

Importance of Loop Effects in Explaining the Accumulated Evidence for New Physics in B Decays with a Vector Leptoquark

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In recent years experiments revealed intriguing hints for new physics (NP) in B decays involving $b \rightarrow c\tau\nu$ and $b \rightarrow s\ell^+\ell^-$ transitions at the 4σ and 5σ level, respectively. In addition, there are slight disagreements in $b \rightarrow u\tau\nu$ and $b \rightarrow d\mu^+\mu^-$ observables. While not significant on their own, they point in the same direction. Furthermore, V_{us} extracted from τ decays shows a slight tension ($\approx 2.5\sigma$) with its value determined from Cabibbo-Kobayashi-Maskawa unitarity, and an analysis of BELLE data found an excess in $B_d \rightarrow \tau^+\tau^-$. Concerning NP explanations, the vector leptoquark SU(2) singlet is of special interest since it is the only single particle extension of the standard model which can (in principle) address all the anomalies described above. For this purpose, large couplings to τ leptons are necessary and loop effects, which we calculate herein, become important. Including them in our phenomenological analysis, we find that neither the tension in V_{us} nor the excess in $B_d \rightarrow \tau^+\tau^-$ can be fully explained without violating bounds from $K \rightarrow \pi\nu\bar{\nu}$. However, one can account for $b \rightarrow c\tau\nu$ and $b \rightarrow u\tau\nu$ data finding intriguing correlations with $B_q \rightarrow \tau^+\tau^-$ and $K \rightarrow \pi\nu\bar{\nu}$. Furthermore, the explanation of $b \rightarrow c\tau\nu$ predicts a positive shift in C_7 and a negative one in C_9 , being nicely in agreement with the global fit to $b \rightarrow s\ell^+\ell^-$ data. Finally, we point out that one can fully account for $b \rightarrow c\tau\nu$ and $b \rightarrow s\ell^+\ell^-$ without violating bounds from $\tau \rightarrow \phi\mu$, $\Upsilon \rightarrow \tau\mu$, or $b \rightarrow s\tau\mu$ processes.

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Introduction.—So far, the LHC has not directly observed any particles beyond the standard model (SM). However, intriguing hints for lepton flavor universality (LFU) violating new physics (NP) have been acquired.

$b \rightarrow s(d)\ell^+\ell^-$.—The ratios

$$R(K^{(*)}) = \frac{\text{Br}[B \rightarrow K^{(*)}\mu^+\mu^-]}{\text{Br}[B \rightarrow K^{(*)}e^+e^-]} \quad (1)$$

[1,2] indicate LFU violation with a combined significance of $\approx 4\sigma$ [3–8]. Taking also into account all other $b \rightarrow s\mu^+\mu^-$ observables, like the angular observable P'_5 [9] in the decay $B \rightarrow K^*\mu^+\mu^-$, the global fit of the Wilson coefficients to all

available data even shows compelling evidence [10] for NP ($>5\sigma$).

Concerning $b \rightarrow d\ell^+\ell^-$ transitions, the theoretical analysis of Ref. [11] shows that the LHCb measurement of $B \rightarrow \pi\mu^+\mu^-$ [12] slightly differs from the theory expectation. Even though this is not significant on its own, the central value is very well in agreement with the expectation from $b \rightarrow s\ell^+\ell^-$ under the assumption of a V_{td}/V_{ts} -like scaling of the NP effect [13]. In other words, an effect of the same order and sign as in $b \rightarrow s\ell^+\ell^-$, relative to the SM, is preferred. Furthermore, an (unpublished) analysis of BELLE data found an excess in $B_d \rightarrow \tau^+\tau^-$ [14].

$b \rightarrow c(u)\tau\nu$.—The ratios

$$R(D^{(*)}) = \frac{\text{Br}[B \rightarrow D^{(*)}\tau\nu]}{\text{Br}[B \rightarrow D^{(*)}\ell\nu]} \quad \text{with } \ell = \{e, \mu\}, \quad (2)$$

which measure LFU violation in the charged current by comparing τ modes with light leptons ($\ell = e, \mu$), differ in combination from their SM predictions by $\approx 4\sigma$ [15]. Also, the ratio

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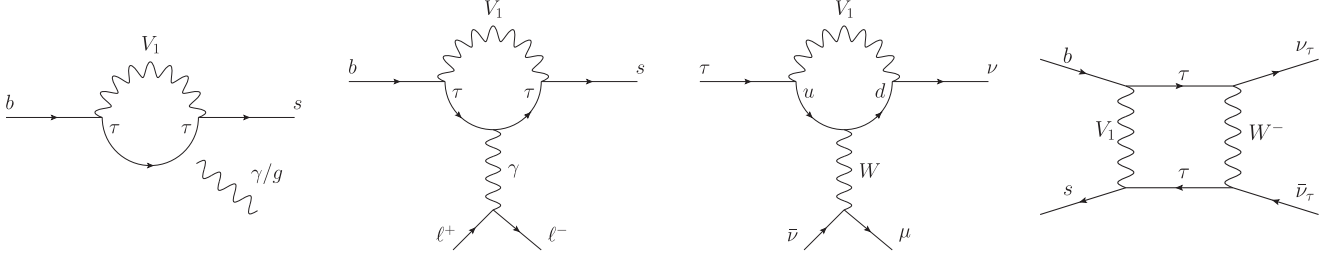


