

Measurement of the Top-Quark Mass in All-Hadronic Decays in $p\bar{p}$ Collisions at CDF II

T. Aaltonen,²³ A. Abulencia,²⁴ J. Adelman,¹³ T. Affolder,¹⁰ T. Akimoto,⁵⁶ M.G. Albrow,¹⁷ D. Ambrose,¹⁷ S. Amerio,⁴⁴ D. Amidei,³⁵ A. Anastassov,⁵³ K. Anikeev,¹⁷ A. Annovi,¹⁹ J. Antos,¹⁴ M. Aoki,⁵⁶ G. Apollinari,¹⁷ J.-F. Arguin,³⁴ T. Arisawa,⁵⁸ A. Artikov,¹⁵ W. Ashmanskas,¹⁷ A. Attal,⁸ F. Azfar,⁴³ P. Azzi-Bacchetta,⁴⁴ P. Azzurri,⁴⁷ N. Bacchetta,⁴⁴ W. Badgett,¹⁷ A. Barbaro-Galtieri,²⁹ V.E. Barnes,⁴⁹ B.A. Barnett,²⁵ S. Baroiant,⁷ V. Bartsch,³¹ G. Bauer,³³ F. Bedeschi,⁴⁷ S. Behari,²⁵ S. Belforte,⁵⁵ G. Bellettini,⁴⁷ J. Bellinger,⁶⁰ A. Belloni,³³ D. Benjamin,¹⁶ A. Beretvas,¹⁷ J. Beringer,²⁹ T. Berry,³⁰ A. Bhatti,⁵¹ M. Binkley,¹⁷ D. Bisello,⁴⁴ R.E. Blair,² C. Blocker,⁶ B. Blumenfeld,²⁵ A. Bocci,¹⁶ A. Bodek,⁵⁰ V. Boisvert,⁵⁰ G. Bolla,⁴⁹ A. Bolshov,³³ D. Bortoletto,⁴⁹ J. Boudreau,⁴⁸ A. Boveia,¹⁰ B. Brau,¹⁰ L. Brigliadori,⁵ C. Bromberg,³⁶ E. Brubaker,¹³ J. Budagov,¹⁵ H.S. Budd,⁵⁰ S. Budd,²⁴ S. Budroni,⁴⁷ K. Burkett,¹⁷ G. Busetto,⁴⁴ P. Bussey,²¹ K. L. Byrum,² S. Cabrera,¹⁶ M. Campanelli,²⁰ M. Campbell,³⁵ F. Canelli,¹⁷ A. Canepa,⁴⁹ S. Carillo,¹⁸ D. Carlsmith,⁶⁰ R. Carosi,⁴⁷ M. Casarsa,⁵⁵ A. Castro,⁵ P. Catastini,⁴⁷ D. Cauz,⁵⁵ M. Cavalli-Sforza,³ A. Cerri,²⁹ L. Cerrito,⁴³ S.H. Chang,²⁸ Y.C. Chen,¹ M. Chertok,⁷ G. Chiarelli,⁴⁷ G. Chlachidze,¹⁵ F. Chlebana,¹⁷ I. Cho,²⁸ K. Cho,²⁸ D. Chokheli,¹⁵ J.P. Chou,²² G. Choudalakis,³³ S.H. Chuang,⁶⁰ K. Chung,¹² W.H. Chung,⁶⁰ Y.S. Chung,⁵⁰ M. Ciljak,⁴⁷ C.I. Ciobanu,²⁴ M.A. Ciocci,⁴⁷ A. Clark,²⁰ D. Clark,⁶ M. Coca,¹⁶ G. Compostella,⁴⁴ M.E. Convery,⁵¹ J. Conway,⁷ B. Cooper,³⁶ K. Copic,³⁵ M. Cordelli,¹⁹ G. Cortiana,⁴⁴ F. Crescioli,⁴⁷ C. Cuenca Almenar,⁷ J. Cuevas,¹¹ R. Culbertson,¹⁷ J.C. Cully,³⁵ D. Cyr,⁶⁰ S. DaRonco,⁴⁴ M. Datta,¹⁷ S. D'Auria,²¹ T. Davies,²¹ M. D'Onofrio,³ D. Dagenhart,⁶ P. de Barbaro,⁵⁰ S. De Cecco,⁵² A. Deisher,²⁹ G. De Lentdecker,⁵⁰ M. Dell'Orso,⁴⁷ F. Delli Paoli,⁴⁴ L. Demortier,⁵¹ J. Deng,¹⁶ M. Deninno,⁵ D. De Pedis,⁵² P.F. Derwent,¹⁷ G.P. Di Giovanni,⁴⁵ C. Dionisi,⁵² B. Di Ruzza,⁵⁵ J.R. Dittmann,⁴ P. DiTuro,⁵³ C. Dörr,²⁶ S. Donati,⁴⁷ M. Donega,²⁰ P. Dong,⁸ J. Donini,⁴⁴ T. Dorigo,⁴⁴ S. Dube,⁵³ J. Efron,⁴⁰ R. Erbacher,⁷ D. Errede,²⁴ S. Errede,²⁴ R. Eusebi,¹⁷ H.C. Fang,²⁹ S. Farrington,³⁰ I. Fedorko,⁴⁷ W.T. Fedorko,¹³ R.G. Feild,⁶¹ M. Feindt,²⁶ J.P. Fernandez,³² R. Field,¹⁸ G. Flanagan,⁴⁹ A. Foland,²² S. Forrester,⁷ G.W. Foster,¹⁷ M. Franklin,²² J.C. Freeman,²⁹ I. Furic,¹³ M. Gallinaro,⁵¹ J. Galyardt,¹² J.E. Garcia,⁴⁷ F. Garbersson,¹⁰ A.F. Garfinkel,⁴⁹ C. Gay,⁶¹ H. Gerberich,²⁴ D. Gerdes,³⁵ S. Giagu,⁵² P. Giannetti,⁴⁷ A. Gibson,²⁹ K. Gibson,⁴⁸ J.L. Gimmell,⁵⁰ C. Ginsburg,¹⁷ N. Giokaris,¹⁵ M. Giordani,⁵⁵ P. Giromini,¹⁹ M. Giunta,⁴⁷ G. Giurgiu,¹² V. Glagolev,¹⁵ D. Glenzinski,¹⁷ M. Gold,³⁸ N. Goldschmidt,¹⁸ J. Goldstein,⁴³ A. Golossanov,¹⁷ G. Gomez,¹¹ G. Gomez-Ceballos,¹¹ M. Goncharov,⁵⁴ O. González,³² I. Gorelov,³⁸ A.T. Goshaw,¹⁶ K. Goulianos,⁵¹ A. Gresele,⁴⁴ M. Griffiths,³⁰ S. Grinstein,²² C. Grosso-Pilcher,¹³ R.C. Group,¹⁸ U. Grundler,²⁴ J. Guimaraes da Costa,²² Z. Gunay-Unalan,³⁶ C. Haber,²⁹ K. Hahn,³³ S.R. Hahn,¹⁷ E. Halkiadakis,⁵³ A. Hamilton,³⁴ B.-Y. Han,⁵⁰ J.Y. Han,⁵⁰ R. Handler,⁶⁰ F. Happacher,¹⁹ K. Hara,⁵⁶ M. Hare,⁵⁷ S. Harper,⁴³ R.F. Harr,⁵⁹ R.M. Harris,¹⁷ M. Hartz,⁴⁸ K. Hatakeyama,⁵¹ J. Hauser,⁸ A. Heijboer,⁴⁶ B. Heinemann,³⁰ J. Heinrich,⁴⁶ C. Henderson,³³ M. Herndon,⁶⁰ J. Heuser,²⁶ D. Hidas,¹⁶ C.S. Hill,¹⁰ D. Hirschbuehl,²⁶ A. Hocker,¹⁷ A. Holloway,²² S. Hou,¹ M. Houlden,³⁰ S.-C. Hsu,⁹ B.T. Huffman,⁴³ R.E. Hughes,⁴⁰ U. Husemann,⁶¹ J. Huston,³⁶ J. Incandela,¹⁰ G. Introzzi,⁴⁷ M. Iori,⁵² Y. Ishizawa,⁵⁶ A. Ivanov,⁷ B. Iyutin,³³ E. James,¹⁷ D. Jang,⁵³ B. Jayatilaka,³⁵ D. Jeans,⁵² H. Jensen,¹⁷ E.J. Jeon,²⁸ S. Jindariani,¹⁸ M. Jones,⁴⁹ K.K. Joo,²⁸ S.Y. Jun,¹² J.E. Jung,²⁸ T.R. Junk,²⁴ T. Kamon,⁵⁴ P.E. Karchin,⁵⁹ Y. Kato,⁴² Y. Kemp,²⁶ R. Kephart,¹⁷ U. Kerzel,²⁶ V. Khotilovich,⁵⁴ B. Kilminster,⁴⁰ D.H. Kim,²⁸ H.S. Kim,²⁸ J.E. Kim,²⁸ M.J. Kim,¹² S.B. Kim,²⁸ S.H. Kim,⁵⁶ Y.K. Kim,¹³ N. Kimura,⁵⁶ L. Kirsch,⁶ S. Klimentenko,¹⁸ M. Klute,³³ B. Knuteson,³³ B.R. Ko,¹⁶ K. Kondo,⁵⁸ D.J. Kong,²⁸ J. Konigsberg,¹⁸ A. Korytov,¹⁸ A.V. Kotwal,¹⁶ A. Kovalev,⁴⁶ A.C. Kraan,⁴⁶ J. Kraus,²⁴ I. Kravchenko,³³ M. Krepis,²⁶ J. Kroll,⁴⁶ N. Krumnack,⁴ M. Kruse,¹⁶ V. Krutelyov,¹⁰ T. Kubo,⁵⁶ S. E. Kuhlmann,² T. Kuhr,²⁶ Y. Kusakabe,⁵⁸ S. Kwang,¹³ A.T. Laasanen,⁴⁹ S. Lai,³⁴ S. Lami,⁴⁷ S. Lammel,¹⁷ M. Lancaster,³¹ R.L. Lander,⁷ K. Lannon,⁴⁰ A. Lath,⁵³ G. Latino,⁴⁷ I. Lazzizzera,⁴⁴ T. LeCompte,² J. Lee,⁵⁰ J. Lee,²⁸ Y.J. Lee,²⁸ S.W. Lee,⁵⁴ R. Lefèvre,³ N. Leonardo,³³ S. Leone,⁴⁷ S. Levy,¹³ J.D. Lewis,¹⁷ C. Lin,⁶¹ C.S. Lin,¹⁷ M. Lindgren,¹⁷ E. Lipeles,⁹ A. Lister,⁷ D.O. Litvintsev,¹⁷ T. Liu,¹⁷ N.S. Lockyer,⁴⁶ A. Loginov,⁶¹ M. Loreti,⁴⁴ P. Loverre,⁵² R.-S. Lu,¹ D. Lucchesi,⁴⁴ P. Lujan,²⁹ P. Lukens,¹⁷ G. Lungu,¹⁸ L. Lyons,⁴³ J. Lys,²⁹ R. Lysak,¹⁴ E. Lytken,⁴⁹ P. Mack,²⁶ D. MacQueen,³⁴ R. Madrak,¹⁷ K. Maeshima,¹⁷ K. Makhoul,³³ T. Maki,²³ P. Maksimovic,²⁵ S. Malde,⁴³ G. Manca,³⁰ F. Margaroli,⁵ R. Marginean,¹⁷ C. Marino,²⁶ C.P. Marino,²⁴ A. Martin,⁶¹ M. Martin,²¹ V. Martin,²¹ M. Martínez,³ T. Maruyama,⁵⁶ P. Mastrandrea,⁵² T. Masubuchi,⁵⁶ H. Matsunaga,⁵⁶ M.E. Mattson,⁵⁹ R. Mazini,³⁴ P. Mazzanti,⁵ K.S. McFarland,⁵⁰ P. McIntyre,⁵⁴ R. McNulty,³⁰ A. Mehta,³⁰ P. Mehtala,²³ S. Menzemer,¹¹ A. Menzione,⁴⁷

