# Harmonic Analysis of Tempered Distributions on Semisimple Lie Groups of Real Rank One 

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#### Abstract

SUMMARY. Let $G$ be a real semisimple Lie group. Harish-Chandra has defined the Schwartz space, $\mathscr{E}(G)$, on $G$. A tempered distribution on $G$ is a continuous linear functional on $\mathscr{E}(G)$.

If the real rank of $G$ equals one, Harish-Chandra has published a version of the Plancherel formula for $L^{2}(G)[3(k), \S 24]$. We restrict the Fourier transform map to $\mathscr{C}(G)$, and we compute the image of the space $\mathscr{C}(G)$ [Theorem 3]. This permits us to develop the theory of harmonic analysis for tempered distributions on $G$ [Theorem 5].


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## Summary

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[^0]$\S 1$. Introduction. Let $G$ be a real semisimple Lie group. The Fourier transform map, $\mathscr{F}$, can be regarded as an isometry form $L^{2}(G)$ onto $L^{2}(\hat{G}) . L^{2}(\hat{G})$ is a Hilbert space defined with the help of the discrete series, $\mathscr{E}_{d}$, and the various continuous series, $\mathscr{E}_{c}$, of irreducible unitary representations of $G . L^{2}(\hat{G})$ consists of certain functions whose domain is $\mathscr{E}_{d} \cup \mathscr{E}_{c}$, and whose range is the space of Hilbert-Schmidt operators on the Hilbert spaces on which the representations in $\mathscr{E}_{d} \cup \mathscr{E}_{c}$ act.

In [3(1)] Harish-Chandra introduces the Schwartz space, $\mathscr{E}(G)$, of functions on $G$. It is analogous to the space, $\mathscr{S}(\mathbf{R})$, of rapidly decreasing functions on the real line. $\mathscr{C}(G)$ is a Fréchet space. It is dense in $L^{2}(G)$, and its injection into $L^{2}(G)$ is continuous. It is of interest to ask about the image of $\mathscr{C}(G)$ in $L^{2}(\hat{G})$ under $\mathscr{F}$. There is a candidate, $\mathscr{E}(\hat{G})$, for this image space. $\mathscr{C}(\hat{G})$ is a Fréchet space defined by a natural family of seminorms on $L^{2}(\hat{G})$.

A tempered distribution on $G$ is a continuous linear functional on $\mathscr{C}(G)$. If we can prove that the Fourier transform gives a topological isomorphism from $\mathscr{C}(G)$ onto $\mathscr{C}(\hat{G})$, we could define the Fourier transform of a tempered distribution as a continuous linear functional on $\mathscr{C}(\hat{G})$. This would include as a special case the theory of Fourier transforms on $L^{2}(G)$.

We confine ourselves to the case in which the real rank of $G$ equals one. In this case Harish-Chandra has published a version of the Plancherel formula for $L^{2}(G)[\mathbf{3}(\mathrm{k}), \S 24]$. Our main result is Theorem 3, which asserts the bijectivity between $\mathscr{C}(G)$ and $\mathscr{C}(\hat{G})$ of the Fourier transform, $\mathscr{F}$.

The most difficult part of this theorem is to prove surjectivity. We have to show that the inverse Fourier transform of an element in $\mathscr{C}(\hat{G})$, which is a priori in $L^{2}(G)$, is actually in $\mathscr{C}(G)$. We use some estimates which Harish-Chandra develops from the study of a differential equation on $G$ [3(1), §27]. In $\S 9$ we review his work and show that his estimates are actually uniform, in a sense which will become clear. In $\S 10$ we use these estimates to prove that $\mathscr{F}(\mathscr{C}(G))$ contains $\mathscr{C}_{0}(\hat{G})$, a subspace of $\mathscr{C}(\hat{G})$ associated with the discrete series.

To prove that $\mathscr{F}(\mathscr{C}(G))$ contains $\mathscr{C}_{1}(\hat{G})$, the subspace of $\mathscr{C}(\hat{G})$ associated with the continuous series, requires more work. It is necessary to derive a formula (Lemma 41) for the norms of certain linear transformations, $c^{+}(\Lambda)$ and $c^{-}(\Lambda)$, which arise in $\S 12$. This we do in $\S 13$ by studying a second-order symmetric differential operator on $\mathfrak{a}_{\mathfrak{p}}$, a one-dimensional subspace of the Lie algebra of $G$. As a biproduct of this formula we obtain in $\S 14$ a condition for irreducibility of certain representations in the continuous series.

For convenience we work with generalized spherical functions. We develop the pertinent information in $\S 5$ and then use it in $\S 6$ to prove the injectivity of the Fourier transform.

In $\S 16$ we define the Fourier transform of a tempered distribution on $G$. Theorem 6 proves that any continuous linear functional on $\mathscr{C}(\hat{G})$ is a certain sum of tempered distributions on the real line.

It seems likely that some of our methods can be used for proving the analogue of Theorem 3 for arbitrary $G$. The injectivity of the Fourier transform should
carry over quite easily. Harish-Chandra's estimates are proved in [3(1), §27] for arbitrary $G$. That these estimates are uniform can also be shown, although the proof of this is somewhat more complicated than in the real rank 1 case. Our proof of Lemma 27 does not carry over in general. However, it gives a good start toward a general proof.

The general Plancherel formula will be complicated by the existence of more than one continuous series of representations. However, in each continuous series linear transformations $c(\Lambda)$ can be defined. The formulae in Lemma 41 can probably be proved, although perhaps not by our methods. In general, Lemma 44 would be proved by induction on the real rank of $G$. Harish-Chandra does this for ordinary spherical functions in [3(h), Theorem 3].
2. Preliminaries. Let $G$ be a connected real semisimple Lie group with Lie algebra g. Let

$$
\mathfrak{g}=\mathfrak{k}+\mathfrak{p}
$$

be a fixed Cartan decomposition with Cartan involution $\theta$. Let $a_{p}$ be a fixed maximal abelian subspace of $\mathfrak{p}$. The dimension of $\mathfrak{a}_{\mathfrak{p}}$ is called the real rank of $G$. We shall assume that $\operatorname{dim} \mathfrak{a}_{\mathfrak{p}}=1$.

Let $\mathfrak{a}_{\mathfrak{k}}$ be a subspace of $\mathfrak{k}$ such that

$$
\mathfrak{a}=\mathfrak{a} \mathfrak{k}+\mathfrak{a}_{\mathfrak{p}}
$$

is a Cartan subalgebra of $\mathfrak{g}$. Let $K$ be the analytic subgroup of $G$ corresponding to $k$. We assume that $G$ has finite center. This implies that $K$ is compact.

We can make further technical assumptions on $G$ without losing generality. In order to do this we state some definitions of Harish-Chandra.

If $L$ is a connected reductive Lie group over the reals, $\mathbf{R}$, with Lie algebra $\mathfrak{l}$, let

$$
j: \mathfrak{l} \mathfrak{l}_{c}
$$

be inclusion into the complexification of $\mathfrak{l}$. (From now on, if $\mathfrak{h}$ is any real Lie algebra we write $\mathfrak{h}_{c}$ for its complexification.) Then if $L_{c}$ is a complex analytic group with Lie algebra $l_{\mathbf{c}}, L_{\mathbf{c}}$ is called a complexification of $L$ if $j$ extends to a homomorphism of $L$ into $L_{\mathbf{c}}$. Let $\mathfrak{l}=\mathfrak{l}_{1}+\mathfrak{c}$, where $\mathfrak{l}_{1}$ is semisimple and $\mathfrak{c}$ is abelian. Let $\mathfrak{l}_{1 \mathbf{c}}$ and $\mathfrak{c}_{\boldsymbol{c}}$ be the respective complexifications of $\mathfrak{l}_{1}$ and $\mathfrak{c}$. Let $L_{1}, C\left(L_{1 \mathbf{c}}, C_{\mathbf{c}}\right)$ be the analytic subgroups of $L\left(L_{\mathbf{c}}\right)$ corresponding to $\mathfrak{l}_{1}, \mathfrak{c}\left(l_{1 c}, \mathfrak{c}_{\mathbf{c}}\right)$ respectively. We call $L_{\mathrm{c}}$ quasi-simply connected (q.s.c.) if $L_{1 \mathrm{c}} \cap C_{\mathrm{c}}=\{1\}$ and if $L_{1 \mathrm{c}}$ is simply connected. We say that $L$ is q.s.c. if it has a q.s.c. complexification.

Fix a complexification $j: L \rightarrow L_{\mathbf{c}}$ and let $\mathfrak{h}$ be a Cartan subalgebra of $\mathfrak{l}$. Let $A$ and $A_{\mathbf{c}}$ be the Cartan subgroups of $L$ and $L_{\mathbf{c}}$ corresponding to $\mathfrak{h}$ and $\mathfrak{h}_{\mathbf{c}}$ (that is, the centralizers of $\mathfrak{h}$ and $\mathfrak{h}_{\mathbf{c}}$ in $G$ and $G_{\mathbf{c}}$ respectively). Clearly. $j(A) \subseteq A_{\mathbf{c}}$. It is known that $A_{c}$ is connected [ $\mathbf{3}(\mathrm{j})$, corollary to Lemma 27]. If $\lambda$ is a linear functional on $\mathfrak{h}_{\mathbf{c}}$, there exists at most one complex analytic homomorphism

$$
\xi_{\lambda}: A_{\mathbf{c}} \rightarrow \mathbf{C}
$$

such that for every $H$ in $\mathfrak{h}_{\mathbf{c}}$

$$
\xi_{\lambda}(\exp H)=e^{\lambda(H)}
$$

We also write $\xi_{\lambda}$ for the homomorphism

$$
\xi_{\lambda} \circ j: A \rightarrow \mathbf{C}
$$

$\xi_{\lambda}$ can be seen to be independent of the complexification $L_{c}$ used, provided that $\xi_{\lambda}$ is defined on that complexification.

Clearly $\xi_{\alpha}$ exists for any root $\alpha$ of $\left(\mathfrak{l}_{\mathbf{c}}, \mathfrak{h}_{\mathrm{c}}\right)$. If $P_{\mathfrak{h}}$ is the set of positive roots relative to some ordering, let

$$
\rho=\frac{1}{2} \sum_{\alpha \in P_{\natural}} \alpha
$$

It is easy to see that the question of the existence of $\xi_{\rho}$ is independent of the ordering of the roots of $\left(l_{\mathbf{c}}, \mathfrak{h}_{\mathbf{c}}\right)$ and of the choice of Cartan subalgebra $\mathfrak{h}$. If $\xi_{\rho}$ exists we call $L_{c}$ acceptable. We say that $L$ is acceptable if it has an acceptable complexification.

If $L_{\mathbf{c}}$ is q.s.c., it is known that it is acceptable [ $\mathbf{3}(\mathbf{j})$, Lemma 29]. If $L_{1} \cap C$ is finite, it is clear that $L$ has a finite, and hence acceptable, cover.

Suppose $L$ is a compact, connected acceptable Lie group with Lie algebra $\mathfrak{l}$. Let $\mathfrak{h}, P_{\mathfrak{h}}, A$, and $\rho$ be defined as above. For each $\alpha$ define an element $H_{\alpha}$ in $\mathfrak{h}_{\mathfrak{c}}$ by

$$
B\left(H_{\alpha}, H\right)=\alpha(H)
$$

for all $H$ in $\mathfrak{h}_{\mathbf{c}}$, where $B$ is the Killing form of $\mathfrak{h}_{\mathbf{c}}$ restricted to $\mathfrak{h}_{\mathbf{c}}$. Put

$$
\tilde{\omega}=\prod_{\alpha \in P} H_{\alpha}
$$

$\tilde{\omega}$ is in $S\left(\mathfrak{h}_{c}\right)$, the symmetric algebra on $\mathfrak{h}_{c}$. Let $\Pi$ be the lattice of linear functionals

$$
\lambda:(-1)^{1 / 2} \mathfrak{h} \rightarrow \mathbf{R}
$$

for which $\xi_{\lambda}$ exists. Let $\Pi^{\prime}=\{\lambda \in \Pi: \tilde{\omega}(\lambda) \neq 0\}$. If $W$ is the Weyl group of $\left(l_{c} / \mathfrak{h}_{\mathrm{c}}\right), W$ acts on $(-1)^{1 / 2} \mathfrak{h}$. Then $W$ acts on $\Pi$ as follows:

$$
s \mu(H)=\mu\left(s^{-1} H\right)
$$

for $\mu$ in $\Pi, s$ in $W$, and $H$ in $(-1)^{1 / 2} \mathfrak{h}$. For $s$ in $W$, put $\varepsilon(s)=(-1)^{n(s)}$, where $n(s)$ is the number of positive roots that are mapped by $s$ into negative roots. For $h$ a regular element of $A$, put

$$
\Delta(h)=\xi_{\rho}(h) \prod_{\alpha \in P_{\mathfrak{k}}}\left(1-\xi_{\alpha}\left(h^{-1}\right)\right) .
$$

LEMMA 1. There is a map $\mu \rightarrow \sigma(\mu)$ from $\Pi^{\prime}$ onto the set of unitary equivalence classes of irreducible representations of $L . \sigma\left(\mu_{1}\right)=\sigma\left(\mu_{2}\right)$ if and only if $\mu_{1}=s \mu_{2}$ for some $s$ in $W$. Furthermore, if $h$ is a regular element of $A$,

$$
\operatorname{tr} \sigma(\mu)(h)=(\operatorname{sign} \tilde{\omega}(\mu)) \cdot \Delta(h)^{-1} \cdot\left(\sum_{s \in w} \varepsilon(s) \xi_{s \mu}(h)\right)
$$

Also there exists a constant $c_{L}$, independent of $\mu$, such that

$$
\operatorname{dim} \sigma(\mu)=c_{L}|\tilde{\omega}(\mu)|
$$

Finally, if $\mu$ is in $\Pi^{\prime}$, and $B(\mu, \alpha)>0$ for each $\alpha$ in $P_{\mathfrak{h}}$, then $\mu-\rho$ is the highest weight of the representation of the Lie algebra $\mathfrak{l}_{c}$ corresponding to $\sigma(\lambda)$.

Proof. Since $L$ has a finite q.s.c. cover, we will assume without loss of generality that $L$ is q.s.c. We can assume further that $L$ is semisimple. Then $L$ is simply connected, so $\Pi$ is precisely the lattice of weights of $\mathfrak{h}[3(j)$, Lemma 29]. If $\mu^{\prime}$ is a dominant integral function (in the terminology of [ $\left.5, \mathrm{p} .215\right]$ ), and if $\mu=\mu^{\prime}+\rho$, then $B(\mu, \alpha)>0$ for any $\alpha$ in $P_{\mathfrak{h}}$ so $\mu$ is in $\Pi^{\prime}$. Conversely, if $\mu$ is in $\Pi^{\prime}$, there exists a unique $s$ such that $B(s \mu, \alpha)>0$ for each $\alpha$ in $P_{\mathfrak{h}}$. Then $\mu^{\prime}=\mu-\rho$ is a dominant integral function on $\mathfrak{h}$. This demonstrates the relation between $\mu$ and the highest weight of $\sigma(\mu)$. The correspondence between representations and dominant integral functions is well known (see [5, Chapter VII]).

The other two statements of the lemma follow from the Weyl character formula [5, p. 255] and the Weyl dimension formula [5, p. 257].

Now let us return to our group $G$. By going to a finite cover we can assume that $G$ is q.s.c. and hence acceptable. Thus, if $j: \mathfrak{g} \bar{\subset} \mathfrak{g}_{\mathbf{c}}$ and $G_{\mathbf{c}}$ is a simply connected analytic group with Lie algebra $g_{c}$ then $j$ extends to a homomorphism

$$
j: G \rightarrow G_{\mathbf{c}} .
$$

Now $K$ is reductive. Therefore, by going to a further finite cover of $G$, we may also assume that $K$ is acceptable.

If we understand the harmonic analysis of a finite cover, $\tilde{G}$, of $G$ then we understand the theory for $G$. We merely throw out those unitary representations of $\tilde{G}$ which are nontrivial on the kernel of the covering projection. Therefore, the above two assumptions can be made with no loss of generality.

There are two possibilities for $G$.
Case I. There exists a Cartan subalgebra $\mathfrak{b}$ of $\mathfrak{g}$ such that $\mathfrak{b}$ is contained in $\mathfrak{k}$. We can assume that $\mathfrak{b}$ has been chosen so that it contains $\mathfrak{a}_{\mathfrak{f}}$. Then it is known that $\{\mathfrak{b}, \mathfrak{a}\}$ is a set of representatives of conjugacy classes of Cartan subalgebras of $\mathfrak{g}$.

Case II. Such a $\mathfrak{b}$ does not exist. Then there is only one conjugacy class of Cartan subalgebras and it is represented by $\mathfrak{a}$.

We shall try as far as possible to deal with these two cases together. Whenever we speak of $\mathfrak{b}$, we shall be implicitly referring to Case I. However, any mention of $\mathfrak{a}$, unless otherwise stated, will refer to either case.

Let $B$ be the Cartan subgroup of $G$ corresponding to $\mathfrak{b}$. Since it is a maximal abelian subgroup in the compact connected Lie group $K$, it is connected [4, Corollary 2.7, p. 247].

Let $A$ be the Cartan subgroup of $G$ corresponding to $\mathfrak{a}$. Then

$$
A=A_{I} A_{\mathfrak{p}}
$$

where $A_{\mathfrak{p}}=\exp \mathfrak{a}_{\mathfrak{p}}$, and $A_{I}$ is contained in $K$. In Case II, $A_{I}$ is a Cartan subgroup of $K$ and is connected. Otherwise, $A_{I}$ may not be connected. In any case, let $\mathfrak{m}$ and $M$ be the centralizers of $\mathfrak{a}_{\mathfrak{p}}$ in $\mathfrak{k}$ and $K$, respectively. Then $M$ is compact with a finite number of connected components.

Fix compatible orders on the real dual spaces of $\mathfrak{a}_{\mathfrak{p}}$ and $\mathfrak{a}_{\mathfrak{p}}+(-1)^{1 / 2} \mathfrak{a}_{\mathfrak{p}}$. Let $P$ be the set of positive roots of $\left(\mathfrak{g}_{c}, a_{c}\right)$ relative to this order. Let $P_{+}$be the set of roots in $P$ which do not vanish on $\mathfrak{a}_{\mathfrak{p}}$ and let $P_{M}$ equal $P-P_{+}$. $\mathfrak{a}_{\mathfrak{f}}$ is a Cartan subalgebra of the reductive Lie algebra $\mathfrak{m}$ and we can regard $P_{M}$ as the set of positive roots of ( $\mathfrak{m}, \mathfrak{a}_{\mathfrak{k}}$ ).

Let $M^{0}$ and $A_{I}^{0}$ be the connected components of $M$ and $A_{I}$. Let $W$ and $W_{1}$ be the Weyl groups of ( $\mathfrak{g} / \mathfrak{a}$ ) and ( $\mathfrak{m} / \mathfrak{a}_{\mathfrak{k}}$ ), respectively. Now in any connected component of $M$, it is possible to choose an element $\gamma_{1}$ such that

$$
\operatorname{Ad} \gamma_{1} \cdot \mathfrak{a}_{\mathfrak{l}}=\mathfrak{a}_{\mathfrak{e}}
$$

But Ad $\gamma_{1}$ leaves $\mathfrak{a}_{\mathfrak{p}}$ pointwise fixed. Therefore, the action of $\gamma_{1}$ on $\mathfrak{a}_{\mathfrak{p}}$ can be regarded as coming from an element of the subgroup of $W$ generated by those roots in $P$ which vanish on $a_{p}$. That is, the action of $\operatorname{Ad} \gamma_{1}$ on $\mathfrak{a}_{\mathfrak{t}}$ is the same as for some element in $W_{1}$. Therefore, we can choose a new element $\gamma$, in the same component of $M$, that leaves $\mathfrak{a}_{\mathfrak{t}}$ pointwise fixed. This means that $\gamma$ is in $A_{I}$. Therefore, $A_{I}$ has the same number of connected components as $M$.

As usual, let

$$
\rho=\frac{1}{2} \sum_{\alpha \in P} \alpha, \quad \rho_{M}=\frac{1}{2} \sum_{\alpha \in P_{M}} \alpha .
$$

Then since $G$ is acceptable, it is known that $M^{0}$ is also acceptable and that for any $a_{1}$ in $A_{I}^{0}$

$$
\begin{equation*}
\xi_{\rho}\left(a_{1}\right)=\xi_{\rho_{M}}\left(a_{1}\right) \tag{2.1}
\end{equation*}
$$

[see 3(j), Lemma 30].
Let $\mathscr{E}_{M}$ be the set of equivalence classes of irreducible unitary representations of $M$. Let $C$ be the set of irreducible characters of the group $A_{I}$ (the set of characters coming from irreducible representations of $A_{I}$ ). For $\zeta$ in $C$ and $a$ in $A_{I}$ write $\langle\zeta, a\rangle$ for the value of $\varsigma$ at $a$. It is clear that $W_{1}$ operates on $C$.

Put $\tilde{\omega}^{\mathfrak{m}}=\prod_{\alpha \in P_{M}} H_{\alpha}$ and let $L_{1}$ be the lattice of real linear functionals, $\mu$, on $(-1)^{1 / 2} \mathfrak{a}_{k}$ such that $\xi_{\mu}$ exists. Let $L_{1}^{\prime}=\left\{\mu \in L_{1}: \tilde{\omega}^{m}(\mu) \neq 0\right\}$. Let $Z(A)=$ $\left\{\gamma \in A_{I}: j(\gamma) \in \exp (-1)^{1 / 2} \mathfrak{a}_{\mathrm{R}}\right\}$. Then $Z(A)$ is a finite subgroup of $A_{I}$. It is known that if $\gamma$ is in $Z(A)$ and $m$ is in $M^{0}$, then $\gamma$ and $m$ commute [ $3(\mathrm{j})$, Lemma 51]. Also $Z(A) A_{I}^{0}=A_{I}$, by [3(k), Lemma 20], so $Z(A) M^{0}=M$. Let $Z(A)^{0}=$ $Z(A) \cap A_{I}^{0} . Z(A)^{0}$ is a central subgroup of both $Z(A)$ and $M^{0}$. Then $M$ is the central product of $M^{0}$ and $Z(A)$ with respect to $Z(A)^{0}$ (see [2, p. 29]). Thus if $\bar{M}=M^{0} \times Z(A)$ and $\overline{Z(A)}=\left\{\left(\gamma, \gamma^{-1}\right): \gamma \in Z(A)^{0}\right\}$ then $\overline{Z(A)}{ }^{0}$ is a discrete normal subgroup of $\bar{M} . M$ is isomorphic to $\bar{M} / \overline{Z(A)}$. Similarly, if $\overline{A_{I}}=A_{I}^{0} \times$ $Z(A)$, then $A_{I}$ is isomorphic to $\left.\overline{A_{I}} / \overline{Z(A}\right)^{0}$. Therefore, irreducible representations of $M$ (or $A_{I}$ ) are in one-to-one correspondence with representations of $\bar{M}$ (or
$A_{I}$ ) which are trivial on $\overline{Z(A)}^{0}$. An irreducible representation of $\overline{A_{I}}$ is of the form $\xi_{\mu} \otimes \delta$, where $\mu$ is in $L_{1}$ and $\delta$ is an irreducible representation of $Z(A)$. Let $C^{\prime}$ be the set of irreducible characters $\zeta$ in $C$ that come from representations $\xi_{\mu} \otimes \delta$ of $\bar{M}$ for which $\mu$ is actually in $L_{1}^{\prime}$. If $\mu$ and $\varsigma$ are so related, we shall write $\mu=\mu_{\varsigma}$. We would like to prove a lemma which will relate the representations in $\mathscr{E}_{M}$ with characters in $C^{\prime}$.

Let $\sigma$ be an arbitrary representation of $M$. Then

$$
\begin{equation*}
\sigma=\sigma_{0} \times \varepsilon \tag{2.2}
\end{equation*}
$$

where $\sigma_{0}$ and $\varepsilon$ are irreducible representations of $M^{0}$ and $Z(A)$, respectively, such that for any $\gamma_{0}$ in $Z(A)^{0}, \sigma_{0}\left(\gamma_{0}\right) \otimes \varepsilon\left(\gamma_{0}^{-1}\right)$ is the identity. $Z(A)^{0}$ is in the center of both $M^{0}$ and $Z(A)$ so $\sigma_{0}\left(\gamma_{0}\right)$ and $\varepsilon\left(\gamma_{0}\right)$ are both scalars. Therefore

$$
\sigma_{0}\left(\gamma_{0}\right)=\varepsilon\left(\gamma_{0}\right)
$$

Suppose that $\sigma_{0}=\sigma_{0}(\mu)$ in the notation of Lemma 1. $\mu$ is a linear functional in $L_{1}^{\prime}$. Then there exists an $s$ in $W_{1}$ such that $s \mu-\rho$ is the highest weight for $\sigma_{0}$. Let $\gamma_{0}$ be an element in $Z(A)^{0}$. By looking at the action of $\xi_{s \mu}-\rho\left(\gamma_{0}\right)$ on a highest weight vector for $\sigma_{0}$ we see that the scalar $\sigma_{0}\left(\gamma_{0}\right)$ is equal to $\xi_{s \mu}-\rho\left(\gamma_{0}\right)$. Therefore

$$
\varepsilon\left(\gamma_{0}\right)=\sigma_{0}\left(\gamma_{0}\right)=\xi_{s \mu}\left(\gamma_{0}\right) \xi_{\rho}\left(\gamma_{0}^{-1}\right)=\xi_{\mu}\left(s^{-1} \gamma_{0}\right) \xi_{0}\left(\gamma_{0}^{-1}\right)
$$

However, $\gamma_{0}$ is in the center of $M$ so $s^{-1} \gamma_{0}=\gamma_{0}$. Therefore

$$
\begin{equation*}
\varepsilon\left(\gamma_{0}\right) \xi_{\rho}\left(\gamma_{0}\right)=\xi_{\mu}\left(\gamma_{0}\right) \tag{2.3}
\end{equation*}
$$

for any $\gamma_{0}$ in $Z(A)^{0}$.
For any $\gamma$ in $Z(A)$, define

$$
\begin{equation*}
\delta(\gamma)=\varepsilon(\gamma) \xi_{\rho}(\gamma) \tag{2.4}
\end{equation*}
$$

This is an irreducible representation of $Z(A)$ and by (2.3), $\xi_{\mu} \otimes \delta$ can be regarded as in irreducible representation of $A_{I}$. Let

$$
\begin{equation*}
\left\langle s, \gamma a_{0}\right\rangle=s_{\mu}\left(a_{0}\right) \cdot \operatorname{tr} \delta(\gamma) \tag{2.5}
\end{equation*}
$$

for $a_{0}$ in $A_{I}^{0}, \gamma$ in $Z(A)$. $\varsigma$ is an element in $C^{\prime}$ and $\mu=\mu_{\varsigma}$. Therefore, given a $\sigma$ in $\mathscr{E}_{M}$, we have constructed an element $\zeta$ in $C^{\prime}$. We write $\sigma=\sigma(\varsigma)$.

Conversely, let us start with an element in $C^{\prime}$. By working backward we can show that there is a unique element $\sigma$ in $\mathscr{E}_{M}$ such that $\sigma=\sigma(\varsigma)$.

Suppose that $a_{0}$ is a regular element of $A_{I}^{0}$ and $\gamma$ is in $Z(A)$. We wish to compute the trace of $\sigma\left(a_{0} \gamma\right)$. Define

$$
\Delta_{M}\left(a_{0}\right)=\xi_{\rho}\left(a_{0}\right) \cdot \prod_{\alpha \in P_{M}}\left(1-\xi_{\alpha}\left(a_{0}^{-1}\right)\right)
$$

In the above notation

$$
\operatorname{tr} \sigma\left(a_{0} \gamma\right)=\operatorname{tr} \sigma_{0}\left(a_{0}\right) \cdot \operatorname{tr} \varepsilon(\gamma)
$$

But from Lemma 1,

$$
\operatorname{tr} \sigma\left(a_{0}\right)=\left(\operatorname{sign} \tilde{\omega}^{m}(\mu)\right) \cdot \Delta_{M}\left(a_{0}\right)^{-1} \cdot\left(\sum_{s \in W_{1}} \varepsilon(s) \xi_{s \mu}\left(a_{0}\right)\right)
$$

Therefore, the trace of $\sigma\left(a_{0} \gamma\right)$ is equal to

$$
\left(\operatorname{sign} \tilde{\omega}^{\mathfrak{m}}(\mu)\right) \cdot \Delta_{M}\left(a_{0}\right)^{-1} \cdot\left(\sum_{s \in W_{1}} \varepsilon(s)\left\langle s \varsigma, a_{0} \gamma\right\rangle\right) \xi_{\rho}\left(\gamma^{-1}\right)
$$

Now it is easy to show that if $\gamma_{c}$ is in $j(Z(A))$ then $\left(\gamma_{c}\right)^{2}=1$. Therefore if $\gamma$ is in $Z(A), \xi_{\rho}(\gamma)=\xi_{\rho}(\gamma)^{-1}$. For future convenience, we rewrite the trace of $\sigma\left(a_{0} \gamma\right)$ as

$$
\begin{equation*}
\left(\operatorname{sign} \tilde{\omega}^{\mathrm{m}}(\mu)\right) \cdot \Delta_{M}\left(a_{0}\right)^{-1} \cdot\left(\sum_{s \in W_{1}} \varepsilon(s)\left\langle s \varsigma, a_{0} \gamma\right\rangle\right) \xi_{\rho}(\gamma) \tag{2.6}
\end{equation*}
$$

Lemma 2. There is a $\operatorname{map} \varsigma \rightarrow \sigma(\varsigma)$ from $C^{\prime}$ onto $\mathscr{E}_{M} . \sigma\left(\varsigma_{1}\right)=\sigma\left(\varsigma_{2}\right)$ if and only if $s \varsigma_{1}=\varsigma_{2}$ for some $s$ in $W_{1}$. If $a_{0}$ is a regular element in $A_{I}^{0}$ and $\gamma$ is in $Z(A)$ then the trace of $\sigma(\zeta)\left(a_{0} \gamma\right)$ equals

$$
\left(\operatorname{sign} \tilde{\omega}^{\mathfrak{m}}(\mu)\right) \cdot \Delta_{M}\left(a_{0}\right)^{-1} \cdot\left(\sum_{s \in W_{1}} \varepsilon(s)\left\langle s \zeta, a_{0} \gamma\right\rangle\right) \cdot \xi_{\rho}(\gamma)
$$

Also, there exists a constant $C_{M}$, independent of $\zeta$, such that

$$
\operatorname{dim} \sigma(\varsigma)=C_{M} \cdot\left|\tilde{\omega}^{\mathfrak{m}}\left(\mu_{\varsigma}\right)\right| \cdot \operatorname{dim} \varsigma .
$$

$\left(\operatorname{dim} \varsigma\right.$ means the dimension of the representation of $A_{I}$ of which $\varsigma$ is the character.)

Proof. The dimension formula follows from Lemma 1. All other statements in the lemma follow from the above discussion.

Let us say that the linear functional $\mu_{\varsigma}$ is associated with $\sigma$ if $\sigma=\sigma(\varsigma)$, in the above notation. For any $\sigma$ in $\mathscr{E}_{M}$ there are exactly [ $W_{1}$ ] associated real linear functionals on $\mathfrak{a}_{\mathrm{f}}$.

Now with $B$ there is associated a discrete series of unitary representations of $G$. With $A$ there is associated a continuous series. We shall describe these.

For the discrete series there is a formal analogy with Lemma 1. Let $\Sigma$ be the set of positive roots of ( $\mathfrak{g}_{\mathbf{c}}, \mathfrak{b}_{\mathbf{c}}$ ) relative to some order. For any $\alpha$ in $\Sigma$ define $H_{\alpha}$ in $(-1)^{1 / 2} \mathfrak{b}$ by the formula

$$
B\left(H_{\alpha}, H\right)=\alpha(H)
$$

for any $H$ in $\mathfrak{b}_{\mathbf{c}}$. Put $\tilde{\omega}^{\mathfrak{b}}=\prod_{\alpha \in \Sigma} H_{\alpha}$. Let $L$ be the lattice of real linear functionals, $\lambda$, on $(-1)^{1 / 2} \mathfrak{b}$ such that $\xi_{\lambda}$ exists. Let $L^{\prime}==\left\{\lambda \in L: \tilde{\omega}^{\mathfrak{b}}(\lambda) \neq 0\right\}$. Let $N(B)$ be the normalizer of $B$ in $G$. Define

$$
W_{G}=W_{G, \mathfrak{b}}=N(B) / B
$$

This is a finite group. It acts on $B$ and therefore on $L$.

An irreducible representation $\pi$ of $G$ on a Hilbert space $\mathscr{H}$ is said to be squareintegrable if there exist nonzero vectors $\Phi_{1}, \Phi_{2}$ if $\mathscr{H}$ such that $\left(\Phi_{1}, \pi(x) \Phi_{2}\right)$ is a square-integrable function of $x$. If $\pi$ and $\pi^{\prime}$ are square-integrable representations on $\mathscr{H}$ and $\mathscr{H}^{\prime}$ and if $\pi$ and $\pi^{\prime}$ are not unitarily equivalent, then for $\Phi_{1}, \Phi_{2}$ in $\mathscr{H}$ and $\Phi_{1}^{\prime}, \Phi_{2}^{\prime}$ in $\mathscr{H}^{\prime}$,

$$
\begin{equation*}
\int_{G}\left(\Phi_{1}, \pi(x) \Phi_{2}\right)\left(\pi^{\prime}(x) \Phi_{2}^{\prime}, \Phi_{1}^{\prime}\right) d x=0 \tag{2.7}
\end{equation*}
$$

On the other hand, there is a number $\beta(\pi)$, the formal degree of $\pi$, such that for every $\Phi_{1}, \Phi_{2}, \Psi_{1}, \Psi_{2}$ and $\mathscr{H}$,

$$
\begin{equation*}
\int_{G}\left(\Phi_{1}, \pi(x) \Phi_{2}\right)\left(\pi(x) \Psi_{2}, \Psi_{1}\right) d x=\beta(\pi)^{-1}\left(\Phi_{1}, \Psi_{1}\right)\left(\Psi_{2}, \Phi_{2}\right) . \tag{2.8}
\end{equation*}
$$

These are the Schur orthogonality relations on $G$. They are proved in $[\mathbf{3}(\mathrm{d})$, Theorem 1].

Let $\mathscr{E}_{d}$ be the set of unitary equivalence classes of square-integrable representations of $G$. Harish-Chandra gives a map $\lambda \rightarrow \omega(\lambda)$ from $L^{\prime}$ onto $\mathscr{E}_{d}$ [see $3(\mathrm{l})$, Theorem 16]. $\omega\left(\lambda_{1}\right)=\omega\left(\lambda_{2}\right)$ if and only if there is an $s$ in $W_{G}$ such that $s \lambda_{1}=\lambda_{2}$. Finally, there is a constant $C_{G}$, independent of $\lambda$, such that

$$
\beta(\omega(\lambda))=C_{G}\left|\tilde{\omega}^{\mathbf{b}}(\lambda)\right|
$$

Lemma 3. $\left\{\beta(\omega): \omega \in \mathscr{E}_{d}\right\}$ is bounded away from zero.
Proof. It is clearly enough to show that for any $\alpha$ in $\Sigma,\left\{\lambda\left(H_{\alpha}\right)\right\}_{\lambda \in L^{\prime}}$ is bounded away from zero. Let $\tilde{L}$ be the lattice of real linear functionals on $(-1)^{1 / 2} \mathfrak{b}$ generated by the roots. Then it is known that $L / \tilde{L}$ is isomorphic to the center of $G$, which is finite. It is also known that $\left\{\tilde{\lambda}\left(H_{\alpha}\right)\right\}_{\tilde{\lambda} \in \tilde{L}}$ is a lattice in R. Therefore $\left\{\lambda\left(H_{\alpha}\right)\right\}_{\lambda \in L}$ is also a lattice in $\mathbf{R}$. But if $\lambda$ is in $L^{\prime}, \lambda\left(H_{\alpha}\right) \neq 0$, so the lemma follows.

Now we shall-describe the continuous series. There is a linear functional $\mu_{0}$ from $\mathfrak{a}_{p}$ to $\mathbf{R}$ such that the restriction of any root in $P_{+}$to $a_{p}$ is either $\mu_{0}$ or $2 \mu_{0}$. Fix $H_{0}$ in $\mathfrak{a}_{\mathfrak{p}}$ so that $\mu_{0}\left(H_{0}\right)=1$. Extend the definition of $\mu_{0}$ to $\mathfrak{a}$ by letting it equal zero on $\mathfrak{a}_{\mathfrak{k}}$.

Let $\mathfrak{n}_{\mathbf{c}}=\sum_{\alpha \in P_{+}} \mathbf{C} X_{\alpha}$, where for any $\alpha$ in $P, X_{\alpha}$ is a fixed root vector. Let $\mathfrak{n}=\mathfrak{n}_{\mathrm{c}} \cap \mathfrak{g}$. Let $N$ be the analytic subgroup of $G$ corresponding to $\mathfrak{n}$. It is well known (see [4, p. 373]) that the map

$$
(k, a, n) \rightarrow k a n, \quad k \in K, a \in A_{\mathfrak{p}}, n \in N
$$

is a diffeomorphism of $K \times A_{\mathfrak{p}} \times N$ with $G$. For $f$ in $C_{0}^{\infty}(G)$,

$$
\begin{equation*}
\int_{G} f(x) d x=\int_{K \times A_{\mathfrak{p}} \times N} f(k a n) e^{2 \rho(\log a)} d k d a d n \tag{2.9}
\end{equation*}
$$

for a suitable normalization of the Haar measure $d x$. If $x=k a n$, write $K(x)=k$ and $H(x)=\log a$.

It is clear that $P=M A_{\mathfrak{p}} N$ is a subgroup of $G$. If $\sigma$ in $\mathscr{E}_{M}$ acts on a finite dimensional Hilbert space $V_{\sigma}$, and if $\Lambda$ is in $\mathbf{R}$, then the map $\sigma_{\Lambda}$ from $P$ into $\operatorname{End}\left(V_{\sigma}\right)$ given by

$$
\sigma_{\Lambda}\left(m \cdot \exp t H_{0} \cdot n\right)=\sigma(m) e^{-i \Lambda t}, \quad m \in M, n \in N, t \in \mathbf{R}
$$

is an irreducible unitary representation of $P$. (We shall sometimes write $i$ instead of $(-1)^{1 / 2}$.) Let $\pi_{\sigma, \Lambda}$ be the unitary representation of $G$ on the Hilbert space $\mathscr{H}_{\sigma, \Lambda}$ obtained by inducing $\sigma_{\Lambda}$ from $P$ to $G$.

Then $\mathscr{H}_{\sigma, \Lambda}$ is the set of functions $\Phi$ from $G$ into $V_{\sigma}$ such that

$$
\begin{equation*}
\Phi\left(x \xi^{-1}\right)=\sigma_{\Lambda}(\xi) \Phi(x), \quad x \in G, \xi \in P \tag{2.10}
\end{equation*}
$$

$$
\begin{equation*}
\int_{K}\|\Phi(k)\|^{2} d k<\infty \tag{2.11}
\end{equation*}
$$

The inner product on $\mathscr{H}_{\sigma, \Lambda}$ is given by

$$
(\Phi, \Psi)=\int_{K}(\Phi(k), \Psi(k))_{V_{\sigma}} d k, \quad \Phi, \Psi \in \mathscr{H}_{\sigma, \Lambda}
$$

where $(,)_{V_{\sigma}}$ is the inner product in $V_{\sigma}$. If $\Phi$ is in $\mathscr{H}_{\sigma, \Lambda}, \pi_{\sigma, \Lambda}(y) \Phi$ is given by

$$
\begin{equation*}
\left(\pi_{\sigma, \Lambda}(y) \Phi\right)(x)=\Phi\left(y^{-1} x\right) e^{-\rho\left(H\left(y^{-1} x\right)\right)+\rho(H(x))}, \quad x, y \in G \tag{2.13}
\end{equation*}
$$

For any real $\Lambda$, and any $\Phi$ in $\mathscr{H}_{\sigma, \Lambda}$ we can define a function $\tilde{\Phi}$ from $K$ to $V_{\sigma}$ by restricting $\Phi$ to $K$. This identifies $\mathscr{H}_{\sigma, \Lambda}$ with a Hilbert space, $\mathscr{H}_{\sigma}$, of square-integrable functions from $K$ into $V_{\sigma} . \mathscr{H}_{\sigma}$ is independent of $\Lambda$. In fact, if $\pi_{\sigma}$ is the representation of $K$ obtained by inducing $\sigma$ to $K, \mathscr{H}_{\sigma}$ is the Hilbert space on which $\pi_{\sigma}$ acts. The above equivalence between $\mathscr{H}_{\sigma}$ and $\mathscr{H}_{\sigma, \Lambda}$ gives an intertwining operator between $\pi_{\sigma}$ and $\left.\pi_{\sigma, \Lambda}\right|_{K}$, the restriction of $\pi_{\sigma, \Lambda}$ to $K$.

Let $M^{\prime}$ be the normalizer of $\mathfrak{a}_{\mathfrak{p}}$ in $K . M$ is a normal subgroup of $M^{\prime} . M^{\prime} / M$ is a group consisting of two elements, $\{1, \delta\}$ say. $\delta$ acts on $\mathfrak{a}_{\mathfrak{p}}$ by reflection. $\delta$ also induces an automorphism of $M$, modulo the group of inner automorphisms. Therefore $\delta$ defines a bijection.

$$
\delta: \sigma \rightarrow \sigma^{\prime}
$$

of $\mathscr{E}_{M}$ onto itself. If we let $\delta$ act on $P$, we can transform the representation $\sigma_{\Lambda}$ into the representation $\left(\sigma^{\prime}\right)_{-\Lambda}$. Now, if $\Lambda$ is real and $\sigma$ is in $\mathscr{E}_{M}$, it is known that the representation $\pi_{\sigma, \Lambda}$ is equivalent to $\pi_{\sigma^{\prime},-\Lambda}$. Furthermore, the representations $\left\{\pi_{\sigma, \Lambda}\right\}, \sigma \in \mathscr{E}_{M}, \Lambda>0$ are all irreducible and inequivalent [see 1, Theorem 7; 2].

For each $\sigma$ in $\mathscr{E}_{M}$ and $\Lambda \neq 0$, let $N_{\sigma}(\Lambda)$ be a fixed unitary intertwining operator between the representations $\pi_{\sigma, \Lambda}$ and $\pi_{\sigma^{\prime},-\Lambda}$. Then

$$
N_{\sigma}(\Lambda) \pi_{\sigma, \Lambda}(x) N_{\sigma}(\Lambda)^{-1}=\pi_{\sigma^{\prime},-\Lambda}(x), \quad x \in G
$$

Notice that since $\pi_{\sigma, \Lambda}$ is irreducible,

$$
\begin{equation*}
N_{\sigma^{\prime}}(-\Lambda)=N_{\sigma}(\Lambda)^{-1} \tag{2.14}
\end{equation*}
$$

It will be convenient to assign a positive real number to any equivalence class of representations in either $\mathscr{E}_{d}$ or $\mathscr{E}_{M}$. If $\omega$ is in $\mathscr{E}_{d}$, choose $\lambda$ in $L^{\prime}$ such that $\omega=\omega(\lambda)$. The Killing form, $B$, of $\mathfrak{g}_{\mathrm{c}}$ can be regarded as a positive definite form on either $(-1)^{1 / 2} \mathfrak{b}$ or its real dual space. Then put

$$
|\omega|^{2}=B(\lambda, \lambda)
$$

Since $W_{G}$ acts on $(-1)^{1 / 2} \mathfrak{b}$ as a group of isometries under the Killing form, $|\omega|$ is well defined. Similarly, for $\sigma$ in $\mathscr{E}_{M}$, we define

$$
|\sigma|^{2}=B\left(\mu_{\sigma}, \mu_{\sigma}\right)
$$

where $\mu_{\sigma}$ is any real linear functional on $(-1)^{1 / 2} \mathfrak{a}_{\mathfrak{k}}$ associated with $\sigma$. $|\sigma|$ is well defined by the above argument.

Let $\mathscr{E}_{K}$ be the set of unitary equivalence classes of irreducible representations of $K$. Let $\mathfrak{h}$ be the subspace of $\mathfrak{k}$ which is equal to either $\mathfrak{b}$ or $\mathfrak{a}_{\mathfrak{l}}$, depending on whether we are in Case I or Case II. $\mathfrak{h}$ is a Cartan subalgebra of $\mathfrak{k}$. In either case, we have already ordered the dual space of $\mathfrak{h}$. $K$ is acceptable by assumption, so the representations in $\mathscr{E}_{K}$ can be indexed by certain real linear functionals on $(-1)^{1 / 2} \mathfrak{h}$ as in Lemma 1. If $\tau$ is in $\mathscr{E}_{K}$ and $\tau=\tau(\mu)$ for some real linear functional $\mu$ on $(-1)^{1 / 2} \mathfrak{h}$, then we write

$$
|\tau|^{2}=B(\mu, \mu)
$$

$|\tau|$ is well defined.
3. Plancherel formula for $L^{2}(G)$. In order to put the Plancherel formula for $G$ into the form we want, we must discuss characters of unitary representations of $G$. To do this we must introduce some more notation of Harish-Chandra.

For $t$ in $\mathbf{R}$, put $h_{t}=\exp t H_{0}$. For $g$ in $C_{0}^{\infty}\left(M^{0} A_{\mathfrak{p}}\right)$, write

$$
F_{g}^{M}\left(a_{0} h_{t}\right)=\Delta_{M}\left(a_{0}\right) \cdot \int_{M^{0} / A_{I}^{0}} g\left(m^{*-1} a_{0} h_{t} m^{*}\right) d m^{*}
$$

for $a_{0}$ in $A_{I}^{0}$ and $a_{0} h_{t}$ a regular element in $A$. Here $d m^{*}$ is the invariant measure on the homogeneous space $M^{0} / A_{I}^{0}$. It is known that there exists a constant $c_{1}>0$ such that for any $g$ in $C_{0}^{\infty}\left(M^{0} A_{\mathfrak{p}}\right)$

$$
\begin{equation*}
\int_{M^{0} \times \mathbf{R}} g\left(m_{0} h_{t}\right) d m_{0} d t=c_{1} \int_{A_{i}^{0} \times \mathbf{R}} \overline{\Delta_{M}\left(a_{0}\right)} \cdot F_{g}^{M}\left(a_{0} h_{t}\right) d a_{0} d t \tag{3.1}
\end{equation*}
$$

(see [3(j), Lemma 41]).
For $a$ in $A_{I}$ and $a h_{t}$ a regular element in $A$ write

$$
\begin{gathered}
\Delta\left(a h_{t}\right)=\xi_{\rho}\left(a h_{t}\right) \cdot \prod_{\alpha \in P}\left(1-\xi_{\alpha}\left(a h_{t}\right)^{-1}\right) \\
\varepsilon_{\mathbf{R}}\left(a h_{t}\right)=1 \quad \text { if } t>0, \quad=-1 \quad \text { if } t<0
\end{gathered}
$$

If $f$ is in $C_{0}^{\infty}(G)$ write

$$
F_{f}\left(a h_{t}\right)=\varepsilon_{R}\left(a h_{t}\right) \cdot \Delta\left(a h_{t}\right) \cdot \int_{G^{*}} f\left(x^{*-1} a h_{t} x^{*}\right) d x^{*}
$$

( $G^{*}$ is the homogeneous space $G / A_{0}$ where $A_{0}$ is the center of $A$. Let $d x^{*}$ be the $G$ invariant measure on $G^{*}$ ). It is clear that

$$
\begin{equation*}
F_{f}\left(s a h_{t}\right)=\varepsilon(s) F_{f}\left(a h_{t}\right), \quad s \in W_{1} . \tag{3.2}
\end{equation*}
$$

It is known that if $f$ is in $C_{0}^{\infty}(G)$, then $F_{f}$ extends to an infinitely differentiable function on $A$ (see [3(f), Lemma 40]). Furthermore, $F_{f}$ has compact support in A [3(f), Theorem 2].

Let $\tilde{\omega}=\tilde{\omega}^{\mathfrak{a}}=\prod_{\alpha \in P} H_{\alpha}$. Let

$$
q=\frac{1}{2}(\operatorname{dim} \mathfrak{g}-\operatorname{dim} \mathfrak{k}-\operatorname{rank} \mathfrak{g}+\operatorname{rank} \mathfrak{k}) .
$$

It is known that $q$ is an integer. If we are in Case II, the Cartan subgroup $A$ is fundamental, in the terminology of [3(f), p. 759]. Then Harish-Chandra's limit formula applies to $F_{f}[\mathbf{3}(\mathbf{f})$, Theorem 4]. Namely, there is a positive constant $c$ such that for any $f$ in $C_{0}^{\infty}(G)$

$$
\begin{equation*}
c f(1)=(-1)^{q} F_{f}(1 ; \tilde{\omega}) \tag{3.3}
\end{equation*}
$$

Here $\tilde{\omega}$ is to be regarded as a differential operator on $A$.
For $f$ in $C_{0}^{\infty}(G)$ define a function $g_{f}$ in $C_{0}^{\infty}\left(M A_{\mathfrak{p}}\right)$ by

$$
g_{f}\left(m h_{t}\right)=e^{\rho\left(t H_{0}\right)} \int_{N} \int_{K} f\left(k m h_{t} n k^{-1}\right) d k d n, \quad m \in M, t \in \mathbf{R}
$$

For $\gamma$ in $Z(A)$ and $m_{0}$ in $M^{0}$ put

$$
g_{f, \gamma}\left(m_{0} h_{t}\right)=g_{f}\left(\gamma m_{0} h_{t}\right)
$$

Then in $[\mathbf{3}(\mathbf{j})$, Lemma 52$]$ it is shown that there is a constant $c_{2}>0$ such that

$$
\begin{equation*}
F_{f}\left(\gamma a_{0} h_{t}\right)=c_{2} \cdot \xi_{\rho}(\gamma) F_{g_{f, \gamma}}^{M}\left(a_{0} h_{t}\right) . \tag{3.4}
\end{equation*}
$$

While we are at it, we shall state another Jacobian formula which we shall need later in the paper (see [4, p. 381, Proposition 1.17]). The map from $K \times \mathfrak{a}_{\mathfrak{p}}^{+} \times K$ into $G$ given by

$$
\left(k_{1}, t H_{0}, k_{2}\right) \rightarrow k_{1} \cdot \exp t H_{0} \cdot k_{2}
$$

is a diffeomorphism onto $G$. (We write $a_{p}^{+}=\left\{t H_{0}: t>0\right\}$.) Furthermore, there is a constant $c>0$ such that for any $f$ in $C_{0}^{\infty}(G)$

$$
\begin{align*}
\int_{G} f(x) d x & =c \int_{0}^{\infty} \int_{K \times K} f\left(k_{1} \cdot \exp t H_{0} \cdot k_{2}\right)|D(t)| d k_{1} d k_{2} d t \\
& =\frac{c}{2} \int_{-\infty}^{\infty} \int_{K \times K} f\left(k_{1} \cdot \exp t H_{0} \cdot k_{2}\right)|D(t)| d k_{1} d k_{2} d t \tag{3.5}
\end{align*}
$$

Here $D(t)=\left(e^{t}-e^{-t}\right)^{r_{1}} \cdot\left(e^{2 t}-e^{-2 t}\right)^{r_{2}}$, where $r_{1}$ and $r_{2}$ are the number of roots in $P_{+}$which, when restricted to $\mathfrak{a}_{\mathfrak{p}}$, are respectively equal to $\mu_{0}$ and $2 \mu_{0}$.

Let $\pi$ be an irreducible unitary representation of $G$ on the Hilbert space $\mathscr{H}$. Let $f$ be a function in $C_{0}^{\infty}(G)$. It is known that the operator

$$
\pi(f)=\int_{G} f(x) \pi(x) d x
$$

is of trace class. The map

$$
f \rightarrow \operatorname{tr} \pi(f)
$$

is a distribution on $C_{0}^{\infty}(G)$ (see $[3(\mathbf{c}), \S 5]$ ). This distribution is called the character of $\pi$.

If $\sigma$ is in $\mathscr{E}_{M}$, and $\Lambda$ is in $\mathbf{R}$, let $\Theta_{\sigma, \Lambda}$ be the character of the representation $\pi_{\sigma, \Lambda}$. Let $m=\frac{1}{2}(\operatorname{dimg}-\operatorname{rank} \mathfrak{g})$. Choose $\varsigma$ in $C^{\prime}$ and the associated $\mu_{\varsigma}$ in $L_{1}$ such that $\sigma=\sigma(\varsigma)$ as in Lemma 2.

