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## Keck Observations of the Galactic Center Source G2: Gas Cloud or Star?

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### ABSTRACT

We present new observations and analysis of G2 – the intriguing red emission-line object which is quickly approaching the Galaxy’s central black hole. The observations were obtained with the laser guide star adaptive optics systems on the W. M. Keck I and II telescopes and include spectroscopy ( $R \sim 3600$ ) centered on the Hydrogen Br- $\gamma$  line as well as  $K'$  ( $2.1 \mu\text{m}$ ) and  $L'$  ( $3.8 \mu\text{m}$ ) imaging. Analysis of these observations shows the Br- $\gamma$  line emission has a positional offset from the  $L'$  continuum. This offset is likely due to background source confusion at  $L'$ . We therefore present the first orbital solution derived from Br- $\gamma$  line astrometry, which when coupled with radial velocity measurements, results in a later time of closest approach ( $2014.21 \pm 0.14$ ), closer periastron ( $130 \text{ AU}$ ,  $1900R_s$ ), and higher eccentricity ( $0.9814 \pm 0.0060$ ) compared to a solution using  $L'$  astrometry. The new orbit casts doubt on previous associations of G2 and a low surface brightness “tail.” It is shown that G2 has no  $K'$  counterpart down to  $K' \sim 20$  mag. G2’s  $L'$  continuum and the Br- $\gamma$  line-emission is unresolved in almost all epochs; however it is marginally extended in our highest quality Br- $\gamma$  data set from 2006 and exhibits a clear velocity gradient at that time. While the observations altogether suggest that G2 has a gaseous component which is tidally interacting with the central black hole, there is likely a central star providing the self-gravity necessary to sustain the compact nature of this object.

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## 1. Introduction

Recently, Gillessen et al. (2012, 2013) reported the discovery of G2, an extremely red object with spatially resolved Br- $\gamma$  emission. The object was interpreted as a  $\sim 3$  Earth mass gas cloud based upon an inferred low dust temperature, observed elongation of the Br- $\gamma$  emission along the object’s direction of motion, and a claimed tail along the same orbital trajectory as G2. This interpretation is particularly interesting because G2 is on a highly eccentric orbit with closest approach to our Galaxy’s central black hole within the next year, potentially allowing us to observe an unprecedented accretion event onto a supermassive black hole and offer insight into the region surrounding the black hole (e.g. Morris 2012; Mościbrodzka et al. 2012; Anninos et al. 2012; Saitoh et al. 2012; Bartos et al. 2013; Yusef-Zadeh & Wardle 2013).

The interpretation of G2 as a gas cloud, however, is not definitive. One challenge for the pure gas cloud scenario is that, given the strong tidal fields in this region and G2’s low self-gravity, G2 needs to have formed quite recently ( $\sim 1995$ , just prior to the initial observations; Burkert et al. 2012; Schartmann et al. 2012). Since such a gas cloud would be tidally disrupted during its periape passage in the upcoming year, the gas cloud model implies G2 will be observed over almost the exact extent of its existence. Therefore several alternative scenarios invoking an underlying star have been proposed (Miralda-Escudé 2012; Murray-Clay & Loeb 2012; Scoville & Burkert 2013). In these scenarios, the observed heated gas is posited to be circumstellar, either intrinsic or a consequence of interaction of the star and surrounding ambient gas. In these stellar scenarios, G2 not only has existed for timescales much longer than the observed time baseline, but will also survive periape passage.

Regardless of its nature, G2’s properties and possible origin depend critically on its orbital parameters. These parameters have been estimated from observations with a short time baseline compared to the orbital period ( $\sim 10$  vs.  $\sim 200$  years; Gillessen et al. 2013) and in a very crowded region, making the orbital solution susceptible to biases (Hartkopf et al. 2001). We therefore present new measurements and analysis of G2 that minimize the effects of source confusion on estimates of G2’s orbital parameters and examine the temporal evolution of G2’s properties.

## 2. Observations

Two types of new data were collected for this study using the laser-guide-star adaptive optics (LGS AO) systems at the W.M. Keck Observatory (Wizinowich et al. 2006; van Dam et al. 2006;

Chin et al. 2012). Spectra were obtained using the OSIRIS integral field spectrograph (Larkin et al. 2006) through the narrow-band Kn3 filter, which is centered on the Br- $\gamma$  hydrogen line (2.1661  $\mu\text{m}$ ), at a spectral resolution of  $R \sim 3600$ . Imaging data were obtained in the  $K'$ -band filter ( $\lambda_0 = 2.124 \mu\text{m}$ ) and the  $L'$ -band filter ( $\lambda_0 = 3.776 \mu\text{m}$ ) using the Keck II near-infrared camera (NIRC2, PI: K. Matthews). These data were obtained and reduced in a similar manner as in our previous publications (Lu et al. 2009; Yelda et al. 2010; Meyer et al. 2012) and the specific spectroscopic and  $L'$  observations utilized in this paper are described in table 1.

