

ARTICLE

Received 27 Jun 2014 | Accepted 20 Nov 2014 | Published 8 Jan 2015

DOI: 10.1038/ncomms6910

OPEN

Generation of spin currents by surface plasmon resonance

K. Uchida^{1,2}, H. Adachi^{3,4}, D. Kikuchi^{1,5}, S. Ito¹, Z. Qiu⁵, S. Maekawa^{3,4} & E. Saitoh^{1,3,4,5}

Surface plasmons, free-electron collective oscillations in metallic nanostructures, provide abundant routes to manipulate light–electron interactions that can localize light energy and alter electromagnetic field distributions at subwavelength scales. The research field of plasmonics thus integrates nano-photonics with electronics. In contrast, electronics is also entering a new era of spintronics, where spin currents play a central role in driving devices. However, plasmonics and spin-current physics have so far been developed independently. Here we report the generation of spin currents by surface plasmon resonance. Using Au nanoparticles embedded in Pt/BiY₂Fe₅O₁₂ bilayer films, we show that, when the Au nanoparticles fulfill the surface-plasmon-resonance conditions, spin currents are generated across the Pt/BiY₂Fe₅O₁₂ interface. This spin-current generation cannot be explained by conventional heating effects, requiring us to introduce nonequilibrium magnons excited by surface-plasmon-induced evanescent electromagnetic fields in BiY₂Fe₅O₁₂. This plasmonic spin pumping integrates surface plasmons with spin-current physics, opening the door to plasmonic spintronics.

¹Institute for Materials Research, Tohoku University, Sendai 980-8577, Japan. ²PRESTO, Japan Science and Technology Agency, Saitama 332-0012, Japan. ³Advanced Science Research Center, Japan Atomic Energy Agency, Tokai 319-1195, Japan. ⁴CREST, Japan Science and Technology Agency, Tokyo 102-0075, Japan. ⁵WPI Advanced Institute for Materials Research, Tohoku University, Sendai 980-8577, Japan. Correspondence and requests for materials should be addressed to K.U. (email: kuchida@imr.tohoku.ac.jp).

Surface plasmons enable subdiffraction-limit localization of light and strong enhancement of electromagnetic fields. Owing to this feature, plasmonics^{1–5} has vast potential in solar cells, light generation, microscopy, data storage and biosensors. One of the other emerging developments in this field is magneto-plasmonics^{6,7}, which offers unique possibilities to manipulate light by the use of external magnetic fields or magnetic materials. Important topics in magneto-plasmonics include the enhancement of magneto-optical effects in plasmonic nanostructures⁶, influence of magnetic fields on surface plasmon resonance (SPR) and observation of SPR in ferromagnetic metals⁷. However, the integration of plasmonics into spintronics^{8,9} is yet to be realized, and the interaction between surface plasmons and spin currents^{10,11} contains unexplored physics.

In the field of spintronics, for spin-current generation, we often rely on a spin pumping effect^{12–15}, which refers to the transfer of spin-angular momentum from magnetization dynamics in a ferromagnet to conduction-electron spins in an attached paramagnet; when a motion of magnetic moments in the ferromagnet is excited, a spin current is pumped out of the ferromagnet into the paramagnet.

Recent studies in spintronics have revealed that there are two types of spin pumping effects. One is coherent spin pumping, which is due to coherent phase precession of the magnetization, that is, ferromagnetic or spin-wave resonances in ferromagnets typically excited by microwave irradiation^{13–15}. The other is incoherent spin pumping, which is due to a nonequilibrium state of spins at a ferromagnet/paramagnet interface excited by incoherent external perturbations; when magnons in the ferromagnet and/or electrons in the paramagnet deviate from thermal equilibrium, a net spin current is generated across the interface. A typical example of the incoherent spin pumping is the spin Seebeck effect^{16–24}, where the magnon/electron distribution functions are modulated by external temperature gradients.

In this work, we report the observation of plasmonic spin pumping; we show that the incoherent spin pumping can be driven by surface plasmons under visible light illumination. This plasmonic spin pumping will invigorate research in both the fields of spintronics and plasmonics, paving the way to create new devices, such as solar spin-current generators and spin-based plasmonic sensors.

Results

Device structure. Figure 1a shows a schematic illustration of the sample system used in the present study. The sample consists of a paramagnetic Pt/ferromagnetic BiY₂Fe₅O₁₂ bilayer film with Au nanoparticles (NPs) embedded in the BiY₂Fe₅O₁₂ layer. The Au NPs with an in-plane diameter of 30–90 nm and a height of 20–50 nm were formed by heating a Au thin film⁶ on a single-crystalline Gd₃Ga₅O₁₂ (111) substrate (Fig. 2b,c). The 110-nm-thick BiY₂Fe₅O₁₂ film was then coated on them by means of a metal-organic decomposition (MOD) method²² (see Methods). Finally, a 5-nm-thick Pt film was sputtered on to the BiY₂Fe₅O₁₂ film. To excite surface plasmons, the Pt/BiY₂Fe₅O₁₂/Au-NP sample was illuminated with unpolarized monochromatic light with wavelength λ (400–770 nm) and power P from the Pt-film side at normal incidence (Fig. 1a). As the Pt and BiY₂Fe₅O₁₂ layers are very thin, the light passes through the layers and interacts with the Au NPs. We checked that the electric resistance between the ends of the BiY₂Fe₅O₁₂ layer with Au NPs under the light illumination is much greater than the measurement limit of our electrometer ($> 2 \times 10^{11} \Omega$) at all the λ values; the possibility of optical charge-carrier generation in BiY₂Fe₅O₁₂ is excluded.

In Fig. 1c, we show the light-transmittance spectrum of the Pt/BiY₂Fe₅O₁₂/Au-NP sample (see the blue circle data points). We found that the sample exhibits a clear dip structure in the transmittance spectrum around $\lambda = 690$ nm in comparison with the spectra of a Gd₃Ga₅O₁₂ substrate, a plain Pt film and

