



Mapping the performance of a versatile water-based condensation particle counter (vWCPC) with COMSOL simulation and experimental study

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15 Abstract.

16 Accurate airborne aerosol instrumentation is required to determine the spatial distribution of ambient aerosol particles, 17 particularly when dealing with the complex vertical profiles and horizontal variations of atmospheric aerosols. A 18 versatile water-based condensation particle counter (vWCPC) has been developed to provide aerosol concentration 19 measurements under various environments with the advantage of reducing the health and safety concerns associated 20 with using butanol or other chemicals as the working fluid. However, the airborne deployment of vWCPCs is relatively 21 limited due to the lack of characterization of vWCPC performance at reduced pressures. Given the complex 22 combinations of operating parameters in vWCPCs, modeling studies have advantages in mapping vWCPC 23 performance. 24 25 In this work, we thoroughly investigated the performance of a laminar flow vWCPC using COMSOL Multiphysics® 26 simulation coupled with MATLAB. We compared it against a modified commercial vWCPC (vWCPC Model 3789, 27 TSI, Shoreview, MN, USA). Our simulation determined the performance of particle activation and droplet growth in 28 the vWCPC growth tube, including the supersaturation, $D_{p,kel,0}$ (smallest size of particle that can be activated), 29 $D_{\rm p,kel,50}$ (particle size activated with 50% efficiency) profile, and final growth particle size $D_{\rm d}$ under wide operating temperatures, inlet pressures P(0.3 - 1 atm), and growth tube geometry (diameter D and initiator length L_{ini}). The 30 31 effect of inlet pressure and conditioner temperature on vWCPC 3789 performance was also examined and compared 32 with laboratory experiments. The COMSOL simulation result showed that increasing the temperature difference (ΔT)

between conditioner temperature T_{con} and initiator T_{ini} will reduce $D_{p,kel,0}$ and the cut-off size $D_{p,kel,50}$ of the vWCPC. In addition, lowering the temperature midpoint $(T_{mid} = \frac{T_{con} + T_{ini}}{2})$ increases the supersaturation and slightly

decreases the $D_{p,kel}$. The droplet size at the end of the growth tube is not significantly dependent on raising or lowering the temperature midpoint but significantly decreases at reduced inlet pressure, which indirectly alters the vWCPC

empirical cut-off size. Our study shows that the current simulated growth tube geometry (D = 6.3 mm and $L_{ini} = 30$

38 mm) is an optimized choice for current vWCPC flow and temperature settings. The current simulation can more

39 realistically represent the $D_{p,kel}$ for 7 nm vWCPC and also achieved a good agreement with the 2 nm setting. Using

40 the new simulation approach, we provide an optimized operation setting for the 7 nm setting. This study will guide

41 further vWCPC performance optimization for applications requiring precise particle detection and atmospheric aerosol

42 monitoring.





43 1 Introduction

Aerosols, defined as any solid or liquid particles suspended in air, are one of the fundamental components of the atmosphere and have a significant impact on air quality, climate change and human health (Seinfeld et al., 2016; Anderson et al., 2020; Lighty et al., 2000; Pöschl, 2005; Prather et al., 2020; Li et al., 2017; Paasonen et al., 2013; Darquenne, 2012). However, accurate and comprehensive monitoring of aerosol particles is challenging because aerosol particle sizes and number concentrations vary widely both spatially and temporally (Davidson et al., 2005; Yu and Luo, 2009; Krudysz et al., 2009). Airborne measurements and characterization, therefore, are often required to capture the vertical profiles and horizontal variability of atmospheric aerosols.

52 In understanding the variability of atmospheric aerosol and determining the size distribution and number concentration of aerosols, laminar-flow, butanol-based Condensation Particle Counters (CPCs) used in conjunction with differential 53 54 mobility analyzers (DMAs) can provide real-time measurements of airborne particles and are widely used by 55 effectively exploiting the working principle of condensation growth (Hermann et al., 2007; Kangasluoma and Attoui, 56 2019; Mordas et al., 2008; Sem, 2002; Wiedensohlet et al., 1997). However, conventional butanol CPCs face 57 difficulties in characterizing particles below 3 nm in size. In addition, health and safety risks, such as the odor, 58 flammability and toxicity of the butanol, are an issue for many deployments in offices, homes, aircraft, and other 59 inhabited locations. These limitations have led directly to researchers designing advanced aerosol instruments that can 60 be more widely used in both atmospheric environments and laboratory studies.

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62 In 2005, Hering and Stolzenburg (2005) developed a continuous-flow, water-based laminar condensation particle 63 counter (WCPC) (Hering et al., 2005; Hering and Stolzenburg, 2005), which uses distilled water as the working fluid 64 to avoid the health and safety concerns. It was also found to have comparable performance to butanol-based CPCs in previous studies (Biswas et al., 2005; Franklin et al., 2010; Iida et al., 2008; Kupc et al., 2013; Liu et al., 2006; Mordas 65 et al., 2008). A modified version of the WCPC featuring an additional new moderator section has been developed 66 67 (Hering et al., 2014). With this new moderated approach, the initiator provides water vapor for particle activation 68 while the moderator provides distance and time for particle growth. This improved water CPC achieves the same peak 69 supersaturation and similar droplet growth while reducing the water vapor, particle loss, and side effects of heating 70 the flow in the earlier version of WCPC. Furthermore, a versatile WCPC was then developed capable of particle 71 detection near 1 nm without using a filtered sheath flow. The operating temperatures can also be adjusted in accordance 72 with the cut-point desired (Hering et al., 2017).

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Since water-based CPCs have comparable performance to butanol-based CPCs while also offering the advantage of avoiding health and safety risks, it is desirable to explore advanced water-based CPCs in a broader range of environmental and energy applications. To improve the detection performance of the vWCPC, we need to identify the effects of operational factors and geometry. However, the limited analysis of relevant temperature and geometric parameters in the vWCPC makes it challenging to control condensational growth conditions. In addition, the inlet pressure effect is another critical factor affecting the detection efficiency of vWCPCs. The potential of using vWCPC





for airborne deployment or other lower pressure monitoring has not been fully explored. Mei et al. (2021) found that the counting efficiency of the vWCPC 3789 operated at the factory settings decreased with decreasing the operating pressure, particularly at operating pressures below 700 hPa. However, determining how to reduce the lower detection limit under various ambient pressures also needs to be investigated.

