# MONOTONICITY OF ENTROPY FOR REAL MULTIMODAL MAPS 

HENK BRUIN AND SEBASTIAN VAN STRIEN


#### Abstract

In 1992, Milnor [28] posed the Monotonicity Conjecture that within a family of real multimodal polynomial interval maps with only real critical points, the isentropes, i.e., the sets of parameters for which the topological entropy is constant, are connected. This conjecture was already proved in the mid-1980s for quadratic maps by a number of different methods, see 30, 10, 9, 26, 42. In 2000, Milnor \& Tresser 31, provided a proof for the case of cubic maps. In this paper we will prove the general case of this 20 year old conjecture.


## 1. Introduction and Statement of Results.

Given a family of continuous maps $f$ from an interval $I$ to itself, one can ask how its 'dynamical complexity' depends on $f$. Let us assume that $I$ can be decomposed into finitely many subintervals $I_{0}, \ldots, I_{b}$ on which $f$ is monotone. The smallest number $b+1$ of such intervals is called the lap number $\ell(f)$ of $f$. Note that $b$ is the number of extrema of $f$, and is often called the modality of $f$. Maps for which $b$ is equal to one or two are called unimodal and bimodal respectively. One natural measurement of the dynamical complexity of $f$ is the rate of exponential growth of the lap numbers $\ell\left(f^{n}\right)$ where $f^{n}$ denotes the $n$-th iterate of $f$. This growth rate $\lim _{n \rightarrow \infty} \frac{1}{n} \log \ell\left(f^{n}\right)$ exists and is equal to the usual notion of topological entropy $h_{\text {top }}(f)$ of $f$, see [33] and also [26, Section II.7]. Topological entropy can be used to classify maps with finite modality, up to semi-conjugacy, see [30] - a bit like the rotation number enables a classification up to semi-conjugacy of degree one circle maps. Conventionally, a map is called chaotic if and only if $h_{\text {top }}(f)>0$.

For continuous interval maps, $h_{\text {top }}(f)$ coincides with the exponential growth rate of the number of $n$-periodic orbits. Therefore if we consider a family of continuous interval maps $f_{t}, t \in[0,1]$ and $h_{\text {top }}\left(f_{1}\right)>h_{\text {top }}\left(f_{0}\right)$, then many periodic orbits are created as the parameter $t$ increases from 0 to 1 . However, it should be noted that entropy is only a coarse indicator of the birth of periodic orbits, because periodic

[^0]

Figure 1. Monotonicity of entropy for the maps $f_{a}(x)=a x(1-x)$, $a \in[3.5,4]$.
orbits can both appear and disappear in parameter ranges of constant entropy. This is clearly true in modality $\geq 2$ (since entropy is only a one-dimensional observable in higher dimensional parameter space), but already in modality one, entropy is constant on any parameter interval within a period doubling cascade. In order to obtain a complete picture on the emergence and disappearance of periodic orbits, it is therefore necessary to combine monotonicity of entropy with so-called Thurston Rigidity, which we explain later on in this introduction.

In this paper we consider the space of $b$-modal continuous interval maps. Obviously, $h_{\text {top }}(f) \in[0, \log (b+1)]$ for a $b$-modal map. It turns out that $f \mapsto h_{\text {top }}(f)$ is not continuous on the entire space of $b$-modal maps, but if we restrict to $C^{1}$-smooth maps then it is, see [30, 43, 32, 26].

The question whether $h_{\text {top }}(f)$ 'increases' with $f$ goes at least back to the early 70s, see [27]. In the unimodal situation, one of the simplest ways of asking this question is as follows. Let $I=[0,1]$ and consider a smooth unimodal map $f:[0,1] \rightarrow[0,1]$ with $f(0)=f(1)=0, f(1 / 2)=1$ and the family $f_{a}(x)=a f(x), a \in[0,1]$.
(1) Does the topological entropy of $f_{a}(x)=a f(x)$ increase with $a \in[0,1]$ ?

As mentioned before, entropy cannot be strictly increasing with $a$. It has been conjectured in the 90 's that if a $C^{3}$ unimodal convex map $f:[0,1] \rightarrow[0,1]$ as above has negative Schwarzian and is symmetric around the critical point, then the answer to (1) is positive. This conjecture is subtle: there are $C^{3}$ close maps $f, g:[0,1] \rightarrow[0,1]$ of this type for which $f \leq g$ yet $h_{\text {top }}(f)>h_{\text {top }}(g)$, see [3]. Moreover, none of the assumptions can be dropped, see the examples in [45, 19, 34] and also [26, Section II.10].

It was shown in the mid 1980's by Douady \& Hubbard [9, 10] and Milnor \& Thurston [30] that for the quadratic family $f_{a}(x)=4 a x(1-x)$ the entropy $h_{\text {top }}\left(f_{a}\right)$ depends monotonically on $a \in[0,1]$. All known proofs of this use that the quadratic map can be extended to the complex plane and require tools from complex analysis, see the above references and also Tsujii's proof [42] and [26, Section II.10]. To show how subtle this question is, let us note that it was only very recently shown that the topological entropy of $f_{a}(x)=a \sin (\pi x)$ depends monotonically on $a$, see [37]. In fact, that paper shows that (11) holds for each unimodal $f:[0,1] \rightarrow[0,1]$ with $f(0)=f(1)=0$ which extends to an entire transcendental map on the complex plane, with a finite number of singular values and satisfying the so-called sector condition.

The above proofs not only show that the topological entropy of $f_{a}(x)=4 a x(1-x)$ increases with $a$, but also that periodic orbits do not disappear when $a$ increases. In other words, as $a$ increases, only new periodic orbits are created (by period doubling and saddle-node bifurcations). That this is true, follows essentially from the following property:

> Thurston Rigidity (combinatorially equivalent critically finite maps are unique): Consider $f_{a}$ and $f_{a^{\prime}}$ for which their critical points $c=1 / 2$ have finite orbits $O$ and $O^{\prime}$. If there exists an order preserving homeomorphism $h: I \rightarrow I$ with $h(O)=O^{\prime}$ and $h \circ f_{a}=f_{a^{\prime}} \circ h$, then $a=a^{\prime}$.

In fact, much more is known: whenever $f_{a}$ and $f_{a^{\prime}}$ have the same 'kneading invariant' and $f_{a}$ has no periodic attractor, then $a=a^{\prime}$. This result was proven in [13, 23] and is usually called the density of hyperbolicity for real quadratic maps. This result implies that there exists a dense set $H \subset[0, \log 2]$ so that for any $h_{0} \in H$ there exists precisely one $a \in[0,1]$ with $h_{\text {top }}\left(f_{a}\right)=h_{0}$. It follows that $a \mapsto h_{\text {top }}\left(f_{a}\right)$ is a devil's staircase, the plateaus of which correspond to intervals of parameters containing a periodic attractor and the subsequent period doubling cascade. By density of hyperbolicity, such parameters form a dense set, and so there exists no interval of parameters on which $a \mapsto h_{\text {top }}\left(f_{a}\right)$ is strictly increasing.

Douady \& Hubbard, see [10], showed the following monotonicity result:
Bifurcations are monotone in the quadratic family: Assume that $\left(a_{-}, a_{+}\right)$is a parameter range so that the quadratic family $f_{a}(x)=$ $4 a x(1-x)$ has a hyperbolic periodic attractor $p_{a}$ of period $n$ for each $a \in\left(a_{-}, a_{+}\right)$, then $a \mapsto D f_{a}^{n}\left(p_{a}\right) \in(-1,1)$ is differentiable and strictly decreasing.

The corresponding parameter space for higher degree polynomials is parametrized by Blaschke products, see [29] and Theorems 3.4 and 3.6 below. Combining Thurston

Rigidity with the previous property shows that period doubling cascades are traversed monotonically in the quadratic family. (We should point out that there are additional results on the transversality of bifurcations in polynomial families, see [39, 11, 22].)

Let us turn to real cubic maps. Take $I=[-1,1]$ and cubic maps $f: I \rightarrow I$ with exactly two critical points, both in the interior of $I$. Since this space consists of two connected components, it makes sense to separate the cases where $f(-1)=$ $-1, f(1)=1$ and where $f(-1)=1, f(1)=-1$. In the former case, such cubic maps can be written in the form $f_{\mathrm{a}, \mathrm{b}}(x)=\mathrm{a} x^{3}+\mathrm{b} x^{2}+(1-\mathrm{a}) x-\mathrm{b}$ where $(\mathrm{a}, \mathrm{b}) \in \mathbb{R}^{2}$ are contained in a simply-connected region bounded four algebraic curves (this follows as in [28]). It is not hard to show that for a smooth one-parameter family $f_{t}$ of such cubic maps, $t \mapsto h_{\text {top }}\left(f_{t}\right)$ need not be monotone, see Remark 4.2. Perhaps this is not too surprising, as the level sets of $(\mathrm{a}, \mathrm{b}) \mapsto h_{\text {top }}\left(f_{\mathrm{a}, \mathrm{b}}\right)$ are very complicated fractallike sets. Related to this is the result by Kan, Koçak \& Yorke [18] that within the Hénon family $F_{\mathrm{a}, \mathrm{b}}(x, y)=\left(1-\mathrm{a} x^{2}+\mathrm{b} y, x\right)$, the entropy of $F_{\mathrm{a}, \mathrm{b}}$ does not depend monotonically on a for fixed b.

Yet a compelling question is whether, within the space of all real cubic maps, those with a given topological entropy form a connected set. In the early 1990's Milnor made this question precise, by defining the following space. Consider the space $P_{\epsilon}^{b}$ of real polynomials $f$ with
(1) precisely $b$ distinct critical points, all of which are real, non-degenerate and contained in the interior of $I$;
(2) $f(\partial I) \subset \partial I$;
(3) with shape $\epsilon=\epsilon(f)$, where

$$
\epsilon(f)= \begin{cases}+1 & \text { if } f \text { is increasing at the left endpoint of } I, \\ -1 & \text { otherwise. }\end{cases}
$$

Note that $P_{\epsilon}^{b}$ consists of polynomials of degree $d=b+1$.
Milnor's conjecture essentially asserts that within this space, bifurcations are 'as efficient as possible':

Milnor's Monotonicity of Entropy Conjecture [28]: For each $\epsilon \in\{-,+\}, b \in \mathbb{N}$ and $h_{0} \geq 0$, the isentrope

$$
\left\{f \in P_{\epsilon}^{b} ; h_{\text {top }}(f)=h_{0}\right\}
$$

is connected.

This conjecture was motivated by numerical experiments, made for the family of real cubic maps, of the isentropes and also by considering the 'bones' for this family. These bones are curves within the parameter space such that one critical point is periodic with a specified order type, and were introduced by MacKay and Tresser in


Figure 2. Isentropes for cubic maps $f_{\mathrm{a}, \mathrm{b}}(x)=\mathrm{a} x^{3}+\mathrm{b} x^{2}+(1-\mathrm{a}) x-\mathrm{b}$. The horizontal axis gives $a \in[2.5,4]$ and the vertical axis $\mathrm{b} \in[0, \sqrt{4 \mathrm{a}}-$ a]. The maps $f_{\mathrm{a},-\mathrm{b}}$ and $f_{\mathrm{a}, \mathrm{b}}$ are conjugate, and $f_{\mathrm{a}, \mathrm{b}}([-1,1]) \not \subset[-1,1]$ for $|\mathrm{b}|>\sqrt{4 \mathrm{a}}-\mathrm{a}$.
a study of the boundary of chaos for bimodal maps of the interval [24]. A few years later it was shown in [8 that in the cubic case $b=2$, this conjecture follows from another conjecture (density of hyperbolicity for cubic maps). We should emphasize that although Thurston Rigidity holds for polynomials of any degree, this is not sufficient for proving the monotonicity conjecture for degree $d>2$. In 2000, Milnor \& Tresser [31] showed that one does not quite need density of hyperbolicity for real cubic maps. More precisely, they showed that on some curves in the parameter space, the bimodal family behaves essentially like a one-parameter family of unimodal maps. Combining this with an extension of density of hyperbolicity in the quadratic case, due to Heckman [16] and using the Jordan theorem, they were able to conclude the cubic case. Using similar techniques as Milnor \& Tresser, Radulescu [36] proved monotonicity of entropy for a two-parameter family of quartic polynomials made up of the composition of two quadratic maps.