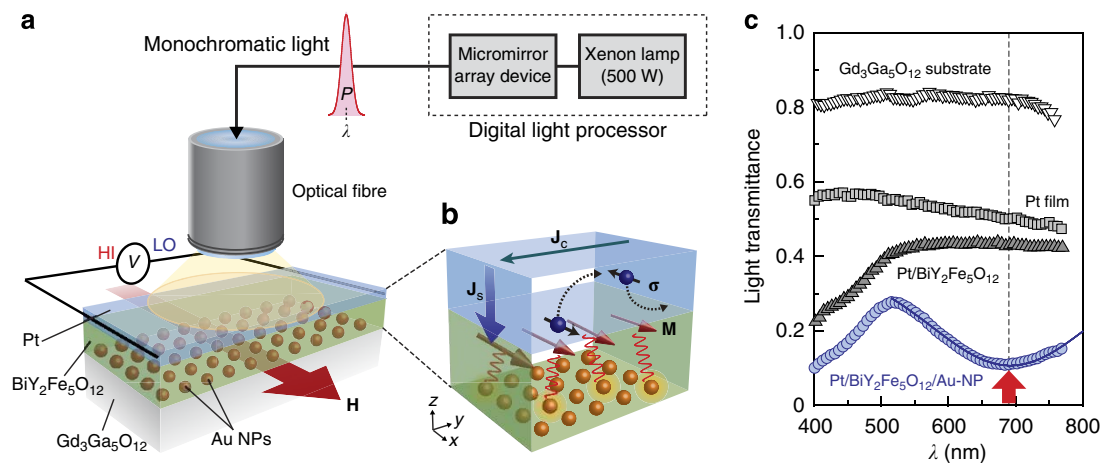


Figure 1 | Experimental set-up and device structure. (a) A schematic illustration of the Pt/BiY₂Fe₅O₁₂/Au-nanoparticle (NP) sample and experimental configuration for measuring the plasmonic spin pumping. The Pt/BiY₂Fe₅O₁₂/Au-NP sample is of a 10 mm × 5 mm rectangular shape. The sample was illuminated with monochromatic light with the wavelength λ (400–770 nm), power P , bandwidth of 20 nm and spot diameter of ~ 5 mm from the Pt-film side at normal incidence by using a digital light processor comprising a 500 W Xenon lamp and a micromirror array device (Gooch & Housego, OL490). HI and LO represent the inputs of a voltmeter. H denotes the external magnetic field vector (red arrow) with the magnitude H . (b) Plasmonic spin pumping and inverse spin Hall effect (ISHE) in the Pt/BiY₂Fe₅O₁₂/Au-NP sample. Here, M , J_c , J_s and σ denote the magnetization vector (dark red arrows), charge current generated by the ISHE (green arrow), spatial direction of the spin current generated by the plasmonic spin pumping (blue arrow) and spin-polarization vector of the spin current (black arrows), respectively. Blue (orange) spheres schematically show electron charge (Au NPs). Black dotted arrows and red wavy lines schematically show the trajectory of the electrons and the interaction between magnetic moments in the BiY₂Fe₅O₁₂ layer and evanescent electromagnetic fields induced by the surface plasmon resonance (SPR), respectively. (c) λ dependence of the light transmittance of the 0.5-mm-thick Gd₃Ga₅O₁₂ substrate, 5-nm-thick Pt film on the Gd₃Ga₅O₁₂ substrate, Pt/BiY₂Fe₅O₁₂ sample and Pt/BiY₂Fe₅O₁₂/Au-NP sample. In the Pt/BiY₂Fe₅O₁₂ sample, the Pt and BiY₂Fe₅O₁₂ films were formed directly on the Gd₃Ga₅O₁₂ substrate without Au NPs. A blue solid curve was obtained by fitting the observed dip structure using a Gaussian function. A red arrow shows the wavelength at which the SPR is excited in the Pt/BiY₂Fe₅O₁₂/Au-NP sample.

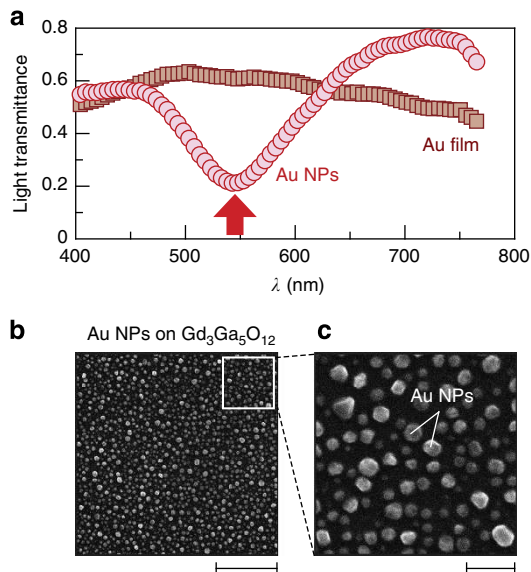


Figure 2 | Au nanoparticles. (a) λ dependence of the light transmittance of a 15-nm-thick Au film and Au NPs on the $\text{Gd}_3\text{Ga}_5\text{O}_{12}$ substrate. The light-transmittance spectrum of the Au-NP sample exhibits a clear dip structure due to the SPR, while that of the Au film does not exhibit strong λ dependence. The position of the SPR wavelength for the Au-NP sample is marked with a red arrow. (b) Scanning electron microscope image of Au NPs on the $\text{Gd}_3\text{Ga}_5\text{O}_{12}$ substrate, measured before forming the $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ and Pt layers. (c) Magnified scanning electron microscope image of the Au NPs on the $\text{Gd}_3\text{Ga}_5\text{O}_{12}$ substrate. The scale bars in (b) and (c) represent 1 μm and 200 nm, respectively.

Pt/ $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ sample without Au NPs. This dip structure for the Pt/ $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ /Au-NP sample is attributed to the excitation of localized SPR^{1,2} in the Au NPs embedded in the $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ layer (see also Fig. 2a).

Electromagnetic field distributions. To visualize the SPR in the $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ /Au-NP structure, we performed electromagnetic field simulations by means of a finite-difference time-domain (FDTD) method²⁵. Figure 3a shows a schematic illustration of a $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ /Au-NP model used for the FDTD simulations (see Methods for details). As shown in Fig. 3b, under the SPR condition ($\lambda = 690$ nm), strong evanescent electromagnetic fields, that is, near fields^{1,2}, are generated in the $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ film in the vicinity of the $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ /Au-NP interface, while the enhancement and localization of electromagnetic fields do not appear at $\lambda = 500$ nm (Fig. 3c), consistent with the light-transmittance spectrum of the Pt/ $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ /Au-NP sample in Fig. 1c. The FDTD simulations also show that the evanescent electromagnetic fields in the $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ layer due to the SPR are not affected by the presence of the Pt layer irrespective of the position and size of Au NPs (Supplementary Figs 1 and 2).

Measurement mechanism of plasmonic spin pumping. In the Pt/ $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ /Au-NP sample, if the evanescent electromagnetic fields concomitant with surface plasmons in Au NPs excite magnons in the $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ film and the excited magnons drive the incoherent spin pumping, a spin current is generated in the Pt layer with the spatial direction \mathbf{J}_S and the spin-polarization vector $\boldsymbol{\sigma}$ parallel to the magnetization \mathbf{M} of the $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ film (Fig. 1b). This plasmon-induced spin current is converted into a d.c. charge current \mathbf{J}_C due to the inverse spin Hall effect (ISHE)^{14,26,27} in the Pt layer owing to the strong spin-orbit interaction in Pt. When \mathbf{M} of the $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ film is along the x

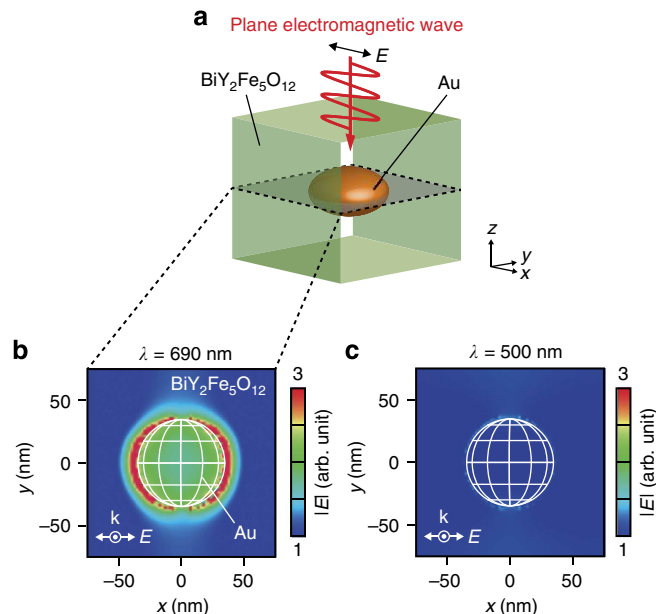


Figure 3 | Numerical calculation. (a) A schematic illustration of the $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ /Au-NP model used for the numerical calculation based on a finite-difference time-domain method. Plane electromagnetic waves with the polarization along the x direction and the wave vector \mathbf{k} along the z direction are applied to the $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ /Au-NP model. \mathbf{E} denotes an electric field of the electromagnetic waves. (b,c) Simulated distributions of the electric field intensity $|E|$ in the $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ /Au-NP model in the x - y plane across the centre of the Au NP at $\lambda = 690$ nm (b) and 500 nm (c).