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85 Due to the complex matrix of geometry, operating temperature, and inlet pressure parameters in vWCPCs, modeling 86 studies are advantageous in determining and optimizing the detection efficiency of vWCPCs. The Graetz model was 87 first used by Stolzenburg (1988) to examine the detection of ultrafine particles in CPCs. In recent years, COMSOL 88 Multiphysics® has been widely used to simulate coupled heat, mass, and momentum transfer problems associated with complex geometries in CPCs. Moreover, COMSOL has advantages in interfacing with post-processing software 89 such as MATLABTM. A series of parametric analyses for butanol CPCs were simulated using COMSOL to investigate 90 91 the performance of particle activation and droplet growth (Hao et al., 2021; Attoui, 2018; Kangasluoma et al., 2015; 92 Barmpounis et al., 2018; Thomas et al., 2018). Our previous work (Hao et al., 2021) first demonstrated that the 93 COMSOL results neglecting the temperature dependence of vapor thermodynamic properties and axial diffusion, agree with the Graetz solution used by Stolzenburg (1988). Considering temperature dependence of vapor 94 95 thermodynamic properties and axial diffusion can generate more accurate results that can guide the optimization of CPC designs. Previous research using COMSOL to examine vWCPC performance has been limited. Bian et al. (2020) 96 97 compared two-stage and three-stage operating temperature methods for growth tubes and parameters such as flow rate 98 and temperature difference to obtain the ideal activation and final growth sizes. Mei et al. (2021) used COMSOL 99 aiding to examine how inlet pressure affects particle activation in the vWCPCs. However, the lack of thorough and 100 systematic examination of vWCPC performance using COMSOL leaves it unclear on how well the vWCPC will 101 perform in multiple complex research areas and applications, such as at reduced atmospheric pressure levels.

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In this study, we thoroughly determined the saturation profile, activation efficiency, and droplet growth for varying 103 airborne operations through numerical simulation of laminar flow vWCPC in COMSOL and experimental validation 104 of a commercial vWCPC (TSI Model 3789). By mapping vWCPC performance in the modeling, the effects of various 105 operational factors, such as inlet pressures (0.3 - 1 atm), growth tube diameter and initiator length, and temperature 106 107 gradients on particle activation and droplet growth, were investigated. In addition, detailed modeling methods are outlined below. The detection efficiency was also examined in the experimental and COMSOL modeling work. The 108 109 results of this study will guide further optimization of the performance of vWCPCs for accurate detection of particles 110 and atmospheric aerosol measurement applications.





112 2 Methods

113 2.1 Numerical simulation

114 2.1.1 COMSOL setup

115 The finite element COMSOL simulation software (COMSOL Multiphysics 5.3a, COMSOL Inc, Stockholm, Sweden) 116 can handle a variety of fields including, but not limited to, electromagnetics, fluid dynamics, heat transfer, chemical 117 reactions, and structural mechanics. Here, a two-dimensional axisymmetric model is developed to simulate fluid flow 118 in a cylindrical tube. The heat, momentum, and mass transfer equations are solved for incompressible parabolic flow. 119 This COMSOL model follows the three-stage tube of the versatile water-based CPC (TSI Inc, Shoreview, MN, USA) 120 described by Hering et al. (2017), which consists of a fully developed laminar flow tube that can be separated into a 121 cool-wall conditioner region, a warm-wall initiator region, and a cool-wall moderator region, where r is the radial 122 coordinate of the tube diameter and z is the axial coordinate of the tube length (Fig. 1a). At the inlet of the conditioner 123 tube, sampled aerosols are fed and saturated with water vapor before entering the initiator region. The manufacturer 124 provides two default cut-off diameter settings: 2 and 7 nm configurations based on the characteristics introduced by 125 Kangasluoma et al. (2017). This study used the 7 nm configuration as an example to demonstrate the simulation 126 mapping effort. The methodology can also be used for other targeted cut-off sizes. The tube diameter (D) is 6.3 mm, 127 the conditioner length (L_{con}) is 73 mm, the initiator length (L_{ini}) is 30 mm, and the moderator length (L_{mod}) is 73 mm. The default settings are (Table 1): the conditioner temperature (T_{con}) is 30 °C, the initiator temperature (T_{ini}) is 59 °C, 128 and the moderator temperature (T_{mod}) is 10 °C. The aerosol flow rate (Q_v) is 0.3 L min⁻¹. The relative humidity (RH) 129 of inlet flow is set at 20%, and the water vapor is assumed to be saturated at the wall. The inlet pressure (P) is 1 atm. 130 To investigate how the vWCPC performance depends on these parameters, we simulate a wide range of values for 131 132 mapping the vWCPC geometry, working temperature conditions, and inlet pressure, as discussed in Section 2.1.4 for the tasks in this study. 133

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135 In this COMSOL model, the coupled heat transfer and fluid flow are first solved by the conjugate heat transfer module, and then the mass transfer of the water vapor is solved based on the obtained temperature and flow field. Lastly, the 136 137 particles are introduced from the inlet of the vWCPC and are considered diluted species that follow the convective 138 diffusion equation in the simulation, which is numerically solved to calculate the temperature, supersaturation, and 139 water vapor concentration profiles (Fig. 1). Regarding the initial and boundary value settings, the inflow temperature 140 is set to an ambient temperature of 25 °C, and wall temperature is set to the default factory settings described above 141 and changed as the experimental conditions changed for mapping performance. The simulated temperature and concentration fields are then imported into MATLABTM R2022b (The MathWorks, Inc., Natick, MA, USA) and 142 143 interpolated at any given point in the r-z plane. COMSOL Multiphysics 5.3a with MATLAB allows us to adjust 144 geometry, modify physics settings, perform parametric studies, control solvers, and post-process the results. 145





146 **2.1.2 Theory of particle activation**

Particle activation is key to the evaluation of CPC performance. Particle activation within the vWCPC depends on the degree of supersaturation, namely the saturation ratio (S) of water vapor, which is the ratio of the partial pressure of the water vapor (p) to the saturation vapor pressure of the water vapor (p_s) for the given flow temperature (T),

$$S = \frac{p}{p_s}$$

(1)

151 The spatial profile of *S* within the vWCPC allows us to calculate the Kelvin effect, the homogeneous nucleation, and 152 further condensational particle growth, as discussed in Section 2.1.3. The Kelvin effect is dependent upon 153 thermodynamic principles and described as the Kelvin equivalent size $(D_{p,kel})$, the minimum diameter of a particle 154 that can be activated for condensation growth. It is determined by water surface tension (σ), water molecular volume 155 (v_m) , the Boltzmann constant (k), temperature (T) and the saturation ratio (S) calculated at each location within the 156 initiator,

$$D_{\rm p,kel} = \frac{4\sigma v_{\rm m}}{kT ln(S)} \tag{2}$$

157 The Kelvin equivalent size is inversely proportional to the distribution of saturation ratio, where the greater the 158 saturation ratio, the smaller the size of the particles that can be activated. In other words, the smaller the particle, the 159 higher the degree of supersaturation required to activate growth. When particle size (D_p) is above $D_{p,kel}$, the particle 160 can be successfully activated and grown by water vapor condensation, while when D_p is below $D_{p,kel}$, the particles 161 cannot be activated. **Fig. 1b - 1e** show examples of saturation ratio, Kelvin equivalent size, water vapor concentration, 162 and temperature profiles within the simulated geometry at the default temperature condition of $T_{con} = 30$ °C, $T_{ini} = 59$ 163 °C, $T_{mod} = 10$ °C, respectively.

