

Hierarchical stripe phases in IrTe₂ driven by competition between Ir dimerization and Te bondingJixia Dai,^{1,2} Kristjan Haule,¹ J. J. Yang,³ Y. S. Oh,^{1,2} S-W. Cheong,^{1,2,3} and Weida Wu^{1,2,*}¹*Department of Physics and Astronomy, Rutgers University, Piscataway, New Jersey 08854, USA*²*Rutgers Center for Emergent Materials, Rutgers University, Piscataway, New Jersey 08854, USA*³*Laboratory for Pohang Emergent Materials and Department of Physics, Pohang University of Science and Technology, Pohang 790-784, Korea*

(Received 18 August 2014; published 11 December 2014)

Layered 5d transition metal dichalcogenide (TMD) IrTe₂ is distinguished from the traditional TMDs (such as NbSe₂) by the existence of multiple charge-density wave (CDW)-like stripe phases and superconductivity at low temperatures. Despite intensive studies, there is still no consensus on the physical origin of the stripe phases or even the ground state modulation for this 5d material. Here we present atomic-scale evidence from scanning tunneling microscopy and spectroscopy (STM/STS) that the ground state of IrTe₂ is a $q = 1/6$ stripe phase, identical to that of the Se-doped compound. Furthermore, our data suggest that the multiple transitions and stripe phases are driven by the intralayer Ir-Ir dimerization that competes against the interlayer Te-Te bonding. The competition results in a unified phase diagram with a series of hierarchical modulated stripe phases, strikingly similar to the renowned “devil’s staircase” phenomena.

DOI: [10.1103/PhysRevB.90.235121](https://doi.org/10.1103/PhysRevB.90.235121)

PACS number(s): 71.45.Lr, 71.27.+a, 74.55.+v

One major challenge in condensed matter physics is to understand the intriguing interplay between unconventional superconductivity and other modulated electronic ground states, e.g., antiferromagnetism, charge/orbital ordering, stripes, structure distortion, spin/charge-density wave (SDW/CDW), etc. Clarifying the microscopic nature of these phases is crucial for understanding the mechanism of unconventional superconductivity [1,2]. Layered TMD have been an active playground for exploring the interplay between CDW and superconductivity in the last several decades [3–6]. Yet it is still under debate what is responsible for the CDW transitions in these materials [7–10]. This exploration is further fueled by the recently discovered superconductivity [11–14] emerging from the suppression of a CDW-like stripe phase transition ($T_C \approx 280$ K) with $q = 1/5$ along the (1,0,−1) direction in IrTe₂ [11,12,15]. The qualitative similar phase diagram between IrTe₂ and other unconventional superconductors suggests that the driving mechanism of the stripe phase may hold the key to understand the superconductivity. Recent STM and transport studies suggest that the CDW-like transition in IrTe₂ is followed by another one near $T_S \approx 180$ K (when cooled) with further increase in electric resistivity [16], a phenomenon similar to the multistep CDW transitions in the quasi-1D system NbSe₃ [17] and quasi-2D TaS₂ [6]. Unlike these two systems, IrTe₂ is unique in terms of single- q (i.e., stripe) modulations with unusually large hysteretic transitions [16], implying a distinctive but yet to be unveiled driving mechanism behind them.

Although many candidate mechanisms, ranging from Fermi surface nesting [11], Ir bonding instability [12,18], crystal field effect, or interlayer bonding of Te p orbitals [13,14,19], have been proposed based on various experimental observations, none of them explains all the experimental results satisfactorily. For example, the absence of CDW gap in angle-resolved photoemission spectroscopy (ARPES) and optical spectroscopy measurements excludes the Fermi surface instability mechanism [13,20]. The central question of the

heated debate is whether Ir or Te plays the key role in the driving mechanism. This is partially due to the complication that the single- q modulation of the stripe phases results in multidomain states below the transition, which washes out the intrinsic properties of spatially averaged experimental observation. Indeed, the atomic position of the $q = 1/5$ stripe phase was unambiguously resolved with single crystal x-ray refinement on a single domain sample [21]. Yet this technique cannot address the emergence of short-range multiple modulations below T_S . Therefore, real-space resolving techniques, especially at low temperature, are necessary for addressing the fundamental mechanism.

