

High-performance photonic reservoir computer based on a coherently driven passive cavity

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Reservoir computing is a recent bio-inspired approach for processing time-dependent signals. It has enabled a breakthrough in analog information processing, with several experiments, both electronic and optical, demonstrating state-of-the-art performance for hard tasks such as speech recognition, time series prediction, and nonlinear channel equalization. A proof-of-principle experiment using a linear optical circuit on a photonic chip to process digital signals was recently reported. Here we present a photonic implementation of a reservoir computer based on a coherently driven passive fiber cavity processing analog signals. Our experiment has error rate as low as or lower than previous experiments on a wide variety of tasks, and also has lower power consumption. Furthermore, the analytical model describing our experiment is also of interest, as it constitutes a very simple high-performance reservoir computer algorithm. The present experiment, given its good performance, low energy consumption, and conceptual simplicity, confirms the great potential of photonic reservoir computing for information processing applications ranging from artificial intelligence to telecommunications. © 2015 Optical Society of America

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1. INTRODUCTION

Reservoir computing is a recent bio-inspired method for processing information [1–3] that uses a recurrent dynamical system called a “reservoir” to process time-dependent signals. The internal variables of the dynamical system, also called “reservoir states,” provide a nonlinear mapping of the input into a high-dimensional space. The time-dependent output of the reservoir is then given by a linear combination of the internal variables. The readout weights used to compute this linear combination are optimized so as to minimize the mean square error between the target and the output signal, leading to a simple and easy training process. Moreover, some global parameters can be tuned to get the best performance, depending on the reservoir architecture and on the task. The simplicity and the flexibility of such systems make them powerful to solve a very large range of tasks, such as speech recognition [4], nonlinear channel equalization [3,5–7], detection of epileptic seizures [8], robot control [9], time series prediction [1,3,7,10], financial forecasting, and handwriting recognition; see [10,11] for recent reviews.

Recently, experimental implementations of reservoir computing have provided a breakthrough in analog information processing, and in particular in optical information processing. Several

experiments [12–22] (see also the simulations in [23]) report analog information processing, often with error rate comparable to the best digital algorithms. These works all use as a nonlinear dynamical system a delay line with a single nonlinear node.

Integrated optical reservoir computers (RCs) have also been investigated in simulation [24–28]. Very recently, a proof-of-principle experiment that carried out simple processing of digital signals was demonstrated [29]. This experiment uses a linear optical network made from interconnected delay lines and works with coherent light, thereby encoding information in both the amplitude and the phase of the electromagnetic field. This can improve performance [24,25] compared to systems that use only the light intensity. The necessary nonlinearity was provided by the readout photodiodes that, since they produce a current proportional to the intensity of the light, provide quadratic nonlinearity. The use of nonlinearity in the readout layer of reservoir computers has been previously investigated theoretically [5,6].

Here we combine the advantages of the approach based on a delay dynamical system and of the approach based on a linear optical circuit with quadratic nonlinearity provided by the readout photodiode. We present an experimental implementation of a photonic reservoir computer based on a coherently driven passive

fiber cavity that processes analog signals. Our system exhibits several advantages compared to previous experiments. First, it is much more flexible than the architecture presented in [29], in which many parameters, such as the number of internal variables or the strength of the internal connections, are fixed by the hardware and cannot be tuned to improve performance on specific tasks. In addition, we are able to process analog signals experimentally (contrary to the experiment described in [29], which only processes digital signals, due to the limitations of very high-speed electronics). Nevertheless, the reservoir presented in [29] is able to process information faster thanks to the integration on chip. Second, our reservoir is a passive optical cavity, with very low intracavity losses. We therefore do not need any amplifier or active element in the cavity as in previous experiments based on a delay line [12–22,30] (for instance, in [17], passive nonlinearity was used, but because of high intracavity losses an amplifier was placed inside the cavity). This absence of active elements in the cavity removes a major source of noise, and therefore improves performance. It also decreases the energy consumption of the reservoir. In fact the total optical power injected in the reservoir layer (0.57 mW peak power at the entrance of the cavity) is to our knowledge the lowest used so far.

Most importantly, in all tasks on which we have tested it, our experimental reservoir computer has error rate as low as or lower than all previous experiments [12–17,21,22,29,31] that studied the same tasks, and the same also holds true for many results obtained previously in simulation or using digital algorithms [7,12,22,24,29,32]. Note that the experiment can also be translated into a digital algorithm. This algorithm is of interest by itself as it constitutes a very simple high-performance reservoir computer algorithm.

2. OPERATION PRINCIPLE

In the experiment reported here, the input signal is coded in the amplitude $A_{in}(t)$ of the electromagnetic field. This signal is then sent to the reservoir, which consists of a passive fiber cavity (see Fig. 1). The evolution equation of the amplitude $A(t)$ inside the cavity is given by

$$A(t) = \alpha \exp(j\Delta\varphi)A(t - T) + \beta A_{in}(t), \quad (1)$$

where T is the roundtrip time, α is the feedback gain that is tuned through an intracavity optical attenuator, β is the transmission coefficient of the input fiber coupler, $j = (-1)^{1/2}$, and $\Delta\varphi$ is the phase detuning of the cavity.

