Graphene based all-optical spatial terahertz modulator

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Abstract

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Graphene based All-Optical Spatial Terahertz Modulator

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We demonstrate an all-optical terahertz modulator based on single-layer graphene on germanium (GOG), which can be driven by a 1.55 μm CW laser with a low-level photodoping power. Both the static and dynamic THz transmission modulation experiments were carried out. A spectrally wide-band modulation of the THz transmission is obtained in a frequency range from 0.25 to 1 THz, and a modulation depth of 94% can be achieved if proper pump power is applied. The modulation speed of the modulator was measured to be ~200 KHz using a 340 GHz carrier. A theoretical model is proposed for the modulator and the calculation results indicate that the enhanced THz modulation is mainly due to the third order nonlinear effect in the optical conductivity of the graphene monolayer.

Terahertz (THz) technology offers a variety of applications including spectroscopy, imaging, and communications. Much effort has been devoted to developing THz sources and detectors, which has promoted THz research into one of the most rapidly growing fields[1]. However, substantive progression of THz applications also depends on the realization of active components for wave manipulation and modulation. One of the key components mostly desired for advanced THz system is a spatial light modulator that actively controls the spatial transmission (reflection) of an incident THz wave[2]. This kind of device enables one to encode information on THz waves thus provides potential applications in imaging, telecommunication, beam shaping, etc. Demonstrations of THz modulators have been reported based on 2DEG[3,4], metamaterials[2,4,5], superconductors[6,7], and phase-transition materials[8]. An alternative means of controlling and modulating THz wave is based on optically induced carriers in semiconductors[9,10]. Very recently, an optical driven spatial THz modulator (STM), based on a plain silicon wafer, has been proposed to realize photo-designed THz devices[11] or reconfigurable quasi-optical THz components[12]. The capability of these devices to manipulate THz wave front also enables one to generate special intensity patterns and vortex THz beams[13]. Unfortunately, silicon wafers exhibits a weak modulation depth of 19.9% and a limit of modulation speed less than 100 KHz even when a large laser fluence is applied, which greatly limits the achievable tunability and versatility[13,14]. Therefore, optical STM with large modulation depth under moderate laser power is highly desired in producing efficient and high speed THz devices.

Graphene, a single layer of carbon atom with honeycomb structure, has proven useful for THz and optical modulation due to its unique band structure and extremely high carrier mobility[15–19]. With graphene, light modulators were realized with desirable features such as broadband or high-speed[20,21]. Recently, a graphene-based THz modulator was demonstrated by electrically tuning the density of states available for intraband transitions[15,19]. However, its modulation depth and speed are limited to 15% and 20 KHz by the electrical device prototype. Striking progresses were made by Rahm et al[22]. They observed an enhanced THz modulation in a graphene/silicon structure over a wide frequency band. When a pulsed femtosecond laser with wavelength of 808 nm is incident upon the graphene/silicon, a great number of free carrier are produced in silicon and these photo-generated carries diffuse into the graphene layer, giving rise to a stronger change of conductivity than would occur in pure Si since graphene has higher carrier mobility. This phenomenon results in a THz modulation with a depth of 99% in the frequency range from 0.2 to 2 THz. However, experiments performed so far were primarily focused on the static properties and the modulation depth and speed have not been systematically studied. Furthermore, Si modulators are not efficient at wavelengths of 1.3–1.55 μm due to the inherent large bandgap of Si thus has the problem of integrating with existing optical fiber communications[23].
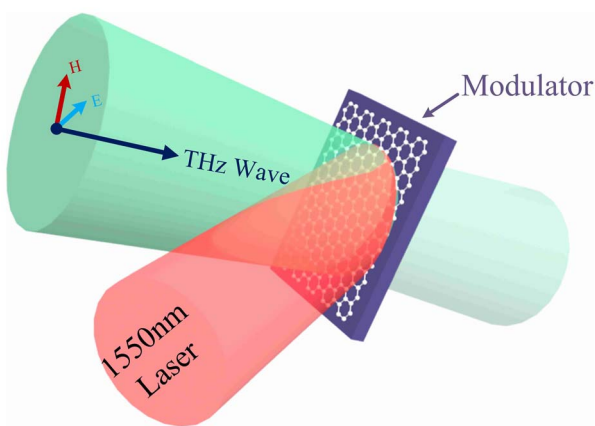
Figure 1 | Prototype and spatial configuration of the all-optical spatial terahertz modulator. The modulator consists of a single-layer graphene sheet on a germanium substrate, which is driven by a 1550 nm laser. The beam of the THz wave is completely overlapped by the laser beam.