Previous STM studies revealed that the second transition at T_S is a hysteretic transition from a periodic soliton lattice ($q = 1/5$) to a partially melted soliton state consisting of emergent stripe modulations of $q_n = (3n + 2)^{-1}$, with gradual reduction of soliton density down to $T = 50$ K [16]. A natural extrapolation ($n \rightarrow \infty$) of this behavior indicates the ground state of IrTe₂ is a soliton-free $q = 1/6$ modulated phase, in contrast to previous studies [11,19,22]. Thus, a clarification of the IrTe₂ ground state is imperative for understanding the fundamental mechanism that drives the transition. In this work we performed low-temperature STM/STS studies of undoped, selenium- and platinum-doped IrTe₂. Our high resolution topographic and spectroscopic data provide strong evidence that the low temperature phase of IrTe₂ is a $q = 1/6$ phase, the same as that in the Se-doped IrTe₂ [19]. Combining the STM/STS measurements and first-principle calculations, our results support that the competition between Ir-Ir dimerization and interlayer Te-Te bonding is the fundamental driving force of the hierarchical CDW-like phase transitions in IrTe₂. Such intriguing phenomena are closely related to the emergence of hierarchical modulations due to competing microscopic interactions in the famous “devil’s staircase” phenomena [23].

In this study we used single crystals of IrTe₂, IrTe_{2-x}Se_x ($x = 0.34$), and Pt_xIr_{1-x}Te₂ ($x = 0.04$ and 0.1) grown using Te flux, as described in Ref. [19]. For STM/STS measurements, an Omicron LT-STM with base pressure of 1×10^{-11} mbar was used. Electrochemically etched tungsten tips were

*Corresponding author: wdu@physics.rutgers.edu

calibrated on clean Au (111) surface before STM experiments on IrTe₂ (see Supplemental Material Fig. S5 [24]). The single crystal samples were mounted on the sample plates using conducting silver epoxy (H20E from EPO-TEK). After being introduced into the STM chamber, the samples were cleaved *in situ* and loaded into the STM scan head within 1 min. The temperature of the scan head was held at 4.5 K by liquid helium. The differential conductance measurements were performed with the standard lock-in technique with amplifier gain $R_{\text{gain}} = 3 \text{ G}\Omega$, modulation frequency $f = 455 \text{ Hz}$, and modulation amplitude $V_{\text{mod}} = 5\text{--}10 \text{ mV}$. To model the LDOS measured by STS, we have carried out the density functional theory calculation, using the WIEN2K package [25]. The p_z partial density of states was obtained by the DMFT-WIEN2K package [26], where spin rotations (in the presence of large spin-orbit coupling) are properly implemented. The local coordinate system on Te atoms is chosen such that p_z orbital points perpendicular to the surface Te layer, which maximizes its coupling to the tunneling tip. We used the crystal structure for IrTe₂ in the 1/6 phase determined in Ref. [27].

Figures 1(a) and 1(b) show the atomically resolved STM images of IrTe₂ and IrTe_{2-x}Se_x ($x = 0.34$), respectively. The contribution from electronic (charge) modulations was removed to reveal atomic positions of Te atoms (see Supplemental Material Fig. S1 [24]). At 4.5 K, the surface of IrTe₂

is dominated by uniform 1/6 modulation (Fig. S2) except for a small fraction of areas containing dilute solitons [Fig. 3(a)]. The dimer stripes pattern is clearly visible. Similarly, long range 1/6 modulation was observed [Fig. 1(b)] on the surface of IrTe_{2-x}Se_x, even in the presence of Te/Se atomic disorder. The atomic scale inhomogeneity is removed by unit-cell averaging, resulting in images [Figs. 1(c) and 1(e)] showing practically identical dimer stripe pattern. This strongly suggests that the 1/6 modulation in these two compounds are the same. Furthermore, the “bond” length of Te dimers (the red dashed curves) is about 3.2 Å, which is ~18% shorter than that in the high temperature phase. This is in excellent agreement with the results of single crystal x-ray refinement of the Te dimers in both 1/5 and 1/6 phases [21,27], further indicating the dimers in the 1/5 and the 1/6 states are of the same origin. (See Supplemental Material Fig. S3 for STM images of the 1/5 state [24].) As shown in the simplified atomic structure model of 1/6 supercell [Fig. 1(d)], the Te atoms shift collectively to generate the 3 + 3' periodic structure. The second half of supercell (3') is related to the first half (3) through symmetry operations of a mirror plus a translation along the horizontal plane. Therefore, the primary 1/6 peak is absent due to a structural factor cancellation, while the satellite 1/6 peak is observed near (010) as shown in the Fourier map [Fig. 1(f)] and line profiles [Fig. 1(g)], consistent with previous TEM

