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MEMORANDUM REPORT NO. 2796

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(Supersedes IMR No. 233)

MISSILE WARHEAD MODELING: COMPUTATIONS

AND EXPERIMENTS

William W. Predebon Walter G. Smothers Charles E. Anderson

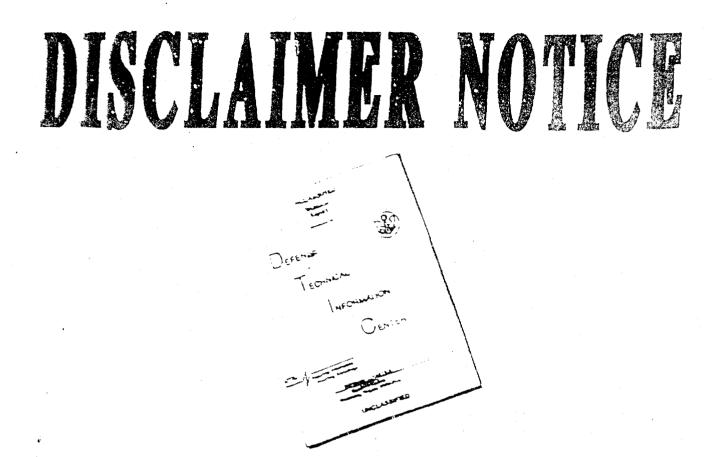
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October 1977

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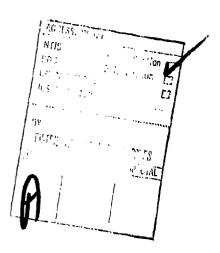
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## I. INTRODUCTION

Missilè warheads are generally quite different in design and construction than warheads used in other military applications, such as artillery shells. This difference mainly is a consequence of the low launch loads experienced by missile warheads. Warhead design in applications where high launch loads are experienced (artillery shells) are usually of the natural fragmentation or controlled fragmentation type. By way of background, natural fragmentation warheads are continuous motal casing warheads with no preselected control sites of failure. Controlled fragmentation warheads, on the other hand, are continuous warheads with preselected control sites or lines of failure. These failure sites are usually arrived at mechanically, metallurgically or explosively. However, missile warheads are generally discrete-fragment warheads and not continuous metal-casing systems. Discrete-fragment warheads are preformed-fragment warheads in which the fragments of the desired shape and mass are placed on a liner (motal or plastic) which encases the high explosive. The warhead is then placed inside a missile casing (skin). These proformed-fragments are attached to the liner and neighboring fragments by an epoxy filler. Due to the very low launch load environment of a missile, this evidently is sufficient for structural integrity.

In the past the main analytical method for predicting fragmentation warhead performance\*, i.e., fragment speed and direction, has been the Gurney<sup>1</sup> and Taylor<sup>2</sup> formulae. These formulae are one-dimensional by assumption and the Gurney formula requires an empirically determined constant. This constant, called the Gurney constant, is obtained experimentally via warhead tests for each different situation. Due to the one-dimensional nature of these formulae, their utility for predicting warhead performance near the ends is questionable. It has been

 $\mathrm{shown}^3$  that for cases when the flow is essentially one-dimensional, such

\*Warhead performance as used in this report will mean fragment velocity distribution, which, it can be shown, yields fragment mass distribution.

<sup>1</sup>Gurney, R. W., "The Initial Velocities of Fragments from Bombs, Shells and Grenades", BRL Report 405, September 1943, AD#36218.

<sup>2</sup>Birkhoff, G., MacDougall, P., Pugh, E. M. and Taylor, G. I., Sir., "Fxplosives with Lined Cavities", J. Appl. Phys., Vol. 19, p. 563, June 1948.

<sup>3</sup>Karpp, R. R. and Predebon, W. W., "Calculations of Fragment Velocities from Naturally Fragmenting Munitions", BRL Memorandum Report No. 2509, July 1975. (AD #B007377L)

as long artillery projectiles (L/D > 2), these formulae are adequate for predicting fragment velocity. However, for cases where the flow is two-dimensional, such as near the warhead ends for any size projectile, and for axisymmetric warheads with length to diameter ratios less than or equal to two, these formulae are inadequate<sup>3</sup>. Lastly, these formulae were initially derived for continuous warheads (i.e., natural fragmentation warheads) and for one specific initiation posture. Consequently, for discrete-fragment warheads, these formulae are not applicable without significant modification and experimental recalibration.

In an effort to find a method to predict warhead performance not only over its midsection but over its entire length including the ends, a time-dependent, two-dimensional, Lagrangian computer code has been applied to naturally fragmenting warheads. In Reference 3, finitedifference calculations of warheads with varying explosive fills and casing materials were compared with experimental results; good agreement was demonstrated. However, the computer code used to obtain the calculations in Reference 3 assumes a continuum. Warheads eventually fragment, and after fragmentation the continuum assumption needs to be modified. A method to model the effects of fragmentation and subsequent explosive gas leakage was presented in Reference 3 as an integral part of the computational scheme.

In this report the earlier work by Karpp and Predebon<sup>3</sup> on natural fragmenting warhead is extended to missile warheads, i.e., discretefragment warheads. The objective of the work reported here was to develop a warhead modeling capability to predict the performance of missile warheads in general, as well as warhead designs pertinent to the "Focused-Blast-Fragment" advanced development program.\* Computational modeling was emphasized with limited experimental verification as needed.

In the next section the experimental program is discussed in detail. In Section III the computational modeling is explained. Comparisons between the computational and experimental results are discussed in Section IV. The conclusions and recommendations follow in Section V.

\*This work was supported by the BRL, Picatinny Arsenal and AMC, HQ.

### 11. EXPERIMENTAL PROGRAM

Figure 1 is a schematic drawing of part of a typical test set-up used in the warhead test firings for this program. Shown is the warhead on an angled\* wood stand which rests on metal plates on the floor of the second level of the test facility. The two metal plates form a preset aperture through which passes the column of fragments to be radiographed. The lower level of the test facility contains two orthogonal banks (each bank has four flash x-ray tubes) which provide time-sequential, orthogonal radiography.

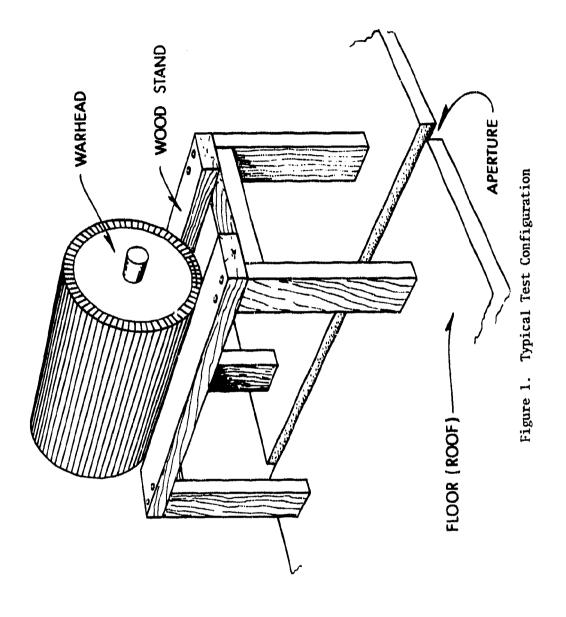
Due to the funding and time constraints on part of this program, it became obvious at the very beginning that cubical fragments\*\* could not be placed completely around each test warhead and still complete the minimum number of warhead tests planned. Also, from a long term point of view, it is desirable to have a more economical method for testing preformed-fragment warheads. The method used determined the minimum azimuthal-sector angle containing preformed-fragments which adequately represented the situation where the entire cylindrical warhead contained preformed-fragments. All the warhoads tested had an overall cylindrical shape. Figure 2 illustrates a test warhead with an azimuthal fragmentbearing sector angle of 22.5° for a typical length to diameter (L/D)ratio of 2.0. In all tests where an azimuthal sector of the warhead contained the preformed-fragmonts, the remainder of the warhead contained strips (rods) of the same material. The width of the strips was equal to the width of the fragments (for proper radial explosive gas leakage), and the strip lengths were equal to the explosive charge length.

Table I contains all the characteristics of the test warheads by their respective round number. As indicated in Table I, the warhead length to diameter ratio was varied from 1.0 to 2.0. The explosive charge in all cases was Octol (25% TNT, 75% HMX). The fragments were steel (AISI 1018 or 1020) cubes and, in all but one test, were placed directly on the bare explosive charge. This type of configuration was adopted in order to determine with greater accuracy the effect which explosive gas leakage between fragments has on the fragment velocity distribution, particularly for the computational modeling portion of the program.

