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# **Linearity and Mobility Degradation in Strained Si MOSFETs with Thin Gate Dielectrics**

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**Abstract-**As gate dielectrics are scaled to a few atomic layers and the channel doping is increased to mitigate short channel effects, high vertical electric fields cause considerable mobility degradation through surface roughness scattering in silicon MOSFETs. This high field mobility degradation is known to influence the harmonic distortion through higher order drain current derivatives. Failure to take these higher order derivatives into account can cause significant error in the predictive evaluation of linearity ( $V_{IP3}$ ) in MOSFETs. Electrical measurements are used to extract the 2<sup>nd</sup> order mobility degradation factor ( $\theta_2$ ) from strained silicon MOSFETs with different germanium contents. Linearity and high-field mobility degradation are shown to be independent of strain in spite of atomic force microscopy measurements showing that the surface roughness root-mean-square amplitude increases with the germanium content. It is also shown that  $\theta_2$  is required for the accurate modelling of linearity. The impact of oxide thickness on linearity is also investigated through  $\theta_2$ . In this paper, an analytical relationship between  $\theta_2$  and the effective oxide thickness is developed and validated by electrical measurements on MOSFETs with different oxide thicknesses and  $\theta_2$  values from the literature. Using the extracted  $\theta_2$  values as inputs to analytical MOSFET models, a correlation between the oxide thickness and linearity is analyzed.

***Index Terms*** – Distortion, Linearity, Mobility degradation, Strained Silicon.

## INTRODUCTION

The miniaturization of the metal oxide on semiconductor field effect transistor (MOSFET) has made complimentary metal oxide on semiconductor (CMOS) devices considerable radio frequency (RF) contenders where bipolars and high-electron-mobility-transistors are traditionally dominant [1, 2]. Strain engineering in deep submicrometer CMOS devices has further improved the high speed performance required for RF implementation [3]. However in analog/mixed-signal applications, other metrics like noise and linearity are important [4-6]. Linearity is particularly of interest in CMOS devices since MOSFETs exhibit better linearity than bipolars [2]. In a perfectly linear MOSFET, a signal with a single harmonic at the input would yield an output signal at the same frequency. However, since MOSFETs are not perfectly linear, the output signal usually contains higher order harmonics that may interfere with the fundamental. These higher order harmonics are related to the higher order derivatives of the drain current with respect the terminal voltages according to

$$i_{DS} \approx g_m v_{GS} + \frac{\partial g_m}{\partial v_{GS}} v_{GS}^2 + \frac{\partial^2 g_m}{\partial v_{GS}^2} v_{GS}^3 + \dots \dots \frac{\partial^{n-1} g_m}{\partial v_{GS}^{n-1}} v_{GS}^n \quad (1)$$

where  $i_{DS}$  is the drain-source current,  $g_m$  is the transconductance and  $v_{GS}$  is the small signal gate-source voltage. Distortion analysis is important for analog MOSFETs since higher order harmonics at the output can cause interference with fundamental harmonic thereby resulting in inter-modulation distortion and degrading the signal integrity of the system. Of the higher order harmonics, the most important is the 3<sup>rd</sup> order harmonic due to the fact that a signal close in frequency to the fundamental signal at the input of a MOSFET will have a third order intermodulation harmonic at

the output of the MOSFET that will be close in frequency to the fundamental. Various figures of merit have been developed to quantify the linearity of RF transistors. The third order intercept point ( $IP3$ ) is defined as the input power level at which the fundamental harmonic and the third order harmonic have equal power levels at the output. This is usually determined from extrapolations of RF power measurements. Another indicator of linearity is the  $V_{IP3}$  which is defined as the extrapolated gate voltage ( $V_{GS}$ ) bias at which the amplitudes of the 1<sup>st</sup> and 3<sup>rd</sup> order derivatives of  $I_{DS}$  are equal [2, 7-11]. The  $V_{IP3}$  can be measured from MOSFET DC  $I_{DS}$  vs.  $V_{GS}$  characteristics and is given by

$$V_{IP3} = \sqrt{24 \left[ \frac{\partial g_m}{\partial V_{GS}} \right] \times \left[ \frac{\partial^2 g_m}{\partial V_{GS}^2} \right]^{-1}} \quad (2)$$

It has been shown that accounting for mobility degradation accurately is essential for understanding the distortion characteristics of MOSFETs [7]. Since applications that require low distortion use strongly inverted MOSFETs (where linearity is highest), surface roughness scattering becomes the important mobility limiting mechanism in the determination of linearity. Gate dielectric scaling and higher substrate doping has increased the transverse electric fields and hence, mobility degradation from surface roughness scattering. A 2<sup>nd</sup> order mobility degradation model was developed to account for mobility reduction related to  $(V_{GS} - V_{TH})^2$  at high vertical fields where  $V_{TH}$  is the threshold voltage [12]. The effective mobility,  $\mu_{EFF}$  can be expressed as

$$\mu_{EFF} = \frac{\mu_O}{1 + \theta_1(V_{GS} - V_{TH}) + \theta_2(V_{GS} - V_{TH})^2} \quad (3)$$

where  $\theta_1$  is the 1<sup>st</sup> order mobility degradation factor (MDF),  $\theta_2$  is the 2<sup>nd</sup> order MDF and  $\mu_O$  is the low field mobility. Previous accounts of modelling mobility degradation from surface roughness scattering used only the 1<sup>st</sup> order MDF model which can be derived from equation 3 by assuming  $\theta_2=0$  [13-15]. However, as the gate dielectric is thinned, the amplitude of the semiconductor surface asperities at the dielectric interface becomes a larger proportion of the total dielectric thickness hence 2<sup>nd</sup> order effects become non-negligible. The model in equation 3 was shown to account more accurately for mobility degradation in thin gate dielectric MOSFETs and has been used in other studies [16-18]. In this paper, the linearity of strained Si nMOSFETs is assessed, the 2<sup>nd</sup> order MDF ( $\theta_2$ ) is extracted for 2.5 nm thick oxides and a MOSFET model is used to assess the importance of  $\theta_2$  in the predictive modelling of linearity for CMOS devices. Also, an analytical model is developed relating  $\theta_2$  to the oxide thickness and is validated by experimental measurements of  $\theta_2$  together with values taken from literature.

## DEVICE FABRICATION

Si control and strained Si nMOSFETs on Si<sub>1-x</sub>Ge<sub>x</sub> strain relaxed buffers (SRBs) with x=0.15, 0.20 and 0.25 are co-fabricated in a CMOS process flow. The SiGe SRBs were grown by low-pressure chemical vapour deposition with SiH<sub>4</sub> and GeH<sub>4</sub> as pre-cursors. The total thickness of the SiGe SRBs was 4 µm with the compositional grading done to a thickness of 2.5 µm and a 1.5 µm layer of constant composition SiGe upon which 10 nm of tensile strained Si layers are deposited. The

tensile strain is confirmed by Raman spectroscopy. Atomic force microscopy (AFM) measurements with different scan sizes were performed on the strained Si wafers. Table 1 shows the results of the AFM scan. The rms roughness amplitude is proportional to the germanium content. It is known that cross-hatching resulting from strain relaxation in the underlying virtual substrate increases the surface roughness amplitude in strained Si MOSFETs hence, these results are not surprising [19, 20]. The gate oxide is grown by thermal oxidation to an effective oxide thickness (EOT) of 2.5 nm which is confirmed by gate-bulk capacitance measurements. The gate polysilicon was deposited; after which source-drain arsenic LDD and boron anti-punch-through pocket implantation was performed. The next process steps were the silicon nitride side wall spacer formation, source-drain HDD implantation and cobalt silicidation done at 1000 °C for 30s. Contact was made with the MOSFET through tungsten plugs with titanium nitride barriers lining the contact via. The back end metallisation consisted of aluminium and copper.

## **RESULTS AND DISCUSSION**

Using the split CV technique with series resistance corrections, the effective mobility is extracted for the Si control and strained Si nMOSFETs [21]. Fig. 1 shows the effective mobility as a function of the vertical effective field for all the devices. It can be seen from Fig. 1 that the effective mobility is proportional to the Ge content (strain) over the entire range of vertical field. The gate-channel capacitance and gate-bulk capacitance measurements used in the calculation of the inversion charge density and vertical effective fields were done on 10  $\mu\text{m}$  ( $W$ ) by 10  $\mu\text{m}$  ( $L_G$ ) nMOSFETs biased at 25 mV drain voltage.



