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Galaxy And Mass Assembly (GAMA): the signatures of galaxy interactions as viewed from small scale galaxy clustering

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ABSTRACT

Statistical studies of galaxy-galaxy interactions often utilise net change in physical properties of progenitors as a function of the separation between their nuclei to trace both the strength and the observable timescale of their interaction. In this study, we use two-point auto, cross and mark correlation functions to investigate the extent to which small-scale clustering properties of star forming galaxies can be used to gain physical insight into galaxy-galaxy interactions between galaxies of similar optical brightness and stellar mass. The H α star formers, drawn from the highly spatially complete Galaxy And Mass Assembly (GAMA) survey, show an increase in clustering on small separations. Moreover, the clustering strength shows a strong dependence on optical brightness and stellar mass, where (1) the clustering amplitude of optically brighter galaxies at a given separation is larger than that of optically fainter systems, (2) the small scale clustering properties (e.g. the strength, the scale at which the signal relative to the fiducial power law plateaus) of star forming galaxies appear to differ as a function of increasing optical brightness of galaxies. According to cross and mark correlation analyses, the former result is largely driven by the increased dust content in optically bright star forming galaxies. The latter could be interpreted as evidence of a correlation between interaction-scale and optical brightness of galaxies, where physical evidence of interactions between optically bright star formers, likely hosted within relatively massive halos, persist over larger separations than those between optically faint star formers.

Key words: galaxies: interactions – galaxies: starburst – galaxies: star formation – galaxies: haloes – galaxies: statistics – galaxies: distances and redshifts

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1 INTRODUCTION

Historically, the field of galaxy interactions dates as far back as the 1940s, however, it was not until 1970s that the concept of tidal

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forces being the underlying drivers of morphological distortions in galaxies was fully accepted. It was the pioneering works by Toomre and Toomre (1972) on numerically generating "galactic bridges and tails" from galaxy interactions, and by Larson and Tinsley (1978) on broadband optical observations of discrepancies in "star formation rates in normal and peculiar galaxies" that essentially solidified this concept. Since then, the progress that followed revealed that interacting galaxies often show enhancements in H α emission (e.g. Keel *et al.* 1985; Kennicutt *et al.* 1987), infrared emission (e.g. Lonsdale *et al.* 1984; Soifer *et al.* 1984; Sanders *et al.* 1986; Solomon and Sage 1988), in radio continuum emission (e.g. Condon *et al.* 1982), and in molecular (CO) emission (e.g. Young *et al.* 1996) compared to isolated galaxies.

Over the past decade or so, numerous studies based on large sky survey datasets have provided ubiquitous evidence for, and signatures of tidal interactions. The enhancement of star formation is perhaps the most important and direct signature of a gravitational interaction (Kennicutt 1998; Wong et al. 2011), however, not all starbursts are interaction driven, and not all interactions trigger starbursts. Starbursts, by definition, are short-lived intense periods of concentrated star formation confined within the galaxy and are expected to be triggered only by the increase in molecular gas surface density in the inner regions over a short timescale. The tidal torques generated during the interactions of gas-rich galaxies are, therefore, one of the most efficient ways of funnelling gas to the centre of a galaxy (Smith et al. 2007; Di Matteo et al. 2007; Cox et al. 2008). In the absence of an interaction, however, bars of galaxies, which are prominent in spiral galaxies, can effectively facilitate both gas inflows and outflows (Regan and Teuben 2004; Owers et al. 2007; Ellison et al. 2011a; Martel et al. 2013), and trigger starbursts. Nuclear starbursts appear to be a common occurrence of interactions and mergers, however, there are cases where starbursts have been observed to occur, for example, in the overlapping regions between two galaxies (e.g. the Antennae galaxies; Snijders et al. 2007).

In the local Universe, most interacting galaxies have been observed to have higher than average central star formation (e.g. Lambas et al. 2003; Smith et al. 2007; Ellison et al. 2008; Xu et al. 2010; Scott and Kaviraj 2014; Robotham et al. 2014; Knapen and Cisternas 2015), though in a handful of cases, depending on the nature of the progenitors, moderate (e.g. Rogers et al. 2009; Darg et al. 2010; Knapen and Cisternas 2015) to no enhancements (e.g. Bergvall et al. 2003; Lambas et al. 2003) have also been reported. Likewise, interactions have been observed to impact circumnuclear gas-phase metallicities. In most cases, interactions appear to dilute nuclear gas-phase metallicities (e.g. Kewley et al. 2006b; Scudder et al. 2012; Ellison et al. 2013) and flatten metallicity gradients (e.g. Kewley et al. 2006a; Ellison et al. 2008). There are also cases where an enhancement in central gas-phase metallicities (e.g. Barrera-Ballesteros et al. 2015) has also been observed. The other observational signatures of galaxy-galaxy interactions include enhancements in optical colours, with enhancements in bluer colours (e.g. De Propris et al. 2005; Darg et al. 2010; Patton et al. 2011) observed to be tied to gas-rich and redder colours to gas-poor interactions (e.g. Rogers et al. 2009; Darg et al. 2010), increased Active Galactic Nuclei activities (AGNs, e.g.; Rogers et al. 2009; Ellison et al. 2011b; Kaviraj et al. 2015; Sabater et al. 2015) and substantially distorted galaxy morphologies (e.g. Casteels et al. 2013).

The strength and the duration of a physical change triggered in

an interaction can potentially shed light on to the nature of that interaction, progenitors and the roles of their galaxy- and halo-scale environments in driving and sustaining that change. In this regard, the projected separation between galaxies, R_p , can essentially be used as a clock for dating an interaction, measuring either the time elapsed since or time to the pericentric passage.

One of the more widely used approaches to understanding the effects of galaxy-galaxy interactions involves directly quantifying net enhancement or decrement of a physical property as a function of Rp. For example, the strongest enhancements in SFR have typically been observed over $< 30 h_{70}^{-1}$ kpc (e.g. Ellison *et al.* 2008; Li et al. 2008a; Wong et al. 2011; Scudder et al. 2012; Patton et al. 2013). The lower-level enhancements, on the other hand, have been observed to persist for relatively longer timescales. Ellison et al. (2008) report a net enhancement in SFRs and a decrement in metallicity of ~ 0.05 – 0.1 dex out to separations of ~ 30 – 40 h_{70}^{-1} kpc, and an enhancement in SFR out to wider separations for galaxy pairs of equal mass. Wong et al. (2011) report observations of SFR enhancements out to ~ $50 h_{70}^{-1}$ kpc based galaxy pair sample drawn from PRIMUS, Scudder et al. (2012) find that net changes in both SFR and metallicity persist out to at least ~ $80 h_{70}^{-1}$ kpc, Patton et al. (2013) find a clear enhancement in SFR out to \sim 150 kpc with no net enhancement beyond, while Patton et al. (2011) report enhancement in colours out to ~ 80 h_{70}^{-1} kpc, and Nikolic *et al.* (2004) report an enhancement in SFR out to \sim 300 kpc for their sample of actively star forming late-type galaxy pairs.

Even though the direct measure of a net change is advantageous as it can provide insight into dissipation rates and observable timescales of interaction-driven alterations (Lotz et al. 2011; Robotham et al. 2014), as highlighted above, the reported values of R_p out to which a given change persists often varies. The strength and the scale out to which a physical change is observable is expected to be influenced by orbital parameters and properties of progenitors (Nikolic et al. 2004; Owers et al. 2007; Ellison et al. 2010; Patton et al. 2011), as well as by the differences in dynamical timescales associated with short and long duration star formation events (Davies et al. 2015). Furthermore, galaxy-galaxy interactions do not always lead to observable changes. In particular, the subtle physical changes on R_ps at which progenitors are just starting to experience the effects of an interaction can be too weak to be observed. A further caveat is that this method fails to provide any physical insights into potential causes for the observed changes, i.e. whether the change is a result of the first pericentric passage, second or environment.

Another approach to studying the effects of galaxy-galaxy interactions involves two-point and higher order correlation statistics. The correlation statistics are often used in the interpretation of clustering properties of galaxies within one- and two-halo terms, and can to be utilised with or without incorporating physical information of galaxies. In this study, we aim to investigate whether large-scale environment plays any role in driving and sustaining interactiondriven changes in star forming galaxies with the aid of two-point correlation statistics.

In the local Universe, correlation functions have been ubiquitously used to study the clustering strength of galaxies with respect to galaxy properties like stellar mass, galaxy luminosities, and optical colours. Norberg *et al.* (2002) and Madgwick *et al.* (2003), for example, find clustering strength to be dependent strongly upon galaxy luminosity. Zehavi *et al.* (2005b); Li *et al.* (2006, 2009);

Zehavi *et al.* (2011); Ross *et al.* (2014); Favole *et al.* (2016) and Loh *et al.* (2010) report that galaxies with optically redder colours, which tend to be characterised with bulge dominated morphologies and higher surface brightnesses, correlate strongest with the strength of clustering than those residing in the green valley or in the blue cloud.

Even though much work has been done in this area, very few of those studies have focussed on investigating clustering of galaxies with respect to their star forming properties such as star formation rate (SFR), specific SFR (sSFR) and dust. The Sloan Digital Sky Survey (SDSS) based analysis of Li et al. (2008a) reports a strong dependence of the amplitude of the correlation function on specific star formation rate (sSFR) of galaxies on $R_p\,\,\lesssim\,\,100\,\,kpc.$ They find a dependence between clustering amplitude and sSFR, where the amplitude is observed to increase smoothly with increasing sSFR such that galaxies with high specific SFRs are clustered more strongly than those with low specific SFRs. The strongest enhancements in amplitude are found to be associated with the lowest mass galaxies and over the smallest R_p . They interpret this behaviour as being due to tidal interactions. Using GALEX imaging data of SDSS galaxies, Heinis et al. (2009) investigate the clustering dependence with respect to both (NUV - r) and sSFR. Over 0.01 $< R_p [h^{-1}Mpc] < 10$, they find a smooth transition in clustering strength from weak-to-strong as a function of the blueto-red change in (NUV - r) and the low-to-high change in sSFR. It must be noted, however, that on the smallest scales the clustering of the bluest (NUV - r) galaxies shows an enhancement.

Coil *et al.* (2016) use the PRIMUS and DEEP2 galaxy surveys spanning 0.2 < z < 1.2 to measure the stellar mass and sSFR dependence of the clustering of galaxies. They find that clustering dependence is as strong of a function of sSFR as of stellar mass, such that clustering smoothly increases with increasing stellar mass and decreasing sSFR, and find no significant dependence on stellar mass a fixed sSFR. This same trend is also found within the quiescent population. The DEEP2 survey based study of Mostek *et al.* (2013) too finds that within the star forming population the clustering amplitude increases as a function of increasing SFR and decreasing sSFR. Their analysis of small scale clustering of both star forming and quiescent populations, however, shows a clustering excess for high sSFR galaxies, which they attribute to galaxy-galaxy interactions.

The spatial and redshift completenesses of a galaxy survey largely determine the smallest R_p that can be reliably probed by two-point correlation statistics, thus the ability to trace galaxy-galaxy interactions reliably. The lack of sufficient overlap between pointings to ensure the full coverage of all sources can significantly impact the spatial completeness of a fibre-based spectroscopic survey. The resulting spatial incompleteness can considerably decrease the clustering signal on $R_p \leq 0.2$ [Mpc], especially for non-projected statistics (Yoon *et al.* 2008), and can have non-negligible effects even on larger scales (Zehavi *et al.* 2005b). Therefore many of the aforementioned studies are generally limited to probing clustering on $R_p \gtrsim 0.1$ [Mpc h⁻¹].

For this study, we draw a star forming sample of galaxies from the Galaxy And Mass Assembly (GAMA) survey (Driver *et al.* 2011; Liske *et al.* 2015), which has very high spatial and redshift completenesses (> 98.5%). The GAMA achieves this very high spatial completeness by surveying the same field over and over (~ 8 - 10 times) until all targets have been observed (Robotham *et al.* 2010, see the subsequent section for a discussion on the characteristics of the survey). Galaxy surveys like SDSS are limited both by the finite size of individual fibre heads as well as by the number of overlaps (~ 1.3 times). Therefore GAMA survey is ideal for a study, such as ours, that investigates the small-scale clustering properties of star forming galaxies as a function of the star forming properties.

This paper is structured as follows. In § 2, we describe the characteristics of the GAMA survey and the different GAMA catalogues that have been used in this study. This section also details the spectroscopic completeness of the GAMA survey, the selection of a reliable star forming galaxy sample from GAMA and the construction of galaxy samples for the clustering analyses. The different clustering techniques and definitions used in this analyses, as well as the modelling of the selection function associated with random galaxies, are described in § 3. Subsequently, in § 4, we present the trends of star forming galaxies with respect to different potential indicators of galaxy-galaxy interactions, and the correlation functions of star forming based on auto, cross and mark correlation statistics. Finally, in § 5 and 6, we discuss and compare the results of this study with the results reported in other published studies of star forming galaxies in the local Universe. This paper also includes four appendices, which are structured as follows. A discussion on sample selection and systematics is given in Appendix A. In Appendices B and C, we present a volume limited analysis involving auto and cross correlation functions, and further correlation results involving different galaxy samples introduced in §2. Finally, in Appendix D, we present the mark correlation analyses as we chose to show only the rank ordered mark correlation analysis in the main paper.

The assumed cosmological parameters are $H_{\circ} = 70 \text{ km s}^{-1} \text{Mpc}^{-1}$, $\Omega_{\text{M}} = 0.3$ and $\Omega_{\Lambda} = 0.7$. All magnitudes are presented in the AB system, and a Chabrier (2003) IMF is assumed throughout.

2 GALAXY AND MASS ASSEMBLY (GAMA) SURVEY

We utilise the GAMA (Driver *et al.* 2011; Liske *et al.* 2015) survey data for the analysis presented in this paper. In the subsequent sections, we briefly describe the characteristics of the GAMA survey and the workings of the GAMA spectroscopic pipeline.

2.1 GAMA survey characteristics

2.1.1 GAMA imaging

GAMA is a comprehensive multi-wavelength photometric and spectroscopic survey of the nearby Universe. GAMA brings together several independent imaging campaigns to provide a near-complete sampling of the UV to far-IR (0.15–500 μ m) wavelength range, through 21 broad-band filters; FUV, NUV (*GALEX*; Martin *et al.* 2005), *ugriz* (Sloan Digital Sky Survey data release 7, i.e. SDSS DR7; Fukugita *et al.* 1996; Gunn *et al.* 1998; Abazajian *et al.* 2009), Z, Y, J, H, K (VIsta Kilo-degree INfrared Galaxy survey, i.e. VIKING; Edge *et al.* 2013), W1, W2, W3, W4 (Wide-field Infrared Survey Explorer, i.e. *WISE*; Wright *et al.* 2010), 100 μ m, 160 μ m, 250 μ m, 350 μ m, and 500 μ m (*Herschel*-ATLAS; Eales *et al.* 2010). A complete analysis of the multi-wavelength successes of GAMA is presented in the end of survey report of Liske

et al. (2015) and in the panchromatic data release of Driver *et al.* (2015).

2.1.2 GAMA redshifts

GAMA's independent spectroscopic campaign was primarily conducted with the 2dF/AAOmega multi-object instrument (Sharp et al. 2006) on the 3.9m Anglo-Australian Telescope (AAT). Between 2008 and 2014, GAMA has surveyed a total sky area of $\sim 286 \ deg^2$ split into five independent regions; three equatorial (called GAMA-09hr or G09, G12, and G15) and two southern (G02 and G23) fields of $12 \times 5 \text{ deg}^2$ each. The GAMA equatorial targets are drawn primarily from SDSS DR7 (Abazajian et al. 2009). We refer the readers to the paper by Baldry et al. (2010) for detailed discussions on target selection strategies and input catalogues. The equatorial fields have been surveyed to an extinction corrected Petrosian r-band magnitude depth of 19.8. A key strength of GAMA is its high spatial completeness, both in terms of the overall completeness and completeness on small spatial scales. This is also advantageous for the present study aimed at investigating SFR enhancement due to galaxy interactions via small scale galaxy clustering. The tiling and observing strategies of the survey are discussed in detail in Robotham et al. (2010) and Driver et al. (2011). At the conclusion of the spectroscopic survey, GAMA has achieved a high redshift completeness of 98.5% for the equatorial regions, and we discuss in detail the spectroscopic completeness of the survey in § 2.3.

2.1.3 GAMA spectroscopic pipeline

A detailed summary of the GAMA redshift assignment, reassignment, and quality control procedure is given in Liske *et al.* (2015), according to which galaxy redshifts with normalised redshift qualities (NQ) \geq 3 are secure redshifts. GAMA does not reobserve galaxies with high-quality spectra originating from other surveys, such that the GAMA spectroscopic catalogues comprise spectra from a number of other sources, e.g. SDSS, 2-degree Field Galaxy Redshift Survey (2dFGRS; Colless *et al.* 2001), Millennium Galaxy Catalogue (MGC; Driver *et al.* 2007), see § 2.3 for a discussion on the contribution of non-GAMA spectral measures to our analysis. Finally, given the exceptionally high redshift completeness of the GAMA equatorial fields, we restrict our analysis to the equatorial data.

The GAMA spectroscopic analysis procedure, including data reduction, flux calibration, and spectral line measurements, is presented in Hopkins *et al.* (2013). The GAMA emission line catalogue (SpecLineSFR) provides line fluxes and equivalent width measurements for all strong emission line measurements. A more detailed description of the spectral line measurement procedure and SpecLineSFR catalogue, in general, can be found in Gordon *et al.* (2017). Additionally, the strength of the λ 4000-Å break (D₄₀₀₀) is measured over the D₄₀₀₀ bandpasses (i.e. 3850 – 3950Å and 4000 – 4100Å) defined in Balogh *et al.* (1999) following the method of Cardiel *et al.* (1998). SpecLineSFR also provides a continuum (6383 – 6538Å) signal-to-noise per pixel measurement, which is representative of the red-end of the spectrum.

2.2 Galaxy properties

The two main intrinsic galaxy properties used in this investigation are $H\alpha$ SFRs and galaxy stellar masses. Below, we briefly overview the derivation of these properties and discuss their uncertainties.

2.2.1 Ha Star Formation Rates

The GAMA intrinsic H α SFRs are derived following the prescription of Hopkins *et al.* (2003), using the Balmer emission line fluxes provided in SpecLineSFR. The spectroscopic redshifts used in the calculation are corrected for the effects of local and large-scale flows using the parametric multi-attractor model of Tonry *et al.* (2000) as described in Baldry *et al.* (2012), and the application of stellar absorption, dust obscuration and fibre aperture corrections to SFRs is described in detail in Gunawardhana *et al.* (2013).

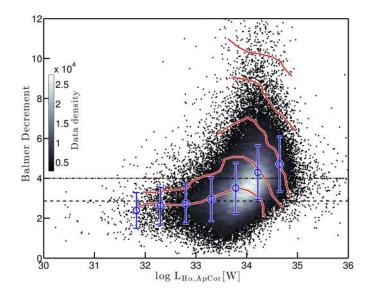


Figure 1. The distribution of Balmer decrement in aperture corrected H α luminosity (L_{H α , ApCor}, i.e. H α luminosity before correcting for dust obscuration) illustrating the luminosity dependence of dust obscuration. The grey colour scale shows the data density distribution of all star forming galaxies. The black dashed and dot-dashed lines indicate the theoretical Case B recombination ratio of 2.86, and the Balmer decrement corresponding to the assumption of one magnitude extinction at the wavelength of H α . The blue points denote the mean variation and one-sigma error in dust obscuration as a function of L_{H α , ApCor}. The constant log sSFR contours, shown in red, are defined in steps of 0.3 dex, where log sSFR increases from $-10.2[yr^{-1}]$ at low Balmer decrements to $-9[yr^{-1}]$ at high Balmer decrements.

The luminosity (or SFR) dependent dust obscuration, reflecting that massive star forming galaxies also contain large amounts of dust relative to their low-SFR counterparts, is observationally well established in the local Universe (e.g. Hopkins *et al.* 2003; Brinchmann *et al.* 2004; Garn and Best 2010; Ly *et al.* 2012; Zahid *et al.* 2013; Jimmy *et al.* 2016). The mean variation in Balmer decrement with aperture corrected H α luminosity for our sample is shown as blue points in Figure 1, with red contours indicating the dependence of Balmer decrement on specific SFR. The dot-dashed line denotes the Balmer decrement approximately corresponding to the assumption of an extinction of one magnitude at the wavelength of H α for all galaxy luminosities (Kennicutt 1992). In this study, for

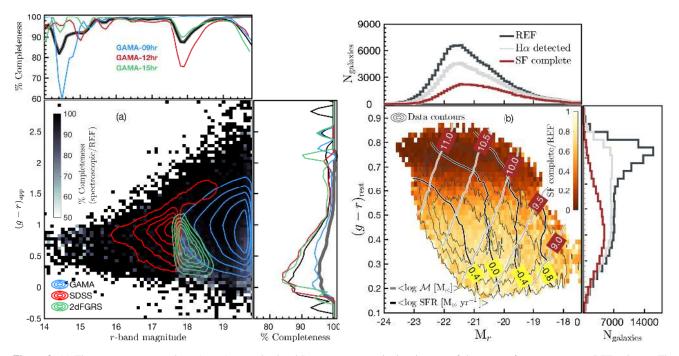


Figure 2. (a) The apparent g - r colour, $(g - r)_{app}$, and *r*-band Petrosian magnitude distributions of the ratios of spectroscopic-to-REF galaxies. The colour code corresponds to the percentage completeness with lighter colours indicating the deviation of the ratios from unity. The coloured contours show the approximate distribution of galaxies in our sample originating from GAMA, SDSS, 2dFGRS surveys. The top and side panels show completeness as a function of *r*-band Petrosian magnitude and $(g - r)_{app}$, respectively, with black and thick grey lines showing the overall completeness across the three equatorial fields with (black) and without (grey) a spectral signal-to-noise cut, and the coloured lines showing the completenesses for individual GAMA fields. (b) The $(g - r)_{rest}$ and M_r distribution of the ratio of SF complete-to-REF galaxies. The closed contours from inwards-to-outwards enclose ~ 25, 50, 75 and 90% of the SF complete data. Also shown are the constant mean log stellar mass ($< \log M/M_{\odot} >$) and mean log SFR ($< \log SFR [M_{\odot} yr^{-1}] >$) contours corresponding to SF complete galaxies. The top and side panels show the univariate M_r and $(g - r)_{rest}$ distributions of REF (black) and SF complete (brown) galaxies, as well as the distribution all SF galaxies with reliably measured H α emission line fluxes (grey).

galaxies without reliable H β flux measurements, we approximate a Balmer decrement based on the relation shown in blue in Figure 1.

