# Interaction Between Colloidal Quantum Dots and Halide Perovskites: Looking for Constructive Synergies

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#### **Abstract**

Colloidal Quantum Dots (QDs) have received an extensive attention during the last few decades due to their amazing properties emerging from the quantum confinement. In parallel, halide perovskites have attracted attention by the demonstration of very high performances especially in solar cell and LEDs, and other optoelectronic devices. Both families of materials can be prepared in a relatively simple way, facilitating their integration. There are several examples of their interaction enhancing the properties of the final nanocomposite. Perovskite can effectively passivate QDs, or act as efficient charge transporters. While QDs can be used to modify the selective contacts in perovskite devices or can be used as efficient light emitters or absorbers for enhanced LEDs and Photodetectors, respectively. Moreover, QD can seed the perovskite crystal growth, improving the morphology and ultimately the solar cell performance. In addition, new advanced devices can appear by the constructive synergy between both families of materials.

# TOC Quantum Dots Perovskite Synergy 1 + 1 = 3

#### **Quotes**

- Combinations of materials with different nature has been a successful strategy in order to develop new materials with enhanced properties.
- The combination of halide perovskites and colloidal quantum is a promising strategy looking for constructive synergies where the whole is more than the sum of the parts.
- Combination of halide perovskites and colloidal quantum have improved the performance of solar cells, LEDs and photodetectors.
- Advanced configurations as tunable LEDs and Intermediate Band Gap Solar cells could be potentially implemented if non-radiative recombination can be avoided.

Combinations of materials with different nature has been a successful strategy in order to develop new materials with enhanced properties. Adobe, concrete, stained glass, stainless steel are just some examples of how this strategy has provided some of the most employed materials that has been used for the humankind for centuries even for millennia. This strategy is still useful nowadays when the development of chemistry, quantum mechanics and nanotechnology has created a revolution of material science. Two paradigmatic relatively new materials with a clear applicability in optoelectronics are halide perovskites (HPs) and semiconductor quantum dots (QDs). Both families of materials are exciting by themselves and they are somewhat easy to combine, consequently there is a huge interest in building up constructive synergies in order to obtain properties and/or devices with enhanced functionalities and performance by their combined use.

Halide Perovskite (HP) is a family of materials with outstanding versatility not only for the preparation of high efficiency solar cells, with certified efficiencies higher than 23%, but also photodetectors, light emitting diodes (LEDs), 4-6 light amplifier and lasers. Although HPs has been studied in the past, the interest for these materials has reached its maximum level in the last few years. This interest has been mainly originated by the impressive photovoltaic conversion efficiencies reported in perovskite solar cells (PSCs) and how this efficiency has increased quickly from 10% to 23% in less than six years. In PSCs it has been shown that HP not only acts as excellent light harvester but also as Electron Transporting Material (ETM), Hole Transporting (HTM), or both in planar devices with no mesoporous scaffold. Perovskite is able to accumulate charge by itself, constituting a new class of photovoltaic device differentiated from its ancestor, the dye sensitized solar cells. Additionally HPs can be prepared from solution methods at low temperatures and, consequently, using low cost fabrication techniques. Solution processes also allows HPs to combine easily with other materials such as organic compounds and QDs.

On the other hand, colloidal QDs offer a high versatility with tunable bandgap just by size or shape control, due to quantum confinement effect, pure color and bright emission, and generally better stability than organic chromophores.<sup>5, 27-28</sup> Consequently, QDs have attracted considerable attention due to their unique and suitable electro-optical properties for the application of QDs in LEDs, <sup>27, 29-31</sup> solar cells <sup>32-38</sup> and other optoelectronic devices. <sup>27</sup> Besides, solar cells prepared with QD thin films have demonstrated certified photovoltaic conversion efficiencies as high as 13.4 %, <sup>39</sup> while for sensitized configuration efficiencies close to 13% have been reached. <sup>40</sup>

Considering all these outstanding properties for HPs and colloidal QDs the combination of both systems is a promising strategy looking for constructive synergies where the whole is more than the sum of the parts. There are previous examples of constructive combination of different family materials as molecular dyes and colloidal QDs that can provide enhanced charge photoinjection into TiO<sub>2</sub>, <sup>41</sup> and consequently increase the performance of the sensitized solar cells. <sup>42</sup> In this case, the obtained photocurrent is not just the consequence of the addition of

materials absorbing at different light wavelengths but higher. The internal recombination inside the QDs of the photogenerated carriers before being injected is highly suppressed due to a fast hole extraction towards the attached dye.<sup>42</sup> This is a clear example of a constructive synergy.

