

Low temperature silicon dioxide by thermal atomic layer deposition: Investigation of material properties

D. Hiller,^{1,2,a)} R. Zierold,³ J. Bachmann,^{3,2} M. Alexe,² Y. Yang,^{1,2} J. W. Gerlach,⁴ A. Stesmans,⁵ M. Jivanescu,⁵ U. Müller,⁶ J. Vogt,⁷ H. Hilmer,⁷ P. Löper,⁸ M. Künle,⁸ F. Munnik,⁹ K. Nielsch,³ and M. Zacharias¹

¹Faculty of Engineering, IMTEK, Albert-Ludwigs-University Freiburg, Georges-Köhler-Allee 103, 79110 Freiburg, Germany

²Max Planck Institute of Microstructure Physics, Weinberg 2, 06120 Halle, Germany

³Institute of Applied Physics, Hamburg University, Jungiusstrasse 11, 20355 Hamburg, Germany

⁴Leibniz Institute of Surface Modification (IOM), Permoserstrasse 15, 04318 Leipzig, Germany

⁵Department of Physics, University of Leuven, Celestijnenlaan 200D, 3001 Leuven, Belgium

⁶Swiss Federal Laboratories for Materials Testing and Research (EMPA), Überlandstrasse 129, CH-8600 Dübendorf, Switzerland

⁷Faculty of Physics and Earth Sciences, University of Leipzig, Linnéstrasse 5, 04103 Leipzig, Germany

⁸Fraunhofer ISE, Heidenhofstrasse 2, 79110 Freiburg, Germany

⁹Forschungszentrum Dresden-Rossendorf e.V., Bautzner Landstrasse 400, 01328 Dresden, Germany

(Received 17 December 2009; accepted 25 January 2010; published online 29 March 2010)

SiO₂ is the most widely used dielectric material but its growth or deposition involves high thermal budgets or suffers from shadowing effects. The low-temperature method presented here (150 °C) for the preparation of SiO₂ by thermal atomic layer deposition (ALD) provides perfect uniformity and surface coverage even into nanoscale pores, which may well suit recent demands in nanoelectronics and nanotechnology. The ALD reaction based on 3-aminopropyltriethoxysilane, water, and ozone provides outstanding SiO₂ quality and is free of catalysts or corrosive by-products. A variety of optical, structural, and electrical properties are investigated by means of infrared spectroscopy, UV-Vis spectroscopy, secondary ion mass spectrometry, capacitance-voltage and current-voltage measurements, electron spin resonance, Rutherford backscattering, elastic recoil detection analysis, atomic force microscopy, and variable angle spectroscopic ellipsometry. Many features, such as the optical constants (n, k) and optical transmission and surface roughness (1.5 Å), are found to be similar to thermal oxide quality. Rapid thermal annealing (RTA) at 1000 °C is demonstrated to significantly improve certain properties, in particular by reducing the etch rate in hydrofluoric acid, oxide charges, and interface defects. Besides a small amount of OH groups and a few atomic per mille of nitrogen in the oxide remaining from the growth and curable by RTA no impurities could be traced. Altogether, the data point to a first reliable low temperature ALD-growth process for silicon dioxide. © 2010 American Institute of Physics. [doi:10.1063/1.3327430]

I. INTRODUCTION

Silicon dioxide is the most widely used dielectric material for optical and electronic applications. However, conventional deposition techniques for SiO₂ (e.g., chemical vapor deposition, sputtering, electron beam evaporation, etc.) have two main disadvantages: the necessity of a high growth temperature and the shadowing effect, causing a poor surface coverage in the case of structured substrates. Atomic layer deposition (ALD) provides a technique to circumvent these problems.¹ To date, atomic layer deposited SiO₂ has been rarely processed and had involved either high temperatures (>300 °C),² catalysts,^{3,4} or corrosive by-products.⁵ Besides these thermal ALD approaches, one report on plasma-enhanced ALD was published, however, with rather insufficient material purity (several atom percent of carbon and nitrogen contamination).⁶ Recently, the deposition of SiO₂ has been reported by thermal ALD using a self-catalytic reaction based on 3-aminopropyltriethoxysilane (APTES), wa-

ter, and ozone as precursors.⁷ It was shown that a high degree of thickness control and perfect surface coverage even into nanoscale pores can be achieved at low growth temperatures (120–200 °C).

An ultrathin layer of SiO₂ is of advantage as interlayer between the Si substrate and high- κ dielectrics (with κ as the dielectric constant).⁸ It would enable one to establish the standard high Si/SiO₂ interface quality, yet its thickness should be limited to the bare minimum to meet the low equivalent-oxide-thickness requirement. SiO₂ also possesses the possibility to prevent chemical reactions of the high- κ material with the silicon substrate.⁹ ALD is the favorite deposition technique for high- κ materials but an incubation period for the growth of the initial layers can result in a delay of several ALD cycles before the onset of linear growth.¹⁰ Chemical oxide is the most suitable surface for linear deposition due to its high concentration of OH groups, as demonstrated by studies on the initial growth of hafnium oxide on silicon substrates with various pretreatments, while ultrathin thermal oxide and H-terminated silicon yielded a much worse performance.¹¹ Since the ALD-grown oxide presented

^{a)}Electronic mail: daniel.hiller@imte.uni-freiburg.de.

here is terminated by OH groups, it might subsequently enable linear ALD growth of high- κ materials from the very beginning but with a higher degree of thickness control in the low nanometer range. Therefore, an *in situ* gate stack deposition of SiO₂ and high- κ material could be realized with the help of ALD-SiO₂. In addition, the increasing interest in three-dimensional (3D) transistor structures (e.g., nanowire or FinFETs) requires ultrathin SiO₂ as gate insulator, diffusion barrier, or sacrificial layer covering 3D nanostructures homogeneously during processing. Hence, new and superior SiO₂ deposition techniques are required since thermal oxidation is only applicable to planar silicon and cannot form homogeneous and uniform layers on convex and concave curvatures of structured silicon due to stress.¹² The investigation of the electronic properties of SiO₂ prepared by ALD is mandatory for its future use in 3D nanoelectronics. Furthermore, various applications in nanobiotechnology also favor SiO₂ due to its biocompatibility and chemical stability.¹³⁻¹⁵ Recently, the applicability of ALD-SiO₂ as a protective and inert material was demonstrated by deposition of a SiO₂/Fe₂O₃/SiO₂ layer stack into anodic aluminum oxide (AAO) templates and subsequent thermal reduction in hydrogen to ferrimagnetic Fe₃O₄.¹⁶ In Ref. 17 the improvement by ALD-SiO₂ coverage of AAO-based pH sensors has been demonstrated.

In this publication we present detailed material properties of this novel low temperature ALD-SiO₂ layers obtained by extensive characterization using Fourier transform infrared spectroscopy (FTIR), UV-Vis spectroscopy, photoluminescence (PL), time-of-flight secondary ion mass spectrometry (TOF-SIMS), capacitance-voltage (CV) and current-voltage measurements (IV), electron spin resonance (ESR), Rutherford backscattering spectrometry (RBS), elastic recoil detection analysis (ERDA), transmission electron microscopy (TEM), atomic force microscopy (AFM), variable angle spectroscopic ellipsometry (VASE) to determine optical constants, and laser ellipsometry for wet and dry chemical etch rates. Besides the quality achieved already after deposition (labeled *as-prepared*) we will evidence how rapid thermal annealing (RTA) can still improve material properties significantly.