2.2.2 Stellar masses

The GAMA stellar masses and absolute magnitudes¹ provided in STELLARMASSESV16 (Taylor *et al.* 2011; Kelvin *et al.* 2012) catalogue are used for this study. A Bayesian approach is used in the derivation of the stellar masses, and are based on *u*, *g*, *r*, *i*, *z* spectral distributions and Bruzual and Charlot (2003) population synthesis models. Furthermore, the derivation assumes a Chabrier (2003) stellar IMF and Calzetti *et al.* (2000) dust law. The stellar mass uncertainties, modulo any uncertainties associated with stellar population synthesis models, are determined to be ~ 0.1 dex. A detailed discussion on the estimation of GAMA stellar masses and the associated uncertainties can be found in Taylor *et al.* (2011).

2.3 Sample selection and spectroscopic completeness

We select a *reference* sample of galaxies, henceforth REF, consisting only of equatorial objects that satisfy both the GAMA main survey selection criteria (Baldry *et al.* 2010), and have spectroscopic redshifts, z_{spec} , in the range $0.002 \leq z_{\text{spec}} < 0.35$, representing the *z* window over which the H α spectral feature is observable in the GAMA spectra (Driver *et al.* 2011). The REF sample consists of 157 079 objects in total.

Out of the REF galaxies, those observed either as a part of GAMA and/or SDSS spectroscopic surveys with spectral signal-to-noise >3 form the *spectroscopic* sample. Objects with other survey spectra (e.g. 2dFGRS, MGC) are excluded as they lack the necessary information needed to reliably flux calibrate their spectra, and the objects with duplicate spectra² are removed on the basis of their spectral signal-to-noise, leaving 148 834 galaxies in the spectroscopic sample.

We assess the spectroscopic completeness of the survey by comparing the bivariate colour-magnitude distributions of REF and spectroscopic samples. Figure 2(a) shows the colour-magnitude distribution of the ratio of spectroscopic-to-REF galaxies in a given *r*-band magnitude and apparent g - r colour, hereafter $(g - r)_{app}$, cell, and the top and right-side panels show the completeness as a function of the *r*-band magnitude and $(g - r)_{app}$. The exclusion of 2dFGRS spectra, in particular, leads to an overall incompleteness of ~ 20% across the three equatorial regions over the magnitude range probed by the 2dFGRS (green contours in Figure 2(a) highlight the colour and magnitude range corresponding

¹ The rest-frame colours used in this analysis are based on these absolute magnitudes.

 $^{^2}$ In cases where an object has an independent GAMA and a SDSS spectrum, the SDSS spectrum is generally found to have the highest spectral signal-to-noise, and is selected to be part of the sample.

to the 2dFGRS galaxy distribution). The incompleteness present in each field, however, varies considerably, with G12 being the most incomplete (i.e. relatively a larger number of 2dFGRS galaxies reside in this region) and G09 being the most complete (i.e. no 2dFGRS galaxies reside in this region) as shown in the top panel of Figure 2(a). Additionally, recall that GAMA spectral signal-tonoise measures are representative of the red end of the spectrum, therefore, the application of a signal-to-noise cut results in the incompleteness evident at fainter magnitudes and bluer colours in the same figure. The implication being that the spectroscopic sample is biased against optically faint bluer galaxies (the thin and thick black lines shown in the side panels of Figure 2(a) clearly demonstrate this bias). Note that the variations in completeness seen at optically redder colours is largely driven by small number statistics. See § A2 for discussion on the impact of spectroscopic incompleteness on the results and conclusions of this study.

Out of the galaxies with detected H α emission in the spectroscopic sample, those dominated by active galactic nuclei (AGN) emission are removed using the standard optical emission line ([N II] λ 6584/H α and [O III] λ 5007/H β) diagnostics (BPT; Baldwin et al. 1981) and the Kauffmann et al. (2003a) pure star forming (SF) and AGN discrimination prescription. If all four emission lines needed for a BPT diagnostic are not detected for a given galaxy, then the two line diagnostics based on the Kauffmann et al. (2003a) method (e.g. log [N II] λ 6584/H α > 0.2 and log [O III] λ 5007/H β > 1.0) are used for the classification. The galaxies that were unable to be classified this way are retained in our sample as a galaxy with measured H α flux but without an [N II] λ 6584 or $[O_{III}] \lambda 5007$ measurement are more likely to be SF galaxies than AGNs (Cid Fernandes et al. 2011). Overall, ~16% of objects are classified either as an AGN or as an AGN-SF composite and are removed from the sample, and the ${\sim}28\%$ unable to be classified are retained in the sample.

As a consequence of the bivariate magnitude and H α flux selection that is applied to our sample, our sample is biased against optically faint SF galaxies. This is a bias that not only affects any star forming galaxy sample drawn a broadband magnitude survey, but it becomes progressively more significant with increasing *z* (Gunawardhana *et al.* 2015). Therefore to select an approximately complete SF galaxy sample, henceforth *SF complete*, we impose an additional flux cut of 1×10^{-18} Wm⁻², which roughly corresponds to the turnover in the observed H α flux distribution of GAMA H α detected galaxies (Gunawardhana *et al.* 2013).

A comparison between the SF complete sample and REF galaxies in rest-frame g - r colour, hereafter $(g - r)_{rest}$, and M_r space is shown in Figure 2(b). The closed contours denote the fraction of the data enclosed, while the open black and grey contours denote constant $\langle \log SFR [M_{\odot} yr^{-1}] \rangle$ and $\langle \log \mathcal{M}/M_{\odot} \rangle$ lines, respectively. Even though the SF complete galaxies are dominated by optically bluer systems, a significant fraction of galaxies with optically redder colours have reliably measured H α SFRs, indicating on-going star formation, albeit at lower rates. Also shown are the univariate M_r and $(g - r)_{rest}$ distributions of REF (black), SF complete (brown), and of galaxies with reliable H α emission detections that are classified as SF following the removal of AGNs (grey) to illustrate how the H α flux cut of 1×10^{-18} Wm⁻² act to largely exclude optically redder systems from our sample.

2.4 REF and SF complete samples for clustering analysis

In order to investigate the clustering properties of star forming galaxies with respect to optical luminosity and stellar mass (§ 4.2 to § 4.4), we use REF and SF complete samples to further define three disjoint luminosity selected, three disjoint stellar mass selected, and several volume limited samples, for which all selection effects are carefully modelled.

The three disjoint luminosity selected samples, called M_f , M_* and M_b , together cover the $-23.5 \leq M_r < -19.5$ range, and the three disjoint stellar mass selected samples, called $\mathcal{M}_{\mathcal{L}}$, $\mathcal{M}_{\mathcal{I}}$ and $\mathcal{M}_{\mathcal{H}}$, together span the 9.5 $\leq \log \mathcal{M}/M_{\odot} < 11$. See Tables 1 and 2 for individual magnitude and stellar mass coverages of each luminosity and stellar mass selected sample, as well as for a description of their key characteristics. We also define two redshift samples for each M_b , M_* , and M_f , and for each $\mathcal{M}_{\mathcal{H}}$, $\mathcal{M}_{\mathcal{I}}$ and $\mathcal{M}_{\mathcal{L}}$, where one set covers the full redshift range of the SF complete galaxies, and the second spans only the 0.001 $\leq z \leq 0.24$ range.

Out of the two redshift samples mentioned above, the former (i.e. the samples covering the full redshift range) is used for the auto correlation analysis, and the latter for the cross and mark correlation analyses (§ 4.3 and 4.4). The main reason for restricting the redshift coverage of galaxy samples in the latter case is to overcome the effects of the equivalent width bias³ (Liang et al. 2004; Groves et al. 2012, see also Appendix A). In this study, we find that the cross correlation functions, hereafter CCFs, of low sSFR galaxies spanning $0.24 \leq z < 0.34$ in redshift computed using two different clustering estimators, the Landy and Szalay (1993) and Hamilton (1993) estimators, differ systematically from each other, suggesting a failure in the modelling of the selection function of low sSFR galaxies over the $0.24 \leq z < 0.34$ range. The respective results for the low sSFR galaxies in the 0.01 $\leq z \leq$ 0.24 range, on the other hand, are consistent with each other. Therefore we limit the redshift range of all galaxy samples used for the cross and mark correlation analyses to $0.01 \leq z \leq 0.24$.

The log sSFR and $(g-r)_{rest}$ distributions of the three disjoint luminosity selected samples are shown in Figures 3 and 4. In Figure 3, with increasing optical luminosity, the peak of the distribution of log sSFRs moves progressively towards lower sSFRs. The notably broader peak of the M_b distribution arises as a result of the bimodality present in the bivariate SFR (or sSFR) and M distribution (see, for example, Figure 10). Similarly, the $(g-r)_{rest}$ distributions show a progressive shift towards redder colours with increasing optical luminosity. From each disjoint luminosity (stellar mass) selected sample, we select the 30% highest and lowest sSFR (SFR), $(g-r)_{rest}$, Balmer decrement and D₄₀₀₀ (i.e. the strength of the 4000 Å break, Kauffmann *et al.* 2003b) galaxies to be used in the cross correlation analysis (§ 4.3). The red and blue arrows in Figures 3 and 4 show these 30% selections.

As none of the samples defined so far are truly volume limited, we define a series of volume limited luminosity and stellar mass samples, which are described in Table B1. The volume limited SF

³ Emission line samples drawn from a broadband survey, like GAMA, can be biased against low SFR and weak-line systems. This can become significant with increasing redshift and apparent magnitude, and the differences in clustering results obtained from different clustering estimators can be used to quantify the significance of such biases.

Table 1. The key characteristics of the three disjoint luminosity selected sub-samples ($M_b: -23.5 \le M_r < -21.5$; $M_*: -21.5 \le M_r < -20.5$; $M_f: -20.5 \le M_r < -19.5$) drawn from the SF complete and REF samples are given. For each sample, we provide the size of the sample, the average redshift and central ~ 50% redshift range, median log sSFR [yr⁻¹], (g - r)_{rest} and log \mathcal{M} [M_{\odot}] along with their central ~ 50% ranges. We define two redshift samples for each M_b , M_* and M_f , where one sample covers the full redshift range over which the H α feature is visible in GAMA spectra (i.e. 0.001 < z < 0.34), and the second covers a narrower 0.001 < $z \le 0.24$ range (see § 4.3). Using both the r-band magnitude selection of the GAMA survey and the H α flux selection of our sample, we estimate a completeness for each disjoint luminosity selected sub-sample, which is shown within brackets under N_{galaxies}.

subset	Ngalaxies	$\langle z \rangle$	z	log sSFR	log sSFR	$\langle (g-r)_{\rm rest} \rangle$	$(g-r)_{\rm rest}$	$\langle \log \mathcal{M} \rangle$	$\log \mathcal{M}$
	-		σ=25%,75%	[yr ⁻¹]	σ=25%,75%		σ=25%,75%	$[M_{\odot}]$	σ=25%,75%
SF complete									
M_b	8 100 (53%) ^a	0.24	(0.19, 0.29)	-10.28	(-10.70, -9.87)	0.55	(0.47, 0.63)	10.9	(10.8, 11.1)
	3 749 (68%)	0.17	(0.13, 0.21)	-10.67	(-11.08, -10.13)	0.60	(0.51, 0.69)	10.8	(10.68, 10.99)
M_*	20976 (12%)	0.21	(0.18, 0.27)	-9.90	(-10.20, -9.61)	0.48	(0.39, 0.56)	10.46	(10.31, 10.65)
	12308 (62%)	0.17	(0.13, 0.21)	-10.11	(-10.52, -9.79)	0.50	(0.41, 0.59)	10.32	(10.15, 10.50)
M_f	14000 (< 1%)	0.14	(0.11, 0.18)	-9.84	(-10.14, -9.54)	0.42	(0.32, 0.51)	9.98	(9.81, 10.16)
5	13650 (< 1%)	0.14	(0.11, 0.18)	-9.94	(-10.24, -9.64)	0.42	(0.33, 0.51)	9.83	(9.66, 10.02)
					REF				
M_b	33 406	0.25	(0.20, 0.30)	_	-	0.67	(0.59, 0.75)	10.95	(10.83, 11.09)
M_*	64618	0.22	(0.18, 0.27)	_	-	0.59	(0.48, 0.72)	10.50	(10.34, 10.69)
M_f	34 868	0.15	(0.13, 0.19)	-	-	0.51	(0.37, 0.67)	9.98	(9.76, 10.20)

^a The sample completeness

Table 2. The key characteristics of the three disjoint stellar mass selected sub-samples ($\mathcal{M}_{\mathcal{H}}$: 10.5 $\leq \log \mathcal{M}/M_{\odot} \leq 11.0$; $\mathcal{M}_{\mathcal{I}}$: 10.0 $\leq \log \mathcal{M}/M_{\odot} \leq 10.5$; $\mathcal{M}_{\mathcal{L}}$: 9.5 $\leq \log \mathcal{M}/M_{\odot} \leq 10.0$) drawn from the SF complete and REF samples are given. For each sample, we provide the size of the sample, the average redshift and central ~ 50% range, median log sSFR, (g - r)_{rest} and M_r along with their central ~ 50% ranges. As described in the caption of Table 1 above, we define two redshift samples for each $\mathcal{M}_{\mathcal{H}}$, $\mathcal{M}_{\mathcal{I}}$ and $\mathcal{M}_{\mathcal{L}}$. The completeness of each sample due to the dual *r*-band magnitude and H α flux is indicated within brackets in the second column (after N_{galaxies}), which is approximately the fraction of galaxies seen over the full volume. This value does not take into account the maximum volume out to which a galaxy of a given stellar mass would be detected.

subset	Ngalaxies	$\langle z \rangle$	z	log sSFR	log sSFR	$\langle (g-r)_{\rm rest} \rangle$	$(g-r)_{\rm rest}$	$\langle \mathrm{M}_r \rangle$	M_r
			σ=25%,75%	[yr ⁻¹]	σ=25%,75%		σ=25%,75%		σ=25%,75%
SF complete									
$\mathcal{M}_{\mathcal{H}}$	11 600 (36%)	0.23	(0.18, 0.30)	-10.35	(-10.72, -9.98)	0.57	(0.50, 0.64)	-21.53	(-21.78, -21.29)
	5 597 (61%)	0.16	(0.12, 0.21)	-10.57	(-10.99, -10.17)	0.61	(0.54, 0.68)	-21.46	(-21.72, -21.20)
\mathcal{M}_I	18 103 (11%)	0.20	(0.14, 0.26)	-10.01	(-10.29, -9.71)	0.47	(0.40, 0.54)	-20.82	(-21.10, -20.55)
	12135 (47%)	0.16	(0.12, 0.21)	-10.12	(-10.43, -9.81)	0.51	(0.43, 0.58)	-20.69	(-20.96, -20.43)
$\mathcal{M}_{\mathcal{L}}$	12647 (< 1%)	0.15	(0.11, 0.19)	-9.86	(-10.16, -9.57)	0.39	(0.31, 0.45)	-20.01	(-20.34, -19.69)
	11 648 (~ 14%)	0.14	(0.10, 0.18)	-9.90	(-10.18, -9.62)	0.40	(0.32, 0.46)	-19.95	(-20.27, -19.66)
					REF				
$\mathcal{M}_{\mathcal{H}}$	54 681	0.24	(0.19, 0.29)	-	-	0.67	(0.60, 0.74)	-21.36	(-21.61, -21.10)
\mathcal{M}_I	44 146	0.19	(0.15, 0.24)	-	-	0.55	(0.44, 0.67)	-21.64	(-20.95, -20.33)
$\mathcal{M}_{\mathcal{L}}$	23 615	0.15	(0.11, 0.18)	-	-	0.42	(0.33, 0.50)	-19.91	(-20.26, -19.57)

complete samples are defined to be at least 95% complete⁴ with respect to the bivariate *r*-band magnitude and H α flux selections. While this implies, by definition, that each volume limited luminosity sample is at least 95% volume limited, the same cannot be said about the volume limited stellar mass samples. To achieve a 95% completeness in volume limited stellar mass samples would require the additional consideration of the detectability of a galaxy of a given stellar mass within the survey volume. It is, however, reasonable to assume that the "volume limited stellar mass" samples are close to 95% volume limited given the strong correlation between stellar mass and optical luminosity. For our sample, the 1 σ scatter in stellar mass-luminosity correlation is ~ 0.4 dex. The volume limited REF samples have the same redshift coverage as their SF counterparts, and as such, they are 100% complete with respect to their univariate magnitude selection.

3 CLUSTERING METHODS

In this section, we describe the modelling of the galaxy selection function using GAMA random galaxy catalogues, and introduce two-point galaxy correlation function estimators used in the analysis.

3.1 Modelling of the selection function

To model the selection function, we use the GAMA random galaxy catalogues (Random DMU) introduced in Farrow *et al.* (2015). Briefly, Farrow *et al.* (2015) employ the method of Cole (2011) to generate clones of observed galaxies, where the number of clones generated per galaxy is proportional to the ratio of the maximum

⁴ This completeness is achieved through excluding very low-SFR sources as they can significantly limit the redshift coverage of a volume limited sample, resulting in samples with small number statistics.

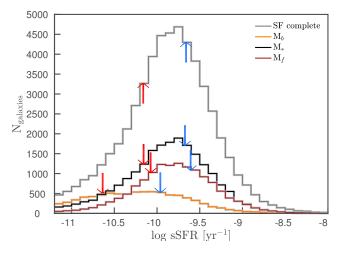


Figure 3. The log sSFR distributions of all SF complete galaxies (grey), as well as M_b , M_* , and M_f galaxies of SF complete sample. The redshift range considered is $0.001 < z \le 0.24$, and the arrows indicate the sSFR cuts used to select the 30% highest (blue arrows) and the 30% lowest (red arrows) sSFR galaxies from each distribution.

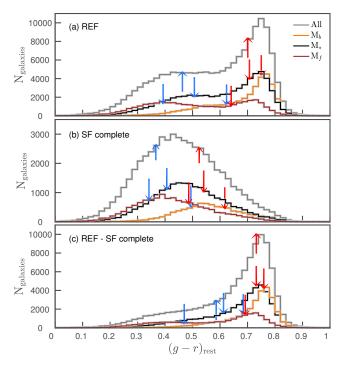


Figure 4. The $(g-r)_{\text{rest}}$ distributions of (a) all REF and (b) all SF complete galaxies, as well as the distributions of their respective M_b , M_* and M_f sub-samples. For completeness, we also show in panel (c) the distributions of REF - SF complete galaxies. The redshift range considered is $0.001 < z \leq 0.24$, and the arrows indicate the colour cuts used to select the 30% bluest (blue arrows) and the 30% reddest (red arrows) colour galaxies from each distribution. The arrows show a clear change in position with luminosity (i.e. arrows move towards redder colours with increasing optical brightness), that is not seen with log sSFR (Figure 3).

volume out to which that galaxy is visible given the magnitude constraints of the survey $(V_{max,r})$ to the same volume weighted by the number density with redshift, taking into account targeting and redshift incompletenesses.

In effect, Random DMU provides N_r , with $\langle N_r\rangle\approx 400$, clones per GAMA galaxy in TilingCatv43. The clones share all intrinsic physical properties (e.g. SFR, stellar mass, etc.) as well as the unique galaxy identification (i.e. CATAID) of the parent GAMA galaxy, and are randomly distributed within the parent's $V_{max,r}$, while ensuring that the angular selection function of the clones matches that of GAMA. Therefore for any galaxy sample drawn from TilingCatv43 based on galaxy intrinsic properties, an equivalent sample of randomly distributed clones can be selected from Random DMU by applying the same selection. If however, a selection involves observed properties, then the clones need to be tagged with "observed" properties before applying the same selection.

In order to select a sample of clones representative of galaxies in SF complete sample, firstly we exclude the clones of GAMA galaxies not part of SF complete sample. Secondly, each clone is assigned an "observed" H α flux based on their redshift and their parent's intrinsic H α luminosity. Finally, the clones with H α fluxes $> 1 \times 10^{-18}$ W/m² and with redshifts outside the wavelength range dominated by the O2 atmospheric band but within the detection range of H α (i.e. SF complete selection criteria) are selected for the analysis. The redshift distribution of the selected clones, hereafter random SF complete, normalised by the approximate number of replications (i.e. $\langle N_r \rangle$) is shown in Figure 6 (green line). Also shown for reference is the redshift distribution of the GAMA SF complete sample (red line). The clear disagreement between the two distributions is a result of the differences in the selections. Recall that only the *r*-band selection of the survey is considered in the generation of clones, i.e. the clones are distributed within their parent's $V_{max,r}$, whereas we also impose an H α flux cut to select the SF complete sample. In essence, we require the clones to be distributed within their parent's $min(V_{max,r}, V_{max,H\alpha})$, where $V_{max,H\alpha}$ is the maximum volume given the H α flux limit, in order to resolve the disagreement between the two distributions.

Instead of regenerating the random DMU with a bivariate selection, we adopt a weighting scheme for the clones, where the original distribution of clones within a given parent's $V_{max,r}$ is altered to a distribution within min($V_{max, r}, V_{max, H\alpha}, V_{zlim}$), where V_{zlim} is the volume out to the detection limit of the H α spectral line in GAMA spectra. The weight of a galaxy, *i*, is defined as

$$N_{\text{weight}}^{i} = \frac{N_{V_{\text{max, r}}}^{i}}{N_{\min(V_{\text{max, H}\alpha}, V_{\text{max, r}}, V_{\text{zlim}})},$$
(1)

where $N_{V_{\text{max, r}}}^{i} \equiv N_{\text{r}}$ is the total number of clones originally generated for the galaxy *i* and distributed within its $V_{\text{max, r}}$, and $N_{\min(V_{\text{max, H}\alpha}, V_{\text{max, r}}, V_{\text{zlim}})}^{i}$ is the number of clones within $\min(V_{\text{max, H}\alpha}, V_{\text{zlim}})$ of the *i*th galaxy.