direction, \mathbf{J}_C is generated along the y direction because of the following ISHE symmetry:

$$\mathbf{J}_C \propto \mathbf{J}_S \times \boldsymbol{\sigma} \quad (1)$$

where the x , y and z directions are defined in Fig. 1b. We note that extrinsic artefacts induced by a static magnetic proximity effect²⁸ at the Pt/ $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ interface are negligibly small in the present structure^{23,24}. To detect the ISHE induced by the plasmonic spin pumping, we measured an electric voltage V between the ends of the Pt layer under open circuit condition while illuminating with monochromatic light and applying an external magnetic field \mathbf{H} (with the magnitude H) to the Pt/ $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ /Au-NP sample along the x - y plane (Fig. 1a). All of the V measurements were carried out at room temperature and atmospheric pressure.

Observation of plasmonic spin pumping. To observe the plasmonic spin pumping, it is important to separate the plasmon-induced signals from extrinsic heating effects. First of all, to check the heating effect due to the light illumination, we measured the voltage in the Pt/ $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ sample without Au NPs. The grey triangle data points in Fig. 4b show the λ dependence of the voltage normalized by the incident light power, V/P , in the Pt/ $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ sample at $H = 200$ Oe. When \mathbf{H} is applied along the x direction, a finite voltage signal appears in the Pt layer at all the λ values. Because of the absence of Au NPs and surface plasmons, the signal in the Pt/ $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ sample is attributed to the heating of the sample by the light illumination, that is, the longitudinal spin Seebeck effect^{20,22–24} (see also Supplementary Fig. 3). Here, the sign of the V/P signal shows that the temperature of the Pt layer is higher than that of the $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ layer under the light illumination (Supplementary Fig. 4 and Supplementary Note 1). Importantly, this heating signal in the Pt/ $\text{BiY}_2\text{Fe}_5\text{O}_{12}$ sample does not exhibit strong λ dependence (Fig. 4b).

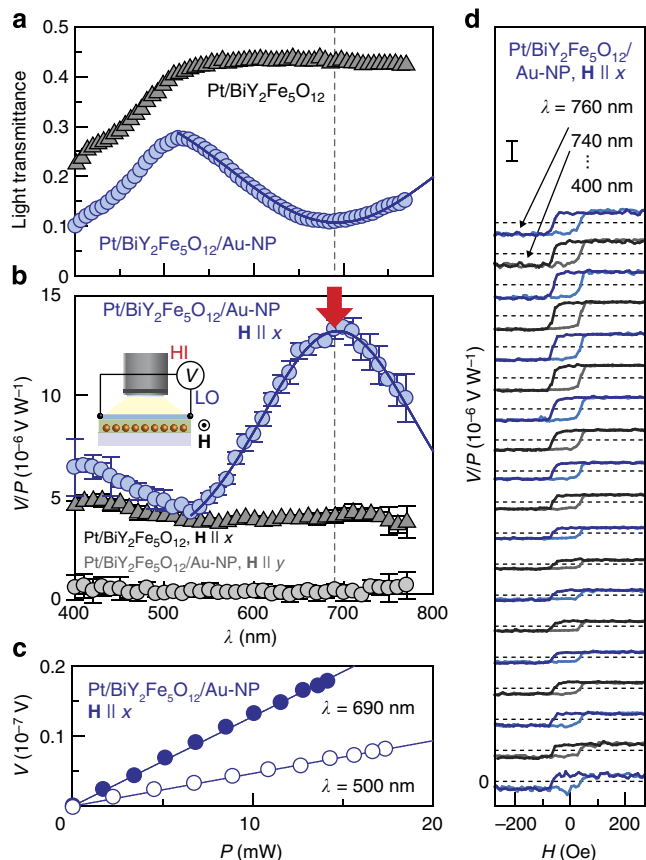


Figure 4 | Observation of plasmonic spin pumping. (a) λ dependence of the light transmittance of the Pt/BiY₂Fe₅O₁₂/Au-NP and Pt/BiY₂Fe₅O₁₂ samples. The samples were prepared at the same time. (b) λ dependence of the electric voltage between the ends of the Pt layer normalized by the incident light power, V/P , in the Pt/BiY₂Fe₅O₁₂/Au-NP and Pt/BiY₂Fe₅O₁₂ samples, measured when \mathbf{H} of $H = 200$ Oe was applied along the x or y direction. The in-plane coercive force of the BiY₂Fe₅O₁₂ layer is around 30 Oe; the magnetization of BiY₂Fe₅O₁₂ is aligned along the \mathbf{H} direction at $H = 200$ Oe. The position of the SPR wavelength is marked with a red arrow. A blue solid curve was obtained by fitting the observed peak voltage structure using a Gaussian function. The error bars represent the s.d. of the measurements. (c) P dependence of V in the Pt/BiY₂Fe₅O₁₂/Au-NP sample at $\lambda = 690$ nm and 500 nm, measured when \mathbf{H} of $H = 200$ Oe was applied along the x direction. (d) H dependence of V/P in the Pt/BiY₂Fe₅O₁₂/Au-NP sample for various values of λ ($= 760, 740, \dots, 400$ nm), measured when \mathbf{H} was along the x direction. The scale bar represents $V/P = 20 \times 10^{-6}$ V W⁻¹. The dark blue and black lines (light blue and grey lines) show the V/P signals measured when \mathbf{H} was swept from positive to negative (from negative to positive).