### 3. Analysis

#### 3.1. OSIRIS IFU Measurements

Because G2 has no detectable  $K$  continuum and is fainter than any object we have previously extracted, some analysis steps differ from the spectral extraction presented in Ghez et al. (2008) and Do et al. (2009). For all epochs, we created a combined data cube before extracting G2's spectrum rather than extracting spectra from individual cubes. The OH sky lines in the data are subtracted using sky frames scaled to the strength of families of OH lines in the observed frames to account for temporal variations. In 2012, the Br- $\gamma$  emission line from G2 is coincident with the 2.180  $\mu\text{m}$  OH line, so to minimize the systematic effects associated with OH line subtraction, we scale the sky only to this line.

An iterative process was required to estimate the position and spectral properties of G2 in the OSIRIS data cubes (see figures 1 and 2) since G2's position is needed to place the aperture for spectral extraction, and G2's spectral properties are needed to determine which OSIRIS channels should be used to measure its position. G2's position was first obtained by visual inspection of the data cube. Then, an initial spectrum was extracted at this position using an aperture with a radius of 35 mas. Emission from local ambient gas was subtracted using a region free of stellar halos within  $\sim 0''.5$ . A Gaussian was fit to the resulting emission line. In order to refine the position, the three-dimensional data cube was median collapsed over the wavelength range corresponding to twice the standard deviation of the Gaussian fit to the emission line, centered on the line peak. Continuum emission was subtracted by averaging the median of the 25 nearest channels on either side of the line. G2's position was further refined with a two-dimensional Gaussian fit to the continuum-subtracted image. The iterative extraction process was repeated again to obtain a final position and spectrum for G2. Further iterations produced no significant changes. Measuring G2's position on a continuum-subtracted frame (see figure 2(b)) removes the effect of source confusion, thereby avoiding astrometric biases.

G2's spectral properties were obtained using the final Gaussian fits to the Br- $\gamma$  emission line. The radial velocities (RVs) for G2 were calculated from the offset of the line from the rest wavelength ( $\lambda_{Br\gamma} = 2.1661 \mu\text{m}$ ) and corrected to the local standard of rest. The reported FWHM measurements were corrected for instrumental broadening (FWHM=85km s $^{-1}$ ). Line flux measurements were

Table 1. Summary of Observations and Measurements of G2

UT date	Fractional Date	AO type/ Telescope	pix scale (mas)	$N_{\text{frames}}$ observed	$N_{\text{frames}}$ used	FWHM (mas)	Orig Pub. <sup>a</sup>	vlsr (km s <sup>-1</sup> )	Br- $\gamma$ FWHM (km s <sup>-1</sup> )	G2 $\Delta$ RA <sup>b</sup> (mas)	G2 $\Delta$ Dec (mas)
OSIRIS, Kn3											
2006 Jun 18,30; Jul 1	2006.495	Keck II LGS	35	28	27	74	3	1125 $\pm$ 9	135 $\pm$ 30	222.97 $\pm$ 5.05	-105.40 $\pm$ 2.40
2008 May 16; Jul 25	2008.487	Keck II LGS	35	22	21	78	4,0	<sup>e</sup>	<sup>e</sup>	181.82 $\pm$ 6.90	-64.97 $\pm$ 2.51
2009 May 5,6	2009.344	Keck II LGS	35	24	19	79	0	1352 $\pm$ 32	171 $\pm$ 121	176.26 $\pm$ 2.74	-65.76 $\pm$ 2.04
2010 May 5,8	2010.349	Keck II LGS	35	17	16	82	0	1479 $\pm$ 29	262 $\pm$ 165	166.29 $\pm$ 9.11	-43.95 $\pm$ 11.05
2012 Jun 9,11; Aug 11,12	2012.613	Keck I LGS	20 <sup>c</sup>	27	21 <sup>d</sup>	68	0	2071 $\pm$ 146	706 $\pm$ 250	103.16 $\pm$ 4.07	-8.79 $\pm$ 11.36
NIRC2, L'											
2003 Jun 10	2003.440	Keck II NGS	10	12	12	86	0			260.1 $\pm$ 4.5	-154.6 $\pm$ 5.1
2004 Jul 26	2004.567	Keck II LGS	10	11	11	80	1			262.6 $\pm$ 4.8	-140.4 $\pm$ 4.2
2005 Jul 30	2005.580	Keck II LGS	10	62	56	87	2			251.6 $\pm$ 3.4	-129.0 $\pm$ 2.6
2006 May 21	2006.385	Keck II LGS	10	19	19	81	0			226.2 $\pm$ 1.9	-99.4 $\pm$ 0.8
2009 Jul 22	2009.561	Keck II LGS	10	4	4	85	0			175.4 $\pm$ 4.9	-70.7 $\pm$ 2.1
2012 Jul 20 - Jul 23	2012.562	Keck II LGS	10	1316	1314	93	0			105.9 $\pm$ 0.8	-21.1 $\pm$ 0.6