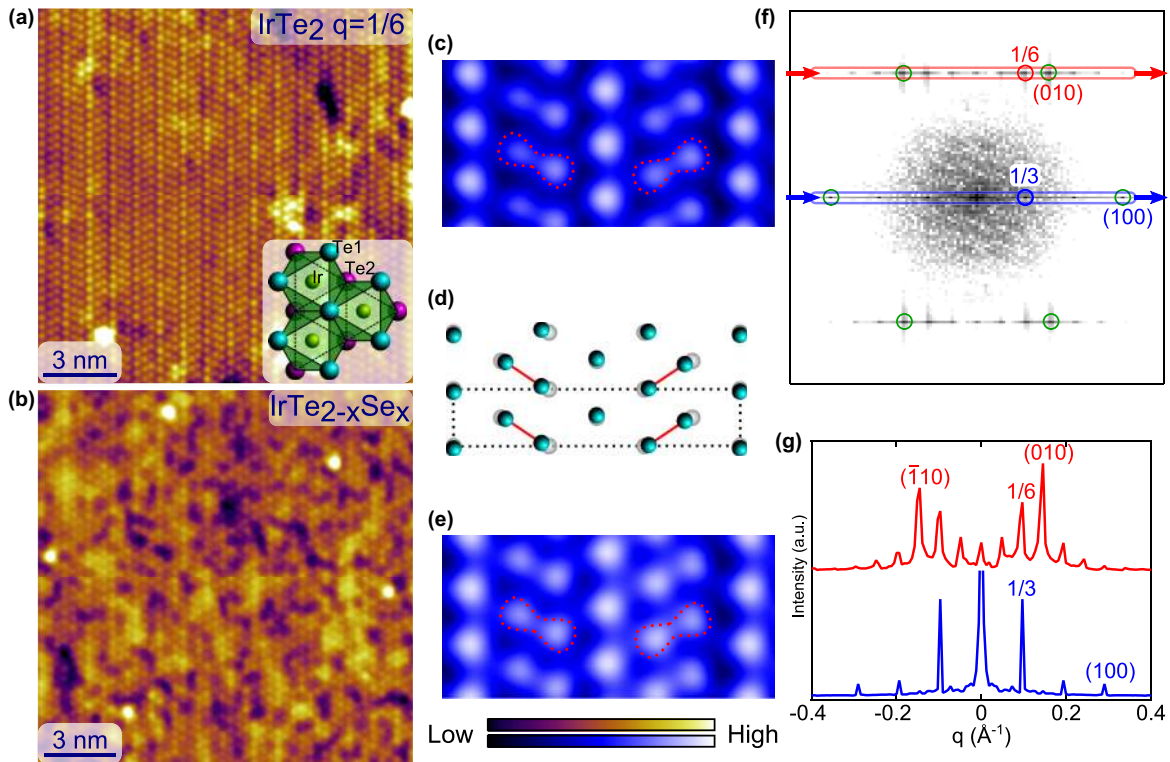


FIG. 1. (Color online) Visualizing the Te-Te dimer stripes on the surface of IrTe₂ and IrTe_{2-x}Se_x. (a) and (b) Topographic images of the $q = 1/6$ states in IrTe₂ and IrTe_{2-x}Se_x at 4.5 K without the 1/3 charge density modulation. The inset of (a) shows the crystal structure of IrTe₂. The atoms in the STM images correspond to the top Te layer (Te1). (c) and (e) The unit-cell-averaged images of (a) and (b), respectively. Here the alternating Te-Te dimers are clearly visible (indicated by the red dashed curves). (d) Schematic diagram of the unit cell for the 1×6 structure, illustrating individual atomic shifts and the Te-Te dimers. The gray/cyan spheres represent the undistorted/distorted positions of top Te atoms. (f) The Fourier transform of the topographic image (including the 1/3 modulation) of IrTe₂ in the area shown in (a). A 1/6 peak near (010) is clearly seen. (g) The line profiles along the red and blue arrows show the peaks (1/6 and 1/3) associated with the $q = 1/6$ modulation. The tunneling set up conditions for (a) and (b) are $V_{\text{bias}} = 5 \text{ mV}$ and $I_{\text{set}} = 5/15 \text{ nA}$.

studies of IrTe_{2-x}Se_x [19]. Accordingly, the primary electronic (charge) modulation is 1/3 (Supplemental Material Fig. S1 [24]).