The 2<sup>nd</sup> order mobility degradation factor is extracted using a technique introduced by McLarty et al [12]. The starting point of the technique is the drain current linear model (low  $V_{DS}$  and high  $V_{GS}$ ) for strongly inverted MOSFETs

$$I_{DS} = \frac{W\mu_{EFF}C_{OX}}{L_G}(V_{GS} - V_{TH})V_{DS} \quad (4)$$

where  $W$  is the gate width,  $L_G$  is the gate length,  $C_{OX}$  is the gate dielectric capacitance density and  $V_{DS}$  is the drain voltage. Substituting equation (3) into (4), the 1<sup>st</sup> order and 2<sup>nd</sup> order derivatives of the inverse of the drain current ( $I_{DS}^{-1}$ ) with respect to  $V_{GS}$  is calculated as

$$\frac{\partial}{\partial V_{GS}} \frac{1}{I_{DS}} = \frac{L_G}{W\mu_0 C_{OX} V_{DS}} \left( \theta_2 - \frac{1}{(V_{GS} - V_{TH})^2} \right) \quad (5)$$

$$\left( \frac{\partial^2}{\partial V_{GS}^2} \frac{1}{I_{DS}} \right)^{-\frac{1}{3}} = \left( \frac{W\mu_0 C_{OX} V_{DS}}{2L_G} \right)^{\frac{1}{3}} (V_{GS} - V_{TH}) \quad (6)$$

$\theta_2$  is extracted from equation 5 as the x-axis intercept of the linear plot of  $\partial(I_{DS}^{-1})/\partial V_{GS}$  against  $(V_{GS}-V_{TH})^{-2}$  whereas  $\mu_0$  is extracted from the slope of the straight line. The threshold voltage can be extracted by calculating the 2<sup>nd</sup> derivative of the inverse of the drain current ( $I_{DS}^{-1}$ ) with respect to  $V_{GS}$  and plotting  $(\partial^2(I_{DS}^{-1})/\partial V_{GS}^2)^{-1/3}$  as a linear function of  $V_{GS}$ . The threshold voltage is the intercept of the extrapolated linear plots with the  $V_{GS}$  axis.

Fig. 2(a) and Fig. 2(b) show the plots of  $\partial(I_{DS}^{-1})/\partial V_{GS}$  against  $(V_{GS}-V_{TH})^{-2}$  for 0.5  $\mu\text{m}$  ( $L_G$ ) Si control and strained Si nMOSFETs biased at 50 mV  $V_{DS}$ . A low drain

voltage is used because the MOSFET channel must be biased in strong inversion (linear mode) for the model in equation 4 to be applicable. It can be seen from Fig. 2(a) that the slopes of the lines reduce as the Ge composition (or strain content) increases. This observation correlates with equation 5 where it can be seen that the slope  $L_G/(W\mu_O C_{OX} V_{DS})$  is inversely proportional to the low field mobility. The extracted low field mobility, shown in table 2, is proportional to strain and is on average 100% higher than the effective mobility. The 2<sup>nd</sup> order MDF and low field mobility are respectively extracted from the x-axis intercepts and slopes in Fig. 2 (a). Fig. 2(b) is an enlarged version of Fig. 2(a), showing how the values of  $\theta_2$  are read off from the x-axis intercepts. When calculating the slope and intercept from Fig. 2(a), it is important to only consider measurement points at least 500 mV above the threshold voltage. This corresponds to  $(V_{GS}-V_{TH})^2 < 4 V^{-2}$  in Fig. 2. This is because the MOSFET model in (4) applies to only MOSFETs in strong inversion and mobility degradation from the vertical effective field applies only at high  $V_{GS}$ . It can be seen from Fig. 2(b) that the values of  $\theta_2$  vary from 0.28 V<sup>-2</sup> to 0.29 V<sup>-2</sup> for the Si control and strained Si nMOSFETs.  $\theta_1$  is also calculated by substituting the experimentally extracted values of  $\mu_{EFF}$ ,  $\mu_O$  and  $V_{GS}-V_{TH}$  (which is approximately 1.5 V at an  $E_{EFF}$  of 1.8 MVcm<sup>-1</sup>) into equation 3. The values of  $\theta_1$ ,  $\theta_2$  and  $\mu_{EFF}$  can also be seen in table 2.

Fig. 2(b) and Table 2 show that the values of  $\theta_2$  for the Si control and strained Si nMOSFETs are within 2% of each other, hence it can be assumed that  $\theta_2$  is independent of strain. Since the MOSFETs under investigation here are co-processed (the same EOT, pocket implant and substrate doping), the only factor that may influence mobility degradation differently would be the different surface morphological properties (rms roughness amplitude and correlation length). However, in spite of the increasing AFM measured surface roughness with Ge content (strain),

the mobility enhancement is maintained in the strained Si devices. High mobility enhancement in strained Si MOSFETs at high vertical fields (in the surface-roughness scattering regime) has led researchers to propose that strained Si MOSFETs have “smoother” surfaces compared with Si control MOSFETs [22, 23]. Lower surface roughness amplitudes and longer correlation lengths were required to match simulated high field mobilities with measured high field mobilities in strained Si MOSFETs [22]. Smoother surfaces in strained silicon are assumed because the mobility enhancement mechanism (reduced phonon scattering from  $\Delta_2$ - $\Delta_4$  conduction band splitting) can explain mobility enhancement at medium vertical fields (where phonon scattering is the dominant mobility limiting mechanism) but not at high vertical fields (where surface roughness scattering is the dominant mobility limiting mechanism). Furthermore, the band-splitting due to quantum confinement at high vertical fields renders the band splitting from tensile strain redundant [23]. Although the measurements here indicate that mobility degradation is independent of strain, more advanced MOSFET models than (1) and more advanced mobility models than (2) would be needed to understand the impact of strain/germanium-related surface morphological properties on the high field mobility in strained Si layers.

Fig. 3 shows the extraction of  $V_{TH}$  for the Si control and strained Si nMOSFETs based on equation 6. Since, the devices are process matched, it is expected that the  $V_{TH}$  will reduce as the Ge content in the  $\text{Si}_{1-x}\text{Ge}_x$  SRB increases. This is due to the increase in the electron affinity with tensile strain. Table 2 also shows the  $V_{TH}$  values extracted from Fig. 3. The  $V_{TH}$  values extracted from Fig. 3 agree well with  $V_{TH}$  values extracted using the linear transconductance method.

The linearity ( $V_{IP3}$ ) of the Si control and strained Si nMOSFETs is measured according to (2). The transconductance ( $g_m$ ) and the 2<sup>nd</sup> derivative of the

transconductance ( $g_{m3}$ ) are used to calculate the linearity. Fig. 4 shows the  $V_{IP3}^2$  as functions of the gate voltage overdrive for the 0.5  $\mu\text{m}$  gate length Si control and strained Si nMOSFETs with 50 mV  $V_{DS}$ . The characteristic peak in the  $V_{IP3}^2$  plot of Fig. 4 indicates when the 2<sup>nd</sup> order derivative of the transconductance ( $g_{m3}$ ) crosses the zero mark [8, 9].  $V_{IP3}$  is low in weak inversion (left side of the peak) due to the exponential relationship between  $I_{DS}$  and  $V_{GS}$  ( $I_{DS}$  is a diffusion current).  $V_{IP3}^2$  increases in strong inversion because the exponential relationship between  $I_{DS}$  and  $V_{GS}$  becomes parabolic ( $I_{DS}$  is a drift current). The peaks in the  $V_{IP3}^2$  characteristics of the measured devices in Fig. 4 occurs at  $V_{GS}=V_{TH}$ , thereby indicating the change from a diffusion current to a drift current i.e. the onset of moderate inversion or the triode part of the  $I_{DS}$  vs.  $V_{GS}$  characteristic. As can be seen from Fig. 4, there is no difference in the  $V_{IP3}^2$  characteristics hence  $V_{IP3}$  is independent of strain. This is not surprising since mobility is proportional to both  $g_m$  and  $g_{m3}$ , hence, it is removed by the ratio in equation 2. Also, it has already been shown in Fig. 2 that mobility degradation is independent of strain. It has been shown that  $V_{IP3}^2$  depends on the oxide thickness and substrate doping, (i.e. the body factor) both of which affect the vertical electric field [5]. Fig. 5 shows the measured  $g_m$  and  $g_{m3}$  for the Si control MOSFET illustrating the positive and negative peak of  $g_{m3}$ . The region to the right of the maximum transconductance in Fig. 5 is dominated by surface roughness scattering hence, the magnitude of mobility degradation determines the slope of the negative transconductance.