P. Merkel,⁴⁹ C. Mesropian,⁵¹ A. Messina,³⁶ T. Miao,¹⁷ N. Miladinovic,⁶ J. Miles,³³ R. Miller,³⁶ C. Mills,¹⁰ M. Milnik,²⁶ A. Mitra,¹ G. Mitselmakher,¹⁸ A. Miyamoto,²⁷ S. Moed,²⁰ N. Moggi,⁵ B. Mohr,⁸ R. Moore,¹⁷ M. Morello,⁴⁷ P. Movilla Fernandez,²⁹ J. Mülmenstädt,²⁹ A. Mukherjee,¹⁷ Th. Muller,²⁶ R. Mumford,²⁵ P. Murat,¹⁷ J. Nachtman,¹⁷ A. Nagano,⁵⁶ J. Naganoma,⁵⁸ I. Nakano,⁴¹ A. Napier,⁵⁷ V. Necula,¹⁸ C. Neu,⁴⁶ M.S. Neubauer,⁹ J. Nielsen,²⁹ T. Nigmanov,⁴⁸ L. Nodulman,² O. Norriella,³ E. Nurse,³¹ S.H. Oh,¹⁶ Y.D. Oh,²⁸ I. Oksuzian,¹⁸ T. Okusawa,⁴² R. Oldeman,³⁰ R. Orava,²³ K. Osterberg,²³ C. Pagliarone,⁴⁷ E. Palencia,¹¹ V. Papadimitriou,¹⁷ A.A. Paramonov,¹³ B. Parks,⁴⁰ S. Pashapour,³⁴ J. Patrick,¹⁷ G. Pauletta,⁵⁵ M. Paulini,¹² C. Paus,³³ D.E. Pellett,⁷ A. Penzo,⁵⁵ T.J. Phillips,¹⁶ G. Piacentino,⁴⁷ J. Piedra,⁴⁵ L. Pinera,¹⁸ K. Pitts,²⁴ C. Plager,⁸ L. Pondrom,⁶⁰ X. Portell,³ O. Poukhov,¹⁵ N. Pounder,⁴³ F. Prakoshyn,¹⁵ A. Pronko,¹⁷ J. Proudfoot,² F. Ptohos^e,¹⁹ G. Punzi,⁴⁷ J. Pursley,²⁵ J. Rademacker^b,⁴³ A. Rahaman,⁴⁸ N. Ranjan,⁴⁹ S. Rappoccio,²² B. Reisert,¹⁷ V. Rekovic,³⁸ P. Renton,⁴³ M. Rescigno,⁵² S. Richter,²⁶ F. Rimondi,⁵ L. Ristori,⁴⁷ A. Robson,²¹ T. Rodrigo,¹¹ E. Rogers,²⁴ S. Rolli,⁵⁷ R. Roser,¹⁷ M. Rossi,⁵⁵ R. Rossin,¹⁸ A. Ruiz,¹¹ J. Russ,¹² V. Rusu,¹³ H. Saarikko,²³ S. Sabik,³⁴ A. Safonov,⁵⁴ W.K. Sakumoto,⁵⁰ G. Salamanna,⁵² O. Saltó,³ D. Saltzberg,⁸ C. Sánchez,³ L. Santi,⁵⁵ S. Sarkar,⁵² L. Sartori,⁴⁷ K. Sato,¹⁷ P. Savard,³⁴ A. Savoy-Navarro,⁴⁵ T. Scheidle,²⁶ P. Schlabach,¹⁷ E.E. Schmidt,¹⁷ M.P. Schmidt,⁶¹ M. Schmitt,³⁹ T. Schwarz,⁷ L. Scodellaro,¹¹ A.L. Scott,¹⁰ A. Scribano,⁴⁷ F. Scuri,⁴⁷ A. Sedov,⁴⁹ S. Seidel,³⁸ Y. Seiya,⁴² A. Semenov,¹⁵ L. Sexton-Kennedy,¹⁷ A. Sfyrla,²⁰ M.D. Shapiro,²⁹ T. Shears,³⁰ P.F. Shepard,⁴⁸ D. Sherman,²² M. Shimojima^k,⁵⁶ M. Shochet,¹³ Y. Shon,⁶⁰ I. Shreyber,³⁷ A. Sidoti,⁴⁷ P. Sinervo,³⁴ A. Sisakyan,¹⁵ J. Sjolin,⁴³ A.J. Slaughter,¹⁷ J. Slaunwhite,⁴⁰ K. Sliwa,⁵⁷ J.R. Smith,⁷ F.D. Snider,¹⁷ R. Snihur,³⁴ M. Soderberg,³⁵ A. Soha,⁷ S. Somalwar,⁵³ V. Sorin,³⁶ J. Spalding,¹⁷ F. Spinella,⁴⁷ T. Spreitzer,³⁴ P. Squillacioti,⁴⁷ M. Stanitzki,⁶¹ A. Staveris-Polykalas,⁴⁷ R. St. Denis,²¹ B. Stelzer,⁸ O. Stelzer-Chilton,⁴³ D. Stentz,³⁹ J. Strologas,³⁸ D. Stuart,¹⁰ J.S. Suh,²⁸ A. Sukhanov,¹⁸ H. Sun,⁵⁷ T. Suzuki,⁵⁶ A. Taffard,²⁴ R. Takashima,⁴¹ Y. Takeuchi,⁵⁶ K. Takikawa,⁵⁶ M. Tanaka,² R. Tanaka,⁴¹ M. Tecchio,³⁵ P.K. Teng,¹ K. Terashi,⁵¹ J. Thom^d,¹⁷ A.S. Thompson,²¹ E. Thomson,⁴⁶ P. Tipton,⁶¹ V. Tiwari,¹² S. Tkaczyk,¹⁷ D. Toback,⁵⁴ S. Tokar,¹⁴ K. Tollefson,³⁶ T. Tomura,⁵⁶ D. Tonelli,⁴⁷ S. Torre,¹⁹ D. Torretta,¹⁷ S. Tourneur,⁴⁵ W. Trischuk,³⁴ R. Tsuchiya,⁵⁸ S. Tsuno,⁴¹ N. Turini,⁴⁷ F. Ukegawa,⁵⁶ T. Unverhau,²¹ S. Uozumi,⁵⁶ D. Usynin,⁴⁶ S. Vallecorsa,²⁰ N. van Remortel,²³ A. Varganov,³⁵ E. Vataga,³⁸ F. Vázquezⁱ,¹⁸ G. Velev,¹⁷ G. Veramendi,²⁴ V. Veszpremi,⁴⁹ R. Vidal,¹⁷ I. Vila,¹¹ R. Vilar,¹¹ T. Vine,³¹ I. Vollrath,³⁴ I. Volobouevⁿ,²⁹ G. Volpi,⁴⁷ F. Würthwein,⁹ P. Wagner,⁵⁴ R.G. Wagner,² R.L. Wagner,¹⁷ J. Wagner,²⁶ W. Wagner,²⁶ R. Wallny,⁸ S.M. Wang,¹ A. Warburton,³⁴ S. Waschke,²¹ D. Waters,³¹ W.C. Wester III,¹⁷ B. Whitehouse,⁵⁷ D. Whiteson,⁴⁶ A.B. Wicklund,² E. Wicklund,¹⁷ G. Williams,³⁴ H.H. Williams,⁴⁶ P. Wilson,¹⁷ B.L. Winer,⁴⁰ P. Wittich^d,¹⁷ S. Wolbers,¹⁷ C. Wolfe,¹³ T. Wright,³⁵ X. Wu,²⁰ S.M. Wynne,³⁰ A. Yagil,¹⁷ K. Yamamoto,⁴² J. Yamaoka,⁵³ T. Yamashita,⁴¹ C. Yang,⁶¹ U.K. Yang^j,¹³ Y.C. Yang,²⁸ W.M. Yao,²⁹ G.P. Yeh,¹⁷ J. Yoh,¹⁷ K. Yorita,¹³ T. Yoshida,⁴² G.B. Yu,⁵⁰ I. Yu,²⁸ S.S. Yu,¹⁷ J.C. Yun,¹⁷ L. Zanello,⁵² A. Zanetti,⁵⁵ I. Zaw,²² X. Zhang,²⁴ J. Zhou,⁵³ and S. Zucchelli⁵

