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1. Introduction

Electromagnetic (EM) interference shielding materials are currently in high demand (for both commercial and defense purposes) because of electromagnetic pollution which is generated due to extensive use of electrical and electronic equipment. EM radiation produced from an external source can disturb electrical circuits, cell phones, televisions, satellite communications, global positioning systems (GPS), etc. The disturbances across communication channels, automation and process control lead to the loss of precious time, energy, resources and even life. Therefore, shielding must be provided to protect today's society from electromagnetic pollution. In general, there are two practical ways for shielding against electromagnetic interference (EMI): the first is to protect certain components from the radiation by reflecting or absorbing the waves; another is to reduce the reflection and increase the absorption by incorporating dielectric or magnetic particles into it.

In the present study, we focus on the second type of shielding, called absorption dominated microwave shielding. The complex

MnO₂ decorated graphene nanoribbons with superior permittivity and excellent microwave shielding properties[†]

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Microwave shielding properties of chemically synthesized MnO_2 decorated graphene nanoribbons (GNRs) are reported for the first time. The nature of MnO_2 decoration on the GNRs has been investigated using scanning electron microscopy, X-ray diffraction, Raman spectroscopy and high resolution transmission electron microscopy. The electromagnetic interference (EMI) shielding effectiveness of this material was investigated in the microwave region (Ku-band, 12.4–18 GHz). The presence of MnO_2 on GNR enhances the interfacial polarization, multiple scattering, natural resonances and the effective anisotropy energy, which leads to absorption dominated high shielding effectiveness of -57 dB (blocking >99.9999% radiation) by a 3 mm thick sample. Dielectric attributes (ε' and ε'') were evaluated to understand the mechanism of the excellent shielding effectiveness. The material will be an excellent choice for radar absorbing applications.

dielectric permittivity is an important parameter for describing the properties of absorption dominated microwave shielding or radar absorbing materials (RAM). Excellent dielectric properties are needed for making good RAM. Therefore, lightweight EMI shielding materials1-3 are sought in order to shield the workspace and environment from incoming radiation from computers and telecommunications as well as for protection of sensitive circuits. Previously, metal based microwave shielding materials were widely explored and utilized but the high specific gravity, corrosion proneness and cumbersome processing methods made these materials unsuitable for both the researchers and users. Furthermore, metals mainly reflect the radiation and cannot be used for applications where absorption is mandatory. Nowadays, microwave shielding materials⁴ with comparatively light weight, good flexibility and strong absorption in the gigahertz range are in high demand. Recently, carbon nanotubes (CNTs),5 magnetic-particle-doped CNTs and nanostructured ZnO were intensively investigated and found to be promising shielding materials. Liu et al.6 achieved a maximum absorption value of -22 dB at 8.8 GHz by 5 wt% single-walled carbon nanotubes (SWCNTs) loaded polyurethane composites. Fan et al.7 reported a maximum absorption of about 24.3 dB at 15.3 GHZ by 8 wt% multi-walled carbon nanotubes (MWCNTs) loaded polymer (or varnish) composites. Singh et al.8 reported a maximum shielding effectiveness due to absorption of -53 dB in the X-band (8.2-12.4 GHz) by 20.4 wt% MWCNTs loaded epoxy composites. Che et al.9 have observed the microwave absorption of about -25 dB at 11 GHz by Fe-CNTs-epoxy composites. However, the complex fabrication process of CNTs



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or magnetic-particle-doped CNTs is a challenging task (for using such absorbing nanocomposites in practical applications). Therefore, there is a need to look for new materials with relatively light weight, good flexibility, long lasting and especially exceptional microwave shielding properties.

Nanoscale materials based on single-layered 2-D graphene sheets have recently attracted great interest due to their many extraordinary properties.¹⁰ At present, carbon based composite materials employing CNTs are dominating the research and application fields. But the poor dispersion of CNTs due to bundling, the presence of impurities, lack of interfacial adhesion with other materials and high cost of production have marred their applications. It has been proposed that these issues could be reduced by incorporating graphene sheets into composite materials.^{11,12} Graphene nanoribbons (GNRs) have also been found an important material in the carbon allotrope family. GNRs have been synthesized by different routes previously.^{13–18}

On the other hand, manganese oxide (MnO₂) offers a new path for the development of advance microwave shielding materials because of its low cost and environment friendly nature. It is well known that the pseudo-capacitive reaction of MnO₂ is a surface reaction and therefore only a very thin layer of MnO₂ is sufficient for high microwave shielding.¹⁹ Thus, ultrathin MnO₂ can deliver very high microwave shielding.

