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Luminescence, Chiroptical, Magnetic and Ab-initio Crystal-Field Characterizations of an Enantiopure Helicoidal Yb(III) Complex

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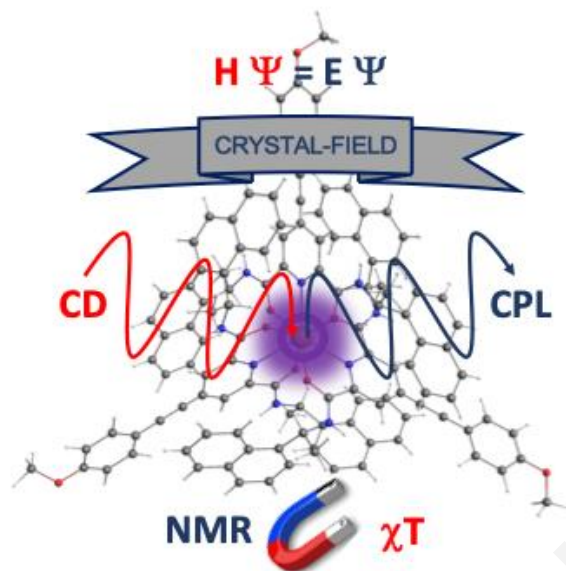
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Abstract

The electronic structure of a chiral Yb(III)-based complex is fully determined by taking advantage of experimental magnetic, luminescence, and chiroptical (NIR-ECD and CPL) characterizations in combination with ab-initio wavefunction calculations. The combined use of these techniques allows determining with high resolution the electronic structure diagram as well as the nature of the different states involved in the magnetic and chiroptical properties of the investigated complex. The crystal-field picture deduced from spectroscopic measurements (absorption and emission) is used to reproduce the magnetic properties. Subsequently, advanced ab-initio calculations demonstrate that global chiroptical spectra correspond to the sum of entangled transitions with similar or opposite polarizations.

INTRODUCTION

The presence of a partially filled 4f valence shell in lanthanide(III) ions has been the origin of an incredible rise in the Ln-based chemistry over the past decades. Among the large panel of applications, the most striking use of such complexes can be found in the recent development of new molecular magnets,¹ named single-molecule magnets (SMMs), which aim at elaborating high density data storage devices, or in the development of Ln(III) luminescent probes for bio-medical purposes.² Additionally, chiral Ln(III)-based complexes have been used lately to achieve circularly polarized luminescence (CPL)³ for applications in biological sensing,⁴ anti-counterfeiting devices⁵ and organic light-emitting diodes (CP-OLEDs) for their appealing use in display technology.⁶ In all of these cases, the magnetic, luminescence and chiroptical properties arise from the same origin: the crystal-field (CF) interaction generated by the ligand sphere around the Ln(III) ion that lifts the degeneracy of the $^{2S+1}L_J$ ground and excited terms.

Indeed, for an isolated Ln(III) ion the 4f transitions are Laporte forbidden, preventing any optical activity due to the vanishing electric transition dipole moments. The presence of a coordination sphere that does not contain an inversion center, allows to (i) break the parity rule and (ii) lift the degeneracy of the $^{2S+1}L_J$ terms, giving rise to optical properties.⁷ Similarly, the CF interaction is mandatory to generate the large magnetic anisotropies required for SMMs.

In fact, an isolated Ln(III) ion only exhibits an isotropic magnetic moment of magnitude $g_J\sqrt{J(J+1)}\mu_B$ corresponding to the magnetic moments average of the different M_J states of the ground term. Depending on the symmetry and magnitude of the CF, the M_J states are allowed to mix with each other's and the degeneracy of these states is lifted, inducing potential magnetic anisotropy. Therefore, in order to be able to tune the optical, chiroptical and magnetic properties in Ln(III)-based complexes, it is necessary to perfectly characterize and rationalize the CF splitting.

Over the last few years, magneto-luminescence correlations have been used in combination with ab-initio multi-configurational wavefunction calculations to probe in detail the electronic structure of several Ln(III)-based SMMs⁸ involving mainly terbium (III),⁹ dysprosium (III),¹⁰ and ytterbium(III).¹¹ In these examples, the luminescence of the Ln(III) ions was used as a snapshot of the energetic splitting of the ground $^{2S+1}L_J$ terms, allowing (i) a direct comparison with the ab-initio electronic structure calculations and (ii) a rationalization of the magnetic properties. The proof of concept was then extended with the use of low-temperature luminescence measurements, allowing an improvement of the spectra's resolution by removing a sizable number of extra transitions corresponding to "hot bands".¹² However, the assignment of these hot bands is not straightforward and can be easily over-interpreted by the presence of additional effects such as vibronic ones.^{11d} Despite its wide use, ab-initio calculations might fail as well to provide accurate electronic structure information, particularly for strongly correlated systems such as ytterbium-based complexes, where electronic correlation and covalent effects are important and remain challenging to tackle with ab-initio methods.

To circumvent such limitations in the high-resolution characterization of the Ln(III) CF interactions, one step further would be to take advantage of the presence of potential chirality. Indeed, the presence of chiral ligands, or of achiral ligands wrapped around the Ln(III) ion with a Λ - or Δ -helicity, generates enantiomers that interact differently with linearly and circularly polarized light. These enantiomers will absorb or emit photons preferentially with one sense of circular polarization. The presence of additional transition selection rules improves the discrimination of each contribution in the absorption and emission spectra, and in combination with the magnetic and luminescence measurements, it should become possible to determine without ambiguities the electronic structure diagram of the

investigated Ln(III) complexes. Such an enhanced analysis would ideally require that transitions do not overlap. However, in the general case, the detected signals consist of a sum of individual transitions with identical or opposite signs, which may lead to erroneous interpretations. For instance, a lack of chiroptical signal at a certain energy may be due to opposite components cancelling each other.

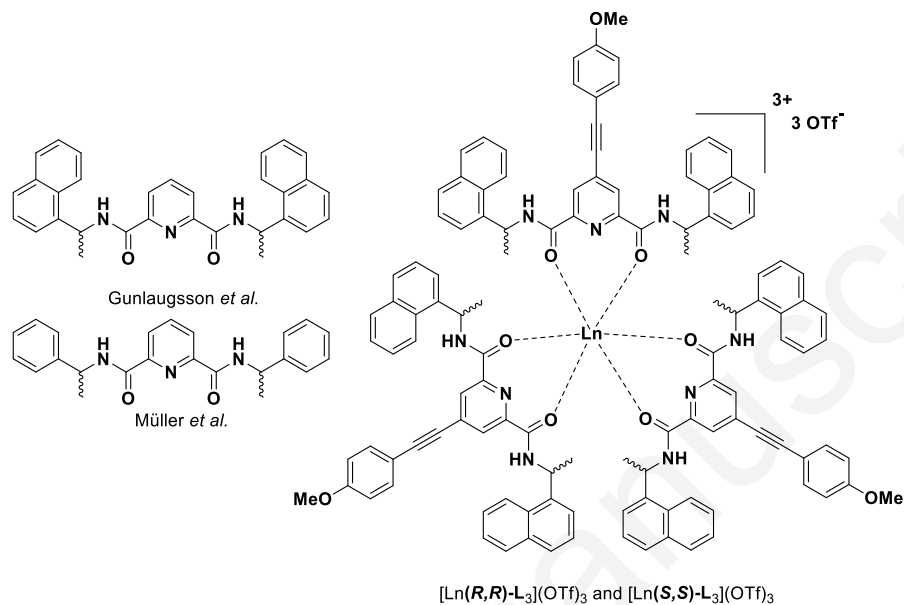


Chart 1. Tridentate ligand leading to the formation of 3-fold symmetric lanthanide complexes featuring an helicoidal chirality.¹³