THEOREM 1. There exists a constant $c_{0}>0$ such that for every $f$ in $C_{0}^{\infty}(G)$, $\Theta_{\sigma, \Lambda}(f)$ is equal to

$$
c_{0}(-1)^{m+\iota}\left(\operatorname{sign} \tilde{\omega}^{\mathrm{m}}\left(\mu_{\varsigma}\right)\right) \cdot \int_{A_{I} \times \mathbf{R}} F_{f}\left(a h_{t}\right)\langle\varsigma, a\rangle e^{-i \Lambda t} d a d t
$$

where $\iota$ equals 1 or 0 depending on whether we are in Case I or Case II.
Proof. Let $A$ be the operator

$$
\int_{G} f(x) \pi_{\sigma, \Lambda}(x) d x
$$

on $\mathscr{H}_{\sigma, \Lambda}$. We want to compute the trace of $A . \operatorname{tr} A$ is equal to $\overline{\operatorname{tr} A^{*}}$ where $A^{*}$ is the adjoint operator of $A$ and the bar denotes complex conjugation. If $\Phi$ is in $\mathscr{H}_{\sigma, \Lambda}$ and $k_{1}$ is in $K$

$$
\begin{aligned}
\left(A^{*} \Phi\right)\left(k_{1}\right) & =\left(\int_{G} \overline{f(x)} \pi\left(x^{-1}\right) d x \cdot \Phi\right)\left(k_{1}\right) \\
& =\int_{G} \overline{f(x)} \Phi\left(x k_{1}\right) e^{-\rho\left(H\left(x k_{1}\right)\right)} d x \\
& =\int_{G} \overline{f\left(x k_{1}^{-1}\right)} \Phi(x) e^{-\rho(H(x))} d x
\end{aligned}
$$

Assume that the Haar measure on $A_{\mathfrak{p}}$ has been normalized so that $d h_{t}=d t$. Then by (2.9) the above integral equals

$$
\int_{K \times \mathbf{R} \times N} \overline{f\left(k h_{t} n k_{1}^{-1}\right)} e^{i \Lambda t} e^{\rho\left(t H_{0}\right)} \Phi(k) d k d t d n
$$

In this integral, substitute $k m$ for $k$ and integrate with respect to $M$. Then $\left(A^{*} \Phi\right)\left(k_{1}\right)$ equals

$$
\int_{K \times M \times \mathbf{R} \times N} \overline{f\left(k m h_{t} n k_{1}^{-1}\right)} \sigma\left(m^{-1}\right) e^{i \Lambda t} e^{\rho\left(t H_{0}\right)} \Phi(k) d m d t d n d k
$$

Now to deal further with this expression we consider the principal fiber bundle

$$
M \rightarrow K \rightarrow K / M
$$

The map $m \rightarrow \sigma\left(m^{-1}\right)$ defines a complex vector bundle $E_{\sigma}$ over $K / M$ with fiber $V_{\sigma}$, the space on which $\sigma$ acts. Let $F\left(k_{1}, k\right)$ be the function

$$
\int_{M \times \mathbf{R} \times N} \overline{f\left(k m h_{t} n k_{1}^{-1}\right)} \sigma\left(m^{-1}\right) e^{i \Lambda t} e^{\rho\left(t H_{0}\right)} d m d t d n
$$

Now it is easy to check that $M$ normalizes $N$ and that for fixed $m$ in $M$ the measures $d n$ and $d\left(m n m^{-1}\right)$ on $N$ are equal. Then for $\overline{m_{1}}, \bar{m}$ in $M$,

$$
F\left(k_{1} \overline{m_{1}}, k \bar{m}\right)=\sigma\left(\bar{m}_{1}^{-1}\right) F\left(k_{1}, k\right) \sigma(\bar{m})
$$

Therefore $F\left(k_{1}, k\right)$ can be regarded as a section of $E_{\sigma} \boxtimes E_{\sigma}^{*}$, where $E_{\sigma}^{*}$ is the adjoint bundle of $E_{\sigma}$ and $E_{\sigma} \boxtimes E_{\sigma}^{*}$ is the exterior tensor product of $E_{\sigma}$ and $E_{\sigma}^{*}$, a bundle with base space $K / M \times K / M$ and fiber $V_{\sigma} \otimes V_{\sigma}^{*}$.

Now there is a natural equivalence between $\mathscr{H}_{\sigma, \Lambda}$ and the space $\mathscr{H}_{\sigma}$ defined earlier. However, $\mathscr{H}_{\sigma}$ is the space of square-integrable sections of $E_{\sigma}$ with respect to a $K$-invariant measure on $K / M . F\left(k_{1}, k\right)$ can be regarded as the kernel of the linear operator $A^{*}$ on this space. Then for any $\Phi$ on $\mathscr{H}_{\sigma}$

$$
\left(A^{*} \Phi\right)\left(k_{1}\right)=\int_{K} F\left(k_{1}, k\right) \Phi(k) d k
$$

To evaluate the trace of $A^{*}$ we appeal to the following lemma.
Lemma 4. Let $X$ be a compact infinitely differentiable manifold of dimension $n$. Let $d x$ be a positive nowhere-vanishing differentiable $n$-form on $X$. If $E \rightarrow X$ is a differentiable Hilbert bundle of fiber dimension s, let $L^{2}(E)$ be the Hilbert space of square-integrable sections of $E$. If $F\left(x_{1}, x\right)$ is a continuous section of $E \boxtimes E^{*}, F\left(x_{1}, x\right)$ defines a bounded linear operator $F$ on $L^{2}(E)$ in the obvious manner. Then if $F\left(x_{1}, x\right)$ is differentiable in both variables, $F$ is of trace class. Furthermore

$$
\operatorname{tr} F=\int_{X}(\operatorname{tr} F(x, x)) d x
$$

Proof. Let $T$ be the closed unit $n$-cube with opposite sides identified. $T$ is an $n$-torus and there is a canonical $n$-form $d t$ on $T$. Let

$$
S=\left\{t \in \mathbf{R}^{n}:|t|<1\right\}
$$

$S$ is an open subset of $T$.
Choose a finite differentiable partition of unity $\left\{\Psi_{\alpha}\right\}_{\alpha \in I}$ and a collection $\left\{U_{\alpha}\right\}_{\alpha \in I}$ of open subsets of $X$ such that the support of $\Psi_{\alpha}$ is contained in $U_{\alpha}$. We assume that for every $(\alpha, \beta)$ in $I \times I$ there is a diffeomorphism $\lambda_{\alpha \beta}$ from $U_{\alpha} \cup U_{\beta}$ onto $S_{\alpha \beta}$, an open subset of $S$. It can be seen that with no loss of generality we may also assume that
(i) $\lambda_{\alpha \beta}^{*}(d t)=d x$.
(ii) If $E_{\alpha \beta}$ is the restriction of $E$ to $U_{\alpha} \cup U_{\beta}$, then $E_{\alpha \beta}$ is trivial.
(iii) The map $\lambda_{\alpha \beta}$ lifts to a bundle map

$$
\Lambda_{\alpha \beta}: E_{\alpha \beta} \rightarrow S_{\alpha \beta} \times \mathbf{R}^{s}
$$

which is an isomorphism between Hilbert bundles preserving the inner product on each fiber. (We assume that $\mathbf{R}^{s}$ is equipped with the natural scalar product.)

Let $F_{\alpha \beta}$ be the integral operator on $L^{2}(E)$ with kernel

$$
F_{\alpha \beta}\left(x_{1}, x\right)=\Psi_{\alpha}\left(x_{1}\right) F\left(x_{1}, x\right) \Psi_{\beta}(x)
$$

It is clear that if each $F_{\alpha \beta}$ is of trace class then so is $F$. In that case $\operatorname{tr} F=$ $\Sigma_{\alpha \beta} \operatorname{tr} F_{\alpha \beta}$. Furthermore

$$
\begin{aligned}
\int_{X}(\operatorname{tr} F(x, x)) d x & =\int_{X}\left(\sum_{\alpha} \Psi_{\alpha}(x)\right)(\operatorname{tr} F(x, x))\left(\sum_{\beta} \Psi_{\beta}(x)\right) d x \\
& =\sum_{\alpha \beta} \int_{X}\left(\operatorname{tr} F_{\alpha \beta}(x, x)\right) d x
\end{aligned}
$$

Therefore it is enough to prove our lemma for the operators $F_{\alpha \beta}$.
$L^{2}\left(E_{\alpha \beta}\right)$ is a closed subspace of $L^{2}(E)$. It is an invariant subspace for the operator $F_{\alpha \beta}$. $\quad F_{\alpha \beta}$ equals zero on the complement of $L^{2}\left(E_{\alpha \beta}\right)$ in $L^{2}(E)$, so the trace of $F_{\alpha \beta}$ is equal to the trace of the restriction of $F_{\alpha \beta}$ to $L^{2}\left(E_{\alpha \beta}\right)$. Let $\mathscr{E}\left(\mathbf{R}^{s}\right)$ be the space of linear transformations of $\mathbf{R}^{s}$. Use the map $\Lambda_{\alpha \beta}$ to transform $F_{\alpha \beta}\left(x_{1}, x\right)$ into a section $R\left(t_{1}, t\right)$ of $\left(S_{\alpha \beta} \times S_{\alpha \beta}\right) \times \mathscr{E}\left(\mathbf{R}^{s}\right)$. Then we can regard $R\left(t_{1}, t\right)$ as an element in $C^{\infty}(T \times T) \times \mathscr{E}\left(\mathbf{R}^{s}\right)$. We have reduced our lemma to the case where $X=T, d x=d t, E=T \times \mathbf{R}^{s}$ and $F=R$.

Let $\left\{\phi_{1}(t), \phi_{2}(t), \ldots\right\}$ be an orthonormal basis $L^{2}(T) \otimes \mathbf{R}^{s}$, consisting of functions of the form

$$
e^{2 \pi i(\nu, t)} \otimes v
$$

Here $\nu$ will be an $n$-tuple of integers and $v$ will be a unit vector in $\mathbf{R}^{s}$. Let

$$
\begin{equation*}
r_{i j}=\int_{T \times T}\left(R\left(t_{1}, t\right) \phi_{i}(t), \phi_{j}\left(t_{1}\right)\right) d t d t_{1} \tag{3.6}
\end{equation*}
$$

The above inner product is of course in $\mathbf{R}^{s}$. Since $R\left(t_{1}, t\right)$ is differentiable, we can show from the harmonic analysis of the group $T \times T$ that if $m_{1}, m_{2}$ are any positive integers,

$$
\begin{equation*}
\sup _{i j}\left|r_{i j}\right|(1+i)^{m_{1}}(1+j)^{m_{2}}<\infty . \tag{3.7}
\end{equation*}
$$

This shows that $R$ is of trace class.
If $v$ is in $\mathbf{R}^{s}$, then from (3.6) we can show that for any $t_{1}, t$ in $T$

$$
R\left(t_{1}, t\right) v=\sum_{i j} r_{i j}\left(\phi_{i}(t), v\right) \phi_{j}\left(t_{1}\right)
$$

Therefore

$$
\begin{equation*}
\operatorname{tr} R\left(t_{1}, t\right)=\sum_{i j} r_{i j}\left(\phi_{i}(t), \phi_{j}\left(t_{1}\right)\right) \tag{3.8}
\end{equation*}
$$

We now compute the trace of the operator $R$.

$$
\begin{aligned}
\operatorname{tr} R & =\sum_{i} r_{i i}=\sum_{i} \int_{T} r_{i i}\left(\phi_{i}(t), \phi_{i}(t)\right) d t \\
& =\sum_{i j} \int_{T} r_{i j}\left(\phi_{i}(t), \phi_{j}(t)\right) d t
\end{aligned}
$$

This last expression is absolutely convergent by (3.7). Therefore

$$
\operatorname{tr} R=\int_{T} \sum_{i j} r_{i j}\left(\phi_{i}(t), \phi_{j}(t)\right) d t
$$

By (3.8) this expression is equal to

$$
\int_{T}(\operatorname{tr} R(t, t)) d t .
$$

This completes the proof of Lemma 4.
Let us return to the proof of the theorem. By the lemma, $\operatorname{tr} A^{*}$ equals

$$
\int_{K \times M \times \mathbf{R} \times N} \overline{f\left(k m h_{t} n k_{1}^{-1}\right)} \cdot \operatorname{tr} \sigma\left(m^{-1}\right) e^{i \Lambda t} e^{\rho\left(t H_{0}\right)} d k d m d t d n
$$

Therefore

$$
\begin{aligned}
\operatorname{tr} A & =\overline{\operatorname{tr} A^{*}} \\
& =\int_{K \times M \times \mathbf{R} \times N} f\left(k m h_{t} n k^{-1}\right) \overline{\operatorname{tr} \sigma\left(m^{-1}\right)} e^{-i \Lambda t} e^{\rho\left(t H_{0}\right)} d k d m d t d n \\
& =\int_{M \times \mathbf{R}} g_{f}\left(m h_{t}\right) \cdot \operatorname{tr} \sigma(m) \cdot e^{-i \Lambda t} d m d t,
\end{aligned}
$$

since

$$
\overline{\operatorname{tr} \sigma\left(m^{-1}\right)}=\overline{\operatorname{tr} \sigma(m)^{*}}=\operatorname{tr} \sigma(m)
$$

Let $Z_{A}$ be a set of representatives of cosets of $Z(A) / Z(A)^{0}$. Then $M$ is diffeomorphic with $Z_{A} \times M^{0}$. Therefore the trace of $A$ equals

$$
\sum_{\gamma \in Z_{A}} \int_{M^{0} \times \mathbf{R}} g_{f}\left(\gamma m_{0} h_{t}\right) \cdot \operatorname{tr} \sigma\left(\gamma m_{0}\right) e^{-i \Lambda t} d m_{0} d t
$$

For any finite set $S$ let $[S]$ denote the number of elements in $S$. Then $[P]=m$. Recall that

$$
q=\frac{1}{2}(\operatorname{dim} \mathfrak{g}-\operatorname{dim} \mathfrak{k}-\operatorname{rank} \mathfrak{g}+\operatorname{rank} \mathfrak{k}) .
$$

Then

$$
\begin{aligned}
& q=\frac{1}{2}\left(\left[P_{+}\right]+1\right) \quad \text { in Case I }, \\
& q=\frac{1}{2}\left[P_{+}\right] \quad \text { in Case II. }
\end{aligned}
$$

Therefore

$$
\left[P_{M}\right]=[P]-\left[P_{+}\right]=m-2 q+\iota .
$$

If $a_{0}$ is in $A_{I}^{0}$ then

$$
\overline{\Delta_{M}\left(a_{0}\right)}=\Delta_{M}\left(a_{0}\right)(-1)^{\left[P_{M}\right]}=\Delta_{M}\left(a_{0}\right)(-1)^{m+\iota}
$$

Now for any $m_{0}$ in $M$

$$
\operatorname{tr} \sigma\left(m_{0}^{-1} \gamma a_{0} m_{0}\right)=\operatorname{tr} \sigma\left(\gamma a_{0}\right)
$$

Therefore, from (3.1) we see that $\operatorname{tr} A$ equals

$$
c_{1} \sum_{\gamma \in Z_{A}} \int_{A_{I}^{0} \times \mathbf{R}} F_{g_{f, \gamma}}^{M}\left(a_{0} h_{t}\right) \cdot \overline{\Delta_{M}\left(a_{0}\right)} \cdot \operatorname{tr} \sigma\left(\gamma a_{0}\right) \cdot e^{-i \Lambda t} d a_{0} d t
$$

By Lemma 2, this equals

$$
\begin{aligned}
&\left(\operatorname{sign} \tilde{\omega}^{\mathrm{m}}\left(\mu_{\mathrm{S}}\right)\right)(-1)^{m+\iota} c_{1} \sum_{\gamma \in Z_{A}} \int_{A_{I}^{0} \times \mathbf{R}} \xi_{\rho}(\gamma) F_{g_{f, \gamma}}^{M}\left(a_{0} h_{t}\right) \\
& \times e^{-i \Lambda t}\left(\sum_{s \in W_{1}} \varepsilon(s)\left\langle s \varsigma, \gamma a_{0}\right\rangle\right) d a_{0} d t .
\end{aligned}
$$

By formula (3.4) this expression then equals

$$
\begin{aligned}
&\left(\operatorname{sign} \tilde{\omega}^{\mathfrak{m}}\left(\mu_{\varsigma}\right)\right)(-1)^{m+\iota}\left(\frac{c_{1}}{c_{2}}\right) \sum_{\gamma \in Z_{A}} \int_{A_{1}^{0} \times \mathbf{R}} F_{f}\left(\gamma a_{0} h_{t}\right) e^{-i \Lambda t} \\
& \times\left(\sum_{s \in W_{1}} \varepsilon(s)\left\langle s \zeta, \gamma a_{0}\right\rangle\right) d a_{0} d t \\
&=\left(\operatorname{sign} \tilde{\omega}^{\mathfrak{m}}\left(\mu_{\varsigma}\right)\right)(-1)^{m+\iota}\left(\frac{c_{1}}{c_{2}}\right) \int_{A_{I} \times \mathbf{R}} F_{f}\left(a h_{t}\right) e^{-i \Lambda t} \\
& \times\left(\sum_{s \in W_{1}} \varepsilon(s)\langle s \zeta, a\rangle\right) d a d t
\end{aligned}
$$

Now if $s$ is in $W_{1}$, substitute $s a$ for $a$ in the above expression. From (3.2) we obtain the formula

$$
\operatorname{tr} A=\left(\operatorname{sign} \tilde{\omega}^{\mathfrak{m}}\left(\mu_{\varsigma}\right)\right) \cdot(-1)^{m+\iota} \cdot\left(\frac{c_{1}}{c_{2}}\right) \cdot\left[W_{1}\right] \cdot \int_{A_{I} \times \mathbf{R}} F_{f}\left(a h_{t}\right) e^{-i \Lambda t}\langle\varsigma, a\rangle d a d t
$$

This proves the theorem if we let $c_{0}=\left(c_{1} / c_{2}\right)\left[W_{1}\right]$.
For every $\varsigma$ in $C^{\prime}$ there is associated a unique $\mu_{\varsigma}$ in $L_{1}^{\prime}$. For any real $\Lambda$ we write

$$
\tilde{\omega}(\varsigma, \Lambda)=\tilde{\omega}\left(\mu_{\zeta}+i \Lambda \mu_{0}\right)
$$

For our discussion of the Plancherel formula it is necessary to examine this expression separately for Cases I and II. We have the formula

$$
\begin{equation*}
\tilde{\omega}(\varsigma, \Lambda)=\tilde{\omega}^{\mathfrak{m}}\left(\mu_{\varsigma}\right) \cdot \prod_{\alpha \in P_{+}}\left\langle\mu_{\varsigma}+i \Lambda \mu_{0}, H_{\alpha}\right\rangle . \tag{3.9}
\end{equation*}
$$

Now $P_{+}$is the union of the positive real roots, $P_{\mathbf{R}}$, and the positive complex roots $P_{\mathbf{c}}$. Let $\eta$ be the conjugation of $\mathfrak{g}_{\mathbf{c}}$ with respect to the real form $\mathfrak{g} . \eta$ acts as a permutation of period 2 on $P_{+}$. A root in $P_{+}$is fixed by $\eta$ if and only if it is a real root, so the positive complex roots occur in pairs. Since dim $\mathfrak{a}_{\mathfrak{p}}=1$, there can be at most one positive real root. Now it is known that $\mathfrak{a}$ has a real root if and only if $\mathfrak{a}$ is not fundamental. Therefore there exists a real root if and only if we are in Case I.

If $\alpha$ is a complex root and $\alpha^{\eta}$ is its conjugate root, then

$$
\left\langle\mu_{\varsigma}+i \Lambda \mu_{0}, H_{\alpha}\right\rangle \cdot\left\langle\mu_{\varsigma}+i \Lambda \mu_{0}, H_{\alpha}^{\eta}\right\rangle=-\left(\mu_{\varsigma}\left(H_{\alpha}\right)^{2}+\Lambda^{2} \mu_{0}\left(H_{\alpha}\right)^{2}\right)
$$

Therefore the sign of the real number

$$
\prod_{\alpha \in P_{\mathrm{c}}}\left\langle\mu_{\varsigma}+i \Lambda \mu_{0}, H_{\alpha}\right\rangle
$$

is equal to $(-1)^{\left[P_{\mathrm{c}}\right] / 2}$, which equals $(-1)^{q+\iota}$. Therefore

$$
\begin{align*}
& \tilde{\omega}(\varsigma, \Lambda) \cdot|\tilde{\omega}(\varsigma, \Lambda)|^{-1}=i(-1)^{q+\iota} \cdot \operatorname{sign} \Lambda \cdot \operatorname{sign} \tilde{\omega}^{m}\left(\mu_{\varsigma}\right) \quad \text { in Case I, } \\
& \tilde{\omega}(\varsigma, \Lambda) \cdot|\tilde{\omega}(\varsigma, \Lambda)|^{-1}=(-1)^{q} \cdot \operatorname{sign} \tilde{\omega}^{\mathfrak{m}}\left(\mu_{\varsigma}\right) \quad \text { in Case II. } \tag{3.10}
\end{align*}
$$

It is also clear that

$$
\begin{equation*}
\tilde{\omega}(\varsigma,-\Lambda)=(-1)^{\iota} \tilde{\omega}(\varsigma, \Lambda) . \tag{3.11}
\end{equation*}
$$

If $\varsigma$ is in $C^{\prime}$, choose $\sigma$ in $\mathscr{E}_{M}$ such that $\sigma=\sigma(\varsigma)$. In $\S 2$ we defined the representation $\sigma^{\prime}$. Choose $\zeta^{\prime}$ in $C^{\prime}$ such that $\sigma^{\prime}=\sigma^{\prime}\left(\varsigma^{\prime}\right)$. Given $\varsigma, \varsigma^{\prime}$ is not uniquely defined. However the expression

$$
\operatorname{sign} \tilde{\omega}^{\mathfrak{m}}\left(\mu_{\varsigma^{\prime}}\right) \cdot \tilde{\omega}\left(\zeta^{\prime}, \Lambda\right)
$$

is well defined for any real $\Lambda$. Furthermore

$$
\begin{equation*}
\operatorname{sign} \tilde{\omega}^{\mathfrak{m}}\left(\mu_{\varsigma^{\prime}}\right) \cdot \tilde{\omega}\left(\varsigma^{\prime}, \Lambda\right)=\operatorname{sign} \tilde{\omega}^{\mathfrak{m}}\left(\mu_{\varsigma}\right) \cdot \tilde{\omega}(\zeta, \Lambda) \tag{3.12}
\end{equation*}
$$

For any $\omega$ in $\mathscr{E}_{d}$, let $\Theta_{\omega}$ and $\beta(\omega)$ be the character and formal degree of $\omega$. A formula for $\beta(\omega)$ was quoted in $\S 2$. It is clear that there is a polynomial $p$ such that

$$
\begin{equation*}
\beta(\omega) \leq p(|\omega|), \quad \omega \in \mathscr{E}_{d} \tag{3.13}
\end{equation*}
$$

LEMMA 5. There exists a nonnegative function $\beta(\sigma, \Lambda)$ on $\mathscr{E}_{M} \times \mathbf{R}$ such that for any $f$ in $C_{0}^{\infty}(G)$,

$$
f(1)=\sum_{\omega \in \mathscr{E}_{d}} \beta(\omega) \Theta_{\omega}(f)+\sum_{\sigma \in \mathscr{E}_{M}} \int_{0}^{\infty} \beta(\sigma, \Lambda) \Theta_{\sigma, \Lambda}(f) d \Lambda .
$$

In addition $\beta(\sigma, \Lambda)$ has the following properties.
(i) $\beta(\sigma, \Lambda)=\beta(\sigma,-\Lambda)=\beta\left(\sigma^{\prime}, \Lambda\right)$.
(ii) For any $\sigma$ in $\mathscr{E}_{M}, \beta(\sigma, \Lambda)$ is the restriction to $\mathbf{R}$ of a meromorphic function on $\mathbf{C}$ with no real poles.
(iii) If $\sigma$ is in $\mathscr{E}_{M}$ and $\Lambda \neq 0$, then $\beta(\sigma, \Lambda) \neq 0$.
(iv) For every $r>0$, there are polynomials $p_{1}, p_{2}$ such that for $\sigma$ in $\mathscr{E}_{M}, \Lambda$ in R ,

$$
\left|\left(\frac{d}{d \Lambda}\right)^{r} \beta(\sigma, \Lambda)\right| \leq p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|)
$$

Proof. We deal with Case I first. Although the lemma is true in general for Case I, Harish-Chandra in $[\mathbf{3}(\mathbf{k}), \S 24]$ proves it only in case $j$ is one-to-one; that is, $G \subset G_{\mathbf{c}}$. We shall content ourselves with dealing with this situation.

Let $C^{ \pm}=\left\{\zeta \in C:\left\langle\zeta, \exp (-1)^{1 / 2} \pi H_{0}\right\rangle= \pm 1\right\}$. Since $\left(\exp (-1)^{1 / 2} \pi H_{0}\right)^{2}=1$, any element in $C$ is in either $C^{+}$or $C^{-}$. Then in [3(k), Lemma 56], HarishChandra shows that there is a constant, which he writes as $c_{A} / c c_{B}$, such that
for every $f$ in $C_{0}^{\infty}(G)$

$$
\begin{aligned}
& f(1)= \sum_{\omega \in \mathscr{E}_{d}} \beta(\omega) \Theta_{\omega}(f)-\left(c_{A} / c c_{B}\right) \cdot i \\
& \cdot(-1)^{m+q}\left\{\sum_{\varsigma \in C^{+}} \int_{0}^{\infty} \operatorname{coth} \frac{\pi \Lambda}{2} \cdot \tilde{\omega}(\varsigma, \Lambda)\left(\int_{A_{1} \times \mathbf{R}} F_{f}\left(a h_{t}\right)\langle\varsigma, a\rangle \cdot e^{i \Lambda t} d a d t\right) d \Lambda\right. \\
&\left.+\sum_{\varsigma \in C^{-}} \int_{0}^{\infty} \tanh \frac{\pi \Lambda}{2} \cdot \tilde{\omega}(\varsigma, \Lambda)\left(\int_{A_{I} \times \mathbf{R}} F_{f}\left(a h_{t}\right)\langle\varsigma, a\rangle e^{i \Lambda t} d a d t\right) d \Lambda\right\}
\end{aligned}
$$

This equals

$$
\begin{aligned}
& \sum_{\omega \in \mathscr{E}_{d}} \beta(\omega) \Theta_{\omega}(f)-\left(c_{A} / c_{0} c c_{B}\right) \cdot i \\
& \cdot(-1)^{q-1}\left\{\sum_{\varsigma \in C^{+}} \int_{0}^{\infty} \operatorname{coth} \frac{\pi \Lambda}{2} \cdot \tilde{\omega}(\varsigma, \Lambda) \cdot \operatorname{sign} \tilde{\omega}^{\mathrm{m}}\left(\mu_{\varsigma}\right) \cdot \Theta_{\sigma(\varsigma),-\Lambda}(f)\right. \\
&\left.+\sum_{\varsigma \in C^{-}} \int_{0}^{\infty} \tanh \frac{\pi \Lambda}{2} \cdot \tilde{\omega}(\varsigma, \Lambda) \cdot \operatorname{sign} \tilde{\omega}^{\mathrm{m}}\left(\mu_{\varsigma}\right) \cdot \Theta_{\sigma(\varsigma),-\Lambda}(f)\right\}
\end{aligned}
$$

We then define

$$
\beta(\varsigma, \Lambda)=\left(c_{A} / c_{0} c c_{B}\right)(-i)(-1)^{q-1}\left\{\begin{array}{l}
\operatorname{coth} \frac{\pi \Lambda}{2}  \tag{3.14}\\
\tanh \frac{\pi \Lambda}{2}
\end{array} \cdot \operatorname{sign} \tilde{\omega}^{\mathrm{m}}\left(\mu_{\varsigma}\right)\right.
$$

where $\operatorname{coth}(\pi \Lambda / 2)$ or $\tanh (\pi \Lambda / 2)$ is used depending on whether $\varsigma$ is in $C^{+}$or $C^{-}$.

Now let us deal with Case II. Then $\mathscr{E}_{d}$ is empty. We can use the limit formula (3.3). Therefore

$$
f(1)=(1 / c)(-1)^{q} F_{f}(1 ; \tilde{\omega})
$$

By the Fourier inversion formula on the connected abelian group $A_{I} \times \mathbf{R}$,

$$
f(1)=\left(\frac{1}{c}\right)(-1)^{q} \sum_{\zeta \in C} \int_{-\infty}^{\infty}\left[\int_{A_{I} \times \mathbf{R}} F_{f}\left(a h_{t} ; \tilde{\omega}\right)\langle\zeta, a\rangle e^{i \Lambda t} d a d t\right] d \Lambda
$$

Since $m=[P]$ and since $F_{f}$ is in $C_{0}^{\infty}\left(A_{I} \times \mathbf{R}\right)$, we see by integration by parts that $f(1)$ is equal to

$$
\left(\frac{1}{c}\right)(-1)^{q+m} \sum_{\varsigma \in C} \int_{-\infty}^{\infty} \tilde{\omega}(\varsigma, \Lambda)\left[\int_{A_{I} \times \mathbf{R}} F_{f}\left(a h_{t}\right)\langle\zeta, a\rangle e^{i \Lambda t} d a d t\right] d \Lambda .
$$

By Theorem 1 this expression equals

$$
\left(\frac{1}{c_{0} c}\right)(-1)^{q} \sum_{\varsigma \in C} \int_{-\infty}^{\infty} \tilde{\omega}(\varsigma, \Lambda) \cdot \operatorname{sign} \tilde{\omega}^{\mathfrak{m}}\left(\mu_{\varsigma}\right) \cdot \Theta_{\sigma(\varsigma),-\Lambda}(f) d \Lambda
$$

We then define

$$
\begin{equation*}
\beta(\varsigma, \Lambda)=\left(2 / c_{0} c\right)(-1)^{q} \cdot \tilde{\omega}(\varsigma, \Lambda) \cdot \operatorname{sign} \tilde{\omega}^{\mathrm{m}}\left(\mu_{\varsigma}\right) \tag{3.15}
\end{equation*}
$$

In either case, we see from (3.10) that $\beta(\varsigma, \Lambda)$ is nonnegative. Also from (3.11) we see that

$$
\beta(\varsigma,-\Lambda)=\beta(\varsigma, \Lambda)
$$

It is clear that the expression $\beta\left(s^{\prime}, \Lambda\right)$ is well defined. (3.12) implies the formula

$$
\beta\left(\varsigma^{\prime}, \Lambda\right)=\beta(\varsigma, \Lambda)
$$

Since for any $\Lambda \neq 0$ the representations $\pi_{\sigma^{\prime}, \Lambda}$ and $\pi_{\sigma,-\Lambda}$ are equivalent,

$$
\Theta_{\sigma^{\prime}, \Lambda}=\Theta_{\sigma,-\Lambda}
$$

Therefore in either Case I or Case II we obtain the formula

$$
f(1)=\sum_{\omega \in \mathscr{E}_{d}} \beta(\omega) \Theta_{\omega}+\sum_{\zeta \in C} \int_{0}^{\infty} \beta(\varsigma, \Lambda) \Theta_{\sigma(s), \Lambda}(f) d \Lambda
$$

It is clear in either case that

$$
\beta(s \zeta, \Lambda)=\beta(\zeta, \Lambda), \quad s \in W_{1} .
$$

Then if $\sigma$ is in $\mathscr{E}_{M}$, choose any $\varsigma$ in $C^{\prime}$ such that $\sigma=\sigma(\varsigma)$. Define

$$
\beta(\sigma, \gamma)=\left[W_{1}\right] \beta(\varsigma, \Lambda)
$$

Then $\beta(\sigma, \Lambda)$ is well defined and $\beta(\sigma, \Lambda)$ satisfies the formula of the lemma.
Property (i) of the lemma follows from the above discussion. Properties (ii), (iii), and (iv) follow easily from formulae (3.14) and (3.15).

For $\omega$ in $\mathscr{E}_{d}$, let $\pi_{\omega}$ be a representation in the class of $\omega$, acting on the Hilbert space $\mathscr{H}_{\omega}$. Let $\mathscr{H}_{2}(\omega)$ be the space of Hilbert-Schmidt operators on $\mathscr{H}_{\omega}$ with the Hilbert-Schmidt norm $\|\cdot\|_{2}$. Similarly, for $\sigma$ in $\mathscr{E}_{M}$, write $\mathscr{H}_{2}(\sigma)$ as the space of Hilbert-Schmidt operators on $\mathscr{H}_{\sigma}$.

Let $L_{0}^{2}(\hat{G})$ be the set of functions

$$
a_{0}: \mathscr{E}_{d} \rightarrow \bigoplus_{\omega \in \mathscr{E}_{d}} \mathscr{H}_{2}(\omega)
$$

such that
(i) $a_{0}(\omega)$ is in $\mathscr{H}_{2}(\omega)$ for each $\omega$ in $\mathscr{E}_{d}$.
(ii) $\left\|a_{0}\right\|^{2}=\sum_{\omega \in \mathscr{Z}_{d}}\left\|a_{0}(\omega)\right\|_{2}^{2} \beta(\omega)<\infty$.

Notice that if we are in Case II, $\mathscr{E}_{d}$ is empty so that $L_{0}^{2}(\hat{G})$ is empty.
Let $L_{1}^{2}(\hat{G})$ be the set of functions

$$
a_{1}: \mathscr{E}_{M} \times \mathbf{R} \rightarrow \bigoplus_{\sigma \in \mathscr{E}_{M}} \mathscr{H}_{2}(\sigma)
$$

such that
(i) $a_{1}(\sigma, \Lambda)$ is in $\mathscr{H}_{2}(\sigma)$ for each $\sigma$ in $\mathscr{E}_{M}$ and $\Lambda$ in $\mathbf{R}$.
(ii) $a_{1}\left(\sigma^{\prime},-\Lambda\right)=N_{\sigma}(\Lambda) a_{1}(\sigma, \Lambda) N_{\sigma}(\Lambda)^{-1}, \sigma \in \mathscr{E}_{M}, \Lambda \neq 0$.
(iii) For any $\sigma$ in $\mathscr{E}_{M}, a_{1}(\sigma, \Lambda)$ is a Borel function of $\Lambda$.
(iv)

$$
\left\|a_{1}\right\|^{2}=\frac{1}{2} \sum_{\sigma \in \mathscr{F}_{M}} \int_{-\infty}^{\infty}\left\|a_{1}(\sigma, \Lambda)\right\|_{2}^{2} \beta(\sigma, \Lambda) d \Lambda<\infty
$$

(In (ii) we can regard the operators $N_{\sigma}(\Lambda)$ as maps from $\mathscr{H}_{\sigma}$ to $\mathscr{K}_{\sigma^{\prime}}$ if we recall the canonical isomorphisms $\mathscr{H}_{\sigma, \Lambda} \leftrightarrow \mathscr{H}_{\sigma}, \mathscr{H}_{\sigma^{\prime},-\Lambda} \leftrightarrow \mathscr{H}_{\sigma^{\prime}}$.)

Notice that since $N_{\sigma}(\Lambda)$ is unitary, condition (ii) implies that

$$
\left\|a_{1}\left(\sigma^{\prime},-\Lambda\right)\right\|_{2}^{2}=\left\|a_{1}(\sigma, \Lambda)\right\|_{2}^{2}
$$

Therefore

$$
\begin{equation*}
\left\|a_{1}\right\|^{2}=\sum_{\sigma \in \mathscr{E}_{M}} \int_{0}^{\infty}\left\|a_{1}(\sigma, \Lambda)\right\|_{2}^{2} \beta(\sigma, \Lambda) d \Lambda . \tag{3.16}
\end{equation*}
$$

$L_{0}^{2}(\hat{G})$ and $L_{1}^{2}(\hat{G})$ are Hilbert spaces. Let $L^{2}(\hat{G})=L_{0}^{2}(\hat{G}) \oplus L_{1}^{2}(\hat{G})$. If $f$ is in $C_{0}^{\infty}(G)$, define $\hat{f}$ in $L^{2}(\hat{G})$ by

$$
\begin{gathered}
\hat{f}=\left(\hat{f}_{0}(\omega), \hat{f}_{1}(\sigma, \Lambda)\right), \\
\hat{f}_{0}(\omega)=\int_{G} f(x) \pi_{\omega}(x) d x, \quad \omega \in \mathscr{E}_{d}, \\
\hat{f}_{1}(\sigma, \Lambda)=\int_{G} f(x) \pi_{\sigma, \Lambda}(x) d x, \quad \sigma \in \mathscr{E}_{M}, \Lambda \in \mathbf{R}
\end{gathered}
$$

(We can regard $\hat{f}_{1}(\sigma, \Lambda)$ as an operator on $\mathscr{H}_{\sigma}$.)
Theorem 2 (Plancherel Formula). The map

$$
f \rightarrow \hat{f}, \quad f \subset C_{0}^{\infty}(G)
$$

extends uniquely to an isometry from $L^{2}(G)$ onto $L^{2}(\hat{G})$.
Proof. Fix $f$ in $C_{0}^{\infty}(G)$. Define

$$
g(x)=\int_{G} f(y) \overline{f\left(x^{-1} y\right)} d y, \quad x \in G
$$

Clearly $g$ is in $C_{0}^{\infty}(G)$ and $g(1)$ equals $\|f\|_{2}^{2}$. If $\pi$ is an irreducible unitary representation of $G$,

$$
\begin{aligned}
\pi(g) & =\int_{G \times G} f(y) \overline{f\left(x^{-1} y\right)} d y \cdot \pi(x) d x \\
& =\int_{G \times G} f(y) \overline{f\left(x^{-1}\right)} \cdot \pi(y x) d y d x \\
& =\left(\int_{G} f(y) \cdot \pi(y) d y\right)\left(\int_{G} f(x) \cdot \pi(x) d x\right)^{*} \\
& =\pi(f) \cdot \pi(f)^{*}
\end{aligned}
$$

where $\pi(f)^{*}$ is the adjoint of $\pi(f)$. Therefore

$$
\operatorname{tr} \pi(g)=\|\pi(f)\|_{2}^{2}=\|\hat{f}(\pi)\|_{2}^{2}
$$

Therefore, applying Lemma 5 to $g(x)$ we see that

$$
\|f\|_{2}^{2}=\|\hat{f}\|^{2}
$$

Thus, the map $f \rightarrow \hat{f}$ is an isometry. We need only show that it is surjective.

By the Schur orthogonality relations (2.7) and (2.8), the map is onto $L_{0}^{2}(\hat{G})$. We must show that it is onto $L_{1}^{2}(\hat{G})$.

Let $\hat{\rho}_{1}$ be the representation of $G \times G$ on $L_{1}^{2}(\hat{G})$ given by

$$
\hat{\rho}_{1}(x, y) a_{1}(\sigma, \Lambda)=\pi_{\sigma, \Lambda}(x) a_{1}(\sigma, \Lambda) \pi_{\sigma, \Lambda}\left(y^{-1}\right)
$$

for $\sigma \in \mathscr{E}_{M}, \Lambda \in \mathbf{R}$, and $(x, y) \subset G \times G . G \times G$, being semisimple, is of type I [3(a), p. 30], so $\hat{\rho}_{1}$ is of type I. Let $\mathbf{R}^{+}=\{\Lambda \in \mathbf{R}: \Lambda>0\}, S=\mathscr{E}_{M} \times \mathbf{R}^{+}$, and let $C$ be the measure class on $S$ defined by the discrete measure on $\mathscr{E}_{M}$ and Lebesgue measure on $\mathbf{R}^{+} . \beta(\sigma, \Lambda)$ does not vanish for any $(\sigma, \Lambda)$ in $S$, and the representations $\left\{\pi_{\sigma, \Lambda} \times \pi_{\sigma, \Lambda}:(\sigma, \Lambda) \in S\right\}$ of $G \times G$ are all irreducible and inequivalent. $\hat{\rho}_{1}$ is clearly the direct integral of these representations of $G \times G$ with respect to the measure class $C$. Therefore $\hat{\rho}_{1}$ is multiplicity-free by [ $6(\mathrm{~b})$, Theorem 5]. This means that the algebra $R\left(\hat{\rho}_{1}, \hat{\rho}_{1}\right)$ of intertwining operators of $\hat{\rho}_{1}$ is commutative.

Let $\rho$ be the two-sided regular representation of $G \times G$ on $L^{2}(G)$. Then the map

$$
f \rightarrow \hat{f}_{1}, \quad f \in L^{2}(G)
$$

is an intertwining operator between $\rho$ and $\hat{\rho}_{1}$. Thus if $L$ is the closure of the set $\left\{\hat{f}_{1}: f \in L^{2}(G)\right\}$, and $P$ is the orthogonal projection of $L_{1}^{2}(\hat{G})$ onto $L$, then $P$ is in $R\left(\hat{\rho}_{1}, \hat{\rho}_{1}\right)$. But since $R\left(\hat{\rho}_{1}, \hat{\rho}_{1}\right)$ is commutative, it is well known that $P$ is of the form $P_{E}$, where $E$ is a Borel subset of $S$ and

$$
P_{E}=\left\{a_{1} \in L_{1}^{2}(\hat{G}): a_{1} \text { vanishes outside } E\right\}
$$

To complete the proof of the surjectivity of the map $f \rightarrow \hat{f}_{1}$, we need only show that the complement of $E$ in $S$ is a null set with respect to $C$.

Let us assume the contrary. Then there is a $\sigma$ in $\mathscr{E}_{M}$ and a subset $R_{\mathbf{1}}$ of $\mathbf{R}^{+}$ of positive Lebesgue measure such that for any $f$ in $C_{0}^{\infty}(G)$,

$$
\hat{f}_{1}(\sigma, \Lambda)=0 \text { for almost all } \Lambda \text { in } R_{1}
$$

Choose a $\tau$ in $\mathscr{E}_{K}$ for which there is a nonzero intertwining operator $T$ between the restriction of $\tau$ to $M$ and $\sigma$. Choose a vector $\xi$ in the space on which $\tau$ acts such that $T \xi \neq 0$. Define

$$
\Phi(k)=T\left(\tau\left(k^{-1}\right) \xi\right), \quad k \in K
$$

Then $\Phi$ is in $\mathscr{H}_{\sigma}$. For any $f$ in $C_{0}^{\infty}(G)$,

$$
\begin{aligned}
(\hat{f}(\sigma, \Lambda) \Phi)(1) & =\left(\int_{G} f(x) \pi_{\sigma, \Lambda}(x) d x \cdot \Phi\right) \\
& =\int_{G} f\left(x^{-1}\right) \Phi(x) e^{-\rho(H(x))} d x
\end{aligned}
$$

Then by $(2.9),(\hat{f}(\sigma, \Lambda) \Phi)(1)$ is equal to

$$
\int_{K \times \mathbf{R} \times N} f\left(n^{-1} \cdot \exp \left(-t H_{0}\right) \cdot k^{-1}\right) e^{\left(i \Lambda+\rho\left(H_{0}\right)\right) t} \Phi(k) d k d t d n
$$

Let $f\left(n^{-1} \cdot \exp \left(-t H_{0}\right) \cdot k^{-1}\right)$ equal

$$
\chi(k) \cdot \alpha(t) \cdot \nu(n)
$$

where $\chi(k)=(\tau(k) \xi, \xi)$ and $\nu$ is any function in $C_{0}^{\infty}(N)$ such that $\int_{N} \nu(n) d n=$ 1. $\alpha$ is some function in $C_{0}^{\infty}(\mathbf{R})$ such that $\int_{-\infty}^{\infty} \alpha(t) e^{\left(i \Lambda+\rho\left(H_{0}\right)\right) t} d t$ is not equal to zero for any $\Lambda$ belonging to a subset $R_{2}$ of $R_{1}$ of positive measure. Clearly such an $\alpha$ exists.

For a fixed $\Lambda_{0}$ in. $R_{2}$,

$$
\left(\hat{f}\left(\sigma, \Lambda_{0}\right) \Phi\right)(1)=T(\xi) \cdot \int_{-\infty}^{\infty} \alpha(t) e^{\left(i \Lambda_{0}+\rho\left(H_{0}\right)\right) t} d t
$$

This is a nonzero vector in the space on which $\sigma$ acts. However, $\left(\hat{f}\left(\sigma, \Lambda_{0}\right) \Phi\right)(k)$ is a continuous function of $k$, so $\left(\hat{f}\left(\sigma, \Lambda_{0}\right) \Phi\right)(k)$ is nonzero for a subset of $K$ of positive measure. Therefore $\hat{f}\left(\sigma, \Lambda_{0}\right) \Phi$ is a nonzero vector in $\mathscr{H}_{\sigma}$. This means that the operators $\hat{f}_{1}(\sigma, \Lambda)$ do not vanish for any $\Lambda$ in $R_{2}$. We have a contradiction. The proof of Theorem 2 is now complete.
4. Statement of Theorem 3. For $x$ in $G$, define

$$
\Xi(x)=\int_{K} e^{-\rho(H(x k))} d k
$$

Define a norm on $\mathfrak{g}$ by putting

$$
\|X\|^{2}=-B(X, \theta X), \quad X \in \mathfrak{g}
$$

where $B$ is the Killing form on $\mathfrak{g}$. Since $G=K A_{\mathfrak{p}} K$ there exist a unique function $\sigma$ on $G$ such that
(i) $\sigma\left(k_{1} x k_{2}\right)=\sigma(x), k_{1}, k_{2} \in K, x \in G$;
(ii) $\sigma(\exp H)=\|H\|, H \in \mathfrak{a}_{\mathfrak{p}}$.

It is known that there exist numbers $c, d$ such that for any $a$ in $A_{p}^{+}$ $\left(=\left\{\exp t H_{0}: t,>0\right\}\right)$,

$$
\begin{equation*}
1<\Xi(a) e^{\rho(\log a)}<c(1+\sigma(a))^{d} \tag{4.1}
\end{equation*}
$$

(see [3(g), Theorem 3 and Lemma 36]). Also there is an $r_{0}>0$ such that

$$
\begin{equation*}
\int_{G} \Xi(x)^{2}(1+\sigma(x))^{-r_{0}} d x=N\left(r_{0}\right)<\infty \tag{4.2}
\end{equation*}
$$

(see [3(1), Lemma 11]).
Choose $\delta$ in $K$ such that $\delta^{-1} a \delta=a^{-1}$ for any $a$ in $A_{\mathfrak{p}}$. We obtain the formulae

$$
\begin{align*}
& \Xi\left(a^{-1}\right)=\Xi\left(\delta^{-1} a \delta\right)=\Xi(a) \\
& \sigma\left(a^{-1}\right)=\sigma\left(\delta^{-1} a \delta\right)=\sigma(a) \tag{4.3}
\end{align*}
$$

Let $\mathscr{B}$ be the universal enveloping algebra of $\mathfrak{g}_{\mathrm{c}}$. We can identify $\mathscr{B}$ with the algebra of left invariant differential operators on $G$. Let $\rho$ be the canonical antiisomorphism with $\mathscr{B}$ and the algebra of right invariant differential operators on $G$. If $g_{1}$ and $g_{2}$ are in $\mathscr{B}$ and $f$ is a differentiable function on $G$, then the actions
of $\rho\left(g_{1}\right)$ and $g_{2}$ on $f$ commute. We denote the resultant of this action at any $x$ in $G$ by $f\left(g_{1} ; x ; g_{2}\right)$.

Now for every $g_{1}, g_{2}$, in $\mathscr{B}$ and $s$ in R, we define a seminorm on $C^{\infty}(G)$ by

$$
\|f\|_{g_{1}, g_{2}, s}=\sup _{x \in G}\left|f\left(g_{1} ; x ; g_{2}\right)\right| \Xi(x)^{-1}(1+\sigma(x))^{s}, \quad f \in C^{\infty}(G)
$$

Let $\mathscr{C}(G)=\left\{f \subset C^{\infty}(G):\|f\|_{g_{1}, g_{2}, s}<\infty\right.$, for any $g_{1}, g_{2}$ in $\mathscr{B}$ and $s$ in $\left.\mathbf{R}\right\}$. These seminorms make $\mathscr{C}(G)$ into a Fréchet space.

Clearly

$$
C_{0}^{\infty}(G) \subset \mathscr{C}(G)
$$

is a continuous inclusion, and it is known that $C_{0}^{\infty}(G)$ is dense in $\mathscr{C}(G)$ [3(1), Theorem 2]. Also from (4.2) we see that there is a continuous inclusion of $\mathscr{E}(G)$ into $L^{2}(G) . \mathscr{C}(G)$ is called the Schwartz space of $G$.

We wish to define a subspace of $L^{2}(\hat{G})$ which will ultimately turn out to be the image of $\mathscr{C}(G)$ under the Fourier transform map, $f \rightarrow \hat{f}$. We shall need to fix appropriate bases for the Hilbert spaces $\mathscr{H}_{\omega}$ and $\mathscr{H}_{\sigma, \Lambda}$.

For each $\omega$ in $\mathscr{E}_{d}$ let $\pi_{\omega}$ be a representation in the class of $\omega$ acting on the Hilbert space $\mathscr{H}_{\omega}$. We can choose an orthonormal basis

$$
\begin{equation*}
\left\{\Phi_{\tau, i}=\Phi_{\tau, i}(\omega)\right\}_{\tau \in \mathscr{E}_{K}} \tag{4.4}
\end{equation*}
$$

of $\mathscr{H}_{\omega}$ such that $\Phi_{\tau, i}$ transforms under $\left.\pi_{\omega}\right|_{K}$, the restriction of $\pi_{\omega}$ to $K$, according to the irreducible representation $\tau$ of $K$. The second subscript, $i$, ranges from 1 to $[\omega: \tau] \cdot \operatorname{dim} \tau$, where $[\omega: \tau]$ is the multiplicity of $\tau$ in $\left.\pi_{\omega}\right|_{K}$. It is known that $[\omega: \tau]<\operatorname{dim} \tau$ (see [3(b), Theorem 4]).

We shall construct explicit bases for the Hilbert spaces $\mathscr{H}_{\sigma, \Lambda}$. As we remarked earlier, there is a canonical intertwining operator between the representations $\left.\pi_{\sigma, \Lambda}\right|_{K}$ and $\pi_{\sigma}$ of $K$. Therefore we shall choose a fixed orthonormal basis for the Hilbert space $\mathscr{H}_{\sigma}$.

The multiplicity of $\tau$ in $\left.\pi_{\sigma, \Lambda}\right|_{K}$ equals the multiplicity of $\tau$ in $\pi_{\sigma}$. But $\pi_{\sigma}$ is just the representation $\sigma$ induced to $K$. Therefore by the Frobenius reciprocity theorem for compact groups [ $\mathbf{6}(\mathbf{a}$ ), Theorem 8.2], these multiplicities are just equal to $[\tau: \sigma]$, the multiplicity of $\sigma$ in $\left.\tau\right|_{M}\left(\left.\tau\right|_{M}\right.$ is the restriction of $\tau$ to $\left.M\right)$.

Fix $\tau$ in $\mathscr{E}_{K}$ and $\sigma$ in $\mathscr{E}_{M}$ acting on the Hilbert spaces $V_{\tau}$ and $V_{\sigma}$ of dimension $t$ and $s$ respectively. Let $R(\tau, \sigma)$ be the set of intertwining operators from $V_{\tau}$ to $V_{\sigma}$ for $\left.\tau\right|_{M}$ and $\sigma$. The Hilbert-Schmidt norm makes $R(\tau, \sigma)$ into a Hilbert space of dimension $[\tau: \sigma]$.