II. EXPERIMENTAL DETAILS

The ALD process used here involves APTES, water, and ozone as precursors. The APTES precursor with a purity of 99% was purchased from Sigma-Aldrich. The substrate temperature for all samples was 150 °C. Details of the deposition process were published before.⁷ To account for various measurement techniques used here a variety of layer thicknesses resulting from 200 to 1500 ALD cycles were deposited onto different substrates such as n- and p-type silicon and quartz glass. Details are described in Table I. All substrates were thoroughly RCA cleaned prior to ALD of SiO₂. The average growth per cycle was determined to be 0.85 Å. RTA was carried out in a JETFIRST 200C rapid thermal processor (RTP) from Jipelec at 1000 °C for 1 min under N₂ atmosphere with ramp rates of about 100 K/s.

FTIR measurements were carried out with a Bruker

TABLE I. Overview of the ALD-SiO₂ samples and substrates used for the different measurements presented in this study. The growth rate is about 0.85 Å per cycle, hence the layer thicknesses were between approximately 20 and 130 nm.

Measurement technique	Substrate	No. of ALD cycles
FTIR	FZ-Si	1000
UV-Vis spectroscopy	Quartz glass	1500
SIMS, CV, TEM	n-Si	200
ESR, AFM, ERDA	p-Si	1000
IV	Cr on quartz glass	500
RBS, VASE	n-Si	1000

IFS113V spectrometer from 400 to 4000 cm⁻¹ with a resolution of 6 cm⁻¹. UV-Vis transmission and reflection were measured in a Varian Cary-500i in the wavelength range from 250 to 2500 nm in steps of 1 nm using an integrating sphere. PL was measured with a single grating monochromator, LN₂-cooled charge coupled device, and a HeCd laser (3.8 eV line) as excitation source. A TOF-SIMS (ION-TOF GmbH) was used to investigate elemental composition and depth profile of the SiO₂ layer. The measurements were performed using a Cs⁺ ion sputter source operated at 0.5 keV and 30 nA. The analysis was carried out in negative polarity using a pulsed Ga⁺ ion beam at 15 keV and 1 pA. An electron flood gun was used to compensate for sample charging. High-frequency (1 MHz) capacitance-voltage (C-V) and conductance-voltage (G-V) curves were measured using a HP4194A impedance analyzer. Current-voltage (I-V) measurements were carried out using a Keithley 6517 electrometer. Investigation of paramagnetic defects was performed using conventional cw locally constructed K-band (~20.5 GHz) ESR spectrometer driven in the adiabatic absorption mode and operated in the 4.2–300 K temperature range. All samples were investigated before and after extended irradiation by vacuum UV photons ($h\nu \sim 10$ eV) using a Kr resonant-discharge lamp equipped with a MgF₂ window (flux $\sim 10^{15}$ cm⁻² s⁻¹) in order to maximally activate the defects. RBS was measured with 2 MeV He⁺ projectiles under grazing incidence (76°) and a detection angle of 170°. The measurements were modeled by RUMP (RBS analysis package). ERDA was carried out using 35 MeV ³⁵Cl⁷⁺ projectiles from a 5 MV tandem accelerator. The scattering angle was 31°, the angle between beam and sample normal 75°. Recoiled target ions were detected by a Bragg-ionization chamber. Only the hydrogen ions were detected separately under 38° by a surface barrier detector with an 18 μm Al-foil entrance window. WINDF was used to simulate the measurements.¹⁸ A TEM cross section micrograph was obtained with a JEM-1010 operated at 100 kV. Surface roughness was measured under ambient conditions on a MobileS AFM from Nanosurf with the 10 μm scan head. Images were acquired in the noncontact mode (phase contrast operation mode) using a SuperSharpSilicon™ (SSS-NCLR) tip from Nanosensors. 1×1 μm² images with 512×512 pixel resolution were measured on an area with no larger dust particles and a linear line fit background subtraction was applied line by line prior to the roughness calculation. A spectral ellipsometer from J. A. Woollam Co., Inc.

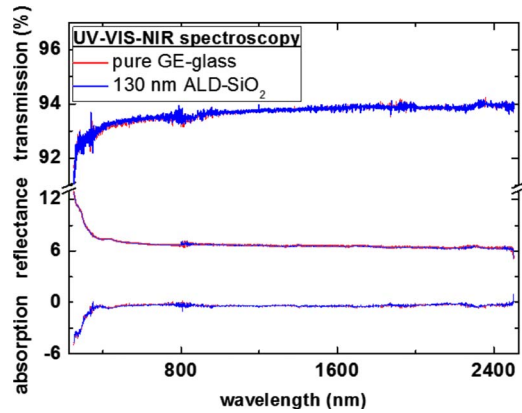


FIG. 1. (Color online) UV-Vis-NIR spectra of a 130 nm thick ALD-SiO₂ film on quartz substrate. No measurable difference appears between the ALD-coated sample and the pure quartz, indicating no absorption of the ALD-SiO₂ over the whole spectral range investigated.

was utilized for VASE measurements at 60° and 70° from 1.2 to 4.5 eV in steps of 20 meV.

III. RESULTS AND DISCUSSION

A. Optical properties

Quartz glass, the amorphous form of bulk silicon dioxide, is the most prominent material for optical lenses, mirrors, and devices. Due to the high band gap of approximately 9 eV SiO₂ has an excellent transmission behavior over a wide spectral range. Furthermore, it is possible to produce glass with a very high degree of purity to avoid luminescence centers such as alkali metals. Optical spectroscopy was performed from 250 nm (UV) to 2500 nm (NIR) to determine the transmission (T) and reflectance (R) behavior of ALD-SiO₂. The absorption A was calculated by $A=1-R-T$. The films were deposited onto 2 mm thick quartz glass substrates (GE-214) with 1500 ALD cycles, resulting in a layer thickness of about 130 nm. The results are shown in Fig. 1. Over the whole spectral range there is no significant difference in either transmission or reflectance between the ALD-coated and the virgin GE glass. Please note that the negative absorption in the UV range originates from an artifact of the white standard used to calibrate the spectrometer. Thus, the evidence that ALD-SiO₂ shows outstanding optical properties identical to that of quartz glass is not affected.

PL was performed using 90 nm thick ALD-SiO₂ layers on silicon substrate since several possible luminescence centers are known in silicon dioxide, e.g., nonbridging oxygen hole centers at around 2 eV and self-trapped excitons (2.6–2.8 eV).¹⁹ Another typical SiO₂ related defect is the E'-center; however, since its luminescence energy is 4.3–4.6 eV we cannot excite it with the 3.8 eV line of the HeCd laser used in this setup. The presence of E' centers will be addressed by ESR. The PL measurements show no luminescence for ALD-SiO₂ in the as-prepared state. Thus, the quartz glass-like optical quality of as-prepared ALD-SiO₂ is confirmed.