<sup>a</sup>References: (0) This Work; (1) Ghez et al. (2005); (2) Hornstein et al. (2007); (3) Ghez et al. (2008); (4) Do et al. (2009)

<sup>b</sup>Offset from SgrA\*-radio.

<sup>c</sup>A smaller square dither pattern of  $0''.5 \times 0''.5$  was used.

<sup>d</sup>In 2012, a more stringent quality cut (FWHM < 68mas) was used because G2 is closer to stars than in other epochs and because the improved LGS performance on Keck I allowed this more stringent cut.

<sup>e</sup>In 2008, low SNR and changes in the local standard of rest velocity between the two observation dates prevent reliable line measurements.

made by comparing the Gaussian fit to the non-variable star S0-2 between 2.17-2.18  $\mu\text{m}$ . Statistical uncertainties were taken from the root mean square (rms) error of Gaussian fits to three independent data subsets; however, we note that systematic errors (e.g. OH sky lines) are not included.

The final Br- $\gamma$  line maps were used to evaluate G2’s physical size and absolute astrometric position. The size of G2’s Br- $\gamma$  emission was estimated by comparing the FWHM of G2 to that of nearby stars. To track the motions of G2 in the plane of the sky, its OSIRIS position was transformed into an absolute coordinate system. To construct an absolute reference frame, we used the precise locations of well characterized stars in the field (see figure 2(a)), as measured with *StarFinder* (Diolaiti et al. 2000). The position of G2 from the two-dimensional Gaussian fit was added to the list of positions for each epoch. Then, these positions were matched to those from LGS AO  $K'$  observations of the same region at a nearby epoch using a second-order polynomial transformation to account for translation, rotation, and pixel scale differences between the images<sup>1</sup>. Once the positions in the OSIRIS frames were transformed into a  $K'$  epoch, they were transformed into an absolute reference frame in which SgrA\*-radio is at rest at the origin, as originally described in Yelda et al. (2010) and updated in Yelda et al. *in prep*. We investigated possible systematic astrometric effects in OSIRIS by cross-checking the OSIRIS astrometry of S0-2 to our standard NIRC2 analysis. The results are consistent, validating the use of OSIRIS astrometry.

### 3.2. NIRC2 Imaging Measurements

The  $L'$  position of G2 was estimated from calibrated near-infrared images in a similar fashion to that described above for the Br- $\gamma$  line maps. One difference was that we deconvolved the individual images using the Lucy-Richardson algorithm (Richardson 1972; Lucy 1974; Lucy & Hook 1992) in order to help isolate the point sources from the extended  $L'$  dust emission. Both the large-scale background and the PSF were estimated using *StarFinder*. After beam-restoring the individual frames with a Gaussian having FWHM of half the theoretical resolution limit (40 mas), the frames were averaged to create a final image for each epoch (see figure 2(d)). The results of this analysis do not vary if the image is restored using a wider Gaussian beam. To locate stars used as astrometric reference sources, *StarFinder* was run on the combined (deconvolved) frame using the beam-restoring Gaussian as the reference PSF. G2’s position was determined from an elliptical two-dimensional Gaussian fit<sup>2</sup>. The fitting of an elliptical Gaussian allows for extended structure and is analogous to the method used for the Br- $\gamma$  detection. The positions for G2 from this method are consistent with those produced via *StarFinder*. The size of G2 was measured by comparing to the PSF of stars with a similar flux. The observed  $L'$  positions were transformed into an absolute coordinate system as described for the Br- $\gamma$  detections.

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<sup>1</sup>In 2012, only a linear transformation was used because fewer reference stars were available due to the smaller field of view (7 versus 19-27 in other epochs).

