Measurement-device-independent Randomness from local entangled states

ISCQI-2016 Manik Banik

Post Doctoral Fellow
Optics & Quantum Information Group
The Institute of Mathematical Sciences

1 / 30

Colaborator..



Figure: Anubhav Chaturvedi

Colaborator...



Figure: Anubhav Chaturvedi

Link: Euro. Phys. Lett. 112, 30003 (2015)

(doi: 10.1209/0295-5075/112/30003)

Manik Banik (PDF)

Random Numbers...

- ★ Random numbers have many practical uses in modern science.
 - Cryptography
 - Statistical research
 - Numerical Simulation (eg. Monte Carlo method)
 - Lotteries and gambling
 - PIN number generation
 - Mobile prepaid systems

★ Consider a sequence of two numbers '0' and '1':

with
$$p(0) = p(1) = \frac{1}{2}$$

★ Consider a sequence of two numbers '0' and '1':

with
$$p(0) = p(1) = \frac{1}{2}$$

★ Consider another sequence:

010011101001010011101001010011101001....

with
$$p(0) = p(1) = \frac{1}{2}$$

★ Consider a sequence of two numbers '0' and '1':

with
$$p(0) = p(1) = \frac{1}{2}$$

★ Consider another sequence:

010011101001010011101001010011101001....

with
$$p(0) = p(1) = \frac{1}{2}$$

 $\{0100111010010\}\{0100111010010\}\{0100111010010\}.....$

Certification: Mathematically impossible...

★ Using algorithmic information theory it can be shown that true randomness can not exist from a mathematical point of view. {Chaitin G. J., IBM J. Res. Dev., 21 (1977) 350; Knuth D., The Art of Computer Programming, Semi-numerical Algorithms}



Certification: Mathematically impossible...

★ Using algorithmic information theory it can be shown that true randomness can not exist from a mathematical point of view. {Chaitin G. J., IBM J. Res. Dev., 21 (1977) 350; Knuth D., The Art of Computer Programming, Semi-numerical Algorithms}



★ Any one who considers arithmetical methods of producing random digits is, of course, in a state of sin.

——John von Neumann

★ Therefore generation of randomness must rely on unpredictability of physical phenomena, like

★ Coin tossing...



6 / 30

★ Dice rolling...



- ★ However all classical processes ('coin tossing'/'dice rolling') are deterministic from fundamental point of view.
- ★ This is because, fate of any classical object at any time is completely predictable in Newtonian dynamics.
- ★ Therefore no classical process can be a source of "true" randomness.



- ★ However all classical processes ('coin tossing'/'dice rolling') are deterministic from fundamental point of view.
- ★ This is because, fate of any classical object at any time is completely predictable in Newtonian dynamics.
- ★ Therefore no classical process can be a source of "true" randomness.



★ So we shift our attention from classical world (CW) to quantum world (QW).

Certification: Possible in QW (?)...

★ Quantum theory:

- ullet State of a system: Vectors, $|\psi
 angle\in\mathcal{H}$
- Observables: Hermitian operator, $\mathcal{A} \in \mathcal{B}(\mathcal{H})$, acting on \mathcal{H} $\mathcal{A} = \sum_i a_i |\alpha_i\rangle\langle\alpha_i|$
- Possible measurement results: Eigenvalues of the given observable
- Outcome probability: $p(a_1) = |\langle \psi | \alpha_i \rangle|^2$ (Born rule)
- ★ Due to the Born's rule, in QW we can obtain our desired randomness.

Certification: Possible in QW (?) $\{an\ example\}...$

- ★ Consider a spin-1/2 system ($\mathcal{H} \equiv \mathbb{C}^2$)
 - State: 'up' eigenstate $(|\uparrow\rangle)$ of the Pauli σ_z
 - Measurement:Pauli σ_x observable, whose eigenstate are denoted as $|\rightleftarrows\rangle=\frac{1}{\sqrt{2}}(|\uparrow\rangle\pm|\downarrow\rangle)$.
 - Outcome probability: $p(| \rightarrow \rangle) = \frac{1}{2}$ and also $p(| \leftarrow \rangle) = \frac{1}{2}$

Manik Banik (PDF)

Certification: Possible in QW (?) {an example}...

