

# Method for decoupling error correction from privacy amplification

Hoi-Kwong Lo\*

*MagiQ Technologies, Inc., 275 Seventh Avenue, 26th Floor, New York, NY 10001-6708*  
(September 17, 2004)

## Abstract

Entanglement purification provides a unifying framework for proving the security of quantum key distribution schemes. Nonetheless, up till now, a *local* commutability constraint in the CSS code construction means that the error correction and privacy amplification procedures of BB84 are not fully decoupled. Here, I provide a method to decouple the two processes completely. The method requires Alice and Bob to share some initial secret string and use it to encrypt the bit-flip error syndrome using one-time-pad encryption. As an application, I prove the unconditional security of the interactive Cascade protocol, proposed by Brassard and Salvail for error correction, modified by one-time-pad encryption of the error syndrome, and followed by the random matrix protocol for privacy amplification. This is an efficient protocol in terms of both computational power and key generation rate.

Keywords: Quantum Cryptography, Quantum Key Distribution, Unconditional Security

Typeset using REVTeX

---

\*email: hoi.kwong@magiqtech.com

## I. INTRODUCTION

An important application of quantum information processing is quantum key distribution (QKD) [1,2]. The goal of QKD is to allow two communicating parties to detect any eavesdropper. Unlike conventional key distribution scheme, QKD makes no assumptions on the eavesdropper's computing power. Rather, the security of QKD is supposed to be based on the fundamental laws of quantum mechanics.

Security proofs of QKD is an important but difficult problem in quantum information theory. Recently, entanglement purification [3,4] has become a fruitful avenue of studying the security of QKD. Roughly speaking, entanglement purification is a generalized form of quantum error correction for a quantum *communication* channel, rather than quantum storage which is dealt with by standard quantum error correction. It was first suggested by Deutsch *et al.* [5] that entanglement purification protocols (EPPs) can correct errors introduced by the eavesdroppers and allow the two communicating parties, Alice and Bob, to obtain perfectly entangled (i.e., quantum-mechanically correlated) quantum systems, so-called EPR pairs, from which they can generate a secure key.

A proof of security by Mayers [6] applies to a standard QKD scheme, BB84 [1], published by Bennett and Brassard in 1984. Mayers' proof makes no explicit reference to entanglement purification, but is rather complex. A proof of security of QKD based on entanglement purification has been provided by Lo and Chau [7]. It has the advantage of being intuitive and conceptually simple, but it requires that Alice and Bob possess quantum computers for its implementation. Recently, Shor and Preskill [8] has removed this requirement and applied the approach of entanglement purification to prove the security of BB84 [1]. Other proofs of security of QKD that make no explicit reference to entanglement purification include [9,10]. Recently, a security proof with a practical set-up (weak coherent states, lossy channels and inefficient detectors, etc) has been presented by Inamori, Lütkenhaus and Mayers [11].

Recall that error correction and privacy amplification are necessary in the generation of the final secure key from the raw quantum transmission data. Error correction ensures that Alice and Bob will share a common string and, roughly speaking, privacy amplification ensures that Eve most likely knows almost nothing about the key. Unfortunately, so far the application of entanglement purification approach to QKD implies a non-trivial constraint between the two processes, namely the corresponding measurement operators employed by Alice and Bob must be *locally commuting*. Such a local commutability constraint means that the two processes are not totally decoupled from each other. Therefore, it is not entirely obvious how to study error correction and privacy amplification independently.

In this paper, I propose a novel method to remove this local commutability constraint, thus allowing us to decouple the error correction process from the privacy amplification process. This amounts to much simplification in the study of both processes. In the EPP picture, the proposed method requires Alice and Bob to share some ancillary pre-distributed pure EPR pairs. Instead of measuring the (bit-flip) error syndrome directly, each user collects the output into those ancillary EPR pairs and measures those pairs. (A specific instance, so-called breeding method, of such a general method, was used in [3].) In the BB84 picture, the proposed method requires Alice and Bob to share initially some common secret ancillary binary string,  $a$ . Instead of announcing the bit-flip error syndrome, which is a binary string  $x$ , each user encrypts the error syndrome bit-wise using  $a$  as a one-time pad and announces

the encrypted version,  $y = x + a \pmod{2}$ , bit-wise.

As an application of the proposed method, I consider a rather general class of classical error correction methods—the so-called *symmetric* stabilizer-based<sup>1</sup> schemes—which may involve either one-way or two-way classical communications. I show that any symmetric stabilizer-based scheme can be modified and subsequently combined with *any symmetric* stabilizer-based privacy amplification procedure into an unconditionally secure protocol for QKD. This means that one can study the two processes—error correction and privacy amplification—independently. Such a decoupling of error correction from privacy amplification allows one to simplify the analysis of security of a general error correction scheme.