As explained in detail in [12,13], successful use of the delay dynamical system architecture as a reservoir computer depends crucially on the method used to encode the input signal. Here we use the method introduced in [13]. The input $u(t)$ is held constant for duration T' using a sample and hold procedure; then this input is multiplied by a mask function $m(t)$, which is periodic with period T' . The interconnections between the internal variables are obtained by desynchronizing the masked input $m(t)u(t)$ and the cavity so $T' \neq T$. In the present case, it is further essential for some tasks to add a bias A_0 to the masked input. The input signal amplitude is therefore equal to $A_{in}(t) = m(t)u(t) + A_0$. The mask function $m(t)$ is taken to be a step function, constant over intervals of duration $\theta = T'/N$, where N is the number of internal variables of the reservoir. The desynchronization is defined by $T' = TN/(N + k)$, where $1 \leq k < N$. All our simulation and exper-

imental results were obtained with $k = 1$; i.e., each internal variable is thus coupled with its nearest neighbor. In this way, we get the richest dynamics possible for this unsynchronized regime (e.g., if k and N have a common divisor, then there would be several independent reservoirs, each with a fraction of the number of neurons. The dynamics would then be less rich, and the performance worse).

In order to discuss in more detail the operation principle of the reservoir, it is useful to introduce a discretized time $t(n, i)$ corresponding to the middle of each interval of duration θ : $t(n, i) = nT' + (i + 1/2)\theta$, where $n \in \mathbb{Z}$ and $i \in [0, N - 1]$. With this notation, the sample and hold input $u(n)$ is only a function of n , and the mask $m(i)$ is only a function of i . We denote $x_i(n) = A(t(n, i))$ the state of the internal variables. The continuous time evolution Eq. (1) can thus be approximated by the discrete time evolution equations:

$$\begin{cases} x_i(n) = \alpha x_{i-k}(n-1) \exp(j\Delta\varphi) + \beta[m_i u(n) + A_0] & k \leq i \leq N \\ x_i(n) = \alpha x_{N+i-k}(n-2) \exp(j\Delta\varphi) + \beta[m_i u(n) + A_0] & 0 \leq i \leq k \end{cases}, \quad (2)$$

with $1 \leq k < N$. Note that when $\Delta\varphi \neq 0$, x_i are complex values and the dynamics of the system is richer. As discussed in the next section, it is often experimentally convenient to carry out a non-linear preprocessing of the input. This small change does not modify the performance of the reservoir, except for the evaluation of the memory capacities for which the results depend on this input signal preprocessing.

All the reservoir states denoted $x_i(n) = A(t(n, i))$ are recovered by a photodiode, which performs a quadratic transformation on $A(t)$, since the photodiode output is proportional to $|x_i(n)|^2$. The output $y(n)$ of the reservoir is taken to be a linear combination of $|x_i(n)|^2$:

$$y(n) = \sum_{i=0}^{N-1} W_i |x_i(n)|^2. \quad (3)$$

Note that Eqs. (2) and (3) are the numerical models we will refer to later in the text.

In the operation of a reservoir computer we distinguish two phases. In the training phase we send inputs for which the target output $y^*(n)$ is known. We then record the values of $|x_i(n)|^2$ and use them to compute the readout weights W_i , by minimizing the mean square error $\langle [y^*(n) - y(n)]^2 \rangle_n$ using Tikhonov regularization [13] (also called ridge regression). In the test phase, the readout weights W_i are kept fixed, and the output signal $y(n)$ is computed using Eq. (3), and compared to the target output $y^*(n)$. The performance of the reservoir computer can be optimized by adjusting the values of the parameters α , $\Delta\varphi$, A_0 , and the amplitude of m_i (equivalent to changing β , which in the experiment is fixed by the coupling ratio of the cavity input coupler and cannot be modified). The input mask m_i , randomly chosen from a uniform distribution, is kept fixed.

As we show in the next sections, experimental and numerical simulations based on Eqs. (2) and (3) perform very well on a number of benchmark tasks. This good performance can be explained by the following points.

First, using complex variables (when $\Delta\varphi \neq 0$) makes the internal dynamics richer. For some tasks this is crucial, and contributes importantly to the improvement of performance (a fact already pointed out in [24,25,29]). Indeed using complex variables doubles the number of internal variables ($2N$), even

if we can only use N readout weights W_i to compute the output y .

Then, another contributing fact is that there is very little noise in our experiment. This is confirmed by the total memory capacity being close to its maximum value (see Section 4). This weak noise level is due to the absence of active (and therefore noise-inducing) elements in the reservoir.

Finally, a third favorable factor has more to do with the tasks investigated than with the reservoir itself. Indeed a reservoir computer can be seen as a map from the input time series $u(n)$ to the variables used to compute the output (here $|x_i(n)|^2$). The output of the reservoir is the projection of the desired output $y^*(n)$ onto the space spanned by $|x_i(n)|^2$ (see [33] for a detailed discussion). In the present case the reservoir states $x_i(n)$ are given by a linear combination of the previous inputs $u(n-l)$ ($l = 0, 1, 2, \dots$) with exponentially decaying dependence on l (the fading memory property essential for good reservoir performance). The variables $|x_i(n)|^2$ used to compute the output are therefore given by the sum of a constant, a linear, and a quadratic fading memory function of previous inputs (note the role of the bias A_0 that ensures that $|x_i(n)|^2$ contain linear terms). The fact that a reservoir based on Eqs. (2) and (3) performs well on a task therefore means that for this task, the desired output $y^*(n)$ can be well approximated by quadratic functions of previous inputs. The higher-order nonlinearities present in many reservoir algorithms are therefore not necessary for these tasks, and may even be detrimental. There are of course some tasks that require nonlinearities of higher order than quadratic, and for these tasks our architecture would not give good performance. However, by cascading reservoirs of the type reported here, nonlinearities of arbitrarily high order could be obtained, so this is not a fundamental problem.