Herein we present the synthesis, the structural determination, and the spectroscopic characterizations of a helical Yb(III)-based complex containing a charge transfer antenna, functionalized version of an analogous chiral complex developed by Muller and Gunlaugsson (see Chart 1).¹³ Our objectives are first to fully characterize the CF splitting of this chiral Yb(III) complex by combining magnetic measurements in solid state, with room and low-temperature absorption and luminescence measurements.¹⁴ Secondly, room temperature circular dichroism (CD) and CPL measurements have been carried out using an innovative CPL NIR setup for samples in solution. Hence ab-initio calculations rationalize the electronic structure and the transitions between the CF states and shed light on the current limitation of chiroptical measurements at room temperature.

Synthesis of Ligands and Complexes

The synthesis of the target ligand consists of the reaction of chelidamic acid (4-hydroxypyridine-2,6-dicarboxylic acid dihydrate) with thionyl chloride to form the transient acid chloride that further reacted with 2,6-bis-*N*-substituted chiral amide (*R*)-(+)- or (*S*)-(-)- 1-(1-naphthyl)ethylamine (see Scheme S1). During this step, the simultaneous aromatic nucleophilic substitution in para-position with a chloro atom was achieved. After replacement of the chloro with a iodo atom using already reported procedure,¹⁵ the palladium-catalyzed Sonogashira cross-coupling with 4-ethynyl anisole was performed to yield the two enantiopure antennas (*R,R*)- and (*S,S*)-L, obtained with an overall yield of ca. 50% for both the enantiomers (Scheme S1, for the detailed procedure see the experimental section and the ESI). The enantiopure complexes are synthesized in a 1:1 mixture of methanol and dichloromethane by adding the lanthanide salt (LnOTf₃, Ln = Nd, Pr, Eu, Dy, Tm, Er, Yb and Lu) to 3 equivalents of each enantiomer of the ligands (*R,R*)-L or

(S,S)-L (Scheme S1): the complexation process is instantaneous, accompanied by a slight color change of the solution from pale to more intense yellow. The pure complexes are isolated by filtration, after a precipitation adding diethyl ether, and fully characterized by ^1H , ^{13}C NMR and HRMS (see ESI).

Solid-State and Solution Structures

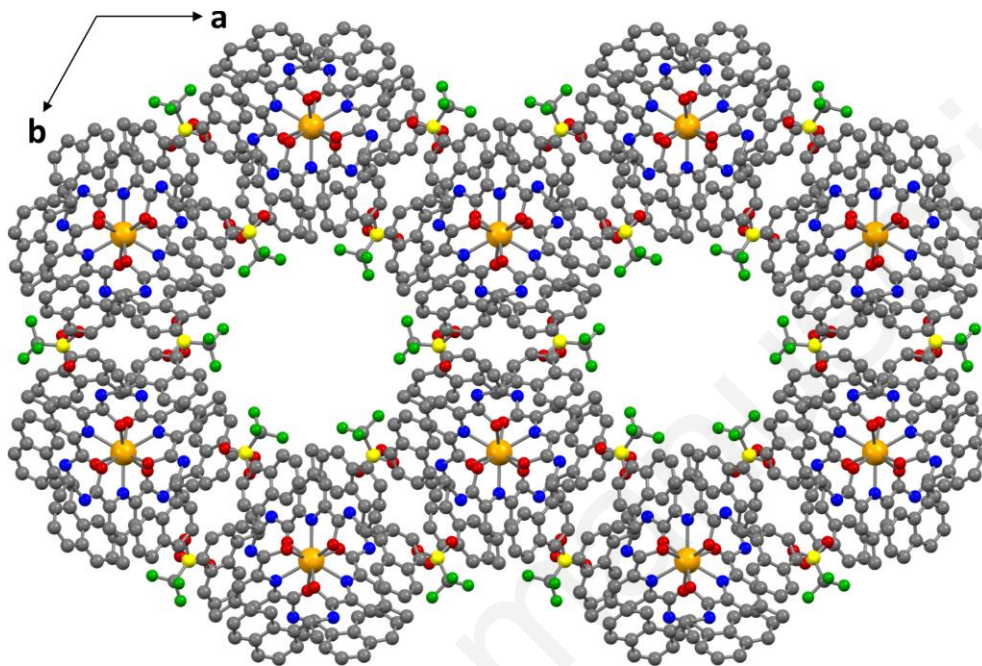


Figure 1. Partial X-ray structure of the $[\text{Er}(\mathbf{R,R})\text{-L}_3](\text{OTf})_3$ representing the packing as viewed along the c crystallographic axis. The 4-ethynylanisole moiety have not been determined and were replaced by methyl fragments. Protons have been omitted for clarity. Er, O, N, C, F and S atoms are in orange, red, blue, grey, green and yellow, respectively.

Single crystals for X-ray diffraction were grown as yellow needles by slow diffusion of diethyl ether in a methanol solution of the complexes. The best data were collected with crystals of the $[\text{Er}(\mathbf{R,R})\text{-L}_3](\text{OTf})_3$ complex. Unfortunately, only a partial structure was determined (see Figure 1), which misses the 4-ethynyl-anisole moiety on each ligand unit. Only the alkyne carbon directly linked to the pyridine ring could be unambiguously located in the electronic density map. However, a comparison can be made with the parent system reported previously by Gunnlaugsson *et al.* and shown in Figure S1.^{13a} The three $(\mathbf{R,R})\text{-L}$ ligands wrap around the lanthanide ion with a Λ -helicity promoting the π -interactions in-between each pyridine and naphthalene rings from the two other ligands (Figure S2). Surprisingly, in spite of the presence of the additional π -extended antennas, the complex crystals are almost isostructural to what previously observed within a $P6_3$ chiral space group (Figure S1). An identical hexagonal packing of the complex, promoted by intermolecular π - π stacking in-between naphthalenes, led to the formation of solvent channels where the ethynyl-anisole moieties stretched (Figure 1). Considering the ca. 8 Å radius of these cavities, a high degree of structural disorder of the antennas is expected, induced both by a possible free rotation of the terminal phenyl rings and steric clash of the three methoxy groups at the center of the channel.

