

## KECK ECHELLETTE SPECTROGRAPH AND IMAGER OBSERVATIONS OF METAL-POOR DAMPED Ly $\alpha$ SYSTEMS\*

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### ABSTRACT

We present the first results from a survey of SDSS quasars selected for strong H I damped Ly $\alpha$  (DLA) absorption with corresponding low equivalent width absorption from strong low-ion transitions (e.g., C II  $\lambda$ 1334 and Si II  $\lambda$ 1260). These metal-poor DLA candidates were selected from the SDSS fifth release quasar spectroscopic database, and comprise a large new sample for probing low-metallicity galaxies. Medium-resolution echellette spectra from the Keck Echellette Spectrograph and Imager spectrograph for an initial sample of 35 systems were obtained to explore the metal-poor tail of the DLA distribution and to investigate the nucleosynthetic patterns at these metallicities. We have estimated saturation corrections for the moderately underresolved spectra, and systems with very narrow Doppler parameters ( $b \lesssim 5 \text{ km s}^{-1}$ ) will likely have underestimated abundances. For those systems with Doppler parameters  $b > 5 \text{ km s}^{-1}$ , we have measured low-metallicity DLA gas with  $[X/H] < -2.4$  for at least one of C, O, Si, or Fe. Assuming non-saturated components, we estimate that several DLA systems have  $[X/H] < -2.8$ , including five DLA systems with both low equivalent widths and low metallicity in transitions of both C II and O I. All of the measured DLA metallicities, however, exceed or are consistent with a metallicity of at least 1/1000 of solar, regardless of the effects of saturation in our spectra. Our results indicate that the metal-poor tail of galaxies at  $z \sim 3$  drops exponentially at  $[X/H] \lesssim -3$ . If the distribution of metallicity is Gaussian, the probability of identifying interstellar medium gas with lower abundance is extremely small, and our results suggest that DLA systems with  $[X/H] < -4.0$  are extremely rare, and could comprise only  $8 \times 10^{-7}$  of DLA systems. The relative abundances of species within these low-metallicity DLA systems are compared with stellar nucleosynthesis models, and are consistent with stars having masses of  $30 M_{\odot} < M_{*} < 100 M_{\odot}$ . The observed ratio of [C/O] for values of  $[O/H] < -2.5$  exceeds values seen in moderate metallicity DLA systems, and also exceeds theoretical nucleosynthesis predictions for higher mass Population III stars. We also have observed a correlation between the column density  $N(\text{C IV})$  with [Si/H] metallicity, suggestive of a trend between mass of the DLA system and its metallicity.

*Key words:* galaxies: ISM – quasars: absorption lines – quasars: individual (SDSS0814+5029, SDSS1001+0343)

*Online-only material:* color figures, machine-readable table

### 1. INTRODUCTION

A key probe of nucleosynthesis in the early universe has been low-metallicity stars discovered in the Galaxy and its nearest neighbors. New samples of stars developed for this purpose have detected iron abundances of  $[\text{Fe}/\text{H}] < -5.0$ , and similarly low abundances of other elements (Frebel et al. 2006, 2007; Cohen et al. 2007). The abundance patterns of these metal-poor stars have been compared with stellar nucleosynthesis models, and have provided useful constraints on the initial mass function (IMF) and other properties of the Population III stars (Matteucci 2003; Meynet & Maeder 2002; Woosley & Weaver 1995). Despite the success in finding many low-metallicity stars, the abundances derived from these stars have uncertainty related to the unknown mixing and convection in the star, the possible separation of dust and gas in star formation, and enrichment of the stellar atmosphere from nearby companion stars.

The damped Ly $\alpha$  (DLA) systems have been a useful complementary probe of the early universe, and of metallic-

ity in the intergalactic medium (IGM). The H I threshold of  $N(\text{H I}) = 2 \times 10^{20} \text{ cm}^{-2}$  defines the DLAs. Above this threshold, the gas may safely be assumed to be predominantly neutral due to shielding by hydrogen, allowing the determination of metallicity from coexisting species such as C II, O I, Si II, Al II, and Fe II compared to  $N(\text{H I})$  with minimal corrections for ionization. Since early surveys of DLA systems (Wolfe et al. 1995; Sargent et al. 1989), the number of DLA systems has continued to rise into the hundreds, allowing a statistically significant sample (Prochaska et al. 2005). The larger surveys of DLAs show increasing metallicity as a function of cosmic time (Prochaska et al. 2003a, 2003b), but with significant scatter, and some evidence of dust depletion in the medium (e.g., Pettini et al. 1994). Using the largest samples of DLA systems, it is possible to probe the low-metallicity tail of DLAs. Previous works have reported a “floor” to the metallicity distribution at  $[X/H] \approx -2.7$  (Prochaska et al. 2003a, 2003b), with only a small fraction of only 10% or less of DLA systems showing  $[X/H] < -2.5$ .

Recent work has focused on the abundance patterns of these very low metallicity systems. Pettini et al. (2008) provide high-resolution spectra of a sample of three low-metallicity DLA systems at redshifts ranging from  $2.07 < z(\text{DLA}) < 2.81$ , and detect values of  $[\text{O}/\text{H}]$  in the range of  $-2.04 < [\text{O}/\text{H}] < -2.42$ , with lower metallicities in  $[\text{Fe}/\text{H}]$ , which ranges from

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$-2.55 < [\text{Fe}/\text{H}] < -2.80$ . The UVES spectra also provide useful information on the line widths of these low-metallicity DLAs and indicate that metal-poor DLAs tend to have quiescent kinematics with  $b$ -values in the range of  $3.7 \text{ km s}^{-1} < b < 25.1 \text{ km s}^{-1}$ , with an average/median  $b$ -values of  $9.24 \text{ km s}^{-1}$  and  $6.5 \text{ km s}^{-1}$ , respectively.

In order to study the lowest metallicity DLA systems, we began a survey to examine a very large sample of quasars with DLA systems of low metallicity using the Sloan Digital Sky Survey (SDSS) fifth release (DR5), which contains 77,229 newly detected quasars (Adelman-McCarthy et al. 2007; Schneider et al. 2007). The DLA sample was assembled from fits to the SDSS spectra in which the estimated H I column density satisfied  $\log(N(\text{H I})) \geq 20.3$  (Prochaska et al. 2005). From the sample of 968 DLA systems, a subset was chosen that was toward relatively bright ( $r < 20$ ) quasars, and which had DLAs with unusually weak absorption from generally strong low-ion transitions (e.g., C II  $\lambda$  1334, Si II  $\lambda$  1260, and O I 1302) based on equivalent width analysis of the SDSS spectra. Note that follow-up observations and re-analysis of the Ly $\alpha$  transition may reveal new estimates for  $N(\text{H I})$  that fall below the DLA threshold. In this manuscript, we will consider the metallicities and abundance patterns of these systems even if they are now formally referred to as sub-DLAs (or super-Lyman limit systems; SLLS).

A Keck observing program with the Echellette Spectrograph and Imager (ESI) echellette spectrograph was then conducted to provide more accurate measurements of column densities of C II, O I, Si II, and other species, to determine the actual metallicities of the DLA systems. By targeting the best newly discovered DLA systems from such a large sample of quasars, we aim to efficiently identify the lowest metallicity DLA systems and, in turn, constrain the low-metallicity distribution function of high- $z$  galaxies and enable detailed comparison with nucleosynthetic models. Our sample probes a wide redshift range,  $2.3 < z(\text{DLA}) < 4.12$ , addressing the metallicity distribution at higher redshifts than previous DLA surveys.

Models of Population III stars have gained in sophistication and predictive ability. The current generation of models (Woosley et al. 2002; Woosley & Heger 2007) can provide exact nucleosynthetic yields as a function of stellar mass, to enable comparison with stellar or DLA observations. In most of the Population III stars, a strong ‘‘odd–even effect’’ is seen in which C, O, Si, and other even atomic number elements are as much as 1 dex more abundant than odd atomic number elements such as N, Al, and Mn. The odd–even effect is well known as the result of rapid nucleosynthesis in C and O core burning stages through the triple alpha process. The end point of Population III stars differs from other stars and is thought to involve either Type II or pair–pair annihilation supernovae, depending on the mass of the progenitor stars. The yields of high-mass stars have been modeled by a number of groups, including the effects of rotation, mass loss through winds, and turbulent mixing within the stars (Chieffi & Limongi 2002; Heger & Woosley 2002; Umeda & Nomoto 2002).

The latter effects, combined with an unknown IMF, lead to a large range of predictions for the integrated yields of Population III stars. New observations of our low-metallicity DLA sample can complement the work of theoretical models, and stellar studies to help better determine the nature of both the high-mass Population III stars and star formation in the early universe.

Despite the importance of low-metallicity systems to constraining nucleosynthesis within the early universe, very few DLA systems of low metallicity have been discovered to date.

**Table 1**  
SDSS-DR5 Metal-poor DLA Candidates

QSO	$z_{\text{abs}}$	$\log N_{\text{HI}}$
J001115.23+144601	3.612	20.65
J001134.52+155137	4.317	20.50
J001328.21+135828.0	2.612	20.10
J003749.19+155208	3.479	20.35
J012211.11+150914.3	3.519	20.45
J012403.77+004432	3.077	20.30
J014609.33–092918	3.680	20.25
J032226.09–055824	3.763	20.30
J033119.66–074143	4.192	20.65
J034300.88–062229.9	3.511	20.15
J073146.99+364346	3.591	20.80
J073149.51+285448	2.686	20.55
J073823.94+340303.8	3.817	20.15
J073938.85+305951	3.355	20.10
J074145.00+215932.9	2.578	20.80

**Notes.** List of all candidate metal-poor DLAs flagged by Prochaska and collaborators from their analysis of the SDSS DR5. We caution that many of these will not be bona-fide DLAs.

(This table is available in its entirety in a machine-readable form in the online journal. A portion is shown here for guidance regarding its form and content.)

Previous work has compared DLA derived metallicities with nucleosynthesis yields. Akerman et al. (2004) adopted yields from Meynet & Maeder (2002) for stars of mass  $8 M_{\odot} < M_{*} < 80 M_{\odot}$ , and for lower mass stars with  $0.8 M_{\odot} < M_{*} < 8 M_{\odot}$  those of van den Hoek & Groenewegen (1997). Models of Heger & Woosley (2002), Chieffi & Limongi (2002), and Umeda & Nomoto (2002) agree in their predictions of yields of some elements, but provide divergent estimates of the ratios of [C/O] for high-mass nucleosynthesis.

In this paper, we present a large set of Keck ESI observations that demonstrate the efficiency of pre-selection, and provide estimates of metallicity for 35 DLA systems. We present the first results on the low-metallicity distribution function of high- $z$  galaxies and compare these results to earlier studies and to nucleosynthesis models.

## 2. OBSERVATIONS AND DATA REDUCTION

### 2.1. Sample Selection

The DLA systems studied in this paper were drawn from a much larger list of metal-poor DLA candidates that was generated as follows. We started first with the complete set of DLA candidates discovered by Prochaska et al. (2005) and Prochaska & Wolfe (2009) in the SDSS, Data Releases 1–5. The candidates were identified using an automated algorithm that focused solely on Ly $\alpha$  absorption. These authors also searched independently for metal-line systems in the same quasar spectra, restricting the search to the spectral region redward of the Ly $\alpha$  forest (see Herbert-Fort et al. 2006, for a description). All of the metal-poor DLA candidates had fewer than three significantly ( $4\sigma$ ) detected metal-line transitions. Furthermore, when analyzing the Ly $\alpha$  line, Prochaska and collaborators searched again for metal-line absorption using the estimated centroid of Ly $\alpha$  as a strong prior. Those without any obvious metal-line absorption (corresponding to  $\lesssim 0.5 \text{ \AA}$ ) were flagged, in part to allow for an additional systematic uncertainty when estimating the H I column density. Altogether, we culled 405 metal-poor DLA candidates as listed in Table 1.

Note that these systems are not required to satisfy the selection criteria used by Prochaska et al. (2005) and Prochaska & Wolfe (2009) in their analysis. Many have poorer S/N spectra and many have  $N_{\text{HI}}$  values estimated to be lower than the standard DLA threshold of  $2 \times 10^{20} \text{ cm}^{-2}$ . We also emphasize that the overwhelming majority of DLAs detected in the SDSS survey do exhibit significant metal-line absorption indicating a metallicity in excess of  $\approx 1/100$  solar. The sample studied in this paper therefore represents the most promising candidates for systems with very low metallicity.

From Table 1, we chose a subset for follow-up observations at higher spectral resolution and signal-to-noise ratio (S/N). Our principal goal was to discover the most metal-poor DLAs, not to choose targets that broadly sample the metal-poor candidate list. To this end, we selected targets with the following characteristics: (1) the background quasar's R.A. coincided with our Winter/Spring observing allocation, (2) brighter quasars to maximize the efficiency of our observing time, (3) candidates with larger  $N_{\text{HI}}$  values, (4) candidates with absorption redshifts near the quasar emission redshift such that key transitions (e.g., O I 1302, and C II 1334) lay outside the Ly $\alpha$  forest, and (5) quasars with more than one metal-poor DLA candidates. The targets eventually observed represent a balance between these various factors. Note that the combined factors generally imply that the DLA candidates satisfy the statistical criteria of Prochaska et al. (2005).

## 2.2. Data Acquisition and Reduction

The data for this paper were acquired during two observing runs at the Keck II telescope in 2007 March and 2008 May. We describe the observations of both observing runs below. Observations of a sample of 10 quasars selected for low metallicity were taken at the Keck II telescope for  $\sim 5.5$  hr on 2007 March 16. The ESI (Sheinis et al. 2002) was used in echellette mode, which provides a free spectral range of 3900–10900 Å at a dispersion ranging from 0.16 Å pixel $^{-1}$  to 0.30 Å pixel $^{-1}$ , corresponding to a constant velocity dispersion of 11.5 km s $^{-1}$  pixel $^{-1}$ . The slit width remained constant at 0".75 for all of the observations, providing an FWHM resolution of  $\approx 57$  km s $^{-1}$ . The Clear S filter was used for the observations. Exposure times on the target quasars ranged from 16 to 42 minutes, and resulted in a median S/N in the observed spectra of S/N = 29 per resolution element, with a range of values of  $23 < \text{S/N} < 64$ . A second sample of 23 quasars was observed in a second observing run at the Keck II telescope over three nights during 2008 May 7–9. The ESI spectrograph and instrument parameters were the same as on the first observing run, and greatly improved atmospheric conditions enabled a very high S/N to be obtained for most of the 23 quasars, resulting in a large sample of spectra with an average S/N of approximately S/N = 50.

The data from both observing runs were reduced using the IDL-based ESIRedux<sup>4</sup> pipeline (Prochaska et al. 2003b), and further processed using IDL to fit a high order spline to remove the quasar continuum, and extract column densities. The observed sample of quasars, apparent magnitudes, exposure times, the DLA redshifts, and other quantities are presented below in Table 2.

For each of the quasars, the H I profile was fit using XIDL and IDL routines, and the damping wings of the H I profiles were visually compared to the observed continuum in the quasar

spectrum. The H I profiles were fit independently by two of the team (J.X.P. and B.E.P.) and our H I columns agreed in all cases within the quoted uncertainty, in most cases to within 0.1 dex in  $N(\text{H I})$ . In a number of cases, our new estimates for the  $N(\text{H I})$  values indicate that the absorber does not meet the DLA threshold criterion of  $2 \times 10^{20} \text{ cm}^{-2}$ . The effects of these sub-DLA systems and their ionization corrections are discussed in Section 4.3.

Column densities of a wide range of species were derived using both the weak-line limit and apparent optical depth (AOD) technique (Savage & Sembach 1991). In some cases, a local continuum fit was performed during the AOD column density measurement to improve accuracy. Where multiple transitions existed, the transitions were combined with a weighted average among selected transitions to create a single optical depth profile for which the column density can be derived. For species with multiple transitions, we rejected strongly blended transitions and transitions in which complete saturation was observed (i.e., pixel values  $< 0.1$ ) either due to blending or saturation. A lower limit of column density was adopted for strongly saturated lines (such as Mg II) based on the corresponding column density from the equivalent width of the transition with the smallest oscillator strength in the weak-line limit. The column densities derived from low-ion transitions of the elements C, O, Si, Al, Fe, Mg, and Mn were compiled and converted to logarithmic abundances using the solar abundances of Lodders (2003), and the derived H I column density.

Concern about saturation of some species resulted in a separate analysis of the effects of saturation within the ESI spectra using standard COG analysis. These procedures are detailed in the following section.

## 2.3. Saturation Criteria

Since the limited resolution of the ESI spectrograph precludes a definitive measurement of component  $b$ -values, even when a large number of transitions allow a determination from the curve of growth, we have adopted the conservative assumption that many of our DLA systems will have narrow  $b$ -values, which would decrease the equivalent width limit at which saturation effects would become significant. In the case of such narrow  $b$ -values, lines which appear to be weak in our ESI spectra could begin to become saturated, implying some of the AOD derived column densities underestimate the actual values.

Within Pettini et al. (2008), three low-metallicity DLA systems were observed with high resolution and included seven separate velocity components with  $b$ -values ranging from 3.7 to 25.1 km s $^{-1}$ , with a mean  $b$ -value of 9.24 km s $^{-1}$ , and a median  $b$ -value of 6.5 km s $^{-1}$ . For our ESI sample, we have also examined the resulting  $b$ -values from the best fit to a normalized curve of growth for multiple transitions. When unblended absorption profiles exist over multiple transitions spanning a range of  $f$ -values and equivalent widths, we have been able to determine  $b$ -values for our ESI spectra, which lie within the range of values reported in Pettini et al. (2008). Figure 1 presents a histogram of measured  $b$ -values within our ESI sample (solid line), compared to the histogram of component  $b$ -values within the low-metallicity DLAs reported in Pettini et al. (2008). Our curve of growth fits showed  $b$ -values consistent with the range of reported  $b$ -values from Pettini et al. (2008), with a mean value of 12.5 km s $^{-1}$ , a median value of  $b = 7.5$  km s $^{-1}$ , and a range of values between  $5 \text{ km s}^{-1} < b < 25 \text{ km s}^{-1}$ .

