Wright State University CORE Scholar

Physics Faculty Publications

Physics

1-1-2003

Hydrogen Incorporation and Diffusivity in Plasma-Exposed Bulk ZnO

K. Ip

M. E. Overberg

Y. W. Heo

D. P. Norton

S. J. Pearton

See next page for additional authors

Follow this and additional works at: https://corescholar.libraries.wright.edu/physics

Part of the Physics Commons

Repository Citation

Ip, K., Overberg, M. E., Heo, Y. W., Norton, D. P., Pearton, S. J., Stutz, C. E., Luo, B., Ren, F., Look, D. C., & Zavada, J. M. (2003). Hydrogen Incorporation and Diffusivity in Plasma-Exposed Bulk ZnO. *Applied Physics Letters*, *82* (3), 385-387. https://corescholar.libraries.wright.edu/physics/81

This Article is brought to you for free and open access by the Physics at CORE Scholar. It has been accepted for inclusion in Physics Faculty Publications by an authorized administrator of CORE Scholar. For more information, please contact library-corescholar@wright.edu.

Authors

K. Ip, M. E. Overberg, Y. W. Heo, D. P. Norton, S. J. Pearton, C. E. Stutz, B. Luo, F. Ren, David C. Look, and J. M. Zavada

Hydrogen incorporation and diffusivity in plasma-exposed bulk ZnO

K. Ip, M. E. Overberg, Y. W. Heo, D. P. Norton, and S. J. Pearton^{a)} Department of Materials Science and Engineering, University of Florida, Gainesville, Florida 32611

C. E. Stutz

AFRL/MLPS, Wright-Patterson AFB, Ohio 45433

B. Luo and F. Ren

Department of Chemical Engineering, University of Florida, Gainesville, Florida 32611

D. C. Look

Semiconductor Research Center, Wright State University, Dayton, Ohio 45435

J. M. Zavada US Army Research Office, Research Triangle Park, North Carolina 27709

(Received 30 October 2002; accepted 3 December 2002)

Hydrogen incorporation depths of >25 μ m were obtained in bulk, single-crystal ZnO during exposure to ²H plasmas for 0.5 h at 300 °C, producing an estimated diffusivity of ~8 ×10⁻¹⁰ cm²/V·s at this temperature. The activation energy for diffusion was 0.17±0.12 eV, indicating an interstitial mechanism. Subsequent annealing at 500–600 °C was sufficient to evolve all of the hydrogen out of the ZnO, at least to the sensitivity of secondary ion mass spectrometry (<5×10¹⁵ cm⁻³). The thermal stability of hydrogen retention is slightly greater when the hydrogen is incorporated by direct implantation relative to plasma exposure, due to trapping at residual damage in the former case. © 2003 American Institute of Physics. [DOI: 10.1063/1.1539927]

Recently, there has been much interest in the properties of hydrogen in ZnO and related oxides.^{1–18} Theoretical predictions suggest that interstitial atomic hydrogen will introduce a shallow donor state in ZnO,^{1–4} and this assertion is supported by some experimental data for its muonium counterpart,^{5,6} by electron paramagnetic resonance data of bulk single crystals⁷ and by the fact that hydrogen can lead to electron conduction in other wide-band-gap oxides such as SnO.¹⁹ In addition, the presence of hydrogen in the growth or processing ambient can significantly affect the electrical and optical properties of ZnO,^{9–18} but there is little systematic data available on its diffusivity and thermal stability when introduced by different methods.

In this letter, we report on an investigation on the diffusion of hydrogen in ZnO during exposure to ²H plasmas, a comparison of the thermal stability of hydrogen incorporated by either direct implantation or by plasma exposure and finally on changes in the electrical and optical properties of the ZnO or a result of hydrogen incorporation. The hydrogen is found to diffuse very rapidly even at temperatures of 100 °C, but can be evolved from the ZnO by subsequent annealing under N₂ at ≤ 600 °C.