The excellent agreement between the two 1/6 states in the undoped and the Se-doped IrTe₂ is further corroborated by STS measurements (i.e., differential tunneling conductance dI/dV), which probe local density of states (LDOS) [28]. Figure 2(a) [2(c)] shows the unprocessed STM topography of IrTe₂ (IrTe_{2-x}Se_x). The spatially averaged STS data taken on Fig. 2(a) [2(c)] is shown in Fig. 2(b) [2(d)]. Over a wide energy range (± 2 eV), the dI/dV spectra from these two compounds are qualitatively the same: sharing both a clifflike particle-hole asymmetry near Fermi level and a pronounced peak near 1.5 eV. Even on the areas of dilute solitons in IrTe₂, a similar spectrum is observed between the solitons [red box in Fig. 2(c) and the corresponding dI/dV in Fig. 2(g)], confirming previous conjecture that 1/6 phase is the soliton free limit of $q_n = (3n + 2)^{-1}$ modulations [16]. High spatial

resolution STS measurements reveal further variation of LDOS at different atomic rows. As shown in Fig. 2(e), the STS data of individual atomic columns [defined in Fig. 2(f)] show LDOS variation, indicating that the LDOS of column 2 (which is above Ir dimers) is lower than those of column 1 and 3 between Fermi energy (E_F) and ~ 1.6 eV. These results are in qualitative agreement with the calculated LDOS of Te p_z orbitals shown in Fig. 2(g), because they dominate the tunneling process due to orbital orientation. The difference between theory and experimental data may originate from energy dependence of a tunneling matrix element.

However, the LDOS of soliton crest [averaged over the blue box in Fig. 3(a)] is markedly different: showing a pronounced “hump” at the Fermi level accompanied by spectral-weight suppression from 0.5 eV and above [Fig. 3(b)]. The observed reduction of density of states near E_F in the dimer phase agrees with previous first-principle calculation [21]. Consistently, resistivity jumps were observed at the two transitions during