FIG. 1. Feynman diagrams depicting the one-loop contributions of the vector LQ singlet to $C_{7/8}^{sb}$, $b \rightarrow s\ell^+\ell^-$, $\tau \rightarrow \mu\nu\bar{\nu}$, and $b \rightarrow s\nu\bar{\nu}$ (from left to right).

$$R(J/\psi) = \frac{\text{Br}[B_c \rightarrow J/\psi\tau\nu]}{\text{Br}[B_c \rightarrow J/\psi\mu\nu]} \quad (3)$$

[16] exceeds the SM prediction in agreement with the expectations from $R(D^{(*)})$ [17,18].

Concerning $b \rightarrow u\tau\nu$ transitions, the theory prediction for $B \rightarrow \tau\nu$ crucially depends on V_{ub} . While previous lattice calculations resulted in rather small values of V_{ub} , recent calculations give a larger value (see Ref. [19] for an overview). However, the measurement is still above the SM prediction by more than 1σ , as can be seen from the global fit [20]. In

$$R(\pi) = \frac{\text{Br}[B \rightarrow \pi\tau\nu]}{\text{Br}[B \rightarrow \pi\ell\nu]} \quad (4)$$

there is also a small disagreement between theory [21] and experiment [22] which does not depend on V_{ub} . These results are not significant on their own but lie again above the SM predictions like in the case of $b \rightarrow c\tau\nu$.

V_{us}^τ .— V_{us} extracted from τ lepton decays (V_{us}^τ) shows a tension of 2.5σ compared to the value of V_{us} determined from Cabibbo-Kobayashi-Maskawa (CKM) unitarity (V_{us}^{uni}) [15,23].

The only possible single particle explanation, which can (at least in principle) address all these anomalies is the vector leptoquark (LQ) $SU(2)_L$ singlet V_1 with hypercharge [24] $-4/3$ [25–31] arising in the famous Pati-Salam model [32]: This LQ can explain $b \rightarrow c\tau\nu$ data without violating bounds from $b \rightarrow s\nu\bar{\nu}$ and/or direct searches, provides (at tree level) a left-handed solution to $b \rightarrow s\ell^+\ell^-$ data, and does not lead to proton decay. Therefore, a sizable effect in $b \rightarrow u\tau\nu$ and $b \rightarrow d\ell^+\ell^-$ is straightforward, and also an explanation of V_{us}^τ could be possible. A huge enhancement of $b \rightarrow s\tau^+\tau^-$ rates is predicted as well [33], making an amplification of $B_d \rightarrow \tau^+\tau^-$ possible.

Several attempts to construct a UV completion for this LQ to address the anomalies have been made [34–44]. In order to fully account for the $b \rightarrow c\tau\nu$ data (while respecting perturbativity), one needs sizable couplings to third generation leptons and V_1 generates, via $SU(2)_L$ invariance, also large contributions to the operators $d_i d_j \tau \tau$ and $u_i u_j \nu_\tau \nu_\tau$ at tree level. These operators give rise to couplings

of down quarks to neutrinos or light charged leptons at loop level (see Fig. 1).

In this Letter we will calculate these loop effects [45], which turn out to be not only numerically important but also give rise to additional correlations among observables. Even though a theory with a massive vector boson without an explicit Higgs sector is not renormalizable, we still identify several phenomenologically important loop effects which are gauge independent and finite and can therefore be calculated reliably (in analogy to flavor observables within the SM).

Model and one-loop effects.—We work in a simplified model extending the SM by a vector LQ $SU(2)_L$ singlet with hypercharge $-4/3$, mass M , and interactions with fermions determined by

$$\mathcal{L}_{V^\mu} = (\kappa_{fi}^L \bar{Q}_f \gamma_\mu L_i + \kappa_{fi}^R \bar{d}_f \gamma_\mu e_i) V_1^{\mu\dagger} + \text{H.c.} \quad (5)$$

Here, Q (L) are quark (lepton) $SU(2)_L$ doublets, d (e) are down quark (charged lepton) singlets, and f, i are flavor indices. In the following, we will neglect the right-handed couplings, which are not necessary to explain the anomalies. This then generates the effective four-fermion interactions encoded in

$$\mathcal{L}_{\text{eff}} = -\frac{\kappa_{il}^L \kappa_{jk}^{L*}}{M^2} \bar{Q}_i^\alpha \gamma^\mu Q_j^\beta \bar{L}_k^\beta \gamma_\mu L_l^\alpha, \quad (6)$$

where α and β label the $SU(2)$ components. After electroweak symmetry breaking, we work in the down basis; i.e., no CKM elements appear in flavor changing neutral currents of down quarks. We recall our definitions and the tree-level results in the Supplemental Material [49], which includes Refs. [50–66].