Ellipsometry provides for ALD a well suited technique to measure *in situ* or *ex situ* film thickness and properties. In this study we present VASE data from ALD-SiO₂ films

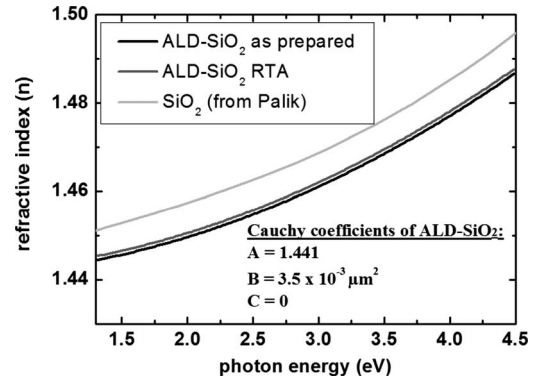


FIG. 2. Refractive index n modeled by a Cauchy approximation from variable angle spectral ellipsometry measurements. The refractive index of as-prepared and RTP-annealed ALD-SiO₂ is virtually identical. SiO₂-data from the literature (Refs. 20 and 21) reveal an offset of only about 0.5%. The modeled Cauchy parameters for ALD-SiO₂ are given in the figure.

grown by 1000 cycles on silicon wafers. The measured ellipsometric angles Ψ and Δ were modeled by means of layer stack analysis where the dielectric functions of silicon oxide were obtained numerically to get the optical constants n (refractive index) and k (extinction coefficient) as well as the layer thickness. In the model only the silicon substrate and the oxide layer were considered. A contribution of a surface roughness layer to the model was found to be negligible (see AFM section). Since the extinction coefficient k is almost zero over the whole spectral range we can use Cauchy's equation $n(\lambda)=A+B/\lambda^2+C/\lambda^4$ to model the refractive index (Fig. 2). The Cauchy coefficients of the fit of as-prepared ALD-SiO₂ are also given in Fig. 2. At first sight it is evident that no significant differences exist between ALD-SiO₂ and data known from the literature.^{20,21} The almost constant offset in the refractive indices between the ALD and the thermal oxide is only about 0.5%. A dry thermal oxide was measured as well and its data were found to be similar to the values tabulated in Ref. 20 (data not shown here). Furthermore, as-prepared state and RTP-annealed sample show almost an identical optical behavior. Therefore, in terms of optical properties we can hardly distinguish between quartz glass or thermal oxide and our ALD-grown silicon oxide. The modeled thickness of the as-prepared sample was 86.6 nm, corresponding to a growth per cycle of 0.85 Å when 2 nm chemical oxide from the RCA cleaning procedure is taken into account. However, the annealed sample could be modeled by a thickness of only 78.4 nm, unveiling a material compaction of 10% due to the 1000 °C-RTA step.

B. Structural and chemical properties

In order to obtain further insight into the structural and chemical properties of ALD-SiO₂ we performed various measurements. First, the question of chemical stability and etch rates was investigated by laser ellipsometry measurements (633 nm) before and after various treatments. The as-prepared material was subjected to RCA-1 (NH₄OH+H₂O₂+H₂O at 75 °C), RCA-2 (HCl+H₂O₂+H₂O at 75 °C), and piranha solution (H₂SO₄+H₂O₂ at 75 °C), respectively, for

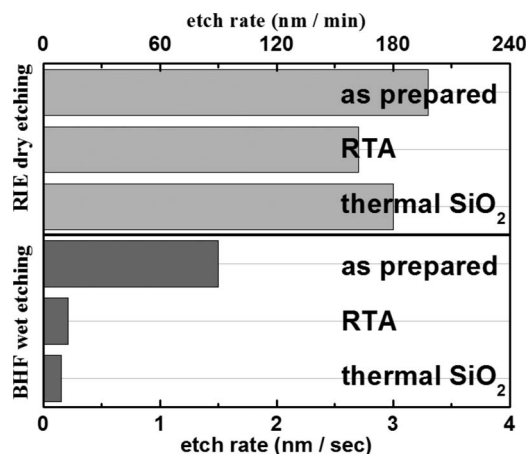


FIG. 3. Comparison of ALD-SiO₂ etch rates measured by laser ellipsometry. The upper part shows the almost identical etch rates of ALD-SiO₂ and thermal oxide achieved by RIE. The lower part indicates a similar BHF acid etch rate for both RTP-annealed ALD-SiO₂ and dry thermal oxide. In the as-prepared state the ALD material is wet etched seven times faster, demonstrating the improvement by RTA.

10 min each. No thickness decrease could be measured indicating complete material stability against these common cleaning solutions.

Buffered hydrofluoric (BHF) acid diluted 1:12 with deionized water at room temperature was used as etchant for the ALD-SiO₂ samples (as-prepared and after RTA) and a dry thermal oxide as reference. The average etch rates determined from six different etch times are presented in the lower part of Fig. 3. In the as-prepared state of ALD-SiO₂ a BHF-etch rate of 1.5 nm/s was measured; however, after RTA the etch rate is seven times lower (0.2 nm/s). This value almost reaches that of dry thermal oxide (0.15 nm/s). Hence substantial improvements in the material structure were induced by RTA at 1000 °C. Since wet etch rates are typically accelerated by a loose material composition or pin holes, this result is in good agreement with the material densification of 10% measured by VASE. However, pin holes do not play a role in this context, as has been confirmed by AFM measurements (see AFM section and Fig. 8). Therefore, the variation in etch properties of our material upon RTA may have a chemical origin: OH groups present in the as-prepared material may assist in the heterolysis of HF, thereby accelerating the etch rate of OH containing SiO₂. In contrast to HF etching, dry etching by reactive ion etching (RIE) is most efficient if the material purity is high. In this work a RIE etcher (Surface Technology Systems STS) was used with a standard oxide recipe based on CF₄, CHF₃, and Ar chemistry. It was found that dry thermal oxide is etched with that recipe at rate of 3 nm/s while, e.g., sheet glass wafers (i.e., impure silica) could be etched at a rate of only 1 nm/s. However, the ALD-SiO₂ was etched at a rate of 3.3 and 2.7 nm/s (as-prepared and RTA, respectively) indicating within measurement uncertainties similar RIE dry etch rates as thermal oxide.

FTIR was measured on float zone silicon (FZ-Si) substrates covered by 1000 ALD cycles of SiO₂. Again, a dry thermal oxide of equal thickness was used as a reference. The FTIR spectra shown in Fig. 4 are baseline corrected and

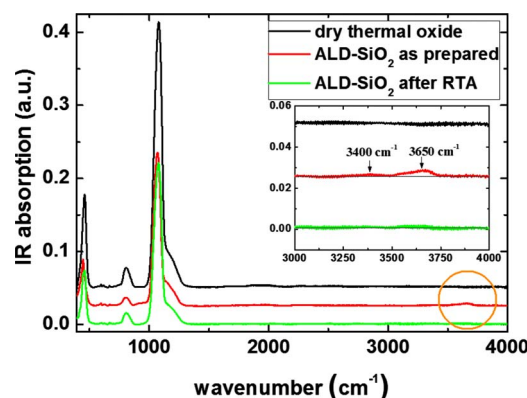


FIG. 4. (Color online) FTIR spectra of ALD-SiO₂ (as-prepared and rapid thermal annealed) and dry thermal oxide. Besides the Si-O related bonds only the as-prepared sample (red curve) shows a small peak around 3650 cm⁻¹ (magnified in the inset) corresponding to the Si-O-H stretching mode. The elimination of the OH groups by RTA is demonstrated by the green curve.

the contribution from the Si substrate is subtracted. First of all, three prominent peaks are present for all samples: the Si-O-Si out-of-plane rocking mode (460 cm⁻¹), the Si-O-Si bending mode (810 cm⁻¹), and the Si-O-Si stretching mode (1075 cm⁻¹).²² In addition, only the as-prepared sample reveals a weak peak at 3650 cm⁻¹ corresponding to the Si-O-H stretching mode. For better visibility the respective spectral region is magnified in the inset of Fig. 4. The baseline drawn in the inset for all three measurements is merely a guide for the eye. An extremely weak peak is also visible at 3400 cm⁻¹ which could be the H-O-H stretching mode. In contrast, the Si-O-H-related peak at 3650 cm⁻¹ is a clear evidence of the incorporation of OH groups into the SiO₂ material during the ALD growth. Interestingly, it vanishes after RTP annealing which results in a FTIR spectrum identical to that of dry thermal oxide. Besides the OH-related signal no other signals such as Si-H-, N-H-, N-O-, or Si-N-related modes could be evidenced by FTIR indicating no significant impurity incorporation from the APTES precursor.