<sup>2</sup>In 2012, a two-component elliptical Gaussian was fit since G2 is blended with another point source.

One complication for measuring the  $L'$  position of G2 in the plane of the sky is that G2 was superimposed on a filament of dust emission during the early- to mid- 2000's, which caused significant astrometric bias (see figure 2(d)). To quantify this effect, we performed a series of star-planting simulations using a 2012 May  $L'$  image, in which G2 is well off the filament. In each of the four simulations, an extra point source ( $L' = 14$ ) was planted at the expected G2 position at the time of earlier  $L'$  observations (2003.440, 2004.567, 2005.580, and 2006.385) based upon the orbit fit to the Br- $\gamma$  measurements. The background dust emission does not have noticeable motions in the plane of the sky, thus the astrometric bias at these locations in the 2012 image should be representative of the bias at earlier epochs. The position of the artificial point source was extracted using *StarFinder*.

To constrain the spectral energy distribution of the object, we measured the integrated flux of G2 in the  $L'$  frame (e.g. Witzel et al. 2012). In contrast with Eckart et al. (2013), but in agreement with Gillessen et al. (2012), we detect no  $K'$  counterpart for G2; therefore, we performed star planting simulations to determine the  $K'$  upper limit. We used the Br- $\gamma$  astrometry of G2 to predict the positions of G2 in the  $K'$  data. Then, an artificial point source was planted at G2's predicted position in the May 2010 image, when G2 is most isolated (see figure 2(c)). Using a modified version of the *StarFinder* algorithm (Boehle et al. *in prep*), we searched for a point source within a 3 pixel box centered on the planted star, given the PSF originally extracted from the May 2010 image.

### 3.3. Orbital Modeling

We fit an orbit to G2's astrometric positions and RV measurements.  $L'$ -astrometric and RV measurements reported by Gillessen et al. (2012, 2013) were also included; however, we fit the Br- $\gamma$  and  $L'$  astrometry separately. We assumed G2 follows a purely Keplerian orbit around the central black hole.

The orbit inference problem aims to find the most likely orbital parameters (eccentricity, period, time of periape passage, inclination, position angle of the ascending node, and the longitude of periape) given the data. In addition to G2's orbital parameters, the parameters of the gravitational potential need to be inferred (mass and distance to black hole, and its perceived 2D position and 3D velocity). Since the observed motion of G2 does not entail enough information to constrain all these parameters, the orbit of S0-2, which has undergone a full orbit, was fit simultaneously and effectively determines the gravitational potential (see Meyer et al. 2012). Compared to earlier publications, we amended our Keplerian orbit fitting code to use the Bayesian sampler MultiNest (Feroz & Hobson 2008). We verified the results are equivalent with our previous Monte Carlo approach (e.g. Ghez et al. 2008).