There are only few reports on using MnO_2 as a microwave shielding material. Duan *et al.*²⁰ studied the microwave absorbing properties of α -MnO₂-carbon black composites and showed that these composites can absorb microwaves in a wide band range at about 8.7 GHz with lower than -10 dB (90% absorption) due to dielectric loss. Microwave absorption properties of MnO₂ composites were also studied by Guan *et al.*²¹ In their study, MnO₂–SiO₂–PVA composites had the strongest absorbing peak of about 25 dB at 8.5 GHz for 30 vol% MnO₂.

In the present study, we demonstrate the successful preparation of MnO_2 decorated GNRs *via* a chemical method in the presence of concentrated sulfuric acid which plays an important role in increasing the interlayer distance of graphene layers. The EMI shielding and dielectric properties of MnO_2 decorated GNRs are studied and it is observed that the incorporation of MnO_2 improves the electromagnetic absorption properties which is due to the enhancement of dielectric loss. We also explored the effect of varying thickness of the composites on the EMI shielding performance.

2. Experimental

2.1. Materials

Natural graphite with purity > 99% and BS mesh 500 was purchased from Pune Carbon, Pune, India. KMnO₄ (purity > 99%) and H₂SO₄ (purity > 98%) were purchased from Qualigens fine chemicals, Mumbai, India and RANKEM, New Delhi, India, respectively.

2.2. Process for decoration of GNRs with MnO₂

In a typical procedure as shown in Scheme 1, graphite flakes (50 mg) were mixed with $KMnO_4$ crystallites (300 mg) using planetary ball milling (2 h, 250 rpm) and a powder mixture was

obtained. This mixture was suspended in 50 ml distilled water and stirred using a magnetic stirrer (speed 600 rpm) at room temperature for 12 h. 5 ml of concentrated H₂SO₄ was introduced into the solution while continuing stirring (at 80 °C) for an additional 2 h. The reaction mixture was removed from the heat source, and poured into 500 ml distilled water to cool it down and dilute it. The solution was then filtered using a 0.22 μ m pore size Whatman filter followed by several washings with de-ionized water till the pH became neutral.

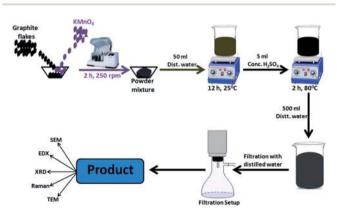
2.3. Characterization

Surface morphology of MnO2 decorated GNRs was investigated using SEM (Leo 440S, UK) and microstructural properties were investigated using HRTEM (Technai G30-stwin, 300 kV instrument). Raman studies were carried out using a Renishaw inVia reflex Raman spectrometer, UK with an excitation source of 514.5 nm. The XRD patterns were recorded using a Bruker D8 Advance X-ray diffractometer in the diffraction (2 θ) range of 10– 70° (slit width of 0.1 mm) using the Cu K α line ($\lambda = 1.5405$ A) as radiation source. EMI shielding effectiveness (SE) and electromagnetic attributes (real and imaginary permittivity) were measured using a vector network analyzer (VNA E8263B Agilent Technologies) by placing pellet having dimensions 15.8 \times 7.9 mm² inside a copper sample holder connected between the wave-guide flanges of the network analyzer in the Ku-band (12.4-18.0 GHz) waveguide. These obtained complex parameters were estimated from experimental scattering parameters $(S_{11} \text{ and } S_{21})$ by standard Nicholson-Ross and Weir theoretical calculations.²² Furthermore, reflection and transmission coefficients were also evaluated using scattering parameters:, R = $|S_{11}|^2 = |S_{22}|^2$ and $T = |S_{21}|^2 = |S_{12}|^2$. The absorbance was calculated as A = 1 - R - T.

3. Results and discussion

3.1. Morphological and structural characterizations

The surface morphology of the as-synthesized MnO_2 decorated GNR was studied using SEM. In Fig. 1(a), some layered structures are observed in MnO_2 decorated GNRs, and the GNRs are not clearly visible due to the coating of MnO_2 over the surface of GNRs. EDX analysis (Fig. 1(b)) shows the presence of Mn (53%),



Scheme 1 Synthesis process of decoration of GNRs with MnO₂.

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O (40%), and C (7%). XRD analysis was used to determine the structure of the hybrid. The peak at a 2θ value of 12.73° in Fig. 1(c) is due to the (110) plane of α -MnO₂. Other peaks are assigned to the tetragonal phase of α -MnO₂ (JCPDS 44-0141) indicating better crystallinity in strong acidic medium.^{23,24} Moreover, a low intensity peak corresponding to the (002) plane (inset of Fig. 1(c)) of graphitic structure was observed in MnO₂ decorated GNR samples. This might be due to the presence of excess MnO₂ which prevented the stacking of graphene sheets over each other to form crystalline structures.²⁵ This is evident

from the EDX results which show that the carbon content in $\rm MnO_2$ decorated GNRs is only 7 wt%.