Now suppose $T$ is in $R(\tau, \sigma)$. Since $\sigma$ is irreducible, we can assume that there are orthonormal bases $\left\{\xi_{1}, \ldots, \xi_{t}\right\}$ and $\left\{\eta_{1}, \ldots, \eta_{s}\right\}$ of $V_{\tau}$ and $V_{\sigma}$ respectively such that there is a constant $c$ for which

$$
\begin{array}{cc}
T \xi_{i}=c \eta_{i}, & i \gtrless s \\
T \xi_{i}=0, & i>s
\end{array}
$$

Suppose $T$ has been normalized such that $c=(t / s)^{1 / 2}$. Then

$$
\begin{align*}
& T \xi_{i}=(t / s)^{1 / 2} \eta_{i}, \quad i \gtrless s \\
& \|T\|=t^{1 / 2} \tag{4.5}
\end{align*}
$$

Fix an element $\xi$ of norm 1 in $V_{\tau}$. Write $\tau^{*}(k)$ for $\tau\left(k^{-1}\right)$ if $k$ is in $K$. Define

$$
\Phi(k)=T\left(\tau^{*}(k) \xi\right), \quad k \in K
$$

Then

$$
\begin{align*}
\Phi\left(k m^{-1}\right) & =T\left(\tau(m) \tau^{*}(k) \xi\right)=\sigma(m) T\left(\tau^{*}(k) \xi\right)  \tag{i}\\
& =\sigma(m) \Phi(k), \quad m \in M, k \in K .
\end{align*}
$$

Therefore $\Phi$ is an element in $\mathscr{H}_{\sigma}$.
(ii) $\|\Phi\|=1$, because

$$
\begin{aligned}
(\Phi, \Phi) & =\int_{K}\left(T\left(\tau^{*}(k) \xi\right), T\left(\tau^{*}(k) \xi\right)\right) d k \\
& =\int_{K} \sum_{i j}\left(T\left[\left(\tau^{*}(k) \xi, \xi_{i}\right) \xi_{i}\right], T\left[\left(\tau^{*}(k) \xi, \xi_{j}\right) \xi_{j}\right]\right) d k \\
& =\sum_{i=1}^{s} \int_{K}\left(\left(\tau^{*}(k) \xi, \xi_{i}\right) \eta_{i},\left(\tau^{*}(k) \xi, \xi_{i}\right) \eta_{i}\right) d k \cdot\left(\frac{t}{s}\right) \\
& =\sum_{i=1}^{s} \int_{K}\left(\tau^{*}(k) \xi, \xi_{i}\right)\left(\overline{\tau^{*}(k) \xi, \xi_{i}}\right) d k \cdot\left(\frac{t}{s}\right) \\
& =\left(\frac{t}{s}\right) \cdot\left(\frac{\operatorname{dim} \sigma}{\operatorname{dim} \tau}\right)(\text { by the Schur orthogonality relations on } K) \\
& =1
\end{aligned}
$$

Conversely, let $\Phi$ be any unit vector in $\mathscr{H}_{\sigma}$ such that $\Phi$ transforms under $\pi_{\sigma}$ according to $\tau$. Then there exists a unit vector $\xi$ in $V_{\tau}$ and a $T$ in $R(\tau, \sigma)$ with $\|T\|=(\operatorname{dim} \tau)^{1 / 2}$ such that

$$
\Phi(k)=T\left(\tau^{*}(k) \xi\right), \quad k \in K
$$

For $\Phi$ defined as above, the vector $N_{\sigma}(\Lambda) \Phi$ is in $\mathscr{H}_{\sigma^{\prime}}$. Clearly $N_{\sigma}(\Lambda) \Phi$ transforms under $\pi_{\sigma^{\prime}}$ according to $\tau$. Then there exists a unique $T^{\prime}$ in $R\left(\tau, \sigma^{\prime}\right)$ with $\left\|T^{\prime}\right\|=(\operatorname{dim} \tau)^{1 / 2}$ such that

$$
\left(N_{\sigma}(\Lambda) \Phi\right)(k)=T^{\prime}\left(\tau^{*}(k) \xi\right), \quad k \in K
$$

The map $T \rightarrow T^{\prime}$ from $R(\tau, \sigma)$ into $R\left(\tau, \sigma^{\prime}\right)$ will be denoted $n_{\sigma}(\Lambda)$, so $T^{\prime}=$ $n_{\sigma}(\Lambda) T . n_{\sigma}(\Lambda)$ is norm-preserving and hence unitary.

Fix an orthonormal base $\left\{T_{1}, \ldots, T_{r}\right\}$ of $R(\tau, \sigma)$ of elements of norm equal to $(\operatorname{dim} \tau)^{1 / 2}$. For $1<l \gtrless r, 1<j \gtrless t$, and $k$ in $K$, define

$$
\begin{equation*}
\Phi_{\tau,(l-1) t+j}(k)=T_{l}\left(\tau^{*}(k) \xi_{j}\right) \tag{4.6}
\end{equation*}
$$

Then $\left\{\Phi_{\tau, i}: \tau \in \mathscr{E}_{K}, 1 ₹ i ₹[\tau: \sigma] \operatorname{dim} \tau\right\}$ is an orthonormal base for $\mathscr{H}_{\sigma}$.
The bases (4.4) and (4.6) can be used to define a collection of seminorms on $L_{0}^{2}(\hat{G})$ and $L_{1}^{2}(\hat{G})$ respectively. For each triplet ( $p, q_{1}, q_{2}$ ) of polynomials we define a seminorm on $L_{0}^{2}(\hat{G})$ by letting $\left\|a_{0}\right\|_{p, q_{1}, q_{2}}$ be the supremum over $\omega$, ( $\left.\tau_{1}, i_{1}\right),\left(\tau_{2}, i_{2}\right)$ of the expressions

$$
\begin{equation*}
\left|\left(\Phi_{\tau_{1}, i_{1}}, a_{0}(\omega) \Phi_{\tau_{2}, i_{2}}\right)\right| p(|\omega|) q_{1}\left(\left|\tau_{1}\right|\right) q_{2}\left(\left|\tau_{2}\right|\right), \quad a_{0} \in L_{0}^{2}(\hat{G}) \tag{4.7}
\end{equation*}
$$

Let $\mathscr{C}_{0}(\hat{G})$ be the set of all $a_{0}$ in $L_{0}^{2}(\hat{G})$ for which $\left\|a_{0}\right\|_{p, q_{1}, q_{2}}<\infty$ for every triplet $\left(p, q_{1}, q_{2}\right)$.

For each set of polynomials ( $p_{1}, p_{2}, q_{1}, q_{2}$ ) and each integer $n$ define a seminorm on $L_{1}^{2}(\hat{G})$ as follows: put $\left\|a_{1}\right\|_{\left(p_{1}, p_{2}, q_{1}, q_{2}: n\right)}=\infty$ if for some $\sigma$ in $\mathscr{E}_{M}$ and some $\Phi_{\tau_{1}, i_{1}}$ and $\Phi_{\tau_{2}, i_{2}}$ the function $\left(\Phi_{\tau_{1}, i_{1}}, a_{1}(\sigma, \Lambda) \Phi_{\tau_{2}, i_{2}}\right)$ is not $n$ times continuously differentiable in $\Lambda$. Otherwise, let $\left\|a_{1}\right\|_{\left(p_{1}, p_{2}, q_{1}, q_{2}: n\right)}$ equal the supremum over $(\sigma, \Lambda),\left(\tau_{1}, i_{1}\right),\left(\tau_{2}, i_{2}\right)$ of the expressions

$$
\begin{equation*}
\left|\left(\frac{d}{d \Lambda}\right)^{n}\left(\Phi_{\tau_{1}, i_{1}}, a(\sigma, \Lambda) \Phi_{\tau_{2}, i_{2}}\right)\right| p_{1}(|\sigma|) p_{2}(|\Lambda|) q_{1}\left(\left|\tau_{1}\right|\right) q_{2}\left(\left|\tau_{2}\right|\right) \tag{4.8}
\end{equation*}
$$

Let $\mathscr{C}_{1}(\hat{G})$ be the set of all $a_{1}$ in $L_{1}^{2}(\hat{G})$ for which $\left\|a_{1}\right\|_{\left(p_{1}, p_{2}, q_{1}, q_{2}: n\right)}<\infty$ for every set ( $p_{1}, p_{2}, q_{1}, q_{2}: n$ ).

The above seminorms define topologies on $\mathscr{C}_{0}(\hat{G})$ and $\mathscr{C}_{1}(\hat{G})$. Define

$$
\mathscr{C}(\hat{G})=\mathscr{C}_{0}(\hat{G}) \oplus \mathscr{C}_{1}(\hat{G})
$$

$\mathscr{C}(\hat{G})$ is a Fréchet space.
THEOREM 3. The map $f \rightarrow \hat{f}$ gives a topological isomorphism of $\mathscr{C}(G)$ onto $\mathscr{E}(\hat{G})$.

We shall spend most of the rest of this paper proving this theorem.
5. Spherical functions. In this section we shall define $\tau$-spherical functions on $G$ and develop some of their elementary properties.

A unitary double representation $\tau$ of the compact group $K$ is a Hilbert space on which there is both a left and a right unitary $K$ action. In addition, these actions are required to commute with each other. We denote both the left and the right action of $K$ by $\tau$. If $\tau$ is a unitary double representation of $K$ on the vector space $V_{\tau}$, define a representation $\tau^{\prime}$ of $K \times K$ on $V_{\tau}$ by

$$
\tau^{\prime}\left(k_{1}, k_{2}\right) v=\tau\left(k_{1}\right) v \tau\left(k_{2}^{-1}\right), \quad v \in V_{\tau}, k_{1}, k_{2} \in K
$$

There is a one-to-one correspondence between double representations of $K$ and representations of $K \times K$.

Suppose $\tau$ is a unitary double representation of $K$ on the vector space $V_{\tau}$. A function $\phi$ from $G$ to $V_{\tau}$ is said to be $\tau$-spherical if for every $k_{1}, k_{2}$ in $K$ and $x$ in $G$,

$$
\phi\left(k_{1} x k_{2}\right)=\tau_{1}\left(k_{1}\right) \phi(x) \tau_{2}\left(k_{2}\right)
$$

We shall write $|\phi(x)|$ to indicate the norm of $\phi(x)$ in $V_{\tau}$.
Suppose $f(x)$ is a continuous complex-valued function on $G$ such that the left and right translates of $f$ by elements in $K$ span a finite-dimensional space of functions on $G$. We shall use $f$ to define a spherical function.

Let $\phi$ be the function from $G$ into $L^{2}(K \times K)$ defined by

$$
\phi(x)\left(k_{1}, k_{2}\right)=f\left(k_{1}^{-1} x k_{2}^{-1}\right), \quad x \in G, k_{1}, k_{2} \in K
$$

Define a double $K$ representation $\mu$ on $L^{2}(K \times K)$ by

$$
\begin{aligned}
& {\left[\mu\left(\bar{k}_{1}\right) u\right]\left(k_{1}, k_{2}\right)=u\left(\bar{k}_{1}^{-1} k_{1}, k_{2}\right),} \\
& {\left[u \mu\left(\bar{k}_{2}\right)\right]\left(k_{1}, k_{2}\right)=u\left(k_{1}, k_{2} \bar{k}_{2}^{-1}\right),}
\end{aligned}
$$

for $u$ in $L^{2}(K \times K)$ and $k_{1}, k_{2}, \bar{k}_{1}, \bar{k}_{2}$ in $K$. Let $V_{\mu}$ equal $s p_{x \in G}\{\phi(x)\}$, the finite-dimensional subspace of $L^{2}(K \times K)$ spanned by $\{\phi(x): x \in G\}$. Then for any $x$ in $G$, and $\bar{k}_{1}, \bar{k}_{2}, k_{1}, k_{2}$ in $K$,

$$
\phi\left(\bar{k}_{1} x \bar{k}_{2}\right)\left(k_{1}, k_{2}\right)=f\left(k_{1}^{-1} \bar{k}_{1} x \bar{k}_{2} k_{2}^{-1}\right)=f\left(\left(\bar{k}_{1}^{-1} k_{1}\right)^{-1} x\left(k_{2} \bar{k}_{2}^{-1}\right)^{-1} .\right.
$$

This expression equals

$$
\left(\mu\left(\bar{k}_{1}\right) \phi(x) \mu\left(\bar{k}_{2}\right)\right)\left(k_{1}, k_{2}\right)
$$

Therefore $\phi$ is a $\mu$-spherical function, which we shall call the $\mu$-spherical function associated with $f$.

Notice that if $\tau$ is an irreducible unitary double representation of $K$ on the finite-dimensional Hilbert space $V_{\tau}$, then $\tau$ can be regarded as an irreducible representation $\tau_{1} \otimes \tau_{2}^{*}$ of $K \times K$ on $V_{1} \otimes V_{2}^{*}$. Here $\tau_{1}$ and $\tau_{2}$ are irreducible representations of $K$ on the spaces $V_{1}$ and $V_{2}$, and $\tau_{2}^{*}$ is the dual representation of $\tau_{2}$ acting on $V_{2}^{*}$, the dual space of $V_{2}$. We write $\tau$ as $\left(\tau_{1}, \tau_{2}\right)$ and $|\tau|$ as $\left|\tau_{1}\right|+\left|\tau_{2}\right|$. Let $\mathscr{E}_{K}^{2}$ be the set of equivalence classes of irreducible unitary double representations of $K$.

Suppose that $f(x)$ is the function ( $\Phi_{1}, \pi(x) \Phi_{2}$ ) where $\pi$ is a unitary representation of $G$ on a Hilbert space $\mathscr{H}$. We assume that for $\alpha=1$ or $2, \Phi_{\alpha}$ is a unit vector in $\mathscr{H}$ that transforms under $\left.\pi\right|_{K}$ according to the irreducible unitary representation $\tau_{\alpha}$ of $K$, acting on the Hilbert space $V_{\alpha}$. Let $\tau=\left(\tau_{1}, \tau_{2}\right) \cdot \tau$ is in $\mathscr{E}_{K}^{2}$ and acts on the Hilbert space $V_{\tau}=V_{1} \otimes V_{2}^{*}$. We shall find a formula for the spherical function $\phi$ associated with $f$. Then we shall specialize to the case where $\pi$ is one of the induced representations $\pi_{\sigma, \Lambda}$ defined in $\S 2$.

Let $\tau_{\alpha}$ have dimension $t_{\alpha}$ and let $\left\{\xi_{\alpha 1}, \ldots, \xi_{\alpha t_{\alpha}}\right\}$ be an orthonormal base for $V_{\alpha}$, for $\alpha=1$ or 2 . Let $V_{\alpha}^{\prime}$ be the subspace of $\mathscr{\mathscr { L }}$ spanned by $\left\{\pi(k) \Phi_{\alpha}: k \in K\right\}$. Choose an orthonormal base $\left\{\Phi_{\alpha 1}, \ldots, \Phi_{\alpha t_{\alpha}}\right\}$ of $V_{\alpha}^{\prime}$ such that the correspondence

$$
\xi_{\alpha i} \leftrightarrow \Phi_{\alpha i}, \quad i=1,2, \ldots, t_{\alpha},
$$

gives an intertwining operator between $\tau_{\alpha}$ and $\left.\pi\right|_{K}$ acting on the space $V_{\alpha}^{\prime}$. Define functions $e_{1 i}\left(k_{1}\right)$ and $e_{2 j}\left(k_{2}\right)$ as follows:

$$
\begin{gathered}
e_{1 i}\left(k_{1}\right)=\left(\pi\left(k_{1}\right) \Phi_{1}, \Phi_{1 i}\right), \quad k_{1} \in K, \\
e_{2 j}\left(k_{2}\right)=\left(\pi\left(k_{2}^{-1}\right) \Phi_{2}, \Phi_{2 j}\right)=\left(\pi\left(k_{2}\right) \Phi_{2 j}, \Phi_{2}\right), \quad k_{2} \in K .
\end{gathered}
$$

Then

$$
\phi(x)\left(k_{1}, k_{2}\right)=f\left(k_{1}^{-1} x k_{2}^{-1}\right)=\left(\pi\left(k_{1}\right) \Phi_{1}, \pi(x) \pi\left(k_{2}^{-1}\right) \Phi_{2}\right) .
$$

This is equal to the expression

$$
\sum_{i j} e_{1 i}\left(k_{1}\right) e_{2 j}\left(k_{2}\right)\left(\Phi_{1 i}, \pi(x) \Phi_{2 j}\right) .
$$

$V_{\mu}$ is the subspace of $L^{2}(K \times K)$ spanned by the functions $e_{1 i}\left(k_{1}\right) e_{2 j}\left(k_{2}\right)$. Let $\left\{\xi_{21}^{*}, \ldots, \xi_{2 t_{2}}^{*}\right\}$ be the dual basis in $V_{2}^{*}$ to $\left\{\xi_{21}, \ldots, \xi_{2 t_{2}}\right\}$. Then for $1<i<t_{1}$, $1 \gtrless j<t_{2}$, identify $e_{1 i}\left(k_{1}\right) e_{2 j}\left(k_{2}\right)$ with $\left(t_{1} t_{2}\right)^{-1 / 2} \xi_{1 i} \otimes \xi_{2 j}^{*}$. This gives an intertwining operator between the double representations $\mu$ and $\tau$. Therefore, we can regard $\phi$ as a $\tau$-spherical function from $G$ to $V_{\tau}$. We have the formula

$$
\begin{equation*}
\phi(x)=\left(t_{1} t_{2}\right)^{-1 / 2} \sum_{i j} \xi_{1 i} \otimes \xi_{2 j}^{*}\left(\Phi_{1 i}, \pi(x) \Phi_{2 j}\right), \quad x \in G \tag{5.1}
\end{equation*}
$$

Now suppose that $\pi$ is one of the representations $\pi_{\sigma, \Lambda}$, for $\sigma$ in $\mathscr{E}_{M}$ and $\Lambda$ in R. Recall that $R\left(\tau_{\alpha}, \sigma\right)$ was the space of intertwining operators between $\tau_{\alpha}$ and $\sigma$. Fix $T_{\alpha}$ in $R\left(\tau_{\alpha}, \sigma\right)$ such that

$$
\left\|T_{\alpha}\right\|^{2}=\operatorname{dim} \tau_{\alpha}=t_{\alpha}, \quad \alpha=1,2
$$

Suppose that $\xi_{1}$ and $\xi_{2}$ are unit vectors in $V_{1}$ and $V_{2}$ respectively. Let

$$
\Phi_{\alpha}(k)=T_{\alpha}\left(\tau_{\alpha}^{*}(k) \xi_{\alpha}\right), \quad k \in K, \alpha=1,2
$$

Then $\Phi_{1}$ and $\Phi_{2}$ are unit vectors in $\mathscr{H}_{\sigma}$. Define

$$
\Phi_{\alpha i}(k)=T_{\alpha}\left(\tau_{\alpha}^{*}(k) \xi_{\alpha i}\right), \quad k \in K, \alpha=1,2, \text { and } 1 \gtrless i \gtrless t_{\alpha}
$$

Then $\left\{\Phi_{\alpha i}\right\}$ is an orthonormal basis of $V_{\alpha}^{\prime}$. Also

$$
f(x)=\left(\Phi_{1}, \pi_{\sigma, \Lambda}(x) \Phi_{2}\right)=\left(\pi_{\sigma, \Lambda}\left(x^{-1}\right) \Phi_{1}, \Phi_{2}\right)
$$

This is equal to the expression

$$
\begin{equation*}
\int_{K}\left(T_{1}\left[\tau_{1}^{*}(K(x k)) \xi_{1}\right], T_{2}\left[\tau_{2}^{*}(k) \xi_{2}\right]\right) e^{\left(i \Lambda \mu_{0}-\rho\right)(H(x k))} d k \tag{5.2}
\end{equation*}
$$

The inner product in this integrand is on $V_{\sigma}$, the space on which $\sigma$ acts. $H(x k)$ and $K(x k)$ were defined in $\S 2$. Combining the formulae (5.1) and (5.2) we obtain the following formula

$$
\begin{gather*}
\phi(x)=\left(t_{1} t_{2}\right)^{-1 / 2} \sum_{i j} \xi_{1 i} \otimes \xi_{2 j}^{*} \cdot \int_{K}\left(T_{1}\left[\tau_{1}^{*}(K(x k)) \xi_{1 i}\right]\right.  \tag{5.3}\\
\left.T_{2}\left[\tau_{2}^{*}(k) \xi_{2 j}\right]\right) e^{\left(i \Lambda \mu_{0}-\rho\right)(H(x k))} d k
\end{gather*}
$$

Let $L=L^{\tau}$ be the following set of functions on $M$ :

$$
\left\{\psi: M \rightarrow V_{\tau}: \psi\left(m_{1} m m_{2}\right)=\tau\left(m_{1}\right) \psi(m) \tau\left(m_{2}\right), m, m_{1}, m_{2} \in M\right\} .
$$

$L^{\tau}$ is a Hilbert space with inner product

$$
\begin{aligned}
\left(\psi_{1}, \psi_{2}\right)_{M} & =\int_{M}\left(\psi_{1}(m), \psi_{2}(m)\right) d m \\
& =\int_{M}\left(\tau_{1}(m) \psi_{1}(1), \tau_{1}(m) \psi_{2}(1)\right) d m=\left(\psi_{1}(1), \psi_{2}(1)\right)
\end{aligned}
$$

If $\psi$ is in $L^{\tau}$, then $\psi(1)$ is in $V_{1} \otimes V_{2}^{*}$ and it can be regarded as an intertwining operator from $V_{2}$ to $V_{1}$ for $\left.\tau_{2}\right|_{M}$ and $\left.\tau_{1}\right|_{M}$. Conversely, if $S$ is such an intertwining operator, then

$$
\psi(m)=\tau_{1}(m) S=S \tau_{2}(m)
$$

is a function in $L^{\tau}$.

If $\sigma$ is in $\mathscr{E}_{M}$, let $L_{\sigma}^{\tau}$ be the set of functions $\psi$ in $L^{\tau}$ such that $\psi(m)$ transforms under left and right translates of $M$ according to the representation $\sigma$ of $M$. Then there exists a finite number of representations $\left\{\sigma_{1}, \ldots, \sigma_{r}\right\}$ in $\mathscr{E}_{M}$ such that

$$
L^{\tau}=L_{\sigma_{1}}^{\tau} \oplus \cdots \oplus L_{\sigma_{r}}^{\tau}
$$

For any $\psi$ in $L^{\tau}$ let us extend the domain of $\psi$ to all of $G$ by defining

$$
\psi(k a n)=\tau_{1}(k) \psi(1), \quad k \in K, a \in A_{\mathfrak{p}}, n \in N
$$

Let us return to our function $\phi(x)$ above. Define

$$
T_{1}^{*}: V_{\sigma} \rightarrow V_{1}
$$

as the adjoint of $T_{1}$. Let

$$
S=T_{1}^{*} T_{2}: V_{2} \rightarrow V_{1}
$$

$S$ is an intertwining operator for $\left.\tau_{2}\right|_{M}$ and $\left.\tau_{1}\right|_{M} . S$ is also canonically an element in $V_{1} \otimes V_{2}^{*}$. Note that

$$
S=\left(t_{1} t_{2}\right)^{-1 / 2} \sum_{i j} \xi_{1 i} \otimes \xi_{2 j}^{*}\left(\xi_{1 i}, S \xi_{2 j}\right)_{V_{1}}
$$

Therefore

$$
\begin{equation*}
S=\left(t_{1} t_{2}\right)^{-1 / 2} \sum_{i j} \xi_{1 i} \otimes \xi_{2 j}^{*}\left(T_{1} \xi_{1 i}, T_{2} \xi_{2 j}\right)_{V_{\sigma}} \tag{5.4}
\end{equation*}
$$

The subscripts $V_{1}$ and $V_{\sigma}$ indicate in what space the inner product is taken. Then the function

$$
\psi(m)=\tau_{1}(m) S=S \tau_{2}(m), \quad m \in M
$$

is in $L_{\sigma}^{\tau}$. Also

$$
\begin{aligned}
\|\psi\|_{M}^{2} & =(\psi(1), \psi(1))=(S, S) \\
& =\left(t_{1} t_{2}\right)^{-1 / 2} \sum_{i j}\left|\left(T_{1} \xi_{1 i}, T_{2} \xi_{2 j}\right)\right|^{2}
\end{aligned}
$$

Since $\left\|T_{\alpha}\right\|=\left(t_{\alpha}\right)^{1 / 2}$, it can be shown from (4.5) that this last expression is equal to $\operatorname{dim} \sigma$.

From (5.3) and (5.4) we obtain the formula

$$
\begin{equation*}
\phi(x)=\int_{K} \psi(x k) \tau\left(k^{-1}\right) e^{\left(i \Lambda \mu_{0}-\rho\right)(H(x k))} d k \tag{5.5}
\end{equation*}
$$

For any $\psi$ in $L^{\tau}$ we write

$$
E_{\Lambda}(\psi: x)=\int_{K} \psi(x k) \tau\left(k^{-1}\right) e^{\left(i \Lambda \mu_{0}-\rho\right)(H(x k))} d k
$$

$E_{\Lambda}(\psi: x)$ is called the Einsenstein integral of $\psi$ and $\Lambda$.
Suppose, conversely, that we were given $\psi$ in $L_{\sigma}^{\tau}$ such that $\|\psi\|_{M}^{2}=\operatorname{dim} \sigma$. Then we could choose $T_{\alpha}$ in $R\left(\tau_{\alpha}, \sigma\right)$ with $\left\|T_{\alpha}\right\|^{2}=\operatorname{dim} \tau_{\alpha}$ for $\alpha=1$ or 2 such that

$$
\psi(1)=T_{1}^{*} T_{2}
$$

Again we can define

$$
\Phi_{\alpha i}(k)=T_{\alpha}\left(\tau_{\alpha}^{*}(k) \xi_{\alpha i}\right), \quad k \in K, \alpha=1,2
$$

Then $\Phi_{\alpha i}$ is a unit vector in $\mathscr{H}_{\sigma}$. Working backward we can obtain the formula

$$
\begin{equation*}
E_{\Lambda}(\psi: x)=\left(t_{1} t_{2}\right)^{-1 / 2} \sum_{i j} \xi_{1 i} \otimes \xi_{2 j}^{*}\left(\Phi_{1 i}, \pi_{\sigma, \Lambda}(x) \Phi_{2 j}\right) \tag{5.6}
\end{equation*}
$$

Now, if $\psi(1)=T_{1}^{*} T_{2}$ as above, and $\Lambda \neq 0$, then $n_{\sigma}(\Lambda) T_{\alpha}$ is in $R\left(\tau_{\alpha}, \sigma^{\prime}\right)$ and $\left\|T_{\alpha}\right\|^{2}=t_{\alpha}$ for $\alpha$ equal to 1 or 2 . Define

$$
\begin{equation*}
\left(M_{\sigma}(\Lambda) \psi\right)(1)=\left(n_{\sigma}(\Lambda) T_{1}\right)^{*}\left(n_{\sigma}(\Lambda) T_{2}\right) \tag{5.7}
\end{equation*}
$$

Then $M_{\sigma}(\Lambda) \psi$ can be regarded as a function in $L_{\sigma^{\prime}}^{\tau}$. It has the same norm as $\psi$. Therefore, $M_{\sigma}(\Lambda)$ is a unitary map of $L_{\sigma}^{\tau}$ onto $L_{\sigma^{\prime}}^{\tau}$. We can then define a unitary linear transformation $M(\Lambda)$ of $L^{\tau}$ by defining it to be $M_{\sigma}(\Lambda)$ on each of the orthogonal subspaces $L_{\sigma}^{\tau}$ of $L^{\tau}$.

If $\Lambda \neq 0$ we have the equation

$$
\left(\Phi_{1}, \pi_{\sigma, \Lambda}(x) \Phi_{2}\right)=\left(N_{\sigma}(\Lambda) \Phi_{1}, \pi_{\sigma^{\prime},-\Lambda}(x) N_{\sigma}(\Lambda) \Phi_{2}\right)
$$

Then from (5.6) we obtain the formula

$$
\begin{equation*}
E_{\Lambda}(\psi: x)=E_{-\Lambda}(M(\Lambda) \psi: x) . \tag{5.8}
\end{equation*}
$$

This is the functional equation for the Eisenstein integral.
From (5.7) and (2.14) we obtain the formula

$$
\begin{equation*}
M(\Lambda)^{-1}=M(-\Lambda) \tag{5.9}
\end{equation*}
$$

Since $M(\Lambda)$ is unitary, it is clear that

$$
\begin{equation*}
M(\Lambda)^{*}=M(\Lambda)^{-1}=M(-\Lambda) \tag{5.10}
\end{equation*}
$$

We make a final remark about the irreducibility of the representations $\pi_{\sigma, 0}$, for $\sigma$ in $\mathscr{E}_{M}$. These representations may or may not be irreducible. If $\pi_{\sigma, 0}$ is irreducible, then for any nonzero vectors $\Phi_{1}$ and $\Phi_{2}$ in $\mathscr{H}_{\sigma, 0}$ the function ( $\Phi_{1}, \pi_{\sigma, 0}(x) \Phi_{2}$ ) cannot vanish identically in $x$. Therefore if $\tau$ is in $\mathscr{E}_{K}^{2}$ and $\psi$ is any nonzero vector in $L_{\sigma}^{\tau}, E_{0}(\psi: x)$ does not vanish identically in $x$.

On the other hand, suppose that

$$
\mathscr{H}_{\sigma, 0}=\mathscr{H}_{1} \oplus \mathscr{H}_{2}
$$

where $\mathscr{H}_{1}$ and $\mathscr{H}_{2}$ are nonzero closed subspaces of $\mathscr{H}_{\sigma, 0}$ which are invariant under $\pi_{\sigma, 0}$. It is possible to choose nonzero vectors $\Phi_{\alpha}$ in $\mathscr{H}_{\alpha}$ such that $\Phi_{\alpha}$ transforms under $\left.\pi_{\sigma, 0}\right|_{K}$ according to some irreducible representation $\tau_{\alpha}$ of $K$, for $\alpha$ equal to 1 or 2 . Let $\tau$ be the double representation ( $\tau_{1}, \tau_{2}$ ). Then there exists a nonzero element $\psi$ in $L_{\sigma}^{\tau}$ such that $E_{0}(\psi: x)$ vanishes identically in $x$.
6. Proof that the map is injective. Let $\pi$ be a unitary representation of $G$ on a Hilbert space $\mathscr{H}$. If $v$ is a vector in $\mathscr{H}$ such that the map from $G$ to $\mathscr{H}$ given by

$$
x \rightarrow \pi(x) v, \quad x \in G,
$$

is infinitely differentiable, $v$ is called a differentiable vector. Let $\mathscr{H}^{\infty}$ be the set of differentiable vectors in $\mathscr{H}$. If $v$ is in $\mathscr{H}^{\infty}$ and $X$ is in $\mathfrak{g}$, define

$$
\pi(X) v=\lim _{t \rightarrow 0} \frac{1}{t}(\pi(\exp t X) v-v)
$$

It can be checked that this gives a representation of the Lie algebra $\mathfrak{g}$ on the vector space $\mathscr{H}^{\infty}$. It extends to a representation, again denoted $\pi$, of the universal enveloping algebra, $\mathscr{B}$, of $\mathfrak{g}_{\mathrm{c}}$.

Let $\mathcal{Z}$ be the center of $\mathscr{B}$. If the restriction of $\pi$ to $\mathfrak{Z}$ is one-dimensional, we obtain a homomorphism

$$
\chi: \mathcal{Z} \rightarrow \mathbf{C}
$$

In this case we say that $\pi$ is quasisimple, and we call $\chi$ the infinitesimal character of $\pi$. It is known that any irreducible unitary representation of $G$ is quasisimple.

Let $\eta$ be the conjugation of $\mathfrak{g}_{c}$ with respect to the real form $\mathfrak{g}$. We define three involutions on $\mathfrak{g}_{\mathrm{c}}$ by

$$
\begin{aligned}
X^{*} & =-\eta X, & & X \in \mathfrak{g}_{\mathbf{c}} \\
X^{+} & =-X, & & X \in \mathfrak{g}_{\mathbf{c}} \\
\bar{X} & =\eta X, & & X \in \mathfrak{g}_{\mathbf{c}}
\end{aligned}
$$

If $X$ and $Y$ are in $\mathfrak{g}_{\mathrm{c}}$ and $c$ is a complex number, it is easy to show that

$$
\begin{aligned}
& {[\bar{X}, \bar{Y}]=[\overline{X, Y}], \quad(\overline{c X})=\bar{c} \bar{X}} \\
& {\left[X^{+}, Y^{+}\right]=-[X, Y]^{+}, \quad(c X)^{+}=c X^{+}} \\
& {\left[X^{*}, Y^{*}\right]=-[X, Y]^{*}, \quad(c X)^{*}=\bar{c} X^{*}}
\end{aligned}
$$

All three involutions extend to involutions of $\mathscr{B}$.
If $\pi$ is a unitary representation of $G$, then for $g$ in $\mathscr{B}$,

$$
\pi\left(g^{*}\right)=\pi(g)^{*}
$$

where $\pi(g)^{*}$ is the adjoint operator of $\pi(g)$.
LEMMA 6. Suppose that $\pi$ is quasisimple. Assume $\Phi_{1}$ and $\Phi_{2}$ are vectors in $\mathscr{H}$ such that the vector spaces

$$
\operatorname{sp}\left\{\pi(k) \Phi_{2}: k \circledast K\right\}, \quad \alpha=1,2
$$

are both finite-dimensional. Then $\Phi_{1}$ and $\Phi_{2}$ are in $\mathscr{H} \infty$. Furthermore if $g_{1}$ and $g_{2}$ are in $\mathscr{B}$, and

$$
f(x)=\left(\Phi_{1}, \pi(x) \Phi_{2}\right)
$$

then

$$
f\left(g_{1} ; x ; g_{2}\right)=\left(\pi\left(g_{1}^{+}\right) \Phi_{1}, \pi(x) \pi\left(\overline{g_{2}}\right) \Phi_{2}\right)
$$

Proof. By [3(a), Theorem 6], $\Phi_{1}$ and $\Phi_{2}$ are actually "analytic vectors" for the representation $\pi$ so in particular they are differentiable. The other statement of the lemma is easy to check.

Now, in order to discuss the infinitesimal characters of the representations $\pi_{\omega}$ and $\pi_{\sigma, \Lambda}$ we shall quickly review how Harish-Chandra classifies homomorphisms from 3 to $\mathbf{C}$.

Let $\mathfrak{M}, \mathfrak{X}_{\mathfrak{p}}, \mathfrak{X}_{\mathfrak{f}}$, and $\mathfrak{A}$ be the universal enveloping algebras of $\mathfrak{m}_{c}, \mathfrak{a}_{\mathfrak{p}, \mathbf{c}}, \mathfrak{a}_{\mathfrak{p}, \mathbf{c}}$, and $\mathfrak{a}_{c}$ respectively. Let $\mathfrak{Z}_{M}$ be the center of $\mathfrak{M}$. Then $\mathfrak{M} \mathfrak{A}_{\mathfrak{p}}$ is the universal enveloping algebra of $m_{c}+\mathfrak{a}_{p, c}$, and its center is $\mathcal{Z}_{M} \mathfrak{A}_{p}$.

If $z$ is in $\mathfrak{Z}$, there exists a unique element $\gamma_{0}^{\prime}(z)$ in $\mathfrak{Z}_{M} \mathfrak{A}_{\mathfrak{p}}$ such that $z-\gamma_{0}^{\prime}(z)$ is in $\sum_{\alpha \in P_{+}} \mathscr{B} X_{\alpha}$ [see $\mathbf{3}(\mathbf{k})$, Lemma 13].

If $z_{1}$ is in $\mathfrak{Z}_{M} \mathfrak{A}_{\mathfrak{p}}$, there exists a unique element $\gamma_{1}^{\prime}\left(z_{1}\right)$ in $\mathfrak{A}$ such that $z_{1}-\gamma_{1}^{\prime}\left(z_{1}\right)$ is in $\sum_{\alpha \in P-P+} \mathfrak{W A}_{\mathfrak{p}} X_{\alpha}$ [see 3(e), Lemma 13].

If $z$ is in $\mathfrak{Z}$, there exists a unique element $\gamma^{\prime}(z)$ in $\mathfrak{A}$ such that $z-\gamma^{\prime}(z)$ is in $\sum_{\alpha \in P} \mathscr{A} X_{\alpha}$ [see 3(e), Lemma 18].

Notice that if $z$ is in 3 ,

$$
z-\gamma_{1}^{\prime}\left(\gamma_{0}^{\prime}(z)\right)=\left(z-\gamma_{0}^{\prime}(z)\right)+\left(\gamma_{0}^{\prime}(z)-\gamma_{1}^{\prime}\left(\gamma_{0}^{\prime}(z)\right)\right.
$$

The right-hand sum is an element in $\sum_{\alpha \in P} \mathscr{B} X_{\alpha}$. Therefore

$$
\begin{equation*}
\gamma_{1}^{\prime} \circ \gamma_{0}^{\prime}=\gamma^{\prime} \tag{6.1}
\end{equation*}
$$

Define automorphisms $\beta$ and $\beta_{1}$ of $\mathfrak{A}$ by

$$
\begin{array}{lr}
\beta(H)=H+\rho(H), & H \in \mathfrak{a}_{\mathbf{c}} \\
\beta_{1}(H)=H+\rho_{\mathfrak{m}}(H), & H \in \mathfrak{a}_{\mathbf{c}}
\end{array}
$$

Let $\gamma=\beta^{-1} \circ \gamma^{\prime}$ and let $\gamma_{1}=\beta_{1}^{-1} \circ \gamma_{1}^{\prime}$. It is known that the maps

$$
\gamma: \mathfrak{Z} \rightarrow \mathfrak{A}, \quad \gamma_{1}: \mathfrak{Z}_{M} \mathfrak{A}_{\mathfrak{p}} \rightarrow \mathfrak{A}
$$

are algebraic isomorphisms onto those elements in $\mathfrak{A}$ which are invariant under $W$ and $W_{1}$ respectively [ $\mathbf{3}(\mathbf{e})$, Lemma 19]. $\mathfrak{A}$ can be regarded as the algebra of polynomial functions from the dual space, $\mathfrak{a}_{\mathbf{c}}^{*}$, of $\mathfrak{a}$, into $\mathbf{C}$. If $\lambda$ is in $\mathfrak{a}_{\mathbf{c}}^{*}$, denote the evaluation of $p$ in $S\left(\mathfrak{a}_{\mathbf{c}}\right)$ at $\lambda$ by $\langle p, \lambda\rangle$. Then for any $\lambda$ in $\mathfrak{a}_{\mathbf{c}}^{*}$ define the homomorphism $\chi_{\lambda}: \mathcal{Z} \rightarrow \mathbf{C}$ by

$$
\chi_{\lambda}(z)=\langle\gamma(z), \lambda\rangle, \quad z \in \mathfrak{Z}
$$

Any homomorphism from 3 into $\mathbf{C}$ is of this form and $\chi_{\lambda_{1}}=\chi_{\lambda_{2}}$ if and only if $\lambda_{1}=s \lambda_{2}$ for some $s$ in $W$. We shall sometimes call $\chi_{\lambda}$ the homomorphism corresponding to the linear functional $\lambda$. Similarly, we can define $\chi_{\lambda}^{M}: \mathfrak{Z}_{M} \mathfrak{A}_{p} \rightarrow$ C by

$$
\chi_{\lambda}^{M}\left(z_{1}\right)=\left\langle\gamma_{1}\left(z_{1}\right), \lambda\right\rangle, \quad z_{1} \in \mathcal{Z}_{M} \mathfrak{A}_{\mathfrak{p}}
$$

Define an automorphism $\beta_{0}$ as follows:

$$
\begin{aligned}
& \beta_{0}(X)=X, \quad X \in \mathfrak{m}_{\mathbf{c}} \\
& \beta_{0}(H)=H+\rho(H), \quad H \in \mathfrak{a}_{\mathfrak{p}}
\end{aligned}
$$

$\beta_{0}$ takes $3_{M} \mathfrak{A}_{\mathfrak{p}}$ onto itself. Put $\gamma_{0}=\beta_{0}^{-1} \circ \gamma_{0}^{\prime}$. By (2.5) and (6.1) we see that

$$
\begin{equation*}
\gamma=\gamma_{1} \circ \gamma_{0} \tag{6.2}
\end{equation*}
$$

Now, what are infinitesimal characters of the representations in $\mathscr{E}_{d}$ ? It is known that all Cartan subalgebras of the complex Lie algebra $\mathfrak{g}_{c}$ are conjugate under an element in $G_{\mathbf{c}}$. Fix $y$ in $G_{c}$ such that $\operatorname{Ad} y \cdot \mathfrak{b}_{\mathbf{c}}=\mathfrak{a}_{\mathbf{c}}$. Then $\operatorname{Ad} y$ preserves the Killing form on $\mathfrak{g}_{c}$, and Ad $y$ maps the roots of ( $\mathfrak{g}_{c}, \mathfrak{b}_{c}$ ) onto the roots of ( $\mathfrak{g}_{c}, \mathfrak{a}_{c}$ ). Therefore Ad $y$ maps $(-1)^{1 / 2} \mathfrak{b}$ onto $(-1)^{1 / 2} \mathfrak{a}_{\mathfrak{t}}+\mathfrak{a}_{p}$. We can also assume that Ad $y$ maps the positive roots of $\left(g_{c}, b_{c}\right)$ onto the positive roots of $\left(g_{c}, a_{c}\right)$. If $\lambda$ is in $L$, define

$$
\lambda_{y}(H)=\lambda(\operatorname{Ad} y \cdot H), \quad H \in \mathfrak{a}_{c}
$$

Then

$$
\begin{equation*}
B(\lambda, \lambda)=B\left(\lambda_{y}, \lambda_{y}\right), \quad \lambda \in L \tag{6.3}
\end{equation*}
$$

Also

$$
\begin{equation*}
\tilde{\omega}^{\mathfrak{b}}(\lambda)=\tilde{\omega}^{\mathfrak{a}}\left(\lambda_{y}\right), \quad \lambda \in L \tag{6.4}
\end{equation*}
$$

If $\lambda$ is in $L^{\prime}$ and $\omega=\omega(\lambda)$ is in $\mathscr{E}_{d}$, then by [3(1), Theorems 15 and 16], $\chi_{\lambda_{y}}$ is the infinitesimal character of any representation in the equivalence class of $\omega$. Write $L_{y}$ and $L_{y}^{\prime}$ as the image of $L$ and $L^{\prime}$ under $y$.

Now if $\sigma$ is in $\mathscr{E}_{M}$, let $\mu$ be a real linear functional on $(-1)^{1 / 2} \mathfrak{a}_{t}$ associated with $\sigma$. Regard $\mu$ as a linear functional on $(-1)^{1 / 2} \mathfrak{a}_{\mathfrak{t}}+\mathfrak{a}_{p}$ by making it equal zero on $\mathfrak{a}_{\mathfrak{p}}$. By looking at a highest weight vector for $\sigma$, we can easily check that for any $z_{1}$ in $\mathfrak{Z}_{M}$

$$
\begin{equation*}
\sigma\left(z_{1}\right)=\chi_{\mu}^{M}\left(z_{1}\right)=\left\langle\gamma_{1}\left(z_{1}\right), \mu\right\rangle \tag{6.5}
\end{equation*}
$$

Lemma 7. Fix $\sigma$ in $\mathscr{E}_{M}$ and $\Lambda$ in $\mathbf{R}$. Then for any $z$ in $\mathfrak{Z}$,

$$
\pi_{\sigma, \Lambda}(z)=\chi-\mu-i \Lambda \mu_{0}(z)
$$

Proof. It is known that the representation $\pi_{\sigma, \Lambda}$ is quasi-simple [3(a), p. 243]. Therefore, there exists a complex linear functional $\lambda$ on $a_{c}$ such that

$$
\pi_{\sigma, \Lambda}(z)=\chi_{\lambda}(z), \quad z \in \mathfrak{Z}
$$

Our job is to evaluate $\lambda$.
Choose a $\tau$ in $\mathscr{E}_{K}^{2}$ such that $L_{\sigma}^{\tau} \neq 0$. Fix $\psi$ in $L_{\sigma}^{\tau}$ so that $\|\psi\|_{M}^{2}=\operatorname{dim} \sigma$. Then by (5.6) (using the notation in that formula) and Lemma 6, we obtain the formula

$$
\begin{aligned}
\overline{\chi_{\lambda}(\bar{z})} \cdot E_{\Lambda}(\psi: x) & =\sum_{i j}\left(t_{1} t_{2}\right)^{-1 / 2} \cdot \xi_{1 i} \otimes \xi_{2 j}^{*} \cdot\left(\Phi_{1 i}, \pi_{\sigma, \Lambda}(x) \pi_{\sigma, \Lambda}(\bar{z}) \Phi_{2 j}\right) \\
& =E_{\Lambda}(\psi: x ; z), \quad z \in \mathcal{Z}
\end{aligned}
$$

Let $F(x)=\psi(x) e^{\left(i \Lambda \mu_{0}-\rho\right)(H(x))}$, and define

$$
F(x: k)=F(x k) \tau\left(k^{-1}\right), \quad x \in G, k \in K
$$

Then by (5.5),

$$
E_{\Lambda}(\psi: x)=\int_{K} F(x: k) d k
$$

Let $z$ be an arbitrary element in $\mathfrak{3}$. It can be regarded as a left and right invariant differential operator, so

$$
F(x ; z: k)=F(x k ; z) \tau\left(k^{-1}\right)
$$

Therefore

$$
\begin{aligned}
E_{\Lambda}(\psi: x ; z) & =\int_{K} F(x ; z: k) d k \\
& =\int_{K} F(x k ; z) \tau\left(k^{-1}\right) d k
\end{aligned}
$$

Clearly $F(x n)=F(x)$ for any $n$ in $N$, so if $g$ is in $\mathscr{B} \mathfrak{n}, F(x ; g)=0$. Therefore

$$
\begin{aligned}
E_{\Lambda}(\psi: x ; z) & =\int_{K} F\left(x k ; \gamma_{0}^{\prime}(z)\right) \tau\left(k^{-1}\right) d k \\
& =\int_{K} F\left(x k ; \beta_{0} \gamma_{0}(z)\right) \tau\left(k^{-1}\right) d k
\end{aligned}
$$

Suppose that

$$
\gamma_{0}(z)=\sum_{i} z_{i} h_{i}, \quad z_{i} \in \mathfrak{Z}_{M}, h_{i} \in \mathfrak{A}_{\mathfrak{p}}
$$

Then $\beta_{0} \gamma_{0}(z)=\sum_{i} \beta_{0}\left(h_{i}\right)$.
Now $F\left(y \exp t H_{0}\right)=F(y)\left(i \Lambda t-\rho\left(t H_{0}\right)\right)$, so for any $h$ in $\mathfrak{A}_{p}$,

$$
\begin{aligned}
F\left(y ; \beta_{0}(h)\right) & =F(y)\left\langle\beta_{0}(h), i \Lambda \mu_{0}-\rho\right\rangle \\
& =F(y)\left\langle h, i \Lambda \mu_{0}\right\rangle
\end{aligned}
$$

On the other hand, if $m$ is in $M$,

$$
\begin{aligned}
F(y m) & =\psi(y m) e^{\left(i \Lambda \mu_{0}-\rho\right)(H(y))} \\
& =\tau_{1}(K(y)) \cdot T_{1}^{*} \cdot \sigma(m) \cdot T_{2} \cdot e^{\left(i \Lambda \mu_{0}-\rho\right)(H(y))}
\end{aligned}
$$

where $\psi(1)=T_{1}^{*} T_{2}$ in the notation of $\S 5$. Therefore, if $z_{M}$ is in $\mathcal{Z}_{M}$

$$
\begin{aligned}
F\left(y ; z_{M}\right) & =\tau_{1}(K(y)) \cdot T_{1}^{*} \cdot \sigma\left(z_{M}\right) \cdot T_{2} \cdot e^{\left(i \Lambda \mu_{0}-\rho\right)(H(y))} \\
& =\tau_{1}(K(y)) \cdot T_{1}^{*} \cdot\left\langle\gamma_{1}\left(z_{M}\right), \mu\right\rangle \cdot T_{2} \cdot e^{\left(i \Lambda \mu_{0}-\rho\right)(H(y))} \\
& =\left\langle\gamma_{1}\left(z_{M}\right), \mu\right\rangle F(y)
\end{aligned}
$$

Therefore we see that

$$
\begin{aligned}
F\left(y ; \beta_{0} \gamma_{0}(z)\right) & =F(y)\left\langle\sum_{i} \gamma_{1}\left(z_{i}\right) h_{i}, \mu+i \Lambda \mu_{0}\right\rangle \\
& =F(y)\left\langle\gamma_{1} \gamma_{0}(z), \mu+i \Lambda \mu_{0}\right\rangle \\
& =F(y)\left\langle\gamma(z), \mu+i \Lambda \mu_{0}\right\rangle
\end{aligned}
$$

Therefore,

$$
E_{\Lambda}(\psi: x ; z)=E_{\Lambda}(\psi: x)\left\langle\gamma(z), \mu+i \Lambda \mu_{0}\right\rangle
$$

It follows that

$$
\overline{\chi_{\lambda}(\bar{z})}=\left\langle\gamma(z), \mu+i \Lambda \mu_{0}\right\rangle
$$

However, $z$ was an arbitrary element in 3. It is easy to show, then, that $\lambda=$ $-\mu-i \Lambda \mu_{0}$. The proof of Lemma 7 is complete.

The linear transformation $\theta$ is orthogonal with respect to the bilinear form $B$. $\theta$ has eigenvalues of +1 and -1 on $(-1)^{1 / 2} a_{t}$ and $a_{t}$ respectively. This implies that $\mathfrak{a}_{\mathfrak{p}}$ is orthogonal to $(-1)^{1 / 2} \mathfrak{a}_{\mathfrak{p}}$ with respect to $B$.

Suppose the rank of g equals $n$. Then let $H_{1}$ and $H_{2}, \ldots, H_{n}$ be orthonormal bases of $\mathfrak{a}_{p}$ and $(-1)^{1 / 2} \mathfrak{a}_{t}$ respectively. For any $\alpha$ in $P$, fix root vectors $X_{\alpha}$ and $X_{-\alpha}$ in such a way that $B\left(X_{\alpha}, X_{-\alpha}\right)=1$. Then $\left[X_{\alpha}, X_{-\alpha}\right]=H_{\alpha}$. (If $\lambda$ is any linear functional on $\mathfrak{a}_{c}$, define $H_{\lambda}$ to satisfy the property

$$
\left.B\left(H_{\lambda}, H\right)=\lambda(H), \quad H \in \mathfrak{a}_{\mathbf{c}} .\right)
$$

The Casimir element of $\mathfrak{B}$ is given by

$$
\begin{equation*}
\omega_{\mathfrak{g}}=H_{1}^{2}+\cdots+H_{n}^{2}+\sum_{\alpha \in P}\left(X_{\alpha} X_{-\alpha}+X_{-\alpha} X_{\alpha}\right) \tag{6.6}
\end{equation*}
$$

$\omega_{g}$ is in 3 and is equal to the expression

$$
H_{1}^{2}+\cdots+H_{n}^{2}+\sum_{\alpha \in P}\left(\left[X_{\alpha}, X_{-\alpha}\right]+2 X_{-\alpha} X_{\alpha}\right)
$$

so we see that

$$
\begin{equation*}
\gamma^{\prime}\left(\omega_{\mathfrak{g}}\right)=H_{1}^{2}+\cdots+H_{n}^{2}+\sum_{\alpha \in P} H_{\alpha}=\sum_{i} H_{i}^{2}+2 H_{\rho} \tag{6.7}
\end{equation*}
$$

Therefore,

$$
\begin{aligned}
\gamma\left(\omega_{\mathfrak{g}}\right) & =\sum_{i}\left(H_{i}-\rho\left(H_{i}\right)\right)^{2}+2 H_{\rho}-2 \rho\left(H_{\rho}\right) \\
& =\sum_{i} H_{i}^{2}-\sum_{i} 2 \rho\left(H_{i}\right) H_{i}+\sum_{i} \rho\left(H_{i}\right) \cdot \rho\left(H_{i}\right)+2 H_{\rho}-2 B(\rho, \rho) \\
& =\sum_{i} H_{i}^{2}-B(\rho, \rho)
\end{aligned}
$$

Then if $\lambda$ is any linear functional on $\mathfrak{a}_{c}$, we have the formula

$$
\begin{equation*}
\chi_{\lambda}\left(\omega_{\mathfrak{g}}\right)=B(\lambda, \lambda)-B(\rho, \rho) \tag{6.8}
\end{equation*}
$$

Suppose that $\pi_{\omega}$ is a representation in the class of some $\omega$ in $\mathscr{E}_{d}$. Then from (6.3) and (6.8) we obtain the formula

$$
\begin{equation*}
\pi_{\omega}\left(\omega_{\mathfrak{g}}\right)=|\omega|^{2}-B(\rho, \rho) \tag{6.9}
\end{equation*}
$$

The restriction of any root in $P_{+}$to $\mathfrak{a}_{\mathfrak{p}}$ is equal to either $\mu_{0}$ or $2 \mu_{0}$. Suppose there are $r_{1}$ and $r_{2}$ roots in $P_{+}$of each type. Then

$$
B\left(H_{0}, H_{0}\right)=2 \sum_{\beta \in P_{+}} \beta\left(H_{0}\right)^{2}=2\left(r_{1}+4 r_{2}\right)
$$

Let $r^{2}=2\left(r_{1}+4 r_{2}\right)$. Then $B\left(\mu_{0}, \mu_{0}\right)=r^{-2}$. If $\mu$ is any linear functional on $\mathfrak{a}_{\mathfrak{t}}$,

$$
B\left(-\mu-i \Lambda \mu_{0},-\mu-i \Lambda \mu_{0}\right)=B(\mu, \mu)-r^{-2} \Lambda^{2}
$$

Therefore if $\sigma$ is in $\mathscr{E}_{M}$, and $\Lambda$ is in $\mathbf{R}$,

$$
\begin{equation*}
\pi_{\sigma, \Lambda}\left(\omega_{\mathfrak{g}}\right)=|\sigma|^{2}-r^{-2} \Lambda^{2}-B(\rho, \rho) \tag{6.10}
\end{equation*}
$$

The Lie algebra $\mathfrak{k}$ is reductive, so $\mathfrak{k}=\mathfrak{k}_{1}+\mathfrak{k}_{2}$, where $\mathfrak{k}_{1}$ is semisimple and $\mathfrak{k}_{2}$ is abelian. Let $\mathscr{K}$ be the universal enveloping algebra of $\mathfrak{t}_{\mathbf{c}}$, and let $\mathcal{Z}_{K}$ be its center. Let $\mathfrak{b}$ be the Cartan subalgebra of $\mathfrak{k}$ defined at the end of $\S 2$. For linear functionals $\nu$ on $\mathfrak{h}_{c}$ we can define the homomorphism

$$
\chi_{\nu}^{K}: \mathfrak{Z}_{K} \rightarrow \mathbf{C}
$$

Since the Killing form of $\mathfrak{g}_{c}$ when restricted to $\mathfrak{k}_{1}$ is $K$-invariant, this Killing form is a linear combination of the Killing forms of $\mathfrak{k}_{1, i}$ on $\mathfrak{k}_{1}$, where $\left\{\mathfrak{k}_{1, i}\right\}$ are the simple ideals of $\mathfrak{k}_{1}$. Then it is clear that we can choose an element $\omega_{\mathfrak{k}}$ in $\mathfrak{Z}_{K}$ such that

$$
\begin{equation*}
\chi_{\nu}^{K}\left(\omega_{\mathfrak{l}}\right)=B(\nu, \nu)-B\left(\rho_{K}, \rho_{K}\right) \tag{6.11}
\end{equation*}
$$

for any linear functional $\nu$ on $\mathfrak{h}_{\mathbf{c}}$.
Notice that

$$
\begin{array}{cll}
\omega_{\mathfrak{g}}^{*}=\omega_{\mathfrak{g}}, & \bar{\omega}_{\mathfrak{g}}=\omega_{\mathfrak{g}}, & \omega_{\mathfrak{g}}^{+}=\omega_{\mathfrak{g}} \\
\omega_{\mathfrak{k}}^{*-}=\omega_{\mathfrak{k}}, & \bar{\omega}_{\mathfrak{k}}=\omega_{\mathfrak{k}}, & \omega_{\mathfrak{k}}^{+}=\omega_{\mathfrak{k}} \tag{6.12}
\end{array}
$$

LEMMA 8. The map $f \rightarrow \hat{f}_{0}$ is a continuous map from $\mathscr{C}(G)$ into $\mathscr{C}_{0}(\hat{G})$.
Proof. Fix polynomials $p, q_{1}, q_{2}$. For $f$ in $C_{0}^{\infty}(G)$, we have

$$
\left\|\hat{f}_{0}\right\|_{p, q_{1}, q_{2}}=\sup _{0}\left|\left(\Phi_{\tau_{1}, i_{1}}, \hat{f}_{0}(\omega) \Phi_{\tau_{2}, i_{2}}\right)\right| \cdot p(|\omega|) \cdot q_{1}\left(\left|\tau_{1}\right|\right) \cdot q_{2}\left(\left|\tau_{2}\right|\right)
$$

By sup we mean the supremum over all $\omega,\left(\tau_{1}, i_{1}\right),\left(\tau_{2}, i_{2}\right)$. This is an arbitrary continuous seminorm on $\mathscr{C}_{0}(\hat{G})$.