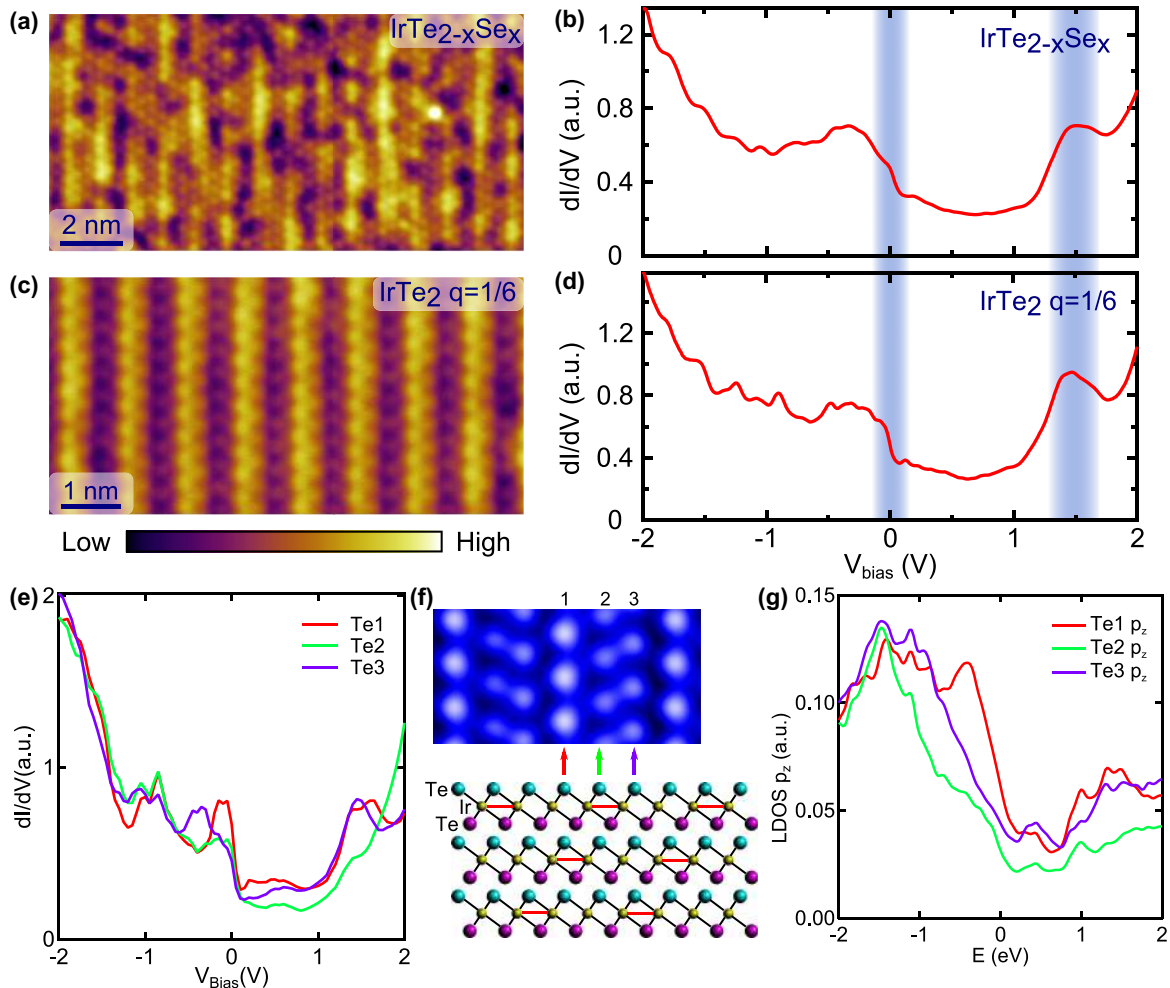


FIG. 2. (Color online) LDOS of IrTe_{2-x}Se_x and IrTe₂ at 4.5 K. (a) and (b) Typical topographic image (4.5 K) of IrTe_{2-x}Se_x and IrTe₂, respectively. These suggest that the ground state of IrTe_{2-x}Se_x is a long range ordered $q = 1/6$ phase, even in the presence of Te/Se disorder. (c) and (d) Averaged dI/dV spectroscopy data (LDOS) with large energy range of IrTe_{2-x}Se_x and IrTe₂, respectively, indicating qualitatively identical LDOS. (e) STS of individual atomic columns [defined in (f)] within half period of the 1/6 phase. (f) Unit-cell-averaged STM image with labels of atomic columns (1, 2, 3) and a cartoon of IrTe₂ lattice (side view). The crystal orientation is determined by tracing the shift of the stripes across a step edge on the same surface (see Fig. S4 [24]). (g) Calculated LDOS of Te p_z orbitals on the top Te layer [defined in (f)], while the p_z orbital is pointing out of plane. Here the spectroscopic data were measured over a very wide high energy range (± 2 V).