We show the mean variation of N_{weight} in SFR and M_r space in Figure 5 for three different redshift bins. At a fixed M_r , N_{weight} declines with increasing SFR and redshift, and at a fixed SFR, N_{weight} decreases with increasing optical brightness and decreasing redshift. The implication being that the maximum volume out to which a high-SFR galaxy would be detectable is limited only by the *r*-band magnitude selection of the survey (i.e. no weighting is

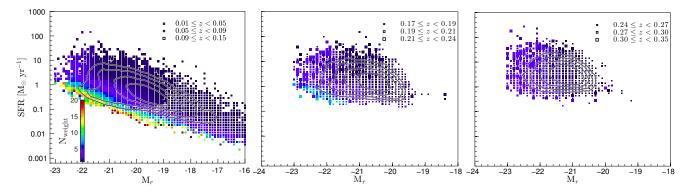


Figure 5. Mean weight applied to the random SF complete sample as a function of their intrinsic SFR and M_r . The size of the markers indicates the mean redshift of GAMA SF complete galaxies with a given SFR and M_r . The closed contours from inwards-to-outwards enclose 25, 50, 75 and 90% of the data in the 0.01 $\leq z < 0.15$, 0.17 < z < 0.24 and 0.24 $\leq z < 0.35$ ranges (left-to-right panels). Only the lowest-redshift sample (left panel) contains galaxies with large N_{weight} measures.

required), and vice versa. For example, a (low SFR) galaxy with $N_{weight} \approx 20$ has ~ 20 clones out of ~ 400 within its $V_{max,\,H\alpha}$. While low SFR galaxies can have larger values of N_{weight} , we demonstrate in Figure 6 that the modelling of the redshift distribution is only very marginally affected by cutting the sample on N_{weight} . Moreover, in Appendix A, we show that the differences between the redshift distributions of clones weighted by N_{weight} with and without removing large values of N_{weight} are minimal. The differences are largely confined to lower redshifts, where most low-SFR systems reside. The impact of galaxies with large values of N_{weight} on the clustering results is, again, minimal, and is not surprising as most of the low-SFR systems with large N_{weight} lie outside the 90% data contour (Figure 5).

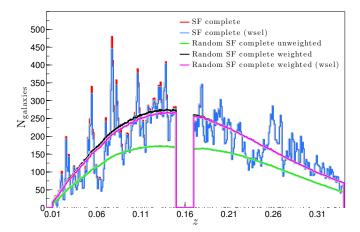


Figure 6. The redshift $(0.01 \leq z < 0.35)$ distribution of the SF complete sample in comparison to the weighted (black and magenta lines) and not-weighted (green) distributions of the random SF complete sample. The weights are determined according to Eq. 1, and the gap in the distributions centred around $z \sim 0.16$ indicates the redshift range where the redshifted H α line overlaps with the atmospheric Oxygen-A band. The galaxies, both GAMA and random, with redshifts in this range are excluded from the analysis as described in § 2.3. Shown also are the weight-selected (wsel) distributions of the SF complete sample and the equivalent weighted random SF complete sample. These distributions exclude all galaxies (and their random clones) with $N_{\text{weights}} > 10$.

A comparison between the redshift distribution of the clones

weighted by N_{weight}, called *random SF complete weighted*, and the distributions of the unweighted clones and GAMA star forming galaxies is presented in Figure 6. We also illustrate the relatively small effect on the weighted distribution if objects with N_{weight} > 10 (i.e. wsel selection in Figure 6) are removed from the analysis. Consequently, the impact on the results of the correlation analyses are also minimal as demonstrated in Appendix A1.

Alternatively, N_{weight} can also be calculated in redshift slices. We refer readers to Appendix A for a discussion on the resulting redshift distributions, mean N_{weight} variations with respect to SFR, M_r , and redshift, as well as on the clustering analysis. The main caveat in calculating N_{weight} in (smaller) redshift slices is that a relatively higher fraction of clones will require larger weights as V_{zlim} now defines the volume of a given redshift slice. For this reason we choose to use N_{weight} calculated assuming a V_{zlim} defined by the detection limit of H α spectral line in GAMA spectra as described above for the clustering analysis presented in subsequent sections.

In summary, in this section, we presented a technique with which the available random clones of GAMA galaxies can be used, without the need to recompute them to take into account any additional constraints resulting from star formation selections.

3.2 Two-point galaxy correlation function

The spatial two-point correlation function, $\xi(r)$, is defined as the *excess* probability dP, relative to that expected for a random distribution, of finding a galaxy in a volume element dV at a distance r from another galaxy (Peebles 1980), i.e.,

$$dP = n[1 + \xi(r)] \, dV, \tag{2}$$

where n is the galaxy number density determined from a given galaxy catalogue.

To disentangle the effects of redshift space distortions from intrinsic spatial clustering, the galaxy CF is often estimated in a two-dimensional grid of pair separations parallel (π) and perpendicular (r_p) to the line of sight, where $r = \sqrt{\pi^2 + r_p^2}$. Using the notation of Fisher *et al.* (1994), for a pair of galaxies with redshift positions \mathbf{v}_1 and \mathbf{v}_2 , we define the redshift separation vector $\mathbf{s} \equiv$ \mathbf{v}_1 - \mathbf{v}_2 and the line of sight vector $\boldsymbol{\ell} \equiv \frac{1}{2}(\mathbf{v}_1+\mathbf{v}_2)$. The parallel and perpendicular separations are then,

$$\pi \equiv |\mathbf{s} \cdot \boldsymbol{\ell}| / |\boldsymbol{\ell}| \quad \text{and} \quad r_p^2 \equiv \mathbf{s} \cdot \mathbf{s} - \pi^2.$$
(3)

The projected two point CF, $\omega_p(r_p)$, obtained by integrating the two-point CF over the line-of-sight (π) direction, then allows the real space $\xi(r)$ to be recovered devoid of redshift distortion effects (Davis and Peebles 1983). The $\omega_p(r_p)$ is defined as,

$$\omega_p(r_p) = 2 \int_{0}^{\pi_{\max}} \xi(r_p, \pi) d\pi = 2 \sum_i \xi(r_p, \pi_i) \Delta \pi_i.$$
(4)

We integrate to $\pi_{\text{max}} \approx 40 h^{-1}$ Mpc, which is determined to be large enough to include all the correlated pairs, and suppress the noise in the estimator (Skibba *et al.* 2009; Farrow *et al.* 2015).

The statistical errors on clustering measures are generally estimated using jackknife resampling (e.g. Zehavi *et al.* 2005a, 2011), using several spatially contiguous subsets of the full sample omitting each of the subsets in turn. The uncertainties are estimated from the error covariance matrix,

$$C_{ij} = \frac{N_{JK}}{N_{JK} - 1} \times \sum_{n=1}^{N_{JK}} [\omega_p^n(r_{p_i}) - \omega_p(r_{p_i})] [\omega_p^n(r_{p_j}) - \omega_p(r_{p_j})],$$
(5)

where N_{JK} is the number of jackknife samples used. We use 18 spatially contiguous subsets (i.e. $N_{JK} = 18$), each covering 16 deg² of the full area, and the results are robust to the number of samples considered (e.g. from 12 to 24).

There are several two-point galaxy CF estimators widely used in the literature (e.g. Hamilton 1993; Landy and Szalay 1993; Davis and Peebles 1983; Peebles and Hauser 1974). Here we adopt the Landy and Szalay (1993) estimator to perform; (i) two-point auto correlation, (ii) two-point cross correlation, and (iii) mark two-point cross correlation analyses, as explained in the subsequent subsections. In Appendix A, we compare the results of Landy and Szalay (1993) with that obtained from the Hamilton (1993) estimator to check whether our results are in fact independent of the estimator used.

3.2.1 Two-point auto correlation function

The two-point auto CF, ξ_a , estimator by Landy and Szalay (1993) is,

$$\xi_a(r_p, \pi)_{\rm LS} = \frac{DD(r_p, \pi)}{RR(r_p, \pi)} - 2\frac{DR(r_p, \pi)}{RR(r_p, \pi)} + 1.$$
 (6)

The DD(r_p , π), RR(r_p , π) and DR(r_p , π) are normalised data-data, random–random and galaxy–random pair counts, and randoms are weighted by N_{weight} (Eq. 1).

3.2.2 Two-point cross correlation function

The estimators given in Eq. 6 and A1 are adapted for the two-point galaxy cross CF, ξ_c , respectively, as follows;

$$\xi_c(r_p,\pi)_{\rm LS} = \frac{D_1 D_2(r_p,\pi) - D_1 R_2(r_p,\pi) - D_2 R_1(r_p,\pi)}{R_1 R_2(r_p,\pi)} + 1,$$
(7)

The $D_1D_2(r_p, \pi)$ is the normalised galaxy–galaxy pair count between data samples 1 and 2, and $R_1R_2(r_p, \pi)$ is the normalised random–random pair count between random clone samples 1 and 2, and the randoms are weighted by N_{weight} as defined in Eq. 1.

The projected cross CFs and their uncertainties are estimated following the same principles as the auto CFs (§ 3.2.1).

Finally, in most cases below, we present GAMA auto and cross correlation functional results relative to the Zehavi *et al.* (2011) power law fit to their $-21 \leq M_r^{0.1} - 5 \log h \leq -20$ sample, hereafter ω_p^{Z11} , given by,

$$\omega_p^{Z11} = \frac{5.33^{\gamma}}{r_p} \Gamma(0.5) \Gamma[0.5(\gamma - 1)] \Gamma(0.5\gamma), \tag{8}$$

where $\gamma = 1.81$.

3.2.3 Two-point mark correlation function

Over the last few decades, numerous clustering studies based on auto and cross correlation techniques have quantitatively characterised the galaxy clustering dependence on galaxy properties in the low-to-moderate redshift Universe. While these studies use the physical information to define galaxy samples for auto and cross correlation analyses, that specific information is not considered in the analysis itself. In other words, galaxies are weighted as "ones" or "zeros" regardless of their physical properties, leading to a potential loss of valuable information. The *mark* clustering statistics, on the other hand, allow physical properties or "marks" of galaxies to be used in the clustering estimation.

The two-point mark CF relates the conventional galaxy clustering to clustering in which each galaxy in a pair is weighted by its mark, therefore, allowing not only clustering as a function of galaxy properties to be measured, but also the spatial distribution of galaxy properties themselves and their correlation with the environment to be efficiently quantified (Sheth *et al.* 2005). As it is the difference between weighted to unweighted clustering at a particular scale that is considered, the mark CF has serval advantages over conventional clustering statistics; (1) it essentially quantifies the degree to which a galaxy mark is correlated with the environment at that scale, and (2) it is less affected by issues related survey/sample selection and incompleteness than conventional methods (Skibba *et al.* 2009). The two-point mark CF is defined as,

$$M(r_p, \pi) = \frac{1 + W(r_p, \pi)}{1 + \xi(r_p, \pi)},$$
(9)

where $\xi(r_p, \pi)$ is the galaxy two-point CF defined above, and $W(r_p, \pi)$ is the weighted CF in which the product of the weights of each galaxy pair taken into account.

For the galaxy pair weighting, we adopt a multiplicative scheme, i.e.,

$$DD(r_p, \pi) = \sum_{ij} \omega_i \times \omega_j, \tag{10}$$

where ω_i is the weight of the *i*th galaxy given by the ratio of its mark to the mean mark across the whole sample. Thus $\frac{1}{N_D} \sum \omega_d^i = 1$ by construction.

The projected mark two-point CF is defined in a similar fashion:

$$E_m(r_p) = \frac{1 + W_p(r_p)/r_p}{1 + \omega_p(r_p)/r_p}.$$
(11)

On large scales, M(r) and $E_m(r_p)$ approach unity (Skibba *et al.* 2009).

Again, we adopt the Landy and Szalay (1993) and Hamilton (1993) clustering estimators for this analysis.

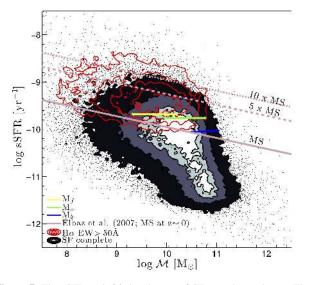


Figure 7. The sSFR and \mathcal{M} distribution of SF complete galaxies. The filled-in and red contours enclose 25, 50, 75, and 94% of SF complete galaxies and SF complete galaxies with H α EW> 50Å (i.e. the "starburst" definition of Rodighiero *et al.* 2011), respectively. The dark pink lines denote the $z \sim 0$ star formation main sequence (solid line, Elbaz *et al.* 2007), and two starburst selections, 5× (dashed line) and 10× the main sequence (dotted line), generally used in the literature (e.g. Rodighiero *et al.* 2011; Silverman *et al.* 2015). The rest of the lines (yellow, green and blue) show the 30% highest sSFR selections applied to the three disjoint luminosity selected galaxy samples used in this analysis (see Table 1).

4 SIGNATURES OF INTERACTION DRIVEN STAR FORMATION

In this study, we consider several different physical properties of galaxies, such as sSFR, colour, dust obscuration and the strength of the 4000 Å break (D₄₀₀₀), that are most likely to be altered in a galaxy-galaxy interaction. A discussion of these properties is given in § 4.1, followed by the results of the auto and cross correlation analyses in § 4.2 and 4.3, respectively. Finally, in § 4.4, we present the results of the mark correlation analysis, where sSFRs and $(g-r)_{\text{rest}}$ of galaxies are used as marks to investigate the spatial correlations of star forming galaxies.

4.1 Characteristics of GAMA star forming galaxies

The enhancement of star formation, or starburst, is perhaps the most important and direct signature of a gravitational interaction (Kennicutt 1998; Wong *et al.* 2011). There are several definitions of "starburst" galaxies. Bolton *et al.* (2012), for example, define "starburst" as SF galaxies with $H\alpha$ equivalent widths (EW), a proxy for sSFR, larger than 50Å. Rodighiero *et al.* (2011); Luo *et al.* (2014), and Knapen and Cisternas (2015) use enhancement of SFR as a

function of stellar mass to identify starbursts. Additionally, the evidence of certain ionised species (e.g. $[Ne III] \lambda 3869 \text{Å}$) indicative of the high ionisation state of gas, as well as the overall enhancement of emission features in galaxy spectra (e.g. [O II], [O III], $H\alpha$, $H\beta$) are other signatures of starbursts (Wild *et al.* 2014). Despite the differences, most "starburst" definitions rely on spectral and/or physical properties of galaxies that are powerful tracers of SFR per unit mass.

The sSFR and \mathcal{M} distribution of star forming galaxies used in this analysis (filled contours) is presented in Figure 7. Over-plotted are several well known "starburst" definitions in the literature; red open contours show the distribution of starbursts (H α EW> 50Å, Rodighiero *et al.* 2011), and the dotted and dashed dark pink lines denote the star formation main sequence (solid dark pink line, Elbaz *et al.* 2007) based starburst definitions (e.g. Rodighiero *et al.* 2011; Silverman *et al.* 2015). The rest of the lines indicate the selection limits of the 30% highest sSFR galaxies of M_b, M_{*} and M_f samples. Note that most of the galaxies selected based on the 30% highest sSFR criterion are in fact those that qualify as starbursts according to the different starburst definitions discussed above.

The signatures of interaction driven star formation that we consider for this analysis are sSFR, SFR, colour, D_{4000} and Balmer decrements, and we use the BPT diagnostics to show (average) variations of these properties in star forming galaxies (Figures 8 and 9). The BPT diagnostics themselves are indicators of gas-phase metallicities (i.e. oxygen abundances) in galaxies (Pettini and Pagel 2004) that can be heavily affected by pristine gas inflows and enriched gas outflows triggered during an interaction. Overall, relatively more massive and lower sSFR galaxies in our star forming sample have higher metallicites (Figure 8) and are characterised by redder optical colours and D_{4000} indices (Figure 9)

Galaxy interactions impact dust to a lesser extent than metallicities as inflowing pristine gas cannot dilute the line-of-sight dust obscuration, though, outflows can remove dust from the interstellar medium. The dust is thought to rapidly build up during a burst of star formation (da Cunha *et al.* 2010; Hjorth *et al.* 2014), giving rise to the observed relationship between dust obscuration and host galaxy SFR (Garn and Best 2010; Zahid *et al.* 2013). This relationship between dust obscuration and SFR is evident in Figure 9 (right panels), where the increment in Balmer decrement approximately mirrors the increase in SFR.

The observed bimodality in optical colours (Baldry *et al.* 2004) can also be used to assess the level of star formation in galaxies. A sudden influx of new stars alters the colour of a galaxy that lasts on time scales that are considerably longer than the parent starburst itself. The trends evident in the distributions of $(g-r)_{rest}$ and D₄₀₀₀ indices (left and middle panels of Figure 9) are such that high sSFR galaxies, including starbursts, are typically characterised with bluer colours.

Overall, SFR or stellar mass alone cannot effectively discriminate a low mass galaxy undergoing a burst of star formation from a quiescently star forming high mass galaxy (see Figure 10). Likewise, optical colour, while indicative of the state of star formation within galaxies, taken alone is insufficient to discriminate starbursts from post starburst and/or dusty starburst systems.

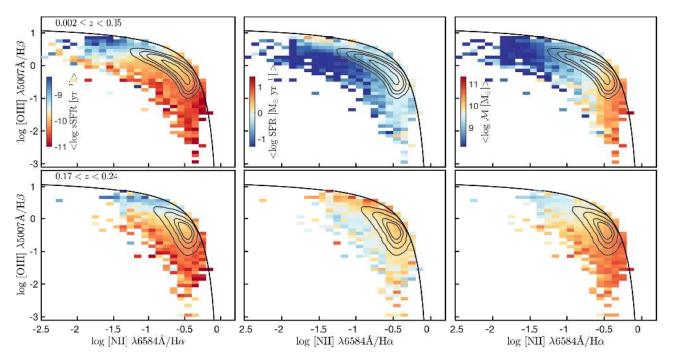


Figure 8. The mean variation in sSFR, SFR, stellar mass (i.e. log \mathcal{M}) of star forming galaxies across the BPT plane in two redshift bins (from top-to-bottom, with the key shown in left-most panels). The mean value of each property in a given $[O_{III}]/H\beta$ and $[N_{II}]/H\alpha$ (i.e. the BPT diagnostics) bin is shown in colour, with the black line denoting the Kauffmann *et al.* (2003a) AGN/SF discrimination criterion. The contours enclose ~25, 50, 75, 90% of the data in each redshift range.

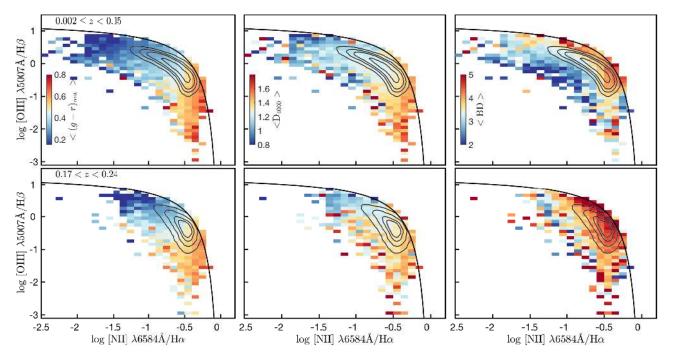


Figure 9. Same as above, now showing the mean variation of $(g - r)_{rest}$, D₄₀₀₀ and Balmer decrement (i.e. BD) of star forming galaxies across the BPT plane in two redshift bins.

4.2 Auto correlation functions of star forming galaxies

The projected auto correlation functions (ACF) of the disjoint luminosity selected samples (Table 1) are presented in the main panels of Figure 11, and the ACFs relative to Zehavi *et al.* (2011) power law fit (ω_p^{z11} , Eq. 8), hereafter ACF $_{\omega_p^{z11}}$, are shown in the top panels. The ACFs of REF versus SF complete galaxies differ significantly over most scales, reflecting the differences in the clustering of the two sets of galaxy populations. These differences are in agreement with the previous clustering studies of the local Universe that find galaxies with bluer optical colours, representative of star forming systems, tend to cluster less strongly than optically redder galaxies

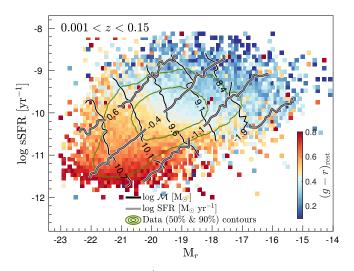


Figure 10. The log sSFR $[yr^{-1}]$ and M_r distribution of z < 0.15 SF complete galaxies, colour-coded by the mean $(g - r)_{rest}$ of galaxies at a given log sSFR and M_r . The thin black and thick grey lines denote the constant stellar mass (in log \mathcal{M} [M_{\odot}]) and log SFR [M_{\odot} yr⁻¹] contours, respectively, that span a relatively large range in both M_r and log sSFR. The green contours enclose 50 and 90% of the data.

(Zehavi *et al.* 2005b; Skibba *et al.* 2009; Zehavi *et al.* 2011; Bray *et al.* 2015; Farrow *et al.* 2015). In our case, REF galaxies comprise of both optically bluer and redder galaxies. Likewise, the ACFs of disjoint stellar mass selected samples (Figure 12) show a qualitative agreement with the ACFs of luminosity selected samples introduced in Figure 11.

On the $-0.15 \leq \log R_p$ [Mpc] ≤ 1.3 range, we find that the ACFs of REF and SF complete galaxies, on average, are consistent with a power-law. On smaller scales (log $R_p \leq -0.15 \text{ Mpc}^5$), however, both sets of functions show varying levels of increase in the strength of clustering with decreasing R_p and optical brightness. This is most clearly evident in the ACF_{$\omega_n^{z_{11}}$}s (i.e. top panels of Figure 11) that demonstrate that at a fixed R_p , the amplitude of ACF_{$\omega_n^{z_{11}}$}s increase with increasing optical brightness. This increase in amplitude appears to be stronger in the ACF_{$\omega_p^{z_{11}}$}s of SF complete galaxies than in REF functions on smaller scales, and vice versa on larger scales. Overall, the behaviour we see on larger scales (log $R_p \gtrsim -0.15$ Mpc) is consistent with other studies that report stronger clustering of massive and luminous galaxies than less massive, low luminosity systems (e.g. Norberg et al. 2001; Marulli et al. 2013; Guo et al. 2014; Zehavi et al. 2005b; Skibba et al. 2009; Zehavi et al. 2011; Bray et al. 2015), and on smaller scales, is mostly consistent with the results of another GAMA study by Farrow et al. (2015).