Now we are in a position to investigate the plasmonic generation of spin currents. The blue circle data points in Fig. 4b are the V/P spectrum in the Pt/BiY₂Fe₅O₁₂ film that contains Au NPs. Also in this Pt/BiY₂Fe₅O₁₂/Au-NP sample, when $\mathbf{H} \parallel x$, the V/P signal appears with the same sign (Fig. 4b) and the magnitude of V is proportional to the light power P (Fig. 4c). The sign of the signal at each λ was observed to be reversed by reversing H with a hysteresis loop, indicating that the signal is affected by the \mathbf{M} direction of the BiY₂Fe₅O₁₂ layer (Fig. 4d). As also shown in Fig. 4b, we confirmed that the V/P signal disappears when \mathbf{H} is applied along the y direction, consistent with equation (1). These results indicate that the observed V/P signal in the Pt/BiY₂Fe₅O₁₂/Au-NP sample is due to the ISHE induced by spin currents in the Pt layer. Significantly,

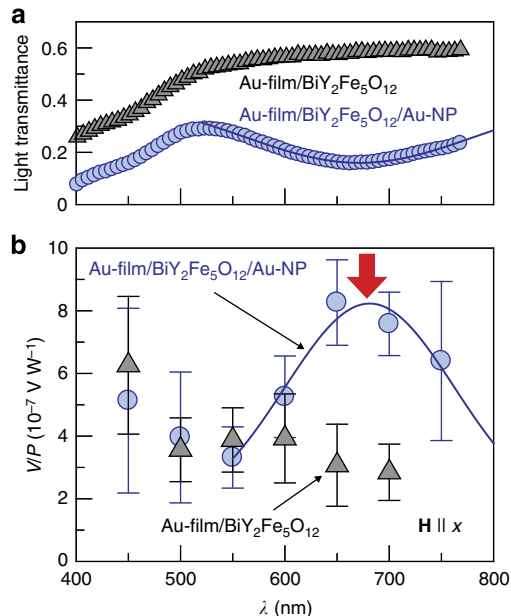


Figure 5 | Measurements using Au film. (a) λ dependence of the light transmittance of the Au-film/BiY₂Fe₅O₁₂/Au-NP and Au-film/BiY₂Fe₅O₁₂ samples. (b) λ dependence of V/P in the Au-film/BiY₂Fe₅O₁₂/Au-NP and Au-film/BiY₂Fe₅O₁₂ samples, measured when \mathbf{H} of $H = 200$ Oe was applied along the x direction. The Au-film/BiY₂Fe₅O₁₂/Au-NP sample was illuminated with monochromatic light with the wavelength λ , power P , bandwidth of 50 nm and spot diameter of ~ 5 mm from the Au-film side at normal incidence.

Fig. 4b shows that the V/P signal in the Pt/BiY₂Fe₅O₁₂/Au-NP sample exhibits peak structure and is markedly enhanced around $\lambda = 690$ nm. This peak position in the V/P spectrum coincides with the SPR wavelength of this Pt/BiY₂Fe₅O₁₂/Au-NP sample (Fig. 4a,b). As both the Pt/BiY₂Fe₅O₁₂/Au-NP and Pt/BiY₂Fe₅O₁₂ samples are illuminated from the Pt-film side, the temperature rise of the Pt layer is almost the same for both, confirming that the heating of the Pt layer is irrelevant to the ISHE enhancement under the SPR condition in the Pt/BiY₂Fe₅O₁₂/Au-NP sample (see Supplementary Fig. 4 and Supplementary Note 1), where we note again that the background heating signal does not exhibit strong λ dependence. This peak V/P structure in the Pt/BiY₂Fe₅O₁₂/Au-NP sample also cannot be explained by the heating of Au NPs or the BiY₂Fe₅O₁₂ layer under the SPR condition because of the different sign; if the BiY₂Fe₅O₁₂ layer was heated by electromagnetic loss in Au NPs, the ISHE voltage should decrease under the SPR condition as the temperature gradient generated by the Au-NP heating across the Pt/BiY₂Fe₅O₁₂ interface is of an opposite sign to the case when the Pt layer is heated (see Supplementary Fig. 4 and Supplementary Note 1). A possible slight heating of Au NPs in the off-resonance region might reduce the magnitude of the background signal slightly, but is also irrelevant to the V/P signal with the positive peak structure. Our numerical calculations of electromagnetic field distributions based on the FDTD method show that the electromagnetic coupling between the Pt film and Au NPs under the SPR condition is too weak to explain the enhancement of the ISHE signal in the Pt/BiY₂Fe₅O₁₂/Au-NP sample; as no SPR mode due to the coupling appears in the configuration used in this study, the temperature rise of the Pt layer is not changed by the SPR in Au NPs (Supplementary Figs 1 and 2).

We also performed control experiments by using different materials. We found that similar V/P signals appear in a

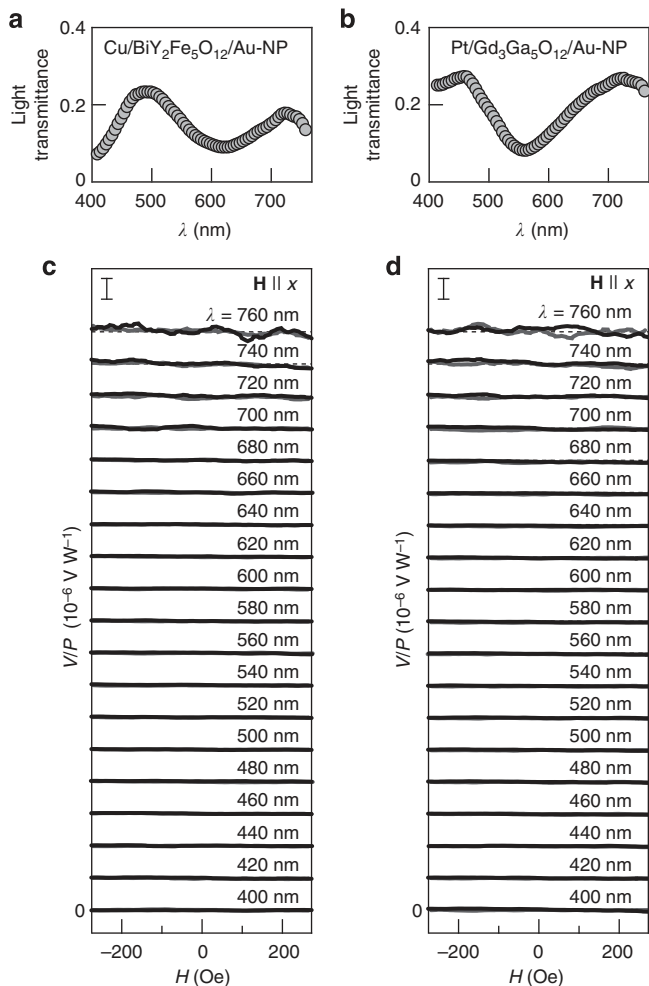


Figure 6 | Control experiments using different materials. (a,b) λ dependence of the light transmittance of the Cu/BiY₂Fe₅O₁₂/Au-NP (a) and Pt/Gd₃Ga₅O₁₂/Au-NP (b) samples. (c,d) H dependence of V/P in the Cu/BiY₂Fe₅O₁₂/Au-NP (c) and Pt/Gd₃Ga₅O₁₂/Au-NP (d) samples for various values of λ (= 760, 740, ..., 400 nm), measured when H was along the x direction. The scale bars represent $V/P = 20 \times 10^{-6} \text{ V W}^{-1}$. All the control experiments were performed by using the samples of a 10 mm \times 5 mm rectangular shape.

