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# Generation of Broadband HighPurity Dual-Mode OAM Beams Using A Four-Feed Patch Antenna: Theory and Implementation 

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#### Abstract

A broadband four-feed circular patch antenna is demonstrated to generate high-purity dual-mode vortex beams carrying orbital angular momentum (OAM) for the first time. The electric field formula of the circularly polarized (CP) component is derived for analyzing radiation characteristics of different OAM modes, such as the CP property, phase and amplitude distributions. Theoretical calculation and numerical simulation results show that the four-feed patch antenna operating in the $\mathrm{TM}_{n 1}$ mode can generate $\pm n$ OAM waves. A prototype of a four-feed circular $\mathrm{TM}_{21}$-mode patch antenna is fabricated and measured to verify the effectiveness of the theoretical analysis. The proposed antenna consists of a four-feed patch with an air-coupled parasitic patch and a four-way phase-shifting feeding network. Experimental results show that the proposed antenna can generate a left-hand CP OAM with mode $l=+2$ and a right-hand CP OAM with mode $l=-2$. The OAM beams are generated in a wide band from 5.50 to 6.50 GHz with a mode purity over $90 \%$. The proposed method is very suitable for the generation of wideband, dual-mode and high-purity OAM beams in the microwave and millimeter wave bands.


The vortex beam carrying orbital angular momentum (OAM) has received much attention recently for its potential in the field of optical communications ${ }^{1}$, high-capacity communications ${ }^{2}$, the radar field of detection and imaging ${ }^{3}$, etc. The OAM vortex waves feature a helical phase front of $e^{-j l \phi}$, where $l$ is the OAM mode order and $\phi$ is the azimuthal angle. The OAM modes are orthogonal to each other, which enables multiple signals to be transmitted simultaneously at the same frequency. Therefore, OAM can be a useful degree of freedom for increasing the channel capacity and the spectral efficiency ${ }^{4}$. Although the OAM-based radio systems are considered equivalent to conventional multiple-input-multiple-output (MIMO) systems ${ }^{5}$, an experiment demonstrates that OAM-mode division multiplexing systems have low receiver complexity and high capacity compared with conventional line-of-sight MIMO systems ${ }^{6}$. Furthermore, an OAM-based synthetic aperture radar imaging method is experimentally verified, whose azimuth resolution is higher than that of a conventional one ${ }^{7}$. Thus, the potential of OAM in various fields remains to be discussed.

Different methods for generating OAM beams have been reported untill now, yet the generation of broadband and high-purity OAM beams is still considered as a big challenge. Spiral phase plates (SPPs) ${ }^{8,9}$ are widely used in the field of optics due to their simple structures. The SPPs can only generate a single mode of OAM at a time. Also, in the low frequency radio domain, SPPs are of bulky structure which increases the cost. The metallic traveling-wave ring resonant cavity antennas ${ }^{10,11}$ are proposed to generate linearly polarized OAM waves carrying $\pm 3$ modes with the purity of $87 \%$ and the attenuation coefficient of $0.036^{12}$. The sizes of these models are large and the bandwidth is quite narrow. Several small-size microstrip antennas have been reported as well, such as the circular patch antenna ${ }^{13}$, the concentric patches antenna ${ }^{14}$, and the half-mode substrate integrated waveguide (HMSIW) antenna ${ }^{15}$. These microstrip antennas are limited to a narrow impedance bandwidth of about $2 \%$ and

[^0]