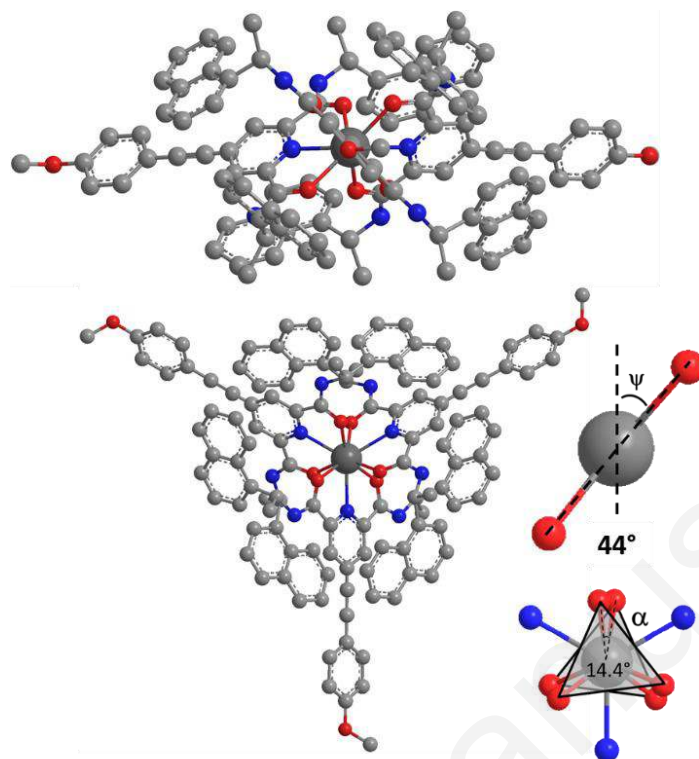


Figure 2. Chem3D model of $[\text{Yb}(\mathbf{R,R})\text{-L}_3]^{3+}$ PERSEUS structure along two views (hydrogen atoms omitted for clarity) with highlighted two important angles characterizing the polyhedron. Color code: C in gray, N in blue, O in red, Yb in dark gray.

To tackle this drawback, the solution structure of the $[\text{Yb}(\mathbf{R,R})\text{-L}_3](\text{OTf})_3$ complex was determined with the PERSEUS (Paramagnetic Enhanced Relaxation and Shift for Eliciting the Ultimate Structure) program¹⁶ and the resulting structure is given in Figure 2. Such procedure has been successfully applied to determine the solution structures of a large number of Ln-based complexes.¹⁷ It takes advantage of the high sensitivity of the NMR paramagnetic shifts to the complex geometry.¹⁸ Indeed, if one assumes that the spin density is localized at the Ln(III) nucleus, the paramagnetic NMR shift of a ligand nuclei i is then dominated by the pseudo-contact term ($\delta^{\text{PC}}(i)$),¹⁹ which is strongly dependent on the relative position of this nuclei with respect to the principal magnetic susceptibility anisotropy axes. Details about the full procedure are given in ESI. As expected, the coordination polyhedron configuration of the solution structure is Λ with the $(\mathbf{R,R})\text{-L}$ enantiomer and its shape is better described as a slightly distorted tricapped trigonal prism: the α angle, which quantifies the distortion from the regular prism, has a value of 14° with the polyhedron close to be achiral (Figure 2). The skewed angle ψ of the pyridine 2,6-bis-amide plane (which involves also the conjugated anisole moiety) with respect to the D_3 axis is 44° and confirms the moderate $\delta^{\text{PC}}(i)$ observed. Moreover, the dihedral angles $\text{C11-N-C12-C13} = 85^\circ$ and $\text{Me-C12-C13-C21} = 79^\circ$ (see Figure S15 for the numbering) describe the orientation of the naphthyl rings in space: the rings of one ligand unit are perfectly oriented for a π -stacking interaction with the pyridine rings of the others (distance around 4-4.5 Å). The strength of this interaction is a key point of the stability of the system and it can be interpreted as a case of a multi π -stacking network between electron rich rings naphthalenes and an electron poor pyridine. Moreover, the ethynyl anisole moiety is perfectly

coplanar with the 2,6-bis amido pyridine plane. All these structural parameters obtained from the solution structure determination compared well with those obtained by X-ray diffraction (Table S1).

Magnetic Characterizations

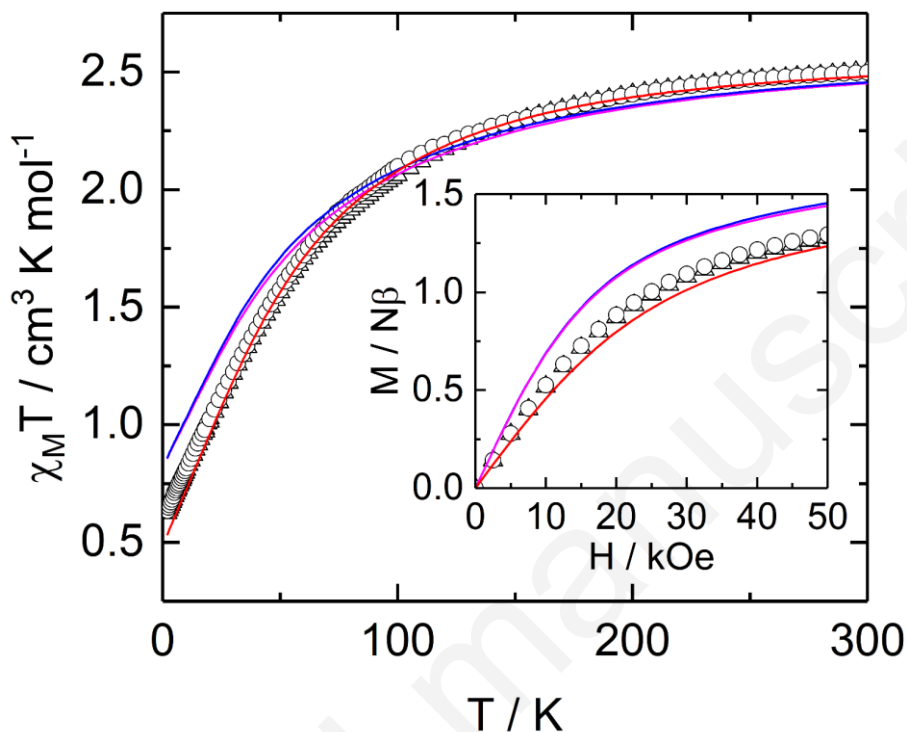


Figure 3. Experimental thermal variations of $\chi_M T$ (in $\text{cm}^3 \text{K mol}^{-1}$) for the $[\text{Yb}(\mathbf{R},\mathbf{R})\text{-L}_3](\text{OTf})_3$ (black circles) and $[\text{Yb}(\mathbf{S},\mathbf{S})\text{-L}_3](\text{OTf})_3$ (black triangles) with simulated curves: ab-initio PT2-SO (red line) and crystal field (see text for details) analyses (solution luminescence: magenta, solid state luminescence: blue). Inset: magnetization curves at 2 K with identical color code.