164

165 It is worth noting that $D_{n,kel}$ varies at different locations of the initiator due to the spatial variation of temperature, 166 surface tension and saturation ratio. We observed the potential particle activation in the moderator by simulation 167 results. However, only the Kelvin equivalent size in the initiator was considered in this work. Although the particles were activated in the moderator, particle detection was unlikely due to droplet growth being the dominant water vapor 168 169 sink in the initiator. Thus, in reality, the actual supersaturation in the moderator may be different, and the activation 170 of smaller particles will be hindered by droplet growth (Hering et al., 2014; Hering et al., 2017). Note that the $D_{p,kel}$ in 171 the conditioner region is blank in color due to no particles being activated in this region ($S \leq 1$). Particles near the 172 wall of the initiator, where there is a lower S, are more difficult to activate due to the larger $D_{p,kel}$. This difference can 173 be explained by the water mass and heat diffusivity differences. As the colder, water-saturated flow passes through 174 the growth tube, the mass transport of water vapor is faster than the heating of the flow from the wall because the mass 175 diffusivity of water vapor is higher than the thermal diffusivity of air, producing a maximum supersaturation of water 176 vapor at the centerline of the tube. As a result, the seed particles entering near the centerline of the growth tube are 177 activated in the warmer initiator. One comparison of the saturation ratio and Kelvin equivalent size along the centerline 178 (r = 0) is shown in Fig. S1a and S1b. We observed the appearance of a double-peaked saturation ratio curve. Again,





- 179 only the Kelvin equivalent size in the initiator was considered in this work due to insufficient water vapor and droplet 180 growth in the moderator. Note that our calculations do not include the solute effects, and we assume wettable insoluble 181 particles in the modeling.
- 182
- 183 The activation efficiency of particles with a size of $D_{\rm p}$ in vWCPC is derived using an approach similar to our previous
- 184 work (Hao et al., 2021), which is calculated by the ratio of the number concentration of the activated particles over
- 185 the total particle number concentration in the growth tube. The activation efficiency is calculated as

$$\eta_{\rm act} = \frac{\int_0^{\kappa_{\rm act}} 2\pi r w N dr}{Q_{\rm v} N_0} \tag{3}$$

186 where w is the velocity along the axial direction, N is the concentration of particles, both at the axial location of z =187 $Z_{\text{act.}} R_{\text{act}}$ is the maximum radius of the contour corresponding to $D_{\text{p,kel}} = D_{\text{p}} N_0$ is the particle concentration at the inlet of the conditioner. Q_v is the flow rate through the vWCPC. An example of activation efficiency as a function of 188 particle diameter D_p can be found in Fig. S1c. Note that the calculation of the activation efficiency in Eq. (3) does not 189 190 consider the diffusion loss of the particles in the conditioner. Since this model does not have sheath flow that 191 minimizes the diffusion losses and constrains the aerosols to the high supersaturation region, the activation efficiency cannot reach 100%. On the activation curve, there are two points of interest: minimum activated size, D_{p,kel,0} (the 192 193 smallest size of particle that can be activated in the initiator), and 50% cut-off size, D_{p,kel,50} (the size of a particle with 194 50% activation efficiency extracted from the activation efficiency curve). $D_{p,kel,50}$ is essential to the performance of CPCs because it determines the general particle size range in which the CPC can confidently measure. D_{p,kel,50} can 195 196 be used as the representative of particle activation efficiency performance. Furthermore, note that negligible 197 homogeneous nucleation occurs in the growth tube of the initiator and moderator under all tested conditions in this 198 study, which means the total nucleation rate is equal to or less than one particle per second (1 s^{-1}) .

199

200 2.1.3 Theory of droplet growth

201 Once the particles are activated, their condensational growth along their trajectories in the initiator region was 202 simulated by numerically solving two coupled differential equations in MATLABTM. First, the evolution of droplet 203 diameter (D_p) can be governed by (Seinfeld and Pandis, 2008; Wang et al., 2017)

$$\frac{dD_{\rm p}}{dt} = \frac{4D_{\rm v}'M\left(C - C_{\rm d}\right)}{\rho} \tag{5}$$

204 where M and ρ represent the molecular weight and density of water, D'_v represents the modified diffusivity of the

205 water vapor accounting for the non-continuum effect of the particles and is given by $D'_{\rm v} = D_{\rm v} \left[1 + \frac{2D_{\rm v}}{\alpha_{\rm c} D_{\rm p}} \left(\frac{2\pi M}{RT} \right)^{1/2} \right]^{-1}$,

- 206 where D_v is the diffusivity of the water vapor, and α_c is the mass accommodation coefficient of water and is assumed
- 207 as 1. C_d represents the equilibrium water concentration at the surface of the growth droplets and is given by $C_d =$





 $C_{\rm s}(T_{\rm d})exp(\frac{4\sigma M}{\rho RT_{\rm d}D_{\rm p}})$, where $C_{\rm s}$ is saturation water concentration, $T_{\rm d}$ is the droplet surface temperature, which is 208 ed by

$$\frac{dT_{\rm d}}{dt} = \frac{3}{c_{\rm p}\rho D_{\rm p}} \left(H_{\rm vap}\rho \frac{dD_{\rm p}}{dt} - 4k'_{\rm g} \frac{(T_{\rm d} - T)}{D_{\rm p}} \right) \tag{6}$$

210 where c_p , ρ , and H_{vap} are the heat capacity, density, and heat of vaporization of water. k'_g represents the modified thermal conductivity of air accounting for the non-continuum effects in heat transfer and is calculated as k'_g 211 $k_{\rm g} \left[1 + \frac{2k_{\rm g}}{\alpha_{\rm T} D_{\rm p} \rho_{\rm g} c_{\rm p,g}} \left(\frac{2\pi M_{\rm g}}{RT}\right)^{1/2}\right]^{-1}$, where $M_{\rm g}$, $\rho_{\rm g}$, $c_{\rm p,g}$, and $k_{\rm g}$ are the molecular weight, density, heat capacity, and 212 thermal conductivity of air. α_T is the thermal accommodation coefficient and was assumed to be 1 (Seinfeld and 213 214 Pandis, 2008; Wang et al., 2017).

215

216 In the simulation of droplet growth, the Brownian motion of the particles inside the conditioner before activation is 217 neglected. Since there is also no electric field inside the conditioner, we assume that the particles move axially along 218 the vWCPC with a velocity of w. Therefore, Eqs. (5) and (6) can be converted to a function of axial location using 219 w = dz/dt, and the droplet size and droplet surface temperature at the end of the moderator ($z = L_{ini} + L_{mod}$) can be 220 calculated. To determine the final droplet growth size at the outlet of the moderator, the condensational growth of 8 221 nm particles as seed particles was studied along the centerline (r = 0) of the growth tube in this work. An example of 222 droplet growth size as a function of distance along the axis of the tube for the default temperature condition of T_{con} = 223 30 °C, T_{ini} = 59 °C, T_{mod} = 10 °C can be found in Fig. S1d. Note that in this simulation, we do not consider the 224 increase in the equilibrium vapor pressure due to warming of the flow from condensational heat release, which would 225 further reduce the droplet growth.

