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Half-life of ²²¹Fr in Si and Au at 4 K and at millikelyin temperatures

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The half-life of the α -decaying nucleus ²²¹Fr was determined in different environments, that is, embedded in Si at 4 K, and embedded in Au at 4 K and about 20 mK. No differences in half-life for these different conditions were observed within 0.1%. Furthermore, we quote a value for the absolute half-life of ²²¹Fr of $t_{1/2} = 286.1(10)$ s that is of comparable precision to the most precise value available in the literature.

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I. INTRODUCTION

The half-lives of α - and β -decaying nuclei have always been considered constant, independent of their surroundings. Recently, however, it has been claimed that the half-life of a radioactive isotope would change if embedded in a metallic host. Using the Debye plasma model applied to quasifree metallic electrons [1–3], a $1/\sqrt{T}$ dependence for the screening energy U_D was predicted. When cooled down to a few kelvins, the half-lives of $\beta^-(\beta^+)$ -decaying nuclei would then increase (decrease) by several tens of percent and the half-lives of α -decaying nuclei would even be shortened by several orders of magnitude. This hypothetical mechanism was proposed as a possible solution for long-lived transuranic waste produced by fission reactors [1].

A series of subsequent experiments claimed to have observed changes in the half-lives of both β^{\mp} and α decays, although the effects were less profound than predicted by the Debye plasma model. The half-life of α -decaying ²¹⁰Po implanted in copper was reported to shorten by 6.3(14)% when cooled down to 12 K [4], the half-life of β^+ -decaying ²²Na embedded in Pd was observed to be 1.2(2)% shorter at 12 K [5], the half-life of β^- -decaying ¹⁹⁸Au embedded in Au was observed to be 4(2)% longer at 12 K [6], and for ⁷Be [electron capture (EC)] decay), an increase of about 1% was reported [7]. Note that electron capture decay rates may depend on the material hosting the radioactive isotope via small modifications of the electron density around the EC-decaying nucleus (see, e.g., [8] and references therein). In contrast to the above results, several experiments carried out at a later stage on the β -decaying isotopes ¹⁹⁸Au [9–12], ²²Na [11,12], ⁶⁴Cu [13], and ⁷⁴As [14], and on the EC-decaying ⁷Be [10], did not observe any changes of the half-lives up to the permille level when these isotopes were embedded in a metallic environment and cooled down to 10-20 K.

As to α decay, it was demonstrated theoretically that taking into account not only the electron screening effect on the α particle's tunneling potential, but also on its binding

energy inside the nucleus, the half-life would not significantly change [15]. Further, it was argued [16,17] that if the Debye plasma model would be applicable to α decays, the effect should have been observed previously in Low Temperature Nuclear Orientation (LTNO) experiments [18]. In this type of experiment, radioactive isotopes are implanted in a metallic foil, typically Fe, and cooled to millikelvin temperatures. It was claimed [4,6] that no appreciable effect on nuclear half-lives can be observed in LTNO experiments, as the radioactive ions implanted at typically 60 keV would end up in the oxidized surface layer of the sample foil, which acts as an insulator. This claim is clearly incorrect as the LTNO technique is precisely based on the fact that the radioactive nuclei end up at substitutional sites in a pure bcc Fe lattice, where they experience a unique hyperfine interaction. Nuclei that end up in an oxide surface layer cannot be oriented and thus also do not show anisotropic emission. To avoid a surface oxide layer, the foils used in LTNO measurements are carefully prepared by polishing and annealing procedures prior to inserting them in the vacuum of the setup [19,20]. Substitutional fractions of about 70% to 95% are generally observed for 60-keV implantations into a cold (i.e., 4 K or lower) Fe foil (e.g., Refs. [21–24]). In LTNO experiments on α -decaying isotopes implanted in Fe, no half-life changes at the percent level were observed between room temperature and 1 K [17], and between 4 K and 50 mK [16].

Here, we report a dedicated half-life measurement of the α -decaying nucleus ²²¹Fr implanted in Si at 4 K and in Au at 4 K and at 20 mK. The three different half-life values obtained agree with each other within 0.1% and also agree with the literature room temperature value. Previously, it was shown that the half-life of ²²¹Fr is constant to less than a percent at room temperature, regardless of its chemical environment [25].