Magnetic studies were performed on microcrystalline powders of the Yb(III) derivatives and as expected, the two enantiomers behave similarly since the two $\chi_M T$ vs. T curves are perfectly superimposed (Figure 3). At room temperature, the $\chi_M T$ values ($2.5 \text{ cm}^3 \text{K mol}^{-1}$) agree with the expected value for a $^2F_{7/2}$ multiplet ground state with $g_J = 8/7$ ($2.57 \text{ cm}^3 \text{K mol}^{-1}$). On cooling, $\chi_M T$'s decrease down to $0.63 \text{ cm}^3 \text{K mol}^{-1}$ due to the splitting of the multiplet ground state. The M vs H curves are also perfectly superimposed (inset of Figure 3). Since these two systems do not show any out-of-phase component of the ac susceptibility in zero external dc field, they are not zero field SMM. The example of $[\text{Yb}(\mathbf{S},\mathbf{S})\text{-L}_3](\text{OTf})_3$ is given on Figure S3 and only the data for this compound will be presented here below. The application of a moderate external dc field induces an out-of-phase signal in the frequency window (1-1500 Hz) that does not significantly shift with the applied field but grows in amplitude (Figure S3). One can estimate that the optimum field for which the signal appears at the lowest frequency with the maximum amplitude is close to 1 kOe (Figure S3). At this field, the variation of the ac susceptibility with the oscillation frequency of the magnetic field can be quantitatively analyzed in the framework of the extended Debye model (see Figure 4 and Figure S4). The non-relaxing fraction of the magnetic susceptibility remains relatively small (~17-20%) that means that the vast majority of the magnetic moments are involved in the relaxation process (Table S2). The distribution of the relaxation time

remains small with α below 0.1 whatever the temperature (Table S2). α close to 1 means an infinite distribution of the relaxation time while α close to 0 means a single relaxation time. The thermal variation of the relaxation time (Figure 5) is reproduced with only Raman and direct relaxation processes ($\tau^{-1} = CT^n + AH^4T$).²⁰ The best-fitted curve is obtained with $C = 75.22 \text{ s}^{-1} \text{ K}^{-n}$, $n = 5.37$, $A = 2.21 \times 10^{-9} \text{ s}^{-1} \text{ Oe}^4 \text{ K}^{-1}$ (Figure 5). The Orbach process ($\tau^{-1} = \tau_0^{-1} \exp(\Delta/T)$) that guides the relaxation pathway through Kramers doublets excited states via Arrhenius law, appears to be too slow with reasonable τ_0 values and $\Delta = 124 \text{ K}$ (ab initio calculations) to interfere with the relaxation process.

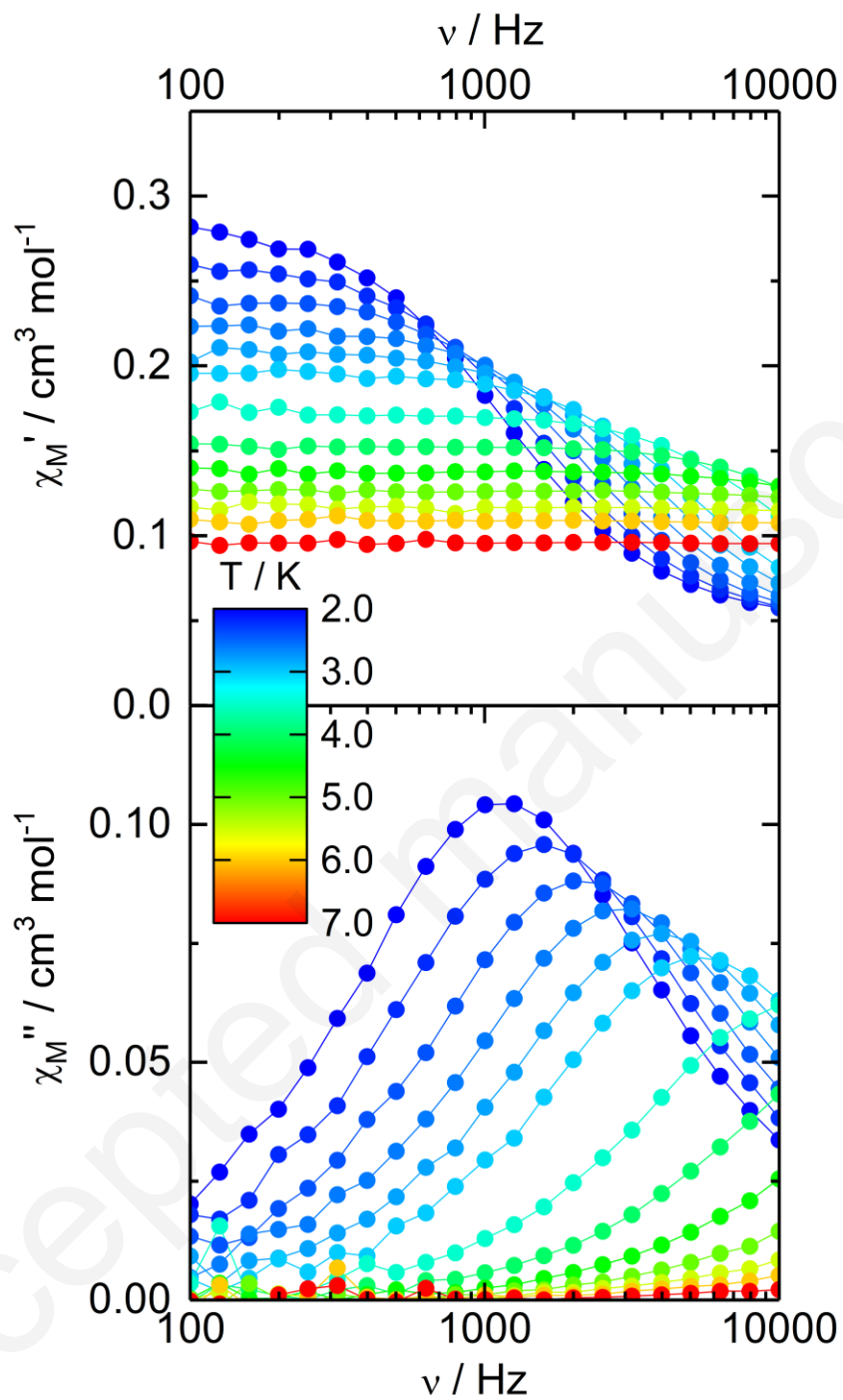


Figure 4. Frequency dependences of the two components, in-phase (χ_M' : top) and out-of-phase (χ_M'' : bottom), of the ac magnetic susceptibility measured at 1 kOe as a function of the temperature for $[\text{Yb}(\text{S,S})\text{-L}_3](\text{OTf})_3$.

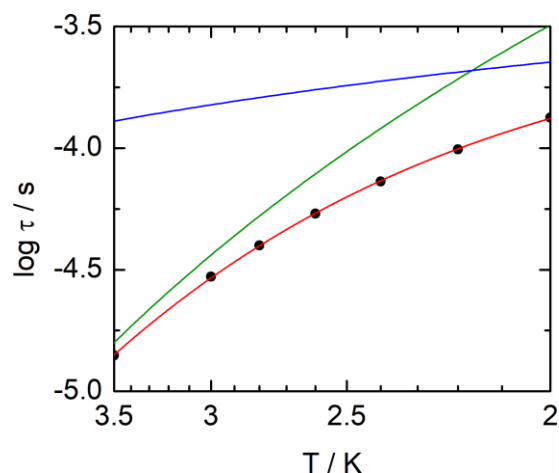


Figure 5. Thermal variation of the relaxation time (black dots) with the best fitted curve in red line and the separated contributions of the Raman (green line) and direct (blue line) processes for $[\text{Yb}(\text{S,S})\text{-L}_3](\text{OTf})_3$.