226

227 2.1.4 Simulation plan

228 Table 1 summarizes the operating temperatures, inlet pressures, and geometric parameters for each simulation task 229 characterizing the vWCPC in this study. In Task 1, we first conducted two matrix combinations of absolute conditioner 230 and initiator operating temperatures, with an interval of 5 °C in each region, for a total of nine different combinations. Task 2 investigates how raising or lowering the temperature midpoints (T_{mid}) , the average value between the 231 conditioner temperature and initiator temperature $\left(\frac{T_{con}+T_{ini}}{2}\right)$ affects particle activation and droplet growth. In 232 233 addition, different inlet pressures are also included when comparing different temperature midpoints. Task 3 examines the effect of inlet pressure by comparing the default temperature of 30 °C and the conditioner temperature of 27 °C. 234 Tasks 4 and 5 further test how the vWCPC geometry, including tube diameter D, and initiator length Lini, affects the 235 performance of the vWCPC. These simulations reveal optimal working conditions and effects for the influence of 236 237 each parameter.





239 2.2 Experimental measurement

240 The modified vWCPC 3789 (TSI Inc, Shoreview, MN, USA) was tested in this study. Operating flow, temperatures 241 and geometry can be found in Section 2.1.1 and our previous study (Mei et al., 2021). Two methods were used to 242 generate the test aerosol: an atomizer coupled with a furnace; and a glowing wire generator (GWG). Ammonium 243 sulfate (AS) has been commonly used for CPC characterization and was the tested material used in this study (Hering 244 et al., 2014; Kangasluoma et al., 2017). It was dissolved into deionized water for aerosol generation using atomization 245 techniques. To increase the aerosol number concentration for particles less than 30 nm, polydisperse AS aerosols were 246 also passed through a tube furnace generator (Lindberg/Blue, Thermal Scientific, TX, USA) to shift the size 247 distribution to a smaller size. A lab-built GWG was also used to generate aerosol particles in size range between 2.5 248 -16 nm. More details about the generator can be found in Attoui (2022). Using the low-pressure testing setup shown in Fig. S2, the counting efficiency of a vWCPC 3789 was measured between 0.5-0.9 atm for AS particles of 3 - 20249 nm (mobility diameter) and NiCr oxidants of 2.5 -16 nm. The aerosol concentrations in this test were maintained in 250 the range of $2 \times 10^4 - 4 \times 10^4$ cm⁻³. During the testing, the temperature variations in the conditioner and moderator were 251 less than 0.5 °C, and the initiator temperature had a variation of 1 °C. The y-axis error bar indicates the standard 252 253 deviation of the counting efficiency averaged over ~ 5 min of sampling time at a 1 Hz sampling rate.

254

255 3 Results and discussion

256 **3.1 Comparisons of temperature-dependent particle activation and droplet growth performance**

Selection of appropriate operating temperatures in CPCs is essential because the supersaturation is significantly 257 258 temperature dependent, which affects particle activation and further droplet growth. In addition, the temperature 259 difference between different regions in CPCs is an important factor in controlling supersaturation. For this reason, the minimum activation size for butanol-based CPCs is significantly impacted by the temperature difference between the 260 saturator and condenser and the raising or lowering of the temperature midpoints, as has been demonstrated by many 261 previous studies (Hermann and Wiedensohler, 2001; Kangasluoma and Attoui, 2019; Barmpounis et al., 2018; Kuang 262 263 et al., 2012). The results showed that in the butanol-based CPCs, the greater the temperature difference between the saturator and condenser, the higher the degree of supersaturation, and the smaller particle could be activated. 264

265

The numerical COMSOL model was used to compare operating temperature-dependent particle activation and droplet 266 267 growth performance in the vWCPC, including minimum activated size $(D_{p,kel,0})$, 50% cut-off size $(D_{p,kel,50})$, and final growth particle size at the outlet of the moderator along the centerline $(r = 0) (D_d)$, as shown in Fig. 2. Previous studies 268 269 confirmed that the centerline saturation rate is insensitive to the moderator wall temperature (Hering et al., 2014; Bian 270 et al., 2020). Thus, this study investigated moderator temperature ($T_{\rm mod}$) at the constant of 10 °C, conditioner temperature (T_{con}) at the range of 25 – 35 °C, initiator temperature (T_{ini}) at the range of 55 – 65 °C. Note that conditions 271 272 that can lead to a lower $D_{p,kel,50}$ value and larger droplet growth size are favored for improving the performance of 273 the vWCPC.





275 Firstly, in order to compare the effect of the conditioner temperature T_{con} , we increased T_{con} from 25 °C to 35 °C while maintaining the same initiator temperature T_{ini} and moderator temperature T_{mod} , $D_{p,kel,0}$ increased significantly 276 by 5.21, 3.32, and 2.27 nm at the initiator temperature T_{ini} of 55, 60, and 65 °C, respectively, and D_{p,kel,50} increased 277 278 significantly by 6.65, 4.16, and 2.75 nm at the initiator temperature T_{ini} of 55, 60, and 65 °C, respectively. The final 279 droplet size D_d decreased by approximately 1 µm at all the initiator temperature T_{ini} of 55, 60, 65 °C. The lower 280 conditioner temperature provided higher saturation ratios in the initiator and more water vapor for particle growth, 281 which is also consistent with the previous growth tube simulation (Bian et al., 2020; Mei et al., 2021). Secondly, the initiator temperature T_{ini} was increased from 55 °C to 65 °C while maintaining the same conditioner temperature T_{con} 282 283 and moderator temperature T_{mod} , $D_{\text{p,kel,0}}$ was decreased by 1.90, 3.74, and 4.87 nm at the conditioner temperature 284 $T_{\rm con}$ of 25, 30, and 35 °C, respectively, $D_{\rm p,kel,50}$ was decreased significantly by 2.32, 3.74, and 6.21 nm at the 285 conditioner temperature T_{con} of 25, 30, and 35 °C, respectively, and D_d was increased by 2.9 μ m at all the conditioner 286 temperature T_{con} of 25, 30, and 35 °C.

287

288 By comparing all combinations, we can find that the activated size becomes smaller as the temperature difference between $T_{\rm con}$ and $T_{\rm ini}$ increases, indicating that the temperature differences between the conditioner and initiator 289 290 dominate the particle activation. After comparing the temperature differences, we conclude that the higher the 291 temperature between these two regions, the better the particle activation. However, in the actual operation of the CPC, one also needs to ensure that the self-nucleation in the growth tube is minimized (<1 s⁻¹) so that the CPC does not 292 293 report false particle counting. The homogeneous nucleation rate is less than 10-8 s-1 at all tested conditions, meaning 294 that the temperatures can be further adjusted to optimize particle activation and droplet growth. Moreover, D_d is the 295 greatest, with a maximum size of 12.20 µm, at the temperature setting of 25-65-10 °C among all these temperature conditions. We also found that the effect of the initiator temperature on droplet growth was greater than that of the 296 297 conditioner temperature. Thus, the following section examines the effect of temperature midpoint on the vWCPC 298 performance.