In this work, we demonstrate a high speed, broadband and cost-effective THz all-optical modulator. The modulator is based on graphene on germanium (GOG), which allow a pumping light with 1.3–1.55 µm wavelength since Ge has a small bandgap of 0.66 eV. Another advantage of Ge compared to silicon is that Ge has higher bulk mobility for both electrons and holes, which should ideally correspond to an increase in surface mobility and ultimately an increase in the device performance such as modulation depth and speed. In the present work, both the static and dynamic THz transmission modulation experiments were carried out. A wide-band modulation of the THz transmission was obtained in a frequency range from 0.25 to 1 THz under a 1550 nm pumping laser. A maximum modulation depth of 94% and a modulation speed of ~200 KHz were measured by using a 340 GHz carrier. These features make GOG a viable candidate for integration with telecommunication fiber for broadband, high speed and low-cost spatial THz modulator.

The prototype and the spatial configuration of the GOG based STM is shown in Fig. 1. The THz wave is overlapped by the 1550 nm modulation laser beam, and both beams are incident from the graphene side. The spectral transmission of the modulator was measured by a home-made THz time domain spectroscopy (TDS) with and without photoexcitation. In this TDS setup, an Er:fiber laser produces 1550 nm pulses of 100 fs duration with 100 MHz

repetition rate. The source beam is split into three portions, corresponding to terahertz generation, probe, and pump beams, respectively. The pump beam has a maximum average power of 400 mW and a spot diameter of 5 mm so that it can overlap the THz beam (~3 mm).

In Fig. 2, the amplitude transmittances of the THz modulator and the Ge substrate, with and without laser pumping, are plotted in both the time- and frequency-domain. The frequency-domain spectra are normalized to a reference spectrum measured without any sample. Upon pumping at a light fluence of 400 mW, both GOG and pure Ge substrate show considerable transmission attenuations of THz wave. It is well known that free carriers can be generated in a semiconductor through photoexcitation, which forms a thin conductive layer at the surface of the sample[10]. The reflection and absorption by this conductive layer of a THz wave causes the broadband attenuation of the transmitted THz wave[10,22]. Since the conductivity of this thin layer, σ, is proportional to the density and the mobility of the photo-induced carriers, the larger mobility of graphene gives rise to the stronger THz attenuation of GOG as compared to pure Ge. It is worth to point out that the modulation effect of pure Ge substrate is also quite notable since pure Ge itself has high mobility for both electrons (typically 3900 $cm^2V^{-1}s^{-1}$) and holes (typically 1900 $cm^2V^{-1}s^{-1}$). The peak to peak amplitude of the main pulse of the time domain spectrum is specially plotted as inset in Fig. 2 (a) as a function of the pump power. It is clear that the THz transmission for both samples decreases with the increase of the pump fluence, confirming that it is the increase of the conductivity that causes the attenuation of the transmitted THz wave[22].

The frequency-domain spectra shown in Fig. 2(b) provide more detailed information for the broadband modulation properties. Without photoexcitation, both samples show a moderate transmission of ~50% of the THz wave due to the reflection and partial absorption from the carriers since the Ge substrate is slightly p-doped. Similar result has been reported in pure Ge in THz frequency range[24]. It is interesting to see that GOG exhibits a higher transmission than Ge at frequency beyond 0.4 THz. This phenomenon is well reproducible. The possible reason is that the graphene acts as an ultrathin anti-reflection coating to Ge, just like a parylene thin film on Si substrate[24], thus the transmission through GOG is slightly enhanced. Upon pumping by 1.55 µm laser with a power of 400 mW, the modulation depth averaged from 0.25 THz to 1 THz is 83% for GOG, while that for pure Ge is 68%. The modulation depth mentioned here is defined as $(T_{\text{no pumping}} - T_{\text{pumping}})/T_{\text{no pumping}}$, where $T_{\text{no pumping}}$ and T_{pumping} represent the intensity of THz transmission without and with light pumping respectively. The
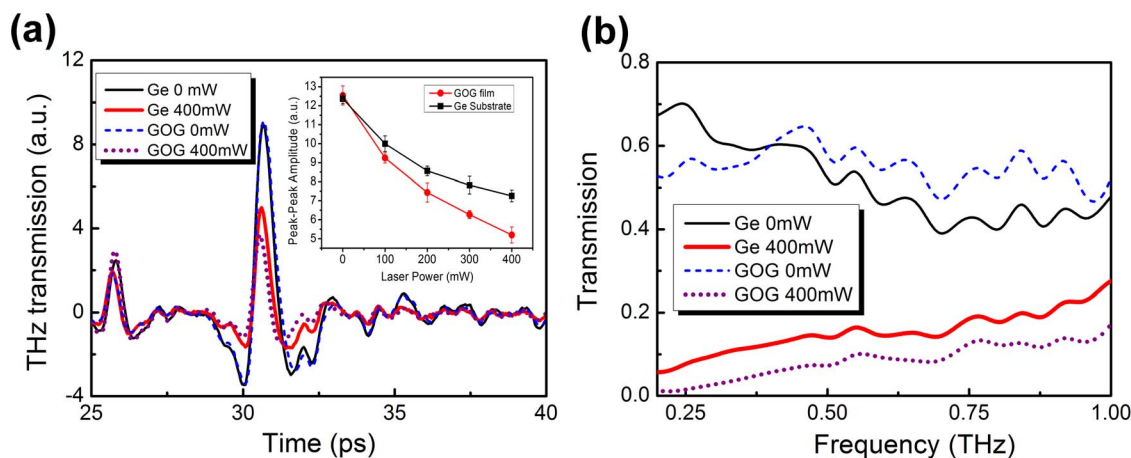

