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Krylov subspace methods for linear systems with tensor product structure

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Seminar for Applied Mathematics, ETH Zürich

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Outline



- 2 Basic Algorithm
- 3 Convergence bounds
- Solving the compressed equation
- 5 Numerical experiments

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Numerical experiments

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Linear system with tensor product structure

We consider the linear system

$$\mathcal{A}x = b$$

with

$$\mathcal{A} = \sum_{s=1}^{d} I_{n_1} \otimes \cdots \otimes I_{n_{s-1}} \otimes A_s \otimes I_{n_{s+1}} \otimes \cdots \otimes I_{n_d},$$
$$b = b_1 \otimes \cdots \otimes b_d,$$
$$A_s \in \mathbb{R}^{n_s \times n_s} \text{ positive definite, } b_s \in \mathbb{R}^{n_s}.$$

Example for 3 dimensions:

 $(A_1 \otimes I \otimes I + I \otimes A_2 \otimes I + I \otimes I \otimes A_3)x = b_1 \otimes b_2 \otimes b_3$

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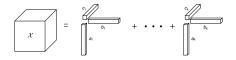
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Tensor decompositions

CP decomposition:

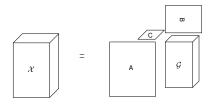
$$\mathsf{vec}(\mathcal{X}) = \sum_{r=1}^{k} a_r \otimes b_r \otimes c_r$$
 $a_r \in \mathbb{R}^m, b_r \in \mathbb{R}^n, c_r \in \mathbb{R}^p$



Tucker decomposition:

$$\mathsf{vec}(\mathcal{X}) = \sum_{i=1}^{\tilde{m}} \sum_{s=1}^{\tilde{n}} \sum_{l=1}^{\tilde{p}} \mathcal{G}_{ijl} \ \mathbf{a}_i \otimes \mathbf{b}_j \otimes \mathbf{c}_l$$

 $= (\mathbf{A} \otimes \mathbf{B} \otimes \mathbf{C}) \operatorname{vec}(\mathcal{G})$
 $\mathcal{G} \in \mathbb{R}^{\tilde{m} imes \tilde{n} imes \tilde{p}},$
 $\mathbf{a}_i \in \mathbb{R}^m, \mathbf{b}_j \in \mathbb{R}^n, \mathbf{c}_l \in \mathbb{R}^p$



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About this system

The eigenvalues of the matrix \mathcal{A} are given by all possible sums

$$\lambda_{i_1}^{(1)} + \lambda_{i_2}^{(2)} + ... + \lambda_{i_d}^{(d)}$$

where $\lambda_{i_s}^{(s)}$ denotes an eigenvalue of A_s . For A_s positive definite, the system has a unique solution.

Note that *x* and *b* are vector representations of tensors in $\mathbb{R}^{n_1 \times \cdots \times n_d}$, and *b* represents a rank-one tensor.

A tensor arising from the discretization of a sufficiently smooth function *f* can be approximated by a short sum of rank-one tensors. By superposition, we can insert such a tensor as right-hand side into our system.

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Tensorized Krylov subspaces

A Krylov subspace is defined as

$$\mathcal{K}_k(A,b) = \operatorname{span}\{b, Ab, \ldots, A^{k-1}b\}.$$

We define a tensorized Krylov subspace as

$$\mathcal{K}^{\otimes}_{\mathfrak{K}}(\mathcal{A}, b) := \operatorname{span} \big(\mathcal{K}_{k_1}(A_1, b_1) \otimes \cdots \otimes \mathcal{K}_{k_d}(A_d, b_d) \big).$$

Note that

$$\mathcal{K}_{k_0}(\mathcal{A}, b) \subset \mathcal{K}_{\mathfrak{K}}^{\otimes}(\mathcal{A}, b)$$

for $\Re = (k_0, \dots, k_0)$. \Rightarrow Tensorized Krylov subspaces are richer than standard Krylov subspaces.