Choose integers $m, n_{1}, n_{2}$ such that

$$
p(|\omega|) \cdot q_{1}\left(\left|\tau_{1}\right|\right) \cdot q_{2}\left(\left|\tau_{2}\right|\right) \leq\left(1+|\omega|^{2}\right)^{m} \cdot\left(1+\left|\tau_{1}\right|^{2}\right)^{n_{1}} \cdot\left(1+\left|\tau_{2}\right|^{2}\right)^{n_{2}}
$$

Define elements $g_{1}$ and $g_{2}$ in $\mathfrak{B}$ by

$$
\begin{gathered}
g_{1}=\left(\omega_{\mathfrak{k}}+B\left(\rho_{K}, \rho_{K}\right)+1\right)^{n_{1}} \\
g_{2}=\left(\omega_{\mathfrak{g}}+B(\rho, \rho)+I\right)^{m} \cdot\left(\omega_{\mathfrak{k}}+B\left(\rho_{K}, \rho_{K}\right)+1\right)^{n_{2}}
\end{gathered}
$$

By (6.12), $g_{1}^{+}=g_{1}$ and $\bar{g}_{2}=g_{2}$. Since $\Phi_{\tau_{1}, i_{1}}$ and $\Phi_{\tau_{2}, i_{2}}$ transform under $\pi_{\omega}$ according to $\tau_{1}$ and $\tau_{2}$ respectively, we see that

$$
\begin{gathered}
\pi_{\omega}\left(g_{1}\right) \Phi_{\tau_{1}, i_{1}}=\left(1+\left|\tau_{1}\right|^{2}\right)^{n_{1}} \Phi_{\tau_{1}, i_{1}} \\
\pi_{\omega}\left(g_{2}\right) \Phi_{\tau_{2}, i_{2}}=\left(1+\left|\tau_{2}\right|^{2}\right)^{n_{2}} \cdot\left(1+|\omega|^{2}\right)^{m} \Phi_{\tau_{2}, i_{2}}
\end{gathered}
$$

Therefore,

$$
\left\|\hat{f}_{0}\right\|_{p, q_{1}, q_{2}} \leq \sup _{0}\left|\left(\pi_{\omega}\left(g_{1}\right) \Phi_{\tau_{1}, i_{1}}, \hat{f}_{0}(\omega) \pi_{\omega}\left(g_{2}\right) \Phi_{\tau_{2}, i_{2}}\right)\right|
$$

If $h(x)$ is the function $\left(\Phi_{\tau_{1}, i_{1}}, \pi_{\omega}(x) \Phi_{\tau_{2}, i_{2}}\right)$, we see from Lemma 6 that

$$
\begin{aligned}
\left\|\hat{f}_{0}\right\|_{p, q_{1}, q_{2}} & \leq \sup _{0}\left|\int_{G} \overline{f(x)} \cdot h\left(g_{1} ; x ; g_{2}\right) d x\right| \\
& =\sup _{0}\left|\int_{G} \overline{f\left(g_{1} ; x ; g_{2}\right)} \cdot h(x) d x\right| \\
& \leq \sup _{0}\left(\int_{G}\left|f\left(g_{1} ; x ; g_{2}\right)\right|^{2} d x\right)^{1 / 2} \cdot\left(\int_{G}|h(x)|^{2} d x\right)^{1 / 2} .
\end{aligned}
$$

By (2.8) this last expression equals

$$
\sup _{0}\left(\int_{G}\left|f\left(g_{1} ; x ; g_{2}\right)\right|^{2} d x\right)^{1 / 2} \cdot \beta(\omega)^{-1 / 2}
$$

But $\left\{\beta(\omega)^{-1 / 2}: \omega \in \mathscr{E}_{d}\right\}$ is bounded by Lemma 3. Also, by (4.2),

$$
\begin{aligned}
& \left(\int_{G}\left|f\left(g_{1} ; x ; g_{2}\right)\right|^{2} d x\right)^{1 / 2} \\
& \quad \leq \sup _{x \in G}\left(\Xi(x)^{-1} \cdot(1+\sigma(x))^{r_{0} / 2} \cdot\left|f\left(g_{1} ; x ; g_{2}\right)\right|\right) \cdot N\left(r_{0}\right)^{1 / 2}
\end{aligned}
$$

We have bounded $\left\|\hat{f}_{0}\right\|_{p, q_{1}, q_{2}}$ by a continuous seminorm for $\mathscr{C}(G)$ on $f$. Since $C_{0}^{\infty}(G)$ is dense in $\mathscr{C}(G)$, we have proved the lemma.

Lemma 9. Let sup denote the supremum over all $(\sigma, \Lambda),\left(\tau_{1}, i_{1}\right),\left(\tau_{2}, i_{2}\right)$. Then for nonnegative integers $m, m_{1}, m_{2}$, and $n$, the continuous seminorms

$$
\begin{aligned}
\left\|a_{1}\right\|_{\left(m, m_{1}, m_{2}, n\right)}= & \sup _{1}\left|\left(\frac{d}{d \Lambda}\right)^{n}\left[\left(\Phi_{\tau_{1}, i_{1}}, a_{1}(\sigma, \Lambda) \Phi_{\tau_{2}, i_{2}}\right) \cdot\left(1+|\sigma|^{2}-r^{-2} \Lambda^{2}\right)^{m}\right]\right| \\
& \cdot\left(1+\left|\tau_{1}\right|^{2}\right)^{m_{1}} \cdot\left(1+\left|\tau_{2}\right|^{2}\right)^{m_{2}}, \quad a_{1} \in \mathscr{C}_{1}(\hat{G})
\end{aligned}
$$

form a base for the topology of $\mathscr{C}_{1}(\hat{G})$.
Proof. By Leibnitz' rule and induction on $n$, we can see that it is enough to prove the lemma for the seminorms

$$
\begin{gather*}
\sup _{1}\left|\left(\frac{d}{d \Lambda}\right)^{n}\left(\Phi_{\tau_{1}, i_{1}}, a_{1}(\sigma, \Lambda) \Phi_{\tau_{2}, i_{2}}\right)\right| \cdot\left|\left(1+|\sigma|^{2}-r^{2} \Lambda^{2}\right)\right|^{m}  \tag{6.13}\\
\cdot\left(1+\left|\tau_{1}\right|^{2}\right)^{m_{1}} \cdot\left(1+\left|\tau_{2}\right|^{2}\right)^{m_{2}}, \quad a_{1} \in \mathscr{C}_{1}(\hat{G})
\end{gather*}
$$

Fix $\sigma$ in $\mathscr{E}_{M}$ and $\tau$ in $\mathscr{E}_{K}$. There is a nonzero vector of the form $\Phi_{\tau, i}$ in our basis for $\mathscr{H}_{\sigma}$ if and only if the representation $\sigma$ occurs in $\left.\tau\right|_{M}$. Suppose that this is the case. Then if $\tau$ acts on the finite-dimensional vector space $V_{\tau}, \sigma$ acts on a subspace of $V_{\tau}$. The Cartan subalgebras of $\mathfrak{m}$ and $\mathfrak{k}$ are $\mathfrak{a}_{\mathfrak{k}}$ and $\mathfrak{h}$ respectively, and we have already ordered their dual spaces. Regard $\sigma$ as a representation $\sigma_{0} \times \varepsilon$ of $M^{0} \times Z(A)$ as in $\S 2$, and let $\nu$ and $\mu$ be the highest weights of the representations $\tau$ and $\sigma_{0}$ respectively. Then

$$
|\sigma|^{2}=B\left(\mu+\rho_{M}, \mu+\rho_{M}\right), \quad|\tau|^{2}=B\left(\nu+\rho_{K}, \nu+\rho_{K}\right)
$$

Write $\mathfrak{h}=\mathfrak{a}_{\mathfrak{k}} \oplus \mathfrak{b}_{2}$ where $\mathfrak{b}_{2}$ is a Lie algebra of dimension 1 or 0 , depending on whether we are in Case I or II. Assume that $B\left(\mathfrak{a}_{\mathfrak{f}}, \mathfrak{b}_{2}\right)=0$. Extend $\mu$ to a linear functional on $\mathfrak{h}$ by letting it equal zero on $\mathfrak{b}_{2}$.

Let $\xi$ be a highest weight vector in $V_{\tau}$ for $\sigma_{0} . V_{\tau}$ is a direct sum of weight spaces for $\tau$. Examine the action of $\tau(\mathfrak{h})$ on $\xi$. It is clear that there is a linear functional $\mu_{2}$ on $\mathfrak{h}$ which is zero on $\mathfrak{a}_{\mathfrak{k}}$ and such that $\nu_{1}=\mu+\mu_{2}$ is a weight for $\tau$. Therefore

$$
B(\mu, \mu)<B\left(\nu_{1}, \nu_{1}\right)
$$

However,

$$
B\left(\nu_{1}, \nu_{1}\right)^{1 / 2} \gtrless B\left(\nu_{1}+\rho_{K}, \nu_{1}+\rho_{K}\right)^{1 / 2}+B\left(\rho_{K}, \rho_{K}\right)^{1 / 2}
$$

Since $\nu$ is a highest weight for $\tau$, we see by [5, Lemma 3, p. 248] that

$$
B\left(\nu_{1}+\rho_{K}, \nu_{1}+\rho_{K}\right) \gtrless B\left(\nu+\rho_{K}, \nu+\rho_{K}\right)
$$

We have shown that there is a constant $C$, independent of $\tau$ and $\sigma$, such that

$$
\begin{equation*}
|\sigma| ₹|\tau|+C . \tag{6.14}
\end{equation*}
$$

From (6.14) we obtain the additional formula

$$
\begin{equation*}
|\Lambda|^{2}<r^{2}\left|\left(1+|\sigma|^{2}-r^{2}|\Lambda|^{2}\right)\right|+r^{2}\left(1+(C+|\tau|)^{2}\right) \tag{6.15}
\end{equation*}
$$

Formulas (6.14) and (6.15) show that any seminorm of the form (4.8) is dominated by a seminorm of the form (6.13). Since the seminorms (4.8) form a base for the topology of $\mathscr{C}_{1}(\hat{G})$, our lemma is proved.

Lemma 10. The map $f \rightarrow \hat{f}_{1}$ is a continuous map from $\mathscr{C}(G)$ into $\mathscr{C}_{1}(\hat{G})$.
Proof. Let $\|\cdot\|_{1}$ be an arbitrary continuous seminorm on $\mathscr{E}_{1}(\hat{G})$. Since $\operatorname{dim} \sigma$, $\operatorname{dim} \tau_{1}$, and $\operatorname{dim} \tau_{2}$ are bounded by polynomials in $|\sigma|,\left|\tau_{1}\right|$, and $\left|\tau_{2}\right|$ respectively, we can use Lemma 9 to choose integers $m, m_{1}, m_{2}, n$ such that for any $a_{1}$ in $L_{1}^{2}(\hat{G})$,

$$
\begin{aligned}
\left\|a_{1}\right\|_{1} ₹ & \sup _{1}\left(\operatorname{dim} \sigma \cdot \operatorname{dim} \tau_{1} \cdot \operatorname{dim} \tau_{2}\right)^{-1 / 2} \\
& \cdot\left|\left(\frac{d}{d \Lambda}\right)^{n}\left[\left(\Phi_{\tau_{1}, i_{1}}, a_{1}(\sigma, \Lambda) \Phi_{\tau_{2}, i_{2}}\right)\left(|\sigma|^{2}-r^{-2} \Lambda^{2}\right)^{m}\right]\right| \\
& \cdot\left(1+\left|\tau_{1}\right|^{2}\right)^{m_{1}} \cdot\left(1+\left|\tau_{2}\right|^{2}\right)^{m_{2}}
\end{aligned}
$$

Therefore, for any $f$ in $C_{0}^{\infty}(G)$

$$
\left\|\hat{f}_{1}\right\|_{1} \leq \sup _{1}\left(\operatorname{dim} \sigma \cdot \operatorname{dim} \tau_{1} \cdot \operatorname{dim} \tau_{2}\right)^{1 / 2} \cdot\left|\left(\frac{d}{d \Lambda}\right)^{n} \int_{G} h\left(g_{1} ; x ; g_{2}\right) \cdot \overline{f(x)} d x\right|
$$

by (6.10). Here

$$
\begin{gathered}
h(x)=\left(\Phi_{\tau_{1}, i_{1}}, \pi_{\sigma, \Lambda}(x) \Phi_{\tau_{2}, i_{2}}\right) \\
g_{1}=\left(\omega_{\mathfrak{k}}+B\left(\rho_{K}, \rho_{K}\right)+1\right)^{m_{1}} \\
g_{2}=\left(\omega_{\mathfrak{g}}+B(\rho, \rho)+1\right)^{m} \cdot\left(\omega_{\mathfrak{k}}+B\left(\rho_{K}, \rho_{K}\right)+1\right)^{m_{2}} .
\end{gathered}
$$

Therefore
$\left\|\hat{f}_{1}\right\|_{1}<\sup _{i}\left(\operatorname{dim} \sigma \cdot \operatorname{dim} \tau_{1} \cdot \operatorname{dim} \tau_{2}\right)^{-1 / 2} \cdot\left|\left(\frac{d}{d \Lambda}\right)^{n} \int_{G} h(x) \cdot \overline{f\left(g_{1} ; x ; g_{2}\right)} d x\right|$.
Let $\tau$ be the double representation $\left(\tau_{1}, \tau_{2}\right)$ of $K$. Then by (5.6) there is a $\psi$ in $L_{\sigma}^{\tau}$, with $\|\psi\|_{M}^{2}=\operatorname{dim} \sigma$, such that the last expression in the above inequality is bounded by

$$
\sup _{1}(\operatorname{dim} \sigma)^{-1 / 2} \cdot \int_{G}\left|f\left(g_{1} ; x ; g_{2}\right)\right| \cdot\left|\left(\frac{d}{d \Lambda}\right)^{n} E_{\Lambda}(\psi: x)\right| d x
$$

Now if $k$ is in $K$,

$$
|\psi(k)|=|\psi(1)|=\|\psi\|_{M}=(\operatorname{dim} \sigma)^{1 / 2}
$$

Therefore, by (5.5),

$$
\begin{aligned}
& (\operatorname{dim} \sigma)^{-1 / 2}\left|\left(\frac{d}{d \Lambda}\right)^{n} E_{\Lambda}(\psi: x)\right| \\
& \quad=(\operatorname{dim} \sigma)^{-1 / 2}\left|\int_{K} \psi(x k) \tau\left(k^{-1}\right) \cdot\left[i \mu_{0}(H(x k))\right]^{n} \cdot e^{\left(i \Lambda \mu_{0}-\rho\right)(H(x k))} d k\right| \\
& \quad \leq \int_{K}\left|\mu_{0}(H(x k))\right|^{n} \cdot e^{-\rho(H(x k))} d x
\end{aligned}
$$

Lemma 11. For any $x$ in $G$,

$$
\left|\mu_{0}(H(x))\right| ₹ r^{-1} \sigma(x)
$$

Assuming the proof of Lemma 11 for the moment, we see that

$$
\begin{aligned}
(\operatorname{dim} \sigma)^{-1 / 2}\left|\left(\frac{d}{d \Lambda}\right)^{n} E_{\Lambda}(\psi: x)\right| & \gtrless \int_{K} r^{-n} \sigma(x k)^{n} e^{-\rho(H(x k))} d k \\
& =r^{-n} \sigma(x)^{n} \cdot \int_{K} e^{-\rho(H(x k))} d k \\
& =r^{-n} \sigma(x)^{n} \cdot \Xi(x)
\end{aligned}
$$

Therefore

$$
\left\|\hat{f}_{1}\right\|_{1} \gtrless r^{-n} \int_{G}\left|f\left(g_{1} ; x ; g_{2}\right)\right| \cdot \sigma(x)^{n} \cdot \Xi(x) d x
$$

By (4.2), the right-hand side of this inequality is bounded by

$$
r^{-n} \cdot N\left(r_{0}\right) \cdot \sup _{x \in G}\left(\Xi(x)^{-1}(1+\sigma(x))^{n+r_{0}} \cdot\left|f\left(g_{1} ; x ; g_{2}\right)\right|\right)
$$

We have dominated $\left\|\hat{f}_{1}\right\|_{1}$ by a continuous seminorm for $\mathscr{C}(G)$ on $f$. Since $C_{0}^{\infty}(G)$ is dense in $\mathscr{C}(G)$, this is enough to prove Lemma 10.

We still must prove Lemma 11. If $x$ is in $G$, we can find $k^{\prime}, k$ in $K$, and $t$ in $\mathbf{R}$ such that $x=k^{\prime} \cdot \exp t H_{0} \cdot k$. There exists a real number $t_{k}$ such that

$$
H(x)=H\left(\exp t H_{0} \cdot k\right)=t_{k} H_{0}
$$

Now [3(g), Lemma 35] and [3(g), Lemma 35, Corollary 2] establish precisely that $-t ₹ t_{k} ₹ t$. Therefore,

$$
\left|\mu_{0}(H(x))\right|=\left|t_{k}\right|<|t|=r^{-1} \sigma(x)
$$

Lemmas 8 and 10 show that the Fourier transform is a continuous map from $\mathscr{C}(G)$ into $\mathscr{C}(\hat{G})$. It will be more difficult to prove surjectivity.
7. Theorem $3^{\prime}$ and some elementary formulae. Let $a=\left(a_{0}, a_{1}\right)$ be an element in $L^{2}(\hat{G})$. There exists a unique function $f$ in $L^{2}(G)$ such that $\hat{f}=a$. We wish to find a formula for $f . f$ is the unique function in $L^{2}(G)$ such that for every $g$ in $C_{0}^{\infty}(G)$,

$$
\int_{G} f(x) \overline{g(x)} d x=(a, \hat{g})
$$

(The latter inner product is that of $L^{2}(\hat{G})$.) We shall write $\sum_{0}$ to denote summation over all $\omega,\left(\tau_{1}, i_{1}\right),\left(\tau_{2}, i_{2}\right)$. We write $\sum_{1}$ to denote summation over all $\sigma$, $\left(\tau_{1}, i_{1}\right),\left(\tau_{2}, i_{2}\right)$.

Then ( $a, \hat{g}$ ) equals

$$
\begin{aligned}
& \sum_{0}\left(a_{0}(\omega) \Phi_{\tau_{2}, i_{2}}, \Phi_{\tau_{1}, i_{1}}\right)\left(\Phi_{\tau_{1}, i_{1}}, \hat{g}_{0}(\omega) \Phi_{\tau_{2}, i_{2}}\right) \beta(\omega) \\
&+\sum_{1} \int_{-\infty}^{\infty}\left(a_{1}(\sigma, \Lambda) \Phi_{\tau_{2}, i_{2}}, \Phi_{\tau_{1}, i_{1}}\right)\left(\Phi_{\tau_{1}, i_{1}}, \hat{g}_{1}(\sigma, \Lambda) \Phi_{\tau_{2}, i_{2}}\right) \cdot \beta(\sigma, \Lambda) d \Lambda \\
&= \sum_{0} \int_{G}\left(a_{0}(\omega) \Phi_{\tau_{2}, i_{2}}, \Phi_{\tau_{1}, i_{1}}\right)\left(\Phi_{\tau_{1}, i_{1}}, \pi_{\omega}(x) \Phi_{\tau_{2}, i_{2}}\right) \cdot \beta(\omega) \overline{g(x)} d x \\
& \quad+\sum_{1} \int_{-\infty}^{\infty} \int_{G}\left(a_{1}(\sigma, \Lambda) \Phi_{\tau_{2}, i_{2}}, \Phi_{\tau_{1}, i_{1}}\right)\left(\Phi_{\tau_{1}, i_{1}}, \pi_{\sigma, \Lambda}(x) \Phi_{\tau_{2}, i_{2}}\right) \\
& \quad \cdot \beta(\sigma, \Lambda) \cdot \overline{g(x)} d x d \Lambda .
\end{aligned}
$$

Let us assume that the integrals in the last expression are absolutely convergent. This is true for example if $a$ is in $\mathscr{C}(\hat{G})$. Then we may take the integration on $G$ outside. Define functions $\check{a}_{0}(x)$ and $\check{a}_{1}(x)$ by

$$
\begin{align*}
& \check{a}_{0}(x)=\sum_{0}\left(a_{0}(\omega) \Phi_{\tau_{2}, i_{2}}, \Phi_{\tau_{1}, i_{1}}\right)\left(\Phi_{\tau_{1}, i_{1}}, \pi_{\omega}(x) \Phi_{\tau_{2}, i_{2}}\right) \beta(\omega) \\
& 㐅_{1}(x)=\sum_{1} \int_{-\infty}^{\infty}\left(a_{1}(\sigma, \Lambda) \Phi_{\tau_{2}, i_{2}}, \Phi_{\tau_{1}, i_{1}}\right)\left(\Phi_{\tau_{1}, i_{1}}, \pi_{\sigma, \Lambda}(x) \Phi_{\tau_{2}, i_{2}}\right) \beta(\sigma, \Lambda) d \Lambda \tag{7.1}
\end{align*}
$$

Then we have the formula

$$
\int_{G} f(x) \overline{g(x)} d x=\int_{G}\left(\check{a}_{0}(x)+\check{a}_{1}(x)\right) \overline{g(x)} d x
$$

Since the function $g(x)$ is arbitrary we see that

$$
f(x)=\check{a}_{0}(x)+x_{1}(x)
$$

By definition of $f(x)$, we have the formulae

$$
\begin{equation*}
\hat{a}_{0}=a_{0}, \quad \hat{a}_{1}=a_{1} \tag{7.2}
\end{equation*}
$$

To prove the surjectivity in Theorem 3, we have to show that both $\chi_{0}(x)$ and $\check{a}_{1}(x)$ are in $\mathscr{C}(G)$. If $h$ is in $\mathscr{S}(\mathbf{R})$, the Schwartz space of $\mathbf{R}$, we shall be
studying the functions

$$
\begin{aligned}
& f_{0}(x)=\left(\Phi_{\tau_{1}, i_{1}}, \pi_{\omega}(x) \Phi_{\tau_{2}, i_{2}}\right) \beta(\omega) \\
& f_{1}(x)=\int_{-\infty}^{\infty} h(\Lambda)\left(\Phi_{\tau_{1}, i_{1}}, \pi_{\sigma, \Lambda}(x) \Phi_{\tau_{2}, i_{2}}\right) \beta(\sigma, \Lambda) d \Lambda
\end{aligned}
$$

If $\nu$ is a continuous seminorm on $\mathscr{C}(G)$, we will need to know that both $\nu\left(f_{0}\right)$ and $\nu\left(f_{1}\right)$ are finite. Furthermore, we will have to determine the dependence of $\nu\left(f_{0}\right)$ and $\nu\left(f_{1}\right)$ on the variables $\tau_{1}, \tau_{2}, \sigma, \omega$, and $h$.

It follows from the remarks of $\S 5$ that the study of the functions $f_{0}(x)$ and $f_{1}(x)$ is equivalent to the study of certain $\left(\tau_{1}, \tau_{2}\right)$ spherical functions. To complete the proof of Theorem 3 it is enough to prove the following

THEOREM 3'. Let $\tau$ be in $\mathscr{E}_{K}^{2}$.
(a) Let $\phi_{\omega}^{\tau}$ be a $\tau$-spherical function of the form (5.1) corresponding to a representation $\pi_{\omega}$ with $\omega$ in $\mathscr{E}_{d}$. Then for each $g_{1}, g_{2}$ in $\mathscr{B}$, and real $s$, there are polynomials $p, q$ such that

$$
\sup _{x \in G}\left|\phi_{\omega}^{\tau}\left(g_{1} ; x ; g_{2}\right) \beta(\omega) \Xi(x)^{-1}(1+\sigma(x))^{s}\right| \gtrless p(|\omega|) q(|\tau|) .
$$

(b) If $\sigma$ is in $\mathscr{E}_{M}$, and $\psi$ is in $L_{\sigma}^{\tau}$, with $\|\psi\|_{M}=1$, let $E_{\Lambda}(\psi: x)$ be the Eisenstein integral as in (5.6). Then for each $g_{1}, g_{2}$ in $\mathscr{B}$, and real $s$, there exist polynomials $p_{1}, p_{1}, q$, and an integer $N$, such that whenever $h$ is in $\mathscr{S}(\mathbf{R})$,

$$
\begin{aligned}
& \sup _{x \in G}\left|\int_{-\infty}^{\infty} h(\Lambda) E_{\Lambda}\left(\psi: g_{1} ; x ; g_{2}\right) \beta(\sigma, \Lambda) d \Lambda \cdot \Xi(x)^{-1}(1+\sigma(x))^{s}\right| \\
& \gtrless p_{1}(|\sigma|) q(|\tau|) \cdot \sup _{1 \gtrless t<N} \cdot \sup _{\Lambda \in \mathbf{R}}\left(p_{2}(|\Lambda|)\right)\left|\left(\frac{d}{d \Lambda}\right)^{t} h(\Lambda)\right| .
\end{aligned}
$$

We shall devote most of our remaining work to proving Theorem $3^{\prime}$.
Let $\mathscr{L}_{1}^{2}(\hat{G})$ and $\mathscr{S}_{1}(\hat{G})$ be defined analogously to $L_{1}^{2}(\hat{G})$ and $\mathscr{C}_{1}(\hat{G})$ respectively, but without the symmetry condition

$$
a_{1}\left(\sigma^{\prime},-\Lambda\right)=N_{\sigma}(\Lambda) a_{1}(\sigma, \Lambda) N_{\sigma}(\Lambda)^{-1}
$$

Now suppose that $a_{1}$ is in $\mathscr{L}_{1}^{2}(\hat{G})$, and that

$$
\sum_{1} \int_{-\infty}^{\infty}\left(a_{1}(\sigma, \Lambda) \Phi_{\tau_{2}, i_{2}}, \Phi_{\tau_{1}, i_{1}}\right) \beta(\sigma, \Lambda) d \Lambda
$$

is absolutely convergent. This is true in particular if $a_{1}$ is in $\mathscr{S}_{1}(\hat{G})$. Then we define $x_{1}(x)$ by the equation in (7.1).

Suppose that for each $\sigma$ in $\mathscr{E}_{M},\left\{\Phi_{\tau, i}^{\prime}\right\}$ is another orthonomal basis for $\mathscr{H}_{\sigma}$ such that each $\Phi_{\tau, i}^{\prime}$ transforms under $\pi_{\sigma}$ according to the irreducible representation $\tau$ of $K$. Then it is not hard to show that the expression

$$
\sum_{1} \int_{-\infty}^{\infty}\left(a_{1}(\sigma, \Lambda) \Phi_{\tau_{2}, i_{2}}^{\prime}, \Phi_{\tau_{1}, i_{1}}^{\prime}\right)\left(\Phi_{\tau_{1}, i_{1}}^{\prime}, \pi_{\sigma, \Lambda}(x) \Phi_{\tau_{2}, i_{2}}^{\prime}\right) \beta(\sigma, \Lambda) d \Lambda
$$

is also absolutely convergent and equal to $\partial_{1}(x)$. Define

$$
b_{1}(\sigma, \Lambda)=\frac{1}{2}\left(a_{1}(\sigma, \Lambda)+N_{\sigma}(\Lambda)^{-1} a_{1}\left(\sigma^{\prime},-\Lambda\right) N_{\sigma}(\Lambda)\right)
$$

Then, using the above remark we can show that

$$
\check{b}_{1}(x)=\check{a}_{1}(x)
$$

Then by (7.2), $\hat{a}_{1}(\sigma, \Lambda)=b(\sigma, \Lambda)$. Therefore, we have the formula

$$
\begin{equation*}
\hat{a}_{1}(\sigma, \Lambda)=\frac{1}{2}\left[a_{1}(\sigma, \Lambda)+N_{\sigma}(\Lambda)^{-1} a_{1}\left(\sigma^{\prime},-\Lambda\right) N_{\sigma}(\Lambda)\right] . \tag{7.3}
\end{equation*}
$$

Now let $\tau=\left(\tau_{1}, \tau_{2}\right)$ be an irreducible double representation of $K$. It is of interest to define Fourier transforms for $\tau$-spherical functions on $G$.

Choose $\sigma_{1}, \ldots, \sigma_{r}$ in $\mathscr{E}_{M}$ such that

$$
L^{\tau}=L_{\sigma_{1}}^{\tau} \oplus \cdots \oplus L_{\sigma_{r}}^{\tau}
$$

Suppose $u$ is a function which maps $\mathbf{R}$ into $L^{\tau}$. Then if $\Lambda$ is in $\mathbf{R}$, write

$$
u(\Lambda)=u_{1}(\Lambda)+\cdots+u_{r}(\Lambda)
$$

where $u_{i}(\Lambda)$ is the projection of $u(\Lambda)$ onto $L_{\sigma_{i}}^{\tau}$. Let us say that $u$ is in $L^{2}(\mathbf{R} ; \tau)$ if $u_{i}$ is in $L^{2}\left(\mathbf{R}, \beta\left(\sigma_{i}, \Lambda\right) d \Lambda\right) \otimes L_{\sigma_{i}}^{\tau}$ for each $i, 1<i<r$. Then if $u$ is in $L^{2}(\mathbf{R} ; \tau)$, define

$$
\begin{equation*}
\check{u}(x)=\sum_{i=1}^{r} \operatorname{dim} \sigma_{i} \cdot \lim _{R \rightarrow \infty} \int_{-R}^{R} E_{\Lambda}\left(u_{i}(\Lambda): x\right) \beta(\sigma, \Lambda) d \Lambda . \tag{7.4}
\end{equation*}
$$

The limit in this formula is taken in the topology of $L^{2}(G) \otimes V_{\tau}$, where $V_{\tau}$ is the space on which $\tau$ acts. Then $\check{u}$ is a $\tau$-spherical function which is square-integrable.

On the other hand, if $\phi$ is a square-integrable $\tau$-spherical function, define a function $\hat{\phi}$ in $L^{2}(\mathbf{R} ; \tau)$ as follows. If $\psi$ is an arbitrary element in $L^{\tau}$ then

$$
\begin{equation*}
(\hat{\phi}(\Lambda), \psi)_{M}=\lim _{N \rightarrow \infty} \int_{G_{N}}\left(\phi(x), E_{\Lambda}(\psi: x)\right) d x, \quad \Lambda \in \mathbf{R} \tag{7.5}
\end{equation*}
$$

Here $G_{N}$ is any increasing sequence of compact sets whose union is $G$, and the limit is taken in the topology of $L^{2}(\mathbf{R} ; \tau)$.

Using the formulae (5.6) and (7.3), and recalling the definition of $M(\Lambda)$, it is possible to derive the following formula for any function $u$ in $L^{2}(\mathbf{R} ; \tau)$ :

$$
\begin{equation*}
\hat{\tilde{u}}(\Lambda)=\frac{1}{2}[u(\Lambda)+M(-\Lambda) u(-\Lambda)], \quad \Lambda \neq 0 \tag{7.6}
\end{equation*}
$$

We shall use this formula in $\S 13$.
8. Basic estimates for derivatives. We open this chapter by stating a lemma of E. Nelson.

Lemma 12. Let $\pi$ be an irreducible unitary representation of $G$ on a Hilbert space $\mathscr{H}$. Fix $g$ in $\mathscr{B}$. Then there are an integer $m_{g}>0$ and a constant $C$, both independent of $\pi$, such that for every $\Phi$ in $\mathscr{H}^{\infty}$

$$
\|\pi(g) \Phi\|<C\left\|\pi\left(\left(\omega_{\mathfrak{g}}+\omega_{\mathfrak{p}}-1\right)^{m}\right) \Phi\right\|
$$

Here $\omega_{\mathfrak{g}}$ and $\omega_{\mathfrak{k}}$ are the elements in $\mathscr{Z}$ and $\mathscr{Z}_{K}$ respectively, defined in $\S 6$.
For a proof of this lemma see [7, Lemma 6.3]. It turns out that $m_{g}$ can be taken to equal the order of $g$.

Suppose $\rho$ is the conjugation of $\mathfrak{g}_{c}$ with respect to the compact real form $\mathfrak{u}=\mathfrak{k}+i \mathfrak{p}$. We define a Hermitian scalar product on $\mathfrak{g}_{\mathbf{c}}$ by

$$
(X, Y)=-B(X, \rho Y), \quad X, Y \subset g_{c}
$$

We obtain a Hermitian inner product on $\mathfrak{a}_{c}$ by restriction. This permits us to define an inner product on the dual space of $\mathfrak{a}_{\mathrm{c}}$. If $\lambda$ is a complex valued linear functional on $a_{c}$, we write $|\lambda|$ as the norm of $\lambda$ with respect to this inner product.

Suppose $\tau_{1}$ and $\tau_{2}$ are in $\mathscr{E}_{K}$. Write $\tau$ as the double representation $\left(\tau_{1}, \tau_{2}\right)$ of $K$. Then

$$
|\tau|=\left|\tau_{1}\right|+\left|\tau_{2}\right|, \quad \operatorname{dim} \tau=\operatorname{dim} \tau_{1} \cdot \operatorname{dim} \tau_{2}
$$

Lemma 13. Let $\pi$ be an irreducible unitary representation of $G$ on the Hilbert space $\mathscr{H}$, with infinitesimal character $\chi_{\lambda}$, for $\lambda$ a linear functional on $\mathfrak{a}_{\mathrm{c}}$. Suppose $g_{1}$ and $g_{2}$ are in $\mathscr{B}$. Then there exist polynomials $p$ and $q$, independent of $\lambda$, such that the following (somewhat complicated) property is satisfied:

Whenever $\Phi_{1}$ and $\Phi_{2}$ are unit vectors in $\mathscr{H}$ that transform under $\left.\pi\right|_{K}$ according to the representations $\tau_{1}$ and $\tau_{2}$ in $\mathscr{E}_{K}$, set $f(x)=\left(\Phi_{1}, \pi(x) \Phi_{2}\right)$. Then there are two sets, $\left\{\Psi_{1 \alpha}: 1 \gtrless \alpha \gtrless t_{1}\right\},\left\{\Psi_{2 \beta}: 1 \gtrless \beta<t_{2}\right\}$, of orthogonal vectors in $\mathscr{H}$, and two sets $\left\{\tau_{1 \alpha}: 1 \gtrless \alpha<t_{1}\right\},\left\{\tau_{2 \beta}: 1 \gtrless \beta \gtrless t_{2}\right\}$ of representations in $\mathscr{E}_{K}$ such that
(i) $\left(\left\|\Psi_{1 \alpha}\right\|+\left\|\Psi_{2 \beta}\right\|\right)<p(|\lambda|) \cdot q(|\tau|)$.
(ii) $\Psi_{1 \alpha}$ and $\Psi_{2 \beta}$ transform under $\left.\pi\right|_{K}$ according to the representations $\tau_{1 \alpha}$ and $\tau_{2 \beta}$ respectively.
(iii) $\left(\left|\tau_{1 \alpha}\right|+\left|\tau_{2 \beta}\right|\right)<q(|\tau|)$.
(iv) $t_{1}+t_{2}<q(|\tau|)$.
(v) $f\left(g_{1} ; x ; g_{2}\right)=\sum_{\alpha=1}^{t_{1}} \sum_{\beta=1}^{t_{2}}\left(\Psi_{1 \alpha}, \pi(x) \Psi_{2 \beta}\right)$.

Proof. By Lemma 6, $f(x)$ is differentiable. If $\Psi_{1}=\pi\left(g_{1}^{+}\right) \Phi_{1}$ and $\Psi_{2}=$ $\left(\bar{g}_{2}\right) \Phi_{2}$, we have the formula

$$
f\left(g_{1} ; x ; g_{2}\right)=\left(\Psi_{1}, \pi(x) \Psi_{2}\right)
$$

By the last lemma, we can choose $C$ and $m$ such that

$$
\left\|\pi\left(g_{1}^{+}\right) \Phi_{1}\right\|<C\left\|\pi\left(\left(\omega_{\mathfrak{g}}+\omega_{\mathfrak{k}}-1\right)^{m}\right) \Phi_{1}\right\| .
$$

However, $\Phi_{1}$ transforms under $\left.\pi\right|_{K}$ according to $\tau_{1}$. Therefore, by (6.8) and (6.11) we see that

$$
\left\|\pi\left(\left(\omega_{\mathfrak{g}}+\omega_{\mathfrak{k}}-1\right)^{m}\right) \Phi_{1}\right\|=\left\|\left(B(\lambda, \lambda)-B(\rho, \rho)+\left|\tau_{1}\right|^{2}-B(\rho, \rho)-1\right)^{m} \Psi_{1}\right\|
$$

This last expression is bounded by polynomials in $|\lambda|$ and $\left|\tau_{1}\right|$. Similarly, $\left\|\Psi_{2}\right\|$ is bounded by a product of polynomials in $|\lambda|$ and $|\tau|$.

Suppose that $g_{1}$ has order $m$. Let $\mathscr{B}_{m}$ be the set of elements in $\mathscr{B}$ of order less than or equal to $m . \mathscr{B}_{m}$ is a finite-dimensional vector space, and there is a natural representation $\nu$ of $K$ on $\mathscr{B}_{m}$ given by the adjoint map.

Define a linear map from the vector space $\mathscr{B}_{m} \otimes \mathscr{H}^{\infty}$ into $\mathscr{H}^{\infty}$ by

$$
g \otimes \Phi \rightarrow \pi(g) \Phi, \quad g \in \mathscr{B}_{m}, \Phi \in \mathscr{H}{ }^{\infty}
$$

$\mathscr{H}^{\infty}$ is a $K$-module under $\left.\pi\right|_{K}$, so $\mathscr{B}_{m} \otimes \mathscr{H}^{\infty}$ is a $K$-module. If $X$ is in $\mathfrak{g}, k$ is in $K$ and $\Phi$ is in $\mathscr{H}^{\infty}$, then

$$
\pi(k) \pi(X) \Phi=\lim _{t \rightarrow \infty} \pi(k) \cdot t^{-1}(\pi(\exp t X)-I) \cdot \Phi
$$

This last expression equals

$$
\pi(\operatorname{Ad} k \cdot X) \cdot \pi(k) \Phi
$$

Therefore, the following diagram is commutative:


Therefore, since $\Phi$ transforms under $\left.\pi\right|_{K}$ according to $\tau_{1}$, the vector $\Psi_{1}=$ $\pi\left(g_{1}^{+}\right) \Phi_{1}$ transforms under $\left.\pi\right|_{K}$ as a vector in the space on which $\tau_{1} \otimes \nu$ acts. Then

$$
\begin{equation*}
\Psi_{1}=\Psi_{11}+\cdots+\Psi_{1 t_{1}} \tag{8.1}
\end{equation*}
$$

where the vectors $\left\{\Psi_{1 \alpha}\right\}$ are orthogonal to each other, and each $\Psi_{1 \alpha}$ transforms under $\left.\pi\right|_{K}$ according to some irreducible representation $\pi_{1 \alpha}$ of $K$ that occurs in the decomposition of $\tau_{1} \otimes \nu$ into irreducible representations of $K$.

Since (8.1) is an orthogonal decomposition,

$$
\left\|\Psi_{1 \alpha}\right\|<\left\|\Psi_{1}\right\|<p(|\lambda|) \cdot q(|\tau|), \quad 1 \gtrless \alpha<t_{1}
$$

Also, we have the formula

$$
t_{1} ₹ \operatorname{dim}\left(\tau_{1} \otimes \nu\right)=\operatorname{dim} \tau_{1} \cdot \operatorname{dim} \nu
$$

This expression is bounded by a polynomial in $\left|\tau_{1}\right|$ by the Weyl dimension formula.

Now if $\Psi_{1 \alpha}$ is not zero, $\tau_{1 \alpha}$ occurs in the decomposition of $\tau_{1} \otimes \nu_{0}$ into irreducible representations, for some irreducible representation $\nu_{0}$ of $K$ occurring in $\nu$. Let $\lambda_{1}, \lambda_{0}$, and $\lambda_{1 \alpha}$ be the highest weights of $\tau_{1}, \nu_{0}$, and $\tau_{1 \alpha}$ respectively. By examining the formula for the multiplicity of $\tau_{1 \alpha}$ in $\tau_{1} \otimes \nu_{0}[5, \mathrm{Pg} .262]$, we see that there is a nonnegative sum, $\mu$, of positive roots, such that

$$
\left(\lambda_{1}+\rho_{K}\right)+\left(\lambda_{0}+\rho_{K}\right)-\left(\lambda_{1 \alpha}+2 \rho_{K}\right)=\mu
$$

Therefore

$$
\left(\lambda_{\mathbf{1}}+\rho_{K}\right)+\left(\lambda_{0}+\rho_{K}\right)=\left(\lambda_{\mathbf{1} \alpha}+\rho_{K}\right)+\left(\mu+\rho_{K}\right)
$$

Now $B\left(\lambda_{1 \alpha}+\rho_{K}, \mu+\rho_{K}\right)>0$ since $\lambda_{1 \alpha}+\rho_{K}$ is in the positive Weyl chamber. It is then easy to show that

$$
\left|\tau_{1 \alpha}\right| ₹\left|\tau_{1}\right|+\left|\nu_{0}\right|
$$

if we recall the definitions of $\left|\tau_{1 \alpha}\right|,\left|\tau_{1}\right|$, and $\left|\nu_{0}\right|$. We have now verified conditions (i)-(iv) for $\left\{\Psi_{1 \alpha}\right\}$ and $\left\{\tau_{1 \alpha}\right\}$.

We define collections $\left\{\Psi_{2 \beta}\right\}$ and $\left\{\tau_{2 \beta}\right\}$ the same way. They satisfy conditions (i)-(iv). It is clear that

$$
f\left(g_{1} ; x ; g_{2}\right)=\sum_{\alpha=1}^{t_{1}} \sum_{\beta=1}^{t_{2}}\left(\Psi_{1 \alpha}, \pi(x) \Psi_{2 \beta}\right)
$$

The proof of Lemma 13 is complete.
Suppose that $\pi$ equals $\pi_{\omega}$, a representation in the equivalence class of some $\omega$ in $\mathscr{E}_{d}$, acting on the Hilbert space $\mathscr{H}_{\omega}$. Suppose $\omega=\omega(\lambda)$ for some $\lambda$ in $L^{\prime}$. The infinitesimal character of $\pi_{\omega}$ is $\chi_{\lambda_{y}}$, where $y$ is the element in $G_{c}$ defined in $\S 6$. Furthermore, since $\operatorname{Ad} y$ maps $(-1)^{1 / 2} \mathfrak{b}$ onto $(-1)^{1 / 2} \mathfrak{a}_{\mathfrak{t}}+\mathfrak{a}_{\mathfrak{p}}, \rho\left(\lambda_{y}\right)$ equals $-\lambda_{y}$. Therefore

$$
\begin{equation*}
\left|\lambda_{y}\right|^{2}=-B\left(\lambda_{y}, \rho\left(\lambda_{y}\right)\right)=B\left(\lambda_{y}, \lambda_{y}\right)=B(\lambda, \lambda)=|\omega|^{2} \tag{8.2}
\end{equation*}
$$

If $\pi=\pi_{\sigma, \Lambda}$, for $\sigma$ in $\mathscr{E}_{M}$, choose a real linear functional $\mu$ on $(-1)^{1 / 2} \mathfrak{a}_{\mathfrak{k}}$ associated with $\sigma$. The infinitesimal character of $\pi_{\sigma, \Lambda}$ is $\chi_{-\mu-i \Lambda \mu_{0}}$, while $\rho(\mu)=$ $-\mu$ and $\rho\left(i \Lambda \mu_{0}\right)=i \Lambda \mu_{0}$. Therefore

$$
\begin{align*}
\left|-\mu-i \Lambda \mu_{0}\right|^{2}=- & B\left(\mu+i \Lambda \mu_{0}, \rho\left(\mu+i \Lambda \mu_{0}\right)\right) \\
& =B(\mu, \mu)-B\left(i \Lambda \mu_{0}, i \Lambda \mu_{0}\right)=|\sigma|^{2}+r^{-2} \Lambda^{2} \tag{8.3}
\end{align*}
$$

9. Preparation for the main estimates. Suppose $\left\{\phi_{\lambda}^{\tau}: \lambda \in \mathscr{E}, \tau \in T_{\lambda}\right\}$ is a collection of infinitely differentiable $\tau$-spherical functions. $\tau$ indexes certain irreducible unitary double representations $\tau=\left(\tau_{1}, \tau_{2}\right)$ of $K$ on the finitedimensional Hilbert spaces $V_{\tau}=V_{\tau_{1}} \otimes V_{\tau_{2}}^{*} \cdot \lambda$ indexes linear functions from $\mathfrak{a}_{\mathrm{c}}$ to C. We have the homomorphisms $\chi_{\lambda}: \mathscr{Z} \rightarrow \mathbf{C}$ defined in $\S 6$. We assume that

$$
z \phi_{\lambda}^{\tau}=\chi_{\lambda}(z) \phi_{\lambda}^{\tau}, \quad z \in \mathcal{Z}
$$

By $|\lambda|$ we shall mean the real number $[-B(\lambda, \rho(\lambda))]^{1 / 2}$ as in the last section. We assume that for any $g_{1}, g_{2}$ in $\mathscr{B}$ there are polynomials $p$ and $q$ such that for each $x$ in $G$

$$
\begin{equation*}
\left|\phi_{\lambda}^{\tau}\left(g_{1} ; x ; g_{2}\right)\right| \gtrless p(|\lambda|) q(|\tau|) \Xi(x) \tag{9.1}
\end{equation*}
$$

In [3(l), §27], Harish-Chandra has defined for each $\phi_{\lambda}^{\tau}$ an infinitely differentiable function $\theta=\theta_{\lambda}^{\tau}$ mapping $M A_{\mathfrak{p}}$ into $V_{\tau} . \theta$ is $\bar{\tau}$-spherical, where $\bar{\tau}$ is the restriction of $\tau$ to $M$. Harish-Chandra shows that $\theta$ vanishes if $\phi_{\lambda}^{\tau}$ is in $L^{2}(G) \otimes V_{\tau}$. We shall make two assumptions on the collection of linear functionals $\{\lambda: \lambda \bullet \mathscr{E}\}$. It then turns out that there exist polynomials $p$ and $q$, and a number $\varepsilon>0$, independent of $(\lambda, \tau)$, such that for each $t>0$,

$$
\begin{equation*}
\left|e^{\rho\left(t H_{0}\right)} \phi_{\lambda}^{\tau}\left(\exp t H_{0}\right)-\theta_{\lambda}^{\tau}\left(\exp t H_{0}\right)\right| \gtrless p(|\lambda|) q(|\tau|) e^{-\varepsilon t} \tag{9.2}
\end{equation*}
$$

We shall review Harish-Chandra's work and prove the estimate (9.2).
Recall that $\mathfrak{M} \mathfrak{X}_{p}$ is the universal enveloping algebra of $\boldsymbol{m}_{c}+a_{p c} ; \operatorname{let} \mathfrak{Z}_{1}=\mathfrak{Z}_{M} \mathfrak{A}_{p}$ be its center. In $\S 6$ we defined isomorphisms

$$
\gamma: \mathfrak{Z} \rightarrow J, \quad \gamma_{1}: \mathfrak{Z}_{1} \rightarrow J_{1}
$$

where $J$ and $J_{1}$ are the elements in $S\left(\mathfrak{a}_{\mathbf{c}}\right)$, the symmetric algebra on $\mathfrak{a}_{\mathbf{c}}$, which are invariant under $W$ and $W_{1}$ respectively. (We can regard $W_{1}$ as the Weyl group of $\mathfrak{m}_{\mathrm{c}}+\mathfrak{a}_{\mathrm{pc}}$ acting on $\mathfrak{a}_{\mathrm{c}}$.) $J$ is contained in $J_{1}$. We have defined the map

$$
\gamma_{0}=\gamma_{1}^{-1} \circ \gamma: \mathfrak{Z} \rightarrow \mathfrak{Z}_{1}
$$

Let $\mathfrak{U}_{\lambda}$ be the annihilator of $\phi_{\lambda}^{\tau}$ in $\mathfrak{Z}$, and let $\mathfrak{U}_{1 \lambda}=\mathfrak{Z}_{1} \cdot \gamma_{0}\left(\mathfrak{U}_{\lambda}\right)$. Let $\mathfrak{Z}_{1}^{*}$ be the quotient algebra $\mathfrak{Z}_{1} / \mathfrak{U}_{1 \lambda}$. We can regard $\mathfrak{Z}_{1}^{*}$ as a complex vector space on which there is a natural $\mathcal{Z}_{1}$ action.

If $\varsigma$ is in $\mathcal{Z}_{1}$, let $\varsigma^{*}$ be the projection of $\varsigma$ onto $\mathcal{Z}_{1}^{*}$. Let $\mathfrak{Z}_{1}^{* *}$ be the vector space dual of $\mathcal{Z}_{1}^{*}$. Let $\mathscr{V}_{\tau}=V_{\tau} \otimes \mathcal{Z}_{1}^{* *}$. Make $\mathscr{V}_{\tau}$ a double $K$-module by letting $K$ act trivially on $\mathfrak{Z}_{1}^{* *}$. Make it a $\mathfrak{Z}_{1}$-module by defining

$$
\Gamma(z)\left(v \otimes \varsigma^{* *}\right)=v \otimes z \varsigma^{* *}, \quad z \in \mathfrak{Z}_{1}, v \in V_{\tau}, \varsigma^{* *} \in \mathfrak{Z}_{1}^{* *}
$$

(Since $\mathfrak{Z}_{1}^{*}$ is a $\mathfrak{Z}_{1}$-module, there is a natural action of $\mathfrak{Z}_{1}$ on $\mathfrak{Z}_{1}^{* *}$ obtained by taking transposes.)

