



Homogeneous Thermal Cloak with Constant Conductivity and Tunable Heat Localization

Tiancheng Han^{1,2}, Tao Yuan¹, Baowen Li^{3,4} & Cheng-Wei Qiu¹

¹Department of Electrical and Computer Engineering, National University of Singapore, Kent Ridge 119620, Republic of Singapore, ²School of Physical Science and Technology, Southwest University, Chongqing 400715, China, ³Department of Physics and Centre for Computational Science and Engineering, National University of Singapore, Singapore 117546, Republic of Singapore, ⁴Center for Phononics and Thermal Energy Science, School of Physical Science and Engineering, Tongji University, 200092, Shanghai, China.

Invisible cloak has long captivated the popular conjecture and attracted intensive research in various communities of wave dynamics, e.g., optics, electromagnetics, acoustics, etc. However, their inhomogeneous and extreme parameters imposed by transformation-optic method will usually require challenging realization with metamaterials, resulting in narrow bandwidth, loss, polarization-dependence, etc. In this paper, we demonstrate that thermodynamic cloak can be achieved with homogeneous and finite conductivity only employing naturally available materials. It is demonstrated that the thermal localization inside the coating layer can be tuned and controlled robustly by anisotropy, which enables an incomplete cloak to function perfectly. Practical realization of such homogeneous thermal cloak has been suggested by using two naturally occurring conductive materials, which provides an unprecedentedly plausible way to flexibly realize thermal cloak and manipulate heat flow with phonons.

Recently, many significant achievements of invisibility cloaking have been achieved hitherto, owing to pioneering theoretical proposals of Pendry¹ and Leohardt². Since the ideal electromagnetic cloaks generally come with extremely complex constitutive parameters (inhomogeneous, anisotropic, and singular), many simplified strategies have been adopted, including reduced cloaks^{3–6} and carpet cloaks^{7–13}. To overcome the lateral shift problem of the carpet cloak designed by quasi-conformal mapping¹⁴, ground-plane cloak constructed with calcite crystals^{15,16} and one dimensional full-parameter cloak constructed with metamaterials¹⁷ have been experimentally demonstrated recently. Decoupling electric and magnetic effects, static magnetic cloak¹⁸ and static electric cloak¹⁹ have been experimentally realized with anisotropic but homogeneous permeability and permittivity, respectively. In addition to manipulation of electromagnetic wave^{1–19}, the theoretical tool of coordinate transformation has been extended to acoustic waves^{20,21}, matter waves^{22,23}, elastic waves²⁴, and heat flux^{25–30}.

On the basis of the invariance of heat conduction equation under coordinate transformations, transformation thermodynamics has provided a new method to manipulate heat flux at will²⁵. The most attractive application is thermal cloaking: making the temperature of a certain region invariant. However, in an analog to electromagnetic cloaking^{1,3}, conventional thermal cloaking is dependent on its geometry, whose materials are usually inhomogeneous and singular^{26–29}, and in turn the practical applications of thermal cloak may be limited. More recently, an experiment has been reported to shield, concentrate, and invert heat current³⁰ by utilizing latex rubber and processed silicone.

Inspired and motivated by the pioneering work³⁰, we establish the theoretical account and general design roadmap for creating realizable thermal cloak and manipulating the heat localization, by using only homogeneous, non-singular, and natural conduction materials. The proposed novel thermal cloak is homogeneous, non-singular, independent on its geometrical size, and dominated by only anisotropy, which is distinguished from thermal cloaks^{25–30} reported so far. More interesting is that, by judiciously selecting natural materials, a partially constructed cloak can perform perfectly. This unique functionality is enabled by the controllable thermal localization, e.g., most heat is confined in the vicinity of the cloak's outer boundary. Therefore an ultra-thin thermal cloak could be created, if a natural material with strong conduction anisotropy can be found or an effective material of highly anisotropic conduction can be constructed. Due to homogeneity and non-singularity, the

SUBJECT AREAS:

THERMODYNAMICS

MECHANICAL ENGINEERING

COMPUTATIONAL METHODS

COMPOSITES

Received

5 February 2013

Accepted

19 March 2013

Published

3 April 2013

Correspondence and requests for materials should be addressed to C.-W.Q. (eleqc@nus.edu.sg)



proposed cloak may be fabricated by multilayer composition approach exploiting two naturally occurring materials throughout.