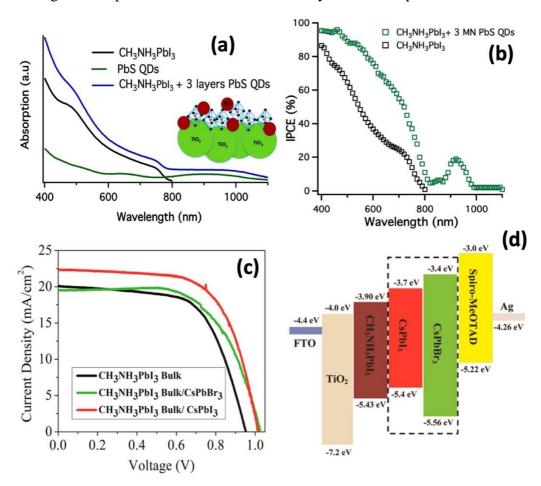
As both HPs and QDs are promising for optoelectronic devices, their combination could benefit of the advantages of these two materials such as the strong luminescent efficiency of QDs and long range charge carrier transport of HPs. Here we summarize important studies on the interaction between HPs and QDs which has been applied for the development of optoelectronic devices such as solar cells<sup>43-46</sup>, LEDs<sup>25-26, 47-48</sup> and photodetectors,<sup>49-51</sup> and discuss the perspectives of the synergistic interaction between them.

HPs and QDs can be fabricated using relatively simple synthetic routes allowing an easy integration of both materials and consequently favored the apparition of different works which explored the integration of both materials after the breakthrough of the first reports on all-solid PSCs. Different QDs have been investigated in order to combine with HPs. The PbS QDs have been likely the most employed ones due to their bandgap that can be tuned deeper in the IR than HPs and also due to the relatively low lattice mismatch.<sup>24</sup> In fact the lattice mismatch can be further reduced by the use of mixed halide I-Br perovskite.,<sup>48</sup> or the use of inorganic perovskites.<sup>52</sup> Researchers have combined HPs with colloidal QDs following different objectives as the use of the selective and transport properties of one of the compounds, attain a broader light absorption range, the passivation of QD surface or the implementation of advanced solar cell or LED configurations.

QDs have been used as HTM in PSCs. Having a proper band energy alignment is a key parameter to achieve a good performance of optoelectronic devices in general and solar cells in particular. Hu et al. synthesized PbS QDs with tuned bandgap from 0.7 to 2.1 eV and used them as HTM in PSCs with a structure of ITO/PbS QDs/MAPbI<sub>3</sub>/PCBM/Al.<sup>43</sup> By tuning the size of PbS colloidal QDs, they engineered the energy alignment between PbS QD acting as HTM and HP active layer. The use of PbS QDs HTM allowed to harvest photons up to 1000 nm wavelength and increased significantly the performance respect cells with no HTM but far of the efficiency obtained with the most common HTMs.<sup>43</sup>

Co-sensitization of HPs and QDs is one of the possibilities that have awaken more interest regarding the combination of the two material families. In this combination the purpose is to use the broader absorption range of the PbS QDs in the IR with the good transport properties of HPs. The combination of HPs and QDs has been used as a co-sensitizer in a heterojunction solar cells. CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> (MAPbI<sub>3</sub>) has an excellent absorption of the solar spectrum in the visible range but not in the near infrared range because of its bandgap (1.57 eV). The use of cosensitizer HP and QDs can overcome this absorption limitation as suggested the first time by Etgar et al.<sup>44</sup> QDs act as an absorber in NIR region, see Figure 1a, resulting a higher short circuit current density,  $J_{sc}$ , than for single MAPbI<sub>3</sub> devices, <sup>44, 53-54</sup> with contribution to the photocurrent of photons harvested by the two co-absorbers as it is observed from IPCE measurements, see Figure 1b. On the

other hand, the study of Vinh et al. pointed out a synergistic effect of the addition of PbS QDs as it induced hydrophobic modification of the perovskite surface, leading to an improvement of the device stability in air atmosphere. <sup>54</sup>



**Figure 1:** (a) Absorption spectra on mesoporous TiO<sub>2</sub> substrate of the PbS QDs, lead iodide perovskite and lead iodide perovskite with 3 layers of PbS QDs. (b) Incident photon to current efficiency (IPCE) spectra of the PbS QDs, lead iodide perovskite and lead iodide perovskite with 3 layers of PbS QDs. Reproduced with permission of the Royal Society of Chemistry.<sup>44</sup> (c) J–V curves spectra of perovskite solar cells incorporating CsPbX<sub>3</sub> QDs between the MAPbI<sub>3</sub> active layer and the HTM (d) Energy band diagram of each material within the perovskite/CsPbX<sub>3</sub> structure, with energy levels given in eV. Reproduced with permission of the American Chemical Society.<sup>55</sup>

