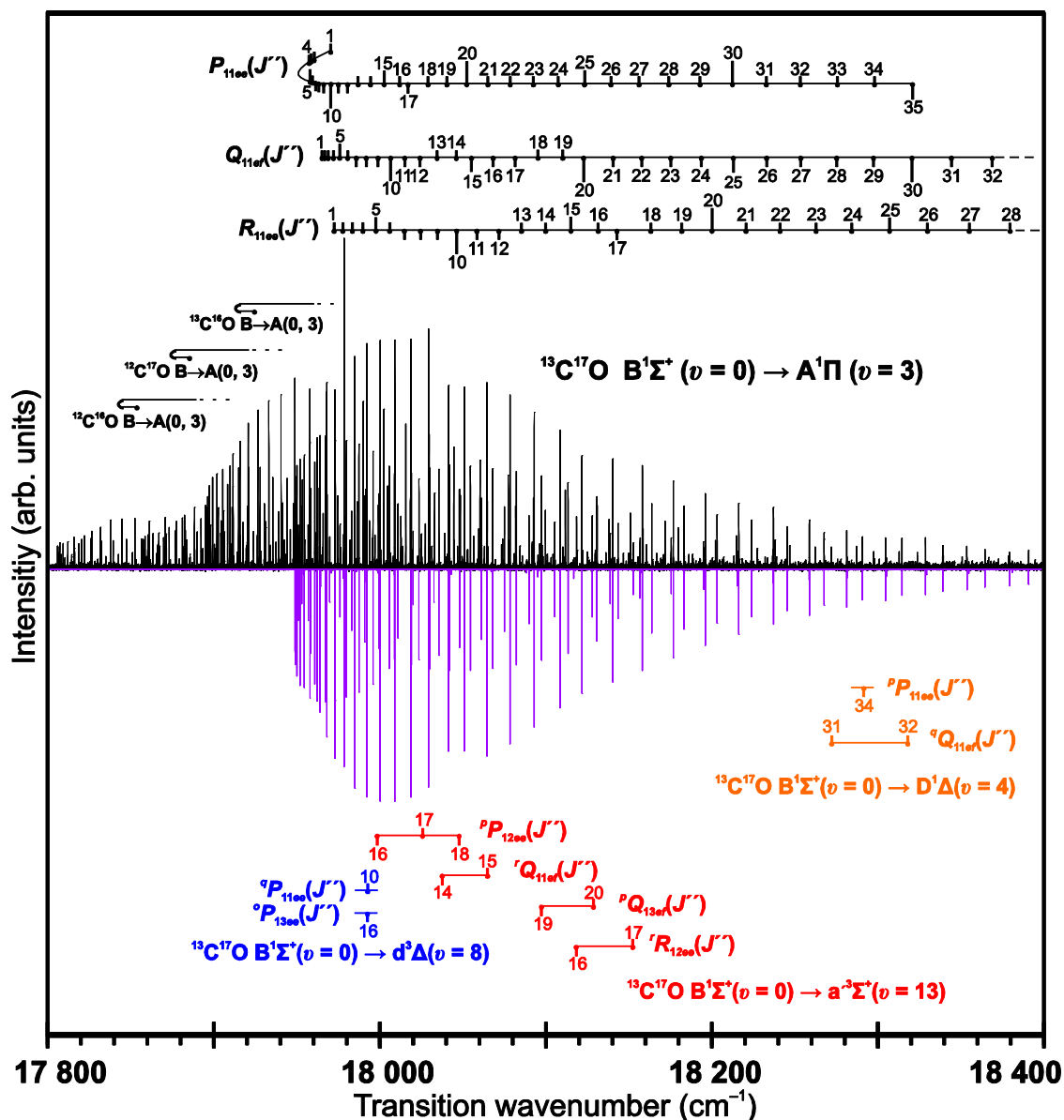


Fig. 2. High-resolution emission spectrum (upper trace) of the (0, 3) band of the Ångström ($B^1\Sigma^+ \rightarrow A^1\Pi$) system in the less-abundant $^{13}\text{C}^{17}\text{O}$ isotopologue recorded in the visible region with the Fourier-transform spectrometer at an instrumental resolution of 0.018 cm^{-1} . The extra-lines belong to the $B^1\Sigma^+ \rightarrow D^1\Delta$ (0, 4) system as well as nominally forbidden $B^1\Sigma^+ \rightarrow d^3\Delta_i$ (0, 8) and $B^1\Sigma^+ \rightarrow a^3\Sigma^+$ (0, 13) intercombination systems, that are a result of the $A^1\Pi$ ($v = 3$) $\sim D^1\Delta$ ($v = 4$) as well as $A^1\Pi$ ($v = 3$) $\sim d^3\Delta_i$ ($v = 8$), and $A^1\Pi$ ($v = 3$) $\sim a^3\Sigma^+$ ($v = 13$) interactions, respectively. Other details same as in Fig. 1.