Results

Heat flows spontaneously from a high temperature region toward a low temperature region. For a steady state and without heat source, the thermal conduction equation can be written as $\nabla \cdot (\kappa \nabla T) = 0$, where κ is the thermal conductivity, and T is the temperature. Transformation thermodynamics has demonstrated that the conduction equation is invariant in its form under the coordinate transformation²⁵. Specifically, in the transformed space, the thermal conduction equation can be written as $\nabla' \cdot (\bar{\kappa} \nabla' T') = 0$. We can obtain

$$\bar{\kappa} = \frac{\mathbf{A}\kappa\mathbf{A}^T}{\det(\mathbf{A})}, \quad \text{with } \mathbf{A} = \frac{\partial(x', y', z')}{\partial(x, y, z)} \quad (1)$$

Considering two dimensional case, where a circular region ($r \leq b$) in original space (r, θ, z) is compressed into an annular region ($a \leq r' \leq b$) in physical space (r', θ', z'). The transformation equation can be expressed as

$$r' = \frac{b-a}{b}r + a, \quad \theta' = \theta, \quad z' = z \quad (2)$$

Submitting Eq. (2) into Eq. (1) and assuming $\kappa = 1$, the conductivity of ideal cloak can be expressed as

$$\begin{aligned} \kappa'_r &= \frac{r'-a}{r'} \\ \kappa'_\theta &= \frac{r'}{r'-a} \end{aligned} \quad (3)$$

Clearly, the conductivity in Eq. (3) is spatially variable and has singularity at $r' = a$ ($\kappa'_r \rightarrow 0$ and $\kappa'_\theta \rightarrow \infty$), which is extremely difficult, if not impossible, to realize it. Apart from the ideal cloak, can we construct an advanced cloak with finite constant conductivity under the premise of maintaining perfect cloaking performance? Fortunately, this is feasible in thermodynamics. Examining Eq. (3),

we can obtain the relationship $\kappa'_r \kappa'_\theta = 1$ and $0 \leq \kappa'_r = \frac{1}{\kappa'_\theta} \leq \frac{b-a}{b}$. By making $\kappa'_r = \frac{1}{\kappa'_\theta} = C$, where C is a constant of $0 < C < \frac{b-a}{b}$, a homogeneous cloak without singularity is achieved. The detailed theoretical analysis is demonstrated in Methods. The inner and outer radii are $a = 1$ m and $b = 2$ m, respectively.

Full-wave simulations are carried out based on finite element method (FEM). Fig. 1(a) corresponds to the ideal cylindrical cloak described in Eq. (3), and Fig. 1(b)–(d) correspond to proposed homogeneous cloak with $C=0.1$, $C=0.2$, and $C=0.3$, respectively. As we can see, in Fig. 1(b)–(c), the heat fluxes travel around the inner domain and eventually returns to their original pathway. Therefore, the object inside the inner domain is protected from the invasion of