The TOF-SIMS depth profiles of selected ions (¹H⁻, ²⁸Si ¹⁴N⁻, ²⁸Si₃⁻) from the ALD-SiO₂ sample grown by 200 ALD cycles (~20 nm) on silicon substrate and a dry thermal oxide of similar thickness grown by rapid thermal oxidation (RTO) are presented in Fig. 5. All signals have been normalized to the Si₃⁻ signal in bulk silicon. Besides the signals shown and other Si-O compounds no other impurity masses were detected. Four regions of the depth profile can be distinguished: the sample surface with the usual superficial contaminations (hydrocarbons, water, etc.), the SiO₂ layer, the interface, and the silicon substrate. A similar ¹²C⁻ signal of around 100 counts could be measured for all three samples in the oxide layers and in the substrate (not shown in Fig. 5). Only surface and interface exhibit a little higher carbon signal intensity. Therefore we can rule out significant carbon residues from the APTES precursor in the ALD silicon oxide. We used the Si₃⁻ signal to determine the interface since silicon atoms bound to more than one silicon atom (Si_n, n > 2) are only present in the substrate, i.e., in contrast to the situation in the SiO₂ layer.²³ Notably, we ob-

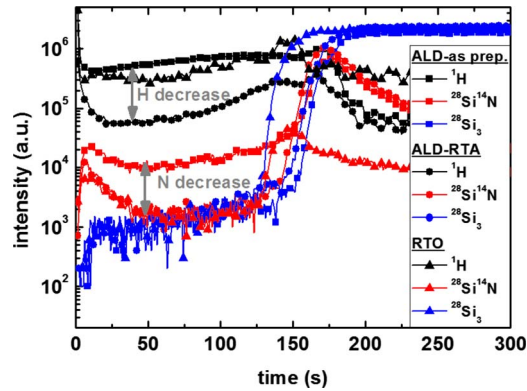


FIG. 5. (Color online) Comparison of TOF-SIMS depth profiles of selected ions ($^1\text{H}^-$, $^{28}\text{Si}^-^{14}\text{N}^-$, $^{28}\text{Si}_3^-$) for ALD- SiO_2 as-prepared (■), after RTA (●), and a rapid thermal oxide (RTO, symbol: ▲). The increase in the Si_3 signal (blue) indicates the layer-substrate interface. The as-prepared ALD oxide contains a small amount of hydrogen and nitrogen which can be decreased to thermal oxide values by RTA. Besides, no other impurities such as carbon could be detected.

serve in the films a constant increase in the H^- and Si_3^- signals toward the interface which cannot be attributed to a material inherent property, as has been checked in comparison to the homogeneous thermal oxide where similar effects were obtained, i.e., this might be a measurement artifact. The most obvious difference between the as-prepared sample (■) and the RTA sample (●) is the intensity of the hydrogen signal, as shown in Fig. 5. We observe a reduction in one order of magnitude by RTA while the ^1H -level in the substrate of these two samples is identical. Furthermore, we only have a slightly higher ^1H -signal in the middle of the RTA sample (●) compared with the Si substrate indicating a very low H content. In contrast to this, the H level in the as-prepared sample (■) is one order of magnitude higher compared with its substrate. Comparison to the thermal oxide (▲) shows a constant albeit higher ^1H -level in the layer and the substrate, except for distortions at the surface and the interface. The same situation applies for the nitrogen level represented by the $^{28}\text{Si}^-^{14}\text{N}^-$ signal. We also observe a decrease in one order of magnitude by RTA to an equal level as in the RTO sample. Hence, we can conclude from the TOF-SIMS depth profile of the ALD- SiO_2 material that (1) the as-prepared oxide is undoubtedly of a high purity with the exception for hydrogen and nitrogen residues from the APTES precursor and (2) the RTA substantially decreases the amount of hydrogen or hydrogen related moieties (such as $-\text{OH}$) as well as nitrogen in the material.

RBS was measured on an 87 nm thick as-prepared sample (1000 ALD cycles). In Fig. 6 the RBS spectrum and

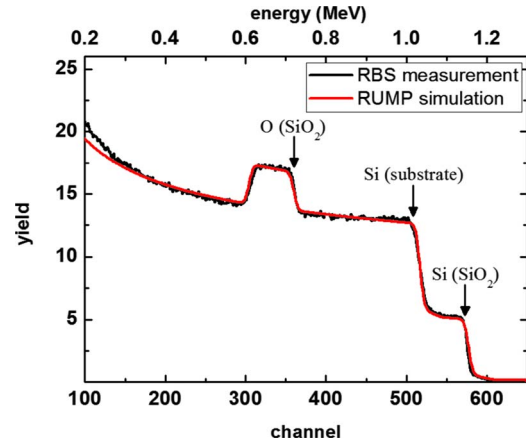


FIG. 6. (Color online) Rutherford backscattering spectrum from a SiO_2 sample grown by 1000 ALD cycles. The black curve is the measurement; the red line is the fit modeled by RUMP. Best fit parameters correspond to a Si:O stoichiometry of 1:2 and a layer thickness of 83 nm. No heavy impurities could be measured.

the RUMP simulation are presented. The results of the best-fit simulation are a layer composition O:Si=2 and a thickness of (83 ± 2) nm, i.e., in agreement with the VASE results within uncertainty. It is noteworthy to mention that in terms of RBS thickness is calculated by the measured atomic areal densities. An agreement with the thickness measured by VASE indicates a material density close to the table values for SiO_2 . Due to absence of any peaks beyond the silicon edge (1.15 MeV) we can exclude heavy impurities in our material down to the ppm range.

Complementary to RBS ERDA was carried out. First of all no carbon could be detected; however, a small amount of around (0.55 ± 0.05) at. % of nitrogen was measured for the as-prepared sample but not for the RTA sample. The origin of the N signal is probably the amino groups of the APTES precursor. The mechanism of nitrogen reduction by RTA is not obvious. ERDA also enabled the quantification of the hydrogen content. As depicted in Table II it is around 6.6 at. % for the as-prepared ALD- SiO_2 sample and more than one order of magnitude less (i.e., close to detection limit) after RTA. A simple estimation of the stoichiometry of the as-prepared material fails when only the ratio O:Si is taken into account, as represented by O:Si=2.21 in Table II. From the FTIR and ESR data (see below) we have to assume a Si-O-H bonding configuration for the hydrogen, hence not every oxygen atom is able to bond to two silicon atoms. The last column in Table II represents the stoichiometry when the amount of hydrogen is subtracted from the oxygen prior to

TABLE II. Quantitative composition of ALD- SiO_2 from ERDA. O:Si represents the ratio of the measured oxygen to silicon concentration. Since Si-O-H configuration is expected from FTIR the ratio between the difference of oxygen and hydrogen (i.e., O-H) to silicon is supposed to be more representative in terms of stoichiometry. The trace amounts of nitrogen are neglected in this consideration.

