

HADAMARD'S VARIATIONAL FORMULA FOR THE SZEGÖ KERNEL

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§1. A variational formula. The present note is concerned with the Hadamard's (first) variational formula for the Szegö kernel associated with a strictly pseudo-convex domain in \mathbf{C}^n with $n \geq 3$. A similar formula for the Bergman kernel has been given in [7].

Let $\Omega^0 \subset \mathbf{C}^n$ with $n \geq 1$ be a bounded domain with smooth boundary $\partial\Omega^0$. Every smoothly perturbed domain of Ω^0 can be parametrized by a small function $\rho \in C^\infty(\partial\Omega^0; \mathbf{R})$ in such a way that the boundary of that domain Ω^ρ is given by

$$(1) \quad \partial\Omega^\rho = \{\zeta + \rho(\zeta)\nu(\zeta); \zeta \in \partial\Omega^0\},$$

where $\nu(\zeta) = \partial/\partial\nu_\zeta$ denotes the unit exterior normal vector to Ω^0 at $\zeta \in \partial\Omega^0$ identified with an element of \mathbf{C}^n .

Let $S^\rho(z, w)$ for $z, w \in \Omega^\rho$ denote the Szegö kernel associated with Ω^ρ , which is the reproducing kernel associated with the space $L^2_\delta H(\Omega^\rho)$ of holomorphic functions in Ω^ρ with L^2 boundary values equipped with the $L^2(\partial\Omega^\rho)$ scalar product. With $\delta\rho \in C^\infty(\partial\Omega^0; \mathbf{R})$ and $z, w \in \Omega^\rho$ fixed arbitrarily, we set

$$(2) \quad \delta S^\rho(z, w) = \frac{d}{d\varepsilon} S^{\rho+\varepsilon\delta\rho}(z, w)|_{\varepsilon=0},$$

which is the Hadamard's first variation of $S^\rho(z, w)$ at ρ in the direction $\delta\rho$. Our purpose is to show that, for a certain class of domains Ω^0 , the variation (2) at $\rho=0$ exists and is given by

$$(3) \quad -\delta S^0(z, w) = \int_{\partial\Omega^0} \frac{\partial}{\partial\nu_\zeta} [S^0(z, \zeta)S^0(\zeta, w)] \cdot \delta\rho(\zeta) d\sigma^0(\zeta) \\
 + \int_{\partial\Omega^0} S^0(z, \zeta)S^0(\zeta, w)H^0(\zeta)\delta\rho(\zeta) d\sigma^0(\zeta),$$

where $d\sigma^0(\zeta)$ denotes the induced surface measure of $\partial\Omega^0 \subset \mathbf{C}^n$ at ζ , and $H^0(\zeta)$ stands for the mean curvature of $\partial\Omega^0$ at ζ multiplied by $2n-1$. A concrete statement of our result is given as follows:

THEOREM. *If $\Omega^0 \subset \mathbf{C}^n$ is strictly pseudo-convex with $n \geq 3$, then the variation (2) at $\rho=0$ exists and is given by (3).*

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Note that the right hand side of (3) makes sense, for if $\Omega^0 \subset \mathbb{C}^n$ is strictly pseudo-convex then $S^0(\cdot, \cdot)$ extends smoothly to $(\overline{\Omega^0} \times \overline{\Omega^0}) \setminus \Delta(\partial\Omega^0)$, where $\Delta(\partial\Omega^0)$ denotes the diagonal of $\partial\Omega^0 \times \partial\Omega^0$ (see Boutet de Monvel and Sjöstrand [1], see also Kerzman and Stein [5]).

Remark 1. As will be seen in Section 3, the variational formula (3) is valid whenever the Szegö kernel associated with Ω^ρ depends smoothly on ρ in the sense of (6) in Section 2.

It is plausible that (3) holds if $n=1$. In fact, if $n=1$, then the Szegö kernel is expressed in terms of the Bergman kernel and the harmonic measures (see Garabedian [3]). The smooth dependence of the Bergman kernel on ρ has been established (cf. [7], Remark 2), while the harmonic measures are expressed in terms of the Poisson kernel and thus depend smoothly on ρ , cf. Section 2.