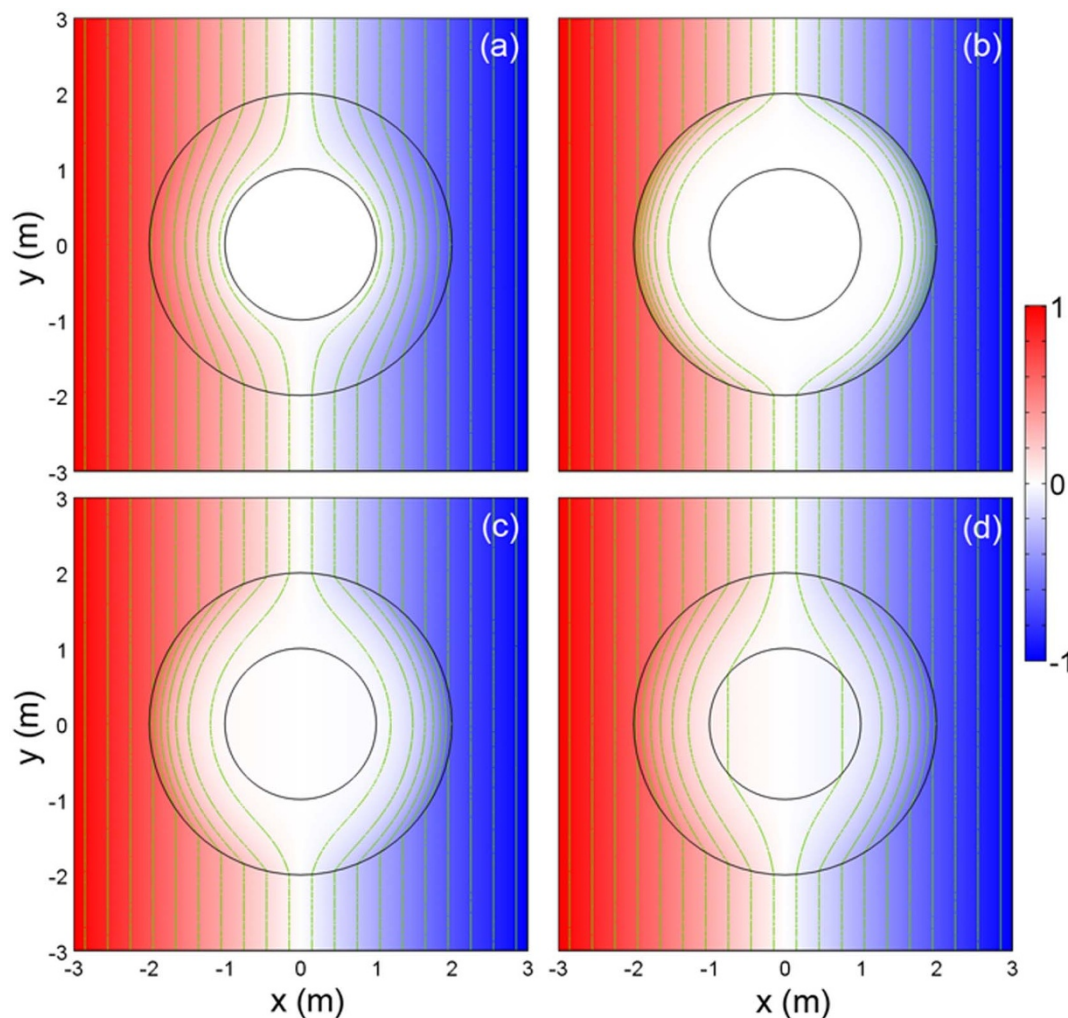


Figure 1 | Temperature profile for a thermal cloak with $a = 1$ m and $b = 2$ m. (a) Ideal conductivity described in Eq. (3). (b) $\kappa'_r = 0.1$ and $\kappa'_\theta = 10$. (c) $\kappa'_r = 0.2$ and $\kappa'_\theta = 5$. (d) $\kappa'_r = 0.3$ and $\kappa'_\theta = 3.3$. Isothermal lines are also represented with green color in panel.



external heat flux. Clearly, we have achieved advanced cloak with extremely simple parameters (finite constant conductivity), having performance as perfect as the ideal case.

When C increases to 0.3, shown in Fig. 1(d), a small portion of thermal energy goes into the inner domain, which leads to an imperfect invisibility cloaking. Obviously, smaller C (i.e. larger anisotropy) corresponds to better performance. To determine the maximal value of C , we may set $T_1/T_3 = (a/b)^{l-1} \leq 0.1$, which means that the temperature potential in inner domain is negligible. Then we can obtain $\text{Max}(C) = \frac{\lg(a/b)}{\lg(a/b) - 1}$. When $a = 1$ m and $b = 2$ m, we can obtain $\text{Max}(C) = 0.23$. Clearly, as long as C is smaller than $\text{Max}(C)$, nearly perfect performance can be achieved, as shown in Fig. 1(b) and (c).