It is worth noting that even though the ACFs $\omega_p^{z_{11}}$ of SF complete galaxies show lower clustering amplitudes than their respective REF functions on most scales, the change in the strength of the ACFs $\omega_p^{z_{11}}$ of SF complete galaxies with decreasing R_p is greater than that of REF functions. In other words, the ACFs $\omega_p^{z_{11}}$ of SF complete galaxies show a steeper decline (increase) in strength on log R_p $\gtrsim -0.15$ Mpc (log R_p $\lesssim -0.15$ Mpc) with decreasing R_p than REF functions. This rapid increase in the clustering strength of the ACFs $\omega_p^{z_{11}}$ of SF complete galaxies on smaller scales (i.e. excess clustering) suggests increased galaxy-galaxy interactions. The same behaviour is also apparent in the ACFs $\omega_p^{z_{11}}$ of disjoint stellar mass samples of SF complete galaxies (Figure 12).

Interestingly, the R_p at which the ACFs_{$\omega_p^{z_{11}}$} of SF-complete galaxies begin to show an increase in strength seems to also be optical brightness dependent, such that higher optical luminosities correspond to larger R_p and vice versa. For instance, the SF ACF_{$\omega_p^{z_{11}}$} of M_f galaxies show a turn-over in the signal at ~ 0.1 Mpc, though the signal appears to plateau⁶ at a R_p of ~ 0.4 Mpc (or log R_p of -0.4). The SF ACFs_{$\omega_p^{z_{11}}$} of M_{*} and M_b show turn-overs at larger R_p of ~ 0.31 Mpc and ~ 0.5 Mpc (i.e. log R_p of -0.51 and 0.3), respectively. This is in the sense that optically luminous star forming galaxies show an enhancement in clustering at relatively larger separations than their low luminosity counterparts.

As mentioned earlier, R_p provides an alternative metric to assess the interaction-phase of a galaxy pair through the association of large R_p with time elapsed since or time to pericentric passage and small R_p with galaxies currently undergoing a close encounter. One of the advantages of using ACFs to trace interaction-phase is that, aside from the initial sample selection, ACFs are not affected by the properties of galaxies. As such, it is not the net change in a property with R_p that is being assessed, but the change in the clustering strength with R_p within the one- and two-halo terms. Interpreting the change in the strength of the clustering of $ACFs_{\omega_n^{z_{11}}}$ of SF complete galaxies as a signature of increased interactions between galaxies, then any correlation between optical brightness (or stellar mass) and R_p in which a change in the clustering signal takes place can be taken as a signature of a halo size-interaction scale dependence. This suggests that the physical evidence of interactions between star forming galaxies within massive halos are (or ought to be) visible out to larger radii than those between star formers residing in less massive halos. This is also supported by the fact that optically bright star forming galaxies are likely hosted within massive halos.

4.3 Cross correlation functions of star forming galaxies

In this section, we extend the above analysis to further investigate the clustering properties of star formers with respect to different galaxy properties. For this, from each disjoint luminosity (and stellar mass) selected sample, we draw subsamples containing the 30% highest and the 30% lowest sSFRs, $(g - r)_{rest}$, D₄₀₀₀ and Balmer decrements. This selection is detailed in § 2.4. The smaller 30% samples increase the susceptibility of auto correlation results to the effects of small number statistics, hence we utilise cross correlation techniques for the analyses presented in the subsequent sections. Note that all the CCF results shown in the main paper correspond to cross correlations between a given 30% sample and its parent SF complete sample. As part of this analysis, we also investigated the cross correlations between a given 30% sample and its parent REF sample, and we refer readers to Appendix C for a discussion of that investigation.

The CCFs of the 30% highest and the lowest sSFR M_{*} galaxies, and the 30% bluest and the reddest $(g-r)_{rest}$ M_{*} galaxies are presented

⁵ Corresponds to a R_p of ≤ 0.7 [Mpc]

⁶ Plateau here implies that the ACF has the same gradient as $\omega_p^{z_{11}}$.

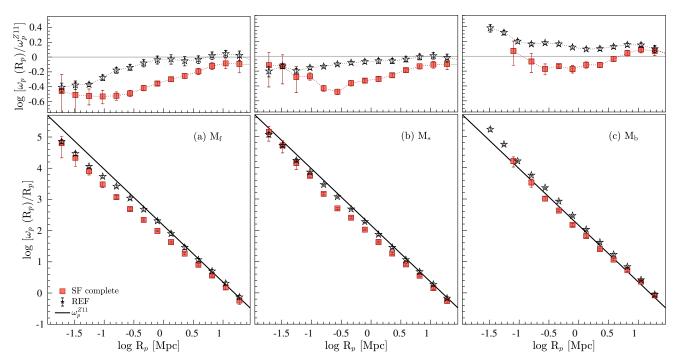


Figure 11. Main panels: The GAMA projected ACFs of luminosity selected (i.e. M_f , M_* and M_b , from left-to-right) REF (open black stars) and SF complete (orange filled squares) samples covering the 0.01 $\leq z \leq 0.34$. The black solid line denote the empirical relation given in Eq. 8 (i.e. ω_p^{Z11}). Top panels: GAMA projected ACFs relative to ω_p^{Z11} . The key is the same as that shown in the left main panel.

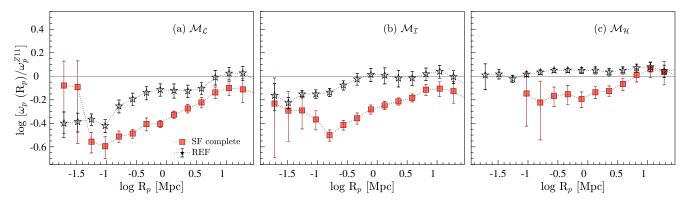


Figure 12. The GAMA projected ACFs of REF (black open symbols) and SF complete (orange filled symbols) stellar mass selected samples (i.e. \mathcal{M}_L , \mathcal{M}_I and \mathcal{M}_H , from left-to-right) relative to ω_p^{Z11} (the key is shown in the left panel).

in the left and right panels of Figure 13, respectively, where each 30% sample is cross correlated with its parent SF complete sample. Also shown in the top panels of Figure 13 are the CCFs relative to ω_p^{Z11} , hereafter CCF ω_p^{Z11} .

Most notable in Figure 13 is, perhaps, the similar clustering excesses on small scales observed for the 30% M_* samples of high sSFR, optically blue, and optically red galaxies. The overlap in clustering amplitudes between high sSFR and optically blue galaxy populations is expected given the correlation between sSFR and optical colour (Figure 10). The overlap between the optically reddest and the highest star forming populations suggests again that a significant fraction of the reddest M_* galaxies in SF complete sample are in fact likely highly dust obscured high sSFR galaxies or starbursts.

In Figures 14 and 15, we compare the $CCFs_{\omega_p^{Z11}}$ of all 30% subsamples drawn from the three disjoint luminosity selected SF complete samples. The top panels of Figure 14 present the $CCFs_{\omega_p^{Z11}}$ of high sSFR galaxies and optically blue galaxies (blue and black symbols), and those of low sSFR and optically red galaxies (red and black symbols) are presented in the bottom panels. To emphasise the degree of the enhancement of the clustering signal on small scales, we fit a linear relation to the log $R_p < -0.64$ Mpc data, where the dashed lines of the same colour denote the bestfitting linear relations to the data of the same colour. Likewise, the $CCFs_{\omega_p^{Z11}}$ of galaxies with low (high) D_{4000} indices and low (high) dust obscurations are presented in the top (bottom) panels of Figure 15. In this figure, for reference, we over plot the best fitting linear relations to the log $R_p < -0.64$ Mpc data shown in Figure 14 as dashed lines.

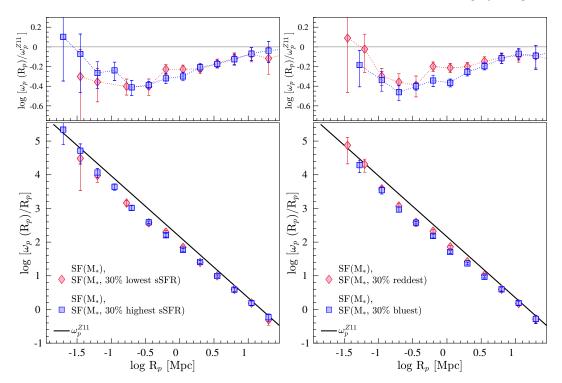


Figure 13. The projected CCFs of the 30% subsets of M_{*} SF complete galaxies. Left panels: The projected CCFs of the 30% highest (blue squares) and the 30% lowest (red diamonds) sSFR galaxies (main panel), and the same functions relative to ω_p^{Z11} (top panel). Right panels: The projected CCFs of the 30% bluest (blue squares) and the 30% reddest (red diamonds) galaxies in $(g - r)_{rest}$ (main panel), and the same functions relative to ω_p^{Z11} (top panel).

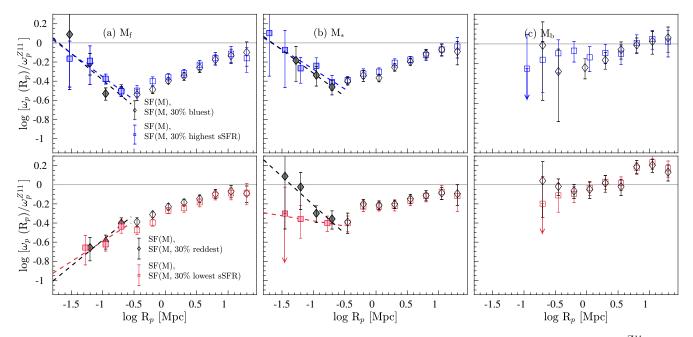


Figure 14. The projected CCFs of high (low) sSFR and optically blue (red) galaxies of luminosity selected SF complete samples relative to ω_p^{Z11} (optical luminosity increases left-to-right). Top panels: the CCFs of optically blue (black diamonds) and high sSFR (blue squares) galaxies. Bottom panels: the CCFs of optically red (black diamonds) and low sSFR (red squares) galaxies. The dashed lines of the same colours denote the best-fitting linear relations to the $R_p < 0.23$ (log $R_p < -0.64$) Mpc data of the same colour. The data points used for the fitting are shown as filled squares, and the arrows denote the data with significant uncertainties.

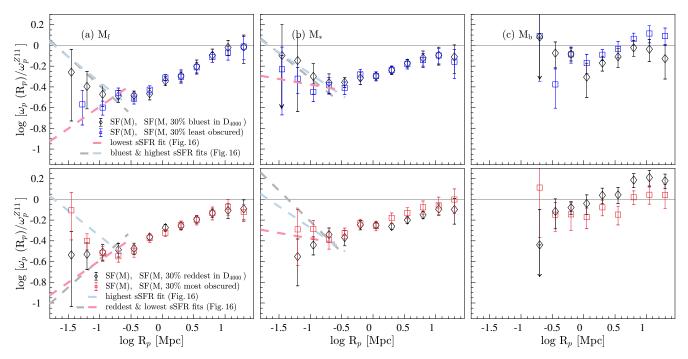


Figure 15. The projected CCFs of least (most) dust obscured and spectroscopically blue (red) galaxies of luminosity selected SF samples relative to ω_p^{Z11} (optical luminosity increases from left to right). Top panels: the CCFs of least dust obscured (the 30% of galaxies with the lowest Balmer decrement measures; blue squares) and spectroscopically blue (the 30% with the lowest D₄₀₀₀ indices; black diamonds) galaxies. Bottom panels: the CCFs of most dust obscured (high Balmer decrement measures; red squares) and spectroscopically red (high D₄₀₀₀ indices; black diamonds) galaxies. For reference, we show the best-fitting linear relations shown in Figure 14 as dashed lines.

For $-0.52 \leq \log R_p$ [Mpc] ≤ 1.3 , all CCFs_{ω_p^{Z11}} show a progressive decline in the strength of clustering with decreasing R_p . The clustering amplitudes of low sSFR and optically red galaxies over this range are, on average, higher than that of their respective high sSFR and optically blue counterparts, in agreement with the studies that find high sSFR galaxies are less clustered than their low sSFR counterparts (e.g. Mostek *et al.* 2013; Coil *et al.* 2016), as well as with the studies that find higher clustering strengths for optically redder galaxies versus optically bluer systems (e.g. Zehavi *et al.* 2011; Ross *et al.* 2014; Favole *et al.* 2016).

Interestingly, on small scales (i.e. $\log R_p \leq -0.52$ Mpc), we see a discrepant behaviour between the $CCF_{\omega_p^{Z_{11}}}$ of optically red galaxies of different luminosity selected samples. The most notable is the $CCF_{\omega_p^{Z_{11}}}$ of M_f galaxies that show a continuous decline in clustering strength with decreasing R_p , whereas the respective $CCF_{\omega_p^{Z_{11}}}$ of M_s and M_b galaxies suggest otherwise. These differences can shed light into dust build up and destruction mechanisms in optically faint (low mass) versus bright (massive) star forming galaxies. At a fixed SFR, an optically faint galaxy would be classified as a starburst, while a luminous system would appear as a normal (or a low) star former (see the distribution of the constant log SFR contours in Figure 10). Therefore to gain further insights into these differences, we add an analysis based on D_{4000} and Balmer decrements (Figure 15), which are complementary to sSFR and $(g-r)_{rest}$, to this study.

The D₄₀₀₀ spectral index is a diagnostic of cumulative star formation history of a galaxy, where lower D₄₀₀₀ indices are indicative of younger stellar populations and vice versa, and is therefore considered a proxy for $(g - r)_{rest}$ (Figure 9). For ease of comparison with the optical colour based analysis discussed above, we hereafter refer to galaxies with lower (higher) D_{4000} indices as spectroscopically blue (red).

The principal advantage of using D_{4000} is that it is less sensitive to dust reddening than $(g-r)_{\text{rest}}^7$. Secondly, it is a spectroscopy-based quantity. In the case of single fibre spectroscopy, the spectrum of a galaxy represents the central region where interaction triggered starbursts are likely to occur (Mihos and Hernquist 1996; Di Matteo et al. 2007; Montuori et al. 2010), whereas photometry based colours represent the light from the whole galaxy. Therefore in galaxies undergoing interactions, where the SFR of centrally triggered starburst, that may also be highly dust obscured, likely mostly contribute to the total SFR, the correlation between SFR and D₄₀₀₀ can be stronger than that between SFR and $(g - r)_{rest}$. Indeed this is evident in Figure 15. The $CCF_{\omega_p^{Z11}}$ of spectroscopically red M_{*} galaxies shows a continuous decline in strength on log $R_p \leq -0.64$ Mpc, whereas the opposite is observed for optically red M* galaxies. In comparison to the $CCF_{\omega_{p}^{Z_{11}}}$ of optically red M_f galaxies, the CCF_{$\omega_n^{Z^{11}}$} of spectroscopically red M_f galaxies too show some differences, though within uncertainties the two $CCFs_{\omega_{p_1}^{Z_{11}}}$ are in agreement. According to these results, the D4000 index appears to be more useful in discriminating starbursts than optical colours. Even though fibre colours are still more susceptible to dust effects than D₄₀₀₀, the correlation between fibre colour and SFR can be stronger than that between global colour and SFR.

The dust obscuration in star forming galaxies has been observed to

⁷ The D_{4000} measures used for this study are based on the Balogh *et al.* (1999) definition, which samples a very narrow range in wavelength.

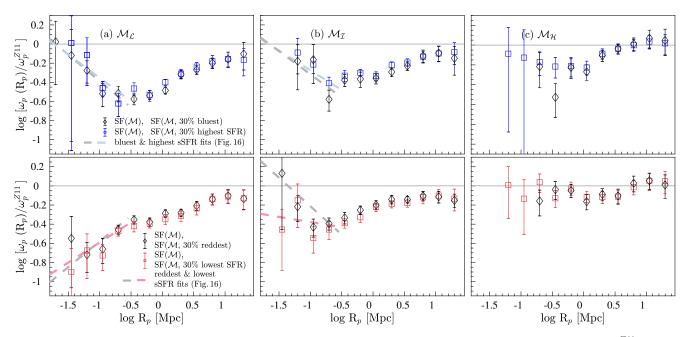


Figure 16. The projected CCFs of high (low) SFR and optically blue (red) galaxies of stellar mass selected star forming galaxies relative to ω_p^{Z11} (Table 2; stellar mass increases left-to-right). Top panels: the CCFs of high SFR (blue squares) and optically blue (black diamonds) galaxies. Bottom panels: the CCFs of low (red squares) and optically red (black diamonds) galaxies. For reference, we show the best-fitting linear relations shown in Figure 14 as dashed lines.

depend on both galaxy SFR and stellar mass (e.g. Brinchmann et al. 2004; Garn and Best 2010; Zahid et al. 2013). Dust is theorised to build up rapidly during a starburst (Hjorth et al. 2014), while a quiescently star forming galaxy experiences a simultaneous decline in dust and SFR as a result of dust destruction and diminishing gas supply (da Cunha *et al.* 2010). The CCFs ω_p^{Z11} of both most and least dust obscured galaxies (Figure 15) show enhancements in clustering amplitudes with decreasing log R_p on ≤ -0.64 Mpc. As shown in Figure 1, the star forming galaxies can have a range of dust obscurations, which can explain the similar enhancements in clustering observed for most and least dust obscured star forming populations of the luminosity selected samples. The clustering excess observed for most dust obscured M* galaxies further supports our earlier assertion that the increase in clustering amplitude of optically red M_{*} galaxies (Figure 13) is, at least in part, caused by the presence of dusty starbursts.

For $-0.52 \leq R_p$ [Mpc] ≤ 1.3 , the CCFs_{ω_p^{Z11}} of the most (least) dust obscured M_f and M_{*} galaxies agree qualitatively with the high (low) sSFR counterparts, as well as with the CCFs_{ω_p^{Z11}} of spectroscopically blue (red) galaxies. The M_b CCFs_{ω_p^{Z11}} of the most and least dust obscured star formers, on the other hand, show an agreement with that of high sSFR M_b galaxies.

For completeness, we present the CCFs ω_p^{Z11} of high (low) SFR and optically blue (red) galaxies of the three disjoint stellar mass selected samples (Table 2) in Figure 16. The dashed lines are the same as in previous figures. These results are, as expected, largely comparable to that observed for high (low) sSFR and optically blue (red) galaxies of luminosity selected samples, and as such, we do not discuss them separately here.

Finally, we also perform a volume limited cross correlation analysis, the results of which are presented and discussed in Appendix B2. Briefly, the $\text{CCFs}_{\omega_R^{Z^{11}}}$ of volume limited samples show

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a qualitative agreement with their respective non-volume limited counterparts on most scales. There are some quantitive differences between the two sets of $\text{CCFs}_{\omega_p^{Z11}}$ on smaller scales that rise as a result of small number statistics.

4.4 The rank-ordered mark correlation functions of star forming galaxies

The mark clustering statistics are different than the auto- and crosscorrelation techniques discussed in the previous sections. The mark statistics can shed light on the dependence of a given physical property on the separation of a galaxy pair by weighting each galaxy in that pair by that physical property. Given this sensitivity of mark correlation functions (MCF) to environmental effects, they form a useful tool in identifying and quantifying underlying correlations of various galaxy properties with environment.

In conventional mark two-point clustering statistics, the correlation function is directly weighted by a given mark, i.e. a physical property (e.g. SFR, sSFR). Consequently, the amplitude of a MCF depends not only on the distribution of marks (Skibba et al. 2006, 2009), but also on the differences in the formulation of a mark (e.g. log or linear; Skibba et al. 2006, 2009). Therefore, unless the distributions of different marks are similar, different MCFs cannot be compared with each other to understand the dependence of different galaxy properties with galaxy separation. In our case, the SFR, sSFR, and $(g - r)_{rest}$ distributions of SF complete samples used differ in shape, in magnitude, and in range. As such, in order to compare the SFR, sSFR and $(g - r)_{rest}$ MCFs, we rank-order the marks and use the rank as the mark. This method, introduced in Skibba et al. (2013), allows the effects of the shape of the distribution on the strength of the MCF to be removed, such that mark correlation signal can be compared between different marks. The

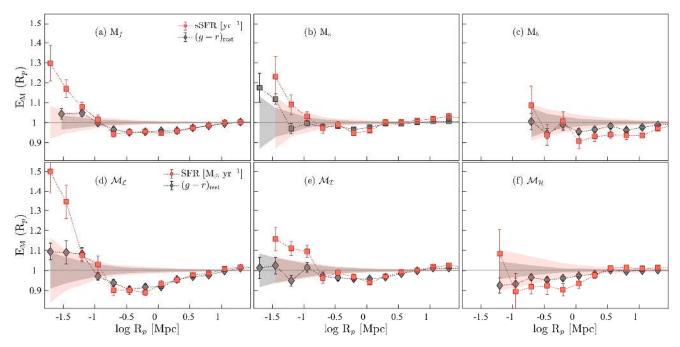


Figure 17. The rank-ordered MCFs of luminosity (top row; optical luminosity increases left-to-right) and stellar mass (bottom row; stellar mass increases left-to-right) selected SF complete samples. The orange and grey filled symbols on the top panels denote the rank-ordered sSFR and $(g - r)_{rest}$ MCFs, respectively, and on bottom panels SFR and $(g - r)_{rest}$ MCFs, respectively. The shaded regions indicate the scatter from randomising the marks.

caveat of this method is that any information contained in the shape of a distribution will be lost.

We present the rank-ordered sSFR and $(g - r)_{rest}$ MCFs of luminosity selected (top panels), and rank-ordered SFR and $(g - r)_{rest}$ MCFs of stellar mass selected (bottom panels) SF complete samples in Figure 17. On small scales, the rank-ordered sSFR and SFR MCFs indicate a clear enhancement in amplitude compared to that of $(g - r)_{rest}$ MCFs. This suggests that sSFR and SFR correlation signals indeed correlate more strongly with the environment than optical colour. The decrement in sSFR, SFR and $(g-r)_{rest}$ mark correlation signals between the $-0.82 \leq \log R_p$ [Mpc] ≤ 0.6 , which is more strongly evident in rank-ordered MCFs of M_f and $M_{\mathcal{L}}$ galaxies, likely demonstrates the effects of post-starbursts, where certain physical properties of a galaxy, e.g. SFR and colour, are affected by the increased presence of now ageing stellar population produced during a starburst.

For completeness, we also present and discuss the conventional MCFs in Appendix D (Figures D1 and D2). The most notable in the conventional case is the strengthening in clustering amplitude with increasing optical brightness observed for sSFR and SFR populations, which mirrors that observed in auto- and cross-correlation functions presented in previous sections.