(CDF Collaboration*)

¹*Institute of Physics, Academia Sinica, Taipei, Taiwan 11529, Republic of China*

²*Argonne National Laboratory, Argonne, Illinois 60439*

³*Institut de Física d'Altes Energies, Universitat Autònoma de Barcelona, E-08193, Bellaterra (Barcelona), Spain*

⁴*Baylor University, Waco, Texas 76798*

⁵*Istituto Nazionale di Fisica Nucleare, University of Bologna, I-40127 Bologna, Italy*

⁶*Brandeis University, Waltham, Massachusetts 02254*

⁷*University of California, Davis, Davis, California 95616*

⁸*University of California, Los Angeles, Los Angeles, California 90024*

⁹*University of California, San Diego, La Jolla, California 92093*

¹⁰*University of California, Santa Barbara, Santa Barbara, California 93106*

¹¹*Instituto de Física de Cantabria, CSIC-University of Cantabria, 39005 Santander, Spain*

¹²*Carnegie Mellon University, Pittsburgh, PA 15213*

¹³*Enrico Fermi Institute, University of Chicago, Chicago, Illinois 60637*

¹⁴*Comenius University, 842 48 Bratislava, Slovakia; Institute of Experimental Physics, 040 01 Kosice, Slovakia*

¹⁵*Joint Institute for Nuclear Research, RU-141980 Dubna, Russia*

¹⁶*Duke University, Durham, North Carolina 27708*

¹⁷*Fermi National Accelerator Laboratory, Batavia, Illinois 60510*

¹⁸*University of Florida, Gainesville, Florida 32611*

¹⁹*Laboratori Nazionali di Frascati, Istituto Nazionale di Fisica Nucleare, I-00044 Frascati, Italy*