Let us examine the algebras $J$ and $J_{1}$ more closely, in order to obtain a basis of $\mathfrak{Z}_{1}^{*}$. Such results appear in $[\mathbf{3}(\mathbf{g}), \S 3]$. We identify $S=S\left(\mathfrak{a}_{\mathrm{c}}\right)$ with the algebra of polynomial functions on $\mathfrak{a}_{\mathrm{c}}^{*}$, the dual space of $\mathfrak{a}_{\mathrm{c}}$. Let $C(S), C\left(J_{1}\right)$, and $C(J)$ be the quotient fields of $S, J_{1}$, and $J$ respectively.

Lemma 14. If $\left[W: W_{1}\right]=r$, then there are homogeneous elements $v_{1}=1$, $v_{2}, \ldots, v_{r}$ in $J_{1}$ such that $J_{1}=\sum_{1 \gtrless i \gtrless r} J v_{i}$. Moreover, the elements $v_{1}, \ldots, v_{r}$ are linearly independent over $C(J)$.

For a proof of this lemma, see $\left[\mathbf{3}(\mathbf{g})\right.$, Lemma 8]. $C(S) / C\left(J_{1}\right)$ and $C(S) / C(J)$ turn out to be normal extensions with Galois groups $W_{1}$ and $W$ respectively.

Now suppose that $\lambda$ is in $a_{\mathbf{c}}^{*}$. Let $S_{\lambda}$ be the ideal of polynomial functions in $S$ that vanish at $\lambda$. Let $J_{\lambda}=J \cap S_{\lambda}$ and let $J_{1 \lambda}=J_{1} \cap S_{\lambda} . J=\mathbf{C} \oplus J_{\lambda}$ is a vector space decomposition of $J$ and the projection from $J$ onto $J / J_{\lambda}=\mathbf{C}$ is given by

$$
w \rightarrow w(\lambda), \quad w \in J
$$

$J_{1} J_{\lambda}$ is an ideal in $J_{1}$ and it is clear that $\gamma_{1}$ defines an isomorphism from $\mathcal{Z}_{1}^{*}$ onto $J_{1} / J_{1} J_{\lambda}$. We shall obtain a basis of $J_{1} / J_{1} J_{\lambda}$ over the complex numbers.

We have the formula

$$
J_{1}=\sum_{1 \gtrless i \gtrless r} J v_{i}=\sum_{i} \mathbf{C} v_{i}+\sum_{i} J_{\lambda} v_{i}
$$

But $\sum_{i} J_{\lambda} v_{i}$ is in $J_{1} J_{\lambda}$ so $\left\{v_{i}: 1 \gtrless i<r\right\}$ spans $J_{1} / J_{1} J_{\lambda}$. On the other hand, suppose $\left\{c_{i}\right\}$ is a set of complex numbers such that the vector $\sum_{i} c_{i} v_{i}$ is in $J_{1} J_{\lambda}$. Now

$$
J_{1} J_{\lambda}=\left(\sum_{i} J v_{i}\right) J_{\lambda}=\sum_{i} J_{\lambda} v_{i}
$$

However, $\left\{v_{i}: 1 \gtrless i \gtrless r\right\}$ are linearly independent over $C(J)$ so each $c_{i}$ is in $J_{\lambda}$. This implies that each $c_{i}$ equals zero. Therefore, $\left\{v_{i}: 1 \gtrless i \gtrless r\right\}$ is a basis for the vector space $J_{1} / J_{1} J_{\lambda}$. Let us regard $J_{1} / J_{1} J_{\lambda}$ as a Hilbert space with orthonormal basis $\left\{v_{1}, \ldots, v_{r}\right\}$.

Define elements $\eta_{1}=1, \eta_{2}, \ldots, \eta_{r}$ in $\mathfrak{Z}_{1}$ by

$$
\gamma_{1}\left(\eta_{i}\right)=v_{i}, \quad 1 \gtrless i \gtrless r
$$

Then $\left\{\eta_{1}^{*}, \ldots, \eta_{r}^{*}\right\}$ is a basis for $\mathfrak{Z}_{1}^{*}$. Let $\left\{\eta_{1}^{* *}, \ldots, \eta_{r}^{* *}\right\}$ be the dual basis $\mathfrak{Z}_{1}^{* *}$. If we make $\mathcal{Z}_{1}^{* *}$ into a Hilbert space with orthonormal basis $\left\{\eta_{1}^{* *}, \ldots, \eta_{r}^{* *}\right\}$, we can regard $\mathscr{V}_{\tau}$ as a Hilbert space.

Define a function on $M A_{p}$ by

$$
d(m a)=e^{\rho(\log a)}, \quad a \in A_{\mathfrak{p}}, m \in M
$$

We define an automorphism $\varsigma \rightarrow \varsigma^{\prime}$ of $\mathfrak{Z}_{1}$ by

$$
\varsigma^{\prime}=d^{-1} \varsigma \circ d, \quad \varsigma \in \mathfrak{Z}_{1}
$$

(We are regarding $\mathfrak{Z}_{1}$ as an algebra of differential operators on $M^{0} A_{\mathfrak{p}}$.)
If $f$ is in $C^{\infty}\left(M A_{\mathfrak{p}}\right), v$ is in $\mathfrak{M} \mathscr{A}_{\mathfrak{p}}, m$ is in $M$, and $a$ is in $A_{\mathfrak{p}}$, we can define $f(m a ; v)$ even though $M A_{\mathfrak{p}}$ may not be connected. If $X$ is in $\mathfrak{m}+\mathfrak{a}_{\mathfrak{p}}$, we just write

$$
f(m a ; X)=\left.\frac{d}{d t} f(m a \exp t X)\right|_{t=0}
$$

We extend the definition to all $v$ in the universal enveloping algebra $\mathfrak{M P}_{\mathfrak{p}}$ in the usual way. Since $M A_{\mathfrak{p}}=Z(A) M^{0} A_{\mathfrak{p}}$, and since $Z(A)$ and $M^{0} A_{\mathfrak{p}}$ commute, any $\varsigma$ in $\mathcal{Z}_{1}$ can be regarded as a left and right invariant differential operator on $C^{\infty}\left(M A_{\mathfrak{p}}\right)$.

For $m$ in $M$ and $a$ in $A_{\mathfrak{p}}$, define

$$
\phi_{i}(m a)=d(m a) \phi_{\lambda}^{\tau}\left(m a ; \eta_{i}^{\prime}\right)
$$

Define

$$
\Phi(m a)=\sum_{1 \gtrless i<r} \phi_{i}(m a) \otimes \eta_{i}^{* *}
$$

If $\varsigma$ is in $\mathfrak{Z}_{1}$, choose complex numbers $\left\{c_{i j}: 1<i<r\right\}$ such that the differential operators

$$
\varsigma \eta_{i}-\sum_{j} c_{i j} \eta_{j}=u_{i}(\zeta), \quad 1 \gtrless i<r
$$

are in $\mathfrak{U}_{1 \lambda}$. Then define

$$
\Psi_{\zeta}(m a)=\sum_{1 \gtrless i \gtrless r} d(m a) \phi_{\lambda}^{\tau}\left(m a ; u_{i}(\varsigma)^{\prime}\right) \otimes \eta_{i}^{* *}
$$

$\Phi$ and $\Psi_{\zeta}$ are functions from $M A_{\mathfrak{p}}$ to the vector space $\mathscr{V}_{\tau}$. Also, $\Phi$ and $\Psi_{\zeta}$ are both $\bar{\tau}$-spherical functions on $M A_{\mathfrak{p}}$, since elements in $\mathfrak{Z}_{1}$ act on $C^{\infty}\left(M A_{\mathfrak{p}}\right)$ as left and right invariant differential operators.

Lemma 15. Let $\varsigma$ be in $\mathfrak{Z}_{1}$. Then for $m$ in $M$ and $a$ in $A_{\mathfrak{p}}$,

$$
\Phi(m a ; \zeta)=\Gamma(\varsigma) \Phi(m a)+\Psi_{\varsigma}(m a)
$$

Proof. We have the equation

$$
\Phi(m a ; \varsigma)=\sum_{i} d(m a) \phi_{\lambda}^{\tau}\left(m a ; \varsigma^{\prime} \eta_{i}^{\prime}\right) \otimes \eta_{i}^{* *}
$$

Also,

$$
\varsigma^{\prime} \eta_{i}^{\prime}=\left(\varsigma \eta_{i}\right)^{\prime}=\sum_{j} c_{i j} \eta_{j}^{\prime}+u_{i}(\varsigma)^{\prime}
$$

Therefore, $\Phi(m a ; \varsigma)$ is equal to

$$
\sum_{i} \sum_{j} d(m a) \phi_{\lambda}^{\tau}\left(m a ; \eta_{j}^{\prime}\right) \otimes \eta_{i}^{* *}+\sum_{i} d(m a) \phi_{\lambda}^{\tau}\left(m a ; u_{i}(\zeta)^{\prime}\right)
$$

which in turn equals

$$
\sum_{i j} c_{i j} \phi_{j}(m a) \otimes \eta_{i}^{* *}+\Psi_{s}(m a)
$$

Since $u_{i}(\varsigma)^{*}=0$, we have the formula

$$
\left(\varsigma \eta_{i}\right)^{*}=\sum_{j} c_{i j} \eta_{j}^{*}
$$

But $\left\{\eta_{i}^{* *}\right\}$ is a dual basis of $\left\{\eta_{i}^{*}\right\}$, so the matrix of the linear transformation $\varsigma$ acting on $\mathfrak{Z}_{1}^{* *}$ is the transpose matrix of its action on $\mathfrak{Z}_{1}^{*}$, with respect to these bases. Therefore,

$$
\Gamma\left(H_{0}\right) \Phi(m a)=\sum_{i j} c_{j i} \phi_{i}(m a) \otimes \eta_{j}^{* *}=\sum_{i j} c_{i j} \phi_{j}(m a) \otimes \eta_{i}^{* *}
$$

This proves the lemma.
Clearly $H_{0}$ is in $\mathfrak{Z}_{1}$. Write $\Psi_{H_{0}}$ as $\Psi$.
COROLLARY. For any $m$ in $M$, a in $A_{\mathfrak{p}}$, and $T$ in $\mathbf{R}$, we have the integral equation

$$
\Phi\left(m a \cdot \exp T H_{0}\right)=e^{T \Gamma\left(H_{0}\right)} \Phi(m a)+\int_{0}^{T} e^{(T-t) \Gamma\left(H_{0}\right)} \Psi\left(m a \cdot \exp t H_{0}\right) d t
$$

Proof. If $t$ is in $R$, we see that

$$
\begin{aligned}
& \frac{d}{d t}\left(e^{-t \Gamma\left(H_{0}\right)} \Phi\left(m a \cdot \exp t H_{0}\right)\right) \\
&=e^{-t \Gamma\left(H_{0}\right)} \Phi\left(m a \cdot \exp t H_{0} ; H_{0}\right)-e^{-t \Gamma\left(H_{0}\right)} \Gamma\left(H_{0}\right) \Phi\left(m a \cdot \exp t H_{0}\right) \\
&=e^{-t \Gamma\left(H_{0}\right)} \Psi\left(m a \cdot \exp t H_{0}\right)
\end{aligned}
$$

Integrate the above equation by $\int_{0}^{T} d t$. We obtain the formula

$$
e^{-T \Gamma\left(H_{0}\right)} \Phi\left(m a \exp T H_{0}\right)-\Phi(m a)=\int_{0}^{T} e^{-t \Gamma\left(H_{0}\right)} \Psi\left(m a \cdot \exp t H_{0}\right) d t
$$

Multiply by $e^{T \Gamma\left(H_{0}\right)}$. This proves the corollary.
Now, fix $\varsigma$ in $\mathfrak{Z}_{1}$. We would like to obtain an estimate for $\Psi_{\varsigma}(m a)$. To do this, we must first express the differential operators $u_{i}(\varsigma)$ in another form.

By Lemma 14, there exist elements $w_{i j}$ in $J$ such that for $1<i<r$,

$$
\gamma_{1}(\varsigma) v_{i}=\sum_{j} w_{i j} v_{j}
$$

Now, the coordinates of $\gamma_{1}(\varsigma) v_{i}$ relative to the basis $\left\{v_{j}\right\}$ of $J_{1} / J_{1} J_{\lambda}$ are clearly $\left\{w_{i j}(\lambda): 1 \gtrless j \gtrless r\right\}$. Then the element

$$
\gamma_{1}\left(u_{i}(\zeta)\right)=\gamma_{1}(\zeta) v_{i}-\sum_{j} w_{i j}(\lambda) v_{j}
$$

is in $J_{1} J_{\lambda} . \gamma_{1}\left(u_{i}(\varsigma)\right)$ is equal to

$$
\sum_{j}\left(w_{i j}-w_{i j}(\lambda)\right) v_{j}
$$

For each $i$ and $j,\left(w_{i j}-w_{i j}(\lambda)\right)$ is in $J_{\lambda}$. Let $z_{i j}^{\lambda}$ be the differential operator

$$
\gamma^{-1}\left(w_{i j}-w_{i j}(\lambda)\right)
$$

Then

$$
u_{i}(\varsigma)=\sum_{j} \gamma_{0}\left(z_{i j}^{\lambda}\right) \gamma_{1}^{-1}\left(v_{j}\right)=\sum_{j} \gamma_{0}\left(z_{i j}^{\lambda}\right) \eta_{j}
$$

Let $u_{i j}^{\lambda}$ be the differential operator

$$
z_{i j}^{\lambda}-\gamma_{0}\left(z_{i j}^{\lambda}\right)^{\prime}
$$

$u_{i j}$ is an element of $\mathfrak{B}$, and it is independent of $\lambda$. Recall that $\theta$ is the Cartan involution of $g_{c}$. In the appendix of [3(1)] it is shown that there exist elements $N_{i j}$ in $n_{c}$ and $g_{i j}$ in $\mathscr{B}$, both independent of $\lambda$, such that

$$
u_{i j}=\theta\left(N_{i j}\right) g_{i j}, \quad 1 \gtrless i, j \gtrless r
$$

Lemma 16. For fixed $\varsigma$ in $\mathfrak{Z}_{1}$ and $h$ in $\mathfrak{A}_{p}$, there are polynomials $p$ and $q$ and an integer $d$, independent of $(\lambda, \tau)$, such that for every $t>0$

$$
\left|\Psi_{\varsigma}\left(\exp t H_{0} ; h\right)\right|<p(|\lambda|) q(|\tau|) e^{-t}(1+t)^{d} .
$$

Proof. We have the equations

$$
\begin{aligned}
\Psi_{\zeta}\left(\exp t H_{0} ; h\right) & =\sum_{i} e^{t \rho\left(H_{0}\right)} \phi_{\lambda}^{\top}\left(\exp t H_{0} ; h^{\prime} u_{i}(\zeta)^{\prime}\right) \otimes \eta_{i}^{* *} \\
& =\sum_{i j} e^{t \rho\left(H_{0}\right)} \phi_{\lambda}^{\tau}\left(\exp t H_{0} ; h^{\prime} \eta_{j}^{\prime} \gamma_{0}\left(z_{i j}^{\lambda}\right)^{\prime}\right) \otimes \eta_{i}^{* *}
\end{aligned}
$$

Since $z_{i j}^{\lambda}$ is in $\mathfrak{U}_{\lambda}$, the annihilator of $\phi_{\lambda}^{\tau}$, this last expression is equal to

$$
\begin{aligned}
& \sum_{i j} e^{t \rho\left(H_{0}\right)} \phi_{\lambda}^{\tau}\left(\exp t H_{0} ; h^{\prime} \eta_{j}^{\prime} u_{i j}\right) \otimes \eta_{i}^{* *} \\
&=\sum_{i j} e^{t \rho\left(H_{0}\right)} \phi_{\lambda}^{\tau}\left(\exp t H_{0} ; u_{i j} h^{\prime} \eta_{j}^{\prime}\right) \otimes \eta_{i}^{* *} \\
&=\sum_{i j} e^{t \rho\left(H_{0}\right)} \phi_{\lambda}^{\tau}\left(\exp t H_{0} ; \theta\left(N_{i j}\right) g_{i j} \eta_{j}^{\prime} h^{\prime}\right) \otimes \eta_{i}^{* *}
\end{aligned}
$$

If $N$ is any vector in $\mathfrak{n}_{\mathrm{c}}, N$ is equal to $\sum_{\alpha \in P_{+}} X_{\alpha}$, where $\left\{X_{\alpha}\right\}$ are vectors in $\mathfrak{g}_{\mathrm{c}}$ such that $\left[H_{0}, X_{\alpha}\right]=\alpha\left(H_{0}\right) X_{\alpha}$. But if $\alpha$ is in $P_{+}$,

$$
\phi_{\lambda}^{\tau}\left(\exp t H_{0} ; \theta\left(X_{\alpha}\right) g_{i j} \eta_{j}^{\prime} h^{\prime}\right)=e^{-t \alpha\left(H_{0}\right)} \phi_{\lambda}^{\tau}\left(\theta\left(X_{\alpha}\right) ; \exp t H_{0} ; g_{i j} \eta_{j}^{\prime} h^{\prime}\right)
$$

But $e^{-t \alpha\left(H_{0}\right)}$ equals either $e^{-t}$ or $e^{-2 t}$. The lemma then follows from (9.1) and (4.1).

Recall that $\tilde{\omega}=\prod_{\alpha \in P} H_{\alpha}$. In particular, $\tilde{\omega}$ is in $S$. For the rest of this section we make the following assumption.

Assumption 1. For each $\lambda$ in $\mathscr{E}, \tilde{\omega}(\lambda)$ is not equal to zero.
We would like to find the eigenvalues of the linear transformation $H_{0}$ acting on the vector space $J_{1} / J_{1} J_{\lambda}$. We also want to find the norms of the projections of $J_{1} / J_{1} J_{\lambda}$ onto these eigenspaces. (These are norms as operators on the Hilbert space $J_{1} / J_{1} J_{\lambda}$; these projections are not necessarily selfadjoint.)

The field $C\left(J_{1}\right)$ is an extension of degree $r$ of the field $C(J)$. Therefore the trace, $\operatorname{tr}_{C\left(J_{1}\right) / C(J)}$, is a function from $C\left(J_{1}\right)$ into $C(J)$. Define an element $v^{i}$ in $C\left(J_{1}\right)$ by

$$
\operatorname{tr}_{C\left(J_{1}\right) / C(J)}\left(v^{i} v_{j}\right)=\delta_{j}^{i}, \quad 1<i, j<r
$$

( $\delta_{j}^{i}=0$ if $i \neq j, \delta_{i}^{i}=1$.) Recall that $\tilde{\omega}^{\mathfrak{m}}=\prod_{\alpha \in P_{M}} H_{\alpha}$. Define

$$
D=\tilde{\omega} / \tilde{\omega}^{\mathfrak{m}}, \quad \tau^{i}=D v^{i}
$$

Then $D$ and $\tau^{i}$ are both in $C(S)$. In [3(g), Lemma 12] Harish-Chandra shows that $\tau^{i}$ is actually in $S$.

Now let $\left\{s_{1}=1, s_{2}, \ldots, s_{r}\right\}$ be a set of representatives of right cosets of $W_{1}$ in $W$. If $v$ is in $C\left(J_{1}\right)$, then

$$
\operatorname{tr}_{C\left(J_{1}\right) / C(J)}(v)=\sum_{1 \gtrless i<r} s_{i}^{-1}(v)
$$

However,

$$
\sum_{1 \geqslant k \gtrless r} \tau^{i}\left(s_{k} \lambda\right) \cdot v_{j}\left(s_{k} \lambda\right) / D\left(s_{k} \lambda\right)=\delta_{j}^{i}
$$

Define $A_{i k}=\tau^{i}\left(s_{k} \lambda\right)$ and $B_{k j}=v_{j}\left(s_{k} \lambda\right) / D\left(s_{k} \lambda\right)$. Then $A=\left(A_{i k}\right)$ and $B=$ $\left(B_{k j}\right)$ are $r \times r$ matrices, and $A B=I$. Define

$$
\begin{equation*}
f_{s_{i} \lambda}=\sum_{k} \tau^{k}\left(s_{i} \lambda\right) v_{k} \tag{9.3}
\end{equation*}
$$

Then $\left\{f_{s_{i} \lambda}: 1 ₹ i ₹ r\right\}$ is a base for $J_{1} / J_{1} J_{\lambda}$. Also

$$
\begin{equation*}
v_{j}=\sum_{k}\left(v_{j}\left(s_{k} \lambda\right) / D\left(s_{k} \lambda\right)\right) f_{s_{k} \lambda} \tag{9.4}
\end{equation*}
$$

Lemma 17. If $p$ is in $J_{1}$

$$
p f_{s_{i} \lambda} \equiv p\left(s_{i} \lambda\right) f_{s_{i} \lambda} \quad\left(\bmod J_{1} J_{\lambda}\right)
$$

The proof of this lemma is in $[\mathbf{3}(\mathbf{g})$, Lemma 15]. $\square$
In particular, the operator $p$ on $J_{1} / J_{1} J_{\lambda}$ is semisimple. Since $H_{0}$ is in $J_{1}$, the lemma tells us that the set of eigenvalues of $H_{0}$ is

$$
\begin{equation*}
\left\{\lambda\left(H_{0}\right), \lambda\left(s_{2}^{-1} H_{0}\right), \ldots, \lambda\left(s_{r}^{-1} H_{0}\right)\right\} \tag{9.5}
\end{equation*}
$$

We can lift $J_{1}$ to $\mathfrak{Z}_{1}$ by $\gamma_{1}^{-1}$. Then if $\zeta$ is in $\mathfrak{Z}_{1}$, we have an analogous statement to Lemma 17 for the eigenvalues of the operator $\Gamma(\varsigma)$ on $\mathscr{V}_{\tau}$. In particular, the eigenvalues of $\Gamma\left(H_{0}\right)$ are also given by (9.5).

Let $\left\{E_{s_{1}}^{\prime}, \ldots, E_{s_{r}}^{\prime}\right\}$ be the projections in $J_{1} / J_{1} J_{\lambda}$ relative to the direct decomposition

$$
J_{1} / J_{1} J_{\lambda}=\mathbf{C} f_{s_{1} \lambda} \oplus \cdots \oplus \mathbf{C} f_{s_{r} \lambda}
$$

Define

$$
E_{s_{l}}=(\mathrm{Id}) \otimes\left(\gamma_{1}^{-1} E_{s_{l}}^{\prime} \gamma_{1}\right)^{*}, \quad 1<l \gtrless r
$$

where Id stands for the identity operator on $V_{\tau}$, and the star denotes the vector space transpose operator. Then the operators $\left\{E_{s_{l}}\right\}$ are the projections of $\mathscr{V}_{\tau}$ onto the eigenspaces of $\Gamma\left(\mathcal{Z}_{1}\right)$.

If $u$ is a vector in $\mathscr{V}_{\tau}$, there are elements $u_{1}, \ldots, u_{r}$ in $V_{\tau}$ such that

$$
u=\left(u_{1} \otimes \eta_{1}^{* *}\right)+\cdots+\left(u_{r} \otimes \eta_{r}^{* *}\right)
$$

Denote $u_{i}$ by $t_{i}(u)$.
Lemma 18. There is a fixed set of elements $\left\{p_{i j}^{l}: 1 \gtrless i, j, l \gtrless r\right\}$ in $S$, independent of $\tau$ and $\lambda$, such that for any $u$ in $\mathscr{V}_{\tau}$

$$
t_{i}\left(E_{s_{l}} u\right)=\sum_{j=1}^{r}\left(p_{i j}^{l}(\lambda) / \tilde{\omega}(\lambda)\right) t_{j}(u)
$$

Proof. (9.4) implies the formula

$$
E_{s_{l}}^{\prime} v_{i}=E_{s_{l}}^{\prime}\left(\sum_{j=1}^{\tau}\left(v_{i}\left(s_{j} \lambda\right) / D\left(s_{j} \lambda\right)\right) f_{s_{j} \lambda}\right)=\left(v_{i}\left(s_{l} \lambda\right) / D\left(s_{l} \lambda\right)\right) \cdot f_{s_{l} \lambda}
$$

By (9.3) this last expression is equal to

$$
\left(v_{i}\left(s_{l} \lambda\right) / D\left(s_{l} \lambda\right)\right) \cdot \sum_{j} \tau^{j}\left(s_{l} \lambda\right) v_{j}
$$

Now $D\left(s_{l} \lambda\right)$ equals $\left(\tilde{\omega}(\lambda) / \tilde{\omega}^{m}(\lambda)\right) \varepsilon$, where $\varepsilon$ equals either 1 or -1 . This is enough to prove the lemma.

The set of eigenvalues of either the linear transformation $H_{0}$ on $J_{1} / J_{1} J_{\lambda}$ or $\Gamma\left(H_{0}\right)$ on $\mathscr{V}_{\tau}$ is given by (9.5). Let $Q^{+}, Q^{-}, Q^{0}$ be the subsets of these eigenvalues with real parts greater than, less than, or equal to zero, respectively. Let $E^{\prime+}$, $E^{\prime-}, E^{\prime 0}$, and $E^{+}, E^{-}, E^{0}$ be the corresponding projections in $J_{1} / J_{1} J_{\lambda}$ and $\mathscr{V}_{\tau}$ respectively.

Let $\varepsilon^{\prime}$ be the minimum of the absolute values of the real parts of the numbers in the set $Q^{+} \cup Q^{-}$. (Set $\varepsilon^{\prime}=1$ if $Q^{+} \cup Q^{-}$is empty.)

For the remainder of this chapter we make the following assumption.
Assumption 2. The real parts of all eigenvalues (9.5), as $\lambda$ ranges in $\mathscr{E}$, generate a lattice in $\mathbf{R}$. In particular, $\varepsilon^{\prime}$ is bounded away from zero independently of $\lambda$ in $\mathscr{E}$.

LEMMA 19. The norm of $H_{0}$ as an operator on $J_{1} / J_{1} J_{\lambda}$ is bounded by a polynomial in $|\lambda|$.

Proof. $\left\{v_{1}, \ldots, v_{r}\right\}$ is an orthonormal base for $J_{1} / J_{1} J_{\lambda}$. By Lemma 14 there exist elements $w_{i j}$ in $J$, independent of $\lambda$, such that

$$
H_{0} v_{i}=w_{i 1} v_{1}+\cdots+w_{i r} v_{r}, \quad 1<i \gtrless r
$$

This means that

$$
H_{0} v_{i} \equiv w_{i 1}(\lambda) v_{1}+\cdots+w_{i r}(\lambda) v_{r} \quad\left(\bmod J_{1} J_{\lambda}\right)
$$

Now $w_{i j}(\lambda)$ is a polynomial function in $\lambda$, so it is bounded by a polynomial in $|\lambda|$. Since $\left\{v_{i}\right\}$ is an orthonormal basis for $J_{1} / J_{1} J_{\lambda}$, the lemma follows.

Lemma 20. Choose a complex number $\xi$ in the resolvent set of the operator $H_{0}$. Let d be the distance from $\xi$ to the spectrum of $H_{0}$. Then there are polynomials $p_{1}$ and $p_{2}$, independent of $\xi$ and $\lambda$, such that the norm of the operator $\left(\xi-H_{0}\right)^{-1}$ is bounded by $d^{-r} p_{1}(|\xi|) p_{2}(|\lambda|)$.

Proof. By Lemma 19 the matrix $W(\lambda)=\left\{w_{i j}(\lambda)\right\}$ of the linear transformation $H_{0}$ with respect to the orthonormal basis $\left\{v_{1}, \ldots, v_{r}\right\}$ of $J_{1} / J_{1} J_{\lambda}$ has entries which are bounded by a polynomial in $|\lambda|$. Let $u_{i}=\left(\xi-H_{0}\right) v_{i}$. Then for 1 ₹ $i$ ₹ $r$

$$
\begin{equation*}
u_{i}=\xi v_{i}-\sum_{j} w_{i j}(\lambda) v_{j} \tag{9.6}
\end{equation*}
$$

We can solve the equations (9.6) to obtain the formula

$$
v_{i}=\operatorname{det}(\xi I-W(\lambda))^{-1} \cdot \sum_{j} p_{i j}(\xi, \lambda) u_{j}
$$

where $p_{i j}(\xi, \lambda)$ are polynomial functions of $\xi$ and $\lambda$. In particular, each $p_{i j}(\xi, \lambda)$ is bounded by a polynomial in $|\xi|$ and $|\lambda|$. Now

$$
\operatorname{det}(\xi I-W(\lambda))=\prod_{j=1}^{r}\left(\xi-\lambda\left(s_{j}^{-1} H_{0}\right)\right)
$$

Therefore

$$
|\operatorname{det}(\xi I-W(\lambda))|^{-1}<d^{-r}
$$

Since $\left(\xi-H_{0}\right)^{-1} u_{j}=v_{j}$, our lemma is proved.
Let $\varepsilon$ be the minimum of $\varepsilon^{\prime} / 3$ and $1 / 3$.
Lemma 21. There exists a polynomial $p$ such that

$$
\begin{gathered}
\left|e^{-t \Gamma\left(H_{0}\right)} E^{+}\right|+\left|e^{t \Gamma\left(H_{0}\right)} E^{-}\right|<p(|\lambda|) e^{-2 \varepsilon t}, \quad t>0 \\
\left|e^{t \Gamma\left(H_{0}\right)} E^{0}\right|<p(|\lambda|) e^{\varepsilon t}, \quad t \in \mathbf{R} .
\end{gathered}
$$

Proof. It is clearly enough to prove the same statements for the linear transformations $e^{-t H_{0}} E^{\prime+}, e^{t H_{0}} E^{\prime-}$, and $e^{t H_{0}} E^{\prime 0}$ on the Hilbert space $J_{1} / J_{1} J_{\lambda}$.

Let $\Gamma^{+}, \Gamma^{-}$, and $\Gamma^{0}$ be closed curves in the complex plane that wind around the corresponding sets of eigenvalues $Q^{+}, Q^{-}$, and $Q^{0}$ in a positive sense, but which contain no other eigenvalues. By looking at Assumption 2 and the eigenvalues (9.5) we see that the curves can be chosen to satisfy the following conditions.
(i) $|\xi|$ is bounded by a polynomial in $|\lambda|$ for any $\xi$ on one of the curves.
(ii) The arc length of each of the curves is bounded by a polynomial in $|\lambda|$.
(iii) If $\xi$ is on one of the curves, the distance from $\xi$ to the spectrum of $H_{0}$ is not less than $\varepsilon$.
(iv) If $\xi$ is a complex number, let $\xi_{\mathrm{R}}$ be its real part. Then

$$
\begin{aligned}
\xi_{\mathbf{R}} & >2 \varepsilon \text { for } \xi \text { on } \Gamma^{+}, \\
\xi_{\mathbf{R}} & \gtrless-2 \varepsilon \text { for } \xi \text { on } \Gamma^{-}, \\
\left|\xi_{\mathbf{R}}\right| & <\varepsilon \text { for } \xi \text { on } \Gamma^{0} .
\end{aligned}
$$

From the spectral theory for a linear transformation on a finite-dimensional vector space we have the formulae

$$
\begin{array}{rlrl}
e^{-t H_{0}} E^{\prime+} & =\int_{\Gamma^{+}} e^{-t \xi}\left(\xi-H_{0}\right)^{-1} d \xi, & & t>0 \\
e^{t H_{0}} E^{\prime-} & =\int_{\Gamma^{-}} e^{t \xi}\left(\xi-H_{0}\right)^{-1} d \xi, & & t>0 \\
e^{t H_{0}} E^{\prime 0} & =\int_{\Gamma^{0}} e^{t \xi}\left(\xi-H_{0}\right)^{-1} d \xi, & t \in \mathbf{R}
\end{array}
$$

Therefore if $t>0$

$$
\left|e^{-t H_{0}} E^{t+}\right| \gtrless \int_{\Gamma^{+}}\left|e^{-t \xi}\right| \cdot\left|\left(\xi-H_{0}\right)^{-1}\right| d \xi
$$

By Lemma 20 and conditions (iii) and (iv), this last expression is bounded by

$$
e^{-2 \varepsilon t} e^{-r} p_{2}(|\lambda|) \int_{\Gamma^{+}} p_{1}(|\xi|) d \xi
$$

Therefore by conditions (i) and (ii) there exists a polynomial $p$ such that

$$
\left|e^{-t H_{0}} E^{\prime+}\right|<e^{-2 \varepsilon t} \cdot p(|\lambda|), \quad t>0 .
$$

The inequalities for $\left|e^{t H_{0}} E^{\prime-}\right|$ and $\left|e^{t H_{0}} E^{\prime 0}\right|$ follow in the same way.
Write $E^{ \pm} \Phi$ as $\Phi^{ \pm}$and $E^{0} \Phi$ as $\Phi^{0}$.
Lemma 22. For any fixed $h$ in $\mathfrak{A}_{\mathfrak{p}}$ there are polynomials $p$ and $q$ such that for any $T>0$

$$
\left|\Phi^{-}\left(\exp T H_{0} ; h\right)\right| \gtrless p(|\lambda|) q(|\tau|) e^{-\varepsilon t}
$$

Proof. $\Gamma\left(H_{0}\right)$ and $E^{-}$commute. Therefore by the corollary to Lemma 15, $\Phi^{-}\left(\exp T H_{0} ; h\right)$ is equal to

$$
e^{T \Gamma\left(H_{0}\right)} \Phi^{-}(1 ; h)+\int_{0}^{T} e^{(T-t) \Gamma\left(H_{0}\right)} \cdot E^{-} \Psi\left(\exp T H_{0} ; h\right) d t
$$

The first term of this expression is easily handled with the help of (9.1) and Lemma 21. On the other hand, Lemma 21 tells us that

$$
\begin{aligned}
& \left|\int_{0}^{T} e^{(T-t) \Gamma\left(H_{0}\right)} E^{-} \Psi\left(\exp t H_{0} ; h\right) d t\right| \\
& \quad<p(|\lambda|) \int_{0}^{T} e^{-2 \varepsilon(T-t)}\left|\Psi\left(\exp t H_{0} ; h\right)\right| d t \\
& \quad<p(|\lambda|) e^{-\varepsilon t} \int_{0}^{T / 2}\left|\Psi\left(\exp t H_{0} ; h\right)\right| d t+p(|\lambda|) \int_{T / 2}^{T}\left|\Psi\left(\exp t H_{0} ; h\right)\right| d t
\end{aligned}
$$

Our lemma then follows from Lemma 16.

LEMMA 23. For any fixed $h$ in $\mathfrak{A}_{\mathfrak{p}}$ there are polynomials $p$ and $q$ such that for any $T>0$

$$
\left|\Phi^{+}\left(\exp T H_{0} ; h\right)\right|<p(|\lambda|) q(|\tau|) e^{-\varepsilon T}
$$

Proof. By means of a change of variables we can rewrite the integral equation of the corollary of Lemma 15. Then for $a$ in $\mathfrak{A}_{p}$ and $t$ in $\mathbf{R}$

$$
\Phi(a ; h)=e^{-T \Gamma\left(H_{0}\right)} \Phi\left(a \cdot \exp T H_{0} ; h\right)-\int_{0}^{T} e^{-t \Gamma\left(H_{0}\right)} \Psi\left(a \exp t H_{0} ; h\right) d t
$$

Operate on both sides of this equation by $E^{+}$, and let $T$ approach $\infty$. Now $\left|e^{-T \Gamma\left(H_{0}\right)} E^{+}\right|$decreases exponentially in $T$. However, by (9.1) and (4.1), $\left|\Phi\left(a \exp T H_{0} ; h\right)\right|$ is bounded by a polynomial in $T$. Therefore, the first term of right-hand side of the above equation approaches zero. We have the formula

$$
\Phi^{+}(a ; h)=-\int_{0}^{\infty} e^{-t \Gamma\left(H_{0}\right)} E^{+} \Psi\left(a \exp t H_{0} ; h\right) d t
$$

Let $a$ equal $\exp T H_{0}$. We obtain the equation

$$
\Phi^{+}\left(\exp T H_{0} ; h\right)=-\int_{T}^{\infty} e^{-(t-T) \Gamma\left(H_{0}\right)} E^{+} \Psi\left(\exp t H_{0} ; h\right) d t
$$

Now $\left|e^{-(t-T) \Gamma\left(H_{0}\right)} E^{+}\right|$is bounded by a polynomial in $|\lambda|$ if $t>T$. Our lemma then follows from Lemma 16.

For $a$ in $A_{\mathfrak{p}}, m$ in $M$, define

$$
\Theta(m a)=\Phi^{0}(m a)+\int_{0}^{\infty} e^{-t \Gamma\left(H_{0}\right)} E^{0} \Psi\left(m a \exp t H_{0}\right) d t
$$

(The integral converges absolutely by Lemmas 16 and 21.) It is clear that for any $h$ in $\mathfrak{A}_{\mathrm{p}}$

$$
\Theta(m a ; h)=\Phi^{0}(m a ; h)+\int_{0}^{\infty} e^{-t \Gamma\left(H_{0}\right)} E^{0} \Psi\left(m a \exp t H_{0} ; h\right) d t
$$

From the corollary of Lemma 15 we obtain the formula

$$
\Theta(m a)=\lim _{T \rightarrow \infty} e^{-T \Gamma\left(H_{0}\right)} \Phi^{0}\left(m a \exp T H_{0}\right)
$$

Therefore

$$
\begin{equation*}
\Theta\left(m a \exp t H_{0}\right)=e^{t \Gamma\left(H_{0}\right)} \boldsymbol{\Theta}(m a), \quad m \in M, a \in A_{\mathfrak{p}}, t \in \mathbf{R} \tag{9.7}
\end{equation*}
$$

Since $\Psi$ is $\bar{\tau}$-spherical, and since both the left and right actions of $\bar{\tau}(m)$ on $\mathscr{V}_{\tau}$ commute with $E^{\circ}$, we have the formula

$$
\begin{equation*}
\Theta\left(m_{1} m a m_{2}\right)=\bar{\tau}\left(m_{1}\right) \Theta(m a) \bar{\tau}\left(m_{2}\right), \quad m_{1}, m_{2}, m \in M, a \in A_{\mathfrak{p}} \tag{9.8}
\end{equation*}
$$

If $\varsigma$ is in $\mathfrak{Z}_{1}$

$$
\Theta(m a ; \varsigma)=\lim _{t \rightarrow \infty} e^{-t \Gamma\left(H_{0}\right)} \Phi^{0}\left(m a \exp t H_{0} ; \varsigma\right)
$$

By Lemma 15 this expression is equal to

$$
\Gamma(\varsigma) \Theta(m a)+\lim _{t \rightarrow \infty} e^{-t \Gamma\left(H_{0}\right)} E^{0} \Psi_{\varsigma}\left(m a \cdot \exp t H_{0}\right)
$$

Now, the last term in this formula approaches 0 as $t$ approaches $\infty$ by Lemmas 16 and 21 . Therefore

$$
\begin{equation*}
\Theta(m a ; \varsigma)=\Gamma(\varsigma) \Theta(m a) \tag{9.9}
\end{equation*}
$$

LEMMA 24. For any fixed $h$ in $\mathfrak{A}_{\mathfrak{p}}$ there are polynomials $p$ and $q$ such that for any $t>0$,

$$
\left|e^{-t \Gamma\left(H_{0}\right)} \Phi^{0}\left(\exp t H_{0} ; h\right)-\Theta(1 ; h)\right|<p(|\lambda|) q(|\tau|) e^{-2 \varepsilon t}
$$

Proof. Using the definition of $\theta$ and the formula (9.7) we see that

$$
\begin{aligned}
\Theta(1 ; h) & =e^{-T \Gamma\left(H_{0}\right)} \Theta\left(\exp T H_{0} ; h\right) \\
& =e^{-T \Gamma\left(H_{0}\right)} \Phi^{0}\left(\exp T H_{0} ; h\right)+\int_{T}^{\infty} e^{-t \Gamma\left(H_{0}\right)} E^{0} \Psi\left(\exp t H_{0} ; h\right) d t
\end{aligned}
$$

The lemma then follows from Lemmas 16 and 21.
COROLLARY. For any fixed $h$ in $\mathfrak{A}_{\mathfrak{p}}$ there are polynomials $p$ and $q$ such that for any $t>0$

$$
\left|\Phi\left(\exp t H_{0} ; h\right)-\Theta\left(\exp t H_{0} ; h\right)\right| \gtrless p(|\lambda|) q(|\tau|) e^{-\varepsilon t}
$$

PROOF. We see that

$$
\begin{aligned}
& \left|\Phi\left(\exp t H_{0} ; h\right)-\Theta\left(\exp t H_{0} ; h\right)\right| \\
& \quad<\left|\Phi^{+}\left(\exp t H_{0} ; h\right)\right|+\left|\Phi^{-}\left(\exp t H_{0} ; h\right)\right| \\
& \quad+\left|e^{t \Gamma\left(H_{0}\right)} E^{0}\right| \cdot\left|e^{-t \Gamma\left(H_{0}\right)} \Phi^{0}\left(\exp t H_{0} ; h\right)-\Theta(1 ; h)\right| .
\end{aligned}
$$

The corollary then follows from Lemmas 21, 22, 23, and 24.
It is clear from the definition of $\Theta$ that for any $a$ in $A_{\mathfrak{p}}$

$$
\begin{equation*}
E^{0} \Theta(a)=\Theta(a) \tag{9.10}
\end{equation*}
$$

For $a$ in $A_{\mathfrak{p}}$, and $1<j \gtrless r$, let $\theta_{j}(a)$ equal the vector $t_{j}(\Theta(a))$. Then

$$
\Theta(a)=\sum_{1 \gtrless j \gtrless r} \theta_{j}(a) \otimes \eta_{j}^{* *}
$$

Write $\theta_{1}$ as $\theta$.
LEMMA 25. For any nonnegative integer $n$ there are polynomials $p$ and $q$ such that for any $t>0$

$$
\left|\left(\frac{d}{d t}\right)^{n}\left[e^{t \rho\left(H_{0}\right)} \phi_{\lambda}^{\tau}\left(\exp t H_{0}\right)-\theta\left(\exp t H_{0}\right)\right]\right| \gtrless p(|\lambda|) q(|\tau|) e^{-\varepsilon t} .
$$

Proof. The expression

$$
\left(\frac{d}{d t}\right)^{n}\left[e^{t \rho\left(H_{0}\right)} \phi_{\lambda}^{\tau}\left(\exp t H_{0}\right)-\theta\left(\exp t H_{0}\right)\right]
$$

is equal to

$$
t_{1}\left(\Phi\left(\exp t H_{0} ; H_{0}^{n}\right)-\Theta\left(\exp t H_{0} ; H_{0}^{n}\right)\right)
$$

Now for any $u$ in $\mathscr{V}_{\tau},\left|t_{1}(u)\right| \gtrless|u|$. Our lemma then follows from the corollary to Lemma 24.

Lemma 26. If $\phi_{\lambda}^{\tau}$ is in $L^{2}(G) \otimes V_{\tau}, \theta$ is equal to zero.
Proof. The function $D(t)$ was defined in $\S 3$. Let

$$
S^{+}=\left\{t>0: D(t)<\frac{1}{2} e^{2 t \rho\left(H_{0}\right)}\right\} .
$$

$S^{+}$is compact. Let $\mathbf{R}^{+}=\{t>0\}$. Break up the integral

$$
\int_{0}^{\infty}\left|e^{t \rho\left(H_{0}\right)} \phi_{\lambda}^{\tau}\left(\exp t H_{0}\right)\right|^{2} d t
$$

into the sum of the integral over $S^{+}$and the integral over $\mathbf{R}^{+}-S^{+}$. The integral over $S^{+}$is finite since $\phi_{\lambda}^{\tau}$ is continuous. Now

$$
\begin{aligned}
& \int_{\mathbf{R}^{+}-s^{+}}\left|e^{t \rho\left(H_{0}\right)} \phi_{\lambda}^{\tau}\left(\exp t H_{0}\right)\right|^{2} d t<2 \cdot \int_{\mathbf{R}^{+}-s^{+}}|D(t)| \cdot\left|\phi_{\lambda}^{\tau}\left(\exp t H_{0}\right)\right|^{2} d t \\
& \leq 2 \int_{0}^{\infty} D(t) \cdot\left|\phi_{\lambda}^{\tau}\left(\exp t H_{0}\right)\right|^{2} d t .
\end{aligned}
$$

By (3.5) this last integral equals $(2 / c) \int_{G}\left|\phi_{\lambda}^{\tau}(x)\right|^{2} d x$, which is finite. Therefore,

$$
\int_{\underline{0}}^{\infty} e^{t \rho\left(H_{0}\right)}\left|\phi_{\lambda}^{\tau}\left(\exp t H_{0}\right)\right|^{2} d t<\infty .
$$

Therefore, by Lemma 25

$$
\int_{0}^{\infty}\left|\theta\left(\exp t H_{0}\right)\right|^{2} d t<\infty
$$

Now $\Theta\left(\exp t H_{0}\right)=e^{t \Gamma\left(H_{0}\right)} \Theta(1)$. But by (9.10), $\Theta(1)$ is nonzero only on the subspace of $\mathscr{V}_{\tau}$ spanned by eigenvectors of $\Gamma\left(H_{0}\right)$ that are associated with purely imaginary eigenvalues. Therefore, $\theta$ must be zero.
10. Completion of the proof of Theorem $3^{\prime}(\mathrm{a})$. We would like to apply the results of $\S 9$ to the collection $\left\{\phi_{\omega}^{\tau}(x): \omega \in \mathscr{E}_{d}, \tau \in \mathscr{E}_{K}^{2}\right\}$ of $\tau$-spherical functions which are derived from the matrix elements of square-integrable representations. That is,

$$
\begin{equation*}
\phi(x)=\phi_{\omega}^{\tau}(x)=(\operatorname{dim} \tau)^{-1 / 2} \sum_{i j} \xi_{1 j} \otimes \xi_{2 j}^{*}\left(\Phi_{1 i}, \pi_{\omega}(x) \Phi_{2 j}\right) \tag{10.1}
\end{equation*}
$$

in the notation of (5.1). Here $\pi_{\omega}$ is a square-integrable representation in the class of $\omega$, acting on the Hilbert space $\mathscr{H}_{\omega}$. If $\tau=\left(\tau_{1}, \tau_{2}\right)$ then $\Phi_{1 i}$ and $\Phi_{2 j}$ are unit vectors in $\mathscr{H}_{\omega}$ that transform under $\left.\pi_{\omega}\right|_{K}$ according to $\tau_{1}$ and $\tau_{2}$ respectively. $\left\{\xi_{1 i}\right\}$ and $\left\{\xi_{2 j}\right\}$ are orthonormal bases for the spaces on which $\tau_{1}$ and $\tau_{2}$ act, so $\left\{\xi_{1 i} \otimes \xi_{2 j}^{*}\right\}$ is an orthonormal base for $V_{\tau}$, the space on which $\tau$ acts. Therefore

$$
\begin{equation*}
\left|\phi_{\omega}^{\tau}(x)\right|^{2}=(\operatorname{dim} \tau)^{-1} \sum_{i j}\left|\left(\Phi_{1 i}, \pi_{\omega}(x) \Phi_{2 j}\right)\right|^{2} . \tag{10.2}
\end{equation*}
$$

It follows that

$$
\begin{equation*}
\left|\phi_{\omega}^{\tau}(x)\right|<1, \quad x \in G . \tag{10.3}
\end{equation*}
$$

From the Schur orthogonality relations on $G$, (2.7) and (2.8), we see that

$$
\begin{equation*}
\int_{G}\left|\phi_{\omega}^{\tau}(x)\right|^{2} d x=\beta(\omega)^{-1} \tag{10.4}
\end{equation*}
$$

Now we need to establish the estimate (9.1) for our collection $\left\{\phi_{\omega}^{\tau}\right\}$. Fix $g_{1}, g_{2}$ in $\mathscr{B}$. We use (10.1) and Lemma 13 to obtain an expression for $\phi_{\omega}^{\tau}\left(g_{1} ; x ; g_{2}\right)$. We obtain polynomials $p$ and $q$, orthogonal sets of vectors $\left\{\Psi_{1 \alpha}\right\}_{1 \gtrless \alpha ₹ t_{1}}$, $\left\{\Psi_{2 \beta}\right\}_{1<\beta<t_{2}}$ in $\mathscr{H}_{\omega}$, and representations $\left\{\tau_{1 \alpha}\right\}_{1<\alpha<t_{1}},\left\{\tau_{2 \beta}\right\}_{1<\beta<t_{2}}$ in $\mathscr{E}_{K}$ that satisfy the conditions of Lemma 13 and such that

$$
\begin{equation*}
\left|\phi_{\omega}^{\tau}\left(g_{1} ; x ; g_{2}\right)\right|^{2}=\sum_{\alpha=1}^{t_{1}} \sum_{\beta=1}^{t_{2}}\left|\left(\Psi_{1 \alpha}, \pi_{\omega}(x) \Psi_{2 \beta}\right)\right|^{2} \tag{10.5}
\end{equation*}
$$

Suppose $\omega=\omega(\lambda)$ for some $\lambda$ in $L^{\prime}$. Then the infinitesimal character of $\pi_{\omega}$ is $\chi_{\lambda_{y}}$. (8.2) tells us that $\left|\lambda_{y}\right|^{2}=|\omega|^{2}$. Therefore from (10.5) and the conditions of Lemma 13 we obtain the inequality

$$
\begin{equation*}
\left|\phi_{\omega}^{\tau}\left(g_{1} ; x ; g_{2}\right)\right|^{2} \gtrless t_{1} t_{2} \cdot p(|\omega|) \cdot q(|\tau|)<p(|\omega|) \cdot q(|\tau|)^{3} \tag{10.6}
\end{equation*}
$$

for any $x$ in $G$. From (2.8) and (10.5) we also see that

$$
\begin{equation*}
\int_{G}\left|\phi_{\omega}^{\tau}\left(g_{1} ; x ; g_{2}\right)\right|^{2} d x<\beta(\omega)^{-1} \cdot p(|\omega|) \cdot q(|\tau|)^{3}, \quad x \in G . \tag{10.7}
\end{equation*}
$$

Lemma 27. There are polynomials $p$ and $q$ such that for any $x$ in $G$,

$$
\left|\phi_{\omega}^{\tau}(x)\right|<p(|\omega|) \cdot q(|\tau|) \Xi(x)
$$

Proof. Recall that $r_{1}$ and $r_{2}$ denote the number of roots in $P_{+}$which when restricted to $\mathfrak{a}_{\mathfrak{p}}$ are respectively equal to $\mu_{0}$ and $2 \mu_{0}$. Then $\rho\left(H_{0}\right)=\frac{1}{2}\left(r_{1}+2 r_{2}\right)$. Also

$$
B\left(H_{0}, H_{0}\right)=2\left(r_{1}+4 r_{2}\right)=r^{2}
$$

Define $D(t)$ as.in (3.5). Define a set $S$ by

$$
S=\left\{t \in \mathbf{R}:|D(t)| \gtrless \frac{1}{2} e^{2|t| \rho\left(H_{0}\right)}\right\}
$$

$S$ is bounded. Define a positive infinitely differentiable function $\eta$ on $\mathbf{R}$ such that for $t$ not in $S$

$$
\eta(t)=\frac{1}{2} e^{2|t| \rho\left(H_{0}\right)} .
$$

Let $h(t)$ be the function

$$
\phi_{\omega}^{\tau}\left(\exp t H_{0}\right) \cdot \eta^{1 / 2}(t) .
$$

Clearly $h(t)$ is infinitely differentiable. We wish to show that $h(t)$ is in $\mathscr{S}(\mathbf{R}) \otimes V_{\tau}$, the tensor product of the Schwartz space of $\mathbf{R}$ with $V_{\tau}$, the Hilbert space on which $\tau$ acts.