3. EXPERIMENTAL IMPLEMENTATION

Our fiber optics experiment is depicted in Fig. 1. The optical signal is generated by a continuous wave (CW) laser at 1550 nm with a coherence time much greater than the inverse linewidth of the cavity, with output power adjusted between 11 and 41 mW. Note that below we always quote the optical power at the entrance of the cavity, as this is the experimentally relevant quantity, and the setup was not optimized to minimize the losses between the laser and the cavity.

The input signal $A_{in}(t)$ is encoded by modulating the amplitude of the laser signal with a lithium niobate Mach–Zehnder (M-Z) interferometer in a push–pull configuration driven by an arbitrary waveform generator (AWG). Consider first that no bias A_0 is needed. As the M-Z transfer function is sinusoidal, we first tried to precompensate the voltage signal V generated by the AWG, so that $A_{in}(t)$ is proportional to $m(t)u(t)$. Then we removed the precompensation of the masked input signal, so that the optical M-Z output signal $A_{in}(t)$ is proportional to $\sin\{V(t)[\pi/(2V_\pi)]\}$ with $V(t) = \gamma m(t)u(t) \in [-\gamma V_\pi, \gamma V_\pi]$ (V_π being the characteristic voltage of the M-Z modulator and $\gamma \in [0, 1]$ an adjustable parameter). We have checked that both codings give the same performance of the reservoir, except for the evaluation of the memory capacities for which the results depend on this input signal preprocessing.

For some tasks, it is necessary to bias the input. Experimentally, the bias is introduced by applying a direct current (DC) voltage $V_0 \in [-V_{\pi,DC}, V_{\pi,DC}]$ to the “DC electrode” of the

M-Z. When a bias is applied, the scaling of the input amplitude is important, and the voltage driving the radio frequency (RF) electrode of the M-Z is taken to be $V = \gamma m(t)u(t) \in [-\gamma V_{\pi,RF}, \gamma V_{\pi,RF}]$. The input signal therefore has the form

$$A_{in}(t) = \sin \left[\frac{\pi}{2} \left(\frac{\gamma m(t)u(t)}{V_{\pi,RF}} + \frac{V_0}{V_{\pi,DC}} \right) \right]. \quad (4)$$

In our case, we measured $V_{\pi,DC} = 7.54$ V and $V_{\pi,RF} = 7.27$ V. The AWG followed by RF amplifiers provides a maximum voltage amplitude V of 3.27 V on both RF electrodes of the M-Z, corresponding to the possibility to tune γ in the range of $[0, 0.45]$. All the results presented here were obtained with a maximum power at the output of the M-Z (i.e., at the entrance of the cavity) of 2.11 mW for unbiased masked input signal (respectively 5 mW for biased masked input signal). This value was chosen to provide a sufficient signal-to-noise ratio (SNR) on the readout photodiode, while avoiding Brillouin backscattering issues. However, we experimentally showed that we can decrease the power at the output of the M-Z to as low as 0.57 mW for unbiased masked input signal (respectively 1.35 mW for biased masked input signal) without affecting performance. Below this level, the SNR on the photodiode was too small.

The reservoir itself consists of a ~ 230 m long passive fiber cavity (made of single-mode fiber SMF-28e), corresponding to a roundtrip time of $T = 1.13209$ μ s. With the desynchronization parameter $k = 1$, and using $N = 50$ internal variables (which we use for most tasks), we therefore have $T' = 1.10989$ μ s, and the temporal length of each internal variable is equal to $\theta = 22.2$ ns. The output refresh rate $1/T' \sim 0.9$ MHz can be seen as the processing speed of our RC. It is quite low due to the use of rather slow electronics (AWG and photodiode) and a correspondingly long fiber cavity. However, this is not a fundamental limit since the refresh rate can in principle be easily increased by using a smaller cavity and faster electronics.

The input signal $A_{in}(t)$ is injected into the passive fiber cavity using a 90/10 coupler, corresponding to $\beta = (0.1)^{1/2} = 0.316$. A tunable optical attenuator is used to adjust the feedback gain in the range $0 < \alpha < 0.806$.

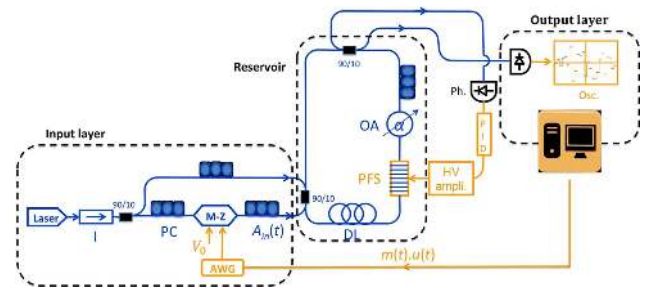


Fig. 1. Experimental setup: Blue lines (dark gray) correspond to fiber optic and orange lines (light gray) to electrical connections. Laser, long coherence telecom wavelength laser; I, isolator; PC, polarization controller; M-Z, amplitude modulator (lithium niobate Mach–Zehnder interferometer in push–pull configuration); OA, optical attenuator; PFS, piezoelectric fiber stretcher; AWG, arbitrary waveform generator; DL, delay loop; HV ampli., high-voltage amplifier; PID, proportional–integral–derivative regulator; Ph., photodiode; Osc., oscilloscope. The PC inside the cavity is used to control the Jones matrix of the cavity, and the PCs before the cavity ensure that $A(t)$ and the counterpropagating control signal on the two different polarization eigenmodes of the cavity.