Figure 2 shows the results of a simulation comparing derived AOD column densities for an absorption profile of an O I line at

<sup>4</sup> <http://www2.keck.hawaii.edu/inst/esi/ESIRedux/index.html>

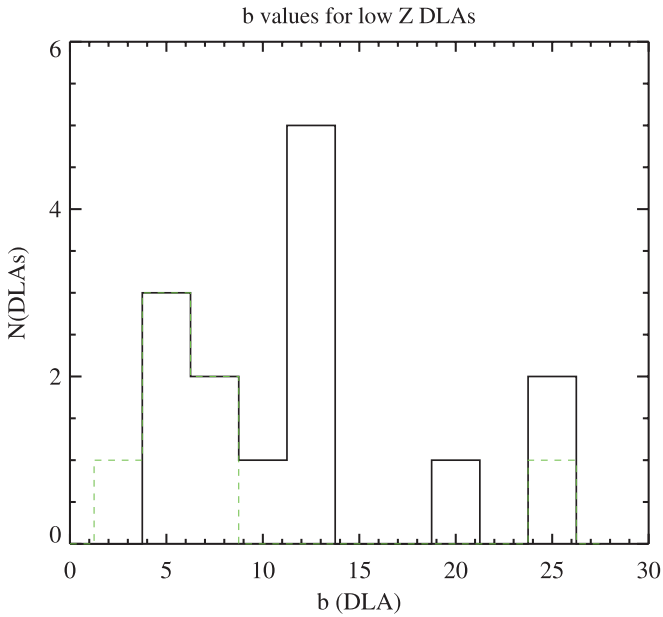
**Table 2**  
Targets Observed, Magnitudes, and Exposure Times

Quasar	$z_{\text{DLA}}$	R.A.(2000)	Decl.(2000)	$R$	Exp.(2008) (s)	Exp.(2007) (s)
SDSS0751+4516	3.046	07:51:13.05	+31:20:37	19.90	...	900.
SDSS0759+3129	3.035	07:59:51.83	+31:29:41	18.78	600.	...
SDSS0811+2838	2.434	08:11:10.32	+28:38:14	18.87	...	1200.
SDSS0814+5029	3.708	08:14:35.18	+50:29:46	18.34	...	1800.
SDSS0831+3358	2.304	08:31:02.55	+33:58:03	18.76	...	1000.
SDSS0844+4624	3.336	08:44:38.63	+46:24:25	19.54	...	1800.
SDSS0910+1026	2.398	09:24:59.91	+09:51:03	19.75	...	1200.
SDSS0924+0951	3.338	09:24:59.91	+09:51:03	19.75	2400.	...
SDSS0928+0939	2.910	09:28:14.93	+09:39:55	19.44	1800.	...
SDSS0940+0549	2.577	09:40:54.56	+05:49:03	18.90	1100.	1200.
****	2.578	****	****	****	****	****
SDSS0955+4116	3.280	09:55:42.12	+41:16:55	19.36	1800	...
SDSS1001+0343	3.078	10:01:51.45	+03:43:01	19.02	1222.	1200.
SDSS1003+5520	2.502	10:03:21.11	+55:20:59	19.40	2000.	...
SDSS1031+4055	2.569	10:31:26.13	+40:55:32	19.43	2000.	...
SDSS1037+0139	2.705	10:37:24.40	+01:39:33	19.14	...	1400.
SDSS1048+3911	2.296	10:48:26.03	+39:11:10	18.82	1200.	...
SDSS1108+1209	3.396	11:08:55.47	+12:09:53	18.63	...	1200.
SDSS1156+5513	2.481	11:56:59.59	+55:13:08	18.84	1200.	1200.
****	2.481	****	****	****	****	****
****	2.498	****	****	****	****	****
SDSS1219+1603	3.004	12:19:28.92	+16:03:57	19.47	1800	...
SDSS1251+4120	2.730	12:51:25.36	+41:20:00	18.95	1300.	...
SDSS1305+2902	2.386	13:05:48.92	+29:02:28	18.77	971.	...
SDSS1325+1255	3.550	13:25:54.12	+12:55:46	19.19	1430.	...
SDSS1327+4845	2.447	13:27:29.75	+48:45:00	19.30	1800.	...
****	2.612	****	****	****	****	...
SDSS1331+4838	3.692	13:31:46.21	+48:38:26	19.26	1800.	...
SDSS1349+1242	3.743	13:49:39.78	+12:42:30	19.61	2200.	...
SDSS1350+5952	2.756	13:50:44.21	+59:52:05	19.26	1800.	...
SDSS1358+0349	2.853	13:58:03.97	+03:49:36	18.65	930.	...
SDSS1358+6522	3.067	13:58:42.92	+65:22:36	18.71	1000.	...
SDSS1402+5909	3.774	14:02:43.97	+59:09:58	19.80	2500.	...
SDSS1440+0637	2.518	14:40:51.89	+06:37:09	19.33	1800.	...
****	2.825	****	****	****	****	...
SDSS1456+0407	2.320	14:56:24.66	+04:07:41	19.23	1600.	...
****	2.674	****	****	****	****	...
SDSS1557+2320	3.540	15:57:38.39	+23:20:57	19.46	1800.	...
SDSS1637+2901	3.496	16:37:47.58	+29:01:35	19.91	2400.	...
SDSS1654+3509	2.811	16:54:08.17	+35:09:54	18.92	1200.	...
SDSS1709+3417	2.530	17:09:31.01	+34:17:31	19.32	1611.	...
****	3.010	****	****	****	****	...
SDSS1717+5802	3.046	17:17:19.45	+58:02:18	19.61	2100.	...
SDSS2114v0632	4.126	21:14:50.33	-06:32:57	19.75	2400.	...

the simulated ESI resolution with the “true” column densities, for a DLA system with low  $b$ -values consistent with some of the low-metallicity DLA systems as measured by Pettini et al. (2008). The three lines represent the correction needed to derive the “true” column densities as a function of equivalent widths for  $b$ -values of 6.5, 7.5, and 8.5 km s<sup>-1</sup>. This plot graphically displays the difficulty in using underresolved spectra for strong absorption lines, as the actual column densities diverge from the AOD column densities as saturation becomes an issue for lines with equivalent widths greater than 50 mÅ, and by equivalent widths of 130 mÅ, the saturation corrections become too large to allow any reliable determination of column densities (see also Prochaska 2006). Fortunately, many of the DLA absorption lines in this work are weak, and have equivalent widths which show that for  $b$ -values in the range of 6.5 km s<sup>-1</sup> <  $b$  < 8.5 km s<sup>-1</sup>, the lines are either unsaturated or mildly saturated.

We do caution, however, that some of the DLAs could have multiple components with very low Doppler parameter (e.g.,  $b = 3$  km s<sup>-1</sup>); such systems would require very large corrections to the reported column densities. Indeed, the highly conservative reader may wish to consider our measurements as lower limits to the column densities until higher resolution spectroscopy is obtained. We wish to emphasize, however, that such low Doppler parameters have not been observed for the majority of DLA systems, nor are they expected. The line-profile fitting of individual DLAs generally report  $b$ -values of 5–10 km s<sup>-1</sup> (e.g., Prochaska & Wolfe 1996; Dessauges-Zavadsky et al. 2006). The only obvious exceptions are the subset of DLAs which show molecular (i.e., cold) gas (e.g., Noterdaeme et al. 2008; Jorgenson et al. 2009). Second, AOD analyses have not indicated any significant “hidden saturation” in DLA profiles (Prochaska & Wolfe 1996). Third, if the gas is a





**Figure 1.** Histogram of best-fit  $b$ -values for our ESI sample (solid black line), based on fitting of a normalized curve of growth for multiple species of varying  $f$ -values for those DLAs with non-blended weak lines in several species. The DLA components within the low-metallicity sample of Pettini et al. (2008) are also included (dashed green line). The median  $b$ -value for our ESI DLA sample is determined to be  $b = 7.5 \text{ km s}^{-1}$ , and agrees well with the median  $b$ -value for the Pettini et al. sample of  $b = 6.5 \text{ km s}^{-1}$ .

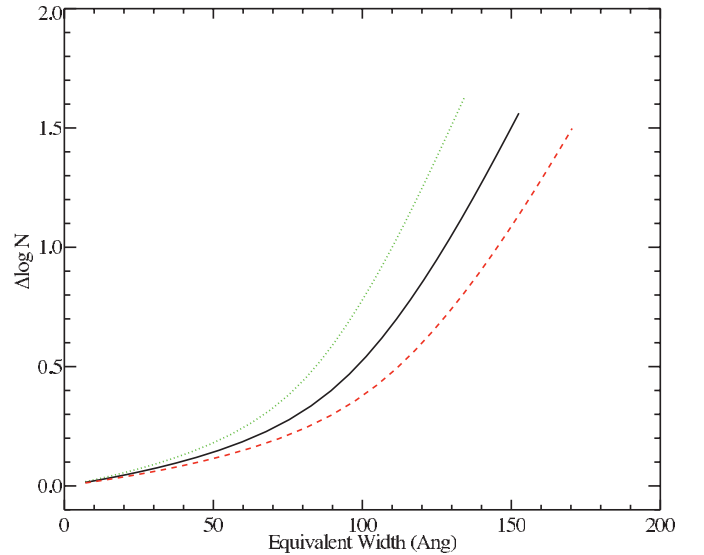
(A color version of this figure is available in the online journal.)

warm neutral medium ( $T \approx 7000 \text{ K}$ ; Kanekar et al. 2009) then even ignoring macroscopic motions one predicts  $b > 3 \text{ km s}^{-1}$  for carbon. Therefore, we expect few of the systems presented here to exhibit such low Doppler parameters.

From the results of our simulation shown in Figure 2, we have developed criteria for determining whether our absorption lines should be considered saturated, mildly saturated, or non-saturated, and for those transitions which are mildly saturated, we have calculated “saturation corrections” to enable an estimation of the range of possible column densities in the case of narrow lines. Transitions were considered mildly saturated when their equivalent widths exceed  $50 \text{ m}\text{\AA}$ , but were less than  $100 \text{ m}\text{\AA}$ . To determine column densities, we have applied corrections to the AOD column density  $\delta N$  in the regime from  $50$  to  $130 \text{ m}\text{\AA}$ , while in all cases we consider column densities from species exceeding  $130 \text{ m}\text{\AA}$  to be lower limits due to saturation effects.

To correct the column density when the equivalent width is between  $50$  and  $130 \text{ m}\text{\AA}$ , we adopt the correction shown in Figure 2 for an adopted  $b$ -value of  $7.5 \text{ km s}^{-1}$ . Our correction involves adjusting the derived AOD column densities upward to a value midway between the non-adjusted (weak-line limit) column density and the fully corrected column density, with error estimates increased to span the range of possible column densities. By choosing to adjust our adopted column densities upward (with corresponding increases in our error estimates), we are able to account for the probable range of column densities if the DLA sightline has narrow absorption lines in the range of  $b = 7.5 \text{ km s}^{-1}$ , but we also preserve as a lower limit of the error bar the estimated column density for larger  $b$ -values.

The result of this procedure is not entirely satisfactory, but represents our best attempt to provide the most information available from our moderate resolution ESI spectra, while



**Figure 2.** Plot of modeled column density offset ( $\Delta \log N \equiv \log N_{\text{tru}} - \log N_{\text{model}}$ ) obtained using the AOD technique for  $b$ -values of  $6.5$ ,  $7.5$ , and  $8.5 \text{ km s}^{-1}$  (top to bottom), using the ESI instrument parameters from our observing run, with a  $0.75 \text{ arcsec}$  slit, on a single  $\text{O I } 1302 \text{ \AA}$  line. The possibility of low  $b$ -values of low-metallicity DLA systems requires caution in deriving column densities from the moderate resolution spectra of the ESI spectrograph. (A color version of this figure is available in the online journal.)

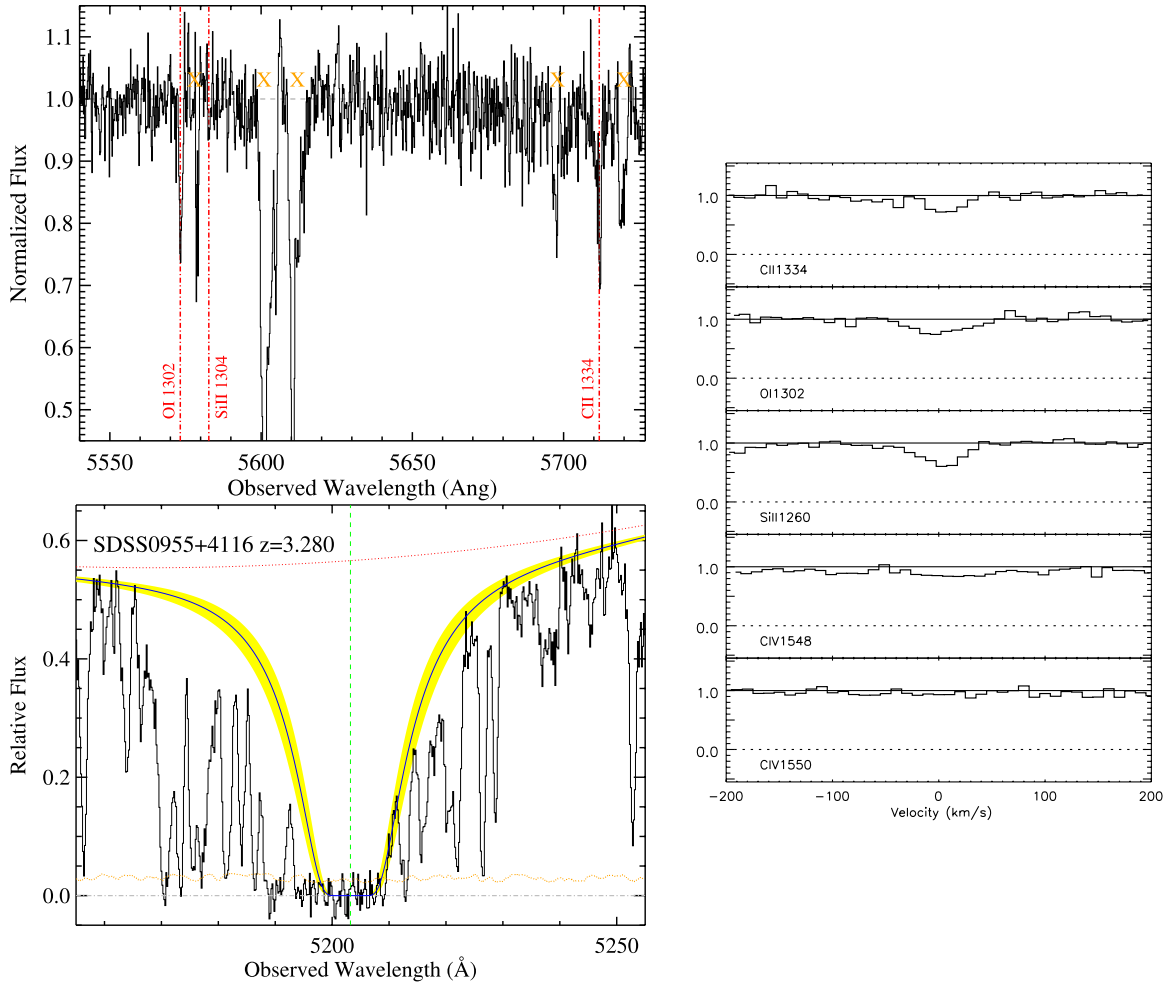
recognizing the limitations imposed by the limited resolution. A future study with higher resolution would enable more accurate determination of column densities for the stronger lines, but with this compromise we provide a range of possible values for these column densities.

#### 2.4. Sensitivity Limit for Extremely Low-metallicity Systems

We also computed the lowest metallicities detectable within our program, as it is important to confirm that our search was capable of detecting extremely low-metallicity systems. For a typical value of  $\text{H I} = 20.5$ , and redshift  $z = 3.0$ , we can calculate the lowest limiting metallicity for various transitions, given a  $4$  pixel wide absorption profile close to the instrumental resolution of  $\text{FWHM} \approx 57 \text{ km s}^{-1}$ , and a  $4\sigma$  detection. Table 3 provides an estimate of the lowest possible observable values of metallicity for a number of species, for spectra with  $\text{S/N} = 10$  and  $\text{S/N} = 20$ . These limits were calculated using parameters for a typical DLA where we adopted values for  $N(\text{H I}) = 20.5$ ,  $z = 3.0$ , and a metal line profile of  $4$  pixels, and a  $4\sigma$  detection. Since our spectra all had values of  $\text{S/N} \geq 20$  for most of the free spectral range, the right column of Table 3 describes the limiting sensitivity of our survey to low-metallicity systems. All of the spectra have  $\text{O I}$  and  $\text{C II}$  equivalent widths larger than our minimum limiting sensitivity, but in some spectra species such as  $\text{Fe II}$  and  $\text{Si II}$  approached this limiting minimum value, and therefore for some DLA systems we report upper limits only in these species. The following section describes our results in detail for the complete sample of DLA systems.

### 3. ANALYSIS

We have conducted a uniform analysis of the entire sample of  $38$  quasars using AOD code and apply corrections for line saturation as discussed in Section 2.3. In cases where the metallicity and elemental abundances are very low, we provide complete details about the DLA system, its detected absorption lines, and the resulting column densities and abundances



**Figure 3.** Section of our spectrum SDSS0955+4116 showing the lines of C II  $\lambda$ 1334, Si II  $\lambda$ 1304, and O I  $\lambda$ 1302 for the DLA system at  $z = 3.28102$  (left), and the H I damped Ly $\alpha$  profile at the same redshift (right), along with our best-fit model (solid line) for  $\log(N(\text{H I})) = 20.10$ , and  $1\sigma$  limits to the H I fit (yellow/shaded region). (A color version of this figure is available in the online journal.)

**Table 3**

Calculated Metallicity Limits for Different Ions, for a Typical DLA with  $N(\text{H I}) = 20.5$ ,  $z = 3.0$ , and with an FWHM of 4 Pixels<sup>a</sup>

Transition	S/N = 10	S/N = 20
O I 1302	-3.5	-3.8
Si II 1304	-2.6	-2.9
C II 1334	-3.8	-4.1
Si II 1526	-2.9	-3.2
Fe II 1608	-2.6	-2.9
Al II 1670	-3.1	-3.4

**Note.**<sup>a</sup> The resulting calculated limits assume a  $4\sigma$  detection for spectra with the two values of S/N listed.

in individual subsections. For those DLAs not discussed below, the results are included in the summary tables and included in our considerations of the statistics of metal abundances in the later sections of the paper.

### 3.1. Results for Individual DLA Systems with Low Metallicity

#### 3.1.1. SDSS0955+4116

We present in the left panel of Figure 3 a plot of our Keck/ESI spectra for the DLA system toward SDSS0955+4116, in the wavelength range which includes important lines of C II  $\lambda$ 1334, Si II  $\lambda$ 1304, and O I  $\lambda$ 1302. The right panel of Figure 3 shows

the H I DLA profile, and we indicate our best-fit model for  $N(\text{H I}) = 20.10$ , along with the one  $\sigma$  range of  $N(\text{H I})$  in the shaded region. A curve of growth analysis of the sightline suggests a  $b$ -value of  $7.5 \text{ km s}^{-1}$ , although all of the lines are weak, and an accurate measurement of the  $b$ -value is difficult. Since all species have equivalent widths less than  $100 \text{ m}\text{\AA}$ , we have corrected the derived AOD column densities according to the prescription in Section 2.3. The results of our adopted column densities for the various species are presented in Table 4. This DLA also is observed to have relatively weak absorption lines for both C II and O I, and is one of five DLA systems from our survey where both equivalent widths of these species are less than  $130 \text{ m}\text{\AA}$ , which forms the weak-line subsample used in Figure 17. In Table 13, we summarize the abundances for this DLA, for which we derive values of  $[\text{C}/\text{H}] = -2.82 \pm 0.14$ ,  $[\text{O}/\text{H}] = -2.82 \pm 0.10$ , and  $[\text{Si}/\text{H}] = -2.75 \pm 0.18$ . The observed value of  $[\text{C}/\text{H}]$  is the second lowest of our sample, while the observed value of  $[\text{O}/\text{H}]$  is the sixth lowest of our sample. Non-detections of the species Al II and Fe II give upper limits of  $[\text{Al}/\text{H}] < -2.74$ , and  $[\text{Fe}/\text{H}] < -2.30$ .

#### 3.1.2. SDSS1001+0343

The DLA system toward SDSS1001+0343 is one of our lowest metallicity systems, and we present in the left panel of Figure 4 a plot of our spectra for the DLA system of

**Table 4**  
Ionic Column Densities for J0955+4116  $z = 3.280$

Ion	$\lambda$ ( $\text{\AA}$ )	$\log f$	$v_{\text{int}}^{\text{a}}$ ( $\text{km s}^{-1}$ )	$W_{\lambda}^{\text{b}}$ ( $\text{m\AA}$ )	$\log N$	$\log N_{\text{adopt}}$
C II	1334.5323	-0.8935	[-100, 60]	$77.8 \pm 10.4$	13.58	$13.72 \pm 0.14$
C II*	1335.7077	-0.9397	[-100, 50]	<21.2	<13.26	
C IV	1548.1950	-0.7194	[-100, 60]	$81.3 \pm 13.3$	$13.33 \pm 0.07$	$13.33 \pm 0.07$
	1550.7700	-1.0213	[-100, 60]	<27.3	<13.32	
O I	1302.1685	-1.3110	[-100, 60]	$65.1 \pm 9.0$	$13.94 \pm 0.05$	$14.04 \pm 0.10$
Al II	1670.7874	0.2742	[-100, 60]	<25.2	<11.92	<11.92
Si II	1260.4221	0.0030	[-100, 60]	$87.0 \pm 6.1$	$12.79 \pm 0.03$	$12.96 \pm 0.18$
	1304.3702	-1.0269	[-100, 60]	<19.2	<13.31	
	1526.7066	-0.8962	[-100, 60]	<29.3	<13.24	
	1808.0130	-2.6603	[-100, 60]	<27.2	<14.79	
Si IV	1393.7550	-0.2774	[-100, 60]	$99.4 \pm 11.6$	$13.08 \pm 0.05$	$13.07 \pm 0.05$
	1402.7700	-0.5817	[-100, 60]	$47.2 \pm 12.9$	$13.04 \pm 0.12$	
Fe II	1608.4511	-1.2366	[-100, 60]	<20.0	<13.34	<13.34

**Notes.**

<sup>a</sup> Velocity interval over which the equivalent width and column density are measured.

<sup>b</sup> Rest equivalent width.

SDSS1001+0343, as in Section 3.1.1, where the right panel of Figure 4 shows the H I DLA profile, with the best-fit model for  $N(\text{H I}) = 20.15$ . A curve of growth analysis of the available transitions for this sightline suggests a  $b$ -value of  $7.5 \text{ km s}^{-1}$ , and all transitions have equivalent widths less than  $80 \text{ m\AA}$ , where possible saturation effects are limited. We have adjusted the stronger lines for possible effects of saturation as described above, typically at a level of 0.1–0.2 dex. The results of our adopted column densities for the various species are presented in Table 5. In Table 13, we summarize the abundances for this DLA, for which we derive values of  $[\text{C}/\text{H}] = -2.85 \pm 0.13$ ,  $[\text{O}/\text{H}] = -2.93 \pm 0.08$ ,  $[\text{Si}/\text{H}] = -2.94 \pm 0.12$ ,  $[\text{Al}/\text{H}] < -2.82 \pm 0.13$ , and an upper limit of  $[\text{Fe}/\text{H}] < -2.32$ . This DLA also is observed to have very weak absorption lines for both C II and O I, and is one of five DLA systems from our survey where both equivalent widths of these species are less than  $130 \text{ m\AA}$ , which forms the weak-line subsample used in Figure 17. The observed value of  $[\text{C}/\text{H}]$  is the lowest in our sample, while the  $[\text{O}/\text{H}]$  is the third lowest of the sample.