For any nonnegative integer $n$

$$
\left(\frac{d}{d t}\right)^{n} h(t)=\sum_{i} \frac{n!}{(n-i)!(i)!}\left(\frac{d}{d t}\right)^{i} \phi_{\omega}^{\tau}\left(\exp t H_{0}\right) \cdot\left(\frac{d}{d t}\right)^{n-i} \eta^{1 / 2}(t)
$$

There exist constants $C^{i}$ such that

$$
\left|\left(\frac{d}{d t}\right)^{n-i} \eta^{1 / 2}(t)\right|<C^{i} e^{|t| \rho\left(H_{0}\right)}
$$

for all $t$ not in $S$ and for $0<i<n$. Therefore since $\eta^{1 / 2}(t)$ is differentiable, there is a constant $C_{n}$ such that

$$
\left|\left(\frac{d}{d t}\right)^{n-i} \eta^{1 / 2}(t)\right| \gtrless C_{n} e^{|t| \rho\left(H_{0}\right)}
$$

for all real $t$ and for $0<i<n$. Now $\phi_{\omega}^{\tau}$ is in $\mathscr{C}(G) \otimes V_{\tau}$ by [3(1), Lemma 65, Corollary 1], so for each $s$ there is a constant $C_{\omega}^{\tau}$ such that

$$
\left|\phi_{\omega}^{\tau}\left(\exp t H_{0} ; H_{0}^{i}\right)\right| \gtrless C_{\omega}^{\tau} \cdot \Xi\left(\exp t H_{0}\right) \cdot\left(1+\sigma\left(t H_{0}\right)\right)^{-s}
$$

for $t$ in $\mathbf{R}$ and $0<i<n$. But

$$
1+\sigma\left(\exp t H_{0}\right)=1+r|t|, \quad t \in \mathbf{R}
$$

By (4.1)

$$
\Xi\left(\exp t H_{0}\right)<c(1+r|t|)^{d} e^{-|t| \rho\left(H_{0}\right)}, \quad t \in \mathbf{R} .
$$

Therefore for all real $t$,

$$
\left.\left|\left(\frac{d}{d t}\right)^{n} h(t)\right| \stackrel{\gtrless}{\sum} \frac{n!}{(n-i)!(i)!}\right) \cdot c \cdot C_{\omega}^{\tau} \cdot C_{n}(1+r|t|)^{-s+d}
$$

This proves that $h$ is in $\mathscr{S}(\mathbf{R}) \otimes V_{\tau}$.
We shall take the Fourier transform of $h(t)$ on $\mathbf{R}$. Write

$$
\hat{h}(s)=\int_{-\infty}^{\infty} h(t) e^{-i s t} d t, \quad s \in \mathbf{R}
$$

For any real $t$

$$
\begin{aligned}
|h(t)| & =(2 \pi)^{-1}\left|\int_{-\infty}^{\infty} \hat{h}(s) e^{i s t} d s\right| \\
& \gtrless(2 \pi)^{-1}\left|\int_{-\infty}^{\infty} \hat{h}(s) e^{i s t}(1+i s) \cdot(1+i s)^{-1} d s\right| \\
& \gtrless(2 \pi)^{-1}\left(\int_{-\infty}^{\infty}|\hat{h}(s)(1+i s)|^{2} d s\right)^{1 / 2} \cdot\left(\int_{-\infty}^{\infty} \frac{d s}{1+s^{2}}\right)^{1 / 2}
\end{aligned}
$$

If we write $h^{\prime}$ for the derivative of $h$, this last expression equals $(\pi)^{-1 / 2}\left\|\hat{h}+\hat{h}^{\prime}\right\|_{2}$, which in turn is equal to $(\pi)^{-1 / 2}\left\|h+h^{\prime}\right\|_{2}$. Therefore for any real $t$

$$
|h(t)| ₹ \pi^{-1 / 2}\left(\|h\|_{2}+\left\|h^{\prime}\right\|_{2}\right)
$$

We have used the fact that all the above integrals are absolutely convergent. This is true because both $h$ and $\hat{h}$ are in $\mathscr{S}(\mathbf{R}) \otimes V_{\tau}$.

We see that it is necessary to estimate both $\|h\|_{2}$ and $\left\|h^{\prime}\right\|_{2}$. Clearly $\|h\|_{2}$ is equal to

$$
\int_{S}\left|\phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|^{2} \eta(t) d t+\int_{\mathbf{R}-S}\left|\phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|^{2} \eta(t) d t
$$

Now $\eta(t)$ is bounded on $S$, so by (10.3) there is a constant $C_{1}$ independent of $\omega$ and $\tau$, such that

$$
\int_{S}\left|\phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|^{2} \eta(t) d t \gtrless C_{1}
$$

On the other hand,

$$
\begin{aligned}
\int_{\mathbf{R}-S} & \left|\phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|^{2} \eta(t) d t \\
& <2 \int_{\mathbf{R}-S}\left|\phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|^{2} \cdot|D(t)| d t \\
& <2 \int_{\mathbf{R}}\left|\phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|^{2} \cdot|D(t)| d t \\
& <2 \int_{K \times K} \int_{\mathbf{R}}\left|\phi_{\omega}^{\tau}\left(k_{1} \cdot \exp t H_{0} \cdot k_{2}\right)\right|^{2} \cdot|D(t)| d t d k_{1} d k_{2}
\end{aligned}
$$

By (3.5) this last integral equals

$$
\frac{4}{c} \int_{G}\left|\phi_{\omega}^{\tau}(x)\right|^{2} d x
$$

This expression equals $\left(\frac{4}{c}\right) \beta(\omega)^{-1}$ by (10.4). Since $\beta(\omega)^{-1}$ is bounded independently of $\omega$, by Lemma $3,\|h\|_{2}$ is bounded independently of $\omega$ and $\tau$.

Now we shall estimate $\left\|h^{\prime}\right\|_{2}$. First of all we need a bound on the expression

$$
\int_{\mathbf{R}}\left|\phi_{\omega}^{\tau}\left(\exp t H_{0} ; H_{0}\right)\right|^{2} \cdot|D(t)| d t
$$

For any $x$ in $G$ we have the following formula by (10.5)

$$
\left|\phi_{\omega}^{\tau}\left(x ; H_{0}\right)\right|^{2}=\sum_{\alpha=1}^{t_{1}} \sum_{\beta=1}^{t_{2}}\left|\left(\Psi_{1 \alpha}, \pi_{\omega}(x) \Psi_{2 \beta}\right)\right|^{2}
$$

where $\left\{\Psi_{1 \alpha}\right\},\left\{\Psi_{2 \beta}\right\}, t_{1}$ and $t_{2}$ satisfy the conditions of Lemma 13. Define

$$
f_{\alpha, \beta}(x)=\left(\Psi_{1 \alpha}, \pi_{\omega}(x) \Psi_{2 \beta}\right)
$$

By (10.7) and Lemma 3, there are polynomials $p_{1}$ and $q_{1}$ such that

$$
\int_{G} \sum_{\alpha, \beta}\left|f_{\alpha, \beta}(x)\right|^{2} d x<p_{1}(|\omega|) \cdot q_{1}(|\tau|)
$$

Let $\tau_{\alpha, \beta}$ be the double representation $\left(\tau_{1 \alpha}, \tau_{2 \beta}\right)$ of $K$. As we did in $\S 5$, we can associate a $\tau_{\alpha, \beta}$-spherical function $\phi_{\alpha, \beta}$ to each $f_{\alpha, \beta}$. (Unlike the situation in $\S 5$, $\Psi_{1 \alpha}$ and $\Psi_{2 \beta}$ are not of norm 1.) Formula (10.2) and the Schur orthogonality relations, (2.7) and (2.8), can be used to derive

$$
\int_{G}\left|f_{\alpha, \beta}(x)\right|^{2} d x=\int_{G}\left|\phi_{\alpha, \beta}(x)\right|^{2} d x
$$

Using this and formulas (10.2) and (3.5), we obtain the following inequality:

$$
\begin{aligned}
& \int_{\mathbf{R}}\left|f_{\alpha, \beta}\left(\exp t H_{0}\right)\right|^{2} \cdot|D(t)| d t \\
& ₹ \operatorname{dim} \tau_{1 \alpha} \cdot \operatorname{dim} \tau_{2 \beta} \cdot \int_{\mathbf{R}}\left|\phi_{\alpha, \beta}\left(\exp t H_{0}\right)\right|^{2} \cdot|D(t)| d t \\
&=\operatorname{dim} \tau_{1 \alpha} \cdot \operatorname{dim} \tau_{2 \beta} \cdot \int_{K \times K} \int_{\mathbf{R}}\left|\phi_{\alpha, \beta}\left(k_{1} \cdot \exp t H_{0} \cdot k_{2}\right)\right|^{2} \cdot|D(t)| d t d k_{1} d k_{2} \\
&=(2 / c) \operatorname{dim} \tau_{1 \alpha} \cdot \operatorname{dim} \tau_{2 \beta} \cdot \int_{G}\left|\phi_{\alpha, \beta}(x)\right|^{2} d x \\
&=(2 / c) \operatorname{dim} \tau_{1 \alpha} \cdot \operatorname{dim} \tau_{2 \beta} \cdot \int_{G}\left|f_{\alpha, \beta}(x)\right|^{2} d x
\end{aligned}
$$

Now the dimensions of $\tau_{1 \alpha}$ and $\tau_{2 \beta}$ are bounded by polynomials in $\left|\tau_{1 \alpha}\right|$ and $\left|\tau_{2 \beta}\right|$, which in turn are bounded by polynomials in $|\tau|$ by Lemma 13. Therefore there are polynomials $p_{2}$ and $q_{2}$ such that

$$
\begin{aligned}
& \int_{\mathbf{R}}\left|\phi_{\omega}^{\tau}\left(\exp t H_{0} ; H_{0}\right)\right|^{2} \cdot|D(t)| d t \\
& \quad<(2 / c) p_{2}(|\omega|) \cdot p_{2}(|\tau|) \cdot \int_{G} \sum_{\alpha, \beta}\left|f_{\alpha, \beta}(x)\right|^{2} d x
\end{aligned}
$$

Therefore

$$
\begin{align*}
& \int_{\mathbf{R}}\left|\phi_{\omega}^{\tau}\left(\exp t H_{0} ; H_{0}\right)\right|^{2}|D(t)| d t  \tag{10.8}\\
& \quad<(2 / c) p_{1}(|\omega|) \cdot p_{2}(|\omega|) \cdot q_{1}(|\tau|) \cdot q_{2}(|\tau|)
\end{align*}
$$

We can now estimate $\left\|h^{\prime}\right\|_{2}$. This norm is the sum of the following two expressions:

$$
\begin{equation*}
\left(\int_{\mathbf{R}}\left|\frac{d}{d t} \phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|^{2} \cdot \eta(t) d t\right)^{1 / 2} \tag{i}
\end{equation*}
$$

(ii)

$$
\left(\int_{\mathbf{R}}\left|\phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|^{2} \cdot\left|\frac{d}{d t} \eta^{1 / 2}(t)\right|^{2} d t\right)^{1 / 2}
$$

There is a constant $c_{2}$ such that

$$
\left|\frac{d}{d t} \eta^{1 / 2}(t)\right|^{2}<c_{2} \eta(t)
$$

Therefore the second expression is bounded by $c_{2}\|h\|_{2}$. We break up the integral in the first expression into integrals over $\mathbf{R}-S$ and $S$. The integral over $S$ causes no problem because $\eta(t)$ is bounded on $S$, and by (10.6), |(d/dt) $\phi_{\omega}^{\tau}\left(\exp t H_{0}\right) \mid$ is bounded by a polynomial in $|\omega|$ and $|\tau|$. On the other hand,

$$
\begin{aligned}
\int_{\mathbf{R}-S}\left|\frac{d}{d t} \phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|^{2} \cdot \eta(t) d t & <2 \int_{\mathbf{R}-S}\left|\frac{d}{d t} \phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|^{2} \cdot|D(t)| d t \\
& <2 \int_{\mathbf{R}}\left|\frac{d}{d t} \phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|^{2} \cdot|D(t)| d t
\end{aligned}
$$

By (10.8) this last expression is bounded by a polynomial in $|\omega|$ and $|\tau|$. Therefore there are polynomials $p_{3}$ and $q_{3}$ such that

$$
\left\|h^{\prime}\right\|_{2}<p_{3}(|\omega|) \cdot q_{3}(|\tau|)
$$

We have now shown that there are polynomials $p_{4}$ and $q_{4}$ such that

$$
|h(t)| ₹ p_{4}(|\omega|) \cdot q_{4}(|\tau|), \quad t \in \mathbf{R}
$$

This implies the inequality

$$
\left|\phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|<\eta^{-1 / 2}(t) p_{4}(|\omega|) \cdot q_{4}(|\tau|)
$$

There is a constant $c_{3}$ such that

$$
\eta^{-1 / 2}(t) \gtrless c_{3} e^{-|t| \rho\left(H_{0}\right)}, \quad t \in \mathbf{R}
$$

Also, by (4.1) and (4.3),

$$
e^{-|t| \rho\left(H_{0}\right)} ₹ \Xi\left(\exp t H_{0}\right), \quad t \in \mathbf{R} .
$$

Therefore, we can choose polynomials $p$ and $q$ such that

$$
\left|\phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|<p(|\omega|) \cdot q(|\tau|) \cdot \Xi\left(\exp t H_{0}\right), \quad t \in \mathbf{R}
$$

However, each $x$ in $G$ is of the form $k_{1} \cdot \exp t H_{0} \cdot k_{2}$ for $k_{1}$ and $k_{2}$ in $K$ and $t$ in R. Therefore

$$
\begin{aligned}
\left|\phi_{\omega}^{\tau}(x)\right| & =\left|\tau_{1}\left(k_{1}\right) \phi_{\omega}^{\tau}\left(\exp t H_{0}\right) \tau_{2}\left(k_{2}\right)\right|=\left|\phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right| \\
& <p(|\omega|) \cdot q(|\tau|) \cdot \Xi\left(\exp t H_{0}\right) \\
& =p(|\omega|) \cdot q(|\tau|) \cdot \Xi(x) .
\end{aligned}
$$

This proves our lemma.
COROLLARY. For any $g_{1}$ and $g_{2}$ in $\mathscr{B}$ there are polynomials $p$ and $q$ such that

$$
\left|\phi_{\omega}^{\tau}\left(g_{1} ; x ; g_{2}\right)\right| \gtrless p(|\omega|) \cdot q(|\tau|) \cdot \Xi(x), \quad x \in G .
$$

Proof. By (10.5) we have the formula

$$
\left|\phi_{\omega}^{\tau}\left(g_{1} ; x ; g_{2}\right)\right|^{2}=\sum_{\alpha=1}^{t_{1}} \sum_{\beta=1}^{t_{2}}\left|\left(\Psi_{1 \alpha}, \pi_{\omega}(x) \Psi_{2 \beta}\right)\right|^{2}
$$

Let $\phi_{\alpha, \beta}$ be the ( $\tau_{1 \alpha}, \tau_{2 \beta}$ )-spherical function associated with $\left(\Psi_{1 \alpha}, \pi_{\omega}(x) \Psi_{2 \beta}\right)$. Apply Lemma 27 to $\phi_{\alpha, \beta}$. The proof of the corollary then follows from the conditions in Lemma 13.

This corollary verifies the inequality (9.1) for our collection $\left\{\phi_{\omega}^{\tau}\right\}$.
Each $\lambda$ in $L_{y}^{\prime}$ is regular, so Assumption 1 of $\S 9$ holds. To see that Assumption 2 is valid, we must look at the eigenvalues of the linear transformation $\Gamma\left(H_{0}\right)$ defined in $\S 9$. ( $\Gamma\left(H_{0}\right)$, of course, depends on $\omega$.) We see from Lemma 17, using the notation there, that the set of eigenvalues of $\Gamma\left(H_{0}\right)$ is equal to

$$
\left\{\left\langle s_{i} \lambda, H_{0}\right\rangle: \lambda \subset L_{y}^{\prime}, 1<i<r\right\} .
$$

Assumption 2 then is established by the following lemma.

LEMMA 28. $\left\{\left\langle s_{i} \lambda, H_{0}\right\rangle: \lambda \in L_{y}^{\prime}, 1 ₹ i ₹ r\right\}$ is a set of real numbers which generates a lattice in $\mathbf{R}$.

Proof. Any $\lambda$ in $L_{y}^{\prime}$ assumes real values on the vector space $\mathfrak{a}_{\mathfrak{p}}+(-1)^{1 / 2} \mathfrak{a}_{\mathfrak{e}}$. Furthermore, any element in $W$ maps the real vector space $\mathfrak{a}_{\mathfrak{p}}+(-1)^{1 / 2} \mathfrak{a}_{\mathfrak{k}}$ onto itself. $\left\langle s_{i} \lambda, H_{0}\right\rangle$ is equal to $\lambda\left(s_{i}^{-1} H_{0}\right)$, which is then a real number.

Let $\widetilde{L_{y}}$ be the lattice of functions from $a_{c}$ into $\mathbf{C}$ generated by the positive roots $P . \widetilde{L_{y}}$ is invariant under the action of $W$. If $\mu$ is in $\widetilde{L_{y}}$, the real number $\left\langle s_{i} \mu, H_{0}\right\rangle$ is an integer because for any $\alpha$ in $P, \alpha\left(H_{0}\right)$ is either 0,1 , or 2 . But $\widetilde{L}_{y}$ is of finite index in $L_{y}$ since $G$ has finite center. This proves the lemma.

We have shown that we can apply the results of $\S 9$ to our collection $\left\{\phi_{\omega}^{\tau}\right\}$. By Lemma 26, the function $\theta$ defined in $\S 9$ is zero. Then by Lemma 25 there are polynomials $p$ and $q$ such that

$$
\left|\phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|<p(|\omega|) \cdot q(|\tau|) \cdot e^{-t \rho\left(H_{0}\right)} e^{-\varepsilon t}, \quad t \gg
$$

Therefore, for any $s>0$ there are polynomials $p$ and $q$ such that

$$
\left|\phi_{\omega}^{\tau}\left(\exp t H_{0}\right)\right|<p(|\omega|) \cdot q(|\tau|) \cdot e^{-t \rho\left(H_{0}\right)}(1+r t)^{-s}, \quad t>0
$$

But for $t>0$ we have the formulae

$$
e^{-t \rho\left(H_{0}\right)}<\Xi\left(\exp t H_{0}\right), \quad 1+\sigma\left(\exp t H_{0}\right)=1+r t
$$

Now each $x$ in $G$ is of the form $k_{1} \cdot \exp t H_{0} \cdot k_{2}$ for $t>0$ and $k_{1}$ and $k_{2}$ in $K$. But $\phi_{\omega}^{\tau}$ is $\tau$-spherical and $\Xi$ and $\sigma$ are bi-invariant under $K$. We have proved the following lemma

Lemma 29. For any $x$ in $G$

$$
\left|\phi_{\omega}^{\tau}(x)\right| \gtrless p(|\omega|) \cdot q(|\tau|) \cdot \Xi(x)(1+\sigma(x))^{-s} .
$$

COROLLARY. For any $g_{1}$ and $g_{2}$ in $\mathscr{B}$ there are polynomials $p$ and $q$ such that for any $x$ in $G$

$$
\left|\phi_{\omega}^{\tau}\left(g_{1} ; x ; g_{2}\right)\right| \gtrless p(|\omega|) \cdot q(|\tau|) \cdot \Xi(x)(1+\sigma(x))^{-s}
$$

Proof. The proof follows from Lemma 29 in the same way as the corollary to Lemma 27 followed from Lemma 27.

We have proved Theorem $3^{\prime}(\mathrm{a})$.
11. Application to the continuous series. Let $\tau$ be an irreducible unitary double representation of $K$ on the Hilbert space $V_{\tau}$. Recall from $\S 5$ that

$$
L^{\tau}=\bigoplus_{\sigma \in \mathscr{E}_{M}} L_{\sigma}^{\tau}
$$

is an orthogonal direct sum. For any fixed $\tau, L_{\sigma}^{\tau}=0$ for all but a finite number of $\sigma$.

We would like to apply the results of $\S 9$ to the collection $\left\{E_{\Lambda}\left(\psi_{\sigma}^{\tau}: x\right)\right\}$ of Eisenstein integrals. $\Lambda$ is to range over the nonzero reals, and $\tau$ and $\sigma$ will range over $\mathscr{E}_{K}^{2}$ and $\mathscr{E}_{M}$ respectively. $\psi_{\sigma}^{\tau}$ will be any unit vector in $L_{\sigma}^{\tau}$.

We will have to check the estimate (9.1) and Assumptions 1 and 2 of $\S 9$ for our collection $\left\{E_{\Lambda}\left(\psi_{\sigma}^{\tau}: x\right)\right\}$. Then we will examine the functions $\theta=\theta_{\Lambda}\left(\psi_{\sigma}^{\tau}: m a\right)$ associated with $E_{\Lambda}\left(\psi_{\sigma}^{\tau}: x\right)$ in $\S 9$.

LEMMA 30. Fix $g_{1}, g_{2}$ in $\mathscr{B}$. Then for fixed $x$ in $G, E_{\Lambda}\left(\psi_{\sigma}^{\tau}: g_{1} ; x ; g_{2}\right)$ can be regarded as an entire function of $\Lambda$. Also, there are polynomials $p_{1}, p_{2}$, and $q$, dependent only on $g_{1}$ and $g_{2}$, such that for every integer $n>0$, for $t>0$, and for $\Lambda$ in $\mathbf{C}$,

$$
\begin{aligned}
&\left|\left(\frac{d}{d \Lambda}\right)^{n} E_{\Lambda}\left(\psi_{\sigma}^{\tau}: g_{1} ; \exp t H_{0} ; g_{2}\right)\right| \\
& \quad<e^{\left|\Lambda_{I}\right| t} \cdot p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|) \mid \cdot q(|\tau|) \cdot \Xi\left(\exp t H_{0}\right) \cdot t^{n}
\end{aligned}
$$

(We write $\Lambda_{I}$ for the imaginary part of the complex number $\Lambda$.)
Proof. We have the formula (5.5):

$$
E_{\Lambda}\left(\psi_{\sigma}^{\tau}: x\right)=\int_{K} \psi_{\sigma}^{\tau}(x k) \tau\left(k^{-1}\right) e^{\left(i \Lambda \mu_{0}-\rho\right)(H(x k))} d k
$$

For fixed $x$, this is clearly an entire function of $\Lambda$. Derivatives of $E_{\Lambda}\left(\psi_{\sigma}^{\tau}: x\right)$ by means of left or right invariant differential operators are entire functions in $\Lambda$.

From (5.6) we obtain the formula

$$
\left(\frac{d}{d \Lambda}\right)^{n} E_{\Lambda}\left(\psi_{\sigma}^{\tau}: x\right)=\left(t_{1} t_{2}\right)^{-1 / 2} \sum_{i j} \xi_{1 i} \otimes \xi_{2 j}^{*} \cdot\left(\frac{d}{d \Lambda}\right)^{n}\left(\Phi_{1 i}, \pi_{\sigma, \Lambda}(x) \Phi_{2 j}\right)
$$

where $t_{1} t_{2}$ is the dimension of the representation $\tau$. We apply Lemma 13 to each of the functions ( $\left.\Phi_{1 i}, \pi_{\sigma, \Lambda}(x) \Phi_{2 j}\right)$. As a result we obtain polynomials $p, q$, orthogonal sets of vectors

$$
\left\{\Psi_{1 \alpha}: 1 \gtrless \alpha<t_{1}\right\}, \quad\left\{\Psi_{2 \beta}: 1 \gtrless \beta \gtrless t_{2}\right\}
$$

in $\mathscr{H}_{\sigma, \Lambda}$, and the representations

$$
\left\{\tau_{1 \alpha}: 1 \gtrless \alpha<t_{1}\right\}, \quad\left\{\tau_{2 \beta}: 1 \gtrless \beta<t_{2}\right\}
$$

in $\mathscr{E}_{K}$ that satisfy the conditions of Lemma 13 . In addition

$$
\begin{equation*}
\left|\left(\frac{d}{d \Lambda}\right)^{n} E_{\Lambda}\left(\psi_{\sigma}^{\tau}: g_{1} ; x ; g_{2}\right)\right|^{2}=\sum_{\alpha=1}^{t_{1}} \sum_{\beta=1}^{t_{2}}\left|\left(\frac{d}{d \Lambda}\right)^{n}\left(\Psi_{1 \alpha}, \pi_{\sigma, \Lambda}(x) \Psi_{2 \beta}\right)\right|^{2} \tag{11.1}
\end{equation*}
$$

For any ( $\alpha, \beta$ ), the vectors $\Psi_{1 \alpha}$ and $\Psi_{2 \beta}$ transform under $\left.\pi_{\sigma, \Lambda}\right|_{K}$ according to the representations $\tau_{1 \alpha}$ and $\tau_{2 \beta}$ respectively. Let $\tau_{\alpha, \beta}$ be the double representation $\left(\tau_{1 \alpha}, \tau_{2 \beta}\right)$ of $K$. Fix a vector $\psi_{\alpha, \beta}$ in $L_{\sigma}^{\tau_{\alpha, \beta}}$ such that $E_{\Lambda}\left(\psi_{\alpha \beta}: x\right)$ is the $\tau_{\alpha, \beta}$-spherical function associated with the function $\left(\Psi_{1 \alpha}, \pi_{\sigma, \Lambda}(x) \Psi_{2 \beta}\right)$. By (5.6) we have the inequality

$$
\begin{aligned}
& \left|\left(\frac{d}{d \Lambda}\right)^{n}\left(\Psi_{1 \alpha}, \pi_{\sigma, \Lambda}\left(\exp t H_{0}\right) \Psi_{2 \beta}\right)\right| \\
& \quad<\left(\operatorname{dim} \tau_{\alpha, \beta}\right)^{1 / 2} \cdot\left|\left(\frac{d}{d \Lambda}\right)^{n} E_{\Lambda}\left(\psi_{\alpha, \beta}: \exp t H_{0}\right)\right|
\end{aligned}
$$

By (5.5), the right-hand side of this inequality equals $\left(\operatorname{dim} \tau_{\alpha, \beta}\right)^{1 / 2} \cdot\left|\left(\frac{d}{d \Lambda}\right)^{n} \int_{K} \psi_{\alpha, \beta}\left(\exp t H_{0} \cdot k\right) \tau_{\alpha, \beta}\left(k^{-1}\right) \cdot e^{\left(i \Lambda \mu_{0}-\rho\right)\left(H\left(\exp t H_{0} \cdot k\right)\right)} d k\right|$,
and this expression equals

$$
\begin{aligned}
\left(\operatorname{dim} \tau_{\alpha, \beta}\right)^{1 / 2} \cdot \mid \int_{K} \psi_{\alpha, \beta}\left(\exp t H_{0} \cdot k\right) \tau_{\alpha, \beta}\left(k^{-1}\right) \cdot & \mu_{0}\left(H\left(\exp t H_{0} \cdot k\right)\right)^{n} \\
\cdot & e^{\left(i \Lambda \mu_{0}-\rho\right)\left(H\left(\exp t H_{0} \cdot k\right)\right)} d k \mid
\end{aligned}
$$

Now there is a real $t_{k}$ such that

$$
H\left(\exp t H_{0} \cdot k\right)=t_{k} H_{0}
$$

Then $[\mathbf{3}(\mathbf{g})$, Lemma 35] and [ $\mathbf{3}(\mathrm{g})$, Lemma 35, Corollary 2] tell us that

$$
\left|t_{k}\right|<t, \quad t>0, k \in K
$$

Therefore, if $t>0$, we have the two inequalities

$$
\left|\mu_{0}\left(H\left(\exp t H_{0} \cdot k\right)\right)\right|^{n}=\left|t_{k}\right|^{n} \gtrless t^{n}
$$

and

$$
\left|e^{\left(i \Lambda \mu_{0}-\rho\right)\left(H\left(\exp t H_{0} \cdot k\right)\right)}\right|<e^{\left|\Lambda_{I}\right| t}
$$

Notice also that

$$
\left|\psi_{\alpha, \beta}\left(\exp t H_{0} \cdot k\right) \tau_{\alpha, \beta}\left(k^{-1}\right)\right|=\left|\psi_{\alpha, \beta}(1)\right| .
$$

Therefore, for $t>0$, we have the inequality

$$
\begin{aligned}
& \left|\left(\frac{d}{d \Lambda}\right)^{n}\left(\Psi_{1 \alpha}, \pi_{\sigma, \Lambda}\left(\exp t H_{0}\right) \Psi_{2 \beta}\right)\right| \\
& \quad<\left(\operatorname{dim} \tau_{\alpha, \beta}\right)^{1 / 2} \cdot\left|\psi_{\alpha, \beta}(1)\right| \cdot e^{\left|\Lambda_{I}\right| t} \cdot t^{n} \cdot \int_{K} e^{-\rho\left(H\left(\exp t H_{0} \cdot k\right)\right)} d k \\
& \quad=\left(\operatorname{dim} \tau_{\alpha, \beta}\right)^{1 / 2} \cdot\left|\psi_{\alpha, \beta}(1)\right| \cdot e^{\left|\Lambda_{I}\right| t} \cdot t^{n} \cdot \Xi\left(\exp t H_{0}\right)
\end{aligned}
$$

But $\left|\psi_{\alpha, \beta}(1)\right|=\left\|\psi_{\alpha, \beta}\right\|_{M}$. From the remarks in $\S 5$ we see that

$$
\left\|\psi_{\alpha, \beta}\right\|_{M}=(\operatorname{dim} \sigma)^{1 / 2} \cdot\left\|\Psi_{1 \alpha}\right\| \cdot\left\|\Psi_{2 \beta}\right\| .
$$

However, by Lemma 13 (i) and formula (8.3),

$$
\left\|\Psi_{1 \alpha}\right\|^{2} \cdot\left\|\Psi_{2 \beta}\right\|^{2}<p\left(\left(|\sigma|^{2}+r^{-2}|\Lambda|^{2}\right)^{1 / 2}\right)^{2} \cdot q(|\tau|)^{2}
$$

Now $\operatorname{dim} \tau_{\alpha, \beta}$ and $\operatorname{dim} \sigma$ are bounded by polynomials in $\left|\tau_{\alpha, \beta}\right|$ and $|\sigma|$ respectively, by the Weyl dimension formula. But by Lemma $13,\left|\tau_{\alpha, \beta}\right|, t_{1}$, and $t_{2}$ are all bounded by polynomials in $|\tau|$. The inequality in our lemma then follows from formula (11.1).

Recall the definitions of $\tilde{\omega}^{\mathrm{m}}, L_{1}$, and $L_{1}^{\prime}$ from $\S 2$. Let $\widetilde{L_{1}}$ be the lattice of real linear functionals on $(-1)^{1 / 2} a_{e}$ which is generated by the restrictions of roots of $(\mathfrak{g}, \mathfrak{a})$. Then $\widetilde{L_{1}} \bar{\subset} L_{1}$ and $\widetilde{L_{1}}$ is a lattice of finite index in $L_{1}$. Recall that

$$
q=\frac{1}{2}(\operatorname{dim} \mathfrak{g}-\operatorname{dim} \mathfrak{k}-\operatorname{rank} \mathfrak{g}+\operatorname{rank} \mathfrak{k})=\left\{\begin{array}{l}
\frac{1}{2}\left[P_{+}\right]+1 \quad \text { in Case I, } \\
\frac{1}{2}\left[P_{+}\right] \quad \text { in Case II. }
\end{array}\right.
$$

Lemma 31. There is a $\delta_{1}>0$ such that for any $\mu$ in $L_{1}^{\prime}$ the function

$$
\Lambda^{2 q} \cdot \tilde{\omega}\left(-\mu-i \Lambda \mu_{0}\right)^{-1}
$$

is holomorphic in the region $\left|\Lambda_{I}\right|<\delta_{1}$. In addition, there exists a polynomial $p$, independent of $\mu$, such that for $\mu$ in $L_{1}^{\prime}$ and $\left|\Lambda_{I}\right|<\delta_{1}$,

$$
\left|\Lambda^{2 q} \cdot \tilde{\omega}\left(-\mu-i \Lambda \mu_{0}\right)^{-1}\right| ₹ p(|\Lambda|) .
$$

Proof. We have the formula

$$
\tilde{\omega}\left(-\mu-i \Lambda \mu_{0}\right)^{-1}=\tilde{\omega}^{\mathrm{m}}(-\mu)^{-1} \cdot \prod_{\alpha \in P_{+}}\left\langle-\mu-i \Lambda \mu_{0}, H_{\alpha}\right\rangle^{-1} .
$$

Note that $\Lambda^{2 q} \cdot \tilde{\omega}\left(-\mathcal{C}^{\prime} \mu-i \Lambda \mu_{0}\right)^{-1}$ is a meromorphic function of $\Lambda$.
Since $\mu$ is in $L_{1}^{\prime}, \tilde{\omega}^{\mathrm{m}}(\mu)$ is not equal to zero. The numbers $\left\{\left|\tilde{\omega}^{\mathrm{m}}(-\mu)\right|: \mu \in L_{1}^{\prime}\right\}$ are actually bounded away from zero. This can be seen by an argument similar to that used in the proof of Lemma 3.

If $\alpha$ is in $P_{+}$, the number $\left\langle-\mu-i \Lambda \mu_{0}, H_{\alpha}\right\rangle$ equals zero only if

$$
\Lambda=i \mu\left(H_{\alpha}\right) \cdot \mu_{0}\left(H_{\alpha}\right)^{-1}
$$

It is well known that the numbers $\left\{\tilde{\mu}\left(H_{\alpha}\right): \tilde{\mu} \in \widetilde{L_{1}}\right\}$ generate a lattice in $\mathbf{R}$. Therefore, since $\widetilde{L_{1}}$ is of finite index in $L_{1}$, the numbers

$$
\left\{\mu\left(H_{\alpha}\right) \cdot \mu_{0}\left(H_{\alpha}\right)^{-1}: \mu \in L_{1}\right\}
$$

generate a lattice in $\mathbf{R}$. Let $\varepsilon_{\alpha}$ be the positive generator of this lattice. Put

$$
\delta_{1}=\inf _{\alpha \in P_{+}}\left(\frac{1}{2} \varepsilon_{\alpha}\right)
$$

Then for any $\alpha$ in $P_{+}$and any $\mu$ in $L_{1}^{\prime}$, either

$$
\left|\mu\left(H_{\alpha}\right) \cdot \mu_{0}\left(H_{\alpha}\right)^{-1}\right|=0
$$

or

$$
\left|\mu\left(H_{\alpha}\right) \cdot \mu_{0}\left(H_{\alpha}\right)^{-1}\right| \gg 2 \delta_{1} .
$$

In either case, the function

$$
\Lambda \cdot\left\langle-\mu-i \Lambda \mu_{0}, H_{\alpha}\right\rangle^{-1}
$$

is holomorphic in the region $\left|\Lambda_{I}\right|<\delta_{1}$, and it is bounded independently of $\mu$ by a polynomial in $|\Lambda|$. Our lemma follows from the fact that $\left.2 q \gg P_{+}\right]$.

COROLLARY. If $\mu$ is in $L_{1}^{\prime}$ and $\Lambda$ is not equal to zero, then $\tilde{\omega}\left(-\mu-i \Lambda \mu_{0}\right)$ is not equal to zero.

LEMMA 32. The real parts of the set

$$
\left\{\left\langle-\mu-i \Lambda \mu_{0}, s H_{0}\right\rangle: \mu \in L_{1}, \Lambda \in \mathbf{R}, s \in W\right\}
$$

of complex numbers form a lattice in $\mathbf{R}$.
Proof. The real part of $\left\langle-\mu-i \Lambda \mu_{0}, s H_{0}\right\rangle$ is equal to $-\left\langle s^{-1} \mu, H_{0}\right\rangle$. If $\mu$ is in $\widetilde{L_{1}}$, we can regard $\mu$ as an integral sum of roots of $(\mathfrak{g}, \mathfrak{a})$. Then $s^{-1} \mu$ is also an integral sum of roots of $(\mathfrak{g}, \mathfrak{a})$. If $\alpha$ is a root of $(\mathfrak{g}, \mathfrak{a}), \alpha\left(H_{0}\right)$ is an integer. Therefore $-\left\langle s^{-1} \mu, H_{0}\right\rangle$ is an integer for any $\mu$ in $\widetilde{L_{1}}$. Our lemma follows from the fact that $\widetilde{L_{1}}$ is of finite index in $L_{1}$.

Lemma 30, the corollary to Lemma 31, and Lemma 32 verify the estimate (9.1) and Assumptions 1 and 2 of $\S 9$ for our collection $\left\{E_{\Lambda}\left(\psi_{\sigma}^{\tau}: x\right)\right\}$. We define the function $\theta_{\Lambda}=\theta_{\Lambda}\left(\psi_{\sigma}^{\tau}: m a\right)$ from $M A_{\mathfrak{p}}$ to $V_{\tau}$ as in §9. By Lemma 25, there exists, for every nonnegative integer $n$, polynomials $p_{1}, p_{2}, q$, and a number $\varepsilon>0$, all independent of $\Lambda, \sigma$, and $\tau$, such that for any $t>0$,

$$
\begin{align*}
& \left|\left(\frac{d}{d t}\right)^{n}\left[e^{t \rho\left(H_{0}\right)} E_{\Lambda}\left(\psi_{\sigma}^{\tau}: \exp t H_{0}\right)-\theta_{\Lambda}\left(\psi_{\sigma}^{\tau}: \exp t H_{0}\right)\right]\right|  \tag{11.2}\\
& \quad<p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|) \cdot q(|\tau|) \cdot e^{-\varepsilon t}
\end{align*}
$$

12. The linear transformations $c^{+}(\Lambda)$ and $c^{-}(\Lambda)$. Let us fix $\tau$ in $\mathscr{E}_{K}^{2}$ and $\sigma$ in $\mathscr{E}_{M}$. Let $\mu$ be a real linear functional on $(-1)^{1 / 2} a_{\ell}$ associated with $\sigma$. Fix $\psi$ in $L_{\sigma}^{\tau}$. We shall let $\Lambda \neq 0$ vary in R. Put $\phi_{\Lambda}(x)=E_{\Lambda}(\psi: x)$. For $m$ in $M$, and $a$ in $A_{\mathfrak{p}}$, define $\Phi_{\Lambda}(m a)=\Phi_{\Lambda}(\psi: m a), \Psi_{\Lambda}(m a)=\Psi_{\Lambda}(\psi: m a)$, and $\Theta_{\Lambda}(m a)=\Theta_{\Lambda}(\psi: m a)$ corresponding to the function $\phi_{\Lambda}$ as in $\S 9$.

The projections $\left\{E_{s_{i}}\right\}$ were defined in $\S 9$. Then

$$
\Theta_{\Lambda}(m a)=E_{s_{1}} \Theta_{\Lambda}(m a)+\cdots+E_{s_{r}} \Theta_{\Lambda}(m a)
$$

Let $\lambda$ be the linear functional $-\mu-i \Lambda \mu_{0}$. If $\varsigma$ is in $\mathcal{Z}_{1}$, then by Lemma 17

$$
\begin{equation*}
\Gamma(\zeta)\left(E_{s_{i}} \Theta_{\Lambda}\right)(m a)=\left\langle\gamma_{1}(\varsigma), s_{i} \lambda\right\rangle\left(E_{s_{i}} \Theta_{\Lambda}\right)(m a) \tag{12.1}
\end{equation*}
$$

Recall that $\bar{\tau}$ was the restriction of $\tau$ to $M$. By (9.8), $\Theta_{\Lambda}$ is $\bar{\tau}$-spherical. Let $\Theta_{\Lambda, i}=E_{s_{i}} \Theta_{\Lambda}$. Since the actions of $\bar{\tau}(M)$ and $\Gamma\left(\mathcal{Z}_{1}\right)$ on the vector space $\mathscr{V}_{\tau}$ commute, and since the $\Theta_{\Lambda, i}$ are eigenvectors of $\Gamma(\varsigma)$, each $\Theta_{\Lambda, i}$ is a $\bar{\tau}$-spherical function.

Now suppose for some $s_{i}$ that the linear functional $s_{i} \mu_{0}$ does not vanish on $\mathfrak{a}_{\mathfrak{t}}$. We will show that $\Theta_{\Lambda, i}=0$. If $\zeta_{M}$ is in $\mathcal{Z}_{M}$, and $m$ is in $M$, by (9.9) and (12.1) we see that

$$
\begin{equation*}
\Theta_{\Lambda, i}\left(m ; s_{M}\right)=\Gamma\left(\varsigma_{M}\right) \Theta_{\Lambda, i}(m)=\left\langle\gamma_{1}\left(s_{M}\right), s_{i} \lambda\right\rangle \Theta_{\Lambda, i}(m) \tag{12.2}
\end{equation*}
$$

Thus if $\Theta_{\Lambda, i}$ is regarded as a function on the compact group $M$, it is an eigenfunction of $3_{M}$. The infinitesimal character corresponds to the restriction of the linear functional $s_{i} \lambda=-s_{i}\left(\mu+i \Lambda \mu_{0}\right)$ to $\mathfrak{a}_{\mathfrak{f}}$. Now $i \Lambda\left(s_{i} \mu_{0}\right)$ is real-valued on $\mathfrak{a}_{\mathfrak{f}}$, so that $s_{i} \lambda$ is not purely imaginary on $\mathfrak{a}_{\mathfrak{k}}$. However, it is well known that the eigenfunctions of a compact Lie group have eigenvalues corresponding to purely imaginary linear functionals on a Cartan subalgebra. Therefore $\Theta_{\Lambda, i}$ must be zero.

Then we can write

$$
\Theta_{\Lambda}=\Theta_{\Lambda}^{+}+\Theta_{\Lambda}^{-} .
$$

$\Theta_{\Lambda}^{+}$and $\Theta_{\Lambda}^{-}$are the sums of those $\Theta_{\Lambda, i}$ for which the Weyl group element $s_{i}$, when restricted to $\mathfrak{a}_{\mathfrak{p}}$, is respectively the identity or reflection about 0 . Now by (9.7),

$$
\Theta_{\Lambda, i}\left(m \exp t H_{0}\right)=e^{t \Gamma\left(H_{0}\right)} \Theta_{\Lambda, i}(m)
$$

Therefore

$$
\begin{aligned}
& \Theta_{\Lambda}^{+}\left(m \exp t H_{0}\right)=e^{i \Lambda t} \Theta_{\Lambda}^{+}(m) \\
& \Theta_{\Lambda}^{-}\left(m \exp t H_{0}\right)=e^{-i \Lambda t} \Theta_{\Lambda}^{-}(m)
\end{aligned}
$$

In the notation of $\S 9, \theta_{\Lambda}(m a)=t_{1}\left(\Theta_{\Lambda}(m a)\right)$. Let $\theta_{\Lambda}^{+}(m a)=t_{1}\left(\Theta_{\Lambda}^{+}(m a)\right)$, $\theta_{\Lambda}^{-}(m a)=t_{1}\left(\Theta_{\Lambda}^{-}(m a)\right)$. Then

$$
\theta_{\Lambda}\left(m \exp t H_{0}\right)=e^{i \Lambda t} \theta_{\Lambda}^{+}(m)+e^{-i \Lambda t} \theta_{\Lambda}^{-}(m)
$$

Since the functions $\theta_{\Lambda}^{ \pm}(m)$ are $\bar{T}$-spherical, so are the functions $\theta_{\Lambda}^{ \pm}(m)$. This means that the functions $\theta^{ \pm}(m)$ are in $L^{\tau}$. $\theta_{\Lambda}^{+}$depend linearly on the $\psi$ in $L_{\sigma}^{\tau}$ that we fixed at the beginning of the section. We write this as $\theta_{\Lambda}^{ \pm}=c^{ \pm}(\Lambda) \psi$. $c^{+}(\Lambda)$ and $c^{-}(\Lambda)$ are linear transformations from $L_{\sigma}^{\tau}$ into $L^{\tau}$. Since $L^{\top}$ is an orthogonal direct sum of spaces $L_{\sigma}^{\tau}$, we can extend the definition of $c^{+}(\Lambda)$ and $c^{-}(\Lambda)$ to all of $L^{\tau}$. Then $c^{ \pm}(\Lambda)$ are linear transformations of $L^{\tau}$ into itself.

For $m$ in $M$ and $t$ in $\mathbf{R}$ we have the formula

$$
\begin{equation*}
\theta_{\Lambda}\left(\psi: m \exp t H_{0}\right)=e^{i \Lambda t}\left(c^{+}(\Lambda) \psi\right)(m)+e^{-i \Lambda t}\left(c^{-}(\Lambda) \psi\right)(m) . \tag{12.3}
\end{equation*}
$$

Suppose that the restriction of some $s_{i}$ to $\mathfrak{a}_{\mathfrak{p}}$ is the identity. Then $\mathfrak{a}_{\mathfrak{e}}$ is an invariant subspace of $s_{i}$. It is known that $s_{i}$ is in $W_{1}$. Since $\left\{s_{1}=1, s_{2}, \ldots, s_{r}\right\}$ is a set of representatives of cosets of $W_{1}$ in $W, s_{i}=s_{1}=1$. On the other hand, if the restriction of an element $s_{j}$ to $\mathfrak{a}_{\mathfrak{p}}$ is a reflection, we can represent $s_{j}$ as the nontrivial element in the group $M^{\prime} / M$ defined in $\S 2$. Therefore there is only one such $s_{j}$, which we shall denote by $s_{2}$. Therefore, if $\psi$ is any vector in $L_{\sigma}^{\tau}$,

$$
\begin{aligned}
& \left(c^{+}(\Lambda) \psi\right)(m)=t_{1}\left(\Theta_{\Lambda, 1}(m)\right), \\
& \left(c^{-}(\Lambda) \psi\right)(m)=t_{1}\left(\Theta_{\Lambda, 2}(m)\right) .
\end{aligned}
$$

Suppose that $\psi$ is an eigenfunction of $3_{M}$ in $L^{\tau}$ with infinitesimal character $\chi_{\mu}^{M}$, for some $\mu$ in $L_{1}^{\prime}$. Then from (12.2), we see that the infinitesimal characters of $c^{+}(\Lambda) \psi$ and $c^{-}(\Lambda) \psi$ are $\chi_{\mu}^{M}$ and $\chi_{\mu^{\prime}}^{M}$, respectively. Here $\mu^{\prime}=s_{2} \mu$.

For any $\mu$ in $L_{1}^{\prime}$ let $\mathscr{E}_{M}(\mu)$ be the set of all $\sigma$ in $\mathscr{E}_{M}$ such that the linear functional associated to either $\sigma$ or $\sigma^{\prime}$ is $\mu$. It is clear that $|\sigma|=\left|\sigma^{\prime}\right|$ for any $\sigma$ in $\mathscr{E}_{M}$. Therefore, for any $\sigma_{1}, \sigma_{2}$ in $\mathscr{E}_{M}(\mu)$,

$$
\begin{equation*}
\left|\sigma_{1}\right|=\left|\sigma_{2}\right| \tag{12.4}
\end{equation*}
$$

Let $L_{\mu}^{\tau}$ be the direct sum of all $L_{\sigma}^{\tau}$ for which $\sigma$ is in $\mathscr{C}_{M}(\mu)$. From the above discussion we see that $L_{\mu}^{\tau}$ is an invariant subspace of the linear transformations $c^{ \pm}(\Lambda)$.

We wish to prove some estimates for $c^{+}(\Lambda)$ and $c^{-}(\Lambda)$. We need to examine the functions

$$
t_{i}\left(\Psi_{\Lambda}\right): M A_{\mathfrak{p}} \rightarrow V_{\tau} .
$$

From the definition in $\S 9$,

$$
\begin{equation*}
t_{i}\left(\Psi_{\Lambda}(\psi: m a)\right)=d(m a) E_{\Lambda}\left(\psi: m a ; u_{i}\left(H_{0}\right)^{\prime}\right) . \tag{12.5}
\end{equation*}
$$

We shall allow $\Lambda$ to assume complex values.
Lemma 33. Suppose $\psi$ is in $L_{\sigma}^{\tau}$, and $\|\psi\|_{M}=1$. Then for fixed $t>0$, $t_{i}\left(\Psi_{\Lambda}\left(\psi: \exp t H_{0}\right)\right)$ is an entire function in $\Lambda$. Furthermore, there exist polynomials $p_{1}, p_{2}$, and $q$ such that for $\Lambda$ in $\mathbf{C}$ and $t>0$,

$$
\left|t_{i}\left(\Psi_{\Lambda}\left(\psi: \exp t H_{0}\right)\right)\right| \gtrless e^{-t\left(1-\left|\Lambda_{I}\right|\right)} \cdot(1+r t)^{d} \cdot p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|) \cdot q(|\tau|) .
$$

Here $r^{2}=B\left(H_{0}, H_{0}\right)$.

Proof. It is clear from (12.5) and Lemma 30 that $t_{i}\left(\Psi_{\Lambda}\left(\psi: \exp t H_{0}\right)\right)$ is an entire function of $\Lambda$.

Now, using the notation of the proof of Lemma 16 , we see $\left|t_{i}\left(\Psi_{\Lambda}\left(\exp t H_{0}\right)\right)\right|$ equals

$$
\left|\sum_{j} e^{t \rho\left(H_{0}\right)} E_{\Lambda}\left(\psi: \exp t H_{0} ; \theta\left(N_{i j}\right) g_{i j} \eta_{j}^{\prime}\right)\right|
$$

If $t>0$, this expression is bounded by

$$
e^{-t}\left|e^{\rho\left(t H_{0}\right)} \sum_{j} E_{\Lambda}\left(\psi: \theta\left(N_{i j}\right) ; \exp t H_{0} ; g_{i j} \eta_{j}^{\prime}\right)\right|
$$

Now by Lemma 30 there are polynomials $p_{1}, p_{2}$, and $q$ such that for any $t>0$

$$
\begin{aligned}
& \left|\sum_{j} E_{\Lambda}\left(\psi: \theta\left(N_{i j}\right) ; \exp t H_{0} ; g_{i j} \eta_{j}^{\prime}\right)\right| \\
& \quad ₹ p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|) \cdot q(|\tau|) \cdot e^{\left|\Lambda_{I}\right| t} \cdot \Xi\left(\exp t H_{0}\right) \\
& \quad<p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|) \cdot q(|\tau|) \cdot e^{\left|\Lambda_{I}\right| t} \cdot e^{-t \rho\left(H_{0}\right)} \cdot(1+r t)^{d}
\end{aligned}
$$

(this last inequality follows from (4.1)). Therefore $\left|t_{i}\left(\Psi_{\Lambda}\left(\exp t H_{0}\right)\right)\right|$ is bounded by

$$
e^{-\left(1-\left|\Lambda_{I}\right| t\right)} \cdot(1+r t)^{d} \cdot p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|) \cdot q(|\tau|)
$$

Let $q$ be the integer in Lemma 31.
Lemma 34. Let $\delta$ be the minimum of $\delta_{1}$ and $\frac{1}{4}$, where $\delta_{1}$ is the positive constant in Lemma 31. Choose $\psi$ in $L_{\sigma}^{\tau}$ with $\|\psi\|_{M}=1$. Then if $1 \gtrless i, j ₹ r$, the function $\Lambda^{2 q} \cdot t_{i}\left(\Theta_{\Lambda, j}(\psi: 1)\right)$ is analytic in the region $\left|\Lambda_{I}\right|<\delta$. In addition, there exist polynomials $p_{1}, p_{2}$, and $q$ such that in the region $\left|\Lambda_{I}\right|<\delta$,

$$
\left|\Lambda^{2 q} \cdot t_{i}\left(\Theta_{\Lambda, j}(\psi: 1)\right)\right|<p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|) \cdot q(|\tau|)
$$

Proof. We have the formula

$$
\Theta_{\Lambda}(1)=E^{0} \Phi_{\Lambda}(1)+\int_{0}^{\infty} e^{-t \Gamma\left(H_{0}\right)} E^{0} \Psi_{\Lambda}\left(\exp t H_{0}\right) d t
$$

We assume that $\Theta_{\Lambda, j}$ is not zero, so by our earlier remarks, $j=1$ or 2 . Then $s_{j} H_{0}=\varepsilon_{j} H_{0}$ where $\varepsilon_{j}=1$ or $-1 . E_{s_{j}}$ is the projection of $\mathscr{V}_{\tau}$ onto the eigenspace of $\Gamma\left(H_{0}\right)$ corresponding to the eigenvalue $\left\langle s_{j}\left(\mu+i \Lambda \mu_{0}\right), H_{0}\right\rangle$ where $\mu$ is a linear functional in $L_{1}^{\prime}$ associated with $\sigma$. This eigenvalue equals $i \varepsilon_{j} \Lambda$, which is purely imaginary. Therefore $E_{s_{j}} E^{0}=E_{s_{j}}$, so that

$$
\Lambda^{2 q} \Theta_{\Lambda, j}(1)=\Lambda^{2 q} E_{s_{j}} \Phi_{\Lambda}(1)+\Lambda^{2 q} \int_{0}^{\infty} e^{-i \varepsilon_{j} \Lambda t} E_{s_{j}} \Psi_{\Lambda}\left(\exp t H_{0}\right) d t
$$

Now if $1 \gtrless l$ ₹ $r$

$$
t_{l}\left(\Phi_{\Lambda}(1)\right)=E_{\Lambda}\left(\chi: 1 ; \eta_{l}^{\prime}\right)
$$

By Lemma 30 there are polynomials $p_{1}, p_{2}, q$ such that

$$
\begin{equation*}
\left|E_{\Lambda}\left(\psi: 1 ; \eta_{l}^{\prime}\right)\right|<p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|) \cdot q(|\tau|) \tag{12.6}
\end{equation*}
$$

By the last lemma we can also choose polynomials $p_{1}, p_{2}, q$ such that for any $t>0$,

$$
\begin{align*}
& \left|e^{-i \varepsilon_{j} \Lambda t} t_{l}\left(\Psi_{\Lambda}\left(\exp t H_{0}\right)\right)\right| \\
& \quad<p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|) \cdot q(|\tau|) \cdot e^{-t\left(1-2\left|\Lambda_{I}\right|\right)} \cdot(1+r t)^{d} \tag{12.7}
\end{align*}
$$

Therefore, the integral

$$
\int_{0}^{\infty} e^{-i \varepsilon_{j} \Lambda t} t_{l}\left(\Psi_{\Lambda}\left(\exp t H_{0}\right)\right) d t
$$

converges uniformly for $\Lambda$ in compact subsets of $\left|\Lambda_{I}\right|<\frac{1}{4}$ and so is an analytic function of $\Lambda$ in this region. Furthermore, by Lemmas 18 and 31, the matrices of the linear transformations $\Lambda^{2 q} E_{s_{j}}$ relative to the basis $\left\{\eta_{1}^{* *}, \ldots, \eta_{r}^{* *}\right\}$ have components which are analytic in $\left|\Lambda_{I}\right|<\delta_{1}$ and bounded by a polynomial in $\Lambda$ in this region. Therefore, the functions $\Lambda^{2 q} t_{i}\left(\Theta_{\Lambda, j}(1)\right)$ are analytic in the region $\left|\Lambda_{I}\right|<\delta$. We see from (12.6) and (12.7) that it is possible to choose polynomials $p_{1}, p_{2}$, and $q$ such that for $\left|\Lambda_{I}\right|<\delta$

$$
\left|\Lambda^{2 q} t_{i}\left(\Theta_{\Lambda, j}(1)\right)\right| \gtrless p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|) \cdot q(|\tau|)
$$

Corollary. Choose $\psi$ in $L_{\sigma}^{\tau}$ with $\|\psi\|_{M}=1$. Then the functions $\Lambda^{2 q} c^{ \pm}(\Lambda) \psi(1)$ are analytic in the region $\left|\Lambda_{I}\right|<\delta$. Furthermore, there are polynomials $p_{1}, p_{2}$, and $q$, independent of $\sigma, \Lambda$, and $\tau$ such that for $\left|\Lambda_{I}\right|<\delta$

$$
\left|\Lambda^{2 q} c^{ \pm}(\Lambda) \psi(1)\right|<p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|) \cdot q(|\tau|)
$$

Proof. The corollary follows from the lemma if we recall that the functions $c^{+}(\Lambda) \psi(1)$ and $c^{-}(\Lambda) \psi(1)$ equal $t_{1}\left(\Theta_{\Lambda, 1}(1)\right)$ and $t_{1}\left(\Theta_{\Lambda, 2}(1)\right)$ respectively.