In order to stabilize the fiber cavity, 10% of the laser light intensity is coupled into it in the counterpropagating direction using a 90/10 coupler (see Fig. 1), and its intensity is used as a control signal processed by a proportional–integral–derivative (PID) regulator that drives a piezoelectric fiber stretcher inserted in the cavity. In this way the phase detuning $\Delta\varphi$ of the cavity can be precisely controlled. A 90/10 coupler is finally used to send 10% of the optical intracavity power to the readout photodiode. A digital oscilloscope records the photodiode signal. The oscilloscope record is used in a digital postprocessing stage realized by a computer to compute the readout weights W_i , to produce the output $y(n)$, and to estimate the performance on specific tasks. Note that no time averaging is used in the oscilloscope recording, contrary to some earlier experiments.

The results of the experiment were carefully compared with discrete time simulations. Further details on the experimental setup and simulations are provided in Appendix A.

4. RESULTS

Here we present the results we obtained both experimentally and through simulations for different benchmark tasks widely used in the reservoir computing community. Except where indicated, we used $N = 50$ internal variables and a ridge parameter of 10^{-4} . As there are many parameters to scan (α , $\Delta\varphi$, $V_0/V_{\pi,DC}$, γ , and the ridge parameter), the reliability of our simulations allowed us to experimentally scan only the regions in parameter space that provide good performance. The experimental results are then compared with the simulation results within these regions. The experimentally reachable parameter ranges with our setup are the following: $\alpha \in [0.366, 0.806]$, $\Delta\varphi \in [-\pi, \pi]$ rad, $V_0/V_{\pi,DC} \in [-1, 1]$, and $\gamma \in [0, 0.45]$. When $V_0/V_{\pi,DC} = 0$, γ was always set to its maximum experimental reachable value.

For each choice of these parameters, we repeated the experiment 10 times with 10 different input series, averaging the performance obtained, except for the speech recognition task. This performance is then compared with the best published results obtained to date with different reservoir computer architectures [7,12–17,21,22,24,29,31,32].

In the following subsections, we provide a small description for each of these tasks. For a more in-depth description, see [13,15].

A. Memory Capacity Evaluation

In this basic task we test the ability of the reservoir to recall simple linear or nonlinear functions P of previous inputs $u(n-l)$. Three different memory capacities (MCs) are considered [33,34], i.e., the linear, quadratic, and cross-memory capacities. The memory function $MF(l)$ is computed as $MF = 1 - NMSE \in [0, 1]$ (NMSE, normalized mean square error) and is obtained by computing the correlation between P and the reservoir states, as is done in [33]. The total memory is the sum of the three different memory capacities. Theoretically it cannot exceed the number of internal variables [33]. Table 1 reports the best values obtained for each memory capacity, as well as for the total memory capacity, compared with three other architectures [13,15,17]. The results were obtained using 10 datasets of 2200 input samples randomly drawn from a uniform distribution in the interval $[-1, 1]$ (except for the total memory in simulation, in which case 10 datasets of 60,000 input samples were used, to lower statistical noise, given that the total memory in our particular reservoir architecture is very close to the maximum).

Note that all the memory capacity results obtained experimentally for the three architectures [13,15,17] selected for comparison and reported in Table 1 are taken from [17]. The reasons are that memory capacity results for the optoelectronic architecture [13] were first reported in [17], and that the memory capacity results for the semiconductor optical amplifier (SOA) based architecture published in [15] contain mistakes that were corrected in [17].

The linear memory we obtained is comparable to the one obtained for the SOA-based RC [15], but lower than for the optoelectronic [13] and saturable absorber (SA) based [17] RCs. The quadratic memory is about three times larger than the values reported for the optoelectronic [13] and SOA-based [15] RCs, and more than five times larger than for the SA-based RC [17]. The cross-memory is slightly larger than that obtained for the optoelectronic architecture [13], but more than six times larger than with the SOA-based RC [15], and two times larger compared with the SA-based RC [17]. Finally, the total memory capacity is comparable to that obtained for the optoelectronic RC [13], and larger than those obtained for the SOA [15] and SA [17] based RCs.