The importance of taking the 2<sup>nd</sup> order MDF into account in the predictive modelling of  $V_{IP3}^2$  is shown by using a single-piece semi-empirical MOSFET model which is valid from weak inversion to strong inversion [9, 24].

$$I_{DS} = I_{D0} \left[ \left( \ln \left( 1 + \exp \left( \frac{V_{GS} - V_{TH}}{2m\Phi_{TH}} \right) \right) \right)^2 - \left( \ln \left( 1 + \exp \left( \frac{V_{GS} - V_{TH} - mV_{DS}}{2m\Phi_{TH}} \right) \right) \right)^2 \right] \quad (7)$$

$$\text{where } I_{D0} = \frac{2W\mu_o C_{ox} \Phi_{TH}^2 m}{L_G (1 + \theta_1 (V_{GS} - V_{TH}) + \theta_2 (V_{GS} - V_{TH})^2)}$$

$m$  is the body factor ( $m=1.16$  for an oxide thickness of 2.5 nm and for a substrate doping of  $5 \times 10^{17} \text{ cm}^{-3}$ ),  $L_G$  is 500 nm,  $V_{DS}=50$  mV and  $\Phi_{TH}$  is the 300 K thermal voltage (26 mV). Fig. 6 shows  $g_m$  and  $g_{m3}$  calculated using the model in equation 7 with  $\theta_1=0.6 \text{ V}^{-1}$  and  $\theta_2=0.3 \text{ V}^{-2}$ . The similarity between the calculated characteristics of Fig. 6 and the measured characteristics of Fig. 5 was obtained by tuning  $\theta_2$  to the measured value. Fig. 7 shows a comparison between the measured  $V_{IP3}^2$  characteristic and the calculated  $V_{IP3}^2$  characteristics with  $\theta_2=0 \text{ V}^{-2}$  and  $\theta_2=0.3 \text{ V}^{-2}$ . It can be seen in Fig. 7 that the calculated characteristic matches the measured characteristic only with  $\theta_2$  taken into account ( $\theta_2=0.3 \text{ V}^{-2}$ ). The calculated  $V_{IP3}$  with  $\theta_2=0 \text{ V}^{-2}$  exhibits high  $V_{IP3}$  in strong inversion which is not representative of the measured characteristic. Fig. 8 shows the calculated  $V_{IP3}^2$  as a function of  $V_{GS}$  for different values of  $\theta_2$ . It can be seen in Fig. 8 that  $V_{IP3}^2$  reduces significantly as  $\theta_2$  increases. Fig. 7 and Fig. 8 show that it is necessary, in the predictive modelling of linearity, to take  $\theta_2$  into account as linearity can easily be over-estimated.

The effective mobility of a carrier in a MOSFET is determined by coulomb scattering from ionized impurities, phonon scattering from lattice vibrations and surface roughness scattering from surface asperities at the Si/SiO<sub>2</sub> interface [25]. This can be expressed, according to the Matthiessen' rule, as

$$\mu_{EFF} = \frac{\mu'_O}{1 + \theta_{PH} E_{EFF}^{1/3} + \theta_{SR} E_{EFF}^n} \quad (8)$$

where  $\theta_{PH}$  and  $\theta_{SR}$  are empirical parameters [7]. The interface between the silicon channel and the gate dielectric is not atomically smooth; hence, carriers scatter against surface asperities in strong inversion conditions. These Si channel surface asperities can be characterised by the root-mean-square roughness and the correlation length [26]. Surface roughness can be modelled as variations in oxide thickness, hence carriers transiting along the channel will be perturbed by a change in potential that is proportional to the average roughness amplitude [26, 27]. According to the Fermi golden rule, the scattering rate is proportional to the square of the perturbation potential resulting from the surface roughness [28]. Under high vertical fields,  $n$  in equation 8 is usually 2 for electrons and 1 for holes [7, 25, 26, 28, 29]. Under such conditions, equation 8 can be re-written for electrons as

$$\mu_{EFF} = \frac{\mu'_O}{1 + \theta_{SR} E_{EFF}^2} \quad (9)$$

The vertical effective field for a MOSFET in inversion can be expressed as

$$E_{EFF} = \frac{1}{\epsilon_{Si}} (Q_B + \eta C_{OX} (V_{GS} - V_{TH})) \quad (10)$$

$$\text{where } Q_B = \sqrt{4q\epsilon_{Si}N_A \left( \frac{k_B T}{q} \ln \left( \frac{n_i}{N_A} \right) \right)}$$

$Q_B$  is the depletion charge density ( $0.39 \mu\text{C.cm}^{-3}$  for a substrate doping,  $N_A$ , of  $5 \times 10^{17} \text{cm}^{-3}$ ),  $\epsilon_{Si}$  is the dielectric constant of silicon,  $C_{OX}$  is the gate dielectric capacitance density ( $1.42 \mu\text{F.cm}^{-2}$  for an EOT of 2.5 nm) and  $\eta$  is 0.5. Substituting equation 10 into 9, the effective mobility can be re-written in the form of equation 3 with the following expressions for the low field mobility, the 2<sup>nd</sup> and 1<sup>st</sup> order MDF.

$$\mu_o = \mu_o' \left( 1 + \frac{\theta_{SR}}{\epsilon_{Si}^2} Q_B^2 \right)^{-1} \quad (11)$$

$$\theta_2 = \frac{\theta_{SR}}{\epsilon_{Si}^2} \eta^2 C_{OX}^2 \left( 1 + \frac{\theta_{SR}}{\epsilon_{Si}^2} Q_B^2 \right)^{-1} \quad (12)$$

$$\theta_1 = 2 \frac{\theta_{SR}}{\epsilon_{Si}^2} Q_B \eta C_{OX} \left( 1 + \frac{\theta_{SR}}{\epsilon_{Si}^2} Q_B^2 \right)^{-1} \quad (13)$$

In the pioneering work of mobility degradation analysis by Fu [13], a similar expression to equation 13 was developed for  $\theta_I$  as shown below

$$\theta_1 = \frac{B}{4\epsilon_{Si}} C_{OX} \left( 1 + \frac{B}{\epsilon_{Si}} Q_B \right)^{-1}$$

where  $B = 2bql/3KT$ .

$l$  is the mean free path,  $T$  is the temperature,  $K$  is Boltzmann's constant and  $b$  is a constant that depends on the ratio of the mean free path to the inversion layer depth.  $B$  in [13] was found to be  $6.41 \times 10^{-6} \text{ cm.V}^{-1}$  assuming a mean free path and inversion layer depth of 5 nm. Comparing equation 13 with the formulation in [13], the following expression can be derived for  $\theta_{SR}$

$$\theta_{SR} = N \left( \frac{2bql}{3\eta KT} \right) \frac{\epsilon_{Si}}{Q_B} \quad (14)$$