²⁰*University of Geneva, CH-1211 Geneva 4, Switzerland*

²¹*Glasgow University, Glasgow G12 8QQ, United Kingdom*

- ²²Harvard University, Cambridge, Massachusetts 02138
²³Division of High Energy Physics, Department of Physics,
 University of Helsinki and Helsinki Institute of Physics, FIN-00014, Helsinki, Finland
²⁴University of Illinois, Urbana, Illinois 61801
²⁵The Johns Hopkins University, Baltimore, Maryland 21218
²⁶Institut für Experimentelle Kernphysik, Universität Karlsruhe, 76128 Karlsruhe, Germany
²⁷High Energy Accelerator Research Organization (KEK), Tsukuba, Ibaraki 305, Japan
²⁸Center for High Energy Physics: Kyungpook National University,
 Taegu 702-701, Korea; Seoul National University, Seoul 151-742,
 Korea; and SungKyunKwan University, Suwon 440-746, Korea
²⁹Ernest Orlando Lawrence Berkeley National Laboratory, Berkeley, California 94720
³⁰University of Liverpool, Liverpool L69 7ZE, United Kingdom
³¹University College London, London WC1E 6BT, United Kingdom
³²Centro de Investigaciones Energeticas Medioambientales y Tecnologicas, E-28040 Madrid, Spain
³³Massachusetts Institute of Technology, Cambridge, Massachusetts 02139
³⁴Institute of Particle Physics: McGill University, Montréal,
 Canada H3A 2T8; and University of Toronto, Toronto, Canada M5S 1A7
³⁵University of Michigan, Ann Arbor, Michigan 48109
³⁶Michigan State University, East Lansing, Michigan 48824
³⁷Institution for Theoretical and Experimental Physics, ITEP, Moscow 117259, Russia
³⁸University of New Mexico, Albuquerque, New Mexico 87131
³⁹Northwestern University, Evanston, Illinois 60208
⁴⁰The Ohio State University, Columbus, Ohio 43210
⁴¹Okayama University, Okayama 700-8530, Japan
⁴²Osaka City University, Osaka 588, Japan
⁴³University of Oxford, Oxford OX1 3RH, United Kingdom
⁴⁴University of Padova, Istituto Nazionale di Fisica Nucleare,
 Sezione di Padova-Trento, I-35131 Padova, Italy
⁴⁵LPNHE, Universite Pierre et Marie Curie/IN2P3-CNRS, UMR7585, Paris, F-75252 France
⁴⁶University of Pennsylvania, Philadelphia, Pennsylvania 19104
⁴⁷Istituto Nazionale di Fisica Nucleare Pisa, Universities of Pisa,
 Siena and Scuola Normale Superiore, I-56127 Pisa, Italy
⁴⁸University of Pittsburgh, Pittsburgh, Pennsylvania 15260
⁴⁹Purdue University, West Lafayette, Indiana 47907
⁵⁰University of Rochester, Rochester, New York 14627
⁵¹The Rockefeller University, New York, New York 10021
⁵²Istituto Nazionale di Fisica Nucleare, Sezione di Roma 1,
 University of Rome “La Sapienza,” I-00185 Roma, Italy
⁵³Rutgers University, Piscataway, New Jersey 08855
⁵⁴Texas A&M University, College Station, Texas 77843
⁵⁵Istituto Nazionale di Fisica Nucleare, University of Trieste/ Udine, Italy
⁵⁶University of Tsukuba, Tsukuba, Ibaraki 305, Japan
⁵⁷Tufts University, Medford, Massachusetts 02155
⁵⁸Waseda University, Tokyo 169, Japan
⁵⁹Wayne State University, Detroit, Michigan 48201
⁶⁰University of Wisconsin, Madison, Wisconsin 53706
⁶¹Yale University, New Haven, Connecticut 06520
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We present a measurement of the top-quark mass, M_{top} , in the all-hadronic decay channel $t\bar{t} \rightarrow W^+bW^-\bar{b} \rightarrow q_1\bar{q}_2bq_3\bar{q}_4\bar{b}$. The analysis is performed using 310 pb⁻¹ of $\sqrt{s}=1.96$ TeV $p\bar{p}$ collisions collected with the CDF II detector using a multi-jet trigger. The mass measurement is based on an event-by-event likelihood which depends on both the sample purity and the value of the top-quark mass, using 90 possible jet-to-parton assignments in the six-jet final state. The joint likelihood of 290 selected events yields a value of $M_{\text{top}}=177.1 \pm 4.9$ (stat.) ± 4.7 (syst.) GeV/ c^2 .

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*With visitors from ^aUniversity of Athens, ^bUniversity of Bristol, ^cUniversity Libre de Bruxelles, ^dCornell University,

^eUniversity of Cyprus, ^fUniversity of Dublin, ^gUniversity of

The mass of the top quark, M_{top} , is an important free parameter of the standard model (SM) and is of the order of the electroweak symmetry breaking scale. Since virtual top quarks are involved in higher-order electroweak processes, by measuring the top-quark mass one can constrain the mass of the SM Higgs boson [1] and particles predicted in extensions of the SM [2]. At the Tevatron, $t\bar{t}$ pairs are produced by QCD processes and decay according to the CKM matrix [3] with a branching ratio of almost 100% into a W boson and a b quark. The final state of the event is then defined by the decay of the W bosons. All-hadronic $t\bar{t}$ events, where both of the W bosons decay into quarks, have a multi-jet final state and no missing energy due to neutrinos. The top-quark mass measurement in this decay channel is motivated by the large branching fraction ($\approx 44\%$) and the complete reconstruction of both top and antitop quarks, relying only on hadronic jets. It is the first top-quark mass measurement in this channel using Tevatron Run II data, and serves as a consistency check of the measurements in the two other $t\bar{t}$ decay modes [4, 5]. The major experimental challenge is the presence of a large amount of multi-jet background events from QCD processes, which dominates $t\bar{t}$ production by three orders of magnitude, even after applying a trigger dedicated to events with a multi-jet final state.