Figures 3 and 4 illustrate typical radiographic results from 1/D = 1.0 and 1/D = 2.0 warheads, respectively. Both test warheads were initiated on the axis at the left end. Figures 3 and 4 show the longitudinal and orthogonal views of the fragment distributions in their actual spacial positions at two different times. In each of the test cases the fragments in the radiographed column were premarked for identification and were recovered in Celotex. The spacial location of the recovered premarked

\*\*In all the warhead tests the fragment shape was oubical.

<sup>\*</sup>The angled wood stand is used to rotate the main fragment distribution into the available radiographic view.



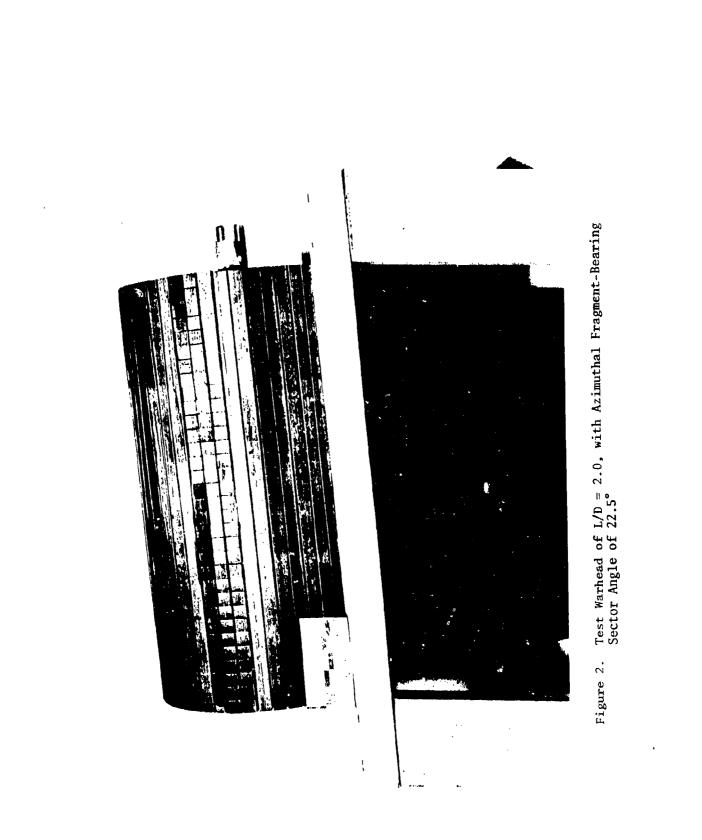


TABLE I. WARHALD CHARACTERISTICS AND TEST DATA

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	ł	9.5A	9,37	9.75	67-6	9.53	9.95	9.72	6. 6	9.55	4.67	4.75	9.97	
 *	ē	3.64	3.77	3.65	3.69	3.67	3.72	3.80	3.78	3.70	0.455	0.464	3.80	
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R <sup>1</sup>	ю	;	60. R	61.4	122.7	122.0	123.0	124.8	124.8	123.9	62.3	63.7	127.8	
Ē	ē	3.84	3.87	3.84	3.84	3.81	3.87	<u>5.91</u>	3.92	3.89	0.486	0.485	3,93	
ر ت	Ĭ	7.89	7.92	7.59	7.89	7.87	7.92	7.95	7.95	7.93	3.96	3.96	7.96	
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<sup>1</sup>These data represent the average value of the variable.

 $^{\sf Z}_{\sf RC}$  represents the measured cube dimension in the hoop direction for the recovered cubes.

 ${}^3 R_{\rm h}/R_0$  is the expansion ratio at fragment separation which is equivalent to  $R_{\rm NC} L_{\rm C}$ 

<sup>4</sup>SEA refers to single end axial initiation and "MA refers to dual end axial initiation.

<sup>5</sup>The fragment layers in this test were aligned (See Section II).

<sup>6</sup>Inc fragment layers in this test were hricked (See Section [1].

<sup>7</sup> The liner was 606J-T4 aluminum with a nominal walf thickness of L.60 mm and rotal mass of 441 (Mr. The endplates were 2024-F5 aluminum with a nominal thickness of 6.35 mm and with total masses of 275 (M (detenator endplate) and 279 (M for the other end plate.

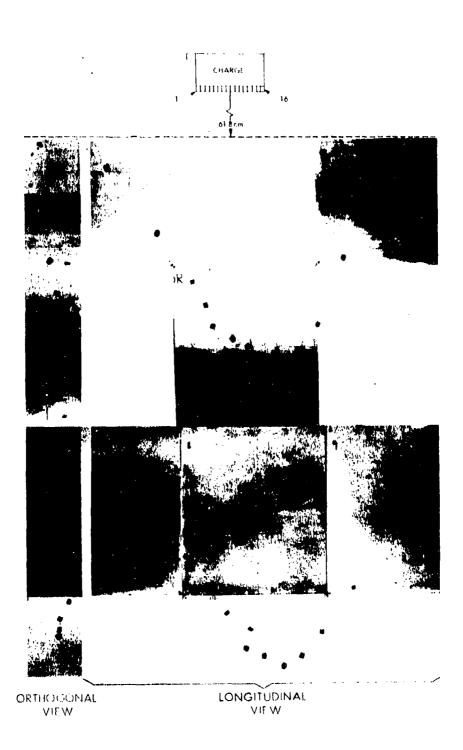
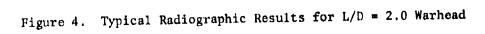


Figure 3. Typical Radiographic Results for L/D = 1.0 Warhead



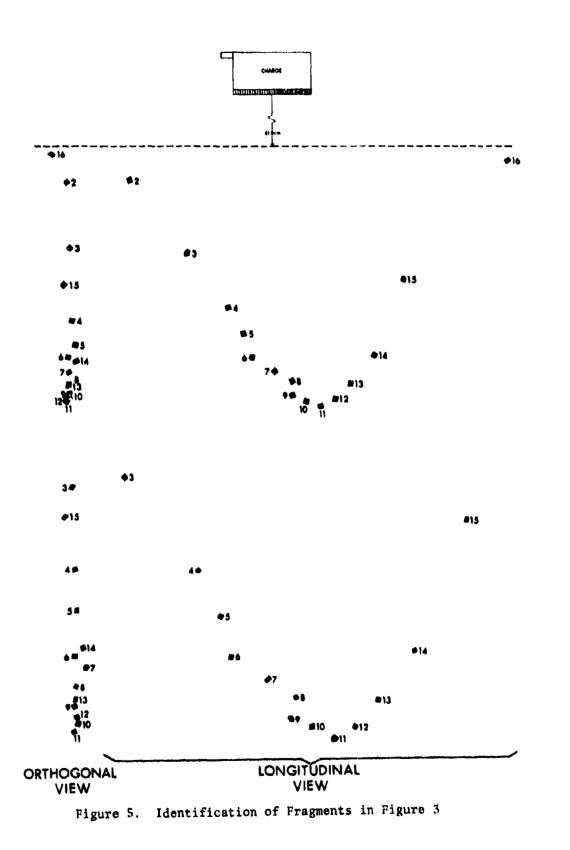


fragments, in conjunction with the time-sequential, orthogonal radiography of the fragments, permitted development of the results shown in Figures 5 and 6 for the two tests illustrated in Figures 3 and 4. In Figures 5 and 6, each of the fragments in flight are labeled according to their original position on the warhead (shown at the top of the figures). Two different times are shown for both longitudinal and orthogonal views.

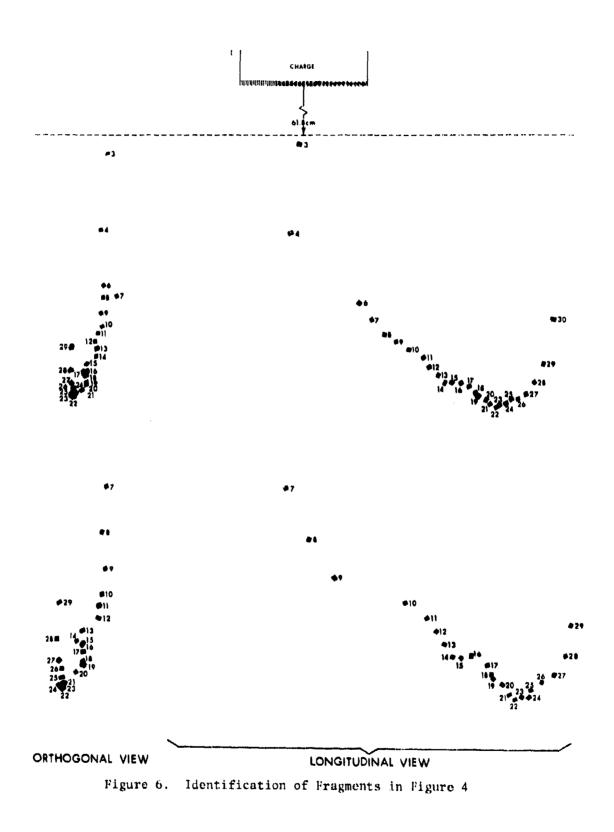
To illustrate radiographically the fragment focusing that can be achieved, Figure 7 shows (for two different times) the longitudinal and orthogonal views of an L/D = 2.0 warhead which was initiated simultaneously at both ends on the axis of symmetry. The fragments of Figure 7 are labeled in Figure 8 according to their original position on the warhead. Again two different times are shown for the longitudinal and orthogonal views.