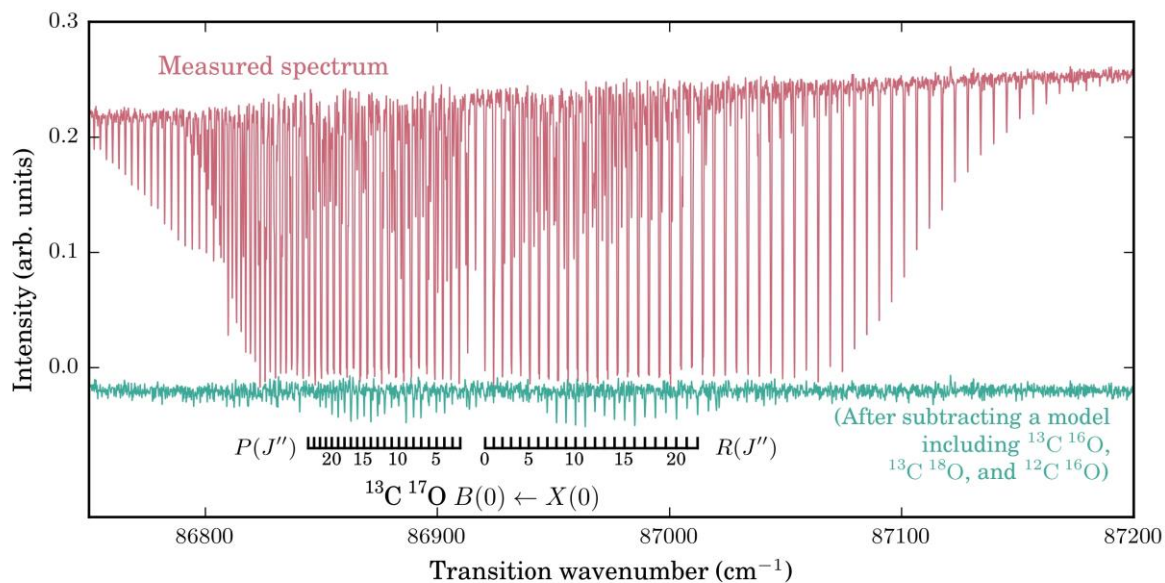


Fig. 3. Modelling of the $^{13}\text{C}^{17}\text{O}$ absorption spectrum (lower trace) of the ($\text{B}^1\Sigma^+ \rightarrow \text{X}^1\Sigma^+$) (0, 0) band of the Hopfield–Birge system superimposed on the high–resolution absorption spectrum (upper trace) present as impurity in a $^{13}\text{C}^{16}\text{O}$ gas sample recorded with the SOLEIL VUV–FTS at an instrumental resolution of 0.20 cm^{-1} . The ratio of the gases used in the experiment, sorted by concentrations, was $^{13}\text{C}^{16}\text{O} : ^{13}\text{C}^{18}\text{O} : ^{13}\text{C}^{17}\text{O} : ^{12}\text{C}^{16}\text{O} = 1 : 0.041 : 0.073 : 0.0045$. The estimated absolute calibration 1σ uncertainty was 0.01 cm^{-1} . The statistical uncertainty of fitted line centers was about 0.003 cm^{-1} for the unblended, medium–strength lines and about 0.09 cm^{-1} for the weakest and/or most blended lines.