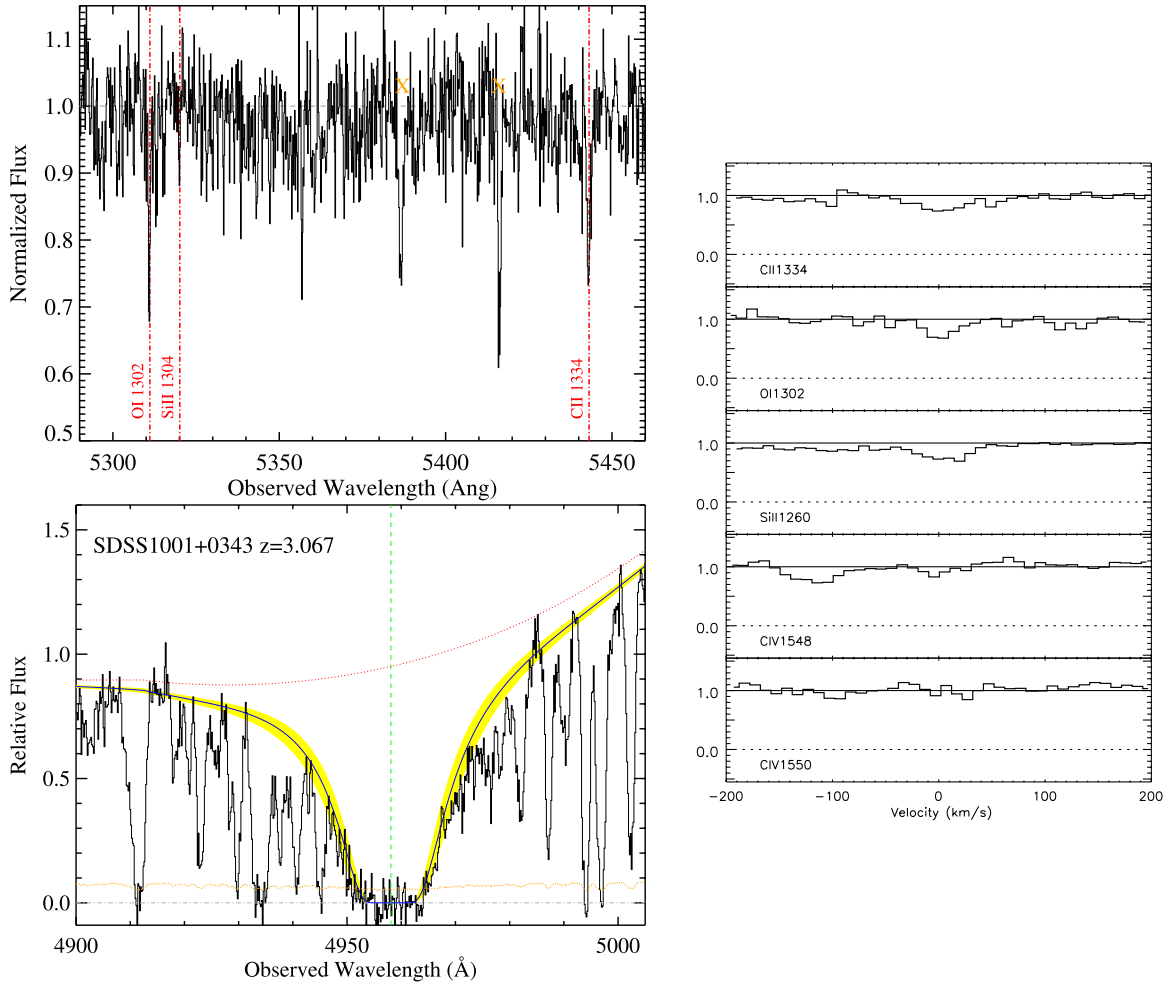
### 3.1.3. SDSS1219+1603

Figure 5 shows details of the spectrum of the DLA system of SDSS1219+1603, with two panels showing the spectrum near key transitions of C II, Si II, and O I (left), and the damped H I profile (right) as in Section 3.1.1. Our best-fit model for  $N(\text{H I}) = 20.35$ , and the SDSS1219+1603 DLA at  $z = 3.00372$ , appears adjacent to a sub-DLA with  $N(\text{H I}) = 20.15$  at  $z = 3.032$ . The results of our adopted column densities for the various species are presented in Table 6. In Table 13, we summarize the abundances for this DLA, for which we derive values of  $[\text{O}/\text{H}] = -2.59 \pm 0.34$ ,  $[\text{Si}/\text{H}] = -2.08 \pm 0.15$ , and  $[\text{Fe}/\text{H}] = -2.09 \pm 0.10$ , and we report only lower limits for the metallicities of  $[\text{C}/\text{H}] > -2.59$ , and  $[\text{Al}/\text{H}] > -2.29$  due to strong absorption of both species with equivalent widths exceeding  $130 \text{ m\AA}$ . The

SDSS1219+1603 DLA shows metal abundances at the high end of our sample, although much less than the mean  $[\text{X}/\text{H}] = -1.52$  from the survey of Prochaska et al. (2003b).

### 3.1.4. SDSS1305+2902

Figure 6 shows details of the spectrum of the DLA system of SDSS1305+2902, with two panels showing the spectrum near key transitions of C II, Si II, and O I (left), and the damped H I profile (right) as in Section 3.1.1. Our best-fit model for  $N(\text{H I}) = 20.25$ , which places this system very close to the lower limit of  $\log(N(\text{H I})) = 20.3$  that provides self-shielding in a DLA system. A curve of growth analysis for the SDSS1305+2902 DLA suggests a  $b$ -value of approximately  $10.0 \text{ km s}^{-1}$ , and some saturation corrections were applied for lines with strengths in the 50–130  $\text{m\AA}$  range. The resulting correction for the C II column density was quite large, and gives us only a very uncertain estimate of its column to within 0.45 dex, but we include the C II column density to allow for comparison with the O I, which had a smaller correction of 0.13 dex corresponding to its weaker absorption of  $70 \text{ m\AA}$  equivalent width. Absorption lines of the species Si II and Fe II both included detections of absorption with weak transitions with equivalent width less than  $50 \text{ m\AA}$  which were not saturated, and therefore we have the AOD adopted column densities for these species using only very small corrections for possible saturation. The results of our adopted column densities for the various species are presented in Table 7. In Table 13, we summarize the abundances for this DLA, for which we derive values of  $[\text{C}/\text{H}] = -2.46 \pm 0.45$ ,  $[\text{O}/\text{H}] = -2.90 \pm 0.12$ ,  $[\text{Si}/\text{H}] = -2.54 \pm 0.10$ ,  $[\text{Al}/\text{H}] = -2.83 \pm 0.10$ , and  $[\text{Fe}/\text{H}] = -2.79 \pm 0.13$ . This DLA has reported column densities for both C II and O I, and is one of five DLA systems from our survey where both equivalent widths of these species are less than  $130 \text{ m\AA}$ , which forms the “weak-line” subsample used in Figure 17, and has the fourth lowest measured value of



**Figure 4.** Section of our spectrum SDSS1001+0343 showing the lines of C II  $\lambda$ 1334, Si II  $\lambda$ 1304, and O I  $\lambda$ 1302 for the DLA system at  $z = 3.07867$  (left), and the H I damped Ly $\alpha$  profile at the same redshift (right), along with our best-fit model (solid line) for  $\log(N(\text{H I})) = 20.15$ , and  $1\sigma$  limits to the H I fit (yellow/shaded region). (A color version of this figure is available in the online journal.)

[O/H] in our sample. The SDSS1305+2902 DLA is also unique in that it has a complete sample of metallicities for the elements C, O, Si, Fe, and Al, which despite the very large uncertainty in the estimated [C/H] metallicity, helps constrain the relative metal abundances in the low-metallicity DLA gas.

### 3.1.5. SDSS1358+0349

Figure 7 shows details of the spectrum of the DLA system of SDSS1358+0349, with two panels showing the spectrum near key transitions of C II, Si II, and O I (left), and the damped H I profile (right) as in Section 3.1.1. Our best-fit model for  $N(\text{H I}) = 20.50$ , placing this system well above the DLA threshold of  $\log(N(\text{H I})) = 20.3$ . A curve of growth fit to the transitions for the SDSS1358+0349 DLA suggests an underlying  $b$ -value of approximately  $12.5 \text{ km s}^{-1}$ , and we have made some allowance for saturation within the regime of 50–100 mÅ, which includes the transitions for Si II and O I. The column densities of the stronger transitions Si II and O I have been corrected to account for the effects of saturation at the level of 0.17 and 0.25 dex, respectively. Sufficient transitions exist with weaker absorption at equivalent widths  $< 70 \text{ mÅ}$  for the species of Fe II, Si II, and Al II, and therefore these species have small estimated saturation, which are close to the experimental errors. The results of our adopted column densities for the various species are presented in Table 8. In Table 13, we summarize the

abundances for this DLA, for which we derive values of [O/H] =  $-2.88 \pm 0.25$ , [Si/H] =  $-2.81 \pm 0.17$ , and [Al/H] =  $-2.87 \pm 0.08$ , and [Fe/H] =  $-3.03 \pm 0.05$ . We report only a lower limit for [C/H]  $> -2.58$ , due to the very strong absorption of C II which has an equivalent width of 194 mÅ. The observed value of [O/H] =  $-2.88$  for the SDSS1358+0349 DLA, even allowing for possible saturation, is the fifth lowest of our sample.

### 3.1.6. SDSS1358+6522

Figure 8 shows details of the spectrum of the DLA system of SDSS1358+6522, with two panels showing the spectrum near key transitions of C II, Si II, and O I (left), and the damped H I profile (right) as in Section 3.1.1. Our best-fit model for  $N(\text{H I}) = 20.35$ , and a curve of growth fit to the absorption lines of the SDSS1358+6522 DLA, indicates small  $b$ -values in the range of 5–10  $\text{km s}^{-1}$ . We have adopted saturation corrections for the absorption lines with equivalent widths between 50 and 130 mÅ using the rubric described in Section 2.3. The adjustment for our C II transition is significant, at 0.43 dex, and greatly increases the uncertainty of the 122 mÅ C II line, but we have included this transition to enable comparison with the O I absorption, which is weaker at only 60 mÅ equivalent width. The absorption from Si II is detected in the transitions at 1526 and 1260 Å, but with different resulting AOD column densities, suggesting that some saturation is occurring within the stronger



**Table 5**  
Ionic Column Densities for J1001+0343  $z = 3.078$

Ion	$\lambda$ (Å)	$\log f$	$v_{\text{int}}^{\text{a}}$ (km s $^{-1}$ )	$W_{\lambda}^{\text{b}}$ (mÅ)	$\log N$	$\log N_{\text{adopt}}$
C II	1334.5323	-0.8935	[-50, 50]	76.7 $\pm$ 9.5	13.63 $\pm$ 0.06	13.76 $\pm$ 0.13
C II*	1335.7077	-0.9397	[-50, 50]	<19.5	<13.20	
C IV	1548.1950	-0.7194	[-200, 50]	136.6 $\pm$ 16.3	13.58 $\pm$ 0.05	13.56 $\pm$ 0.05
	1550.7700	-1.0213	[-200, 50]	59.0 $\pm$ 17.0	13.50 $\pm$ 0.12	
O I	1302.1685	-1.3110	[-50, 50]	57.6 $\pm$ 9.0	13.90 $\pm$ 0.07	13.98 $\pm$ 0.08
Al II	1670.7874	0.2742	[-50, 50]	32.6 $\pm$ 9.4	11.87 $\pm$ 0.13	11.87 $\pm$ 0.13
Si II	1260.4221	0.0030	[-50, 50]	72.5 $\pm$ 4.3	12.70 $\pm$ 0.05	12.82 $\pm$ 0.12
	1304.3702	-1.0269	[-50, 50]	<19.3	<13.32	
	1526.7066	-0.8962	[-50, 50]	<19.0	<13.04	
	1808.0130	-2.6603	[-50, 50]	<26.9	<14.81	
Si IV	1393.7550	-0.2774	[-50, 50]	47.7 $\pm$ 9.8	12.75 $\pm$ 0.09	12.75 $\pm$ 0.09
Fe II	1608.4511	-1.2366	[-50, 50]	<38.7	<13.49	<13.31
	2344.2140	-0.9431	[-50, 50]	<67.5	<13.31	

**Notes.**<sup>a</sup> Velocity interval over which the equivalent width and column density are measured.<sup>b</sup> Rest equivalent width.

**Table 6**  
Ionic Column Densities for J1219+1603  $z = 3.003$

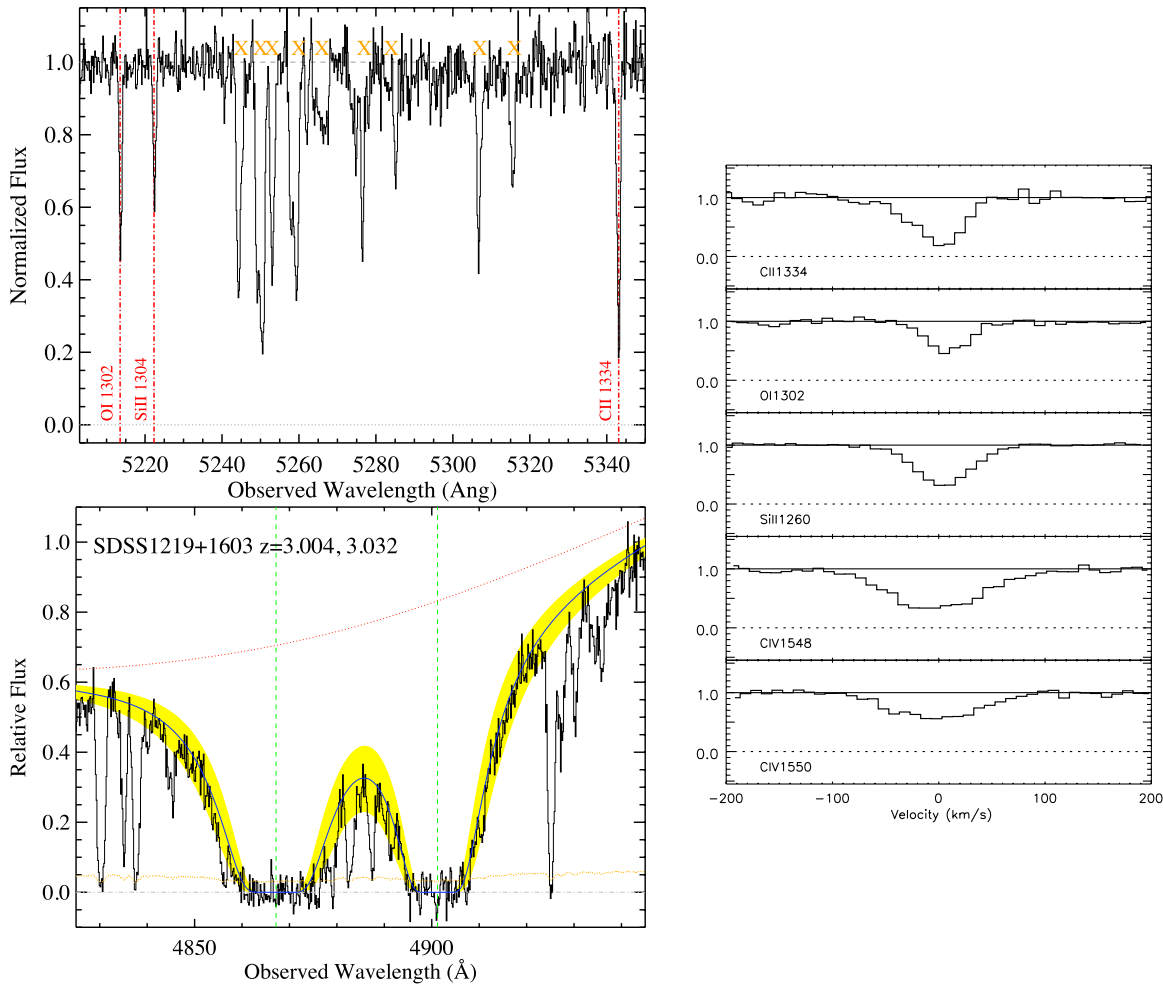
Ion	$\lambda$ (Å)	$\log f$	$v_{\text{int}}^{\text{a}}$ (km s $^{-1}$ )	$W_{\lambda}^{\text{b}}$ (mÅ)	$\log N$	$\log N_{\text{adopt}}$
C II	1334.5323	-0.8935	[-70, 70]	208.0 $\pm$ 10.0	14.02 $\pm$ 0.06	>14.22
C IV	1548.1950	-0.7194	[-150, 100]	391.0 $\pm$ 16.3	13.98 $\pm$ 0.05	14.05 $\pm$ 0.05
	1550.7700	-1.0213	[-150, 100]	231.0 $\pm$ 17.0	14.05 $\pm$ 0.12	
O I	1302.1685	-1.3110	[-70, 70]	112.0 $\pm$ 9.0	14.18 $\pm$ 0.07	14.52 $\pm$ 0.34
Al II	1670.7874	0.2742	[-70, 70]	149.0 $\pm$ 9.4	12.50 $\pm$ 0.10	>12.60
Si II	1260.4221	0.0030	[-70, 70]	186.7 $\pm$ 9.3	>13.11	13.88 $\pm$ 0.15
	1304.3702	-1.0269	[-70, 70]	79.1 $\pm$ 7.0	13.74 $\pm$ 0.05	
	1526.7066	-0.8962	[-70, 70]	135.0 $\pm$ 7.0	>13.70	
Si IV	1393.7550	-0.2774	[-150, 100]	256.0 $\pm$ 9.8	13.45 $\pm$ 0.10	13.45 $\pm$ 0.10
Fe II	1608.4511	-1.2366	[-70, 70]	66.8 $\pm$ 6.0	13.70 $\pm$ 0.05	13.80 $\pm$ 0.10

**Notes.**<sup>a</sup> Velocity interval over which the equivalent width and column density are measured.<sup>b</sup> Rest equivalent width.

line of the DLA. We have adopted the column density from the weaker of the two lines, and increased the uncertainty estimate for  $N(\text{Si II})$  to encompass the possible effects of saturation.

The absorption from the species of Fe II is not convincingly detected in any of several transitions, and so we present only upper limits for this ion. The results of our adopted column densities for the various species are presented in Table 9. In Table 13, we summarize the abundances for this DLA, for which we derive values of  $[\text{C}/\text{H}] = -2.59 \pm 0.43$ ,  $[\text{O}/\text{H}] = -3.08 \pm 0.18$ ,  $[\text{Si}/\text{H}] = -2.96 \pm 0.20$ , and  $[\text{Al}/\text{H}] = 2.78 \pm$

0.16, with an upper limit to  $[\text{Fe}/\text{H}] < -3.09$ . This DLA also is observed to have measurable column densities for both C II and O I, and is one of five DLA systems from our survey where both equivalent widths of these species are less than 130 mÅ, which forms the weak-line subsample used in Figure 17. The measured value of  $[\text{O}/\text{H}]$  is the second lowest of our sample, while the measured value of  $[\text{C}/\text{H}]$  is the fourth lowest of the sample. Further observations of this system at higher resolution would improve our estimates, since the uncertainty from possible saturation effects dominate our reported values of  $[\text{C}/\text{H}]$  and



**Figure 5.** Section of our spectrum SDSS1219+1603 showing the lines of C II  $\lambda$ 1334, Si II  $\lambda$ 1304, and O I  $\lambda$ 1302 for the DLA system at  $z = 3.00372$  (left), and the H I damped Ly $\alpha$  profile at the same redshift (right), along with our best-fit model (solid line) for  $\log(N(\text{H I})) = 20.35$ , and  $1\sigma$  limits to the H I fit (yellow/shaded region). A smaller sub-DLA with  $\log(N(\text{H I})) = 20.15$  appears adjacent to this DLA at  $z = 3.032$ .

(A color version of this figure is available in the online journal.)