Fig. 1(d) shows the room temperature Raman spectrum of MnO_2 decorated GNRs. The Raman spectrum was recorded using a low laser power of 5 mW in order to avoid any decomposition of the sample. The Raman spectrum of MnO_2 decorated GNRs is dominated by a sharp and strong band at approximately 645 cm⁻¹ with a shoulder at 576 cm⁻¹ which can be assigned to α -MnO₂. The Raman band at 645 cm⁻¹ of the A_{1g} mode is attributed to symmetric Mn–O stretching vibrations of

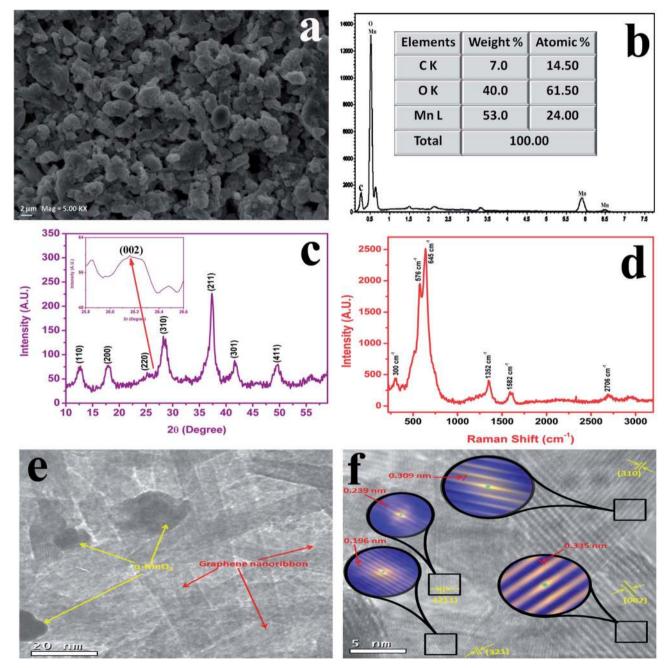


Fig. 1 (a) SEM image, (b) EDX results, (c) XRD pattern, (d) Raman spectrum of MnO₂ decorated GNRs, (e) TEM image showing the presence of MnO₂ (yellow arrows) on the surface of GNR (red arrows) and (f) showing the inter-planar spacing (yellow lines) and an enlarged view of the lattice fringes (black circles).

tetragonal MnO₂. The shoulder peak at 576 cm⁻¹ of the F_{2g} mode is not well distinguishable because of its low intensity. This shoulder peak is closely related to the oxidation state of manganese and appears due to the vibration of Mn–O bonds.^{26,27}

The two most intense features of the Raman spectrum are the presence of a G peak at ~1582 cm⁻¹ and a 2D peak at ~2706 cm⁻¹, as usually observed in graphitic structures.²⁸ The G and D bands are due to the doubly degenerate zone centre E_{2g} mode²⁹ at ~1582 cm⁻¹ and the breathing modes of six atom rings which appears at 1352 cm⁻¹ due to the presence of defects in the graphite,²⁸ respectively. The 2D peak is due to the second order defects. The 2D peak in bulk graphite consists of two components, 2D₁ (shoulder) and 2D₂,^{29,30} with heights roughly 1/4 and 1/2 of the height of the G peak, respectively. But the Raman spectrum of graphene has a single sharp 2D peak. The number of layers in graphene is estimated based on the intensity of the 2D peak. In the above figure, the D, G and 2D peak position values satisfy the criteria for the presence of graphene type structure in the MnO₂ decorated GNRs.

The TEM micrograph (Fig. 1(e)) shows that MnO_2 nanoparticles are spherical in shape (indicated by yellow arrows) and the particle size is in the range of 10–20 nm. They cover the surface of GNRs. The ribbon-like structure (red arrows) is a clear indication of the presence of nanocrystalline GNRs. The HRTEM micrograph (Fig. 1(f)) shows well-defined inter-planar spacings of 0.309 nm, 0.239 nm and 0.196 nm, which correspond to the (310), (211) and (321) planes of α -MnO₂. The inter-planar spacing of 0.354 nm corresponds to the (002) plane of graphene. The detailed microscopic evidence of MnO₂ decoration on GNRs is shown in HRTEM images (Fig. S1, ESI SP1†). All the above characterizations confirm the decoration of MnO₂ on GNRs.