299

300 **3.2 Effect of temperature midpoint on particle activation and droplet growth performance**

301 In addition to temperature difference, lowering the temperature midpoint was also found to cause higher 302 supersaturation. However, there is limited research on how the performance of the vWCPC changes under various 303 temperature midpoints and especially under different inlet pressures, which will be important for applications such as 304 atmospheric airborne deployment and environmental monitoring at elevated locations. Here, we compared the particle activation and droplet growth performance for three different temperature midpoints (40 °C, 43 °C, and 46 °C) of 305 306 conditioner temperature (from 24 °C to 30 °C) and initiator temperature (56 °C to 62 °C) at a wide range of inlet 307 pressures from 0.3 atm to 1 atm, as shown in Fig. 3. The temperature difference ΔT between the conditioner and the 308 initiator was kept constant at 32 °C. The moderator temperature remained constant at 10 °C in all simulations. 309





310 Results show that the minimum activated size $D_{p,kel,0}$ decreases from 5.15 nm to 4.96 nm, and the 50% cut-off size 311 $D_{p,kel,50}$ decreases from 5.88 nm to 5.65 nm as the temperature midpoint decreases from 46 °C to 40 °C, as shown in 312 Fig. 3a and 3b. Thus, a slight control of the minimum activation size can be achieved by lowering the temperature 313 midpoint. Higher supersaturation can explain this slight decrease in the initiator, which also agrees with the previous 314 growth tube WCPC simulation (Bian et al., 2020). On the other hand, a slight increase of 0.07, 0.1, and 0.14 nm occurs 315 in $D_{p,kel,0}$, and negligible change in $D_{p,kel,50}$ by reducing the inlet pressure from 1 to 0.3 atm under three temperature midpoints of 40 °C, 43 °C, and 46 °C. This slight increase is due to a low peak supersaturation caused by the decrease 316 317 in inlet pressure. Since water vapor transport is faster than heat transport, the decrease in pressure affects the location 318 of the peak supersaturation, whereas the degree of the supersaturation does not change significantly. 319 320 In Fig. 3c, we show that the droplet growth is not significantly dependent on raising or lowering the temperature 321 midpoint. By lowering the temperature midpoint by 6 °C, D_d becomes smaller by approximately 14%. When studying 322 the effect of inlet pressure on the D_d , unlike $D_{p,kel,0}$ and $D_{p,kel,50}$, D_d decreases substantially from 1 to 0.3 atm, by approximately 45%. Limited by the optical chamber design of the commercial vWCPC, the droplets smaller than 8 323 µm may not gain sufficient pulse signal to get counted. Thus, when operating under lower inlet pressure, the apparent 324 325 cut-off size of vWCPC may increase and needs to be further determined. The reduced pressure strongly affects the 326 final droplet growth size, likely due to the faster water vapor and heat transport at reduced pressure. The thermal and 327 mass diffusivity is inversely proportional to the pressure in the growth tube, resulting in insufficient time for droplet 328 growth. In addition, we found that with the lower inlet pressure, the final droplet size reduced more notably. For 329 example, D_d decreased from 10.6 to 10.4 µm (by 0.2 µm) as pressure reduced from 1.0 to 0.9 atm, while D_d decreased 330 from 7.2 to 6.1 µm (by 1.1 µm) as pressure reduced from 0.4 to 0.3 atm. The difference can be explained by the 331 competition between heat and water vapor transport. The mass transport of water vapor is faster than the heating flow 332 from the wall because the mass diffusivity of water vapor is higher than the thermal diffusivity of air. Therefore, by 333 reducing the inlet pressure, water vapor transport becomes even faster than heat transfer due to the water vapor 334 diffusivity and air thermal diffusivity being inversely proportional to the pressure, further shortening the time for 335 particle growth at high supersaturation. This observation demonstrates for the first time how the final droplet size is 336 affected by raising or lowering temperature midpoints at standard and various reduced inlet pressure conditions in the vWCPC. 337

338

339 **3.3 Effect of inlet operation pressure on particle activation and droplet growth performance**

With the advantages of safe, eco-friendly and readily available distilled water as working fluid in the vWCPC, applying the vWCPC in various inlet pressures will expand broader applications such as atmospheric airborne aerosol measurements. Here, we examined the effect of inlet pressure on minimum activated size, $D_{p,kel,0}$, 50% cut-off size, $D_{p,kel,50}$, and final growth particle size at the outlet of the moderator along the centerline (r = 0), D_d from 0.3 to 1 atm for two different temperature settings: the conditioner, initiator, and moderator temperatures were 30, 59, and 10 °C and 27, 59, and 10 °C in Fig. 4.





346

347 Figs. 4a and 4b show $D_{p,kel,0}$ and $D_{p,kel,50}$ as a function of inlet pressure, relatively greater (2 - 3%) $D_{p,kel,0}$ was observed at reduced inlet pressures at both conditioner temperatures of 27 °C and 30 °C. This increase is because the 348 supersaturation value at reduced pressure is lower than the saturation profile under standard conditions. We also found 349 350 that the saturation profile peaked earlier, closer to the entrance of the initiator in the low-pressure condition. In 351 addition, greater $D_{p,kel,50}$ is observed at reduced inlet pressures due to the reduction of the saturation peak at both 352 conditioner temperatures of 27 °C and 30 °C. Again, the difference at reduced inlet pressure can be explained by the 353 competition from heat transfer and water vapor transport, as discussed in Section 3.2. For this reason, greater $D_{p,kel,50}$ 354 was observed at reduced inlet pressures. This reduction of saturation peaks is also associated with the growing droplet 355 size decreasing with the decrease in the operating pressure. Again, lowering the conditioner temperature while 356 maintaining the same temperature difference between the initiator and the moderator provided higher saturation ratios 357 in the initiator over all pressure ranges. 358

Fig. 4c shows the final droplet size as a function of inlet pressure. When the conditioner temperature is 27 or 30 °C, a lower final droplet size (~ 40 % reduction in the droplet size) was observed at a reduced inlet pressure of 0.3 atm, indicating insufficient droplet growth happens at low-pressure conditions, which is consistent with the previous study that insufficient droplet growth becomes more significant under low-pressure operation (Mei et al., 2021).

363

Furthermore, in addition to showing consistent results with the previous study (Mei et al., 2021), our simulations enhance guidance for aircraft applications under extreme conditions, which can be achieved by simulating low atmospheric pressure at 0.3 atm. As shown in Section 3.5, by comparing with experimental results, our simulations can provide more accurate estimates of particle activation and droplet growth to guide vWCPC for low-pressure applications.