II. THE EXPERIMENT

The 221 Fr activity was produced at the ISOLDE/CERN facility by spallation reactions induced by a 1.4-GeV proton beam impinging on a UC_x target [26]. After diffusing out of the target, the 221 Fr nuclei were ionized in a W surface

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ion source [27], accelerated to 60 keV, and mass selected by the General Purpose Separator [28]. The radioactive ions were subsequently transported to the NICOLE LTNO setup [24,29], where they were implanted into the cooled Au or Si sample. A SRIM calculation [30] showed that the implantation depth of ²²¹Fr ions with an energy of 60 keV is 80 Å for Au and 330 Å for Si. In previous LTNO experiments, it was shown that under these implantation conditions about 80% of the ²²¹Fr ions end up in a fully metallic environment in a well-prepared Fe sample foil [23,31].

The decay α particles were detected with three 500- μ m-thick Si PIN diode detectors that were positioned at an angle of 15° relative to the surface of the sample foil and at an angle of 90° relative to each other. This type of detector has previously been tested under the conditions of this experiment and has shown good behavior [32]. A high-efficiency coaxial germanium detector, placed outside the cryostat, was used for γ -ray detection. Further, standard NIM electronics and a PC-based data acquisition system were used.

The measurements consisted of a succession of implantation and data-taking periods, wherein the latter a spectrum was recorded every 30 s during at least 10^{-221} Fr half-lives. First, nine such measurements were performed with an Au foil at millikelvin temperatures, after which the temperature of the foil was raised to 4 K and another five measurements were carried out. Thereafter, the Au sample foil was removed and a Si sample was loaded into the system to carry out five measuring cycles at 4 K. During all measurements, a 1-kHz pulser was used to monitor the dead time. Only data for which the dead time was below 5% and the shorter-lived beam contaminant 221 Ra [$t_{1/2} = 28(2)$ s] [33] had decayed were considered for analysis.

A typical α spectrum is shown in Fig. 1. After the 221 Ra activity has decayed, the only α lines appearing in the spectrum originate from the decay of 221 Fr and its daughter nuclei. The α line at 6.34 MeV, which will be used to determine the half-life of 221 Fr, showed an energy resolution of 32 keV (full width at half maximum). The α lines of 217 At and 213 Po were broadened to, respectively, 43 keV and 57 keV because of the relocation of the recoiling daughter nuclei following the α decay of 221 Fr and 217 At, respectively, leading to a larger scattering for the

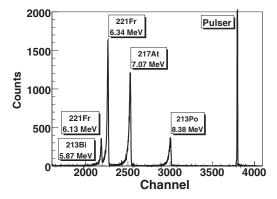


FIG. 1. A typical α spectrum, taken a few minutes after the end of the implantation of the mass 221 beam. The integral of the pulser peak was used to correct for dead time.

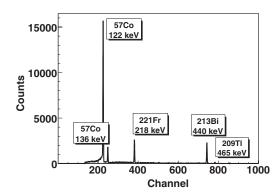


FIG. 2. Typical γ -ray spectrum. Apart form the two γ lines of 57 Co, only γ lines originating from the decay of 221 Fr or its daughter nuclei were observed.

decay α particles and even allowing part of the nuclei to leave the host foil. Therefore, only the α lines of ²²¹Fr, which was directly implanted, were considered for analysis.

The γ spectrum shown in Fig. 2 confirmed that, after the isotope 221 Ra had decayed, 221 Fr and its daughters were the sole activities present in the sample foil. The two γ lines at 122 keV and 136 keV originate from the 57 CoFe nuclear thermometer [34] that was soldered on to the back side of the sample holder so as to monitor the temperature in the millikelvin range.

III. ANALYSIS AND RESULTS

In order not to be sensitive to possible small gain shifts during the measurements, the 221 Fr 6.34-MeV α peak was integrated using reasonably wide markers, including, in fact, both α lines of 221 Fr (at 6.34 MeV and 6.13 MeV [35]), as well as the weak 5.87-MeV α line of 213 Bi [36]. The background contribution to this integral originates from this 213 Bi α line and from the long tail of the 8.38-MeV α line from the decay of 213 Po [36], which has the same effective half-life as 213 Bi. After correcting the integrals for the dead time by normalizing to the pulser peak (see Fig. 1), it was checked that the background indeed decays according to the literature value of the half-life of 213 Bi, that is, 2735(4) s [36] (Fig. 3).

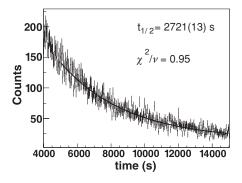


FIG. 3. A fit with a single exponential of the background of the 221 Fr α peak integral. The resulting half-life coincides with the literature value of the half-life of 213 Bi.