As an application, I prove the unconditional security of a modified version of the Cascade scheme [12] for error correction invented by Brassard and Salvail, (followed by, for example, a random hashing procedure for privacy amplification [6]). This is the first time such a computationally efficient scheme has been proven to be secure. Therefore, the result is of practical interest.

Finally, note that the proposed method can be employed as a sub-routine in concatenated entanglement purification procedures, including those involving two-way classical communications, as studied by [13] and those involving degenerate codes [14].

## II. MOTIVATION

A key motivation of this work is to provide a rigorous proof of security of interactive protocols for error correction in QKD. Let me explain in detail. In QKD, one often has to perform error correction at a rather high bit error rate of say a few percents, which is much higher than the typical value of say  $10^{-5}$  in conventional communications. Moreover, one would like the key generation rate to remain high. As a rule of thumb, the fewer bits are exchanged between Alice and Bob, the higher the key generation rate. Furthermore, one would like to implement a QKD scheme efficiently. That is to say with a minimal amount of computational power. In a general implementation of QKD, it is a highly complex question what the trade-off between the various parameters—tolerable error rate, key generation rate, computational power—would be the best.

Forward error correction is commonly employed in conventional communications and works efficiently at low error rates. Unfortunately, QKD has a high bit error rate. If forward error correction is employed in QKD, a very large block size of order  $10^5$  would probably be needed. This translates to a large amount of computing power.<sup>2</sup>

Two-way communications between Alice and Bob are useful in reducing the required computing power for error correction. In the literature, several interactive protocols such as “BBSS” [15] and “Cascade” [12] have been proposed for error correction in QKD.<sup>3</sup> The

---

<sup>1</sup>By stabilizer-based, I only mean that each operator that Alice measures is a Pauli operator. The various operators are *not* required to commute.

<sup>2</sup>I thank enlightening discussions with Tsz-Mei Ko and Norbert Lütkenhaus on this point.

<sup>3</sup>I thank Norbert Lütkenhaus for providing the references.

Cascade protocol, invented by Brassard and Salvail, for instance, has the advantages of being computationally highly efficient and also being one of the best methods in minimizing the number of exchanged bits between Alice and Bob. It works very well in a few percents bit error rate. Therefore, Cascade is well suited for implementations. Unfortunately, up till now, a proof of unconditional security of a QKD scheme based Cascade (and followed by, for example, standard Shor-Prekill [8] or Mayers [6] privacy amplification procedure) has been missing. A key contribution of this paper is to provide such a proof. The proof of security applies not only to Cascade, but to *any* (interactive or non-interactive) protocols for error correction that are based on parity computations in QKD.

Another motivation for this work is to demonstrate the decoupling of error correction from privacy amplification. On the conceptual level, a QKD scheme consists of several steps—“advantage distillation” [16], error correction and privacy amplification. Entanglement purification has recently been proposed by Shor and Preskill [8] as a useful framework for dealing with BB84. The work of Shor and Preskill built on earlier work in [7] and has been subsequently extended in [13] to protocols involving two-way communications and in [14] to the six-state [17] QKD scheme.

Nonetheless, an important constraint remains in those works: The measurement operators employed by Alice and Bob must commute *locally*. This local commutability constraint ensures that those observables are simultaneous observables. Therefore, the measurement of one observable does not introduce any “back-reaction” to the measurement of any other observables. Such a local commutability constraint means that in analyzing QKD, one has to study both error correction and privacy amplification together and ensure that the observables that Alice and Bob measure *do* commute locally. Therefore, this constraint complicates the analysis.

Analysis of protocols of QKD would be greatly simplified if one could divide up its procedure into different components and analyze each component *independently*. A main contribution of this paper is to show that such a decoupling is, in fact, possible for error correction and privacy amplification. The upshot is that, one can study error correction and pick the best that one can find. Then, one studies privacy amplification and pick the best that one can find. Finally, one puts the two together and the composite will remain good. This result is reminiscent of the decoupling of source coding from error correction in classical coding theory.<sup>4</sup>

### III. BB84

The best-known QKD scheme is BB84, in which the sender, Alice, prepares and sends to the receiver, Bob, a sequence of single photons randomly in one of the four polarizations, horizontal, vertical, 45-degrees and 135-degrees. Bob then performs a measurement ran-

---

<sup>4</sup>Actually, the decoupling result in classical coding theory is stronger than what I have stated here. It shows that the combined protocol is *optimal*, even in the case of a finite block size. In contrast, no claim of optimality for the decoupling result for QKD is claimed in the present paper. The issue of optimality is beyond the scope of the current paper.

domly one of the two polarization bases—rectilinear and diagonal. BB84 is an example of standard “prepare-and-measure” protocols, which can be executed without quantum computers. Proving the security of BB84 against the most general attack by the eavesdropper, Eve, turned out to be a hard problem.