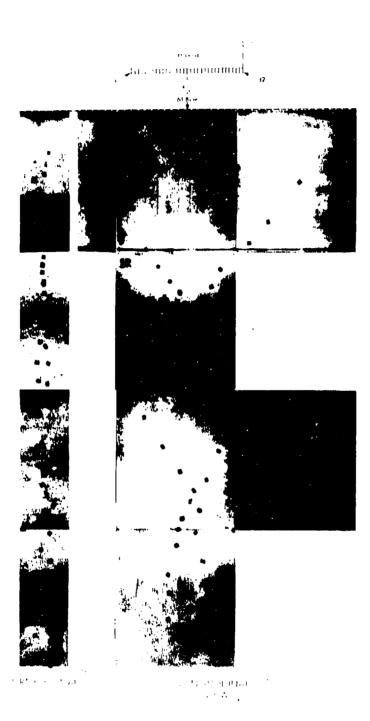
To determine the accuracy of the experimental procedure, a simple error analysis of the measurements was conducted and the round to round variation was quantified. Using a conservative approach to error analysis wherein only maximum differences are considered, for two timesequential observations of fragments the measurement uncertainty in the fragment speed is less than  $\pm$  1% and is less than  $\pm$  0.004 for the tangent of the fragment projection angle. The round to round variation in the warhoad tests is shown in Figure 9. Plotted in Figure 9 is the fragment speed and projection angle distributions versus fragment number or relative initial axial position from two tests of nominally identical (See Table I)  $L/D \approx 2.0$  cylindrical warhcads. The round to round variation in fragment speed is  $\pm$  3% for two time-sequential observations of the fragments and  $\pm$  13% for single time observations of fragments (which occurs for a few fragments at both ends). The overall round to round variation in the tangent of the fragment projection angle is  $\pm$  0.02. Therefore, it appears that the variations from one round to another (nominally the same) are larger than the measurement uncertainty. Consequently, it was inadvisable to attempt to improve the accuracy of the measurement procedure.

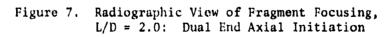
As previously mentioned, the cubical fragments were placed on the warhead in varying azimuthal sector angles, with strips in the remaining azimuthal angle, to determine the minimum azimuthal fragment-bearing sector angle required to approximate the situation in which fragments covered the entire round. Plotted in Figures 10 and 11 are the fragment speed and projection angle distributions for varying azimuthal fragmentbearing sector angles  $\alpha$  for L/D = 1.0 and 2.0 warheads, respectively. Since the percent difference between the  $\alpha = 45^{\circ}$  and  $\alpha = 90^{\circ}$  cases in Figure 10 is within round to round variation shown in Figure 9, it is felt that a 45° azimuthal fragment-bearing sector angle is adequate for the L/D = 1.0 case. For the L/D = 2.0 case, Figure 11 indicates a large difference between  $\alpha = 22.5^{\circ}$  and  $\alpha = 45^{\circ}$ , 90° and 135°. However the differences are much smaller between the  $\alpha = 45^{\circ}$ , 90° and 135° cases, although the fragment speeds for  $\alpha = 45^{\circ}$  tend to be greater than or