Therefore, $c^{ \pm}(\Lambda) \psi(1)$ is meromorphic in $\left|\Lambda_{I}\right|<\delta$, with the only possible pole being at $\Lambda=0$. Assume it has a pole of order $N^{ \pm}(\psi)$ at $\Lambda=0$. Let us agree to write $N^{ \pm}(\psi)=0$ if $c^{ \pm}(\Lambda) \psi$ has no poles at $\Lambda=0$. Then

$$
0 \gtrless N^{ \pm}(\psi) \gtrless 2 q
$$

The function

$$
g^{ \pm}(\psi: \Lambda)=\Lambda^{N^{ \pm}(\psi)} c^{ \pm}(\Lambda) \psi(1)
$$

is holomorphic in the region $\left|\Lambda_{I}\right|<\delta$. For any $\Lambda$ in this region we can write

$$
\begin{equation*}
g^{ \pm}(\psi: \varsigma)=a_{0}^{ \pm}+a_{1}^{ \pm}(\varsigma-\Lambda)+a_{2}^{ \pm}(\varsigma-\Lambda)^{2}+\cdots \tag{12.8}
\end{equation*}
$$

where $\varsigma$ remains in some neighborhood of $\Lambda$, and $a_{0}^{ \pm}, a_{1}^{ \pm}, a_{2}^{ \pm}, \ldots$ are vectors in $V_{\tau}$.

LEMMA 35. For every nonnegative integer $n$ there are polynomials $p_{1}, p_{2}, q$ such that in the region $\left|\Lambda_{I}\right|<\delta / 2$,

$$
\left|\left(\frac{d}{d \Lambda}\right)^{n} g^{ \pm}(\psi: \Lambda)\right|<p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|) \cdot q(|\tau|)
$$

Proof. From (12.8), it is easy to show that

$$
\begin{aligned}
& \left|\left(\frac{d}{d \Lambda}\right)^{n} g^{ \pm}(\psi: \Lambda)\right|=\left|n!(2 \pi i)^{-1} \cdot \int_{\Gamma} g^{ \pm}(\psi: \xi)(\xi-\Lambda)^{-(n+1)} d \xi\right| \\
& =n!(2 \pi)^{-1}\left|\int_{\Gamma} \xi^{2 q} c^{ \pm}(\xi) \psi(1) \cdot(\xi-\Lambda)^{-(n+1)} \cdot \xi^{-2 q+N^{ \pm}(x)} d \xi\right|
\end{aligned}
$$

where $\Gamma$ is any curve in $\left|\xi_{I}\right|<\delta$ that winds around $\Lambda$ once. Choose $\Gamma$ such that for any $\xi$ on $\Gamma$
(i) $\delta / 2 \gtrless|\xi-\Lambda|<1$,
(ii) $\delta / 2<|\xi|$.

This is clearly possible since $\delta<\frac{1}{4}$. Lemma 35 then follows from the corollary to Lemma 34. 미
13. Relation between $c^{ \pm}(\Lambda)$ and $\beta(\sigma, \Lambda)$. Let $\tau=\left(\tau_{1}, \tau_{2}\right)$ be an irreducible unitary double representation of $K$ on the finite-dimensional Hilbert space $V_{\tau}$. Let $\mathscr{E}$ be the finite-dimensional vector space of endomorphisms of $V_{\tau}$. Let $\mathscr{K}$ and $\mathscr{A}_{p}$ be the universal enveloping algebras of $\mathfrak{k}_{c}$ and $\mathfrak{a}_{p c}$ respectively.

Let $\overline{\mathscr{B}}=\mathscr{B} \otimes \mathscr{E} . \mathscr{B}$ acts on $C^{\infty}(G)$ by left invariant differentiation, so there is a natural action of $\overline{\mathscr{B}}$ on $C^{\infty}(G) \otimes V_{\tau}$.

Fix $h$ in $A_{\mathfrak{p}}$ such that $h$ is not equal to the identity. Define

$$
\begin{aligned}
& \overline{\mathfrak{k}}_{2}=\left\{X \otimes 1-1 \otimes \tau_{2}(X): X \in \mathfrak{k}_{\mathbf{c}}\right\} \\
& \overline{\mathfrak{k}}_{1}=\left\{\operatorname{Ad}\left(h^{-1}\right) X \otimes 1-1 \otimes \tau_{1}(X): X \in \mathfrak{k}_{c}\right\} .
\end{aligned}
$$

$\mathfrak{k}_{1}$ and $\mathfrak{k}_{2}$ are vector subspaces of $\overline{\mathscr{B}}$.
The following lemma generalizes [ $\mathbf{3}(\mathrm{g})$, Lemma 21, Corollary 2].
Lemma 36. $\overline{\mathscr{B}}$ is the direct sum of $\overline{\mathfrak{k}}_{1} \overline{\mathscr{B}}+\overline{\mathscr{B}} \overline{\mathfrak{k}}_{2}$ with $\mathfrak{A}_{\mathfrak{p}} \otimes \mathscr{E}$.
PROOF. If $\rho$ denotes conjugation of $\mathfrak{g}_{c}$ with respect to the compact real form $\mathfrak{u}=\mathfrak{k}+i \mathfrak{p}$, then the inner product

$$
-B(X, \rho Y), \quad X, Y \in \mathfrak{g}_{\mathbf{c}}
$$

converts $\mathfrak{g}_{\mathrm{c}}$ into a Hilbert space. Let $\mathfrak{q}$ be the orthogonal complement in $\mathfrak{g}_{\mathrm{c}}$ of $\mathfrak{m}_{\mathfrak{c}}+\mathfrak{a}_{\mathfrak{p c}}$ with respect to this inner product. Let $\mathfrak{q}_{\mathfrak{t}}=\mathfrak{q} \cap \mathfrak{k}_{c}$, and let $S\left(\mathfrak{q}_{\mathfrak{l}}\right)$ be the symmetric algebra on $\mathfrak{q}_{\mathfrak{k}}$. Let $\mathfrak{Q}_{\mathfrak{k}}$ be the image of $S\left(\mathfrak{q}^{\mathfrak{k}}\right)$ in $\mathscr{B}$ under the canonical mapping. (The canonical mapping is defined in [3(a), p. 192]). It is proved in [ $\mathbf{3}(\mathbf{g})$, Lemma 21, Corollary 1] that every element in $\mathscr{B}$ has a unique representation as an element of the form

$$
\sum_{i} q_{i} a_{i} k_{i}, \quad q_{i} \in \mathfrak{O}_{\mathfrak{k}}^{h^{-1}}, a_{i} \in \mathfrak{X}_{\mathfrak{p}}, k_{i} \in \mathscr{K}
$$

(The map $b \rightarrow b^{h}$, for $b$ in $\mathscr{B}$, is the automorphism of $\mathscr{B}$ which coincides with $\operatorname{Ad}(h)$ on $\mathfrak{g}$.) We shall then write $\mathscr{B}=\mathfrak{Q}_{\mathfrak{k}}^{h^{-1}} \mathfrak{A}_{\mathfrak{p}} \mathscr{K}$. Therefore

$$
\overline{\mathscr{B}}=\mathfrak{Q}_{\mathrm{z}}^{h^{-1}} \mathfrak{A}_{\mathfrak{p}} \mathscr{K} \otimes \mathscr{E}
$$

Let $\bar{q}_{\mathfrak{k}}=\left\{\operatorname{Ad}\left(h^{-1}\right) X \otimes 1-1 \otimes \tau_{1}(X): X \in \mathfrak{q}_{\mathfrak{k}}\right\}$. Then it is easy to see that $\overline{\mathscr{B}}$ is the direct sum of $\overline{\mathfrak{q}}_{\mathrm{k}} \overline{\mathscr{B}}+\overline{\mathscr{B}} \mathfrak{l}_{2}$ with $\mathfrak{A}_{\mathfrak{p}} \otimes \mathscr{E}$. However, in [ $\mathbf{3}(\mathrm{g})$, Lemma 21, Corollary 2] it is shown that

$$
\mathfrak{k}^{h^{-1}} \mathscr{B}+\mathscr{B} \mathfrak{k}=\mathfrak{q}^{h^{-1}} \mathscr{B}+\mathscr{B} \mathfrak{k}
$$

Therefore

$$
\left(\mathfrak{k}^{h^{-1}} \mathscr{B}+\mathscr{B} \mathfrak{k}\right) \otimes \mathscr{E}=\left(\mathfrak{q}_{\mathfrak{k}}^{h^{-1}} \mathscr{B}+\mathscr{B} \mathfrak{k}\right) \otimes \mathscr{E}
$$

It then follows that

$$
\overline{\mathfrak{q}}_{\mathfrak{k}} \overline{\mathscr{B}}+\overline{\mathscr{B}} \overline{\mathfrak{k}}_{2}=\overline{\mathfrak{k}}_{1} \overline{\mathscr{B}}+\overline{\mathscr{B}} \overline{\mathfrak{k}}_{2}
$$

This proves the lemma.
COROLLARY. For any $b$ in $\mathscr{B}$ there is a unique element $\delta_{h}^{\prime}(b)$ in $\mathscr{A}_{p} \otimes \mathscr{E}$ such that $b-\delta_{h}^{\prime}(b)$ is in the vector space $\overline{\mathfrak{k}}_{1} \overline{\mathscr{B}}+\overline{\mathscr{B}}_{2}$.

LEMMA 37. Let $\phi: G \rightarrow V_{\tau}$ be an infinitely differentiable $\tau$-spherical function. Then for any $b$ in $\overline{\mathscr{B}}$ and any $h$ in $A_{\mathfrak{p}}, h$ not equal to the identity, we have the formula

$$
\phi(h ; b)=\phi\left(h ; \delta_{h}^{\prime}(b)\right)
$$

(It is clear what the notation $\phi\left(h ; \delta_{h}^{\prime}(b)\right)$ means.)
PROOF. If $b_{2}$ is in $\overline{\mathscr{B}}_{2}$, then it is clear that $\phi\left(h ; b_{2}\right)=0$. If $b_{1}$ is in $\overline{\mathfrak{k}}_{1} \overline{\mathscr{F}}$, we shall assume that

$$
b_{1}=\left(\operatorname{Ad}\left(h^{-1}\right) X \otimes 1-1 \otimes \tau_{1}(X)\right) B_{1}
$$

for some $X$ in $\mathfrak{k}$ and some $B_{1}$ in $\overline{\mathfrak{B}}$. Then

$$
\phi\left(h ; b_{1}\right)=\phi\left(X ; h ; B_{1}\right)-\tau_{1}(X) \phi\left(h ; B_{1}\right)
$$

This last term is equal to zero, since $\phi(k h)=\tau_{1}(k) \phi(h)$ for any $k$ in $K$. This is enough to prove the lemma.

If $g$ is in $\mathscr{B}$, identify $g$ with the element $g \otimes 1$ in $\overline{\mathscr{B}}$. If $T$ is in $\mathscr{E}$, identify $T$ with the element $1 \otimes T$ in $\overline{\mathscr{B}}$.

If $\lambda$ is a real linear functional on the vector space $(-1)^{1 / 2} \mathfrak{a}_{\mathfrak{t}}+\mathfrak{a}_{\mathrm{p}}$, define an element $H_{\lambda}$ in $(-1)^{1 / 2} \mathfrak{a}_{\mathfrak{p}}+\mathfrak{a}_{\mathfrak{p}}$ by the property

$$
B\left(H_{\lambda}, H\right)=\lambda(H), \quad H \in(-1)^{1 / 2} \mathfrak{a}_{\mathfrak{f}}+\mathfrak{a}_{\mathfrak{p}}
$$

Also, write $\bar{\lambda}$ for that linear functional which equals $\lambda$ on $\mathfrak{a}_{\mathfrak{p}}$ and equals zero on $(-1)^{1 / 2} \mathfrak{a}_{\mathrm{f}}$.

Define the vectors $H_{1}, \ldots, H_{n}, X_{\alpha}, X_{-\alpha}$ in $\mathfrak{g}_{c}$ and the Casimir operator $\omega_{\mathfrak{g}}$ as in $\S 5$. Since $\theta$ is an automorphism of $\mathfrak{g}_{\mathbf{c}}, \theta\left(\omega_{\mathfrak{g}}\right)$ equals $\omega_{\mathfrak{g}}$.

Fix a nonzero $H$ in $\mathfrak{a}_{\mathfrak{p}}$ and let $h=\exp H$. We would like to compute $\delta_{h}^{\prime}\left(\omega_{\mathfrak{g}}\right)$. The following lemma will express $\delta_{h}^{\prime}\left(\omega_{\mathfrak{g}}\right)$ as a linear second-order differentiable operator on $A_{\mathfrak{p}}^{\prime}$. The coefficients of this operator will be selfadjoint operators in $\mathscr{E}$ which depend on $h$. The lemma is a generalization of $[\mathbf{3}(\mathrm{g})$, Lemma 27].

Lemma 38. There is a selfadjoint operator $Q(H)$ in $\mathscr{E}$ such that

$$
\delta_{h}^{\prime}\left(\omega_{\mathfrak{g}}\right)=H_{1}^{2}+\sum_{\alpha \in P_{+}} \operatorname{coth} \alpha(H) \cdot H_{\bar{\alpha}}+Q(H) \cdot I
$$

Proof. $\omega_{\mathfrak{g}}$ equals $\frac{1}{2}\left(\omega_{\mathfrak{g}}+\theta\left(\omega_{\mathfrak{g}}\right)\right)$. This is then equal to the expression

$$
H_{1}^{2}+\cdots+H_{n}^{2}+\frac{1}{2} \sum_{\alpha \in P}\left(X_{\alpha} X_{-\alpha}+X_{-\alpha} X_{\alpha}\right)+\frac{1}{2} \sum_{\alpha \in P} \theta\left(X_{\alpha} X_{-\alpha}+X_{-\alpha} X_{\alpha}\right)
$$

If $A$ and $B$ are in $\overline{\mathscr{B}}$ we shall say $A$ is congruent to $B$ and write $A \equiv B$ if $A-B$ is in the vector space $\overline{\mathscr{B}}_{2}+\overline{\mathfrak{k}}_{1} \overline{\mathscr{B}}$. If $i$ is greater than $1, H_{i}$ is in $(-1)^{1 / 2} \mathfrak{k}$ so that $H_{i}^{2} \equiv \tau_{2}\left(H_{i}\right)^{2}$. This implies that $\tau_{2}\left(H_{i}\right)$ is a selfadjoint operator on $V_{\tau}$, since $\tau_{2}$ is a unitary representation. For any $\alpha$ in $P$ let $R_{\alpha}$ be the expression

$$
\frac{1}{2}\left(X_{\alpha} X_{-\alpha}+\theta\left(X_{-\alpha} X_{\alpha}\right)\right)+\frac{1}{2}\left(X_{-\alpha} X_{\alpha}+\theta\left(X_{\alpha} X_{-\alpha}\right)\right)
$$

Fix $\alpha$ such that either $\alpha$ or $-\alpha$ is in $P_{+}$. Let $X_{\alpha}=Y_{\alpha}+Z_{\alpha}$, where $Y_{\alpha}$ equals $\frac{1}{2}\left(X_{\alpha}-\theta\left(X_{\alpha}\right)\right)$, a vector in $\mathfrak{p}_{\mathbf{c}}$, and $Z_{\alpha}$ equals $\frac{1}{2}\left(X_{\alpha}+\theta\left(X_{\alpha}\right)\right)$, a vector in $\mathfrak{k}_{\mathbf{c}}$. Then

$$
\begin{aligned}
Z_{\alpha}^{h^{-1}} & =\frac{1}{2}\left(X_{\alpha}^{h^{-1}}+\theta\left(X_{\alpha}^{h}\right)\right)=X_{\alpha}\left(\frac{1}{2} e^{-\alpha(H)}\right)+\theta\left(X_{\alpha}\right)\left(\frac{1}{2} e^{\alpha(H)}\right) \\
& =\frac{1}{2} Y_{\alpha}\left(e^{-\alpha(H)}-e^{\alpha(H)}\right)+\frac{1}{2} X_{\alpha}\left(e^{-\alpha(H)}+e^{\alpha(H)}\right) .
\end{aligned}
$$

Therefore we have the formula

$$
\begin{equation*}
Y_{\alpha}=\operatorname{coth} \alpha(H) \cdot Z_{\alpha}-\operatorname{csch} \alpha(H) \cdot Z_{\alpha}^{h^{-1}} \tag{13.1}
\end{equation*}
$$

It follows that

$$
\begin{aligned}
X_{\alpha} X_{-\alpha} & =\left(Y_{\alpha}+Z_{\alpha}\right)\left(Y_{-\alpha}+Z_{-\alpha}\right) \\
= & {\left[(1+\operatorname{coth} \alpha(H)) Z_{\alpha}-\operatorname{csch} \alpha(H) \cdot Z_{\alpha}^{h^{-1}}\right]\left[Y_{-\alpha}+Z_{-\alpha}\right] } \\
\equiv & {\left[(1+\operatorname{coth} \alpha(H)) Z_{\alpha}-\operatorname{csch} \alpha(H) \cdot \tau_{1}\left(Z_{\alpha}\right)\right]\left[Y_{-\alpha}+\tau_{2}\left(Z_{-\alpha}\right)\right] } \\
= & (1+\operatorname{coth} \alpha(H)) Z_{\alpha} Y_{-\alpha}+(1+\operatorname{coth} \alpha(H)) Z_{\alpha} \cdot \tau_{2}\left(Z_{-\alpha}\right) \\
& -\operatorname{csch} \alpha(H) \cdot Y_{-\alpha} \cdot \tau_{1}\left(Z_{\alpha}\right)-\operatorname{csch} \alpha(H) \cdot \tau_{1}\left(Z_{\alpha}\right) \tau_{2}\left(Z_{-\alpha}\right) \\
= & (1+\operatorname{coth} \alpha(H))\left[Z_{\alpha}, Y_{-\alpha}\right]+(1+\operatorname{coth} \alpha(H)) Y_{-\alpha} \cdot \tau_{2}\left(Z_{\alpha}\right) \\
& +(1+\operatorname{coth} \alpha(H)) \tau_{2}\left(Z_{-\alpha}\right) \tau_{2}\left(Z_{\alpha}\right)-\operatorname{csch} \alpha(H) \cdot Y_{-\alpha} \cdot \tau_{1}\left(Z_{\alpha}\right) \\
& -\operatorname{csch} \alpha(H) \tau_{1}\left(Z_{\alpha}\right) \tau_{2}\left(Z_{-\alpha}\right) \\
= & (1+\operatorname{coth} \alpha(H))\left[Z_{\alpha}, Y_{-\alpha}\right] \\
& +Y_{-\alpha}\left[(1+\operatorname{coth} \alpha(H)) \tau_{2}\left(Z_{\alpha}\right)-\operatorname{csch} \alpha(H) \cdot \tau_{1}\left(Z_{\alpha}\right)\right] \\
& +(1+\operatorname{coth} \alpha(H)) \tau_{2}\left(Z_{-\alpha}\right) \tau_{2}\left(Z_{\alpha}\right)-\operatorname{csch} \alpha(H) \cdot \tau_{1}\left(Z_{\alpha}\right) \tau_{2}\left(Z_{-\alpha}\right) .
\end{aligned}
$$

We let the Cartan involution $\theta$ act on $\overline{\mathscr{B}}$ by making it act on $\mathscr{B}$ in the usual way and letting it act on $\mathscr{E}$ trivially. Notice that $\theta\left[Z_{\alpha}, Y_{-\alpha}\right]=-\left[Z_{\alpha}, Y_{-\alpha}\right]$, and $\theta\left(Y_{-\alpha}\right)=-Y_{-\alpha}$. If we let $\theta$ act on the above congruence, we obtain a new congruence with respect to the space
where

$$
\overline{\mathfrak{k}}_{1}^{\prime}=\left\{X^{h} \otimes 1-1 \otimes \tau_{1}(X): X \in \mathfrak{k}_{c}\right\}
$$

It follows that

$$
\begin{aligned}
\theta\left(X_{\alpha} X_{-\alpha}\right) \equiv & -(1+\operatorname{coth} \alpha(H))\left[Z_{\alpha}, Y_{-\alpha}\right] \\
& -Y_{-\alpha}\left[(1+\operatorname{coth} \alpha(H)) \tau_{2}\left(Z_{\alpha}\right)-\operatorname{csch} \alpha(H) \cdot \tau_{1}\left(Z_{\alpha}\right)\right] \\
& +(1+\operatorname{coth} \alpha(H)) \tau_{2}\left(Z_{-\alpha}\right) \tau_{2}\left(Z_{\alpha}\right)-\operatorname{csch} \alpha(H) \cdot \tau_{1}\left(Z_{\alpha}\right) \tau_{2}\left(Z_{-\alpha}\right) \\
& \bmod \left(\overline{\mathscr{B}} \overline{\mathfrak{k}}_{2}+\overline{\mathfrak{e}}_{1}^{\prime} \overline{\mathfrak{\varepsilon}}\right)
\end{aligned}
$$

In this congruence, replace $\alpha$ by $-\alpha$ and $H$ by $-H$. Then $h=\exp H$ is replaced by $h^{-1}$ and ${\overline{\mathscr{F}} \bar{k}_{2}}+\overline{\mathfrak{k}}_{1}^{\prime} \overline{\mathscr{B}}$ becomes $\overline{\mathscr{B}}_{2}+\overline{\mathrm{k}}_{1} \overline{\mathscr{B}}$. Then

$$
\begin{aligned}
\theta\left(X_{-\alpha} X_{\alpha}\right) \equiv & -(1+\operatorname{coth} \alpha(H))\left[Z_{-\alpha}, Y_{\alpha}\right] \\
& -Y_{\alpha}\left[(1+\operatorname{coth} \alpha(H)) \tau_{2}\left(Z_{-\alpha}\right)-\operatorname{csch} \alpha(H) \tau_{1}\left(Z_{-\alpha}\right)\right] \\
& +(1+\operatorname{coth} \alpha(H)) \tau_{2}\left(Z_{\alpha}\right) \tau_{2}\left(Z_{-\alpha}\right)-\operatorname{csch} \alpha(H) \cdot \tau_{1}\left(Z_{-\alpha}\right) \tau_{2}\left(Z_{\alpha}\right)
\end{aligned}
$$

Now $\left[X_{\alpha}, X_{-\alpha}\right]=H_{\alpha}$. Therefore

$$
\begin{aligned}
& {\left[Z_{\alpha}, Y_{-\alpha}\right]-\left[Z_{-\alpha}, Y_{\alpha}\right]=\left[Z_{\alpha}, Y_{-\alpha}\right]+\left[Y_{\alpha}, Z_{-\alpha}\right]} \\
& \quad=\frac{1}{4}\left[X_{\alpha}+\theta\left(X_{\alpha}\right), X_{-\alpha}-\theta\left(X_{-\alpha}\right)\right]+\frac{1}{4}\left[X_{\alpha}-\theta\left(X_{\alpha}\right), X_{-\alpha}+\theta\left(X_{-\alpha}\right)\right] \\
& \quad=\frac{1}{2}\left(H_{\alpha}-\theta\left(H_{\alpha}\right)\right)=H_{\bar{\alpha}}
\end{aligned}
$$

We obtain the formula

$$
\begin{aligned}
X_{\alpha} X_{-\alpha}+\theta\left(X_{-\alpha} X_{\alpha}\right) \equiv & (1+\operatorname{coth} \alpha(H)) H_{\alpha} \\
& +Y_{-\alpha}\left[(1+\operatorname{coth} \alpha(H)) \tau_{2}\left(Z_{\alpha}\right)-\operatorname{csch} \alpha(H) \cdot \tau_{1}\left(Z_{\alpha}\right)\right] \\
& -Y_{\alpha}\left[(1+\operatorname{coth} \alpha(H)) \tau_{2}\left(Z_{-\alpha}\right)-\operatorname{csch} \alpha(H) \cdot \tau_{1}\left(Z_{-\alpha}\right)\right] \\
& +(1+\operatorname{coth} \alpha(H))\left[\tau_{2}\left(Z_{\alpha}\right) \tau_{2}\left(Z_{-\alpha}\right)+\tau_{2}\left(Z_{-\alpha}\right) \tau_{2}\left(Z_{\alpha}\right)\right] \\
& -\operatorname{csch} \alpha(H)\left[\tau_{1}\left(Z_{\alpha}\right) \tau_{2}\left(Z_{-\alpha}\right)+\tau_{1}\left(Z_{-\alpha}\right) \tau_{2}\left(Z_{\alpha}\right)\right] .
\end{aligned}
$$

By substituting $-\alpha$ for $\alpha$ in this formula we obtain a similar expression for $X_{-\alpha} X_{\alpha}+\theta\left(X_{\alpha} X_{-\alpha}\right)$. Adding the two expressions together we obtain the formula

$$
\begin{aligned}
R_{\alpha}= & \frac{1}{2}\left[\dot{X_{\alpha}} X_{-\alpha}+\theta\left(X_{-\alpha} X_{\alpha}\right)\right]+\frac{1}{2}\left[X_{-\alpha} X_{\alpha}+\theta\left(X_{\alpha} X_{-\alpha}\right)\right] \\
\equiv & \operatorname{coth} \alpha(H) \cdot H_{\alpha}+Y_{-\alpha}\left[\operatorname{coth} \alpha(H) \cdot \tau_{2}\left(Z_{\alpha}\right)-\operatorname{csch} \alpha(H) \cdot \tau_{1}\left(Z_{\alpha}\right)\right] \\
& -Y_{\alpha}\left[\operatorname{coth} \alpha(H) \cdot \tau_{2}\left(Z_{-\alpha}\right)-\operatorname{csch} \alpha(H) \cdot \tau_{1}\left(Z_{-\alpha}\right)\right] \\
& +\left[\tau_{2}\left(Z_{\alpha}\right) \tau_{2}\left(Z_{-\alpha}\right)+\tau_{2}\left(Z_{-\alpha}\right) \tau_{2}\left(Z_{\alpha}\right)\right] .
\end{aligned}
$$

Using formula (13.1) for $Y_{\alpha}$ and $Y_{-\alpha}$ we find that $R_{\alpha}$ is congruent to the expression

$$
\begin{aligned}
\operatorname{coth} \alpha(H) \cdot H_{\bar{\alpha}} & -(\operatorname{coth} \alpha(H))^{2}\left[\tau_{2}\left(Z_{\alpha}\right) \tau_{2}\left(Z_{-\alpha}\right)+\tau_{2}\left(Z_{-\alpha}\right) \tau_{2}\left(Z_{\alpha}\right)\right] \\
& -(\operatorname{csch} \alpha(H))^{2}\left[\tau_{1}\left(Z_{\alpha}\right) \tau_{1}\left(Z_{-\alpha}\right)+\tau_{1}\left(Z_{-\alpha}\right) \tau_{1}\left(Z_{\alpha}\right)\right] \\
& +2 \operatorname{csch} \alpha(H) \operatorname{coth} \alpha(H)\left[\tau_{2}\left(Z_{\alpha}\right) \tau_{1}\left(Z_{-\alpha}\right)+\tau_{2}\left(Z_{-\alpha}\right) \tau_{1}\left(Z_{\alpha}\right)\right] \\
& +\tau_{2}\left(Z_{\alpha}\right) \tau_{2}\left(Z_{-\alpha}\right)+\tau_{2}\left(Z_{-\alpha}\right) \tau_{2}\left(Z_{\alpha}\right)
\end{aligned}
$$

Using the fact that $1-(\operatorname{coth} \alpha(H))^{2}=-(\operatorname{csch} \alpha(H))^{2}$ we obtain the formula

$$
R_{\alpha} \equiv \operatorname{coth} \alpha(H) \cdot H_{\bar{\alpha}}-F_{\alpha}
$$

where $F_{\alpha}$ is the following operator on $\mathscr{E}$ :

$$
\begin{aligned}
(\operatorname{csch} \alpha(H))^{2}\left[\tau_{2}\left(Z_{\alpha}\right)\right. & \tau_{2}\left(Z_{-\alpha}\right)+\tau_{2}\left(Z_{-\alpha}\right) \tau_{2}\left(Z_{\alpha}\right) \\
& \left.+\tau_{1}\left(Z_{\alpha}\right) \tau_{1}\left(Z_{-\alpha}\right)+\tau_{1}\left(Z_{-\alpha}\right) \tau_{1}\left(Z_{\alpha}\right)\right] \\
& -\operatorname{csch} \alpha(H) \cdot \operatorname{coth} \alpha(H)\left[\tau_{2}\left(Z_{\alpha}\right) \tau_{1}\left(Z_{-\alpha}\right)+\tau_{2}\left(Z_{-\alpha}\right) \tau_{1}\left(Z_{\alpha}\right)\right]
\end{aligned}
$$

We want to show that $\sum_{\alpha \in P_{+}} F_{\alpha}$ is a selfadjoint operator on $V_{\tau}$. Let $\eta$ denote conjugation of $\mathfrak{g}_{\mathrm{c}}$ with respect to $\mathfrak{g}$. Then since $\tau_{1}$ and $\tau_{2}$ are unitary

$$
\tau_{1}(X)^{*}=-\tau_{1}(\eta X), \quad \tau_{2}(X)^{*}=-\tau_{2}(\eta X)
$$

for any $X$ in $\mathfrak{k}_{\mathbf{c}}$. It is easy to see that there exists a constant $c_{\alpha} \neq 0$ such that

$$
\eta X_{\alpha}=c_{\alpha} X_{\alpha^{\eta}}, \quad \eta X_{-\alpha}=c_{\alpha}^{-1} X_{-\alpha^{\eta}}
$$

where $\alpha^{\eta}$ is the conjugate root of $\alpha$. Since $\eta$ commutes with $\theta$, we see that

$$
\eta Z_{\alpha}=c_{\alpha} Z_{\alpha^{\eta}}, \quad \eta Z_{-\alpha}=c_{\alpha}^{-1} Z_{-\alpha^{\eta}}
$$

Therefore $\left(F_{\alpha}\right)^{*}=F_{\alpha^{\eta}}$ for any $\alpha$ in $P_{+}$, since $\alpha^{\eta}=\alpha$ when restricted to $\mathfrak{a}_{\mathfrak{p}}$, and because the functions $\operatorname{csch} \alpha(H)$ and $\operatorname{coth} \alpha(H)$ are real-valued. Therefore

$$
\left(\sum_{\alpha \in P_{+}} F_{\alpha}\right)^{*}=\sum_{\alpha \in P_{+}} F_{\alpha}
$$

We have shown that

$$
\omega_{\mathfrak{g}} \equiv H_{1}^{2}+\sum_{\alpha \in P_{+}} \operatorname{coth} \alpha(H) \cdot H_{\bar{\alpha}}+\sum_{i=1}^{n} \tau_{2}\left(H_{i}\right)^{2}-\sum_{\alpha \in P_{+}} F_{\alpha}+\sum_{\alpha \in P_{M}} R_{\alpha}
$$

We need only show that for any $\alpha$ in $P_{M}, R_{\alpha}$ is congruent to a selfadjoint operator on $V_{\tau}$. In this case

$$
R_{\alpha}=X_{\alpha} X_{-\alpha}+X_{-\alpha} X_{\alpha}
$$

Therefore

$$
R_{\alpha} \equiv \tau_{2}\left(X_{-\alpha}\right) \tau_{2}\left(X_{\alpha}\right)+\tau_{2}\left(X_{\alpha}\right) \tau_{2}\left(X_{-\alpha}\right)
$$

But $\eta R_{\alpha}=R_{\alpha}$. Therefore $\tau_{2}\left(X_{-\alpha}\right) \tau_{2}\left(X_{\alpha}\right)+\tau_{2}\left(X_{\alpha}\right) \tau_{2}\left(X_{-\alpha}\right)$ is a selfadjoint operator on $V_{\tau}$. This completes the proof of Lemma 38 .

Recall that $r_{1}$ and $r_{2}$ stand for the number of roots in $P_{+}$whose restrictions to $a_{p}$ are $\mu_{0}$ and $2 \mu_{0}$ respectively. As before, write

$$
r^{2}=2\left(r_{1}+4 r_{2}\right)=B\left(H_{0}, H\right)
$$

Since $B\left(H_{1}, H_{1}\right)=1, H_{1}^{2}$ equals $r^{-2} H_{0}^{2}$. Since $B\left(H_{\mu_{0}}, H_{0}\right)=\mu_{0}\left(H_{0}\right)=1, H_{\mu_{0}}$ equals $r^{-2} H_{0}$. Therefore, if we write $\delta_{t}^{\prime}\left(\omega_{\mathfrak{g}}\right)$ for $\delta_{\exp t H_{0}}^{\prime}\left(\omega_{\mathfrak{g}}\right)$ and $Q(t)$ for $Q\left(t H_{0}\right)$, we obtain the following formula from Lemma 38:

$$
\begin{equation*}
r^{2} \delta_{t}^{\prime}\left(\omega_{\mathfrak{g}}\right)=\left(\frac{d}{d t}\right)^{2}+\left(r_{1} \operatorname{coth} t+2 r_{2} \operatorname{coth} 2 t\right) \frac{d}{d t}+r^{2} Q(t) \tag{13.2}
\end{equation*}
$$

Recall that

$$
D(t)=(2 \sinh t)^{r_{1}}(2 \sinh 2 t)^{r_{2}}, \quad t \in \mathbf{R} .
$$

If $t$ is greater than zero, $D(t)$ does not vanish. Define a new differential operator $\delta_{t}\left(\omega_{\mathfrak{g}}\right)$ on $C^{\infty}\left(\mathbf{R}^{+}\right) \otimes V_{\tau}$ by

$$
\delta_{t}\left(\omega_{\mathfrak{g}}\right)=D(t)^{1 / 2} \delta_{t}^{\prime}\left(\omega_{\mathfrak{g}}\right) \circ D(t)^{-1 / 2}
$$

Lemma 39. For any $t>0$ there is a selfadjoint operator $q(t)$ on $V_{\tau}$ such that

$$
r^{2} \delta_{t}\left(\omega_{\mathfrak{g}}\right)=\left(\frac{d}{d t}\right)^{2}+q(t)
$$

Proof. From (13.2) we obtain the formula

$$
\begin{aligned}
& r^{2} \delta_{t}\left(\omega_{\mathfrak{g}}\right)=r^{2} D(t)^{1 / 2} \delta_{t}^{\prime}\left(\omega_{\mathfrak{g}}\right) \circ D(t)^{-1 / 2} \\
& \quad=\left(\frac{d}{d t}\right)^{2}+\left(r_{1} \operatorname{coth} t+2 r_{2} \operatorname{coth} 2 t+2 D(t)^{1 / 2} \frac{d}{d t} D(t)^{-1 / 2}\right) \frac{d}{d t}+q(t)
\end{aligned}
$$

$q(t)$ is some operator obtained by adding real-valued scalar functions of $t$ to the operator $Q(t)$. In particular $q(t)$ is selfadjoint. Now

$$
2 D(t)^{1 / 2} \cdot \frac{d}{d t} D(t)^{-1 / 2}=-D(t)^{-1} \cdot \frac{d}{d t} D(t)
$$

If we differentiate $D(t)$ we obtain the formula

$$
-D(t)^{-1} \cdot \frac{d}{d t} D(t)=-\left[r_{1} \operatorname{coth} t+2 r_{2} \operatorname{coth} 2 t\right]
$$

Therefore

$$
r^{2} \delta_{t}\left(\omega_{\mathfrak{g}}\right)=\left(\frac{d}{d t}\right)^{2}+q(t)
$$

We are now in a position to relate the linear transformations $c^{+}(\Lambda)$ and $c^{-}(\Lambda)$ of $L^{\tau}$ with the Plancherel measure $\beta(\sigma, \Lambda)$.

If $\psi$ is in $L^{\tau}, \Lambda \neq 0$, and $t>0$, define

$$
f_{\Lambda}(\psi: t)=D(t)^{1 / 2} E_{\Lambda}\left(\psi: \exp t H_{0}\right)
$$

Recall that $L^{\tau}$ is a direct sum of a finite number of orthogonal subspaces of the form $L_{\sigma}^{\tau}$. If $\psi_{\sigma}$ is in $L_{\sigma}^{\tau}$ for some $\sigma$ in $\mathscr{E}_{M}$, we have the formula

$$
E_{\Lambda}\left(\psi_{\sigma}: \exp t H_{0} ; \omega_{\mathfrak{g}}\right)=\pi_{\sigma, \Lambda}\left(\omega_{\mathfrak{g}}\right) \cdot E_{\Lambda}\left(\psi_{\sigma}: \exp t H_{0}\right)
$$

Fix $\mu$ in $L_{1}^{\prime}$. Choose $\sigma_{1}$ and $\sigma_{2}$ in $\mathscr{E}_{M}(\mu)$ and fix $\psi_{1}$ and $\psi_{2}$ in $L_{\sigma_{1}}^{\tau}$ and $L_{\sigma_{2}}^{\tau}$ respectively. Using Lemma 37 we see that for $\alpha$ equal to 1 or 2 ,

$$
\begin{aligned}
\delta_{t}\left(\omega_{\mathfrak{g}}\right) f_{\Lambda}\left(\psi_{\alpha}: t\right) & =D(t)^{1 / 2} \delta_{t}^{\prime}\left(\omega_{\mathfrak{g}}\right) E_{\Lambda}\left(\psi_{\alpha}: \exp t H_{0}\right) \\
& =\pi_{\sigma_{\alpha}, \Lambda}\left(\omega_{\mathfrak{g}}\right) \cdot f_{\Lambda}\left(\psi_{\alpha}: t\right)
\end{aligned}
$$

By Lemma 39 we obtain the formula

$$
\left[\left(\frac{d}{d t}\right)^{2}+q(t)\right] f_{\Lambda}\left(\psi_{\alpha}: t\right)=r^{2} \pi_{\sigma_{\alpha}, \Lambda}\left(\omega_{\mathfrak{g}}\right) f_{\Lambda}\left(\psi_{\alpha}: t\right), \quad \alpha=1,2
$$

Write $\frac{d}{d t} f_{\Lambda}\left(\psi_{\alpha}: t\right)$ as $f_{\Lambda}^{\prime}\left(\psi_{\alpha}: t\right)$. Fix $\Lambda_{0} \neq 0$. Then

$$
\begin{aligned}
\frac{d}{d t}[ & \left.\left(f_{\Lambda}\left(\psi_{1}: t\right), f_{\Lambda_{0}}^{\prime}\left(\psi_{2}: t\right)\right)-\left(f_{\Lambda}^{\prime}\left(\psi_{1}: t\right), f_{\Lambda_{0}}\left(\psi_{2}: t\right)\right)\right] \\
& =\left(f_{\Lambda}\left(\psi_{1}: t\right), f_{\Lambda_{0}}^{\prime \prime}\left(\psi_{2}: t\right)\right)-\left(f_{\Lambda}^{\prime \prime}\left(\psi_{1}: t\right), f_{\Lambda_{0}}\left(\psi_{2}: t\right)\right) \\
& =r^{2}\left(\pi_{\sigma_{2}, \Lambda}\left(\omega_{\mathfrak{g}}\right)-\pi_{\sigma_{1}, \Lambda}\left(\omega_{\mathfrak{g}}\right)\right) \cdot\left(f_{\Lambda}\left(\psi_{1}: t\right), f_{\Lambda_{0}}\left(\psi_{2}: t\right)\right)
\end{aligned}
$$

We are using the fact that $q(t)$ is a selfadjoint operator on $V_{\tau}$ and that $\pi_{\sigma_{2}, \Lambda}\left(\omega_{\mathfrak{g}}\right)$ is a real number. The scalar product above is of course on the space $V_{\tau}$. Now by (12.4), $\left|\sigma_{1}\right|=\left|\sigma_{2}\right|$. Therefore, by (6.10)

$$
\left(\pi_{\sigma_{2}, \Lambda}\left(\omega_{\mathfrak{g}}\right)-\pi_{\sigma_{1}, \Lambda}\left(\omega_{\mathfrak{g}}\right)\right)=r^{-2}\left(\Lambda^{2}-\Lambda_{0}^{2}\right) .
$$

For $T$ a positive real number, define the number $V_{T}\left(\Lambda, \Lambda_{0}\right)$ by

$$
\begin{equation*}
V_{T}\left(\Lambda, \Lambda_{0}\right)=\int_{0}^{T}\left(f_{\Lambda}\left(\psi_{1}: t\right), f_{\Lambda_{0}}\left(\psi_{2}: t\right)\right) d t \tag{13.3}
\end{equation*}
$$

Then $V_{T}\left(\Lambda, \Lambda_{0}\right)$ is equal to the expression

$$
\left.\left(\Lambda^{2}-\Lambda_{0}^{2}\right)^{-1}\left\{\left(f_{\Lambda}\left(\psi_{1}: t\right), f_{\Lambda_{0}}^{\prime}\left(\psi_{2}: t\right)\right)-\left(f_{\Lambda}^{\prime}\left(\psi_{1}: t\right), f_{\Lambda_{0}}\left(\psi_{2}: t\right)\right)\right\}\right|_{0} ^{T}
$$

Now $D(t)^{1 / 2} \rightarrow 0$ as $t \rightarrow 0$, so $f_{\Lambda}\left(\psi_{2}: t\right) \rightarrow 0$ as $t \rightarrow 0$. Therefore the evaluation of the above expression at $t=0$ is equal to zero. We then have the formula
$V_{T}\left(\Lambda, \Lambda_{0}\right)=\left(\Lambda^{2}-\Lambda_{0}^{2}\right)^{-1}\left\{\left(f_{\Lambda}\left(\psi_{1}: T\right), f_{\Lambda_{0}}^{\prime}\left(\psi_{2}: T\right)\right)-\left(f_{\Lambda}^{\prime}\left(\psi_{1}: T\right), f_{\Lambda_{0}}\left(\psi_{2}: T\right)\right)\right\}$.
Define the integer $q$ as in $\S 3$. Define the distribution $\Delta_{\Lambda_{0}}^{T}$ on $C_{0}^{\infty}(\mathbf{R})$ by

$$
\Delta_{\Lambda_{0}}^{T}(h)=\int_{-\infty}^{\infty} V_{T}\left(\Lambda, \Lambda_{0}\right) \cdot \Lambda^{2 q} \cdot h(\Lambda) d \Lambda, \quad h \in C_{0}^{\infty}(\mathbf{R})
$$

LEMMA 40. If $h$ is in $C_{0}^{\infty}(\mathbf{R})$ and $\Lambda_{0} \neq 0$, the limit as $T$ approaches infinity of $\Delta_{\Lambda_{0}}^{T}(h)$ is equal to each of the following expressions:
(i)

$$
\begin{aligned}
& 2 \pi\left(c^{+}\left(\Lambda_{0}\right) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot \Lambda_{0}^{2 q} \cdot h\left(\Lambda_{0}\right) \\
& \quad+\pi\left(c^{-}\left(-\Lambda_{0}\right) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot\left(-\Lambda_{0}\right)^{2 q} \cdot h\left(-\Lambda_{0}\right)
\end{aligned}
$$

(ii)

$$
\begin{aligned}
& 2 \pi\left(c^{-}\left(\Lambda_{0}\right) \psi_{1}, c^{-}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot \Lambda_{0}^{2 q} \cdot h\left(\Lambda_{0}\right) \\
& \quad+\pi\left(c^{+}\left(-\Lambda_{0}\right) \psi_{1}, c^{-}\left(\Lambda_{0}\right) \psi_{2}\right)_{M}\left(-\Lambda_{0}\right)^{2 q} \cdot h\left(-\Lambda_{0}\right)
\end{aligned}
$$

Proof. Write $\theta\left(\psi_{\alpha}: t\right)$ as the function

$$
c^{+}(\Lambda) \psi_{\alpha}(1) e^{i \Lambda t}+c^{-}(\Lambda) \psi_{\alpha}(1) e^{-i \Lambda t}
$$

Write $\frac{d}{d t} \theta_{\Lambda}\left(\psi_{\alpha}: t\right)$ as $\theta_{\Lambda}^{\prime}\left(\psi_{\alpha}: t\right)$. Write $\frac{d}{d t} E_{\Lambda}\left(\psi_{\alpha}: \exp t H_{0}\right)$ as $E_{\Lambda}^{\prime}\left(\psi_{\alpha}: \exp t H_{0}\right)$. Let

$$
\begin{aligned}
& q_{\alpha}(\Lambda, t)=f_{\Lambda}\left(\psi_{\alpha}: t\right)-\theta_{\Lambda}\left(\psi_{\alpha}: t\right) \\
& q_{\alpha}^{\prime}(\Lambda, t)=f_{\Lambda}^{\prime}\left(\psi_{\alpha}: t\right)-\theta_{\Lambda}^{\prime}\left(\psi_{\alpha}: t\right)
\end{aligned}
$$

for $\alpha=1,2, \Lambda \neq 0$, and $t>0$.

It is clear that we can find a number $\delta_{1}>0$ and a constant $C$ such that for any $t>0$

$$
\left|e^{-t \rho\left(H_{0}\right)} D(t)^{1 / 2}-1\right|+\left|e^{-t \rho\left(H_{0}\right)} \cdot \frac{d}{d t}\left(D(t)^{1 / 2}\right)-\rho\left(H_{0}\right)\right|<C e^{-\delta_{1} t}
$$

It is also clear that

$$
\begin{aligned}
& \left|f_{\Lambda}^{\prime}\left(\psi_{\alpha}: t\right)-\frac{d}{d t}\left[e^{t \rho\left(H_{0}\right)} E_{\Lambda}\left(\psi_{\alpha}: \exp t H_{0}\right)\right]\right| \\
& \quad ₹\left|D(t)^{1 / 2} \cdot E_{\Lambda}^{\prime}\left(\psi_{\alpha}: \exp t H_{0}\right)-e^{t \rho\left(H_{0}\right)} E_{\Lambda}^{\prime}\left(\psi_{\alpha}: \exp t H_{0}\right)\right| \\
& \quad+\left|\frac{d}{d t}(D(t))^{1 / 2} \cdot E_{\Lambda}\left(\psi_{\alpha}: \exp t H_{0}\right)-\frac{d}{d t}\left(e^{t \rho\left(H_{0}\right)}\right) \cdot E_{\Lambda}\left(\psi_{\alpha}: \exp t H_{0}\right)\right| \\
& \quad=\left|e^{t \rho\left(H_{0}\right)} E_{\Lambda}^{\prime}\left(\psi_{\alpha}: \exp t H_{0}\right)\right| \cdot\left|e^{-t \rho\left(H_{0}\right)} D(t)^{1 / 2}-1\right| \\
& \quad+\left|e^{t \rho\left(H_{0}\right)} E_{\Lambda}\left(\psi_{\alpha}: \exp t H_{0}\right)\right| \cdot\left|e^{-t \rho\left(H_{0}\right)} \frac{d}{d t}\left(D(t)^{1 / 2}\right)-\rho\left(H_{0}\right)\right|
\end{aligned}
$$

Now by (4.1) and Lemma 30, we observe that both $\left|e^{t \rho\left(H_{0}\right)} E_{\Lambda}\left(\psi_{\alpha}: \exp t H_{0}\right)\right|$ and $\left|e^{t \rho\left(H_{0}\right)} E_{\Lambda}^{\prime}\left(\psi_{\alpha}: \exp t H_{0}\right)\right|$ are bounded by the product of $(1+r t)^{d}$ and a polynomial in $|\Lambda|$. Therefore, there exists a $\delta>0$ and a polynomial $p$ such that for $\Lambda$ in $\mathbf{R}$ and $t>0$

$$
\begin{equation*}
\left|f_{\Lambda}^{\prime}\left(\psi_{\alpha}: t\right)-\frac{d}{d t}\left(e^{t \rho\left(H_{0}\right)} E_{\Lambda}\left(\psi_{\alpha}: \exp t H_{0}\right)\right)\right|<p(|\Lambda|) e^{-\delta t} \tag{13.5}
\end{equation*}
$$

It then follows from (11.2) that there exists a polynomial $p$ and an $\varepsilon>0$ such that for $\alpha=1,2, \Lambda \neq 0$, and $t>0$

$$
\begin{equation*}
\left|q_{\alpha}^{\prime}(\Lambda, t)\right|<p(|\Lambda|) e^{-\varepsilon t} \tag{13.6}
\end{equation*}
$$

By the same argument we obtain the inequality

$$
\begin{equation*}
\left|q_{\alpha}(\Lambda, t)\right|<p(|\Lambda|) e^{-\varepsilon t} \tag{13.7}
\end{equation*}
$$

For $\Lambda \neq 0$, and $T \gg 0$, let $Q(\Lambda, T)$ be the sum of the following three terms:
(i) $\left(q_{1}(\Lambda, T), q_{2}^{\prime}\left(\Lambda_{0}, T\right)\right)-\left(q_{1}^{\prime}(\Lambda, T), q_{2}\left(\Lambda_{0}, T\right)\right)$,
(ii) $\left(q_{1}(\Lambda, T), \theta_{\Lambda_{0}}^{\prime}\left(\psi_{2}: T\right)\right)-\left(q_{1}^{\prime}(\Lambda, T), \theta_{\Lambda_{0}}\left(\psi_{2}: T\right)\right)$,
(iii) $\left(\theta_{\Lambda}\left(\psi_{1}: T\right), q_{2}^{\prime}\left(\Lambda_{0}, T\right)\right)-\left(\theta_{\Lambda}^{\prime}\left(\psi_{1}: T\right), q_{2}\left(\Lambda_{0}, T\right)\right)$.