### A. entanglement purification based QKD

Entanglement purification [3] has become a useful proof technique. Consider the following entanglement purification based QKD scheme. Alice prepares a sequence of say  $2N$  EPR pairs and sends half of each pair to Bob. Owing to channel noises and eavesdropping attacks, those pairs will be corrupted. Alice and Bob randomly sample say  $N$  of their pairs to estimate the error rates in the two bases. If the error rates are too high, they abort. Otherwise, they now apply a so-called entanglement purification protocol (EPP)  $C$ , which distills from the  $N$  remaining impure pairs a smaller number, say  $m$ , of almost perfectly entangled EPR pairs. They then measure those pairs to generate a secure key.

First of all, suppose Alice and Bob share  $m$  nearly perfect EPR pairs and generate a key by measuring them. The following theorem shows that Eve cannot have much information on the key.

**Theorem 1** ([7]) *If a density matrix  $\rho$  has high fidelity  $F$  to a state of  $m$  perfect EPR pairs, and Alice and Bob produce their key by measuring individual qubits of  $\rho$ , then with high probability, Alice and Bob have identical  $m$ -bit strings  $k$  with a uniform distribution, and Eve has essentially no information about  $k$ . In fact, if  $F \rightarrow 1$  exponentially with  $m$ , then Eve’s information approaches 0 exponentially with  $m$  as well.<sup>5</sup>*

*Definition: Bell-basis.* Given a pair of qubits, a convenient basis to use is the Bell-basis, which has Bell states as its basis vectors. The Bell states are of the form:

$$\Psi^\pm = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle \pm |\downarrow\uparrow\rangle) \quad (1)$$

and

$$\Phi^\pm = \frac{1}{\sqrt{2}}(|\uparrow\uparrow\rangle \pm |\downarrow\downarrow\rangle). \quad (2)$$

It is convenient to label them by two bits such that:

$$\begin{aligned} \Phi^+ &= 00 \\ \Psi^+ &= 01 \\ \Phi^- &= 10 \\ \Psi^- &= 11. \end{aligned}$$

---

<sup>5</sup>As discussed in [18], if we demand that Eve’s information is bounded by some small number independent of  $k$ , then the number of test particles only scales logarithmically with  $k$ .

*Definition: N-Bell basis and BDSW notations.* Suppose Alice and Bob share  $N$  pairs of qubits. A convenient basis to use is the  $N$ -Bell basis. That is to say, each basis vector is the tensor product state of  $N$  Bell basis vectors. Following Eq. 3, it is convenient to label an  $N$ -basis vector by  $2N$  bits. This is the notation employed by Bennett, DiVincenzo, Smolin and Wootters (BDSW) [3].

*Definition: Pauli operator.* A Pauli operator,  $\mathcal{P}$ , is defined as a tensor product of single-qubit operators of the form  $I$  (the identity),  $X$ ,  $Y$  and  $Z$  where  $X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$ ,  $Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$  and  $Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$ .

*Definition: Stabilizer.* An Abelian group whose generators are Pauli operators is called a stabilizer group.

*Definition: Correlated Pauli strategies.* An eavesdropper, Eve, is said to be employing a correlated Pauli strategy if she applies a Pauli operator,  $\mathcal{P}_i$ , to the quantum signals with some probability  $p_i$ .

*Definition: Symmetric stabilizer-based EPP.* An EPP is called symmetric, stabilizer-based if it involves Alice and Bob measuring operators that are the generators of some stabilizer group.

While Eve may use any eavesdropping strategy, the following theorem states essentially that, to consider security, one only needs to consider correlated Pauli strategies.

**Theorem 2** (Adapted from [7]) *Suppose Alice creates  $M$  EPR pairs and sends half of each to Bob. Alice and Bob then test the error rates,  $p_X$  and  $p_Z$ , along the  $X$  and  $Z$  bases for randomly chosen disjoint subsets,  $s_1$  and  $s_2$ , each of  $m \ll M$  objects respectively. If the error rate is too high, they abort. Otherwise, they perform an EPP  $\mathcal{C}$  on the remaining  $N = M - 2m$  pairs to try to distill out  $k$  EPR pairs of high fidelity. Suppose, the EPP  $\mathcal{C}$  can correct up to  $N(p_X + \varepsilon)$  phase errors and up to  $N(p_Z + \varepsilon)$  bit-flip errors. Define a Hilbert subspace  $\mathcal{H}_{good}$  of the  $N$  EPR pairs to be the subspace spanned by  $N$ -Bell-states with good error patterns. (i.e., with up to  $N(p_X + \varepsilon)$  phase errors and up to  $N(p_Z + \varepsilon)$  bit-flip errors). Let us denote the projection operator into  $\mathcal{H}_{good}$  by  $\Pi$ . Then, we have the following:*

*Given any eavesdropping strategy,  $\mathcal{S}_1$ , by Eve, there exists a correlated Pauli strategy,  $\mathcal{S}_2$ , by Eve that will yield exactly the same values to the following two important quantities:*

(i)  $P(\text{verification test is passed by the test sample} \mid s_1, s_2)$

and

(ii)  $\text{tr}(\Pi\rho)$ ,

for all choices of  $s_1$ , and  $s_2$ .