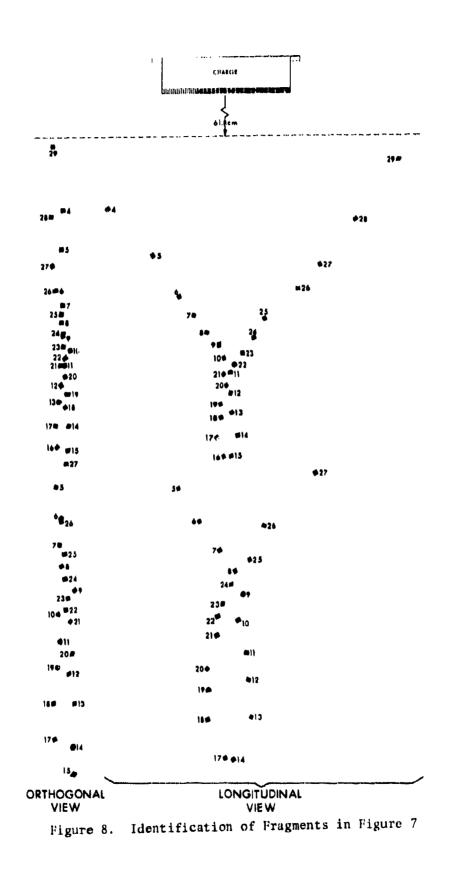


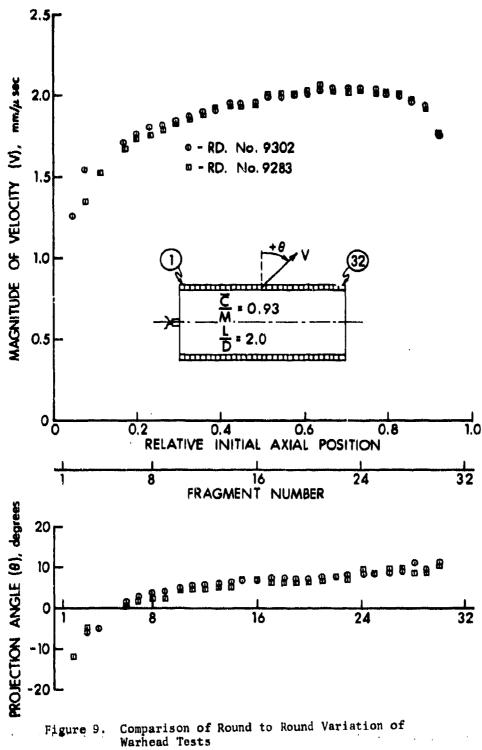




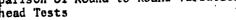




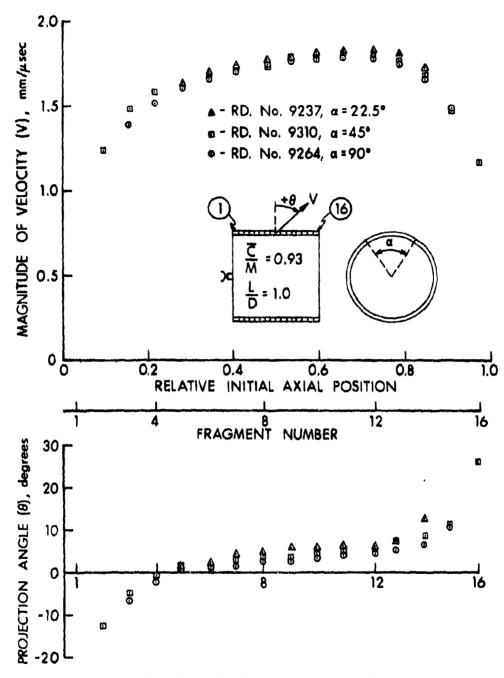


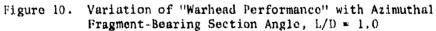


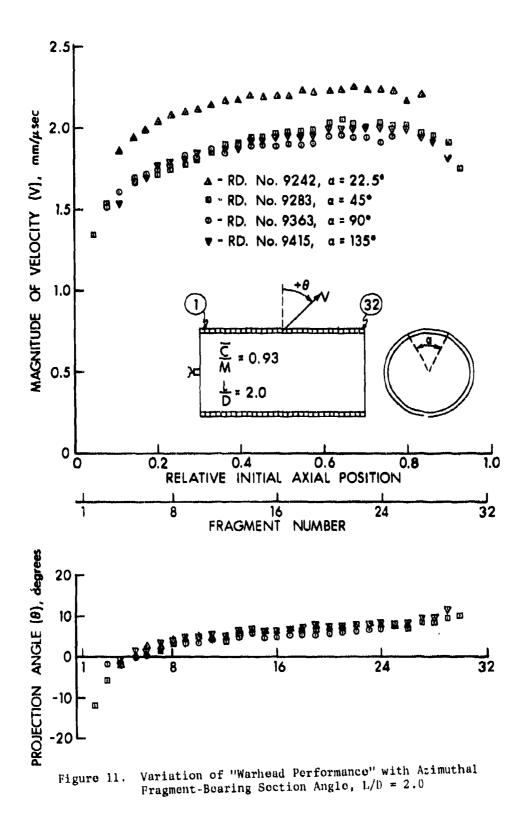










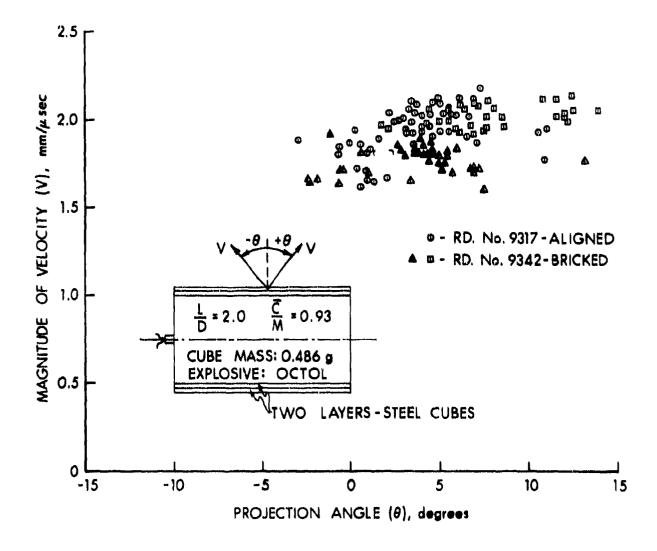


equal to  $\alpha = 90^{\circ}$  or  $\alpha = 135^{\circ}$ , especially in areas remote from end rarefactions. Mindful of the latter point and the round to round variation shown in Figure 9, the variation in fragment speed between the  $\alpha = 90^{\circ}$ and 135° is within the round to round variation; whereas the same variation between the  $\alpha = 45^{\circ}$  and 90° is generally not. In view of the above it is felt that  $\alpha = 90^{\circ}$  is desirable for the L/D = 2.0 case.

Returning to the  $\alpha = 22.5^{\circ}$  case in Figure 11, due to the unexpected large differences between it and  $\alpha = 45^{\circ}$ , 90° and 135°, particularly near the ends where one would expect the fragment speeds to converge to a common value for most practical  $\alpha$ 's, it is questionable whether these large differences are due to increased confinement for  $\alpha = 22.5^{\circ}$  or perhaps some uncovered error in the data. An additional test at  $\alpha =$  $22.5^{\circ}$  will be required to clarify these questions; however, it was judged that the results of such a test would not significantly change the main thrust and purpose of this study and therefore was not completed at this time.

The experimental results of cases in which computational results were obtained are discussed in Section IV. There are a few test cases for which the computations were not performed during this program and these experimental results will be presented at this time.

The modeling of multi-fragment layering is immensely complicated due to the many combinations of layer patterns that can be placed on the warhead. In order to bracket the effects of the simplest multi-fragment layering case, i.e., two layers of fragments, two double-layered discretefragment warheads were fired. The fragment layers were aligned with respect to each other for the first warhead; the other warhead had the fragments bricked (that is, the fragments on the top layer were displaced by a distance equal to half of a cube side relative to the fragments on the bottom layer). In Figure 12 are plotted the results of the aligned and bricked double-layered warheads. In order to compare these results to the single layered test, the charge mass to metal mass (C/M), the type of explosive, explosive initiation and explosive dimensions were kept the same as the single layered L/D = 2.0 case (shown in Figure 16). Further, the fragment shape and fragment material were kept the same. These constraints thus imposed that the double-layered fragment masses be 1/8 the single layered masses. Due to the resulting lower cube masses (0.486 gm), and therefore much smaller dimensions, the previously used cube numbering system was not effective in identifying the cubes with respect to their original position on the warhead. When the cubes were recovered they had generally collided with each other during flight and the punched numbers on each fragment were unidentifiable. Consequently the fragment speed is plotted with respect to its projection angle in Figure 12 rather than its initial position on the warhead.



# Figure 12. Double-Layered Discrete-Fragment Warheads: Fragment Speed Versus Projection Angle

In Figure 12 it is quite evident that there exists a significant variation of the magnitude of the velocity at any projection angle; this observation applies both to the aligned and bricked warheads. By comparison, if the data from the single fragment layer case of Figure 16 were plotted versus projection angle, as in Figure 12, only one fragment would appear at any particular projection angle. It is important to note, for reasons previously discussed, it was very difficult to identify the same fragment in two time-sequential, radiographic views for many of the fragments in the double-layered cases. Consequently, the only fragments that are plotted are those in which this identification was confidently made. Therefore, although the single fragment layer case cannot be compared with the double-layered case over the entire projection angle distribution, meaningful comparisons can be made over the projection angle range for which data were obtained. In the two examples depicted in Figure 12, the spread in the speeds ranges from 0.15 to 0.40 mm/usec over the projection angle distribution. The most probable contributing cause of this spread in speed at a given projection angle is slippage between the layer fragments. Thus, instead of fragments leaving the warhead at a given angle having one particular velocity, collective effects result in a number of fragments with varying velocities at the specified angle.

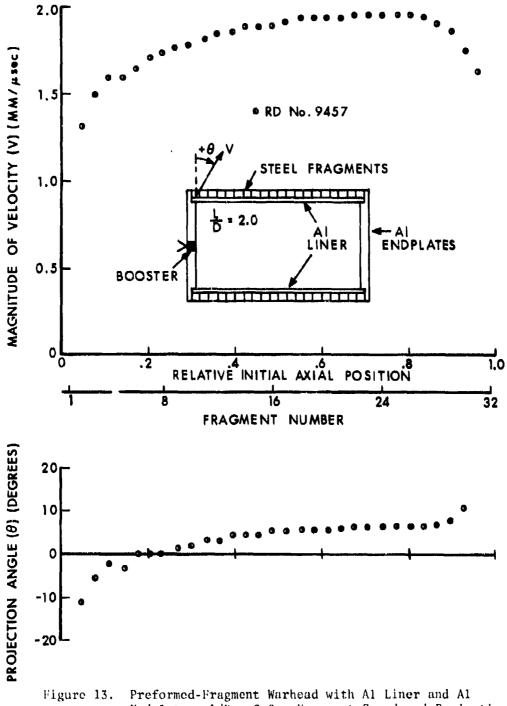
It is lastly noted that the plotted speeds in Figure 12 have not been corrected for drag, which from the data was estimated to be causing a 4%-7% decrease in fragment speeds. This decrease was obtained by comparing single and double flash calculations of identical cubes. This calculation was also completed in the single layered cases; however, in these cases the drag was always negligible. Because of the lower cube masses (0.486 gm) in the double-layered warheads, drag is no longer negligible.

Shown in Figure 13 are the experimental results for single-layered fragments on an aluminum liner with aluminum endplates, L/D = 2.0 cylindrical warhead. As shown, initiation was on the axis on the left side of the warhead. All the necessary data are given in Table I. This is the only test case reported with a liner and endplates. These results will be discussed in some detail in Section IV along with a similar test case without a liner and endplates.

### III. COMPUTATIONAL MODELING

The computer code that is being used for the computation modeling is a modified version of the HEMP<sup>4</sup> computer code. The HEMP code, which was originally developed at the Lawrence Livermore Laboratories, California, is a time-dependent, two-dimensional, Lagrangian finite-difference code.

<sup>&</sup>quot;Wilkins, M. L., "Calculations of Elastic-Plastic Flow," Methods of Computational Physics, Vol. 3, edited by Alder, B., Fernback, S. and Rotenburg, M., Academic Press, NY, 1964.



Lgure 13. Preformed-Fragment Warhead with Al Liner and Al Endplates, L/D = 2.0: Fragment Speed and Projection Angle Distribution

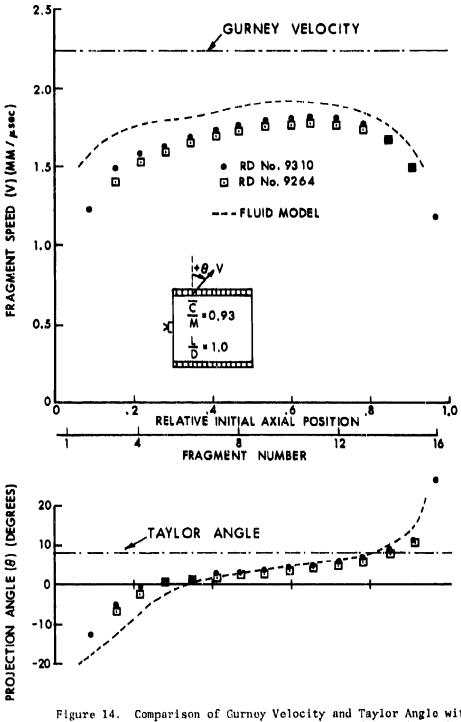
If the standard form of the HEMP code is used to predict the performance of discrote-fragment warhcads without including the effects of fragment separation and subsequent explosive gas leakage between fragments, the dashed line results shown in Figure 14 are obtained. Shown in Figure 14 is a comparison between these HEMP calculations and experimental data of fragment speed and projection angle distributions for an L/D = 1.0 cylindrical warhead. In the HEMP calculations the fragments are modeled as a fluid, i.e., the fragments cannot accommodate a tensile force but they can be recompressed. This would be a logical approach with a continuum code since the fragments are not a continuum but discrete particles which will separate with the application of a small tensile force. It is obvious that this modeling is superior in predictive ability than the prediction of velocity and angle by Gurney and Taylor, also shown in Figure 14. However, the agreement between the HEMP computations and experiment is not as good as desired. Because of this and other objections<sup>5</sup>, modeling of the effects of fragment separation and explosive gas leakage between fragments was initiated to improve the overall predictive capability of the computations.