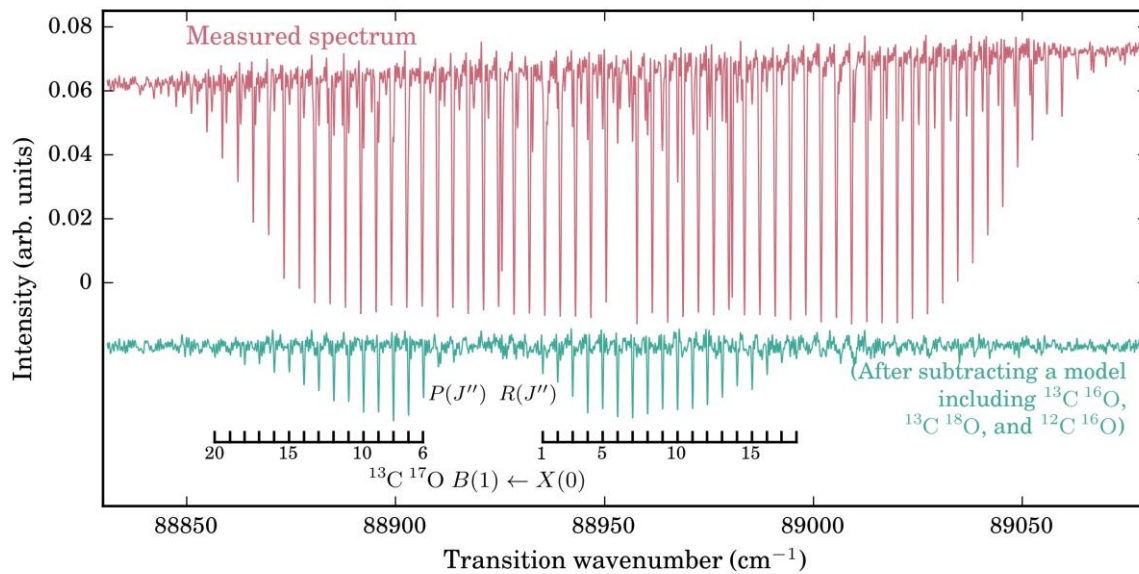


Fig. 4. Same caption as for Fig. 3 for the ($B^1\Sigma^+ \rightarrow X^1\Sigma^+$) (1, 0) band.

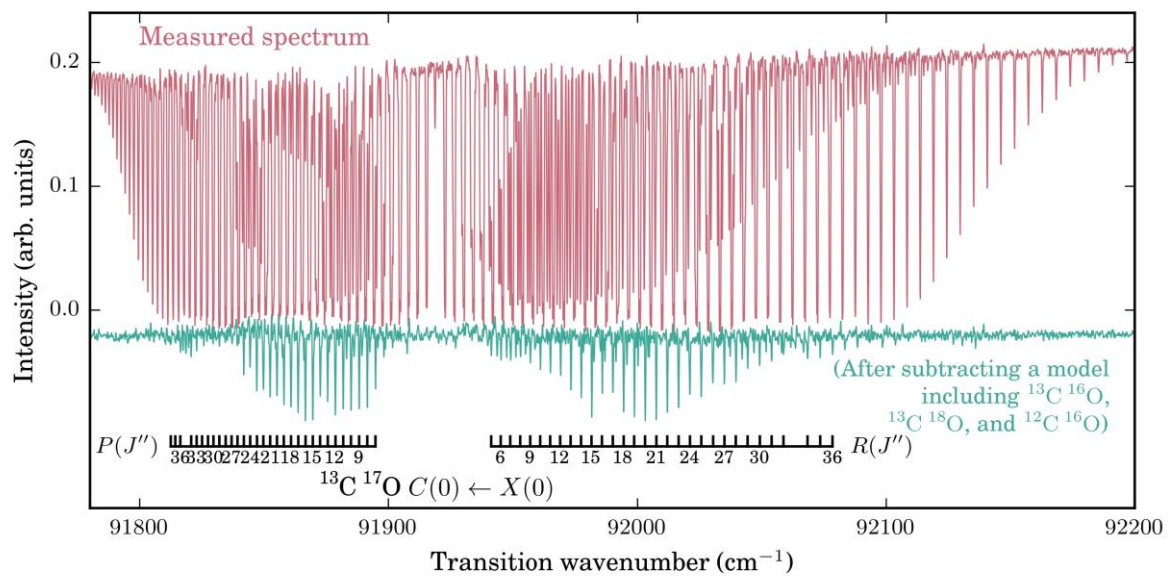


Fig. 5. Same caption as for Fig. 3 for the $(C^1\Sigma^+ \rightarrow X^1\Sigma^+)$ (0, 0) band.