*Sketch of Proof:* The ‘‘commuting observables’’ idea in [7] is employed. An eavesdropping strategy is defined by the choice of an ancilla and the unitary transformation between the combined system of the ancilla and the  $N$  EPR pairs. Given any eavesdropping strategy  $\mathcal{S}_1$  by Eve, let us consider a fixed but arbitrary choice of sampling subsets,  $s_1$ , and  $s_2$ . Let  $O_{s_1, s_2}$  be the observable that determines whether the verification test is passed. Recall that  $\Pi$  is defined as the projection operator into the good (i.e., correctable) Hilbert space. Consider also  $W$ , the observable that gives the  $2M$ -bit string representing the state  $w$  in the BDSW notation. Since all the observables,  $O_{s_1, s_2}$ ’s,  $\Pi$ ’s and  $W$  are simultaneously diagonalizable in the  $M$ -Bell basis, they all commute with each other. Therefore, it is mathematically

consistent to assign probabilities to the simultaneous eigenvalues of those observables, thus giving rise to the two quantities  $P(\text{verification test is passed by the test sample} \mid s_1, s_2)$  and  $\text{tr}(\Pi\rho)$  for all possible choices of  $s_1$ , and  $s_2$ .

Now, imagine applying a hypothetical measurement  $W$  to Alice and Bob's state before the measurements of  $O_{s_1, s_2}$ 's and  $\Pi$ 's. Given that  $W$  commutes with  $O_{s_1, s_2}$ 's and  $\Pi$ 's, a prior measurement of  $W$  in no way effects the outcomes of measurements of  $O_{s_1, s_2}$ 's and  $\Pi$ 's. In other words, if Eve pre-measures the state in the  $N$ -Bell-basis (i.e., measures  $W$ ), neither the probability of passing the verification test, nor the probability of being in the good Hilbert space will be affected by such a prior measurement. However, with such a prior measurement, Eve has reduced her eavesdropping strategy  $S_1$  to a correlated Pauli's strategy,  $S_2$ .

*Remark:* This commuting observables idea applies to all symmetric stabilizer-based EPPs including ones that involve two-way classical communications.

Theorem 2 is telling us that one can treat the two important quantities—i) the probability of passing a verification test and ii) the probability of being in a good Hilbert space,  $\text{tr}(\Pi\rho)$ —as *classical*. In essence, one can apply classical sampling theory to a quantum problem. Furthermore,  $\text{tr}(\Pi\rho)$  provides a bound to the fidelity of the corrected EPR pairs:

**Theorem 3** ([8,19]) *Consider a stabilizer-based EPP  $\mathcal{C}$  which distills  $m$  EPR pairs from  $n$  impure pairs. Suppose  $\mathcal{C}$  works perfectly in a Hilbert subspace  $\mathcal{H}_{\text{good}}$ , which is spanned by Bell-states with good error patterns (i.e., correctable by  $\mathcal{C}$ ). Denote the projection operator onto  $\mathcal{H}_{\text{good}}$  by  $\Pi$ . If we apply the EPP  $\mathcal{C}$  to an initial state  $\rho$ , then the fidelity of the recovered state as  $m$  EPR pairs is bounded below by*

$$F \equiv \langle \bar{\Phi}^{(m)} \mid \rho_{\text{rec.}} \mid \bar{\Phi}^{(m)} \rangle \geq \text{tr}(\Pi\rho). \quad (3)$$

Here,  $\rho_{\text{rec.}}$  is the recovered state after error correction,  $\bar{\Phi}^{(m)}$  is the  $m$ -EPR pair state.

*Proof:* This Theorem follows from standard stabilizer quantum error correcting code (QECC) theory. An explicit proof of essentially the same result can be found in [19]. Q.E.D.

## B. reduction to BB84 via CSS codes

Because of Theorem 3, EPP based QKD schemes are particularly convenient to analyze. Unfortunately, they are difficult to implement because they generally require Alice and Bob to possess quantum computers. A key insight of Shor and Preskill is to remove the requirement of quantum computers by showing that, in fact, the security of a special class of EPP based QKD schemes implies the security of BB84. More concretely, they considered a special class of quantum error-correcting codes, called Calderbank-Shor-Steane (CSS) [20,21] codes (see below for properties of CSS codes) and proved the following theorem:

**Theorem 4** ([8]) *Given an EPP-based QKD scheme that is based on a CSS code and a verification procedure that involves only two bases, its security implies the security of a BB84 scheme.*

*Remark:* Similarly, when the verification procedure involves three bases, an analogous Theorem shows that the security of an EPP-based QKD scheme that is based on a CSS code implies the security of the six-state scheme.