In summary, for most of the memory capacity evaluations, our reservoir architecture exhibits comparable or better performance

Table 1. Memory Capacity Evaluation^a

	Linear Memory	Quadratic Memory	Cross-Memory	Total Memory
MF	$LMF(l):P(u(n-l)) = u(n-l)$	$QMF(l):P(u(n-l)) = 3u^2(n-l) - 1$	$XMF(l,l'):P(u(n-l), u(n-l')) = u(n-l)u(n-l')$	
MC	$LC = \sum_{l=0}^{l_{\max}} LMF(l)$	$QC = \sum_{l=0}^{l_{\max}} QMF(l)$	$XC = \sum_{l=0}^{l_{\max}} \sum_{l'=l+1}^{l'_{\max}} XMF(l,l')$	$C = LC + QC + XC$
Sim.	$LC = 22.61 \pm 0.07$	$QC = 13.25 \pm 0.23$	$XC = 33.90 \pm 0.82$	$C = 49.99 \pm 0.13$
Exp.	$LC = 21.14 \pm 0.34$	$QC = 12.07 \pm 0.10$	$XC = 30.20 \pm 0.46$	$C = 48.37 \pm 0.47$
[13]	$LC = 31.9$	$QC = 4$	$XC = 27.3$	$C = 48.6$
[15]	$LC = 20.84$	$QC = 4.16$	$XC = 4.71$	$C = 25.20$
[17]	$LC = 36.8$	$QC = 2.23$	$XC = 14.32$	$C = 37.05$

^aMF, memory function; MC, memory capacity; Sim., simulation results of our architecture; Exp., experimental results of our architecture; LMF, linear memory function; LC, linear capacity; QMF, quadratic memory function; QC, quadratic capacity; XMF, cross-memory function; XC, cross capacity; C, (total) capacity. The conditions under which these results were obtained in both simulation and experiment are as follows: LC was obtained for $\alpha = 0.806$ and by scanning $\Delta\varphi$ in the range $[0, 1.42]$ rad, $V_0/V_{\pi,DC}$ in the range $[-0.6, -0.53]$, and γ in the range $[0.33, 0.45]$ (without any precompensation of the masked input signal); QC was obtained for $\alpha = 0.806$, $\Delta\varphi = 0$ and by scanning $V_0/V_{\pi,DC}$ in the range $[-0.22, 0.85]$ and γ in the range $[0.3, 0.45]$ (without any precompensation of the masked input signal); and XC and C were obtained for $\alpha = 0.806$, $\Delta\varphi = 0$, and $V_0/V_{\pi,DC} = 0$ (with precompensation of the masked input signal).

than the three experimental RC setups [13,15,17] taken for comparison, sometimes several times larger. The fact that the experimental total capacity C is very close to the maximum value (50) indicates that very little noise is affecting the experiment.

Note that the simulation model is in good agreement with the experimental results. As there is no noise in the simulation, the total memory capacity in simulation is found to be extremely close to the maximum value of 50, in agreement with theoretical expectations.

B. NARMA10

The nonlinear auto-regressive moving average (NARMA10) task is one of the most-used benchmark tasks in the field of reservoir computing [1,7,11–13,32]. The aim is to reproduce the behavior of a nonlinear, 10th-order system with random input $u(n)$ drawn from a uniform distribution over the interval [0,0.5]. The target output is the following:

$$y^*(n+1) = 0.3y^*(n) + 0.05y^*(n) \left[\sum_{i=0}^9 y^*(n-i) \right] + 1.5u(n-9)u(n) + 0.1. \quad (5)$$

We trained the reservoir over 1000 steps (1000 values of $u(n)$), and tested its performance for the subsequent 1000 steps. The standard deviation was evaluated by repeating this procedure 10 times. The performance obtained was measured using the NMSE, defined by

$$\text{NMSE} = \frac{\sum_n [y^*(n) - y(n)]^2}{\sum_n \left\{ y^*(n) - \frac{\sum_n [y^*(n)]}{n} \right\}^2}. \quad (6)$$

For this task, all the results below were obtained with a pre-compensated masked input signal and no bias $V_0/V_{\pi,DC} = 0$. Using $N = 50$ internal variables, we obtained a NMSE of 0.104 ± 0.02 for the simulation and an experimental NMSE of 0.107 ± 0.012 . This performance was obtained with $\alpha = 0.806$, and $\Delta\varphi$ was scanned in the ranges of [0,1.81] rad and [2.61,3.67] rad. This surpasses the best result reported in experiment to date with $N = 50$. For instance, in [13] a NMSE = 0.168 ± 0.015 was obtained with the same number of variables. The value NMSE = 0.16 in fact corresponds to the best that can be obtained with a linear shift register [12].

Our simulation model predicts very well our experimental results. Simulations also showed that our results are strongly dependent on the oscilloscope resolution: by increasing the acquisition resolution from 8 to 14 bits in our simulation model (thus decreasing the quantization noise), we reach a NMSE of 0.062 ± 0.008 , which is a very good result. This simulation result was obtained with $\alpha = 0.806$ and $\Delta\varphi$ scanned in the range [0, π] rad. The optimum was reached for $\Delta\varphi = 0.38$ rad. This is in the range already scanned for the results mentioned above with a resolution acquisition of 8 bits, which implies that the increase in performance is only due to the increase of the oscilloscope resolution. This result is, to our knowledge, better than any result published up to now with $N = 50$ internal variables (note that using 50 variables, the algorithm reported in [7] obtained NMSE = 0.152 ± 0.0138). We attribute this very good performance to our architecture in which the reservoir is linear and the readout quadratic. Indeed the nonlinearities exhibited by the NARMA10 equation are mainly quadratic.