Bulk wurtzite (0001) ZnO crystals from Eagle–Picher (grade I quality) were used in all cases. The samples were nominally undoped with as-received *n*-type carrier concentrations of $\sim 10^{17}$ cm⁻³ and a room temperature mobility of 190 cm²/V·s. The samples were exposed to ²H plasmas at temperatures of 100–300 °C in a Plasma Therm 720 series reactor operating at 900 mTorr with 50 W of 13.56 MHz power. Some of these samples were subsequently annealed at

temperatures up to 600 °C under flowing N₂ ambients for 5 mins cm⁻². Secondary ion mass spectrometry (SIMS) measurements were used to obtain the deuterium profiles as a function of plasma exposure or annealing temperature.²⁰ The electrical properties of some of the samples were examined by electrochemical capacitance–voltage (C-V) measurements using a 0.2 M NaOH/0.1 M EDTA electrolyte as the rectifying contact. Finally, optical properties were measured using photoluminescence (PL) spectroscopy at variable temperatures, with a He–Cd laser as the excitation source.

Figure 1 shows SIMS profiles of ²H in plasma exposed ZnO, for different sample temperatures during the plasma treatment. The profiles follow those expected for diffusion from a constant or semi-infinite source.²¹ The incorporation depths of ²H are very large compared to those in GaN or GaAs under similar conditions, where depths of $1-2 \mu m$ are observed.^{22,23} It is clear that hydrogen must diffuse as an interstitial, with little trapping by the lattice elements or by

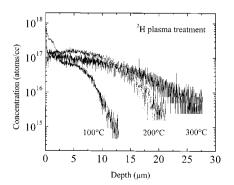


FIG. 1. SIMS profiles of 2 H in ZnO exposed to deuterium plasmas for 0.5 h at different temperatures.

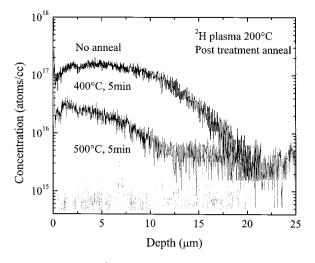


FIG. 2. SIMS profiles of ²H in ZnO exposed to deuterium plasma for 0.5 h at 200 °C and then annealed at 400 or 500 °C for 5 min under flowing N_2 .

defects or impurities. The position of H in the lattice after immobilization has not yet been determined experimentally, but from theory the lowest energy states for H⁺ is at a bondcentered position forming an O-H bond, while for H₂ the antibonding Zn site is most stable.¹

Using a simple estimate of the diffusivity D, from D $=X^{2}/4t$, and where X is taken to be the distance at which ²H concentration has fallen to 5×10^{15} cm⁻³ in Fig. 1, we can estimate the activation energy for diffusion. The extracted activation energy, E_a , is $0.17 \pm 12 \text{ eV}$ for ²H in ZnO. Note that the absolute diffusivities of ¹H would be $\sim 40\%$ larger because of the relationship for diffusivities of isotopes, i.e.,²¹

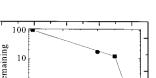
$$\frac{D_{1_H}}{D_{2_H}} = \left(\frac{M_{2_H}}{M_{1_H}}\right)^{1/2}$$

The small activation energy is consistent with the notion that the atomic hydrogen diffuses in interstitial form.

Figure 2 shows SIMS profiles of a ZnO sample exposed to a ²H plasma of 0.5 h at 200 °C, then annealed for 5 min under N₂ at different temperatures. There is significant loss of ²H even after a short anneal at 400 °C, with virtually all of it evolved out of the crystal by 500 °C. This is in sharp contrast to ²H in GaN, where much higher temperatures $(\geq 800 \,^{\circ}\text{C})$ are needed to evolve the deuterium out of the sample.^{22,23}

To compare these data to the thermal stability of ²H incorporated by direct implantation,¹⁸ Fig. 3 shows the percentage of ²H remaining (measured by SIMS) as a function of annealing temperature for incorporation by either plasma exposure or implantation. The ²H is slightly more thermally stable in the latter case, most likely due to trapping at residual damage in the ZnO carried by the nuclear stopping process. Lavrov et al.²⁴ have identified two hydrogen-related defects in ZnO, by using local vibrational mode spectroscopy. The H-I center consists of a hydrogen atom at the bond centered site, while the H-II center contains two inequivalent hydrogen atoms bound primarily to two oxygen atoms.24

Figure 4 shows donor concentration profiles in the ZnO before and after plasma exposure and following subsequent annealing. The ²H plasma treatment causes an increase in



lp et al.