Figure 1. Geometries, current and electric field distributions of the circular patch antennas excited by different feeds. ( $\mathbf{a}, \mathbf{c}, \mathbf{f}$ ) Two feeds. ( $\mathbf{b}, \mathbf{d}, \mathbf{g}$ ) Four feeds. (e) Simulated normalized radiation patterns. The black arrow represents the direction of the current of the $\mathrm{TM}_{21}$ mode. The pink arrow signifies the direction of the clutter current.
a low OAM purity of about $72 \%^{16}$. In addition, the array antennas ${ }^{17-21}$ are proposed to generate OAM beams by controlling a phase shift of the successive element. But the number of the array elements limits the order of the OAM mode and more elements are needed to ensure that the main lobe is in the pure mode region ${ }^{17}$. The OAM purity of $80 \%$ with the mode +1 can be achieved at the maximum gain angle by the uniform circular array with 16 elements. Consequently, formerly reported $2 \times 2$ array antennas ${ }^{18-21}$ can only generate $\pm 1$ OAM modes with the main lobe not being in the pure mode region. To increase the channel capacity, a dual linearly polarized (LP) microstrip antenna array is proposed with an impedance bandwidth of $3.6 \%{ }^{18}$. A high-profile dual LP bow-tie dipole array antenna achieves a wide bandwidth of $25 \%{ }^{19}$. Some dual circularly polarized (CP) patch array antennas ${ }^{20,21}$ are proposed to mitigate the multipath interference, but their feeding networks are quite complex. Several metasurfaces without complex feeding networks are proposed to generate dual-mode OAM beams ${ }^{22,23}$. However, most of these metasurfaces are excited by additional feed horn which makes them difficult to integrate and increase the profile of the antenna. Recently, an innovative OAM generation method based on transformation optic using all-dielectric lens is presented ${ }^{24}$. The lens can generate an OAM mode of +1 from 8 GHz to 16 GHz . In this method, the physical realization of the device is a challenge and radiation patterns of the proposed lens have a slight divergence effect. A phase-modulation converging lens ${ }^{25}$ based on the concept of the optical converging axicon excited by a patch array source can achieve the OAM beam converging from 9.3 GHz to 10 GHz , which is advantageous for improving the propagation performance of the OAM beam.

In this paper, we propose a new strategy to generate broadband high-purity dual-mode OAM beams using the four-feed circular patch antenna. The theoretical formula of the radiation field shows that a four-feed patch antenna at $\mathrm{TM}_{n 1}$ mode can produce $\pm n$ OAM modes. A prototype of the four-feed circular $\mathrm{TM}_{21}$-mode patch antenna is fabricated and measured to verify the effectiveness of the theoretical analysis. The simulated and measured results are well matched with the theoretical ones. The measured $3-\mathrm{dB}$ axial ratio (AR) bandwidth is about $28.33 \%$ from 5.00 to 6.70 GHz for both ports. The OAM beams with $\pm 2$ modes are generated in a wide band from 5.50 to 6.50 GHz with high mode purity.

## Results

Theoretical analysis of a four-feed circular patch antenna with CP OAM modes. The geometries of the circular patch antennas excited by two feeds and four feeds are shown in Fig. 1(a,b). Circular antennas are fed with equal amplitude, a relative $90^{\circ}$ phase shift, and the proper angular spacing to obtain CP radiation. In the case of $\mathrm{TM}_{21}$-mode patch antenna, the feed angular spacing $\alpha$ is set to $135^{\circ}$. We employ two additional feeds 3,4 located diametrically across from the feeds 1,2 . The four-feed antenna is fed with equal amplitude and a phase shift of $\left[0^{\circ}, 90^{\circ}, 0^{\circ}, 90^{\circ}\right]$ for the $\mathrm{TM}_{21}$ mode. As shown in Fig. 1(d), a more symmetric current distribution is obtained by using four-feed structure. This is because most of the cluttered current due to the mutual coupling between the feeds is suppressed by the two opposing feeds. Moreover, the effect of the remaining cluttered current is significantly eliminated by the opposite current. Thus, the four-feed structure achieves a lower cross-polarization level compared to the two-feed structure, as shown in Fig. 1(e). Additionally, Fig. 1(f,g) compares the electric field distributions for the two-feed antenna and the four-feed antenna. The four-feed structure has a more uniform and symmetrical distribution of the $\mathrm{TM}_{21}$ mode because of the low cross polarization. This illustrates that