The performance for the NARMA10 task strongly depends on the number of internal variables. Upon increasing the number of internal variables to $N = 300$, we obtained NMSE = 0.0484 ± 0.0095 in experiment, and NMSE = 0.0463 ± 0.0142 in simulation. These results were obtained with $\alpha = 0.806$ and $\Delta\varphi$ scanned in the range of [0.26,2.35] rad.

Finally it should be noted that using the simple algorithm given by Eqs. (2) and (3) (that is, without the experimental constraints set out in Appendix A and at the end of the first paragraph of this section), we obtained NMSE = 0.0106 ± 0.0030 with $N = 400$ internal variables. This performance was obtained by scanning α in the range [0.7,0.99] and $\Delta\varphi$ in the range [0, π] rad. Using the same number of internal variables, a simulation result of NMSE = 0.022 is reported in [12]. A NMSE = 0.018 is also reported in [32] using $N = 520$ internal variables.

C. Nonlinear Channel Equalization

In this task, first used in the context of RC in [3], the goal is to recover an input symbol sequence $d(n)$ from the signal received at the output $u(n)$ of a standardized nonlinear multipath RF channel, defined as follows:

$$q(n) = 0.08d(n+2) - 0.12d(n+1) + d(n) + 0.18d(n-1) - 0.1d(n-2) + 0.091d(n-3) - 0.05d(n-4) + 0.04d(n-5) + 0.03d(n-6) + 0.01d(n-7), \quad (7)$$

$$u(n) = q(n) + 0.036q^2(n) - 0.011q^3(n) + \nu(n). \quad (8)$$

Here $\nu(n)$ is a Gaussian noise with zero mean, adjusted to yield SNRs ranging from 12 to 32 dB. The symbols $d(t)$ are randomly chosen between four values $\{-3, -1, 1, 3\}$. The performance is evaluated in terms of symbol error rate (SER), which is the fraction of misclassified symbols.

For this task, we used 50 internal variables, 3000 training samples, and 6000 test samples. The standard deviation reported in Fig. 2 was evaluated by repeating this procedure 10 times. The

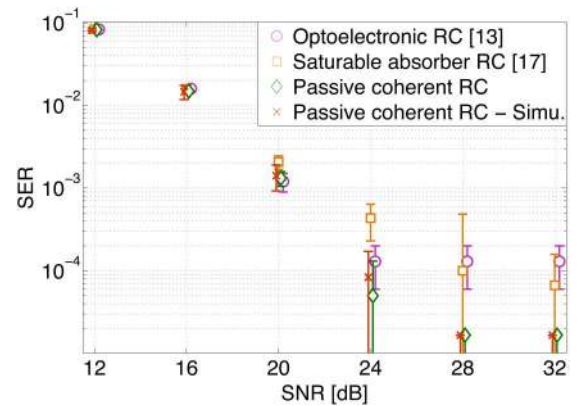


Fig. 2. Experimental and simulation results for the nonlinear channel equalization task. Horizontal axis, signal-to-noise ratio of the nonlinear channel; vertical axis, symbol error rate, i.e., number of misidentified symbols. These results were obtained with $\Delta\varphi = 0$ rad, $\gamma = 0.45$ and without any precompensation of the masked input signal. To get the best performances for each SNR, α was scanned in the range [0.45,0.55], $V_0/V_{\pi,DC}$ in the range [-0.5, 0.5], and the ridge parameter from 10^{-6} to 10^{-4} . Note that at 28 and 32 dB SNR, the SER is zero for the passive coherent RC, both in simulation and in experiment.

results obtained are reported in Fig. 2 and compared to the ones obtained in [13,17].

Again, we can see that our passive coherent RC architecture outperforms the previous architectures, using the same test conditions. In particular, for SNRs of 28 and 32 dB we obtained, both in simulation and experimentally, a SER of 0%, which means that all 60,000 symbols were correctly identified. Such very good performance was never attained in previous experiments.

D. Isolated Spoken Digits Recognition

The goal of this task is to classify 10 digits pronounced 10 times by five different female speakers, first without noise and then with babble noise added to the sound records to reach a SNR of 3 dB, following the procedure explained in [24]. The data were taken from the NIST TI-46 Corpus [35] (National Institute of Standards and Technology Texas Instrument-46 Corpus), and the recording was preprocessed according to the Lyon cochlear ear model [36]. Ten output classifiers were trained, one for each digit. Each classifier was trained to output « + 1 » if a specific digit was sent to the reservoir, and « - 1 » otherwise. A winner-takes-all approach between the classifiers was then used to select the actual response of the reservoir.

For a better estimation of the performance, a cross-validation procedure was applied over five subsets of 100 words, chosen randomly. The reservoir was trained on four of the subsets, and then tested on the fifth one. This procedure was repeated five times, so that each subset was used once as a test subset. The performance values, evaluated in terms of word error rate (WER), which is the fraction of misclassified digits, are reported in Table 2.

Again, we can see that our simulation results are very close to the results obtained experimentally. With no noise on the input signal, we obtained, using 200 internal variables, an experimental WER of 0%, which means that all the 500 digits were well classified. A WER of 0% obtained in simulation is also reported in [22] using 400 internal variables. Our result surpasses the one obtained with the optoelectronic RC [13] (WER of 0.4%) with $N = 200$ internal variables. A WER of $3 \pm (1.2)\%$ is also reported in [15] using $N = 200$.