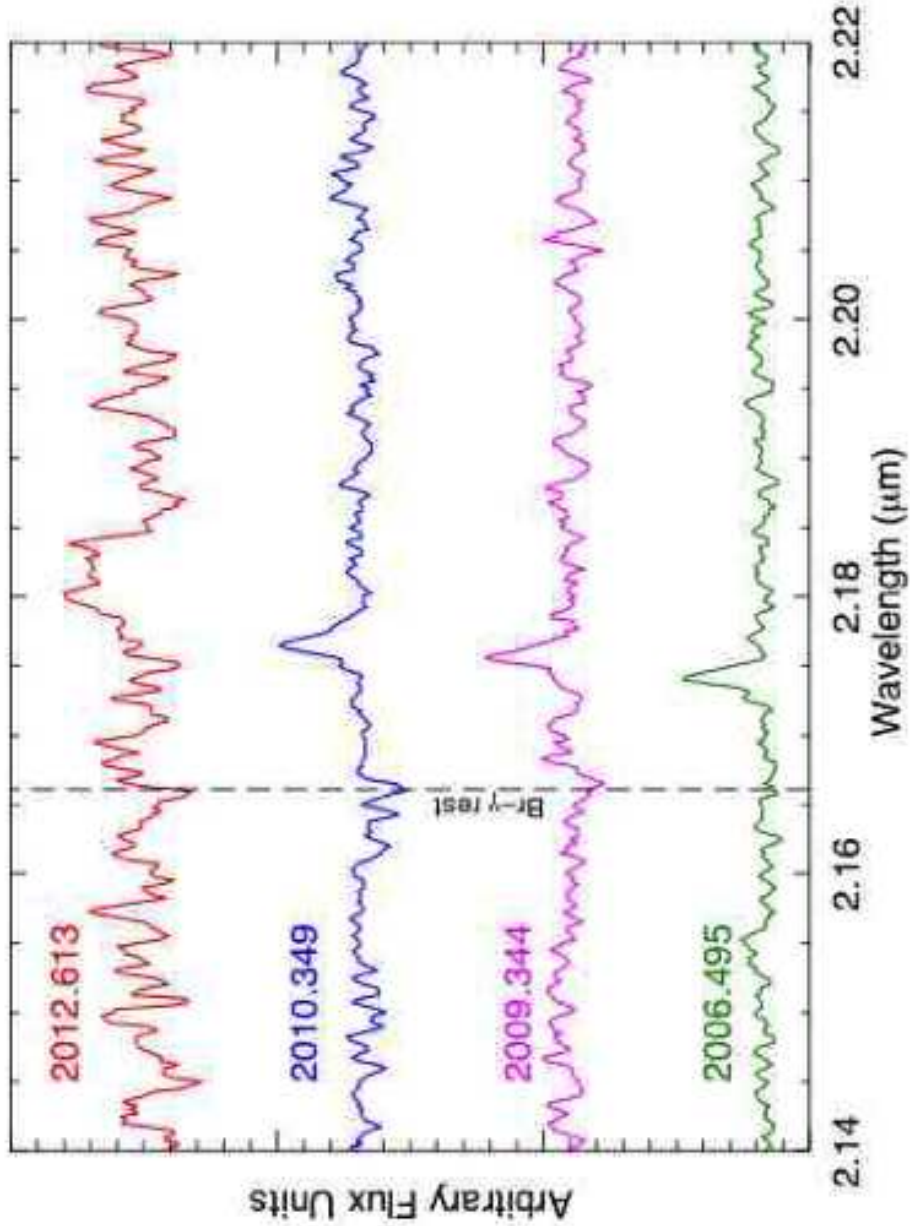


Fig. 1.— The evolving spectra of G2. The highly redshifted Br- $\gamma$  emission line has a FWHM which increases with time (see table 1). The Br- $\gamma$  line peak has been scaled such that it is the same height for each spectrum. Additionally, the spectra have been smoothed using a boxcar average with a width of 3 OSIRIS channels.



## 4. Results

Six key results emerge from our analysis of G2:

- *Apparent positional offset between the Br- $\gamma$  line emission and L' continuum emission of G2:* The L' astrometry is not consistent with our Br- $\gamma$  astrometry, which is reported here for the first time (see figure 3). The positional difference can be explained by astrometric bias in the L' astrometry. In our star planting analysis, we find the extracted position of an artificial L' source on the dust filament where G2 is located at early epochs differs from the input positions by an average of  $\sim 1.2$  pix (12 mas). This is almost an order of magnitude larger than typical centroiding uncertainties at L' and the astrometric bias found in a region isolated from extended background flux. We therefore exclude L' astrometry from our primary orbital fits.
- *Revised Orbital Parameters:* The Keplerian orbit derived using Br- $\gamma$  astrometry differs from the orbit derived using L' astrometry (e.g. Gillessen et al. 2013) by more than can be explained by formal measurement uncertainties (see figure 3). We list the most likely orbital parameters from Br- $\gamma$  in table 2. This new orbit for G2 has pushed the closest approach date to March 2014, has a closer periastron (130 AU,  $1900R_s$ ) and a higher eccentricity ( $0.9814 \pm 0.0060$ ).
- *G2's compact size:* G2's observed size is comparable to the Keck angular resolution in almost all our observations; however, it exhibits marginal spatial extent in our highest quality spectroscopic data set. At this time (2006), the observed half-width at half-maximum is  $20 \pm 4$  mas after accounting for the observed PSF of S0-2 (FWHM =  $74 \pm 3$  mas). G2 also shows a velocity gradient along its orbit, with the most highly redshifted portion of the object closest to the black hole, consistent with what is expected for a tidal interaction event. The observed spatial extent is consistent with that presented in Gillessen et al. (2012), thus the Br- $\gamma$  spatial extent of G2 along the direction of motion is relatively constant and corresponds to a size on the order of  $\sim 100$  AU.
- *Spectral Evolution:* We find an increasing Br- $\gamma$  line width for G2, with FWHM measurements ranging from  $135 \pm 30$  km s $^{-1}$  in 2006 to  $706 \pm 250$  km s $^{-1}$  in 2012 (see table 1). These measurements are consistent with those presented by Gillessen et al. (2012, 2013).
- *Constant Brightness:* Both at L' and Br- $\gamma$  the flux of G2 is constant within uncertainties. The integrated flux of G2 at L' has an average value of  $14.9 \pm 0.3$  mag. The Br- $\gamma$  line fluxes in 2006, 2009, 2010, and 2012 are  $0.884 \pm 0.116\%$ ,  $0.919 \pm 0.098\%$ ,  $1.13 \pm 0.41\%$ ,  $1.14 \pm 0.87\%$  of S0-2's continuum flux between 2.17-2.18  $\mu\text{m}$ , respectively.
- *Source Color:* Our star planting simulations indicate that the upper K' magnitude limit of G2 is  $K' = 20$  and is  $\sim 2$  magnitudes deeper than the earlier upper K' limit imposed on G2 (Gillessen et al. 2012). The deeper K' limit yields a dust temperature below  $\sim 500\text{K}$  if G2 is a pure gas cloud. If G2 has an underlying stellar source, a K' magnitude limit places distinct

