

MATTERS ARISING

Giant solar flares in Antarctic ice

ROOD *et al.*¹ have discovered four prominent 'spikes' in a long time record (circa 1150 to the present) of the NO_3^- concentration inside an Antarctic ice core. These four spikes rise 2–3 times higher than the upper envelope of a fluctuating background level of $0\text{--}20 \mu\text{g l}^{-1}$ that has been plausibly attributed to the action of high-energy solar radiation (photons and particles) impinging on the Earth's upper atmosphere and ionizing N_2 , thereby leading to various chains of chemical reactions that culminate in the formation of NO_3^- , some of which is transported, within a few weeks or months, to Antarctica¹. According to three alternative chronologies provided by Rood *et al.*, the estimated dates of the four NO_3^- spikes lie within the intervals given in Table 1. Three of these dates have been tentatively associated by the same authors with the galactic supernovae of 1604, 1572 and 1181. At the outset, they have rejected energetic particles from the supernova explosion as a possible source of ionization of the terrestrial N_2 because galactic magnetic fields would have greatly delayed and diffused the particles on their way to Earth. Instead, they have shown that photons of energy $\geq 10 \text{ keV}$ are required. Unfortunately, as they admitted, the total energy requirements are difficult to meet, and the matching of dates with historical supernovae is not perfect.

As an alternative explanation, I suggest that the necessary ionizing radiation could have come from unusually powerful solar flares. These flares would be expected to have occurred preferentially during periods when the Sun was generally most active, that is around the times of the largest maxima in the solar cycle. Two good indices of solar activity are available: for the years elapsed since 1700, there are both sunspot numbers² and auroral numbers³; for earlier years auroral statistics^{3–6} are preferred because the sunspot record (mostly from the Far East) is very sporadic⁷. Despite some incompleteness of the record before 1700, the main trends in the statistics are quite unmistakable (see refs 3 and 5).

The intervals of time in which the largest auroral and sunspot maxima occurred are listed in Table 1, where earlier dates are given only to the nearest half-decade. These intervals of time correlate very well with the known epochs of the NO_3^- spikes. Only in one case is it necessary to recognize that episodes of solar flaring need not occur (as they have not always occurred in modern times) precisely at times of maximum auroral or

maximum sunspot numbers. This therefore gives some flexibility in the possible dates of the giant solar flares that is not available in the case of the supernova hypothesis. However, there are no NO_3^- spikes in the years around 1778 and 1957, at which times solar activity was also at a peak. Nevertheless, given the rarity of the proposed flaring events, this absence could simply be a product of statistical fluctuations. Also the Sun may now be somewhat different physically from its state before the long Maunder minimum in 1645–1715 (ref. 2).

According to Table 1, very large solar maxima seem to recur in cycles of ~ 200 yr, as Schöve⁴ originally noted. In the background NO_3^- data, there is also some evidence of the Maunder minimum and of the normal 11-yr solar cycle¹. Bauer⁸ estimated that the background concentration could vary by a factor of two during the 11-yr cycle. To extend this further, I consider the largest solar flares in modern times. These have importance class 3+ or 4 and emit $E_{32} \times 10^{32}$ erg of high-energy radiation (photons and particles), where $E_{32} \sim 1\text{--}2$ (ref. 9). The Earth intercepts $4 \times 10^4 E_{32}$ erg cm^{-2} of this. According to Rood *et al.*¹, the energy flux needed to produce a NO_3^- concentration of $20 C_{20} \mu\text{g l}^{-1}$ near the South Pole is $\sim 1 \times 10^5 C_{20}$ erg $\text{cm}^{-2} \text{yr}^{-1}$. Since $C_{20} \leq 1$, only one or two major flares per year is sufficient to produce all of the background NO_3^- . It is therefore not unreasonable to suppose that, every couple of hundred years or so, a giant flare, perhaps 2–3 times more intense than ordinary major flares, erupts on the Sun. An alternative possibility is that a very rapid succession of major flares of the ordinary type takes place. An event of this type may still develop during the current maximum of solar activity.

The Antarctic ice-core measurements are apparently now being pushed to deeper levels than before¹. Much older NO_3^- spikes may therefore be discovered. One immediate prediction of the solar flare hypothesis is the possibility (though no more than that) of a spike occurring around the year 1000, a time of height-

ened auroral activity^{3–5}. On the other hand, a very bright supernova also appeared in 1006, as Rood *et al.* pointed out. But, fortunately, an *experimentum crucis* to discriminate between the two hypotheses can be made for the middle of the eleventh century, a time of profound auroral quiet but, equally importantly, of a brilliant supernova, the Crab explosion of 1054.