5 DISCUSSION

In this study, we considered several different star forming properties of galaxies (i.e. SFR, sSFR, $(g - r)_{rest}$, D₄₀₀₀ and Balmer decrement), which are most likely to be affected by galaxy-galaxy interactions. We utilised the [O III] λ 5007Å/H β and [N II] λ 6584Å/H α diagnostics (i.e. BPT, Baldwin *et al.* 1981), which can be used as an indicator of gas phase metallicity, to demonstrate the variation of the physical properties considered with metallicity in star forming galaxies (Figures 8-10). In general, the variation in sSFR largely mirrors that of $(g - r)_{rest}$ and D_{4000} , where low sSFR are typically characterised by lower metallicities. Dust obscuration, on the other hand, indicates a variation similar to that seen with SFR, where high SFR galaxies show a higher dust obscuration than low SFR systems.

Below we discuss the main findings of this study and is structured as follows. A discussion of the results of auto, cross and mark correlation analysis of star forming galaxies is presented in §5.1 to §5.4, and in §5.5, we compare the GAMA results of this study with that of SDSS.

5.1 On the potential interaction-scale halo-size dependence of interaction-driven disturbances

The role that large-scale environment plays in driving and sustaining changes induced during a galaxy-galaxy interaction is understood to a lesser extent than the role of the interaction itself. Generally, the net changes in physical properties of galaxies are used as *direct* indicators of interactions and environmental effects. As mentioned before, the focus of our study is to explore the suitability of utilising *two-point correlation statistics* to shed light on any dependence of galaxy-galaxy interactions on their large scale (i.e. halo-scale) environment. For this, we have computed two-point auto-, cross- and mark-correlation functions of star forming galaxies as a function of both optical luminosity and stellar mass, which approximately correlate with halo mass.

In order to quantify the R_p out to which signatures of interactions ought to persist, in § 4.2, we make the assumption that any change in the relative strength of clustering of a given population reflects

its interaction scale⁸. For example, the ACFs of both luminosity and stellar mass selected star forming galaxies are consistent with a power-law on $-0.15 \leq \log R_p$ [Mpc] ≤ 1.3 . On $\log R_p \leq -0.15$ Mpc, they show a significant clustering excess (Figure 11). This is best seen in ACFs ω_n^{Z11} , where this change appears as a turn-over in the signal. It is this 'turn-over' that we consider to approximately correspond to the interaction scale of that galaxy population. The interaction scales estimated this way appear to depend on galaxy luminosity. This is in the sense that the interaction scale of optically brighter star-forming galaxies is greater than that of optically faint galaxies. This could be interpreted as a signature of a halo sizeinteraction scale dependence, where the evidence of interactions between star formers residing in massive halos are visible out to larger radii than those between star formers residing in low mass halos. This can be, in part, due to massive halos playing a greater role in enhancing and sustaining the effects of galaxy interactions than their less massive counterparts. Equally, this could also be an artefact of high-mass inhabitants of massive halos being able to form stars more efficiently than low-mass galaxies in interactions (Ferreras et al. accepted).

The ACFs of both luminosity and stellar mass selected REF galaxies also show similar changes in the small scale clustering. These changes are, however, not as significant as those observed in starforming galaxies. In comparison to the ACFs of REF galaxies, the star formers show lower clustering amplitudes over most scales, except on log R_p \leq -0.15 Mpc. On log R_p \leq -0.15 Mpc, the ACFs of star forming galaxies show a rapid increase in the amplitude of clustering with decreasing R_p . Consequently, over these scales, the clustering of star forming galaxies appears to be similar to that of REF. Both of these results are consistent with the findings of previous studies; the former with the studies that find optically redder galaxies are more strongly clustered than their bluer counterparts (e.g. Zehavi et al. 2005b; Skibba et al. 2009; Zehavi et al. 2011; Bray et al. 2015), and the latter with the Farrow et al. (2015) clustering study of optically selected red and blue galaxies, finding an up-turn in the clustering of the blue systems on small scales, as well as with the results of Heinis et al. (2009) and Mostek et al. (2013).

5.2 On the direct indicators versus two-point correlation statistics tracing interaction scales

Here we discuss the potential reasons for the differences in R_p reported by the studies that utilise direct probes of interactions (see § 1 for a discussion), as well as between those and the predictions of our auto correlation analysis.

A vast number of competing factors can influence both the strength of an interaction-induced physical change and the R_p out to which the net effect is observable. The orbital parameters, for instance, can play a significant role in moderating the SFR response. Both observational and theoretical studies suggest that retrograde encounters lead to higher star formation efficiencies, and thus higher SFR enhancements, than prograde encounters (Di Matteo *et al.* 2007; Mesa *et al.* 2014). The ratio of the stellar masses of the

progenitors and their gas fractions are two other factors that can significantly influence the strengths of direct indicators. Galaxy pairs with mass ratios between 1 - 3 are observed to have the strongest SFR enhancements (e.g. Ellison et al. 2008; Cox et al. 2006). Likewise, lower gas fractions are theorised to lead to lower SFR enhancements (Di Matteo et al. 2007). While starbursts with the shortest durations tend to typically show the strongest enhancements (Di Matteo et al. 2007) and tend to occur over the smallest separations (typically < $30 h_{70}^{-1}$ kpc; e.g. Ellison *et al.* 2008; Li et al. 2008a; Wong et al. 2011; Scudder et al. 2012; Patton et al. 2013), the smallest separations can, also, inhibit SFR if the tidal forces are strong enough to eject molecular gas to tidal tails without allowing the gas to funnel to the centres of galaxies (Di Matteo et al. 2007). Overall, these competing effects can "wash out" the net signal of direct indicators, thus affecting the observability of a physical change.

The differences in dynamical timescales associated with different star formation probes is another factor that must be considered when using physical properties as tracers of interactions. Davies et al. (2015), based on the GAMA survey data, report that short-duration star formation indicators show stronger signs of enhancement/suppression than long-duration tracers. The H α SFR, for example, is a direct tracer of on-going star formation in galaxies, probing on average the star formation over a shorter timescale (i.e. ~ 10 Myr) than broadband photometry, e.g. (g - r)probes star formation over much longer timescales of ~ 1 Gyr. The short-duration indicators are, therefore, expected to be most vulnerable to recent dynamical events. This suggests that the dynamical timescales of processes that likely trigger short-duration star formation events are also shorter than those of processes that likely trigger long-duration events. The implication being that analyses which rely on observations of net changes are susceptible to the differences in the dynamical timescales of physical processes that trigger and sustain different changes. This can, perhaps, further explain the differences in the reported interaction scales. In this sense, auto correlation techniques offer an alternative to trace interaction scales which is almost⁹ independent of the influences of interaction induced direct observables.

A dependence on the interaction scale and the size of a halo suggests that star formation activity evolves differently in different environments. Elbaz et al. (2007) and Ziparo et al. (2014) report a reversal of the SFR-density relationship at $z \sim 1$, from highdensity environments hosting high-SFR galaxies at earlier times to low-density environments hosting high-SFR systems at later times. Popesso et al. (2015b,a) interpret SFR-density relation and "galaxy downsizing" (Cowie et al. 1996) in terms of "halo downsizing", where the SFR contribution of massive halos to the cosmic SFR density becomes progressively less significant with increasing cosmic time. In the local Universe, the bulk of the stellar mass is locked in galaxy-groups (Eke et al. 2005) so that group-sized halos are the most common type of halos for a star forming galaxy to inhabit. Therefore it is likely that most of the aforementioned studies preferentially selected galaxies residing in one type of a halo (i.e. group-sized halos) over the others. In our study, by using disjoint luminosity and stellar mass samples, we attempt to minimise this preferential selection, as well as the overlap between halos of

⁸ We use the term 'interaction scale' to denote the R_p out to which changes in physical properties ought to be evident instead of 'observable timescale' to avoid confusion, as this term is generally used by studies that rely on net changes in physical properties to trace interactions.

⁹ As we have used H α fluxes to select the star forming sample used for this analysis, our results are not completely independent.

different sizes, thereby giving insight into interactions between star formers in relatively low- versus high-mass halos.

5.3 On the use of cross-correlation techniques in the determination of interaction scales with respect to galaxy properties

As mentioned earlier, galaxy-galaxy interactions have been observed to drive many physical changes in galaxies. The best physical tracer of an interaction can, however, differ depending on the progenitors, the environment and the interaction itself. It has been shown, both theoretical and observationally, that SF-SF galaxy pairs largely favour low-to-moderate density environments, which are typical hosts to low mass galaxies with higher gas fractions, whereas non-SF - non-SF and SF - non-SF galaxy pairs are preferentially found in high-density environments (e.g. Lin et al. 2010; Ellison et al. 2010). As such, while interactions still occur in highdensity environments, they may not always lead to an enhancement in star formation (Ellison et al. 2008, 2010), though can, perhaps, lead to a change in another property such as optical colours. Below we discuss the clustering properties of star forming galaxies in different environments with respect to different star forming properties of galaxies obtained from the cross correlation analysis presented in § 4.3.

On small scales, the CCFs $_{\omega_p^{Z11}}$ of M_f star formers of high (low) sSFRs, high (low) dust obscurations, and bluer (redder) optical and spectroscopic colours show enhancements (decrements) of varying degree in clustering amplitudes with decreasing R_p (Figures 14 and 15). In contrast, all CCFs $_{\omega_p^{Z11}}$ of star forming M_{*} galaxies, except spectroscopically red objects, show enhancements. The most notable are the opposing clustering trends observed between optically red M_f versus M_{*} populations, and likewise between spectroscopically red (i.e. higher D₄₀₀₀ indices) and optically red M_{*} populations. As D₄₀₀₀ is less sensitive to dust effects than optical colours, one of the potential drivers of these discrepancies is dust obscuration.

The clustering excess observed for highly dust obscured M_{*} galaxies and the dearth in clustering observed for spectroscopically red galaxies support the assertion that a large fraction of optically red star forming M_{*} galaxies are likely dusty starbursts. To illustrate this further, in Figure 18, we show the distribution of sSFRs as a function of 250 μ m luminosity¹⁰ (L₂₅₀ [W/Hz]) for all optically red M_f and M_{*} star formers detected in HERSCHEL 250 μ m. The colour code denotes mean dust obscuration as measured by Balmer decrement, and black and red contours indicate the distribution of the 30% reddest M_f and M_{*} galaxies, respectively. The sSFR cuts used to select the 30% highest sSFR galaxies from M_f and M_{*} samples are shown by the dashed lines. The significant overlap between the high sSFR and optically red M_{*} populations demonstrates that the redder optical colours of these star forming systems have been enhanced by dust.

Finally, the differences in the environments typically inhabited by optically faint versus bright galaxies provide another explanation for the differences between CCFs. The galaxy-galaxy interactions in higher density environments have been observed to lead to

¹⁰ The HERSCHEL 250 μ m photometry is drawn from HATLASCATV03 (Smith *et al.* 2011).

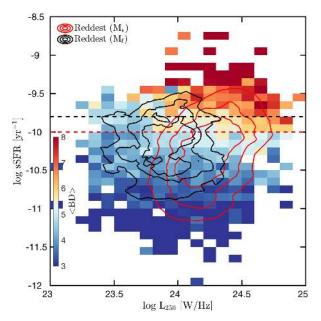


Figure 18. The log L₂₅₀ [W/Hz] and log sSFR [yr⁻¹] distribution of the 30% reddest M_f and M_{*} galaxies of SF complete sample, colour coded by mean dust obscuration (as measured by Balmer decrement, i.e. BD). The contours enclose 25, 50 and 75% of the data, and the dashed lines denote the approximate cuts in log sSFR used to select the 30% highest sSFR M_f (black) and M_{*} (red) galaxies.

quenching of on-going star formation, thus amplifying their redder colours (Patton *et al.* 2011; Ellison *et al.* 2010). This likely also plays a role in enhancing redder colours of M_* galaxies that generally reside in denser environments than fainter systems.

5.4 On the use of mark-correlation techniques in the determination of interaction scales

The mark correlation statistics allow for the dependence of interaction scale on galaxy properties to be investigated. In mark statistics, unlike in cross correlation analysis, the galaxies are either weighted directly by a given physical property or by a rank-order assigned to them based on the distribution of a given physical property. We compute MCFs using both of these methods. The results based on the former (i.e. the conventional) method are presented in Appendix D, and they allow the comparison of MCFs of the same mark between different galaxy samples. Those based on the latter are shown in Figure 17, and permits the comparison of MCFs of different marks between different galaxy samples.

The sSFR and $(g - r)_{rest}$ MCFs based on the conventional method show a strengthening in the mark correlation signal with decreasing R_p , and at a fixed R_p , the strength increases with increasing optical brightness. The same trend is also evident with increasing stellar mass in the SFR and $(g - r)_{rest}$ MCFs of stellar mass selected SF complete samples. The greater enhancement in sSFR observed in the MCF of M_* galaxies than that of M_f galaxies (Figure D1) is in agreement with that expected if the fraction of dusty starbursts with M_* luminosities is higher than those with optically fainter luminosities.

Finally, the comparison of rank-ordered MCFs of sSFR, SFR and $(g-r)_{rest}$ galaxies show that the relative mark correlation strengths

of sSFR MCFs are higher than that of the respective $(g-r)_{rest}$ functions across all luminosity selected SF complete samples. Likewise, the mark correlation strengths of SFR MCFs are higher that of the respective $(g - r)_{rest}$ across all stellar mass selected SF complete samples. This suggests that sSFR and SFR are more sensitive probes of the effects of interactions than optical colours, in agreement with the findings of Heinis *et al.* (2009).

5.5 A comparison between SDSS and GAMA

The CCFs $_{\omega_n^{Z11}}$ of GAMA and SDSS (Li *et al.* 2008a) high and low sSFR galaxies are presented in Figure 19. Li et al. (2008a) define high and low sSFR galaxies as those within the upper and lower 25th percentiles of the sSFR distribution, and they cross correlated with a reference sample containing galaxies in the $0.01 \le z \le 0.3$ and $-23 \leq M_r \leq -17$. Therefore in order to make this comparison as fair as possible, the GAMA $\text{CCFs}_{\omega_p^{Z11}}$ shown in Figure 19 are the CCFs $_{\omega_n^{Z_{11}}}$ obtained from cross correlating SF complete samples with their respective REF samples (see Appendix C for the cross correlation analysis involving REF samples). In general, the SDSS CCFs $\omega_{z^{11}}$ of high sSFR show a good agreement with that of GAMA high sSFR M_{*} galaxies. The low sSFR SDSS function, on the other hand, exhibits a lower clustering strength than the GAMA functions, which is most likely a result of the differences between galaxy samples used for the two studies. For example, even though we show the results of the cross correlation between star forming and reference samples in Figure 19, the redshifts and optical luminosities spanned by the galaxy samples used by Li et al. (2008a) are still larger than the ranges that we considered for our analysis.

In Figure 20, we compare the GAMA sSFR MCFs with the SDSS measures provided in Li *et al.* (2008a). Relative to GAMA, the enhancement in sSFR of SDSS galaxies occurs at a smaller R_p , and the amplitude at a fixed R_p is lower than that of GAMA M_f and M_* sSFR MCFs. Moreover, the GAMA sSFR MCFs of M_f and M_* galaxies show a strengthening in the mark correlation signal with increasing optical luminosity and/or redshift. The enhancement in sSFRs of SDSS galaxies also appears to support this trend, suggesting that interactions between luminous galaxies trigger more intense starbursts than those between faint systems.

We show the GAMA ACFs_{$\omega_p^{Z_{11}}$} (§ 4.2) in the inset of Figure 20 for comparison. On average, the R_p at which the sSFR MCFs of the three luminosity selected SF complete samples show an enhancement in sSFR appear to coincide with the R_p at which the respective ACFs_{$\omega_p^{Z_{11}}$} begin to show a change. This result is not unexpected as MCFs are in a sense ACFs with weights based on the galaxy properties applied.

6 SUMMARY

We have used a sample of galaxies with detected H α emission drawn from the GAMA survey to study the small scale clustering properties of star forming galaxies as a function of both optical luminosity and stellar mass. In the process, we provide a method with which the random clones of galaxies computed by Farrow *et al.* (2015) for the GAMA survey (i.e. computed for a galaxy survey with a univariate primary selection) to be applied to a bivariately selected sample of galaxies (e.g. a star forming sample of galaxies drawn from a broadband survey). The auto, cross and mark twopoint correlation techniques have been used in the computation of galaxy correlation functions for each luminosity and stellar mass selected sample, and below we summarise the main conclusions of this study.

- The strengthening of clustering on small scales observed in GAMA ACFs of star formers (§ 4.2) is a signature of galaxy-galaxy interactions.
- With increasing optical brightness, both the increase in clustering amplitude of star forming population with decreasing R_p at a given R_p (log $R_p \leq -0.15$ [Mpc]) becomes progressively more significant, and the R_p at which the clustering signal of the ACFs of star forming galaxies relative to the fiducial power law show a turn-over becomes progressively larger. This behaviour of star forming galaxies can be interpreted as evidence of an existence of an interaction scale, where physical changes induced in an interaction are, or rather ought to be, evident out to the R_p at which the clustering signal of a given star forming population relative to the fiducial power law starts to alter. This is in the sense that the interactions between optically bright galaxy pairs induce changes that are evident out to larger separations than those between optically faint galaxies.
- The main advantage of utilising auto correlation techniques to map interaction scales is that they are much less susceptible to (1) fluctuations (i.e. enhancements, and decrements) in measured properties, (2) the observability of a change (i.e. the change in a physical property can be too subtle to be observable over some scales), and (3) to the differences arising from the type of star formation indicator used (e.g. short versus long-duration star formation indicators) than methods that employ net changes in properties to trace interactions.
- Out of the different potential signatures of interactions (e.g. sSFR, SFR, optical colour, D_{4000} and Balmer decrement) considered in this study, the clustering with respect to (30%) sSFR and SFR, both based on H α emission, on average show the strongest small scale enhancements across all magnitude and stellar mass ranges considered. Likewise, the 30% lowest sSFR and SFR galaxies show a decrement in clustering across all magnitude and stellar mass ranges.
- The optical colours, i.e. $(g r)_{rest}$, can be affected by the dust obscuration in galaxies. The spectroscopically based D₄₀₀₀ indices, a proxy for colour that less affected by dust and in single-fibre spectroscopy represents the changes in central regions of galaxies, can provide a clearer picture of the effects of interactions than optical colours.
- The comparison between rank ordered sSFR and $(g-r)_{rest}$ MCFs show that the small scale enhancement in sSFR is stronger than that of $(g-r)_{rest}$, supporting the aforementioned conclusion that sSFR is a better tracer of interactions between star forming galaxies than other tracers considered.
- The sSFR MCFs show an increase in small scale clustering, and the amplitude at a given R_p of the MCF of optically bright (e.g. M_*) sSFR galaxies greater than that of optically faint (e.g. M_f) sSFR systems. This suggests that optically brighter star forming systems are characterised by higher SFR than fainter objects. Based on the comparison of dust properties of different star forming populations, it is clear that optically bright high SFR systems contain higher dust contents than their fainter counterparts.

Highly complete datasets with large redshift coverage that will be provided by the future/planned galaxy surveys will allow fur-

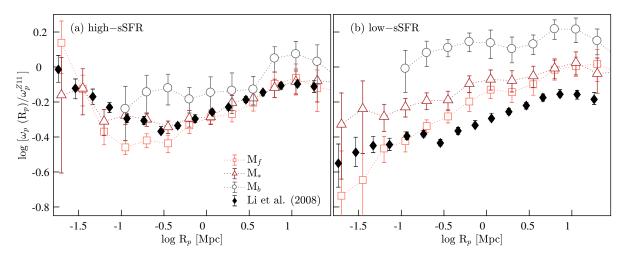


Figure 19. The projected REF CCFs high sSFR (open squares; left panel) and low sSFR (open squares; right panel) M_f , M_* and M_b galaxies relative to ω_p^{Z11} . The filled symbols denote the Li *et al.* (2008a) CCFs of high sSFR (left panel) and low sSFR (right panel) SDSS galaxies. Li *et al.* (2008a) define high and low sSFR galaxies as galaxies contained in the upper and lower 25th percentiles of the distribution of sSFRs of galaxies over the $0.01 \le z \le 0.3$ and $-23 \le M_r \le -17$ ranges.

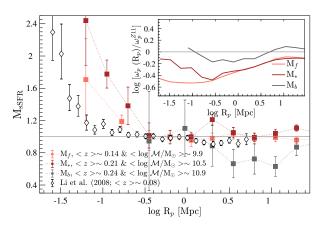


Figure 20. GAMA versus SDSS sSFR MCFs. The GAMA MCFs of M_f , M_* and M_b galaxies (filled symbols) in comparison to the sSFR enhancement functions of SDSS $-23 \le M_{r_{0,1}} \le -17$ galaxies extended over the $0.01 \le z \le 0.3$ range, where $M_{r_{0,1}}$ is *r*-band absolute magnitudes *k*-corrected to z = 0.1 (Li *et al.* 2008a). The mean redshift and stellar mass coverages of each galaxy sample are given in the legend.

ther insights into the relationship between interaction scale and optical brightness, and into underlying physical processes (galaxy and/or cluster-scale) that are responsible for it. Moreover, these datasets will allow any evolution in interaction scales to be tightly constrained, thereby shedding light on the evolution of physical properties and processes of galaxies across time.

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GAMA is a joint European-Australasian project based around a spectroscopic campaign using the Anglo-Australian Telescope. The GAMA input catalogue is based on data taken from the Sloan Digital Sky Survey and the UKIRT Infrared Deep Sky Survey. Complementary imaging of the GAMA regions is being obtained by a number of independent survey programs including GALEX MIS, VST KIDS, VISTA VIKING, WISE, Herschel-ATLAS, GMRT and ASKAP providing UV to radio coverage. GAMA is funded by the STFC (UK), the ARC (Australia), the AAO, and the participating institutions. The GAMA website is http://www.gamasurvey.org/.

Data used in this paper will be available through the GAMA website (http://www.gama-survey.org/) once the associated redshifts are publicly released.

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REFERENCES

- Abazajian, K. N., J. K. Adelman-McCarthy, M. A. Agüeros, S. S. Allam, C. Allende Prieto, D. An, K. S. J. Anderson, S. F. Anderson, J. Annis, N. A. Bahcall, and et al., 2009, ApJS 182, 543-558.
- Adelberger, K. L., C. C. Steidel, M. Pettini, A. E. Shapley, N. A. Reddy, and D. K. Erb, 2005, ApJ **619**, 697.