Interface Engineering in PSCs using perovskite QD layer sandwiched between the perovskite active layer and the HTM has produced a performance enhancement. <sup>55-57</sup> Cha et al. reported an increase of PSC photoconversion efficiency by the deposition of MAPbBr<sub>0.9</sub>I<sub>2.1</sub> between MAPbI<sub>3</sub> and spiro-OMeTAD HTM Quantum Dots. <sup>57</sup> Moreover, inorganic perovskite QDs present higher stability as Sidhik et al. pointed out. <sup>55</sup> The thin layer of HP (CsPbX<sub>3</sub>, X is Br or I) QDs deposited on top a MAPbI<sub>3</sub> layer played acted as the electron blocking of

the perovskite/HTM interface, reducing interfacial recombination and increasing consequently the open circuit potential, see Figure 1c. With an appropriated band alignment, using CsPbI<sub>3</sub> QDs (see Figure 1d), the hole injection is not hindered and the photocurrent can be also enhanced, see Figure 1c.<sup>55</sup> There are some other related studies regarding the interaction of perovskite QD with perovskite thin layers that have been recently summarized in a perspective article by Yang et al.<sup>56</sup>

On the other hand, one of the key issues on the efficiency improvement experienced in the last decade in colloidal QD solar cells, both on depleted heterojunction<sup>58</sup> and in sensitized configuration,<sup>40</sup> has been the passivation of surface trap states of colloidal QDs. HPs also exhibit good performance as passivation agents of QDs surface. Typically, colloidal QDs are synthesized with long-chain organic ligands which stabilize, control size and afford an excellent monodispersity. However, the bulky organic ligands create electronic barriers that compromise efficient charge transport of QD films in optoelectronic devices. To build a high-quality QD-based photovoltaic device, interdot electronic communication needs to be improved. On the other hand, organic ligands also act as surface passivators. QDs possess an inherent large surface-to-volume ratio, resulting in unsaturated dangling bonds, creating undesired electronic trap states within the bandgap of QDs. These trap states increase the chance of carrier recombination, curtailing the efficiency of charge extraction. The ligand exchange procedure opens the possibility to exchange to short ligand, improving the interdot electronic communication, and moreover passivate the trap states. Among various ligand exchange schemes, HPs, including HP precursors, shows a strong passivation and well-protection from oxygen, resulting an improvement in a device performance and stability.

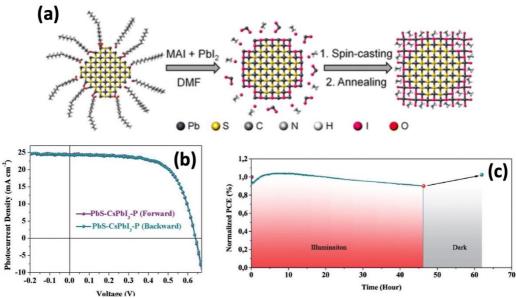
There are several examples of exchange of bulky organic capping layer of colloidal QDs by HPs or HPs precursor shell. In 2014, Dirin et al. presented a broad study of the use of HPs and other metal halide complexes as inorganic capping ligands for colloidal nanocrystals (PbS, CdSe, CdS, CdSe/CdS, InP, Au and Pd), showing that MAPbI<sub>3</sub> was the first example of hybrid capping of PbS able to retain highly efficient IR photoluminescence quantum yield (PLQY) of PbS QDs. In the same year, Ning et al. carried out the solution-phase ligand exchange from the oleic acid capping of PbS QDs to methylammonium iodide (MAI). Short-ligand exchanged PbS QDs could be directly deposited in a single step, with no requirement of solid-state exchange, rendering it compatible with roll-to-roll processes such as screen printing, spray coating, and blade coating. Moreover MAI capped PbS QD films provided a good stability at 100°C annealing temperature. These QD films with MAI short-capping were applied to fabricate QD solar cells, reaching a photovoltaic conversion efficiency of 6.10%.

Next year, taking advantage of the similar crystal structure and lattice match between PbS and MAPbI<sub>3</sub> Yang et al. used MAPbI<sub>3</sub> to exchange the ligand of PbS QDs, see the Figure 2a. High resolution transmission electron microscopy (HRTEM) measurements and the analysis of lattice fringes from scanning tunneling electron microscope (STEM) pointed out the core-shell structure of PbS/MAPbI<sub>3</sub> QDs after ligand exchange. An annealing is necessary to transform HP precursors