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Reminder: CG method (Ax = b)

The best approximation of x in $\mathcal{K}_k(A, b)$ is defined by:

$$\|x_k-x\|_A=\min_{\tilde{x}\in\mathcal{K}_k(A,b)}\|\tilde{x}-x\|_A.$$

Find U_k with (orthonormal) columns that span the Krylov subspace $\mathcal{K}_k(A, b)$. Set $x_k = U_k y$, then y is the solution of the compressed system

$$H_k y = \tilde{b}, \qquad H_k = U_k^\top A U_k, \tilde{b} = U_k^\top b.$$

Convergence bound ($\kappa = \kappa_2(A)$)

$$\|x_k - x\|_A \leq C(A, b, \kappa) \Big(rac{\sqrt{\kappa} - 1}{\sqrt{\kappa} + 1}\Big)^k.$$

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Tensorized Krylov: $\mathcal{A} \operatorname{vec}(\mathcal{X}) = b$

Applying the Arnoldi method to A_s , b_s results in U_s , H_s with

$$U_s^\top A_s U_s = H_s,$$

U_s column-orthogonal, H_s upper Hessenberg matrix.

The columns of U_s span $\mathcal{K}_{k_s}(A_s, b_s)$, and similarly the columns of $\mathcal{U} := U_1 \otimes \cdots \otimes U_d$ span $\mathcal{K}^{\otimes}_{\mathcal{R}}(\mathcal{A}, b)$.

Solve the compressed system

$$\mathcal{H}y = \tilde{b}$$

with $x_{\mathfrak{K}} = \mathcal{U}y$, $\tilde{b} = \mathcal{U}^{\top}b$ and $\mathcal{H} = \mathcal{U}^{\top}\mathcal{A}\mathcal{U}$.

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Compressed system

The structure of $\mathcal{H}y = \tilde{b}$ corresponds to that of $\mathcal{A}x = b$:

$$\mathcal{H} = \sum_{s=1}^{d} I_{k_1} \otimes \cdots \otimes I_{k_{s-1}} \otimes H_s \otimes I_{k_{s+1}} \otimes \cdots \otimes I_{k_d}$$
$$\tilde{b} = \tilde{b}_1 \otimes \cdots \otimes \tilde{b}_d = \|b\|_2 (e_1 \otimes \cdots \otimes e_1)$$

For dense core tensor y, x_{\Re} is a tensor in Tucker decomposition:

$$x_{\mathfrak{K}} = (U_1 \otimes \cdots \otimes U_d)y$$

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Convergence, s.p.d. case (1)

Similarly to CG, we have:

$$\|x_{\mathfrak{K}} - x\|_{\mathcal{A}} = \min_{\tilde{x} \in \mathcal{K}_{\mathfrak{K}}^{\otimes}(\mathcal{A}, b)} \|\tilde{x} - x\|_{\mathcal{A}}$$

Every vector in a Krylov subspace $\mathcal{K}_{k_s}(A_s, b_s)$ can be seen as $p(A_s)b_s$, with p a polynomial of order at most $k_s - 1$. A similar thought reduces the bound calculation to the min-max problem

$$E_{\Omega}(\mathfrak{K}) := \min_{\boldsymbol{\rho} \in \Pi_{\mathfrak{K}}^{\otimes}} \| \boldsymbol{\rho}(\lambda_{1}, \ldots, \lambda_{d}) - \frac{1}{\lambda_{1} + \cdots + \lambda_{d}} \|_{\Omega},$$

where Π_{\Re}^{\otimes} is a space of multivariate polynomials, and Ω contains the eigenvalues $(\lambda_1, \ldots, \lambda_d)$, with $\lambda_i \in [\lambda_{\min}(A_i), \lambda_{\max}(A_i)]$.

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Convergence, s.p.d. case (2)

Inserting an upper bound for $E_{\Omega}(\mathfrak{K})$, we find

$$\|x_{\mathfrak{K}}-x\|_{\mathcal{A}}\leq \sum_{s=1}^{d}C(\mathcal{A},b,\kappa_{s})\Big(rac{\sqrt{\kappa_{s}}-1}{\sqrt{\kappa_{s}}+1}\Big)^{k_{s}},$$

with $\kappa_s = 1 + \frac{\lambda_{\max}(A_s) - \lambda_{\min}(A_s)}{\lambda_{\min}(A)}$.

For the case A_s , k_s constant:

$$\|x_{\mathfrak{K}}-x\|_{\mathcal{A}} \leq C(\mathcal{A},b,d) \Big(rac{\sqrt{\kappa}-1}{\sqrt{\kappa}+1}\Big)^k,$$

with $\kappa = \frac{d-1}{d} + \frac{\kappa_2(\mathcal{A})}{d}$.