In our setup, one-loop effects involving the LQ and third generation leptons (τ 's and τ neutrinos) can be very important, since we aim for large effects in $b \rightarrow c(u)\tau\nu$ and $b \rightarrow s(d)\tau^+\tau^-$ processes. In principle, a massive vector boson, like our LQ, without a Higgs sector is not renormalizable. However, in flavor physics most effects can still be calculated reliably since they are gauge independent and finite (also in unitary gauge) [67]. This is in analogy to the SM, where the contribution of the W to flavor observables

can be correctly calculated in unitary gauge without taking into account the Higgs sector.

We are only interested in effects which are always absent at tree level (like $b \rightarrow s\nu\bar{\nu}$ processes) or are not present at tree level due to a specific coupling structure (like $b \rightarrow s\mu^+\mu^-$ processes in the absence of muon couplings). Furthermore, we neglect tiny dimension-8 effects of the SM Higgs particle. In these cases the loop effects are the leading contributions. We calculate all diagrams at leading order in the external momenta using asymptotic expansion [68].

W boxes contributing to $d_i \rightarrow d_f\nu\bar{\nu}$.—We use the effective Hamiltonian,

$$\begin{aligned} \mathcal{H}_{\text{eff}}^{\nu\nu} &= -\frac{4G_F}{\sqrt{2}} V_{td_k} V_{id_j}^* (C_{L,jk}^{fi} \mathcal{O}_{L,jk}^{fi} + C_{R,jk}^{fi} \mathcal{O}_{R,jk}^{fi}), \\ \mathcal{O}_{L,jk}^{fi} &= \frac{\alpha}{4\pi} [\bar{d}_j \gamma^\mu P_L d_k] [\bar{\nu}_f \gamma_\mu (1 - \gamma_5) \nu_i], \\ \mathcal{O}_{R,jk}^{fi} &= \frac{\alpha}{4\pi} [\bar{d}_j \gamma^\mu P_R d_k] [\bar{\nu}_f \gamma_\mu (1 - \gamma_5) \nu_i], \end{aligned} \quad (7)$$

with $P_{R(L)} = [1 + (-)\gamma_5]/2$ and G_F (α) being the Fermi (electromagnetic fine structure) constant. The result of the box contributions involving a W to $d_i \rightarrow d_f\nu\bar{\nu}$ (an example diagram is shown on the right-hand side of Fig. 1) is gauge invariant in R_ξ gauge and the same finite result is obtained in unitary gauge [with $e = \sqrt{4\pi\alpha}$ and m_t (m_W) the top quark (W boson) mass]:

$$\begin{aligned} C_{L,fa}^{ij} &= \frac{-m_W^2}{2e^2 V_{3a} V_{3f}^* M^2} \left(6\kappa_{fj}^L \kappa_{ai}^{L*} \log\left(\frac{m_W^2}{M^2}\right) \right. \\ &\quad + 3(V_{3a} V_{3k}^* \kappa_{ki}^{L*} \kappa_{fj}^L + V_{3f}^* V_{3k} \kappa_{kj}^L \kappa_{ai}^{L*}) \frac{\log\left(\frac{m_t^2}{m_W^2}\right)}{1 - \frac{m_W^2}{m_t^2}} \\ &\quad \left. + V_{3f}^* V_{3k} \kappa_{kj}^L V_{3a} V_{3l}^* \kappa_{li}^{L*} \frac{m_t^2}{m_W^2} \right). \end{aligned} \quad (8)$$

W off-shell penguins contributing to $\tau \rightarrow \mu\nu\bar{\nu}$.—Here (see third diagram in Fig. 1) we obtain again a finite and gauge independent result for the Wilson coefficient; following the analysis of Ref. [69], we use

$$\mathcal{H}_{\text{eff}}^{\tau\mu\nu\nu_i} = \frac{4G_F}{\sqrt{2}} D_{L,fi}^{\tau\mu} [\bar{\nu}_f \gamma^\sigma P_L \nu_i] [\bar{\mu} \gamma_\sigma P_L \tau], \quad (9)$$

with

$$D_{L,fi}^{\tau\mu} = N_c \delta_{i2} \frac{V_{3k}^* \kappa_{kf}^{L*} \kappa_{33}^L m_t^2}{32\pi^2 M^2} \left[1 + 2 \log\left(\frac{m_t^2}{M^2}\right) \right]. \quad (10)$$

We find, in agreement with Ref. [30], that the effect is small.