- Baldry, I. K., S. P. Driver, J. Loveday, E. N. Taylor, L. S. Kelvin, J. Liske, P. Norberg, A. S. G. Robotham, S. Brough, A. M. Hopkins, S. P. Bamford, J. A. Peacock, et al., 2012, MNRAS 421, 621.
- Baldry, I. K., K. Glazebrook, J. Brinkmann, Ž. Ivezić, R. H. Lupton, R. C. Nichol, and A. S. Szalay, 2004, ApJ **600**, 681.
- Baldry, I. K., A. S. G. Robotham, D. T. Hill, S. P. Driver, J. Liske, P. Norberg, S. P. Bamford, A. M. Hopkins, J. Loveday, J. A. Peacock, E. Cameron, S. M. Croom, *et al.*, 2010, MNRAS **404**, 86.

Baldwin, J. A., M. M. Phillips, and R. Terlevich, 1981, PASP 93, 5.

- Balogh, M. L., S. L. Morris, H. K. C. Yee, R. G. Carlberg, and E. Ellingson, 1999, ApJ **527**, 54.
- Barrera-Ballesteros, J. K., S. F. Sánchez, B. García-Lorenzo, J. Falcón-Barroso, D. Mast, R. García-Benito, B. Husemann, G. van de Ven, J. Iglesias-Páramo, F. F. Rosales-Ortega, M. A. Pérez-Torres, I. Márquez, *et al.*, 2015, A&A **579**, A45.
- Bergvall, N., E. Laurikainen, and S. Aalto, 2003, A&A 405, 31.
- Bolton, A. S., D. J. Schlegel, É. Aubourg, S. Bailey, V. Bhardwaj, J. R. Brownstein, S. Burles, Y.-M. Chen, K. Dawson, D. J. Eisenstein, J. E. Gunn, G. R. Knapp, *et al.*, 2012, AJ **144**, 144.
- Bray, A. D., D. J. Eisenstein, R. A. Skibba, M. R. Blanton, A. L. Coil, B. L. Cool, A. L. Mandag, L. Maustalas, and C. Zhu, 2015, ApJ 811, 00
- R. J. Cool, A. J. Mendez, J. Moustakas, and G. Zhu, 2015, ApJ **811**, 90. Brinchmann, J., S. Charlot, S. D. M. White, C. Tremonti, G. Kauffmann,

T. Heckman, and J. Brinkmann, 2004, MNRAS **351**, 1151. Bruzual, G., and S. Charlot, 2003, MNRAS **344**, 1000.

- Calzetti, D., L. Armus, R. C. Bohlin, A. L. Kinney, J. Koornneef, and
- T. Storchi-Bergmann, 2000, ApJ **533**, 682. Cardiel, N., J. Gorgas, and A. Aragon-Salamanca, 1998, MNRAS **298**, 977.
- Casteels, K. R. V., S. P. Bamford, R. A. Skibba, K. L. Masters, C. J. Lintott, W. C. Keel, K. Schawinski, R. C. Nichol, and A. M. Smith, 2013, MNRAS 429, 1051.
- Chabrier, G., 2003, ApJL 586, L133.
- Cid Fernandes, R., G. Stasińska, A. Mateus, and N. Vale Asari, 2011, MNRAS 413, 1687.
- Coil, A. L., A. J. Mendez, D. J. Eisenstein, and J. Moustakas, 2016, ArXiv e-prints 1609.09090.
- Cole, S., 2011, MNRAS 416, 739.
- Colless, M., G. Dalton, S. Maddox, W. Sutherland, P. Norberg, S. Cole, J. Bland-Hawthorn, T. Bridges, R. Cannon, C. Collins, W. Couch, N. Cross, *et al.*, 2001, MNRAS **328**, 1039.
- Condon, J. J., M. A. Condon, G. Gisler, and J. J. Puschell, 1982, ApJ **252**, 102.
- Cowie, L. L., A. Songaila, E. M. Hu, and J. G. Cohen, 1996, AJ 112, 839.
- Cox, T. J., P. Jonsson, J. R. Primack, and R. S. Somerville, 2006, MNRAS **373**, 1013.
- Cox, T. J., P. Jonsson, R. S. Somerville, J. R. Primack, and A. Dekel, 2008, MNRAS **384**, 386.
- da Cunha, E., C. Eminian, S. Charlot, and J. Blaizot, 2010, MNRAS 403, 1894.
- Darg, D. W., S. Kaviraj, C. J. Lintott, K. Schawinski, M. Sarzi, S. Bamford, J. Silk, D. Andreescu, P. Murray, R. C. Nichol, M. J. Raddick, A. Slosar, *et al.*, 2010, MNRAS **401**, 1552.
- Davies, L. J. M., A. S. G. Robotham, S. P. Driver, M. Alpaslan, I. K. Baldry, J. Bland-Hawthorn, S. Brough, M. J. I. Brown, M. E. Cluver, M. J. Drinkwater, C. Foster, M. W. Grootes, *et al.*, 2015, MNRAS **452**, 616.
- Davis, M., and P. J. E. Peebles, 1983, ApJ 267, 465.
- De Propris, R., J. Liske, S. P. Driver, P. D. Allen, and N. J. G. Cross, 2005, AJ **130**, 1516.
- Di Matteo, P., F. Combes, A.-L. Melchior, and B. Semelin, 2007, A&A 468, 61.
- Driver, S., J. Liske, and A. Graham, 2007, Anglo-Australian Observatory Epping Newsletter **111**, 5.
- Driver, S. P., D. T. Hill, L. S. Kelvin, A. S. G. Robotham, J. Liske, P. Norberg, I. K. Baldry, S. P. Bamford, A. M. Hopkins, J. Loveday, J. A. Peacock, E. Andrae, *et al.*, 2011, MNRAS **413**, 971.
- Driver, S. P., A. H. Wright, S. K. Andrews, L. J. Davies, P. R. Kafle, R. Lange, A. J. Moffett, E. Mannering, A. S. G. Robotham, K. Vinsen,

M. Alpaslan, E. Andrae, et al., 2015, ArXiv e-prints 1508.02076.

- Eales, S., L. Dunne, D. Clements, A. Cooray, G. de Zotti, S. Dye, R. Ivison, M. Jarvis, G. Lagache, S. Maddox, M. Negrello, S. Serjeant, *et al.*, 2010, PASP **122**, 499.
- Edge, A., W. Sutherland, K. Kuijken, S. Driver, R. McMahon, S. Eales, and J. P. Emerson, 2013, The Messenger **154**, 32.
- Eke, V. R., C. M. Baugh, S. Cole, C. S. Frenk, H. M. King, and J. A. Peacock, 2005, MNRAS **362**, 1233.
- Elbaz, D., E. Daddi, D. Le Borgne, M. Dickinson, D. M. Alexander, R.-R. Chary, J.-L. Starck, W. N. Brandt, M. Kitzbichler, E. MacDonald, M. Nonino, P. Popesso, *et al.*, 2007, A&A **468**, 33.
- Ellison, S. L., J. T. Mendel, D. R. Patton, and J. M. Scudder, 2013, MNRAS 435, 3627.
- Ellison, S. L., P. Nair, D. R. Patton, J. M. Scudder, J. T. Mendel, and L. Simard, 2011a, MNRAS **416**, 2182.
- Ellison, S. L., D. R. Patton, J. T. Mendel, and J. M. Scudder, 2011b, MNRAS 418, 2043.
- Ellison, S. L., D. R. Patton, L. Simard, and A. W. McConnachie, 2008, AJ 135, 1877.
- Ellison, S. L., D. R. Patton, L. Simard, A. W. McConnachie, I. K. Baldry, and J. T. Mendel, 2010, MNRAS 407, 1514.
- Farrow, D. J., S. Cole, P. Norberg, N. Metcalfe, I. Baldry, J. Bland-Hawthorn, M. J. I. Brown, A. M. Hopkins, C. G. Lacey, J. Liske, J. Loveday, D. P. Palamara, et al., 2015, ArXiv e-prints 1509.02159.
- Favole, G., C. K. McBride, D. J. Eisenstein, F. Prada, M. E. Swanson, C.-H. Chuang, and D. P. Schneider, 2016, MNRAS **462**, 2218.
- Fisher, K. B., M. Davis, M. A. Strauss, A. Yahil, and J. Huchra, 1994, MNRAS **266**, 50.
- Fukugita, M., T. Ichikawa, J. E. Gunn, M. Doi, K. Shimasaku, and D. P. Schneider, 1996, AJ **111**, 1748.
- Garn, T., and P. N. Best, 2010, MNRAS 409, 421.
- Gordon, Y. A., M. S. Owers, K. A. Pimbblet, S. M. Croom, M. Alpaslan, I. K. Baldry, S. Brough, M. J. I. Brown, M. E. Cluver, C. J. Conselice, L. J. M. Davies, B. W. Holwerda, *et al.*, 2017, MNRAS **465**, 2671.
- Groves, B., J. Brinchmann, and C. J. Walcher, 2012, MNRAS 419, 1402.
- Gunawardhana, M. L. P., A. M. Hopkins, J. Bland-Hawthorn, S. Brough, R. Sharp, J. Loveday, E. Taylor, D. H. Jones, M. A. Lara-López, A. E. Bauer, M. Colless, M. Owers, *et al.*, 2013, MNRAS **433**, 2764.
- Gunawardhana, M. L. P., A. M. Hopkins, E. N. Taylor, J. Bland-Hawthorn, P. Norberg, I. K. Baldry, J. Loveday, M. S. Owers, S. M. Wilkins, M. Colless, M. J. I. Brown, S. P. Driver, et al., 2015, MNRAS 447, 875.
- Gunn, J. E., M. Carr, C. Rockosi, M. Sekiguchi, K. Berry, B. Elms, E. de Haas, Ž. Ivezić, G. Knapp, R. Lupton, G. Pauls, R. Simcoe, *et al.*, 1998, AJ **116**, 3040.
- Guo, H., Z. Zheng, I. Zehavi, H. Xu, D. J. Eisenstein, D. H. Weinberg, N. A. Bahcall, A. A. Berlind, J. Comparat, C. K. McBride, A. J. Ross, D. P. Schneider, *et al.*, 2014, MNRAS 441, 2398.
- Hamilton, A. J. S., 1993, ApJ 417, 19.
- Heinis, S., T. Budavári, A. S. Szalay, S. Arnouts, M. A. Aragón-Calvo, T. K. Wyder, T. A. Barlow, K. Foster, F. G. Peter, D. C. Martin, P. Morrissey, S. G. Neff, *et al.*, 2009, ApJ **698**, 1838.
- Hjorth, J., C. Gall, and M. J. Michałowski, 2014, ApJL 782, L23.
- Hopkins, A. M., S. P. Driver, S. Brough, M. S. Owers, A. E. Bauer, M. L. P. Gunawardhana, M. E. Cluver, M. Colless, C. Foster, M. A. Lara-López, I. Roseboom, R. Sharp, *et al.*, 2013, MNRAS **430**, 2047.
- Hopkins, A. M., C. J. Miller, R. C. Nichol, A. J. Connolly, M. Bernardi, P. L. Gómez, T. Goto, C. A. Tremonti, J. Brinkmann, Z. Ivezić, and D. Q. Lamb, 2003, ApJ 599, 971.
- Jimmy, K.-V. Tran, A. Saintonge, G. Accurso, S. Brough, P. Oliva-Altamirano, B. Salmon, and B. Forrest, 2016, ApJ 825, 34.
- Kauffmann, G., T. M. Heckman, C. Tremonti, J. Brinchmann, S. Charlot, S. D. M. White, S. E. Ridgway, J. Brinkmann, M. Fukugita, P. B. Hall, Ž. Ivezić, G. T. Richards, *et al.*, 2003a, MNRAS **346**, 1055.
- Kauffmann, G., T. M. Heckman, S. D. M. White, S. Charlot, C. Tremonti, J. Brinchmann, G. Bruzual, E. W. Peng, M. Seibert, M. Bernardi, M. Blanton, J. Brinkmann, *et al.*, 2003b, MNRAS **341**, 33.
- Kaviraj, S., S. S. Shabala, A. T. Deller, and E. Middelberg, 2015, MNRAS **452**, 774.

- Keel, W. C., R. C. Kennicutt, Jr., E. Hummel, and J. M. van der Hulst, 1985, AJ **90**, 708.
- Kelvin, L. S., S. P. Driver, A. S. G. Robotham, D. T. Hill, M. Alpaslan, I. K. Baldry, S. P. Bamford, J. Bland-Hawthorn, S. Brough, A. W. Graham, B. Häussler, A. M. Hopkins, *et al.*, 2012, MNRAS **421**, 1007.
- Kennicutt, R. C., Jr., 1992, ApJ 388, 310.
- Kennicutt, R. C., Jr., 1998, ARA&A 36, 189.
- Kennicutt, R. C., Jr., K. A. Roettiger, W. C. Keel, J. M. van der Hulst, and E. Hummel, 1987, AJ **93**, 1011.
- Kewley, L. J., and M. A. Dopita, 2002, ApJS 142, 35.
- Kewley, L. J., M. J. Geller, and E. J. Barton, 2006a, AJ 131, 2004.
- Kewley, L. J., B. Groves, G. Kauffmann, and T. Heckman, 2006b, MNRAS **372**, 961.
- Knapen, J. H., and M. Cisternas, 2015, ApJL 807, L16.
- Lambas, D. G., P. B. Tissera, M. S. Alonso, and G. Coldwell, 2003, MNRAS **346**, 1189.
- Landy, S. D., and A. S. Szalay, 1993, ApJ 412, 64.
- Larson, R. B., and B. M. Tinsley, 1978, ApJ 219, 46.
- Li, C., D. A. Gadotti, S. Mao, and G. Kauffmann, 2009, MNRAS 397, 726.
- Li, C., G. Kauffmann, T. M. Heckman, Y. P. Jing, and S. D. M. White, 2008a, MNRAS 385, 1903.
- Li, C., G. Kauffmann, T. M. Heckman, S. D. M. White, and Y. P. Jing, 2008b, MNRAS 385, 1915.
- Li, C., G. Kauffmann, Y. P. Jing, S. D. M. White, G. Börner, and F. Z. Cheng, 2006, MNRAS **368**, 21.
- Liang, Y. C., F. Hammer, H. Flores, N. Gruel, and F. Assémat, 2004, A&A 417, 905.
- Lin, L., M. C. Cooper, H.-Y. Jian, D. C. Koo, D. R. Patton, R. Yan, C. N. A. Willmer, A. L. Coil, T. Chiueh, D. J. Croton, B. F. Gerke, J. Lotz, *et al.*, 2010, APJ **718**, 1158.
- Liske, J., I. K. Baldry, S. P. Driver, R. J. Tuffs, M. Alpaslan, E. Andrae, S. Brough, M. E. Cluver, M. W. Grootes, M. L. P. Gunawardhana, L. S. Kelvin, J. Loveday, *et al.*, 2015, MNRAS **452**, 2087.
- Loh, Y.-S., R. M. Rich, S. Heinis, R. Scranton, R. P. Mallery, S. Salim, D. C. Martin, T. Wyder, S. Arnouts, T. A. Barlow, K. Forster, P. G. Friedman, *et al.*, 2010, MNRAS **407**, 55.
- Lonsdale, C. J., S. E. Persson, and K. Matthews, 1984, ApJ 287, 95.
- Lotz, J. M., P. Jonsson, T. J. Cox, D. Croton, J. R. Primack, R. S. Somerville, and K. Stewart, 2011, ApJ **742**, 103.
- Luo, W., X. Yang, and Y. Zhang, 2014, ApJL 789, L16.
- Ly, C., M. A. Malkan, N. Kashikawa, K. Ota, K. Shimasaku, M. Iye, and T. Currie, 2012, ApJL **747**, L16.
- Madgwick, D. S., E. Hawkins, O. Lahav, S. Maddox, P. Norberg, J. A. Peacock, I. K. Baldry, C. M. Baugh, J. Bland-Hawthorn, T. Bridges, R. Cannon, S. Cole, *et al.*, 2003, MNRAS **344**, 847.
- Martel, H., D. Kawata, and S. L. Ellison, 2013, MNRAS 431, 2560.
- Martin, D. C., J. Fanson, D. Schiminovich, P. Morrissey, P. G. Friedman, T. A. Barlow, T. Conrow, R. Grange, P. N. Jelinsky, B. Milliard, O. H. W. Siegmund, L. Bianchi, *et al.*, 2005, ApJL **619**, L1.
- Marulli, F., M. Bolzonella, E. Branchini, I. Davidzon, S. de la Torre, B. R. Granett, L. Guzzo, A. Iovino, L. Moscardini, A. Pollo, U. Abbas, C. Adami, *et al.*, 2013, A&A 557, A17.
- Mesa, V., F. Duplancic, S. Alonso, G. Coldwell, and D. G. Lambas, 2014, MNRAS 438, 1784.
- Mihos, J. C., and L. Hernquist, 1996, ApJ 464, 641.
- Montuori, M., P. Di Matteo, M. D. Lehnert, F. Combes, and B. Semelin, 2010, A&A **518**, A56.
- Mostek, N., A. L. Coil, M. Cooper, M. Davis, J. A. Newman, and B. J. Weiner, 2013, ApJ **767**, 89.
- Nikolic, B., H. Cullen, and P. Alexander, 2004, MNRAS 355, 874.
- Norberg, P., C. M. Baugh, E. Hawkins, S. Maddox, D. Madgwick, O. Lahav, S. Cole, C. S. Frenk, I. Baldry, J. Bland-Hawthorn, T. Bridges, R. Cannon, *et al.*, 2002, MNRAS **332**, 827.
- Norberg, P., C. M. Baugh, E. Hawkins, S. Maddox, J. A. Peacock, S. Cole, C. S. Frenk, J. Bland-Hawthorn, T. Bridges, R. Cannon, M. Colless, C. Collins, *et al.*, 2001, MNRAS **328**, 64.

- Owers, M. S., C. Blake, W. J. Couch, M. B. Pracy, and K. Bekki, 2007, MNRAS 381, 494.
- Patton, D. R., S. L. Ellison, L. Simard, A. W. McConnachie, and J. T. Mendel, 2011, MNRAS 412, 591.
- Patton, D. R., P. Torrey, S. L. Ellison, J. T. Mendel, and J. M. Scudder, 2013, MNRAS **433**, L59.
- Peebles, P. J. E., 1980, The large-scale structure of the universe.
- Peebles, P. J. E., and M. G. Hauser, 1974, ApJS 28, 19.
- Pettini, M., and B. E. J. Pagel, 2004, MNRAS 348, L59.
- Popesso, P., A. Biviano, A. Finoguenov, D. Wilman, M. Salvato, B. Magnelli, C. Gruppioni, F. Pozzi, G. Rodighiero, F. Ziparo, S. Berta, D. Elbaz, *et al.*, 2015a, A&A 574, A105.
- Popesso, P., A. Biviano, A. Finoguenov, D. Wilman, M. Salvato, B. Magnelli, C. Gruppioni, F. Pozzi, G. Rodighiero, F. Ziparo, S. Berta, D. Elbaz, *et al.*, 2015b, A&A **579**, A132.
- Regan, M. W., and P. J. Teuben, 2004, ApJ 600, 595.
- Robotham, A., S. P. Driver, P. Norberg, I. K. Baldry, S. P. Bamford, A. M. Hopkins, J. Liske, J. Loveday, J. A. Peacock, E. Cameron, S. M. Croom, I. F. Doyle, *et al.*, 2010, PASA 27, 76.
- Robotham, A. S. G., S. P. Driver, L. J. M. Davies, A. M. Hopkins, I. K. Baldry, N. K. Agius, A. E. Bauer, J. Bland-Hawthorn, S. Brough, M. J. I. Brown, M. Cluver, R. De Propris, *et al.*, 2014, MNRAS **444**, 3986.
- Rodighiero, G., E. Daddi, I. Baronchelli, A. Cimatti, A. Renzini, H. Aussel, P. Popesso, D. Lutz, P. Andreani, S. Berta, A. Cava, D. Elbaz, *et al.*, 2011, ApJL **739**, L40.
- Rogers, B., I. Ferreras, S. Kaviraj, A. Pasquali, and M. Sarzi, 2009, MN-RAS **399**, 2172.
- Ross, A. J., L. Samushia, A. Burden, W. J. Percival, R. Tojeiro, M. Manera, F. Beutler, J. Brinkmann, J. R. Brownstein, A. Carnero, L. A. N. da Costa, D. J. Eisenstein, *et al.*, 2014, MNRAS 437, 1109.
- Sabater, J., P. N. Best, and T. M. Heckman, 2015, MNRAS 447, 110.
- Sanders, D. B., N. Z. Scoville, J. S. Young, B. T. Soifer, F. P. Schloerb, W. L. Rice, and G. E. Danielson, 1986, ApJL **305**, L45.
- Scott, C., and S. Kaviraj, 2014, MNRAS 437, 2137.
- Scudder, J. M., S. L. Ellison, P. Torrey, D. R. Patton, and J. T. Mendel, 2012, MNRAS **426**, 549.
- Sharp, R., W. Saunders, G. Smith, V. Churilov, D. Correll, J. Dawson, T. Farrel, G. Frost, R. Haynes, R. Heald, A. Lankshear, D. Mayfield, et al., 2006, in Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series, volume 6269 of Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series, p. 0, astro-ph/0606137.
- Sheth, R. K., A. J. Connolly, and R. Skibba, 2005, ArXiv Astrophysics e-prints astro-ph/0511773.
- Silverman, J. D., E. Daddi, G. Rodighiero, W. Rujopakarn, M. Sargent, A. Renzini, D. Liu, C. Feruglio, D. Kashino, D. Sanders, J. Kartaltepe, T. Nagao, *et al.*, 2015, ApJL **812**, L23.
- Skibba, R., R. K. Sheth, A. J. Connolly, and R. Scranton, 2006, MNRAS **369**, 68.
- Skibba, R. A., S. P. Bamford, R. C. Nichol, C. J. Lintott, D. Andreescu, E. M. Edmondson, P. Murray, M. J. Raddick, K. Schawinski, A. Slosar, A. S. Szalay, D. Thomas, *et al.*, 2009, MNRAS **399**, 966.
- Skibba, R. A., R. K. Sheth, D. J. Croton, S. I. Muldrew, U. Abbas, F. R. Pearce, and G. M. Shattow, 2013, MNRAS **429**, 458.
- Smith, B. J., C. Struck, M. Hancock, P. N. Appleton, V. Charmandaris, and W. T. Reach, 2007, AJ **133**, 791.
- Smith, D. J. B., L. Dunne, S. J. Maddox, S. Eales, D. G. Bonfield, M. J. Jarvis, W. Sutherland, S. Fleuren, E. E. Rigby, M. A. Thompson, I. K. Baldry, S. Bamford, *et al.*, 2011, MNRAS **416**, 857.
- Snijders, L., L. J. Kewley, and P. P. van der Werf, 2007, ApJ 669, 269.
- Soifer, B. T., M. Rowan-Robinson, J. R. Houck, T. de Jong, G. Neugebauer, H. H. Aumann, C. A. Beichman, N. Boggess, P. E. Clegg, J. P. Emerson, F. C. Gillett, H. J. Habing, *et al.*, 1984, ApJL **278**, L71.
- Solomon, P. M., and L. J. Sage, 1988, ApJ 334, 613.
- Taylor, E. N., A. M. Hopkins, I. K. Baldry, M. J. I. Brown, S. P. Driver, L. S. Kelvin, D. T. Hill, A. S. G. Robotham, J. Bland-Hawthorn, D. H. Jones, R. G. Sharp, D. Thomas, *et al.*, 2011, MNRAS **418**, 1587.
- Tonry, J. L., J. P. Blakeslee, E. A. Ajhar, and A. Dressler, 2000, ApJ **530**, 625.

