

1 **Geologic studies of planetary surfaces using radar polarimetric imaging**
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14 **Abstract:** Radar is a useful remote sensing tool for studying planetary geology because it is
15 sensitive to the composition, structure, and roughness of the surface and can penetrate some
16 materials to reveal buried terrain. The Arecibo Observatory radar system transmits a single
17 sense of circular polarization, and both senses of circular polarization are received, which allows
18 for the construction of the Stokes polarization vector. From the Stokes vector, daughter products
19 such as the circular polarization ratio, the degree of linear polarization, and linear polarization
20 angle are obtained. Recent polarimetric imaging using Arecibo has included Venus and the
21 Moon. These observations can be compared to radar data for terrestrial surfaces to better
22 understand surface physical properties and regional geologic evolution. For example,
23 polarimetric radar studies of volcanic settings on Venus, the Moon and Earth display some
24 similarities, but also illustrate a variety of different emplacement and erosion mechanisms.
25 Polarimetric radar data provides important information about surface properties beyond what can
26 be obtained from single-polarization radar. Future observations using polarimetric synthetic
27 aperture radar will provide information on roughness, composition and stratigraphy that will
28 support a broader interpretation of surface evolution.

29

29 **1.0 Introduction**

30

31 Radar polarimetry has the potential to provide more information about surface physical
32 properties than single-polarization backscatter measurements, and has often been used in remote
33 sensing observations of Solar System objects. Many pioneering dual-polarization studies utilized
34 ground-based radio telescopes and radar systems, in part because early Earth-orbital and
35 planetary sensors were constrained in sensitivity or downlink data rate. For example, the radar
36 on the Magellan mission to Venus measured only horizontally (H) polarized radiation, except for
37 a few orbits where the spacecraft was rotated to measure the vertical (V) polarization. The
38 Cassini radar also images in a single polarization. The recent Mini-RF (radio frequency) radars
39 on Chandrayaan-1 and the Lunar Reconnaissance Orbiter, using a hybrid-polarity architecture to
40 generate the Stokes vector [Raney et al. 2010; Nozette et al. 2010], mark the first time that such
41 imaging radar data has been obtained from a planetary mission.

42 The first radar polarimetric studies of Venus employed a system that transmitted one circular
43 polarization and received two orthogonal circular polarizations [Levy and Schuster, 1964]. This
44 allowed for the first measurements of the ratio of the circular polarizations and enabled scattering
45 law comparisons. Comparing spectra of Venus in both circular polarization channels allowed the
46 identification of comparatively rough surface features that appeared at specific rotational phases
47 [Goldstein, 1965; Evans et al. 1966]. The diffuse scattering behavior appeared similar to that
48 observed for the Moon, but the quasi-specular component was stronger and exhibited a steeper
49 drop with incidence angle, indicative of smaller surface slopes [Evans et al. 1966]. Hagfors and
50 Campbell [1974] observed Venus using the 70-cm wavelength radar at Arecibo and measured
51 cross-section, circular polarization ratio, and the fraction (degree) of linear polarization in the

52 received echo versus incidence angle along the apparent rotation axis (time delay). The degree of
53 linear polarization was much lower than corresponding lunar values, which suggested less
54 surface penetration of the radar wave on Venus.

55 Studies of the Moon have also long employed polarimetry as a means to understand surface
56 composition and structure. Many experiments revealed that the total echo power received in the
57 same sense circular polarization as was transmitted is about 13 dB below that received in the
58 opposite sense circular polarization (for example, see Evans and Pettengill [1963]). Hagfors et al.
59 [1965] used the Millstone Hill 23-cm wavelength radar to investigate the nature and distribution
60 of the lunar regolith covering. A circular polarization was transmitted and two orthogonal linear
61 polarizations were received; the resulting frequency spectra and polarization ratios were
62 consistent with a tenuous tens-of-centimeters thick surface layer [Hagfors et al. 1965]. Evans and
63 Hagfors [1966] measured the backscatter behavior of linear-polarized waves using a specially
64 constructed polarizer at Millstone Hill, also at 23-cm wavelength. They found that in the diffuse
65 region of echoes (incidence angles greater than 40°) only 1/8 of the total power is returned in the
66 linear mode orthogonal to that transmitted, suggesting that either multiple reflections can occur
67 or that the echo is partially reflected from the subsurface. Zisk et al. [1987] acquired dual-
68 circular polarization data of the Moon at 3.0 cm wavelength using Haystack Observatory. These
69 data were used to study the scattering properties of the Apollo 15 landing site [Zisk et al. 1987]
70 and to measure surface properties around lunar impact craters [Campbell et al. 1988; B. A.
71 Campbell et al. 1992]. Stacy [1993] used the 12.6-cm radar system at Arecibo to search for
72 possible evidence of ice at the lunar poles [Stacy et al. 1993; 1997]. Images of Sinus Iridum and
73 Mare Imbrium were used to investigate scattering from the mare, and a comparison of the

74 fraction of linear polarized echo power to models showed that some of the backscattered power
75 must come from subsurface quasi-specular scattering [Stacy, 1993].

76 Over the past decade, improvements to ground-based telescope facilities, including a higher-
77 powered Arecibo Observatory radar, improved receivers, and new fast-sampling instruments,
78 have allowed full Stokes vector radar imaging of the inner Solar System at higher resolution than
79 was previously possible. These new capabilities provide the opportunity to use radar polarimetry
80 to study geology at local to regional scales. In particular, radar polarimetry has been used to
81 study volcanism, impact cratering, and surface properties on Venus and the Moon (e.g. Stacy et
82 al. 1997; Carter et al. 2004; Carter et al. 2006; Campbell et al. 2010; Campbell et al. 2008;
83 Thompson et al. 2006; Ghent et al. 2005; Wells et al. 2010), which are both relatively close to
84 Earth and yield high signal-to-noise data. Planetary radar data can also be compared to imaging
85 radar observations of terrestrial-analog settings to better understand the observed scattering
86 behaviors. Cross-comparisons of the available data sets demonstrate that radar polarimetry is
87 useful for discriminating between different types of geologic surfaces and mantling cover.
88 Below, we use radar data obtained and analyzed over the last several years to compare the
89 scattering behaviors of similar terrain types on Venus, the Moon, and Earth. This is the first time
90 that relatively high-resolution polarimetric radar imaging has been available for such a cross-
91 planet comparison. These comparisons elucidate the range of observed scattering behaviors and
92 identify puzzling cases that have no clear physical explanation based on current models.

93

94 **2.0 Description of polarimetric data products**

95

96 Ground-based radio telescopes are typically able to receive two orthogonal polarizations

97 simultaneously, and from this data it is possible to create the Stokes vector. The data described
98 below were obtained using the Arecibo Observatory 12.6 cm wavelength (S-band) radar system
99 as a transmitter, and either Arecibo or the Robert C. Byrd Green Bank Telescope (GBT) as the
100 receiver. The Arecibo radar transmits a circular polarization, and both Arecibo and the GBT can
101 receive two orthogonal circular polarizations, referred to as the same-sense circular (SC) and
102 opposite-sense circular (OC) polarizations (to that transmitted). Fully polarimetric observations
103 were not used because typically the radar transmits continuously. In the past, when pulsed
104 waveforms were used, suitable switches that would permit transmit polarization switching were
105 not available at the power levels used (typically 100's of kilowatts).

106 In the case of Venus, the beam of the Arecibo telescope at S-band is two arcminutes, which
107 is about twice the angular size of Venus at closest approach to Earth. To reduce the resulting
108 north-south delay-Doppler ambiguity problem, we pointed north and south of the planet on
109 alternating radar runs to allow the central portion of the beam to preferentially illuminate one
110 hemisphere. The resulting images therefore contain echo power from the entire Earth-facing
111 hemisphere of Venus. The Venus data were mapped to a Mercator projection that shows either
112 the north or south (depending on the telescope pointing), with areas near the Doppler equator
113 excluded. The incidence angle variation across the surface is due to both the curvature of the
114 spherical planet and to changes in topography. However, when looking at the global maps at low
115 (~12 km) resolution, the largest changes in incidence angle are from the curvature of the planet.
116 For the lunar observations, the Arecibo beam subtends only a small portion of the lunar surface,
117 and there is no ambiguity except in areas near the Doppler equator. Incidence angle variation
118 within a given lunar image can be dominated by local topography, although this depends on the
119 particular sub-radar point and the amount of local topographic variation.

120 In both the Venus and lunar cases, the received data are converted into complex-valued
 121 delay-Doppler images, and a relative channel balance is applied based on background noise
 122 measured off-planet. The delay-Doppler maps are converted into latitude/longitude format,
 123 which includes focusing for the lunar case [Campbell et al. 2007]. The maps are then converted
 124 into real-valued Stokes vector images, and daughter products such as the circular polarization
 125 ratio are computed. Additional discussion of observational parameters and data processing can
 126 be found in Carter et al. [2004] and Campbell et al. [2010]. Below, we describe the daughter
 127 products in more detail.