369

370 **3.4 Effect of tube diameter and initiator length on particle activation and droplet growth performance**

371 The geometry in CPCs also impacts the CPC activation performance and particle growth due to the changed 372 supersaturation and temperature profile in the tube, as discussed in previous studies (Hao et al., 2021; Hering et al., 373 2014). Here, we examined how the tube diameter D and the length of initiator L_{ini} in the vWCPC may affect the 374 minimum activated size, $D_{p,kel,0}$, 50% cut-off size, $D_{p,kel,50}$, and final growth particle size at the outlet of the moderator along the centerline (r = 0), D_d under default temperature $T_{con} - T_{ini} - T_{mod}$ of 30–59–10 °C, standard inlet pressure 375 376 and reduced pressure of 0.5 atm using the numerical COMSOL model. Again, one needs to note that conditions that 377 can lead to a lower $D_{p,kel,0}$ and $D_{p,kel,50}$ value and larger droplet growth size are favored for improving the performance 378 of the vWCPC. 379

We examined four values of *D* from 4 to 8 mm, and five values of L_{ini} from 10 to 50 mm, shown in **Fig. 5** and **Fig. 6**, respectively. The results indicate that a smaller *D* can slightly decrease in $D_{p,kel,0}$ approximately 0.03 nm, while no





382 noticeable changes on $D_{p,kel,50}$ at the standard pressure (Figs. 5a and 5b). By reducing the tube diameter, the flow 383 speed in the tube increases under the same flow rate, reducing the residence time of the condensed water vapor. This reduction in residence time suppresses homogeneous nucleation in the initiator. Unlike our previous study on CPCs 384 (Hao et al., 2021), the homogeneous nucleation rate is minimal in vWCPC and has no impact on the temperature 385 386 difference compared to butanol-based CPCs. For this reason, this suppressed homogeneous nucleation has limited 387 effects on D_{p,kel,0} and D_{p,kel,50}. However, the increase of the flow speed will significantly limit the time for droplet growth, as will be discussed later. At the reduced pressure of 0.5 atm, a smaller D can slightly decrease in 388 $D_{p,kel,0}$ approximately 0.08 nm, and a slight decrease of approximately 0.03 nm on $D_{p,kel,50}$ (Figs. 5a and 5b). Overall, 389 390 the reduction in pressure plays a more critical role in negatively impacting CPC performance for relatively large tube 391 diameters. Note that buoyancy effects (Roberts and Nenes, 2005) may be critical for large temperature differences if 392 the tube diameter is too large, which is not discussed in the study. 393

394 On the other hand, we found that reducing L_{ini} leads to limited effects on $D_{p,kel,0}$ and $D_{p,kel,50}$, except for the shortest 395 initiator length of 10 mm at the standard pressure (Figs. 6a and 6b). The effect of these relatively long initiator lengths is limited because the degree of supersaturation is determined by the absolute temperature of the tube flow. The 396 temperature difference did not change in the standard pressure and reduced pressure, leaving both $D_{p,kel,0}$ and $D_{p,kel,50}$ 397 398 unchanged. However, at the initiator length of 10 mm, D_{p,kel,0} and D_{p,kel,50} increase significantly due to insufficient 399 water vapor diffusion before passing through the next moderator region, resulting in a lower peak supersaturation 400 along the centerline than for longer initiators operating at the same temperature (Hering et al., 2014). At reduced 401 pressure, $D_{p,kel,0}$ and $D_{p,kel,50}$ have no noticeable changes at all tested initiator lengths, however, this is due to the 402 sufficient diffusion of water vapor, from which the water transport is faster than at the standard pressure. Again, if the initiator is longer, the difference in peak supersaturation will be negligible, while the peak temperature along the 403 404 centerline and the amount of added water vapor will be higher. Thus, for relatively short initiators, such as 20 mm 405 used in the simulation, one can provide all the necessary water vapor to create the same peak supersaturation as for 406 the longer initiators. However, the droplet growth size will be smaller (Fig. 6c), mainly due to the shorter growth time discussed later. 407

408

409 With regard to the performance of particle growth, an increased D and an increased L_{ini} are beneficial for improving 410 the performance of particle growth in vWCPC at both standard and reduced pressure (Figs. 5c and 6c). An increased 411 D implies a decrease in the flow velocity through the high saturation region, greatly increasing the time for particle 412 growth and contributing to the sufficient growth of the particles. Fig. 5c shows that the final droplet sizes increase 413 significantly from 6.72 µm to 13.88 µm when D is increased from 4 mm to 8 mm at the standard pressure and increase 414 from 4.92 to 10.78 μ m at the reduced pressure of 0.5 atm. The final droplet size is found to be 2 – 3 μ m smaller than 415 the standard pressure at the reduced pressure. Similarly, a longer L_{ini} also leads to a larger droplet growth size. The final droplet size increases from 8.66 μ m to 11.26 μ m when L_{ini} is increased from 10 mm to 50 mm at the standard 416 417 pressure and from 7.40 to 8.29 µm at the reduced pressure of 0.5 atm (Fig. 6c). This increase is likely due to the longer 418 growth time of the longer initiator. Also, we found that the final droplet size increases much faster at shorter initiator





- lengths than at lengths above 20 mm, which tells us that the droplet size is more susceptible to the effects of initiator
 length below 20 mm. This difference also means that having a longer length does not further enhance the final size of
- 421 the particle growth.
- 422
- In addition to the performance of particle activation, it is crucial to evaluate the droplet growth performance of complex geometries in the vWCPC. The time that allows the activated particle to grow in the initiator and moderator
- 425 is an important droplet growth kinetics assumption, representing the vWCPC performance of droplet growth. We use
- 426 $t_{\rm g}$ to represent allowed particle growth time, approximated with Eq. (7). $t_{\rm g} \sim D^2 L^* / Q_{\rm w}$

- (7)
- 427 where L^* indicates the length of the initiator and moderator beyond the point of activation. This equation can explain
- 428 that the residence time is impacted more by the change in tube diameter than the initiator length. The allowed particle
- 429 growth time as a function of final growth particle size at the outlet of the moderator along the centerline (r = 0), D_d is
- 430 shown in Fig. S3. The longer the allowed particle growth time, the larger the droplet growth size. Based on this droplet
- 431 growth time shown, the vWCPC geometry of D and L_{ini} are not independent variables if we consider the droplet
- 432 growth for further particle detection.
- 433

434 **3.5 Experimental measurement validation of detection efficiency**

- Experimental validation is essential for simulation work in terms of the accuracy of the simulation model and the correctness of the underlying trends. Furthermore, validation and good agreement will provide well-guided approaches for future applications. Therefore, we compare the experimental and simulation results of the counting efficiency and detection efficiency of vWCPC for two configurations of 2 nm and 7 nm at different conditioner and initiator temperature settings and different low-pressure conditions in **Fig. 7**.
- 440

As the experimental results in a previous study (Mei et al., 2021) are shown in **Fig. 7a**, the counting efficiency of vWCPC 3789 varied with different working pressures (500, 700 and 910 hPa) when the conditioner temperature is 27 °C, the initiator temperature is 59 °C, and the moderator temperature is 10 °C. The results indicate that the counting efficiency slightly decreases with the decrease in the operating pressure of 500, 700 and 900 hPa, which shows the same trend in **Fig. 7b.** In addition, the cut-off size in both experimental and simulation results are in the range of 5 -7 nm, which is also an acceptable range within error when compared to commercial vWCPC detection efficiency.