We shall refer the readers to [8,19] for details of the proof of Theorem 4. A CSS code is a stabilizer-based quantum code with generators that are either i) tensor products of the identities and  $Z$ 's only or ii) tensor products of the identities and  $X$ 's only. It has the advantage that the phase and bit-flip error correction procedures are totally decoupled from each other.<sup>6</sup>

More concretely, a CSS code is defined as follows: Consider a binary linear classical code  $C_1$  and its subcode  $C_2$ . A codeword of a CSS code is an equal superposition of codewords of  $C_1$  that are in the same coset of  $C_2$ :

$$|\phi_u\rangle = \sum_{v \in C_2} |u + v\rangle. \quad (4)$$

Note that, if  $u_1 - u_2 \in C_2$ , then  $|\phi_{u_1}\rangle = |\phi_{u_2}\rangle$ . Therefore, the codeword of a CSS code is in one-one correspondence with the cosets of  $C_2$  in  $C_1$ . Suppose both  $C_1$  and the dual of  $C_2$ ,  $C_2^\perp$ , can correct up to  $t$  errors. Then, the CSS code based on  $C_1$  and  $C_2$  can correct up to  $t$  bit-flip errors and  $t$  phase errors.

On reduction from EPP to BB84, the EPP leaves its mark as an error correction/privacy amplification protocol in the following manner. Alice sends a random quantum state  $|w\rangle$  to Bob. Owing to noises in the channel and eavesdropping actions, Bob receives it as a corrupted string  $w + e$ . Afterwards, Alice picks a random codeword  $u \in C_1$  and broadcasts  $w + u$ . Bob subtracts this from his string to obtain  $u + e$ . He then corrects the error to obtain  $u$ . Finally, he generates the key as the coset  $u + C_2$ . Notice that, the cosets of a code, say  $C_2$ , is in one-one correspondence with the error syndromes. Indeed, the value of the key is given by the error syndrome of the subcode  $C_2$  for a codeword in  $C_1$ .

Using CSS codes and Theorem 4, BB84 is proven to be secure up to an error rate of 11 percents. By using two-way classical communications, BB84 can be made secure at a much higher error rate of about 17 percents.<sup>7</sup> This is due to the following theorem by Gottesman and myself [13], which generalizes Theorem 4.

---

<sup>6</sup>Applying an operator  $X$  to a state will introduce a bit-flip error to the state. Similarly, applying an operator  $Z$  to a state will introduce a phase error. Finally, applying an operator  $Y$  will lead to both a bit-flip and a phase error.

The intuitive reason why an EPP-based QKD can be reduced to BB84 is that Alice and Bob do not need to compute or announce their phase error syndromes. This is because the phase errors do not affect the value of the final key. Roughly speaking, randomizing the state over all possible phase error syndromes, one recovers BB84. In other words, provided that, from Eve's point of view, Alice and Bob *could have* performed the QKD scheme by quantum computers, the resulting BB84 scheme is secure. Alice and Bob do not really have to use quantum computers.

<sup>7</sup>Note that it has been shown that BB84 with only one-way classical communications is necessarily insecure at an error rate of about 15% [22,23]. Therefore, this result in [13] shows clearly that BB84 with two-way classical communications is definitely better than BB84 with only one-way classical communications.



**Theorem 5** ( [13]) *Suppose a two-way EPP satisfies the following conditions:*

1. (*Symmetric*) *It can be described as a series of measurements  $M_i$ , with both Alice and Bob measuring the same  $M_i$ .*
2. (*CSS-like*) *Each of its generators  $M_i$  can be written as either a) a product of  $X$ 's only or b) a product of  $Z$ 's only.*
3. (*Locally-commuting*) *Each pair of  $M_i$  and  $M_j$  commute locally in Alice's (or Bob's) side.*
4. (*Conditional on  $Z$ 's only*) *All conditional operations depend on the result of measuring  $Z$  operators only.*

*Let us call such a protocol a reducible protocol. Claim: a reducible protocol can be converted to a standard "prepare-and-measure" QKD scheme with security equal to the EPP-based QKD scheme.*

*Remark:* Here, the notation has been slightly abused. By a products of  $Z$ 's only, I actually mean a product of the identities and the  $Z$ 's only. Similarly, for  $X$ 's.

*Remark:* If the verification stage involves two bases, then the "prepare-and-measure" QKD scheme is BB84. If it involves three bases, then the "prepare-and-measure" QKD scheme is the six-state scheme.

We will refer the readers to [13] for the details of the proof of Theorem 5.

#### IV. CONSTRAINT ON LOCAL COMMUTABILITY

Theorem 5 is a strong result in QKD. Nonetheless, the constraint 3 in Theorem 5 seriously restricts its applicability. In the EPP picture, the constraint demands that all the local measurement operators that Alice and Bob employ must *commute locally* with each other. Therefore, one is not at liberty to choose the bit-flip and phase error correction measurement operators independently.