To quantitatively examine cloaking performance with variance of anisotropy (denoted by C) and geometrical size (denoted by b/a), Fig. 2 (a) and (b) show the temperature gradient (∇T) of the inner region ($r < a$) as functions of C and b/a . When the geometrical size is fixed ($a = 1$ m and $b = 2$ m), temperature gradient as function of C is demonstrated in Fig. 2(a). Clearly, nearly perfect performance can be achieved as C is smaller than 0.23. When anisotropy is fixed ($\kappa'_r = 0.1$ and $\kappa'_\theta = 10$), i.e. $C = 0.1$, temperature gradient as function of b/a is demonstrated in Fig. 2(b). Obviously, good performance is kept until $b/a = 1.3$, which means the cloaking shell is very thin. This is because $b/a = 1.3$ corresponds to $\text{Max}(C) = 0.1$, and the fixed $C = 0.1$ is not larger than $\text{Max}(C)$, thus good performance could still be achieved.

According to prediction of the theoretical equation $T_2/T_3 = (r/b)^{l-1}$, the temperature distribution in region II is more concentrated near outer boundary with the decrease of C , which has been shown from Fig. 1(b) to Fig. 1(d). To demonstrate this phenomenon clearly, Fig. 2(c) show the isothermal contour with different C . Obviously, the isothermal lines are more concentrated near outer boundary with the decrease of C . When $C = 0.01$, nearly all of the energy in shell region is confined to the inner side of the outer boundary, which means an ultra-thin cloak with homogenous conductivity can be created.

Due to the proposed cloak with finite constant conductivity, it could be easily realized through alternating layered isotropic medium and only two types of isotropic materials (medium A and

medium B) are needed throughout. The conductivities of medium A and medium B are defined as $\kappa_{A,B} = \kappa'_\theta \pm \sqrt{\kappa'^2_\theta - \kappa'_r \kappa'_\theta}$. Considering a thermal cloak with $a = 1$ m, $b = 2$ m, and $C = 0.25$, i.e. $\kappa'_r = 0.25$ and $\kappa'_\theta = 4$, the temperature profile is shown in Fig. 3(a). The meshes formed by streamline and isothermal values illustrate the deformation of the transformed space, which is curved smoothly around the central invisibility region. Fig. 3(b) presents the temperature profile for the multilayered cloak with homogeneous, nonsingular, and isotropic conductivities, which is believed to be quite advanced. The anisotropy in Fig. 3(a) has been removed by replacing the anisotropic material with two isotropic conductivities $\kappa_A = 7.87$ (thermal epoxy) and $\kappa_B = 0.13$ (natural latex rubber), as shown in Fig. 3(c).

To validate the tunable and robust functionality of the proposed cloak, we simulate the same sized cloak with $C = 0.1$, i.e. $\kappa'_r = 0.1$ and $\kappa'_\theta = 10$, as shown in Fig. 3(d). Since the energy is mainly distributed near the outer boundary and is confined to the region ($0.15 \leq r \leq 0.2$ m), we can construct a partially constructed cloak made of multilayered materials in Fig. 3(e), where those multilayers initially situated in no energy region ($0.1 \text{ m} \leq r \leq 0.15 \text{ m}$) can be simply discarded and then removed. The two constitutive materials of the thinner cloak are stainless steel ($\kappa_A = 20$) and wood ($\kappa_B = 0.05$), as shown in Fig. 3(f). Clearly, the partially multilayered cloak behaves as perfectly as the ideally complete case. More importantly, one needs just two kinds of conductivities, in contrast to the singular and inhomogeneous cloak requiring $2N$ kinds of different conductivities²⁹.

Discussion

We have proposed an advanced methodology for the design of thermal cloak with finite constant conductivity (without inhomogeneity and singularity), which drastically facilitates feasible realization and fabrication. The proposed cloak is independent on its geometrical size and is dominated by only anisotropy, which could be easily replaced by periodically alternating isotropic conductivities. Furthermore, given two isotropic conductivities, anisotropy can still be tunable in a large range by adjusting individual filling ratio, empowering many flexible recipes for using naturally occurring materials in thermal cloaking. In this connection, we demonstrate the possibility of