Figure 2 | Measured THz transmission from the Ge and GOG samples. The measurements were carried out on a THz-TDS system without and with photodoping at a power of P = 400 mW. (a) The time domain spectra and (b) the corresponding frequency domain spectra calculated from the time domain data. Inset in (a) plotted the peak-peak amplitude of the main THz pulse as a function of the pumping power.
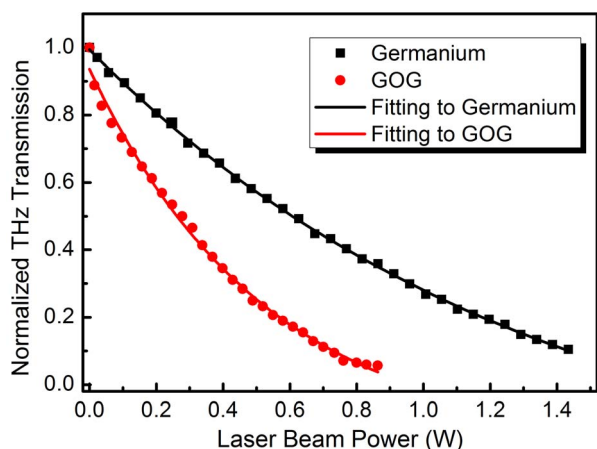
Figure 3 | Measured static modulation characteristics of the STM. Normalized THz transmission at 340 GHz carrier as a function of the laser beam power for the Ge and GOG. The solid lines are fitting curves to the experimental data by an exponential decay function.

modulation depth of GOG is more significant as compared to pure Ge across the whole measurement range from 0.25 THz to 1 THz.

For a practical implementation of all-optical terahertz modulators, continuous wave (CW) lasers are more convenient and cost-effective than pulsed femtosecond lasers. Static and dynamic modulation characteristics of the GOG modulator were further investigated at 340 GHz carrier with a CW laser. The detected THz transmission signals are normalized and plotted in Fig. 3 as a function of the photoexitation power. It can be seen that for both Ge and GOG, the THz transmission decreases exponentially with the increasing beam power and finally the transmission decay approached a saturation point. At equivalent power levels of the photodoping, GOG has larger attenuation to the THz wave when compared with Ge. It is shown that for GOG the modulation saturation occurs at a pump power of 0.8 W and at that power the modulation depth is 94%. At exactly the same pumping power the modulation depth of pure Ge is only 64.2%. This feature of GOG to obtain large modulation depth at low-level photodoing power is very attractive for all-optical THz devices.