$N$  accounts for the numerical constant difference between the derivation of  $\theta_I$  in this study and that in [13]. Solving equation 14 using the value of  $B$  calculated in [13],  $\theta_{SR}$  is calculated to be approximately  $17 \times 10^{-12} \text{ cm}^2 \cdot \text{V}^{-2}$ . Substituting  $\theta_{SR} = 17 \times 10^{-12} \text{ cm}^2 \cdot \text{V}^{-2}$  into equations 12 and 13, solving for  $\theta_I$  and  $\theta_2$  yields  $2.5 \text{ V}^{-1}$  and  $2.3 \text{ V}^{-2}$  respectively. This is much higher than what was measured. Adjusting  $N$  to 0.34 and  $\theta_{SR}$  to  $0.8 \times 10^{-12} \text{ cm}^2 \cdot \text{V}^{-2}$ , solving for  $\theta_I$  and  $\theta_2$  yields  $0.6 \text{ V}^{-1}$  and  $0.3 \text{ V}^{-2}$  respectively (much closer to the measured values). For  $\theta_{SR} = 0.8 \times 10^{-12} \text{ cm}^2 \cdot \text{V}^{-2}$ ,  $Q_B^2 \theta_{SR} / \epsilon_{Si}^2 \ll 1$ , hence, equations 11, 12 and 13 can be approximated as

$$\mu_O = \mu_O' \quad (15)$$

$$\theta_2 = \frac{\theta_{SR}}{\epsilon_{Si}^2} \eta^2 C_{OX}^2 \quad (16)$$

$$\theta_1 = 2 \frac{\theta_{SR}}{\epsilon_{Si}^2} Q_B \eta C_{OX} \quad (17)$$

It can be observed that in equations 12 and 13 that  $\theta_I$  and  $\theta_2$  are inversely related to the gate dielectric thickness. This is expected since the vertical effective field increases as the gate dielectric is scaled. What is also interesting to note is that the dielectric constant of the gate insulator is directly proportional to  $\theta_I$  and  $\theta_2$ . This is important for advanced technology nodes that used high-k gate dielectrics. To experimentally characterise the dependence of  $\theta_2$  on  $t_{OX}$ ,  $\theta_2$  is extracted from a 1.4 nm and a 6 nm thick gate oxide nMOSFET. Although the MOSFETs were not



fabricated in the same process flow, the comparison serves as the first approximation of the  $\theta_2$  vs.  $t_{OX}$  relationship. Fig. 9 shows the measured  $\theta_2$  as a function of  $t_{OX}$  for the 1.4 nm, 2.5 nm and 6 nm nMOSFETs together with measurements taken from 3.5 nm and 10 nm gate oxide MOSFETs in literature [12]. It should be noted that since these MOSFETs are not co-fabricated, they will exhibit different  $\theta_{SR}$  values due to different surface roughness amplitudes and correlation lengths i.e.  $\theta_{SR}$  will be unique for each process. However, a universal trend can be observed with  $\theta_2$  increasing as  $t_{OX}$  is reduced. Also shown in Fig. 9 is the  $\theta_2$  vs.  $t_{OX}$  calculated by equation 12 with  $\theta_{SR} = 1 \times 10^{-12} \text{ cm}^2 \cdot \text{V}^{-2}$  and  $Q_B = 0.39 \text{ } \mu\text{C} \cdot \text{cm}^{-3}$ . It can be seen in Fig. 9 that there is good agreement between the measured and modelled  $\theta_2$  vs.  $t_{OX}$  relationship.

The linearity of the 6 nm and 1.4 nm oxide MOSFETs is also measured. It has been shown in Fig. 7 and Fig. 8 that  $V_{IP3}$  reduces as  $\theta_2$  increases. Fig. 10 shows the  $V_{IP3}^2$  as a function of  $V_{GS} - V_{TH}$  for 1.4 nm, 2.5 nm and 6 nm oxide MOSFETs. It can be seen in Fig. 10 that  $V_{IP3}^2$  reduces with the oxide thickness. At  $V_{GS} - V_{TH} = 1.25 \text{ V}$ , there is a difference of at least an order of magnitude between the  $V_{IP3}^2$  of the MOSFETs.

## CONCLUSIONS

The results in this paper show the importance of the 2<sup>nd</sup> order MDF in the predictive modelling of linearity in ultra-thin gate dielectric MOSFETs. The 2<sup>nd</sup> order MDF and low field mobility has been extracted from strained Si MOSFETs with 2.5 nm silicon dioxide gate dielectrics and varying germanium composition. The results indicate that mobility degradation and linearity is independent of strain and is dominated by oxide thickness. A semi-empirical MOSFET model calibrated with the measured parameters is used to further understand the impact of the 2<sup>nd</sup> order MDF on linearity. It was shown that linearity reduces substantially as the 2<sup>nd</sup> order MDF is

increased. An analytical relationship between the 2<sup>nd</sup> order MDF and the gate dielectric thickness is developed. It is shown that the 2<sup>nd</sup> order MDF is inversely proportional to the oxide thickness. Measurements of the 2<sup>nd</sup> order MDF from MOSFETs with different oxide thicknesses and values taken from literature are used to validate the accuracy of the model. The model also indicates that linearity will be reduced in MOSFETs with high-k dielectrics due to the increased dielectric constant. MOSFET models must take account of the 2<sup>nd</sup> order MDF as these results show that linearity can easily be over-estimated.

### **ACKNOWLEDGEMENT**

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TABLE 1  
AFM MEASUREMENTS

AFM Scan size	Ge=15%		Ge=20%		Ge=25%	
	RMS (nm)		RMS (nm)		RMS (nm)	
	Max	St dev	Max	St dev	Max	St dev
1x1 $\mu\text{m}^2$	0.20	0.01	0.23	0.04	0.23	0.04
3x3 $\mu\text{m}^2$	0.21	0.02	0.40	0.12	0.43	0.06
10x10 $\mu\text{m}^2$	0.28	0.05	0.39	0.09	0.83	0.14
30x30 $\mu\text{m}^2$	0.40	0.09	0.61	0.15	0.89	0.10
100x100 $\mu\text{m}^2$	0.70	0.05	0.90	0.09	1.20	0.22

TABLE 2  
PARAMETERS EXTRACTED

Device	$\mu_O$ ( $\text{cm}^2\text{V}^{-1}\text{S}^{-1}$ )	$\mu_{EFF}$ ( $\text{cm}^2\text{V}^{-1}\text{S}^{-1}$ )	$\theta_2$ ( $V^{-2}$ )	$\theta_I$ ( $V^{-1}$ )	$V_{TH}$ (V)
Si	230	125	0.28	0.63	0.40
Ge=15%	515	250	0.29	0.59	0.34
Ge=20%	569	280	0.28	0.61	0.30
Ge=25%	594	310	0.28	0.57	0.28

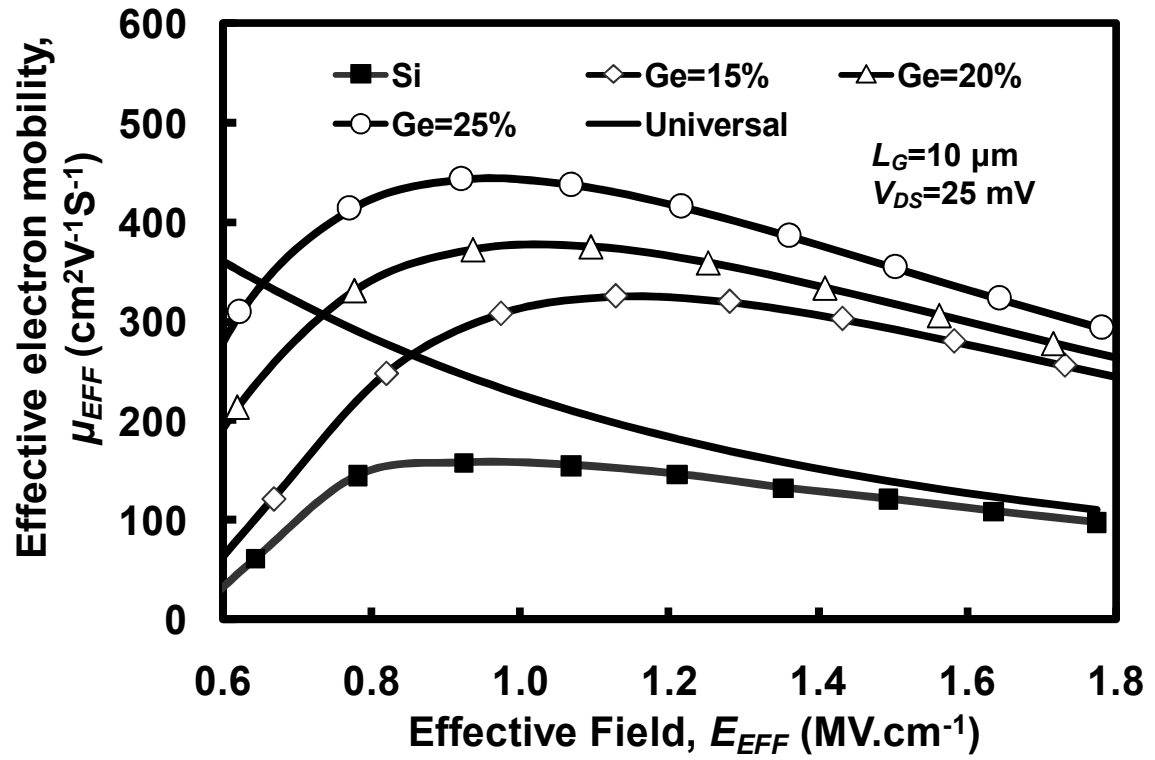


Fig. 1. The effective mobility as a function of the vertical effective field for the Si control and strained Si nMOSFETs on  $\text{Si}_{1-x}\text{Ge}_x$  SRBs for  $x=0.15, 0.20$  and  $0.25$ . The effective mobility is proportional to strain.