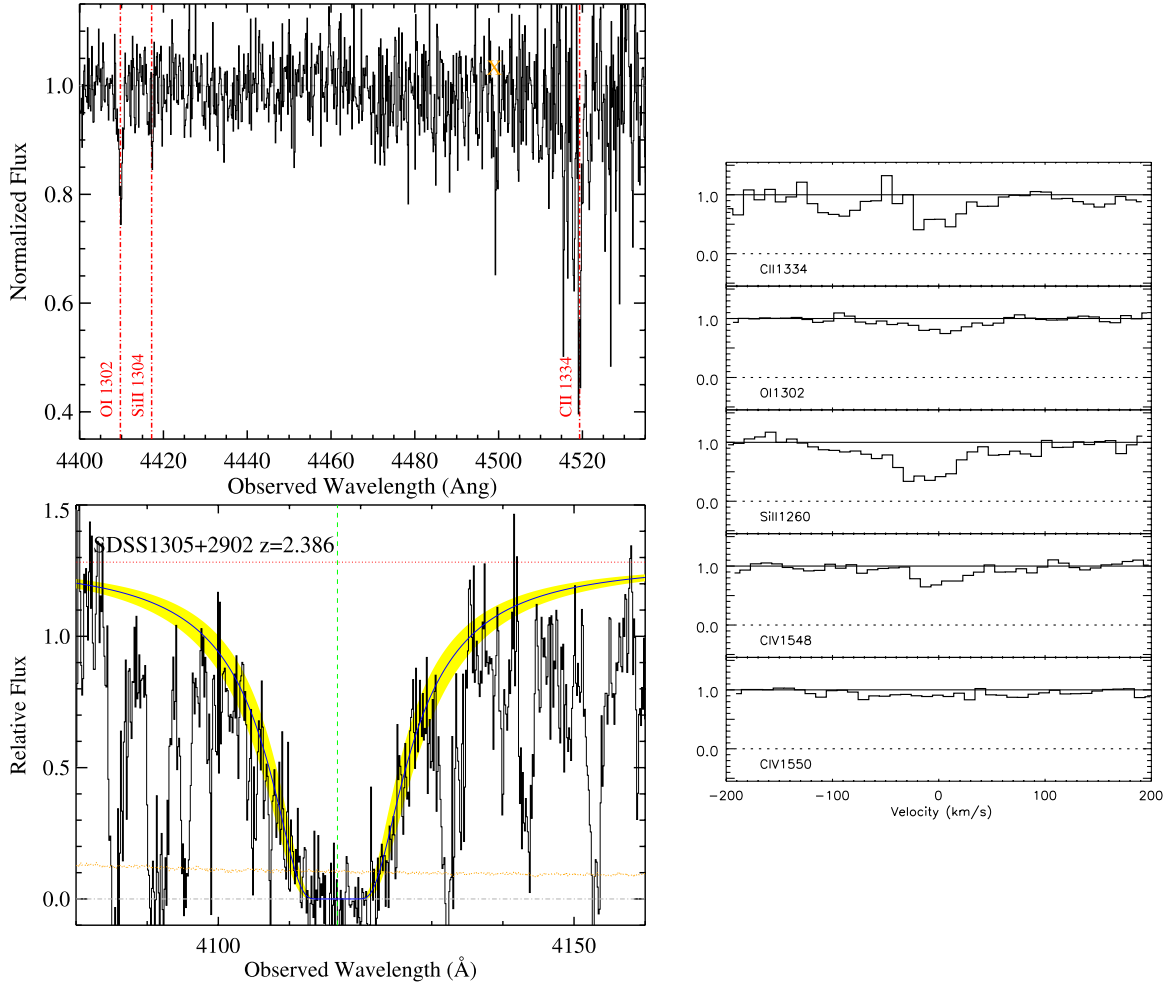
[O/H], and especially limit our ability to provide more than a very approximate estimate of [C/H].

### 3.1.7. SDSS1456+0407

Figure 9 shows details of the spectrum of the DLA system of SDSS1456+0407, with two panels showing the spectrum near key transitions of C II, Si II, and O I (left), and the damped H I profile (right) as in Section 3.1.1. Our best-fit model for  $N(\text{H I}) = 20.35$ , and the SDSS1456+0407 DLA, includes good detections of O I, Al II, Si II, and Fe II, which have equivalent widths in the range of 30–100 mÅ, and a strong absorption line from C II with equivalent width in excess of 190 mÅ, which only provides a lower limit to  $N(\text{C II})$ . A set of weak Fe II lines gives additional data points for our curve of growth, which has a best fit for  $b > 15 \text{ km s}^{-1}$ , but with significant uncertainty in the fit. To be conservative in our estimates of metallicity, we have assumed a  $b$ -value of  $7.5 \text{ km s}^{-1}$ , and have adopted the corrections for saturation described in Section 2.3 for the species O I, Al II, Si II, and Fe II. The results of our adopted column densities for the various species are presented in Table 10. In Table 13, we summarize the abundances for this DLA, for which we derive values of  $[\text{C}/\text{H}] > -2.48$ ,  $[\text{O}/\text{H}] = -2.56 \pm 0.28$ ,  $[\text{Si}/\text{H}] = -2.47 \pm 0.07$ ,  $[\text{Al}/\text{H}] = -1.96 \pm 0.45$ ,  $[\text{Fe}/\text{H}] = -2.89 \pm 0.1$ , and  $[\text{Mg}/\text{H}] = -2.37 \pm 0.1$ .

### 3.1.8. SDSS1637+2901

The SDSS1637+2901 DLA is one of our lowest metallicity quasar DLA systems, and Figure 10 shows details of the spectrum of the DLA system of SDSS1456+0407, with two panels showing the spectrum near key transitions of C II, Si II, and O I (left), and the damped H I profile (right) as in Section 3.1.1. Our best-fit model for  $N(\text{H I}) = 20.70$ , well above the threshold of  $\log(N(\text{H I})) = 20.3$  for DLAs. The SDSS1637+2901 DLA shows evidence of narrow lines, and the lowest  $b$ -value of our sample, with  $b = 5 \text{ km s}^{-1}$  from curve of growth fitting. For this reason, we have adopted the saturation corrections described in Section 2.3. The strong absorption from C II is accompanied by an even stronger absorption line at the position of the C II\* absorption at 1335.7 Å for  $z = 3.496$ , which is evidence for either extremely strong C II\* excitation or blending from other species, perhaps arising from a C II absorption at slightly higher redshift. There also appears to be a pair of absorption lines to the left of the O I and Si II lines from Figure 10 arising from an unrelated C IV doublet. The blending allows only a determination of the upper limit for the C II\* column density. The likelihood of saturation in the C II makes our derived column density of C II a lower limit. However, the absorption from O I has an equivalent width of 91 mÅ, which we have corrected for saturation effects with a correction of 0.18



**Figure 6.** Section of our spectrum SDSS1305+2902 showing the lines of C II  $\lambda$ 1334, Si II  $\lambda$ 1304, and O I  $\lambda$ 1302 for the DLA system at  $z = 2.38645$  (left), and the H I damped Ly $\alpha$  profile at the same redshift (right), along with our best-fit model (solid line) for  $\log(N(\text{H I})) = 20.25$ , and  $1\sigma$  limits to the H I fit (yellow/shaded region). (A color version of this figure is available in the online journal.)

dex, according to the saturation correction for  $b = 7.5 \text{ km s}^{-1}$  described in Section 2.3. Smaller corrections for saturation of 0.14 and 0.11 dex were applied to the column densities of Si II and Al II, respectively. The results of our adopted column densities for the various species are presented in Table 11. In Table 13, we summarize the abundances for this DLA, for which we derive values of  $[\text{O}/\text{H}] = -3.17 \pm 0.20$  (the lowest of our sample),  $[\text{Si}/\text{H}] = -2.90 \pm 0.14$ ,  $[\text{Al}/\text{H}] = -2.95 \pm 0.11$ ,  $[\text{Fe}/\text{H}] = -2.40 \pm 0.10$ , and a lower limit of  $[\text{C}/\text{H}] > -3.20$ , due to the very strong absorption at equivalent width of more than  $140 \text{ m}\text{\AA}$ .

### 3.1.9. SDSS2114–0632

Figure 11 shows details of the spectrum of the DLA system of SDSS2114–0632, with two panels showing the spectrum near key transitions of C II, Si II, and O I (left), and the damped H I profile (right) as in Section 3.1.1. Our best-fit model for  $N(\text{H I}) = 20.40$ , just above the DLA threshold of  $\log(N(\text{H I})) = 20.3$ . The SDSS2114–0632 DLA is at the highest redshift ( $z = 4.126$ ) of our sample, and includes extremely weak detected absorption lines of the species C II, O I, and Si II. The absorption from the species Al II and Fe II is extremely weak, and we are not able to detect these species, and therefore provide only upper limits for both of them. A curve of growth analysis of the detected species suggests a large  $b$ -value in the range of

$20 \text{ km s}^{-1}$ , which would greatly reduce the required correction for saturation. Despite the possibility of large  $b$ -values for the DLA, we account for possible saturation in our reported column densities, as our curve of growth fitting cannot rule out the possibility of multiple narrow components within the DLA. We have included the O I transition in our compilation since the lower bound of the measured equivalent width is within our threshold of  $130 \text{ m}\text{\AA}$ , to enable comparison with C II and other species. The results of our adopted column densities for the various species are presented in Table 12. In Table 13, we summarize the abundances for this DLA, for which we derive values of  $[\text{C}/\text{H}] = -2.63 \pm 0.45$ ,  $[\text{O}/\text{H}] = -2.44 \pm 0.45$ , and  $[\text{Si}/\text{H}] = -2.78 \pm 0.13$ , with upper limits for the metallicities of  $[\text{Al}/\text{H}] < -3.17$ , and  $[\text{Fe}/\text{H}] < -2.43$ . This DLA also is observed to have relatively weak absorption lines for both C II and O I, and is one of five DLA systems from our survey where both equivalent widths of these species are less than  $130 \text{ m}\text{\AA}$ , which forms the weak-line subsample used in Figure 17. The observed value of  $[\text{C}/\text{H}] = -2.63$  is the third lowest of our sample.

### 3.2. Comparison of ESI-derived H I Column Densities with SDSS Spectra

Our sample was selected based on fits to the H I profiles within the SDSS spectra where the derived values of  $\log(N(\text{H I})) \geq 20.3$ ,

**Table 7**  
Ionic Column Densities for J1305+2902  $z = 2.386$

Ion	$\lambda$ (Å)	$\log f$	$v_{\text{int}}^{\text{a}}$ (km s <sup>-1</sup> )	$W_{\lambda}^{\text{b}}$ (mÅ)	$\log N$	$\log N_{\text{adopt}}$
C II	1334.5323	-0.8935	[-50, 70]	130.8 ± 15.4	13.80 ± 0.10	14.25 ± 0.45
C II*	1335.7077	-0.9397	[-90, 70]	<35.0	<13.46	
C IV	1548.1950	-0.7194	[-90, 70]	109.8 ± 15.5	<13.49 ± 0.06	<13.52
	1550.7700	-1.0213	[-90, 70]	77.6 ± 15.8	<13.61 ± 0.09	
O I	1302.1685	-1.3110	[-90, 70]	69.9 ± 10.3	13.99 ± 0.06	14.11 ± 0.12
Mg I	2852.9642	0.2577	[-90, 70]	<109.6	<12.11	<12.11
Mg II	2796.3520	-0.2130	[-90, 70]	254.5 ± 54.4	>12.88	>12.88
	2803.5310	-0.5151	[-90, 70]	115.8	>12.46	
Al II	1670.7874	0.2742	[-90, 70]	26.7 ± 5.0	11.82 ± 0.10	11.96 ± 0.10
Si II	1193.2891	-0.3019	[-90, 70]	180.8 ± 20	>13.45	13.32 ± 0.1
	1260.4221	0.0030	[-90, 70]	220.9 ± 21.4	>13.33	
	1304.3702	-1.0269	[-90, 70]	21.5 ± 5	13.37 ± 0.1	
	1526.7066	-0.8962	[-90, 70]	30.7 ± 5	13.26 ± 0.1	
	1808.0130	-2.6603	[-90, 70]	<56.5	<14.97	
Si IV	1393.7550	-0.2774	[-90, 70]	59.5 ± 17.3	12.86 ± 0.12	12.86 ± 0.12
	1402.7700	-0.5817	[-90, 70]	35.7 ± 12.0	12.89 ± 0.12	
Fe II	1608.4511	-1.2366	[-90, 70]	25.5 ± 10.6	13.24 ± 0.13	13.00 ± 0.13
	2344.2140	-0.9431	[-90, 70]	<42.4	>12.99	
	2374.4612	-1.5045	[-90, 70]	<43.7	<13.63	
	2382.7650	-0.4949	[-90, 70]	118.3 ± 21.4	>12.90	
	2600.1729	-0.6216	[-90, 70]	126.8 ± 20.0	12.98 ± 0.13	

**Notes.**

<sup>a</sup> Velocity interval over which the equivalent width and column density are measured.

<sup>b</sup> Rest equivalent width.

but we have refitted the Ly $\alpha$  profiles in the ESI spectra. Indeed, the SDSS estimates of  $\log(N(\text{H I}))$  are uncertain due to the limited resolution of the SDSS spectra, and the possibility of blending of the DLA system with either other DLA absorbers or strong Ly $\alpha$  forest components. Figure 12 shows a comparison of the two values of  $\log(N(\text{H I}))$ , which in most cases shows that our selection procedure identified DLA systems from the SDSS spectra. As a population, we find that the  $N(\text{H I})$  values from SDSS spectra are systematically lower than those derived from the ESI data. We calculate a mean (median) offset of 0.07 (0.1) dex and expect one of two reasons for the differences. First, if there was no metal-line transition identified then one tends to center the Ly $\alpha$  profile on the absorption line. This will yield the highest value possible for the data and therefore gives a systematic overestimate for the  $N(\text{H I})$  value. Second, we suspect there is a subtle, statistical effect related to our pre-selection of DLAs with weak metal lines. In any case, this offset has no significant implication for this paper.

Two of the systems drawn from the SDSS database (SDSS1156+5513 at  $z \approx 2.49$ ; SDSS1327+4845 at  $z \approx 2.61$ ) are now observed to be a pair of strong Ly $\alpha$  lines that were analyzed as a single absorption system by Prochaska et al. (2005). Each of these pair has  $N(\text{H I}) < 10^{20} \text{ cm}^{-2}$  and is plotted as a single point in the figure. It is likely that by focusing on systems with very weak metal-line absorption, we have preferentially identified systems that suffer from this (rare) systematic effect.

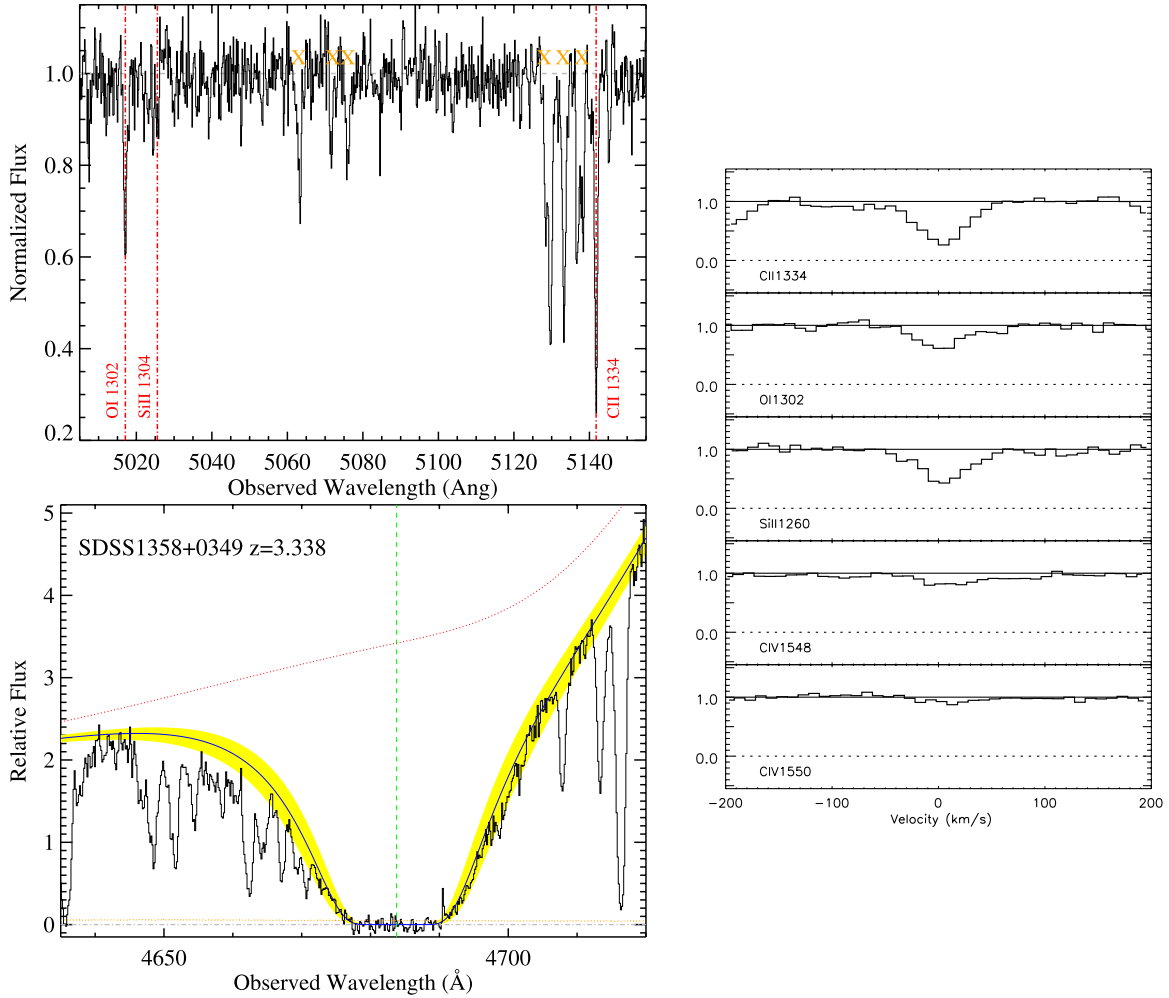
The uncertainty in the SDSS H I determinations also caused a few of our DLA systems to have column densities of  $\log(N(\text{H I})) < 20.3$ , due to the uncertainty of the SDSS H I determination. For this reason, a few of our DLA systems could require small corrections for ionization, but these corrections in most cases were below our systematic uncertainty from the limited resolution of the ESI spectrograph. We discuss some of the implications of the sub-DLA systems in Section 4.5.

#### 4. ELEMENTAL ABUNDANCES AND COMPARISONS WITH LITERATURE AND THEORETICAL MODELS

##### 4.1. Distribution of DLA Metallicities

We present the results of our AOD column density analysis in Table 13 for the entire sample, which shows either measurements or limits for the abundances of the available elements. For each element, we measured the column density of the species, and then adjusted for solar abundances using recently recalibrated values (Lodders 2003), after combining with the measured H I column density. The results in Table 13 then are  $[X/\text{H}]$ , or abundance of element X relative to solar on a logarithmic scale. In some cases, transitions appear to be blended and provide only an upper limit to the column densities. Due to the limited spectral resolution of the ESI spectrograph, we also considered profiles with normalized flux values  $F_{\lambda} < 0.3$  in





**Figure 7.** Section of our spectrum SDSS1358+0349 showing the lines of C II  $\lambda$ 1334, Si II  $\lambda$ 1304, and O I  $\lambda$ 1302 for the DLA system at  $z = 2.8528$  (left), and the H I damped Ly $\alpha$  profile at the same redshift (right), along with our best-fit model (solid line) for  $\log(N(\text{H I})) = 20.50$ , and  $1\sigma$  limits to the H I fit (yellow/shaded region). (A color version of this figure is available in the online journal.)

any pixel to be potentially saturated, and many of these column densities are presented in Table 13 as lower limits.

In most cases, where we have detections of absorption, we can determine values of metallicity to within  $\pm 0.15$  dex, while some of the lower S/N spectra only enable a weaker constraint to within  $\pm 0.25$  dex, and the degree of precision of the abundance is provided with labels within Table 13. Our estimate of the uncertainty is based on comparisons of column densities from multiple transitions of the same species and allows for the combination of uncertainty from continuum fitting, blending, and photon noise within the data.