Au-film/BiY₂Fe₅O₁₂/Au-NP sample, where the Pt film is replaced with a 5-nm-thick Au film (Fig. 5a,b); as Au is a typical metal far from ferromagnetism^{29,30}, proximity ferromagnetism in Pt cannot be the origin of the observed peak voltage structure. In contrast, this voltage signal was found to disappear in a Cu/BiY₂Fe₅O₁₂/Au-NP sample, where the Pt layer is replaced with a 5-nm-thick Cu film in which the spin-orbit interaction is very weak, indicating the important role of spin-orbit interaction, or the ISHE, in the voltage generation (Fig. 6a,c). The voltage signal also disappears in a Pt/Gd₃Ga₅O₁₂/Au-NP sample, where Au NPs are embedded in a paramagnetic insulator Gd₃Ga₅O₁₂ film instead of BiY₂Fe₅O₁₂ (see Fig. 6b,d and Methods), indicating that direct contact between BiY₂Fe₅O₁₂ and Pt is necessary for the observed voltage generation; electromagnetic artefacts are irrelevant. All the control experiments support our interpretation that the enhancement of the ISHE signal under the SPR condition in the Pt/BiY₂Fe₅O₁₂/Au-NP sample results from the plasmonic spin pumping.

Separation of plasmonic spin pumping from heating effects. To further verify the effect of surface plasmons on the spin-current

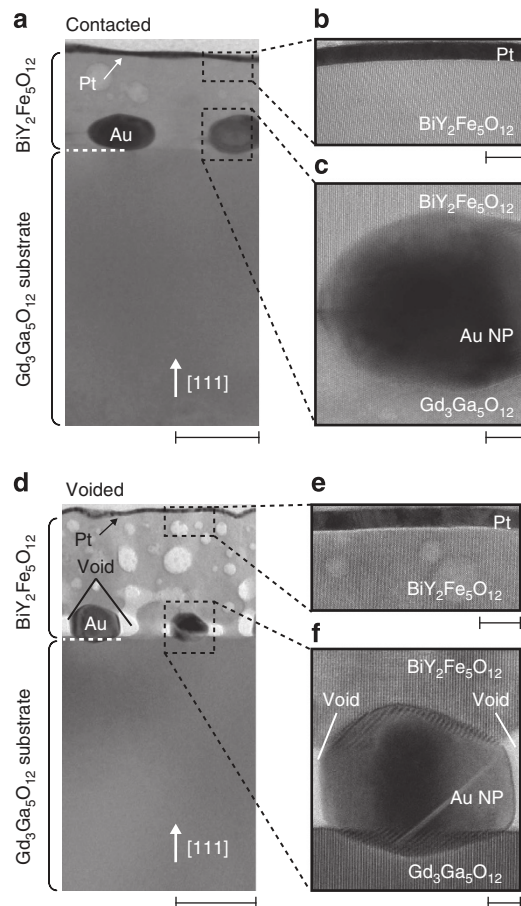


Figure 7 | Cross-sectional images. (a) Cross-sectional transmission electron microscope (TEM) images of the contacted Pt/BiY₂Fe₅O₁₂/Au-NP sample. (b,c) Magnified TEM images of the contacted Pt/BiY₂Fe₅O₁₂/Au-NP sample near the Pt film (b) and the Au NP (c). (d) Cross-sectional TEM images of the voided Pt/BiY₂Fe₅O₁₂/Au-NP sample. (e,f) Magnified TEM images of the voided Pt/BiY₂Fe₅O₁₂/Au-NP sample near the Pt film (e) and the Au NP (f). Scale bars in (a) and (d) (b,c,e and f) represent 100 nm (10 nm). White arrows in (a) and (d) represent the crystal orientation of the Gd₃Ga₅O₁₂ substrate. The voided Pt/BiY₂Fe₅O₁₂/Au-NP sample has small voids with the size of 20–30 nm between the BiY₂Fe₅O₁₂ film and Au NPs, while Au NPs in the contacted sample are densely embedded in BiY₂Fe₅O₁₂.

generation, we measured the V/P spectra in the Pt/BiY₂Fe₅O₁₂/Au-NP samples on changing the contact condition between BiY₂Fe₅O₁₂ and Au NPs. Here, we prepared the Pt/BiY₂Fe₅O₁₂/Au-NP samples by means of two different annealing conditions for the BiY₂Fe₅O₁₂ layer (see Methods for details). In the Pt/BiY₂Fe₅O₁₂/Au-NP sample used for the experiments in Fig. 4, Au NPs are densely embedded in the BiY₂Fe₅O₁₂ film (Fig. 7a–c), where the peak V/P structure appears under the SPR condition (Fig. 8a–d). In contrast, we also fabricate the Pt/BiY₂Fe₅O₁₂/Au-NP sample in which the Au NPs are separated from the BiY₂Fe₅O₁₂ film by small voids (Fig. 7d–f), which interrupt the interaction between magnons in BiY₂Fe₅O₁₂ and surface plasmons in Au NPs as the evanescent electromagnetic fields generated by the SPR are localized in the vicinity of Au NPs (Supplementary Fig. 5). The difference in the SPR wavelength between the ‘contacted’ and ‘voided’ Pt/BiY₂Fe₅O₁₂/Au-NP samples is attributed to the difference in the dielectric constant of the medium around Au NPs; As Au NPs in the voided Pt/BiY₂Fe₅O₁₂/Au-NP sample are surrounded by voids, its SPR