In this paper we solve this conjecture in full generality:
Main Theorem (Milnor's Monotonicity of Entropy Conjecture). For each $\epsilon \in$ $\{+,-\}, b \in \mathbb{N}$ and $h_{0} \geq 0$, the isentrope

$$
I\left(h_{0}\right)=\left\{f \in P_{\epsilon}^{b} \quad ; \quad h_{\text {top }}(f)=h_{0}\right\}
$$

and the set

$$
I\left(h_{0}^{+}\right):=I\left(h_{0}\right) \cap \operatorname{closure}\left(\left\{f \in P_{\epsilon}^{b} \quad ; \quad h_{\text {top }}(f)>h_{0}\right\}\right)
$$

are both connected.

In particular, the boundary of 'chaos', i.e., the boundary of the set of maps in $P_{\epsilon}^{b}$ with positive entropy is connected. In fact, the proof of the theorem also shows that for each $h_{1} \leq h_{0}$, the set

$$
I=\left\{f \in P_{\epsilon}^{b} \quad ; \quad h_{1} \leq h_{\text {top }}(f) \leq h_{0}\right\}
$$

is connected.
The set of maps $P_{\epsilon}^{b} \subset P^{b}$ with shape $\epsilon$ can be parametrized by the coefficients of the polynomial, or more suitably by the critical values of $f$, see [26, Section II.4] or [30]. As mentioned, one should not expect that the entropy depends monotonically on any of these parameters.
Remark 1.1. We should emphasise that we prove that the isentrope $I\left(h_{0}\right)=\{f \in$ $\left.P_{\epsilon}^{b} ; h_{\text {top }}(f)=h_{0}\right\}$, rather than the weaker statement that the closure of this space is connected. (So we prove connectedness within the space of maps with nondegenerate critical points, rather than merely in the closure of this space.)

As in Milnor \& Tresser's paper [31], our proof relies on stunted sawtooth maps. The other important ingredient is density of hyperbolicity, but now for real polynomials of arbitrary degree, see [20] and [21]. More precisely, we use an analogue of Thurston Rigidity proved in [20] which holds for all real polynomials with real critical points regardless of whether the orbits of the critical points are finite or not.

On the way to proving the Main Theorem, we will also prove the following result (see Theorem 5.12 and Theorem 6.1):

Theorem A. Fix $\epsilon \in\{+,-\}, b \in \mathbb{N}$, let $f \in P_{\epsilon}^{b}$ and define

$$
\mathcal{H}(f)=\left\{g \in P_{\epsilon}^{b} \quad ; \quad g \text { has the same kneading invariants as } f\right\} .
$$

Then $\mathcal{H}(f)$ is connected.
1.1. Related results and some conjectures. In this paper we will also consider the space of so-called admissible stunted sawtooth maps $\mathcal{S}_{*}^{b}$, and show that isentropes within this space are contractible, see Theorem 7.1. In view of this, we would like to propose the following
Conjecture. Any isentrope $\left\{f \in P_{\epsilon}^{b} \quad ; \quad h_{t o p}(f)=s\right\}$ is contractible.
Isentropes in $P_{\epsilon}^{b}$ are extremely complicated. Indeed, we prove in [4] the following
Theorem. When $b \geq 4$, there are infinitely many values for $s \geq 0$ for which $\{f \in$ $\left.P_{\epsilon}^{b} ; h_{\text {top }}(f)=s\right\}$ is not locally connected.

In fact, it is not known whether there exists any value $s \in[0, \log (b+1)]$ so that the corresponding isentrope $\left\{f \in P_{\epsilon}^{b} ; h_{\text {top }}(f)=s\right\}$ is locally connected. The methods used in the proof of the previous theorem rely on $b \geq 4$, and it is possible that each isentrope is connected in the cubic case.

In the survey [41] a number of related questions and conjectures are discussed. In particular, the following question due to Tresser: Consider the space Pol ${ }_{\epsilon}^{d}$ of real polynomials $f$ of degree $d$, not necessarily with all critical points on the real line, but still with $f(\{ \pm 1\}) \subset\{ \pm 1\}$ and $\epsilon(f)=\epsilon$ as in the definition of $P^{b}$.
Conjecture 1.2 (Tresser). Fix $\epsilon \in\{-1,1\}$. Isentropes in Pol ${ }_{\epsilon}^{d}$ are connected.
Davoud Cheraghi and the second author have made progress towards this conjecture in the context of real unimodal polynomials of degree 4 with at most one real critical point, see [7].
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## 2. Strategy of the Proof, organization of the paper and notation

It is well-known that any multimodal map (with positive topological entropy) is entropy-preservingly semi-conjugate to a piecewise monotone map of constant slope, [35, 30]. However, such piecewise affine maps do not exhibit all possible combinatorial types which exist for polynomials maps. Instead, one of the ingredients in Milnor \& Tresser's proof is to consider the space of stunted sawtooth maps, all obtained from a single sawtooth map $S_{0}$ as in Figure 3. For example, for each cubic map, there exists a stunted sawtooth map $T$ as in this figure with the same combinatorics. This map $T$ is obtained by moving the two plateaus up or down as appropriate in a way which is made precise in Section 5. In our paper, we will use the space $\mathcal{S}_{\epsilon}^{b}$ of $b$-modal stunted sawtooth maps to "parametrize" the space $P_{\epsilon}^{b}$ of $b$-modal polynomials. Indeed, we introduce a map

$$
\Psi: P_{\epsilon}^{b} \rightarrow \mathcal{S}_{\epsilon}^{b}
$$

which assigns to $f \in P_{\epsilon}^{b}$ the unique map $T \in \mathcal{S}_{\epsilon}^{b}$ which has the same 'kneading invariant' (i.e., symbolic itineraries of critical points) as $f$. We discuss the definition of $\Psi$ in detail in Section 5. An important property of $\Psi$ is that $\Psi(f)$ and $f$ have the same topological entropy.

Since we shall fix the shape $\epsilon$ during the proof, we shall write from now on mostly $P^{b}$ and $\mathcal{S}^{b}$ instead of $P_{\epsilon}^{b}$ and $\mathcal{S}_{\epsilon}^{b}$. One of the crucial benefits of the space $\mathcal{S}^{b}$ is that outside its plateaus, any map $T \in \mathcal{S}^{b}$ agrees with the same map $S_{0}$. This means that all orbits of two stunted sawtooth maps $T$ and $\tilde{T}$ agree except on the preimages of their plateaus. By decreasing the width of a plateau $Z_{i}$ (i.e., by moving the image $T\left(Z_{i}\right)$ up or down depending on whether $T$ assumes a local maximum or minimum at $Z_{i}$ ), we create new orbits without destroying any others, and hence the entropy can only increase. This means that within the space $\mathcal{S}^{b}$, entropy is a monotone


Figure 3. Two bimodal stunted sawtooth maps $T$ and $\tilde{T}$ (drawn in bold lines) constructed from the same sawtooth map $S_{0}$ (drawn in thin lines).
function of each critical value separately, a property which fails for $P^{b}$, see Remark 4.2 and [4]. Using this idea, Milnor \& Tresser showed the following:

Theorem 2.1 (31]). Isentropes in $\mathcal{S}_{\epsilon}^{b}$ are connected and contractible.
This result would imply the main theorem if $\Psi: P_{\epsilon}^{b} \rightarrow \mathcal{S}_{\epsilon}^{b}$ was a homeomorphism, but as we shall see that is unfortunately not the case.
2.1. Non-bijectivity of $\Psi: P^{b} \rightarrow \mathcal{S}^{b}$. One of the reasons the map $\Psi: P^{b} \rightarrow \mathcal{S}^{b}$ is non-bijective is simple to see: if $f \in P^{b}$ is hyperbolic (i.e., if each critical point is in the basin of a hyperbolic periodic attractor), then the itinerary of each critical point of $f$ is eventually periodic. From the definition of $\Psi$ it then follows that the endpoints of each plateau of $T=\Psi(f)$ are also eventually periodic. Since there are uncountably many hyperbolic maps in $P^{b}$ and only countably many such maps $T$, the map $\Psi$ is obviously not injective. Neither is $\Psi$ surjective, see Example 5.2. To overcome this we consider equivalence classes in $P^{b}$ and $\mathcal{S}^{b}$.
2.2. Equivalence classes in $P^{b}$ : partial conjugacy. The first ingredient aimed at overcoming the fact that $\Psi: P^{b} \rightarrow \mathcal{S}^{b}$ is neither injective nor surjective is to define a notion of equivalence classes within these spaces, corresponding to sets of maps which have the same dynamics except inside the basins of their attractors. It turns out that we will need slightly different notions within the spaces $P^{b}$ and $\mathcal{S}^{b}$. For $P^{b}$, two maps $f$ and $\tilde{f}$ will be taken to be equivalent if they are partially conjugate. For this to hold, we require that they are conjugate (on the real line) outside the basins $B(f)$ and $B(\tilde{f})$ of their periodic attractors and that they have the same number of critical points in corresponding components of $B(f)$ and $B(\tilde{f})$. In other words, two interval maps $f, \tilde{f}: I \rightarrow I$ are partially conjugate, if there exists a homeomorphism $h: I \rightarrow I$ so that $h \circ f=\tilde{f} \circ h$ holds outside $B(f)$, so that $h(B(f))=B(\tilde{f})$ and so that $h$ maps critical points of $f$ to critical points of $\tilde{f}$. Given $f \in P^{b}$, we define
$\mathcal{P H}(f)$ to be the set of polynomials $\tilde{f} \in P^{b}$ which are partially conjugate to $f$.