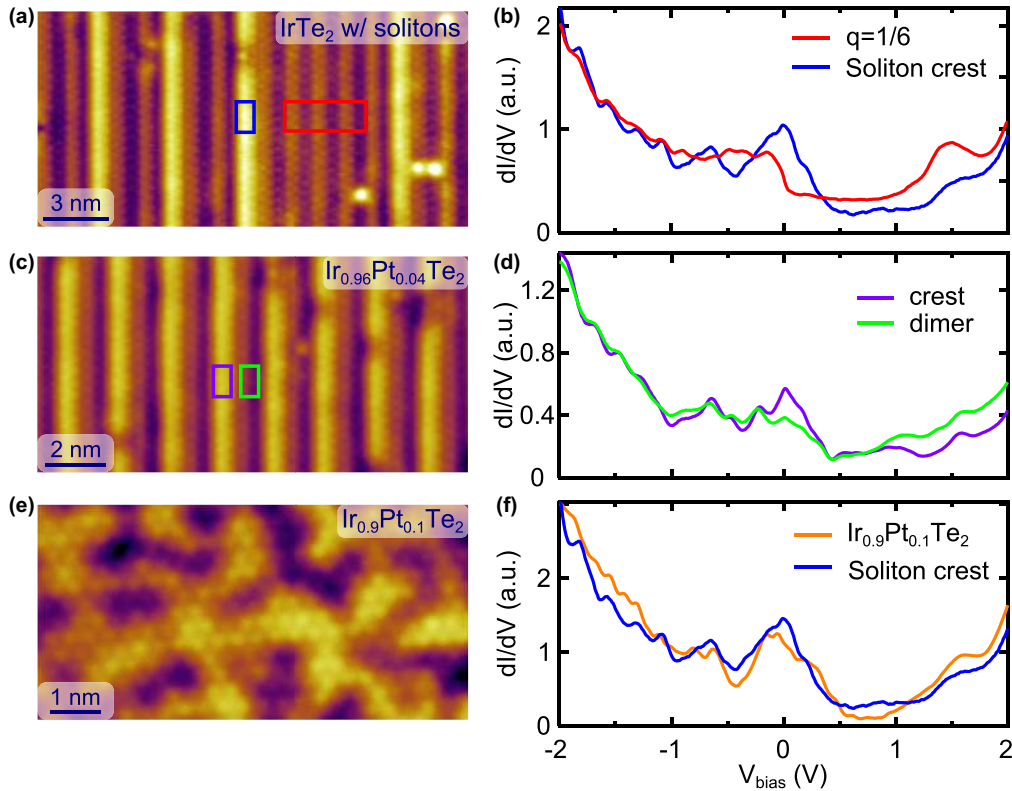


FIG. 3. (Color online) LDOS of soliton crest in IrTe_2 and Pt-doped IrTe_2 at 4.5 K. (a) Diluted solitons occasionally observed in IrTe_2 at 4.5 K. (b) Areal averaged STS of soliton crest (blue) in comparison with that of $1/6$ phase (red). (c) Typical topographic image of slightly Pt doped IrTe_2 , where the $1/5$ phase is stabilized at 4.5 K. (d) Column-averaged STS data of crest (purple) and valley (green) in the $1/5$ phase. The LDOS of soliton (crest) at Fermi energy E_F is higher than that of valley (dimerized stripes). (e) Typical topographic image of $\sim 10\%$ Pt doped IrTe_2 . The stripe modulated phases is completely suppressed with significant Pt doping, resulting in the high temperature phase (1×1). (f) Averaged STS data of $\text{Ir}_{0.9}\text{Pt}_{0.1}\text{Te}_2$ (orange) in comparison with that of soliton crest in (a) (blue). The qualitative similarity between these data suggests that the soliton (crest) is undimerized atomic rows.

cooling as the dimer density increases [16]. Additionally, there is no gap feature around zero bias in any of the above spectra, in agreement with previous angle-resolved photoemission spectroscopy (ARPES) and optical measurements [13,20].

The striking similarity between the low temperature $1/6$ stripe phases between undoped and Se-doped IrTe_2 clearly demonstrate that they share the same origin. Furthermore, the transition temperature (T_C) is enhanced smoothly with Se doping [19]. Together these observations suggest that Te plays a secondary, likely a competing role in the CDW-like transition. Indeed, a completely different scenario emerges in Pt-doped compound $\text{Pt}_x\text{Ir}_{1-x}\text{Te}_2$. At low doping as $x = 0.04$, the $q = 1/5$ stripe phase is stabilized at low temperature [Fig. 3(c)], indicating that the soliton melting transition (T_S) is suppressed. When the Pt concentration is further increased ($x = 0.1$), transport measurements show that the CDW-like transitions is suppressed while superconductivity emerges in the vicinity [12,13,29]. Consequently, the high temperature trigonal symmetry (1×1 phase) is observed in our STM measurements of $x = 0.1$ samples [Fig. 3(e)], demonstrating the absence of the long range order of any stripe modulations. Furthermore, the dI/dV spectrum measured on this 1×1 phase [Fig. 3(f)] is very similar to that of soliton crests (undimerized stripes), confirming previous theoretical prediction that these undimerized stripes are essentially the

high temperature phase [21]. The strong suppression of stripe phases by Pt doping is in sharp contrast to the enhancement of T_C with Se doping, which provides compelling evidence that Ir dimerization plays the dominant role on driving the transition.

Compiled with previous studies, our STM results can be summarized in a unified phase diagram (Fig. 4) showing the antagonistic behaviors of T_C with Se and Pt doping. The salient feature of this phase diagram is the common ground state ($1/6$ phase) of Se doped and undoped IrTe_2 . Previous studies emphasize the importance of Te orbitals, especially interlayer Te-Te bonding (polymerization) in the high temperature 1×1 phase [13,14,9,30]. However, it was not elucidated what breaks the Te-Te bonding below T_C . Using single-crystal x-ray refinement, Pascut *et al.* observed striking intralayer Ir-Ir dimer formation with $>20\%$ bond length reduction, stronger than that of Te-Te dimer (10%–17%) in the $1/5$ stripe phase. This bond length change of Ir-Ir dimer is even larger than that in CuIr_2S_4 [31] and VO_2 [32–34] where dimerization is associated with their metal-insulator transitions. Clearly it is the Ir dimerization that competes against the interlayer Te-Te bonding and drives the system into the dimer stripe phases. Competing microscopic interactions often leads to phenomena of emergent periodicities such as incommensurate density waves and devil's staircase [23,35,36]. Therefore, this