Photon and gluon penguins.—We use the standard Hamiltonian (see, for example, Ref. [70]) also defined in

the Supplemental Material [49]. For on-shell photons and gluons the result of the left-hand diagram in Fig. 1 is finite in unitary gauge and the same result is obtained in R_ξ gauge:

$$C_{7(8)}^{sb} = \frac{-\sqrt{2}}{G_F V_{tb} V_{ts}^* M^2} \frac{11}{72} \left(\frac{5}{48} \right) \kappa_{2i}^L \kappa_{3i}^{L*}. \quad (11)$$

Taking into account the running from the LQ scale $\mu_{\text{LQ}} = M = 1$ TeV down to $\mu_b = 5$ GeV (see, e.g., Refs. [71,72]), we obtain

$$C_7^{sb}(\mu_b) \approx 0.29 \kappa_{2i}^L \kappa_{3i}^{L*}. \quad (12)$$

For off-shell photons the full result (second diagram in Fig. 1) for the amplitude is gauge dependent and, in general, divergent. However, one can calculate the mixing of $C_{9,sb}^{\tau\tau} = -C_{10,sb}^{\tau\tau}$ into the four-fermion operators $O_{9,sb}^{\ell\ell}$ (containing light leptons as well) within the effective theory (i.e., after integrating out the LQ at tree level). In this way, a gauge independent result is obtained and the leading logarithm of the (unknown) full result is recovered. For off-shell photons we thus calculate the effect in the effective field theory (below the LQ scale), generating the following mixing into the four-fermion operators with light leptons:

$$C_{9,sb}^{\ell\ell} = \frac{\sqrt{2}}{G_F V_{tb} V_{ts}^* M^2} \frac{1}{6} \log\left(\frac{M^2}{\mu_b^2}\right) \kappa_{2i}^L \kappa_{3i}^{L*}. \quad (13)$$

Note that this result is model independent (at leading-log accuracy) in the sense that it does not depend on the model which generates $C_{9,sb}^{\tau\tau} = -C_{10,sb}^{\tau\tau}$. In principle, there are also Z penguins generating $C_{9,sb}^{\ell\ell}$ and $C_{10,sb}^{\ell\ell}$. However, this effect is suppressed by light lepton masses (or small momenta) and is therefore of dimension 8. Further, note that there are no box diagram contributions which generate $\bar{s}b\bar{\mu}\mu$ ($\bar{s}b\bar{e}e$) operators if the couplings of the LQ to muons (electrons) are zero at tree level.

Box diagrams with LQs.—What cannot be calculated consistently are box diagrams involving only LQs [35]. Here, the results are divergent in unitary gauge which corresponds to a gauge dependence in R_ξ gauge. However, these effects are suppressed if $|\kappa^L| < g_2$ and can be further suppressed in the presence of vectorlike fermions by a GIM-like mechanism [38] which, in analogy to the SM, would render the result finite.

Phenomenology.—Assuming $\kappa_{33}^L V_{cb} \ll \kappa_{23}^L$, one is safe from LHC bounds, and the effects in $B_s \rightarrow \tau^+\tau^-$, $C_7^{sb}(\mu_b)$ [Eq. (12)], and $C_{9,sb}^{\ell\ell}$ [Eq. (13)] directly depend on $R(X)/R(X)_{\text{SM}}$ (with $X = D^{(*)}, J/\psi$). In Fig. 2 we show these dependences. Intriguingly, the effect generated in $C_7^{sb}(\mu_b)$ and $C_{9,sb}^{\ell\ell}$, within the preferred region from $b \rightarrow c\tau\nu$ data, exactly overlaps with the 1σ ranges of the model independent fit to $b \rightarrow s\mu^+\mu^-$ data excluding LFU violating

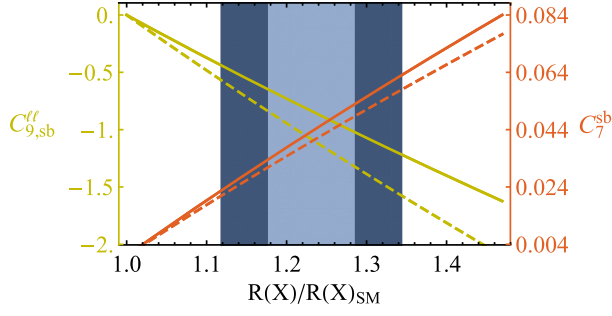


FIG. 2. $C_{9, sb}^{\ell\ell}$ and $C_7^{sb}(\mu_b)$ as functions of $R(X)/R(X)_{\text{SM}}$ with $X = \{D, D^*, J/\psi\}$. The solid lines correspond to $M = 1$ TeV and the dashed ones to $M = 5$ TeV while the (dark) blue region is preferred by $b \rightarrow c\tau\nu$ data at the 1σ (2σ) level. From the global fit, taking into account only lepton flavor conserving observables we have $-1.29 < C_{9, sb}^{\ell\ell} < -0.87$ [70] and $-0.01 < C_7^{sb}(\mu_b) < 0.05$ [10] at the 1σ level. Therefore, our model predicts just the right sign and size of the effect in $C_{9, sb}^{\ell\ell}$ and $C_7^{sb}(\mu_b)$ necessary to explain $b \rightarrow s\ell^+\ell^-$ data, assuming an explanation of $b \rightarrow c\tau\nu$.

observables [70,73] [therefore, only lepton flavor universality conserving observables like P'_5 but not $R(K^{(*)})$ can be explained].