The discrete-fragment warhead (missile warhead) model presented involves the following three essential features: A) modeling of the discrete-fragments, B) modeling of fragment separation and subsequent explosive gas leakage and, C) criteria for fragment separation and method to determine this criteria.

### A. Modeling of Discrete-Fragments

When discrete-fragments are placed on a cylindrical body, whether it be the bare explosive or a liner containing the explosive, (assuming conventional fragment shapes) air gaps will occur between the fragments in the circumferential direction. Consequently, the discrete-fragment casing forms a composite of metal and air. In reality, this is a threedimensional system. However, the approach that was adopted with the twodimensional modeling is similar to the methods used in the modeling of composite materials. The discrete metal-air composite casing (metal preformed-fragments) is modeled as a homogeneous continuous material of the same metal material as the experimental fragments, but with the requirement that the total mass of the new continuous metal casing equal the total miss of all the fragments on the experimental warhead. For example, in the case of cubical fragments, the resulting continuous metal casing wall thickness in the model will be less than one of the sides of

<sup>&</sup>lt;sup>5</sup>Karpp, R. R. and Predebon, W. W., "Calculations of Fragment Velocities from Fragmentation Munitions," First International Symposium on Ballistics, Orlando, Florida, 13-15 November 1974, Section IV, pp. 145-176, Proceedings Published by American Defense Preparedness Association, Union Trust Building, Washington, DC.



ure 14. Comparison of Gurney Velocity and Taylor Angle with HEMP (Fluid Model) Computations and Experimental Data, L/D = 1.0 Cylindrical Preformed.Fragment Warhead

an actual cubical fragment. For a cylindrical warhead, the modeled casing outside diameter for a given explosive diameter is obtained from

$$M_{T} = \sum_{i} m_{i} = p_{f} \left\{ \frac{\pi}{4} L_{W} \left( D_{M}^{2} - D_{HE}^{2} \right) \right\},$$

where  $M_{\rm T}$  is the measured total fragment mass (plus liner mass if the warhoad has a liner),  $m_{\rm i}$  are the measured individual fragment masses,  $\rho_{\rm f}$  is the fragment density,  $L_{\rm W}$  is the total length of the fragments on the warhead,  $D_{\rm HE}$  is the measured diameter of the explosive and  $D_{\rm M}$  is the outside diameter of the continuous metal easing to be computed. If one ignores the air gap between the fragments and uses the fragment cube dimension as the continuous easing wall thickness, then for the cases reported in Sections 11 and 1V, the modeled total mass will be as much as 6% greater than the actual total mass.

#### B. Fragment Separation and Explosive Gas Leakage

The general ideas of the model for fragment separation and explosive gas leakage will be presented here; however, the specific details and equations have been presented in an earlier report<sup>5</sup>.

The HEMP code models continuum behavior in general. However, in reality, when fragment separation occurs a continuous casing circumference ceases to increase. But the standard form of the code causes the casing to remain continuous. This treatment tends to yield a higher radial acceleration and thus a higher final fragmont velocity. This can be corrected by computing the true force acting on an element of the casing. The true force is computed in the code by reducing the prossure acting on the interface of the casing by the factor  $R_{\rm b}/R$ , where  $R_{\rm b}$  is the radius at fragment separation and R is the current radius. At present,  $R_{\rm b}$  must be determined by experiment. After fragment separation, i.e.,  $R > R_{\rm b}$ , a void appears between fragments and gas leakage occurs. This effect is modeled in the code by calculating the rate of offlux from an ideal nozzle passing its maximum rate of flow. The rate of efflux is computed in the code for every cycle, and the mass loaked for a single cycle is the product of the mass rate and the cycle time increment. In this manner, the model accounts for gas leakage between fragments after casing breakup. Gas leakage roduces the prossure accordingly. Lastly, to simulate a loss in circumferential stress upon casing breakup, the yield strongth of the fragment material is set equal to zero when  $R > R_{\rm b}$ .

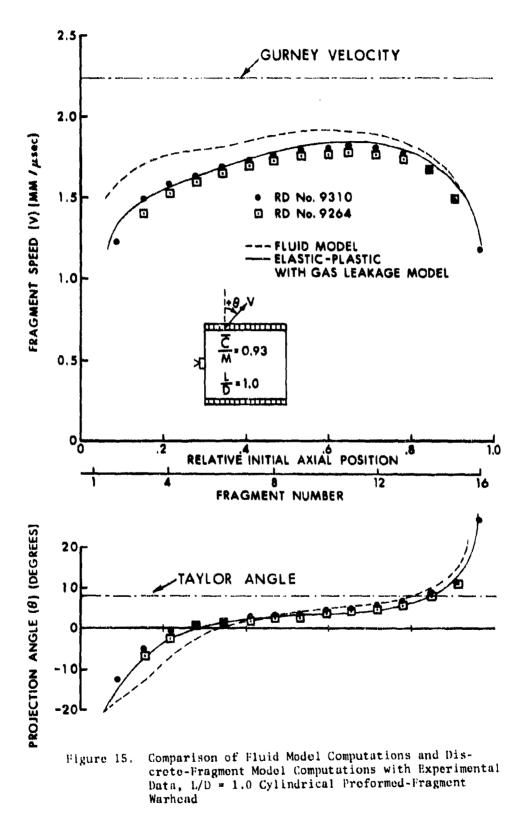
C. Criterion for Fragment Separation

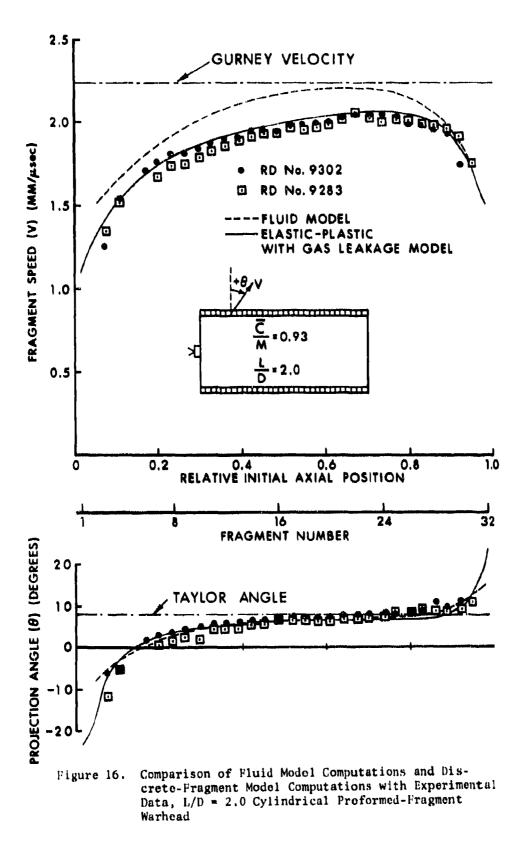
Currently, the criterion for fragment separation,  $R_b$ , is an empirically determined quantity. As discussed in Section II the fragments are all premarked and recovered in a soft recovery medium (Celotex) after firing. The recovered fragments are weighed and measured for their dimensions. Obtained from these measurements is the fragment plastic deformation due to the explosive loading (or final dimensions). Since its original dimensions are known a straight forward calculation determines the corresponding radius of the warhead at which the fragments separated. (The elastic displacements are neglected in this analysis). For the case of cubical fragments on bare explosive for L/D = 1.0 and 2.0 cases, the expansion ratio at fragment separation, i.e.,  $R_b/R_o$  where  $R_o$  is the inside radius of the casing, was found to be 1.2. The recovered fragment measurements and expansion ratios are tabulated in Table I. This criterion for fragment separation is empirical; the criterion will not be applicable in general. Thus, the next refinement in this analysis is to relate this criterion to a fundamental material property.