128 The Stokes vector generated from the received circular polarizations can be used to
 129 completely describe the polarization state of the received wave [Jackson 1999]:

130

$$131 \quad S = \begin{bmatrix} S_1 \\ S_2 \\ S_3 \\ S_4 \end{bmatrix} = \begin{bmatrix} \langle |E_L|^2 \rangle + \langle |E_R|^2 \rangle \\ 2\text{Re}\langle E_L E_R^* \rangle \\ 2\text{Im}\langle E_L E_R^* \rangle \\ \langle |E_L|^2 \rangle - \langle |E_R|^2 \rangle \end{bmatrix} \quad (1)$$

132

133 where E_L and E_R are the electric fields for the left and right circular polarizations, respectively,
 134 and the averages are time or spatial averages. In practice, Eqn. 1 is used to derive the Stokes
 135 parameters using fully processed OC and SC complex-valued images. The first Stokes parameter
 136 (S_1) is a measure of the total average power in the echo. The S_2 and S_3 Stokes parameters
 137 describe the linearly polarized state of the wave. The S_4 Stokes parameter gives the direction
 138 and magnitude of the circularly polarized power.

139 There are two particularly useful daughter products that can be derived from the Stokes
 140 vector: the circular polarization ratio ($\mu_c = \text{SC}/\text{OC}$) and the degree of linear polarization (m_l). The

141 circular polarization ratio can be used as an indicator of surface roughness. Specular echoes
142 from surfaces that are smooth at wavelength scales will lead to low ratios, while diffuse
143 scattering from rough surfaces generates μ_c values approaching one, or even greater than one for
144 extremely rough terrain and low temperature water ice when it is present. The circular
145 polarization ratio can be calculated from:

146

$$147 \quad \mu_c = \frac{S_1 - S_4}{S_1 + S_4} \quad (2)$$

148

149 The circular polarization ratio will tend to increase with incidence angle because near-nadir
150 scattering is primarily quasi-specular (i.e., dominantly OC). At larger incidence angles, diffuse
151 scattering generally contributes a greater fraction of power to the received echo, which leads to a
152 higher SC/OC ratio.

153 For a circular-polarized transmit and receive system such as Arecibo, sources of error in
154 the circular polarization ratio are dominated by uncertainties in the relative channel gains of the
155 OC and SC channels. In most cases, the noise background can be measured from areas of noise
156 in off-planet areas of each image. It is important that the dynamic range of the data is well
157 captured by the quantization; for example 4-8 bit sampling has proved better than 2-bit sampling
158 for planetary targets with a large range in echo power (e.g. Moon). Otherwise it may not be
159 possible to obtain accurate measurements of the noise floor in each channel. During processing,
160 the data is divided by the noise measurements to balance the channels. Typically, and for the data
161 shown below, a very good channel balance can be achieved, and the resultant uncertainty in μ_c
162 values is small (a few percent).

163 The linear-polarized component of the received echo can be used to infer the presence of

164 subsurface scattering. A circularly polarized incident wave can be thought of as a combination of
 165 two orthogonal linear vectors that are vertically (V) and horizontally (H) polarized with respect
 166 to the plane of incidence. These two components have different power transmission coefficients
 167 [Jackson, 1999]:

168

$$169 \quad T_H(\theta, \epsilon') = \frac{4 \cos \theta \sqrt{\epsilon' - \sin^2 \theta}}{(\cos \theta + \sqrt{\epsilon' - \sin^2 \theta})^2} \quad (3)$$

$$170 \quad T_V(\theta, \epsilon') = \frac{4 \epsilon' \cos \theta \sqrt{\epsilon' - \sin^2 \theta}}{(\epsilon' \cos \theta + \sqrt{\epsilon' - \sin^2 \theta})^2} \quad (4)$$

171

172 where θ is the angle between the surface normal and the incoming radar wave (incidence angle)
 173 and ϵ' is the real component of the dielectric constant. If this wave penetrates the surface, the V
 174 polarization will be preferentially transmitted, and the polarization state will change from
 175 circular to elliptical. The reflected wave exiting the surface towards the radar will experience a
 176 similar preferential transmission of the V polarization component. The elliptical polarization can
 177 be thought of as a combination of circularly polarized power and a linear component with a
 178 measurable magnitude and direction. The degree (or percent) of linear polarization is:

179

$$180 \quad m_l = \frac{\sqrt{S_2^2 + S_3^2}}{S_1}. \quad (5)$$

181

182 The long axis of the polarization ellipse indicates the direction of the linear polarization vector.
 183 The angle of this major axis, with respect to H polarization, is given by:

184

$$185 \quad \chi = \frac{1}{2} \arctan\left(\frac{S_3}{S_2}\right). \quad (6)$$

186

187 For subsurface scattering, the direction of the linear polarization will be parallel to the plane of
188 incidence and reflection (a plane that includes the direction of the incident or reflected wave and
189 the normal to the surface). The linear polarization angle will change as the surface tilts in the
190 azimuth direction with respect to the radar. Measurements of the linear polarization angle can
191 assist with interpretation of the degree of linear polarization; scattering from the subsurface
192 should produce linear polarization angles that vary across the scene as the local surface normal
193 varies with topography and the curvature of the planet. Linear polarization angles have also been
194 used to measure local slopes [Stacy 1993; Stacy and Campbell 1993].

195 There are three primary sources of error in measurements of the degree of linear polarization
196 as measured from a circular-polarized transmit and receive system: instrumental cross-coupling
197 between the circular polarizations, statistical fluctuations in the S_2 and S_3 images that create a
198 spurious degree of linear polarization when they are squared and added, and errors in subtracting
199 the noise background in the S_1 image. These errors are discussed in detail in Carter et al. [2004].
200 Instrumental cross-coupling at S-band is low for both the Arecibo and the Green Bank Telescope
201 systems, and contributes a spurious degree of linear polarization of ~ 0.04 . The statistical
202 background in the S_2 and S_3 images depends on the amount of spatial averaging in a given image,
203 but the added spurious degree of linear polarization ranges from 0.02 to 0.04 in the Venus data.
204 Errors in the S_1 background power measurement usually range from a few to several percent,
205 which results in an added spurious linear polarization of 0.02 for a degree of linear polarization
206 value of 0.35.

207 The value of m_l depends on surface physical properties (the dielectric constant and relative
208 amounts of surface vs. subsurface scattering), the radar viewing geometry with respect to the
209 surface topography (incidence angle), and the amount of spatial averaging across surfaces with
210 differing amounts of radar penetration. The degree of linear polarization increases with incidence
211 angle and is zero at normal incidence. The m_l value will also be higher if the surface permittivity
212 (ϵ') is higher, because the transmission coefficients lead to a greater difference in H and V
213 transmitted power. However, an increase in permittivity also increases the fraction of power that
214 is reflected directly from the surface and reduces the amount of power transmitted into the
215 surface. For terrestrial planet surfaces, the viewing geometry is well understood. However, the
216 spatial variation in terrain within a pixel, the dielectric properties, and the amount of power
217 reflected from the surface and subsurface, are generally unknown. Other data sources, such as
218 higher resolution optical images, and laboratory measurements of dielectric properties, can be
219 used to estimate reasonable values in some cases.

220 Changes in μ_c and m_l can be used to infer differences in the surface and subsurface materials
221 and structure. As described by Stacy [1993], the backscattered radar wave can be modeled using
222 a simple single layer model with four components; surface quasi-specular scattering (σ_{qs}),
223 surface diffuse scattering (σ_d), subsurface quasi-specular scattering (σ'_{qs}) and subsurface diffuse
224 scattering (σ'_d). The power returned from the subsurface depends on the depth to any subsurface
225 scatterers, their total cross section, and the loss tangent of the medium. The absolute backscatter
226 values also depend upon the incidence angle, the dielectric constant of the surface material, and
227 the surface roughness properties.

228 It is possible to compare in a relative sense the expected μ_c and m_l values for different values
229 of the four backscatter components, assuming that the viewing geometry and dielectric properties

230 of the materials are similar. This approach leads to some endmember cases shown in Fig. 1,
231 where the left side shows smooth surfaces with various types of subsurface scattering, and the
232 right side shows cases where the surface echo comes primarily from diffuse scattering on
233 surfaces that are rough at the wavelength scale.