	Si (at. %)	O (at. %)	N (at. %)	H (at. %)	O:Si	(O-H):Si
As-prepared	28.9	64.0	0.55	6.63	2.21	1.99
RTA (1000 °C)	33.0	66.5	0.02	0.50	2.02	2.00

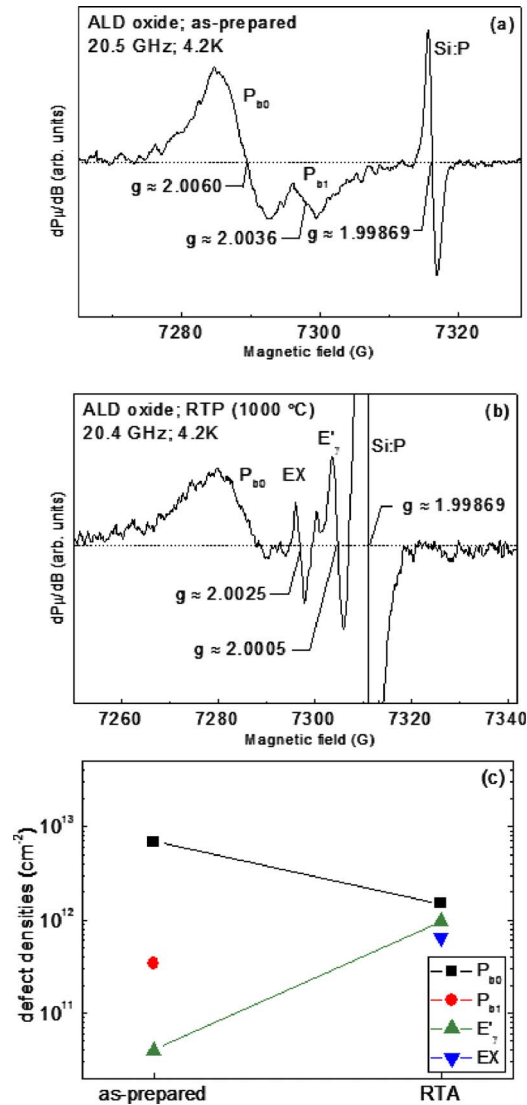


FIG. 7. (Color online) First-derivative K-band ESR spectra observed at 4.2 K on SiO₂(87 nm)/(100)Si entities fabricated by ALD SiO₂ deposition (1000 cycles) (a) in the as-prepared sample for the orientation of the applied magnetic field $\mathbf{B} \parallel \mathbf{n}$, the (100) interface normal, and (b) after subjecting to RTP (1000 °C; 1 min; N₂) followed by VUV irradiation at 300 °C. The applied modulation field amplitude was 0.4 G and the incident microwave power $P_{\mu} \approx 0.3$ nW. The signal at $g = 1.99869$ stems from a comounted Si:P marker sample, serving as g value and defect density reference. Presentation of all paramagnetic defect densities derived from the ESR spectra (c). The densities of interfacial DB defects (i.e., P_{b0} and P_{b1}) of as-prepared ALD-SiO₂ are typical for low temperature grown oxides. After RTA the P_{b0} defects are reduced close to the values of the standard thermal Si/SiO₂ interface and P_{b1} close to the detection limit. Insufficient E'_γ formation by VUV irradiation (~ 10 eV) for the as-prepared sample indicates the presence of OH groups; after RTA the E'_γ-density adjusts close to standard values. In addition, the oxide reorganization by RTA introduces EX centers.

calculating the ratio to silicon [labeled (O–H):Si]. Therefore, a balanced O to Si ratio is proven for ALD-SiO₂ before and after RTP annealing.

Occurring paramagnetic point defects were investigated by K-band ESR at 4.2 K on samples prepared by 1000 ALD cycles on p-type (100)Si substrates, resulting in a SiO₂ layer of 87 nm. Typical ESR spectra are shown in Figs. 7(a) and 7(b), the defect densities determined from the ESR measurements before and after RTA are depicted in Fig. 7(c). On the as-prepared sample [Fig. 7(a)], we observe as prominent sig-

nals the anisotropic spectra from the archetypal intrinsic P_{b0} and P_{b1} interface defects, inherently generated in (100)Si/SiO₂ as a result of interface network/lattice mismatch, with densities inferred as $[P_{b0}] = (6.9 \pm 0.3) \times 10^{12} \text{ cm}^{-2}$ and $[P_{b1}] = (3.4 \pm 0.2) \times 10^{11} \text{ cm}^{-2}$. As to P_{b0} centers (interfacial $\text{*Si} \equiv \text{Si}_3$ centers, where the dot represents an unpaired Si sp₃ hybrid), electrically identified as detrimental interface traps,²⁴ the density is substantially larger compared with standard thermal oxide interfaces grown in the range of 850–950 °C, invariably with inherent density $\approx 1 \times 10^{12} \text{ cm}^{-2}$ for both P_{b0} and P_{b1}.²⁵ Besides, no other defects could be observed, i.e., no SiO₂-specific centers such as the E'_γ center (O vacancy; generic entity $\text{*Si} \equiv \text{O}_3$). Subsequent vacuum ultraviolet (VUV) irradiation (~ 130 min) was found not to affect the P_b-type defect spectra, which may be expected in the absence of any substantial passivation of Si dangling bond (DB) defects by H. The latter was confirmed by the observation that the P_b-type defect status was left unaltered after annealing the samples at 500 °C (1 h, N₂). However, noteworthy is that hardly any E'_γ-centers are observable, with $[E'_{\gamma}] < 4.4 \times 10^{10} \text{ cm}^{-2}$ (volumetric density of $\sim 5 \times 10^{15} \text{ cm}^{-3}$), that is about 100 times less than typically observed in standard thermal oxides after such treatment ($\sim 4 \times 10^{17} \text{ cm}^{-3}$).²⁶ For the current work, this was independently verified by growing a standard thermal (100)Si/SiO₂ structure with oxide thickness of 75 ± 0.5 nm, comparable to the ALD-SiO₂ layer thickness. Subjecting this layer to a similar VUV treatment resulted in a density $[E'_{\gamma}] \sim 3.3 \times 10^{12} \text{ cm}^{-2}$, corresponding to a volume density of $\sim 4.4 \times 10^{17} \text{ cm}^{-3}$. The likely explanation for this reduced E'_γ density behavior is that the ALD oxide, a low-temperature oxide (LTO), is an O-rich silica with substantial content of Si–OH groups (see FTIR and ERDA section), and hence much reduced density of E' precursor sites (O₃ \equiv Si–H bonds). So, on this matter, the ALD oxide would differ from standard thermal SiO₂. A noteworthy observation also is that no carbon signal could be detected: due to the utmost ESR sensitivity (i.e., ppb range) this implies amazingly the absence of any carbon residues from the metalorganic APTES precursor in agreement with the TOF-SIMS data. After RTP annealing (1000 °C, 1 min, N₂) of the as-prepared ALD-SiO₂ the ESR spectrum changes. The P_{b0} defect density is now found to be substantially reduced to $[P_{b0}] = (1.5 \pm 0.5) \times 10^{12} \text{ cm}^{-2}$, nearing the typical value (quality) for standard thermal oxide. In addition, EX centers now appear ($[EX] \sim 6.5 \times 10^{11} \text{ cm}^{-2}$), i.e., another SiO₂-associated defect, which might be related to some kind of oxide damage introduced by the annealing as observed several times before.²⁷ As illustrated in Fig. 7(b), after subsequent VUV treatment, we notice, besides the unaltered P_{b0} signal, the additional appearance of the E'_γ signal now in density of $\sim 9.6 \times 10^{11} \text{ cm}^{-2}$, i.e., in the same order of magnitude as in the standard thermal oxide. This result would indicate that the oxide matrix has altered, i.e., no influence of OH groups left after RTA. Summing up, we conclude that a reasonable Si/SiO₂-interface quality comparable to other low temperature grown interfaces was found in the as-prepared state, with the influential presence of OH groups in the oxide layer. This effect can be suppressed by RTA, where the