| $\mathbf{T M}_{\boldsymbol{n} 1}$ | $\mathbf{T M}_{\mathbf{1 1}}$ | $\mathbf{T M}_{\mathbf{2 1}}$ | $\mathbf{T M}_{\mathbf{3 1}}$ | $\mathbf{T M}_{\mathbf{4 1}}$ | $\mathbf{T M}_{\mathbf{5 1}}$ | $\mathbf{T M}_{\mathbf{6 1}}$ | $\mathbf{T M}_{\mathbf{7 1}}$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $X_{n 1}$ | 1.8412 | 3.0542 | 4.2012 | 5.3176 | 6.4156 | 7.5013 | 8.5578 |

Table 1. Some Values of $X_{n 1}$.
the radiation of the working mode ( $\mathrm{TM}_{21}$ mode) is enhanced and the fields of the unwanted adjacent modes ( $\mathrm{TM}_{11}$ and $\mathrm{TM}_{31}$ modes) are suppressed. As a result, the CP bandwidth is widened and the mode purity is improved.

Next, a theoretical analysis of a four-feed circular microstrip patch antenna generating CP OAM beams is presented as follows. Based on the cavity-model approximation, different $\mathrm{TM}_{n m}$ modes can be excited by different circular patch sizes. The symbols $n$ and $m$ are the number of the angle mode and the radial mode, respectively. The electric field components of the circular patch excited by a single feed in the spherical coordinate are given as follows ${ }^{26}$

$$
\begin{align*}
& E_{\theta n}(\phi, \theta)=j^{n} \frac{e^{-j k_{0} r}}{2 r} k_{0} a_{e} h E_{0} J_{n}\left(a_{e} k_{0} \sqrt{\varepsilon_{r}}\right) \cos n \phi\left[J_{n+1}(\gamma)-J_{n-1}(\gamma)\right] \\
& E_{\phi n}(\phi, \theta)=j^{n} \frac{e^{-j k_{0} r}}{2 r} k_{0} a_{e} h E_{0} J_{n}\left(a_{e} k_{0} \sqrt{\varepsilon_{r}}\right) \sin n \phi\left[J_{n+1}(\gamma)+J_{n-1}(\gamma)\right] \cos \theta \tag{1}
\end{align*}
$$

where $\gamma=k_{0} a_{e} \sin \theta, k_{0}$ is the phase constant in a vacuum, $a_{e}$ is the equivalent circular patch radius, $h$ is the substrate thickness, $E_{0}$ is the electric field amplitude at the edge of the patch, $J_{n}$ is the Bessel function of the first kind for order integer $n$, and $\varepsilon_{r}$ is the relative permittivity. Moreover, the equivalent circular patch radius can be derived by $a_{e}=\lambda_{0} X_{n m} /\left(2 \pi \sqrt{\varepsilon_{r}}\right)^{27}$, where $\lambda_{0}$ denotes the free-space wavelength at the center operating frequency ( $f_{0}=6 \mathrm{GHz}$ ), and $X_{n m}$ is the $m$ th zero of the derivative of $J_{n}$, which is a constant and independent of frequency. Some values of $X_{n m}$ are shown in Table 1. Because of the fringe fields, the physical patch radius $a_{p}$ is slightly smaller than the equivalent one.

For the case of four feeds, the excitation amplitudes are same and mutual coupling is neglected. The feed angular spacing is given by $\alpha=(2 k+1) \pi / 2 n, k=0,1,2, \ldots, n-1^{15}$. Additionally, the four-feed antenna has a phase shift of $\left[0^{\circ}, 90^{\circ}, 0^{\circ}, 90^{\circ}\right]$ for the even-order modes and $\left[0^{\circ}, 90^{\circ}, 180^{\circ}, 270^{\circ}\right]$ for the odd-order modes ${ }^{27}$. For simplicity, only the radiated electric field of the right-hand circular polarization (RHCP) is given. The radiated electric field of the left-hand circular polarization (LHCP) can also be obtained in a similar way. The total electric field components generated by the four orthogonal modes are as follows:

$$
\begin{align*}
E_{\theta n}^{t}= & E_{\theta n}^{1}(\phi, \theta)+j E_{\theta n}^{2}(\phi+\alpha, \theta)+(-1)^{n}\left[E_{\theta n}^{3}\left(\phi+180^{\circ}, \theta\right)\right. \\
& \left.+j E_{\theta n}^{4}\left(\phi+\alpha+180^{\circ}, \theta\right)\right] \\
E_{\phi n}^{t}= & E_{\phi n}^{1}(\phi, \theta)+j E_{\phi n}^{2}(\phi+\alpha, \theta)+(-1)^{n}\left[E_{\phi n}^{3}\left(\phi+180^{\circ}, \theta\right)\right. \\
& \left.+j E_{\phi n}^{4}\left(\phi+\alpha+180^{\circ}, \theta\right)\right] \tag{2}
\end{align*}
$$

where superscripts 1, 2, 3, and 4 correspond to individual electric fields excited by the four feeds. By combining Eqs (1) and (2), the total electric field components are written as

$$
\begin{align*}
E_{\theta n}^{t} & =j^{n} \frac{C}{r} e^{-j k_{0} r} e^{-j n \phi}\left[J_{n+1}(\gamma)-J_{n-1}(\gamma)\right] \\
E_{\phi n}^{t} & =j^{n+1} \frac{C}{r} e^{-j k_{0} r} e^{-j n \phi}\left[J_{n+1}(\gamma)+J_{n-1}(\gamma)\right] \cos \theta \tag{3}
\end{align*}
$$

where $C$ is the amplitude constant. Since the vortex EM waves dynamically rotate around the $z$-axis, the electric field components should be obtained in the cylindrical coordinate system.

The $\rho$-, $\phi$-, and $z$-components of the total electric field in the cylindrical coordinate system are

$$
\begin{align*}
E_{\rho n}^{t} & =\cos \theta E_{\theta n}^{t}=-j^{n} \frac{C}{r} e^{-j k_{0} r} e^{-j n \phi}\left[J_{n-1}(\gamma)-J_{n+1}(\gamma)\right] \cos \theta \\
E_{\phi n}^{t} & =E_{\phi n}^{t}=j^{n+1} \frac{C}{r} e^{-j k_{0} r} e^{-j n \phi}\left[J_{n-1}(\gamma)+J_{n+1}(\gamma)\right] \cos \theta \\
E_{z n}^{t} & =-\sin \theta E_{\theta n}^{t}=j^{n} \frac{C}{r} e^{-j k_{0} r} e^{-j n \phi}\left[J_{n-1}(\gamma)-J_{n+1}(\gamma)\right] \sin \theta . \tag{4}
\end{align*}
$$

The radial electric field can be divided into two orthogonal LP components or CP components, which can be expressed as

$$
\begin{equation*}
\overrightarrow{E_{l n}^{t}}=\overrightarrow{e_{\rho}} E_{\rho n}^{t}+\overrightarrow{e_{\phi}} E_{\phi n}^{t}=\overrightarrow{e_{L}} E_{L n}^{t}+\overrightarrow{e_{R}} E_{R n}^{t} \tag{5}
\end{equation*}
$$

where $E_{L n}^{t}=\left(E_{\rho n}^{t}-j E_{\phi n}^{t}\right) / \sqrt{2}, E_{R n}^{t}=\left(E_{\rho n}^{t}+j E_{\phi n}^{t}\right) / \sqrt{2}$.
By combining Eqs (4) and (5), the LHCP and RHCP components of the radial electric field are


Figure 2. The calculated $\operatorname{AR}$ (a) for integer orders $n=1,2,3, \ldots, 7$ and (b) for $n=2$ with different permittivities.