into thin crystalline matrix to archive good solar cell efficiency. By introducing the coherently lattice-matching thin shell of MAPbI<sub>3</sub>, QDs were passivated, limiting the aggregation and thus achieving a better balance of open circuit potential,  $V_{oc}$ , and  $J_{SC}$ , yielding improved device performance, reaching 8.95% of photoconversion efficiency, using the depleted heterojunction quantum dot solar cell configuration. <sup>46</sup> The main mechanism in the ligand exchange from organic QD capping to MAI or MAPbI<sub>3</sub> shell is based on the binding of iodine anions to lead atoms on the QD surface. 60 However, agglomeration could be produced. With a proper optimization of the ligand exchange process the inclusion of iodine on OD surface was done without the detrimental effect of fusion.<sup>61</sup> A co-solvent system was employed to tune the solvent polarity in order to achieve the solvation of MAI and the dispersion of hydrophobic PbS QDs simultaneously in a homogeneous phase enhancing the performance of QD solar cells to 10.6%. This strategy has been further improved using lead halide passivators reaching a solar cell efficiency of 11.28%. 62 Moreover the passivation of lead halide also improved the stability of devices, 90% of initial efficiency was maintained after 1000 hours stored in ambient conditions.<sup>62</sup>

MAPbI<sub>3</sub> shell showed an excellent passivation of PbS QDs however QD solar cells strongly suffer from hysteresis in photocurrent-voltage curves. 46 An interesting alternative to hybrid halide perovskite is the use of inorganic halide perovskites as CsPbI<sub>3</sub>, that exhibits higher stability. <sup>63</sup> Although, perovskite α-phase is not stable for bulk CsPbI<sub>3</sub>, it can be stabilized in CsPbI<sub>3</sub> QDs.<sup>64</sup> On the other hand, it also possesses similar crystal structure and low lattice mismatch (less than 5%) comparing with those of PbS.<sup>52</sup> Very recently, Zhang et al. used CsPbI<sub>3</sub> passivated PbS ODs. 65 As in the case of MAPbI<sub>3</sub>, after ligand exchange from organic capping to CsPbI<sub>3</sub>, a PbS/CsPbI<sub>3</sub> core-shell QD structure also formed. CsPbI<sub>3</sub> shell passivates the PbS QD surface, diminishing the trap states, hence improving the charge collection and reducing the charge recombination rate, allowing to obtain performances as high as 10.5%. In addition, colloidal OD solar cells based on PbS/CsPbI<sub>3</sub> QDs present extremely low hysteresis, see Figure 2b, not only in the different scanning directions but also in different scan sweep steps. Moreover these devices also show a good stability not only under long illumination, 46 hours at continuously one sun illumination, see Figure 2c, but also under ambient conditions (42 days).<sup>65</sup>

Perovskite passivation of PbS QDs have been also used successfully in QD sensitized solar cell configuration. <sup>45, 66</sup> MAPbI<sub>3</sub> layer coated on PbS QDs in sensitized solar cells, passivates QD surface, reducing significantly the defects as the PbSO<sub>3</sub> oxidation, resulting finally in a better charge extraction with an improvement of the performance of the photovoltaic device.



**Figure 2:** (a) Schematic of solution ligand exchange with MAPbI<sub>3</sub> perovskite precursors and corresponding colloidal QD film preparation. Reproduced with permission of the American Chemical Society. (b) Forward and backward *J–V* curves under 1 sun AM1.5G illumination of PbS/CsPbI<sub>3</sub> based colloidal QD solar cell. (c) Photostability test of the PbS/CsPbI<sub>3</sub> based QD solar cell under illumination and after stored during 16 h in ambient conditions under dark after the illumination test. The photostability test was performed in N<sub>2</sub> atmosphere under one sun equivalent illumination provided by a white LED lamp. Reproduced with permission of Wiley. (65)

Nevertheless, beyond of the use of halide perovskite as mere passivation agent of colloidal QDs, likely the most appealing results have been obtained when a more synergistic interaction was stablished between both materials. In 2015, Ning et al. reported by first time the preparation of thin films with perovskite matrix and embedded colloidal QDs.<sup>24</sup> PbS and MAPbI<sub>3</sub> possess a related crystal structure, each having a six-coordinated Pb atom, and with Pb-Pb distances (MAPbI<sub>3</sub>, 6.26A°; PbS, 5.97A°, ~ 5% lattice mismatch), see Figure 3a-b. Their analysis of density functional theory (DFT) of an interface formation between PbS (100) and MAPbI<sub>3</sub> (110) plane showed an interfacial energy less than 10 meV·Å<sup>-2</sup>, suggesting that growth of HP on PbS at room temperature is nearly as feasible as homoepitaxy of PbS on PbS or HP on HP. DFT further reveals that the epitaxial threedimensional embedding of PbS CQDs inside a HP matrix is achieved without the formation of interfacial defects; that is, the bandgap is predicted to remain open, with no in-gap defects predicted. They concluded structurally PbS QDs lattice has a relative good match with MAPbI<sub>3</sub> HP structure both three dimensionally and two dimensionally.<sup>24</sup> Perovskite matrix with embedded QDs increased significantly the PL emission from PbS QDs as photogenerated electrons and holes in HP were transferred to PbS QDs with a value exceeding 80% of efficiency. Since this first demonstration, HP films with embedded QDs were applied as an active layer in different optoelectronic devices such as solar cells<sup>67-68</sup>, LEDs<sup>26, 48</sup> and photodetectors. 49-51 In these works, embedding QDs inside HPs matrix showed an improvement respect to the reference devices, without embedded QDs.