For more precise definitions see Section 3. Although the space $\mathcal{H}(f)$ (of maps with the same kneading as $f$ ) and $\mathcal{P H}(f)$ are closely related, neither is a subset of the other one, see Example 5.2. Extending the rigidity theorems from [21] we obtain

Theorem 2.2 (See Theorem 3.4). For any $f \in P_{\epsilon}^{b}$, the set $\mathcal{P H}(f)$ is connected.
This result is a non-trivial extension of Douady \& Hubbard's result that hyperbolic components with the space of (complex) quadratic polynomials are topologically discs. We emphasize that an important ingredient in the proof of this theorem is that all critical points of maps in $P^{b}$ are real. The situation when two real polynomials are conjugate on the real line, but have critical points which are outside the real line, is more subtle and the subject of ongoing research, see [7].
2.3. Preplateau equivalence in $\mathcal{S}^{b}$. The same definition can also be used to say when $T, \tilde{T} \in \mathcal{S}^{b}$ are partially conjugate. In Lemma 5.1 we will see that if $f, \tilde{f} \in$ $P^{b}$ are partially conjugate then $T=\Psi(f), \tilde{T}=\Psi(\tilde{f})$ are partially conjugate, but unfortunately the reverse implication does not hold. This is why we also introduce another equivalence class within the space $\mathcal{S}^{b}$. Indeed, in $\mathcal{S}^{b}$ we will define a set $W(T)$ which deviates slightly from $B(T)$ and which is based on the preimages of plateaus, see Section 4.2. We then say that $T, \tilde{T} \in \mathcal{S}^{b}$ are equivalent if $W(\tilde{T})=W(T)$ and define

$$
\langle T\rangle=\{\tilde{T} ; W(\tilde{T})=W(T)\} \quad \text { and } \quad[T]=\operatorname{closure}(\langle T\rangle)
$$

In Theorem 4.17 we show that $\langle T\rangle$ and therefore its closure $[T]$ is connected (in fact, it is a cell). From the definitions it follows that all maps within $\mathcal{P} \mathcal{H}(f)$ (respectively within $[T]$ ) have the same topological entropy.
2.4. The set $\mathcal{A}_{b} \subset P^{b}$ and a useful property of the map $\Psi$. Unfortunately it is not true that $\Psi(\mathcal{P H}(f)) \subset[\Psi(f)]$ for any $f \in P^{b}$. To address this issue and in order to relate $\mathcal{P H}(f)$ and $[T]$, we introduce a subset $\mathcal{A}_{b} \subset P^{b}$, see Definition 5.3 , which enables us to obtain the following property:

Proposition 2.3 (See Proposition 5.4). If $f \in \mathcal{A}_{b}$ then $\Psi(\mathcal{P H}(f)) \subset[\Psi(f)]$. Within any $\mathcal{P H}(f)$ we can find special maps $f_{0} \in \mathcal{P H}(f) \cap \mathcal{A}_{b}$.

Because of this proposition we can morally view $\Psi$ as a map which sends equivalences classes (consisting of partially conjugate maps) in $P^{b}$ to equivalence classes (consisting of preplateau equivalent maps) in $\mathcal{S}^{b}$.
2.5. Non-surjectivity of $\Psi: P^{b} \rightarrow \mathcal{S}^{b}$ because of wandering pairs. There is an additional, more serious way, in which $\Psi: P^{b} \rightarrow \mathcal{S}^{b}$ is not surjective, and this has to do with wandering intervals. An interval $J \subset[-1,1]$ is called wandering for $f$ if all its iterates are disjoint and $f^{n}(J)$ does not converge to an attracting periodic orbit as $n \rightarrow \infty$. It is well-known that polynomial interval maps (indeed $C^{2}$ interval map with non-flat critical points) have no wandering intervals. This implies that, when
$b \geq 3$, there are many stunted sawtooth maps $T$ in $\mathcal{S}^{b}$ for which there is no $f \in P^{b}$ with $T=\Psi(f)$. Indeed, take $T$ with the property that there exists an interval containing two adjacent plateaus which is eventually mapped into a third plateau which is not eventually periodic. Then for any interval map $f$ with $\Psi(f)=T$, the interval connecting the corresponding adjacent critical points would be wandering. Since a polynomial $f$ does not have wandering intervals, these adjacent critical points coincide and so that $f$ has only two critical points. This in turn implies that $T \neq \Psi(f)$. Therefore by using merely the space $\mathcal{S}^{b}$ we could at best prove that isentropes within $P^{b}$ are connected within the larger space of polynomial maps with $\leq b$ critical points.
2.6. The space of non-degenerate stunted sawtooth maps $\mathcal{S}_{*}^{b}$. To overcome this problem we define the notion of wandering pairs of plateaus, in Definition 4.12, and introduce the space $\mathcal{S}_{*}^{b}$ of stunted sawtooth maps without wandering pairs. In turns out that $\Psi$ maps $P^{b}$ into $\mathcal{S}_{*}^{b}$. The topology of $\mathcal{S}_{*}^{b}$ is much more complicated than that of $\mathcal{S}^{b}$, which makes it tricky to show that isentropes in $\mathcal{S}_{*}^{b}$ are contractible. In $\mathcal{S}^{b}$, this is much easier: within $\mathcal{S}^{b}$ one can construct a retract of an isentrope by moving plateaus with relatively great liberty. To construct a retract for isentropes in $\mathcal{S}_{*}^{b}$ we are forced to move plateaus in exactly the right order and with exactly the right speed. The description of this occupies most of Section 7, and leads to
Theorem 2.4 (Connectivity of isentropes in $\mathcal{S}_{\epsilon, *}^{b}$, see Theorem 7.1). Isentropes in $\mathcal{S}_{\epsilon, *}^{b}$ are connected and even contractible.
Remark 2.5. The construction in Theorem 7.1 of this contraction is quite subtle. Nevertheless it would be interesting to explore whether one can use the same methodology to construct deformations within the space $P^{b}$ (and show that isentropes within that space are contractible).
2.7. The main steps in the proof. With these notions in place, we will obtain that $\Psi$ is 'almost' surjective and injective:
Proposition 2.6 ( $\Psi$ is 'almost' surjective, see Proposition 5.9. Take $T \in \mathcal{S}_{*}^{b}$ and let $\tilde{T} \in\langle T\rangle_{\bullet}$. Then there exists a polynomial $f \in P^{b} \cap \mathcal{A}_{b}$ such that $\tilde{T}=\Psi(f)$ and $T \in[\Psi(f)]$.
Proposition 2.7 ( $\Psi$ is almost injective, see Proposition5.10). The map $\Psi: P^{b} \rightarrow \mathcal{S}_{*}^{b}$ is 'almost injective' in the sense that if $f_{1}, f_{2} \in \mathcal{A}_{b}$ and $\left[\Psi\left(f_{1}\right)\right] \cap\left[\Psi\left(f_{2}\right)\right] \neq \emptyset$, then $\overline{\mathcal{P H}\left(f_{1}\right)} \cap \overline{\mathcal{P H}\left(f_{2}\right)} \neq \emptyset$.

To prove the latter proposition, it is important to analyze how two subsets $\mathcal{P} \mathcal{H}(f)$ and $\mathcal{P H}(\tilde{f})$ of the space of polynomials $P_{\epsilon}^{b}$ can intersect. It is for this reason that we require a description of what bifurcations occur at the boundary of these sets, see Theorem 3.6. A corresponding description for the boundary of $\langle T\rangle$ within the space $\mathcal{S}_{\epsilon}^{b}$ is also needed, and is given in Theorem 4.17.

Proposition 2.8 ( $\Psi$ is almost continuous, see Proposition 5.13). $\Psi: P^{b} \rightarrow \mathcal{S}_{*}^{b}$ is 'almost continuous' in the following sense. Assume that $f_{n} \rightarrow f$ where $f_{n}, f \in P^{b}$ and $f_{n} \in \mathcal{A}_{b}$. Then there exists $T \in \mathcal{S}_{*}^{b}$ so that $\Psi\left(f_{n}\right) \rightarrow[T]$ and $\Psi(f) \in[T]$.

Combined this gives the following:
Theorem 2.9. There exists a map $\Psi: P_{\epsilon}^{b} \rightarrow \mathcal{S}_{\epsilon, *}^{b}$ such that

- $\Psi$ is 'almost continuous', 'almost surjective' and 'almost injective' (as described in the previous three propositions);
- There exists a connected set $[\Psi(f)] \ni \Psi(f)$ such that the topological entropy of any map $T \in[\Psi(f)]$ is equal to the topological entropy of $f$;
- If $K$ is closed and connected then $\Psi^{-1}(K)=\{f \quad ; \quad[\Psi(f)] \cap K \neq \emptyset\}$ is connected.

Isentropes in $\mathcal{S}_{\epsilon, *}^{b}$ are contractible (and therefore connected), so Theorem 2.9 implies that isentropes in $P_{\epsilon, *}^{b}$ are connected, proving the Main Theorem.
2.8. Organization of the paper. Section 3 discusses the notion of partial conjugacy and shows that partial conjugacy classes $\mathcal{P H}(f)$ within the space $P^{b}$ of polynomials are connected. It also describes when different sets $\mathcal{P H}(f)$ and $\mathcal{P} \mathcal{H}(\tilde{f})$ have common boundary points. This section relies on methods which use complex analysis and results from the theory of holomorphic dynamics. In particular, this section it relies on a theorem on rigidity of real polynomials. The remainder of the paper only uses real methods. Section 4 discusses the space $\mathcal{S}^{b}$ of stunted sawtooth maps and properties of the equivalence classes $\langle T\rangle$. Section 5 discusses the map $\Psi: P^{b} \rightarrow \mathcal{S}^{b}$ and its properties. In Section 6 the proof of the Main Theorem is provided. The technical result that isentropes in $\mathcal{S}_{*}^{b}$ are connected (and even contractible) is proved in Section 7.

### 2.9. Notation used in this paper.

- $B(f)$ is the union of the basin of periodic attractors, see Definition 3.1.
- $\mathbb{D}$ denotes the open unit disc in $\mathbb{C}$.
- $I$ is an interval in $\mathbb{R}$.
- $Z_{i}$ are (closed) plateaus of the stunted sawtooth map, see Section4.1.
- $P^{b}$ is the space of real polynomials of degree $b+1$ mapping $I$ (and $\partial I$ ) into itself, with precisely $b$ non-degenerate critical points each of which is contained in the interior of $I$.
- $P_{\epsilon}^{b} \subset P^{b}$ is the space of maps which is increasing (respectively decreasing) at the left endpoint of $I$ when $\epsilon=1$ (respectively $\epsilon=-1$ ).
- $\mathcal{H}(f)$ is the space of maps $g \in P^{b}$ with the same kneading invariant as $f$, see Theorem A in the introduction.
- $\mathcal{P H}(f)$ is the space of maps $g \in P_{\epsilon}^{b}$ which are partially conjugate to $f$, see Definition 3.2.
- $\mathcal{K} \mathcal{S}(f)$ is the space of Kupka-Smale maps, see Definition 3.3.
- $S_{0}$ is a sawtooth map of modality $b$, see Section 4.1.
- $\mathcal{S}^{b}$ is the space of stunted sawtooth maps, see Section 4.1.
- $W(T)$ is the set of points which are eventually mapped into the interior of a block of plateaus, see equation (5).
- $\langle T\rangle$ is the set of maps with the same $W(T)$, see equation (6).
- $[T]$ is the closure of $\langle T\rangle$, see equation (6).
- $\min [T]$ and $[T]_{b}$ are certain subsets of the boundary of $\langle T\rangle$, see Definition 4.7.
- $\mathcal{J}:=\left[Z_{i}, Z_{j}\right]$ is the convex hull of plateaus $Z_{i}$ and $Z_{j}$, see Definition 4.12.
- $\mathcal{S}_{*}^{b}$ is the space of non-degenerate maps in $\mathcal{S}^{b}$, see Definition 4.12,
- $\Psi: P^{b} \rightarrow \mathcal{S}^{b}$ is the map which assigns to a polynomial a stunted sawtooth map, see equation (8) in Section 5 .
- $\mathcal{A}_{b}$ is a subset of polynomials with parabolic periodic points, see Definition 5.3.
- $M_{n, \epsilon}, M_{n, \epsilon}^{o}, M_{n, \epsilon}^{\Sigma}$ are spaces of Blaschke products, see Definition 3.9.
- $\Gamma_{t}, \gamma_{t}, \hat{\Gamma}_{t}$ are entropy increasing deformations, see Section 7.2 and 7.8 .
- $\delta_{t}, \hat{\delta}_{t}$ and $\Delta_{t}$ are entropy decreasing deformation, see Sections 7.3 and 7.9 .
- $K_{i}$ and $\hat{K}_{i}$ are periodic intervals related to the $i$-th plateau of $T$, see Section 7.9.
- $\beta_{t}$ is an entropy preserving deformation, see Section 7.7
- $R_{t}$ and $r_{t}$ are retracts, see Sections 7.5, 7.6 and 7.3.