$$
\begin{align*}
E_{L n}^{t} & =2 j^{n} \frac{C}{r} e^{-j k_{0} r} e^{-j n \phi} \cos \theta J_{n+1}(\gamma) / \sqrt{2} \\
E_{R n}^{t} & =-2 j^{n} \frac{C}{r} e^{-j k_{0} r} e^{-j n \phi} \cos \theta J_{n-1}(\gamma) / \sqrt{2} . \tag{6}
\end{align*}
$$

The calculated AR is defined in decibels as

$$
\begin{equation*}
A R=20 \lg \left|\frac{\left|E_{L n}^{t}\right|+\left|E_{R n}^{t}\right|}{\left|E_{L n}^{t}\right|-\left|E_{R n}^{t}\right|}\right|=20 \lg \left|\frac{\left|J_{n-1}(\gamma)\right|-\left|J_{n+1}(\gamma)\right|}{\left|J_{n-1}(\gamma)\right|+\left|J_{n+1}(\gamma)\right|}\right| \tag{7}
\end{equation*}
$$

where $\gamma=k_{0} a_{e} \sin \theta=X_{n m} \sin \theta / \sqrt{\varepsilon_{r}}$. The calculated AR is shown in Fig. 2. For different integer orders with $\varepsilon_{r}=3.66$, the AR value is less than 0.6 indicating that the CP radiation is obtained. For the case of the four-feed antenna working in the $\mathrm{TM}_{21}$ mode, the AR value decreases with the increase in permittivity of the substrate as shown in Fig. 2(b). Because $\left|E_{L n}^{t}\right|<\left|E_{R n}^{t}\right|$, the RHCP component is the co-polarized and the LHCP component is the cross-polarized. Equation (6) shows that the RHCP component of the electric field has the helical phase wavefront of $e^{-j n \phi}$, which indicates the generation of the RHCP vortex waves carrying an OAM mode of $-n$. Similarly, the phase wavefront of $e^{j n \phi}$ can be obtained for LHCP, indicating the four-feed antenna can obtain an OAM mode with a positive number.

To verify the formula of the electric field, the phase and amplitude distributions under ideal conditions operating in the 6 GHz are analysed using MATLAB software, as shown in Fig. 3(a). The square observation window lies at $z=2 \lambda_{0}$ with a side length of $12 \lambda_{0}$. The phase change along a concentric circle is $2 \pi l$ for $l=2,3$, and 4 . An amplitude null is observed in the center of the beam, similar to that of the optical vortex beam carrying OAM. Based on the model shown in Fig. 1(b), the theory is further verified by using HFSS software. Figure 3(b,c) shows the simulated results for different OAM modes for different sizes of the metal ground at 6 GHz . The calculated and simulated results for the metal ground of the $10 \lambda_{0} \times 10 \lambda_{0}$ square are in good agreement. However, Figure 3(c) reveals a difference in the center of the electric field for the OAM mode of +4 for the metal ground of the $1.6 \lambda_{0} \times 1.6 \lambda_{0}$ square, which is mainly due to the effect of the finite metal ground. Both the size of the patch antenna and the divergence angle of the radiation start to rise with an increase in the order of the OAM mode. Therefore, the beam peak of the higher mode antenna is at a lower angle. The finite metal ground cannot maintain the horizontal electric field to propagate well for low-angle radiation, which deteriorates the electric field distribution of the $\mathrm{TM}_{41}$ mode. In conclusion, these features with the helical phase and the donut-shaped amplitude validate that the four-feed circular patch antenna can generate $\pm n \mathrm{OAM}$ states for the $\mathrm{CP}^{\mathrm{TM}}{ }_{n 1}$ mode.