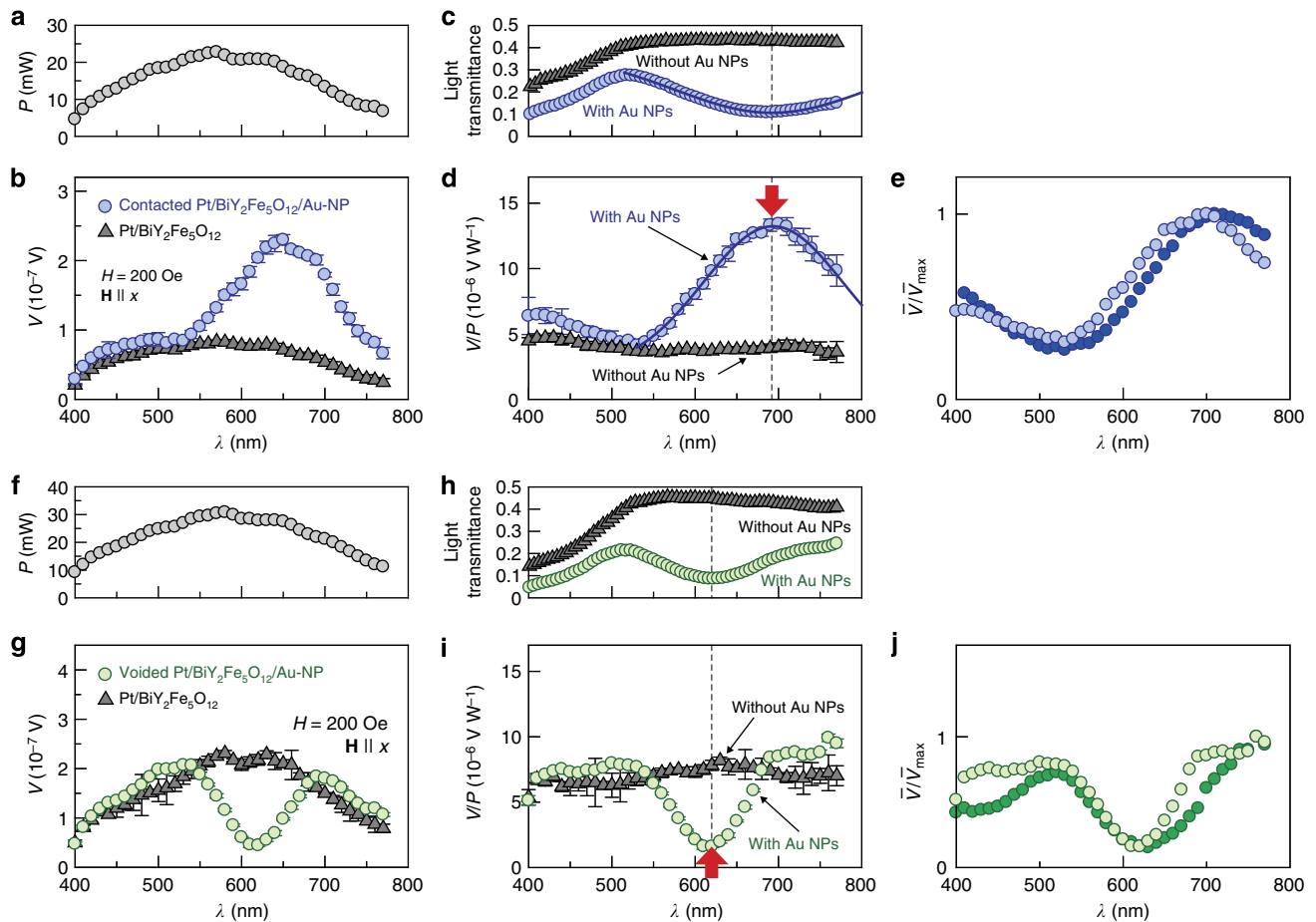


Figure 8 | Evidence for magnon excitation by surface plasmons. (a) λ dependence of P , applied when the data in (b) were measured. (b) λ dependence of V in the contacted Pt/BiY₂Fe₅O₁₂/Au-NP and Pt/BiY₂Fe₅O₁₂ samples, measured when \mathbf{H} of $H = 200$ Oe was applied along the x direction. The error bars represent the s.d. of the measurements. (c) λ dependence of the light transmittance of the contacted Pt/BiY₂Fe₅O₁₂/Au-NP and Pt/BiY₂Fe₅O₁₂ samples. (d) λ dependence of V/P in the contacted Pt/BiY₂Fe₅O₁₂/Au-NP and Pt/BiY₂Fe₅O₁₂ samples. (e) λ dependence of \bar{V}/\bar{V}_{\max} in two different contacted Pt/BiY₂Fe₅O₁₂/Au-NP samples, where $\bar{V} \equiv V/P$ and \bar{V}_{\max} is the maximum value of \bar{V} in the $\lambda - \bar{V}$ spectrum. (f–j) Experimental results for the voided Pt/BiY₂Fe₅O₁₂/Au-NP and Pt/BiY₂Fe₅O₁₂ samples. (j) shows the λ dependence of \bar{V}/\bar{V}_{\max} in two different voided Pt/BiY₂Fe₅O₁₂/Au-NP samples.

wavelength is closer to that of plain Au NPs without BiY₂Fe₅O₁₂ (compare Figs 2a and 8c,h). In the voided Pt/BiY₂Fe₅O₁₂/Au-NP sample, dip structure was found to appear in the V/P spectrum under the SPR condition (Fig. 8f–i). This behaviour was observed not only in one sample but also in our different samples as exemplified in Fig. 8e,j. The dip V/P signal in the voided Pt/BiY₂Fe₅O₁₂/Au-NP sample is attributed to the heating of Au NPs by the SPR, which generates the temperature gradient across the Pt/BiY₂Fe₅O₁₂ interface and the longitudinal spin Seebeck voltage along the direction opposite to the case when the Pt layer is heated (see Supplementary Figs 4,5 and Supplementary Note 1). These results indicate that the enhancement of the ISHE under the SPR condition, observed in the contacted Pt/BiY₂Fe₅O₁₂/Au-NP sample, requires direct contact between the BiY₂Fe₅O₁₂ and the Au NPs and cannot be explained by conventional heating effects, providing evidence for the operation of the plasmonic spin pumping due to nonequilibrium magnons excited by surface plasmons.

Discussion

The above experiments clearly show that the spin current generated by the light illumination in the contacted Pt/BiY₂Fe₅O₁₂/Au-NP sample is strongly enhanced under the

SPR condition, and this enhancement is not connected to extrinsic heating effects. So then, what is the origin of the spin current under the SPR condition? Following previous studies on the incoherent spin pumping^{31,32}, to generate the positive ISHE voltage in the present configuration, magnons in the BiY₂Fe₅O₁₂ layer should be excited (see Supplementary Notes 1 and 2). A potential candidate for the magnon-excitation mechanism in the present sample structure is the strong evanescent electromagnetic fields (near-field photons) concomitant with surface plasmons. If the near-field photons excite magnons via the magnon–photon interaction^{33–36} near the BiY₂Fe₅O₁₂/Au-NP interface, the excited magnons in the BiY₂Fe₅O₁₂ layer produce a spin current owing to the incoherent spin-pumping mechanism and the positive ISHE voltage in the Pt layer: this is the plasmonic spin pumping. Using a standard many-body technique, we formulate theoretically the plasmonic spin pumping based on magnon–photon Raman scattering process (Supplementary Fig. 6). The spin current pumped by light illumination is described as

$$J_S = \Gamma^+ \Gamma^- G_S \frac{\hbar v a_S^4}{c\alpha} |E|^2 \quad (2)$$

within the linear response region (see Supplementary Note 2 for details). Here, $\Gamma^{+(-)}$ is the dimensionless magnon–photon coupling constant, G_S the strength of the magnetic coupling at the

Pt/BiY₂Fe₅O₁₂ interface, \hbar the Planck constant divided by 2π , ν the photon frequency, c the velocity of light, a_s^3 the effective block spin volume, α the Gilbert damping constant and E the electric field amplitude induced by surface plasmons. Equation (2) indicates that the plasmon-induced spin current in the Pt/BiY₂Fe₅O₁₂/Au-NP sample is proportional to the electromagnetic energy in the BiY₂Fe₅O₁₂ layer ($\propto |E|^2$), a situation consistent with the experimental results in Fig. 4c as $|E|^2$ generated by surface plasmons is proportional to the incident light power P . The observed peak ISHE voltage in the contacted Pt/BiY₂Fe₅O₁₂/Au-NP sample is attributed to the strong evanescent electromagnetic fields enhanced by surface plasmons, while the spin current described by Equation (2) is undetectably small in the plain Pt/BiY₂Fe₅O₁₂ film due to the small energy transfer from external photons to magnons (see Supplementary Note 2). Note that the ratio of the voltage enhancement in Fig. 4b does not correspond to that of the $|E|^2$ enhancement due to the presence of the background heating signal.

The plasmonic spin pumping observed here can be applied to the construction of light/spin-current converters for driving spintronic devices. The plasmonic spin pumping also adds a light/voltage conversion function to spin Seebeck thermoelectric devices²², enabling hybrid voltage generation from both light and heat in a single device. The plasmonic spin pumping follows the same scaling law as the spin Seebeck effect: the output power is increased simply by extending the device area. This light/voltage conversion is conceptually different from that based on a PN junction, and is free from the output-voltage limitation caused by built-in potential. Although the spin and electric voltages generated by the plasmonic spin pumping are still small, there is plenty of scope for performance improvement; they can be enhanced by using magnetic insulators with large magnon-photon coupling, improving magnetic-insulator/NP interfaces for efficient magnon-plasmon energy transfer, and introducing plasmonic crystals that exhibit sharp and strong SPR. Plasmonic spintronics based on the interaction between surface plasmons and spin currents is still in its infancy; more detailed experimental and theoretical investigations are necessary to obtain full understanding of the mechanism of the plasmonic spin pumping and to quantitatively evaluate plasmon-magnon-electron conversion efficiency, which might be realized by removing the size, position and shape variations of Au NPs.