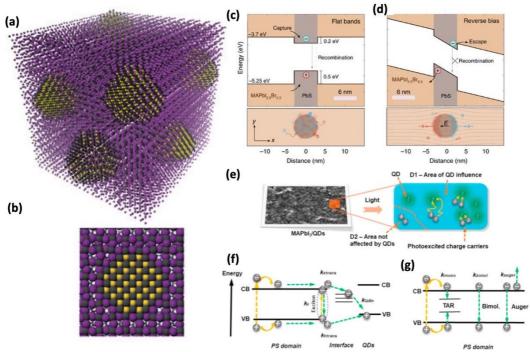


Figure 3: (a) Three-dimensional atomistic model and (b) Cross-section (two-dimensional view) of PbS QDs embedded in a halide perovskite matrix. (a-b) Reproduced with permission of Nature.<sup>24</sup> (c, d) Band alignment of PbS QDs embedded in mixed halide perovskite matrix as calculated with a SCAPS model at (c) no bias and (d) reverse bias conditions. In the absence of an imposed electric field photogenerated charges are uniformly distributed across the dot (bottom panels). At sufficient reverse bias conditions they are spatially separated and can inject into the MAPbI<sub>2.5</sub>Br<sub>0.5</sub> host, favoring carrier recirculation over capturing. (c-d) Reproduced with permission of Nature.<sup>50</sup> (e) SEM of MAPbI<sub>3</sub> layer with embedded PbS/CdS core/shell ODs and the cartoon showing the areas of QDs influence (D1) and without their influence (D2) upon excitation over the band gap of perovskite. (f) Energy diagram showing the dynamics of photoexcited charge carriers in MAPbI<sub>3</sub>-ODs film in D1 area and (g) D2 area upon excitation over the band gap. For clarity, (f) only shows the processes involving the influence of QDs, while those in (g) representing the classical events in pure perovskite. ketrans and khtrans characterize the effective transfer processes of charge carriers from perovskite to the interface or/and QDs. kx is the exciton recombination rate constant. kodin is the recombination rate constant between trapped electrons and holes in QDs. (e-g) Reproduced with permission from Elsevier.<sup>69</sup>

The introduction of PbS QDs into HP matrix has increased photoconversion performance of PSCs by the effect of QDs in the crystallization process of perovskite layer. Li et al. prepared solar cells with MAPbI<sub>3-x</sub>Cl<sub>x</sub> matrix an embedded PbS QDs by dispersing a small quantity of core/shell PbS/MAI QDs in the solution containing the HP precursors. PbS acted as effective seed-like sites to promote the formation of HP grains, resulting in substantial morphological improvements in terms of grain size, surface coverage and uniformity. Similar

effect in the grain size by QD addition have been also observed in other works.<sup>26</sup> The PSC performance observed by the appropriate incorporation of PbS, with concentration 1.0wt%, reached 17.4%, proximately a 25% of improvement respect the reference devices.<sup>67</sup> Recently, Han et al. attained a performance of 18.6% with the QD in perovskite strategy.<sup>68</sup>

In this sense, the maximum photoconversion efficiency with hybrid QD and HP films have been obtained with QDs acting with a passive role, as additives in order to obtain a better perovskite crystallization. This is not the case in hybrid QD-HP LEDs where the ODs are the active light emitters and HPs act as carrier transporting matrix. Gong et al. reduced the lattice mismatch between HP and PbS QDs (with exciton peak at 1391 nm) by incorporating a small quantity (17%) of bromide into MAPbI<sub>3</sub> matrix, forming a mixed halide MAPbI<sub>3-x</sub>Br<sub>x</sub>. 48 The introduction of bromide reduced on one hand the strain at the interface of HP and QDs, and on the other hand the amount of defect on PbS surface. The function of a mixed halide perovskite is to transport generated charger carriers into PbS QDs. The selected HP and PbS QDs presented a type I band alignment, see Figure 3c, and the transported carriers in the perovskite matrix are finally injected into the PbS QDs where they recombine radiatively, resulting in an increase of the LED external quantum efficiency (EQE), reaching 4.9%, two folds higher than pure QDs devices. 48 This approach has been also used by the same group but with layered 2D perovskites, incorporating longer amine molecules, butylamine, pentylamine and hexylamine. The substitution of MA<sup>+</sup> cation by these amine molecules with longer carbon backbones breaks the symmetry of the 3D structure and introduces 2D structuring layered along c-axis. The ODs exchanged were easily deposited in a single step. By using 2D HP passivated PbS QDs as an active layer, the authors could increase the operating voltages of LEDs up to 7.5 V due to better stability of 2D HPs in a comparison with 3D HPs.<sup>47</sup>