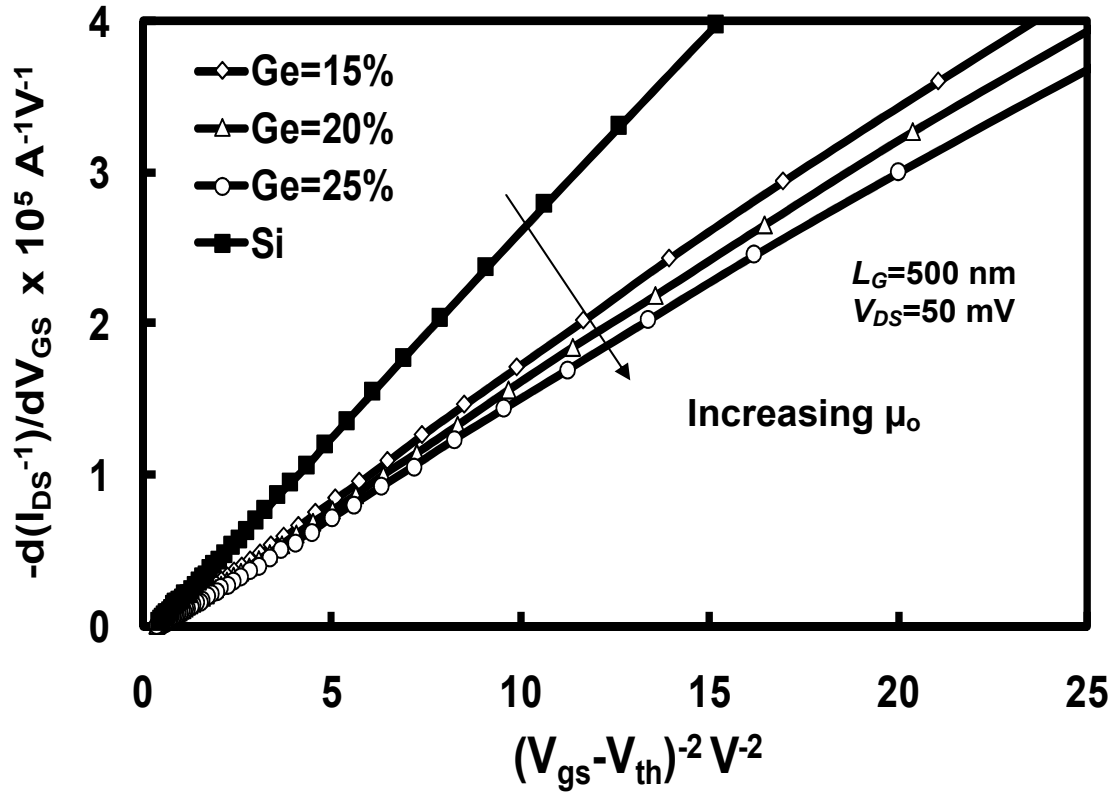


Fig. 2(a).  $\partial(I_{DS}^{-1})/\partial V_{GS}$  plotted as a function of  $(V_{GS}-V_{TH})^{-2}$  for Si control and strained Si nMOSFETs. The low field mobility is inversely proportional to the slope of the line and  $\theta_2$  can be extracted from the x-axis intercepts.

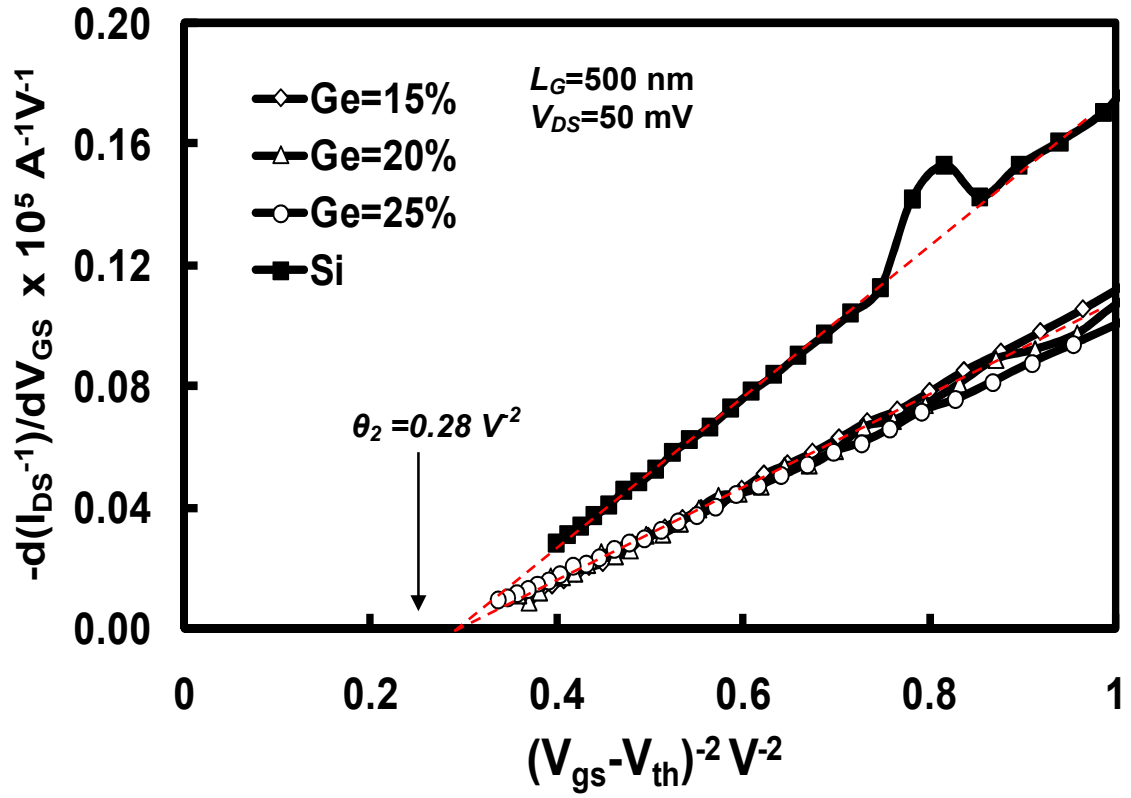


Fig. 2(b). An enlarged version of Fig. 2(a) showing the intercepts of the straight lines with the x axis. The 2<sup>nd</sup> order MDFs extracted for all the MOSFETs is approximately 0.29 V<sup>-2</sup>.