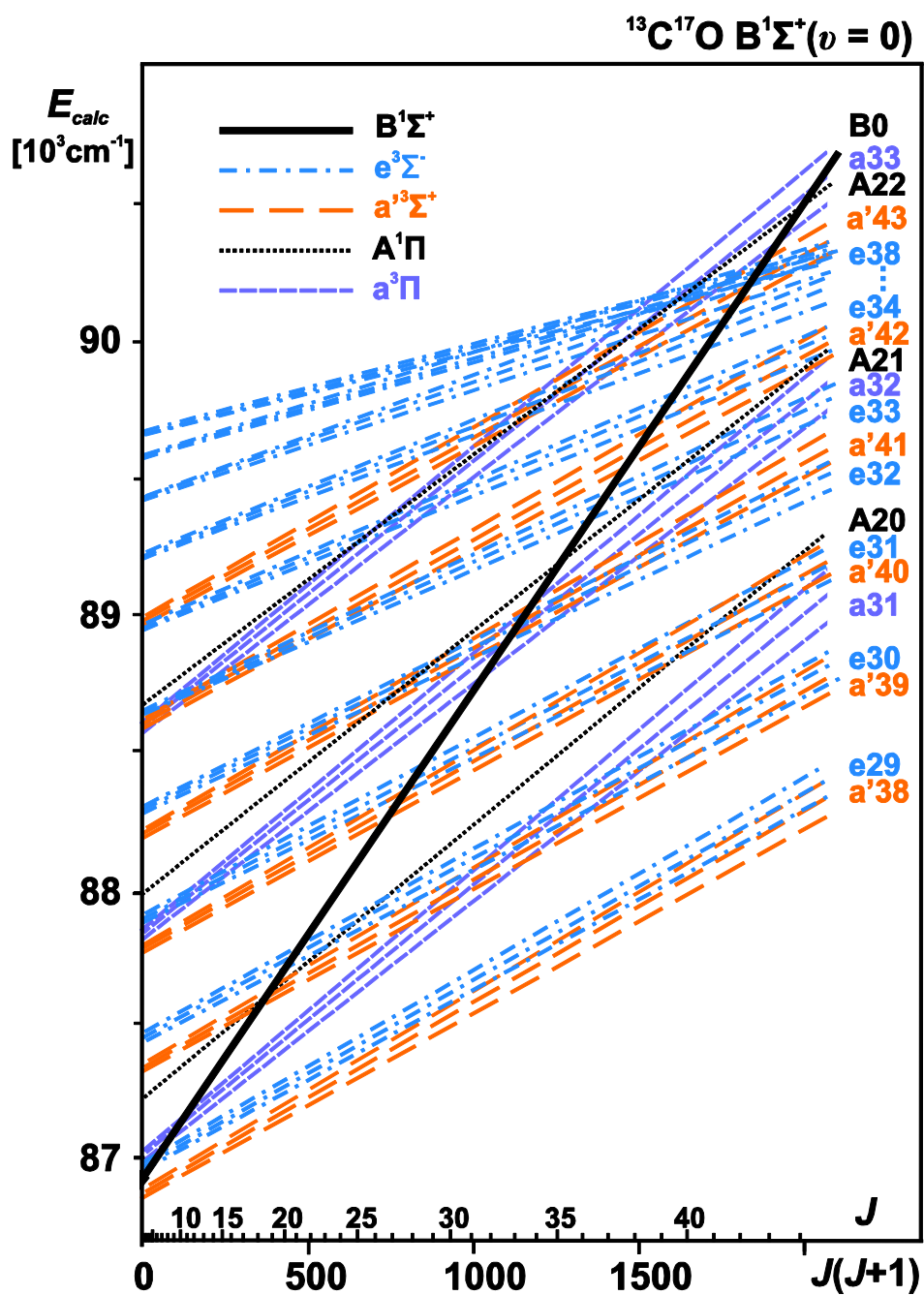


Fig. 6. Rovibronic term crossing diagram for the $^{13}\text{C}^{17}\text{O } B^1\Sigma^+(v=0)$ level. The curves are indicated by an electronic state label and vibrational quantum number to their right.