Luminescence Characterizations

Emission. The luminescence has been studied in solution and in the solid state both at room temperature and at 77 K (Figure 6). Upon excitation at 360 nm (Figures S5-S6), the classical Yb(III) emission profile is observed and corresponds to the ${}^2\text{F}_{5/2} \rightarrow {}^2\text{F}_{7/2}$ transition in the NIR. The emission spectra are similar for both enantiomers (Figure S7). Decreasing the temperature down to 77 K considerably narrowed the emission bands, suppressed the “hot bands” contribution and revealed the crystal field splitting fine structure. In a threefold symmetry, the ${}^2\text{F}_{7/2}$ multiplet ground state splits in four CF states noted 0,1,2,3 and in three for the ${}^2\text{F}_{5/2}$ excited state noted 0',1',2'. **Erreur ! Argument de commutateur inconnu.** In the 77 K solid state spectrum, four signals - three main peaks and one shoulder - can be clearly identified using Gaussian deconvolution (Figure S8, 979.2, 989.8, 998.2, 1013.7 nm) and consequently the energy diagram can be deduced assigning the 979.2 nm (10213 cm^{-1}) as the zero-line transition (ZLT) (0, 109, 194, 348 cm^{-1}). This simple spectroscopic experimental measurement enables a direct reading of the ${}^2\text{F}_{7/2}$ ground state CF. In frozen solution, the similar structure induces a comparable energy diagram obtained from the luminescence spectrum (Figures 6 and S8, 0, 107, 201, 355 cm^{-1}). Note that the experimental error for such measurements is in the nanometer range and corresponds to a precision in the energy splitting determination estimated to $\pm 10 \text{ cm}^{-1}$. The two CF splittings obtained in frozen solution and in the solid state at 77 K are almost identical, indicating that the coordination sphere remains the same in both cases. It is well known that the total CF splitting is a signature of the local symmetry of the Yb(III) coordination polyhedron.²¹ In the present case, the total splitting noted Δ_{CF} is small, about 348 cm^{-1} in the solid state (355 cm^{-1} in frozen solution). This value is perfectly in line with that previously determined in the literature for three-fold complexes like tris-dipicolinate ($\Delta_{\text{CF}} = 348 \text{ cm}^{-1}$),²² helicates ($\Delta_{\text{CF}} = 372 \text{ cm}^{-1}$),²³ or murex complexes ($\Delta_{\text{CF}} = 366 \text{ cm}^{-1}$) **Erreur ! Argument de commutateur inconnu.**^b and much lower than that of C_{2v} ($\Delta_{\text{CF}} = 670 \text{ cm}^{-1}$)²⁴ or even lower symmetry complexes (Δ_{CF} up to 880 cm^{-1}).²⁵ Finally, the luminescence lifetime for the two enantiomers of the Yb(III) complex has been measured (Figure S9) and fitted with a perfect mono-exponential decay giving $\tau = 6.1\text{--}6.2 \text{ }\mu\text{s}$, in the classical range for non-deuterated Yb(III) complexes in organic **solution**.^{11b,11f,22,24,25}

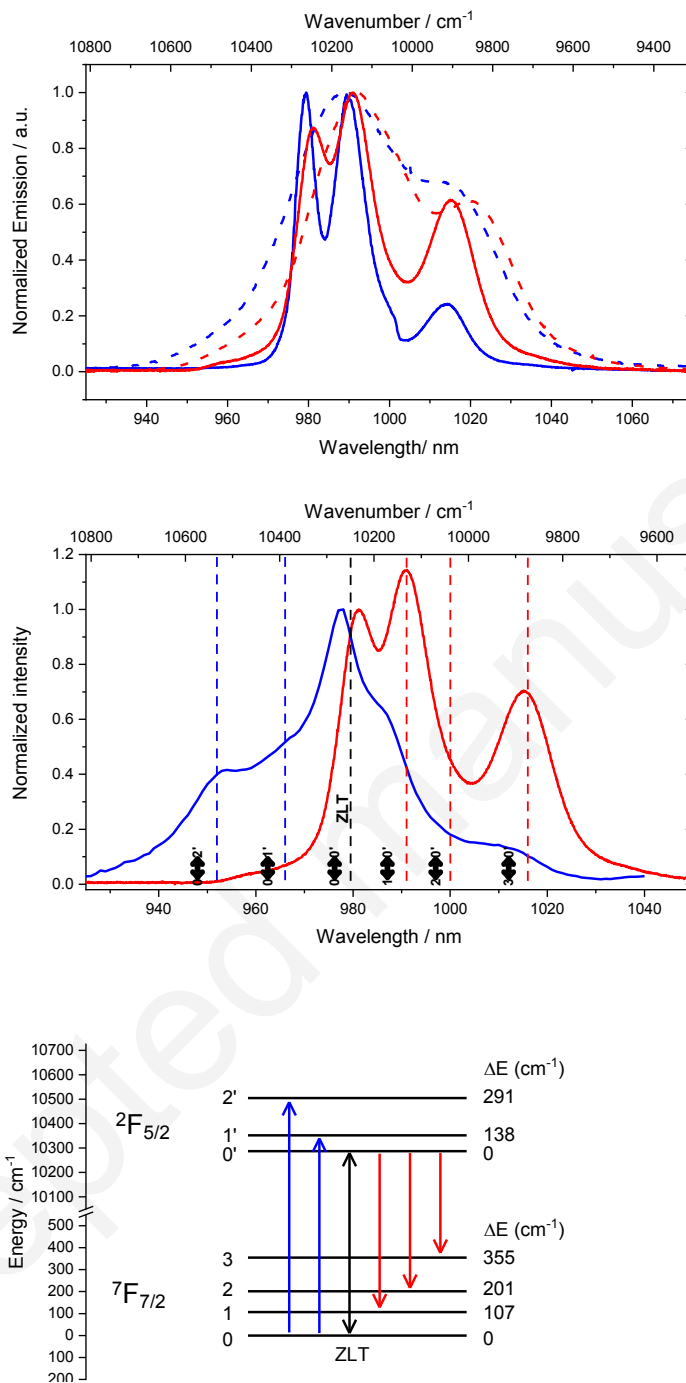


Figure 6. (top) Luminescence spectra of $[\text{Yb}(\text{S,S})\text{-L}_3](\text{OTf})_3$ in solution of ethanol : methanol = 4 : 1 at 298 (red dashed) and at 77 K (red solid); in solid state at 298 (blue dashed) and 77 K (blue solid). (Middle) Normalized absorption (blue) and emission (red) spectra at 298 K in methanol. (Bottom) Experimental energy diagram with related transitions.

Absorption. The NIR-absorption spectrum was measured only in solution and at room temperature. It spreads between 950 and 1050 nm and presents a rather poor resolution. Overlay with the emission spectrum (Figure 6)

confirms the ZLT position at 979.2 nm in the absorption, as well. At room temperature, the maximum in absorption is blue-shifted by 3 nm compared to the ZLT. Five main structures can be observed. The short-wave part of the spectrum corresponds to the $0 \rightarrow i'$ ($i' = 1', 2'$) transitions (shoulders at 952 and 966 nm) while the long wavelength side can be assigned to the transitions from thermally occupied ground state levels (*i.e.* hot bands in absorption). These three hot absorption bands (991, 1000 and 1016 nm) match well the emission spectrum. From this assignment, the CF splitting of the $^2F_{5/2}$ excited state can be deduced as 0, 138, 291 cm^{-1} .