I remark that the local commutability constraint is a big obstacle in the application of Theorem 5 to prove the security of interactive Cascade scheme [12] for error correction proposed by Brassard and Salvail. Recall the Cascade protocol involves a binary search subroutine, "BINARY", by Alice and Bob, which allows them to identify the location of an error. The binary search subroutine, BINARY, involves the computation of the parity of a set and subsequently dividing it into two sets and computing the parity of each subset, etc, until the location of the error is found. Note that at the end of BINARY, the size of a subset is reduced to a single object, which means Alice (and also Bob) has to announce the eigenvalue  $Z_i$  of a *single* qubit at location  $i$  (i.e., the  $i$ -th qubit). Now, any quantum error correcting procedure that corrects the phase error of the announced bit must contain a measurement operator  $M$  with a component  $X_i$  for also the  $i$ -th qubit. This means that  $M$  anti-commutes, rather than commutes with  $Z_i$ . In conclusion, with Cascade protocol, it would be impossible to correct all the phase errors. Therefore, the application of Theorem 5 to the Cascade protocol looks problematic.

### A. Using ancillary EPR pairs

To resolve this problem of local non-commutability, notice that a) all symmetric measurement operators,  $M_i = M_i^A \otimes M_i^B$  do commute globally and b) in many cases, only this *relative* error syndrome between Alice and Bob is of interest. For instance, in BINARY, Alice and Bob are interested in only whether their corresponding parities agree or disagree, but not in the actual values of the individual parities. A simple method to bring two distant quantum systems together and allow a *global* operator to be measured is teleportation. To achieve teleportation, some ancillary EPR pairs must be shared by Alice and Bob. This motivates the basic insight of the current paper—to use ancillary EPR pairs to compute the relative error syndrome.

Instead of teleportation, a more efficient way of measuring the global error syndrome will be employed. Here is a main theorem of the current paper.

**Theorem 6** *Suppose Alice and Bob share a number of impure EPR pairs and they would like to compute  $r$  symmetric global operators each of the form  $M_i = M_i^A \otimes M_i^B$  (As before, by symmetric, it means that  $M_i^A$  is the same as  $M_i^B$  except that they act on Alice's and Bob's Hilbert spaces respectively) and  $M_i^A$  is a Pauli operator. Suppose further that they would like to know only the eigenvalues of  $M_i$ 's, but otherwise leave the state unchanged. The claim is that they can do so with  $r$  ancillary EPR pairs.*

*Sketch of Proof:* The notation is such that an EPR pair is an eigenstate of  $ZZ$  and  $XX$ , with eigenvalue  $+1$  for both. Let us call the two qubits of the  $j$ -th EPR pair shared by Alice and Bob,  $A'_j$  and  $B'_j$  respectively. For each operator,  $M_i$ , Alice measures  $M_i^A \otimes Z_{A'_i}$  and broadcasts her outcome and Bob measures  $M_i^B \otimes Z_{B'_i}$  and broadcasts his outcome. The relative outcome, the product of  $M_i^A \otimes Z_{A'_i} \otimes M_i^B \otimes Z_{B'_i}$  gives the eigenvalue of the operator  $M_i$  (because the state of the ancillary EPR pair gives an eigenvalue  $+1$  for the operator  $Z_{A'_i} \otimes Z_{B'_i}$ ). More importantly, by an explicit calculation analogous to the argument in teleportation, one can show that no disturbance to the state is made except for the determination of the eigenvalue of  $M_i = M_i^A \otimes M_i^B$ . Q.E.D.

The above theorem employs a generalization of the so-called breeding method for EPP, studied in [4] (see also [3]). In [3], the breeding method was only mentioned on passing because it had been superseded by the standard hashing method, which can be performed without ancillary EPR pairs. Let me call a general EPP that involves ancillary EPR pairs a *generalized* breeding protocol/method. In contrast to prior art, here I notice that the generalized breeding protocol is, in general, *not* reducible to a non-breeding protocol. In fact, it is more powerful because it allows the decoupling of error correction from privacy amplification. In summary, the decoupling of error correction from privacy amplification is achieved at the price of introducing ancillary EPR pairs shared by Alice and Bob.

I remark that the calculation of  $M_i^A \otimes Z_{A'_i}$  (and similarly  $M_i^B \otimes Z_{B'_i}$ ) in Theorem 6 can, indeed, be done by local quantum gates. The actual quantum circuit diagram is very similar to the ones discussed in for example, [3] and [4]. Since the actual construction is outside the main theme of this paper, the details will be skipped here.

## B. Reduction to BB84

Using ancillary EPR pairs in a generalized breeding protocol, the above subsection shows that one can decouple error correction from privacy amplification in a QKD scheme. However, such a scheme generally requires a quantum computer to implement. So, the next question is: how to reduce the above protocol to standard BB84? Here is the second main Theorem of the current paper.