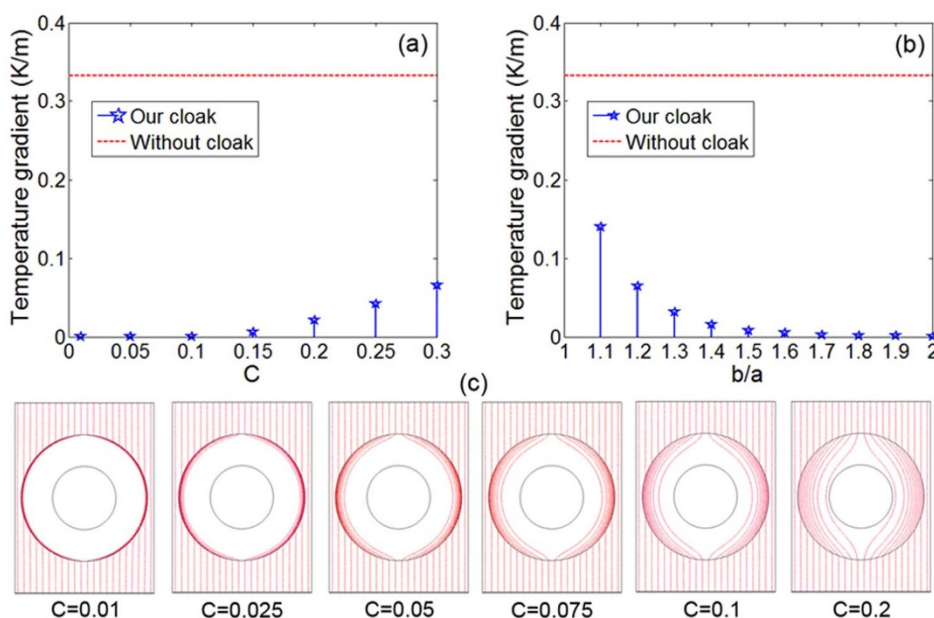


Figure 2 | (a) Temperature gradient of the inner region ($r \leq a$) as function of C with $a = 1$ m and $b = 2$ m. (b) Temperature gradient of the inner region ($r \leq a$) as function of b/a with $\kappa'_r = 0.1$ and $\kappa'_\theta = 10$. (c) Isothermal contour with different C values at $a = 1$ m and $b = 2$ m.