Let us now include the effect of κ_{13}^L . Here, many correlations arise. First of all, $b \rightarrow c(u)\tau\nu$ is already at tree level correlated to $b \rightarrow s(d)\tau^+\tau^-$. In addition, the W boxes in Eq. (8) generate effects in $B \rightarrow K^{(*)}(\pi)\nu\bar{\nu}$ and $K \rightarrow \pi\nu\bar{\nu}$. While the bounds from $B \rightarrow K^{(*)}(\pi)\nu\bar{\nu}$ turn out to be weaker than the ones from $B_q \rightarrow \tau^+\tau^-$, there are striking correlations with $K \rightarrow \pi\nu\bar{\nu}$, as can be seen from Fig. 3. Furthermore, we get an effect

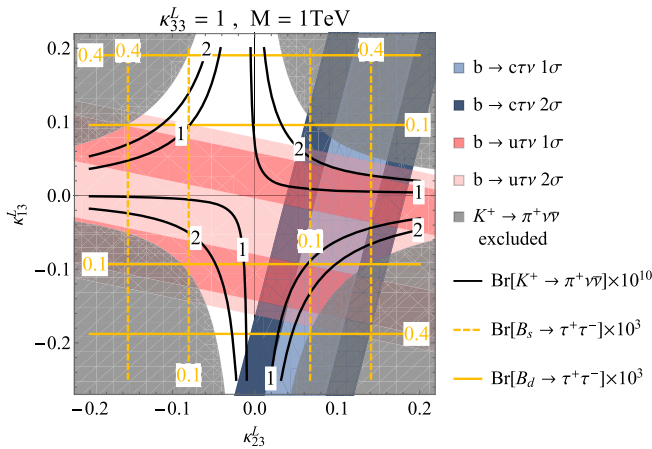


FIG. 3. Predictions for $B_q \rightarrow \tau^+\tau^-$ and $K \rightarrow \pi\nu\bar{\nu}$ (contour lines) in the $\kappa_{13}^L - \kappa_{23}^L$ plane for $M = 1$ TeV and $\kappa_{33}^L = 1$. The colored regions are preferred by $b \rightarrow c(u)\tau\nu$ data, where we naively averaged (i.e., we computed the weighted average of the observables and added their errors in quadrature, disregarding correlations) $R(D^{(*)})$ and $R(J/\psi)$ or $R(\pi)$ and $B \rightarrow \tau\nu$, respectively. The gray region is excluded by $K^+ \rightarrow \pi^+\nu\bar{\nu}$. Here we assumed all couplings κ_{ij}^L to be real.

$$\delta V_{us}^\tau = \frac{V_{us}^\tau - V_{us}^{\tau(0)}}{V_{us}^{\text{uni}}} \approx -C_{us}^{\tau\tau}, \quad (14)$$

where $V_{us}^{\tau(0)}$ is the CKM matrix element extracted from τ decays without NP. However, Eq. (8) generates $K \rightarrow \pi\nu\bar{\nu}$, and respecting these bounds, the relative effect in V_{us}^τ can only be at the per-mill level, $|\delta V_{us}^\tau| \approx 0.05\%$, excluding the possibility to account for the discrepancy of $|V_{us}^{\text{uni}}| = 0.22547 \pm 0.00095$ versus $|V_{us}^\tau| = 0.2212 \pm 0.0014$ [15,23]. The same is true about $B_d \rightarrow \tau^+\tau^-$, where the currently preferred region of analysis using BELLE data [14] of $\text{Br}[B_d \rightarrow \tau^+\tau^-]_{\text{exp}} = (4.39_{-0.83}^{+0.80} \pm 0.45)$ lies outside the plot range.

Now, in addition to the couplings κ_{33}^L and κ_{23}^L , we allow nonvanishing κ_{32}^L and κ_{22}^L . These couplings give rise to tree-level effects in $b \rightarrow s\mu^+\mu^-$. In Fig. 4 we show the allowed (colored) regions from $b \rightarrow s\mu^+\mu^-$ and $b \rightarrow c\tau\nu$ as well as the exclusions from $b \rightarrow s\tau\mu$ and $\tau \rightarrow \phi\mu$. Note that a simultaneous explanation of the anomalies is perfectly possible since the colored regions overlap and do not extend to the parameter space excluded by $b \rightarrow s\tau\mu$ and $\tau \rightarrow \phi\mu$. Interestingly, due to the loop effects originating from the $b \rightarrow c\tau\nu$ explanation, we predict a flavor universal effect in $C_{9, sb}^{\ell\ell}$ and C_7^{sb} which is supplemented by a tree-level effect of the form $C_{9, sb}^{\mu\mu} = -C_{10, sb}^{\mu\mu}$ with muons only. This means that the relative NP effect compared to the SM in lepton flavor conserving observables (like P'_5) should be larger than in $R(K^{(*)})$, which is in perfect agreement with the global fit [74].