The assumption $n \geq 3$ in Theorem above is imposed in order to use an expression of the Szegö kernel in terms of the $\bar{\partial}_b$ -Neumann operator, see (9) in Section 2. It is likely that Theorem above is valid also for strictly pseudo-convex domains in \mathbb{C}^2 .

Remark 2. In case $n=1$, Schiffer [9] has obtained another expression for the variation (2) in terms of the Szegö kernel and the so-called adjoint kernel. It is not difficult to see that his formula follows from (3).

§ 2. Existence of the variation (2). Setting

$$\mathcal{CV}^0(\varepsilon_1) = \{ \rho \in C^\infty(\partial\Omega^0; \mathbf{R}); |\rho(\zeta)| < \varepsilon_1 \text{ for } \zeta \in \partial\Omega^0 \}$$

with $\varepsilon_1 > 0$ small, we begin with constructing a family of diffeomorphisms $e_\rho : \mathbb{C}^n \rightarrow \mathbb{C}^n$ for $\rho \in \mathcal{CV}^0(\varepsilon_1)$ such that

$$(4) \quad \begin{cases} e_\rho(\zeta) = \zeta + \rho(\zeta)\nu(\zeta) & \text{for } \zeta \in \partial\Omega^0 \quad (\text{cf. (1)}), \\ \mathcal{CV}^0(\varepsilon_1) \ni \rho \mapsto e_\rho \in C^\infty(\mathbb{C}^n; \mathbb{C}^n) & \text{is continuous,} \\ e_\rho - e_0 & \text{depends linearly on } \rho \text{ and } e_0 = \text{identity.} \end{cases}$$

In particular, (4) will imply that e_ρ depends smoothly on ρ and that $e_\rho(\partial\Omega^0) = \partial\Omega^\rho$ so that $e_\rho(\Omega^0) = \Omega^\rho$. Several ways of constructing such a family $\{e_\rho; \rho \in \mathcal{CV}^0(\varepsilon_1)\}$ are possible. We shall employ the one as in [7], which will be convenient for our purpose.

Given a small constant $\varepsilon_0 > 0$, we consider a tubular neighborhood $N(\varepsilon_0) = \{z \in \mathbb{C}^n; |r^0(z)| < \varepsilon_0\}$ of $\partial\Omega^0$ in \mathbb{C}^n , where $r^0 \in C^\infty(\mathbb{C}^n; \mathbf{R})$ is a defining function of Ω^0 such that

$$\Omega^0 = \{z \in \mathbb{C}^n; r^0(z) < 0\}, \quad |dr^0(z)| = 1 \text{ for } z \in N(\varepsilon_0).$$

Then, every point $z \in N(\varepsilon_0)$ is uniquely expressed as $z = \zeta_z + r^0(z)\nu(\zeta_z)$, where $\zeta_z \in \partial\Omega^0$ is the nearest point to z . Fixing a constant ε_1 with $0 < \varepsilon_1 < \varepsilon_0/4$, we

choose $\chi_0 \in C_0^\infty(\mathbf{R}; \mathbf{R})$ satisfying

$$\begin{aligned} \chi_0(r) &= 1 & \text{for } |r| \leq \varepsilon_1, & & \text{and } \left| \frac{d}{dr} \chi_0(r) \right| &\leq \frac{3}{4\varepsilon_1} & \text{for } r \in \mathbf{R}. \\ \chi_0(r) &= 0 & \text{for } |r| \geq 3\varepsilon_1, & & & & \end{aligned}$$

For $\rho \in \mathcal{C}\mathcal{V}^0(\varepsilon_1)$, we define a mapping $e_\rho : \mathbf{C}^n \rightarrow \mathbf{C}^n$ by setting

$$(5) \quad \begin{aligned} e_\rho(z) &= z + \chi_0(r^0(z))\rho(\zeta_z)\nu(\zeta_z) & \text{for } z \in N(\varepsilon_0), \\ e_\rho(z) &= z & \text{otherwise.} \end{aligned}$$

Then, $\{e_\rho; \rho \in \mathcal{C}\mathcal{V}^0(\varepsilon_1)\}$ is a family of diffeomorphisms satisfying (4).