Embedding QDs inside HPs matrix was applied in photodetectors as well. Liu et al. pointed out that the incorporation of PbS QDs induced traps-assisted charge injection. <sup>49</sup> PbS QDs produced traps residing on the surface of MAPbI<sub>3</sub> thin films. Under illumination, excitons formed in HP layer rapidly would dissociate into electrons and holes. Then, electrons would be trapped at the HP-PbS QD surface which would shift downward the conduction band (CB) energy level of HP-QD nanocomposites, followed by the re-alignment of its fermi level with the work function of the Al cathode. The trapped electrons can be considered as a heavy-doping on the surface of HP-QDs layer, which dramatically narrow the depletion region formed at the interface. Thus injected holes from Al cathode can be easily tunneled through the hole blocking layers (PC61BM/BCP) at low external biases, resulting an EQE of ~4500%, a photoresponsivity of ~15,000 mA/W and a detectivity of over  $6\times10^{13}$  Jones (1 Jones=1 cm·Hz<sup>1/2</sup>·W<sup>-1</sup>) at a small bias of 2 V, and a photoresponse time of 11.5  $\mu$ s. <sup>49</sup>

Moreover, following the strategy of lattice matching developed by Gong et al.<sup>48</sup> a further study on the interaction between HP and PbS QDs under applied reverse bias condition for the development of photodetectors has been performed

by the same group.<sup>50</sup> The lattice of mixed HP was adjusted, by optimizing the ratio between iodide and bromide MAPbI<sub>2.5</sub>Br<sub>0.5</sub> to suitable to that of PbS QDs (exciton peak at 1240 nm, ~ 1 eV). The photodiode has been prepared using a structure of FTO/compact TiO<sub>2</sub>/PbS QDs embedded in HP matrix/spiro-OMeTAD/Au. Under no external electrical field, generated excitons are confined inside QDs due to the type I junction forming between HP and QDs, and lose their energy via the recombination, radiative or non-radiative, see Figure 3c. However, under the presence of a suitably large electric field, Figure 3d, the excitons photogenerated in the QDs will then separate, and the charges may be ejected from the dots, under the action of the field, into the host perovskite matrix. Escape charge carriers can be potentially re-captured into QDs, as determined by the competition between carrier thermalization and extraction time. This structure allowed the preparation of photodetectors with specific detectivities that exceed 10<sup>12</sup> Jones.<sup>50</sup>

In order to further improve the performance of optoelectronic devices developed with HP-QD systems, the understating of the deep interaction of these materials is mandatory. Recently, we have furthermore elucidated the interaction mechanism between HP and QDs. We studied the charge carrier dynamics of HP-QDs hybrid films as function of QD concentrations (C<sub>ODs</sub>) using ultrafast UV–VIS transient absorption (TA) and THz spectroscopies measurements.<sup>69</sup> By photogenerating charge carriers over the band gap of MAPbI<sub>3</sub> (~ 1.55 eV) in the HP-ODs hybrid material, we observed fast and efficient charge carrier deactivation due to their transfer to the QDs. The TA decay becomes shorter upon increasing C<sub>ODs</sub>. For samples containing C<sub>ODs</sub> lower than 0.18 wt%, the TA decays were governed by the behavior of neat HP while for higher concentrations a progressive faster decay was observed.<sup>69</sup> The observed dynamics were analyzed in terms of two different domains of the HP layer: a domain D1 where photocarriers are influenced by the presence of QDs and a domain D2 where they do not interact with QDs. The events happened in D1 were described in Figure 3f: (i) charge carriers transport and transfer from HP domain to QDs and MAPbI<sub>3</sub>/QDs interface, (ii) exciton recombination, and (iii) recombination of transferred charge carriers. In addition, electron and hole recombination processes in HP is also produced. Using a kinetic model and analyzing the TA signals, we observed that 85–89% of both electrons and holes were transferred to the HP/QDs interface trap states, and to the valence band (VB) of the ODs. Up to 7% charges formed excitons in HP domain in the close vicinity of interface. A small fraction of photoexcited charges (8–15%) were not affected by the presence of QDs, and thus, their deactivation followed the standard processes of neat HP, see Figure 3g. The changes in the values are related to the QDs concentration in the matrix. The analysis of the data shows the important role of trap states and gives clues for the further improvement of this hybrid system. The size of D1 and D2 domains depends on the QD concentration and on the carrier diffusion length in HP, consequently this method can be also used to determine the diffusion length of the HP layer as we showed.<sup>69</sup>