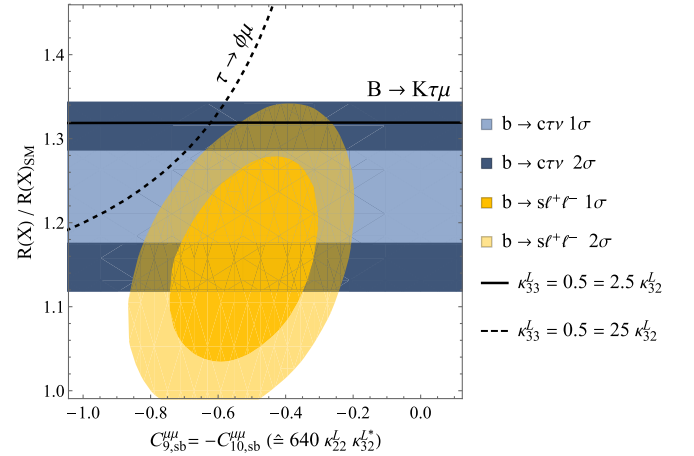


FIG. 4. Allowed (colored) regions in the $C_{9, sb}^{\mu\mu} = -C_{10, sb}^{\mu\mu}$ ($\hat{=} 640\kappa_{22}^L\kappa_{32}^L$)— $R(X)/R(X)_{\text{SM}}$ plane for $M = 1$ TeV and $X = D, D^*, J/\psi$ at the 1σ and 2σ level for $\kappa_{33}^L V_{cb} \ll \kappa_{23}^L$. The region above the black dashed (solid) line is excluded by $\tau \rightarrow \phi\mu$ ($B \rightarrow K\tau\mu$) for $\kappa_{33}^L = 0.5 = 25\kappa_{32}^L$ ($\kappa_{33}^L = 0.5 = 2.5\kappa_{32}^L$). The bound from $\tau \rightarrow \phi\mu$ ($B \rightarrow K\tau\mu$) depends on κ_{33}^L and κ_{32}^L and gets stronger if κ_{32}^L gets smaller (larger). That is, for $\kappa_{33}^L = 0.5$ and $2.7 \lesssim \kappa_{33}^L/\kappa_{32}^L \lesssim 27$, the whole 2σ region preferred by $b \rightarrow c\tau\nu$ and $b \rightarrow s\ell^+\ell^-$ data is consistent with $B \rightarrow K\tau\mu$ and $\tau \rightarrow \phi\mu$.

Conclusions.—The vector leptoquark SU(2) singlet is a prime NP candidate to explain the current hints for LFU violation. In this Letter, we calculated and studied the important loop effects arising within such a model and performed a phenomenological analysis.

We find that an explanation of $b \rightarrow c\tau\nu$ data generates lepton flavor universal effects in $b \rightarrow s\ell^+\ell^-$ transitions which nicely agree with the model independent fit (see Fig. 2). Therefore, the $C_9 = -C_{10}$ -like tree-level effect, which is in general LFU violating, is supplemented by these effects generating a new pattern for the Wilson coefficients. This can be tested with future data. That is, with more precise measurements of lepton flavour universality violating and lepton flavor universality conserving effects, one can test if in fact there is a lepton flavor universality conserving contribution in addition to the lepton flavor universality violating ones [75]. Similar conclusions hold for the correlations between $b \rightarrow u\tau\nu$ data generating lepton flavor universal effects in $b \rightarrow d\ell^+\ell^-$ processes.

We also find that NP in $b \rightarrow c(u)\tau\nu$ generates important effects in $B_{s(d)} \rightarrow \tau^+\tau^-$ which are even correlated to $b \rightarrow s(d)\nu\bar{\nu}$ processes and $K \rightarrow \pi\nu\bar{\nu}$ via W box contributions (see right-hand diagram in Fig. 1). The V_{us}^τ puzzle (like the CP asymmetry in $\tau \rightarrow K_S\pi\nu$ [76]) cannot be solved due to the stringent constraints from $K \rightarrow \pi\nu\bar{\nu}$, and because of $b \rightarrow u\tau\nu$ bounds one cannot fully account for the BELLE excess in $B_d \rightarrow \tau^+\tau^-$ (see Fig. 3).

Additionally, $b \rightarrow c\tau\nu$ and $b \rightarrow s\ell^+\ell^-$ data can be simultaneously explained without violating other bounds like $\tau \rightarrow \phi\mu$ (see Fig. 4). Furthermore, one could at the same time also account for NP effects in $b \rightarrow d\mu^+\mu^-$ without violating $K_L \rightarrow \mu^+\mu^-$ bounds.

