

LETTER TO THE EDITOR

Generation of massive stellar black holes by rapid gas accretion in primordial dense clusters

Zacharias Roupas¹ and Demosthenes Kazanas²

¹ Centre for Theoretical Physics, The British University in Egypt, Sherouk City 11837, Cairo, Egypt
e-mail: zaroupas@gmail.com

² Astrophysics Science Division, NASA Goddard Space Flight Center, Greenbelt, MD 20771, USA

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ABSTRACT

Supernova theory suggests that black holes of a stellar origin cannot attain masses in the range of 50–135 solar masses in isolation. We argue here that this mass gap is filled in by black holes that grow by gas accretion in dense stellar clusters, such as protoglobular clusters. The accretion proceeds rapidly, during the first 10 megayears of the cluster life, before the remnant gas is depleted. We predict that binaries of black holes within the mass gap can be observed by LIGO.

Key words. globular clusters: general – galaxies: star clusters: general – gravitational waves – accretion, accretion disks – black hole physics

1. Introduction

Pair-instability supernova (Heger & Woosley 2002) signifies the presence of an upper mass gap at 50–135 M_{\odot} of population I/II stellar origin black holes (BHs; Belczynski et al. 2016; Farmer et al. 2019). Up to the time this work is published, no BH mass in binary black hole (BBH) gravitational wave (GW) signals detected by The LIGO Scientific Collaboration & the Virgo Collaboration (2019) exceeds 50 M_{\odot} . Nevertheless, 75% of these BHs have masses greater than 20 M_{\odot} , which is far more than expected from stellar evolution and more than is observed in X-ray binaries. Fishbach & Holz (2017) suggested that because of LIGO's higher sensitivity to BBH members $\gtrsim 20 M_{\odot}$, an excess of higher mass BHs should be expected, with a maximum at 40 M_{\odot} .

We suggest here that besides LIGO's bias toward higher BH masses, the BH mass function observed will favor masses that are higher than predicted by stellar evolution because of gas accretion by BHs in dense stellar clusters. This will occur rapidly (Roupas & Kazanas 2019) before the gas is depleted through the first stellar formation event. More importantly, we estimate that this shift is sufficient to fill the upper mass gap. This mechanism is different, but not mutually exclusive, from the repeated mergers scenario (Gerosa & Berti 2017, 2019; Rodriguez et al. 2019; Doctor et al. 2019).

It is plausible that globular clusters (GCs) started life as dense gas clouds, which are referred to as protoglobular or primordial clusters. They underwent prolonged star formation early in their lifetimes (Gratton et al. 2012). A huge gas reservoir, comparable to or higher in mass than that of its stellar component, is available for accretion immediately after the first stellar formation event. Feedback processes from stellar evolution in star-forming regions are believed to clear away the surrounding gas (e.g., Voss et al. 2010; Galván-Madrid et al. 2013; Krumholz et al. 2014). The effectiveness of the process depends on the compactness (Krause et al. 2012, 2016;

Silich & Tenorio-Tagle 2017, 2018), that is, the mass over half-mass radius of the cluster,

$$C = \frac{M_{\text{total}}}{10^5 M_{\odot}} \left(\frac{r_{\text{hm}}}{\text{pc}} \right)^{-1}. \quad (1)$$

For sufficiently compact clouds, feedback processes are expected to become ineffective because they are proportional to the total mass. The gravitational binding energy is proportional to the square of it.

Proposals regarding the precise mechanism for gas depletion in primordial GCs abound in the literature, but a general consensus has not been achieved (Spergel 1991; Thoul et al. 2000; Fender et al. 2005; Moore & Bildsten 2011; Herwig et al. 2012; Krumholz et al. 2014; Krause et al. 2016, 2013, 2012; Silich & Tenorio-Tagle 2017, 2018; Marks et al. 2008; Kruijssen 2012; D'Ercole et al. 2008; Conroy 2012; Renzini et al. 2015; Bagetakos et al. 2011; Jaskot et al. 2011; Fierlinger et al. 2016; Yadav et al. 2017). Leigh et al. (2013) proposed that in any cluster that is able to form massive stars, the primordial gas is depleted exactly due to the accretion onto BHs. They find that accreting BHs can deplete the whole gas reservoir within as short a time as 10 Myr.

