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In(Ga)As self-assembled quantum ring formation by molecular beam epitaxy

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The effect of growth conditions on the morphological properties of InAs/GaAs(001) quantum dots covered by a thin (<3 nm) GaAs cap has been studied by atomic force microscopy. Each dot turns into an elongated nanostructure at 540 °C upon deposition of the cap in As₄ atmosphere, while structures with two humps are obtained when capping at 500 °C. The use of As₂ atmosphere instead of As₄ at 500 °C leads to the formation of quantum rings. Photoluminescence spectroscopy and polarization photoluminescence (PL) at 15 K show dramatic changes due to the different kinds of confinement. This allows the possibility of tailoring PL emission by controlling the size and shape. © 2003 American Institute of Physics. [DOI: 10.1063/1.1566799]

There is increased interest nowadays in methods to control the size and shape of quantum dots (QDs) to tune QD optical properties. In particular, morphological changes of QD grown by molecular beam epitaxy (MBE) due to a thin GaAs cap have been reported. This approach has turned into a powerful technique by which to obtain self-assembled nanostructures such as quantum rings (QRs). Although the first QRs were grown some years ago, only a few groups have reported similar findings. In some cases fee growns used, and in others in which MBE was used, the ring-like structures were related to InAs "giant" QDs, probably associated with dislocated islands.

In this work we present results of QDs covered with a 2 nm GaAs cap layer. We show that the final morphology and optical properties depend strongly on the starting shape of the QD and on detailed growth conditions such as the substrate temperature for the capping process ($T_{\rm cap}$) and the molecular species of the As flux.

The final sample morphology was characterized by *ex situ* contact mode atomic force microscopy (AFM). Their emission characteristics are measured by photoluminescence (PL) spectroscopy and polarization photoluminescence (PPL) with a 514.5 nm Ar laser line and a 0.22 m monochromator with a liquid nitrogen cooled Ge diode. For PPL measurements the $\langle 110 \rangle$ direction of the sample is oriented 45° with respect to the direction of laser polarization. A Glan–Thomson polarizer is placed at the entrance of the monochromator and is used to measure PL intensity along $\langle 110 \rangle$ and $\langle 1\bar{1}0 \rangle$ directions (PL₁₁₀,PL₁₁₀). The polarization degree (PD) is defined as PD=(PL₁₁₀-PL₁₁₀)/(PL₁₁₀+PL₁₁₀).

The substrate temperature (T) is carefully calibrated by observation of the reflection high energy electron diffraction (RHEED) diffraction pattern, the oxide desorption (\sim 630 °C) and surface reconstruction phase transition from $c(4\times4)$ to a clear 2×4. This transition occurs at T=520 °C and at an As₄ beam equivalent pressure (BEP) of 3.5 \times 10⁻⁶ mbar 30 s after shutting down the As source, and is observed by RHEED. The QDs are grown by depositing 1.7

monolayers (MLs) of InAs at 540 °C, using $\sim 3-4\times 10^{-6}$ mbar BEP-As₂ pressure (a incipient 2×4 RHEED pattern reconstruction is observed). InAs deposition takes place in a growth sequence of 0.1 ML InAs (at ~ 0.06 ML/s) plus a 2 s pause under As flux. The QDs are next annealed 1 min to enhance the size distribution and to obtain low density ensembles (10^8-10^9 cm⁻²), which is required for improved AFM analysis. Subsequently, a GaAs thin cap layer is grown (1 ML/s) at $T_{\rm cap}$ and annealed for 1 min. The samples for AFM characterization are cooled down immediately and removed from the chamber. The samples for PL measurements are obtained by burying these transformed nanostructures under a thick GaAs layer where the first 20 nm of GaAs is grown at $T_{\rm cap}$ and then T is increased to 595 °C.

Figure 1(a) shows an AFM image of QDs with a full width at half maximum (FWHM) vertical size distribution of 9% centered at \sim 10 nm. When a 2 nm cap is deposited at 540 °C under $4-5\times10^{-6}$ mbar As₄ BEP each QD is trans-

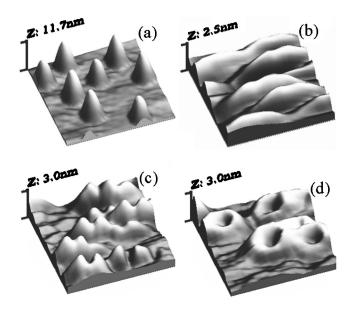


FIG. 1. 250 nm \times 250 nm AFM images of (a) a QD grown at 540 °C, (b) a QD after 2 nm capping at 540 °C, (c) a QD after 2 nm capping at 500 °C and (d) a QD after 2 nm capping at 500 °C with As₂.