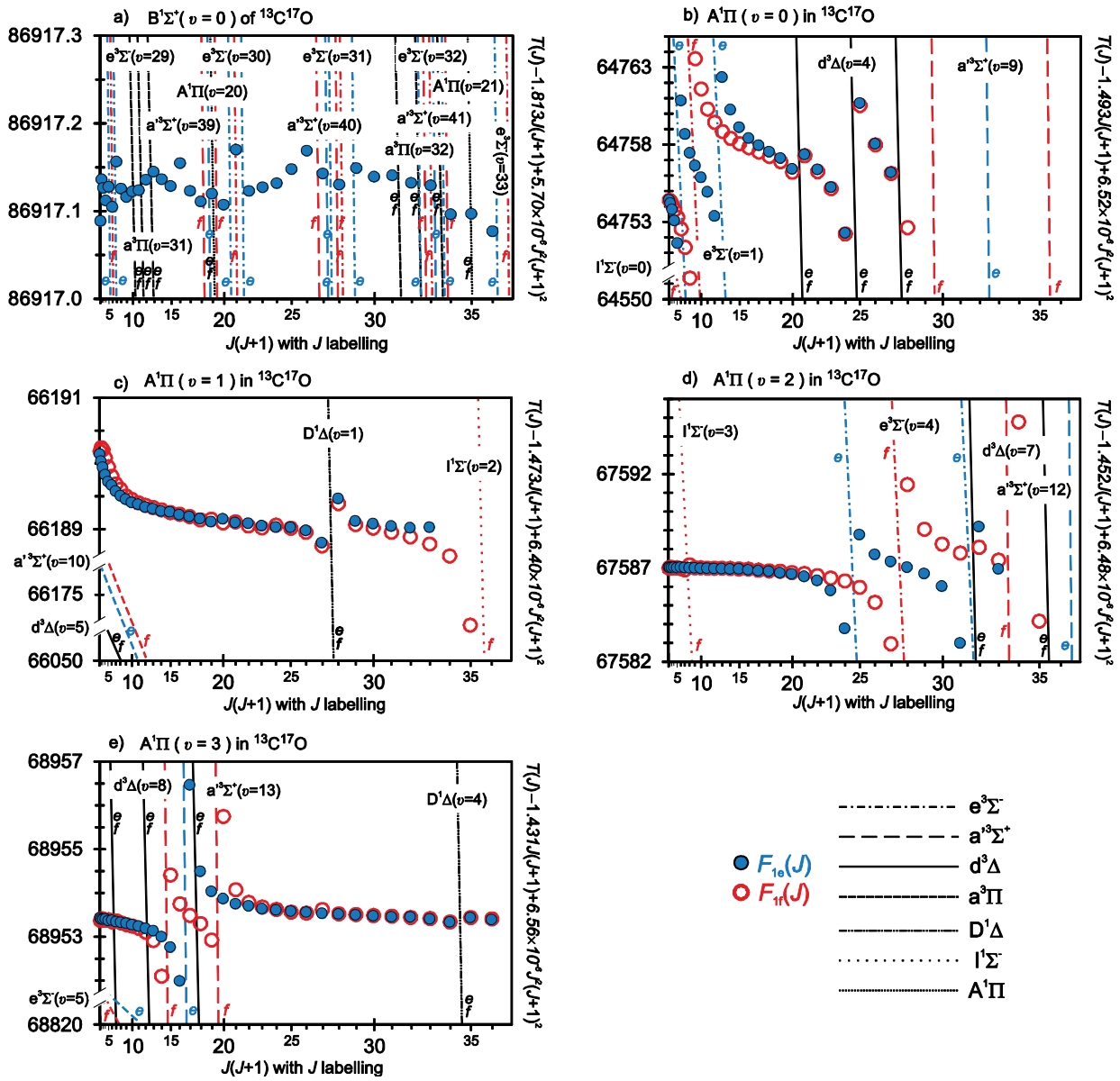


Fig. 7. The reduced term values (in cm^{-1}) plotted for the $^{13}\text{C}^{17}\text{O}$ $B^1\Sigma^+(v=0)$ and $A^1\Pi(v=0, 1, 2, 3)$ levels together with the hypothetical unperturbed crossing rovibronic levels of the possible perturbers. Filled (blue) and open (red) circles indicate e and f electronic symmetry of the B and A states, if applicable. In order to expand the vertical scale, $T(J) - BJ(J+1) + DJ^2(J+1)^2$ is subtracted from the observed term values (zero energy is placed at $v=0, J=0$ of the $X^1\Sigma^+$ ground state). Note that different reduced-energy scales are used for different vibrational levels.

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