234 Although real surfaces will be some combination of these cases, the endmembers provide a
235 basis for understanding how μ_c and m_l can change in different geologic settings. For example,
236 areas that have a low abundance of wavelength-scale subsurface scatterers within the penetration
237 depth of the radar, or that have a gradual change in dielectric constant with depth, will have a
238 low degree of linear polarization (Figs. 1a, 1d). Very rough crater walls and ejecta will have a
239 high μ_c value, but a low m_l value (Fig. 1d). Fine-grained deposits (e.g. impact ejecta and
240 ash) that contain embedded rocks will produce a high m_l value (Fig. 1c), and they may also have
241 a high μ_c value if a substantial amount of the echo power comes from subsurface diffuse
242 scattering off blocks. It is usually not possible to distinguish between these schematic cases
243 solely based on the polarization measurements. For example, it can be difficult to distinguish
244 between different types of subsurface scattering (e.g. between Figs. 1b and 1c, or Figs. 1e and 1f)
245 unless there is other evidence to infer a buried near-planar interface or buried rocks. For
246 example, it may be possible to determine from other remote sensing techniques that there is a
247 buried lava flow that would lead to cases shown in Fig. 1b or 1e. But without this information, it
248 can be impossible to differentiate between similar endmembers, particularly because attenuation
249 in the surface layer is unknown.

250 In cases where the measured μ_c and m_l values can be combined with information about the
251 geologic setting derived from radar, optical, and infrared images, it is sometimes possible to
252 make reasonable assumptions about the types of scattering present and thus derive a physical

253 model. Such models can be used to compute dielectric constants, make depth estimates for
254 mantling deposits, investigate the amount of surface vs. subsurface scattering, and search for
255 areas that are rock poor at the wavelength scale [Stacy 1993; Campbell and Hawke 2005;
256 Campbell et al. 2008]. However, in many cases there are simply too many uncertainties to
257 derive unique quantitative values through modeling. This is particularly true for the Venus data
258 shown below, where the resolution is of order ten kilometers, and it may not be a good
259 assumption that each pixel has a uniform surface type. In these cases, a qualitative comparison
260 between geologic units can still provide valuable information that cannot be obtained using
261 single polarization radar data.

262

263 **3.0 Earth**

264

265 Polarimetric radar observations of planetary analogs are useful for understanding the end-
266 member scenarios described above. The NASA/JPL AIRSAR system collects the full 4x4 Stokes
267 scattering operator [van Zyl et al. 1987; Campbell et al. 2004], which can be used to derive the
268 daughter products discussed above. In this section, we focus on an analysis of the degree of
269 linear polarization because it provides useful context for explaining some of the polarization
270 behavior explored on Venus. All the study areas contain little to no vegetation, since tree trunks
271 and shrubs will change the polarization signature.

272 In many AIRSAR images examined to date, a high degree of linear polarization appears to
273 be produced when the radar wave travels into a smooth surficial layer and reflects from a buried
274 interface that is also smooth at the wavelength scale. One such example is smooth sand deposits
275 in the Stovepipe Wells region of Death Valley, CA [Carter et al. 2004], where the radar wave

276 penetrates a few centimeters of sand to reflect from a buried mud-cracked surface. Another
277 example is the Lunar Lake playa deposit in Nevada. Fig. 2 shows 24-cm AIRSAR data for Lunar
278 Lake. The playa has a silty-clay surface with mud cracks and localized gravel patches [Greeley et
279 al. 1997]. The radar data show high m_i values (up to 0.6) across most of the playa, and very low
280 m_i values from the surrounding rough mountains. In this case, the radar wave is likely
281 penetrating the silt surface coating and reflecting from a buried horizon with a higher dielectric
282 constant than the silt.

283 Although the geologic setting is different, a similar physical scenario (reflection from a
284 smooth buried interface) appears to occur in areas of smooth, ponded flows in and around the
285 Kilauea caldera in Hawaii. In this case, near-surface air gaps beneath the flows provide the
286 smooth and nearly continuous subsurface interface that allows for a strong subsurface return
287 [Carter et al. 2006]. Although the permittivity of the basalt lava flows is high, a 68-cm radar
288 wave is able to penetrate a few centimeters and reflect from the boundary between the flow and
289 internal air gaps. The dielectric contrast between the lava and air is high, and the upper and basal
290 interfaces involved are very smooth at the wavelength scale.

291 One terrestrial example that may involve a scenario more like Fig. 1c is an area near Sunset
292 Crater, AZ [Carter et al. 2006] where a high degree of linear polarization is associated with a few
293 specific radar-dark cinder and ash deposits. Field observations of one such cinder cone reveal
294 centimeter-scale cinders as well as meter-scale blocks covering the ground and partially buried in
295 a soil and ash matrix that slopes smoothly away from the cinder cone. Larger blocks are located
296 closer to the base of the cone, where the degree of linear polarization is higher. In this instance,
297 the radar wave may reflect from the large blocks buried in the soil matrix.

298 These terrestrial examples demonstrate that in all of the currently identified cases where a
299 significant degree of linear polarization is observed, the upper surface is smooth at the
300 wavelength scale. A detectible degree of linear polarization value also appears to require
301 abundant subsurface reflectors with relatively strong dielectric contrast with the covering layer or
302 surrounding medium.

303

304 **4.0 Venus**

305

306 Arecibo observations from 1999 through 2004 produced images of the hemisphere of Venus
307 visible from Earth at inferior conjunction, at resolutions of 12 and 16 km. These regional data
308 show that areas with a high degree of linear polarization are concentrated in discrete areas, or
309 “features”. These features correspond to impact crater ejecta, lava flows, dome fields, and
310 aeolian settings [Carter et al. 2004; Carter et al. 2006]. In this section, we include a new analysis
311 of the circular polarization ratio to better understand the types of surfaces that are present in
312 volcanic and impact cratering settings.

313 Distal impact crater ejecta deposits, including parabola-shaped features surrounding some
314 craters [D. B. Campbell et al. 1992], account for many of the local increases in the degree of
315 linear polarization. In most cases, these correspond to areas of slightly increased radar
316 backscatter, as is the case with the Carson crater parabola [Carter et al. 2004]. In these cases, the
317 radar bright terrain mostly likely reflects subsurface scattering from larger blocks embedded in
318 fine-grained ejecta material (e.g. Fig. 1c, 1e, or 1f).

319 In other cases, such as the radar dark halos surrounding the craters Galina (Fig. 3) and Shih
320 Mai-Yu, high m_1 values correspond to areas with a low backscatter cross section. On the Moon,

321 radar-dark halos with low circular polarization ratios surround young craters, and are caused by a
322 zone of fine-grained pulverized material that falls near the impact site [Ghent et al. 2005; 2010].
323 Over time, these radar-dark zones become mixed with larger blocks, and the distinctive low
324 circular polarization ratio ring disappears [Ghent et al. 2005]. As can be seen in Fig. 3, the
325 Galina halo has a low μ_c , as is the case for the lunar craters. In the case of these dark halo
326 deposits on Venus, the radar wave may travel into a smooth, fine-grained deposit and reflect
327 from a gently undulating, buried interface (e.g. Fig. 1b). If large quantities of buried blocks were
328 present, it is likely that a higher circular polarization ratio would be observed.

329 Some volcanic settings also show an enhanced degree of linear polarization, including both
330 shield fields and lava flows. Fig. 4 shows the degree of linear polarization (red) and the circular
331 polarization ratio (green) overlaid on a Magellan image of Sif and Gula Montes. It is clear from
332 this image that different lava flows have different physical properties. Radar dark lava flows on
333 the flank of Sif Mons show a higher degree of linear polarization than most other flows in the
334 vicinity [Carter et al. 2006]. It is possible that in these cases, gaps within the flows may be
335 causing the increased degree of linear polarization values, similar to the Kilauea example
336 discussed in Section 3.0 (Fig. 1b). Radar bright flow complexes north of Gula Mons have high
337 circular polarization ratio values, indicating that they likely have a surface that is mostly rough
338 with little appreciable surface coating (Fig 1d). Even within these bright flows, however, it is
339 possible to see color variations due to the changing polarization values that indicate changing
340 surface roughness and radar penetration. The flanks of Sif Mons have a greater area of increased
341 linear polarization than the flanks of Gula Mons, suggesting that there is more overall surface
342 mantling across Sif Mons. It is not possible to be certain what causes this difference, but there

343 may be fine volcanic material across much of Sif Mons, or perhaps recent cratering events on or
344 east of Sif Mons deposited impact ejecta across the edifice.