limits on the luminosity of the star. This limit is similar to the expected  $K'$  magnitude for a low-mass T Tauri star in the Galactic Center (e.g. Scoville & Burkert 2013), and eliminates a more massive star unless it is shrouded by dust which self-extincts the star at  $K'$ .

## 5. Discussion

The two orbits derived by Gillessen et al. (2012, 2013) differ from each other by more than can be explained with formal measurement uncertainties. Systematic errors in the L' astrometry naturally explain the discrepancy. Because the offset becomes less significant and the RV measurements play an increasingly dominant role in the orbital modeling, the measurements seem to converge on the solution obtained with the less biased Br- $\gamma$  astrometry and RV measurements.

There are several implications of the new orbital model parameters. First, the revised periapse date of March 2014 and the closer periastron are important for the design and interpretation of SgrA\* monitoring programs designed to test for an increased accretion flow. With a special NRAO call for proposals and tremendous attention called to this object, which was originally anticipated to experience periapse in June 2013, more than 30 programs have been approved for the summer of 2013 covering radio to  $\gamma$ -ray wavelengths. While some models predict that a cloud would generate enhanced radio emission from SgrA\* well in advance of periapse passage (e.g. Narayan et al. 2012; Sadowski et al. 2013a,b), none has yet been detected (Kassim et al. 2013), although this may not yet be expected with the revised periapse passage.

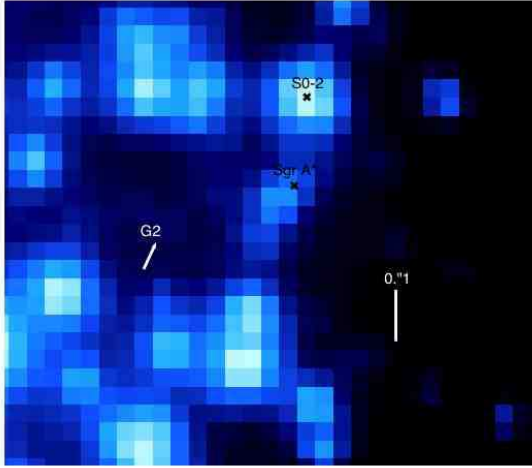
Second, the geometry of the new orbit may call two previously suggested associations into question. Most importantly, the previous claimed “tail” of low surface brightness emission, which was central to the original claim that G2 is a pure gas cloud, is no longer securely affiliated with the compact “head”, which has been analyzed in this work. The new orbit also falls  $3.5\sigma$  off the plane of the disk of young stars orbiting the central black hole at  $\sim 0.05 - 2$  pc (Genzel et al.

Table 2. Orbital elements for G2<sup>a</sup>

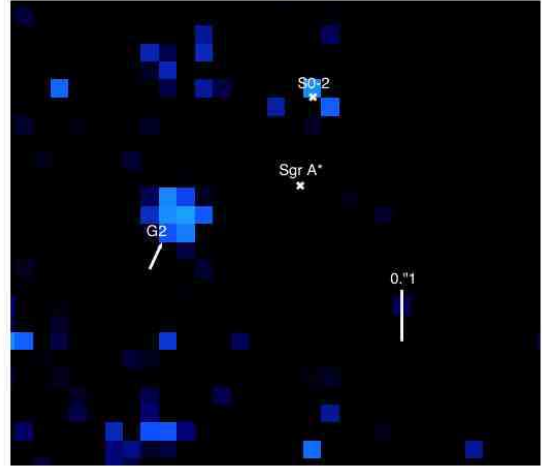
Parameter	Value <sup>b</sup>
Time of closest approach ( $T_0$ )	$2014.21 \pm 0.14$ yrs
Eccentricity ( $e$ )	$0.9814 \pm 0.0060$
Period ( $P$ )	$276 \pm 111$ yrs
Angle to periapse ( $\omega$ )	$88 \pm 6$ deg
Inclination ( $i$ )	$121 \pm 3$ deg
Position angle of the ascending node ( $\Omega$ )	$56 \pm 11$ deg

<sup>a</sup>Parameters describing the gravitational potential are found in Meyer et al. (2012).