Boltzmann distribution. The $^2F_{5/2}$ and $^2F_{7/2}$ CF splitting determined above reveals that the multiplet extension ($\sim 300 \text{ cm}^{-1}$) is in the same order of magnitude as kT at room temperature (208 cm^{-1} at 300 K). This population distribution leads to an overall broadening of the achiral optical features. We have then computed the variation of the Boltzmann's distribution with the temperature using a three and four levels Boltzmann model (see SI for complete calculation, Figure S10). It must be mentioned that all CF states are Kramers doublets and are then all doubly degenerated. A relative distribution of 57, 29 and 14% for the $0', 1', 2'$ CF states, respectively, is clearly determined for the $^2F_{5/2}$ excited state at room temperature. Population of the ground state multiplet is distributed as 46, 28, 18 and 8% for the $0, 1, 2, 3$ CF states. These values clearly show that all the multiplets are populated and, consequently, all the $i \leftrightarrow j$ transitions can have an impact on the spectra.

Chiroptical Characterizations

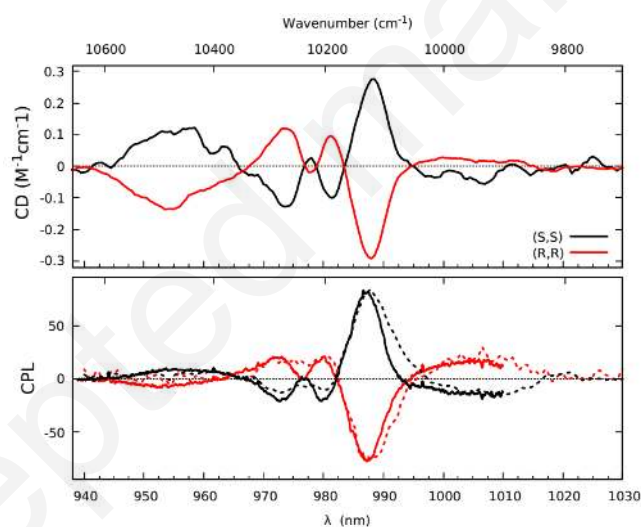


Figure 7. NIR-ECD (top) and NIR-CPL (bottom) spectra of $[\text{Yb}(\mathbf{R},\mathbf{R})\text{-L}_3](\text{OTf})_3$ (red) and $[\text{Yb}(\mathbf{S},\mathbf{S})\text{-L}_3](\text{OTf})_3$ (black) complexes in methanol solution at 298 K. For the NIR-CPL, two set-ups were used either with a CCD Si camera (continuous line) or an IR photomultiplier (dotted lines) as detectors.

The chiroptical super-spectrum²⁶ of the $[\text{YbL}_3](\text{OTf})_3$ complex, combining various chiroptical spectroscopies in different spectral range, visible electronic circular dichroism (ECD), NIR-ECD, circularly polarized luminescence (NIR-CPL) and vibrational circular dichroism (VCD) was determined for the first time for an ytterbium derivative (Figure 7 and Figures S11-S12). The NIR-ECD and NIR-CPL spectra of the two enantiomers of $[\text{YbL}_3](\text{OTf})_3$ were recorded in concentrated methanol solution (4 mM) with our home-made apparatus (see Supporting Information). For the CPL, we have used two set-ups: one is based on a CCD camera and waveplates similar to what is used for chiral contrast imaging.²⁷ It allows recording one-shot high-resolution spectra but is limited by the sensitivity of the Si sensors.²⁸ The

other one that uses a standard photoelastic modulator and a single channel detector, is less resolved but can be used down to 1500 nm.²⁹ It is worth noting that NIR-chiroptical measurements are relatively rare³⁰ compared to visible ones that are widely exemplified for europium(III), terbium(III) and samarium(III) chiral complexes. **Erreur ! Argument de commutateur inconnu.**²⁹

NIR-ECD and CPL spectra for the ${}^2F_{7/2} \leftrightarrow {}^2F_{5/2}$ transition of the $[\text{YbL}_3](\text{OTf})_3$ complexes are displayed in [Figure 7](#).

Erreur ! Signet non défini. This transition is magnetically allowed, highly sensitive to the Yb(III) ion environment and therefore particularly suited for chiroptical spectroscopy. At room temperature, the ECD and CPL spectra are much more structured than the absorption and emission ones. They are comparable in terms of energetic position and are mirror images for the two enantiomers. The strongest band appears in both ECD and CPL at around 987 nm (FWHM = 5 nm). This band is positive for the (S,S) enantiomer with significant absorption and emission dissymmetry factors $g = +0.04$ and $g_{\text{lum}} = +0.11$. On the low energy side of this band, a broad CPL signal of opposite sign extends from 994 to 1016 nm. The corresponding ECD signal is hardly observable as expected for a hot band in absorption. At shorter wavelengths, we clearly count three bands in CPL at 980 (negative), 973 (negative) and 954 nm (positive) about 4 times lower in intensity. This last transition is much broader (FWHM = 20 nm) than the other two (around 5 nm). The ECD spectra get the same overall shape but the short wavelength intensities are stronger and a small peak of opposite sign appears between the two peaks at 973 and 980 nm. The different features observed at room temperature do not match with the energies of the CF states determined from absorption and emission spectroscopies (*vide supra*). This apparent inconsistency has been investigated under the light of ab-initio calculations ([Figure 9](#)). The short wavelength band (around 950 nm) is calculated as a $2' \leftrightarrow 0$ transition which also explains the fact that it is more intense in ECD than CPL. This negative band is followed at higher wavelengths by a complicated structure mixing three positive bands involving $0'$, $1'$ and $2'$ states as well as the negative contribution from the $1' \leftrightarrow 1$ transition. The exact position, as well as the relative intensities of these bands can lead to the experimental shape observed in the 965-980 nm spectral range ([Figure 9](#)). The strong band at 988 nm is the sum of the $1' \leftrightarrow 2$, $2' \leftrightarrow 3$ and $0' \leftrightarrow 1$ transitions. Its intensity will depend strongly on the temperature of measurement as a consequence of the thermal population of the different ${}^2F_{5/2}$ sublevels. The long wavelength CPL signal is related to the $1' \leftrightarrow 3$ transition, it is thus a hot band in absorption leading to a low ECD signal. Based on all these interpretations, it clearly appears that the apparently well-resolved ECD and CPL spectra at room temperature are in fact the combination of all CF states and consequently the experimental determination of the energy diagram from room temperature chiroptical data remains hardly possible.