Methods

Preparation process of Pt/BiY₂Fe₅O₁₂/Au-NP samples. The Au NPs were formed by the following two-step procedures: fabrication of a 5-nm-thick Au thin film on a Gd₃Ga₅O₁₂ (111) substrate by using a d.c. sputtering system and heating of the Au film at 1,000 °C for 30 min (ref. 6). To increase the density of Au NPs, these procedures were repeated three times. The BiY₂Fe₅O₁₂ thin film was then formed by the MOD method²². The MOD solution for the contacted (voided) Pt/BiY₂Fe₅O₁₂/Au-NP sample includes Bi, Y and Fe carboxylates, dissolved in organic solvents with the concentration of 3% (5%). Its chemical composition is Bi:Y:Fe = 1:2:5. The solution for the contacted (voided) sample was spin-coated on the Gd₃Ga₅O₁₂ substrate with Au NPs at 5,000 r.p.m. for 60 s, followed by a drying step at 50 °C (170 °C) for 3 min and pre-annealing at 480 °C (550 °C) for 1 h (5 min). To increase the thickness of the BiY₂Fe₅O₁₂ film, the processes from spin-coating to pre-annealing were repeated three times. Then, the contacted (voided) sample was annealed at 725 °C (680 °C) for 15 h in air to form a crystallized BiY₂Fe₅O₁₂ film. The cross-sectional transmission electron microscope images in Fig. 7 indicate that a crystalline BiY₂Fe₅O₁₂ film was formed on the Gd₃Ga₅O₁₂ substrate. Finally, a 5-nm-thick Pt film was deposited over the whole surface of the BiY₂Fe₅O₁₂ film by using an rf magnetron sputtering.

Preparation process of Pt/Gd₃Ga₅O₁₂/Au-NP sample. The Pt/Gd₃Ga₅O₁₂/Au-NP sample used in Fig. 6b,d consists of a Pt/Gd₃Ga₅O₁₂ bilayer film with Au NPs embedded in the Gd₃Ga₅O₁₂ layer. After forming Au NPs on a Gd₃Ga₅O₁₂ (111) substrate with the procedures described above, an additional Gd₃Ga₅O₁₂ thin film was coated on them by means of the MOD method. The MOD solution for the Gd₃Ga₅O₁₂ film includes Gd and Ga carboxylates, dissolved in organic

solvents with the concentration of 3%. Its chemical composition is Gd:Ga = 3:5. The spin-coating and annealing processes for the Gd₃Ga₅O₁₂ film are the same as those for the BiY₂Fe₅O₁₂ film in the contacted Pt/BiY₂Fe₅O₁₂/Au-NP sample. Finally, a 5-nm-thick Pt film was deposited over the whole surface of the Gd₃Ga₅O₁₂ film. The mean distance between the Pt film and Au NPs in the Pt/Gd₃Ga₅O₁₂/Au-NP sample is \sim 100 nm, which is comparable to that in the Pt/BiY₂Fe₅O₁₂/Au-NP sample.

Numerical calculation. To investigate the electromagnetic field distributions induced by surface plasmons, we performed the numerical calculation based on the FDTD method²⁵ using KeyFDTD (Science Technology Research Institute, Japan). The simulation results shown in Fig. 3 were obtained from a model comprising a BiY₂Fe₅O₁₂ rectangular parallelepiped with the size of 150 nm \times 150 nm \times 140 nm and a Au spheroid with the in-plane diameter d_{Au} of 70 nm and the height h_{Au} of 40 nm embedded at the centre of the BiY₂Fe₅O₁₂, where the distance L between the tops of the BiY₂Fe₅O₁₂ and the Au is 50 nm. The unit cell of the model is 2 nm \times 2 nm \times 1 nm. A flat light source is placed on the top of the BiY₂Fe₅O₁₂ rectangular parallelepiped in the x - y plane, which generates plane electromagnetic waves with the polarization along the x direction and the wave vector \mathbf{k} along the z direction (note that the incident light used in the experiments is unpolarized). Here, the x , y and z directions are defined in Fig. 3a. At the outer boundaries of the BiY₂Fe₅O₁₂, we set a periodic boundary condition. Under this condition, we calculated the distribution of the electric field intensity $|E|$ while fixing the wavelength of the incident electromagnetic waves. In Fig. 3b,c, we plot the $|E|$ distribution in the BiY₂Fe₅O₁₂/Au-NP model in the x - y plane across the centre of the Au spheroid. By using the FDTD method and similar models, we have also investigated the electromagnetic coupling between the Pt film and Au NPs in the Pt/BiY₂Fe₅O₁₂/Au-NP sample, which might affect the voltage signal under the SPR condition, by comparing the simulated distributions of $|E|$ induced by the SPR between the BiY₂Fe₅O₁₂/Au-NP models with and without the Pt layer for various values of L , d_{Au} and h_{Au} (Supplementary Figs 1 and 2). In Supplementary Fig. 5, we compare the simulated $|E|$ distributions between the Pt/BiY₂Fe₅O₁₂/Au-NP models with and without a void. The refraction index and absorption-constant spectra of BiY₂Fe₅O₁₂, Au and Pt in the calculations were obtained from refs 37–39, respectively.