**Theorem 7** *Suppose a two-way EPP satisfies the following conditions:*

1. *(Symmetric) It can be described as a series of measurements  $M_i$ , with both Alice and Bob measuring the same  $M_i$ .*
2. *(CSS-like) Each of its generators  $M_i$  can be written as either a) a product of  $X$ 's only or b) a product of  $Z$ 's only. Let me call them  $M_X$  and  $M_Z$  operators respectively.*
3. *( $r$ -locally-non-commuting) There exists a set of  $r$   $M_Z$  operators, which, thereafter I shall call the non-commuting set such that, after deleting them from the set of measurements, each pair of operators  $M_j$  and  $M_k$  chosen from the remaining set of measurements commute locally in Alice's (or Bob's) side.*
4. *(Conditional on  $Z$ 's only) All conditional operations depend only on the result of measuring  $Z$  operators.*

Then the protocol can be converted to a standard “prepare-and-measure” QKD scheme with security equal to the EPP-based QKD scheme, provided that Alice and Bob initially share an  $r$ -bit secret string and use it to encode the measurement outcome of  $M_Z$ 's of the non-commuting set in Condition 3.<sup>89</sup>

*Sketch of Proof:* Combine the proofs of Theorems 5 and 6. In other words, the proof of Theorem 6 can be used to relax the constraint of local commutability in Theorem 5, thus giving Theorem 7.

We have the following Corollary:

**Corollary 8** *Consider the purification of  $N$  impure EPR pairs. Suppose one is given a symmetric stabilizer-based bit-flip (interactive or one-way) error correction procedure with  $s$  operators  $M_Z$ 's and also a symmetric stabilizer-based phase error correction procedure with  $t$  operators  $M_X$ 's acting on the  $N$  pairs.*

---

<sup>8</sup>Note that the same key is used to encode the measurement outcomes in both Alice and Bob's sides. This is because the relative error syndrome is *allowed* to be disclosed to Eve.

<sup>9</sup>Note that the final key is now a coset of  $C_2$  in  $F_2^n$ , whereas in Shor-Prekill's proof, the key is a coset of  $C_2$  in  $C_1$ . The difference is due to the fact that, in Theorem 9, an ancillary secret is sacrificed. The *net* key generation rate is the same if Theorem 9 is applied in lieu of Shor-Prekill's proof.

*Claim: The combined error correction/privacy amplification protocol can be reduced to a standard prepare-and-measure QKD protocol, provided that Alice and Bob initially share an  $s$ -bit secret string. Having sacrificed the initial  $s$ -bit string, the output of the procedure is an  $N - t$ -bit secret string.<sup>10</sup>*

*Remark:* As an application of the above Corollary, the following protocol for error correction/privacy amplification of QKD is unconditional secure: Step 1: the Cascade scheme for error correction, modified by the one-time-pad encryption of its bit-flip error syndrome, followed by Step 2: a random hashing procedure [6,8]. Notice that this is a rather efficient protocol in terms of both the key generation rate and computational power.

For schemes involving concatenation, there is the following Corollary:

**Corollary 9** *Suppose an EPP,  $\mathcal{C}$  is a concatenation of two subroutines,  $S_1$  and  $S_2$ , where the first subroutine,  $S_1$  satisfies all the conditions in Theorem 5 (i.e., symmetric, CSS-like, locally-commuting and conditional on  $Z$ 's only) and the second subroutine,  $S_2$  satisfies Theorem 9 as an  $r$ -locally-noncommuting (symmetric, CSS-like, conditional on  $Z$ 's only) EPP. Then, the protocol  $\mathcal{C}$  can be converted to a prepare-and-measure QKD protocol with the same security, provided that Alice and Bob initially share an  $r$ -bit secret string and use it for one-time-pad encryption of the measurement outcomes of the  $r$  pairs<sup>11</sup> of measurement outcomes in the non-commuting set.*

The upshot of the above Corollary is that the decoupling result remains valid even when there are two way classical communications [13] and even when concatenated codes are employed.

## V. CONCLUDING REMARKS

In summary, I have considered a rather general class of entanglement purification schemes, more specifically, symmetric, stabilizer-based schemes and their reduction to BB84. It was shown that in those schemes, the procedure for error correction can be decoupled from the procedure for privacy amplification. The decoupling is achieved by requiring Alice and Bob to share a modest initial string and use it for the one-time-pad encryption of the bit-flip error syndrome. This is no change in the net key generation rate because the loss of this initial string will be exactly compensated by the generation of a longer key. (See footnotes 8 and 9.) As a corollary, I prove the security of the Cascade scheme, modified by one-time-pad encryption of error syndrome, followed by a random hashing privacy amplification procedure. This is an efficient scheme in terms of both key generation rate and computational power.

---

<sup>10</sup>See footnote 9.

<sup>11</sup>See footnote 8.

## VI. ACKNOWLEDGEMENT

I particularly thank Norbert Lütkenhaus for bringing to my attention the question of the security proof of the Cascade scheme and for many enlightening discussions. Helpful conversations with colleagues including Daniel Gottesman, Tsz-Mei Ko, John Preskill and Peter Shor are also gratefully acknowledged.