345 The highest degrees of linear polarization are confined to local areas, but we detect a
346 measurable (i.e., above the expected spurious few percent due to calibration uncertainties)
347 amount of linearly polarized echo across most of the venusian surface. Perhaps the best
348 illustration of this is the linear polarization angle maps. Figures 5 and 6 show the northern and
349 southern hemisphere of Venus, respectively, from data acquired in 2001. The degree of linear
350 polarization values are low near the sub-radar point where the incidence angle is small, and
351 increase towards the limbs as the incidence angle increases. Close to the sub-radar point where
352 degree of linear polarization values are expected to be very small (Eqns. 3 and 4), it can be
353 difficult to discern whether subsurface scattering still occurs and whether specific features are
354 present. Non-random values in the polarization angle maps demonstrate that even at low
355 incidence angles, some fraction of the echo power is returned from the subsurface, and that the
356 slight regional increases in the degree of linear polarization are associated with specific geologic
357 structures.

358 The linear polarization angle rotates across the surface of the planet as the plane of incidence
359 and reflection changes across the sphere, as discussed in Section 2. The discrete surface features
360 visible in the degree of linear polarization maps have consistent linear polarization angle values,
361 and a quantitative but relative comparison of the χ values for different features demonstrates that
362 the linear polarization angles rotate in step with the change in the plane of incidence caused by
363 the planetary curvature [Carter et al. 2006].

364 The maps in Figs. 5 and 6 show that large portions of the Venus surface have some
365 component of subsurface scattering at 12.6 cm wavelength. This suggests that while large and

366 continuous surficial deposits occur mostly in local areas on Venus, there may be a broadly
367 distributed background of patchy, penetrable surface materials such as dust or regolith. These
368 deposits could be very thin (cm-scale), so it is not clear that the data imply significant additional
369 fine-grained surface materials. However, it does suggest that patchy coatings may be more
370 common than thought based on analysis of Magellan radar images, which cannot distinguish
371 subtle variations in such thin surface coatings from minor changes in surface roughness.

372 Many of the high degree of linear polarization features, including the high radar reflectivity
373 summits of Tepev and Theia Montes, correspond fully or partially to areas with lower than the
374 Venus-average emissivity value of 0.84 [Pettengill et al. 1992]. Low emissivity areas have
375 generally been attributed to an increase in the dielectric constant [Pettengill et al. 1992], although
376 volume scattering in a low-loss medium could also be responsible [Tryka and Muhleman 1992].
377 A higher surface dielectric constant can lead to a larger m_1 value (Section 2), but it can also
378 increase the surface echo component and decrease the fraction of the echo that comes from
379 subsurface scattering. A small decrease in the Magellan emissivity could be caused by a change
380 in surface roughness, which may explain the low emissivity and high m_1 correlations in the case
381 of distal crater ejecta. If the surface is smooth, then the Magellan data measure the true H-
382 polarized emissivity, which is less than that of a rougher surface [Campbell et al. 1994].

383 In some instances, particularly Stuart crater and the Tepev and Theia Montes summits, the
384 emissivity values are too low to be explained through surface roughness variations. The floor of
385 Stuart crater has an emissivity of 0.69 and a high degree of linear polarization (Fig. 7). Stuart is
386 the only example where high m_1 values are associated with the interior of a crater; the rough
387 surface of the crater floor and walls should generally preclude any strong echo from subsurface
388 scattering. Indeed, the circular polarization ratio image demonstrates that the interior of the crater

389 is rough. The combination of polarization behaviors suggests a scenario like Fig. 1e or 1f.
390 However, it is still not clear how an area with a significant surface cover of low density material
391 could have such a low emissivity. In the case of the high-reflectivity, low-emissivity summit
392 regions, perhaps the radar wave is able to penetrate into high dielectric materials in small areas
393 with a very smooth surface texture. Alternatively, there may be thin, unresolved patches of
394 surface coatings across some parts of the high-dielectric-constant regions. Higher-resolution
395 polarimetric imaging or emissivity, radar images at a different wavelength, or surface images,
396 would all help to differentiate between these cases.

397

398 **5.0 Moon**

399

400 Lunar data at both S- and P-band wavelength (12.6 cm and 70 cm, respectively) have been
401 obtained using the Arecibo Observatory radar transmitter and the Green Bank Telescope as a
402 receiver [Campbell et al. 2010; Campbell et al. 2007]. These data sets have a much higher
403 resolution (20-80 m/pixel at S-band) than the Venus data, and therefore have the potential to
404 show finer details in areas of geologic interest.

405 Preliminary work using the lunar Stokes vector data has focused on volcanic terrains. The
406 Aristarchus region of the Moon is an uplifted plateau that has been the site of extensive
407 volcanism. A large rille emanates from the Cobra Head source vent, and the plateau is mantled in
408 thick pyroclastic deposits that cover an area of ~49,000 km [Gaddis *et al.*, 2003]. S-band images
409 (Fig. 8) show radar-dark, lobate, slumping terrain surrounding the head of the rille. Radar-bright
410 streaks cross some areas of the pyroclastic deposit and mark deposits of blocky material from the
411 Aristarchus impact. In corresponding P-band images, some areas of the deposit have a higher

412 radar backscatter, and Campbell et al. [2008] propose that the pyroclastic material is covering
413 lava flows that lie, at most, 15 m below the surface. The same areas are slightly brighter at S-
414 band as well, probably because impacts have penetrated the deposit and mixed blocky material
415 into the pyroclastics [Campbell et al. 2008].

416 Images of the circular polarization ratio, degree of linear polarization, and linear polarization
417 angle are shown in Fig 8, from data acquired at 40 m/pixel single-look resolution. The degree of
418 linear polarization and linear polarization angle data were averaged to 800 m/pixel resolution to
419 reduce speckle before being overlaid on the total power image (Fig. 8c, d). The pyroclastic
420 materials are radar dark and have a very low circular polarization ratio, most likely because they
421 are smooth, fine-grained, and block free. The circular polarization ratio overlay clearly highlights
422 the difference between fine-grained pyroclastics and the higher circular polarization ratio impact
423 crater ejecta.

424 The degree of linear polarization image, however, shows little correlation with the circular
425 polarization ratio image; the high CPR streaks do not uniformly match with high or low m_l
426 values, and the buried lava flow shown in Campbell et al. [2008] is not visible in the linear
427 polarization data. The crater Aristarchus has a low degree of linear polarization, as might be
428 expected for a rough, blocky surface, but other craters, including the large crater Herodotus, have
429 average m_l values. Aristarchus has a radar-dark halo in 70-cm wavelength ground-based radar
430 images [Ghent et al. 2005] that can also be seen in the S-band (Fig. 8) data, although it is
431 somewhat subdued. However, there is no corresponding increase in the degree of linear
432 polarization in the S-band data (Fig. 8c), as is seen for some of the Venus craters. The
433 pyroclastic deposit itself has a degree of linear polarization value of around 0.14, which is
434 identical to the surrounding mare basalts. Fig. 8c has fairly uniform values across large portions

435 of the scene, with some of the minor fluctuations probably due to statistical errors of ~ 0.04 due
436 to S2 and S3 noise and other errors as discussed in Section 2.

437 The linear polarization angle image shown in Fig. 8d covers a much smaller area than the
438 Venus maps in Figs. 5 and 6, and therefore does not show large angle changes due to the
439 curvature of the Moon. Instead, the image shows that the value of χ changes as the plane of
440 incidence and reflection changes with topography around crater rims. The linear polarization
441 angle image also shows a large area to the northeast with different angle values (colored red),
442 which probably represents a large plateau-forming block that has been tilted along a fault.

443 The uniform value of degree of linear polarization across the Aristarchus pyroclastic suggests
444 that the radar wave may penetrate into the surface and reflect from buried objects almost
445 everywhere in the image. Other lunar data, such as an area near the crater Focas on the western
446 limb of the Moon [Campbell et al. 2010], and the Cauchy dome field in Mare Tranquillitatis, also
447 show a nearly uniform degree of linear polarization across the image despite the presence of
448 various-aged impact craters and volcanic structures. In the case of the Moon, the ubiquitous
449 regolith covering may provide an adequate medium for a significant component of subsurface
450 scattering of a 12.6-cm wave in most situations, with the absolute value of the backscattered
451 power modulated by the volume population of wavelength-scale rocks within the probing depth
452 of the signal.

453 While the change in degree of linear polarization values across the lunar images is less
454 pronounced than for Venus, the minimum and maximum m_1 values at a given incidence angle are
455 roughly similar. For example, parts of the Aristarchus pyroclastic (image center $\theta=53^\circ$) have m_1
456 values of about 0.05, and the radar bright Aristarchus ejecta have m_1 values of 0.09, so there are
457 some localized patches where little echo power is returned from the subsurface. At higher

458 incidence angles, such as the Focas crater example (center $\theta=88^\circ$) [Campbell et al. 2010], the
459 average degree of linear polarization reaches values of up to 0.3, which is similar to some
460 examples on Venus. At S-band wavelength, the major difference between the two bodies
461 appears to be the higher “background” fraction of subsurface scatter from the lunar regolith.