Here, we do not investigate this specific scenario and do not focus on the gas reservoir, but on the effect of accretion onto the BHs of the cluster. Calura et al. (2015) estimated that within ~ 14 Myr the gas is depleted by 99% by star formation feedback processes in a primordial cluster with an initial total mass of about $10^7 M_{\odot}$. Typical globular clusters (GCs) should form about 10^2 – 10^3 BHs within about 3 Myr (Morscher et al. 2013, 2015) and could retain most of them initially if their natal kicks are sufficiently low (see Wong et al. 2012 and references therein). Additional observational evidence suggests that a BH subcluster may be retained even to present-day GCs (Sedda et al. 2019; Abbate et al. 2019).

We assume in our analysis an initial population of 500 BHs, segregated in the core of the cluster (Spitzer 1987; Leigh et al. 2014), and that the primordial gas is depleted by 99% within time $t_f = 10$ Myr, following an exponential law,

$$\rho_{\text{gas}} = \rho_{\text{gas}}(0)e^{-t/\tau}, \quad \tau = \frac{t_f}{2 \ln(10)}. \quad (2)$$

Therefore our analysis applies to any type of depletion mechanism as long as it does not proceed faster than an exponential law and the loss is approximately uniform in the core. To the exponential gas loss, we add the loss by accretion onto the BHs.

In the next section we briefly describe our model, and in Sect. 3 we present the results of our analysis. In the final section we discuss our conclusions.

2. Model

We considered $N_{\bullet} = 500$ BHs moving inside a fixed external potential with gaseous and stellar components. This analysis focuses solely on the effect of accretion onto the BH mass function, and in addition, of accretion within a short timescale $t_f = 10$ Myr. We do not study the general effect on the cluster or BH kinematics. For simplicity and with the aim to provide only statistical estimates, we therefore assumed a spherical distribution of BHs at any t , that is, we did not follow the angular changes of their orbital planes (the BH subsystem dynamics is dominated by the gaseous and stellar components on the timescales considered). We further neglected close encounters. We discuss this further in our conclusions. We intend to include and study both of these effects in a separate more detailed work that will include the estimation of merger rates.

We assumed that the fixed external potential is generated by Plummer density profiles for stellar and gaseous components,

$$\rho_{\star}(r) = \frac{3\epsilon M_{\text{total}}}{4\pi a^3} \left(1 + \frac{r^2}{a^2}\right)^{-5/2}, \quad (3)$$

$$\rho_{\text{gas}}(r, t) = \frac{3(1 - \epsilon)M_{\text{total}}}{4\pi a^3} \left(1 + \frac{r^2}{a^2}\right)^{-5/2} e^{-t/\tau}, \quad (4)$$

where ϵ is the stellar formation efficiency, and τ , given in Eq. (2), corresponds to 99% gas depletion by the final time t_f . Following Silich & Tenorio-Tagle (2018), we assumed that the gas is dominated by turbulence (Johnson et al. 2015; Elmegreen 2017), in which case the equation of state is $P_{\text{gas}} = \rho_{\text{gas}}\sigma_{\text{gas}}^2$.

The BHs were chosen from a Salpeter initial mass function (Perna et al. 2019) with a cutoff

$$f(m_{\bullet}) = \begin{cases} Am_{\bullet}^{-2.35}, & 5 M_{\odot} \leq m_{\bullet} \leq 50 M_{\odot} \\ 0, & m_{\bullet} < 5 M_{\odot} \text{ or } m_{\bullet} > 50 M_{\odot} \end{cases}, \quad (5)$$

with $A = 1.35/(5^{-1.35} - 50^{-1.35})$. They were all assumed to be initially bound inside the core radius $r_c = 0.64a$ of the initial Plummer sphere. The initial radial distribution and velocities of the BHs were chosen randomly from a Plummer distribution with the same softening radius a of Eqs. (3) and (4). Their initial angular distribution was spherical.

Hot-type accretion, whose typical representative is Bondi-Hoyle accretion, is appropriate in the dense gaseous environment we examined, where the speed of sound is greater than or on the order of the accretor velocity (Merritt 2013). The spherically symmetric accretion rate for a given cross section πR_{acc}^2 may be written as (Frank et al. 1985)

$$\dot{m}_{\bullet} = \pi \rho_{\text{gas}} v_{\text{rel}} R_{\text{acc}}^2, \quad (6)$$

where m_{\bullet} is the mass of the BH that accretes the gas with relative velocity

$$v_{\text{rel}} = \sqrt{v^2 + c_s^2}, \quad (7)$$

and v is the BH velocity with respect to the center of mass of the cluster, and c_s is the speed of sound of the gas.