## REFERENCES

- [1] C. H. Bennett and G. Brassard, Quantum cryptography: Public key distribution and coin tossing, *Proceedings of IEEE International Conference on Computers, Systems, and Signal Processing*, IEEE, 1984, pp. 175-179.
- [2] A. K. Ekert, Quantum cryptography based on Bell's theorem, *Phys. Rev. Lett.* vol. 67, (1991), pp. 661-663.
- [3] C. H. Bennett, D. P. DiVincenzo, J. A. Smolin, and W. K. Wootters, *Phys. Rev. A*, vol. 54, (1996), pp. 3824-.
- [4] C. H. Bennett, G. Brassard, S. Popescu, B. Schumacher, J. A. Smolin, and W. K. Wootters, *Phys. Rev. Lett.* vol. 76, (1996), pp. 722-.
- [5] D. Deutsch, A. Ekert, R. Jozsa, C. Macchiavello, S. Popescu, and A. Sanpera, *Phys. Rev. Lett.* **77**, 2818 (1996); **80**, 2022 (1998) (errata).
- [6] D. Mayers, preprint <http://xxx.lanl.gov/abs/quant-ph/9802025>, to appear in J. of Assoc. Comp. Mach. A preliminary version in D. Mayers, Quantum key distribution and string oblivious transfer in noisy channel, *Advances in Cryptology — Proceedings of Crypto' 96*, Lecture Notes in Computer Science, vol. 1109, Springer-Verlag, 1996, pp. 343-357.
- [7] H.-K. Lo and H. F. Chau, Unconditional security of quantum key distribution over arbitrarily long distances, *Science*, vol. 283, (1999), pp. 2050-2056; also available at <http://xxx.lanl.gov/abs/quant-ph/9803006>.
- [8] P. W. Shor and J. Preskill, Simple proof of security of the BB84 quantum key distribution scheme, *Phys. Rev. Lett.* vol. 85, (2000), pp. 441-444.
- [9] E. Biham, M. Boyer, P. O. Boykin, T. Mor, and V. Roychowdhury, A proof of security of quantum key distribution, *Proceedings of the Thirty-Second Annual ACM Symposium on Theory of Computing*, ACM Press, New York, 2000, pp. 715- .
- [10] M. Ben-Or, to appear.
- [11] H. Inamori, N. Lütkenhaus, and D. Mayers, Unconditional Security of Practical Quantum Key Distribution, available at <http://xxx.lanl.gov/abs/quant-ph/0107017>
- [12] G. Brassard and L. Salvail, Secret key reconciliation by public discussion, *Advances in Cryptology, Eurocrypt' 93 Proceedings*, 1993, pp. 410-423.
- [13] D. Gottesman and H.-K. Lo, Proof of security of quantum key distribution with two-way classical communications, available at <http://xxx.lanl.gov/abs/quant-ph/0105121>
- [14] H.-K. Lo, Proof of unconditional security of six-state quantum key distribution scheme, *QIC* (Quantum Information and Computation), vol. 1, No. 2, (2001), pp. 81-94.
- [15] C. H. Bennett, F. Bessette, G. Brassard, L. Salvail, and J. Smolin, Experimental quantum cryptography, *J. Cryptol.* vol. 5 (1992), pp.3-28.
- [16] C. H. Bennett, G. Brassard, C. Crépeau, and U. M. Maurer, Generalized Privacy Amplification, *IEEE Transactions on Information Theory*, vol. IT-41, no. 6, (1995), pp. 1915-1923.
- [17] D. Bruss, Optimal eavesdropping in quantum cryptography with six states, *Phys. Rev. Lett.*, vol. 81, (1998), pp. 3018-3021.
- [18] H.-K. Lo, H. F. Chau and M. Ardehali, Efficient quantum key distribution scheme and proof of its unconditional security, available at <http://xxx.lanl.gov/abs/quant-ph/0011056>

- [19] D. Gottesman and J. Preskill, Secure quantum key distribution using squeezed states, available at <http://xxx.lanl.gov/abs/quant-ph/0008046>
- [20] A. R. Calderbank and P. Shor, Good quantum error correcting codes exist, *Phys. Rev. A*, Vol. 54, (1996), 1098-1105.
- [21] A. M. Steane, Multiple particle interference and error correction, *Proc. R. Soc. London A*, vol. 452, (1996), pp. 2551-2577.
- [22] C. Fuchs, N. Gisin, R. B. Griffiths, C. S. Niu, and A. Peres, "Optimal eavesdropping in quantum cryptography. I," *Phys. Rev.*, vol. A56, p. 1163, 1997.
- [23] I. Cirac and N. Gisin, "Coherent eavesdropping strategies for the 4-state quantum cryptography protocol," *Phys. Lett.*, vol. A229, p. 1, 1997.