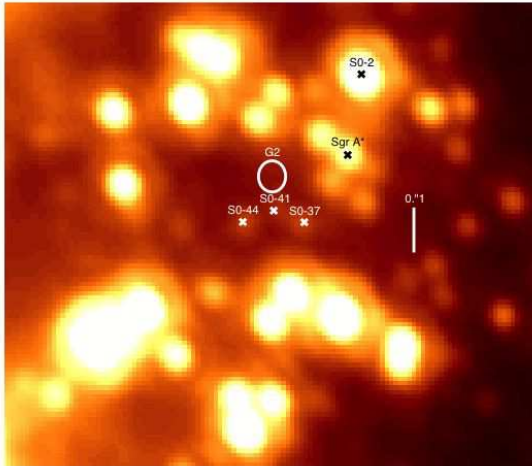
<sup>b</sup>Values provided are the mean and standard deviation of the marginalized one-dimensional distributions.



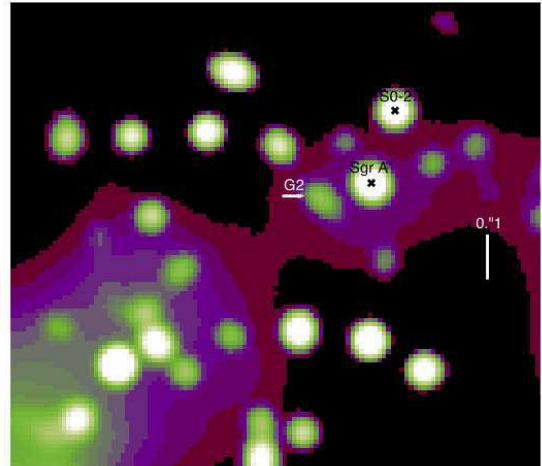
(a) A collapsed OSIRIS cube from May 2006 showing the continuum sources through the narrow-band Kn3 filter. The position of G2, which has no detectable  $K$  continuum, is indicated with an arrow.



(b) The Kn3 image shown in (a) is continuum-subtracted such that only Br- $\gamma$  emission at G2's redshift is shown. G2 is isolated in this continuum-subtracted map, minimizing the effect of stellar confusion on positional measurements.



(c) A NIRC2  $K'$  image from May 2010 with the OSIRIS position of G2 indicated by a  $3\sigma$  contour. For reference, the detections of other  $K'$  sources are shown. This image was used to derive G2's  $K'$  magnitude upper limit since it is most isolated in this imaging epoch.



(d) A NIRC2  $L'$  image from July 2012 which has been deconvolved and restored with a Gaussian Beam. The background estimation from *StarFinder* has also been added back in to show the dust filament which biases G2's  $L'$  position in early epochs. This feature extends upwards from the lower left of the image. G2 is blended with another point source in 2012.

Fig. 2.— Images of the Galactic Center showing the Br- $\gamma$  and  $L'$  detections of G2, as well as the map in which the  $K'$  upper limit was derived.

2000; Levin & Beloborodov 2003; Paumard et al. 2006; Lu et al. 2009; Bartko et al. 2009; Yelda et al. *in prep*) which has been suggested as the possible origin of the gas cloud through colliding stellar winds of young stars (Ozernoy et al. 1997; Cuadra et al. 2006) or a central stellar source (e.g. Murray-Clay & Loeb 2012).

Third, the even higher eccentricity provides strong constraints on G2’s origin. It is notable that a high eccentricity is consistent with the outcome of three body exchanges between a binary system and the central black hole which may explain the dense concentration of B stars (the S-stars) in the innermost regions of the Galaxy (e.g. Alexander & Livio 2004). If a triple star system is invoked rather than a binary, G2 could possibly be the result of a recent merger between two components to produce the observed physical properties (red, compact, and at times marginally resolved) and an exchange with the central black hole with the remaining component to produce the observed orbital parameters (highly eccentric, short period). While we favor a stellar model, this question will be soon addressed, as G2 should remain intact through its periapse if this is indeed correct.