- ★ Consider a spin-1/2 system ($\mathcal{H} \equiv \mathbb{C}^2$)
 - State: 'up' eigenstate $(|\uparrow\rangle)$ of the Pauli σ_z
 - Measurement:Pauli σ_x observable, whose eigenstate are denoted as $|\rightleftarrows\rangle=\frac{1}{\sqrt{2}}(|\uparrow\rangle\pm|\downarrow\rangle)$.
 - Outcome probability: $p(| \rightarrow \rangle) = \frac{1}{2}$ and also $p(| \leftarrow \rangle) = \frac{1}{2}$

★ If we associate '0' ('1') with $|\uparrow\rangle$ ($|\downarrow\rangle$), then we obtain a sequence of '0' and '1' with $p(0)=p(1)=\frac{1}{2}$ and there will be no pattern in the sequence \Longrightarrow Randomness



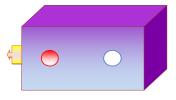
QRNG...

★ Such QRNG already exists:



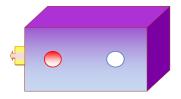
DI Certification...

★ Consider that we are ignored about the internal working of the device, i.e., the device is like a black box with just input and output.



DI Certification...

★ Consider that we are ignored about the internal working of the device, i.e., the device is like a black box with just input and output.



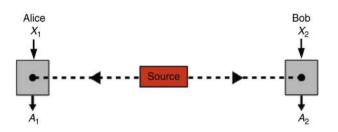
★ In this situation, is it still possible to be certain that the device is producing the outcomes without following any particular pattern?



DI Certification...

★ In a recent result it has been shown that DI randomness certification is possible.

★ "Random Numbers Certified by Bell's Theorem", S. Pironio, A. Acin, S. Massar, A. Boyer de la Giroday, D. N. Matsukevich, P. Maunz, S. Olmschenk, D. Hayes, L. Luo, T. A. Manning, C. Monroe, Nature 464, 1021 (2010)



★ Local Realism (LR):

$$P(A_1, A_2|X_1, X_2) = \sum_{\lambda \in \Lambda} \rho(\lambda) P(A_1|X_1, \lambda) P(A_2|X_2, \lambda)$$

- ★ Bell inequality: $|\langle X_1X_2\rangle + \langle X_1X_2'\rangle + \langle X_1'X_2\rangle \langle X_1'X_2'\rangle| \le 2$
- ★ Bell inequality can also be derived under two operational assumptions, namely 'predictability' and 'signal-locality'



14 / 30

Manik Banik (PDF) IMSc, Chennai 17th February, 2016

★ Determinism
$$\land$$
 Locality \Rightarrow factorizability \Rightarrow Bell's inequality (BI) i.e., $P(A_1, A_2 | X_1, X_2, \psi) = \int_{\lambda \in \Lambda} \mu(\lambda | \psi) P(A_1 | X_1, \psi, \lambda) P(A_2 | X_2, \psi, \lambda) d\lambda$

- **Determinism** (D) $\Rightarrow P(A_1, A_2 | X_1, X_2, \psi, \lambda) \in \{0, 1\}$
- Locality (L) \Rightarrow $P(A_1|X_1, X_2, \psi, \lambda) = P(A_1|X_1, \psi, \lambda)$ $P(A_2|X_1, X_2, \psi, \lambda) = P(A_2|X_2, \psi, \lambda)$
- Proof:

$$P(A_1, A_2|X_1, X_2, \psi, \lambda) = P(A_1|A_2, X_1, X_2, \psi, \lambda)P(A_2|X_1, X_2, \psi, \lambda)$$

$$= P(A_1|X_1, X_2, \psi, \lambda)P(A_2|X_1, X_2, \psi, \lambda); [D]$$

$$= P(A_1|X_1, \psi, \lambda)P(A_2|X_2, \psi, \lambda); [L]$$



15 / 30

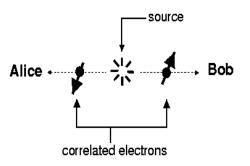
- ★ Predictability \wedge Signal Locality \Rightarrow factorizability \Rightarrow BI i.e., $P(A_1, A_2 | X_1, X_2, \psi) = \int_{\lambda \in \Lambda} \mu(\lambda | \psi) P(A_1 | X_1, \psi, \lambda) P(A_2 | X_2, \psi, \lambda) d\lambda$
 - **Predictability** (P) $\Rightarrow P(A_1, A_2 | X_1, X_2, \psi) \in \{0, 1\}$
 - Signal Locality (SL) $\Rightarrow P(A_1|X_1, X_2, \psi) = P(A_1|X_1, \psi) P(A_2|X_1, X_2, \psi) = P(A_2|X_2, \psi)$