3.2. Proposed reaction mechanism

Based on the above results, a mechanism for the formation of MnO_2 decorated GNRs is proposed and the scheme is shown in Fig. 2. In the preliminary stage, when a large amount of KMnO₄ is mixed with graphite under neutral conditions, it acts as a 'weak oxidizing agent' due to the much lower positive potential of KMnO₄ in a neutral medium than that in an acidic medium.^{31,32} During the slow oxidation process, oxygen-containing groups would attach themselves spontaneously on the graphite sheets.

In the next stage, when a small amount of H_2SO_4 is introduced into the system, $KMnO_4$ in the acidic medium oxidizes graphite faster, which links the oxidized sites to curved epoxy chains.³³ Step by step, MnO_4^- ions are absorbed and precipitated on the oxidized location through the interactions between remaining $KMnO_4$ and hydroxyl groups. When the solution mixture is heated, oxidation of the graphite layers is significantly enhanced. Intercalation of SO_4^{2-} increases the distance between graphite layers. Insoluble MnO_2 precipitates on the expanded graphite and cannot be washed away.

3.3. Microwave shielding effectiveness

The EMI SE is defined as the logarithmic ratio of incoming power (P_i) to outgoing power (P_t) of radiation. In general, the

efficiency of any shielding material is expressed in decibels (dB). For a higher decibel level of EMI shielding effectiveness, less energy must be transmitted through the shield. It depends upon the following three factors:^{3,34}

(a) absorption: the wave is absorbed in the shielding material,

(b) *reflection*: the wave is reflected from the front surface of the shielding material,

(c) *multiple reflections*: reflections of the waves from various interfaces

Therefore, the total EMI shielding effectiveness can be expressed as^{1,2,8,35–37}

$$SE_{T} (dB) = SE_{A} + SE_{R} + SE_{M} = 10 \log_{10} \left(\frac{P_{t}}{P_{i}}\right)$$
 (1)

As $P_{\rm t}$ is always less than $P_{\rm i}$, the shielding effectiveness is a negative quantity. When the total EMI shielding effectiveness of the material is more than -10 dB, then loss due to multiple reflections (SE_M) becomes negligible and can be neglected.^{34,38,39} Thus, the total shielding effectiveness (SE_T) can be expressed as

$$SE_{T} (dB) = SE_{A} + SE_{R}$$
⁽²⁾

A schematic diagram showing the interaction of electromagnetic waves with MnO_2 decorated GNRs is shown in Fig. 3. Reflection (R), absorption (A) and multiple internal reflections (MIR) which contribute towards the overall attenuation (with SE_R , SE_A and SE_M components, respectively) are shown. When the electromagnetic wave falls on MnO_2 decorated GNRs, the dominant loss of the incident electromagnetic wave is through absorption. The energy is consumed mainly due to dielectric loss (because of the presence of a large amount of dielectric material, α -MnO₂). A small part of the incident electromagnetic wave is reflected from the front and back interfaces of the material.

Microwave shielding materials having strong absorption, low density, and good physical and chemical stability are in high demand. Conventional microwave shielding materials (metals and magnetic materials) have limited practical applications due to their high specific gravity, difficulty in formulation and lack of proper response to gigahertz frequencies.⁴⁰

The frequency dependence of microwave absorption is a key criterion in selecting materials for EMI shielding applications. MnO₂ and graphene were selected in our experiments because of their high dielectric and absorption properties, respectively. Microwave shielding properties of these MnO₂ decorated GNRs were examined by making rectangular pellets of different thicknesses (0.5 mm, 1.0 mm, 1.5 mm, 2.0 mm, 2.5 mm and 3.0 mm). The pellets were placed in the Ku-band sample holder and shielding effectiveness was examined in the frequency range of 12.4-18.0 GHz. The plots of microwave shielding effectiveness versus frequency for MnO2 decorated GNRs having six different thicknesses are shown in Fig. 4. The value of SE_T is -11 dB for the sample having a thickness of 0.5 mm and it reaches -57 dB for the 3.0 mm thick sample (Fig. 4(a)). The terms SE_R and SE_A show the shielding contributions due to reflection and absorption, respectively and can be expressed mathematically as:2,34

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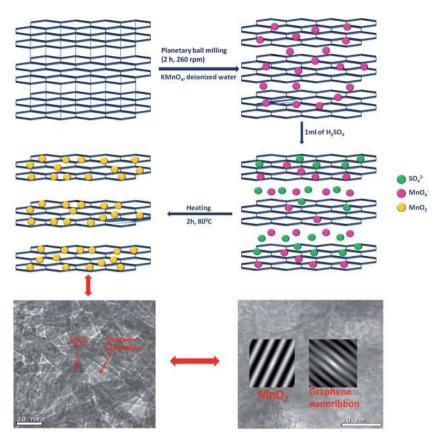


Fig. 2 Proposed mechanism for the formation of MnO₂/graphene nanoribbons.