462

463 **6.0 Discussion and Conclusions**

464

465 Radar data processed using the dual-polarimetry Stokes parameter technique demonstrate
466 some clear differences among planetary surfaces. In the case of the Moon, a thick regolith
467 covering with suspended rocks appears to generate a fairly uniform degree of linear polarization.
468 In contrast, on Venus, and in terrestrial cases like those discussed in Section 3, the degree of
469 linear polarization is typically low for most surfaces, and is enhanced only for certain well-
470 defined features such as layered lava flows and smooth dust- or sand-covered areas. On both the
471 Moon and Venus, the circular polarization ratio varies significantly as the surface roughness and
472 subsurface block abundance change between local geologic units.

473 Since the local behavior of the degree of linear polarization is so different for the Moon and
474 Venus, it is interesting to compare the incidence angle trends of the degree of linear polarization
475 to determine the global behavior of this parameter. A simple model for the degree of linear
476 polarization assumes a single-layer model like that shown in Fig. 1b, with a fine-grained surface
477 coating overlying a quasi-specular subsurface scattering layer. In this case, the degree of linear
478 polarization depends solely on the transmission coefficients in Eqns. 3 and 4 [Stacy 1993]. To
479 some extent, this model provides an upper limit, because power loss through the surface layer,
480 and increased diffuse scattering from the surface or subsurface, will both lower the degree of

481 linear polarization values across all incidence angles. Fig. 9 shows the predicted scattering
482 behavior for this model, which predicts high m_l values at moderate to large angles of incidence.

483 The actual behaviors for the Moon and Venus, also plotted in Fig. 9, shows that the
484 background degree of linear polarization values averaged across Venus are fairly low out to 70°
485 incidence. Areas near the edges of the Venus images have increased statistical and systematic
486 errors due to decreased signal-to-noise and are not plotted, but m_l values averaged over large
487 areas are still less than the ~ 0.3 predicted for low-dielectric constant materials (e.g. see Figs. 5
488 and 6). Average m_l values for individual Venus features are larger than the background,
489 presumably because a larger fraction of the surface is covered in mantling material; however, the
490 m_l values are still low at higher incidence angles. For example, an area near the crater Xantippe
491 (62° incidence) has an average value around 0.25. The highest m_l values measured for individual
492 pixels within specific features are still lower than 0.35. The Moon displays a similar behavior,
493 although the polarimetric analysis is still in early stages and so only points from a few analyzed
494 scenes are available. At high incidence angles, for example near-limb areas near Mare Orientale
495 and Focas crater ($\theta \sim 80\text{-}85^\circ$), the average degree of linear polarization values are ~ 0.25 .

496 It is clear that the simple model shown in Fig. 9 does not entirely explain the scattering
497 regime leading to high m_l values. At moderate incidence angles ($30^\circ\text{-}40^\circ$), the model roughly
498 matches the average measured degree of linear polarization values, but as the incidence angles
499 exceed 50° , the model is less reasonable. The average permittivity of the lunar regolith is 2.8
500 [Carrier et al. 1991], and at high incidence angles, the measured degree of linear polarization
501 values would suggest a surface material with an even lower permittivity. Instead, changes in the
502 relative amounts of surface and subsurface scattering, and incidence angle related changes in the

503 relative amount of quasi-specular and diffuse scattering, likely contribute to the incidence angle
504 behavior of the measured fraction of linear-polarized echo power.

505 In summary, radar polarimetry has the capability to better distinguish between different types
506 of surface and subsurface physical properties than single-polarization radar imagery, and can be
507 useful for comparing scattering regimes across planetary objects. Future modeling using high-
508 resolution data, such as the terrestrial and lunar data sets, will help to better understand the types
509 of surfaces that produce different combinations of polarimetric behaviors. Additional
510 comparisons with radiometric data will be useful to better understand why the high degree of
511 linear polarization features on Venus are often correlated with low microwave emissivity. SAR
512 systems designed for future planetary missions will benefit greatly from including polarimetric
513 capabilities.

514

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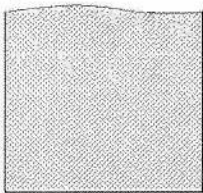
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Surface and Subsurface Scattering Scenarios

Smooth surfaces

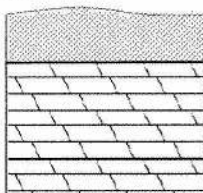
$$(\sigma_d \rightarrow 0)$$

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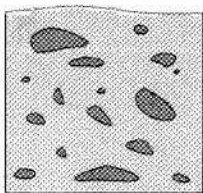
low μ_c ; low m_l
no subsurface
scattering

b



low μ_c ; high m_l
Quasi-specular
scattering from sub-
surface gives high m_l

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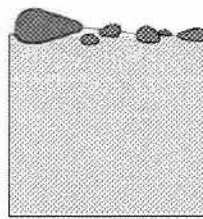


mod.-high μ_c ;
mod.-high m_l
Diffuse subsurface
scattering generates
lower m_l than case b

Rough surfaces

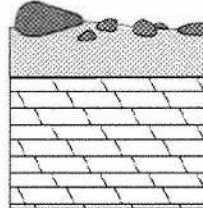
$$(\sigma_{qs} \rightarrow 0)$$

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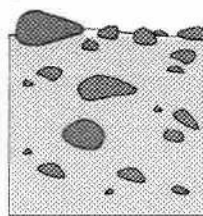
high μ_c ; low m_l
no subsurface
scattering

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high μ_c ; low-mod. m_l
Radar wave may not
penetrate into rough
surface.

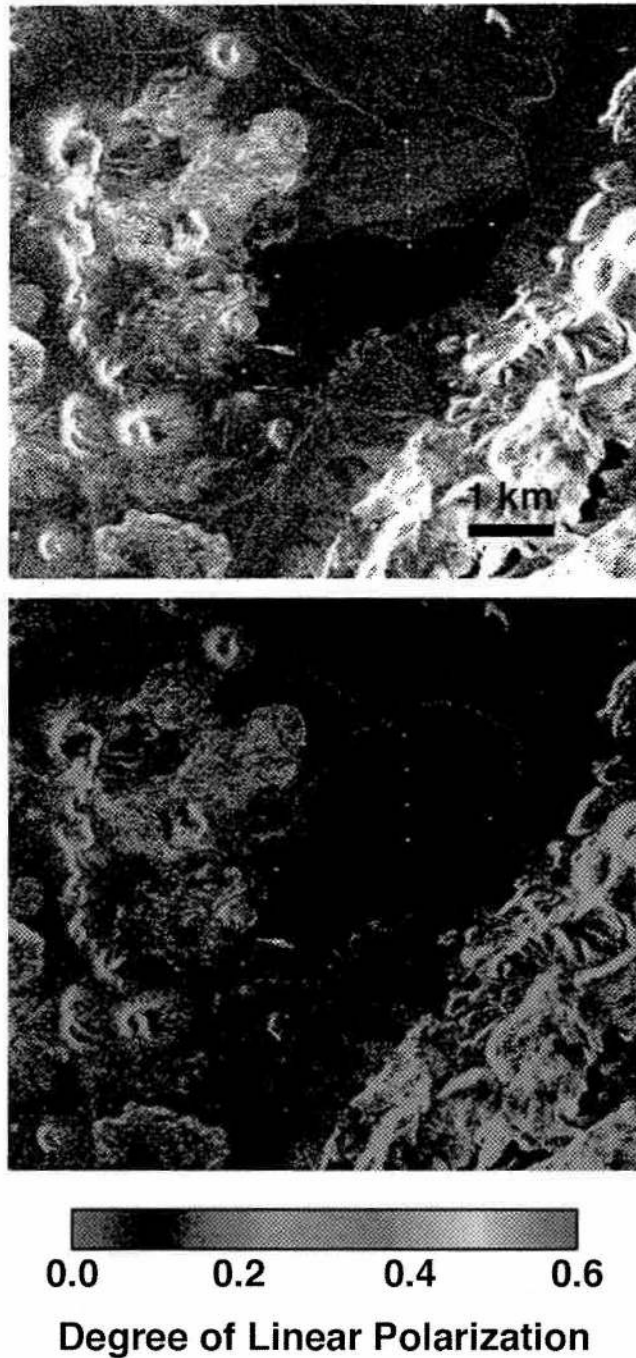
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high μ_c ; low-mod. m_l
Radar wave may not
penetrate into rough
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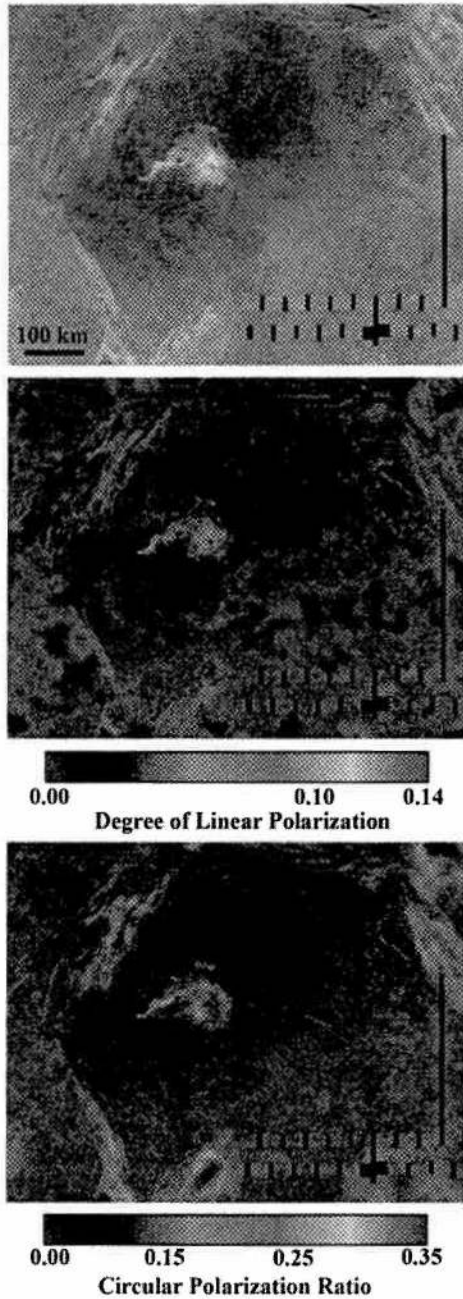
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Fig 1: Scattering endmember cases for smooth and rough surfaces, and smooth and rough subsurface reflectors. Smooth surfaces (left column) are more likely to allow for surface penetration of the radar wave. In the case of rough surfaces, the fraction of the received echo that comes from the subsurface will partially depend on the extent of the wavelength scale surface roughness. In all cases, to produce a linearly polarized echo component, the loss tangent of the mantling layer or medium must be low enough to allow the radar wave to travel far enough to reflect from buried scattering surfaces.