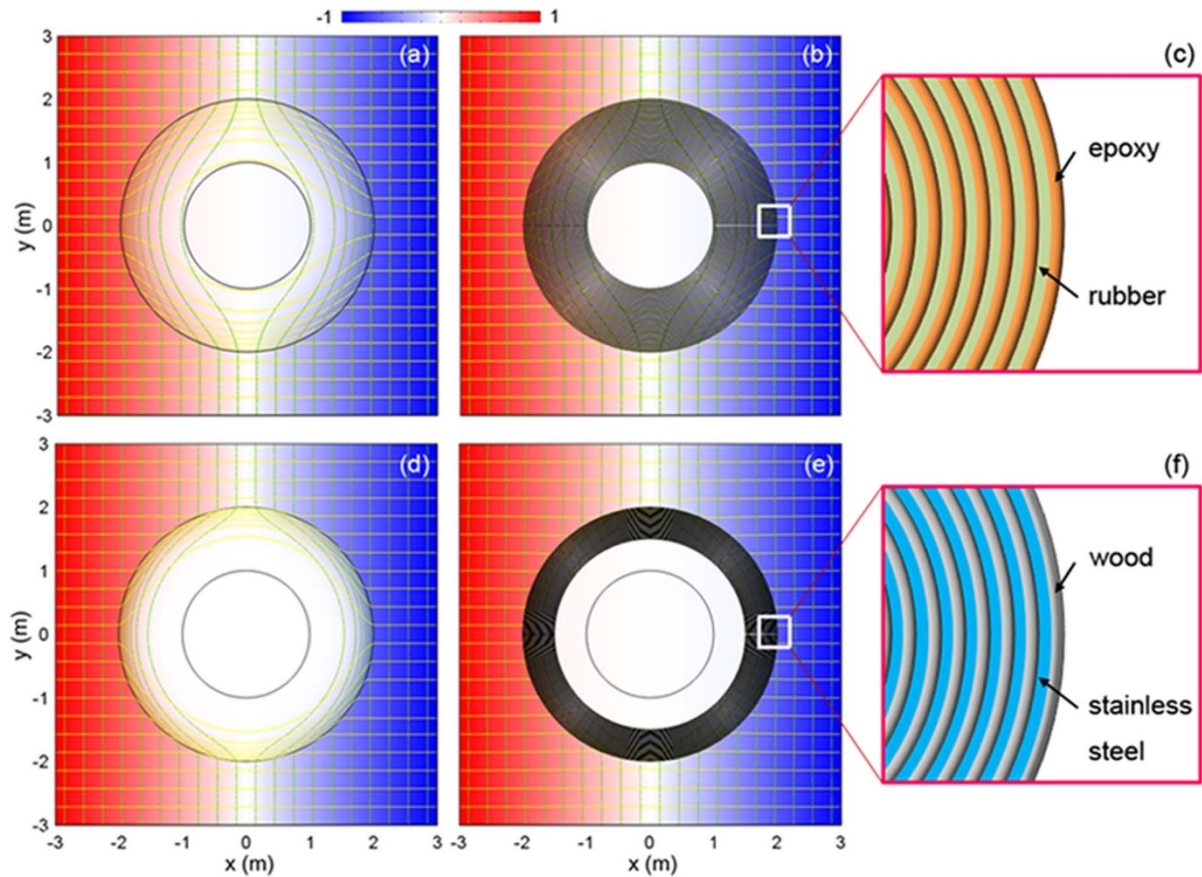


Figure 3 | Temperature profile for the thermal cloak with $a = 1$ m and $b = 2$ m. (a) $C=0.25$. (b) The multilayered composition realization for the cloak in (a). (c) Close-up view of the multilayered cloak in (b) showing the constitutive materials available in nature. (d) $C = 0.1$. (e) The multilayered composition realization for the cloak in (d). (f) Close-up view of the multilayered cloak in (e) showing the constitutive materials. Streamlines of thermal flux and isothermal are also represented with yellow and green colors in panel, respectively.

creating ultra-thin thermal cloak by partial construction while maintaining perfect functionality. Theoretical analysis and full-wave simulations validate the advanced thermal cloak with utmost simple materials.

It should be pointed out that the presented work has put additional dimension to the emerging field of phononics³¹: controlling and manipulating heat flow with phonons. The ideas proposed in this paper might be applicable to another two forms of heat transfer – heat convection and heat radiation, given the fact that

many phononic concepts have been extended to heat radiation and acoustics.

Methods

We introduce the rigorous theoretical analysis for such a homogeneous cloak, whose model is shown in Fig. 4. The conduction equation in cylindrical coordinate can be expanded as

$$\frac{\partial^2 T}{\partial r^2} + \frac{1}{r} \frac{\partial T}{\partial r} + \frac{l^2}{r^2} \frac{\partial^2 T}{\partial \theta^2} = 0 \tag{4}$$

where $l = 1$ for region I ($0 \leq r \leq a$) and III ($r > b$), $l = \sqrt{\kappa'_\theta / \kappa'_r}$ for region II ($a \leq r \leq b$). Considering the symmetry relation of $T(x, y) = T(x, -y)$, the temperature potential of three regions can be respectively expressed as

$$T_1 = \sum_{n=1}^{\infty} A_{2n-1} r^{2n-1} \cos(2n-1)\theta \tag{5a}$$

$$T_2 = \sum_{n=1}^{\infty} [B_{2n-1} r^{(2n-1)l} + C_{2n-1} r^{-(2n-1)l}] \cos(2n-1)\theta \tag{5b}$$

$$T_3 = \sum_{n=1}^{\infty} [D_{2n-1} r^{2n-1} + E_{2n-1} r^{-2n+1}] \cos(2n-1)\theta \tag{5c}$$

Owing to the temperature potential and the normal component of heat flux vector being continuous across the interfaces, we have

$$\begin{cases} T_1|_{r=a} = T_2|_{r=a}, & \nabla T_1|_{r=a} = \kappa'_r \nabla T_2|_{r=a} \\ T_2|_{r=b} = T_3|_{r=b}, & \kappa'_r \nabla T_2|_{r=b} = \nabla T_3|_{r=b} \end{cases} \tag{6}$$

Taking into account the boundary condition $T(x = \pm x_0) = \mp T_0$ (where T_0 is constant), we can obtain