Ab-initio Calculations

In order to gain further insight into the optical and magnetic properties of the $[\text{Yb}(\mathbf{R},\mathbf{R})\text{-L}_3](\text{OTf})_3$ complex, multi-reference calculations were performed at the SCF/PT2-SO level on a model compound $[\text{YbL}'_3]^{3+}$ resulting from the solution structure determined by NMR (see computational details). The nature of the ground state (GS) was first characterized with the help of the natural spin orbitals (NSOs) as shown in [Figure 8](#). These NSOs represent the distribution of the unpaired electron among the 4f spin-orbitals.³¹ The largest positive spin-populations (i.e. spin-up) are calculated for the NSOs $4f_{yz}$ (0.40 e^-) and $4f_{z(x^2-y^2)}$ (0.31 e^-) which are associated to a projected orbital angular momentum $m_l = \pm 2$. Additionally, a sizable negative spin-population (-0.22 e^-) is calculated for the NSO $4f_{x(x^2-3y^2)}$ corresponding to a $m_l = \pm 3$. The sum of these spin-populations gives a spin expectation value $\langle S_{\parallel} \rangle = 0.21$, revealing the importance of the SO coupling in the nature of the GS. Using the total angular momentum J and its projections M_J ,

the distribution of the unpaired electron derives formally from the Kramers doublet $|7/2, \pm 5/2\rangle$ of the ${}^2F_{7/2}$ multiplet. For reference, the composition of this M_J state for a $4f^{13}$ free ion in term of m_l and m_s would be of ca. 86% $|\pm 2, \pm 1/2\rangle$ and 14% $|\pm 3, \pm 1/2\rangle$. However, the almost trigonal environment around the Yb(III) ion in $[\text{YbL}'_3]^{3+}$ leads to a mixing of the $M_J = \pm 5/2$ with the $M_J = \pm 1/2$ and $\pm 7/2$. As seen in Figure 8, the small deviation from a perfectly trigonal environment creates a symmetry breaking in the wave function that is visible with the sizable mixing with the $M_J = -3/2$ state. This mixing of the different M_J states into the SO wave function is responsible for the small magnetic anisotropy of the GS, characterized by the two large perpendicular components of the g-factors and a relatively small parallel component ($g_{\parallel} = 3.38$). The calculated isotropic magnetization is found in really good agreement with the experimental one at very low T , suggesting a proper description of the nature of the GS (Figure 3). Additionally, for the temperature-dependence of $\chi_M T$ a nice agreement with the experiment is obtained at the PT2-SO level. This behavior corresponds to the thermal population of the excited M_J states of ${}^2F_{7/2}$ that are calculated at 86, 138 and 290 cm^{-1} above the GS, respectively.

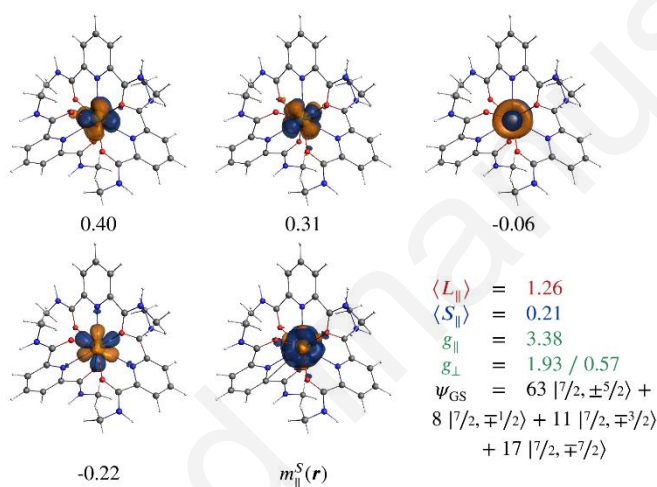


Figure 8. Selected Natural Spin Orbitals (NSOs) (isosurface values = ± 0.03 a.u.) and their corresponding spin populations for the Kramers doublet GS obtained at the PT2-SO level. The isosurface (± 0.001 a.u.) of the spin-magnetization ($m_{\parallel}^S(r)$) for the doublet component $\langle S_{\parallel} \rangle > 0$ is also shown. The orbital and spin expectation values, the EPR g-factors and the nature of the wave function calculated for the GS are given for comparison. Additional data are presented in Tables S3 and S4 of the Supporting Information.

The calculated CPL spectrum is shown in Figure 9. Here, the calculations were performed at the RASPT2 level in order to mix the $4f^{13}$ configurations with the $4f^{12}5d^1$ ones. This mixing of configurations allows gaining electric dipole intensities for the 4f-4f transitions by breaking the parity rule.³² Such a procedure was successfully applied recently to obtain the CPL spectrum of an analogous Eu(III) tris-dipicolinate complex.³³ At 0 K, the calculated CPL spectrum exhibits a strong positive and a strong negative band corresponding to the $0' \rightarrow 0$ and $0' \rightarrow 1$ transitions, respectively. The increase of the temperature leads to the appearance of a second positive band corresponding to the $1' \rightarrow 0$ transition, as observed experimentally. However, the comparison with the experimental spectrum suggests that the excited state $1'$ is calculated too low in energy as the intensity of this hot band becomes too important at room temperature. Interestingly, the breakdown of the calculated CPL spectrum into the different contributions from the three excited states of ${}^2F_{5/2}$ (Figure 9) reveals that the principal negative CPL band does not only correspond to the $0' \rightarrow 1$ transition as suggested experimentally, but also contains contributions from the $1' \rightarrow 2$ and $2' \rightarrow 3$ transitions.

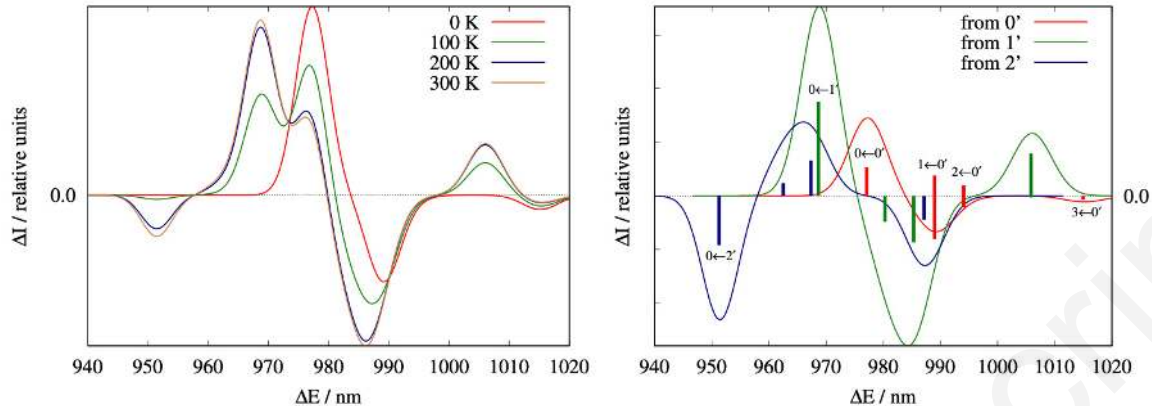


Figure 9. Calculated CPL spectrum of $[\text{Yb}(\text{R,R})\text{-L}'_3]^{3+}$ as a function of the temperature (left) and the different contribution at 300K from the ${}^2\text{F}_{5/2}$ term. Calculated spectra have been shifted by 0.07 eV to match the experimental $0' \rightarrow 0$ transition. Calculated raw data are given in Table S5.