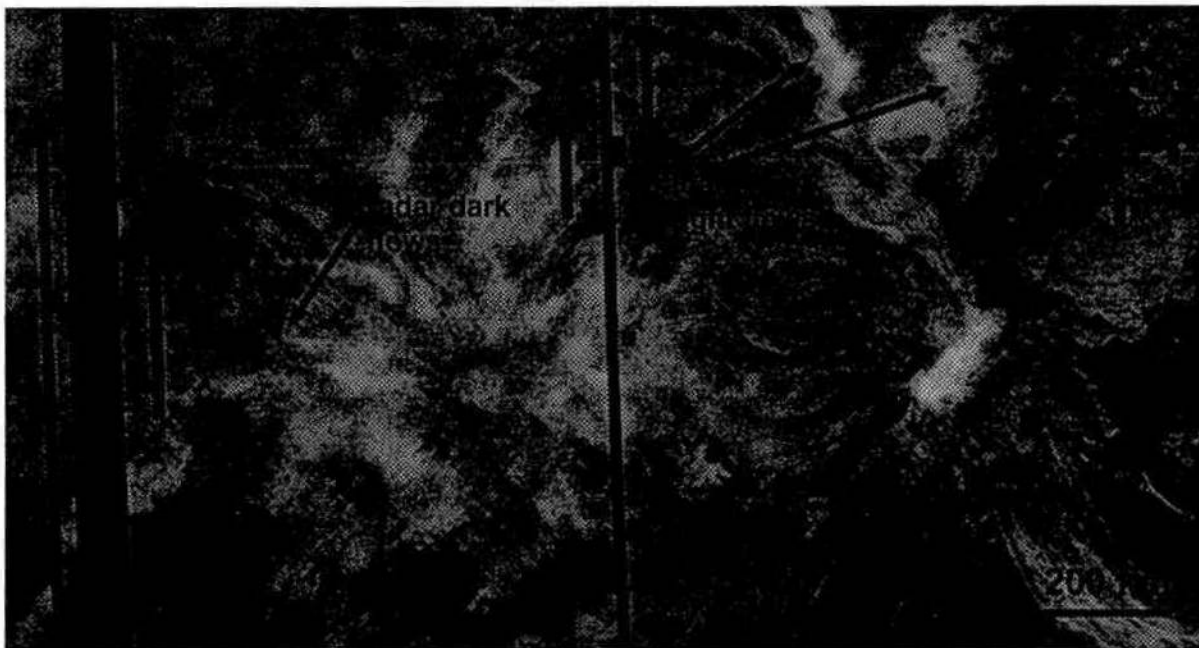
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Fig. 2: AIRSAR L-band data (24 cm wavelength) of the Lunar Lake playa, in the Lunar Crater volcanic field in Nevada. North is toward the top of the images. Top: Total power image. Bottom: Corresponding color overlay of the degree of linear polarization showing high values associated with the playa.



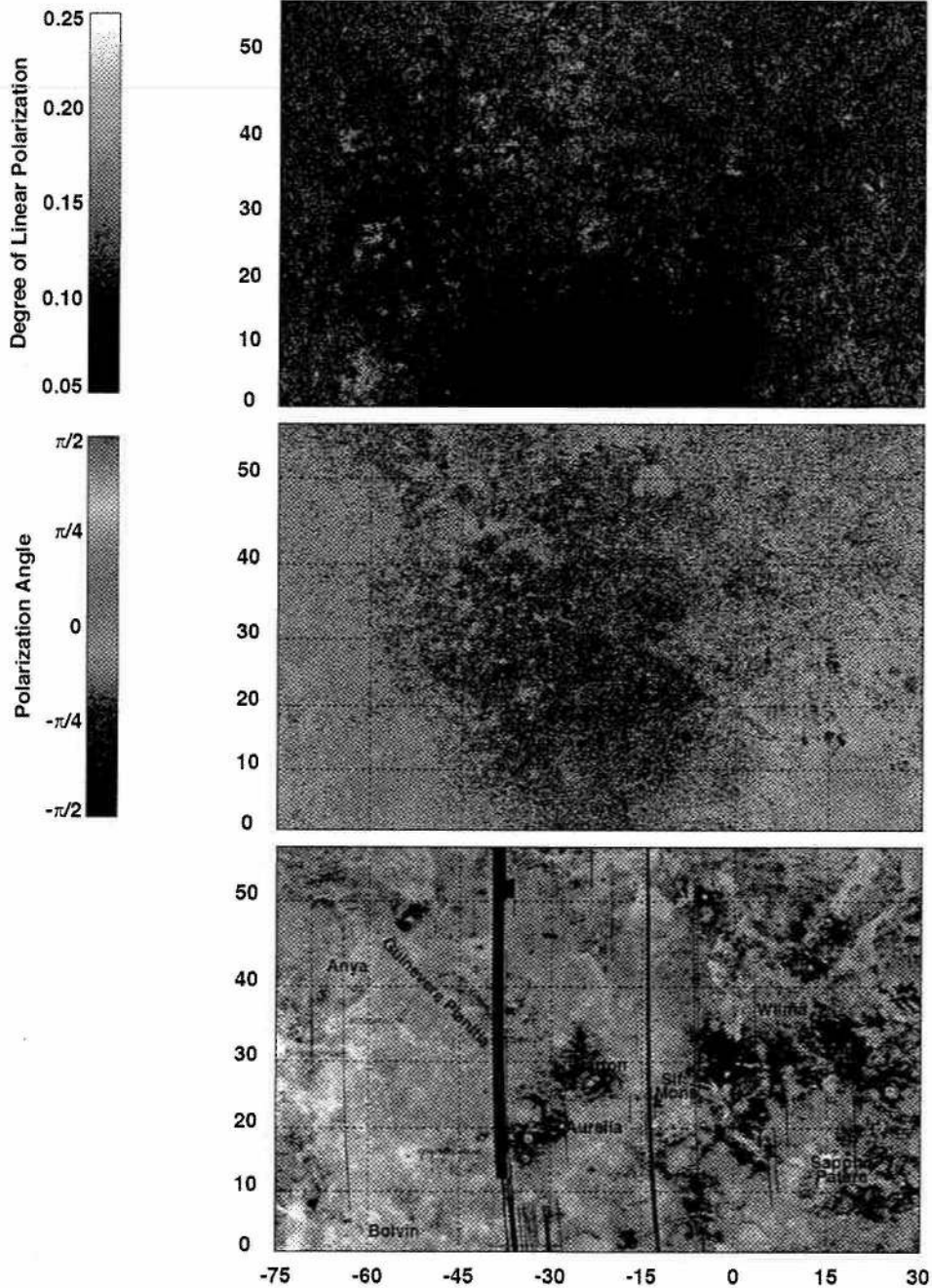
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Fig 3: Images of Galina crater (47.6° N, 307.1° E) on Venus. Top: Magellan SAR image showing the dark halo surrounding the crater. Middle: A degree of linear polarization image stretched to a color scale and overlaid on the Magellan image. The radar dark halo has larger degree of linear polarization values than surrounding areas of radar-bright plains. Bottom: A circular polarization ratio image stretched to a color scale and overlaid on the Magellan image. The radar dark halo surrounding the crater has a low circular polarization ratio, similar to dark halos that surround some lunar craters [Ghent et al. 2005].