The modulation speed of GOG was measured at 340 GHz carrier. In the measurement, a square-wave voltage was used to drive the CW diode laser to produce pump beam with the power alternating

between zero and 200 mW. Here a small laser power of 200 mW was adopted for the purpose of avoiding the possible accumulation of heating effects. The detected voltage signal is plotted in Fig. 4(a) for different modulation frequencies. Generally, the transmission amplitude of THz wave decreases with the increase of modulation frequency. At the modulation frequency of 20 KHz, the waveform of the modulated THz signal is almost square, which gradually changes to a triangle shape when the modulation frequencies increase beyond 40 KHz. The detected voltage at 100 KHz falls to 78.4% of that for 20 KHz, while at 200 KHz this value notably falls to ~50%. At 1 MHz, the device still has a response to the photoexitation but the signal falls to near noise level. The normalized modulation magnitude at different modulation frequency was summarized in Fig. 4(b), which indicate a 3 dB operation bandwidth (f_c) is ~200 kHz. According to the principle of GOG modulator mentioned above, it is the carrier recombination time of Ge, not graphene, that sets the ultimate limit. Previous research on germanium crystals show that the carrier lifetimes depend on the resistivity of Ge and the photo-induced carrier level[25]. Considering the resistivity of Ge used here (about 3 $\Omega\cdot$cm, corresponding to a carrier density of ~10^6 cm^{-3}) and excitation level of 200 mW, the carrier recombination time is in the neighborhood of 2 μs, suggesting an actual limit of 0.5 MHz for real applications. This result is in good agreement with the directly measured modulation speed of the device. It is reported that a carrier lifetime as short as 30 ns can be obtained in Ge[25]. Therefore, the modulation rate of GOG based STM can be tens of megahertz if suitable Ge materials and optical lasers are used.

Discussion

In a similar structure of graphene on silicon (GOS), a semianalytical model has been proposed to calculate the spectral transmission and the modulation depth of THz radiation[20]. However, by that model the derived numerical results only agree with the experimental data at very low beam power of 30 mW. For higher laser power, the calculated transmissions of GOS deviated qualitatively and quantitatively from the experimental results. The effect of graphene on THz transmission cannot be well described by that model especially at high modulation intensities. It is reported that graphene is a strong nonlinear material in THz band and the optical activity of single layer graphene can be significantly enhanced by nonlinear effects[26]. Here we give a qualitative analysis of the observed strong nonlinear transmission in GOG by using the model of nonlinear optical conductivity of monolayer graphene in the terahertz regime[26,27].
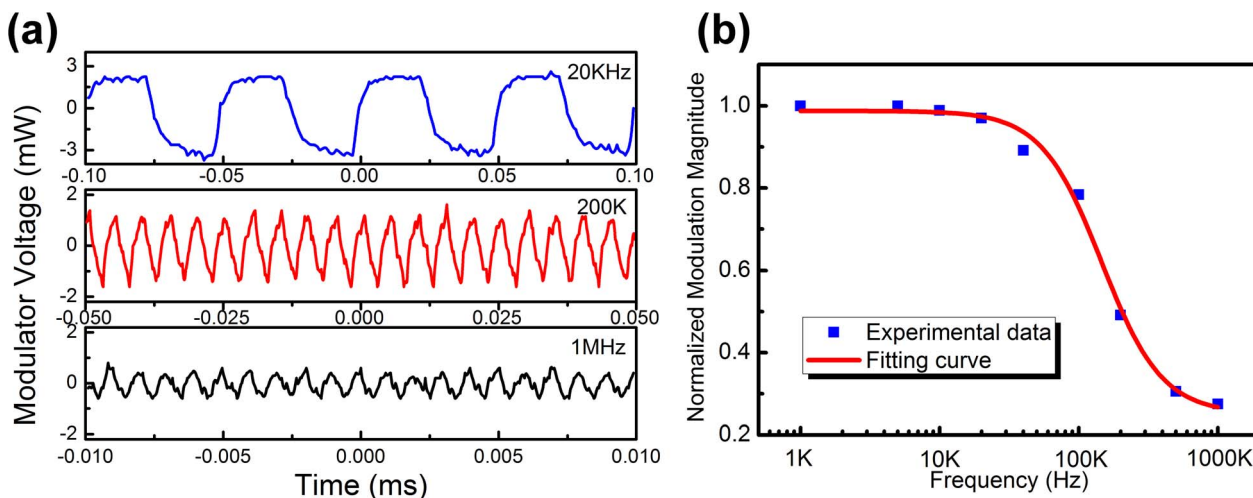