## 3. The partial conjugacy class of maps in $P^{b}$ is connected

As usual, we say that a polynomial $f$ is hyperbolic if each of its periodic orbits is hyperbolic and each of its critical points lies in the basin of a periodic attractor. A well-known result due to Douady \& Hubbard asserts that each connected component of the set $\left\{c \in \mathbb{C} ; q_{c}(z)=z^{2}+c\right.$ is hyperbolic $\}$ is topologically an open disc parametrised by the multiplier of the periodic attractor. The corresponding case for polynomials of higher degree was considered in [29]. In [11, 22] it shown that the multipliers of non-repelling periodic points are independent parameters. In this section we will generalize these results to polynomials of higher degree with the crucial difference that we no longer assume that each critical point is in the basin of hyperbolic periodic attractors and restrict to partial conjugacy classes (defined below). We shall only prove this generalization for real polynomials, because one of the main ingredients we need is a rigidity result which is only known in that context.

Before stating this generalization we will introduce some terminology.
Definition 3.1 (Basin of an interval map). We say that a periodic orbit $O$ of an interval map $f: I \rightarrow I$ is attracting if its basin $B^{s}(O)=\left\{x ; f^{n}(x) \rightarrow O\right.$ as $\left.n \rightarrow \infty\right\}$ contains a (possibly one-sided) neighborhood of $O$. Let $B(f)$ be the union of the basins of periodic attractors of $f$, i.e., $B(f)$ consists of all points $x$ so that $f^{n}(x)$ tends to a (possibly one-sided) periodic attractor. Note that if $f$ has a neutral periodic point (i.e., a periodic point which is non-hyperbolic), $B(f)$ need not be open. When $f$ is a polynomial we also will consider $f$ as acting on the complex plane, and in order to emphasize this we sometimes write $B_{\mathbb{C}}(f)$ to denote the basin of $f$ as a subset of $\mathbb{C}$.

Definition 3.2 (Partially conjugate). We say that two b-modal maps $f, g: I \rightarrow I$ are partially conjugate if there is an orientation preserving homeomorphism $h: I \rightarrow I$ such that

- $h$ maps $B(f)$ onto $B(g)$;
- $h$ maps the $i$-th critical point of $f$ to the $i$-th critical point of $g$;
- $h \circ f(x)=g \circ h(x)$ for all $x \notin B(f)$.

We denote by $\mathcal{P} \mathcal{H}(f)$ the set of maps $g \in P_{\epsilon}^{b}$ which are partially conjugate to $f$.

Note that when a critical point is eventually mapped to the boundary of a component of $B(f)$ that this property persists within $\mathcal{P H}(f)$.

Definition 3.3 (Kupka-Smale maps). Let $\mathcal{K} \mathcal{S}$ be the set of $g \in P_{\epsilon}^{b}$ which are KupkaSmale in the sense that

- $g$ has only hyperbolic periodic points and
- $g$ has no homoclinic orbits, i.e., no critical point of $g$ is mapped to the boundary of a component of $B(\mathrm{~g})$.

The set

$$
\mathcal{P} \mathcal{H}^{o}(f):=\mathcal{P H}(f) \cap \mathcal{K} \mathcal{S} .
$$

generalizes the notion of hyperbolic component for quadratic maps allowing, for the situation that some critical points are not attracted to periodic attractors provided the dynamics of such critical points agrees for all maps within $\mathcal{P} \mathcal{H}^{\circ}(f)$. Note that $f \in \mathcal{K} \mathcal{S}$ does not imply $\mathcal{P} \mathcal{H}^{\circ}(f)=\mathcal{P H}(f)$ because even in this case $\mathcal{P H}(f)$ can contain maps with neutral periodic orbits.

The main result in this section is the following theorem and its more detailed version Theorem 3.6.

Theorem 3.4 (Connectedness of $\mathcal{P H}(f))$. Let $f \in P^{b}$.

- If $f \in \mathcal{K} \mathcal{S}$ then $\mathcal{P} \mathcal{H}^{\circ}(f)$ is homeomorphic to an open ball of dimension equal to the number of critical points in $B(f)$.
- $\mathcal{P H}(f) \subset \overline{\mathcal{P} \mathcal{H}^{o}(f)}$ and therefore the set $\mathcal{P H}(f)$ is connected.

In fact, we shall also need Theorem 3.6 which states that for any $f \in P^{b}$ one can find a continuous family of maps $f_{\mu}, \mu \in[0,1]$ with $f_{\mu} \in \mathcal{K} \mathcal{S}$ for $\mu \in(0,1]$ and $f_{0}=f$ with the crucial additional property that $f_{\mu}$ has the same dynamics as $f_{0}$ outside the basins. Before stating that theorem more formally, let us clarify what the types of non-hyperbolic periodic points can occur for maps within the space $P^{b}$.

Lemma 3.5. Let $f \in P^{b}$. Then each attracting or neutral periodic point of $f$ is real and contains a critical point in its basin. Moreover, if $p$ is a neutral point (say of minimal period $n$ ) then it attracts at least from one side and is of one of the following types:
(pd) $p$ is attracting from both sides with multiplier -1 and up to a change of coordinates $f^{2 n}$ has the form $x \mapsto x-x^{3}+O\left(x^{4}\right)$ near $p$.
( $\mathbf{p f )} p$ is attracting from both sides with multiplier 1 and up to a change of coordinates $f^{n}$ has the form $x \mapsto x-x^{3}+O\left(x^{4}\right)$ near $p$.
(sn) $p$ is one-sided attracting with multiplier 1 and up to a change of coordinates $f^{n}$ has the form $x \mapsto x-x^{2}+O\left(x^{3}\right)$ near $p$.

Proof. Here we use that maps in $P^{b}$ are real and only have real critical points. This condition implies that if $p$ is a periodic orbit of period $n$ with $D f^{n}(p)=1$ then it must be attracting one from side (otherwise the attracting petals will not intersect the real line, but this is impossible since all critical points lie in the real line). Using that each attracting forward invariant petal of a neutral periodic point contains a critical point, the result follows. Alternatively, one can use the fact that maps $f \in P_{\epsilon}^{b}$ have negative Schwarzian derivative, i.e., $S f=\left[f^{\prime} \cdot f^{\prime \prime \prime}-(3 / 2)\left(f^{\prime \prime}\right)^{2}\right] /\left(f^{\prime}\right)^{2}<0$. Since this implies that the Schwarzian derivative of $f^{n}$ and $f^{2 n}$ are negative, see Exercise IV.1.7 in [26], the required statement follows.


Figure 4. A map with a pf (pitchfork), pd (period doubling) respectively sn (saddle-node) fixed point and the corresponding unfoldings constructed in Theorem 3.6.

These case are described in Figure 4. One of the main ingredients in this paper is the following theorem which shows that one can find a deformation $f_{\mu}, \mu \in[0,1]$ of a any map $f \in P^{b} \backslash \mathcal{K} \mathcal{S}$ so that $f_{\mu} \in \mathcal{K} \mathcal{S}$ for $\mu \in(0,1]$ and so that $f_{\mu} \in \mathcal{P} \mathcal{H}\left(f_{1}\right)$ for each $\mu \in(0,1]$. Therefore the only bifurcations of $f_{\mu}$ are that basins can merge or split and that neutral orbits become hyperbolic.
Theorem 3.6 (Bifurcations of maps which are not in $\mathcal{K} \mathcal{S}$ ). Assume that $f \notin \mathcal{K} \mathcal{S}$ and let $O_{1}, \ldots, O_{k}$ be the non-hyperbolic periodic orbits of $f$ and pick a periodic point $p_{i} \in O_{i}$ for each $i=1, \ldots, k$. For each periodic orbit $O_{i}$, choose a sign $\sigma_{i} \in\{-,+\}$.

Then there exists a family $f_{\mu} \in P_{\epsilon}^{b}$ and periodic points $p_{i, \mu}$ all depending continuously on $\mu \in[0,1]$ so that $f_{0}=f, p_{i, 0}=p_{i}$ and

$$
f_{\mu} \in \mathcal{P} \mathcal{H}^{\circ}\left(f_{1}\right) \text { for each } \mu \in(0,1]
$$

If $p_{i}$ has (minimal) period $n_{i}$ and
(1) $p_{i}$ is attracting from both sides and has multiplier -1 (the ( $p d$ )-case), then depending on the sign of $\sigma_{i}$ a period doubling or a period halving bifurcation occurs as $\mu$ becomes positive; i.e., up to a parameter dependent coordinate change $f_{\mu}^{2 n_{i}}$ has for each $\mu \in[0,1]$ near $p_{i, \mu}$ the form $x \mapsto\left(1+\sigma_{i} \mu\right) x-x^{3}+$ h.o.t.;
(2) $p_{i}$ is attracting from both sides with multiplier 1 (the ( $p f$ )-case), then depending on the sign of $\sigma_{i}$ a pitch-fork or a reverse pitch-form bifurcation occurs as $\mu$ becomes positive; i.e., up to a parameter dependent coordinate change $f_{\mu}^{n_{i}}$ has for each $\mu \in[0,1]$ near $p_{i, \mu}$ the form $x \mapsto\left(1+\sigma_{i} \mu\right) x-x^{3}+$ h.o.t.;
(3) $p_{i}$ is one-sided attracting and has multiplier 1 (the (sn)-case), then one can create a saddle-node pair for $f_{\mu}$ as $\mu$ becomes positive; i.e., up to a parameter dependent coordinate change $f_{\mu}^{n_{i}}$ has for each $\mu \in[0,1]$ near $p_{i, \mu}$ the form $x \mapsto(1+\mu) x-x^{2}+$ h.o.t.
(4) If $p_{i}$ is one-sided attracting and has multiplier 1 and there exists a neighborhood $U$ of $p_{i}$ so that $U \backslash\left\{p_{i}\right\}$ is contained in $B(f)$, then one can create or destroy a saddle-node pair for $f_{\mu}$ as $\mu$ becomes positive; i.e., up to a parameter dependent coordinate change $f_{\mu}^{n_{i}}$ has for each $\mu \in[0,1]$ near $p_{i, \mu}$ the form $x \mapsto x-x^{2}+$ $\sigma_{i} \mu+$ h.o.t.
(5) Moreover, if a critical point $c$ is eventually mapped into the boundary of a component of $B(f)$, then $f_{\mu}(c)$ is contained in the interior of $B\left(f_{\mu}\right)$ for each $\mu \in(0,1]$.

This theorem asserts that one can find a family of maps $f_{\mu} \in \mathcal{K} \mathcal{S}$ when $\mu \in(0,1]$ so that the two adjacent parabolic petals of $f=f_{0}$ at a neutral periodic point $p_{i}$ of type (pf) and (pd) as in case (1) and (2) correspond for $\mu>0$ to two adjacent hyperbolic basins when $\sigma_{i}=+$ or to one hyperbolic basin when $\sigma_{i}=-$ (so the sign of $\sigma_{i}$ determines whether or not one has a reverse period doubling of pitchfork bifurcation). In case (4) two basins which touch at a a saddle-node orbit are deformed into a map where the basins touch at a repelling orbit (when $\sigma_{i}=+$ ) or merge into one hyperbolic basin (when $\sigma_{i}=-$ ). This situation corresponds to the fixed point with multiplier $=1$ in Figure 5 .