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formed into an elongated dash-like nanostructure (160 $nm \times 40 \text{ nm}$) $\sim 2 \text{ nm high [Fig. 1(b)]}$. Cap deposition at lower temperatures, $T_{\rm cap}$ =500 °C, produces two camel hump-like nanostructures (100 nm×50 nm) ~2.5 nm high [Fig. 1(c)]. Figure 1(d) shows quantum rings (100 nm \times 90 nm) \sim 1.5 nm high, which formed at $T_{\rm cap}$ =500 °C under As₂ flux. It is worth mentioning that we never obtained QRs for samples with initially low (<7 nm) QD height.

These morphological changes allow a better understanding of the quantum ring formation mechanism. Some authors have suggested^{4,7,8} that a change in the balance of surface free energy due to the thin capping layer creates an outward pointing force. This force brings about material redistribution and results in a ring shaped structure. Although this would imply a liquid InAs QD, there are recent experimental⁹ and theoretical results¹⁰ showing that biaxial epitaxial stress of InAs on GaAs(001) leads to a mixture of stress-free matter including phases of "floating" liquid In and InAs. Moreover, this effect has been observed 11 independent of T in the 170-520 °C range. The driving force for this melting is the elastic energy,⁹ it is expected that when a relaxed QD is partially covered, and consequently compressed, it undergoes a phase transition into liquid phase, and allows a de-wetting process that expels the InAs from the QD. We have to consider also the simultaneous Ga-In alloying process that takes place. The presence of liquid InAs is basically a T independent process, 11 whereas Ga-In alloying is strongly T dependent. 12 The competition between these two processes at different $T_{\rm can}$ can explain our experimental observations. Although liquid InAs must have very high mobility, at large temperatures ($T_{\rm cap}$ =540 °C) the strong In-Ga alloying processes result in the formation of an immobile 13 InGaAs phase at the base of the dot, resulting in a dash-like structure [Fig. 1(b)]. For lower cap temperatures ($T_{\text{cap}}=500\,^{\circ}\text{C}$), the migration of Ga atoms into the center region is reduced and dewetting of the liquid center InAs region leads to the formation of a depleted region surrounded by immobile InGaAs. This causes the camel hump-like shape structures.

The observation of quantum ring formation under As₂ is related to enhancement of the reactivity along the (110) direction and it is discussed next. During GaAs homoepitaxy, it is well known that the migration distance of both principal directions is different due to the anisotropic chemical reactivity of the steps. 14 Arsenic terminated B steps (perpendicular to the $\langle 1\overline{10}\rangle$ direction) provide highly active sites for gallium adsorption whereas gallium terminated A steps are less reactive to group III migrating atoms. The structures elongated along the (110) direction usually found when growing by MBE and step flow growth mode in B type vicinal surfaces are attributed to this chemical reactivity anisotropy. The step flow of A-type facets can be achieved for higher As pressures.¹⁴ If we assume similar chemical behavior of the original InAs island steps, an increase of As pressure will produce enhancement of the reactivity of A steps to Ga atoms. The use of As₂ then has similar effects to employing a larger As₄ flux.¹⁵

Figure 2 (squares) shows PL emission from QDs at 15 K. The interband transitions between the ground and excited first states of electrons and holes are observed at low power excitation (\sim 20 W cm $^{-2}$). The PL spectrum of dash-like the top and around the crater. Downloaded 14 May 2010 to 161.111.235.252. Redistribution subject to AIP license or copyright; see http://apl.aip.org/apl/copyright.jsp

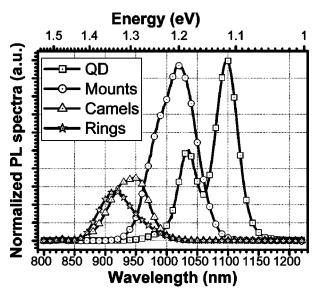


FIG. 2. Low temperature PL spectra at 20 W/cm² of samples with the nanostructures shown in Fig. 1 buried with GaAs. Strong blueshifting of the wavelength is observed from the dot to ring emission.

nanostructures (Fig. 2, circles) is blueshifted ~85 meV with respect to the QD emission, in agreement with the AFM measured reduction in vertical size of the original QDs. Although the camel hump-like nanostructures are slightly taller (0.5 nm) than the dash-like ones, the PL emission (Fig. 2, triangles) shows a larger blueshift than that for the former ones (178 meV). The PL emission of the QR (Fig. 2, stars) shows the largest blueshift (220 meV) with respect to the dots. This larger blueshift for camel hump-like and QR nanostructures suggests a reduction of confinement, probably associated with the formation of the center depleted region.