In this Letter, we present the measurement of the top-quark mass in the all-hadronic decay channel using a sample of $t\bar{t}$ decays corresponding to 310 pb^{-1} of proton-antiproton collisions at $\sqrt{s}=1.96 \text{ TeV}$, collected using the CDF II detector between February, 2002, and August, 2004. We measure the mass using the ideogram method, which was used in the DELPHI experiment at LEP for the W boson mass measurement [6]. The method is based on an event-by-event likelihood that reconstructs the top and antitop quarks using all 90 assignments of jets to quarks from W boson decays and b quarks (jet combinations) from the $t\bar{t}$ decay. Considering all jet combinations enhances the amount of top-quark mass information we can extract from each event. The 90 jet combinations arise from the $6!/8$ permutations of the six jets, where the two possible top antitop permutations and the two possible permutations of each of the two quark pairs coming from the W decay are treated equally.

The sample purity, \mathcal{P} , defined as the fraction of all-hadronic $t\bar{t}$ events contained in the selected data sample, is a free parameter of the likelihood together with the value of M_{top} . The reason for measuring both M_{top} and \mathcal{P} simultaneously is that the cross sections of the QCD

backgrounds are not well known. Even though the background processes do not contain any the top-quark mass information, the sensitivity of the likelihood on M_{top} does depend on the total amount of QCD background events. This is reflected both in the statistical variance of M_{top} , and in corrections applied to M_{top} which take into account the presence of background events.

The CDF II detector [7] is a general-purpose charged and neutral particle detector designed to study $p\bar{p}$ collisions at the Fermilab Tevatron. It consists of an eight-layer silicon microstrip detector array and a cylindrical drift chamber contained inside a 1.4-T solenoid magnet, surrounded by electromagnetic and hadronic sampling calorimeters with a geometrical acceptance up to a pseudorapidity of $|\eta| = 3.6$ [8]. Muon chambers are located outside the calorimeters, allowing the reconstruction of track segments for penetrating particles.

The dataset used for this measurement is selected with a multi-jet trigger that relies solely on calorimetry. This trigger requires at least four calorimeter clusters with a transverse energy $E_{\text{T}} > 15 \text{ GeV}$ and the scalar sum of the transverse jet energies to exceed 125 GeV . To improve the signal-to-background ratio, we impose offline kinematic requirements based on the scalar sum of the transverse jet energies and event shape observables. A full description of the trigger and kinematic selection is given in an earlier publication [9]. For the final top-quark mass measurement we consider only events with exactly six jets, each with transverse energy $E_{\text{T}} > 15 \text{ GeV}$ and pseudorapidity $|\eta| \leq 2$. Jets are identified as clusters of energy deposits in the calorimeter segments (towers) which fall within a cone radius of $\Delta R = \sqrt{\Delta\phi^2 + \Delta\eta^2} \leq 0.4$ [8]. The raw jet energy scale (JES) is the result of a multi-step correction procedure to convert the measured transverse jet energy to the expected transverse energy of the parton corresponding to the jet. The corrections, assessed using data and a simulation of the CDF detector, include corrections for the response inhomogeneity in η , contributions from multiple interactions, the non-linearity of the calorimeter response, the underlying event, and the energy flow out of the jet cone [10]. We apply additional jet-parton energy corrections, specific to $t\bar{t}$ events, which are parameterized independently for b -quark jets and light-flavor jets [11].

In order to further reduce the QCD background, we apply b -quark identification to each event passing the kinematic selection. At least one jet must have a reconstructed displaced vertex (b -tag), consistent with a long-lived bottom hadron [12]. This displaced vertex information is not used to reduce the number of jet combinations in the event likelihood. Instead, we derive a weight factor using the CDF jet probability algorithm [13] which takes into account the probability that a jet originates from a b quark, as explained below. The final data sample contains 290 events. Assuming a theoretical cross section of 6.1 pb [14] for a top-quark mass of $178 \text{ GeV}/c^2$ [15], the

Edinburgh, ^hUniversity of Heidelberg, ⁱUniversidad Iberoamericana, ^jUniversity of Manchester, ^kNagasaki Institute of Applied Science, ^lUniversity de Oviedo, ^mUniversity of London, Queen Mary and Westfield College, ⁿTexas Tech University, ^oIFIC(CSIC-Universitat de Valencia),

expected signal-to-background ratio is about one-to-five.

By applying energy and momentum conservation we fit the momenta of the jets and the two top-quark masses, $m_i^{1,2}$, and determine their estimated uncertainties, $\sigma_i^{1,2}$ using a χ^2 minimization. This kinematic fit is applied to each of the 90 jet combinations, i , where we constrain the masses of the two W bosons to a world average value of $m_W = 80.4 \text{ GeV}/c^2$ [3] within their natural widths. In order to achieve a better background reduction, the two reconstructed top-quark masses, $m_i^{1,2}$, are not constrained to be equal. The $m_i^{1,2}$ distributions thus obtained peak around $160 \text{ GeV}/c^2$ and range from $100\text{--}300 \text{ GeV}/c^2$ for the whole data sample, while the uncertainties $\sigma_i^{1,2}$ are typically of the order of $15 \text{ GeV}/c^2$ and range from $5\text{--}30 \text{ GeV}/c^2$. We measure M_{top} using an event likelihood where we weigh each combination i with a factor w_i , expressing the compatibility with both the $t\bar{t}$ decay kinematics and the presence of two b -quark jets together with four light-flavor jets. The weight factor w_i is calculated as