Note that the convergence rate improves with increasing dimension.

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Convergence, non-symmetric positive definite case

General convergence bound:

$$\|x_{\mathfrak{K}} - x\|_{2} \leq \sum_{s=1}^{d} \int_{0}^{\infty} e^{-\hat{\alpha}_{s}t} \|U_{s}e^{-tH_{s}}e_{1} - e^{-tA_{s}}b_{s}\|_{2} dt$$

with
$$\hat{\alpha}_s := \sum_{j \neq s} \alpha_j$$
 and $\alpha_j = \lambda_{\min}(A_j + A_j^{\top})/2$.

The bound on $||U_s e^{-tH_s} e_1 - e^{-tA_s} b_s||_2$ will depend on additional knowledge on A_s .

For example, when the field of values of each matrix A_s is contained in a known ellipse, an explicit convergence bound can be found.

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Grasedyck's method, system $\mathcal{H}y = ilde{b}$

For ${\mathcal H}$ positive definite:

$$\mathcal{H}^{-1} = \int_0^\infty \exp(-t\mathcal{H})dt$$

The exponential of \mathcal{H} has a Kronecker product structure, too:

$$\exp(-t\mathcal{H}) = \exp(-t\sum_{s=0}^{d}\hat{H}_{s}) = \prod_{s=1}^{d}\exp(-t\hat{H}_{s}) = \bigotimes_{s=1}^{d}\exp(-t\mathcal{H}_{s}),$$

Approximation *y_m*:

$$y_m = \sum_{j=1}^m \omega_j \bigotimes_{s=1}^d \exp\left(-\alpha_j H_s\right) \tilde{b}_s,$$

with certain coefficients α_i, ω_i .

Coefficients of the exponential sum (1)

The coefficients α_j, ω_j should minimize

$$\sup_{z\in\Lambda(\mathcal{H})}\Big|\frac{1}{z}-\sum_{j=1}^m\omega_je^{-\alpha_jz}\Big|$$

Case 1: \mathcal{H} symmetric and condition number known The eigenvalues of \mathcal{H} are real: $\Lambda(\mathcal{H}) \subset [\lambda_{\min}, \lambda_{\max}]$. There are coefficients α_j, ω_j s.t.

$$\|\mathbf{y} - \mathbf{y}_m\|_2 \leq C(\mathcal{H}, \tilde{b}) \exp\Big(\frac{-m\pi^2}{\log(8\kappa_2(\mathcal{H}))}\Big).$$

These coefficients may be found using a variant of the Remez algorithm. Tabellated values have been made available by Hackbusch.

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Coefficients of the exponential sum (2)

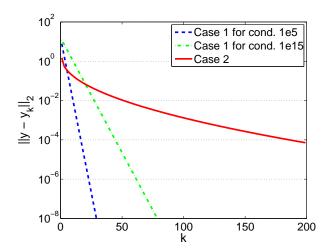
Case 2: Nonsymmetric \mathcal{H} and/or unknown condition number There is an explicit formula for α_j, ω_j , where the eigenvalues of \mathcal{H} only need to have positive real part.

$$\|\boldsymbol{y} - \boldsymbol{y}_{2m+1}\|_2 \le C(\mathcal{H}, \tilde{b}) \exp(\mu/\pi) \exp(-\sqrt{m}),$$

where $\mu = \max\{|\Im m(\Lambda(\mathcal{H}))|\}.$

The convergence is significantly slower than for Case 1.

Theoretical convergence bounds for the coefficients



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Symmetric example: Poisson Equation

Finite difference discretization of the Poisson equation in *d* dimensions:

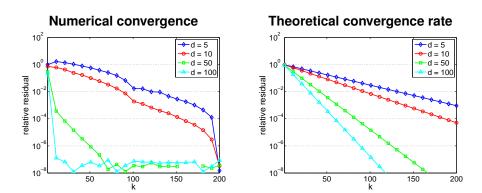
$$\Delta u = f$$
 in $\Omega = [0, 1]^d$
 $u = 0$ on $\Gamma := \partial \Omega$,

where the right-hand side *f* is a separable function.

$$A_{s} = \frac{1}{h^{2}} \begin{pmatrix} 2 & -1 & & \\ -1 & 2 & \ddots & \\ & \ddots & \ddots & -1 \\ & & -1 & 2 \end{pmatrix}, b_{s} : \text{random numbers.}$$

The approximation error is measured by *relative residual*, $\frac{\|Ax_{\bar{x}} - b\|_2}{\|b\|_2}$.