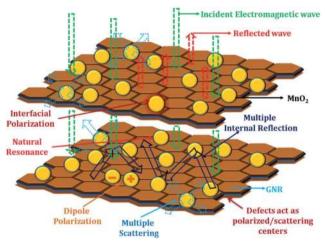


Fig. 3 Schematic presentation of possible microwave shielding mechanism in MnO_2 decorated GNRs.

$$SE_{R} (dB) = 108 + \log_{10} \left(\frac{\sigma}{f\mu} \right)$$
 (3)

$$SE_A (dB) = -10log_{10} \left(\frac{P_t}{P_i} \right) = -8.68 \{ t/\delta \} = -8.68 \alpha t$$
 (4)

where α is the attenuation coefficient which describes the extent to which the intensity of an electromagnetic wave is reduced when it passes through a specific material.

$$\alpha = \frac{4\pi n}{\lambda_0} \tag{5}$$

where *n* is the refractive index of the material and λ_0 is the wavelength in vacuum. Refractive index (*n*) is defined by the following relation:

$$n = (\varepsilon_{\rm r}\mu_{\rm r})^{1/2} \tag{6}$$

where ε_r and μ_r are the relative permittivity and permeability, respectively.

For non-magnetic materials μ_r is very close to 1.⁴¹ Therefore, $n = (\varepsilon_r)^{1/2}$ and we know that for a medium, ε_r is equal to ε' . Therefore,

$$n = (\varepsilon')^{1/2} \tag{7}$$

where ε' is the real part of permittivity.

Wavelength in vacuum (λ_0) is equal to $\lambda_0 = 2\pi c/\omega$ because $\omega = 2\pi n = 2\pi c/\lambda_0$, where *c* is the velocity of light and is described as:

$$c = \left(\varepsilon_{\rm r} \mu_{\rm r}\right)^{-1/2} \tag{8}$$

Therefore,

$$\lambda_0 = \frac{2\pi}{\omega(\varepsilon_r \mu_r)^{1/2}} \tag{9}$$

For non-magnetic materials, $\mu_r = 1$. Thus, high permittivity is very important for the enhancement of SE_A and suppressing

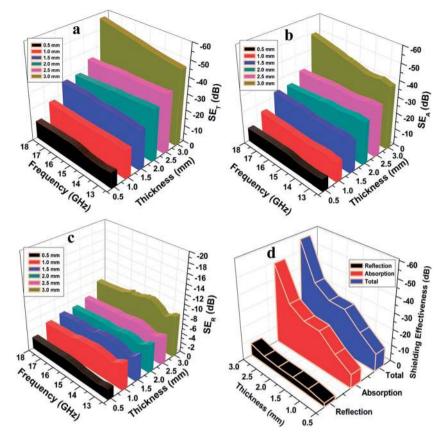


Fig. 4 Variation of (a) SE_T, (b) SE_A, (c) SE_R with frequency for different thicknesses of MnO₂ decorated GNRs and (d) SE_T, SE_A and SE_R for different thicknesses of MnO₂ decorated GNRs at 18.0 GHz.

 SE_R , respectively. Thus, the incorporation of very high permittivity material, MnO_2 nanoparticles, in GNRs enhanced the overall shielding effectiveness.

Fig. 4(b) shows that SE_A increases with increase in thickness and the value increased from -8.63 dB for the 0.5 mm thick sample to -49.34 dB for the 3.0 mm thick sample. This trend in reflection loss (SE_R) is very small (Fig. 4(c)) where it varies from -2.4 dB for 0.5 mm to -7.7 dB for the 3.0 mm thick sample. Fig. 4(d) shows the overall changes in % attenuation with thickness, in which changes in reflection loss with increasing thickness are very small as compared to absorption.

In the study of Hailong *et al.*,⁴² graphene was used as a microwave absorbing material and it showed a microwave absorption of -11 dB for a sample having a thickness of 4.0 mm in the 2.0–18.0 GHz frequency range. In another study, Wang *et al.*⁴³ reported 6.9 dB microwave absorption at 7 GHz using chemically reduced graphene oxide. Therefore an attempt has been made to improve SE_A by incorporating a high dielectric material (MnO₂ nanoparticles).