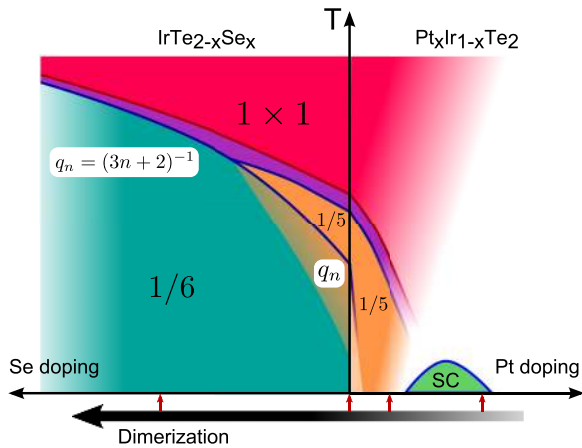


FIG. 4. (Color online) A unified-phase diagram of doped IrTe₂. The first order phase transition T_C is enhanced when Te is partially replaced by Se, suggesting that Se doping strengthens the Ir-Ir dimerization. The large hysteretic transition T_S is suppressed by the enhancement of dimerization. In contrast, Pt doping weakens dimerization tendency so that T_C is suppressed and superconductivity emerges at low temperature (~ 3 K).

mechanism naturally explains the appearance of the $q = 1/5$ phase and the subsequent $q_n = (3n + 2)^{-1}$ stripe phases below 180 K, namely they are the staircase phases preceding the ultimate ground state ($q = 1/6$). Se doping weakens the interlayer Te-Te bonding, resulting in an increase of T_C and likely a suppression of the temperature window of the $q_n = (3n + 2)^{-1}$ stripe phases. In contrast, Pt substitution of Ir directly perturbs the dimerization, resulting in a dramatic suppression of T_C [12,13] and enhanced stability of the

intermediate ($1/5$) phase. Interestingly, superconductivity appears near the suppression of long range ordering of the stripe phase, indicating fluctuating Ir dimers might facilitate Cooper pairing of electrons.

In summary, our results of atomic scale imaging and spectroscopic measurements suggest that the ground state for IrTe₂ is the $q = 1/6$ state with the alternating dimer pattern, while doping into the Ir sites weakens the tendency for the system to form dimers. Nonetheless, comprehensive understanding of the ground states and their relation to the superconductivity would require further experimental and theoretical efforts. Due to the stripe and thus the domain formation, spatially averaged measurements, such as transport, quantum oscillation, ARPES, and scattering experiments, would have extra complication with data interpretation. Measurements within a single domain would provide more straightforward information, such as the Fermi surface, conductivity anisotropy, and the unusual crosslayer quasi-2D electronic behavior. We expect the single domain phase to be achieved by straining the crystal while cooling through the transitions and/or repetitive thermal cycling, similar to the Fe-base superconductors [37]. Furthermore, by low Se doping, the transition temperature could be tuned around room temperature [19], which is highly desirable for potential applications.

We thank G. L. Pascut and V. Kiryukhin for sharing x-ray results and discussions. The work at Rutgers was supported by NSF Grants No. DMREF-1233349 and No. DMR-0844807. The work at POSTECH was supported by the Max Planck POSTECH/KOREA Research Initiative Program (No. 2011-0031558) through the NRF of Korea funded by the Ministry of Education, Science and Technology.