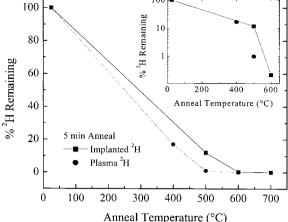


FIG. 3. Percentage of retained ²H incorporated by direct implantation or plasmaexposure, as a function of subsequent annealing temperature. The inset shows the data on a log scale.

donor concentration, consistent with past reports.¹⁰ In that case, the effect was attributed to hydrogen passivation of compensating acceptor impurities present in the as-grown ZnO epitaxial layers.¹⁰ An alternative explanation is that the hydrogen induces a donor state and thereby increases the free electron concentration.¹ Subsequent annealing reduces the carrier density to slightly below the initial value in the asreceived ZnO, which may indicate that it contained hydrogen as a result of the growth process. However, if a 600 °C anneal indeed eliminates hydrogen in as-grown material to the same extent as it does in ion implanted and plasma-treated material, then the data in Fig. 4 must be interpreted to mean that the *n*-type conductivity in the present as-grown ZnO arises from multiple impurity sources 2^{25-27} and not from hvdrogen alone. In other words, although hydrogen appears to be a shallow donor, from this study and others, it is not necessarily the *dominant* shallow donor in as-grown ZnO.

PL spectra from a plasma treated sample was measured as a function of measurement temperature. The sample showed strong band-edge luminescence and a small deeplevel band ($\sim 2.6 \text{ eV}$). Past reports have shown that the efficiency of band-edge emission was increased by plasma hydrogenation of various types of ZnO,¹¹ but that the degree of

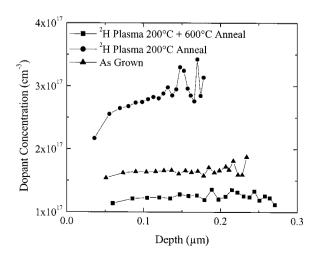


FIG. 4. Donor concentration profiles in ZnO before and after plasma exposure and after subsequent annealing.

improvement depended on the impurity and defect concentration in the original samples.^{11,13} We did not observe any significant difference in the intensity or shape of the PL spectra as a result of plasma hydrogenation of our samples.

In summary, hydrogen is found to exhibit a very rapid diffusion in ZnO when incorporated by plasma exposure, with an estimated value for D of 8.7×10^{-10} cm²/VS at 300 °C. The activation energy for diffusion is indicative of interstitial motion. All of the plasma-incorporated hydrogen is removed from the ZnO by annealing at \geq 500 °C. When the hydrogen is incorporated by direct implantation, the thermal stability is somewhat higher, due to trapping at residual damage. The free electron concentration increases after plasma hydrogenation, consistent with the small ionization energy predicted for H in ZnO¹ and experimentally measured energy as 35 ± 5 meV.⁸ The electrical activity and rapid diffusivity of H in ZnO must be taken into account when designing device fabrication processes such as deposition of dielectrics using SiH₄ as a precursor or dry etching involving use of $CH_4/H_2/Ar$ plasmas since these could lead to significant changes in near-surface conductivity.

The work at University of Florida is partially supported by ARO DAAO 190210420 and NSF (DMR0101438 and CTS 994473). The work at Wright State University was partially supported under AFOSR Grant No. F49620-00-1-0347.

- ¹C. C. van de Walle, Phys. Rev. Lett. **85**, 1012 (2000).
- ²C. C. van de Walle, Phys. Status Solidi B **229**, 221 (2002).
- ³C. Kilic and Z. Zunger, Appl. Phys. Lett. 81, 73 (2002).
- ⁴C. C. van de Walle, Physica B **308-310**, 899 (2001).
- ⁵S. J. F. Cox, E. A. Davis, S. P. Cottrell, P. J. C. King, J. S. Lord, J. M. Gil, H. V. Alberto, R. C. Vilao, DJ. P. Duarte, N. A. de Campos, A. Weidinger,
- R. L. Lichti, and S. J. C. Irving, Phys. Rev. Lett. **86**, 2601 (2001).
- K. L. Lichu, and S. J. C. Hving, Fliys. Rev. Lett. $\mathbf{00}$, 2001 (2001).
- ⁶S. J. F. Cox, E. A. Davis, P. J. C. King, J. M. Gil, H. V. Alberto, R. C. Vilao, J. P. Duarte, N. A. de Campos, and R. L. Lichti, J. Phys. C **13**, 9001 (2001).