By means of e_ρ , one can pull back in general a function f^ρ in Ω^ρ or on $\partial\Omega^\rho$ and a linear operator L^ρ acting on f^ρ as follows:

$$f_\rho = e_\rho^* f^\rho = f^\rho \circ e_\rho, \quad L_\rho f_\rho = (e_\rho^* L^\rho e_\rho^{-1*}) f_\rho = (L^\rho (f_\rho \circ e_\rho^{-1})) \circ e_\rho.$$

Let $S^\rho : L^2(\partial\Omega^\rho) \rightarrow L^2 H_b(\partial\Omega^\rho) \subset L^2(\partial\Omega^\rho)$ denote the Szegő projector associated with Ω^ρ , which is the orthogonal projector onto $L^2 H_b(\partial\Omega^\rho) = L^2 H(\Omega^\rho)|_{\partial\Omega^\rho}$ and is related to $S^\rho(z, w)$ by

$$S^\rho f^\rho(z) = \int_{\partial\Omega^\rho} S^\rho(z, \zeta) f^\rho(\zeta) d\sigma^\rho(\zeta) \quad \text{for } f^\rho \in L^2(\partial\Omega^\rho),$$

where $d\sigma^\rho(\zeta)$ stands for the induced surface measure of $\partial\Omega^\rho \subset \mathbf{C}^n$ at ζ . Then, $S_\rho = e_\rho^* S^\rho e_\rho^{-1*}$ satisfies

$$S_\rho f_\rho(z) = \int_{\partial\Omega^0} S_\rho(z, \zeta) f_\rho(\zeta) d\sigma^\rho(e_\rho(\zeta)) \quad \text{for } f \in L^2(\partial\Omega^0),$$

where we have set

$$S_\rho(z, w) = S^\rho(e_\rho(z), e_\rho(w)) \quad \text{for } (z, w) \in \Omega^0 \times \overline{\Omega^0}.$$

Observe by (5) that $S_\rho(z, w) = S^\rho(z, w)$ for $z, w \in \Omega^0 \setminus N(\varepsilon_0)$. Therefore, the variation (2) exists for $z, w \in \Omega^0 \setminus N(\varepsilon_0)$ fixed, provided that $S_\rho(z, w)$ depends smoothly on ρ as far as ρ is small with respect to the $C^\infty(\partial\Omega^0)$ -topology. For the later use, we shall show that

$$(6) \quad \mathcal{C}\mathcal{V}^2 \ni \rho \mapsto S_\rho(\cdot, w) \in C^\infty(\overline{\Omega^0}) \text{ is smooth}$$

with $w \in \Omega^0 \setminus N(\varepsilon_0)$ fixed, where $\mathcal{C}\mathcal{V}^2$ is a neighborhood of $0 \in C^\infty(\partial\Omega^0; \mathbf{R})$.

In order to prove (6), we first recall that

$$\begin{aligned} S^\rho(z, w) &= \int_{\partial\Omega^\rho} S^\rho(z, \zeta) P^\rho(w, \zeta) d\sigma^\rho(\zeta) \\ &= [S^\rho P^\rho(w, \cdot)](z) \quad \text{for } (z, w) \in \overline{\Omega^\rho} \times \Omega^\rho, \end{aligned}$$

where $P^\rho(w, \zeta)$ denotes the Poisson kernel associated with Ω^ρ , see Kerzman and Stein [5]. Then,

$$(7) \quad S_\rho(\cdot, w) = S_\rho P_\rho(w, \cdot) \quad \text{on } \overline{\Omega^0} \quad \text{for } w \in \Omega^0 \setminus N(\varepsilon_0).$$