#### IV. COMPARISON BETWEEN COMPUTATIONAL AND EXPERIMENTAL RESULTS

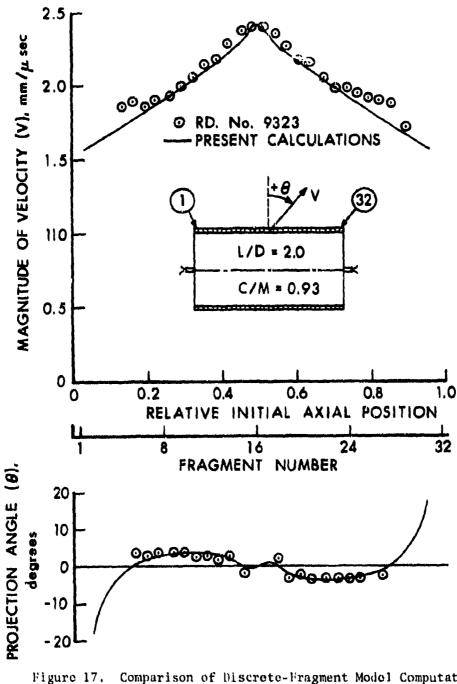
In Figures 15-17 are shown comparisons between the calculations with the discrete-fragment model (labeled Elastic-Plastic with Gas Leakage Model) and experimental results. In Figure 15 and 16 the calculated fragment speed and projection angle distributions are compared with the experimental data for L/D = 1.0 and 2.0 cylindrical warheads, respectively. Data from two test firings are plotted in each figure. In both cases the agreement between computation and experiment is extremely good. Also shown in Figures 15 and 16 are the calculations without the discrete-fragment model, namely, a fluid model, and the Gurney velocity and Taylor angle predictions. Shown in Figure 17 is a comparison between the computations and experimental data for the same case as shown in Figure 16 except with dual-end axial initiation. Again the agreement between computation and experiment is very good. Note that the computations model the region where the two detonation waves collide (higher pressure region) at 0.5 relative initial axial position.

Returning to the test results shown in Figure 13, note that the warhead configuration in Figure 13 is similar to the configuration in Figure 16. Both are L/D = 2.0 cylindrical warheads with the same explosive material, fragment material and size, and with essentially the same explosive diameter. However, the test warhead in Figure 13 has a liner between the explosive and fragments, and endplates. The reasons given for placing the liner between the fragments and the explosive are often varied. It was the intent of this initial experiment, in conjunction with computational modeling (yet to be performed), to help explain the effect of a liner on warhead performance. Due to the additional width of the Al liner, the inside fragment radius increased, which resulted in more fragments on the warhead, and thus a larger total metal weight, and lower C/M. These data are given in Table 1 by Rd. No. 9457. Consequently, the results of Figures 13 and 16 cannot be directly compared. However, it is the authors' judgment at this time, based on these data and other





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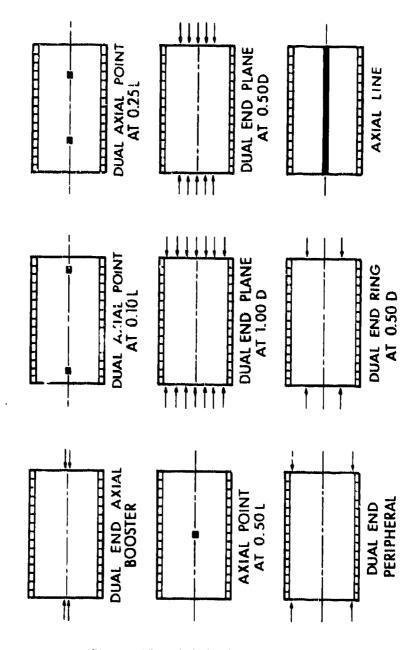
igure 17. Comparison of Discrete-Fragment Model Computations with Experimental Data for L/D = 2.0 Cylindrical Preformed-Fragment Warhead with Dual End Axial Initiation

calculations, that the effect of a liner is to contain the explosive gases for a longer period of time. This in turn causes leakage of explosive gases between fragments to occur at a larger expansion ratio, resulting in a higher final fragment speed, as compared to the case without a liner (assuming all other quantities being equal). This can be illustrated from a plot<sup>6</sup> of calculated fragment speed versus time (from detonation) for these cases. From a plot of this type one observes that fragment separation and leakage occurs during the initial, approximately exponential, acceleration rise. By placing a liner between the fragments and explosive, one can prolong fragment separation and subsequent leakage until the fragment speed versus time profile flattens, which then results in a higher final fragment velocity. From these and other computations, it is estimated that up to approximately a 10% increase in final fragment velocity can be obtained with a liner between the fragments and explosive versus the case without a liner.

llaving demonstrated extremely good agreement between the present calculations and experimental data for both L/D = 1.0 and 2.0 cylindrical warheads with single-end and dual-end axial initiation, the computations were then confidently utilized, without experimental verification, to study the effect of initiation posture on warhead performance. Shown in Figure 18 are the nine initiation postures that were computed. The cylindrical warheads that are drawn in Figure 18 are all with a length (L) to diameter (D) ratio of two and initiated as shown. The first case of dual end axial booster initiation, i.e., initiation with two 19.05 mm diameter explosive boosters, is shown in Figure 17. The computational results of the remaining eight cases are shown respectively in Figures 19-26. Shown in Figure 17 and Figures 19-26 is the effect of a large variety of initiation postures on warhead performance. The general idea of this set of computations is, on the one hand, to demonstrate the versatility and generality of this computational tool, and on the other hand to aid in the understanding of the effect of initiation posture on warhead performance. Furthermore, should the initiation posture of interest lie somewhere between these cases, a simple interpolation of the already computed cases is all that is required for a first approximation of the warhead performance. Similarly, since the computational results are general, i.e., scalable, then if the geometry of interest is not procisely the same as the cases calculated, these results can be scaled up or down assuming proper scaling procedures are followed. Lastly, it is worth while to note that axial line initiation, Figure 26, is the one case of the initiation variations studied in which the fragment projection angle and speed distribution are significantly different from the dual end axial initiation case, Figure 17.

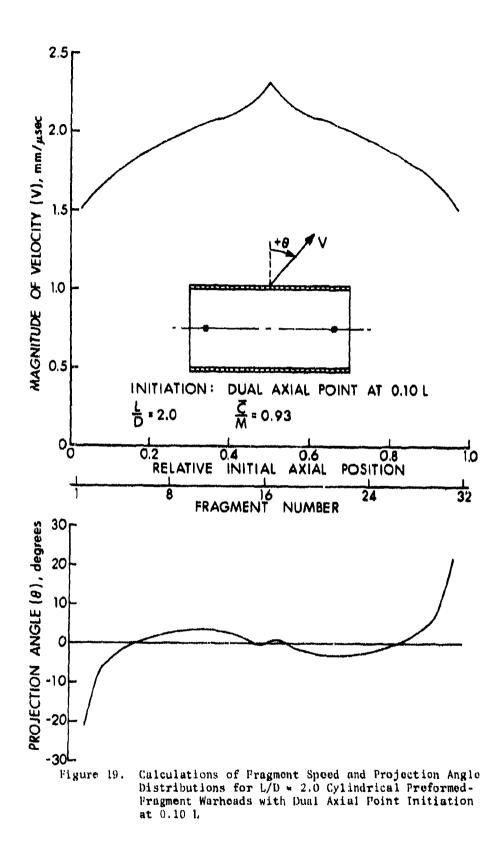
Shown in Figures 27 through 30 are the computational results for two different internal cavity sizes and materials for the same general L/D = 2.0 cylindrical geometry. Shown in Figure 27 are the results for

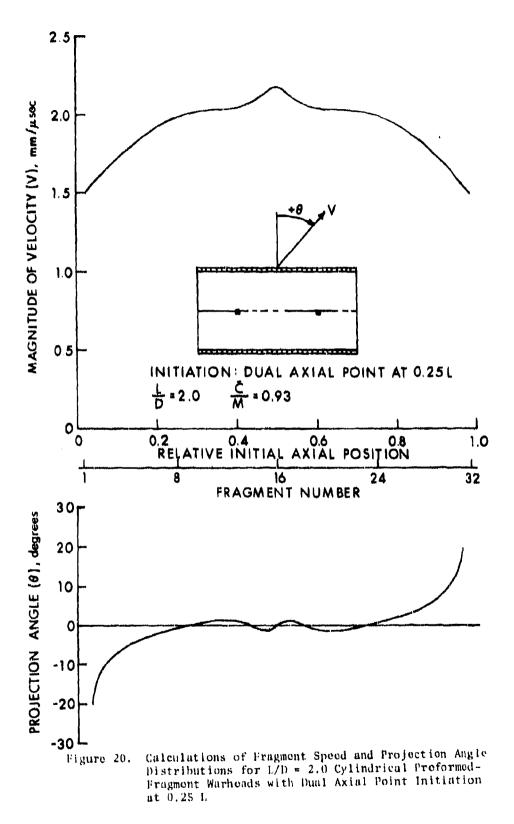
<sup>&</sup>lt;sup>6</sup>Karpp, R. R. and Predebon, W. W., "Calculations of Fragment Velocities from Fragmentation Munitions," p. IV-164.