- Fig. 7c and 7d compare the counting efficiency and detection efficiency versus particle size from experimental and simulation results under initiator temperatures of 75 and 90°C and pressure of 910 and 500 hPa for the 2 nm
- 450 configuration. As expected, the detection efficiency of both experimental and simulated results is lower at the
- 451 temperature $T_{con} T_{ini} T_{mod}$ of 7-75-10 °C at a lower pressure (at 500 hPa). When the temperature $T_{con} T_{ini} T_{mod}$
- 452 is 7-90-10 °C, the higher detection efficiency is seen, and the effect of inlet pressure becomes insignificant. However,
- 453 it is not feasible to maintain 90 °C when operating under lower pressure, such as 500 hPa. Thus, the default 2 nm





454 setting in vWCPC can only be operated near sea level. One should note that although we do not present many 455 simulations for the 2 nm configuration, what we learned from the modeling results with the 7 nm setting will guide

- 456 future simulations with the 2 nm setting.
- 457

458 By comparing with counting efficiency curves, the present simulations can more realistically represent the D_{p,kel} for

- 459 7 nm vWCPC, which also achieved a good agreement with the 2 nm setting. Thus, from the merits of the results of 460 this work, we can find that this work not only provides guidance for 7 nm, but this trend can also help guide one for
- 461 other desired cut-off sizes.
- 462

463 4 Conclusions

464 This study evaluated the particle activation and droplet growth performance of a commercial versatile water CPC 465 using COMSOL in combination with MATLAB data processing. In addition, validation experiments on the detection efficiency of the commercially modified vWCPC (TSI 3789) agreed with the simulation work. Increasing the 466 temperature difference between T_{con} and T_{ini} and lowering the temperature midpoint can enhance particle activation 467 468 at both standard and reduced ambient pressure conditions. However, the lack of droplet growth becomes more significant under low-pressure operations, which might affect the apparent counting efficiency of the vWCPC due to 469 the limited measurable size range of the optical chamber. Additionally, reducing the diameter of the growth tube 470 slightly improved particle activation without enhancing the droplet growth, while increasing the initiator length had a 471 472 limited effect on improving the performance of the vWCPC at both standard and reduced pressure.

473

This simulation realistically represents the $D_{p,kel}$ for 7 nm vWCPC and shows that the current growth tube geometry is an optimized choice for aerosol measurements. This study will guide further vWCPC performance optimization for applications requiring precise particle detection and atmospheric aerosol monitoring. Furthermore, the developed simulation capability provides a vital tool for the aerosol community to understand the effects of temperature and pressure on vWCPC behavior. The knowledge gained will guide the field deployment of vWCPC on the ground level and airborne measurements. Thus, several future experimental studies will be carried out to investigate the performance of the vWCPC.

481

482 Data availability

483 The vWCPC data in the study are available upon request to Fan Mei (fan.mei@pnnl.gov).

484

485 Author contributions

486 WH, FM, and YW designed the research. FM carried out the measurements. WH led the simulation and data analyses.

487 WH led the writing, with significant input from FM and YW as well as further input from all other authors. SH, SS,

- 488 BS, and JT provided suggestions on the revision.
- 489





490	Competing	interests
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- 491 Susanne Hering has a commercial interest in the success of the vWCPC instrument.
- 492

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- 495

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580 Table of Nomenclature

- *c*: Molecular concentration of the water vapor [mol m⁻³]
- C_d : Equilibrium water concentration at the surface of the growth droplets [mol m⁻³]
- $C_{\rm s}$: Saturation water concentration [mol m⁻³]
- c_p : Heat capacity of the water [J K⁻¹ kg⁻¹]
- $c_{p,g}$: Heat capacity of air [J K⁻¹ kg⁻¹]
- $D_{\rm v}$: Diffusivity of the water vapor [m² s⁻¹]
- D'_{v} : Modified diffusivity of the water vapor [m² s⁻¹]
- 588 D: Diameter of the growth tube in vWCPC [m]
- $D_{\rm p}$: Particle size [m]
- D_d : Final growth droplet size in vWCPC [m]
- 591 D_{p,kel}: Size of particle that can be activated according to the Kelvin equation [m]
- $D_{p,kel,0}$: Smallest size of particle that can be activated in the vWCPC [m]
- $D_{p,kel,50}$: Size of particle that has a 50% activation efficiency [m]
- H_{vap} : Heat of vaporization of water [J kg⁻¹]
- 595 k: Boltzmann constant, 1.38×10^{-23} [J K⁻¹]
- k_{g} : Thermal conductivity of air [W m⁻¹ K⁻¹]
- k'_{g} : Modified thermal conductivity of air [W m⁻¹ K⁻¹]
- L_{con} : Length of the conditioner [m]
- L_{ini} : Length of the initiator [m]
- $600 \quad L_{mod}$: Length of the moderator [m]
- 601 L*: Length of the initiator and moderator beyond the point of activation [m]
- *m*: Molecular mass of water [kg]
- M: Molecular weight of water [kg mol⁻¹]
- $M_{\rm g}$: Molecular weight of air [kg mol⁻¹]
- *n*: Molecular concentration of the water vapor [molecules m^{-3}]
- 606 N: Concentration of the particles at the axial location of $z = Z_{act}$ [particles m⁻³]
- N_0 : Concentration of particles at the inlet of the conditioner [particles m⁻³]
- 608 p: Partial pressure of the water vapor [Pa]
- p_s : Saturation vapor pressure of the water vapor [Pa]
- *P*: Inlet pressure in vWCPC [Pa]
- $Q_{\rm v}$: Flow rate through the vWCPC [m³ s⁻¹]
- 612 r: Radial coordinate of the tube diameter of the vWCPC [m or as otherwise explicitly designated]
- R: Gas constant [J mol⁻¹ K⁻¹]
- R_{act} : Maximum radius of the contour corresponding to $D_{p,kel} = D_p [m]$
- 615 S: Saturation ratio [1]





- 616 t: Allowed particle growth time [s]
- 617 T: Flow temperature in the CPC [K]
- *T*_{con}: Conditioner temperature [K]
- *T*_{ini}: Initiator temperature [K]
- T_{mid} : Temperature midpoint corresponding to $\frac{T_{\text{con}}+T_{\text{ini}}}{2}$ [K]
- *T*_{mod}: Moderator temperature [K]
- T_d : Droplet surface temperature [K]
- $v_{\rm m}$: Molecular volume of the water vapor [m³]
- 624 w: Velocity along the axial direction in the vWCPC [m s⁻¹]
- 625 z: Axial coordinate of the tube length of the vWCPC [m or as otherwise explicitly designated]
- Z_{act} : Axial location corresponding to $r = R_{act}$ [m]
- α_c : Mass accommodation coefficient of water [1]
- $\alpha_{\rm T}$: Thermal accommodation coefficient of air [1]
- η_{act} : Activation efficiency [1]
- ρ : Density of water [kg m⁻³]
- ρ_{g} : Density of air [kg m⁻³]
- σ : Surface tension of water [N m⁻¹]
- ΔT : Temperature difference between conditioner temperature and initiator temperature [K]