$$w_i = \exp\left(-\frac{1}{2}\chi_i^2\right)w_i^b, \quad (1)$$

where the exponential term, calculated using the χ^2 value of the kinematic fit, is a measure of the compatibility of jet combination i with the kinematics of a decaying $t\bar{t}$ pair. The second factor, w_i^b , gives the probability of two of the jets to originate from a b (or \bar{b}) quark and four others to originate from light quarks (including c quarks), and is given by

$$w_i^b = \prod_{j=1}^2 p_j^b \prod_{j=3}^6 p_j^q, \quad (2)$$

where we assume quarks with indices 3-6 to be light quarks originating from the W bosons. The probability $p_j^{b(\text{or } q)}$ for a jet j to originate from a b quark (or light-flavor quark, q) is obtained using the CDF jet probability algorithm [13]. This calculates the probability P_j for a given jet j to originate from the primary vertex, based on the impact parameter information from the tracks belonging to that jet. The P_j distribution for b -quark jets (or light-flavor jets) is described by a probability density function, $f(b(\text{or } q)|P_j)$ in order to obtain

$$p_j^{b(\text{or } q)} = \frac{f(b(\text{or } q)|P_j)}{f(b|P_j) + f(q|P_j)}. \quad (3)$$

The event likelihood consists of two terms. The signal likelihood term corresponds to the convolution of two Breit-Wigner distributions, $F_{\text{BW}}(m'_j|M_{\text{top}})$, with two Gaussians, $G(m'_j|m_i^j, \sigma_i^j)$, describing the experimental resolutions, $\sigma_i^{1,2}$, for each of the two reconstructed top-quark masses, $m_i^{1,2}$:

$$\mathcal{L}_i^{\text{sig}}(M_{\text{top}}) = \prod_{j=1,2} \int G(m'_j|m_i^j, \sigma_i^j)$$

$$\times F_{\text{BW}}(m'_j|M_{\text{top}})dm'_j. \quad (4)$$

The M_{top} -independent background likelihood term, $\mathcal{L}_i^{\text{bg}}$, corresponds to the two-dimensional posteriori probability density function (histogram) of $m_i^{1,2}$ obtained from ALPGEN [16] Monte Carlo (MC) multi-jet QCD background. All MC events are passed through the CDF detector simulation and are subjected to the same event selection criteria as the data.

The likelihood for a given event n is derived by summing the signal and background event likelihoods for each jet combination i and is calculated as a function of the top-quark mass, M_{top} , and the sample purity, \mathcal{P} :

$$\mathcal{L}^n(M_{\text{top}}, \mathcal{P}) = \sum_{i=1}^{90} w_i [\mathcal{P}\mathcal{L}_i^{\text{sig}}(M_{\text{top}}) + (1 - \mathcal{P})\mathcal{L}_i^{\text{bg}}]. \quad (5)$$

We obtain a one-dimensional likelihood curve as a function of M_{top} by maximizing the two-dimensional joint likelihood,

$$L(M_{\text{top}}, \mathcal{P}) = \prod_{n=1}^{290} \mathcal{L}^n(M_{\text{top}}, \mathcal{P}), \quad (6)$$

with respect to the sample purity, \mathcal{P} , for each value of M_{top} . The maximum of the total likelihood of the 290 selected data events corresponds to a sample purity value of $\mathcal{P} = 0.21 \pm 0.07(\text{stat.})$, which is compatible with a signal-to-background ratio of about one-to-five expected from SM $t\bar{t}$ production. The value of M_{top} extracted

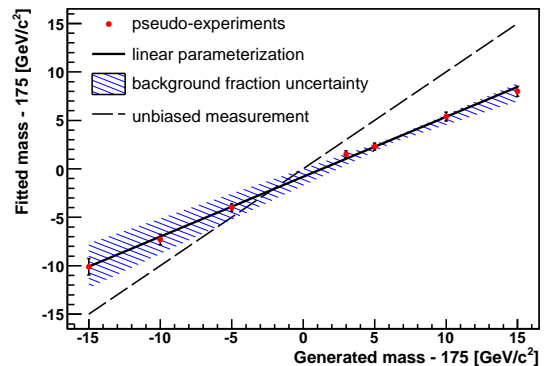


FIG. 1: The correlation between the measured value of M_{top} and the generated top-quark mass obtained with pseudo-experiments using $t\bar{t}$ signal events generated with HERWIG and an admixture of QCD background generated with ALPGEN. The solid line represents the linear parameterization used to correct M_{top} . The shaded band shows the effect of a one-sigma variation of the sample purity, $\mathcal{P} = 0.21 \pm 0.07(\text{stat.})$. The dashed line corresponds to 100% correlation and zero offset.

from the joint likelihood fit in Eq. (6) is biased due to the presence of wrong jet combinations, background events,

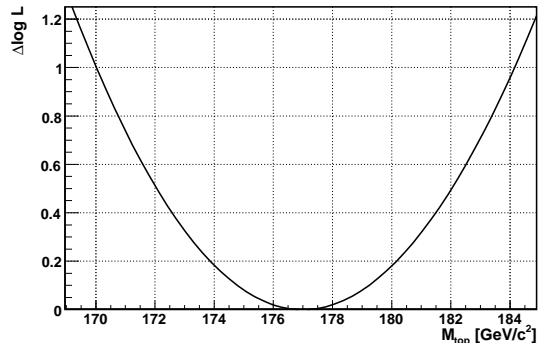


FIG. 2: The one-dimensional negative log-likelihood of M_{top} obtained including corrections using the parameterization shown in Fig. 1