Figure 4 | Measured dynamic characteristics of the STM. (a) The modulated terahertz beam signal for a carrier frequency of 340 GHz with different modulation frequency. (b) Normalized modulation magnitude, showing a 3 dB operation bandwidth of ~200 kHz.
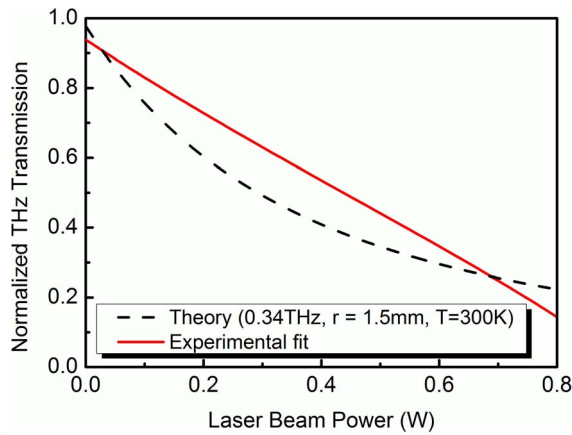
Figure 5 | Comparison with the model of nonlinear optical response. Normalized transmission coefficient of graphene at 340 GHz carrier as a function of the laser beam power. The black dash line is the theoretical result from Eq.(1) with a radius of the incident laser beam r = 1.5 mm. The beam has been modeled as having a Gaussian distribution. The red solid line is the experimental result of $T_{SLG} = T_{eff}/T_{Ge}$.

The transmission coefficient is given as,

$$T(E) = \frac{1}{\left(1 + \frac{1}{2}\alpha\pi\sigma(E)\right)^2} \quad (1)$$

where $\alpha \approx 1/137$ is the fine structure constant. Up to the third order in electric field strength, ignoring the phonon effect and other mechanical effects, the nonlinear conductivity up to the third order nonlinear process can be obtained by using the Floquet expansion[26,28],

$$\sigma(E) = \sqrt{(1 + \sigma_3(E)H_2)^2 + \sigma_3^2(E)H_1^2} \quad (2)$$

where $\sigma_3(E) = \frac{e^2 v_F^2 E_0^2}{\hbar^2 \omega^4}$, $H_1 = \frac{13}{48}N\left(\frac{1}{2}\right) - \frac{2}{3}N(1) + \frac{45}{48}N\left(\frac{3}{2}\right)$, $H_2 = 2N(1)$, $N(x) = \tanh\left(\frac{x\hbar\omega}{2k_B T}\right)$, e is the charge of an electron, $v_F \approx 1 \times 10^6$ m/s is the Fermi velocity of graphene, E_0 is the strength of the incident field, ω is the frequency of incident photons, and T is the temperature of the system.

The H_1 and H_2 terms are a result of taking the time average of the nonlinear currents, $J_3(\omega)$ and $J_3(3\omega)$ respectively, in the graphene sheet. The total current is written as the sum of these terms, i.e. $J_{tot} =$ $J_1(\omega) + J_2(\omega) + J_3(3\omega)$. When we take the time average of $|J_{tot}|^2$, the mixed frequency terms, for example $J_3(\omega) \cdot J_3(3\omega)$ (and complex conjugate), will disappear, leaving just the un-mixed terms. In terms of pumping power P, the third order conductivity of the graphene can be written as,

$$\sigma_3(P) = \frac{e^2 v_F^2}{\hbar^2 \omega^4} \frac{2 \ln(2)}{c\varepsilon_0 \pi r^2} P \quad (3)$$

where r is the approximate radius of the incident laser beam.

For a graphene monolayer on Ge, the effective transmission is the product of the transmission of the monolayer graphene and that of Ge, $T_{eff} = T_{SLG} \times T_{Ge}$. Therefore the measured $T_{SLG} = T_{eff}/T_{Ge}$. In Figure 5, we compare theoretical result based on Eq.(1) and the measurement at the carrier frequency of 340 GHz. In the calculation the diameter of the laser beam is 3 mm and the system temperature is 300 K. We have modeled the beam as having a Gaussian distribution and have set the FWHM as the beam diameter (~3 mm). It can be seen the intensity dependence of the transmission coefficient is mainly due to the third order effect in the optical conductivity in terahertz regime. The qualitative agreement is good and the quantitative difference is also quite small even under high laser beam power of 800 mW. The observed quantitative difference may be due to effects neglected in our model such as the imperfectness of the graphene sheet or the inexact profile and size of the laser beam used in the calculation. Works to refine the theoretical model are in progress. Despite the inaccuracy of the model, the theoretical calculations strongly support that the third order effect occurring in graphene result in the enhanced attenuation of THz transmission in GOG.

In conclusion, we experimentally demonstrated a graphene based all-optical terahertz modulator fabricated on germanium, which can be efficiently driven by a 1.55 μm CW laser at relatively low-level power fluences. We measured a spectrally wide-band modulation of the THz transmission in the frequency range from 0.25 to 1 THz. Dynamic experiments at 340 GHz carrier evidence that GOG provides a 3 dB modulation speed of 200 KHz and an ultimate limit of 0.5 MHz. Theoretical calculations indicates that the laser intensity dependence of the transmission attenuation is mainly due to the third order effect in the optical conductivity of the graphene monolayer. The major advantage of this device is the ability to integrate with low-loss telecommunication fiber thus an all-optical, high speed and low-cost spatial THz modulator can be readily realized.