Toomre, A., and J. Toomre, 1972, ApJ 178, 623.

- Wild, V., O. Almaini, M. Cirasuolo, J. Dunlop, R. McLure, R. Bowler, J. Ferreira, E. Bradshaw, R. Chuter, and W. Hartley, 2014, MNRAS 440, 1880
- Wong, K. C., M. R. Blanton, S. M. Burles, A. L. Coil, R. J. Cool, D. J. Eisenstein, J. Moustakas, G. Zhu, and S. Arnouts, 2011, ApJ 728, 119.
- Wright, E. L., P. R. M. Eisenhardt, A. K. Mainzer, M. E. Ressler, R. M. Cutri, T. Jarrett, J. D. Kirkpatrick, D. Padgett, R. S. McMillan, M. Skrutskie, S. A. Stanford, M. Cohen, et al., 2010, AJ 140, 1868.
- Xu, C. K., D. Domingue, Y.-W. Cheng, N. Lu, J. Huang, Y. Gao, J. M.
- Mazzarella, R. Cutri, W.-H. Sun, and J. Surace, 2010, ApJ 713, 330. Yoon, J. H., K. Schawinski, Y.-K. Sheen, C. H. Ree, and S. K. Yi, 2008,
- ApJS 176, 414-423.
- Young, J. S., L. Allen, J. D. P. Kenney, A. Lesser, and B. Rownd, 1996, AJ 112, 1903.
- Zahid, H. J., R. M. Yates, L. J. Kewley, and R. P. Kudritzki, 2013, ApJ 763. 92.
- Zehavi, I., D. J. Eisenstein, R. C. Nichol, M. R. Blanton, D. W. Hogg, J. Brinkmann, J. Loveday, A. Meiksin, D. P. Schneider, and M. Tegmark, 2005a, ApJ 621, 22.
- Zehavi, I., Z. Zheng, D. H. Weinberg, M. R. Blanton, N. A. Bahcall, A. A. Berlind, J. Brinkmann, J. A. Frieman, J. E. Gunn, R. H. Lupton, R. C. Nichol, W. J. Percival, et al., 2011, ApJ 736, 59.
- Zehavi, I., Z. Zheng, D. H. Weinberg, J. A. Frieman, A. A. Berlind, M. R. Blanton, R. Scoccimarro, R. K. Sheth, M. A. Strauss, I. Kayo, Y. Suto, M. Fukugita, et al., 2005b, ApJ 630, 1.
- Ziparo, F., P. Popesso, A. Finoguenov, A. Biviano, S. Wuyts, D. Wilman, M. Salvato, M. Tanaka, K. Nandra, D. Lutz, D. Elbaz, M. Dickinson, et al., 2014, MNRAS 437, 458.

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APPENDIX A: ON THE MODELLING OF THE SURVEY SELECTION FUNCTION AND ON THE IMPACT OF SAMPLE SYSTEMATICS

Modelling of the survey selection function in redshift A1 bins

We describe the modelling of the selection function used for the analysis presented in the main paper in § 3.1, which is based on Eq. 1 with V_{zlim} set either to the redshift detection limit of 0.34 of the H α spectral feature in GAMA spectra or to $z \sim 0.24$. Alternatively, the Nweights can also be computed in redshift slices such that V_{zlim} defines the volume of a given redshift slice. Figure A1 shows the distributions N_{galaxies} in redshift bins, where N_{weights} computed in redshift slices are used to weight the random galaxies. Note that the redshift ranges are defined such that none include the redshift band centred around $z \sim 0.16$.

Figure A2 shows the mean distribution of N_{weights} with respect to both SFR and optical luminosity. As discussed in the main paper, the V_{max} of each galaxy is used in the computation of its N_{weight} , and V_{max} can either be limited by the galaxy's SFR (i.e. H α flux), or by its r-band magnitude or by the upper limit of the relevant redshift slice. The solid red lines show approximate regions in SFR and M_r plane where a galaxy with a given N_{weight} lie if the V_{max} of that galaxy is limited by its SFR (lower regions), by rband magnitude (upper region) or by the upper redshift limits of the relevant redshift bin (left most region). The main caveat of the calculation of Nweights in redshift slices is that due to the relatively

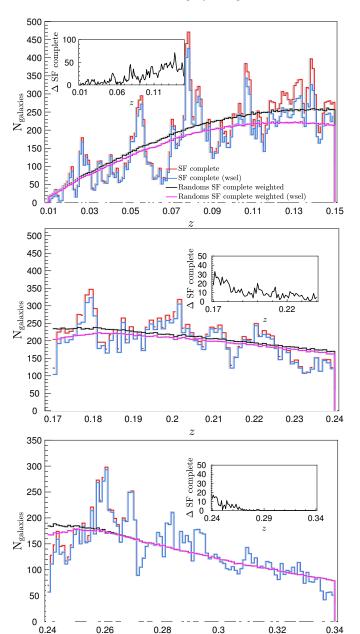


Figure A1. The redshift distributions of SF complete galaxies in comparison to the weighted and weight-selected, where galaxies with $N_{\text{weights}} > 10$ are removed from the sample, distributions of the equivalent random galaxies. The randoms are weighted by N_{weights} computed in redshift slices, and each inset show the difference between SF complete and SF complete (wsel).

z

0.3

0.32

0.24

narrow range in redshift sampled a large fraction of galaxies are assigned $N_{\text{weights}} > 10$.

Figure A3 presents a comparison of the ACFs computed using N_{weight} estimated as described in § 3.1 (right panels) with that computed using N_{weight} estimated in redshift bins (left panels) as described above. In each panel, we compare the ACFs computed from the (Landy and Szalay 1993, open and filled orange squares) estimator with (open orange squares) and without (filled orange squares) with the respective ACFs obtained from the (Hamilton

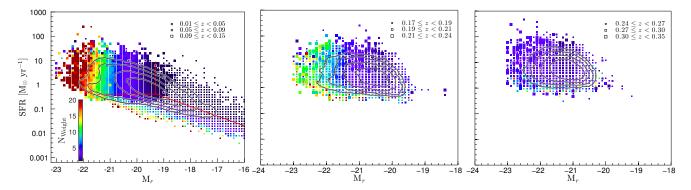


Figure A2. The weights applied to the galaxies in random SF complete as a function of SFR and M_r in 0.01 $\leq z < 0.15$, 0.17 < z < 0.24 and 0.24 $\leq z < 0.35$ ranges. The weights shown are estimated using Eq. 1, where V_{zlim} is the volume of a given redshift slice. From left-to-right, the redshift slices are 0.01 $\leq z \leq 0.15$, 0.17 $< z \leq 0.24$ and 0.24 $< z \leq 0.34$. The marker size is indicative of < z > of SF complete galaxies with a given SFR and M_r measures. The closed contours (from inwards-to-outwards) enclose 25, 50, 75 and 90% of the data, respectively. The red solid lines shown in the left panel approximately indicate the different regions in the SFR and M_r plane where V_{max} of a given galaxy is mostly limited by its measured H α flux (lower region), or by its *r*-band magnitude (upper region) or by zlim (left most region).

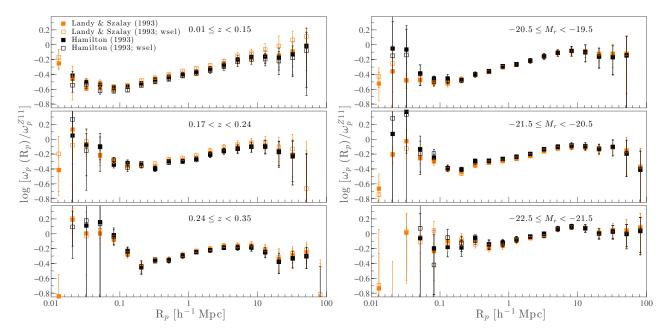


Figure A3. A comparison of ACFs computed from the Landy and Szalay (1993) and Hamilton (1993) estimators (orange and black symbols, respectively) with and without N_{weight} selections (open and filled, respectively). Left panels show the ACFs computed using random galaxies weighted as described in § A, and the right panels show the ACFs computed as described above

1993, open and filled black squares) estimator,

$$\xi_a(r_p, \pi)_{\rm H} = \frac{DD(r_p, \pi) \times RR(r_p, \pi)}{DR(r_p, \pi)^2} - 1.$$
(A1)

The comparison shows that both clustering estimators as well as both methods of estimating N_{weight} yield similar auto correlation results. Given the outcome of this comparison and the caveats associated with estimating N_{weight} in redshift bins (see above), we choose to use a larger redshift range for the analysis presented in the main paper.

Finally, we find that our method of modelling the selection function over the full redshift range over which the H α spectral feature is visible fails to model the low sSFR galaxy population over the 0.24 < z < 0.34 range. This is demonstrated in Figure A4. While the Landy and Szalay (1993) and Hamilton (1993) clustering estimators produce consistent CF results for the high sSFR M_{*} galaxies, the low sSFR M_{*} CFs show a systematic offset. The high versus low sSFR galaxy redshift distributions shown in the insets highlight this issue; modelling the selection function over the full redshift range fails to model the redshift distribution of low sSFR M_{*} galaxies over the 0.24 < z < 0.34 range (i.e. random galaxy redshift distribution of low sSFR M_{*} galaxies is under-predicted). However, limiting the redshift from 0.34 to 0.24 produces consistent results. Therefore for the cross correlation analysis presented in the main paper, we only use galaxies with redshifts in the $0.01 \le z \le 0.24$ range, and re-model the selection function to match this redshift range.

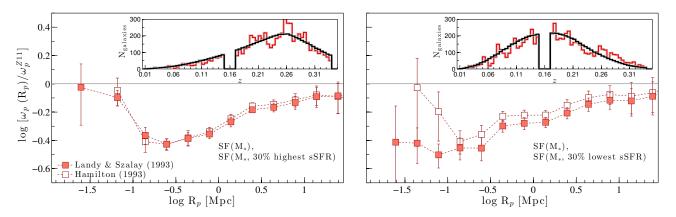


Figure A4. A comparison between Landy and Szalay (1993, filled symbols) and Hamilton (1993, open symbols) cross correlation estimators. Left panel: The projected CCFs of M_* SF galaxies with respect to ω_p^{Z11} (i.e. the reference function introduced in Figure 11), cross correlating all M_* SF galaxies with the 30% highest sSFR galaxies of M_* SF sub-sample. Left panel: The projected cross CF corresponding to the cross correlation between all M_* SF galaxies with the 30% lowest sSFR galaxies of M_* SF sub-sample. Left panel: The projected cross CF corresponding to the cross correlation between all M_* SF galaxies with the 30% lowest sSFR galaxies of M_* SF sub-sample. All galaxy samples used for this figure cover the full redshift range (i.e. z < 0.35) over which Balmer H α feature is visible. As such the weights for the random clones are calculated assuming the maximum redshift of $z \approx 0.35$. The insets show the redshift distributions of the highest (left) and lowest (right) sSFR galaxies of M_* SF sub-sample (red line), and their respective random clones (black line). While the redshift distribution of the random clones of the highest sSFR galaxies match the respective distribution of GAMA galaxies, there is a discrepancy between the two lowest sSFR distributions at higher redshift (i.e. $z \ge 0.24$), which in turn give rise to the systematic discrepancy evident between the cross CFs calculated from the Landy and Szalay (1993, filled symbols) and Hamilton (1993, open symbols) estimators.

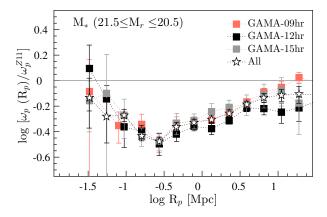


Figure A5. The projected ACFs of M_* galaxies relative to ω_p^{Z11} . The squares denote the CFs corresponding to the three equatorial fields, and stars to all SF M_* galaxies in our sample.

A2 The sample selection & systematics

The selection of the reference (i.e. REF) and SF complete samples is described in detail in § 2.3. Here we investigate how spectroscopic incompleteness of the star forming sample as well as our definition of star forming galaxies impact our results.

A2.1 The lack of 2dFGRS data

One of the main issues discussed in § 2.3, in relation to the selection of the star forming galaxy sample is the incompleteness introduced by the exclusion of 2dFGRS data. Figure 2 demonstrates that our clustering sample is incomplete approximately between 17.7 and 18.8 in apparent *r*-band magnitude, and between $\sim 1 - 0.3$ in $(g - r)_{app}$. The fact that this incompleteness is not randomly distributed over the optical colour and apparent magnitude plane can be problematic for a clustering analysis.

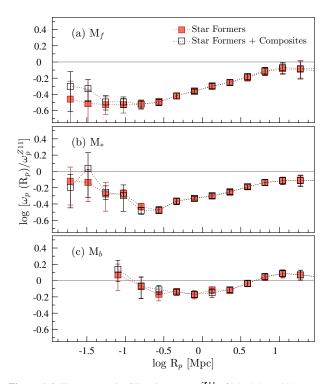


Figure A6. The projected ACFs relative to ω_p^{Z11} of M_f , M_* and M_b pure star forming galaxies (filled squares) in comparison to pure star forming and composites (open squares).

Figure 2 shows the completeness as a function of colour and magnitude for the three GAMA fields individually, as well as the total completeness. The overlap between the 2dFGRS and GAMA surveys is largest in GAMA-12hr field, followed by GAMA-15hr field. The GAMA-09hr field, on the other hand, lies completely outside of the sky regions surveyed by the 2dFGRS survey. Consequently, the spectroscopic incompleteness is significant in GAMA-12hr and relatively insignificant in GAMA-09hr. Therefore to investigate the impact of this incompleteness, we construct the ACFs of M_{*} SF complete galaxies in GAMA-09hr, GAMA-12hr, and GAMA-15hr (Figure A5). Also shown is the ACF of all M_{*} SF complete galaxies. As expected, the ACFs_{ω_p^{Z11}} of M_{*} SF complete galaxies in GAMA-09hr and GAMA-12hr show the largest differences. Despite these differences, however, the two ACFs are in agreement with the ACF of all M_{*} SF complete galaxies to within uncertainties. Note that the differences between the ACFs of individual GAMA fields are not only a result of the differences in spectroscopic completeness between the fields but also reflect sample variance.

Additionally, we have also quantified the impact of excluding the 2dFGRS data on the correlation results by modifying the GAMA redshift completeness to account for the missing 2dFGRS galaxies. The comparison between the correlation functions computed using this modified GAMA redshift completeness mask with that computed using the standard redshift completeness mask (i.e. GAMA main sample of galaxies described in § 2.3) shows that the differences are minimal, and are within the measurement uncertainties.

A2.2 The AGN selections

For the analysis presented in the main paper, we selected the star forming galaxies based on the prescription of Kauffmann *et al.* (2003a). The Kewley and Dopita (2002) prescription is another popular SF/AGN discriminator widely used in the literature. Generally, the Kauffmann *et al.* (2003a) prescription is used to select 'purely' star forming galaxies, while that of Kewley and Dopita (2002) discriminates between galaxies with line emission likely significantly contaminated by the emission from AGNs, and galaxies with line emission likely mostly dominated by massive star formation. The latter class can include objects with some contamination from AGNs (i.e. composites).

In order to understand the impact of the inclusion of composites, we compare the ACFs_{$\omega_p^{Z_{11}}$} of star forming M_f, M_{*} and M_b galaxies selected using the prescription of (Kewley and Dopita 2002, , SF + Composite) with the respective corresponding to star forming galaxies selected using the prescription of Kauffmann *et al.* (2003a). The results of this comparison shown in Figure A6 are qualitatively and quantitatively in agreement with each other. This implies that the composites are galaxies dominated by on-going massive star formation as the AGNs have been observed to have lower clustering amplitudes than star forming galaxies (Li *et al.* 2006, 2008b).

APPENDIX B: VOLUME LIMITED CLUSTERING ANALYSIS

As mentioned in § 2.4, we define several volume limited SF complete samples that are ~ 95% complete with respect to the bivariate r-band magnitude and H α flux selections. In order to achieve this completeness without significantly limiting the redshift coverage of each volume limited SF sample, we impose a low-SFR cut, which excludes very low-SFR galaxies from the sample. The volume limited magnitude samples are, by definition, 95% volume limited, however, the same cannot be said about the volume limited stellar mass samples. In order for the stellar mass samples to be 95% volume limited, we need to consider the maximum volume out which a galaxy of a given stellar mass would be detectable, which has not been taken into account in this analysis. However, given the correlation between stellar mass and optical brightness, the volume limited stellar mass selected samples are likely close to 95% volume limited. Furthermore, we also define several volume limited REF samples have the same redshift coverage as their SF counterpart, which are, therefore, 100% complete with respect the *r*-band magnitude selection of the GAMA survey. Table B1 presents the SF complete and REF volume limited samples used for the clustering analyses.

We present and discuss the ACFs and CCFs constructed using magnitude (Table 1) and stellar mass selected (Table 2) non-volume limited SF complete samples in § 4.2 and § 4.3 of the main paper, respectively. In the subsequent sections, we present and discuss the respective ACFs and CCFs computed using the volume limited samples described in Table B1.

B1 Auto correlation functions of volume limited star forming and REF samples

Figure B1 presents the ACFs $\omega_p^{z_{11}}$ of luminosity selected volume limited samples (Table B1), with the top (bottom) panels showing the results for the low (high) redshift volume samples. The same colour code as in Figure 11 is used, and the shaded black and dark orange regions denote the ACFs of luminosity selected REF and SF complete samples presented in Figure 11.

On small scales, all $ACF_{\omega_p^{z_{11}}}$ of volume limited luminosity selected samples are in quantitative agreement with the respective luminosity selected functions. Compared to the $ACFs_{\omega_p^{z_{11}}}$ of luminosity selected samples, however, the uncertainties associated with the clustering amplitudes of volume limited functions are relatively large, driven by the small number statistics of the volume samples. Given both the agreement between volume limited and non-volume limited ACF results, as well as the importance of sample statistics for studies, such as ours, that aim to investigate small scale clustering properties of star formers, we base the conclusions of this study on the analyses performed using luminosity and stellar mass selected samples.

On large separations, however, the respective ACFs $\omega_p^{z_{11}}$ of volume limited and non-volume limited luminosity samples differ from each other. These disagreements can largely be attributed to the discrepancies between the redshift coverages of the respective volume limited and non-volume limited sample. The redshift coverage of M_f galaxies, for example, is similar to that of M_{f,v1}, and consequently leads to a good agreement between the ACFs based on M_f and M_{f,v1} samples. The M_{*} sample, on the other hand, encompasses both M_{*,v1} and M_{*,v2} galaxies. Therefore the ACF of M_{*} galaxies can be thought of as the average of the ACFs of its respective volume samples. The same trends evident in the ACFs $\omega_p^{z_{11}}$ of volume limited luminosity selected samples are also evident in the ACFs $\omega_p^{z_{11}}$ of volume limited stellar mass selected samples presented in Appendix B1.