Fabricated structure of the antenna working in a CP TM ${ }_{21}$ mode. The configuration of the proposed antenna consists of a parasitic circular patch, a four-feed circular patch, a metal ground, and a feeding network, as shown in Fig. 4. The parasitic circular patch is located on the top layer of Substrate 1. The four-feed circular patch which works as a driven patch is located on the Substrate 2. There is an air gap between Substrate 1 and Substrate 2 in order to achieve the wideband property of CP radiation. The metal ground with four feeding slots is situated on the top layer of Substrate 3. A four-way feeding network with $90^{\circ}$ phase differences is printed on the bottom layer of Substrate 3. An inexpensive FR4 substrate ( $h_{1}=0.8 \mathrm{~mm}, \varepsilon_{r 1}=4.4$, loss tangent $=0.02$ ) is used as the top substrate. The middle and bottom substrates are Rogers 4350B ( $h_{2}=1.524 \mathrm{~mm}, h_{3}=0.508 \mathrm{~mm}$ $\varepsilon_{r 2}=3.66$, loss tangent $=0.004$ ). The detailed dimensions of the proposed antenna are set as follows (unit: mm ): $a_{1}=16, a_{2}=12.5, l_{1}=60, l_{2}=70, h_{a}=4, l_{c 1}=14.48, l_{c 2}=11.08, w_{1}=1.78, w_{2}=1.72, l_{t}=7.2$, and $l_{f}=9$.

Figure 4(d) shows the phase-shifting feeding network consisting of a $3-\mathrm{dB}$ three-branch coupler, two T-junction power dividers, and two quarter-wavelength impedance transformers. The 3-dB three-branch coupler is considered as a cascade connection of two conventional couplers ${ }^{28}$, introducing another resonance point, thus


Figure 3. Calculated and simulated phase and amplitude distributions of CP component of the electric field for $l=+2, l=+3$, and $l=+4$ OAM modes at 6 GHz . (a) Calculated results for the infinite metal ground. (b) Simulated results for the metal ground ( $10 \lambda_{0} \times 10 \lambda_{0}$ square). (b) Simulated results for the metal ground $\left(1.6 \lambda_{0} \times 1.6 \lambda_{0}\right.$ square $)$.


Figure 4. Configuration of the proposed four-feed patch antenna. (a) Exploded view. (b) Side view. (c) Top view without the parasitic patch. (d) Bottom view.


Figure 5. Simulated results of the four-way phase-shifting feeding network. (a) $l=-2$ mode. (b) $l=+2$ mode.


Figure 6. Fabricated antenna. (a) Top view. (b) Bottom view. (c) The near-field test scenario in the anechoic chamber.
providing a broadband $90^{\circ}$ phase shift. We designed the feeding network with six ports: Ports 1,2 are input ports, while other ports are output ports used to feed the patch. When Port 1 is excited and Port 2 is terminated with a $50-\Omega$ load, Ports $3,4,5$, and 6 have equal amplitude and a phase shift of $0^{\circ},-90^{\circ}, 0^{\circ},-90^{\circ}$ respectively. The proposed antenna can generate the OAM mode $l=-2$ for RHCP. Whereas when Port 2 is excited and Port 1 is matched, Ports $3,4,5$, and 6 have equal amplitude and a phase shift of $0^{\circ}, 90^{\circ}, 0^{\circ}, 90^{\circ}$ respectively and the OAM mode $l=+2$ is generated for LHCP. The simulated S-parameters of the broadband feeding network are shown in Fig. 5. The simulated $\left|S_{11}\right|,\left|S_{22}\right|,\left|S_{12}\right|$ and $\left|S_{21}\right|$ are less than -20 dB from 5.5 to 6.5 GHz . For the OAM mode $l=-2$, the variations of output powers and phases at Ports $3,4,5$, and 6 are less than 0.9 dB and $0.9^{\circ}$ from 5.5 to 6.5 GHz , respectively. For the OAM $l=+2$ mode, the variations of output powers and phases are less than 0.8 dB and $0.5^{\circ}$, respectively. Therefore, the four-way phase-shifting feeding network can maintain a stable phase and amplitude response over a wide frequency band, which makes it possible for the proposed antenna with broadband performance.