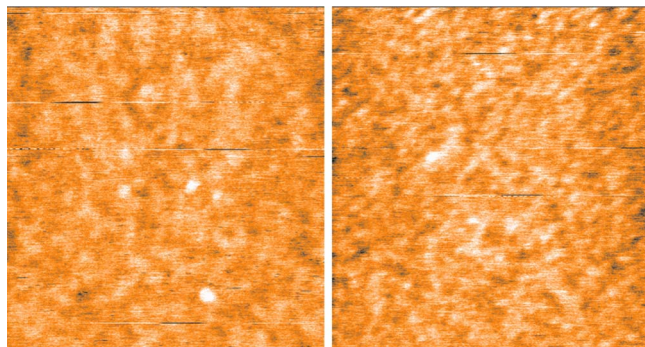


FIG. 8. (Color online) AFM topographies of as-prepared ALD-SiO₂ (left) and ALD-SiO₂ after RTA at 1000 °C (right). The image size is 1 × 1 μm² and the height color scale is 972 pm. The AFM scans prove ideal smoothness of the material with a rms roughness of only 1.5 Å (left: as-prepared) and 1.6 Å (right: after RTA). No deterioration of the surface morphology by RTA is observed.

Si/SiO₂ structure is now taken to an improved quality, closer to that of the more standard thermal Si/SiO₂ entity. On ESR basis, the evidence comes from two sides: both the P_b-type interface defect and E' _γ center densities reach close to standard values.

In order to investigate the surface roughness the ALD-SiO₂ layers were measured by AFM scans. As depicted in Table I these measurements were taken on samples prepared by 1000 ALD cycles. The topography scans of the as-prepared and RTP-annealed state are presented in Fig. 8. Except for a few particles no topographical features are visible. The absence of any pinholes (with widths ≥ 20 nm and depths ≥ 1 nm) is proven as expected for ALD growth. Furthermore, RTA does not introduce any cracks or changes in the surface morphology. Numerically, we find a root-mean-square (rms) roughness of 1.5 and 1.6 Å for the as-prepared and the RTP-annealed state, respectively. Compared with bare silicon wafer roughness of about ≤ 2 Å this demonstrates the ideal layer smoothness of the ALD-SiO₂ process presented here.

C. Electrical properties

Metal oxide semiconductor (MOS) capacitors were used to investigate the electronic properties of ALD-SiO₂. MOS structures were prepared on n-type silicon with aluminum contact pads (area of 0.23 mm²) evaporated onto the 20 nm thick ALD-oxide layer (200 ALD cycles). The backside aluminum contact was evaporated after removal of the backoxide layer by HF etching. Figure 9 shows the high-frequency (1 MHz) capacitance- and conductance-voltage curves of the as-prepared and RTP-annealed samples. A cross-section TEM image is presented in the inset of Fig. 9, demonstrating the smooth surface and interface of our ALD-SiO₂ layers. The bias was swept from inversion to accumulation and back (for the sake of clarity the back sweep is not shown here). The filled symbols (■,▲) denote measurements taken in the initial state, while curves with open symbols (□,△) were measured after a bias stress in accumulation of +10 V for 30 s. As can be seen here, significant differences appear between the as-prepared (Fig. 9 top) and the RTA sample (Fig. 9 bottom). From thermal oxidation it is well known that pos-

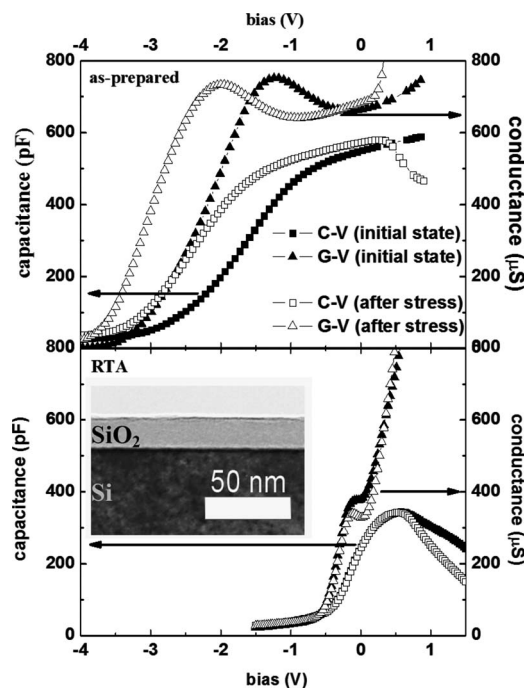


FIG. 9. High-frequency capacitance-voltage (CV) and conductance-voltage (GV) measurements of MOS structures of as-prepared ALD-SiO₂ (top) and after RTP annealing (bottom). Filled symbols (■,▲) represent the measurements taken in the initial state, open symbols (□,△) after a bias stress in accumulation of +10 V for 30 s. Significant changes in terms of oxide capacity and flatband position occur after annealing. The inset shows a TEM cross section image of the as-prepared sample.

toxidation annealing can increase the oxide quality significantly,²⁸ since excess silicon or oxygen atoms from silicon-rich or, respectively, oxygen-rich growth can remain in the layer after oxidation.²⁹ Since the ALD process is based on surface chemistry rather than on mass transport it is stoichiometric, as was proven for this material previously⁷ and by the RBS and ERDA results presented above. To quantify the influence of thermal treatment the capacitance curves are examined first. The dielectric constant can be calculated with the help of the simple plate capacitor model using the oxide layer thickness of about 20 nm and the oxide capacitances in accumulation of the as-prepared sample ($C_{ox}=600$ pF) and the thermally treated ones ($C_{ox}=350$ pF). Table III summarizes the results. The dielectric constant was found to be 6.1 and 3.6 for the as-prepared state and the RTA sample, respectively. Hence, only the sample after annealing is close to the standard literature value for thermal SiO₂ (approximately 3.9). The rather high value (6.1) of the dielectric constant of the as-prepared state might be attributed to additional polar-

TABLE III. Compilation of the electrical parameters derived from high frequency CV-GV measurements. Significant improvements were attained by RTA at 1000 °C.

	As-prepared	RTA (1000 °C)
Dielectric constant ϵ	6.1	3.6
Flatband voltage V_{FB} (V)	-2.27	-0.28
Fixed oxide charge Q_f (cm ⁻²)	$+3.2 \times 10^{12}$	-1.9×10^{10}
Trapped oxide charge Q_t (cm ⁻²)	$+1.0 \times 10^{12}$	0
Interface trap level density D_{it} (cm ⁻² eV ⁻¹)	1.3×10^{13}	1.5×10^{12}