In order to clarify case (5) in the theorem, consider the situation that $f(c)$ is contained in the boundary of a component $B$ of $B(f)$. Then the basin $B_{\mathbb{C}}(f)$ has two components $B_{1}, B_{2}$ which touch at $c$. These components $B_{1}, B_{2}$ can lie symmetrically in the upper and lower half plane or to the left and right of $c$. Which case occurs depends on whether $f(c)$ is a left or right end point of $B \cap \mathbb{R}$ and whether $f$ has a maximum or a minimum at $c$. A situation with the latter case is shown in Figure 5, see also Figure 6 where $U$ plays the role of $B \cap \mathbb{R}$.


Figure 5. For this map, there are three degeneracies: two parabolic fixed points (one with multiplier -1 and another one with multiplier 1) and a homoclinic orbit, i.e., a critical point which is mapped to the boundary of $B(f)$. In this case, both sides of the one-sided attracting fixed point with multiplier 1 are contained in $B(f)$. Theorem 3.6 shows that such a map can be embedded in a family of maps $f_{\mu}$, $\mu \in[0,1]$ so that for $\mu \in(0,1]$ the map $f_{\mu}$ only has hyperbolic periodic points (in the situation shown in the figure, a unique hyperbolic fixed point) and so that the two basins are merged.
3.1. Underlying rigidity theorems. Before proving Theorem 3.4, we should remark that it is related to (and extends) the following result:

Theorem 3.7 (Rigidity Theorem, see [20]). Let $f \in P_{\epsilon}^{b}$. Assume that $f$ has no periodic attractors. Then $\mathcal{P} \mathcal{H}(f)=\{f\}$.

For the quadratic case, Theorem 3.7 was proved independently by Lyubich [23] and Graczyk \& Świa̧tek [13]. Milnor \& Tresser used this result, or rather a version which applies to certain cubic maps due to Heckman, a PhD student of Świątek (see [16]), in their proof of the Main Theorem for the cubic case, see [31.

Let us also note that Theorem 3.7 is related to density of hyperbolicity. We say that an interval map $f$ is hyperbolic if each critical point of $f$ is in the basin of a hyperbolic periodic attractor. Building on Theorem 3.7] it was shown in [21] that one has density of hyperbolicity: each real polynomial can be approximated by a hyperbolic real polynomial of the same degree. In fact, each $C^{\infty}$ interval map can be approximated in the $C^{\infty}$ topology by a hyperbolic $C^{\infty}$ map, see [21]; it follows for example that within generic one-parameter families of interval maps hyperbolic maps are dense, see [40].

The main ingredient in the proof of Theorems 3.4 and 3.6 is the following:
Theorem 3.8 (Generalized Rigidity Theorem, see [20]). Let $f, g \in P_{\epsilon}^{b}$ and assume that $f, g$ are partially conjugate on the real line. Moreover, assume that for each periodic attractor (or parabolic point) of $f$ there exists a conformal map $h: B_{\mathbb{C}}^{f} \rightarrow B_{\mathbb{C}}^{g}$ from the basin $B_{\mathbb{C}}^{f}$ of this periodic attractor to the basin $B_{\mathbb{C}}^{g}$ of the corresponding periodic attractor of $g$ so that $h \circ f=g \circ h$ on $B_{\mathbb{C}}^{f}$. Then $f=g$.

Proof. This follows from the Rigidity Theorem' on page 751 of [20] which states that $f$ and $g$ are quasiconformally conjugate (in this theorem parabolic periodic points are allowed. By the assumption, we can modify this conjugacy away from the Julia set to a quasiconformal homeomorphism which agrees with the conformal conjugacy outside a small neighborhood of the Julia set and with the original conjugacy on the forward orbits of the critical points (and which preserves the real line). Using the usual pullback argument, one then obtains a quasiconformally conjugacy which is conformal on the Fatou set. Since the Julia set of a map in $P_{\epsilon}^{b}$ does not carry an invariant line field, see Theorem 1 in [38, it follows that the quasiconformal conjugacy must be conformal. By the normalization imposed on maps in $P_{\epsilon}^{b}$ it follows that $f=g$.

### 3.2. The set of Blaschke products with real critical points forms a ball.

 To show that $\mathcal{P H}(f)$ is connected, our strategy is to prove first that $\mathcal{P} \mathcal{H}^{o}(f)$ is connected. To this end we shall identify $\mathcal{P} \mathcal{H}^{\circ}(f)$ with a space of Blaschke products by means of quasi-conformal surgery.Definition 3.9. For any $n \geq 1$ and $\epsilon \in\{+,-\}$, let $M_{n, \epsilon}$ denote the set of all proper (i.e., the inverse of a compact set is compact) holomorphic maps $A: \mathbb{D} \rightarrow \mathbb{D}$ of degree $n$ of the open unit disc $\mathbb{D}$, preserving the real axis, such that $A$ has $n-1$ distinct real critical points in $(-1,1)$, and such that the sign of $A^{\prime}(-1)$ is $\epsilon$. Since $A \in M_{n, \epsilon}$ maps $\mathbb{R}$ into itself, it can be written as $z \mapsto \sigma \prod_{i=1}^{n} \frac{z-a_{i}}{1-\bar{a}_{i} z}$, where $\sigma \in\{-1,1\}$ depending on $\epsilon$ and on the parity of $n$, and where $\left\{a_{i}, 1 \leq i \leq n\right\}$ is a subset of $\mathbb{D}$ which is symmetric with respect to the real axis. Therefore $M_{n, \epsilon}$ can be considered as a subset of $\mathbb{D}^{n}$ and is thus supplied with the induced topology.

Not every map $A \in M_{n, \epsilon}$ has a fixed point, but by the Schwarz-Pick lemma any holomorphic map $A: \mathbb{D} \rightarrow \mathbb{D}$ has at most one fixed point. If $A$ has finite degree then it extends to maps $\partial \mathbb{D}$ to itself. If $A: \mathbb{D} \rightarrow \mathbb{D}$ does not have a fixed point in $\mathbb{D}$ then it follows from Denjoy-Wolff that there exists a unique fixed point on $\partial \mathbb{D}$ which attracts all points in $\mathbb{D}$.

Definition 3.10. Let $M_{n, \epsilon}^{o}$ (respectively $M_{n, \epsilon}^{\Sigma}$ ) be the set of maps $A \in M_{n, \epsilon}$ with the additional property that $A(0)=0$ (respectively so that $c_{1}=0$ ) where $-1<c_{1}<$ $\cdots<c_{n-1}<1$ are the critical points of $A$.

When $n \geq 3$, the assumption that $\left\{a_{i}, 1 \leq i \leq n\right\}$ is a subset of $\mathbb{D}$ which is symmetric with respect to the real axis does not imply that all critical points of $A$ are real.

Lemma 3.11. $M_{n, \epsilon}^{o}$ and $M_{n, \epsilon}^{\Sigma}$ are homeomorphic to an open Euclidean ball with (real) dimension equal to the number of critical points of maps in $M_{n, \epsilon}$, i.e., equal to $n-1$. Moreover, $M_{n, \epsilon}$ is homeomorphic to an open Euclidean ball of (real) dimension $n$.

Proof. Using the same argument as the proof of Lemma 3.1 in [30] or of Corollary II.4.1 in [26], one sees that maps in $M_{n, \epsilon}^{o}$ can be reparametrized by their critical values. Since for any map $A \in M_{n, \epsilon}^{o}$ there exists a unique Möbius transformation $M$ so that $A \circ M \in M_{n, \epsilon}^{\Sigma}$ it follows that the space $M_{n, \epsilon}^{\Sigma}$ also has dimension $n-1$. Since any map $A \in M_{n, \epsilon}$ can be written in the form $B \circ M$ where $B \in M_{n, \epsilon}^{o}$ and $M$ is a Möbius transformation, this implies that the space $M_{n, \epsilon}$ is homeomorphic to a Euclidean ball of dimension $n$. Alternatively, this follows from the fact that $M_{n, \epsilon}^{o}$ can be parametrized by their critical points, see [44] and that any $A \in M_{n, \epsilon}$ can be written in the form $M \circ B$ where $B \in M_{n, \epsilon}^{o}$ and $M$ is a Möbius transformation.
3.3. The set $\mathcal{P H}{ }^{o}(f)$ is homeomorphic to a ball. Take $f \in \mathcal{K} \mathcal{S}$ and let us associate spaces $\mathcal{M}(f)$ and $\mathcal{M}^{o}(f)$ to $\mathcal{P} \mathcal{H}^{o}(f)$. For this we will consider $f \in P^{b} \cap \mathcal{K} \mathcal{S}$ as a map acting on the complex plane, and define $B_{\mathbb{C}}(f)$ as the set of points in the complex plane whose iterates converge to periodic attractors (or parabolic points) of $f$. Let $U_{1}, U_{2}, \ldots, U_{m}$ be the components of $B_{\mathbb{C}}(f) \subset \mathbb{C}$ which contain critical points, and let $n_{i}$ be the number of critical points in $U_{i}$. For each $i$ let $s_{i}$ be the minimal positive integer such that $f^{s_{i}}\left(U_{i}\right)=U_{i^{\prime}}$ for some $1 \leq i^{\prime} \leq m$. Note that it is conceivable that some components $U_{i}$ are backward iterates of the immediate basin of the periodic attractor. Let us consider the space $\mathcal{M}(f)=\prod_{i=1}^{m} M_{n_{i}, \epsilon_{i}}$, where $\epsilon_{i}$ denotes the sign of $\left(f^{s_{i}}\right)^{\prime}$ at the left endpoint of $B_{i}:=U_{i} \cap \mathbb{R}$. (If this endpoint is a critical point, then $\epsilon_{i}$ is the sign of the second derivative at this point.) An element $\mathbf{A}=\left(A_{1}, \ldots, A_{m}\right) \in \mathcal{M}(f)$ will be viewed as a dynamical system on the disjoint union of $m$ copies of the unit disk,

$$
\mathbf{A}: \bigcup_{i=1}^{m} \mathbb{D}_{i} \rightarrow \bigcup_{i=1}^{m} \mathbb{D}_{i}, \quad \text { where } \mathbb{D}_{i}=\mathbb{D} \times\{i\}
$$

such that $\mathbf{A}(z, i)=\left(A_{i}(z), i^{\prime}\right)$, where $i^{\prime}$ is as above. Let us say that $\mathbf{A} \sim \tilde{\mathbf{A}}$ if they are conjugate to each other via a component-preserving conformal map $\varphi: \bigcup \mathbb{D}_{i} \rightarrow \bigcup \mathbb{D}_{i}$ such that for each $1 \leq i \leq m, \varphi \mid \mathbb{D}_{i}$ is a real symmetric (i.e., $\left.\varphi(z)=\overline{\varphi(\bar{z})}\right)$ map whose restriction to the real line preserves the orientation.

Let $\mathcal{M}^{o}(f)$ denote the subset of $\mathcal{M}(f)$ consisting of maps $\mathbf{A}=\left(A_{1}, A_{2}, \ldots, A_{m}\right)$ with the following property: if $\mathbf{A}^{k}$ maps $\mathbb{D}_{i}$ onto itself, then $\mathbf{A}^{k}$ has a fixed point in $\mathbb{D}_{i}$. In other words, if $U_{i_{1}}, U_{i_{2}}, \cdots, U_{i_{k}}$ is a cycle of attracting basins of $f$, then we require that $A_{i_{k}} \circ \cdots \circ A_{i_{1}}$ has a fixed point. This means that up to a Möbius transformation we can assume that the periodic points in $U_{i}$ correspond to 0 . It follows that $\mathcal{M}^{o}(f) / \sim=\prod_{i=1}^{m} M_{n_{i}, \varepsilon_{i}}^{\delta_{i}}$ where $M_{n_{i}, \varepsilon_{i}}^{\delta_{i}}=M_{n_{i}, \varepsilon_{i}}^{o}$ when $U_{i}$ contains a periodic attractor and $M_{n_{i}, \varepsilon_{i}}^{\delta_{i}}=M_{n_{i}, \varepsilon_{i}}^{o}$ otherwise.