According to the electromagnetic theory, dielectric losses in a material occur because of complex phenomena such as natural resonances, dipolar relaxation, electronic polarization and its relaxation, and the unique structure of the material. In MnO₂ decorated GNRs, MnO₂ acts as a polarized centre in the presence of microwaves which gives rise to

better microwave shielding due to absorption. The relative complex permittivity ($\varepsilon^* = \varepsilon' - i\varepsilon''$) (at room temperature) has been measured for investigating the microwave shielding properties of MnO₂ decorated GNRs. The real part of permittivity (ε') is mainly associated with the amount of polarization in the material and is related to the storage of electrical energy, whereas the imaginary part of permittivity (ε'') is a measure of loss in energy. Fig. 5(a) shows the real (ε') and imaginary (ε'') parts of the relative complex permittivity of MnO₂ decorated GNRs in the frequency range of 12.4-18 GHz. It can be observed from Fig. 5(a) that ε' shows the highest value of 335 at 12.7 GHz and it decreases with fluctuations (from 320 to 223) in the 12.4-18.0 GHz range, similar to earlier reports.^{36,43,44} The ε'' value is lower than ε' and it fluctuates (from 173 to 100) in the 12.4-18.0 GHz range. The relatively high values of ε'' and $\tan \delta_{\varepsilon}$ (Fig. 5(b)) show that α -MnO₂ decorated GNRs exhibit high dielectric losses. These dielectric losses are attributed to complex phenomena, such as: natural resonance, dominant electronic and dipolar polarization, interfacial polarization and associated relaxation phenomena.36,45,46 The dielectric tangent loss (tan $\delta_{\varepsilon} = \varepsilon''/\varepsilon'$) based on the permittivity of the sample has also been evaluated and the values are in the range of 0.5 to 0.45 as shown in Fig. 5(b). This indicates higher absorption loss (-49.34 dB) which is equivalent to 99.9999% attenuation.

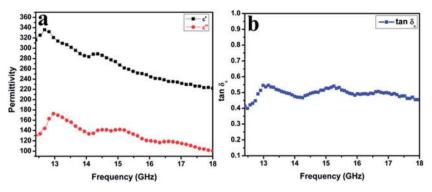


Fig. 5 Frequency dependence of (a) real and imaginary parts of permittivity and (b) corresponding dielectric tangent loss of MnO₂ decorated GNRs for a 3.0 mm thick sample.

Conventionally the relaxation process described by the Cole-Cole semicircles has an important effect on the permittivity behavior of microwave absorbing materials.

According to the Debye dipolar relaxation process,^{43,47} the complex permittivity can be written as

$$(\varepsilon' - \varepsilon_{\infty})^2 + (\varepsilon'')^2 = (\varepsilon_{\rm s} - \varepsilon_{\infty})^2 \tag{10}$$

where ε_s and ε_{∞} are the static dielectric constant and the dielectric constant at infinite frequency, respectively. The $\varepsilon' - \varepsilon''$ curve of MnO₂ decorated GNRs is shown in Fig. 6. The Cole–Cole plot presents a clear segment of three semicircles. The result demonstrates that there are three dielectric relaxation processes in the MnO₂ decorated GNRs with each semicircle corresponding to one Debye relaxation process which enhances the dielectric properties of MnO₂ decorated GNRs. In MnO₂ decorated GNRs, the existence of interfaces between MnO₂–MnO₂, GNR–GNR and MnO₂–GNR gives rise to interfacial polarization or Maxwell–Wagner–Sillars polarization.

The phenomenon of interfacial and space charge relaxation appears because charge carriers are trapped at the interfaces of heterogeneous media. Evidently, the Cole–Cole semicircles are distorted which reveals that except for the dielectric relaxation, other phenomena such as conductance loss, interfacial polarization and oxygen defects contribute to the permittivity spectra. Therefore, in MnO₂ decorated GNRs, ternary dielectric

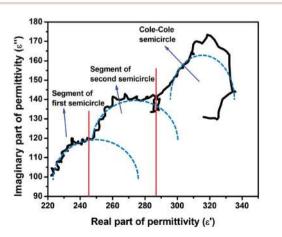


Fig. 6 Typical Cole–Cole plot of MnO₂ decorated GNRs.

relaxation, interfacial polarization and effective anisotropy energy are helpful in improving the level of impedance matching for the excellent microwave shielding due to absorption.

4. Conclusion

 MnO_2 decorated graphene nanoribbons (GNRs) have been synthesized by a chemical route which results in MnO_2 intercalated graphene. The material was further characterized using SEM, EDX, XRD, Raman and HRTEM. The results show that the existence of MnO_2 in GNRs effectively enhances the electronic polarization, interfacial polarization and anisotropy energy in the presence of microwaves. A maximum shielding effectiveness of -57 dB was achieved for a 3.0 mm thick sample and is mainly governed by absorption rather than reflection. Thus, inclusion of MnO_2 in GNRs is highly useful for better microwave absorption in Ku band (12.4–18 GHz) as compared to conventional materials. Additionally, the shielding effectiveness can easily be tailored by varying the thickness of the sample. It is believed that this kind of material will find applications in microwave absorbing area.