Table B1. The volume limited sample definitions corresponding to the three independent magnitude and stellar mass limited samples described in Tables 1
and 2. We note that the SF complete volume limited samples are at least 95% volume limited.

subset	subset definition	Ngalaxies	z coverage					
SF complete								
At least 95% complete with respect to both r-band magnitude and H $lpha$ flux selection								
M_{b,v_1}	$-23.5 \leq M_r < -21.5$, SFR/M _{\odot} yr ⁻¹ ≥ 0.28	1 491	0.01 - 0.15					
M_{*,v_1}	$-21.5 \leq M_r < -20.5$, SFR/M _o yr ⁻¹ ≥ 0.25	4 1 8 8	0.01 - 0.15					
M_{f,v_1}	$-20.5 \leq M_r < -19.5$, SFR/M _{\odot} yr ⁻¹ ≥ 0.33	5 2 9 8	0.01 - 0.14					
		1 5 1 4	0.17 0.00					
M_{b,v_2}	$-23.5 \leq M_r < -21.5$, SFR/M _o yr ⁻¹ ≥ 0.90	1514	0.17 - 0.23					
M_{*,v_2}	$-21.5 \leq M_r < -20.5$, SFR/M _{\odot} yr ⁻¹ ≥ 1.00	4914	0.17 - 0.23					
$\mathcal{M}_{\mathcal{H},v}$	$10.5 \leq \log M/M_{\odot} < 11.0$, SFR/M _{\odot} yr ⁻¹ ≥ 0.29	1 991	0.01 - 0.15					
$\mathcal{M}_{I,v}$	$10.0 \leq \log M/M_{\odot} < 10.5$, SFR/M _{\odot} yr ⁻¹ ≥ 0.35	4 1 6 3	0.01 - 0.15					
$\mathcal{M}_{\mathcal{L},v}$	$9.5 \leqslant \log {\cal M}/M_\odot {<} 10.0, SFR/M_\odot yr^{-1} \geqslant 0.38 \label{eq:second}$	2 906	0.01 - 0.126					
	REF							
1	100% complete with respect to r -band magnitude set	election of the	survey					
M_{b,v_1}	$-23.5 \leq M_r < -21.5$	4 064	0.01 - 0.15					
M_{*,v_1}	$-21.5 \leq M_r < -20.5$	10244	0.01 - 0.15					
M_{f,v_1}	$-20.5 \leq M_r < -19.5$	12751	0.01 - 0.14					
M_{b,v_2}	$-23.5 \leq M_r < -21.5$	6971	0.17 - 0.23					
M _{*, v2}	$-21.5 \leq M_r < -20.5$	17737	0.17 - 0.23					
$\mathcal{M}_{\mathcal{H},v}$	$10.5 \leq \log M/M_{\odot} < 11.0$	7 338	0.01 - 0.15					
$\mathcal{M}_{I,v}$	$10.0 \leq \log M/M_{\odot} < 10.5$	11812	0.01 - 0.15					
$\mathcal{M}_{\mathcal{L},v}$	$9.5 \leqslant \log \mathcal{M}/M_{\odot} < 10.0$	8014	0.01 - 0.126					

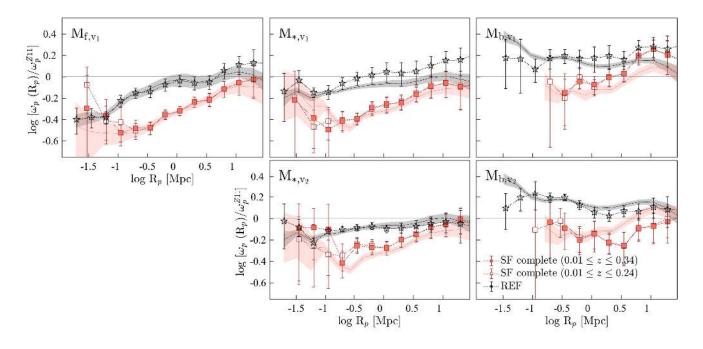


Figure B1. The GAMA projected ACFs of luminosity selected volume limited samples (symbols) compared to the projected ACFs of luminosity selected samples (shaded regions, corresponding to the ACFs shown in Figure 11), all relative to ω_p^{Z11} . The black symbols denote the REF ACFs of luminosity selected volume limited samples (Table B1), and the solid and open orange symbols denote the SF complete ACFs of luminosity selected volume limited samples, where N_{weight} (§ 3.1) is computed based on $0.001 < z \le 0.34$ and $0.001 < z \le 0.24$ galaxy samples, respectively. The black and dark orange shaded regions denote the REF and SF complete ACFs of luminosity selected samples (Table 1) presented in Figure 11. The ACFs of volume limited samples are in qualitative, and in most cases quantitative, agreement with the respective ACFs of luminosity selected samples.

B2 Cross correlation functions of volume limited star forming samples

The cross correlation results of the volume limited magnitude selected samples are presented in Figures B2-B4. The left panels of Figure B2 show the projected CCFs relative to ω_p^{Z11} of high, low and intermediate sSFR galaxies (blue, red and black symbols, respectively), and the right panels show the CCFs of optically blue, red and intermediate colour (blue, red and black symbols, respectively) galaxies. Also shown as shaded regions are the ACFs of respective magnitude selected samples (non-volume limited) with blue and red on left (right) panels denoting high (optically blue) and low (optically red) sSFR galaxies in each magnitude sample, respectively.

The $\text{CCFs}_{\omega_p^{Z_{11}}}$ of volume limited M_f galaxies of SF complete sample show a strong agreement with their non-volume limited counterparts on all scales. Recall that there is a large overlap in redshift between volume limited and non-volume limited magnitude selected M_f galaxies. Therefore the respective $\text{CCFs}_{\omega_p^{Z_{11}}}$ likely mostly probe the clustering properties of similar galaxy populations (as is the case with the respective $\text{ACF}_{\omega_p^{Z_{11}}}$ of M_f galaxies of SF complete sample). In comparison, the $\text{CCFs}_{\omega_p^{Z_{11}}}$ of M_* and M_b galaxies of volume limited luminosity selected samples noticeably differ from their respective non-volume limited counterparts.

On $R_p \leq 0.3$ Mpc, the $CCF_{\omega_p^{Z11}}$ of optically blue M_* galaxies with $0.01 \leq z \leq 0.15$ (i.e. low redshift volume limited M_* sample; Figure B2) show a steady decline in clustering with decreasing R_p . The CCF of optically blue M_* galaxies with $0.17 \leq z \leq 0.24$ (i.e. high redshift volume limited M_* sample; Figure B3), on the other hand, show a steady incline in clustering amplitude with decreasing R_p , in agreement with that seen in the $CCFs_{\omega_p^{Z11}}$ of M_* galaxies over the $0.01 \leq z \leq 0.24$ range. On larger separations, the respective $CCFs_{\omega_p^{Z11}}$ optically blue galaxies are in agreement with each other to within their uncertainties. Finally, the respective volume limited and non-volume limited $CCFs_{\omega_p^{Z11}}$ of optically red and low sSFR M_* galaxies show similar clustering behaviour to within their uncertainties.

On $R_p \lesssim 0.3$ Mpc, the $M_b CCFs_{\omega_p^{Z+1}}$ of volume limited magnitude selected low redshift galaxies with high (low) sSFRs and optically blue (red) colours show on average higher clustering strengths than (similar clustering strength to) their respective magnitude limited $CCFs_{\omega_p^{Z+1}}$. The clustering of the respective high redshift volume limited samples, on the other hand, show the opposite. These differences could be driven by the differences in the redshift distributions of low and high redshift volume samples.

Figure B4 shows the CCFs_{ω_p^{Z11}} of spectroscopically blue and red (left panels) galaxies, and of galaxies with low and high dust obscuration measures (right panels) of volume limited (blue and red symbols, respectively) and non-volume limited (blue and red shaded regions, respectively). Overall, there is a good agreement between the respective CCFs_{ω_p^{Z11}}. Interestingly, the clustering signal on $R_p \leq 0.3$ Mpc of spectroscopically blue M_{*} galaxies does not appear to mirror the decline in amplitude with decreasing R_p evident in the respective CCF_{ω_p^{Z11}} of optically blue M_{*} galaxies, despite being a proxy for optical colour. This is perhaps a result of D₄₀₀₀ index being more sensitive to the colour changes in the cen-

tral regions of galaxies than optical colours (see § 4.3 for detailed discussion).

Finally, the CCFs_{ω_p^{Z11}} of high, low and intermediate SFR galaxies, and of galaxies with optically blue, red and intermediate colours of low redshift volume limited stellar mass selected samples (not shown) indicate a clustering behaviour similar to that observed in Figure B2.

APPENDIX C: CROSS CORRELATING STAR FORMING AND REF SAMPLES

The clustering results presented in the main paper are computed from cross correlating different star forming samples. We have also investigated the clustering of different galaxy populations by cross correlating star forming and reference (REF) galaxy samples. This approach significantly reduces the uncertainties arising from small number statistics as REF samples contain a higher number of galaxies than SF complete samples, however, with the caveat that by mixing star forming and non-star forming populations it becomes difficult to interpret and understand the clustering properties of star forming galaxies. For this reason, we present and discuss the results of this analysis here. Also, in order to differentiate the results of the cross correlation between different SF complete samples of galaxies (i.e. the CCFs presented in § 4.3) from the results of the cross correlations between SF complete and REF samples, we use the labels SF CCFs_w²¹¹ and REF CCFs_w²¹¹, respectively.

Figures B5 and B6 present the REF (symbols) and SF (coloured regions) CCFs $\omega_{z^{11}}$ of volume limited magnitude selected samples. The red and blue colours in Figure B5 denote low and high sSFR galaxies, and in Figure B6, optically red and blue galaxies. The REF CCFs $_{\omega_p^{Z^{11}}}$ are in qualitative agreement with the respective SF $CCFs_{\omega_n^{Z_{11}}}$ on all scales. The $CCFs_{\omega_n^{Z_{11}}}$ of low and high redshift volume limited samples of M_* and M_b galaxies show the evolution of the clustering of optically bright star forming galaxies. These evolutionary effects are present to varying degrees over all R_p probed, however, are particularly notable over the $R_p > 0.5$ Mpc. For instance, the CCFs $_{\omega_n^{Z_{11}}}$ of M_{*} and M_b galaxies show higher clustering amplitudes at low redshift than at high redshift. This result is in agreement with previous studies, in particular with those that investigate the dependence of galaxy clustering on optical luminosity (e.g. Adelberger et al. 2005; Marulli et al. 2013), that report an increase in clustering strength with decreasing redshift.

To summarise, our results based on cross correlating star forming galaxies with all galaxies regardless of star formation (i.e. REF) show that, on most scales, the redder, low-sSFR galaxies at a fixed M_r are clustered more strongly than bluer, high-sSFR systems, in agreement with previous photometric studies of clustering in the local Universe (e.g. Norberg *et al.* 2001; Marulli *et al.* 2013; Guo *et al.* 2014). The environmental effects are likely to be largely responsible for the differences observed in the clustering strengths between the REF CCFs $_{\omega_{P}^{Z_{11}}}$ of low-sSFR and high-sSFR galaxies. Both the lower clustering strengths exhibited by high-sSFR galaxies and the observational evidence of higher SFR enhancements in galaxy pairs of equal mass (Ellison *et al.* 2008) suggest that, at a fixed M_r , it is more likely for an interacting high-sSFR galaxy to have a SF companion than a non-SF companion, and reside in relatively low density environment. The higher clustering of

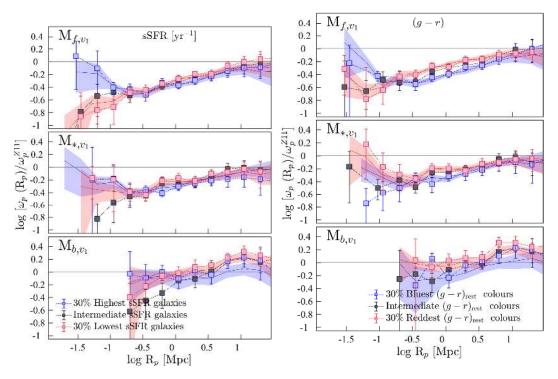


Figure B2. The projected CCFs of volume limited luminosity selected SF complete samples (low redshift samples described in Table B1) relative to ω_p^{211} (luminosity increases down). Right panels: the CCFs of galaxies with optically blue (the 30% bluest in $(g - r)_{rest}$; blue squares), red (the 30% reddest in $(g - r)_{rest}$; red squares) and intermediate (the 40% with intermediate $(g - r)_{rest}$ measures; black squares) colours. Left panels: the CCFs of high (the 30% highest in sSFR; blue squares), low (the 30% lowest in sSFR; red squares) and intermediate (the 40% with intermediate sSFR; black squares) sSFR galaxies. The blue and red shaded regions show the respective CCFs of optically blue and red (right panels), and high and low sSFR (left panels) galaxies of magnitude selected SF complete samples described in Table 1 of the main paper.

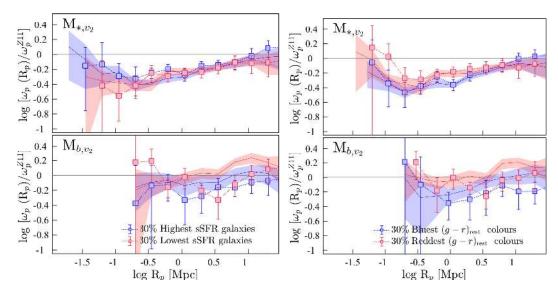


Figure B3. The projected CCFs of volume limited luminosity selected SF complete samples (high redshift samples described in Table B1) relative to ω_p^{211} (luminosity increases down). Right panels: the CCFs of galaxies with optically blue (the 30% bluest in $(g - r)_{rest}$; blue squares) and red (the 30% reddest in $(g - r)_{rest}$; red squares) colours. Left panels: the CCFs of high (the 30% highest in sSFR; blue squares) and low (the 30% lowest in sSFR; red squares) sSFR galaxies. The blue and red shaded regions show the respective CCFs of optically blue and red (right panels), and high and low sSFR (left panels) galaxies of magnitude selected samples described in Table 1 of the main paper.

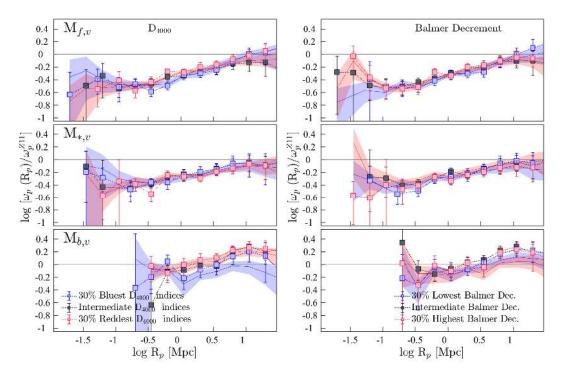


Figure B4. The projected CCFs of volume limited luminosity selected SF complete samples (low redshift samples described in Table B1) relative to ω_p^{Z11} (luminosity increases down). Right panels: the CCFs of galaxies with low (the 30% lowest in Balmer decrement; blue squares), high (the 30% highest in Balmer decrement; red squares) and intermediate (the 40% with intermediate Balmer decrements; black squares) dust obscurations. Left panels: the CCFs of galaxies with spectroscopically blue (the 30% lowest in D₄₀₀₀ indices; blue squares), red (the 30% highest in D₄₀₀₀ indices; red squares) and intermediate (the 40% with intermediate shaded regions show the respective CCFs of magnitude selected samples; galaxies with spectroscopically blue and red colours (right panels), and high and low Balmer decrements (left panels).

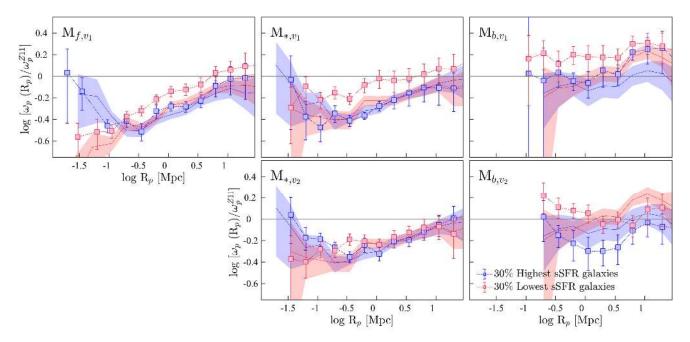


Figure B5. The projected CCFs relative to ω_p^{Z11} computed from cross correlating luminosity selected volume limited SF complete samples (i.e. high and low sSFR galaxies drawn from volume limited SF complete samples) with the respective volume limited REF samples (optical luminosity increases across). Top panels: the CCFs of low redshift high (the 30% highest in sSFR; blue squares) and low (the 30% lowest in SFR; red squares) sSFR galaxies. Bottom panels: the CCFs of high redshift high (the 30% highest in sSFR; blue squares) and low (the 30% lowest in SFR; red squares) sSFR galaxies. The blue and red shaded regions show the CCFs of high and low sSFR galaxies relative to ω_p^{Z11} of magnitude selected volume limited SF complete samples discussed in § B2 and §4.3.

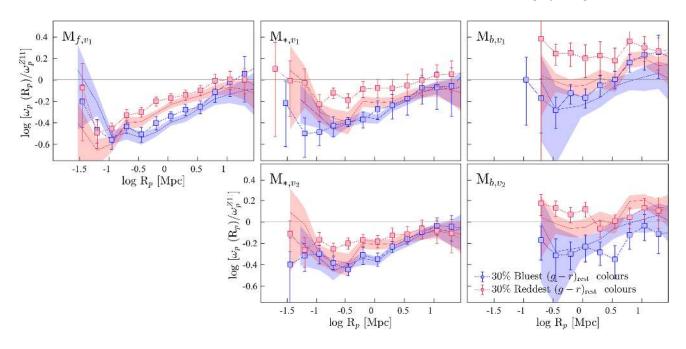


Figure B6. The projected CCFs relative to ω_p^{Z11} computed from cross correlating luminosity selected volume limited SF complete samples (i.e. optically blue and red galaxies drawn from volume limited SF complete samples) with the respective volume limited REF samples (optical luminosity increases across). Top panels: the CCFs of low redshift galaxies with optically blue (the 30% bluest in $(g - r)_{rest}$; blue squares) and red (the 30% reddest in $(g - r)_{rest}$; blue squares) and red (the 30% reddest in $(g - r)_{rest}$; blue squares) and red (the 30% reddest in $(g - r)_{rest}$; blue squares) and red (the 30% reddest in $(g - r)_{rest}$; blue squares) and red (the 30% reddest in $(g - r)_{rest}$; blue squares) and red (the 30% reddest in $(g - r)_{rest}$; blue squares) colours. Bottom panels: the CCFs of high redshift galaxies with optically blue (the 30% bluest in $(g - r)_{rest}$; blue squares) and red (the 30% reddest in $(g - r)_{rest}$; red squares) colours. The blue and red shaded regions show the CCFs of high and low sSFR galaxies relative to ω_p^{Z11} of magnitude selected volume limited SF complete samples discussed in § B2 and § 4.3.

low-sSFR and redder systems at a fixed M_r , on the other hand, suggest that an interacting companion of a low-sSFR galaxy can either be a SF or a non-SF galaxy, both preferentially inhabiting a higher-density environment, and rather than triggering SF in interacting members, their environment has triggered quenching of star formation.

APPENDIX D: THE STANDARD MARK CORRELATION FUNCTIONS OF STAR FORMING GALAXIES

The MCFs of luminosity selected and their equivalent volume limited samples, with sSFR (top panels) and $(g - r)_{rest}$ colour (bottom panels) as marks, are presented in Figure D1. The measurement uncertainties of sSFRs are relatively large compared to those of $(g - r)_{rest}$ colours, and our SF sample likely contains some overestimated sSFR measures. Therefore to limit the impact of sSFR outliers on the MCFs while not removing true starbursts from the clustering samples, we re-adjust sSFRs as follows,

$$sSFR_{adj} = \frac{sSFR \times sSFR_{max}}{sSFR + sSFR_{max}},$$
 (D1)

where $\text{sSFR}_{\text{max}} = 10^{-9.0} [\text{yr}^{-1}]$. The $(g - r)_{\text{rest}}$ and sSFR_{adj} distributions of different clustering samples are shown in the insets of Figure D1.

The $sSFR_{adj}$ MCFs (top panels of Figure D1) not only show a clear dependence of sSFR on the environment, but also show a small-scale dependence of enhancement in the spatial distribution of

sSFR (i.e. E_{sSFR}) on optical brightness. This E_{sSFR} -optical brightness dependence is in the sense that E_{sSFR} shows a strengthening in magnitude at a given R_p on $R_p \leq 0.35$ Mpc scales with increasing optical brightness, a behaviour similar to that seen in the SF ACFs of magnitude-limited samples (§ 4.2). The $(g - r)_{rest}$ MCFs (bottom panels of Figure D1) also show an enhancement in $E_{(g-r)_{rest}}$, in particular in M_f galaxies, however, the strength of $E_{(g-r)_{rest}}$ enhancement does not appear to depend on optical brightness of galaxies.

The MCFs of stellar mass selected SF complete samples are presented in Figure D2. For the same reason mentioned above, we adjust the SFRs as follows,

$$SFR_{adj} = \frac{SFR \times SFR_{max}}{SFR + SFR_{max}},$$
 (D2)

where $SFR_{max} = 50 [M_{\odot} yr^{-1}]$, and SFR_{adj} distributions of different clustering samples are shown in the insets of Figure D2.

The SFR MCFs show an enhancement in E_{SFR} , similar to that observed in sSFR MCFs of luminosity selected samples. The $(g - r)_{rest}$ colour MCFs, other than that of $\mathcal{M}_{\mathcal{L}}$ galaxies, show no enhancement in the signal over small scales.

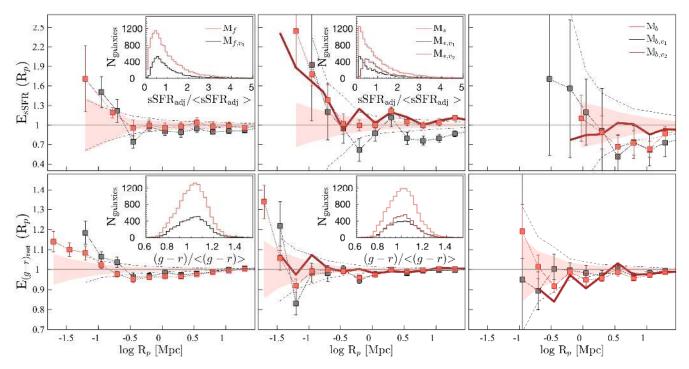


Figure D1. The MCFs of non-volume limited and volume limited luminosity selected samples (with increasing luminosity from left-to-right), with sSFR_{adj} (top panels) and $(g - r)_{rest}$ colour (bottom panels) as marks. From left-to-right: The filled symbols show the MCFs of SF M_f, M_{*} and M_b galaxies, and their respective low-*z* volume samples. We show the MCFs of higher-*z* volume samples as thick solid lines for clarity, noting that they only exist for M_{*} and M_b samples (centre and right panels). The shaded regions and the regions enclosed by black dashed lines denote the 1 σ scatter from scrambling the marks of luminosity selected and corresponding volume limited samples. For clarity, we do not show the scatter on the MCFs of higher-*z* M_{*} and M_b volume samples. The insets on left and middle panels show the distribution of the marks. Again, for clarity, we do not show the distributions of M_b galaxies.

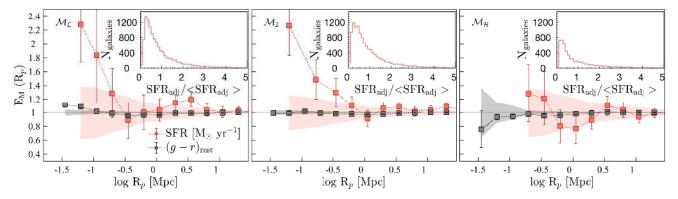


Figure D2. The MCFs of stellar mass selected SF complete samples (with increasing stellar mass from left-to-right), with SFR_{adj} and $(g - r)_{rest}$ colour as marks (orange and grey symbols, respectively). The shaded regions denote the 1σ scatter from scrambling the marks. The insets show the distributions of the SFR_{adj} marks. The distributions of colour marks are similar to that shown in Figure D1.