As Bonnell et al. (2001) pointed out, choosing the proper accretion radius inside a gaseous stellar cluster depends on the relative amount of gas and on the radial position of the accretor. When the gas dominates the cluster potential, the accretion rates are given by a tidal-lobe accretion radius (Paczynski 1971),

$$R_{\text{tid}}(r_i) \sim 0.5 \left(\frac{m_{\bullet}}{M(r < r_i)} \right)^{1/3} r_i, \quad (8)$$

where $M(r < r_i)$ is the total mass of the cluster within the radial position r_i of the i th BH. When the gas has been sufficiently depleted so that stars dominate the potential, the appropriate accretion radius is the Bondi – Hoyle radius,

$$R_{\text{B}} = \frac{2Gm_{\bullet}}{v_{\text{rel}}^2}. \quad (9)$$

Following Bonnell et al. (2001), and with the intention to provide a minimum estimate of the BH mass growth, we chose the accretion radius to be the smaller of the two,

$$R_{\text{acc}}(t) = \min\{R_{\text{B}}, R_{\text{tid}}\}, \quad (10)$$

at any instant of time.

We accounted in the analysis for the dynamical friction generated by the stellar component according to Chandrasekhar's formula

$$F_{\text{df},\star} = \frac{4\pi G^2 m_{\bullet}^2}{v^2} \rho_{\star} \ln \Lambda \left\{ \text{Erf} \left(\frac{v}{\sigma\sqrt{2}} \right) - \frac{2}{\sqrt{\pi}} \frac{v}{\sigma\sqrt{2}} e^{-\frac{v^2}{2\sigma^2}} \right\}, \quad (11)$$

where v is the velocity of the BH, and the Coulomb logarithm is $\ln \Lambda = \ln(b_{\text{max}}/b_{\text{min}})$ with

$$b_{\text{max}} = r_{\text{hm}}, \quad b_{\text{min}} = \frac{Gm_{\bullet}}{3v_{\text{rel}}^2}. \quad (12)$$

Lee & Stahler (2011, 2014) reported that the gaseous dynamical friction on an accretor is

$$\mathbf{a}_{\text{acc}} = -\frac{\dot{m}}{m} \mathbf{v} \quad (13)$$

for both subsonic and supersonic accretors. This result seems to agree with the earlier calculation of Hadjidemetriou (1963). It may be understood simply as manifesting angular momentum preservation (see also Roupas & Kazanas 2019). Lee & Stahler (2011, 2014) argued that Eq. (13) encompasses the entire gaseous dynamical friction. In contrast, Tanaka & Haiman (2009) added to this term the gaseous dynamical friction formula proposed by Ostriker (1999; see also Ruffert 1996; Sánchez-Salcedo & Brandenburg 2001; Kim & Kim 2009; Indulekha 2013; Antoni et al. 2019). This issue seems unresolved. We adopted the former view (only the term in Eq. (13), neglecting Ostriker's term) because it generates the least gaseous dynamical friction. We wish to provide minimum estimates of BH growth, and higher dynamical friction causes BHs to sink deeper into the center where the gas density is higher, leading to more intense BH growth. Including Ostriker's

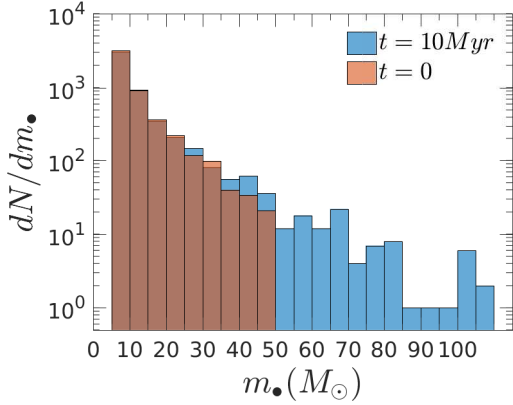


Fig. 1. Total BH mass distribution of ten samples of $N_{\bullet} = 500$ BHs, with an initial mass cutoff at $50 M_{\odot}$, bound in the core of a primordial gaseous stellar cluster with $r_{\text{hm}} = 1$ pc. We assumed the gas to be exponentially depleted by 99% within 10 Myr.

formula would only lead to results that support our conclusions more strongly. We can also report that when Ostriker’s formula is used instead of Eq. (13), we obtained similar results.