Let us define a map

$$
\begin{equation*}
\Theta: \mathcal{P} \mathcal{H}^{o}(f) \rightarrow \mathcal{M}^{o}(f) / \sim \tag{2}
\end{equation*}
$$

as follows. For $g \in \mathcal{P} \mathcal{H}^{o}(f)$, let $U_{i}(g) \subset \mathbb{C}, i=1,2, \ldots, m$, be the components of $B_{\mathbb{C}}(g)$ containing critical points corresponding to the sets $U_{i}$ from above. For each
$i$, let $\varphi_{i}: U_{i}(g) \rightarrow \mathbb{D}$ be some real-symmetric conformal map whose restriction to the real axis is orientation-preserving, and let

$$
\begin{equation*}
A_{i}(g)=\varphi_{i^{\prime}} \circ g^{s_{i}} \circ \varphi_{i}^{-1} \tag{3}
\end{equation*}
$$

Then define

$$
\begin{equation*}
\Theta(g)=\left[\left(A_{1}(g), A_{2}(g), \ldots, A_{m}(g)\right)\right], \tag{4}
\end{equation*}
$$

where $[A]$ denotes the equivalence class of $A$. Note that the space $\mathcal{M}^{\circ}(g)$ associated to any map in $g \in \mathcal{P} \mathcal{H}^{\circ}(f)$ is the same and so this definition makes sense.

Lemma 3.12. The map $\Theta$ defines a homeomorphism between $\mathcal{P} \mathcal{H}^{o}(f)$ and $\mathcal{M}^{o}(f) / \sim$. In particular, $\mathcal{P H}^{\circ}(f)$ is homeomorphic to an open ball of dimension equal to the number of critical points in $B(f)$.

Proof of Lemma 3.12. Since the sets $U_{i}(g)$ move continuously (in the Carathédory topology) with respect to $g \in \mathcal{P} \mathcal{H}^{\circ}(f)$ the map $\Theta$ is continuous. Here we use the continuous dependence of the Riemann mapping from $U_{i}(g)$ to $\mathbb{D}$ as the simply connected domain $U_{i}(g)$ moves continuously with $g$, see the discussion in Section 5.1 in [25]. By the Rigidity Theorem 3.7, $\Theta$ is injective. Indeed, if $\Theta(g)=\Theta(\tilde{g})$ then $g$ and $\tilde{g}$ are topologically conjugate on $\mathbb{R}$, and moreover they are conformally conjugate near the corresponding periodic attractors. Therefore $g$ and $\tilde{g}$ are affinely conjugate.

Because a continuous bijective map between open subsets of Euclidean spaces is a homeomorphism (due to Brouwer's invariance of domain theorem), it remains to prove that $\Theta$ is surjective. Let $\mathbf{A}=\left(A_{1}, A_{2}, \ldots, A_{m}\right)$ be an element in $\mathcal{M}^{\circ}(f)$. Our aim is to construct a map $g \in \mathcal{P} \mathcal{H}^{o}(f)$ so that $\Theta(g)=[\mathbf{A}]$. To do this, one applies quasi-conformal surgery techniques in a standard fashion. Let us therefore be brief, and refer to the exposition given in Theorem VIII.2.1 of [6] for details. Choose $f_{0} \in$ $\mathcal{P} \mathcal{H}^{\circ}(f)$ and let $U_{i}, s_{i}, n_{i}, \varepsilon_{i}$ be the objects associated to $f_{0}$ as above. Let $\varphi_{i}: U_{i} \rightarrow \mathbb{D}$ be a real-symmetric conformal map sending the periodic attractor in $U_{i}$ to 0 . Then $\varphi_{i^{\prime}} \circ f_{0}^{s_{i}} \circ \varphi_{i}^{-1}: \mathbb{D} \rightarrow \mathbb{D}$ is a map $A_{i}^{o}$ in $M_{n_{i}, \varepsilon_{i}}^{o}$ and $\mathbf{A}^{o}=\left(A_{1}^{o}, A_{2}^{o}, \ldots, A_{m}^{o}\right) \in \mathcal{M}^{o}(f)$. Define a new smooth covering map $\tilde{A}_{i}: \mathbb{D} \rightarrow \mathbb{D}$ as follows. Take discs $\Delta\left(r_{i}\right) \subset \mathbb{D}$ with $r_{i}<1$ sufficiently close to 1 so that $\cup_{i} \Delta\left(r_{i}\right) \times\{i\}$ is mapped into itself by $\mathbf{A}$ and also by $\mathbf{A}^{o}$. Let $\mathrm{Ann}_{i}=\left(\left(A_{i}^{o}\right)^{-1} \Delta\left(r_{i}\right)\right) \backslash \Delta\left(r_{i}\right)$ so that $\mathrm{Ann}_{i}$ is a fundamental annulus of $A_{i}^{o}$. Choose $\tilde{A}_{i}: \mathbb{D} \rightarrow \mathbb{D}$ so that it agrees with $A_{i}$ on $\Delta\left(r_{i}\right)$ and with $A_{i}^{o}$ on $\mathbb{D} \backslash\left(\Delta\left(r_{i}\right) \cup \mathrm{Ann}_{i}\right)$, and so that it is a smooth covering map on the fundamental annulus $\mathrm{Ann}_{i}$. Next define a smooth map $\tilde{f}$ which agrees with $f$ outside $\cup U_{i}$ and which is equal to $f^{-(s-1)} \circ \varphi_{i}^{-1} \circ \tilde{A}_{i} \circ \varphi_{i}$ on $U_{i}$, where $f^{-(s-1)}$ stands for inverse of the conformal map $f^{s-1}: f\left(U_{i}\right) \rightarrow U_{i}^{\prime}$. The smooth map $\tilde{f}$ agrees with $f$ outside $\cup U_{i}$ and is conformal outside the annuli $\varphi_{i}^{-1}\left(\operatorname{Ann}_{i}\right)$. Since the $\tilde{f}$ orbit of each point only hits at most once the fundamental annuli we can choose an invariant ellipse field which agrees with the standard linefield in $\varphi_{i}^{-1}\left(\Delta\left(r_{i}\right)\right)$ and on the complement of $B(f)$. Using the Measurable Riemann Mapping Theorem, we obtain a $K$-quasiconformal
homeomorphism $h$ so that $g:=\tilde{h} \circ \tilde{f} \circ \tilde{h}^{-1}$ is again holomorphic and therefore the required polynomial of the same degree as $f$. Since $g$ and $f$ are conjugate outside $B(f)$, we have $g \in \mathcal{P} \mathcal{H}^{\circ}(f)$.

Let us now show that $\Theta(g)=[A]$. Let $U_{i}^{g}:=h\left(U_{i}\right)$ be the components of $B(g)$. Writing

$$
H_{i}:=\varphi_{i} \circ h^{-1}: h\left(U_{i}\right) \rightarrow \mathbb{D} \text { and } H_{i^{\prime}}:=\varphi_{i^{\prime}} \circ h^{-1}: h\left(U_{i^{\prime}}\right) \rightarrow \mathbb{D}
$$

we have that $H_{i^{\prime}} \circ g^{s_{i}} \circ H_{i}^{-1}: \mathbb{D} \rightarrow \mathbb{D}$ agrees with the Blaschke product $A_{i}$ on $H_{i}\left(\Delta\left(r_{i}\right)\right)$. Moreover, $H_{i}^{-1}$ is conformal on this set. Since the forwards iterates of critical points $A$ are contained in this set, by pulling back via the dynamics, one obtains a sequence of $K$-quasiconformal maps $H_{i, n}: h\left(U_{i}\right) \rightarrow \mathbb{D}, H_{i^{\prime}, n}: h\left(U_{i^{\prime}}\right) \rightarrow \mathbb{D}$ which are conformal on larger and larger subsets of $h\left(U_{i}\right)$ and $h\left(U_{i^{\prime}}\right)$ respectively, and so that $H_{i^{\prime}, n} \circ g^{s_{i}} \circ H_{i, n}^{-1}: \mathbb{D} \rightarrow \mathbb{D}$ agrees with $A$ on corresponding increasing subsets of $\mathbb{D}$. By taking limits, one obtains conformal maps $\hat{H}_{i}: h\left(U_{i}\right) \rightarrow \mathbb{D}, \hat{H}_{i^{\prime}}: h\left(U_{i^{\prime}}\right) \rightarrow \mathbb{D}$ so that $\hat{H}_{i^{\prime}} \circ g^{s_{i}} \circ H_{i}^{-1}=A_{i}$.
3.4. The proof of Theorems 3.4 and 3.6. Take $f \in P_{\epsilon}^{b}$. Assertion (i) in Theorem 3.4 was proved in the previous lemma. So let us prove that $\mathcal{P H}(f) \subset \overline{\mathcal{P} \mathcal{H}^{o}(f)}$ and that there exists a family of maps as in Theorem 3.6.

In Steps 1-4 we find a polynomial map $P \in P_{\epsilon}^{b}$ which will essentially play the role of $f_{1}$. To find such a polynomial, we will first approximate $f$ by a suitable continuous map $g$ whose dynamics is the same as that of $f$ except 'on the basins of periodic attractors'.

Step 1. In this step we find a family of piecewise smooth interval maps $g_{t}, t \in[0,1]$ with $g_{0}=f$ which undergoes the required bifurcations at $t=0$, as $t$ becomes positive, for each of the periodic attractors as in case (1)-(4) of the assumption of Theorem 3.6. Here we ensure that $g_{t}$ agrees with $f$ outside a small neighborhood in $B(f) \cap \mathbb{R}$ of the neutral periodic points $p_{1}, \ldots, p_{k}$. Depending on the sign of $\sigma_{i}$ in the assumption of Theorem 3.6, we choose $g_{t}$ so that as $t$ increases, the neutral point $p_{i}$ undergoes a period-doubling or period-halving bifurcation in case (1), a pitchfork or a reverse pitchfork bifurcation in case (2) and a saddle-node or a reverse saddlenode bifurcation in case (4). If $p_{i}$ is as in case (3), we choose $g_{t}$ so that it is merely piecewise smooth at $p_{i}$, and $g_{t}^{n_{i}}(x)=p_{i}+(1+t)\left(x-p_{i}\right)-\left(x-p_{i}\right)^{2}+O\left(x-p_{i}\right)^{3}$ for $x$ in a one-sided attracting neighborhood of $p_{i}$.

Step 2. First assume that $c$ is a critical point as in case (5) so that $c$ is mapped to the boundary of a component $U$ of $B(f)$ and $f^{-1}(U) \cup\{c\}$ contains a (real) neighborhood of $c$, see Figure 6 on the right. Then we can choose the family $g_{t}$ so that $g_{t}(c)$ is in the interior of $B(f)$ for each $t \in(0,1]$. Let us denote $\hat{B}(f)=\cup \bar{J}$ where the union runs over all connected components $J$ of $B(f)$ and $\bar{J}$ is the closure


Figure 6. A critical point which is mapped to the boundary of a component of the basin of a periodic attractor. The cases discussed in step 2 and 3 of the proof of Theorem 3.6 are shown on the left respectively right.
of $J$. Note that the only difference between $g_{t}$ and $f$ is that some basins are merged or split in two and therefore $\hat{B}\left(g_{t}\right)=\hat{B}(f)$ for each $t \in[0,1]$.