Simulated and measuredt results. To verify the effectiveness of the theory and the design, a prototype of the four-feed antenna is fabricated and measured, as shown in Fig. 6. The simulated and measured S-parameters of the four-feed antenna are shown in Fig. 7(a). The measured 10-dB return loss bandwidth is $45 \%$ from 4.75 to 7.45 GHz for Port 1 and $44.2 \%$ from 4.75 to 7.40 GHz for Port 2 . The measured transmission coefficient $\left|S_{12}\right|$ is lower than -10 dB from 5.50 to 6.65 GHz . The simulated and measured gains and AR values of the four-feed antenna versus frequency are depicted in Fig. 7(b). The measured 1-dB gain bandwidth of the antenna is $16.67 \%$ from 5.50 to 6.65 GHz for both ports. The measured 3-dB AR bandwidth is $28.33 \%$ from 5.00 to 6.70 GHz for both ports. The simulated and measured radiation patterns in $x o z$ plane at 6 GHz are shown in Fig. $8(\mathrm{a}, \mathrm{b})$. For the OAM with $l=-2$ mode, the angle of the maximum gain is around $\theta=38^{\circ}$. The measured peak gain is 5.81 dBic


Figure 7. Simulated and measured results of the four-feed antenna. (a) S-parameters. (b) Gains and AR values versus frequency in $x o z$ plane.


Figure 8. Simulated and measured radiation patterns and AR in $x o z$ plane at 6 GHz for $(\mathbf{a}, \mathbf{c}) l=-2$ mode, $(\mathbf{b}, \mathbf{d})$ $l=+2$ mode.
at 6 GHz . The half-power bandwidth (HPBW) ranges from $-46^{\circ}$ to $-12^{\circ}$ and from $20^{\circ}$ to $60^{\circ}$. For the OAM with $l=+2$ mode, the angle of the maximum gain is also around $\theta=38^{\circ}$ with the peak gain of 6.23 dBic at 6 GHz . The HPBW ranges from $-44^{\circ}$ to $-20^{\circ}$ and from $24^{\circ}$ to $56^{\circ}$. The measured cross-polarization levels are substantially lower than -16 dB over the HPBW. Figure $8(\mathrm{c}, \mathrm{d})$ shows that the measured AR is substantially below than 3 dB over the HPBW. Figure 9 represents the simulated and measured near-field distributions of the electric field at different frequencies, which highly coincide with each other in a wide band. The observation window lies in the plane at $z=150 \mathrm{~mm}$ with dimensions of $300 \mathrm{~mm} \times 300 \mathrm{~mm}$ with a step of 20 mm . The measured and simulated


Figure 9. Simulated and measured near-field phase and amplitude distributions of the CP component of the electric field for $l=-2$ and $l=+2$ OAM modes at $z=150 \mathrm{~mm}$ at different frequencies. (a,c) Simulated results. (b,d) Measured results.


Figure 10. OAM spectrum generated by the four-feed antenna at (a) 5.5 GHz , (b) 6 GHz , (c) 6.5 GHz .
results are in good agreement. The purity of different OAM modes is quantitatively calculated based on the Fourier transform. Since the azimuth angle $\phi$ is a periodic function, its Fourier conjugate is the OAM spectrum. The Fourier relationship between the OAM spectrum $A_{l}$ and the sampling phase $\psi(\phi)$ is given by ${ }^{29}$

$$
\begin{equation*}
A_{l}=\frac{1}{2 \pi} \int_{0}^{2 \pi} \psi(\phi) d \phi \exp (-i l \phi) \tag{8}
\end{equation*}
$$