- [1] M. R. Norman, D. Pines, and C. Kallin, *Adv. Phys.* **54**, 715 (2005).
- [2] J. Paglione and R. L. Greene, *Nat. Phys.* **6**, 645 (2010).
- [3] E. Morosan, H. W. Zandbergen, B. S. Dennis, J. W. G. Bos, Y. Onose, T. Klimczuk, A. P. Ramirez, N. P. Ong, and R. J. Cava, *Nat. Phys.* **2**, 544 (2006).
- [4] T. Yokoya, T. Kiss, A. Chainani, S. Shin, M. Nohara, and H. Takagi, *Science* **294**, 2518 (2001).
- [5] Y. I. Joe *et al.*, *Nat. Phys.* **10**, 421 (2014).
- [6] B. Sipoš, A. F. Kusmartseva, A. Akrap, H. Berger, L. Forro, and E. Tutis, *Nat. Mater.* **7**, 960 (2008).
- [7] M. D. Johannes, I. I. Mazin, and C. A. Howells, *Phys. Rev. B* **73**, 205102 (2006).
- [8] D. W. Shen *et al.*, *Phys. Rev. Lett.* **101**, 226406 (2008).
- [9] F. Weber *et al.*, *Phys. Rev. Lett.* **107**, 107403 (2011).
- [10] S. V. Borisenko *et al.*, *Phys. Rev. Lett.* **102**, 166402 (2009).
- [11] J. J. Yang, Y. J. Choi, Y. S. Oh, A. Hogan, Y. Horibe, K. Kim, B. I. Min, and S. W. Cheong, *Phys. Rev. Lett.* **108**, 116402 (2012).
- [12] S. Pyon, K. Kudo, and M. Nohara, *J. Phys. Soc. Jpn.* **81**, 053701 (2012).
- [13] A. F. Fang, G. Xu, T. Dong, P. Zheng, and N. L. Wang, *Sci. Rep.* **3**, 1153 (2013).
- [14] M. Kamitani, M. S. Bahramy, R. Arita, S. Seki, T. Arima, Y. Tokura, and S. Ishiwata, *Phys. Rev. B* **87**, 180501(R) (2013).
- [15] N. Matsumoto, K. Taniguchi, R. Endoh, H. Takano, and S. Nagata, *J. Low Temp. Phys.* **117**, 1129 (1999).
- [16] P.-J. Hsu, T. Maurer, M. Vogt, J. J. Yang, Y. S. Oh, S. W. Cheong, M. Bode, and W. Wu, *Phys. Rev. Lett.* **111**, 266401 (2013).
- [17] P. Monçeau, N. P. Ong, A. M. Portis, A. Meerschaut, and J. Rouxel, *Phys. Rev. Lett.* **37**, 602 (1976).
- [18] D. Ootsuki *et al.*, *Phys. Rev. B* **86**, 014519 (2012).
- [19] Y. S. Oh, J. J. Yang, Y. Horibe, and S. W. Cheong, *Phys. Rev. Lett.* **110**, 127209 (2013).
- [20] D. Ootsuki *et al.*, *J. Phys. Soc. Jpn.* **82**, 093704 (2013).
- [21] G. L. Pascut *et al.*, *Phys. Rev. Lett.* **112**, 086402 (2014).
- [22] T. Machida *et al.*, *Phys. Rev. B* **88**, 245125 (2013).
- [23] P. Bak, *Rep. Prog. Phys.* **45**, 587 (1982).
- [24] See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevB.90.235121>, for details of removal of $1/3$ stripe modulation in the $1/6$ phase, atomic resolving $1/5$ phase, tip calibration process, etc.
- [25] P. Blaha, K. Schwarz, G. K. H. Madsen, D. Kvasnicka, and J. Luitz, *WIEN2k, An Augmented Plane Wave + Local Orbitals*

- Program for Calculating Crystal Properties* (Techn. Universität Wien, Vienna, Austria, 2001).
- [26] K. Haule, C.-H. Yee, and K. Kim, *Phys. Rev. B* **81**, 195107 (2010).
- [27] G. L. Pascut, T. Biroli, M. J. Gutmann, J. J. Yang, S. W. Cheong, K. Haule, and V. Kiryukhin, *Phys. Rev. B* **90**, 195122 (2014).
- [28] J. A. Strosio, R. M. Feenstra, and A. P. Fein, *Phys. Rev. Lett.* **57**, 2579 (1986).
- [29] S. Y. Zhou *et al.*, *Europhys. Lett.* **104**, 27010 (2013).
- [30] S. Jobic, R. Brec, and J. Rouxel, *J. Alloys Compd.* **178**, 253 (1992).
- [31] P. G. Radaelli *et al.*, *Nature (London)* **416**, 155 (2002).
- [32] J. B. Goodenough, *J. Solid State Chem.* **3**, 490 (1971).
- [33] S. Biermann, A. Poteryaev, A. I. Lichtenstein, and A. Georges, *Phys. Rev. Lett.* **94**, 026404 (2005).
- [34] J. H. Park, J. M. Coy, T. S. Kasirga, C. Huang, Z. Fei, S. Hunter, and D. H. Cobden, *Nature (London)* **500**, 431 (2013).
- [35] J. Villain and M. B. Gordon, *J. Phys. C: Solid State Phys.* **13**, 3117 (1980).
- [36] W. L. McMillan, *Phys. Rev. B* **14**, 1496 (1976).
- [37] J.-H. Chu, H.-H. Kuo, J. G. Analytis, and I. R. Fisher, *Science* **337**, 710 (2012).