Step 3. Pick $t_{0}>0$ small and let $\hat{g}=g_{t_{0}}$. Next we need to take special care of case (5) in the situation when there exists a critical point $c$ so that $f(c)$ is mapped into the boundary of a component $U$ of $B(f)$ so that $f^{-1}(U) \cup\{c\}$ contains no (real) neighborhood of $c$, as is drawn in Figure 6 on the right. In this case choose a semi-conjugacy $h: I \rightarrow I$ (i.e., $h$ is continuous, monotone and surjective) and a continuous $b$-modal interval map $g$ so that $h \circ g=\hat{g} \circ h$ where $h^{-1}(x)$ is a point, except if $x \in \cup_{n \geq 0} \hat{g}^{-n}(c)$ for any $c$ as in the previous sentence. Therefore we 'glue-in' intervals in the backward orbit of $c$ allowing us to 'move' $\hat{g}(c)$ into the interior of $B(\hat{g})$. We can do this so that $g(c)$ now is mapped into the interior of the component of $B(g)$ corresponding to $U$. It is possible that a critical point $c^{\prime}$ is eventually mapped to $c$. In this case $\cup_{n>0} \hat{g}^{-n}(c)$ contains $c^{\prime}$, and we proceed in the same way. Also note that there exists an order preserving homeomorphism $h_{0}$ of $I \backslash \hat{B}(f)$ to $I \backslash \hat{B}(g)$ so that $h_{0} \circ f=g \circ h_{0}$ on $I \backslash \hat{B}(f)$.

Step 4. Note that $g$ has no wandering intervals and also no inessential periodic attractors. Hence, by the Fullness Theorem II.4.1 in [26], there exists a polynomial $P \in P_{\epsilon}^{b}$ which is topologically conjugate to $g$ and so that each of its periodic orbits is hyperbolic and so that none of its critical point is mapped to the boundary of a component of the basin of $P$. Note that $f$ and $P$ are conjugate outside their basins: there exists an order preserving homeomorphism $h_{1}$ of $I \backslash \hat{B}(f)$ to $I \backslash \hat{B}(P)$ so that $h_{0} \circ f=P \circ h_{0}$ on $I \backslash \hat{B}(f)$. In other words, $f$ and $P$ are conjugate outside their basin of attractors and components of $B_{\mathbb{C}}(f)$ and $B_{\mathbb{C}}(P)$ correspond to each other in the following manner:
(1) Each component of $B_{\mathbb{C}}(f)$ containing a hyperbolic periodic point corresponds to a unique component of $B_{\mathbb{C}}(P)$.
(2) Each component of $B_{\mathbb{C}}(f)$ of the basin of a neutral periodic point $p_{i}$ of $f$ (i.e., of a petal) corresponds to a component of the basin of a hyperbolic periodic point for $P$. However, two touching adjacent basins (petals) of $B_{\mathbb{C}}(f)$ may
correspond to one component of $B_{\mathbb{C}}(P)$ (so these petals are merged). This happens in case (1), (2) and (4) when the corresponding sign $\sigma_{i}$ is negative.
(3) Each critical point of $f$ which is eventually mapped into the boundary of the basin of $f$, corresponds to a critical point of $P$ which is contained in the interior of $B_{\mathbb{C}}(P) \cap \mathbb{R}$.

We will use $P$ to construct the required family $f_{\mu}$, but we do not claim that the polynomial $P$ is close to $f$, even in the $C^{0}$ topology.

Let $U_{1}, \ldots, U_{m}$ be the components of $B_{\mathbb{C}}(P)$ which contain critical points. We will use the homeomorphism $\Theta: \mathcal{P H}^{o}(P) \rightarrow \mathcal{M}^{o}(P) / \sim$ defined in Subsection 3.3 to construct a family of polynomials through $f$. To do this, we construct in Step 5 a family of polynomials $F_{\mu}$ which may have a much higher degree than $P$. This family $F_{\mu}$ will be used to obtain a one-parameter family of Blaschke products, i.e., a one-parameter family in $\mathcal{M}^{o}(P) / \sim$.

Step 5. Let $U_{1}^{f}, \ldots, U_{k}^{f}$ be the components of $B_{\mathbb{C}}(f)$ which contain critical points and let $U_{k+1}^{f}, \ldots, U_{k^{\prime}}^{f}$ be the components $U$ of $B_{\mathbb{C}}(f)$ for which there exists a critical point $c$ so that $f(c) \in \partial(U \cap \mathbb{R})$. Note that $k^{\prime}=m$ by the construction of $P$. Let $V_{1}^{f}, \ldots, V_{n}^{f}$ be the components of $B_{\mathbb{C}}(f)$ containing the forward iterates of the sets $U_{1}^{f}, \ldots, U_{m}^{f}$. Let $X=\cup_{i} \partial\left(V_{i}^{f} \cap \mathbb{R}\right)$. Moreover, for each critical point $c$ of $f$, let $n(c)$ be the smallest integer for which there exits a component $U$ of $B_{\mathbb{C}}(f)$ so that $f^{n(c)}(c) \in \partial(U \cap \mathbb{R})$ and if there exists no such integer let $n(c)=0$. Let $Y=\cup_{c}\left\{f^{i}(c) ; 0 \leq i \leq n(c)-2\right\}$ where the union is taken over all critical points $c$ of $f$.

Next consider a real polynomial map $Q$ which is zero on the set $X \cup Y$, so that $Q^{\prime}$ is zero on $\operatorname{Crit}(f)$ and let $F_{\mu}=f+\mu Q$. Note that the degree of $Q$ and therefore of $F_{\mu}$ might be much larger than that of $f$. By the choice of $Q$ each periodic point of $f$ on the boundary of $B_{\mathbb{C}}(f)$ is still a periodic point for $F_{\mu}$ and each critical point of $f$ is still a critical point of $F_{\mu}$. By Theorem VI.1.2 in [6] one can choose $Q$ so that as $\mu \in[0,1]$ becomes positive, $F_{\mu}$ undergoes all the bifurcations of neutral periodic points and homoclinic orbits required in cases (1)-(5) of the theorem. In particular, each of the components $U_{1}, \ldots, U_{m}$ of $B_{\mathbb{C}}(P)$ corresponds in a unique way to a component $U_{1, \mu}, \ldots, U_{m, \mu}$ of $B_{\mathbb{C}}\left(F_{\mu}\right)$ when $\mu \in(0,1]$ is small. Note that $F_{\mu}$ can have many more attractors than $P$ and $F_{\mu}$ may not necessarily be conjugate to $f$ and we do not claim that $U_{i, 0}=U_{i}$ either. Each of the attractors in $U_{1, \mu}, \ldots, U_{m, \mu}$ is hyperbolic and $U_{i, \mu}$ depend continuously on $\mu \in\left(0, \mu_{0}\right]$ for $\mu_{0}>0$ small. As below Lemma 3.11, we have that $F_{\mu}^{s_{i}}\left(U_{i, \mu}\right)=U_{i^{\prime}, \mu}$ and we can associate to each of the maps $F_{\mu}^{s_{i}}: U_{i, \mu} \rightarrow U_{i^{\prime}, \mu}$ a Blaschke product $A_{i}$ as in (3). In this way we obtain

$$
\left[\left(A_{1, \mu}, \ldots, A_{2, \mu}\right)\right] \in \mathcal{M}^{o}(P) / \sim
$$

as in Lemma 3.11. Note that $\left[\left(A_{1, \mu}, \ldots, A_{2, \mu}\right)\right]$ depends continuously on $\mu \in\left(0, \mu_{0}\right]$, because the domains $U_{i, \mu}$ vary continuously.

Step 6. Next consider the map

$$
\Theta: \mathcal{P H}^{o}(P) \rightarrow \mathcal{M}^{o}(P) / \sim
$$

from equation (2) and (4) and define

$$
f_{\mu}=\Theta^{-1}\left[\left(A_{1, \mu}, \ldots, A_{m, \mu}\right] .\right.
$$

Therefore $f_{\mu}$ is obtained from $P$ by gluing in the components $U_{i}$ Blaschke products which are obtained from the family $F_{\mu}$. By definition $f_{\mu} \in P_{\epsilon}^{b}$ and by construction, when $\mu>0$, all periodic orbits of $f_{\mu}$ are hyperbolic and no critical point point of $f_{\mu}$ is eventually mapped in the boundary of a component of $f_{\mu}$. Moreover, $f_{\mu} \in$ $\mathcal{P} \mathcal{H}^{o}\left(f_{\mu_{0}}\right)=\mathcal{P} \mathcal{H}^{o}(P)$ for each $\mu \in\left(0, \mu_{0}\right]$.

Step 7. Let us now show that $f_{\mu}$ tends to $f$ as $\mu \rightarrow 0$ (in the sense that the coefficients of $f_{\mu}$ converge to those of $f$ ). To see this, let $V_{1, f_{\mu}}, \ldots, V_{n, f_{\mu}}$ be the components of $B\left(f_{\mu}\right)$ which contain forward iterates of critical points of $f_{\mu}$ and let $V_{1, \mu}, \ldots, V_{n, \mu}$ be the corresponding components of $B\left(F_{\mu}\right)$. By construction there exists a family of conformal homeomorphisms

$$
h_{\mu}: V_{1, f_{\mu}} \cup \cdots \cup V_{n, f_{\mu}} \rightarrow V_{1, \mu} \cup \cdots \cup V_{n, \mu}
$$

so that $h_{\mu} \circ f_{\mu}=F_{\mu} \circ h_{\mu}$ on this set. Moreover, $h_{\mu}$ depends continuously on $\mu \in\left(0, \mu_{0}\right]$. Moreover, even though some of these components pinch (at places where $f$ has a parabolic periodic point), the conformal homeomorphisms $h_{\mu}$ have a well-defined conformal 'limit'

$$
h: V_{1, f_{0}} \cup \cdots \cup V_{n, f_{0}} \rightarrow V_{1} \cup \cdots \cup V_{n}
$$

so that $h \circ f_{0}=F_{0} \circ h$. It follows that $f_{0}$ and $F_{0}=f$ are conformally conjugate on the basin of periodic attractors. Hence by the Generalized Rigidity Theorem 3.8 we obtain that $f_{0}=f$ and that $f_{\mu} \rightarrow f$ as $\mu \rightarrow 0$.

This completes the proof of Theorems 3.4 and 3.6 .
Note that we do not state that the Julia set of $F_{\mu}$ is related to that of $f$. This enables us to avoid using the techniques employed in [14, 15]. Moreover, in general, it is not obvious how to deform a map with an attracting and a repelling orbit to one with a parabolic orbit, or vice versa to deform a map with a parabolic point to a 'subhyperbolic' map in such a way that the Julia set remains topologically the same, see [14, 15]. We are not concerned with this question.