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- [1] R. Aaij *et al.* (LHCb Collaboration), *Phys. Rev. Lett.* **113**, 151601 (2014).
- [2] R. Aaij *et al.* (LHCb Collaboration), *J. High Energy Phys.* **08** (2017) 055.
- [3] W. Altmannshofer, P. Stangl, and D. M. Straub, *Phys. Rev. D* **96**, 055008 (2017).
- [4] G. D’Amico, M. Nardecchia, P. Panci, F. Sannino, A. Strumia, R. Torre, and A. Urbano, *J. High Energy Phys.* **09** (2017) 010.
- [5] M. Ciuchini, A. M. Coutinho, M. Fedele, E. Franco, A. Paul, L. Silvestrini, and M. Valli, *Eur. Phys. J. C* **77**, 688 (2017).
- [6] G. Hiller and I. Nisandzic, *Phys. Rev. D* **96**, 035003 (2017).
- [7] L.-S. Geng, B. Grinstein, S. Jäger, J. Martin Camalich, X.-L. Ren, and R.-X. Shi, *Phys. Rev. D* **96**, 093006 (2017).
- [8] T. Hurth, F. Mahmoudi, D. Martinez Santos, and S. Neshatpour, *Phys. Rev. D* **96**, 095034 (2017).
- [9] R. Aaij *et al.* (LHCb Collaboration), *J. High Energy Phys.* **02** (2016) 104.
- [10] B. Capdevila, A. Crivellin, S. Descotes-Genon, J. Matias, and J. Virto, *J. High Energy Phys.* **01** (2018) 093.
- [11] C. Hambrock, A. Khodjamirian, and A. Rusov, *Phys. Rev. D* **92**, 074020 (2015).
- [12] R. Aaij *et al.* (LHCb Collaboration), *J. High Energy Phys.* **10** (2015) 034.
- [13] Here, V refers to the Cabibbo-Kobayashi-Maskawa (CKM) matrix.
- [14] M. Ziegler, Ph.D. thesis, KIT, Karlsruhe, 2016, <http://belle.kek.jp/belle/theses/doctor/ziegler16.pdf>.
- [15] Y. Amhis *et al.* (HFLAV Collaboration), *Eur. Phys. J. C* **77**, 895 (2017).
- [16] R. Aaij *et al.* (LHCb Collaboration), *Phys. Rev. Lett.* **120**, 121801 (2018).
- [17] R. Watanabe, *Phys. Lett. B* **776**, 5 (2018).
- [18] B. Chauhan and B. Kindra, [arXiv:1709.09989](https://arxiv.org/abs/1709.09989).
- [19] G. Ricciardi, *Mod. Phys. Lett. A* **32**, 1730005 (2017).
- [20] J. Charles, A. Hocker, H. Lacker, S. Laplace, F. R. Le Diberder, J. Malcles, J. Ocariz, M. Pivk, and L. Roos (CKMfitter Group), *Eur. Phys. J. C* **41**, 1 (2005).
- [21] F. U. Bernlochner, *Phys. Rev. D* **92**, 115019 (2015).
- [22] P. Hamer *et al.* (Belle Collaboration), *Phys. Rev. D* **93**, 032007 (2016).
- [23] A. Lusiani, *Nucl. Part. Phys. Proc.* **287–288**, 29 (2017).
- [24] In our conventions, the left-handed lepton doublet has hypercharge -1 .
- [25] R. Alonso, B. Grinstein, and J. Martin Camalich, *J. High Energy Phys.* **10** (2015) 184.
- [26] L. Calibbi, A. Crivellin, and T. Ota, *Phys. Rev. Lett.* **115**, 181801 (2015).
- [27] S. Fajfer and N. Košnik, *Phys. Lett. B* **755**, 270 (2016).
- [28] G. Hiller, D. Loose, and K. Schönwald, *J. High Energy Phys.* **12** (2016) 027.
- [29] B. Bhattacharya, A. Datta, J.-P. Guévin, D. London, and R. Watanabe, *J. High Energy Phys.* **01** (2017) 015.
- [30] D. Buttazzo, A. Greljo, G. Isidori, and D. Marzocca, *J. High Energy Phys.* **11** (2017) 044.
- [31] J. Kumar, D. London, and R. Watanabe, [arXiv:1806.07403](https://arxiv.org/abs/1806.07403).
- [32] J. C. Pati and A. Salam, *Phys. Rev. D* **10**, 275 (1974); **11**, 703(E) (1975).
- [33] B. Capdevila, A. Crivellin, S. Descotes-Genon, L. Hofer, and J. Matias, *Phys. Rev. Lett.* **120**, 181802 (2018).
- [34] R. Barbieri, G. Isidori, A. Pattori, and F. Senia, *Eur. Phys. J. C* **76**, 67 (2016).
- [35] R. Barbieri, C. W. Murphy, and F. Senia, *Eur. Phys. J. C* **77**, 8 (2017).