Manik Banik (PDF)

- ★ Predictability \wedge Signal Locality \Rightarrow factorizability \Rightarrow BI i.e., $P(A_1, A_2 | X_1, X_2, \psi) = \int_{\lambda \in \Lambda} \mu(\lambda | \psi) P(A_1 | X_1, \psi, \lambda) P(A_2 | X_2, \psi, \lambda) d\lambda$
 - **Predictability (P)** \Rightarrow $P(A_1, A_2 | X_1, X_2, \psi) \in \{0, 1\}$
 - Signal Locality (SL) $\Rightarrow P(A_1|X_1, X_2, \psi) = P(A_1|X_1, \psi) P(A_2|X_1, X_2, \psi) = P(A_2|X_2, \psi)$

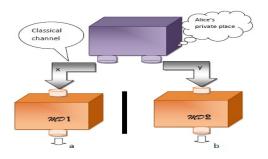
Manik Banik (PDF)

★ Quantum correlation violates BI:



★ Using nonlocal correlation DI randomness certification is possible.

17 / 30



***** Randomness associated with $\{P(ab|xy)\}$, quantified as $H_{\infty} = -\log_2 \max_{a,b} P(ab|xy)$.

★ Which physical correlation shows this nonlocal properties?

★ Which physical correlation shows this nonlocal properties?

★ Entanglement:

- Bipartite quantum system $\to \mathcal{H}_A \otimes \mathcal{H}_B$
- Product state: $|\psi\rangle_A\otimes|\psi\rangle_B\in\mathcal{H}_A\otimes\mathcal{H}_B$
- Non product States are called entangled:

$$|\psi\rangle_{AB} = \sum_{i} c_{i} |\psi^{i}\rangle_{A} \otimes |\psi^{i}\rangle_{B}$$

• Separable states: $\rho_{AB} = \sum_{i} p_{i} \sigma_{A}^{i} \otimes \sigma_{B}^{i}$; with

$$p_i \geq 0 \& \sum \liminf_i p_i = 1$$

★ Example: Werner class of states

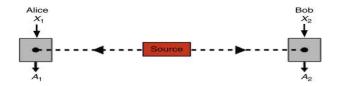
$$W_{
ho}=
ho|\psi^{-}
angle\langle\psi^{-}|+(1-
ho)rac{\mathbb{I}}{2}\otimesrac{\mathbb{I}}{2}$$

with
$$0 \le p \le 1$$
 and $|\psi^-\rangle = \frac{1}{\sqrt{2}} (|\uparrow\rangle|\downarrow\rangle - |\downarrow\rangle|\uparrow\rangle)$

- $p > \frac{1}{\sqrt{2}}$: Violates BI (useful for DI certification)
- $p \le \frac{1}{2}$: LHV for PV
- $p \leq \frac{5}{12}$: LHV for POVM
- $p \ge \frac{1}{3}$: Entangled

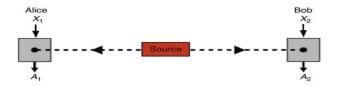
★ Local entangled states are not useful for DI randomness certification

★ Semi-quantum game (F. Buschemi):



- Instead of classical inputs, quantum states $\{|\phi^x\rangle_{\alpha'}\}_{x\in X}$ and $\{|\psi^y\rangle_{\beta'}\}_{y\in Y}$, chosen from Hilbert spaces $\mathcal{H}_{\alpha'}$ and $\mathcal{H}_{\beta'}$, respectively, are sent
- For every entangled state quantum inputs can be chose in such a way that the produced correlation cannot be achieved by local operation and shared randomness

★ Semi-quantum game (F. Buschemi):



- Instead of classical inputs, quantum states $\{|\phi^x\rangle_{\alpha'}\}_{x\in X}$ and $\{|\psi^y\rangle_{\beta'}\}_{y\in Y}$, chosen from Hilbert spaces $\mathcal{H}_{\alpha'}$ and $\mathcal{H}_{\beta'}$, respectively, are sent
- For every entangled state quantum inputs can be chosen in such a way that the produced correlation cannot be achieved by local operation and shared randomness