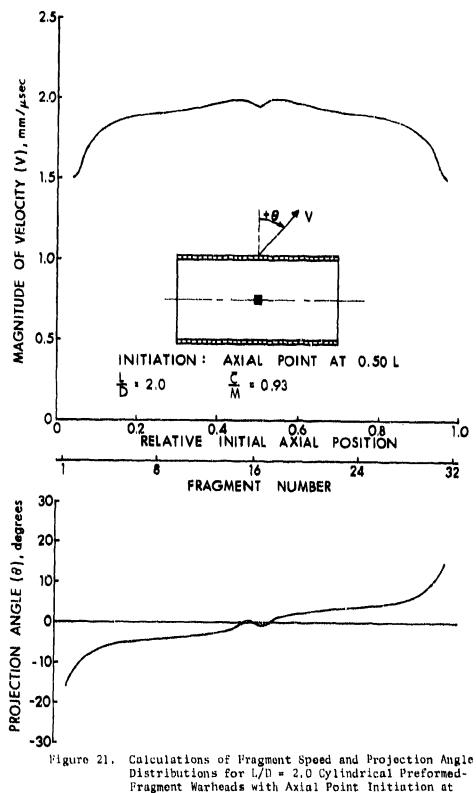


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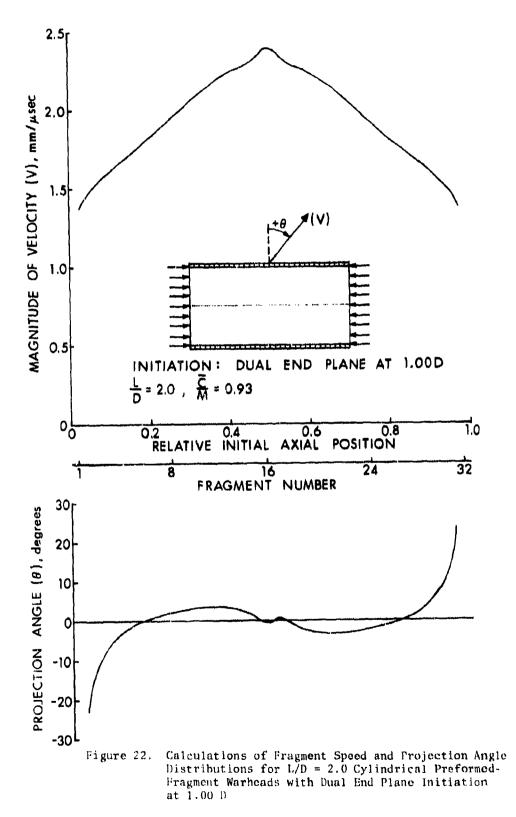
Figure 18. Initiation Postures