Figure 13 shows the histogram of observed values of  $[\text{M}/\text{H}]$  from the survey of Prochaska et al. (2003b), along with a Gaussian fit to the distribution. The Gaussian fit seems to match the data well, and shows a mean value of metallicity for systems with  $2 < z < 4.5$  of  $[\text{M}/\text{H}] = -1.52$ , with a value of  $\sigma([\text{X}/\text{H}]) = 0.52$ . Figure 13 also presents the metallicity distribution for the metal-poor DLAs analyzed in this paper, which shows a significantly lower mean metallicity of  $[\text{M}/\text{H}] = -2.2$ . To estimate an  $[\text{M}/\text{H}]$  values, we have adopted the  $[\text{Si}/\text{H}]$  metallicity in those cases where the system has a reported value. For systems with upper/lower limits for  $[\text{Si}/\text{H}]$ , we have adopted the  $[\text{Fe}/\text{H}]$  value (when measured) incremented by 0.4 dex which represents the typical  $[\text{Si}/\text{Fe}]$  offset observed

in DLAs. Systems with only a lower limit to  $[\text{M}/\text{H}]$  are shown separately in the figure with dotted lines. Histograms of our estimated abundances for individual elements are presented in Section 4.2.

The present work sampled a large number of DLA systems (458) in the SDSS DR5 to provide the small number of low-metallicity DLA systems in this study. If we use the fitted Gaussian distribution we would predict that values of  $[\text{M}/\text{H}] < -3.12$  would comprise 0.1% of the available DLA systems at this redshift, based on the expectation values for the fitted Gaussian metallicity distribution. Likewise a  $3\sigma$  departure of metallicity  $[\text{M}/\text{H}]$  for the distribution corresponds to a value of  $[\text{M}/\text{H}] = -3.08$ , and would comprise 0.00135 of the available distribution of DLA systems. Given that our SDSS sample consisted of 458 DLA systems, we would then predict 0.6 DLA systems with  $[\text{M}/\text{H}] < -3.08$ , and 0.45 systems with  $[\text{M}/\text{H}]$ .

Our results are consistent with this distribution, within the small number statistics of our sample, as our lowest metallicity in  $[\text{C}/\text{H}] = -2.85$ , while our lowest metallicity in  $[\text{O}/\text{H}] = -3.17$ ; both lie close to the  $3\sigma$  limit of metallicity for DLA systems. We observe two systems with  $[\text{O}/\text{H}] \leq -3.08$  (the three  $\sigma$  limit for  $[\text{M}/\text{H}]$ ), and no systems with  $[\text{C}/\text{H}] \leq -3.08$ . It is also interesting to note that extremely low-metallicity

**Table 8**  
Ionic Column Densities for J1358+0349  $z = 2.853$

Ion	$\lambda$ (Å)	$\log f$	$v_{\text{int}}^{\text{a}}$ (km s $^{-1}$ )	$W_{\lambda}^{\text{b}}$ (mÅ)	$\log N$	$\log N_{\text{adopt}}$
C II	1334.5323	-0.8935	[-80, 70]	174.0 $\pm$ 9.0	13.93 $\pm$ 0.15	>14.38
C II*	1335.7077	-0.9397	[-80, 70]	19.9 $\pm$ 8.0	13.24 $\pm$ 0.14	
C IV	1548.1950	-0.7194	[-80, 70]	86.6 $\pm$ 8.9	13.33 $\pm$ 0.04	13.25 $\pm$ 0.04
	1550.7700	-1.0213	[-80, 70]	31.5 $\pm$ 8.2	13.19 $\pm$ 0.10	
O I	1302.1685	-1.3110	[-80, 70]	100.5 $\pm$ 10.7	14.13 $\pm$ 0.05	14.38 $\pm$ 0.25
Al II	1670.7874	0.2742	[-80, 70]	57.7 $\pm$ 11.1	12.09 $\pm$ 0.08	12.17 $\pm$ 0.08
Si II	1193.2891	-0.3019	[-80, 70]	86.0 $\pm$ 8.0	13.13 $\pm$ 0.05	13.30 $\pm$ 0.17
	1260.4221	0.0030	[-80, 70]	141.1 $\pm$ 8.7	13.11 $\pm$ 0.05	
	1304.3702	-1.0269	[-40, 70]	<19.0	<13.32	
	1526.7066	-0.8962	[-80, 70]	64.0 $\pm$ 10.2	13.41 $\pm$ 0.05	
	1808.0130	-2.6603	[-80, 70]	<21.4	<14.70	
Si IV	1393.7550	-0.2774	[-80, 70]	76.2 $\pm$ 9.9	12.96 $\pm$ 0.06	12.96 $\pm$ 0.06
Fe II	1608.4511	-1.2366	[-80, 70]	<19.9	<13.35	13.01 $\pm$ 0.05
	2344.2140	-0.9431	[-80, 70]	51.3 $\pm$ 15.6	13.00 $\pm$ 0.13	
	2374.4612	-1.5045	[-80, 70]	19.15 $\pm$ 8.0	13.08 $\pm$ 0.10	
	2382.7650	-0.4949	[-80, 70]	153.1 $\pm$ 17.4	13.02 $\pm$ 0.05	

**Notes.**

<sup>a</sup> Velocity interval over which the equivalent width and column density are measured.

<sup>b</sup> Rest equivalent width.

**Table 9**  
Ionic Column Densities for J1358+6522  $z = 3.067$

Ion	$\lambda$ (Å)	$\log f$	$v_{\text{int}}^{\text{a}}$ (km s $^{-1}$ )	$W_{\lambda}^{\text{b}}$ (mÅ)	$\log N$	$\log N_{\text{adopt}}$
C II	1334.5323	-0.8935	[-70, 80]	122.0 $\pm$ 12.5	13.78 $\pm$ 0.04	14.22 $\pm$ 0.43
C II*	1335.7077	-0.9397	[-50, 30]	<19.6	<13.23	
C IV	1548.1950	-0.7194	[-50, 50]	<20.8	<12.90	<12.90
O I	1302.1685	-1.3110	[-90, 70]	60.0 $\pm$ 12.2	13.95 $\pm$ 0.04	14.03 $\pm$ 0.08
Al II	1670.7874	0.2742	[-50, 70]	55.5 $\pm$ 10.3	12.11 $\pm$ 0.10	12.11 $\pm$ 0.10
Si II	1260.4221	0.0030	[-60, 70]	55.0 $\pm$ 7.6	12.64 $\pm$ 0.10	13.29 $\pm$ 0.25
	1304.3702	-1.0269	[-60, 70]	<23.0	<13.41	
	1526.7066	-0.8962	[-50, 70]	50.9 $\pm$ 5.0	13.29 $\pm$ 0.10	
	1808.0130	-2.6603	[-50, 70]	<29.4	<14.84	
Si IV	1393.7550	-0.2774	[-50, 50]	<20.2	<12.52	<12.52
Fe II	1608.4511	-1.2366	[-50, 70]	<5.3	<12.8	<12.80

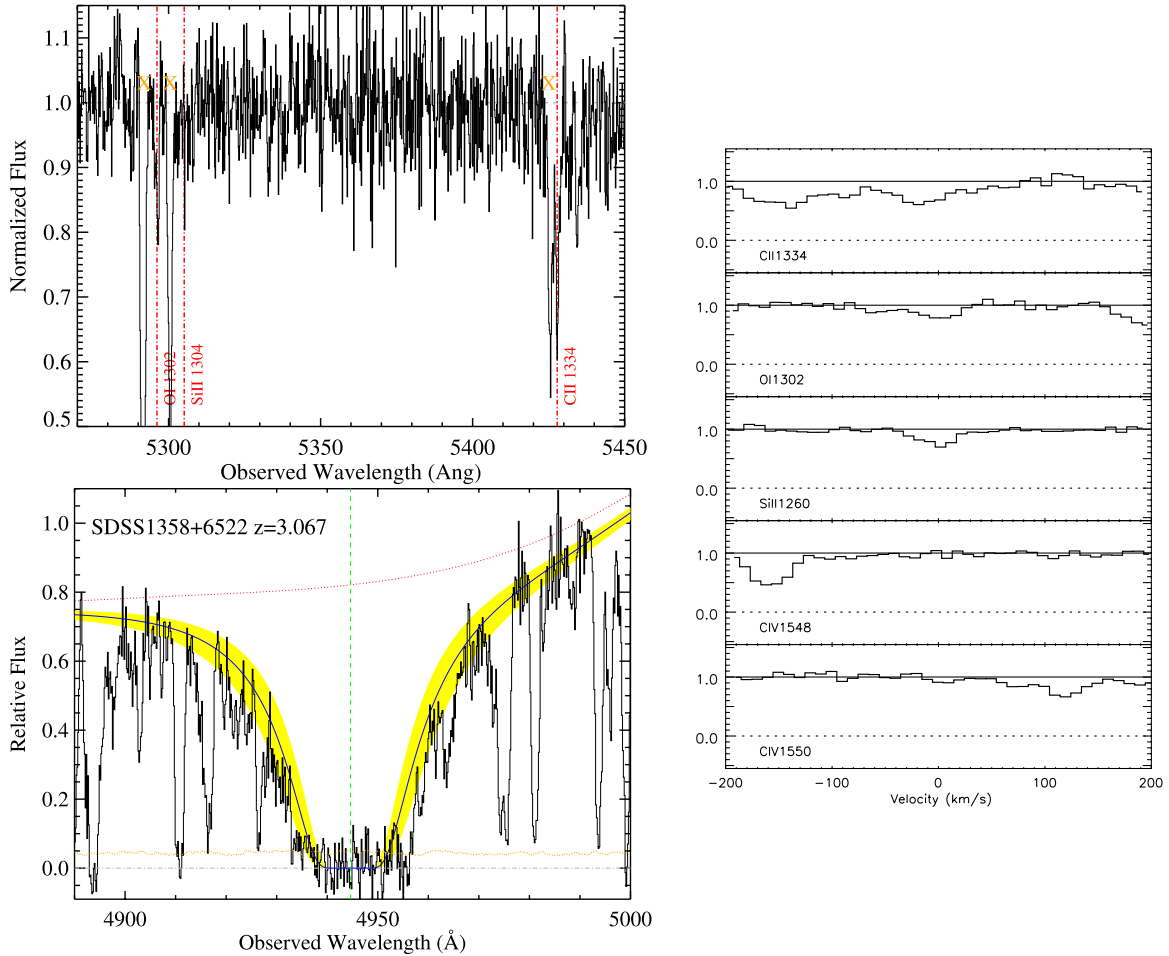
**Notes.**

<sup>a</sup> Velocity interval over which the equivalent width and column density are measured.

<sup>b</sup> Rest equivalent width.

systems with  $[X/H] < -4.0$  are expected to be extremely rare, as they represent a  $4.8\sigma$  deviation in metallicity, and therefore would comprise only  $8 \times 10^{-7}$  of the DLA systems. Put another way, one would need observations of 1.25 million quasar DLA

systems to expect to find one system with  $[M/H] < -4.0$ , making the discovery of such low-metallicity systems highly unlikely, under the assumption that the metallicity distribution obeys a Gaussian distribution.



**Figure 8.** Section of our spectrum SDSS1358+6522 showing the lines of C II  $\lambda$ 1334, Si II  $\lambda$ 1304, and O I  $\lambda$ 1302 for the DLA system at  $z = 3.0674$  (left), and the H I damped  $\text{Ly}\alpha$  profile at the same redshift (right), along with our best-fit model (solid line) for  $\log(N(\text{H I})) = 20.35$ , and  $1\sigma$  limits to the H I fit (yellow/shaded region). (A color version of this figure is available in the online journal.)

#### 4.2. Distribution of Estimated Metallicity for Individual Elements

When we plot the observed elemental abundances for our sample against redshift, we can see that many of the target quasars feature DLA systems of significantly lower metallicity than previously published surveys. Our sample was chosen to have very weak lines of C II  $\lambda$ 1334 and Si II  $\lambda$ 1260 in the SDSS spectra, and the derived abundances from the ESI spectra bear out that the selection process was successful in isolating very low metallicity DLAs.

The top panel of Figure 14 shows the values of  $[\text{C}/\text{H}]$  for our sample, including those sightlines which appear to have only limits to the C abundance, either due to saturation effects (lower limits) or due to a non-detection of the line (upper limits) due to  $\text{Ly}\alpha$  blending. These limits are shown in the figure as arrows. The lower panel of Figure 14 shows the same data points superimposed on the metallicity distribution from Prochaska et al. (2003b), with vertical lines indicating the mean metallicity ( $[\text{M}/\text{H}] = -1.52$ ) and 1, 2, and  $3\sigma$  departures from the mean metallicity.

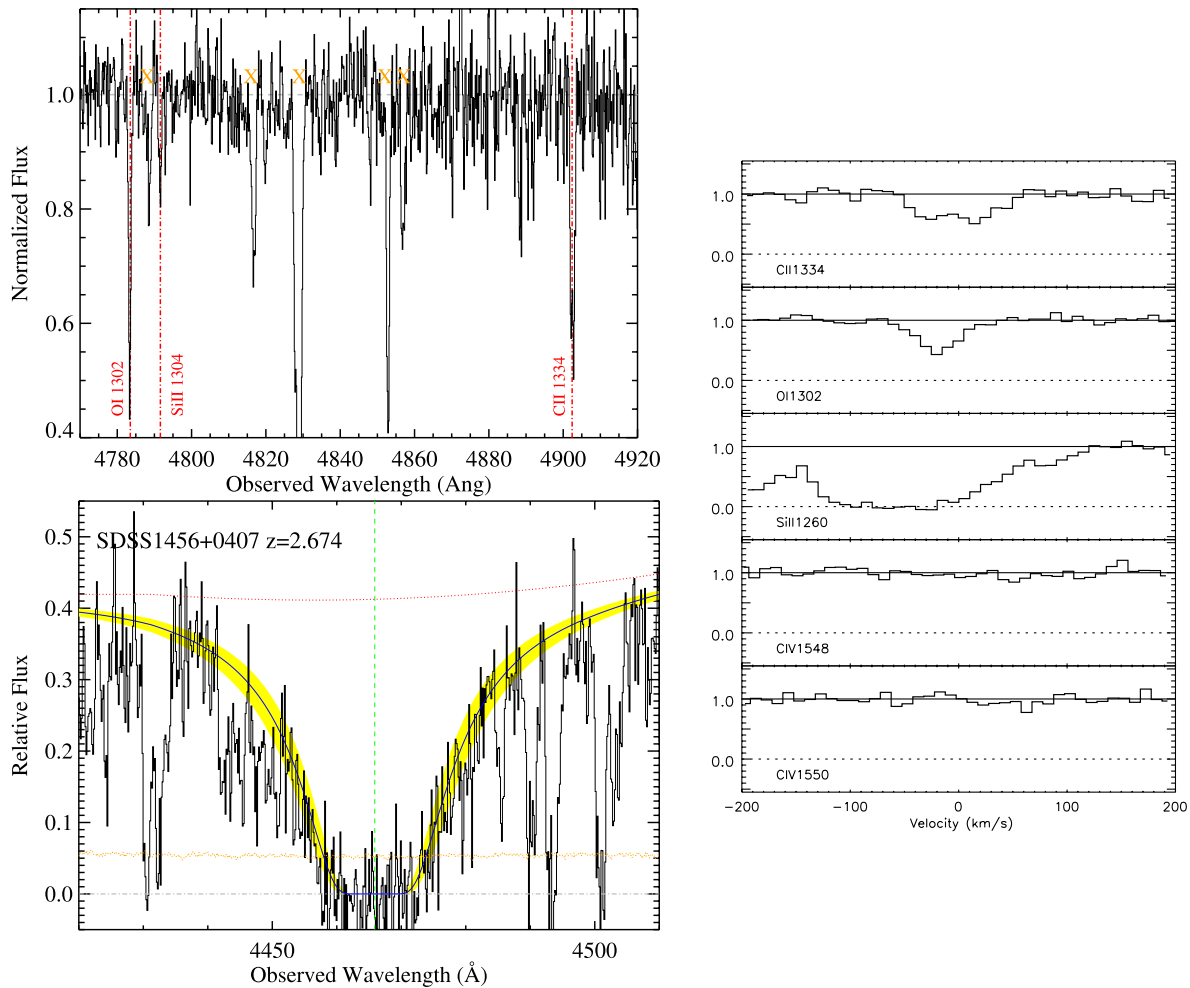
The top panel of Figure 15 compares the results from our survey with the metallicity sample of Prochaska et al. (2003b) for the  $\alpha$ -process elements of O and Si, for those DLA systems with measurements (i.e., excluding upper and lower limits) of the elements O and Si. Our values of  $[\text{O}/\text{H}]$ , and  $[\text{Si}/\text{H}]$

are some of the lowest yet observed, with several values of  $[\alpha/\text{H}] < -2.75$ . The bottom panel of Figure 15 compares these same measurements with the distribution of metallicity within Prochaska et al. (2003b), as in Figure 14.

A similar plot of the metallicity dependence on redshift is shown below in Figure 16, where we present the metallicities of our sample for the elements Fe and Al, against redshift in the top panel, and the comparison of these metallicities against the survey of DLA systems from Prochaska et al. (2003b) in the bottom panel. We discuss the observed abundances of these elements compared with stellar abundance surveys, and in the context of nucleosynthesis models in the following section.

#### 4.3. C and O Abundances and Comparison with Stellar Sample

The lowest abundances of our sample enable a comparison with metal-poor stellar samples, and as such the DLA abundances provide a complementary probe of important nucleosynthetic quantities such as O/C and Si/O yields within the first stars that enrich the DLA with heavy elements. Models of high-mass stellar nucleosynthesis predict small values of  $[\text{C}/\text{O}]$ , due to the prediction of a large oxygen rich core in the final stages of the Population III star (with  $140 M_{\odot} < M_{*} < 260 M_{\odot}$ ), and a much larger yield of O compared to C in massive star explosion, typically with a factor of 10 more O than C produced (Heger & Woosley 2002). However, a great deal of uncertainty exists in reaction rates for the alpha processing of C into O, and other



**Figure 9.** Section of our spectrum SDSS1456+0407 showing the lines of C II  $\lambda$ 1334, Si II  $\lambda$ 1304, and O I  $\lambda$ 1302 for the DLA system at  $z = 2.6736$  (left), and the H I damped Ly $\alpha$  profile at the same redshift (right), along with our best-fit model (solid line) for  $\log(N(\text{H I})) = 20.35$ , and  $1\sigma$  limits to the H I fit (yellow/shaded region). (A color version of this figure is available in the online journal.)

factors such as rotation, mass loss, and convection will have a significant impact on the observed nuclear abundances (e.g., Chieffi & Limongi 2004).

Figure 17 presents the results for [C/O] of our newly measured low-metallicity DLA systems (diamonds), compared to the data of Akerman et al. (2004) (squares are Galactic disk stars; triangles are Galactic halo stars). Our results confirm the upward trend of [C/O] for low metallicity and low [O/H]. We have provided corrections for saturation where relevant to the column densities within Figure 17, as described in the above sections. We also caution that one or more of these values could be compromised if a system has a very low Doppler parameter. We plan to confirm the results presented here with follow-up echelle observations.

In some cases, it is possible that saturation corrections may be a source of the upward trend shown in Figure 17, especially since some of our saturation corrections were adjusted to account for  $b$ -values determined from curve of growth fitting. Recent studies have shown that curve of growth fits of limited resolution spectra for GRB spectra can underestimate column densities, as they occasionally derive larger  $b$ -values (and hence lower column densities) if those spectra contain numerous unresolved components of low  $b$ -value (Prochaska 2006). For those DLA systems with significant corrections for saturation, a lower  $b$ -value would cause us to underestimate the columns

of our absorbing species, especially in the strong lines of C II and O I.

To limit this possibility, we have included in Figure 17 only those DLA systems which have weak absorption in both C II and O I, with equivalent widths less than 130 mÅ in both species. Within this subsample, we have estimated the saturation corrections, using corrections to the weak-line AOD column densities, appropriate to  $b$ -values of 7.5 km s<sup>-1</sup>, as described in Section 2.3. These corrections are conservative, and in most cases, the weak lines result in small corrections of only 0.1 dex or less. We also note that corrections for saturation would likely increase the C/O ratio because the C column density derived from the C II  $\lambda$ 1334 transition is more likely to be underestimated than the O column density from the O I  $\lambda$ 1302 transition, due to its larger oscillator strength. The enhancement of [C/O] is also strongest within the DLA systems with the weakest C II and O I lines, suggesting that the trend is not a result of saturation effects in our spectra.