Finally, we describe the equations of motion of the BH-subcluster. As we noted, the BHs were assumed to be initially spherically distributed, and we are interested only in the amount of gas that they can accrete within 10 Myr. Therefore we did not follow the angular redistribution of their orbital planes and also assumed that the background of stars and gas is spatially fixed, as in Eqs. (3)–(4). We further assumed that the gas is depleted by 99% uniformly and exponentially within $t_f = 10$ Myr due to any process. The BHs accrete gas according to Eq. (6) and are subject to dynamical friction generated by stars, Eq. (11), as well as gas, Eq. (13). The system of equations that we solved numerically for $i = 1, \dots, N_{\text{BH}}$ was therefore

$$\frac{dr_i}{dt} = v_{r,i}, \quad (14)$$

$$\frac{d\phi_i}{dt} = \frac{v_{\phi,i}}{r_i}, \quad (15)$$

$$\frac{dv_{r,i}}{dt} = -\frac{v_{\phi,i}^2}{r_i} - \frac{GM(r < r_i(t))}{r_i^2} - (a_{\text{df},\star,i} + a_{\text{acc},i}) \frac{v_{r,i}}{v_i}, \quad (16)$$

$$\frac{dv_{\phi,i}}{dt} = -\frac{v_{r,i}v_{\phi,i}}{r_i} - (a_{\text{df},\star,i} + a_{\text{acc},i}) \frac{v_{\phi,i}}{v_i}, \quad (17)$$

$$\frac{dm_i}{dt} = \pi \rho_{\text{gas}}(r_i(t)) v_{\text{rel}}(v_i(t), c_s(r_i(t))) R_{\text{acc}}(v_i(t), c_s(r_i(t)))^2. \quad (18)$$

The total mass within r_i , $M(r < r_i(t))$, includes the enclosed mass of field stars, of the remaining enclosed gas at t , and the enclosed mass within r_i of the BH population.

3. Results

We performed the numerical analysis for an initial total cluster mass $M_{\text{tot}} = 10^6 M_{\odot}$, stellar formation efficiency $\epsilon = 0.3$, and several initial compactness values $C_{\text{ini}}(t = 0) = \{5-10\}$, corresponding to half-mass radii $r_{\text{hm}} = \{2-1\}$ pc. For each half-mass radius we performed ten simulations with different initial conditions. We also assumed a depletion time $t_f = 10$ Myr.

In Fig. 1 we show the final BH mass function in the case of a half-mass radius $r_{\text{hm}} = 1$ pc, that is, an initial compactness $C_{\text{ini}} = 10$, which results in a final compactness $C_{\text{fin}} = 3$. It is evident that about 2% of the BHs fill in the theoretical BH mass gap. In addition, the BH mass function increases for mass values

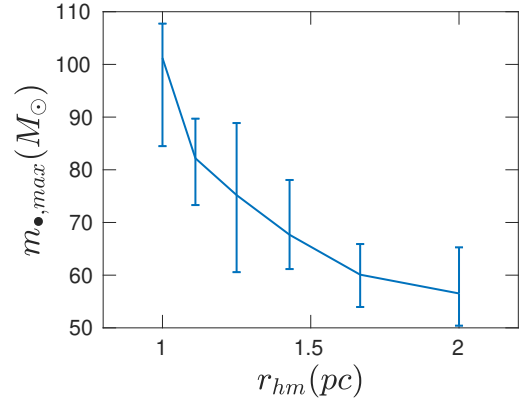


Fig. 2. Maximum BH mass achieved after gas accretion of an initial population of $N_{\bullet} = 500$ BHs, with initial mass cutoff at $50 M_{\odot}$, with respect to the half-mass radius r_{hm} of the stellar cluster for ten samples of the BH population at each r_{hm} . The vertical lines represent the range of maximum mass values achieved in all samples for a certain r_{hm} . The gas is assumed to be exponentially depleted by 99% within $t_f = 10$ Myr.

Table 1. In five out of the ten samples with $r_{\text{hm}} = 1$ pc, six massive BBHs formed within 10 Myr because they sank to the center of the cluster ($\leq 10^{-3}$ pc).