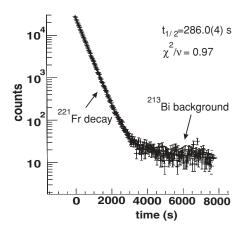


FIG. 4. Fit to the 221 Fr half-life, $t_{1/2}$, of the decay curve of the combined 221 Fr α peaks. The background decays according to the half-life of 213 Bi (see Fig. 3).

To fit the $^{221}\mathrm{Fr}$ decay curve, the following function was used:

$$Ae^{-\frac{(t+B)ln^2}{C}} + D[e^{-\frac{(t+B)ln^2}{E}} - e^{-\frac{(t+B)ln^2}{C}}],$$
 (1)

which takes into account the buildup and decay of the 213 Bi background activity. The free parameters in the fit were the amplitudes of the counts from 221 Fr, A, and from 213 Bi, D, as well as the half-life of 221 Fr, C. The fixed parameter B expresses the start of the ingrowth of the 213 Bi activity and is determined by a fit of the ingrowth curve of the 440-keV γ line in the β decay of 213 Bi. For the parameter E, the previously mentioned literature value of the 213 Bi half-life was used [36]. A typical example of a fit to determine the half-life of 221 Fr using Eq. (1) is shown in Fig. 4.

For the integration of the peaks in the γ spectrum, a linear background subtraction was carried out, taking into account the amplitude of the background to the left and right of the peak. As the integral of the 218-keV γ line in the decay of 221 Fr went to zero, a single exponential could be fitted to extract the half-life (Fig. 5).

All fits were performed using the MINUIT package [37] inside the publicly available ROOT framework. Nearly all fits

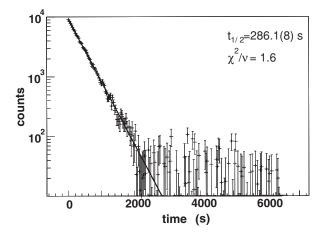


FIG. 5. Single exponential fit to the decay curve of the $^{221}\mbox{Fr}$ 218-keV γ line.

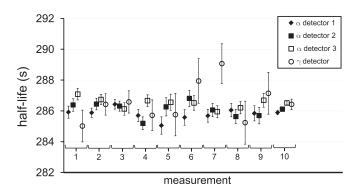


FIG. 6. 221 Fr half-life values obtained in the nine measurements at millikelvin temperatures in the Au host foil, for the three α -particle detectors (using the most intense α lines in the decay of 221 Fr) and the 218-keV γ line. The last block (labeled 10) shows the weighted average values for all nine results for each detector.

resulted in reduced χ^2 values, χ^2/ν , close to unity. This also demonstrates that the integrations and dead-time corrections were being dealt with correctly. Whenever the χ^2/ν of the fit was larger than unity the error on the corresponding $t_{1/2}$ was multiplied by $\sqrt{\chi^2/\nu}$.

The results from all four detectors, three α -particle detectors and one γ detector, as well as the results of successive decay measurements in one given material and for a given temperature, all turned out to be mutually consistent, thus demonstrating the stability of our experiment. This is illustrated in Figs. 6, 7, and 8, where all $t_{1/2}$ values extracted from the three different data sets are shown. The evolution of the temperature of the sample during one of the measurements at millikelvin temperatures is shown in Fig. 9.

Table I summarizes the results for the three different measurement conditions. Within 2.5 standard deviations, no difference is observed between the half-lives of 221 Fr embedded in the nonmetallic Si environment at 4 K, in Au at 4 K, and in Au at millikelvin temperatures. Our results therefore do not support the claimed temperature or host material dependence

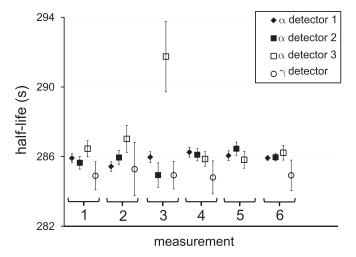


FIG. 7. ²²¹Fr half-life values obtained in the five measurements at 4 K in the Au host foil. The last block (labeled 6) shows the weighted average values for all five results for each detector.