#### 4.4. Fe Abundances Compared with C, O, and Si Abundances

We have also examined trends in the ratio of metal abundances against the [Si/H] and [Fe/H] metallicity, to see if systematic trends in elemental abundances exist at lower metallicities. Figure 18 shows the trend of [C/Si] and [O/Si], measured



**Table 10**  
Ionic Column Densities for J1456+0407  $z = 2.674$

Ion	$\lambda$ (Å)	$\log f$	$v_{\text{int}}^{\text{a}}$ ( $\text{km s}^{-1}$ )	$W_{\lambda}^{\text{b}}$ (mÅ)	$\log N$	$\log N_{\text{adopt}}$
C II	1334.5323	-0.8935	[-80, 60]	$127.0 \pm 14.4$	$13.88 \pm 0.04$	>14.33
C II*	1335.7077	-0.9397	[-80, 60]	<32.0	<13.42	
C IV	1548.1950	-0.7194	[-80, 60]	<35.5	<13.14	<13.14
	1550.7700	-1.0213	[-80, 60]	<36.4	<13.45	
O I	1302.1685	-1.3110	[-80, 60]	$100.3 \pm 8.2$	$14.13 \pm 0.03$	$14.55 \pm 0.28$
Al II	1670.7874	0.2742	[-80, 60]	$87.0 \pm 10.2$	$12.27 \pm 0.05$	$12.93 \pm 0.45$
Si II	1304.3702	-1.0269	[-80, 60]	$32.6 \pm 8.9$	$13.40 \pm 0.12$	$13.49 \pm 0.07$
	1526.7066	-0.8962	[-80, 60]	$84.4 \pm 16.7$	$13.58 \pm 0.08$	
	1808.0130	-2.6603	[-80, 60]	<44.6	<15.02	
Si IV	1393.7550	-0.2774	[-80, 60]	$83.9 \pm 15.8$	$13.01 \pm 0.09$	$13.01 \pm 0.09$
Fe II	1608.4511	-1.2366	[-80, 60]	$30.1 \pm 8.00$	$13.35 \pm 0.10$	$13.00 \pm 0.10$
	2344.2140	-0.9431	[-80, 60]	$41.2 \pm 10.00$	$12.86 \pm 0.10$	
	2374.4612	-1.5045	[-80, 60]	47.4	<13.68	
	2382.7650	-0.4949	[-80, 60]	$167.0 \pm 13.2$	$13.01 \pm 0.06$	

**Notes.**

<sup>a</sup> Velocity interval over which the equivalent width and column density are measured.

<sup>b</sup> Rest equivalent width.

**Table 11**  
Ionic Column Densities for J1637+2901  $z = 3.496$

Ion	$\lambda$ (Å)	$\log f$	$v_{\text{int}}^{\text{a}}$ ( $\text{km s}^{-1}$ )	$W_{\lambda}^{\text{b}}$ (mÅ)	$\log N$	$\log N_{\text{adopt}}$
C II	1334.5323	-0.8935	[-70, 80]	$147.0 \pm 12.4$	>13.86	>14.30
C II	1335.7077	-0.9397	[-70, 80]	$274.8 \pm 11.7$	<14.35	
C IV	1548.1950	-0.7194	[-70, 150]	$130.3 \pm 14.0$	$13.55 \pm 0.05$	$13.55 \pm 0.05$
O I	1302.1685	-1.3110	[-70, 80]	$91.7 \pm 10.6$	$14.09 \pm 0.05$	$14.29 \pm 0.20$
Al II	1670.7874	0.2742	[-70, 80]	$71.0 \pm 12.2$	$12.18 \pm 0.07$	$12.29 \pm 0.11$
Si II	1260.4221	0.0030	[-70, 80]	$132.7 \pm 5.0$	>13.07	$13.41 \pm 0.14$
	1304.3702	-1.0269	[-70, 80]	$42.3 \pm 11.0$	$13.54 \pm 0.11$	
	1526.7066	-0.8962	[-70, 80]	<32.3	$13.30 \pm 0.14$	
	1808.0130	-2.6603	[-70, 80]	<56.4	<14.98	
Si IV	1393.7550	-0.2774	[-70, 80]	$40.7 \pm 13.1$	$12.69 \pm 0.14$	$12.69 \pm 0.14$
	1402.7700	-0.5817	[-70, 80]	<26.3	<12.95	
Fe II	1608.4511	-1.2366	[-70, 80]	$67.8 \pm 9.7$	$13.74 \pm 0.06$	$13.84 \pm 0.10$

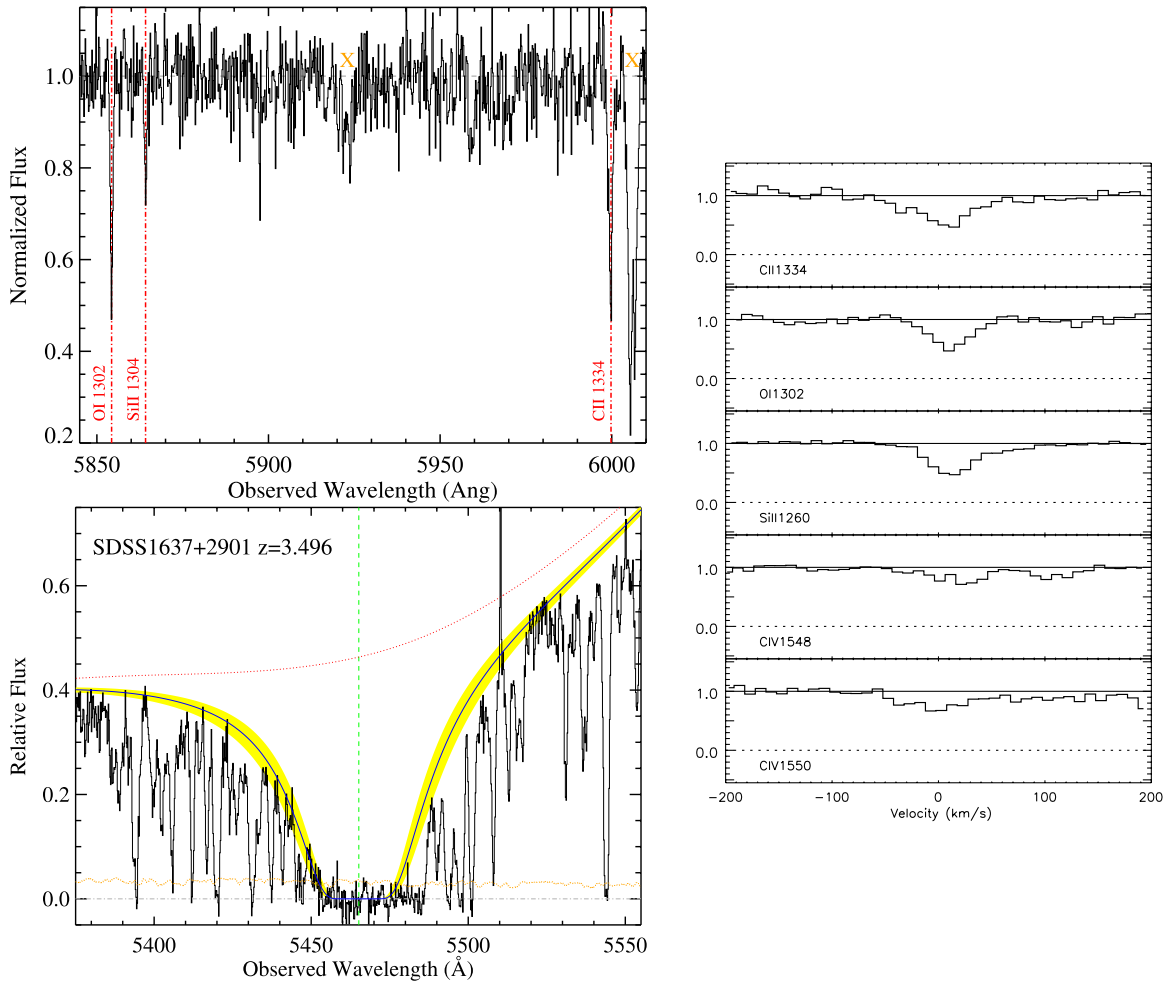
**Notes.**

<sup>a</sup> Velocity interval over which the equivalent width and column density are measured.

<sup>b</sup> Rest equivalent width.

against the [Si/H] metallicity, while Figure 19 shows the element ratios [C/Fe], [O/Fe], and [Si/Fe]. In both cases, we see trends of increasing [C/X] and [O/X] for the sample with decreased metallicity. This suggests an upward trend of  $\alpha$ /Fe, which has implications for the nucleosynthesis which produced the low metal abundances within our DLA sample, perhaps

due to a mix of more high-mass stars within DLA systems which would be expected to produce large amounts of  $\alpha$ -process elements with the lowest values of [X/H]. One might argue that some of the reduced values of  $\alpha$ /Fe within the sample at higher metallicities could arise from depletion effects, and higher spectral resolution observations which include highly



**Figure 10.** Section of our spectrum SDSS1637+2901 showing the lines of C II  $\lambda$ 1334, Si II  $\lambda$ 1304, and O I  $\lambda$ 1302 for the DLA system at  $z = 3.4956$  (left), and the H I damped Ly $\alpha$  profile at the same redshift (right), along with our best-fit model (solid line) for  $\log(N(\text{H I})) = 20.70$ , and  $1\sigma$  limits to the H I fit (yellow/shaded region). (A color version of this figure is available in the online journal.)

depleted species such as Cr are needed to help elucidate this effect. However, since the depletion of Fe relative to Si increases with dust content, with  $0.5 < [\text{Si}/\text{Fe}] < 1.0$  depending on whether the IGM is “halo” or “disk” type interstellar medium (ISM; Savage & Sembach 1996), one would expect the values of  $[\text{Si}/\text{Fe}]$  and  $\alpha/\text{Fe}$  to increase with increasing metallicity, which is the opposite of what is observed. It therefore seems more likely that our trend of increasing values of  $\alpha/\text{Fe}$  for lower metallicity DLA systems is a result of nucleosynthetic effects.

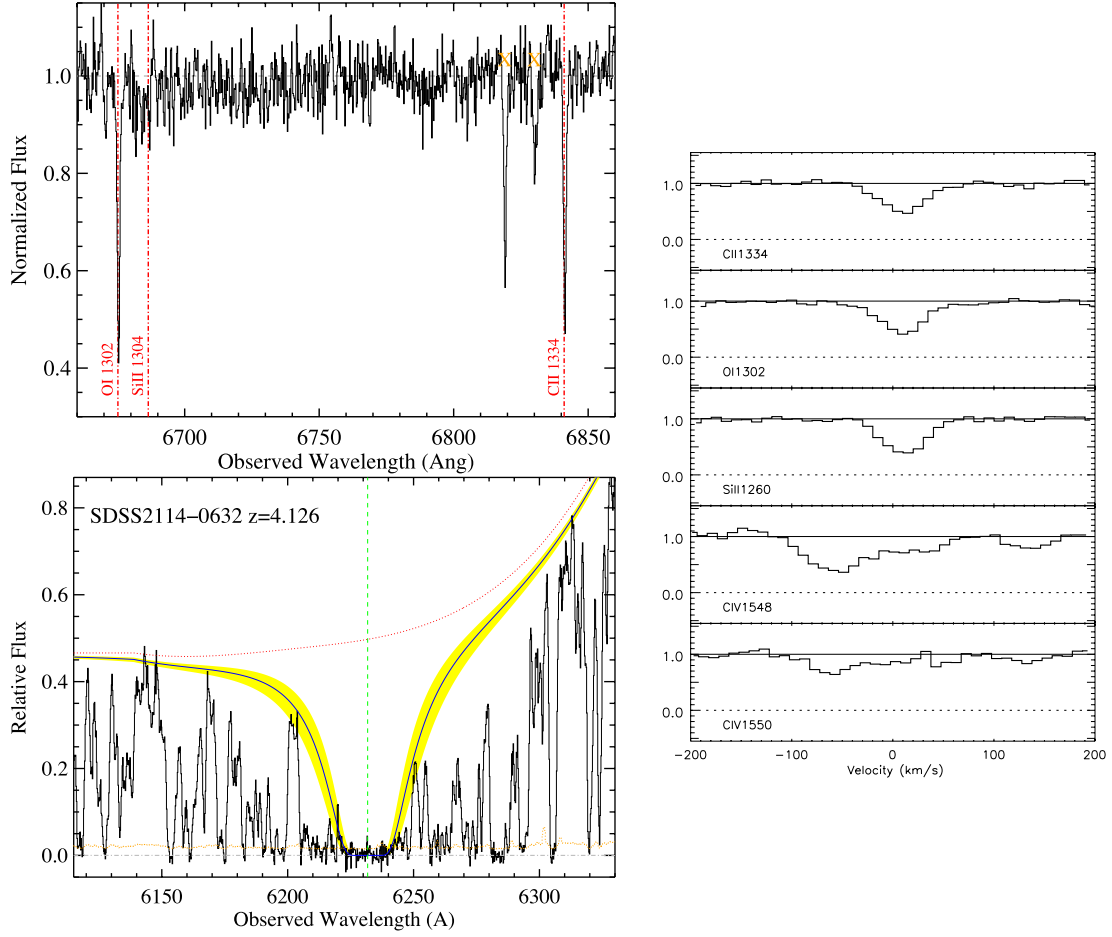
#### 4.5. Comparison of Elemental Abundances with Nucleosynthesis Models

Our results help provide constraints on the nucleosynthetic yields of high-mass stars, since the low values of  $[\text{O}/\text{H}]$  observed in our sample suggest that they were enriched by only one to a few supernovae, depending on the value of  $[\text{O}/\text{H}]$ . The higher yield of  $[\text{C}/\text{O}]$  is not predicted by nucleosynthesis models of Heger & Woosley (2002), and Umeda & Nomoto (2002), but is consistent with the predictions of Chieffi & Limongi (2002, 2004) for stars with masses ranging from  $20 M_{\odot} < M_{*} < 80 M_{\odot}$ . The higher yield of  $[\text{C}/\text{O}]$  is also consistent with those seen in emission line studies of H II regions (Akerman et al. 2004), and is within the range of values of  $[\text{C}/\text{O}]$  estimated for the Ly $\alpha$  forest, where near solar values of  $[\text{C}/\text{O}]$  are predicted using the

most realistic ionizing radiation models with the overdensity  $\delta < 10$  (Aguirre et al. 2008).

We also present a plot in Figure 20 of our elemental abundances against atomic number, for a subset of four of our highest quality low-metallicity DLA systems. We observe that compared to the models of Heger and Woosley, we observe a reduced abundance of O, Si, and Fe compared to the model, and an overabundance of Al. This result also suggests that the ratio of  $[\alpha/\text{Fe}]$  in low-metallicity DLA systems appears closer to a value of  $[\alpha/\text{Fe}] = 0$ , instead of higher values at lower metallicity, as some models have suggested (e.g., Tinsley, 1979). The mean underabundance of the elements O, Si, and Fe for our sample is  $-0.72$  dex relative to C when compared to the models of Heger & Woosley (2002), while the Al abundance we observe is  $+0.52$  dex above the model.

Figure 21 summarizes the observed ratios of  $[\text{C}/\text{O}]$  compared with various nucleosynthesis models as a function of redshift. The horizontal lines indicate the predicted yields from the models of Heger & Woosley (2008), Chieffi & Limongi (2004), and Heger & Woosley (2002). Results from a survey of high-redshift DLA systems by Becker et al. (2006), which may in some ways be similar to our low-metallicity sample, are also indicated on the plot. The results show that the largest values of  $[\text{C}/\text{O}]$  occur for the redshift 2.4 DLA system, and that the higher redshift DLA systems within our sample appear to converge



**Figure 11.** Section of our spectrum SDSS2114–0632 showing the lines of C II  $\lambda$ 1334, Si II  $\lambda$ 1304, and O I  $\lambda$ 1302 for the DLA system at  $z = 4.126$  (left), and the H I damped  $\text{Ly}\alpha$  profile at the same redshift (right), along with our best-fit model (solid line) for  $\log(N(\text{H I})) = 20.40$ , and  $1\sigma$  limits to the H I fit (yellow/shaded region). (A color version of this figure is available in the online journal.)

**Table 12**  
Ionic Column Densities for J2114–0632  $z = 4.126$

Ion	$\lambda$ ( $\text{\AA}$ )	$\log f$	$v_{\text{int}}^{\text{a}}$ ( $\text{km s}^{-1}$ )	$W_{\lambda}^{\text{b}}$ ( $\text{m\AA}$ )	$\log N$	$\log N_{\text{adopt}}$
C II	1334.5323	-0.8935	[-70, 80]	$118.4 \pm 8.0$	$13.80 \pm 0.03$	$14.25 \pm 0.45$
C II*	1335.7077	-0.9397	[-70, 80]	<16.1	<13.13	
C IV	1548.1950	-0.7194	[-120, 80]	$337.7 \pm 12.8$	$14.05 \pm 0.02$	$14.04 \pm 0.02$
	1550.7700	-1.0213	[-120, 80]	$192.3 \pm 14.6$	$14.04 \pm 0.04$	
O I	1302.1685	-1.3110	[-70, 80]	$138.5 \pm 8.6$	$14.27 \pm 0.03$	$14.72 \pm 0.45$
Al II	1670.7874	0.2742	[-70, 80]	<20.2	<11.83	<11.83
Si II	1260.4221	0.0030	[-70, 80]	$135.9 \pm 5.4$	>12.98	$13.23 \pm 0.13$
	1304.3702	-1.0269	[-70, 80]	<18.5	<13.31	
	1526.7066	-0.8962	[-70, 80]	$39.0 \pm 9.7$	$13.17 \pm 0.13$	
	1808.0130	-2.6603	[-70, 80]	<36.4	<14.93	
Si IV	1393.7550	-0.2774	[-100, 80]	$57.5 \pm 10.7$	$12.84 \pm 0.08$	$12.90 \pm 0.06$
	1402.7700	-0.5817	[-100, 80]	$48.8 \pm 10.4$	$13.06 \pm 0.09$	
Fe II	1608.4511	-1.2366	[-40, 80]	<39.6	<13.51	<13.51

**Notes.**

<sup>a</sup> Velocity interval over which the equivalent width and column density are measured.

<sup>b</sup> Rest equivalent width.