$m_{\bullet} (M_{\odot})$	$a_{\text{BBH}} (10^{-3} \text{ pc})$
{108, 45}	0.1
{82, 64}	1.6
{73, 51}	0.8
{90, 63}	0.1
{78, 75}	0.9
{84, 68}	1.6

Notes. We assumed a population of $N_{\bullet} = 500$ BHs, with different initial conditions in each sample, immersed inside the core of the same gaseous primordial cluster. The left column denotes the masses of each BBH member, and the right column lists the semi-major axis of each BBH.

higher than $\gtrsim 35 M_{\odot}$. This is also true for any $r_{\text{hm}} \leq 1.25$ pc. In Fig. 2 we draw the maximum BH mass that is achieved after accretion with respect to the half-mass radius of the cluster.

We report that the probability for a BH to exceed the initial cutoff of $50 M_{\bullet}$ is

$$P(m_{\bullet} > 50 M_{\odot}) \approx 1-2\%, \quad \text{for } r_{\text{hm}} = \{1-2\} \text{ pc}, \quad (19)$$

where the lower half-mass radius corresponds to the upper probability value. The total accreted gas mass lies in the range $\approx \{100-900\} M_{\odot}$ (with an initial mass $5400 M_{\odot}$ of the BH population).

We remark that the more massive BHs sink deeper into the center $r \lesssim 10^{-3}$ pc, where it is most probable that they form hard massive BBHs. About half of the samples at each r_{hm} form a BBH by this process. In particular, for $r_{\text{hm}} = 1$ pc, six hard BBHs were formed in five out of the ten samples, as shown in Table 1. The BBHs will continue to form for $t > 10$ Myr by dynamical processes.

4. Conclusions

We propose that in the early life of dense stellar clusters a rapid accretion process operates. The BHs that have been generated by the more massive stars, most of which are therefore segregated in

the core, rapidly accrete primordial gas that is also accumulated in the core within ~ 10 Myr before it becomes depleted by stellar formation feedback processes. The amount of accreted gas depends sensitively on the mass and compactness of the cluster and on the accretion timescale.

The BH (and BBH) initial mass function are shifted toward higher values. For sufficiently compact clusters the BH mass limit of $50 M_{\odot}$ predicted by supernova theory is exceeded. We estimate that 1–2% of the initial BH population exceeds the theoretical limit for an initial cluster mass $M_{\text{tot}} = 10^6 M_{\odot}$, half-mass radii $r_{\text{hm}} = \{1, 2\}$ pc, depletion time $t_f = 10$ Myr, and stellar formation efficiency $\epsilon = 0.3$.

In addition, the BH mass growth together with dynamical friction causes the more massive BHs to sink deeper into the center of the cluster. There, it is more probable that the BH attains mutual binding energies that are higher than each individual binding energy with the cluster and thus forms hard massive BBHs, as listed in Table 1. This process operates rapidly in the very early life of the cluster, therefore it adds to the dynamical channel of BBH formation and mergers (Abbott et al. 2016), which will continue to operate during the whole life of the cluster (e.g., see Park et al. 2017; Hong et al. 2018). Thus, an observation by LIGO of BBH mergers with member masses above $50 M_{\odot}$ is to be anticipated.

We assumed an initial BH population bound inside the core of the cluster neglecting close encounters. These are not expected to have a strong effect on our results because their timescale are relatively long. The recoil mechanism of BBH-BH three-body encounters, which ejects BBHs out of the cluster (Sigurdsson & Hernquist 1993), operates on gigayear timescales, but our proposed mechanism works on megayear timescales. In addition, it has become evident during the past decade (Maccarone et al. 2007, 2011; Barnard et al. 2011; Shih et al. 2010; Strader et al. 2012) that recoil is not effective in sufficiently dense GCs, which contain a population (≤ 1000) of BHs and BBHs in their core (e.g., Morscher et al. 2015 and references therein).

More sophisticated simulations may be required in order to confirm with higher confidence and detail the effects of the rapid accretion process we proposed here and to provide a better estimation of the resulting BH mass function for a wider range of parameter values $\{M_{\text{tot}}, r_{\text{hm}}, t_f, \epsilon\}$. Nevertheless, our current analysis strongly supports the idea that the BH upper mass gap can be populated by rapid gas accretion onto the BHs of dense primordial stellar clusters.

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