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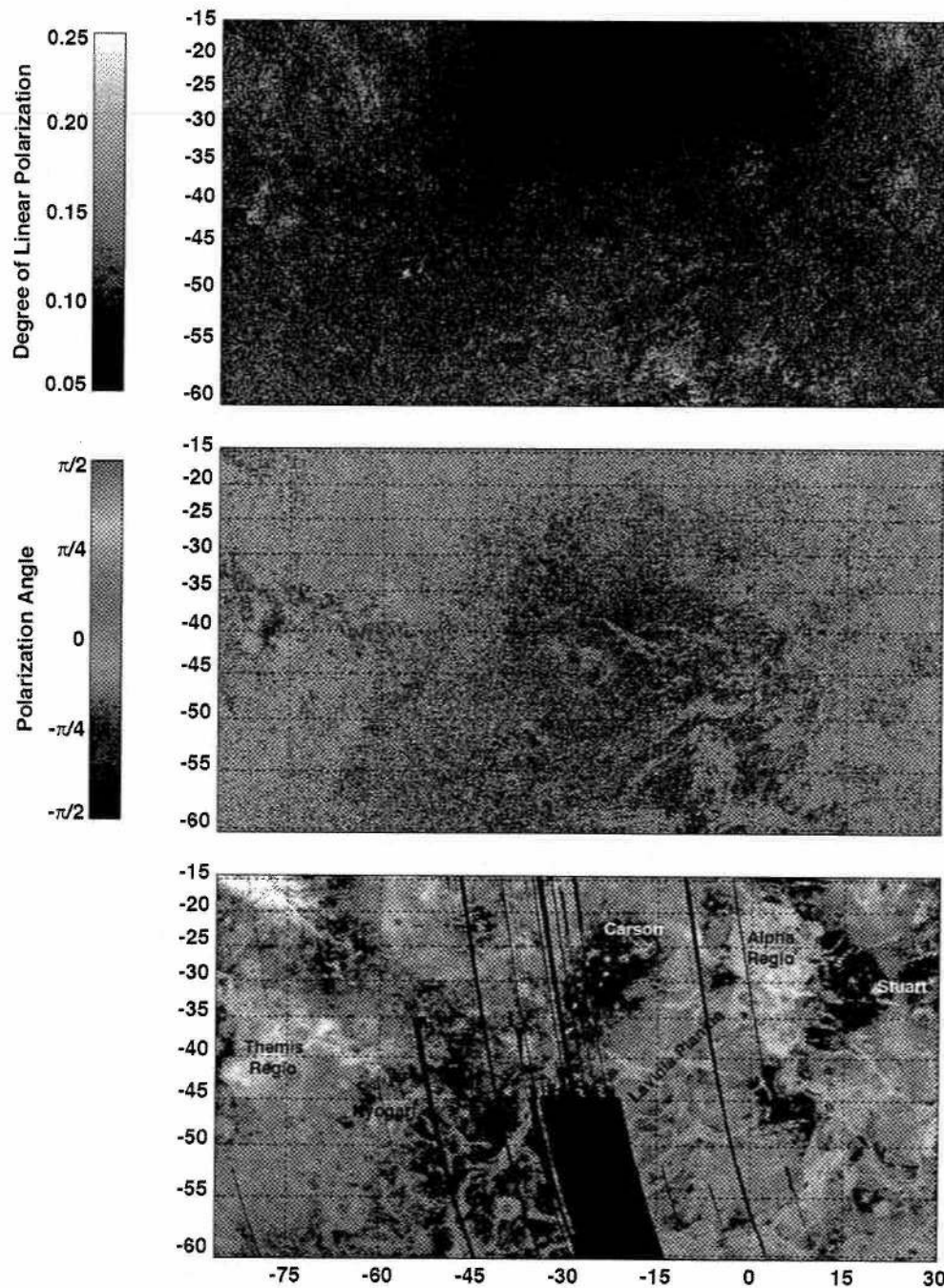
682 **Fig 4:** An image of Sif and Gula Mons (center 40.4° N, 353.7° E) showing both degree of linear
 683 polarization (red) and circular polarization ratio (green) overlaid on a Magellan SAR image.
 684 Both polarizations were stretched with a linear scale such that runs from zero to twice the
 685 average value across the scene. For the degree of linear polarization, the range is 0-0.11. For the
 686 circular polarization ratio the range is 0-0.31. Areas with a higher-than-average degree of linear
 687 polarization, and a lower-than-average circular polarization ratio, are the most red. Areas close to
 688 the average for both polarization products are yellow. Green areas have a lower-than-average
 689 degree of linear polarization and an above-average circular polarization ratio. Displaying both
 690 polarization products simultaneously illustrates the complexity and variety of different surfaces
 691 present in volcanic areas on Venus.

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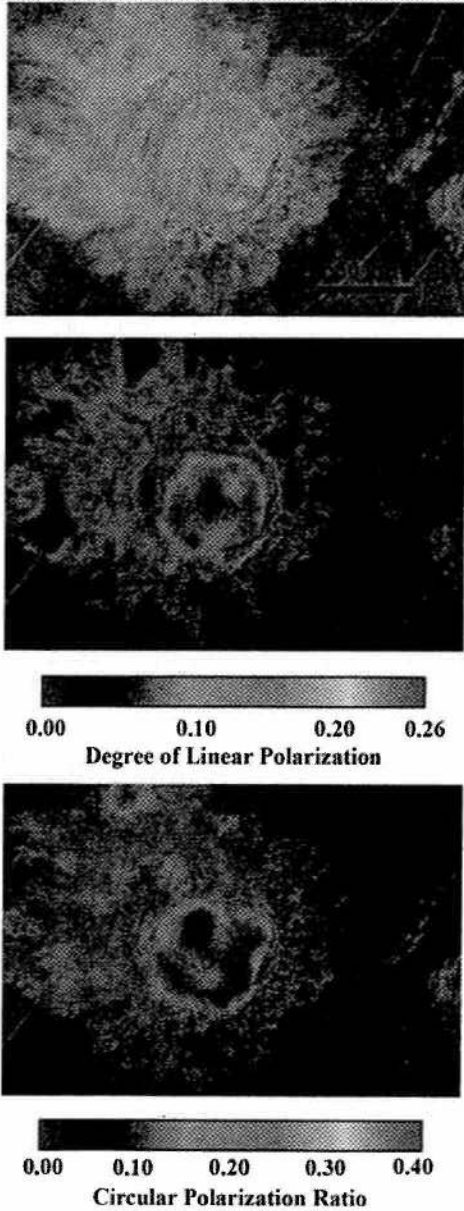
Fig 5: Maps of the northern hemisphere of Venus, from data obtained in 2001. The sub-radar point is in the lower center of the image. Incidence angle increases away from the sub-radar point, towards the top and sides of the image. Top: The degree of linear polarization image. Middle: The linear polarization angle. This angle rotates as the plane of incidence and reflection (perpendicular to the surface) rotates with the curvature of the planet. Bottom: A Magellan SAR image with features labeled for reference.



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Fig 6: Maps of the southern hemisphere of Venus, from data obtained in 2001. The sub-radar point is in the upper center of the image. Incidence angle increases away from the sub-radar point, towards the bottom and sides of the image. Top: The degree of linear polarization image. This angle rotates as the plane of incidence and reflection (perpendicular to the surface) rotates with the curvature of the planet. Middle: The linear polarization angle. Bottom: A Magellan SAR image with features labeled for reference.