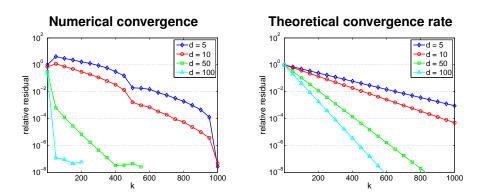
Convergence for system size 200^d



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Numerical experiments

Convergence for system size 1000^d



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Extended Krylov subspaces

Using extended Krylov subspaces

$$\widetilde{\mathcal{K}}_{k_s}(\textit{A}_{s},\textit{b}_{s}) := \operatorname{span}(\mathcal{K}_{k_s}(\textit{A}_{s},\textit{b}_{s}) \cup \mathcal{K}_{k_s+1}(\textit{A}_{s}^{-1},\textit{b}_{s})),$$

the algorithm works analogously. Convergence bound (A_s , k_s constant):

$$\|x_{\mathfrak{K}}-x\|_{2}\leq \mathcal{C}(\mathcal{A},b,d)\Big(rac{\sqrt{ ilde{\kappa}}-1}{\sqrt{ ilde{\kappa}}+1}\Big)^{k}.$$

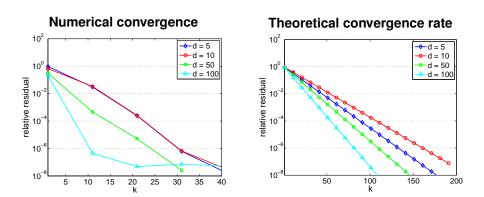
The convergence rate depends on $\tilde{\kappa}$:

$$\tilde{\kappa} \approx \frac{d-1}{d} + \frac{1}{d^2} (d-1)^{\frac{d-1}{d}} \kappa_2(\mathcal{A})^{\frac{d-1}{d}}$$
$$d = 2 : \tilde{\kappa} = \frac{1}{2} + \frac{\sqrt{\kappa_2(\mathcal{A})}}{4}, \qquad d \gg 0 : \tilde{\kappa} \approx \frac{d-1}{d} + \frac{\kappa_2(\mathcal{A})}{d}.$$

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Numerical experiments

Extended Krylov, system size 200^d



Non-symmetric example

Finite difference discretization in *d* dimensions of the convection-diffusion equation

$$\Delta u - (\boldsymbol{c}, \dots, \boldsymbol{c}) \nabla u = f \quad \text{in } \Omega = [0, 1]^d$$
$$u = 0 \quad \text{on } \Gamma := \partial \Omega,$$

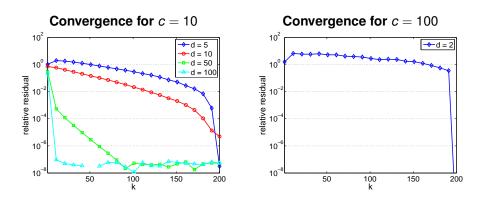
where f is again a separable function.

$$A_{s} = \frac{1}{h^{2}} \begin{pmatrix} 2 & -1 & & \\ -1 & 2 & -1 & & \\ & \ddots & \ddots & \ddots & \\ & & -1 & 2 & -1 \\ & & & & -1 & 2 \end{pmatrix} + \frac{c}{4h} \begin{pmatrix} 3 & -5 & 1 & & \\ 1 & 3 & -5 & \ddots & \\ & \ddots & \ddots & \ddots & 1 \\ & & 1 & 3 & -5 \\ & & & & 1 & 3 \end{pmatrix}$$

For c = 10, all eigenvalues of A are real. For c = 100, this is not the case, and the convergence is significantly worsened.

Numerical experiments

Non-symmetric case, system size 200^d



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Conclusions

- An efficient algorithm to calculate a low-rank approximation of the solution tensor
- The computational complexity is linear in the number of dimensions
- Only matrix-vector operations with the full system matrices are required
- A theoretical convergence bound was found.

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Conclusions

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- Only matrix-vector operations with the full system matrices are required
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Thank you for your attention!

Literature

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