RICHARD STOTHERS

NASA Institute for Space Studies,
Goddard Space Flight Center,
2880 Broadway,
New York, New York 10025

1. Rood, R. T., Sarazin, C. L., Zeller, E. J. & Parker, B. C. *Nature* **282**, 701 (1979).
2. Eddy, J. A. *Climatic Change* **1**, 173 (1977).
3. Gleissberg, W. *Solar Phys.* **21**, 240 (1971).
4. Schöve, D. J. *J. geophys. Res.* **60**, 127 (1955).
5. Link, F. *Planet. Space Sci.* **12**, 333 (1964).
6. Siscoe, G. L. preprint (1980).
7. Clark, D. H. & Stephenson, F. R. Q. *J. R. Astr. Soc.* **19**, 387 (1978).
8. Bauer, E. *Nature* **276**, 96 (1978).
9. Talon, R. *et al.* in *Solar Gamma-, X-, and EUV Radiation* (ed. Kane, S. R.) 315 (Reidel, Dordrecht, 1975).

Magnetic fields and the solar constant

THOMAS¹ has suggested that changes in the magnetic flux content of the convection zone produces changes in radius. However, his calculations did not include the effect of structural changes in the superadiabatic region on the bulk of the convection zone. Theoretical studies^{2–4} have shown that solar luminosity fluctuations can result from small structure adjustments in the convection zone, and can occur on time scales shorter than 1 yr. Such fluctuations are of interest in studies of the terrestrial climate.

The previous calculations have forced these structural changes by assuming a time-dependent mixing length. When the mixing length (physically the convective efficiency) changes, adjustments in the structure occur rapidly in the superadiabatic region, and in turn the bulk of the convection zone adjusts to maintain hydrostatic equilibrium. The result is a temporary luminosity change³. Any mechanism which affects the structure of the superadiabatic region will cause such a luminosity fluctuation. The effect of magnetic pressure changes on the superadiabatic region is described here.

We began by including a global magnetic pressure term in a stellar structure code. If the flux is assumed to be concentrated in vertical flux tubes, the pressure at a given radius (r) is given by

Table 1 Dates of the largest maxima in the Antarctic concentration of NO_3^- and in solar activity indices

Largest NO_3^- maxima	Largest solar maxima
1130–1160	1120–1140
1300–1340	1360–1375
1590–1600	1565–1585
1610–1620	1605–1630
–	1778–1788
–	1947–1959

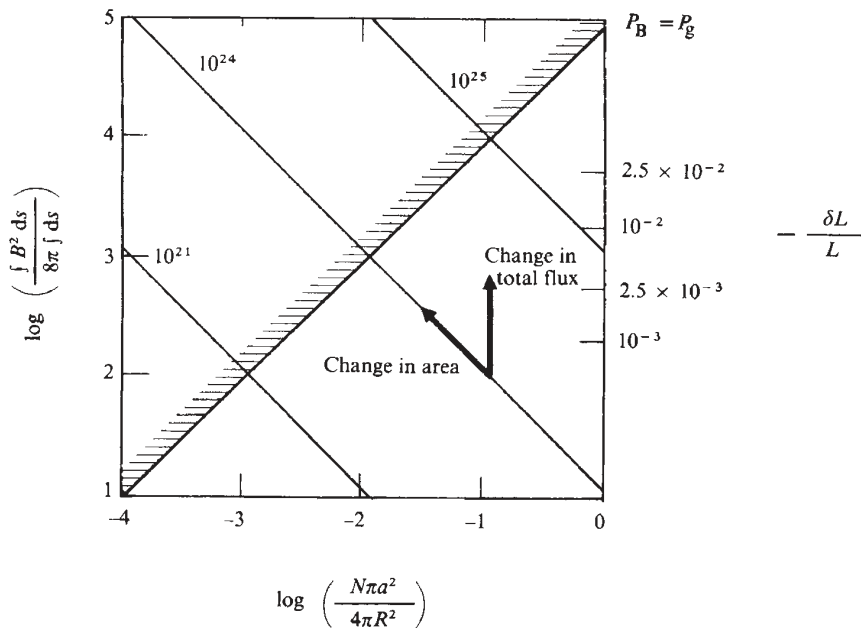


Fig. 1 Global magnetic pressure plotted against the fraction of the solar surface contained in the magnetic flux tubes. The luminosity scale on the right side gives the change resulting from the magnetic pressure change. The oblique lines are contours of constant flux.

$$P_m = \frac{1/8\pi \int B^2 ds}{\int ds} \approx \frac{NB^2 a^2}{32\pi r^2}$$

where N is the number of tubes with flux B and radius a . If the flux in a tube is constant and material is locked to the field lines, then the ratio of magnetic pressure to gas pressure decreases rapidly with depth in agreement with detailed sunspot models by Weiss⁶.