- We consider the following particular semi-quantum game:
- The input quantum states are chosen from a regular tetrahedron on the Bloch sphere i.e.,

$$|\phi^{\mathsf{x}}\rangle\langle\phi^{\mathsf{x}}| = \frac{\mathbb{I} + \vec{\mathsf{v}}_{\mathsf{x}}.\vec{\sigma}}{2}, \quad |\psi^{\mathsf{y}}\rangle\langle\psi^{\mathsf{y}}| = \frac{\mathbb{I} + \vec{\mathsf{v}}_{\mathsf{y}}.\vec{\sigma}}{2},$$

for x,y=1,...,4 we have $\vec{v}_1=\frac{(1,1,1)}{\sqrt{3}}$, $\vec{v}_2=\frac{(1,-1,-1)}{\sqrt{3}}$, $\vec{v}_3=\frac{(-1,1,-1)}{\sqrt{3}}$ and $\vec{v}_4=\frac{(1,-1,-1)}{\sqrt{3}}$; and $\vec{\sigma}=(\sigma_1,\sigma_2,\sigma_3)$ with σ_i (i=1,2,3) being the Pauli matrices

 \bullet The POVM $\{\mathcal{M}_{\mathsf{a}}^{\alpha'\alpha}\}_{\mathsf{a}\in\{0,1\}}$ is given by

$$\mathcal{M}_{1}^{\alpha'\alpha}=|\phi^{+}\rangle\langle\phi^{+}|,\quad \mathcal{M}_{0}^{\alpha'\alpha}=\mathbb{I}-|\phi^{+}\rangle\langle\phi^{+}|,$$

23 / 30

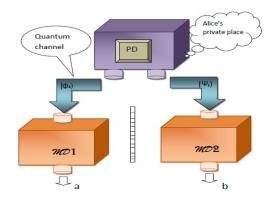
 Using the above quantum-input classical-output statistics one can construct the following MDI-entanglement witness (Branciard et al.):

$$I(P) = \frac{5}{8} \sum_{x=y} p(1,1||\phi^x\rangle, |\psi^y\rangle) - \frac{1}{8} \sum_{x\neq y} p(1,1||\phi^x\rangle, |\psi^y\rangle).$$

Here P denotes the probability distribution $\{p(a,b||\phi^x\rangle,|\psi^y\rangle)|a,b=0,1;x,y=1,..,4\}.$

• For the Werner states the above expression becomes $I(P_{\varrho^{\nu}})=\frac{1-3\nu}{16}$, which is negative for $\nu>\frac{1}{3}$. For any separable state ρ , $I(P_{\rho})=0$

★ Measurement-DI (MDI) randomness certification:



★ MDI min-entropic source

To find the minimum randomness associated with the probability distribution $P = \{p(ab|xy)\}$ one has to solve the following optimization problem,

$$p^*(ab|xy) = \max_{p(ab|xy)} p(ab|xy)$$

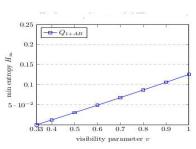
subject to $I(P) = \frac{1-3v}{16}$
 $p(ab|xy) \in Q$,

- where $I(P) = \frac{5}{8} \sum_{x=y} p(1,1||\phi^x\rangle,|\psi^y\rangle) \frac{1}{8} \sum_{x\neq y} p(1,1||\phi^x\rangle,|\psi^y\rangle)$
- the minimum random bits is therefore $H_{\infty}(AB|XY) = -\log_2 \max_{ab} p_q^*(ab|xy)$.
- While the optimization problem is computationally tough, we solve for a relaxed condition $p(ab|xy) \in Q_{1+AB}$ using SDP.

4 D > 4 P > 4 B > 4 B >

★ MDI min-entropic source

- we find zero min-entropy against Q_{1+AB}
- so we put further conditions on the observed statistics: $P(0,0|I,I) = P(0,0|m,m); \ P(0,0|I,I) = P(0,0|m,m); \ \forall \ I,m \in \{1,2,3,4\}$
- interestingly positive min entropy is obtained for I(P) < 0



★ MDI min-entropic source

- two qubits entangled Werner class of states also satisfied the required conditions and hence are useful for MDI min-entropy (randomness) certification
- no seperable state satisfies this condition I(P) < 0 hence no cheating strategy is possible by sharing seperable correlations
- correlations in semi-quantum scenarion cannot even be simulated by local operation and classical correlation (LOCC) (Rosset et al.)

28 / 30

★ Summary of the talk:

- Randomness is a useful resource
- Randomness certification is not possible mathematically and also in classical world
- In quantum world randomness generation is possible
- Bell's theorem: DI certification possible:

 nonlocal entangled states are useful
- local entangled states are useful for MDI randomness certification