We next recall the assumption that Ω^0 is strictly pseudo-convex, so that

$$(8) \quad \sum_{j,k=1}^n \frac{\partial^2 r^0(z)}{\partial z_j \partial \bar{z}_k} \xi_j \bar{\xi}_k \geq C \sum_{j=1}^n |\xi_j|^2 \quad \text{whenever } \sum_{j=1}^n \frac{\partial r^0(z)}{\partial z_j} \xi_j = 0$$

holds for each $z \in \partial\Omega^0$, where $C > 0$ is a constant independent of z . Hence, if $\rho \in \mathcal{C}^0(\varepsilon_1)$ is small with respect to the $C^2(\partial\Omega^0)$ -topology, say,

$$\rho \in \mathcal{C}^2 = \{\rho \in \mathcal{C}^0(\varepsilon_1); |\rho|_{C^2(\partial\Omega^0)} < \varepsilon_2\} \quad \text{with } \varepsilon_2 > 0 \text{ small,}$$

then Ω^ρ is strictly pseudo-convex uniformly in $\rho \in \mathcal{C}^2$ in the sense that (8) holds for each $z \in \partial\Omega^\rho$ with $r^\rho = r^0 \circ e_\rho^{-1}$ in place of r^0 , where $C > 0$ is independent of $\rho \in \mathcal{C}^2$. If moreover $n \geq 3$ then the following formula holds:

$$(9) \quad S^\rho = 1 - \mathcal{D}_b^* N_b^* \bar{\delta}_b^\rho, \quad \text{thus } S_\rho = 1 - (\mathcal{D}_b)_\rho (N_b)_\rho (\bar{\delta}_b)_\rho,$$

where $\bar{\delta}_b^\rho$ and \mathcal{D}_b^ρ denote the tangential Cauchy-Riemann operator acting on $C^\infty(\partial\Omega^\rho)$ and its $L^2(\partial\Omega^\rho)$ adjoint, respectively, and N_b^ρ stands for the $\bar{\delta}_b$ -Neumann operator acting on the space $C_{(0,1)}^\infty(\partial\Omega^\rho)$ of tangential $(0, 1)$ -forms on $\partial\Omega^\rho$ (see Kohn [6], or Folland and Kohn [2]). The definitions of $(\mathcal{D}_b)_\rho$, $(N_b)_\rho$ and $(\bar{\delta}_b)_\rho$ will be clear except for the fact that the space $e_\rho^* C_{(0,1)}^\infty(\partial\Omega^\rho)$ may vary with ρ . However, one may modify it to be independent of ρ by considering the projection: $e_\rho^* C_{(0,1)}^\infty(\partial\Omega^\rho) \rightarrow C_{(0,1)}^\infty(\partial\Omega^0)$ (see Kuranishi [8]). The smooth dependence of this modification of the pull-back of $(N_b)_\rho$ on ρ small is involved in Kuranishi [8]. Therefore, S_ρ depends smoothly on ρ in the sense that

$$\mathcal{C}^2 \times C^\infty(\partial\Omega^0) \ni (\rho, f) \mapsto S_\rho f \in C^\infty(\partial\Omega^0) \quad \text{is smooth.}$$

Since the Poisson kernel $P^\rho(w, \cdot)$ can be expressed in terms of the Green kernel, the smooth dependence of $P_\rho(w, \cdot)$ on ρ can be proved as in Hamilton [4] (the easier case). Hence, by virtue of (7), we have proved (6). In particular, the variation (2) makes sense.