- ⁷S. O. Kucheyev, P. N. K. Deenapannay, C. Jagadish, J. S. Williams, M. Yano, K. Koike, S. Sasa, M. Inoue, and K. Ogata, Appl. Phys. Lett. 81, 3350 (2002).
- ⁸D. M. Hofmann, A. Hofstaetter, F. Leiter, H. Zhou, F. Henecker, B. K. Meyer, S. B. Schmidt, and P. G. Baranov, Phys. Rev. Lett. 88, 045504 (2002).
- ⁹S. J. Baik, J. H. Jang, C. H. Lee, W. Y. Cho, and K. S. Lim, Appl. Phys. Lett. **70**, 3516 (1997).
- ¹⁰B. Theys, V. Sallet, F. Jomard, A. Lusson, J.-F. Rommeluere, and Z. Teukam, J. Appl. Phys. **91**, 3922 (2002).
- ¹¹N. Ohashi, T. Ishigaki, N. Okada, T. Sekiguchi, I. Sakaguchi, and H. Haneda, Appl. Phys. Lett. **80**, 2869 (2002).
- ¹² V. Bogatu, A. Goldenbaum, A. Many, and Y. Goldstein, Phys. Status Solidi B **212**, 89 (1999).
- ¹³T. Sekiguchi, N. Ohashi, and Y. Terada, Jpn. J. Appl. Phys., Part 2 36, L289 (1997).
- ¹⁴S. Tüzemen, G. Xiong, J. Wilkinson, B. Mischuck, K. B. Ucer, and R. T. Williams, Physica B **308-310**, 1197 (2001).
- ¹⁵C. S. Han, J. Jun, and H. Kim, Appl. Surf. Sci. 175/176, 567 (2001).
- ¹⁶Y. Natsume and H. Sakata, J. Mater. Sci.: Mater. Electron. **12**, 87 (2001).
- ¹⁷Y.-S. Kang, H. Y. Kim, and J. Y. Lee, J. Electrochem. Soc. **147**, 4625 (2000).
- ¹⁸K. Ip, M. E. Overberg, Y. W. Heo, D. P. Norton, S. J. Pearton, S. O. Kucheyev, C. Jagadish, J. S. Williams, R. G. Wilson, and J. M. Zavada Appl. Phys. Lett. **81**, 3996 (2002).
- ¹⁹S. Samson and C. G. Fonstad, J. Appl. Phys. 44, 4618 (1973).
- ²⁰SIMS performed at Charles Evans and Associates using a Cameca system.
- ²¹R. J. Borg and C. J. Dienes, An Introduction to Solid State Diffusion (Academic, Boston, 1988).
- ²² R. G. Wilson, S. J. Pearton, C. R. Abernathy, and J. M. Zavada, J. Vac. Sci. Technol. A **13**, 719 (1995).
- ²³S. J. Pearton, J. C. Zolper, R. J. Shul, and F. Ren, J. Appl. Phys. 86, 1 (1999).
- ²⁴ E. V. Lavrov, J. Weber, F. Borrnet, C. G. Van de Walle, and R. Helbig, Phys. Rev. B 66, 165205 (2002).
- ²⁵D. G. Look, Mater. Sci. Eng., B 80, 383 (2001).
- ²⁶D. G. Look, J. W. Hemsky, and J. R. Sizelove, Phys. Rev. Lett. 82, 2552 (1999).
- ²⁷D. C. Look, D. C. Reynolds, J. W. Hemsky, R. L. Jones, and J. R. Sizelove, Appl. Phys. Lett. **75**, 811 (1999).