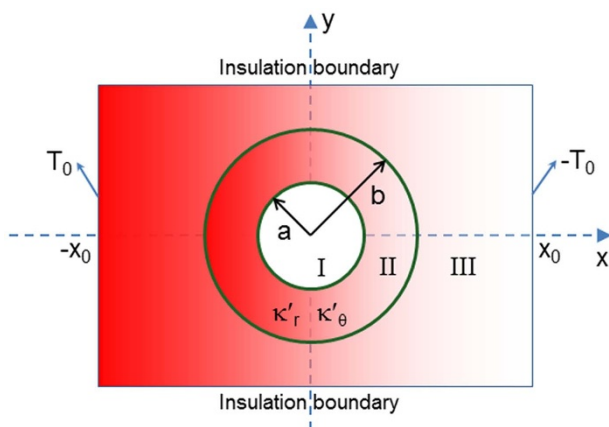


Figure 4 | Schematic of the homogeneous thermal cloak. Red and white colors denote high and low temperature, respectively.



$$\begin{aligned}
 T_1 &= -\frac{T_0}{x_0} \left(\frac{a}{b}\right)^{l-1} r \cos \theta \\
 T_2 &= -\frac{T_0}{x_0} \left(\frac{r}{b}\right)^{l-1} r \cos \theta \\
 T_3 &= -\frac{T_0}{x_0} r \cos \theta
 \end{aligned} \quad (7)$$

Thus the temperature potential of three regions can be fully presented in Eq. (7). It can be seen that $l = 1$ leads to $T_1 = T_2 = T_3$, corresponding to free space. For our homogeneous cloak, we need to control $l = \sqrt{\kappa'_0/\kappa'_r} = 1/C$, with $T_1/T_3 = (a/b)^{l-1}$ and $T_2/T_3 = (r/b)^{l-1}$. When l is large enough, it implies $T_1/T_3 \rightarrow 0$, resulting in that nearly no energy flows into the inner domain. It reveals that larger anisotropy corresponds to better performance, with the price of more difficult fabrication though. We want to have perfect performance ($T_1/T_3 \rightarrow 0$) and small anisotropy ($l \rightarrow 1$) simultaneously. However, it is not possible. In fact, each solution is some compromise between these two quantities, and each quantity can be improved by trading off the other one.

- Pendry, J. B., Schurig, D. & Smith, D. R. Controlling electromagnetic fields. *Science* **312**, 1780–1782 (2006).
- Leonhardt, U. Optical conformal mapping. *Science* **312**, 1777–1780 (2006).
- Schurig, D. *et al.* Metamaterial electromagnetic cloak at microwave frequencies. *Science* **314**, 977–980 (2006).
- Kanté, B., Germain, D. & Lustrac, A. Experimental demonstration of a nonmagnetic metamaterial cloak at microwave frequencies. *Phys. Rev. B* **80**, 201104 (2009).
- Xu, S. *et al.* Experimental demonstration of a free-space cylindrical cloak without superluminal propagation. *Phys. Rev. Lett.* **109**, 223903 (2012).
- Chen, H. & Zheng, B. Broadband polygonal invisibility cloak for visible light. *Scientific Reports* **2**, 255 (2012).
- Li, J. & Pendry, J. B. Hiding under the carpet: a new strategy for cloaking. *Phys. Rev. Lett.* **101**, 203901 (2008).
- Liu, R. *et al.* Broadband ground-plane cloak. *Science* **323**, 366–369 (2009).
- Ma, H. F. & Cui, T. J. Three-dimensional broadband ground-plane cloak made of metamaterials. *Nat. Commun.* **1**, 21 (2010).
- Lee, J. H. *et al.* Direct visualization of optical frequency invisibility cloak based on silicon nanorod array. *Optics Express* **17**, 12922 (2009).
- Valentine, J., Li, J., Zentgraf, T., Bartal, G. & Zhang, X. An optical cloak made of dielectrics. *Nat. Mater.* **8**, 568–571 (2009).
- Gabrielli, L. H., Cardenas, J., Poitras, C. B. & Lipson, M. Silicon nanostructure cloak operating at optical frequencies. *Nat. Photon.* **3**, 461–463 (2009).
- Ergin, T., Stenger, N., Brenner, P., Pendry, J. B. & Wegener, M. Three-dimensional invisibility cloak at optical wavelengths. *Science* **328**, 337–339 (2010).
- Zhang, B., Chan, T. & Wu, B.-I. Lateral shift makes a ground-plane cloak detectable. *Phys. Rev. Lett.* **104**, 233903 (2010).
- Zhang, B., Luo, Y., Liu, X. & Barbastathis, G. Macroscopic invisibility cloak for visible light. *Phys. Rev. Lett.* **106**, 033901 (2011).
- Chen, X., Luo, Y., Zhang, J., Jiang, K., Pendry, J. B. & Zhang, S. Macroscopic invisibility cloaking of visible light. *Nat. Commun.* **2**, 176 (2011).
- Landy, N. & Smith, D. R. A full-parameter unidirectional metamaterial cloak for microwaves. *Nat. Mater.* **12**, 25–28 (2013).