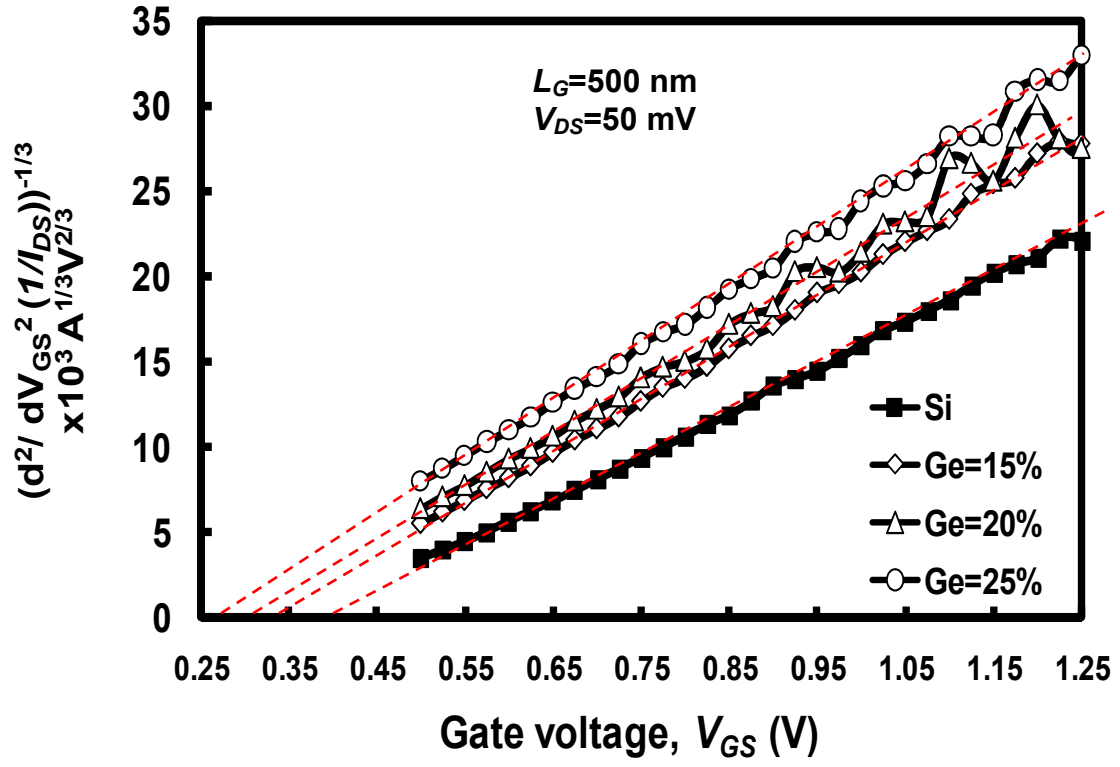


Fig. 3. The extraction of the threshold voltage from the x axis intercepts. The  $V_{TH}$  extracted reduces as the strain content increases.  $V_{TH}$  extracted agrees with values extracted using linear transconductance.

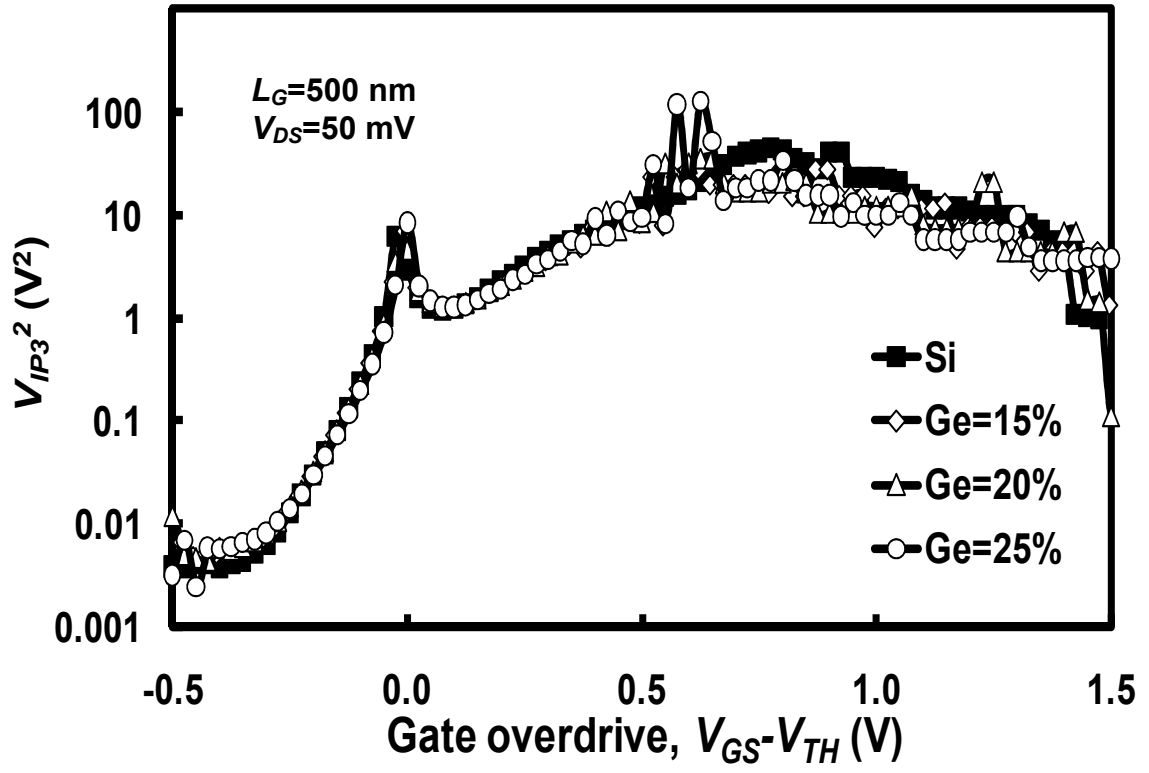


Fig. 4.  $V_{IP3}^2$  as a function of the gate voltage overdrive for strained Si and Si control nMOSFETs. There is no impact of strain on the  $V_{IP3}^2$ .

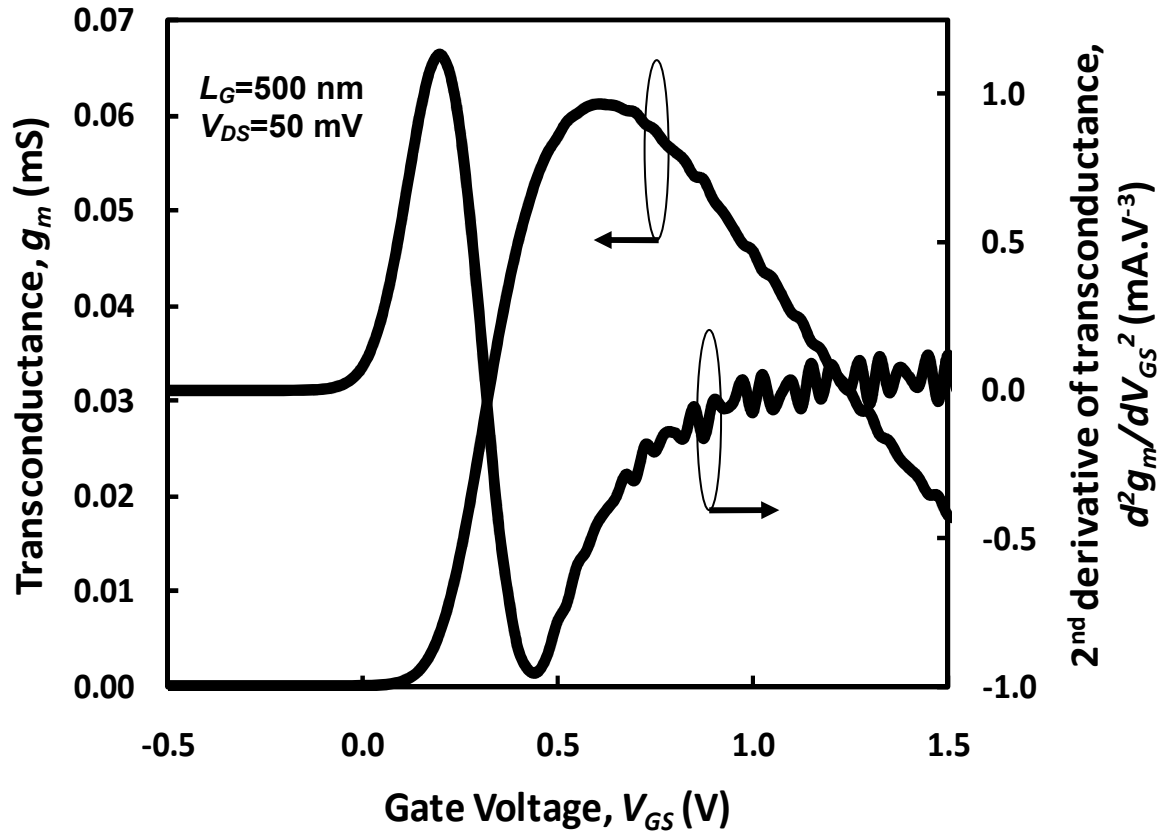


Fig. 5. Measured  $g_m$  and  $\delta^2 g_m / \delta V_{GS}^2$  as functions of  $V_{GS} - V_{TH}$  for the Si control MOSFET.