The interaction of HP and QDs has been studied not only for hybrid organic-inorganic HP but also with a fully inorganic HP. Jung et al. pointed out that PbS

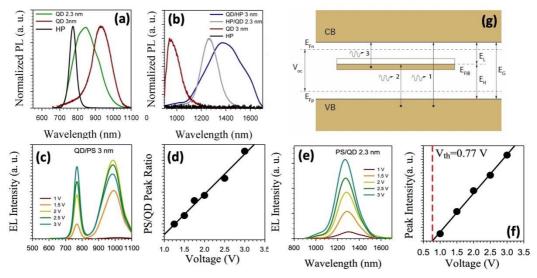
QDs and CsPbBr<sub>3</sub> present just ~ 0.5% lattice mismatch, much lower than the lattice mismatch between PbS and MAPbI<sub>3</sub>, ~ 5%.<sup>52</sup> They studied theoretically the bond formation and carrier confinement at the HP-QD interface. By calculating the surface energy at different facets of PbS and CsPbBr<sub>3</sub>, they could construct the interface slab models along (100) orientation. Using the models and calculating the interface energy taking into account different terminations, the growth mechanism of HP-QD interface is elucidated. (100) facet of CsPbBr<sub>3</sub> with a termination of PbBr<sub>2</sub> spontaneously grows on the (100) facet of PbS by locating Pb on S and Br on Pb with the presence of a strong chemical interaction attributed to heteropolar bond formation. <sup>52</sup>

Finally, combined systems of HP and colloidal QDs have also demonstrated the potentiality for the development of advanced optoelectronic devices, not by the enhancement of the properties already presented in HP or colloidal QDs but by presenting new ones. Single layers and bilayers of MAPbI<sub>3</sub> and colloidal core/shell PbS/CdS QDs were prepared, using two different QD sizes, with diameter 2.3 and 3 nm.<sup>25</sup> Figure 4a shows the normalized photoluminescence (PL) of the single HP and QD layers. In addition, bilayers formed by both combinations of QD/HP and HP/QD were also studied with a type II alignment due to the relative large bandgap of the PbS QDs in comparison with other experiments. Interestingly bilayer samples presents a new feature in the near infrared region (NIR) that it is not observed for single layer samples, HP or QD, as can be observed in Figure 4b. In this region perovskite single layer does not present any PL while for single layer colloidal QD sample just the queue of the band edge emission is observed. However bilayers presents an emission at 1200-1600 range blue shifted for QDs with smaller size, see Figure 4b. The intensity of this emission increases when samples are excited with wavelengths shorter than the perovskite bandgap and it is attributed to the formation of an exciplex state between HP and QD.<sup>25</sup> The exciplex emission has been also observed in perovskite films with embedded QDs.<sup>26</sup>

The use of different emitters produced LEDs with color tuned by the applied bias, see Figure 4c. The color can be easily controlled as the ratio between QD and HP EL intensity emission follows a linear relationship, see Figure 4d. Interestingly, the emission from HP-QD exciplex state is observed not only by optical pumping, Figure 4b, but also by electrical injection, Figure 4e, with a voltage threshold as low as 0.77 V, see Figure 4f.<sup>25</sup> The possibility of band gap tuning of the HP by the use of different cations and anions and of the colloidal QDs can potentially develop dual LEDs working with very different wavelength depending of the applied bias.

Moreover the combination of HP and colloidal QDs can have also an interesting potentiality for the development of advanced photovoltaic devices. Due to the reciprocity principle, if an electron and hole from an exciplex state can recombine producing a photon, the reciprocal process in which a photon can produce an electron-hole pair at the exciplex state is also possible. This possibility opens the door to the preparation of Intermediate Band Gap Solar Cells (IBSCs) taking benefit from the HP and QD synergies.

The concept of IBSC was developed by Luque and Martí. 70 In an IBSC there are one or several intermediate bands at the band gap of a wide gap semiconductor, been E<sub>G</sub> the band gap energy, see Figure 4g, for the case of just one intermediate band. Photons with energy higher than E<sub>G</sub> can be absorbed by the matrix semiconductor, process 1 in Figure 4g. Moreover, intermediate bands would act as stepping stones allowing low energy photons transfer electrons from the VB of the wide band gap semiconductor to its CB by the absorption of a couple of photons with energies E<sub>H</sub> and E<sub>L</sub>, processes 2 and 3 in Figure 4g. In order to take advantage of the intermediate band this level has to be isolated from the extracting contact and three quasi Fermi levels can be defined for VB, CB and intermediate band. 70-71 This fact allow that the IBSC retains the  $V_{oc}$  of the wide band gap semiconductor defined by the splitting of electrons in CB and holes in VB Fermi levels, E<sub>Fn</sub> and E<sub>Fh</sub> respectively, but with higher photocurrent due to the sub-band gap photon absorption, see Figure 4g. In an optimum situation of an IBSC, a half filled intermediate bandgap can either receive electrons from the VB of the matrix compound (process 2 in Figure 4g) and promote electrons from the intermediate band state into the CB of the matrix (process 3 in Figure 4g). Theoretical calculations predicts a maximum photoconversion efficiency of 63% under high solar concentration for an cell with an optimum IB state in contrast with a maximum of 41% predicted for single absorber by the Shockley-Queisser limit at the same conditions, 70 with a performance higher than an optimized tandem cell formed by a couple of absorbers. IBSC will have higher theoretical quantum efficiency than tandem cells at any illumination condition. A tandem cell has a maximum quantum efficiency of 0.5 in all the spectra, needing two photons to extract one electron. In the case of IBSCs the same situations accounts for photons with energy lower than E<sub>G</sub>, whereas for photons with higher energy quantum efficiency of 1 is possible.