- Gömöry, F., Solovyyov, M., Šouc, J., Navau, C., Part-Camps, J. & Sanchez, A. Experimental realization of a magnetic cloak. *Science* **335**, 1466–1468 (2012).
- Jiang, W. X., Luo, C. Y., Mei, Z. L. & Cui, T. J. An ultrathin but nearly perfect direct current electric cloak. *Appl. Phys. Lett.* **102**, 014102 (2013).
- Chen, H. & Chan, C. T. Acoustic cloaking in three dimensions using acoustic metamaterials. *Appl. Phys. Lett.* **91**, 183518 (2007).
- Zhang, S., Xia, C. & Fang, N. Broadband acoustic cloak for ultrasound waves. *Phys. Rev. Lett.* **106**, 024301 (2011).
- Greenleaf, A., Kurylev, Y., Lassas, M. & Uhlmann, G. Isotropic transformation optics: approximate acoustic and quantum cloaking. *New J. Phys.* **10**, 115024 (2008).
- Zhang, S., Genov, D. A., Sun, C. & Zhang, X. Cloaking of matter waves. *Phys. Rev. Lett.* **100**, 123002 (2008).
- Brun, M., Guenneau, S. & Movchan, A. B. Achieving control of in-plane elastic waves. *Appl. Phys. Lett.* **94**, 061903 (2009).
- Milton, G. W., Briane, M. & Willis, J. R. On cloaking for elasticity and physical equations with a transformation invariant form. *New J. Phys.* **8**, 248 (2006).
- Fan, C., Gao, Y. & Huang, J. Shaped graded materials with an apparent negative thermal conductivity. *Appl. Phys. Lett.* **92**, 251907 (2008).
- Chen, T., Weng, C. N. & Chen, J. S. Cloak for curvilinearly anisotropic media in conduction. *Appl. Phys. Lett.* **93**, 114103 (2008).
- Li, J., Gao, Y. & Huang, J. A bifunctional cloak using transformation media. *J. Appl. Phys.* **108**, 074504 (2010).
- Guenneau, S., Amra, C. & Veynante, D. Transformation thermodynamics: cloaking and concentrating heat flux. *Opt. Express* **20**, 8207–8218 (2012).
- Narayana, S. & Sato, V. Heat flux manipulation with engineered thermal materials. *Phys. Rev. Lett.* **108**, 214303 (2012).
- Li, N. B., Ren, J., Wang, L., Huang, P. & Li, B. Colloquium: Phononics: Manipulating heat flow with electronic analogs and beyond. *Rev. Mod. Phys.* **84**, 1045 (2012).

Acknowledgements

C.W.Q. acknowledges the Grant R-263-000-A23-232 administered by National University of Singapore. T.C.H. also acknowledges the support from the Southwest University (SWU112035).

Author contributions

T.H. and C.W.Q. contributed in theoretical analysis and numerical simulation. T.Y. participated in the simulation, discussion and figure preparation. T.H. prepared the manuscript. B.L. and C.W.Q. revised the manuscript. C.W.Q. conceived the idea and supervised this project.

Additional information

Competing financial interests: The authors declare no competing financial interests.

License: This work is licensed under a Creative Commons Attribution-NonCommercial-NoDerivs 3.0 Unported License. To view a copy of this license, visit <http://creativecommons.org/licenses/by-nc-nd/3.0/>

How to cite this article: Han, T., Yuan, T., Li, B. & Qiu, C. Homogeneous Thermal Cloak with Constant Conductivity and Tunable Heat Localization. *Sci. Rep.* **3**, 1593; DOI:10.1038/srep01593 (2013).