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References

- 1 T. K. Gupta, B. P. Singh, S. Teotia, V. Katyal, S. R. Dhakate and R. B. Mathur, *J. Polym. Res.*, 2013, **20**, 1–7.
- 2 B. P. Singh, Prabha, P. Saini, T. Gupta, P. Garg, G. Kumar,
 I. Pande, S. Pande, R. K. Seth, S. K. Dhawan and
 R. B. Mathur, *J. Nanopart. Res.*, 2011, 13, 7065–7074.
- 3 R. Kumar, S. R. Dhakate, T. Gupta, P. Saini, B. P. Singh and R. B. Mathur, *J. Mater. Chem. A*, 2013, **1**, 5727–5735.
- 4 V. M. Petrov and V. V. Gagulin, Inorg. Mater., 2001, 37, 93-98.

- 5 N. Li, Y. Huang, F. Du, X. He, X. Lin, H. Gao, Y. Ma, F. Li,Y. Chen and P. C. Eklund, *Nano Lett.*, 2006, 6, 1141–1145.
- 6 Z. Liu, G. Bai, Y. Huang, F. Li, Y. Ma, T. Guo, X. He, X. Lin, H. Gao and Y. Chen, *J. Phys. Chem. C*, 2007, **111**, 13696–13700.
- 7 Z. Fan, G. Luo, Z. Zhang, L. Zhou and F. Wei, *Mater. Sci. Eng.*, *B*, 2006, **132**, 85–89.
- 8 B. P. Singh, Prasanta, V. Choudhary, P. Saini, S. Pande,
 V. N. Singh and R. B. Mathur, *J. Nanopart. Res.*, 2013, 15, 1–12.
- 9 R. C. Che, L. M. Peng, X. F. Duan, Q. Chen and X. L. Liang, *Adv. Mater.*, 2004, 16, 401–405.
- 10 A. K. Geim and K. S. Novoselov, Nat. Mater., 2007, 6, 183-191.
- S. Stankovich, D. A. Dikin, G. H. B. Dommett, K. M. Kohlhaas, E. J. Zimney, E. A. Stach, R. D. Piner, S. B. T. Nguyen and R. S. Ruoff, *Nature*, 2006, 442, 282–286.
- 12 T. Ramanathan, A. A. Abdala, S. Stankovich, D. A. Dikin, M. Herrera-Alonso, R. D. Piner, D. H. Adamson, H. C. Schniepp, X. Chen, R. S. Ruoff, S. T. Nguyen, I. A. Aksay, R. K. Prud'Homme and L. C. Brinson, *Nat. Nanotechnol.*, 2008, 3, 327–331.
- 13 L. Xie, H. Wang, C. Jin, X. Wang, L. Jiao, K. Suenaga and H. Dai, *J. Am. Chem. Soc.*, 2011, 133, 10394–10397.
- 14 X. Li, X. Wang, L. Zhang, S. Lee and H. Dai, *Science*, 2008, **319**, 1229–1232.
- 15 L. Jiao, L. Zhang, X. Wang, G. Diankov and H. Dai, *Nature*, 2009, **458**, 877–880.
- 16 L. Jiao, L. Zhang, L. Ding, J. Liu and H. Dai, *Nano Res.*, 2010, 3, 387–394.
- 17 L. Jiao, X. Wang, G. Diankov, H. Wang and H. Dai, *Nat. Nanotechnol.*, 2010, 5, 321–325.
- 18 L. Jiao, L. Xie and H. Dai, Nano Res., 2012, 5, 292-296.
- 19 S. B. Ma, K. Y. Ahn, E. S. Lee, K. H. Oh and K. B. Kim, *Carbon*, 2007, **45**, 375–382.
- 20 Y. Duan, Y. Yang, M. He, S. Liu, X. Cui and H. Chen, *J. Phys. D: Appl. Phys.*, 2008, **41**, 125403.
- 21 H. Guan, S. Liu, Y. Zhao and Y. Duan, J. Electron. Mater., 2006, 35, 892–896.
- 22 A. M. Nicolson and G. F. Ross, *IEEE Trans. Instrum. Meas.*, 1970, **19**, 377–382.
- 23 J. Luo, H. T. Zhu, H. M. Fan, J. K. Liang, H. L. Shi, G. H. Rao, J. B. Li, Z. M. Du and Z. X. Shen, *J. Phys. Chem. C*, 2008, **112**, 12594–12598.
- 24 W. Xiao, D. Wang and X. W. Lou, *J. Phys. Chem. C*, 2010, **114**, 1694–1700.
- 25 J. Yan, Z. Fan, T. Wei, W. Qian, M. Zhang and F. Wei, *Carbon*, 2010, **48**, 3825–3833.
- 26 C. Julien, M. Massot, R. Baddour-Hadjean, S. Franger, S. Bach and J. P. Pereira-Ramos, *Solid State Ionics*, 2003, 159, 345–356.