izable moieties in the layer (for instance, $-\text{OH}$). The relationship between dielectric constant and OH content was demonstrated on chemical-vapor deposition (CVD) oxides before.³⁰ The flatband capacitance method was used to investigate the behavior of oxide charges.³¹ Substantial changes in the flatband voltage position were observed: $V_{\text{FB}} = -2.27$ V for the as-prepared sample and -0.28 V for the RTA sample. By using the equation $V_{\text{FB}} = Q_f / C_{\text{ox}}$ the density of fixed oxide charge can be calculated.³¹ The work function difference of approximately -0.3 V between the n-type silicon used here and aluminum (Φ_{ms}) has to be taken into account in order to calculate Q_f accurately. For the as-prepared sample a rather high value of 3.2×10^{12} positive fixed oxide charges per cm^2 (see Table III) was obtained, which is attributed to either hydroxyl groups or charges at the Si/SiO₂ interface. Impurities of the APTES precursor representing fixed positive charges (e.g., metals) can be ruled out, as has been proven by RBS and TOF-SIMS. Amazingly, the RTP annealing at 1000°C reduces fixed oxide charges Q_f to the 10^{10} cm^{-2} range, which is virtually zero within uncertainties. Comparing measurements of the initial state (Fig. 9, filled symbols $\blacksquare, \blacktriangle$) and after bias stress (open symbols \square, \triangle) exhibits a severe voltage shift for the as-prepared sample, which can be attributed to trapped oxide charge (Q_t) or mobile ionic charge (Q_m). Using $\Delta V_{\text{FB}} = Q / C_{\text{ox}}$ we estimated Q to be around 10^{12} cm^{-2} . In our case, the presence of mobile ions such as sodium has been ruled out by the ERDA data. Thus, trapped charges Q_t must be the origin for the shift in the C-V curve. This positive charge in the as-prepared state can be attributed to injection of carriers into the oxide.³¹ The RTA sample has a value of $Q_t = 0$, since within the limits of our measurement accuracy no flatband voltage shift was measured. Therefore, no significant amount of charge is trapped in the oxide under the bias stress used here ($+10$ V, 30 s). Please note that extensive procedures for bias-temperature stress are able to distinguish between trapped and mobile ionic charges but include stress times of up to several hours, elevated temperatures, and much stronger electrical fields.³¹ Our measurement indicates only a substantial improvement due to RTP annealing but cannot reveal that the oxide is absolutely free of charge traps. For the as-prepared sample we have to point out that a hysteresis of several tens of millivolts on the voltage axis exists even between the up and down curves of the C-V sweep (not shown in Fig. 9 for the sake of clarity). Hence, the charges are trapped easily even under low bias in the as-prepared state.

Conductance-voltage (G-V) curves were recorded to investigate the interface properties. The interface trap level densities (D_{it}) were calculated from the peak of the corrected initial-state conductance curves according to Refs. 31 and 32. Again a rather high value of 10^{13} $\text{cm}^{-2} \text{eV}^{-1}$ can be calculated for the as-prepared state (Fig. 9 top). Please note that this amount of interface traps can also cause a contribution of several volts to the value of V_{FB} , hence introducing additional uncertainty about the capacitance curve of the as-prepared state of our ALD-SiO₂.³¹ The conductance peaks are smaller and narrower for the annealed sample (Fig. 9 bottom), resulting in a value of interface trap level density D_{it} in the 10^{12} $\text{cm}^{-2} \text{eV}^{-1}$ range (Table III), i.e., a decrease

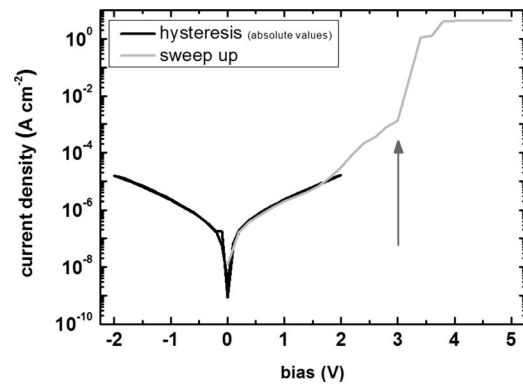


FIG. 10. Current-voltage (IV) measurements of a MOM structure with ALD-SiO₂. The hysteresis curve (black) indicates reproducible electrical behavior under low electrical fields. In the sweep-up measurement (gray curve) already at 3 V (i.e., electrical field of 75 mV/nm) breakdown occurs as indicated by the arrow.

by RTA of one order of magnitude but still approximately one order of magnitude too high for electronic device applications.⁸ However, these values are only indicative since the SiO₂ deposition took place on RCA-cleaned silicon, hence on chemical oxide which is known to have poor electrical quality. The D_{it} -values obtained here are consistent with the P_b-type interface defects from ESR. The sudden increase in the parallel equivalent conductance accompanied by the decrease in capacitance in accumulation is another feature of Fig. 9. Since the admittance is sensitive to electrical loss the measurements are affected by further uncertainty in the capacitance measurement in accumulation, hence distorting the estimation of the oxide capacitance and the calculated dielectric constant. This circumstance could solve the apparent discrepancy between the excellent dielectric behavior measured by VASE and the 50% too high electrically measured ϵ -value (6.1). For both samples this effect is pronounced after bias stress although the annealed sample shows this effect already in the initial state. The jump in conductance is due to the rather high conductivity of the oxide layers and in addition probably due to the creation of more conductive paths by thermal processing or bias stress.

In order to measure the current-voltage behavior a metal-oxide-metal (MOM) structure was prepared by evaporation of 100 nm of chromium onto quartz glass, followed by the deposition of 500 ALD cycles of SiO₂ (~ 40 nm), and the subsequent structuring and evaporation of another 100 nm of Cr as top contact with 0.23 mm^2 . Although IV measurements could be taken with the structure in the as-prepared state, RTP annealing was not successful since cracks occurred and severe material damage was visible. Since this problem does not occur for silicon or quartz substrates it is supposed to be attributed to the chromium layer. Two IV measurements are presented in Fig. 10: a hysteresis loop (from 0 to $+2$ V, then to -2 V and back to 0) and a bias sweep from 0 to $+5$ V. The hysteresis (due to the logarithmic scale on the ordinate the absolute values are shown) indicates reproducible IV behavior under both bias polarities, and at 1 V we can deduce a current density of 10^{-6} A cm^{-2} . Hence, as mentioned above in the CV section the leakage current through the as-prepared ALD-SiO₂ is rather high.

The same applies to the breakdown voltage of 3 V as represented by the bias sweep in Fig. 10. In the current case of an approximately 40 nm thick layer this corresponds to a breakdown field of only 75 mV/nm, while good thermal oxides achieve up to 1–2 V/nm. The origins of the high conductivity and weak breakdown stability as well as ways for improvement, e.g., by RTA, are an object of further investigation.

IV. CONCLUSION

In summary, a detailed investigation of material properties of a novel low temperature thermal-ALD silicon dioxide deposition is presented. The superior aspects of this process are its low deposition temperature (150 °C), the absence of corrosive by-products, no necessity of catalysts, and its ALD-related high degree of thickness control and perfect uniformity. The growth per cycle was determined to be 0.85 Å. Many aspects of this ALD-SiO₂ such as optical transmission and absorption, refractive index, stoichiometry, and chemical stability can be directly compared with thermal oxide or quartz glass already in its as-prepared state, hence representing the highest achievable material quality. Interface quality between ALD-SiO₂ and Si substrate was shown to be comparable to other low temperature deposition methods. SIMS, ESR, RBS, and ERDA verified a high degree of purity: carbon residues can be ruled out with ppb precision and only hydrogen and about 5 atomic per mille of nitrogen could be traced. By means of ERDA the H content was measured to 6.6 at. % and FTIR and ESR measurements revealed that it is bonded in a Si–O–H configuration. This probably affects certain material properties, in particular the fixed and trapped oxide charges (7×10^{12} and 1×10^{12} cm⁻²) or breakdown field (0.75 MV/cm). RTA at 1000 °C was employed to investigate potential improvements. The effect was in most respects outstanding: the N and H contents as well as fixed and trapped oxide charges (Q_f, Q_t) were reduced below detection limits and the interface trap level density was improved by one order of magnitude (1×10^{12} cm⁻² eV⁻¹). In addition, a certainty has been given that the as-prepared ALD-SiO₂ is about 10% less compact compared with the RTP-annealed state, accompanied by an improvement in BHF-etch rate to almost thermal oxide strength. Excellent surface smoothness (1.5 Å) was demonstrated by AFM measurements and no degradation was found after RTP annealing.