| Ref. | Antenna structure | $f_{0}(\mathrm{GHz})$ | Size ( $\lambda_{0}^{3}$ ) | 10-dB BW | Gain (dBi) | Polarization | OAM Mode | Mode Purity | No. of Ports |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 10 | Metallic circular slot + Coupler | 10 | $\pi \times 2.06^{2} \times$ N.A | Narrow | 3.71 | LP | +3/-3 | 87\% | 2 |
| 13 | Circular patch + A coupler | 2.6 | $3.4 \times 3.4 \times 0.12$ | 2.6\% | 6.5 dBic | Dual CP | +1/-1 | >72\% | 2 |
| 14 | Concentric patches | 5.65 | $\pi \times 1.13^{2} \times 0.03$ | Narrow | N.A | CP | +1/-1, +2/-2 | 72\% | 6 |
| 15 | HMSIW | 10 | $3 \times 3 \times 0.07$ | Narrow | N.A | LP | $+2 /-2,+6 /-6$ | 72\% | 4 |
| 18 | $2 \times 2$ Patch array + T-junction power dividers | 5.5 | $2.2 \times 2 \times 0.1$ | 3.6\% | 9.5 | Dual LP | +1/-1 | <80\% | 2 |
| 19 | $2 \times 2$ Dipole array +4 Couplers | 2.4 | $1.76 \times 1.76 \times 0.51$ | 25\% | N.A | Dual LP | +1/-1 | <80\% | 4 |
| 20 | $2 \times 2$ Patch array +4 Couplers | 4.8 | $1.67 \times 1.67 \times 0.1$ | 33.5\% | 9 dBic | Dual CP | $+1 /-1$ | <80\% | 2 |
| 21 | $2 \times 2$ Patch array + P-i-n didoes feed network | 2.5 | $1.28 \times 1.28 \times 0.07$ | 21\% | 5.3 | Dual CP | +1/-1 | <80\% | 1 |
| This work | Four-feed circular patch + A coupler | 6 | $\mathbf{1 . 4} \times \mathbf{1 . 4 \times 0 . 1 4}$ | 44.2\% | 6.2 dBic | Dual CP | +2/-2 | 90\% | 2 |

Table 2. Comparison Between the Proposed and Reported OAM Antennas.

Figure 10 shows the OAM spectrum at maximum gain angle at different frequencies. We can see that the OAM purity of $l= \pm 2$ is greater than $90 \%$ from 5.5 to 6.5 GHz .

## Discussion

In this paper, a four-feed circular patch antenna is proposed for the first time to generate high-purity OAM beams. In principle, arbitrary OAM modes can be obtained by changing the patch radius together with the feed angular spacing. The fabricated prototype of the proposed $\mathrm{TM}_{21}$-mode antenna has an ability to generate the OAM mode of $l=+2$ with LHCP and the OAM mode of $l=-2$ with RHCP. Compared with previously reported OAM antennas shown in Table 2, the proposed antenna has a broader bandwidth of about $44.2 \%$ and a higher mode purity above $90 \%$. Unlike the designs with single polarization ${ }^{10,14,15}$ and dual LP property ${ }^{18,19}$, the proposed antenna has the benefit of providing dual CP property. Moreover, in comparison with the patch antenna ${ }^{13}$ and array antennas ${ }^{18-21}$, the proposed antenna has a higher order of the OAM mode. All of these attractive features illustrate that the proposed four-feed circular patch antenna offers a tremendous potential in high-capacity wireless communications and radar applications of detection and imaging.

## Methods

The theoretical calculation results are processed by MATLAB software. The proposed four-feed circular patch antenna is designed by ANSYS High Frequency Structure Simulator (HFSS), and manufactured on the Rogers 4350B and FR4 substrates using the PCB technology.

The S-parameters are measured by the Keysight N5247A vector network analyzer and the field distributions are obtained by near-field scanning with an open-ended waveguide probe. The radiation pattern and near-field scanning are measured in a microwave anechoic chamber.

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## Author Contributions

Y.H. and X.L. conceived the idea, conducted the design and measurements, and wrote the manuscript. Y.H., Q.L. and Z.Q. conducted the theoretical analyses, Y.H. and Z.Q. conducted the antenna fabrication and measurements. Y.H., H.Z. and Z.A. analysed the results. All authors participated in discussions and reviewed the manuscript.

## Additional Information

Competing Interests: The authors declare no competing interests.
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