The polarization degree for large pumping power of a QD sample is 9% as shown in Fig. 3(a). Similar spectra have been observed by Noda et al., 16 who also showed theoretically that the polarization properties of the transitions of slightly asymmetrical QDs can be assigned to transitions between excited states of electrons and holes with the same quantum number. The PPL signal is positive (stronger PL along (110)) due to overlapping of the emission modes and the predominance of lower energy states due to band-filling effects. 16 For the QR sample, and in spite of the more asymmetrical shape of the rings [Fig. 3(b)], its PPL spectrum is 6%, somewhat smaller than the dot polarization degree. From 950 to 1025 nm the PL signal along (110) is more intense than along (110). Accordingly the polarization degree exhibits a negative character in this wavelength of emission, contrary to what is observed in the QD measurements. This suggests different spatial distributions of the wave function envelopes. The PPL signal of camel hump-like and dashed structures shows no clear trend and is currently under investigation. It has been recently shown¹⁷ that a giant vertical permanent dipole moment opposite in sign to the observed one in QDs is present in similar quantum ring nanostructures. This change in sign may be associated with the distribution of InAs vertically inverted to the one in the QD. Whereas QDs have a high In(Ga)As content at the base, the quantum rings seem to present a high In(Ga)As content on

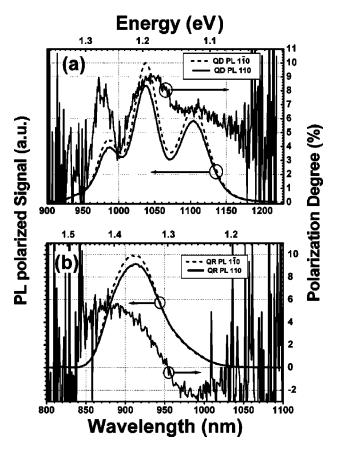


FIG. 3. Polarization PL emission at 20 K using 1000 W/cm² of a sample with (a) buried dots and (b) buried rings.

Several conclusions as to the transformation mechanisms involved during self-assembled nanostructure thin GaAs capping can be suggested. The InAs within the QD is destabilized by the deposition of GaAs, thus driving material redistribution. A high migration rate of liquid phases In and InAs enhances mass transport out toward the center region of the QD, while Ga–In alloying effects reduce the mobility and tend to prevent the formation of a crater at high substrate temperatures. Quantum ring formation is achieved at low

substrate temperatures. The use of As_2 is required for successful isotropic distribution of the InGaAs forming the rim. Large QDs of more than 7 nm in height are necessary to obtain a depleted center region. Understanding the mechanisms by which these nanostructures grow and change shape allows tailoring of their optical properties such as PL wavelength emission and polarization response.

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- ¹ J. M. García, G. Medeiros-Ribeiro, K. Schmidt, T. Ngo, and P. M. Petroff, Appl. Phys. Lett. **71**, 2014 (1997); R. Songmuang, S. Kiravittaya, and O. G. Schmidt, J. Cryst. Growth **249**, 416 (2003).
- ² A. Lorke, R. J. Luyken, A. O. Govorov, J. Kotthaus, J. M. García, and P. M. Petroff, Phys. Rev. Lett. 84, 2223 (2000).
- ³R. J. Warburton, C. Schaffein, D. Haft, F. Bickel, A. Lorke, K. Karrai, J. M. García, W. Schoenfeld, and P. M. Petroff, Nature (London) 405, 926 (2000).
- ⁴T. Raz, D. Ritter, and G. Bahir, Appl. Phys. Lett. **82**, 1706 (2003).
- ⁵J. S. Lee, H. W. Ren, S. Sugou, and Y. Masumoto, J. Appl. Phys. **84**, 6686 (1998).
- ⁶I. Kamiya, I. Tanaka, and I. Sakaki, J. Cryst. Growth **201**, 1146 (1999).
- ⁷ A. Lorke, R. Blossey, J. M. García, M. Bichler, and G. Abstreiter, Mater. Sci. Eng., B 88, 225 (2002).
- ⁸R. Blossey and A. Lorke, Phys. Rev. E 65, 021603 (2002).
- ⁹J. M. García, J. Silveira, and F. Briones, Appl. Phys. Lett. **77**, 409 (2000).
- ¹⁰D. Bottomley, Appl. Phys. Lett. **80**, 4747 (2002).
- ¹¹ J. Silveira, J. M. García, and F. Briones, J. Cryst. Growth **227–228**, 995 (2001).
- ¹² P. B. Joyce, T. J. Krzyzewski, G. R. Bell, B. A. Joyce, and T. S. Jones, Phys. Rev. B 58, R15981 (1998).
- ¹³ A. Lorke, R. J. Luyken, J. M. García, and P. M. Petroff, Jpn. J. Appl. Phys., Part 1 40, 1857 (2001).
- ¹⁴Y. Horikoshi, H. Yamaguchi, F. Briones, and M. Kawashima, J. Cryst. Growth 105, 326 (1990).
- ¹⁵T. Ogura, D. Kishimoto, and T. Nishinaga, J. Cryst. Growth **226**, 179 (2001).
- ¹⁶S. Noda, T. Abe, and M. Tamura, Phys. Lett. B **58**, 7181 (1998).
- ¹⁷ R. J. Warburton, C. Schulhauser, D. Haft, C. Schaflein, K. Karrai, J. M. García, W. Schoenfeld, and P. M. Petroff, Phys. Rev. B 65, 113303 (2002).