Experimental results on this task are also reported in [12,14,16,21,22,31]. The reported WERs are 0.2% [12], $0.04(\pm 0.017)\%$ [14], $0.014(+0.051/-0.014)\%$ [16,22], $0.6(\pm 0.2)\%$ [31], and 0.05% [21]. In these works, $N = 400$ in [12,14,21], $N = 388$ in [16,22], $N = 150$ in [31], and for the cross-validation procedure, the reservoir was trained on 475 spoken digits and tested on 25 digits. Concerning speed,

it should be noted that the setup presented in [16] carried out this task with a refresh rate slightly more than 14 times faster than in the present experiment, while our setup is almost 19 times faster than in [14], 7.2×10^5 times faster than in [12], and more than 200 times faster than in [31].

This task can be made more complicated by adding background babble noise. With a SNR of 3 dB, we experimentally obtained $WER = 0.8(\pm 0.8)\%$ with $N = 500$. Unfortunately, we cannot compare this performance because it has never been tested experimentally until now. In simulation we obtain a $WER = 0.6(\pm 0.9)\%$ with $N = 500$. This can be compared to the simulation results of $WER = 4.5\%$ [24] using $N = 81$ and of $WER = 1\%$ [29] using $N = 500$.

On this task, our algorithm (i.e., the simple algorithm given by Eqs. (2) and (3), without the experimental constraints set out in Appendix A and at the end of the first paragraph of this section) reaches $WER = 0$ with $N = 90$ for noiseless signal, and $WER = 0$ with $N = 350$ for noisy signal. These later results were obtained with precompensation of the masked input signal and without any bias. α was scanned from 0.8 to 0.99, $\Delta\phi$ was scanned in the region of $[0, \pi]$ rad, and the ridge parameter was scanned from 10^{-6} to 10^{-4} .

5. CONCLUSIONS

In this work we have demonstrated a photonic implementation of a passive linear fiber reservoir computer working with coherent light for analog signal processing. This experiment presents many qualities that were either absent or not simultaneously present in previous works: we can perform analog optical signal processing, the easy tunability of each key parameter achieves the best operating point for each task, the system is able to reach a strikingly weak noise floor thanks to the absence of active elements in the reservoir itself, richer dynamics is provided by operating in coherent light, and finally high power efficiency is yielded as a result of the passive nature and simplicity of the setup. Note that at this stage we have only obtained low optical power consumption for the reservoir itself, and have not tried to minimize the overall power consumption, including all control electronics.

The main challenge was to stabilize the system, as our reservoir is a long interferometer made of an optical fibered cavity of approximately 230 m. In the future, faster electronics (photodiode, AWG, and oscilloscope) will enable smaller, possibly even integrated, cavities, and hence much simpler stabilization. In addition, our expertise in interferometric stabilization will allow us in the future to study photonic reservoir computers in which the reservoir states are processed in parallel, rather than sequentially, providing significant further speedup.

On the different tasks we tested, the present experiment has error rate as low as or lower than previous experiments. For instance, the NARMA10 task was up to now considered as a big challenge, and no experiment with performance exceeding a linear shift register had been reported before this work.

Quite remarkably, through its conceptual simplicity, our experimental approach has also contributed to the theory of reservoir computing. Indeed, the discrete time Eqs. (2) and (3) constitute a very simple high-performance reservoir computer algorithm. In particular, it combines three advantages: (1) the simple interconnection matrix first introduced in [7]; (2) the simplicity of a linear reservoir associated with a nonlinear output

Table 2. Experimental and Simulation Results for the Isolated Spoken Digit Recognition Task^a

	No Noise on the Signal, with 200 Internal Variables	SNR = 3 dB, with 500 Internal Variables
Simulation	WER = 0%	WER = 0.6(±0.9)%
Experiment	WER = 0%	WER = 0.8(±0.8)%

^aThese results were obtained with precompensation of the masked input signal, no bias, and $\alpha = 0.806$. The ridge parameter was scanned from 10^{-6} to 10^{-4} , and $\Delta\phi$ was scanned in the range of $[-1.62, 0.19]$ rad.

layer; and (3) very good performance on benchmark tasks even when a small number of internal variables is used. Our algorithm thus seems very computationally efficient and could also find applications in numerical implementations of reservoir computers. This should, however, be confirmed by a more detailed comparison with other RC algorithms.

Given the good performance and conceptual simplicity of the experiment reported here, one can expect that this architecture constitutes an important milestone in the future progress in photonic reservoir computing.

APPENDIX A

1. Numerical Simulations

Our simulation model is based on the discrete time Eqs. (2) and (3), with the nonlinear preprocessing (4) taken into account when relevant. In addition, the simulation includes the effects of the most important features of the components of our experimental setup: the 14 bits of resolution of the AWG, the 8 bits of resolution of the oscilloscope, the saturation transfer function of the RF amplifiers used to amplify the signal sent to the M-Z from the AWG, the low cutoff frequency of these RF amplifiers (30 kHz to 12 GHz bandwidth), and the low cutoff frequency of the photodiode used to read the $|x_i(n)|^2$ (30 kHz to 1 GHz bandwidth).