Methods

Sample fabrication and characteristic. The device studied in this work consists of a single-layer of graphene on top of a p-type Ge substrate, as indicated in Fig. 1. The
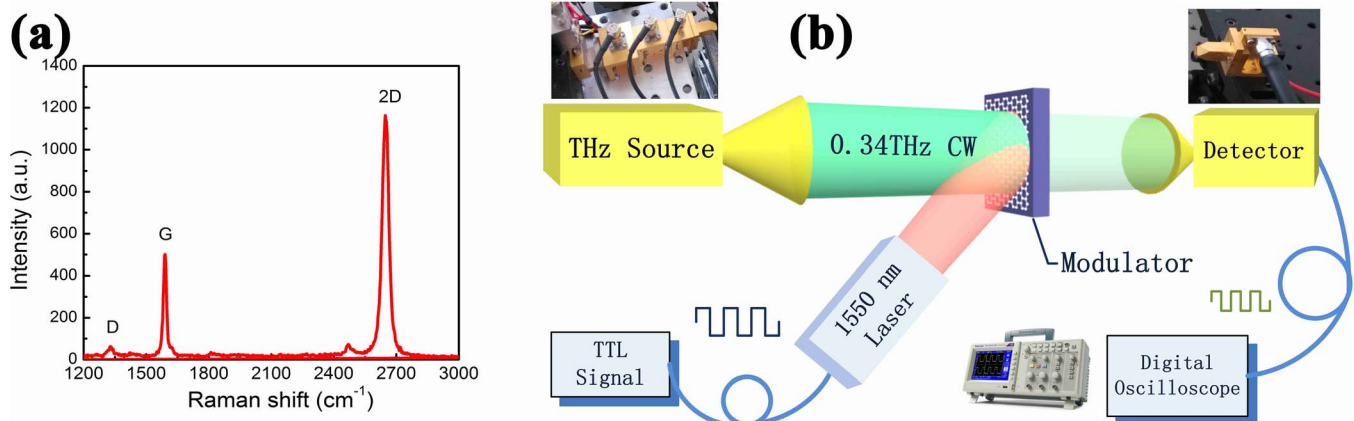

Figure 6 | Graphene quality and terahertz characterization set-up. (a) Representative Raman spectrum of the CVD graphene sample. (b) Schematic of the experimental configuration used for Static and dynamic modulation measurements, which using a 340 GHz THz CW source and a 1550 nm pumping laser. The laser is mounted with a beam incident angle of 30° to the normal of the modulator.

modulator was fabricated using chemical vapor deposition (CVD) graphene grown on copper[29]. The Czochralski Ge substrate (Wafer World, Inc) is p-doped with a moderate resistivity of ~3 $\Omega \cdot$cm. The thickness of Ge is 250 ± 25 µm. Graphene was transferred by means of polymethyl methacrylate (PMMA) and wet etch methods. Fig. 6(a) is the typical Raman spectrum for our graphene obtained using a 442 nm excitation laser. The 2D band is located at 2650 cm^{-1}, and is more than 2 times the height of the corresponding G band, indicating a single layer graphene on Ge substrate. Notice that the film we obtained has some defects since a slight D band was detected.

Terahertz modulator characterization. The measurement set-up for static and dynamic modulation characteristics is shown in Fig. 6(b). The system consists of a VDI (Virginia Diodes) CW terahertz source with a central output in the 340 GHz, and a 240–400 GHz zero-bias Schottky diode intensity detector. The average output power of the terahertz source is ~10 mW. A 1.55 µm CW diode laser was used to generate the pumping beam with a maximum power of 3 W and a spot diameter of d ~ 3 mm (e.g, r ~ 1.5 mm).

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Acknowledgments

Author contributions

Q.Y.W. concerned the devices structure and guided the theoretical and experimental work. W.T. designed and established the 340 GHz measurement system, Q.M. and Z.C. fabricated and characterias the modulator devices. W.W.L. and Q.H.Y. performed the THz-TDS measurements, M.S. developed the theoretical model and did the calculations. H.W.Z. guided the experimental work. All authors discussed the results and co-wrote the manuscript.

Additional information

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