References

- Barnes, W. L., Dereux, A. & Ebbesen, T. W. Surface plasmon subwavelength optics. *Nature* **424**, 824–830 (2003).
- Hutter, E. & Fendler, J. H. Exploitation of localized surface plasmon resonance. *Adv. Mater.* **16**, 1685–1706 (2004).
- Raghu, S., Chung, S. B., Qi, X.-L. & Zhang, S.-C. Collective modes of a helical liquid. *Phys. Rev. Lett.* **104**, 116401 (2010).
- Appelbaum, I., Drew, H. D. & Fuhrer, M. S. Proposal for a topological plasmon spin rectifier. *Appl. Phys. Lett.* **98**, 023103 (2011).
- Di Pietro, P. *et al.* Observation of Dirac plasmons in a topological insulator. *Nat. Nanotechnol.* **8**, 556–560 (2013).
- Uchida, H., Masuda, Y., Fujikawa, R., Baryshev, A. V. & Inoue, M. Large enhancement of Faraday rotation by localized surface plasmon resonance in Au nanoparticles embedded in Bi:YIG film. *J. Magn. Magn. Mat.* **321**, 843–845 (2009).
- Melander, E. *et al.* Influence of the magnetic field on the plasmonic properties of transparent Ni anti-dot arrays. *Appl. Phys. Lett.* **101**, 063107 (2012).
- Wolf, S. A. *et al.* Spintronics: a spin-based electronics vision for the future. *Science* **294**, 1488–1495 (2001).
- Zutić, I., Fabian, J. & Das Sarma, S. Spintronics: fundamentals and applications. *Rev. Mod. Phys.* **76**, 323–410 (2004).
- Slonczewski, J. C. Current-driven excitation of magnetic multilayers. *J. Magn. Magn. Mat.* **159**, L1–L7 (1996).
- Maekawa, S., Saitoh, E., Valenzuela, S. O. & Kimura, T. (Eds). *Spin Current* (Oxford University Press, 2012).
- Tserkovnyak, Y., Brataas, A. & Bauer, G. E. W. Enhanced Gilbert damping in thin ferromagnetic films. *Phys. Rev. Lett.* **88**, 117601 (2002).
- Mizukami, S., Ando, Y. & Miyazaki, T. Effect of spin diffusion on Gilbert damping for a very thin permalloy layer in Cu/permalloy/Cu/Pt films. *Phys. Rev. B* **66**, 104413 (2002).
- Saitoh, E., Ueda, M., Miyajima, H. & Tataru, G. Conversion of spin current into charge current at room temperature: inverse spin-Hall effect. *Appl. Phys. Lett.* **88**, 182509 (2006).
- Kajiwaru, Y. *et al.* Transmission of electrical signals by spin-wave interconversion in a magnetic insulator. *Nature* **464**, 262–266 (2010).
- Uchida, K. *et al.* Observation of the spin Seebeck effect. *Nature* **455**, 778–781 (2008).
- Xiao, J., Bauer, G. E. W., Uchida, K., Saitoh, E. & Maekawa, S. Theory of magnon-driven spin Seebeck effect. *Phys. Rev. B* **81**, 214418 (2010).
- Uchida, K. *et al.* Spin Seebeck insulator. *Nat. Mater.* **9**, 894–897 (2010).
- Jaworski, C. M. *et al.* Observation of the spin-Seebeck effect in a ferromagnetic semiconductor. *Nat. Mater.* **9**, 898–903 (2010).

20. Uchida, K. *et al.* Observation of longitudinal spin-Seebeck effect in magnetic insulators. *Appl. Phys. Lett.* **97**, 172505 (2010).
21. Adachi, H., Ohe, J., Takahashi, S. & Maekawa, S. Linear-response theory of spin Seebeck effect in ferromagnetic insulators. *Phys. Rev. B* **83**, 094410 (2011).
22. Kirihaara, A. *et al.* Spin-current-driven thermoelectric coating. *Nat. Mater.* **11**, 686–689 (2012).
23. Kikkawa, T. *et al.* Longitudinal spin Seebeck effect free from the proximity Nernst effect. *Phys. Rev. Lett.* **110**, 067207 (2013).
24. Kikkawa, T. *et al.* Separation of longitudinal spin Seebeck effect from anomalous Nernst effect: determination of the origin of transverse thermoelectric voltage in metal/insulator junctions. *Phys. Rev. B* **88**, 214403 (2013).
25. Mur, G. Absorbing boundary conditions for the finite-difference approximation of the time-domain electromagnetic-field equations. *IEEE Trans. Electromagn. Compat.* **EMC-23**, 377–382 (1981).
26. Valenzuela, S. O. & Tinkham, M. Direct electronic measurement of the spin Hall effect. *Nature* **442**, 176–179 (2006).
27. Kimura, T., Otani, Y., Sato, T., Takahashi, S. & Maekawa, S. Room-temperature reversible spin Hall effect. *Phys. Rev. Lett.* **98**, 156601 (2007).
28. Huang, S. Y. *et al.* Transport magnetic proximity effects in platinum. *Phys. Rev. Lett.* **109**, 107204 (2012).
29. Ibach, H. & Luth, H. *Solid-State Physics: an Introduction to Principles of Materials Science* (Springer, 2009).
30. Papaconstantopoulos, D. A. *Handbook of the Band Structure of Elemental Solids* (Plenum Press, 1986).
31. Uchida, K. *et al.* Long-range spin Seebeck effect and acoustic spin pumping. *Nat. Mater.* **10**, 737–741 (2011).
32. Uchida, K. *et al.* Acoustic spin pumping: direct generation of spin currents from sound waves in Pt/Y₃Fe₅O₁₂ hybrid structures. *J. Appl. Phys.* **111**, 053903 (2012).
33. Elliott, R. J. & Loudon, R. The possible observation of electronic Raman transitions in crystals. *Phys. Lett.* **3**, 189–191 (1963).
34. Shen, Y. R. & Bloembergen, N. Interaction between light waves and spin waves. *Phys. Rev.* **143**, 372–384 (1966).
35. Moriya, T. Theory of light scattering by magnetic crystals. *J. Phys. Soc. Jpn* **23**, 490–500 (1967).
36. Fleury, P. A. & Loudon, R. Scattering of light by one- and two-magnon excitations. *Phys. Rev.* **166**, 514–530 (1968).
37. Doormann, V., Krumme, J.-P., Klages, C.-P. & Erman, M. Measurement of the refractive index and optical absorption spectra of epitaxial bismuth substituted yttrium iron garnet films at uv to near-ir wavelengths. *Appl. Phys. A* **34**, 223–230 (1984).
38. Johnson, P. B. & Christy, R. W. Optical constants of the noble metals. *Phys. Rev. B* **6**, 4370–4379 (1972).
39. Palik, E. D. *Handbook of Optical Constants of Solids* (Academic Press, 1985).

Acknowledgements

We thank Y. Ohnuma, T. Kikkawa, S. Daimon, A. Kirihaara, M. Ishida, T. Meyer, M. Agrawal, E. Papaioannou and B. Hillebrands for valuable discussions. This work was supported by PRESTO-JST 'Phase Interfaces for Highly Efficient Energy Utilization', CREST-JST 'Creation of Nanosystems with Novel Functions through Process Integration', Grant-in-Aid for Young Scientists (A) (25707029), Grant-in-Aid for Scientific Research (A) (24244051), Grant-in-Aid for Scientific Research on Innovative Areas 'Nano Spin Conversion Science' (26103005) from MEXT, Japan, LC-IMR of Tohoku University, NEC Corporation, the Sumitomo Foundation, the Tanikawa Fund Promotion of Thermal Technology, the Casio Science Promotion Foundation and the Iwatani Naoji Foundation.

Author contributions

K.U. and E.S. planned and supervised the study. K.U. designed the experiments, prepared the samples, collected and analysed all of the data and performed the numerical calculation. H.A. and S.M. developed the explanation of the experiments. D.K. performed the scanning electron microscopy. S.I. performed the transmission electron microscopy. Z.Q. supported the sample preparation. K.U. and H.A. wrote the manuscript with input from E.S. and S.M. All authors discussed the results.

Additional information

Supplementary Information accompanies this paper at <http://www.nature.com/naturecommunications>

Competing financial interests: The authors declare no competing financial interests.

Reprints and permission information is available online at <http://npg.nature.com/reprintsandpermissions/>

How to cite this article: Uchida, K. *et al.* Generation of spin currents by surface plasmon resonance. *Nat. Commun.* **6**:5910 doi: 10.1038/ncomms6910 (2015).



This work is licensed under a Creative Commons Attribution 4.0 International License. The images or other third party material in this article are included in the article's Creative Commons license, unless indicated otherwise in the credit line; if the material is not included under the Creative Commons license, users will need to obtain permission from the license holder to reproduce the material. To view a copy of this license, visit <http://creativecommons.org/licenses/by/4.0/>