- [36] N. Assad, B. Fornal, and B. Grinstein, *Phys. Lett. B* **777**, 324 (2018).
- [37] L. Di Luzio, A. Greljo, and M. Nardecchia, *Phys. Rev. D* **96**, 115011 (2017).
- [38] L. Calibbi, A. Crivellin, and T. Li, *Phys. Rev. D* **98**, 115002 (2018).
- [39] M. Bordone, C. Cornella, J. Fuentes-Martin, and G. Isidori, *Phys. Lett. B* **779**, 317 (2018).
- [40] R. Barbieri and A. Tesi, *Eur. Phys. J. C* **78**, 193 (2018).
- [41] M. Blanke and A. Crivellin, *Phys. Rev. Lett.* **121**, 011801 (2018).
- [42] A. Greljo and B. A. Stefanek, *Phys. Lett. B* **782**, 131 (2018).
- [43] M. Bordone, C. Cornella, J. Fuentes-Martín, and G. Isidori, *J. High Energy Phys.* **10** (2018) 148.
- [44] S. Matsuzaki, K. Nishiwaki, and K. Yamamoto, *J. High Energy Phys.* **11** (2018) 164.
- [45] Similar loop effects for scalar LQs have been calculated in Refs. [46–48].
- [46] C. Bobeth and A. J. Buras, *J. High Energy Phys.* **02** (2018) 101.
- [47] S. Fajfer, N. Košnik, and L. Vale Silva, *Eur. Phys. J. C* **78**, 275 (2018).
- [48] K. Earl and T. Gregoire, *J. High Energy Phys.* **08** (2018) 201.
- [49] See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevLett.122.011805> for the calculation of the observables and a possible UV completion of the EW symmetry breaking mechanism.
- [50] R. Aaij *et al.* (LHCb Collaboration), *Phys. Rev. Lett.* **118**, 251802 (2017).
- [51] C. Bobeth, M. Gorbahn, T. Hermann, M. Misiak, E. Stamou, and M. Steinhauser, *Phys. Rev. Lett.* **112**, 101801 (2014).
- [52] C. Bobeth, in *Proceedings of the 49th Rencontres de Moriond on Electroweak Interactions and Unified Theories, La Thuile, Italy* (2014), pp. 75–80 [arXiv:1405.4907].
- [53] C. Bobeth, M. Misiak, and J. Urban, *Nucl. Phys.* **B574**, 291 (2000).
- [54] T. Huber, E. Lunghi, M. Misiak, and D. Wyler, *Nucl. Phys.* **B740**, 105 (2006).
- [55] A. Crivellin, L. Hofer, J. Matias, U. Nierste, S. Pokorski, and J. Rosiek, *Phys. Rev. D* **92**, 054013 (2015).
- [56] A. J. Buras, D. Buttazzo, and R. Knegjens, *J. High Energy Phys.* **11** (2015) 166.
- [57] M. Blanke, A. J. Buras, B. Duling, K. Gemmler, and S. Gori, *J. High Energy Phys.* **03** (2009) 108.
- [58] G. Isidori and R. Unterdorfer, *J. High Energy Phys.* **01** (2004) 009.
- [59] Y. Miyazaki *et al.* (Belle Collaboration), *Phys. Lett. B* **699**, 251 (2011).
- [60] J. P. Lees *et al.* (BABAR Collaboration), *Phys. Rev. Lett.* **104**, 151802 (2010).
- [61] A. J. Buras, T. Ewerth, S. Jager, and J. Rosiek, *Nucl. Phys.* **B714**, 103 (2005).
- [62] A. J. Buras, D. Buttazzo, J. Girrbach-Noe, and R. Knegjens, *J. High Energy Phys.* **11** (2015) 033.
- [63] A. J. Buras, J. Girrbach-Noe, C. Niehoff, and D. M. Straub, *J. High Energy Phys.* **02** (2015) 184.
- [64] J. Grygier *et al.* (Belle Collaboration), *Phys. Rev. D* **96**, 091101 (2017); **97**, 099902(A) (2018).
- [65] T. Abe *et al.* (Belle-II Collaboration), arXiv:1011.0352.
- [66] C. Patrignani *et al.* (Particle Data Group), *Chin. Phys. C* **40**, 100001 (2016).
- [67] In this Letter, we followed two approaches to check the results. First, we calculated the results in unitary gauge. Then, we derived the couplings of the LQ Goldstones to SM fermions by requiring the tree-level amplitude to be gauge independent. Finally, we calculated its contribution in R_ξ gauge.
- [68] V. A. Smirnov, *Mod. Phys. Lett. A* **10**, 1485 (1995).
- [69] F. Feruglio, P. Paradisi, and A. Pattori, *J. High Energy Phys.* **09** (2017) 061.
- [70] S. Descotes-Genon, L. Hofer, J. Matias, and J. Virto, *J. High Energy Phys.* **06** (2016) 092.
- [71] F. M. Borzumati and C. Greub, *Phys. Rev. D* **58**, 074004 (1998).
- [72] F. Borzumati, C. Greub, T. Hurth, and D. Wyler, *Phys. Rev. D* **62**, 075005 (2000).
- [73] W. Altmannshofer and D. M. Straub, *Eur. Phys. J. C* **75**, 382 (2015).
- [74] See Ref. [75] for a recent analysis of such scenarios.
- [75] M. Algueró, B. Capdevila, S. Descotes-Genon, P. Masjuan, and J. Matias, arXiv:1809.08447.
- [76] V. Cirigliano, A. Crivellin, and M. Hoferichter, *Phys. Rev. Lett.* **120**, 141803 (2018).