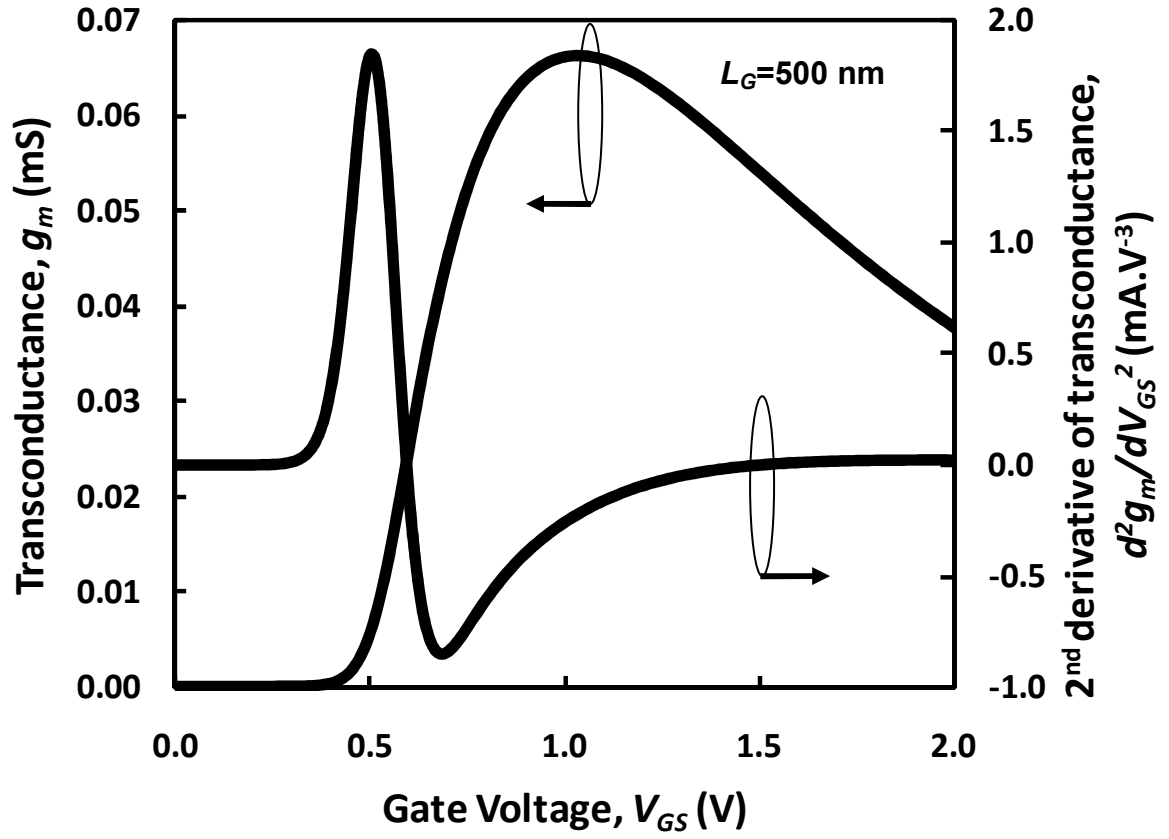


Fig. 6. Calculated  $g_m$  and  $\delta^2 g_m / \delta V_{GS}^2$  as functions of  $V_{GS} - V_{TH}$ . The semi-empirical MOSFET model is calibrated with measured parameters. A  $\theta_2$  of  $0.3 \text{ V}^{-2}$  is required to match the shape of the calculated characteristics with the measured one.

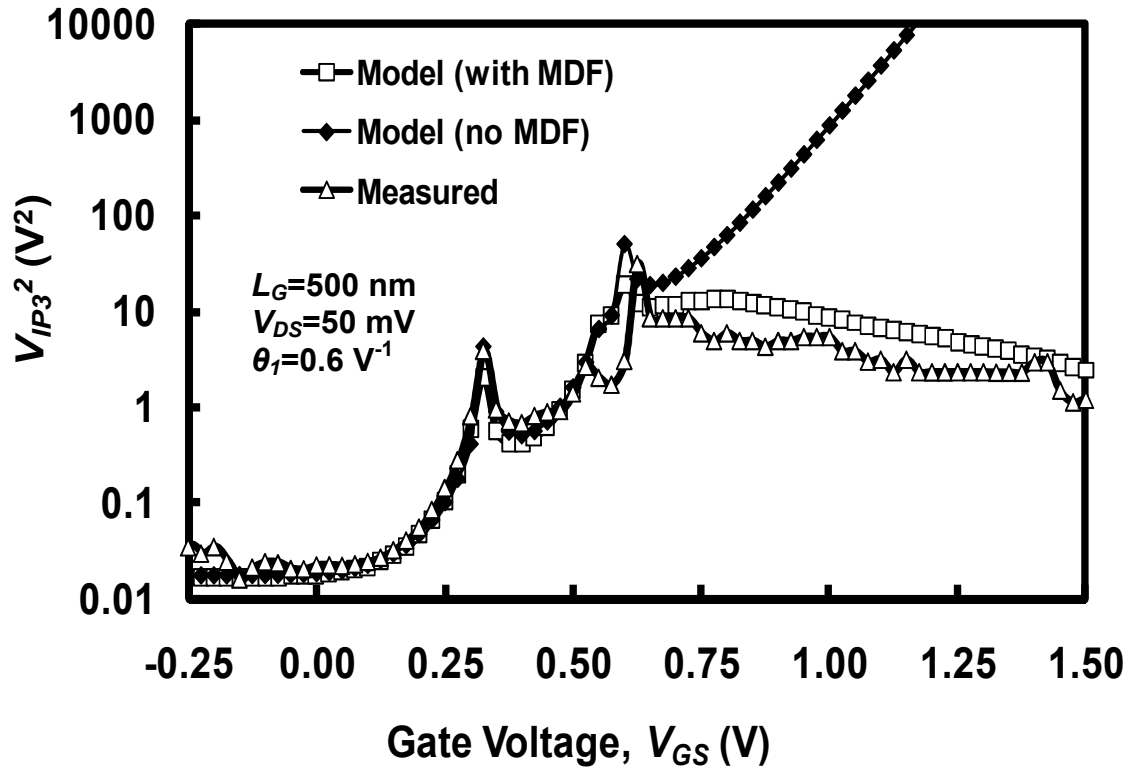


Fig. 7. A comparison of measured  $V_{IP3}^2$  characteristics and calculated  $V_{IP3}^2$  characteristics with  $\theta_2=0.28$  V<sup>-2</sup> and  $\theta_2=0$  V<sup>-2</sup>. There is a match between the calculated  $V_{IP3}^2$  and measured  $V_{IP3}^2$  only when  $\theta_2$  is tuned to the measured value.