and assignment of jets that arise from initial- (ISR) and final-state gluon radiation (FSR). Therefore, we correct the measurement of M_{top} using a linear parameterization of the correlation between the generated top-quark mass value and the mass estimator M_{top} in the range between $160 \text{ GeV}/c^2$ and $190 \text{ GeV}/c^2$, as shown in Fig. 1. Each mass point is obtained from pseudo-experiments with MC samples containing a mixture of 290 Poisson-fluctuated signal plus background events. We use HERWIG [17] to generate the $t\bar{t}$ signal events and ALPGEN to model the QCD background. The tag rate matrix used in [9] models the false identification of b -quarks. The correction depends on the sample purity used to construct the pseudo-experiments as illustrated by the shaded band which corresponds to the one-sigma statistical uncertainty on the sample purity measured from data. Accordingly, we assign a systematic uncertainty due to this purity uncertainty. Using the same pseudo-experiments we verify that the measured sample purity is independent of the generator top-quark mass value and that the statistical uncertainty on M_{top} covers 68% of the sample measurements. We correct the statistical uncertainty from the likelihood curve using the width of the pull distribution $(M_{\text{top}} - m_{\text{gen}})/\sigma_{M_{\text{top}}}$, where m_{gen} is the generated top-quark mass, M_{top} the measured top-quark mass, and $\sigma_{M_{\text{top}}}$ its statistical uncertainty obtained in the pseudo-experiments. The correction factor, 1.17, is independent of the value of M_{top} .

The maximum of the one-dimensional top mass likelihood, after correction with the parameterization shown in Fig. 1, corresponds to the value $M_{\text{top}} = 177.1 \pm 4.9(\text{stat.}) \text{ GeV}/c^2$. Figure. 2 shows the corresponding likelihood.

The uncertainty on the JES which varies between approximately 3% and 8% of the measured jet energy, depending on the η and p_T of the jet [10], dominates the systematic uncertainties on our measurement of M_{top} . We determine its effect using pseudo-experiments in which

the JES is increased (decreased) with one standard deviation. Half of the difference in M_{top} between the positive and negative variation of the JES amounts to $\Delta M_{\text{top}} = 4.3 \text{ GeV}/c^2$.

The corrections to M_{top} , shown in Fig. 1, depend both on the background fraction and shape. The uncertainty due to the background fraction, varied within the uncertainties of the measured sample purity, is $\Delta M_{\text{top}} = 1.1 \text{ GeV}/c^2$ for the measured value of M_{top} . The relative fractions of the light-flavor QCD background and QCD background containing b quarks is not precisely known. We estimate their ratio (light-flavor/ b) to be 2.8 ± 1.0 by performing a binned maximum likelihood fit of the MC expectation to the data based on the jet probability weight distribution in Eq. (2). This result is consistent with a value of 2.7 obtained using the respective QCD cross sections given by the ALPGEN event generator, taking the individual selection efficiencies into account. A systematic uncertainty due to the background composition results in $\Delta M_{\text{top}} = 0.8 \text{ GeV}/c^2$ by varying the background flavor composition between only light-flavor QCD background events and only QCD background containing b quarks.

In accordance with [5], we vary two PYTHIA [18] parton shower parameters, Λ_{QCD} and the ISR/FSR transverse momentum scale K -factor, in order to model the systematic uncertainties due to the amount of ISR and FSR. The largest difference observed in M_{top} between any variation amounts to $0.9 \text{ GeV}/c^2$ ($0.8 \text{ GeV}/c^2$) for ISR (FSR).

To examine the systematic effect due to uncertainties in the parton distribution function (PDF) of the proton and antiproton, we follow the approach used in [5], resulting in $\Delta M_{\text{top}} = 0.8 \text{ GeV}/c^2$.

By applying the parameterized corrections on M_{top} obtained with HERWIG to MC events generated with PYTHIA we estimate MC modeling uncertainties. The resulting difference in M_{top} equals $0.5 \text{ GeV}/c^2$.

Uncertainties in b -quark fragmentation, semi-leptonic branching ratios, and color-flow affect the b -quark JES [5]. This results in an additional contribution to ΔM_{top} of $0.5 \text{ GeV}/c^2$. All systematic uncertainties are summarized in Table I. They result in a total systematic uncertainty of $\Delta M_{\text{top}} = 4.7 \text{ GeV}/c^2$.

In summary, we present a new measurement of the top-quark mass in the all-hadronic channel. The measured value is $M_{\text{top}} = 177.1 \pm 4.9(\text{stat.}) \pm 4.7(\text{syst.}) \text{ GeV}/c^2$. This measurement is the first determination of the top-quark mass in the all-hadronic channel using Run II data and is twice as precise as the Run I measurements in this channel [19, 20]. We expect the statistical uncertainty of this measurement to be reduced to $2 \text{ GeV}/c^2$ after analyzing 1 fb^{-1} of integrated luminosity. Our result can be compared with the latest CDF top-quark mass measurements in the dilepton [4] and lepton+jets [5] channel, which are based on samples with a similar integrated luminosity. The measured value for the dilepton chan-

TABLE I: Systematic uncertainties on the measured top-quark mass.

Systematic uncertainty source	ΔM_{top} [GeV/ c^2]
JES	4.3
Background fraction	1.1
Background shape	0.8
ISR	0.9
FSR	0.8
PDF	0.8
MC modeling (generator)	0.5
b -quark JES	0.5
Total	4.7

nel is $M_{\text{top}} = 167.9 \pm 5.2(\text{stat.}) \pm 3.7(\text{syst.}) \text{ GeV}/c^2$, and the value for the lepton+jets channel is $M_{\text{top}} = 173.5^{+3.7}_{-3.6}(\text{stat.} + \text{JES}) \pm 1.3(\text{syst.}) \text{ GeV}/c^2$, where the first uncertainty includes the uncertainty on the jet energy scale, which is measured simultaneously with M_{top} . At this level of precision, the values of M_{top} are compatible between all three decay channels.

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