§3. Proof of the variational formula (3). The proof is similar to that in [7]. Pick $z, w \in \Omega^0$ arbitrarily and choose $\varepsilon_0 > 0$ so small that $z, w \in \Omega^0 \setminus N(\varepsilon_0)$. Then,

$$S_\rho(z, w) = S^\rho(z, w) \quad \text{for } \rho \in \mathcal{C}^2.$$

By the reproducing property for the Szegő kernel, we have

$$\begin{aligned} S_\rho(z, w) &= S^\rho(z, w) = \int_{\partial\Omega^\rho} S^\rho(z, \zeta) S^\rho(\zeta, w) d\sigma^\rho(\zeta) \\ &= \int_{\partial\Omega^0} S_\rho(z, \zeta) S_\rho(\zeta, w) J_b[e_\rho](\zeta) d\sigma^0(\zeta), \end{aligned}$$

where $J_b[e_\rho]$ stands for the Jacobian determinant of the mapping e_ρ restricted to

$\partial\Omega^0$. Recalling (6), we take the variation at $\rho=0$ in the direction $\delta\rho \in C^\infty(\partial\Omega^0; \mathbf{R})$. Then,

$$\begin{aligned} \delta S^0(z, w) &= \delta S_0(z, w) = \frac{d}{d\varepsilon} S_{\varepsilon\delta\rho}(z, w)|_{\varepsilon=0} \\ &= \int_{\partial\Omega^0} \{(I_1) + (I_2) + (I_3)\} d\sigma^0(\zeta), \end{aligned}$$

where

$$\begin{aligned} (I_1) &= \delta S_0(z, \zeta) S^0(\zeta, w), \quad (I_2) = S^0(z, \zeta) \delta S_0(\zeta, w), \\ (I_3) &= S^0(z, \zeta) S^0(\zeta, w) \delta J_b[e_0](\zeta), \end{aligned}$$

and $\delta J_b[e_0] = \frac{d}{d\varepsilon} J_b[e_{\varepsilon\delta\rho}]|_{\varepsilon=0}$. Setting $\delta X_0 = \frac{d}{d\varepsilon} e_{\varepsilon\delta\rho}|_{\varepsilon=0}$, we get

$$\delta X_0(\zeta) = \delta\rho(\zeta) \frac{\partial}{\partial\nu_\zeta}, \quad \delta J_b[e_0](\zeta) = \operatorname{div} \delta X_0(\zeta) = \delta\rho(\zeta) H^0(\zeta)$$

for $\zeta \in \partial\Omega^0$, and

$$\begin{aligned} (10) \quad \delta S_0(z, \zeta) &= \delta S^0(z, \zeta) + \delta X_0(\zeta) S^0(z, \zeta), \\ \delta S_0(\zeta, w) &= \delta S^0(\zeta, w) + \delta X_0(\zeta) S^0(\zeta, w) \end{aligned}$$

for $\zeta \in \Omega^0$, where the vector field $\delta X_0(\zeta)$ in (10) is acting as a differential operator. Note that $\delta S^0(z, \cdot)$ and $\delta S^0(\cdot, w)$ extend smoothly to $\overline{\Omega^0}$, and that the relations in (10) remain valid for $\zeta \in \overline{\Omega^0}$. Moreover, $\delta S^0(\cdot, w)$ and $\delta S^0(z, \cdot)$ are holomorphic and conjugate holomorphic in Ω^0 , respectively. Since $S^0(\cdot, \cdot)$ is hermitian symmetric with the reproducing property, we have

$$\begin{aligned} \int_{\partial\Omega^0} (I_1) d\sigma^0(\zeta) &= \delta S^0(z, w) + \int_{\partial\Omega^0} \delta X_0(\zeta) S^0(z, \zeta) \cdot S^0(\zeta, w) d\sigma^0(\zeta), \\ \int_{\partial\Omega^0} (I_2) d\sigma^0(\zeta) &= \delta S^0(z, w) + \int_{\partial\Omega^0} S^0(z, \zeta) \cdot \delta X_0(\zeta) S^0(\zeta, w) d\sigma^0(\zeta), \end{aligned}$$

while

$$\int_{\partial\Omega^0} (I_3) d\sigma^0(\zeta) = \int_{\partial\Omega^0} S^0(z, \zeta) S^0(\zeta, w) H^0(\zeta) \delta\rho(\zeta) d\sigma^0(\zeta).$$

Summing them up, we obtain the desired variational formula (3).

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