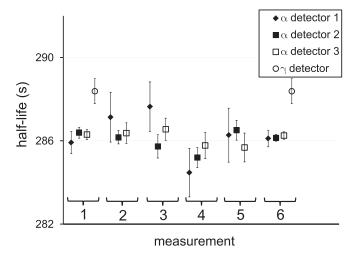


FIG. 8. 221 Fr half-life values obtained in the five measurements in the Si host foil. During the last four measurements of this series, the γ spectra were corrupted. The last block (labeled 6) shows the weighted average values for all five results for each detector.

for α decay. The weighted average value of all three results in Table I is $t_{1/2}(^{221}{\rm Fr})=286.12(9)~(\chi^2/\nu=1.6)$. The quoted error on this value is purely statistical. To quote an absolute half-life of $^{221}{\rm Fr}$, systematic errors have still to be taken into account. The systematic error resulting from the dead-time correction by pulser normalization was estimated by varying the markers that define the integral of the pulser peak, yielding a 0.5-s variation on the half-life of $^{221}{\rm Fr}$. In addition, the integrals during the first half-life of $^{221}{\rm Fr}$ were the most sensitive to the dead-time correction, as at later times the amount of dead time rapidly dropped below 1%. Varying the starting point of the fit reveals another systematic error, related to the dead-time correction, of 0.7 s. The systematic error related to the presence of counts from the decay of $^{213}{\rm Bi}$

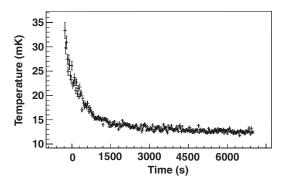


FIG. 9. Temperature evolution of the Au sample during one of the measurements at millikelvin temperatures deduced from the ⁵⁷CoFe nuclear orientation thermometer. During the implantation period, the sample warms to several tens of millikelvins due to the power deposited by the mass 221 beam and the subsequent radioactive decay of the source that is gradually increasing in strength. About 200 s after the end of the implantation period, all of the short lived ²²¹Ra activity is gone and the sample foil quickly cools down to below 20 mK, after which the half-life measurement was started.

was estimated to be 0.5 s by varying the value used for the half-life of 213 Bi, that is, parameter B in Eq. (1), within one standard deviation. Finally, an 0.11-s systematic error comes from the accuracy of our timing system, which was determined by calibrating it with an external clock.

Adding all systematic errors in quadrature, our final value for the half-life of ²²¹Fr becomes

$$t_{1/2}(^{221}\text{Fr}) = 286.1 \pm (0.09)_{\text{stat}} \pm (1.0)_{\text{syst}} \text{ s.}$$
 (2)

This value is in agreement with and more precise than the value of 294(12) s quoted in Refs. [36,38], and is in agreement with and even slightly more precise than the value of 287.4(12) s quoted in Ref. [25].

TABLE I. Half-life values for 221 Fr measured by observing the decay α and γ rays at different temperatures, that is, in the mK region or at 4 K, and in different environments, that is, Au and Si. The values presented in the fifth column are the weighted averages of the values of all four detectors. The last column indicates the difference of the half-life values obtained, compared to the reference Si measurement. The quoted errors are purely statistical. Whenever the χ^2/ν of the weighted average was larger than unity, the error bar was increased by a factor of $\sqrt{\chi^2/\nu}$.

Measurement/condition	Detector	$t_{1/2}(^{221}\text{Fr})$ (s)	$\sqrt{\chi^2/\nu}$	$t_{1/2}(^{221}\text{Fr})$ (s)	$\sqrt{\chi^2/\nu}$	$\Delta t_{1/2} \ (\%)$
mK/Au	α_1	285.89(13)	0.8			
	α_2	286.10(11)	1.2			
	α_3	286.52(13)	0.8			
	γ	286.43(36)	1.1			
	•	, ,		286.17(8)	0.9	0.01(5)
4 K/Au	α_1	285.91(13)	1.2	. ,		
	α_2	285.96(18)	1.4			
	α_3	286.22(43)	3.0			
	γ	284.92(95)	3.6			
	•	, ,		285.93(11)	1.4	0.09(6)
4 K/Si	α_1	286.11(40)	1.2	, ,		
	α_2	286.14(21)	1.5			
	α_3	286.25(25)	0.9			
	γ	288.38(60)	_			
	•			286.19(12)	1.1	_

IV. CONCLUSION

Because of the findings reported in [16] and [17], it comes as no surprise that our results do not support the Debye-Hücker model [1–3], which predicts (based on the calculations of [11,12]) a reduction of the half-life of 221 Fr in a metal at 4 K with a factor of about 50, and at 20 mK with a factor of many orders of magnitude more. No dependency on the solid-state environment and temperature of this α decaying isotope is observed up to a level of 1×10^{-3} , which is at variance with the reported 6% change of the activity of 213 Po nuclei implanted in Cu at 12 K [4]. Our result supports the theoretical calculations of [15], which predicted no significant change of the half-life. These dedicated half-life measurements bring down the level to which no temperature and host-material

effect is observed on nuclear half-lives for α decays to the same level of precision as was already established for β decays and EC decays in Refs. [8–14].

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