**Table 13**  
Elemental Abundances for DLA Systems

Quasar	$z_{\text{DLA}}$	$\log(N(\text{H I}))$ ( $\text{cm}^{-2}$ )	[C/H]	[O/H]	[Si/H]	[Al/H]	[Fe/H]	[Mg/H]
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)
SDSS0751+4516	3.0462	20.35 ± 0.15	> -2.96 <sup>a</sup>	< -2.10 <sup>a</sup>	> -2.37 <sup>a</sup>	> -2.92 <sup>a</sup>	> -2.09 <sup>a</sup>	...
SDSS0759+3129	3.0346	20.60 ± 0.10	> -2.88 <sup>b</sup>	> -2.73 <sup>b</sup>	-2.03 <sup>c</sup>	-2.63 <sup>c</sup>	-2.27 <sup>a</sup>	...
SDSS0811+2838	2.4338	20.50 ± 0.10	> -3.00 <sup>b</sup>	...	...	> -2.92 <sup>b</sup>	> -2.84 <sup>b</sup>	...
SDSS0814+5029	3.7079	21.25 ± 0.15	> -3.46 <sup>a</sup>	> -3.50 <sup>a</sup>	> -2.92 <sup>a</sup>	> -3.43 <sup>a</sup>	> -2.95 <sup>a</sup>	...
SDSS0831+3358	2.3039	20.25 ± 0.15	> -2.72 <sup>a</sup>	...	> -2.76 <sup>a</sup>	> -3.15 <sup>a</sup>	> -2.59 <sup>a</sup>	...
SDSS0844+4624	3.3363	20.70 ± 0.15	< -2.37 <sup>a</sup>	...	> -2.21 <sup>a</sup>	...	> -2.24 <sup>a</sup>	...
SDSS0924+0951	3.3382	20.85 ± 0.10	> -2.65 <sup>c</sup>	> -2.71 <sup>a</sup>	> -1.76 <sup>c</sup>	> -2.19 <sup>c</sup>	> -2.06 <sup>c</sup>	...
SDSS0928+0939	2.9098	20.75 ± 0.15	-2.41 <sup>a</sup>	> -2.78 <sup>a</sup>	> -2.16 <sup>a</sup>	> -2.69 <sup>a</sup>	-2.13 <sup>c</sup>	...
SDSS0955+4116	3.2801	20.10 ± 0.10	-2.82 <sup>a</sup>	-2.82 <sup>a</sup>	-2.75 <sup>a</sup>	< -2.74 <sup>b</sup>	< -2.30 <sup>b</sup>	...
SDSS1001+0343	3.0785	20.15 ± 0.10	-2.85 <sup>a</sup>	-2.93 <sup>a</sup>	-2.94 <sup>a</sup>	-2.82 <sup>a</sup>	< -2.32 <sup>b</sup>	...
SDSS1003+5520	2.5024	20.35 ± 0.15	< -2.59 <sup>a</sup>	< -2.21 <sup>a</sup>	-2.13 <sup>c</sup>	-2.76 <sup>a</sup>	-2.87 <sup>c</sup>	> -2.67 <sup>a</sup>
SDSS1031+4055	2.5686	20.55 ± 0.10	< -2.58 <sup>b</sup>	< -2.21 <sup>b</sup>	< -1.66 <sup>b</sup>	> -2.37 <sup>b</sup>	-2.18 <sup>a</sup>	> -2.51 <sup>b</sup>
SDSS1037+0139	2.7050	20.40 ± 0.25	> -2.31 <sup>c</sup>	> -2.71 <sup>c</sup>	> -2.25 <sup>c</sup>	> -2.77 <sup>c</sup>	> -2.54 <sup>c</sup>	...
SDSS1043+6151	2.7865	20.60 ± 0.15	> -2.71 <sup>a</sup>	> -2.46 <sup>a</sup>	-2.08 <sup>c</sup>	> -2.46 <sup>a</sup>	-2.04 <sup>a</sup>	...
SDSS1048+3911	2.2957	20.70 ± 0.10	-2.79 <sup>b</sup>	> -2.93 <sup>c</sup>	-2.31 <sup>c</sup>	-2.49 <sup>c</sup>	-2.46 <sup>a</sup>	> -2.82 <sup>b</sup>
SDSS1108+1209	3.3964	20.55 ± 0.15	> -2.86 <sup>a</sup>	< -2.61 <sup>a</sup>	> -2.46 <sup>a</sup>	> -3.18 <sup>a</sup>	< -2.29 <sup>a</sup>	...
SDSS1219+1603	3.0037	20.35 ± 0.10	> -2.59 <sup>b</sup>	-2.59 <sup>c</sup>	-2.08 <sup>a</sup>	> -2.29 <sup>b</sup>	-2.09 <sup>a</sup>	...
SDSS1251+4120	2.7296	21.10 ± 0.10	> -3.06 <sup>b</sup>	> -2.92 <sup>b</sup>	-2.73 <sup>c</sup>	-2.86 <sup>c</sup>	-2.35 <sup>c</sup>	...
SDSS1305+2902	2.3865	20.25 ± 0.10	-2.46 <sup>c</sup>	-2.90 <sup>a</sup>	-2.54 <sup>a</sup>	-2.83 <sup>a</sup>	-2.79 <sup>a</sup>	> -2.99 <sup>b</sup>
SDSS1325+1255	3.5507	20.50 ± 0.15	-2.54 <sup>c</sup>	> -2.39 <sup>c</sup>	-2.52 <sup>a</sup>	< -2.06 <sup>a</sup>	< -2.27 <sup>a</sup>	...
SDSS1350+5952	2.7558	20.65 ± 0.10	> -2.77 <sup>b</sup>	> -2.53 <sup>b</sup>	> -2.48 <sup>b</sup>	-2.53 <sup>c</sup>	-2.59 <sup>a</sup>	...
SDSS1358+0349	2.8528	20.50 ± 0.10	> -2.58 <sup>c</sup>	-2.88 <sup>c</sup>	-2.81 <sup>a</sup>	-2.87 <sup>a</sup>	-3.03 <sup>b</sup>	...
SDSS1358+6522	3.0675	20.35 ± 0.15	-2.59 <sup>c</sup>	-3.08 <sup>a</sup>	-2.67 <sup>c</sup>	-2.78 <sup>a</sup>	< -3.09 <sup>a</sup>	...
SDSS1402+5909	3.7745	21.35 ± 0.10	> -3.37 <sup>c</sup>	> -3.17 <sup>c</sup>	> -2.46 <sup>c</sup>	> -2.86 <sup>c</sup>	> -2.34 <sup>c</sup>	...
SDSS1440+0637	2.5177	21.00 ± 0.15	...	...	-2.37 <sup>c</sup>	> -2.69 <sup>c</sup>	-1.99 <sup>c</sup>	> -3.22 <sup>a</sup>
SDSS1440+0637	2.8246	20.20 ± 0.10	< -2.43 <sup>b</sup>	< -1.96 <sup>b</sup>	-2.14 <sup>b</sup>	-2.43 <sup>a</sup>	-2.21 <sup>a</sup>	...
SDSS1456+0407	2.3201	20.15 ± 0.10	< -2.01 <sup>b</sup>	...	> -1.86 <sup>b</sup>	-1.96 <sup>c</sup>	> -2.17 <sup>a</sup>	> -2.37 <sup>b</sup>
SDSS1456+0407	2.6736	20.35 ± 0.10	> -2.48 <sup>c</sup>	-2.56 <sup>c</sup>	-2.47 <sup>a</sup>	-1.96 <sup>c</sup>	-2.89 <sup>a</sup>	...
SDSS1557+2320	3.5383	20.65 ± 0.10	< -2.81 <sup>b</sup>	-2.21 <sup>b</sup>	-2.14 <sup>c</sup>	-2.68 <sup>c</sup>	-2.61 <sup>c</sup>	...
SDSS1637+2901	3.4956	20.70 ± 0.10	> -2.85 <sup>b</sup>	-3.17 <sup>c</sup>	-2.90 <sup>a</sup>	-2.95 <sup>a</sup>	-2.40 <sup>a</sup>	...
SDSS1654+3509	2.8113	20.10 ± 0.10	< -1.70 <sup>b</sup>	> -2.22 <sup>b</sup>	-1.74 <sup>c</sup>	-1.53 <sup>b</sup>	-2.01 <sup>b</sup>	...
SDSS1709+3417	3.0104	20.40 ± 0.10	> -2.80 <sup>b</sup>	...	> -1.91 <sup>b</sup>	-2.22 <sup>c</sup>	-1.99 <sup>a</sup>	...
SDSS1709+3417	2.5303	20.45 ± 0.15	-2.01 <sup>a</sup>	...	-1.51 <sup>a</sup>	-1.87 <sup>a</sup>	-1.65 <sup>a</sup>	> -2.27 <sup>a</sup>
SDSS1717+5802	3.0461	20.25 ± 0.10	> -2.39 <sup>a</sup>	> -2.26 <sup>b</sup>	-2.05 <sup>a</sup>	> -2.29 <sup>b</sup>	-2.37 <sup>b</sup>	...
SDSS2114-0632	4.1262	20.40 ± 0.15	-2.63 <sup>c</sup>	-2.44 <sup>c</sup>	-2.78 <sup>a</sup>	< -3.17 <sup>a</sup>	< -2.43 <sup>a</sup>	...

**Notes.**

<sup>a</sup> Metallicity uncertainty between 0.11 and 0.20 dex.

<sup>b</sup> Uncertainty less than 0.11 dex.

<sup>c</sup> Metallicity uncertainty greater than 0.20 dex.

with the value predicted by most nucleosynthesis models of  $-0.18 < [\text{C}/\text{O}] < -0.30$ , which also is close to the mean value of  $[\text{C}/\text{O}] = -0.30$  for the high-redshift DLA systems observed by Becker et al. (2006). The data are consistent with a trend of increasing  $[\text{C}/\text{O}]$  with decreasing redshift, which could indicate the presence of intermediate-mass stars  $30 M_{\odot} < M_{*} < 60 M_{\odot}$  providing additional enrichment of carbon at later times compared to the more massive stars for which these nucleosynthesis models were calculated.

To help determine whether the observed enhancement of  $[\text{C}/\text{O}]$  could result from ionization effects, we have plotted the observed ratio of  $[\text{C}/\text{O}]$  against  $\text{H I}$  column in Figure 22. The lowest column densities of  $\text{C II}$  and  $\text{O I}$  were observed toward systems which include sub-DLAs, specifically systems with values of  $\text{H I}$  ranging from  $20.1 < N(\text{H I}) < 20.4$ . For these systems, it is possible that a non-zero ionization correction is necessary for these species. Figure 22 does not indicate a trend in  $[\text{C}/\text{O}]$  with decreasing column of  $\text{H I}$ . Ionization corrections for sub-DLAs for  $20.1 < \log(N(\text{H I})) < 20.4$  based on CLOUDY

models are expected to be in the range of  $0.1 < \delta[\text{C}/\text{O}] < 0.3$  and produce an enhancement of  $\text{C II}$  relative to  $\text{O I}$  due to the lower ionization potential of  $\text{O I}$  compared to  $\text{C II}$ . This effect could help explain some of the super-solar values of  $[\text{C}/\text{O}]$  in our sample, but cannot explain all of the observed  $[\text{C}/\text{O}]$  enhancement at low  $[\text{O}/\text{H}]$ .

We also compare the observed ratios of  $[\text{O}/\text{Si}]$  with nucleosynthesis models as a function of redshift in Figure 23. The observed ratios of  $[\text{O}/\text{Si}]$  appear to not be correlated with redshift, and agree well with nucleosynthesis models, and with the high-redshift DLA systems observed by Becker et al. (2006).

#### 4.6. C IV and Si IV Absorption within the DLA Systems

The presence of  $\text{C IV}$  and  $\text{Si IV}$  absorption was detected in most of our DLA systems, but not all of them. These latter examples are among the very few DLAs with undetected high-ion absorption (Lu et al. 1996; Wolfe & Prochaska 2000; Fox et al. 2007). Table 14 shows the derived column densities of these two species, for the entire DLA sample. The fact that several



**Table 14**  
Column Densities for the Species C IV and Si IV

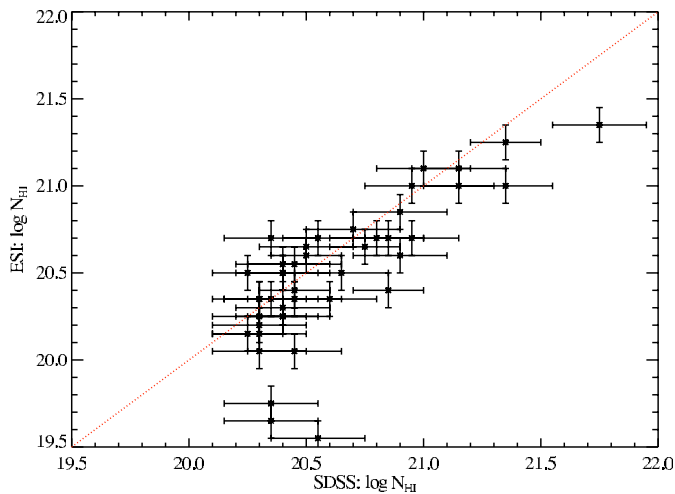
Quasar	$z_{\text{DLA}}$	$\log N(\text{H I})$	$\log N(\text{C IV})$	Err	$\log N(\text{Si IV})$	Err
SDSS0751+4516	3.0462	20.35	13.70	0.10	12.71	0.10
SDSS0759+3129	3.0346	20.6	<12.9	...	<12.0	...
SDSS0811+2838	2.4338	20.50	13.90	0.10	12.40	0.10
SDSS0814+5029	3.7079	21.25	<12.4	...	<12.0	...
SDSS0831+3358	2.3039	20.25	<12.6	...	<12.3	...
SDSS0844+4624	3.3363	20.70	13.31	0.10	12.30	0.10
SDSS0924+0951	3.3382	20.9	12.95	0.043	<12.5	...
SDSS0928+0939	2.9098	20.8	13.60	0.043	<13.8	...
SDSS0955+4116	3.2801	20.1	<13.1	...	13.0	0.043
SDSS1001+0343	3.0785	20.1	13.24	0.043	<12.0	...
SDSS1003+5520	2.5024	20.4	<12.9	...	<13.6	...
SDSS1031+4055	2.5686	20.5	13.20	0.038	13.0	0.043
SDSS1037+0139	2.7050	20.40	13.38	0.10	<12.7	...
SDSS1043+6151	2.7865	20.6	13.74	0.043	13.40	0.039
SDSS1048+3911	2.2957	20.7	13.70	0.042	13.20	0.031
SDSS1108+1209	3.3964	20.55	13.73	0.010	13.50	0.010
SDSS1156+5513	2.4808	19.6	14.28	0.04	...	...
SDSS1156+5513	2.4975	19.6	<13.30	...	...	...
SDSS1219+1603	3.0037	20.4	14.05	0.038	13.50	0.042
SDSS1251+4120	2.7296	21.1	13.93	0.039	13.20	0.042
SDSS1305+2902	2.3865	20.3	<13.4	...	<12.9	...
SDSS1325+1255	3.5507	20.5	<12.9	...	<12.5	...
SDSS1327+4845	2.4468	20.61	13.85	0.040	...	...
SDSS1327+4845	2.6116	19.60	<12.70	...	...	...
SDSS1331+4838	3.6923	21.07	13.90	0.10	...	...
SDSS1349+1242	3.7433	20.00	<13.00	...	...	...
SDSS1350+5952	2.7558	20.6	<13.1	...	<12.5	...
SDSS1358+0349	2.8528	20.5	13.20	0.043	12.80	0.030
SDSS1358+6522	3.0675	20.4	<12.8	...	<12.5	...
SDSS1402+5909	3.7745	21.4	<13.4	...	13.1	0.044
SDSS1440+0637	2.5177	21.0	13.40	0.041	13.60	0.044
SDSS1440+0637	2.8246	20.2	<13.0	...	13.90	...
SDSS1456+0407	2.3201	20.1	<13.2	...	<13.7	...
SDSS1456+0407	2.6736	20.3	<12.6	...	13.10	0.035
SDSS1557+2320	3.5398	20.6	<12.9	...	12.90	0.041
SDSS1637+2901	3.4956	20.7	<13.8	...	11.80	0.043
SDSS1654+3509	2.8113	20.1	14.80	0.041	14.10	0.037
SDSS1709+3417	3.0104	20.4	13.40	0.037	12.80	0.043
SDSS1709+3417	2.5303	20.5	14.80	0.057	14.50	0.041
SDSS1717+5802	3.0461	20.3	14.30	0.041	13.50	0.039
SDSS2114-0632	4.1262	20.4	14.00	0.040	12.90	0.040

of the DLA systems have no detected C IV or Si IV suggests that the low-metallicity DLA systems could be sampling a more neutral region or regions with lower ionizing flux than a typical DLA, where these species are nearly always detected. We also observe within our sample an increasing trend in both  $N(\text{C IV})$  and  $N(\text{Si IV})$  with metallicity (as determined by  $[\text{Si}/\text{H}]$ ), which is consistent with the observations of larger DLA samples (Fox et al. 2007). The equivalent widths of C IV and Si IV could be indicative of a larger velocity dispersion as well as the column densities of these species, due to saturation effects. If so, then the trend in these species with  $[\text{X}/\text{H}]$  is consistent with the observation of larger velocity dispersion and velocity range with increasing metallicity, suggestive of an increasing mass of the DLA system with metallicity (Wolfe & Prochaska 1998; Ledoux et al. 2006). Figure 24 shows the trend of  $N(\text{C IV})$  with increasing metallicity, and further observations at higher resolution can help constrain whether this correlation is caused by increasing velocity dispersion within the higher metallicity DLA systems.

We also have derived estimates of  $N(\text{H II})/N(\text{H I})$  for the low-metallicity DLA sample, using the column densities of C II and C IV and the prescription of Fox et al. 2007, which assumes  $f(\text{C IV}) = 0.3$ . For the two systems which have detected C IV and unsaturated C II, we derive values of  $N(\text{H II})/N(\text{H I})$  in the range  $-0.5 < \log(N(\text{H II})/N(\text{H I})) < -0.1$ , which is slightly higher than the mean value of  $N(\text{H II})/N(\text{H I})$  for the larger DLA sample of Fox et al. (2007).

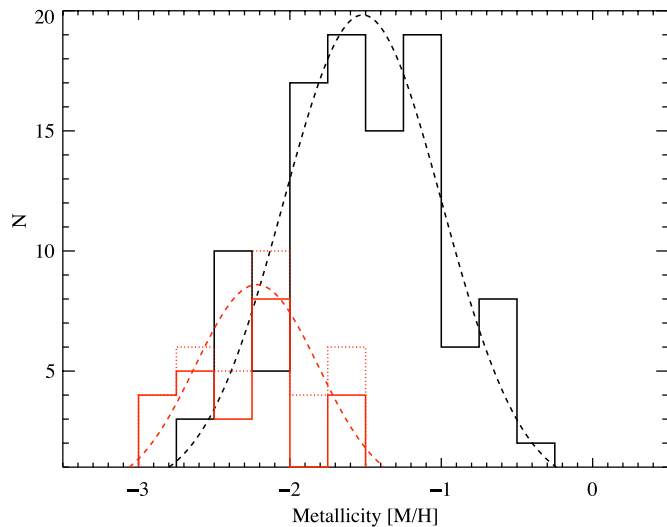
## 5. SUMMARY

Our survey of DLAs pre-selected to have weak metal-line absorption has established a “floor” of metallicity in the ISM of high- $z$  galaxies at  $\approx 1/1000$  solar abundance. Because these systems represent the most likely systems to exhibit very low abundance in the available SDSS sample, it is unlikely that metallicities much lower than these observed values will be found in future observations of DLA systems. The observed floor for the DLA systems with detected C, O, Si, Fe, and Al



**Figure 12.** Derived column density of H I for the ESI spectra, plotted against those for the SDSS spectra. The lower resolution of the SDSS spectra caused some discrepancies due to blending, and continuum placement, but the  $N(\text{H I})$  values are generally within the  $1\sigma$  error estimate.

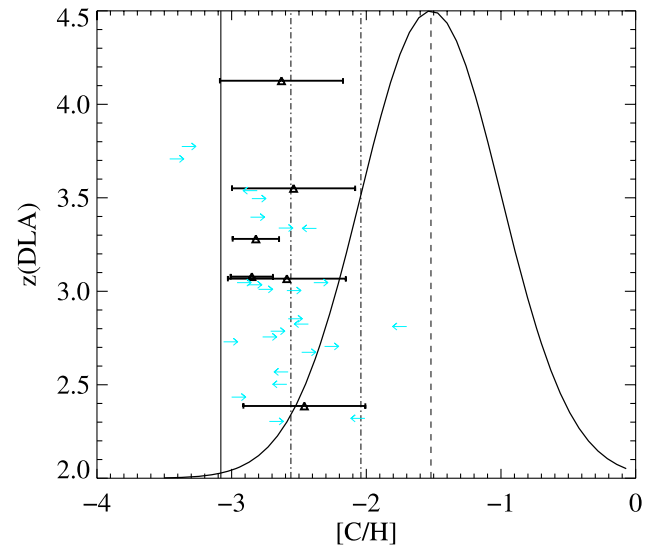
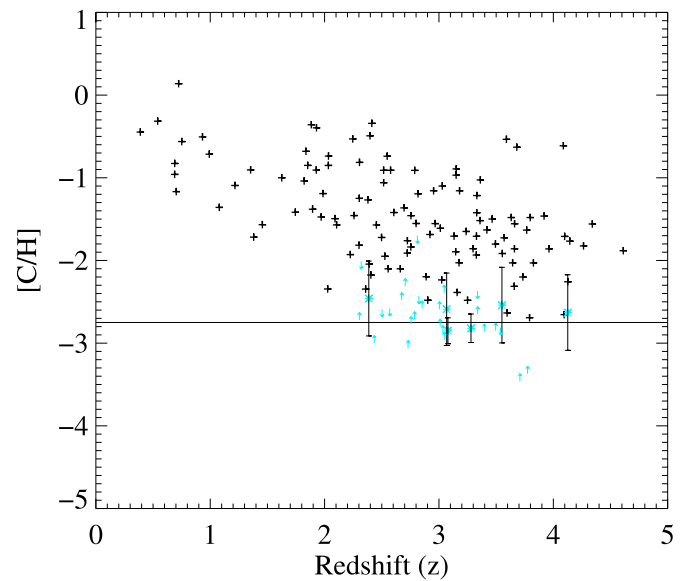
(A color version of this figure is available in the online journal.)



**Figure 13.** Plot of observed values of  $[\text{M}/\text{H}]$  from the survey of Prochaska et al. (2003b), and a Gaussian fit (black dashed line). The mean metallicity for this “control” sample in the redshift range of our sample ( $2 < z < 4.5$ ) is  $[\text{M}/\text{H}] = -1.52$ , with a width of  $\sigma = 0.52$ . We also present a histogram of  $[\text{M}/\text{H}]$  for the low-metallicity sample observed in this paper, for which a Gaussian fit (red dashed line) gives a mean metallicity of  $[\text{M}/\text{H}] = -2.2$ . DLA systems where lower limits of  $[\text{M}/\text{H}]$  have been observed are plotted with the dashed (red) lines on top of the values for the metal-poor DLAs.

(A color version of this figure is available in the online journal.)