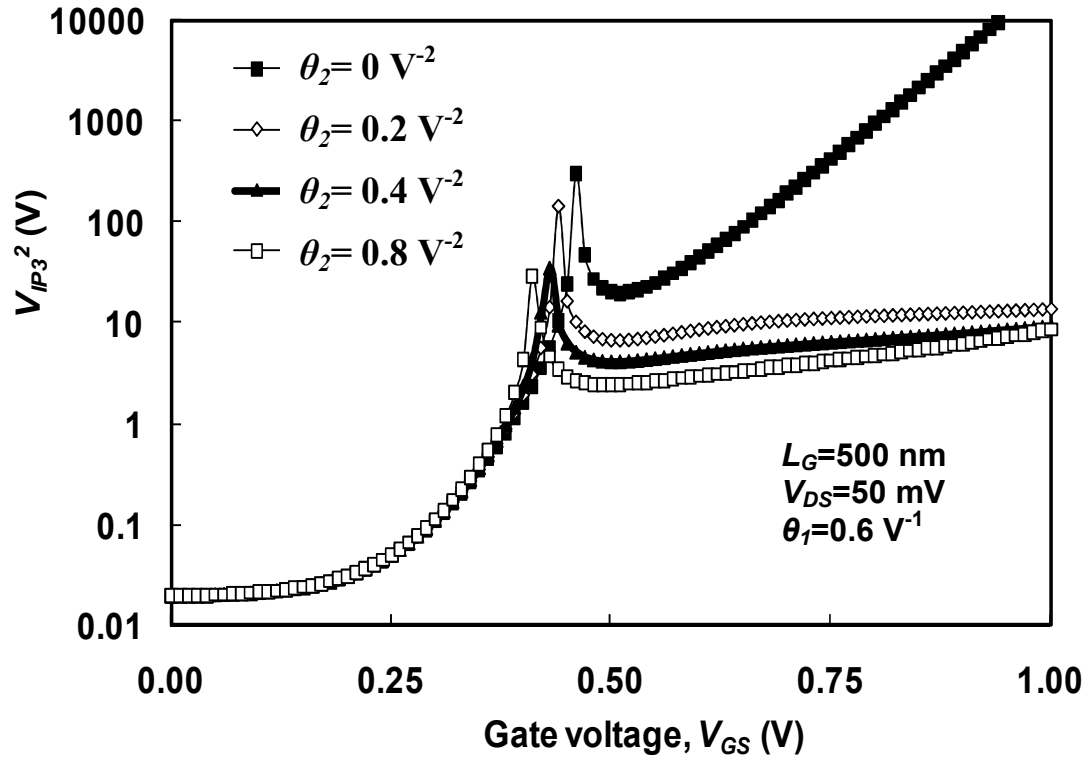


Fig. 8. The calculated  $V_{IP3}^2$  characteristics shown as functions of  $\theta_2$ .  $V_{IP3}^2$  reduces significantly as  $\theta_2$  increases. The linearity can be overestimated without taking  $\theta_2$  into account.

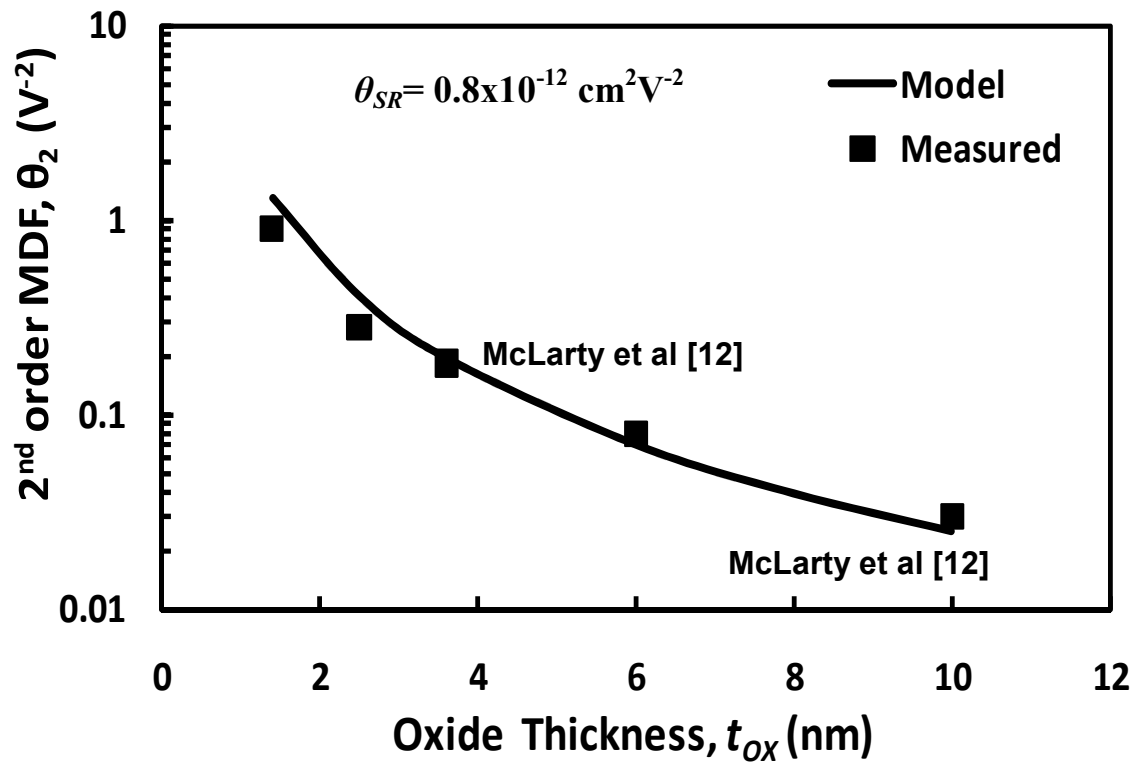


Fig. 9. The calculated and measured  $\theta_2$  values shown as functions of the oxide thickness.  $\theta_2$  increases as the oxide thickness is decreased and there is good matching between the calculated and measured data.

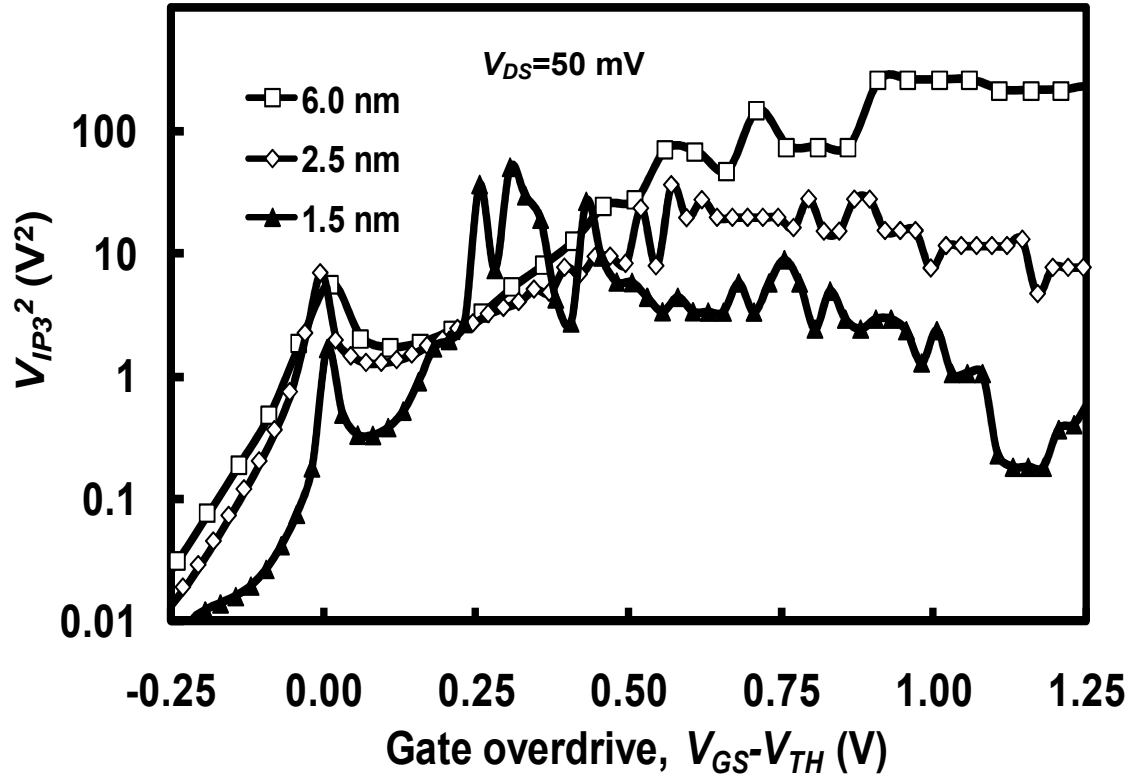


Fig. 10.  $V_{IP3}^2$  as functions of  $V_{GS}-V_{TH}$  for 1.4 nm, 2.5 nm and 6 nm gate oxide MOSFETs.  $V_{IP3}^2$  increases with the oxide thickness because  $\theta_2$  reduces as oxide thickness increases.



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