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*Facilities:* Keck: II (LGS AO, NIRC2, OSIRIS), Keck: I (LGS AO, OSIRIS)

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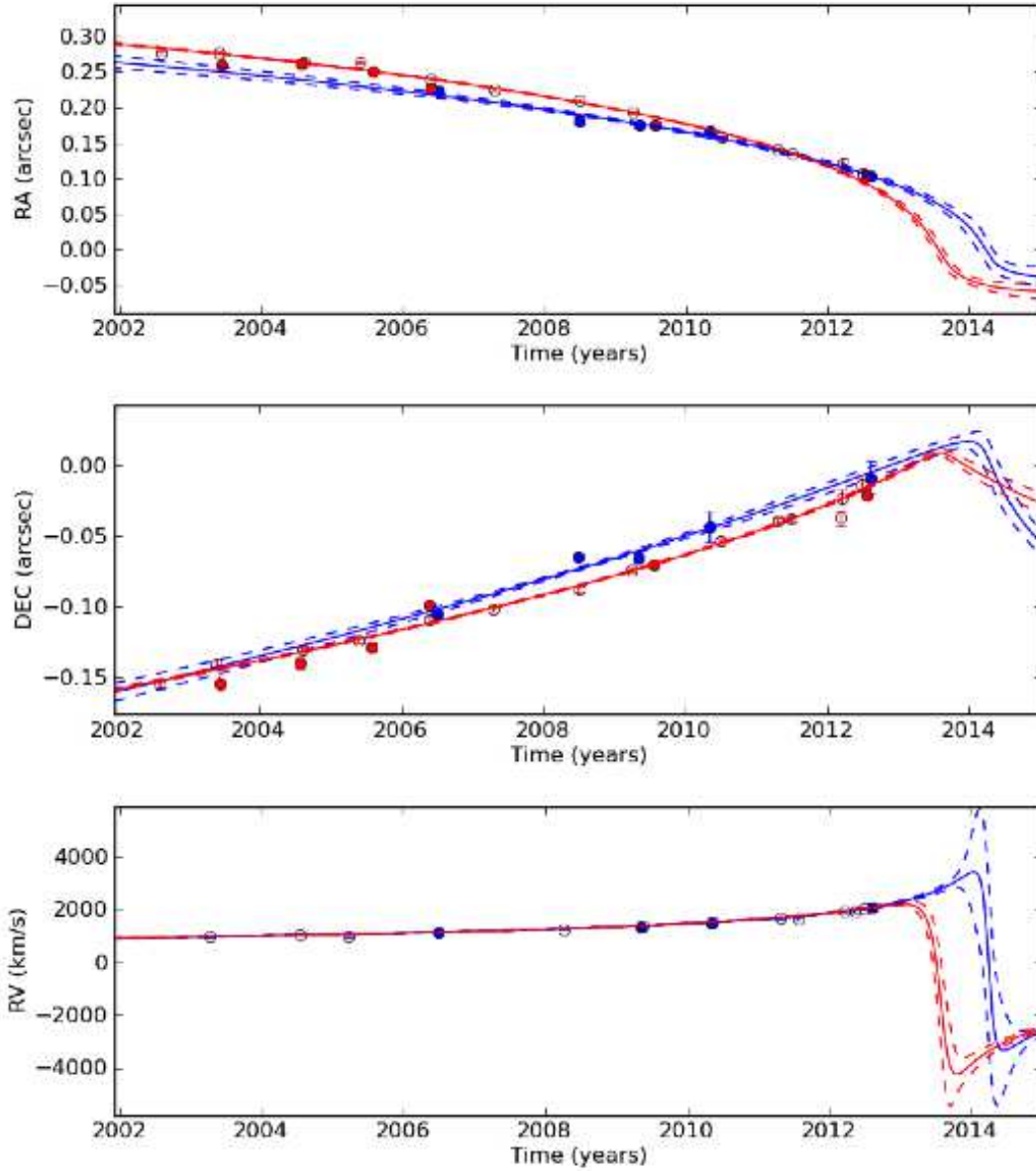


Fig. 3.— The orbital fit using the Br- $\gamma$  astrometry is shown in blue and the fit derived using the  $L'$  astrometry is shown in red. Both solutions use radial velocity fits determined from the Br- $\gamma$  emission line. Astrometric points obtained from data presented in this paper are indicated with filled points, and open circles indicate measurements presented in Gillessen et al. (2012, 2013). Both  $L'$  data sets appear offset from the Br- $\gamma$  astrometry.