Table 1. Parameters of vWCPC for different simulation tasks. Note that the default settings of VWCPC are: the conditioner for the initiator temperature (T_{-1}) is 30 °C the initiator temperature (T_{-1}) is 50 °C and the moderator temperature (T_{-1}) is 10 °C. The

temperature (T_{con}) is 30 °C, the initiator temperature (T_{ini}) is 59 °C, and the moderator temperature (T_{mod}) is 10 °C. The aerosol flow rate (Q_v) is 0.3 L min⁻¹. The relative humidity (RH) of inlet flow is set at 20%, and the water vapor is assumed to be saturated at the wall. The inlet pressure (P) is 1 atm.

$T_{\rm con}$ (°C) - $T_{\rm ini}$ (°C)	$T_{\rm mod}$ (°C)	$T_{\rm mid}$ (°C)	P (atm)	D (mm)	$L_{\rm ini}~(\rm mm)$
(25, 30, 35) - (55, 60, 65)	10	-	1	6.3	30
24 - 56, 27 - 59, 30 - 62	10	40, 43, 46	0.3 - 1	6.3	30
27 - 59, 30 - 59	10	-	0.3 - 1	6.3	30
30 - 59	10	-	0.5, 1	4, 5, 6.3, 8	30
30 - 59	10	-	0.5, 1	6.3	10, 20, 30, 40, 50
	<i>T</i> _{con} (°C) - <i>T</i> _{ini} (°C) (25, 30, 35) - (55, 60, 65) 24 - 56, 27 - 59, 30 - 62 27 - 59, 30 - 59 30 - 59 30 - 59	$T_{\rm con}$ (°C) - $T_{\rm ini}$ (°C) $T_{\rm mod}$ (°C) (25, 30, 35) - (55, 60, 65) 10 24 - 56, 27 - 59, 30 - 62 10 27 - 59, 30 - 59 10 30 - 59 10 30 - 59 10	$T_{\rm con}$ (°C) - $T_{\rm ini}$ (°C) $T_{\rm mod}$ (°C) $T_{\rm mid}$ (°C)(25, 30, 35) - (55, 60, 65)10-24 - 56, 27 - 59, 30 - 621040, 43, 4627 - 59, 30 - 5910-30 - 5910-30 - 5910-	$T_{\rm con}$ (°C) - $T_{\rm ini}$ (°C) $T_{\rm mod}$ (°C) $T_{\rm mid}$ (°C) P (atm)(25, 30, 35) - (55, 60, 65)10-124 - 56, 27 - 59, 30 - 621040, 43, 460.3 - 127 - 59, 30 - 5910-0.3 - 130 - 5910-0.5, 130 - 5910-0.5, 1	$T_{\rm con}$ (°C) - $T_{\rm ini}$ (°C) $T_{\rm mod}$ (°C) $T_{\rm mid}$ (°C) P (atm) D (mm)(25, 30, 35) - (55, 60, 65)10-16.324 - 56, 27 - 59, 30 - 621040, 43, 460.3 - 16.327 - 59, 30 - 5910-0.3 - 16.330 - 5910-0.5, 14, 5, 6.3, 830 - 5910-0.5, 16.3







641 Figure 1. Geometry of vWCPC used in COMSOL simulation and spatial distribution of saturation ratio and Kelvin 642 equivalent size under 30-59-10 °C temperature setting. (a) Geometry of the vWCPC used in COMSOL simulation, (b) 643 Spatial distribution of saturation ratio (S, color contour plot), (c) Spatial distribution of Kelvin equivalent size ($D_{p,kel}$, color 644 contour plot), (d) Spatial distribution of water vapor concentration (c, color contour plot), and (e) Spatial distribution of 645 temperature (T, color contour plot). Note that the color of $D_{p,kel}$ in the conditioner region is blank because no particles are 646 activated in this region.











649 Figure 2. Effect of conditioner (T_{con}) and initiator temperature (T_{ini}) on (a) minimum activated size, $D_{p,kel,0}$, (b) 50%

650 cut-off size, $D_{p,kel,50}$, and (c) final growth particle size at the outlet of the moderator along the centerline (r = 0), D_d .

651 The condensational growth of 15 nm particles was tested as seed particles.





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653 Figure 3. Effect of temperature midpoints at 40 °C, 43 °C, and 46 °C at $T_{\rm con}$ - $T_{\rm ini}$ - $T_{\rm mod}$ of 24–56–10 °C, 27–59–10 °C

and 30–62–10 °C with a constant temperature difference of 32 °C on (a) minimum activated size, $D_{p,kel,0}$, (b) 50% cutoff size, $D_{p,kel,50}$, and (c) final growth particle size at the outlet of the moderator along the centerline (r = 0), D_d . The

656 condensational growth of 8 nm particles was tested as seed particles.







Figure 4. Effect of inlet operation pressure at T_{con} - T_{ini} - T_{mod} of 27–59–10 °C and 30–59–10 °C on (a) minimum activated size, $D_{p,kel,0}$, (b) 50% cut-off size, $D_{p,kel,50}$, and (c) final growth particle size at the outlet of the moderator along the centerline (r = 0), D_d . The condensational growth of 8 nm particles was tested as seed particles.





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663 $D_{p,kel,50}$, and (c) final growth particle size at the outlet of the moderator along the centerline (r = 0), D_d . The

664 condensational growth of 8 nm particles was tested as seed particles.









Figure 6. Effect of initiator length (L_{ini}) at 0.5 atm and 1 atm on (a) minimum activated size, $D_{p,kel,0}$, (b) 50% cut-off

667 size, $D_{p,kel,50}$, and (c) final growth particle size at the outlet of the moderator along the centerline (r = 0), D_d . The 668 condensational growth of 8 nm particles was tested as seed particles.







Figure 7. vWCPC operation validation: (a) the counting efficiency of experimental results as a function of particle size under the conditioner temperature of 27 °C and pressure of 910, 700, and 500 hPa for the 7 nm configuration, (b) the detection efficiency of simulation results as a function of particle size under the conditioner temperature of 27 °C and pressure of 0.9, 0.7, and 0.5 atm for the 7 nm configuration, (c) the detection efficiency of experimental results as a function of particle size under initiator temperatures of 75 and 90°C and pressure of 910 and 500 hPa for the 2 nm configuration, and (d) the detection efficiency of simulation results as a function of particle size under initiator temperatures of 75 and 90°C and pressure of 0.9 and 0.5 atm for the 2 nm configuration.