The ALD-SiO₂ process presented here opens new possibilities for ALD-based depositions in nanotechnology and related fields, since a conventional and well known material of reasonable quality can now be deposited easily at low temperatures and with great uniformity. Thereby it offers an attractive alternative to, e.g., the so-called LTO grown by CVD at around 425 °C or to typical plasma-enhanced chemical-vapor deposition oxides (200–400 °C), especially when low thermal budgets, absence of plasma irradiation, and a homogeneous coverage of 3D nanostructures are essential.

ACKNOWLEDGMENTS

The authors gratefully acknowledge the financial support by DFG under Project No. Za191/14-3 and by the German Ministry of Education and Research (BMBF Grant No. 03N8701). K. P. Meyer is acknowledged for technical support. R. Schmidt-Grund and M. Grundmann are acknowledged for support in ellipsometry.

- ¹M. Knez, K. Nielsch, and L. Niinistö, *Adv. Mater. (Weinheim, Ger.)* **19**, 3425 (2007).
- ²S. W. Lee, K. Park, B. Han, S. H. Son, S. K. Rha, C. O. Park, and W. J. Lee, *Electrochem. Solid-State Lett.* **11**, G23 (2008).
- ³J. D. Ferguson, E. R. Smith, A. W. Weimer, and S. M. George, *J. Electrochem. Soc.* **151**, G528 (2004).
- ⁴B. A. McCool and W. J. DeSisto, *Chem. Vap. Deposition* **10**, 190 (2004).
- ⁵J. W. Klaus, O. Sneh, and S. M. George, *Science* **278**, 1934 (1997).
- ⁶J. W. Lim, S. J. Yun, and J. H. Lee, *ETRI J.* **27**, 118 (2005).
- ⁷J. Bachmann, R. Zierold, Y. T. Chong, R. Hauert, C. Sturm, R. Schmidt-Grund, B. Rheinländer, M. Grundmann, U. Gösele, and K. Nielsch, *Angew. Chem., Int. Ed.* **47**, 6177 (2008).
- ⁸G. D. Wilk, R. M. Wallace, and J. M. Anthony, *J. Appl. Phys.* **89**, 5243 (2001).
- ⁹J. Robertson, *Rep. Prog. Phys.* **69**, 327 (2006).
- ¹⁰M. M. Frank, Y. J. Chabal, M. L. Green, A. Delabie, B. Brijs, G. D. Wilk, M. Y. Ho, E. B. O. da Rosa, I. J. R. Baumvol, and F. C. Stedile, *Appl. Phys. Lett.* **83**, 740 (2003).
- ¹¹M. L. Green, M. Y. Ho, B. Busch, G. D. Wilk, T. Sorsch, T. Conard, B. Brijs, W. Vandervorst, R. I. Räisänen, D. Muller, M. Bude, and J. Graul, *J. Appl. Phys.* **92**, 7168 (2002).
- ¹²R. B. Marcus and T. T. Sheng, *J. Electrochem. Soc.* **129**, 1278 (1982).
- ¹³Y. H. Cho, S. W. Lee, B. J. Kim, and T. Fujii, *Nanotechnology* **18**, 465303 (2007).
- ¹⁴N. Kaji, Y. Tezuka, Y. Takamura, M. Ueda, T. Nishimoto, H. Nakanishi, Y. Horiike, and Y. Baba, *Anal. Chem.* **76**, 15 (2004).
- ¹⁵A. Burns, H. Ow, and U. Wiesner, *Chem. Soc. Rev.* **35**, 1028 (2006).
- ¹⁶K. Pitzschel, J. M. M. Moreno, J. Escrig, O. Albrecht, K. Nielsch, and J. Bachmann, *ACS Nano* **3**, 3463 (2009).
- ¹⁷L. Moreno i Codinachs, C. Birkenstock, T. Garma, R. Zierold, J. Bachmann, K. Nielsch, M. J. Schöning, and A. Fontcuberta i Morral, *Phys. Status Solidi A* **206**, 435 (2009).
- ¹⁸N. P. Barradas, C. Jeynes, and R. P. Webb, *Appl. Phys. Lett.* **71**, 291 (1997).
- ¹⁹M. A. Stevens Kalceff and M. R. Phillips, *Phys. Rev. B* **52**, 5 (1995).
- ²⁰*Handbook of Optical Constants of Solids*, edited by E. D. Palik (Academic, New York, 1997).
- ²¹V. Gottschalch, R. Schmidt, B. Rheinländer, D. Pudis, S. Hardt, J. Kvietkova, G. Wagner, and R. Franzheld, *Thin Solid Films* **416**, 224 (2002).
- ²²V. P. Tolstoy, I. V. Chernyshova, and V. A. Skryshevsky, *Handbook of Infrared Spectroscopy of Ultrathin Films* (Wiley, New York, 2003); M. Zacharias, D. Dimora-Malinovska, and M. Stutzmann, *Philos. Mag. B* **73**, 799 (1996).
- ²³M. Perego, S. Ferrari, S. Spiga, E. Bonera, M. Fanciulli, and V. Soncini, *Appl. Phys. Lett.* **82**, 121 (2003).
- ²⁴E. H. Poindexter, G. J. Gerardi, M.-E. Rueckel, P. J. Caplan, N. M. Johnson, and D. K. Biegelsen, *J. Appl. Phys.* **56**, 2844 (1984); A. Stesmans and V. V. Afanas'ev, *Phys. Rev. B* **57**, 10030 (1998).
- ²⁵A. Stesmans and V. V. Afanas'ev, *J. Vac. Sci. Technol. B* **16**, 3108 (1998).
- ²⁶P. G. Tello, V. V. Afanas'ev, and A. Stesmans, *Microelectron. Eng.* **72**, 81 (2004); A. Stesmans and V. V. Afanas'ev, *J. Appl. Phys.* **97**, 033510 (2005).
- ²⁷W. E. Carlos and S. M. Prokes, *J. Appl. Phys.* **78**, 2129 (1995); H. J. von Bardeleben, J. L. Cantin, J. J. Ganem, and I. Trimaille, in *Defects in High-κ Gate Dielectric Stacks*, NATO Science Series II, edited by E. Gusev (Kluwer, Dordrecht, 2006); B. B. Triplett, P. T. Chen, Y. Nishi, P. H. Kasai, J. J. Chambers, and L. Colombo, *J. Appl. Phys.* **101**, 013703 (2007).

²⁸B. J. O'Sullivan, P. K. Hurley, C. Leveugle, and J. H. Das, *J. Appl. Phys.* **89**, 3811 (2001).

²⁹A. Ludsteck, J. Schulze, I. Eisele, W. Dietl, and Z. Nenyai, *J. Appl. Phys.* **95**, 2827 (2004).

³⁰B. Fowler and E. O'Brian, *J. Vac. Sci. Technol. B* **12**, 441 (1994).

³¹E. H. Nicollian and J. R. Brews, *MOS (Metal Oxide Semiconductor) Physics and Technology* (Wiley, New York, 1982).

³²J. R. Brews, *Solid-State Electron.* **26**, 711 (1983).