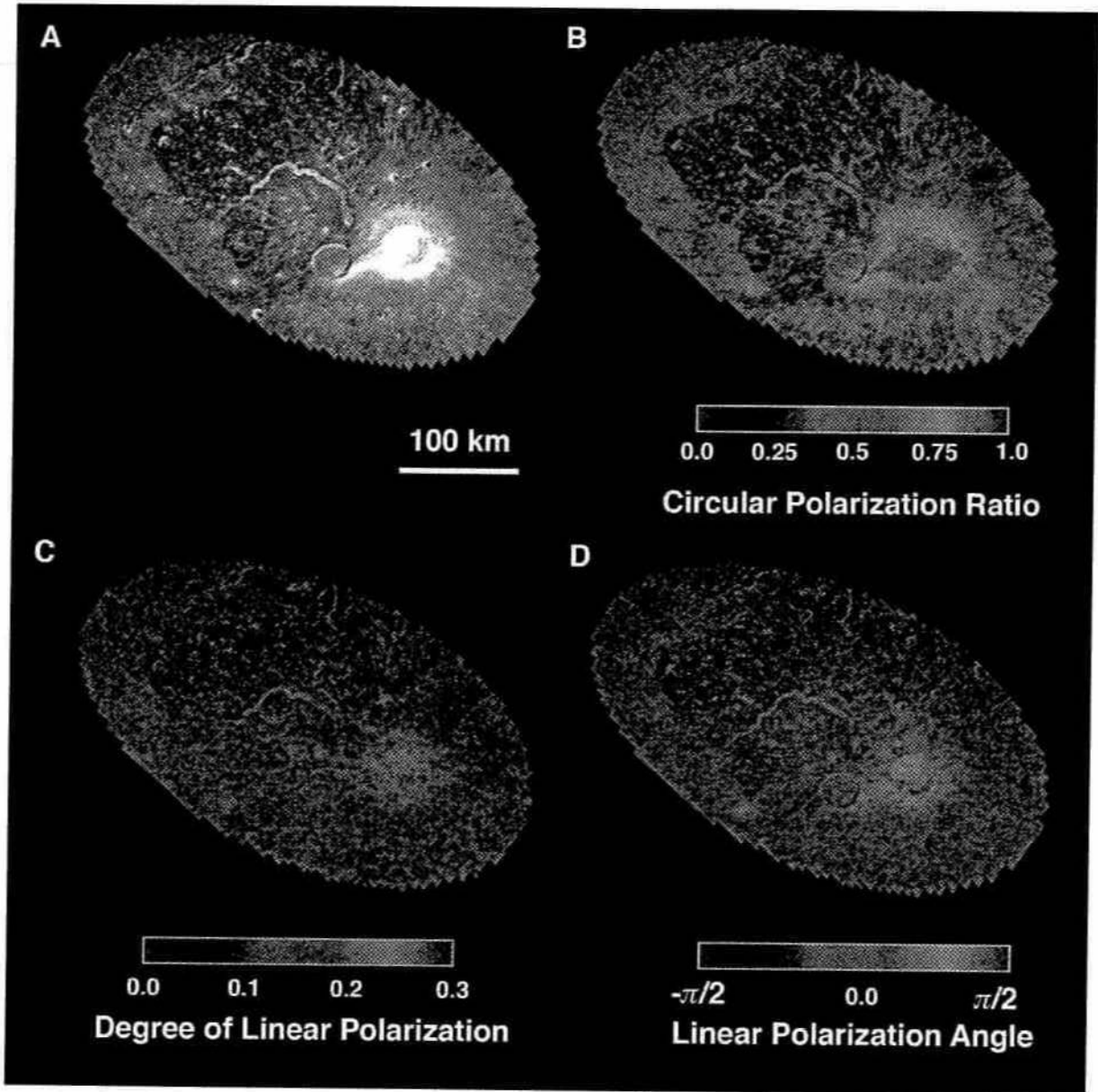
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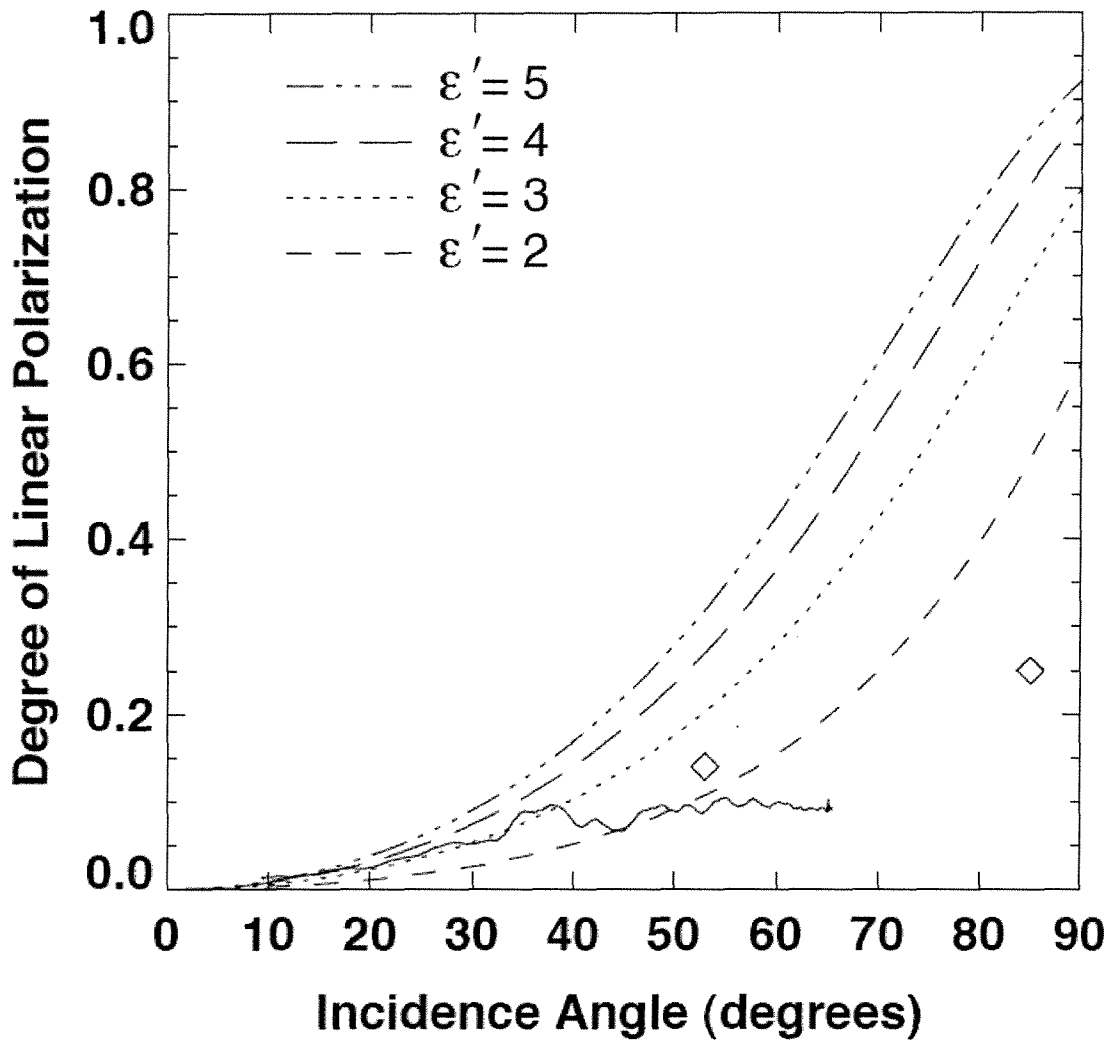
Fig. 7: Images of Stuart crater (30.8° S, 20.2° E), Venus. Top: A Magellan SAR image. Middle: A degree of linear polarization image that has been stretched to a color scale and overlaid on the Magellan image. In this case, a high degree of linear polarization is observed from within the crater walls. Bottom: A circular polarization ratio image that has been stretched to a color scale and overlaid on the Magellan image. The crater floor has a high circular polarization ratio, suggesting a rough surface or subsurface.

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Fig 8: Polarimetric data of Aristarchus plateau. The radar bright Aristarchus crater is located at 23.7° N and 47.4° E. a.) Total power image. b.) Circular polarization ratio stretched to a color scale and overlaid on the total power image. c.) Degree of linear polarization, smoothed, stretched to a color scale, and overlaid on the total power image. d.) Linear polarization angle, smoothed, stretched to a color scale, and overlaid on the total power image.



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Fig 9: Plots of a simple scattering model for the degree of linear polarization shown with Venus data (solid line) and lunar measurements (diamonds). In this model, the radar wave penetrates into a surface that is smooth at the wavelength scale and reflects in a quasi-specular fashion from a subsurface interface. The model is plotted for different values of the mantling (surface) layer dielectric constant. The Venus data line is an average derived from northern hemisphere data acquired in 2001. The increased values around 36° incidence are the result of a polarization feature surrounding Barton crater; individual features typically have a higher than average degree of linear polarization. The lunar data are averages acquired away from large impact craters for the Aristarchus and Focas scenes.