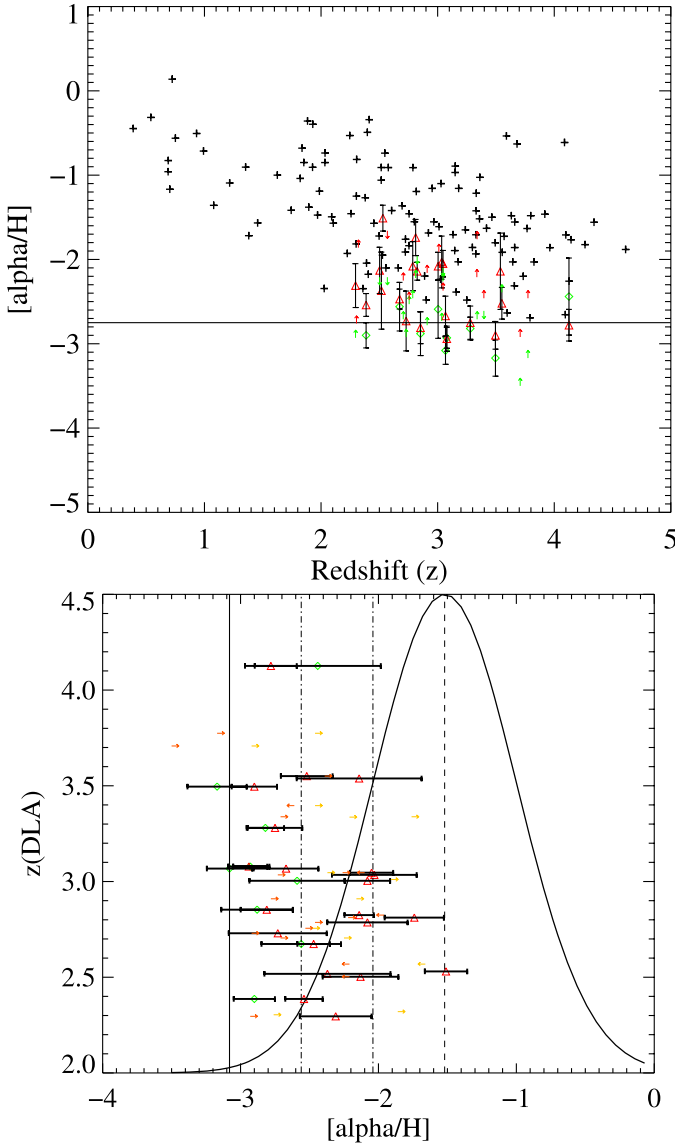
are found to be  $[\text{C}/\text{H}] > -2.9$ ,  $[\text{O}/\text{H}] > -3.2$ ,  $[\text{Si}/\text{H}] > -2.9$ ,  $[\text{Al}/\text{H}] > -2.9$ , and  $[\text{Fe}/\text{H}] > -3.2$ . The existence of a “floor” of metallicity is demonstrated from our observations, regardless of the effects of saturation on our spectra, which could transform some of our reported elemental abundances into lower limits for very low values of Doppler parameter ( $b < 5 \text{ km s}^{-1}$ ). In a very few cases, DLA systems without detected Si or Fe lines could have abundances lower than these “floor” values, but higher resolution spectra or higher signal-to-noise spectra are needed to test this.



**Figure 14.** Plot of DLA element abundance relative to solar ( $[\text{C}/\text{H}]$ ) vs. redshift for our sample (symbols with error bars), compared with that of Prochaska et al. (2003b) (top), and compared to a histogram of metallicity for the same sample (bottom). Vertical lines on the right panel show the location of 1, 2, and  $3\sigma$  departures from the mean DLA metallicity of  $-1.52$  from the Prochaska et al. (2003b) survey.

(A color version of this figure is available in the online journal.)

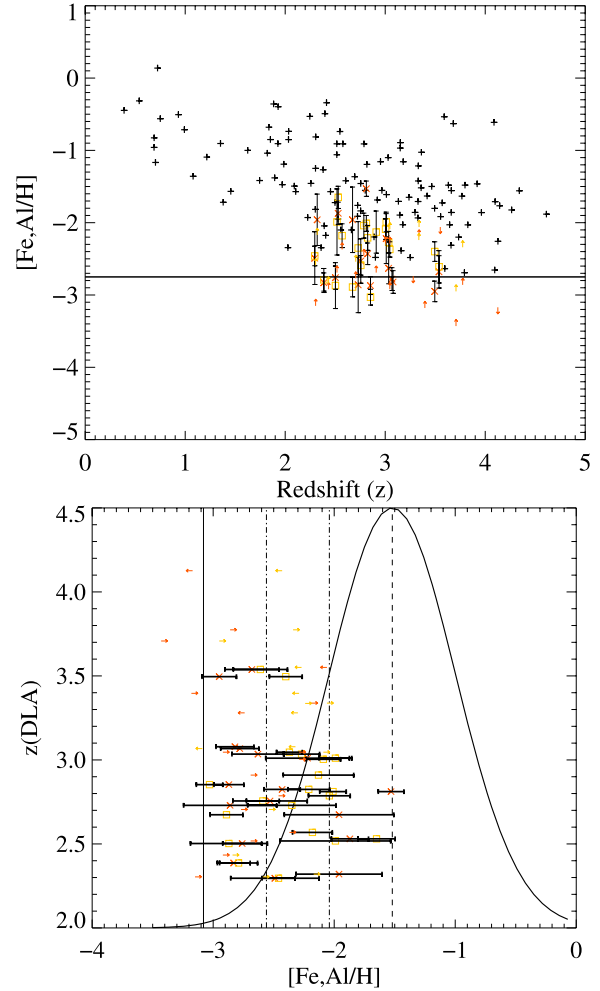
From this study, we observe multiple low-metallicity systems with multiple transitions which have  $[\text{X}/\text{H}] \approx -2.8$ , enabling a comparison with nucleosynthesis models, and providing the first observations of an “ultra-low-metallicity” DLA sample. A comparison of our results with stellar abundance studies is consistent, and extends the  $[\text{C}/\text{O}]$  observations to new low values of  $[\text{O}/\text{H}]$ , in which an upward trend of  $[\text{C}/\text{O}]$  is observed, consistent with nucleosynthetic yields of high-mass stars with masses between  $20$  and  $80 M_{\odot}$ . We observe in the DLA abundances a nucleosynthesis pattern consistent with a Population III population with IMF that includes stars from  $20 M_{\odot} < M_{*} < 80 M_{\odot}$ , and further studies are ongoing to provide higher resolution spectra of these DLA systems to include more elements, and to also extend our sample to provide a larger statistical sample of low-metallicity quasars. The low number



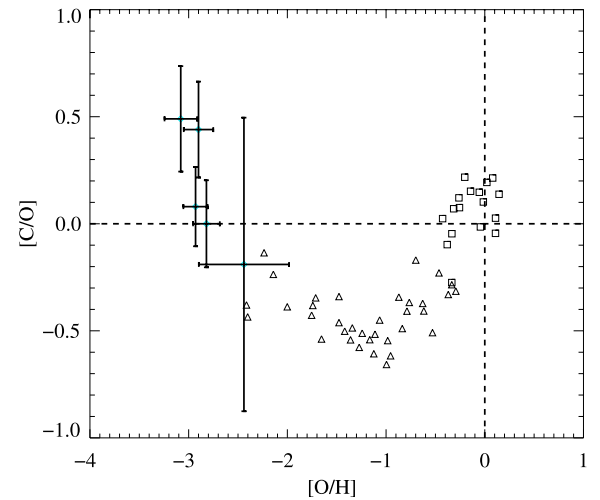
**Figure 15.** Plot of DLA element abundance relative to solar ( $[\alpha/\text{H}]$ ) vs. redshift for our sample (symbols with error bars), compared with that of Prochaska et al. (2003b) (plus signs) for the  $\alpha$ -process elements (top), and compared to a histogram of metallicity for the same sample (bottom). The  $\alpha$ -process elements presented include  $[\text{O}/\text{H}]$  (triangles) and  $[\text{Si}/\text{H}]$  (diamonds). (A color version of this figure is available in the online journal.)

of DLA systems with  $[\alpha/\text{Fe}] < -3.0$  even in our carefully selected sample suggests that the fraction of low-metallicity DLA systems and galaxies with  $[\alpha/\text{Fe}] < -3.0$  may be as low as  $\sim 0.135\%$ , based on a consideration of the statistics of DLA metallicity, and the results from our DLA sample.

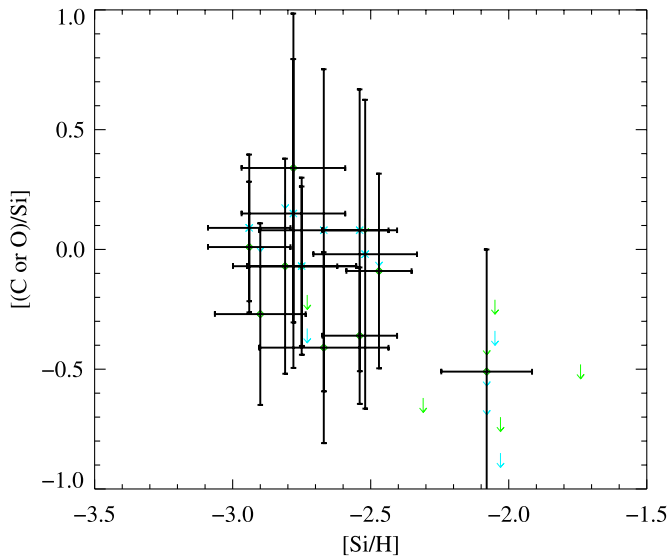
The observations from the Keck telescope were made possible by Pomona College travel grants to B.E.P., I.T.M., and D.J.B and a Pomona SURP grant to I.T.M. The authors also recognize the very significant cultural role and reverence that the summit of Mauna Kea has always had within the indigenous Hawaiian community. J.X.P. acknowledges funding through an NSF CAREER grant (AST-0548180) and an additional NSF grant (AST-0709235). W.W.S and B.E.P. acknowledge support from an NSF grant (AST-0606868).



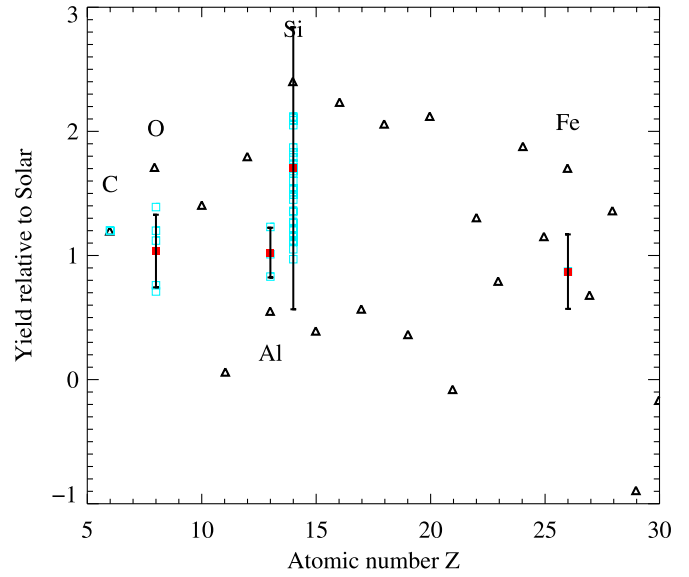
**Figure 16.** Plot of DLA element abundance relative to solar ( $[\text{X}/\text{H}]$ ) vs. redshift for our sample (symbols with error bars), compared with that of Prochaska et al. (2003b) (plus signs) for the Fe-peak elements (top), and compared to a histogram of metallicity for the same sample (bottom). The Fe-peak elements presented include  $[\text{Fe}/\text{H}]$  (squares) and  $[\text{Al}/\text{H}]$  ( $\times$ 's). (A color version of this figure is available in the online journal.)



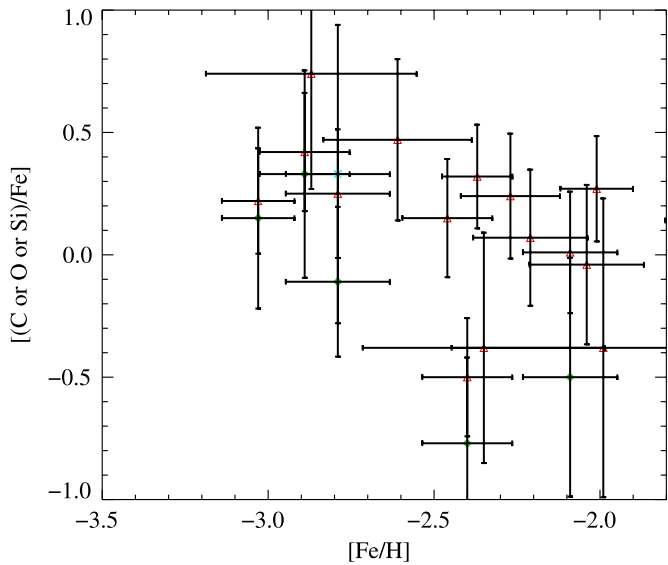
**Figure 17.**  $[\text{C}/\text{O}]$  vs.  $[\text{O}/\text{H}]$  as in the figure observed only for those points with small equivalent widths with equivalent width less than  $130 \text{ m}\text{\AA}$  for both  $\text{C II}$  and  $\text{O I}$ , and corrected for saturation using the conservative procedure described in Section 2.3, which assumes a  $b$ -value of  $b = 7.5 \text{ km s}^{-1}$ . The upward trend in  $[\text{C}/\text{O}]$  persists for this smaller sample, and suggests that the effect is independent of saturation effects. Stellar measurements of  $[\text{C}/\text{O}]$  (squares and triangles) from Akerman et al. (2004) are also plotted for comparison.



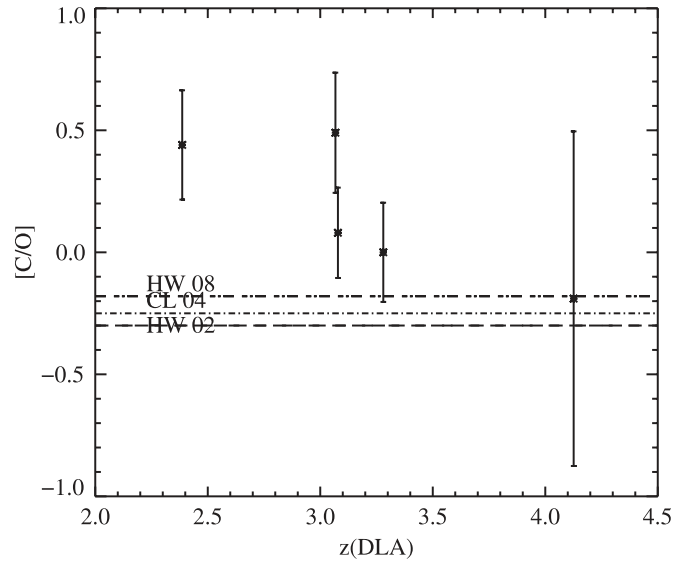
**Figure 18.** Trend in metal abundances of [C/Si] (turquoise X) and [O/Si] (green diamonds) against [Si/H]. The data shows a systematic dependence of [C/Si] and [O/Si] against Si/H.  
(A color version of this figure is available in the online journal.)



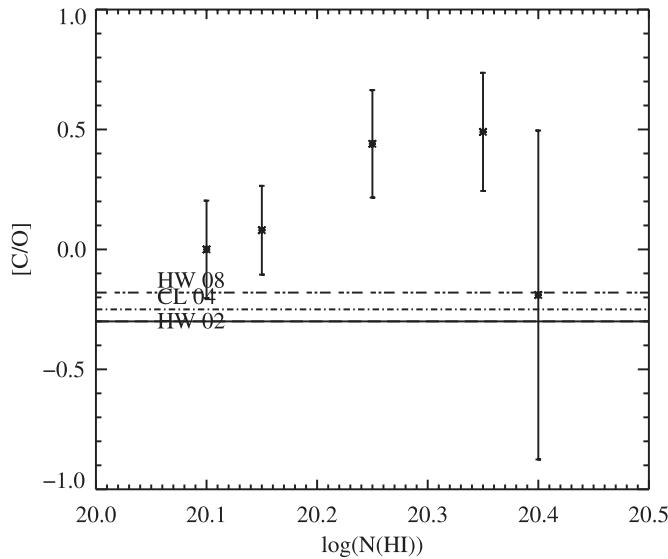
**Figure 20.** Plot of element abundances for the DLA systems shown in Figure 17 against element number, with open squares representing individual DLA systems, and the mean and standard deviation of the sample of four DLAs shown by the filled squares and error bars. The nucleosynthetic yields of Heger & Woosley (2002) for the various elements are shown with triangles. We observe reduced abundances of O, Si, and Fe relative to C compared to the nucleosynthetic model, and a slight enhancement of Al relative to C.  
(A color version of this figure is available in the online journal.)



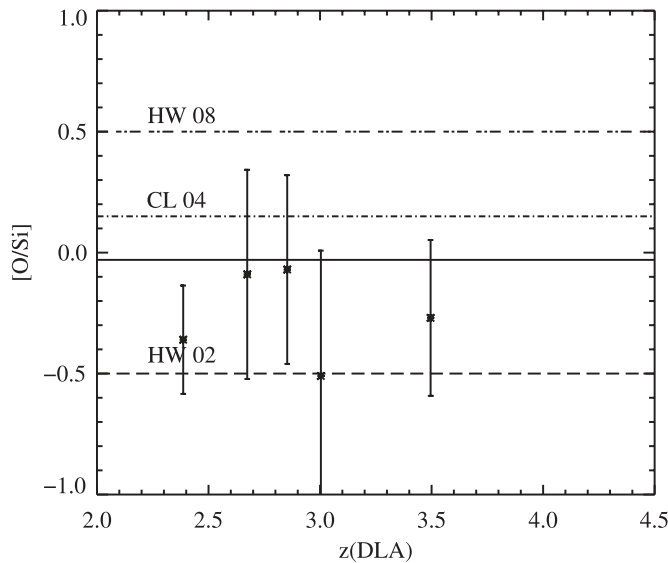
**Figure 19.** Trend in metal abundances of [C/Fe] (turquoise X), [O/Fe] (green diamonds), and [Si/Fe] (red triangles) against [Fe/H]. The data show a systematic increase of lighter element abundances against Si/H.  
(A color version of this figure is available in the online journal.)



**Figure 21.** Plot of our observed values of [C/O] against redshift, for our DLA sample. The predicted yields for nucleosynthesis models of massive stars are indicated by the horizontal lines for the models of Heger & Woosley (2008) (top), Chieffi & Limongi (2004) (middle), and Heger & Woosley (2002) (bottom). Also indicated is the mean value of [C/O] = -0.30 for the high-redshift DLA systems observed by Becker et al. (2006) (solid line).



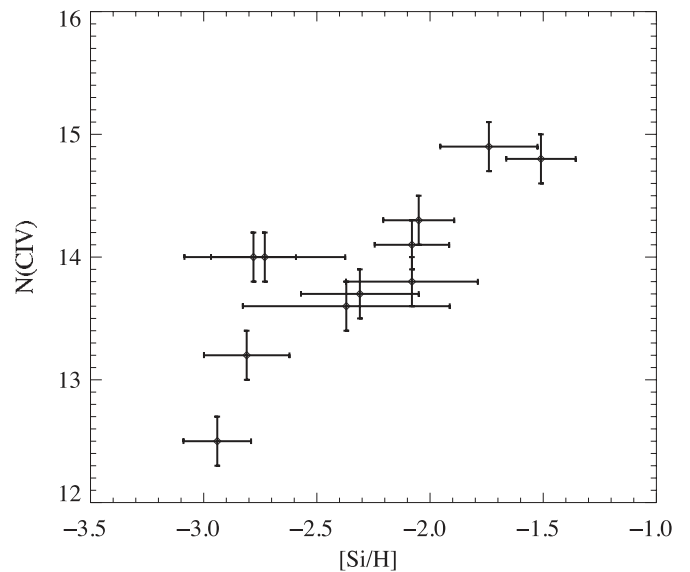
**Figure 22.** Plot of the observed values of  $[\text{C}/\text{O}]$  as in Figure 21, but against the measured column density of  $\text{H I}$ . The lowest column densities of both  $\text{C II}$  and  $\text{O I}$  are observed toward systems which are technically “sub-DLAs,” since  $N(\text{H I}) < 20.3$  for these systems. The observed enhancement of  $[\text{C}/\text{O}]$  for our observed systems is larger than the predicted ionization corrections, which are in the range of 0.1–0.2 dex.



**Figure 23.** Plot of the observed ratios of the elements  $[\text{O}/\text{Si}]$  the same systems as in Figure 21, compared with nucleosynthesis models and the high-redshift DLA systems observed by Becker et al. (2006) (solid line). The values of  $[\text{O}/\text{Si}]$  for our sample seem to agree with the high-redshift DLA sample, and are intermediate between two of the nucleosynthesis models.

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**Figure 24.** Plot of the observed metallicity (in this case  $[\text{Si}/\text{H}]$ ) compared with the  $\text{C IV}$  column density for our sample, showing an obvious correlation between the  $N(\text{C IV})$  and  $[\text{Si}/\text{H}]$ , which could reflect an increase in velocity dispersion for higher metallicity DLA systems.

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