Crystal-field analysis

Based on the luminescence characterizations of the $[\text{YbL}_3](\text{OTf})_3$ complex, it is then possible to establish its electronic structure and hence, determine the CF parameters. At first, one can assume that the Yb(III) ion is located in a trigonal environment (D_3 symmetry), and express the CF Hamiltonian as:

$$\hat{H}^{CF} = \alpha_J[\text{B}_2^0\tilde{\text{O}}_2^0(J)] + \beta_J[\text{B}_4^0\tilde{\text{O}}_4^0(J) + \text{B}_4^3\tilde{\text{O}}_4^3] + \gamma_J[\text{B}_6^0\tilde{\text{O}}_6^0(J) + \text{B}_6^3\tilde{\text{O}}_6^3(J) + \text{B}_6^6\tilde{\text{O}}_6^6(J)] \quad (\text{Eq 1.})$$

where the α_J , β_J and γ_J are parameters obtained using the Wigner-Eckart theorem and depend on the considered J-manifold (here $J = 7/2$), the $\tilde{\text{O}}_k^q(J)$ are the Stevens operators, and the B_k^q terms are the CF parameters. Since the ground term ${}^2\text{F}_{7/2}$ is well separated from the excited one (over 10,000 cm^{-1}), the matrix elements $\langle J, \text{M}_J | \tilde{\text{O}}_k^q | J, \text{M}_J' \rangle$ can be simply determined by using the values originally tabulated by Stevens.²⁰ In this work, the B_2^0 term in Eq. 1 was determined by performing luminescence measurement on the Eu(III) analogue complex and by measuring the energetic splitting of the ${}^7\text{F}_1$ multiplet (see Figure S13).³⁴ Indeed, in a trigonal CF the splitting corresponds to 3 times the term B_2^0 ; the sign being determined by the ordering of the degenerate vs non-degenerate components of the multiplet. The analysis of the emission spectrum in solution leads to J = 1 transition CF splitting of 90 cm^{-1} , and hence to a B_2^0 term equal to ca. -30 cm^{-1} . The value of B_2^0 for the Yb(III) analogue differs probably from the Eu(III) case. However, theoretical investigations have shown that for a given series and assuming small geometrical relaxations,³⁵ the B_2^0 remains almost constant along the series, and therefore B_2^0 was kept fixed at -30 cm^{-1} . The rest of the B_k^q terms were then obtained with least-squares procedures by fitting the energies obtained from luminescence measurements carried out on the investigated complex. The results of each set of CF parameters are given in Table 1 along with the nature of the associated GS. To assess the quality of the fits, the CF parameters were then used to simulate the magnetic properties ($\chi_{\text{M}}T$ product and magnetization) and compared to the experimental data (see Figure 3).

As visible in Table 1, the magnitude of CF parameters remains relatively close between the different luminescence measurements. The fact that the B_2^0 , B_4^0 and B_6^0 terms are of the same order of magnitude and all negative leads to a GS that is mainly of $M_J = \pm 5/2$ character. These small differences in the CF parameters hardly affect χ_{MT} (see Figure 3), which converges to the experimental data at room temperature. However, at low-temperature the experimental χ_{MT} is overestimated by the CF analyses, suggesting a too large magnetic moment for the GS. This result is confirmed with the over-estimation of the magnetization at 2K (see Figure 3). The difference at low temperature between the CF analyses and the experimental data can be rationalized with the help of the ab-initio calculations. As mentioned previously, the $[\text{Yb}(\mathbf{R},\mathbf{R})\text{-L}_3](\text{OTf})_3$ complex is not perfectly trigonal but corresponds to a slightly distorted tricapped trigonal prism (see Figure 2). The lowering of symmetry induces a larger mixing in the GS wavefunction with other M_J states (see Figure 8), which reduces the magnetic moment of the GS. Such an admixture of states cannot be reproduced by a strict trigonal CF. As visible with the calculated ab-initio CF parameters (Table S5), a larger number of CF parameters is considered in order to get the proper wavefunction. It results in a great reproduction of the magnetic susceptibility and magnetization curves (Figure 3).

Table 1. Experimental and calculated energy splitting (in cm^{-1}) of the ground multiplet $^2F_{7/2}$ in $[\text{Yb}(\mathbf{R},\mathbf{R})\text{-L}_3](\text{OTf})_3$. The CF parameters (B_k^q , in cm^{-1}) deduced from these relative energies and the composition of the GS wavefunction are also given for comparison.

	Lum Sol. (77K)	Lum Solid (77K)	Ab-initio* Sol.
$^2F_{7/2}$	0	0	0
	107	109	86
	201	194	138
	355	348	290
B_2^0	-30	-30	-24
B_4^0	-98	-91	-47
B_4^3	-85	-173	-80
B_6^0	-22	-20	-11
B_6^3	1178	1183	-29
B_6^6	262	289	-56
% $M_J = \pm 1/2$	13	14	8
% $M_J = \pm 3/2$	0	0	11
% $M_J = \pm 5/2$	78	75	63
% $M_J = \pm 7/2$	8	11	17

*the additional ab-initio CF parameters are given in Table S6.

Advanced theoretical calculations revealed that chiroptics cannot be analyzed in the same way (*vide supra*). In normal absorption or emission measurements if two different transitions fall at the same wavelength intensities of the transitions gets added: there is no negative absorption or emission. In CD and CPL the signal of each transition at this wavelength can be either positive or negative. The resulting measured signal being the sum it can be positive, negative or even null depending on the respective amplitudes. In other words, the non-observation of a signal does not automatically mean that there are no transitions. It thus makes the analysis of room temperature experiments hazardous especially for lanthanide based complexes in which the crystal field splitting is small and emission lines

close to each other's inside the same multiplet, like in $[\text{YbL}_3](\text{OTf})_3$ system. Indeed, the CPL spectrum is temperature dependent with great modulations and shift of the various lines (as demonstrated with calculations; Figure 9). If one wants chiroptical information to be a source of CFS information one needs low temperature CPL measurements, a task that still needs to be achieved.³⁶

Conclusions

The electronic structure of the chiral $[\text{YbL}_3](\text{OTf})_3$ complex has been characterized with the help of a combination of magnetic, absorption and luminescence characterizations carried out at low and room temperature in solution and in solid-state. From these spectroscopies, we have been able to determine the energetic splitting of the ground $^2F_{7/2}$ and excited $^2F_{5/2}$ terms arising from the crystal-field splitting. The temperature resolved luminescence spectroscopy produced similar electronic structures that with the ab-initio calculations. The CF picture was then achieved by assuming a perfect trigonal environment (D_3 symmetry) that allows reproducing qualitatively temperature variation of the magnetic susceptibility. However, this CF picture was not able to reproduce the proper magnetic moment of ground state because the $[\text{YbL}_3](\text{OTf})_3$ complex exhibits a slight deviation from the ideal D_3 symmetry.

NIR chiroptical measurements (ECD and CPL) at room temperature revealed features that can only be rationalized on the basis of advanced theoretical calculations since the superposition of various transitions at almost the same wavelength in lanthanide-based complexes truncates the analysis in the Boltzmann landscape. This makes the extraction of CF transitions impossible at room temperature. The quantitative description of chiroptical properties thus waits for low-temperature measurements. In this line, g_{lum} values that are considered in the literature as intrinsic observables of each emission line may be experimentally the result of a sum of various contributions and might not be considered definite. In following works, we will combine thus variable-temperature magneto-lumino-chiroptical techniques on a wider range of symmetries and environments to deeper investigate the CF parameters origin.

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