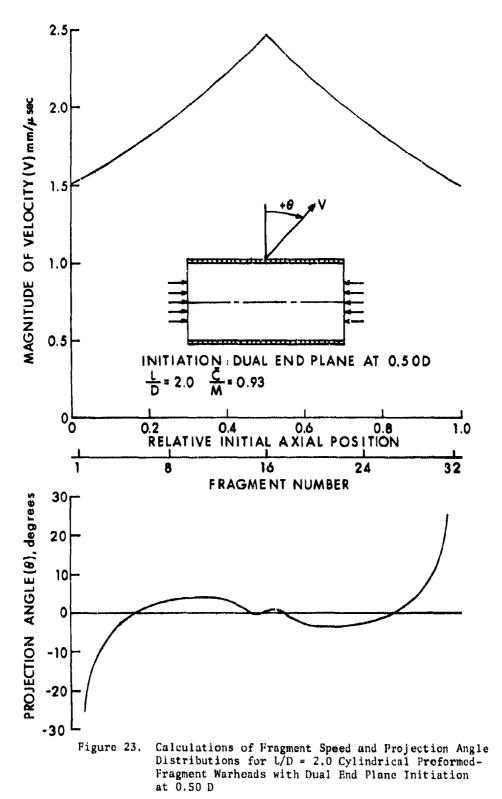






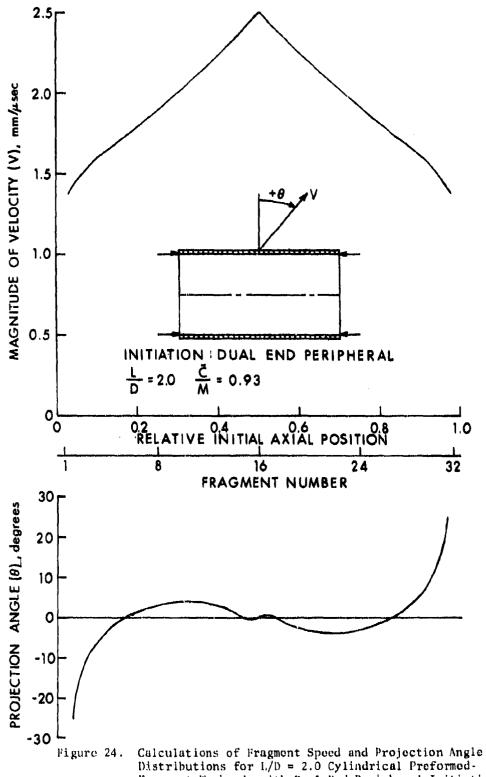




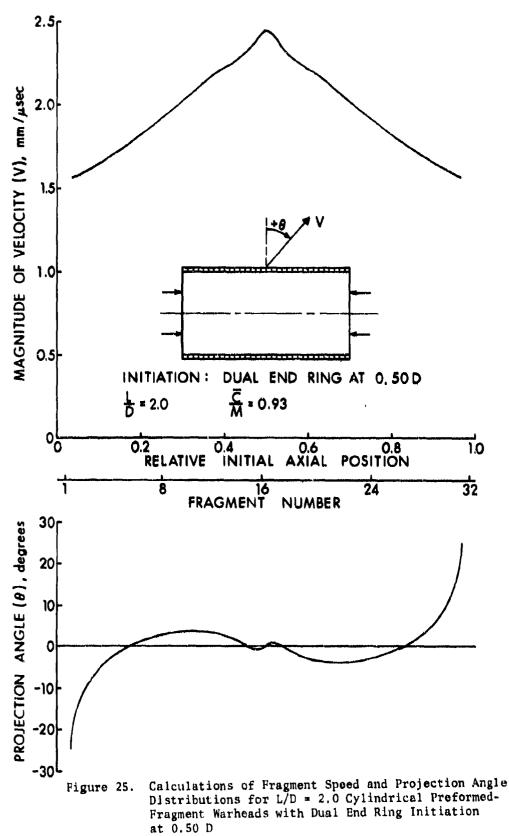


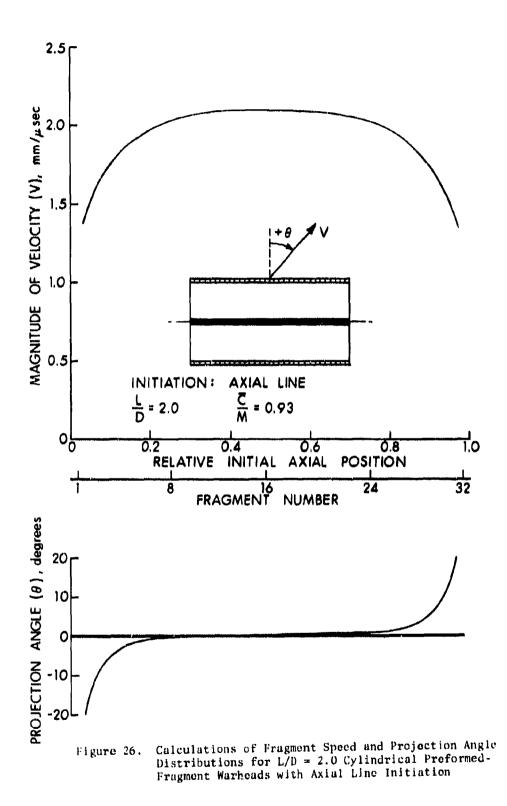
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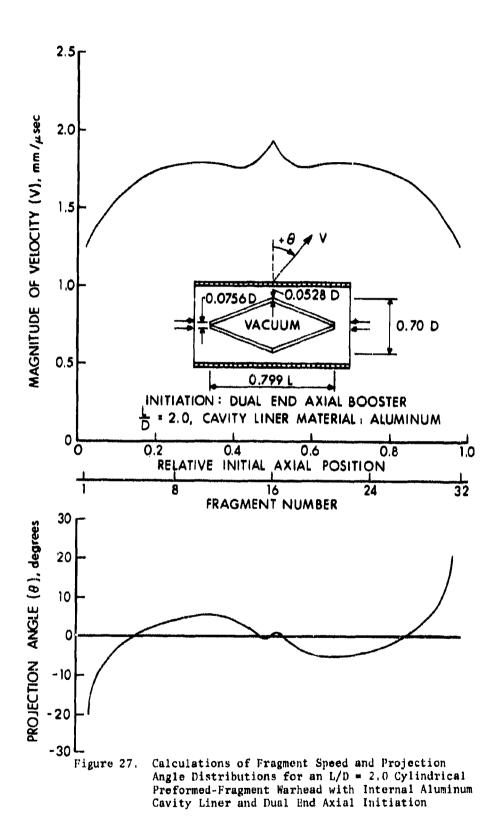
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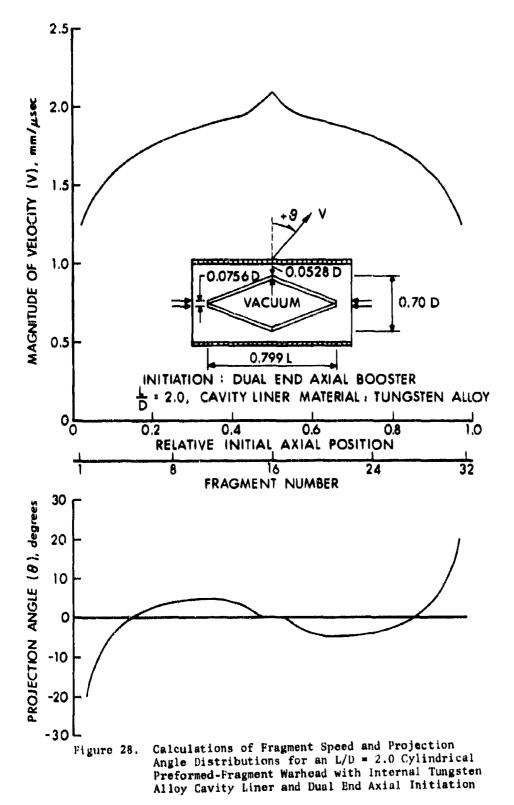


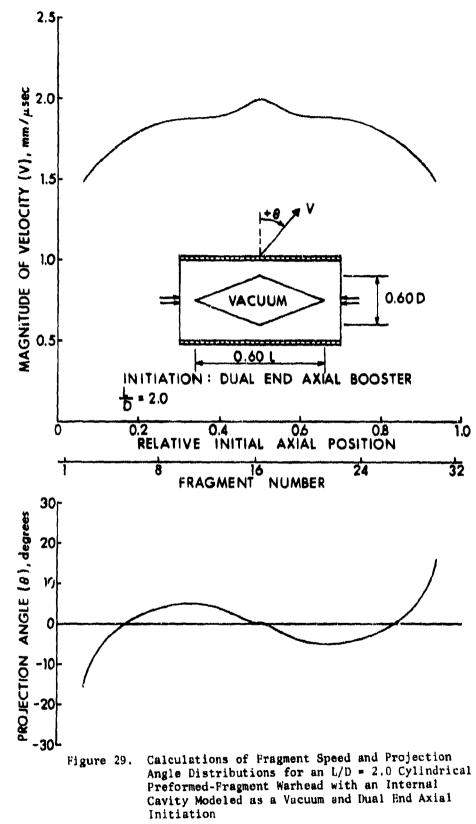
Fragment Warheads with Dual End Peripheral Initiation

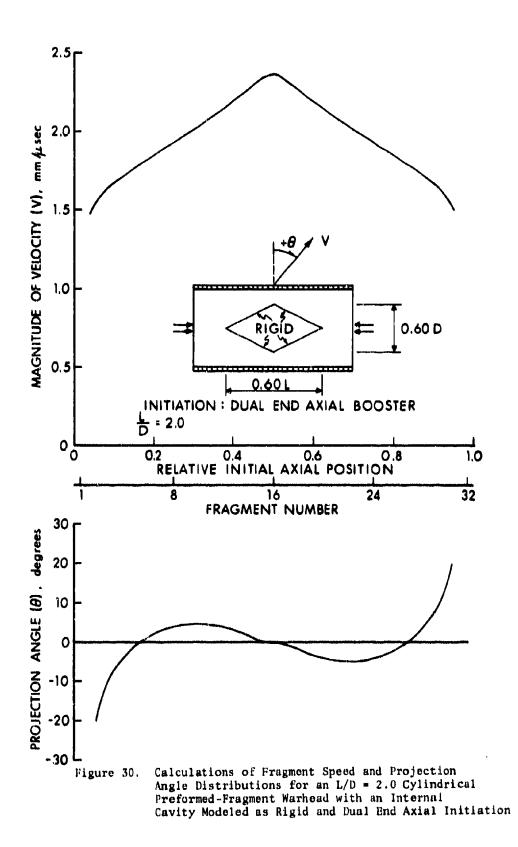












the case with an aluminum cavity liner material with the dimensions as indicated. Comparing the results shown in Figure 27 with those in Figure 17 it can be seen that although the projection angle distribution is only slightly altered, the fragment speed is significantly reduced.

The computational results with a tungsten alloy ( $\rho = 16890 \text{ KG/M}^3$ ) cavity liner material for the same cylindrical geometry and cavity shape given in Figure 27 are shown in Figure 28. As expected, due to the higher density cavity liner material the fragment speed distribution is somewhat higher than with aluminum; however, the fragment projection angle distribution is not noticeably altered. It is important to note that these results are for one cavity shape and therefore are not necessarily true for all cavity shapes and sizes. In Figures 29 and 30 are shown the computational results for a slightly different cavity size where the cavity is modeled either as a vacuum (Figure 29) or as a cavity whose boundaries are rigid (Figure 30) for all time. The purpose of these two computations is to provide a qualitative lower (Figure 29) and upper (Figure 30) bound on fragment speed and projection angle distributions that can be expected for the cylindrical geometry and cavity shape shown.

#### V. CONCLUSIONS AND RECOMMENDATIONS

A discrete-fragment warhead (missile warhead) computational modeling capability has been developed which is capable of modeling with a satisfactory degree of confidence the performance of a large variety of axisymmetric warhead configurations. The computational warhead model includes the ability to model the effects of varying L/D, fragment material and size, explosive material, initiation posture, and internal cavity material and shape. It is felt that the modeling capability can be extended to include two layers of fragments on the warhead, a liner between the fragments and the explosive, and end confinement. The experiments were completed for these cases and are discussed in Section II.

Future efforts should include additional experiments and computational modeling with different liner materials, shapes and sizes, and end confinement to understand fully the influence of these variables on warhead performance. Also, future work should include generalization of the empirical criterion for fragment separation to a more fundamental criterion. For example, material properties, such as the total plastic strain to failure, might be used to predict fragment separation.

It is also believed that the modeling capability can successfully be applied to certain aspects of eccentric warheads (a three dimensional problem). However, while some degree of success may be had, in order to model eccentric warheads properly, a three dimensional computer code is needed.

#### ACKNOWLEDGEMENTS

The authors wish to thank Dr. Robert R. Karpp for the many helpful technical discussions during the course of the work. Also the authors are grateful to Mr. G. Gentle and Mr. Carroll West for their assistance with the experimental program.

## LIST OF SYMBOLS

' **,** ',

С/м	Nominal charge (explosive) mass to metal mass (fragment + liner) ratio
D <sub>HE</sub>	Explosive diameter
D <sub>M</sub>	Calculated outside diameter of continuous metal casing model, from $M_T = \rho_f \{ \pi/4 L_W (D_M^2 \cdot D_{HE}^2) \}$
L <sub>c</sub>	Side of a cubical fragment
L <sub>HE</sub>	Explosive Length
L <sub>W</sub>	Length of fragments on warhead (equal to $L_{HE}$ in all cases)
Mc	Mass of a cubical fragment (also equivalent to m <sub>i</sub> )
M <sub>HE</sub>	Mass of the explosive charge
MR	Mass of a rod
m <sub>i</sub>	Individual fragment mass
MRC	Mass of a recovered cubical fragment
MT	Total mass of fragments and liner on warhead
Ncc	Number of cubical fragments per column (along warhead length)
N <sub>c</sub>	Number of columns of cubical fragments
NR	Number of rods per warhead
R	Current warhead radius in computations
<sup>R</sup> ъ	Radius of warhead at fragment separation in computations
R <sub>b/R</sub> o	Expansion ratio at fragment separation or breakage ratio; equivalent to $W_{\rm RC}/L_{\rm c}$
Ro	Initial inside radius of casing (metal fragments)
WRC	Recovered cube dimension in the circumferential (hoop) direction.
αf	Azimuthal fragment-bearing sector angle
٩f	Fragment mass density

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