2. Interferometric Stabilization and Phase Detuning Setting

Our reservoir is a fiber optic cavity with a roundtrip time $T = 1.13209 \mu\text{s}$ corresponding to approximately 230 m of fiber. This cavity thus needs to be phase stabilized. This is done using a piezoelectric fiber stretcher that compensates phase shifts introduced in the cavity by thermal, vibrational, and phonic noise. To this end, a nonmodulated optical signal, the control signal, is sent in the cavity in the opposite direction compared to the reservoir states signal coded in $A(t)$, as shown in Fig. 1. This signal is collected by a photodiode, which drives a digital PID regulator. The output signal of the PID regulator is then amplified in order to control the piezoelectric fiber stretcher.

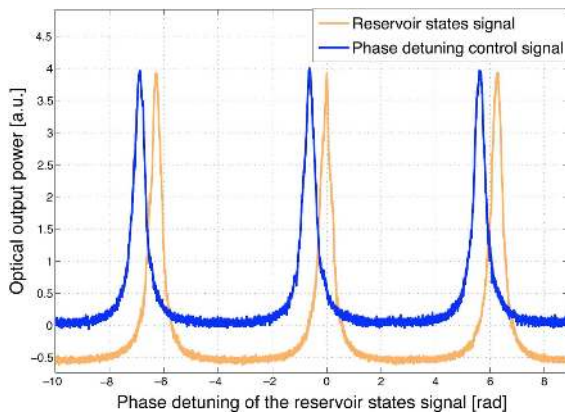


Fig. 3. Resonances of the cavity. Two CW signals are coupled in counterpropagating directions in the cavity, on the two polarization eigenmodes of the cavity. A voltage ramp is applied to the piezoelectric fiber stretcher in order to scan the cavity phase, and the output power is recorded. The optical attenuator was set at its maximum transparency.



Fig. 4. Isolation system of the delay loop, open. It consists of two aluminum boxes: a small box made of 1 cm thick plates inside a bigger box made of 2 cm thick plates. The inner plates of both boxes are covered with stone wool. Each box is isolated from the base on which it rests by sorbothane sheets that absorb vibrations.

The polarization controller inside the cavity (see Fig. 1) is used to tune the Jones matrix of the cavity (the polarization transfer matrix of the cavity), and thereby tune “the phase offset” between the two polarization eigenmodes of the cavity. The reservoir states signal and the counterpropagating control signal are injected on the two different polarization eigenmodes of the cavity. The PID stabilizes the cavity on one of the slopes of the control signal resonance. By changing the Jones matrix of the cavity, the phase distance between the resonances of the reservoir states signal and the control signal can be arbitrarily adjusted, and thus the detuning $\Delta\varphi$ appearing in Eq. (1) can be adjusted to an arbitrary value.

This is illustrated in Fig. 3, which shows the measured cavity transfer function for two nonmodulated input signals coupled in counterpropagating directions in the cavity, on the two different polarization eigenstates.

In order to minimize the phonic, vibrational, and thermal noise, we isolated every cavity component in several boxes made of 1–2 cm thick aluminum plates with stone wool inside, as shown in Fig. 4 for the delay loop.

All the boxes fit together so that the isolation system of the whole cavity is almost airtight. These boxes were designed to isolate the cavity from the phase noise frequencies that the PID regulator is not able to follow, as its output refresh rate of 8.3 kHz is limited by the resonance frequency of the piezoelectric fiber stretcher. Moreover, the whole experiment is placed on an optical table maintained on air cushions.

Note that in this first implementation, the comparatively long fiber cavity (230 m) was designed so as to be able to use rather slow electronics (AWG and photodiode). In future work we intend to use faster electronics, and hence a smaller cavity. A smaller cavity will be much easier to stabilize and isolate from sources of phase noise. An implementation on a chip, as in [29], will not require any stabilization (except for thermal control of the whole chip).

3. Experimental Setup Components

Table 3 presents the relevant technical characteristics of the main components used in our hardware setup.

Table 3. Technical Description of the Most Important Hardware Components

Equipment	Technical Description
Laser	Koheras ajustik KOH1995 laser: - Wavelength: 1550 nm - Full width at half-maximum (FWHM): 1 kHz
Reservoir	Optical fibered cavity made with 2 90/10 couplers and one optical attenuator: - Maximum feedback gain in electromagnetic field amplitude: 0.806 - Minimum FWHM: 60.3 kHz - Maximum finesse: 14.65
Arbitrary waveform generator	Agilent AWG M8190A: - 14 bits resolution - Sample rate from 125 MSamples/s to 8 GSamples/s - DC to 5 GHz analog bandwidth - Up to 128 MSamples arbitrary waveform memory per channel
Acquisition system: oscilloscope	Agilent DSA91204A: - 12 GHz bandwidth - Up to 40 GSamples/s on each of four analog channels - 1 GSamples memory per channel
PID regulator	Homemade PID regulator programmed using mbed NXP LPC1768 development board: - Processor: 32-bits ARM Cortex-M3 core running at 96 MHz - Analog-to-digital converter (ADC): 12 bits resolution - Digital-to-analog converter (DAC): 10 bits resolution - Output refresh rate: 8.3 kHz
Mach-Zehnder interferometer	Lucent electro-optic modulator: - Model: X2624C - 20 GHz analog bandwidth
Reservoir states readout photodiode	New focus low-noise photoreciever, model 1611 IR 1 GHz: - Bandwidth (3 dB): 30 kHz to 1 GHz - Risetime: 400 ps - Current gain: 700 V/A - Responsivity at 1550 nm: 1.04 A/W

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