The effect of changing the value of the magnetic pressure was explored by changing the surface value through the range $\log P_m = 1$ to 3.8 dyn cm^{-2} . Only in the superadiabatic region (upper few thousand kilometres) was the magnetic pressure term significant. Its gradient results in expansion and decreased density. In response to the lowered density in the superadiabatic region, the bulk of the convection zone begins to expand. The energy for expansion is obtained from the luminosity, and thus increasing the magnetic pressure decreases the solar luminosity.

The magnetic pressure depends on the total magnetic flux present and the area containing that flux. However, there is a minimum area over which a given amount of flux may be distributed that is determined by the requirement that the magnetic pressure in the flux tubes should not be greater than the gas pressure outside the tube.

The surface value of the global magnetic pressure is plotted against the fraction of the solar surface contained in the magnetic flux tubes in Fig. 1. Lines of constant total flux are shown as diagonals and the effect on the pressure of changing the area contained in flux tubes can be seen by moving along one of these lines. The effect on the pressure of changing the

total amount of flux contained in the flux tubes is shown by vertical motion in this diagram. The luminosity scale along the right side of Fig. 1 gives the amplitude of the luminosity fluctuation produced by the magnetic pressure change.

The solar budget of magnetic flux is uncertain, but to produce a luminosity change of the order of 0.1% the flux penetrating the superadiabatic region must be $\geq 2 \times 10^{23} \text{ Mx}$. The radius change associated with this flux change was found to be $\Delta R/R \leq 10^{-4}$ which is slightly less than the value found by Thomas¹ for two reasons. First, the vertical flux tubes considered here contain less volume through the superadiabatic region than the horizontal flux tubes considered by Thomas. Second, the radius of the upper parts of the superadiabatic region is affected by the luminosity change because energy transport through the outer 1,000 km is partially by radiation^{3,4}. Thus reducing the radiative flux causes this region to contract.

While the total amount of magnetic flux in the convection zone is uncertain, the amount of flux required for its pressure to produce a significant luminosity change is apparently high. A 1% luminosity change required a flux change of $2 \times 10^{24} \text{ Mx}$ in our model, and resulted in a radius change of $\Delta R/R = 2 \times 10^{-4}$. The action of such large amounts of magnetic flux on convective energy transport may actually dominate the effects discussed here. Unfortunately, although it is well known that strong fields can suppress convection locally (sunspots), the realistic inclusion of the effects of such flux tubes on global convective energy transport is a problem. As the effects of the field on convection have been neglected, the present results should represent the minimum effect of

magnetic fields on the solar luminosity and radius.

This work was supported by The Aerospace Corporation on company-financed funds.

D. S. P. DEARBORN

Steward Observatory,
University of Arizona,
Tucson, Arizona 85721

J. B. BLAKE

Space Sciences Laboratory,
The Aerospace Corporation,
Box 92957
Los Angeles, California 90009

1. Thomas, J. *Nature* **280**, 662 (1979).
2. Ulrich, R. *Science* **190**, 619 (1975).
3. Dearborn, D. & Newman, M. *Science* **201**, 150 (1979).
4. Dearborn, D. & Blake, J. B. *Astrophys. J.* (in the press).
5. Parker, E. *Astrophys. J.* **198**, 205 (1975).
6. Sofia, S., O'Keefe, J., Lesh, J. & Endal, A. *Science* **204**, 1306 (1979).

A tectonic mélange of foreign eclogites and ultramafites in West Norway

MEDARIS¹ suggests that eclogites enclosed within the amphibolite-facies country-rock gneisses in the Basal Gneiss Region, West Norway "developed by *in situ* metamorphism of crustal materials"¹; furthermore he attributes this to an "eclogite-facies metamorphic event"¹. There are, however, some weaknesses in these hypotheses; also the simple *in situ* hypothesis may detract other geologists, especially geodynamicists, from pondering the genesis of these rocks, some of which are amongst the most exceptional on Earth², being the highest-pressure metamorphic rocks in the Caledonide Orogen³.

Many different kinds of eclogite can be distinguished in the region according to mineralogical, textural, chemical, and pressure (P)–temperature (T) criteria, such that the frequent assumption^{1,4–8} of a single origin (the simple *in situ* origin) for all eclogites within country-rock gneiss is difficult to justify. The 'dual' origin of eclogites and ultramafites, with only the ultramafites being foreign, suggested at Lien¹ and elsewhere^{7,9}, is less credible than my alternative multiple foreign origin¹⁰ of the different eclogite types and also of those meta-igneous rocks (ultramafites, anorthosites, dolerites) for which adequate evidence of igneous intrusion into country-rock gneiss is unavailable (most published examples). A gigantic tectonic mélange is thus envisaged having been created by tectonic introduction of rock fragments from diverse foreign sources of diverse ages of equilibration [for example eclogite-facies eclogites + ultramafites (upper mantle); granulite-facies eclogites + anorthosites + gneisses + autometamorphosed dolerites