- 27 C. Julien and M. Massot, Phys. Chem. Chem. Phys., 2002, 4, 4226-4235.
- 28 A. C. Ferrari, J. C. Meyer, V. Scardaci, C. Casiraghi, M. Lazzeri, F. Mauri, S. Piscanec, D. Jiang, K. S. Novoselov, S. Roth and A. K. Geim, *Phys. Rev. Lett.*, 2006, **97**, 187401.
- 29 R. J. Nemanich and S. A. Solin, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1979, 20, 392–401.
- 30 R. P. Vidano, D. B. Fischbach, L. J. Willis and T. M. Loehr, Solid State Commun., 1981, 39, 341–344.
- 31 K.-C. Chang, L. Li and E. F. Gloyna, J. Hazard. Mater., 1993, 33, 51–62.
- 32 Y. Chen, Y. Zhang, D. Geng, R. Li, H. Hong, J. Chen and X. Sun, *Carbon*, 2011, **49**, 4434–4442.
- 33 J.-L. Li, K. N. Kudin, M. J. McAllister, R. K. Prud'homme, I. A. Aksay and R. Car, *Phys. Rev. Lett.*, 2006, **96**, 176101.
- 34 T. K. Gupta, B. P. Singh, S. R. Dhakate, V. N. Singh and R. B. Mathur, *J. Mater. Chem. A*, 2013, 1, 9138–9149.
- 35 H.-B. Zhang, Q. Yan, W.-G. Zheng, Z. He and Z.-Z. Yu, *ACS Appl. Mater. Interfaces*, 2011, 3, 918–924.
- 36 A. P. Singh, M. Mishra, P. Sambyal, B. K. Gupta, B. P. Singh, A. Chandra and S. K. Dhawan, *J. Mater. Chem. A*, 2014, DOI: 10.1039/c3ta14212d.
- 37 A. P. Singh, M. Mishra, A. Chandra and S. K. Dhawan, *Nanotechnology*, 2011, 22, 9.
- 38 K. Singh, A. Ohlan, V. H. Pham, B. R. S. Varshney, J. Jang, S. H. Hur, W. M. Choi, M. Kumar, S. K. Dhawan, B.-S. Kong and J. S. Chung, *Nanoscale*, 2013, 5, 2411–2420.
- 39 A. Ohlan, K. Singh, A. Chandra and S. K. Dhawan, ACS Appl. Mater. Interfaces, 2010, 2, 927–933.
- 40 X. Sun, J. He, G. Li, J. Tang, T. Wang, Y. Guo and H. Xue, *J. Mater. Chem. C*, 2013, **1**, 765–777.
- 41 Y. A. Urzhumov, Sub-wavelength Electromagnetic Phenomena in Plasmonic and Polaritonic Nanostructures: From Optical Magnetism to Super-resolution, The University of Texas at Austin, 2007.
- 42 H. Yu, T. Wang, B. Wen, M. Lu, Z. Xu, C. Zhu, Y. Chen, X. Xue, C. Sun and M. Cao, *J. Mater. Chem.*, 2012, 22, 21679–21685.
- 43 C. Wang, X. Han, P. Xu, X. Zhang, Y. Du, S. Hu, J. Wang and X. Wang, *Appl. Phys. Lett.*, 2011, **98**, 072906–072903.
- 44 T. K. Gupta, B. P. Singh, R. B. Mathur and S. R. Dhakate, Nanoscale, 2014, 6, 842–851.
- 45 X. G. Liu, D. Y. Geng, H. Meng, P. J. Shang and Z. D. Zhang, *Appl. Phys. Lett.*, 2008, **92**, 173117–173113.
- 46 Y.-F. Zhu, L. Zhang, T. Natsuki, Y.-Q. Fu and Q.-Q. Ni, *ACS Appl. Mater. Interfaces*, 2012, 4, 2101–2106.
- 47 M.-S. Cao, J. Yang, W.-L. Song, D.-Q. Zhang, B. Wen, H.-B. Jin, Z.-L. Hou and J. Yuan, ACS Appl. Mater. Interfaces, 2012, 4, 6949–6956.