By Lemma 34 it is clear that for $\alpha=1,2$, both $\left|\Lambda^{2 q} \theta_{\Lambda}\left(\psi_{\alpha}: T\right)\right|$ and $\left|\Lambda^{2 q} \theta_{\Lambda}^{\prime}\left(\psi_{\alpha}: T\right)\right|$ are bounded by a polynomial in $|\Lambda|$, independently of $T$. Therefore, using (13.6) and (13.7), we see that there is a polynomial $p$ and an $\varepsilon>0$ such that for $\Lambda$ in $\mathbf{R}$ and $T>0$

$$
\begin{equation*}
\left|\Lambda^{2 q} Q(\Lambda, T)\right|<p(|\Lambda|) e^{-\varepsilon t} \tag{13.8}
\end{equation*}
$$

We shall also need a weak bound on $\left|\frac{d}{d \Lambda}\left(\Lambda^{2 q} Q(\Lambda, T)\right)\right|$. By using Lemma 30 and Lemma 35 we can show that there exists a polynomial $p$ and an integer $n$ such that for $\Lambda$ in $\mathbf{R}$ and $T>0$

$$
\begin{equation*}
\left|\frac{d}{d \Lambda}\left(\Lambda^{2 q} Q(\Lambda, T)\right)\right|<p(|\Lambda|)(1+T)^{n} \tag{13.9}
\end{equation*}
$$

We shall now break $V_{T}\left(\Lambda, \Lambda_{0}\right)$ up into a sum of terms which we can handle separately with the above estimates.

$$
\begin{aligned}
\Lambda^{2 q} V_{T} & \left(\Lambda, \Lambda_{0}\right) \\
= & \Lambda^{2 q} \cdot\left(\Lambda^{2}-\Lambda_{0}^{2}\right)^{-1} \cdot\left\{\left(f_{1}\left(\psi_{1}: T\right), f_{\Lambda_{0}}^{\prime}\left(\psi_{2}: T\right)\right)-\left(f_{\Lambda}^{\prime}\left(\psi_{1}: T\right), f_{\Lambda_{0}}\left(\psi_{2}: T\right)\right)\right\} \\
= & \Lambda^{2 q} \cdot\left(\Lambda^{2}-\Lambda_{0}^{2}\right)^{-1}\left(\theta_{\Lambda}\left(\psi_{1}: T\right), \theta_{\Lambda_{0}}^{\prime}\left(\psi_{2}: T\right)\right) \\
& -\Lambda^{2 q} \cdot\left(\Lambda^{2}-\Lambda_{0}^{2}\right)^{-1}\left(\theta_{\Lambda}^{\prime}\left(\psi_{1}: T\right), \theta_{\Lambda_{0}}\left(\psi_{2}: T\right)\right) \\
& +\Lambda^{2 q} \cdot\left(\Lambda^{2}-\Lambda_{0}^{2}\right)^{-1} \cdot Q(\Lambda, T)
\end{aligned}
$$

Recall that for any two vectors $\psi$ and $\psi^{\prime}$ in $L^{\tau}$

$$
\left(\psi, \psi^{\prime}\right)_{M}=(\psi(1), \psi(1))
$$

Then we obtain the following formula:

$$
\begin{aligned}
& \Lambda^{2 q} V_{T}\left(\Lambda, \Lambda_{0}\right) \\
&=i \Lambda^{2 q} \cdot\left(\Lambda^{2}-\Lambda_{0}^{2}\right)^{-1}\left\{\left(c^{+}(\Lambda) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot e^{i T\left(\Lambda-\Lambda_{0}\right)}\left(\Lambda_{0}+\Lambda\right)\right. \\
&+\left(c^{-}(\Lambda) \psi_{1}, c^{-}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot e^{i T\left(-\Lambda+\Lambda_{0}\right)}\left(-\Lambda_{0}-\Lambda\right) \\
&+\left(c^{+}(\Lambda) \psi_{1}, c^{-}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot e^{i T\left(\Lambda+\Lambda_{0}\right)}\left(-\Lambda_{0}+\Lambda\right) \\
&\left.+\left(c^{-}(\Lambda) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot e^{i T\left(-\Lambda-\Lambda_{0}\right)}\left(\Lambda_{0}-\Lambda\right)\right\} \\
&+\Lambda^{2 q} \cdot\left(\Lambda^{2}-\Lambda_{0}^{2}\right)^{-1} \cdot Q(\Lambda, T)
\end{aligned}
$$

Now $\left(\Lambda-\Lambda_{0}\right) \cdot \Lambda^{2 q} \cdot V_{T}\left(\Lambda, \Lambda_{0}\right)=0$ if $\Lambda=\Lambda_{0}$. Therefore, for any $T>0$,
$-i \Lambda_{0}^{2 q}\left\{\left(c^{+}\left(\Lambda_{0}\right) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M}-\left(c^{-}\left(\Lambda_{0}\right) \psi_{1}, c^{-}\left(\Lambda_{0}\right) \psi_{2}\right)_{M}\right\}$

$$
+\Lambda_{0}^{2 q} \cdot\left(2 \Lambda_{0}\right)^{-1} \cdot Q\left(\Lambda_{0}, T\right)=0
$$

Now by (13.8), $Q\left(\Lambda_{0}, T\right)$ approaches 0 as $T$ approaches $\infty$, so

$$
\left(c^{+}\left(\Lambda_{0}\right) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M}=\left(c^{-}\left(\Lambda_{0}\right) \psi_{1}, c^{-}\left(\Lambda_{0}\right) \psi_{2}\right)_{M}
$$

This implies that $Q\left(\Lambda_{0}, T\right)=0$ for any $T \gg 0$. In a similar manner, we can show that

$$
\left(c^{-}\left(-\Lambda_{0}\right) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M}=\left(c^{+}\left(-\Lambda_{0}\right) \psi_{1}, c^{-}\left(\Lambda_{0}\right) \psi_{2}\right)_{M}
$$

and that $Q\left(-\Lambda_{0}, T\right)=0$ for any $T \gg 0$.
The number

$$
\Delta_{\Lambda_{0}}^{T}(h)=\int_{\mathbf{R}} h(\Lambda) \cdot \Lambda^{2 q} \cdot V_{T}\left(\Lambda, \Lambda_{0}\right) d \Lambda
$$

is equal to the sum of the following three terms:
(i)

$$
\begin{aligned}
\int_{\mathbf{R}}\{ & \left(c^{+}(\Lambda) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot e^{i T\left(\Lambda-\Lambda_{0}\right)} \\
& \left.-\left(c^{-}(\Lambda) \psi_{1}, c^{-}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot e^{-i T\left(\Lambda-\Lambda_{0}\right)}\right\} \cdot h(\Lambda)(-i) \Lambda^{2 q} \cdot\left(\Lambda-\Lambda_{0}\right)^{-1} d \Lambda
\end{aligned}
$$

(ii)

$$
\begin{aligned}
\int_{\mathbf{R}}\{ & \left(c^{+}(\Lambda) \psi_{1}, c^{-}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot e^{i T\left(\Lambda+\Lambda_{0}\right)} \\
& \left.\quad-\left(c^{-}(\Lambda) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot e^{-i T\left(\Lambda+\Lambda_{0}\right)}\right\} \cdot h(\Lambda)(-i) \Lambda^{2 q} \cdot\left(\Lambda+\Lambda_{0}\right)^{-1} d \Lambda
\end{aligned}
$$

(iii)

$$
\int_{\mathbf{R}} Q(\Lambda, T) \cdot \Lambda^{2 q} \cdot h(\Lambda) \cdot\left(\Lambda^{2}-\Lambda_{0}^{2}\right)^{-1} d \Lambda
$$

Let us assume for the moment that $h$ vanishes in a neighborhood of $\left(-\Lambda_{0}\right)$. We wish to examine the behavior of each of the above three terms as $T$ approaches $\infty$. We deal first with the third term.

Let $n$ be the integer in (13.9). The third term is equal to $I_{1}+I_{2}$ where $I_{1}$ and $I_{2}$ are those parts of the integral in (iii) which are taken over the sets $\left\{\Lambda:\left|\Lambda-\Lambda_{0}\right| ₹ T^{-n-1}\right\}$ and $\left\{\Lambda:\left|\Lambda-\Lambda_{0}\right|>T^{-n-1}\right\}$ respectively. Apply the mean value theorem to the function $\Lambda^{2 q} Q(\Lambda, T)$ of $\Lambda$. Since $Q\left(\Lambda_{0}, T\right)=0$, there is a real $\Lambda_{1}$, with $\left|\Lambda_{1}-\Lambda_{0}\right|<\left|\Lambda-\Lambda_{0}\right|$, such that

$$
\Lambda^{2 q} Q(\Lambda, T)=\left(\Lambda-\Lambda_{0}\right) \cdot \frac{d}{d \Lambda_{1}}\left(\Lambda_{1}^{2 q} Q\left(\Lambda_{1}, T\right)\right)
$$

Then by (13.9) we see that $\left|I_{1}\right|$ is bounded by a constant multiple of $\left(1+T^{n}\right) T^{-n-1}$. Therefore, $\left|I_{1}\right|$ approaches 0 as $T \rightarrow \infty$. On the other hand, there is a constant $C$, independent of $\Lambda$ and $T$, such that if $\left|\Lambda-\Lambda_{0}\right|>T^{-n-1}$,

$$
\left|\Lambda^{2 q} \cdot h(\Lambda) \cdot\left(\Lambda^{2}-\Lambda_{0}^{2}\right)^{-1} \cdot Q(\Lambda, T)\right| ₹ C p(|\Lambda|) \cdot|h(\Lambda)| \cdot T^{n+1} \cdot e^{-\varepsilon t}
$$

by (13.8). (We have used the fact that $h(\Lambda)$ vanishes in a neighborhood of $-\Lambda_{0}$.) This implies that $\left|I_{2}\right|$ approaches 0 as $T$ approaches $\infty$. Therefore, term (iii) goes to 0 as $T$ goes to $\infty$.

To deal with term (ii), we observe that both

$$
(-i) \cdot \Lambda^{2 q} \cdot\left(\Lambda+\Lambda_{0}\right)^{-1} \cdot\left(c^{+}(\Lambda) \psi_{1}, c^{-}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot h(\Lambda)
$$

and

$$
(-i) \cdot \Lambda^{2 q} \cdot\left(\Lambda+\Lambda_{0}\right)^{-1} \cdot\left(c^{-}(\Lambda) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot h(\Lambda)
$$

are continuous functions of $\Lambda$, since $h(\Lambda)$ vanishes in a neighborhood of $\left(-\Lambda_{0}\right)$. Therefore, by the Riemann-Lebesgue lemma, term (ii) goes to 0 as $T$ goes to $\infty$.

Now, since $.\left(c^{+}\left(\Lambda_{0}\right) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M}=\left(c^{-}\left(\Lambda_{0}\right) \psi_{1}, c^{-}\left(\Lambda_{0}\right) \psi_{2}\right)_{M}$, we may rewrite term (i) as

$$
\begin{aligned}
& 2 \int_{\mathbf{R}}\left(c^{+}\left(\Lambda_{0}\right) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot \Lambda^{2 q} \cdot h(\Lambda) \cdot \sin T\left(\Lambda-\Lambda_{0}\right) \cdot\left(\Lambda-\Lambda_{0}\right)^{-1} d \Lambda \\
& \quad+(-i) \int_{\mathbf{R}}\left(c^{+}(\Lambda) \psi_{1}-c^{+}\left(\Lambda_{0}\right) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot\left(\Lambda-\Lambda_{0}\right)^{-1} \cdot \Lambda^{2 q} \cdot h(\Lambda) \\
& \quad \cdot e^{i T\left(\Lambda-\Lambda_{0}\right)} d \Lambda \\
& \quad+(i) \int_{\mathbf{R}}\left(c^{-}(\Lambda) \psi_{1}-c^{-}\left(\Lambda_{0}\right) \psi_{1}, c^{-}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot\left(\Lambda-\Lambda_{0}\right)^{-1} \cdot \Lambda^{2 q} \cdot h(\Lambda) \\
& \quad \cdot e^{-i T\left(\Lambda-\Lambda_{0}\right)} d \Lambda \\
& \quad=J_{1}+J_{2}+J_{3}, \quad \text { say. }
\end{aligned}
$$

Now $\Lambda^{2 q} c^{+}(\Lambda) \psi_{1}$ and $\Lambda^{2 q} c^{-}(\Lambda) \psi_{1}$ are both continuously differentiable in $\Lambda$ by Lemma 35. Therefore, we can use the Riemann-Lebesgue lemma to prove that

$$
\lim _{T \rightarrow \infty}\left(\left|J_{2}\right|+\left|J_{3}\right|\right)=0
$$

It follows that

$$
\begin{aligned}
\lim _{T \rightarrow \infty} & \Delta_{\Lambda_{0}}^{T}(h) \\
= & \lim _{T \rightarrow \infty} 2 \int_{\mathbf{R}}\left(c^{+}\left(\Lambda_{0}\right) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot \Lambda^{2 q} \cdot h(\Lambda) \cdot \sin T\left(\Lambda-\Lambda_{0}\right) \\
& \cdot\left(\Lambda-\Lambda_{0}\right)^{-1} d \Lambda
\end{aligned}
$$

Now it is an easy fact that

$$
\lim _{T \rightarrow \infty} \sin T\left(\Lambda-\Lambda_{0}\right) \cdot\left(\Lambda-\Lambda_{0}\right)^{-1}=\pi \delta_{\Lambda_{0}}
$$

where the limit is taken in the topology of $\mathscr{S}^{\prime}(\mathbf{R})$ and $\delta_{\Lambda_{0}}$ is the Dirac distribution at $\Lambda_{0}$.

We have proved that if $h$ vanishes in a neighborhood of $\left(-\Lambda_{0}\right)$,

$$
\lim _{T \rightarrow \infty} \Delta_{\Lambda_{0}}^{T}(h)=2 \pi\left(c^{+}\left(\Lambda_{0}\right) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot \Lambda_{0}^{2 q} \cdot h\left(\Lambda_{0}\right)
$$

Similarly, if $h$ vanishes in a neighborhood $\Lambda_{0}$, we could show that

$$
\lim _{T \rightarrow \infty} \Delta_{\Lambda_{0}}^{T}(h)=2 \pi\left(c^{-}\left(-\Lambda_{0}\right) \psi_{1}, c^{+}\left(\Lambda_{0}\right) \psi_{2}\right)_{M} \cdot\left(-\Lambda_{0}\right)^{2 q} \cdot h\left(-\Lambda_{0}\right)
$$

This completes the proof of Lemma 40. 四
For any $\sigma$ in $\mathscr{E}_{M}$ and $\Lambda$ in $\mathbf{R}, \beta(\sigma, \Lambda)=\beta\left(\sigma^{\prime}, \Lambda\right)$ by Lemma 5 . Then if $\sigma_{1}$ and $\sigma_{2}$ are in $\mathscr{E}_{M}(\mu)$ for some $\mu$ in $L_{1}^{\prime}, \beta\left(\sigma_{1}, \Lambda\right)=\beta\left(\sigma_{2}, \Lambda\right)$. We write this number as $\beta(\mu, \Lambda)$. It is also easy to see from the Weyl dimension formula and the definition of $\mathscr{E}_{M}(\mu)$ that if $\sigma_{1}$ and $\sigma_{2}$ are in $\mathscr{E}_{M}(\mu), \operatorname{dim} \sigma_{1}=\operatorname{dim} \sigma_{2}$. We write this number as $\operatorname{dim} \mu$.

LEMMA 41. Assume $\Lambda_{0} \neq 0$. Fix $\mu$ in $L_{1}^{\prime}$. Then there is a constant $C$, independent of $\Lambda_{0}$, such that the following operator equations hold on $L_{\mu}^{\tau}$ :
(i) $c^{+}\left(\Lambda_{0}\right)^{*} c^{+}\left(\Lambda_{0}\right)=c^{-}\left(\Lambda_{0}\right)^{*} c^{-}\left(\Lambda_{0}\right)=\left(C \cdot \beta\left(\mu, \Lambda_{0}\right)\right)^{-1}$,
(ii) $c^{+}\left(\Lambda_{0}\right)^{*} c^{-}\left(-\Lambda_{0}\right)=c^{-}\left(\Lambda_{0}\right)^{*} c^{+}\left(-\Lambda_{0}\right)=\left(C \cdot \beta\left(\mu, \Lambda_{0}\right)\right)^{-1} \cdot M\left(-\Lambda_{0}\right)$.

Proof. In §12 we remarked that $L_{\mu}$ was an invariant subspace of $c^{ \pm}\left(\Lambda_{0}\right)$. Then $L_{\mu}$ is also invariant under $c^{ \pm}\left(\Lambda_{0}\right)^{*}$. Then from Lemma 40 we obtain the equations

$$
\begin{aligned}
& c^{+}\left(\Lambda_{0}\right)^{*} c^{+}\left(\Lambda_{0}\right)=c^{-}\left(\Lambda_{0}\right)^{*} c^{-}\left(\Lambda_{0}\right) \\
& c^{+}\left(\Lambda_{0}\right)^{*} c^{-}\left(-\Lambda_{0}\right)=c^{-}\left(\Lambda_{0}\right)^{*} c^{+}\left(-\Lambda_{0}\right)
\end{aligned}
$$

Now fix $\psi_{1}$ in $L_{\sigma_{1}}^{\tau}$ and $\psi_{2}$ in $L_{\sigma_{2}}^{\tau}$ for $\sigma_{1}, \sigma_{2}$ in $\mathscr{E}_{M}(\mu)$. Also fix $h$ in $C_{0}^{\infty}(\mathbf{R})$. Then

$$
\begin{aligned}
\lim _{T \rightarrow \infty} \Delta_{\Lambda_{0}}^{T}(h)= & \lim _{T \rightarrow \infty} \int_{\mathbf{R}} h(\Lambda) \cdot \Lambda^{2 q} \cdot V_{T}\left(\Lambda, \Lambda_{0}\right) d \Lambda \\
= & \lim _{T \rightarrow \infty} \int_{0}^{T} \int_{\mathbf{R}}\left(f_{\Lambda}\left(\psi_{1}: t\right), f_{\Lambda_{0}}\left(\psi_{2}: t\right)\right) \cdot \Lambda^{2 q} \cdot h(\Lambda) d \Lambda d t \\
= & \lim _{T \rightarrow \infty} \int_{0}^{T} D(t) \int_{\mathbf{R}}\left(E_{\Lambda}\left(\psi_{1}: \exp t H_{0}\right), E_{\Lambda_{0}}\left(\psi_{2}: \exp t H_{0}\right)\right) \\
& \cdot \Lambda^{2 q} \cdot h(\Lambda) d \Lambda d t \\
= & \lim _{T \rightarrow \infty}(c)^{-1} \int_{G_{T}} \int_{\mathbf{R}}\left(E_{\Lambda}\left(\psi_{1}: x\right), E_{\Lambda_{0}}\left(\psi_{2}: x\right)\right) \cdot \Lambda^{2 q} \cdot h(\Lambda) d \Lambda d x
\end{aligned}
$$

where $G_{T}=\left\{k_{1} \cdot \exp t H_{0} \cdot k_{2}: k_{1}, k_{2} \in K, 0 \gtrless t<T\right\}$. Let us assume that $h$ vanishes in a neighborhood of 0 . Then we define

$$
u(\Lambda)=\Lambda^{2 q} \cdot h(\Lambda) \cdot \beta(\mu, \Lambda)^{-1} \psi_{1}, \quad \Lambda \in \mathbf{R}
$$

$\beta(\mu, \Lambda)$ is infinitely differentiable and does not vanish if $\Lambda \neq 0$ by Lemma 5. Therefore $u \in C_{0}^{\infty}(\mathbf{R}) \otimes L^{\tau}$.

Then in the notation of (7.4) and (7.5)

$$
\left.\lim _{T \rightarrow \infty} \Delta_{\Lambda_{0}}^{T}(h)=(c \cdot \operatorname{dim} \mu)^{-1} \hat{\tilde{u}}\left(\Lambda_{0}\right), \psi_{2}\right)_{M}
$$

By (7.6) the right-hand side of this formula equals

$$
(2 c \cdot \operatorname{dim} \mu)^{-1}\left(u\left(\Lambda_{0}\right)+M\left(-\Lambda_{0}\right) u\left(-\Lambda_{0}\right), \psi_{2}\right)_{M}
$$

which in turn is equal to

$$
\begin{aligned}
& \left(2 c \cdot \operatorname{dim} \mu \cdot \beta\left(\mu, \Lambda_{0}\right)\right)^{-1} \Lambda_{0}^{2 q} \cdot h\left(\Lambda_{0}\right) \cdot\left(\psi_{1}, \psi_{2}\right)_{M} \\
& \quad+\left(2 c \cdot \operatorname{dim} \mu \cdot \beta\left(\mu, \Lambda_{0}\right)\right)^{-1} \Lambda_{0}^{2 q} \cdot h\left(\Lambda_{0}\right) \cdot\left(M\left(-\Lambda_{0}\right) \psi_{1}, \psi_{2}\right)_{M}
\end{aligned}
$$

However, by the last lemma, $\lim _{T \rightarrow \infty} \Delta_{\Lambda_{0}}^{T}(h)$ equals

$$
\begin{aligned}
& 2 \pi \cdot \Lambda_{0}^{2 q} \cdot h\left(\Lambda_{0}\right) \cdot\left(c^{+}\left(\Lambda_{0}\right)^{*} c^{+}\left(\Lambda_{0}\right) \psi_{1}, \psi_{2}\right)_{M} \\
& \quad+2 \pi \cdot\left(-\Lambda_{0}\right)^{2 q} \cdot h\left(-\Lambda_{0}\right) \cdot\left(c^{+}\left(\Lambda_{0}\right)^{*} c^{-}\left(-\Lambda_{0}\right) \psi_{1}, \psi_{2}\right)_{M}
\end{aligned}
$$

We can choose $h$ so that $h\left(\Lambda_{0}\right)$ and $h\left(-\Lambda_{0}\right)$ are arbitrary. This proves the formulae

$$
\begin{aligned}
& c^{+}\left(\Lambda_{0}\right)^{*} c^{+}\left(\Lambda_{0}\right)=\left(\pi \cdot c \cdot \operatorname{dim} \mu \cdot \beta\left(\mu, \Lambda_{0}\right)\right)^{-1} \\
& c^{+}\left(\Lambda_{0}\right)^{*} c^{-}\left(-\Lambda_{0}\right)=\left(\pi \cdot c \cdot \operatorname{dim} \mu \cdot \beta\left(\mu, \Lambda_{0}\right)\right)^{-1} \cdot M\left(-\Lambda_{0}\right)
\end{aligned}
$$

14. A condition for irreducibility of $\pi_{\sigma, 0}$. We have now done enough work to prove Theorem $3^{\prime}(\mathrm{b})$. However, we shall postpone this until $\S 15$. In this chapter we shall use the inequality (11.2) and Lemma 41 to give a sufficient condition for the irreducibility of the representation $\pi_{\sigma, 0}$.

LEMMA 42. For any $\psi$ in $L_{\sigma}^{\tau}$, where $\tau$ is in $\mathscr{E}_{K}^{2}$ and $\sigma$ is in $\mathscr{E}_{M}$, the meromorphic functions $c^{+}(\Lambda) \psi(1)$ and $c^{-}(\Lambda) \psi(1)$ have a pole of order at most one at $\Lambda=0$.

Proof. If $\psi$ is fixed, then by (11.2) and (12.3) there is a polynomial $p$ and an $\varepsilon>0$ such that

$$
\begin{align*}
& \left|e^{t \rho\left(H_{0}\right)} E_{\Lambda}\left(\psi: \exp t H_{0}\right)-c^{+}(\Lambda) \psi(1) \cdot e^{i \Lambda t}-c^{-}(\Lambda) \psi(1) \cdot e^{-i \Lambda t}\right|  \tag{14.1}\\
& \quad ₹ p(|\Lambda|) e^{-\varepsilon t}, \quad t>0, \Lambda \neq 0 .
\end{align*}
$$

Therefore, for any $t>0$, the function

$$
e^{t \rho\left(H_{0}\right)} E_{\Lambda}\left(\psi: \exp t H_{0}\right)-c^{+}(\Lambda) \psi(1) \cdot e^{i \Lambda t}-c^{-}(\Lambda) \psi(1) \cdot e^{-i \Lambda t}
$$

has no pole at $\Lambda=0$. However, the function

$$
e^{t \rho\left(H_{0}\right)} E_{\Lambda}\left(\psi: \exp t H_{0}\right)
$$

is regular at $\Lambda=0$, so the function

$$
c^{+}(\Lambda) \psi(1) \cdot e^{i \Lambda t}+c^{-}(\Lambda) \psi(1) \cdot e^{-i \Lambda t}
$$

is also regular at $\Lambda=0$.
Consider the Laurent expansions of the functions $c^{+}(\Lambda) \psi(1) \cdot e^{i \Lambda t}$ and $c^{-}(\Lambda) \psi(1) \cdot e^{-i \Lambda t}$. The coefficients will be functions of $t$ with values in $V_{\tau}$, the space on which $\tau$ acts. Allow $t$ to vary. Suppose that one of the two functions has a pole of order greater than one at $\Lambda=0$. Then it is easy to see from the Laurent expansions that this forces the function

$$
c^{+}(\Lambda) \psi(1) \cdot e^{i \Lambda t}+c^{-}(\Lambda) \psi(1) \cdot e^{-i \Lambda t}
$$

to have a pole at $\Lambda=0$ for some $t>0$. We have a contradiction.
THEOREM 4. If $\sigma$ is in $\mathscr{E}_{M}$ and $\beta(\sigma, 0)=0$, then the representation $\pi_{\sigma, 0}$ of $G$ is irreducible.

Proof. Assume the contrary. Then as we saw in $\S 5$, there is a $\tau$ in $\mathscr{E}_{K}^{2}$, acting on the space $V_{\tau}$, and a nonzero $\psi$ in $L_{\sigma}^{\tau}$ such that $E_{0}\left(\psi: \exp t H_{0}\right)$ vanishes for all $t$. We shall use the inequality (14.1) to obtain a contradiction.

Since $E_{0}\left(\psi: \exp t H_{0}\right)$ vanishes for all $t>0$, we can use the mean value theorem to show that for every $\Lambda, t$,

$$
E_{\Lambda}\left(\psi: \exp t H_{0}\right)=\Lambda \cdot \frac{d}{d \Lambda_{1}} E_{\Lambda_{1}}\left(\psi: \exp t H_{0}\right)
$$

$\Lambda_{1}$ is some real number between 0 and $\Lambda$. However, by Lemma 30, there is a polynomial $p_{1}$ such that for any $t>0$,

$$
\left|e^{t \rho\left(H_{0}\right)} \frac{d}{d \Lambda_{1}} E_{\Lambda_{1}}\left(\psi: \exp t H_{0}\right)\right|<p_{1}\left(\left|\Lambda_{1}\right|\right) \cdot \Xi\left(\exp t H_{0}\right) \cdot t
$$

But by (4.1)

$$
\Xi\left(\exp t H_{0}\right) \cdot t ₹ e^{-t \rho\left(H_{0}\right)} \cdot(1+t)^{d+1}, \quad t>0
$$

Therefore there is a polynomial $p_{2}$ such that for $\Lambda \in \mathbf{R}$ and $t>0$,

$$
\begin{equation*}
\left|e^{t \rho\left(H_{0}\right)} E_{\Lambda}\left(\psi: \exp t H_{0}\right)\right|<|\Lambda| \cdot p_{2}(|\Lambda|) \cdot(1+t)^{d+1} \tag{14.2}
\end{equation*}
$$

Since $\beta(\sigma, 0)=0$, Lemma 41 tells us that the functions $c^{+}(\Lambda) \psi(1)$ and $c^{-}(\Lambda) \psi(1)$ must both have a pole at $\Lambda=0$. This pole must be of order 1 by Lemma 42. Let

$$
c^{ \pm}(\Lambda) \psi(1)=c_{-1}^{ \pm} \Lambda^{-1}+c_{0}^{ \pm}+c_{1}^{ \pm} \Lambda+\cdots
$$

be the Laurent expansions about $\Lambda=0$ for these functions. $c_{-1}^{ \pm}, c_{0}^{ \pm}, c_{1}^{ \pm}, \ldots$ are vectors in $V_{\tau}$, and neither $c_{-1}^{+}$nor $c_{-1}^{-}$can equal zero. If $t>0$ is fixed, we apply Taylor's formula with remainder to the functions

$$
g^{ \pm}(t, \Lambda)=\Lambda c^{ \pm}(\Lambda) \psi(1) e^{ \pm i \Lambda t}
$$

Then there is a real number $\Lambda_{1}$ between 0 and $\Lambda$ such that

$$
g^{+}(t, \Lambda)=c_{-1}^{+}+\Lambda\left(c_{0}^{+}+i c_{-1}^{+} t\right)+\Lambda^{2} \cdot\left(\frac{d}{d \Lambda_{1}}\right)^{2} g^{+}\left(t, \Lambda_{1}\right)
$$

The function $\left(d / d \Lambda_{1}\right)^{2} g^{+}\left(t, \Lambda_{1}\right)$ is equal to

$$
\begin{aligned}
& e^{i \Lambda_{1} t}\left(\frac{d}{d \Lambda_{1}}\right)^{2}\left(\Lambda_{1} c^{+}\left(\Lambda_{1}\right) \psi(1)\right)+e^{i \Lambda_{1} t} \cdot 2 i t \cdot \frac{d}{d \Lambda_{1}}\left(\Lambda_{1} c^{+}\left(\Lambda_{1}\right) \psi(1)\right) \\
&-e^{i \Lambda_{1} t} \cdot t^{2} \cdot\left(\Lambda_{1} c^{+}\left(\Lambda_{1}\right) \psi(1)\right)
\end{aligned}
$$

Then by Lemma 35 there exists a polynomial $p_{3}$ such that

$$
\left|\left(\frac{d}{d \Lambda_{1}}\right)^{2} g^{+}\left(t, \Lambda_{1}\right)\right|<p_{3}\left(\left|\Lambda_{1}\right|\right)(1+t)^{2}
$$

This implies that we can choose a polynomial $p_{4}$ such that for $\Lambda \neq 0$ and $t>0$,

$$
\begin{equation*}
\left|c^{+}(\Lambda) \psi(1) e^{i \Lambda t}-\left(c_{-1}^{+} \Lambda^{-1}+c_{0}^{+}+i c_{-1}^{+} t\right)\right|<|\Lambda| \cdot p_{4}(|\Lambda|) \cdot(1+t)^{2} \tag{14.3}
\end{equation*}
$$

Similarly, we can show that

$$
\begin{equation*}
\left|c^{-}(\Lambda) \psi(1) e^{-i \Lambda t}-\left(c_{-1}^{-} \Lambda^{-1}+c_{0}^{-}-i c_{-1}^{-} t\right)\right| ₹|\Lambda| \cdot p_{4}(|\lambda|) \cdot(1+t)^{2} \tag{14.4}
\end{equation*}
$$

Now for $\Lambda \neq 0$ and $t>0$ the expression

$$
\left|\left(c_{-1}^{+}+c_{-1}^{-}\right) \Lambda^{-1}+\left(c_{0}^{+}+c_{0}^{-}\right)+i\left(c_{-1}^{+}-c_{-1}^{-}\right) t\right|
$$

is bounded by the sum of the following four terms:
(i) $\left|c^{+}(\Lambda) \psi(1) e^{i \Lambda t}-\left(c_{-1}^{+} \Lambda^{-1}+c_{0}^{+}+i c_{-1}^{+} t\right)\right|$,
(ii) $\left|c^{-}(\Lambda) \psi(1) e^{-i \Lambda t}-\left(c_{-1}^{-} \Lambda^{-1}+c_{0}^{-}-i c_{-1}^{-} t\right)\right|$,
(iii) $\left|e^{t \rho\left(H_{0}\right)} E_{\Lambda}\left(\psi: \exp t H_{0}\right)\right|$,
(iv) $\left|e^{t \rho\left(H_{0}\right)} E_{\Lambda}\left(\psi: \exp t H_{0}\right)-c^{+}(\Lambda) \psi(1) e^{i \Lambda t}-c^{-}(\Lambda) \psi(1) e^{-i \Lambda t}\right|$.

By (14.1), (14.2), (14.3), and (14.4), there are polynomials $p$ and $p_{5}$ such that this sum is bounded by

$$
p(|\Lambda|) e^{-\varepsilon t}+|\Lambda| \cdot p_{5}(|\Lambda|) \cdot(1+t)^{d+2}
$$

For a fixed $t$ and. for $0<\Lambda ₹ 1$ this expression is clearly bounded. Therefore

$$
\begin{equation*}
c_{-1}^{+}+c_{-1}^{-}=0 \tag{14.5}
\end{equation*}
$$

Let

$$
P=\sup _{0<\Lambda \gtrless 1} p(\Lambda), \quad P_{5}=\sup _{0<\Lambda \gtrless 1} p_{5}(\Lambda)
$$

Then for any $t>0$ and for $0<\Lambda<1$,

$$
\left|\left(c_{0}^{+}+c_{0}^{-}\right)+i\left(c_{-1}^{+}-c_{-1}^{-}\right) t\right| \gtrless P \cdot e^{-\varepsilon t}+P_{5} \cdot \Lambda \cdot(1+t)^{d+2}
$$

Let $\Lambda=t^{-(d+3)}$, and let $t$ approach $\infty$. Then we see that

$$
c_{-1}^{+}-c_{-1}^{-}=0
$$

Therefore, by (14.5)

$$
c_{-1}^{+}=c_{-1}^{-}=0
$$

We have a contradiction, so Theorem 4 is proved.

## 15. Completion of the proof of Theorem $3^{\prime}(b)$.

LEMMA 43. Suppose that $\psi$ is in $L_{\sigma}^{\tau}$ for $\tau$ in $\mathscr{E}_{K}^{2}, \sigma$ in $\mathscr{E}_{M}$, and assume that $\|\psi\|_{M}=1$. Then $c^{ \pm}(\Lambda) \psi(1) \cdot \beta(\sigma, \Lambda)$ is an infinitely differentiable function from $\mathbf{R}$ into $V_{\tau}$, the space on which $\tau$ acts. Furthermore, for every nonnegative integer $n$ there are polynomials $p_{1}, p_{2}, q$ independent of $\sigma, \Lambda, \tau$ such that

$$
\left|\left(\frac{d}{d \Lambda}\right)^{n}\left(c^{ \pm}(\Lambda) \psi(1) \cdot \beta(\sigma, \Lambda)\right)\right| \gtrless p_{1}(|\sigma|) \cdot p_{2}(|\Lambda|) \cdot q(|\tau|)
$$

Proof. $c^{ \pm}(\Lambda) \psi(1) \cdot \beta(\sigma, \Lambda)$ is equal to $\Lambda^{2 q} c^{ \pm}(\Lambda) \psi(1) \cdot \beta(\sigma, \Lambda) \Lambda^{-2 q}$. This is the restriction to $\mathbf{R}$ of a function $h(\Lambda)$ which is meromorphic in the region $\left|\Lambda_{I}\right|<\delta$. (See Lemma 5, (ii), and the corollary to Lemma 34.) The only possible real pole of $h(\Lambda)$ is at $\Lambda=0$. For any real $\Lambda \neq 0$

$$
\begin{aligned}
\left|c^{ \pm}(\Lambda) \psi(1) \cdot \beta(\sigma, \Lambda)\right| & =\beta(\sigma, \Lambda) \cdot\left(c^{ \pm}(\Lambda) \psi, c^{ \pm}(\Lambda) \psi\right)^{1 / 2} \\
& =\left(c^{ \pm}(\Lambda)^{*} c^{ \pm}(\Lambda) \psi, \psi\right)^{1 / 2} \cdot \beta(\sigma, \Lambda)
\end{aligned}
$$

Therefore, by Lemma 41,

$$
\begin{equation*}
\left|c^{ \pm}(\Lambda) \psi(1) \cdot \beta(\sigma, \Lambda)\right|=(C)^{-1} \cdot(\psi, \psi)^{1 / 2} \cdot \beta(\sigma, \Lambda)^{1 / 2} \tag{15.1}
\end{equation*}
$$

But by Lemma $5, \beta(\sigma, \Lambda)$ is regular at $\Lambda=0$, so $|h(\Lambda)|$ is bounded for all real $\Lambda$ in a neighborhood of 0 . Therefore, $h(\Lambda)$ has no real pole and hence is infinitely differentiable at any real $\Lambda$.

Since $\beta(\sigma, \Lambda)$ is regular at $\Lambda=0$, the functions $c^{ \pm}(\Lambda) \psi(1)$ cannot have a zero at $\Lambda=0$. Let $N(\psi)$ be the order of the pole of the functions $c^{+}(\Lambda) \psi(1)$ and $c^{-}(\Lambda) \psi(1)$ at $\Lambda=0 . N(\psi)$ is a nonnegative integer. We see from (15.1) that $\beta(\sigma, \Lambda)$ has a zero of order $2 N(\psi)$ at $\Lambda=0$. Then write

$$
c^{ \pm}(\Lambda) \psi(1) \beta(\sigma, \Lambda)=\Lambda^{N(\psi)} c^{ \pm}(\Lambda) \psi(1) \cdot \Lambda^{-N(\psi)} \beta(\sigma, \Lambda)
$$

The estimate for $(d / d \Lambda)^{n}\left(c^{ \pm}(\Lambda) \psi(1) \beta(\sigma, \Lambda)\right)$ in the lemma then follows from Lemma 5, (iv), Lemma 35, and Leibnitz' rule.

LEMMA 44. Choose $\psi$ in $L_{\sigma}^{\tau}$ with $\|\psi\|_{M}=1$. Then for every nonnegative integer $s$ there are polynomials $p_{1}, p_{2}, q$ such that for every $h$ in $\mathscr{S}(\mathbf{R})$,

$$
\begin{aligned}
& \sup _{x \in G}\left|\int_{\mathbf{R}} h(\Lambda) E_{\Lambda}(\psi: x) \beta(\sigma, \Lambda) d \Lambda \cdot \Xi(x)^{-1}(1+\sigma(x))^{s}\right| \\
& \quad<p_{1}(|\sigma|) \cdot q(|\tau|) \cdot \sup _{0<i<s} \cdot \sup _{\Lambda \in \mathbf{R}}\left(p_{2}(|\Lambda|)\left|\left(\frac{d}{d \Lambda}\right)^{i} h(\Lambda)\right|\right) .
\end{aligned}
$$

Proof. Every $x$ in $G$ is of the form $k_{1} \cdot \exp t H_{0} \cdot k_{2}$, for $k_{1}, k_{2}$ in $K$, and $t>0$. By (4.1) it is enough to prove the lemma when the left-hand side of the inequality is replaced by

$$
\sup _{t>0}\left|e^{t \rho\left(H_{0}\right)}(1+r t)^{s} \int_{-\infty}^{\infty} E_{\Lambda}\left(\psi: \exp t H_{0}\right) \cdot h(\Lambda) \cdot \beta(\sigma, \Lambda) d \Lambda\right|
$$

If $t>0$, the expression

$$
\left|e^{t \rho\left(H_{0}\right)}(1+r t)^{s} \int_{-\infty}^{\infty} E_{\Lambda}\left(\psi: \exp t H_{0}\right) \cdot h(\Lambda) \cdot \beta(\sigma, \Lambda) d \Lambda\right|
$$

is bounded by the sum of the following two expressions:
(i) $I_{1}(t)=(1+r t)^{s} \cdot \int_{-\infty}^{\infty}|h(\Lambda)| \cdot|\beta(\sigma, \Lambda)|$

$$
\left|e^{t \rho\left(H_{0}\right)} E_{\Lambda}\left(\psi: \exp t H_{0}\right)-c^{+}(\Lambda) \psi(1) e^{i \Lambda t}-c^{-}(\Lambda) \psi(1) e^{-i \Lambda t}\right| d \Lambda
$$

(ii) $I_{2}(t)=\left|\int_{-\infty}^{\infty}\left(c^{+}(\Lambda) \psi(1) e^{i \Lambda t}-c^{-}(\Lambda) \psi(1) e^{-i \Lambda t}\right) h(\Lambda) \beta(\sigma, \Lambda) d \Lambda \cdot(1+r t)^{s}\right|$.

By (11.2) and (12.3), there are polynomials $p_{1}, p_{2}, q$, and a number $\varepsilon>0$, all independent of $\sigma, \Lambda, \tau$, such that

$$
I_{1}(t)<p_{1}(|\sigma|) \cdot q(|\tau|) \cdot(1+r t)^{s} \cdot e^{-\varepsilon t} \cdot \int_{\mathbf{R}} p_{2}(|\Lambda|) \cdot|h(\Lambda)||\beta(\sigma, \Lambda)| d \Lambda .
$$

But by Lemma 5 , (iv), there exist polynomials $p_{1}^{\prime}$ and $p_{2}^{\prime}$, such that

$$
\beta(\sigma, \Lambda) ₹ p_{1}^{\prime}(|\sigma|) \cdot p_{2}^{\prime}(|\Lambda|) .
$$

Therefore

$$
\begin{aligned}
\sup _{t>0} I_{1}(t)< & c_{s} \cdot p_{1}(|\sigma|) \cdot p_{1}^{\prime}(|\sigma|) \cdot q(|\tau|) \\
& \cdot \sup _{\Lambda \in R}\left(p_{2}(|\Lambda|) \cdot p_{2}^{\prime}(|\Lambda|) \cdot\left(1+|\Lambda|^{2}\right) \cdot|h(\Lambda)|\right),
\end{aligned}
$$

where

$$
c_{s}=\int_{\mathbf{R}}\left(1+|\Lambda|^{2}\right)^{-1} d \Lambda \cdot \sup _{t>0}(1+r t)^{s} e^{-\varepsilon t}
$$

This takes care of $I_{1}(t)$.
We now obtain a bound for $I_{2}(t)$. Define

$$
\phi^{ \pm}(\Lambda)=c^{ \pm}(\Lambda) \psi(1) \beta(\sigma, \Lambda) \cdot h(\Lambda)
$$

$\phi^{ \pm}$is in $\mathscr{S}(\mathbf{R}) \otimes V_{\tau}$. We see that

$$
\begin{aligned}
I_{2}(t)< & \left|\int_{-\infty}^{\infty} \phi^{+}(\Lambda) e^{i \Lambda t} d \Lambda\right| \cdot(1+r t)^{s}+\left|\int_{-\infty}^{\infty} \phi^{-}(\Lambda) e^{-i \Lambda t} d \Lambda\right| \cdot(1+r t)^{s} \\
= & \left|\int_{-\infty}^{\infty} e^{i \Lambda t}\left(1+i r \cdot \frac{d}{d \Lambda}\right)^{s} \phi^{+}(\Lambda) d \Lambda\right| \\
& +\left|\int_{-\infty}^{\infty} e^{-i \Lambda t}\left(1-i r \cdot \frac{d}{d \Lambda}\right)^{s} \phi^{-}(\Lambda) d \Lambda\right| .
\end{aligned}
$$

Now by Liebnitz' rule and Lemma 43 there are polynomials $p_{1}, p_{2}, q$ such that

$$
\begin{aligned}
& \sup _{\Lambda \in \mathbf{R}}\left|\left(1+|\Lambda|^{2}\right)\left(1 \pm i r \cdot \frac{d}{d \Lambda}\right)^{s} \phi^{ \pm}(\Lambda)\right| \\
& \quad ₹ p_{1}(|\sigma|) \cdot q(|\tau|) \cdot \sup _{1 \gtrless i \gtrless s} \cdot \sup _{\Lambda \in \mathbf{R}}\left(p_{2}(|\Lambda|) \cdot\left|\left(\frac{d}{d \Lambda}\right)^{i} h(\Lambda)\right|\right) .
\end{aligned}
$$

Therefore, we see that

$$
I_{2}(t) \gtrless c \cdot p_{1}(|\sigma|) \cdot q(|\tau|) \cdot \sup _{1 \gtrless i \gtrless s} \cdot \sup _{\Lambda \in \mathbf{R}}\left(p_{2}(|\Lambda|) \cdot\left|\left(\frac{d}{d \Lambda}\right)^{i} h(\Lambda)\right|\right)
$$

where $c=2 \int_{-\infty}^{\infty}\left(1+|\Lambda|^{2}\right)^{-1} d \Lambda$. We have proved Lemma 44 .

Corollary. Suppose $\psi$ is in $L_{\sigma}^{\tau}$ and $\|\psi\|_{M}=1$. Then for every $g_{1}, g_{2}$ in $\mathscr{B}$, and every nonnegative integer $s$, there are polynomials $p_{1}, p_{2}, q$ such that for every $h$ in $\mathscr{S}(\mathbf{R})$,

$$
\begin{aligned}
& \sup _{x \in G}\left|\int_{\mathbf{R}} h(\Lambda) E_{\Lambda}\left(\psi: g_{1} ; x ; g_{2}\right) \cdot \beta(\sigma, \Lambda) d \Lambda \cdot \Xi(x)^{-1}(1+\sigma(x))^{s} d x\right| \\
& \quad<p_{1}(|\sigma|) \cdot \dot{q}(|\tau|) \cdot \sup _{1<i \gtrless s} \cdot \sup _{\Lambda \in \mathbf{R}}\left(p_{2}(|\Lambda|)\left|\left(\frac{d}{d \Lambda}\right)^{i} h(\Lambda)\right|\right)
\end{aligned}
$$

Proof. By (11.1)

$$
\left|E_{\Lambda}\left(\psi: g_{1} ; x ; g_{2}\right)\right|^{2}=\sum_{\alpha=1}^{t_{1}} \sum_{\beta=1}^{t_{2}}\left|\left(\Psi_{1 \alpha}, \pi_{\sigma, \Lambda}(x) \Psi_{2 \beta}\right)\right|^{2}
$$

where $\Psi_{1 \alpha}$ and $\Psi_{2 \beta}$ are the same as in (11.1). Now let $E_{\Lambda}\left(\psi_{\alpha, \beta}: x\right)$ be the ( $\tau_{1 \alpha}, \tau_{2 \beta}$ )-spherical function corresponding to ( $\Psi_{1 \alpha}, \pi_{\sigma, \Lambda}(x) \Psi_{2 \beta}$ ). Apply Lemma 44 to $E_{\Lambda}\left(\psi_{\alpha, \beta}: x\right)$. The corollary follows from the conditions on $\left\{\Psi_{1 \alpha}\right\},\left\{\Psi_{2 \beta}\right\}$ and $\left\{\tau_{1 \alpha}\right\},\left\{\tau_{2 \beta}\right\}$ given in Lemma 13.

This corollary establishes the proof of Theorem $3^{\prime}(\mathrm{b})$. The proof of Theorem 3 is now complete.
16. Tempered distributions. Having proved Theorem 3, we can now extend the definition of Fourier transform to tempered distributions on $G$.

A distribution on $G$ is said to be tempered if it extends to a continuous linear functional from $\mathscr{C}(G)$ to $C$. Since $C_{0}^{\infty}(G)$ is dense in $\mathscr{C}(G)$, and since the inclusion map

$$
C_{0}^{\infty}(G) \bar{\subset} \mathscr{C}(G)
$$

is continuous, we can regard the space of tempered distributions as the dual space of $\mathscr{C}(G)$, that is, the space of continuous linear functionals from $\mathscr{C}(G)$ into $\mathbf{C}$.

Let $\mathscr{C}^{\prime}(G)$ be the set of tempered distributions on $G$. It becomes a locally convex topological vector space when endowed with the weak topology. (A base for the weak topology of $\mathscr{C}^{\prime}(G)$ is given by the seminorms $\left\{\|\cdot\|_{f}: f \in \mathscr{E}(G)\right\}$, where if $T$ is in $\mathscr{C}^{\prime}(G),\|T\|_{f}=|T(f)|$.)

Let $\mathscr{C}^{\prime}(\hat{G})$ be the dual space of $\mathscr{C}(\hat{G})$. Then

$$
\mathscr{C}^{\prime}(\hat{G})=\mathscr{C}_{0}^{\prime}(\hat{G}) \oplus \mathscr{C}_{1}^{\prime}(\hat{G})
$$

where $\mathscr{C}_{0}^{\prime}(\hat{G})$ and $\mathscr{C}_{1}^{\prime}(\hat{G})$ are the dual spaces of $\mathscr{C}_{0}(\hat{G})$ and $\mathscr{C}_{1}(\hat{G})$ respectively. Endow $\mathscr{C}^{\prime}(\hat{G}), \mathscr{C}_{0}^{\prime}(\hat{G})$, and $\mathscr{C}_{1}^{\prime}(\hat{G})$ each with the weak topology induced from $\mathscr{E}(\hat{G}), \mathscr{C}_{0}(\hat{G})$, and $\mathscr{E}_{1}(\hat{G})$ respectively.

THEOREM 5. Denote the map $f \rightarrow \hat{f}$ of $\mathscr{C}(G)$ onto $\mathscr{C}(\hat{G})$ by $\mathscr{F} f$. Then $\mathscr{F}^{*}$, the transpose of $\mathscr{F}$, is a topological isomorphism from $\mathscr{C}^{\prime}(\hat{G})$ onto $\mathscr{C}^{\prime}(G)$.

Proof. The theorem follows directly from the fact that $\mathscr{F}$ is a topological isomorphism.

It is of interest to obtain a slightly different characterization of the space $\mathscr{C}_{1}^{\prime}(\hat{G})$. Define the space $\mathscr{S}_{1}(\hat{G})$ as in $\S 7$. For $a_{1}$ in $\mathscr{S}_{1}(\hat{G})$ define the function

$$
\check{a}_{1}: G \rightarrow \mathbf{C}
$$

as in $\S 7$. Then

$$
\check{a}_{1}(x)=\sum_{1} \int_{-\infty}^{\infty}\left(a_{1}(\sigma, \Lambda) \Phi_{\tau_{2}, i_{2}}, \Phi_{\tau_{1}, i_{1}}\right)\left(\Phi_{\tau_{1}, i_{1}}, \pi_{\sigma, \Lambda}(x) \Phi_{\tau_{2}, i_{2}}\right) \cdot \beta(\sigma, \Lambda) d \Lambda
$$

$\check{a}_{1}(x)$ is in $\mathscr{C}(G)$ by Theorem $3^{\prime}(\mathrm{b})$.
For $a_{1}$ in $\mathscr{S}_{1}(\hat{G})$ define

$$
\left(S a_{1}\right)(\sigma, \Lambda)=\frac{1}{2}\left[a_{1}(\sigma, \Lambda)+N_{\sigma}(\Lambda)^{-1} a_{1}\left(\sigma^{\prime},-\Lambda\right) N_{\sigma}(\Lambda)\right]
$$

Then by (7.3)

$$
\hat{\hat{a}}_{1}(\sigma, \Lambda)=\left(S a_{1}\right)(\sigma, \Lambda)
$$

Therefore, by Theorem $3^{\prime}(\mathrm{b})$ and Lemma 10 , the map

$$
a_{1} \rightarrow S a_{1}, \quad a_{1} \in \mathscr{S}_{1}(\hat{G})
$$

is a continuous transformation from $\mathscr{S}_{1}(\hat{G})$ into $\mathscr{E}_{1}(\hat{G})$. If $a_{1}$ is in $\mathscr{C}_{1}(\hat{G}), S a_{1}=$ $a_{1}$. Therefore $S$ is a continuous projection from $\mathscr{P}_{1}(\hat{G})$ onto $\mathscr{C}_{1}(\hat{G})$. Since $\mathscr{C}_{1}(\hat{G})$ is the kernel of the continuous map $S-1$ on $\mathscr{S}_{1}(\hat{G}), \mathscr{C}_{1}(\hat{G})$ is a closed subspace of $\mathscr{S}_{1}(\hat{G})$.

Let $\mathscr{S}_{1}^{\prime}(\hat{G})$ be the space of distributions on $\mathscr{S}_{1}(\hat{G})$. We can regard the map $S$ as going from $\mathscr{S}_{1}(\hat{G})$ to either $\mathscr{S}_{1}(\hat{G})$ or $\mathscr{E}_{1}(\hat{G})$. In either case let $S^{*}$ be its transpose. Then if $T$ is in either $\mathscr{S}_{1}^{\prime}(\hat{G})$ or $\mathscr{C}_{1}^{\prime}(\hat{G}), S^{*} T$ is a distribution in $\mathscr{S}_{1}^{\prime}(\hat{G})$, and

$$
\left(S^{*} T\right)\left(a_{1}\right)=T\left(S a_{1}\right), \quad a_{1} \in \mathscr{S}_{1}(\hat{G})
$$

THEOREM 6. $S^{*}$ is a canonical isomorphism from $\mathscr{C}_{1}^{\prime}(\hat{G})$ onto the closed subspace

$$
\overline{\mathscr{S}_{1}^{\prime}}(\hat{G})=\left\{T \in \mathscr{S}_{1}^{\prime}(\hat{G}): S^{*} T=T\right\}
$$

of $\mathscr{S}_{1}^{\prime}(\hat{G})$.
Proof. Clearly $S^{*}$ is a one-to-one map of $\mathscr{C}_{1}^{\prime}(G)$ into $\overline{\mathscr{S}_{1}^{\prime}}(\hat{G})$. Now suppose that $T$ is in $\overline{\mathscr{S}}_{1}^{\prime}(\hat{G})$. We can define a new distribution $T_{1}$ in $\mathscr{C}_{1}^{\prime}(\hat{G})$ as the restriction of $T$ to the closed subspace $\mathscr{C}_{1}(\hat{G})$ of $\mathscr{S}_{1}(\hat{G})$. If $a_{1}$ is in $\mathscr{P}_{1}(\hat{G})$,

$$
\left(S^{*} T_{1}\right)\left(a_{1}\right)=T_{1}\left(S a_{1}\right)=T\left(S a_{1}\right)=\left(S^{*} T\right)\left(a_{1}\right)=T a_{1}
$$

Therefore the map

$$
S^{*}: \mathscr{C}_{1}^{\prime}(\hat{G}) \rightarrow \overline{\mathscr{S}_{1}^{\prime}}(\hat{G})
$$

is surjective. Our theorem is proved.
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