**Figure 4:** (a) PL of single layers of MAPbI<sub>3</sub>, and core/shell PbS/CdS QDs 2.3 and 3 nm diameter size corresponding to the band gap transition. (b) PL in the NIR (900-1700 nm) showing a new exciplex transitions in HP-QD bilayers while no new peak is observed in for single HP or QD (for

QD layer just the queue of the band gap transition). (c) EL at different applied bias of a bilayer sample of HP and PbS/CdS QDs of 3 nm. (d) Ratio between the EL intensity peak of PbS/CdS QDs and HP at different applied bias. (e) EL at different applied bias in the NIR region of a bilayer sample of HP and PbS/CdS QDs of 3 nm. (f) EL intensity peak in the NIR of PbS/CdS QDs 3nm and HP at different applied bias. (g) Scheme of energy configuration of an Intermediate Band Gap Solar Cell. (a-g) Reproduced with permission of AAAS.<sup>25</sup>

However, the development of an IBSC with significant efficiency has been elusive. For example the use a monocrystalline GaAs host matrix and InAs QDs grown in the Stranski–Krastanov mode, presented severe limitations in order to produce a practical solar cell. The photocurrent increase is limited by the low QD density. The density of QDs cannot be very high as QDs produce strain in the layer, and very high strain causes the formation of defects that degrades severely the cell performance. However, the use of HP could help to solve these problems as strain is avoided as it is relaxed at the benign perovskite grain boundaries. Moreover, both HPs and QDs present a high versatility allowing tailoring the intermediate band and the band alignment between HP and colloidal QDs to produce an optimum IBSC.

In summary, the interaction between HP and colloidal QDs can produce interesting constructing synergies promoting the properties of each material and generating new ones. Each material obtains benefits from the interaction with the other, HP can increase the passivation of the QD surface, protecting them from the photocorrosion. Moreover, HP can benefit from the appropriate use of QDs in the selective contacts. HP matrix can also work as an excellent carrier transporter into ODs in a view of type I junction for the development of LEDs and photodetectors while QDs acted as effective seed-like effect to promote the proper crystallization of HP lattice structures, resulting in substantial morphological improvements. Furthermore the band alignment between HP and QDs can be tailored to favor the injection of photogenerated electrons and holes from one into the other. Hence this combination could lead to an improvement in performance of optoelectronic devices such as solar cells and LEDs. On the other hand, the interaction between HP and colloidal QDs also rises new properties, which do not exist in singles materials, as the formation of the exciplex state at lower energies than both HP and QD band gap. 25-26 This system presents a high potentiality in order to develop advanced optoelectronic devices as tunable color LED or intermediate band gap solar cells. However, non-radiative recombination has to be avoided as a major challenge of these advanced configurations in order to get a full benefit of this potentiality. Further study on the interaction between these families of materials undoubtedly will contribute to a further optimization of current devices and to explore new and exciting constructive synergies.

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The authors declare no competing financial interest.

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Thi Tuyen Ngo graduated in Physics from Vietnam National University (Vietnam) in 2008 and received her Master from University of the Basque Country (Spain) in 2011. She is currently a Ph.D. student at the Institute of Advanced Materials (INAM), Jaume I University (Spain). Her study is focusing on the interaction of perovskite and other materials such as quantum dots and organic compounds. Before incorporating in a Ph.D, she worked as assistant at Institute of Physics, Vietnam Academy Science and Technology (2008-2009) and worked in the photovoltaic cells at Cidetec (2011-2013) where she learnt about the preparation and characterization of solar cells.



**Iván Mora-Seró** is associated professor at the Institut of Advanced Materials (INAM) at Universitat Jaume I de Castelló. His research has been focused on crystal growth, nanostructured devices, transport and recombination properties, photocatalysis, electrical characterization of photovoltaic, electrochromic, and water splitting systems, making both experimental and theoretical work. Recent

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