## 4. Partial conjugacy classes within the space $\mathcal{S}^{b}$ of stunted SAWTOOTH MAPS

4.1. Definition of the space of stunted sawtooth maps $\mathcal{S}^{b}$. Fix the number of turning points $b$ and the shape $\epsilon$ of the polynomials we will consider. From now on we will drop the symbol $\epsilon$. Following [31], it will be useful to introduce a space of piecewise linear maps with $b$ (possibly touching) plateaus. Fix the slope $\lambda=b+2$
and let $e=b \lambda /(\lambda-1)$. An elementary calculation shows that there exists a unique $b$-modal piecewise linear map $S_{0}$ (a 'sawtooth map') with shape $\epsilon$ and $b$ turning points $c_{1}, \ldots, c_{b}$ at $-b+1,-b+3, \ldots, b-3, b-1$ with

- $b+1$ intervals of monotonicity

$$
I_{0}=\left[-e, c_{1}\right], I_{1}=\left[c_{1}, c_{2}\right], \ldots, I_{b}=\left[c_{b}, e\right] ;
$$

- slope $\pm \lambda$ and extremal values $\pm \lambda$;
- and such that $S_{0}(\{-e, e\}) \subset\{-e, e\}$.


The space of $\mathcal{S}^{b}$ of stunted sawtooth maps consists of continuous maps $T$ with plateaus $Z_{i, T}, i=1, \ldots, b$, which are obtained from $S_{0}$ and satisfying

- $Z_{i, T}$ is a closed symmetric interval around $c_{i}$;
- $T$ and $S_{0}$ agree outside $\bigcup_{i} Z_{i, T}$;
- $T \mid Z_{i, T}$ is constant and $T\left(Z_{i, T}\right) \in[-e, e]$;
- $Z_{i, T}$ have pairwise disjoint interiors.


Maps in $\mathcal{S}^{b}$ are allowed to have touching plateaus (i.e., plateaus with one point in common). We allow plateaus to touch because, as we will see later on, such maps $T$ correspond to polynomial maps for which several critical points lie in one component of the basin of an attracting periodic point.

Note that if $T \in \mathcal{S}^{b}$ has touching plateaus, then the union of these touching plateaus is called a block of plateaus. If $T$ has touching plateaus, then it is constant on at least one lap of $S_{0}$. In this case $T$ is $b$-modal only in a degenerate sense.

It is convenient to use the $b$ 'signed' extremal values $\zeta \in[-e, e]^{b}$ to parametrize $\mathcal{S}^{b}$ :

$$
\zeta_{i}= \begin{cases}T\left(Z_{i, T}\right) & \text { if } S_{0} \text { assumes a maximum at } c_{i}, \\ -T\left(Z_{i, T}\right) & \text { if } S_{0} \text { assumes a minimum at } c_{i}\end{cases}
$$

Sometimes we will denote by $T_{\zeta}$ the map $T$ with parameters $\zeta=\left(\zeta_{1}, \ldots, \zeta_{b}\right)$. Note that decreasing $\zeta_{i}$ results in widening the corresponding plateau of $T_{\zeta}$ and that $\zeta_{i}+\zeta_{i+1}$ is equal to the length of the convex hull of $T\left(Z_{i}\right)$ and $T\left(Z_{i+1}\right)$. Hence

$$
\zeta_{i} \geq-\zeta_{i+1} \quad \text { for } i=1, \ldots, b-1
$$

with equality when the plateaus $Z_{i}$ and $Z_{i+1}$ touch. Let us denote by [ $Z_{i}, Z_{i+1}$ ] the convex hull of the plateaus $Z_{i}$ and $Z_{i+1}$. Thus we can identify $\mathcal{S}^{b}$ with

$$
\left\{\zeta=\left(\zeta_{1}, \ldots, \zeta_{b}\right) ; \zeta_{i} \in[-e, e], \zeta_{i} \geq-\zeta_{i+1}\right\}
$$

We define $T<\tilde{T}$ if for the corresponding parameters $\zeta_{i} \leq \tilde{\zeta}_{i}$ for all $i=1, \ldots, b$ with at least one inequality.

Proposition 4.1. The map $\zeta=\left(\zeta_{1}, \ldots, \zeta_{b}\right) \rightarrow h_{\text {top }}\left(T_{\zeta}\right)$ is non-decreasing in each coordinate.

Remark 4.2. In [4] we prove that the analogous statement is false for $P^{b}$ for $b \geq 2$. That is, if the cubic family is parametrized by its critical values $a, b$, then the map $(a, b) \mapsto h_{\text {top }}\left(f_{a, b}\right)$ is not monotone in each of its parameters separately.

A consequence of this proposition is that $\left\{T \in \mathcal{S}^{b} ; h_{t o p}(T)=s\right\}$ is contractible, see Theorem 6.1 in [31].

Proof of Proposition 4.1. Increasing a parameter $\zeta_{i}$ makes a plateau narrower, and affects none of the orbits that never enter $Z_{i}$. Therefore only new orbits are created, and none destroyed. Hence entropy is non-decreasing in each $\zeta_{i}$.
4.2. The preplateau and the basin of a map $T$. We define the basin $B(T)$ of a map $T \in \mathcal{S}^{b}$ exactly as before, see Definition 3.1. Since maps $T$ have plateaus, we also introduce a related notion: we define the preplateau $W(T)$ of a map $T$ to be the set of points $x$ which eventually map into the interior of the union of the plateaus of $T$, i.e.,

$$
\begin{equation*}
W(T)=\bigcup_{k \geq 0} T^{-k}\left(\operatorname{int}\left(\cup_{i=1}^{b} Z_{i, T}\right)\right) \tag{5}
\end{equation*}
$$

Because we allow the possibility of plateaus touching each other, we take the interior of the union rather than the union of the interiors. We say that a component $W$ of $W(T)$ is periodic of period $s$ if $T^{s}(W) \subset \bar{W}$. A periodic point $p$ of $T$ is called hyperbolic if its orbit enters the interior of a plateau of $T$. The following elementary lemma explains how the sets $W(T)$ and $B(T)$ are related.

Lemma 4.3. Let $T \in \mathcal{S}^{b}$. Then $W(T)$ is open and dense. Moreover, for $W^{\prime}$ and $W^{\prime \prime}$ are components of $W(T)$,
(1) if $T^{n}\left(W^{\prime}\right)$ intersects a boundary point $y$ of $W^{\prime \prime}$, then $T^{n}\left(W^{\prime}\right)=\{y\}$;
(2) if $T^{n}\left(W^{\prime}\right) \cap W^{\prime \prime} \neq \emptyset$, then $T^{n}\left(W^{\prime}\right) \subset W^{\prime \prime}$;
(3) if $W^{\prime}$ is periodic of period s, then $T^{s}\left(\partial W^{\prime}\right) \subset \partial W^{\prime}$ and either $T^{s}\left(W^{\prime}\right) \subset W^{\prime}$ or $T^{s} \mid W^{\prime}$ is constant;
(4) $W^{\prime}$ is either eventually mapped into a periodic component $W^{\prime \prime}$ of $W(T)$ (with $T^{s}\left(W^{\prime \prime}\right) \subset W^{\prime \prime}$ for some s) or there exists $n$ so that $T^{n}\left(W^{\prime}\right)$ is equal to a point;
(5) if $T^{s}\left(W^{\prime}\right) \subset W^{\prime}$, then $W^{\prime}$ contains precisely one periodic point $p$ (so that some forward iterate is contained in the interior of a plateau) and $T^{k s}(x) \rightarrow p$ for every point in $x \in W^{\prime}$;
(6) $B(T)$ is equal to the set of points which are eventually mapped into a periodic plateau. Moreover, for each component $B$ of $B(T)$
(a) there exists a sequence of touching components $W_{i}$ of $W(T)$ so that $\bigcup_{i \in \mathcal{I}} W_{i} \subset B \subset \bigcup_{i \in \mathcal{I}} \bar{W}_{i}$, where $\mathcal{I}$ is an at most countable index set;
(b) if each periodic orbit of $T$ is hyperbolic (i.e., disjoint from $\partial\left(\cup_{i} Z_{i}\right)$ ), then $B$ is equal to a component of $W(T)$;
(7) if $W^{\prime}$ and $W^{\prime \prime}$ touch and one of them intersects a component $B$ of the basin of a periodic attractor, then both of them are contained in $B$.

Proof. Openness of $W(T)$ follows from the definition. Since the complement of $W(T)$ is forward invariant and $T$ is expanding on this complement, the set $W(T)$ is dense. To prove (1), take $x \in W^{\prime}$ so that $y:=T^{n}(x) \in \partial W^{\prime \prime}$. Note that we can assume that $n$ is 'minimal', i.e., there exists no $0<k<n$ such that $T^{k}(x)$ is in the boundary of a component of $W(T)$. Since $y \in \partial W^{\prime \prime}$ no iterate of $y$ is mapped into the interior of a plateau. Since $x \in W$ and $y=T^{k}(x)$ there exists $0 \leq l<n$ so that $f^{l}(x)$ is mapped in the interior of a plateau. It follows that the interior $H$ of the component of $T^{-n}(y)$ containing $x$ is non-empty. Since $x \in W(T)$ and so $x$ is eventually mapped into the $\operatorname{int}\left(\cup_{i=1}^{b} Z_{i, T}\right)$, it follows that $H \subset W(T)$ (here we use the minimality of $n$ ). Let $x^{\prime}$ be an endpoint of $H$. Then $T^{n}$ is not locally constant near $x^{\prime}$ and so $x^{\prime}, T\left(x^{\prime}\right), \ldots, T^{n-1}\left(x^{\prime}\right) \notin \operatorname{int}\left(\cup_{i=1}^{b} Z_{i, T}\right)$. Since $T^{n}\left(x^{\prime}\right)=y \notin W(T)$ and therefore $y, T(y), \ldots \notin \operatorname{int}\left(\cup_{i=1}^{b} Z_{i, T}\right)$, it follows that $T^{k}\left(x^{\prime}\right) \notin \operatorname{int}\left(\cup_{i=1}^{b} Z_{i, T}\right)$ for all $k \geq 0$ and therefore $x^{\prime} \notin W(T)$. Hence $H=W^{\prime}$ and so $T^{n}\left(W^{\prime}\right)=\{y\}$, proving Assertion (1). This implies (2), (3) and (4) because each component of $W(T)$ is open, because $W(T)$ is backward invariant and because $T$ has only finitely many plateaus. Note that if $T^{s}\left(W^{\prime}\right) \subset W^{\prime}$ then by Assertion (1) the one-sided slope of $T^{s}$ at the endpoints of $W$ is $>1$. Therefore $T^{s}$ has a fixed point $p \in W^{\prime}$ at which the map is locally constant. If $T^{2 s} \mid W^{\prime}$ has another fixed point, then $T^{2 s} \mid W^{\prime}$ also has a repelling fixed point which is impossible since $W^{\prime} \subset W(T)$. Since $T^{s}$ is locally constant at $p$, there exists an interval neighborhood $U_{0}$ of $p$ so that $T^{s}\left(U_{0}\right) \subset U_{0}$ and so that $T^{n}(x) \rightarrow p$ for each $x \in U_{0}$. If we denote by $U_{n}$ the component of $T^{-s n}\left(U_{0}\right)$ containing $U_{0}$, we have $U_{n+1} \supset U_{n} \supset \cdots \supset U_{0}$. Therefore $T^{s}$ maps $U=\bigcup U_{n}$ into itself and $T^{s}(\partial U) \subset \partial U$. Since the only fixed point of $T^{2 s} \mid J$ is $p$, it follows that $U=W$. This proves assertion (5). To prove (6) note that each attracting periodic orbit of $T$ necessarily intersects $\cup_{i=1}^{b} Z_{i, T}$ because $T$ and the unstunted sawtooth map $S_{0}$ agree outside this set. If this periodic orbit intersects int $\left(\cup_{i=1}^{b} Z_{i, T}\right)$, then each component $B$ of the basin of this periodic orbit coincides with a component $W^{\prime}$ of $W(T)$. On the other hand, if this periodic orbit does not intersect the interior of $\cup_{i=1}^{b} Z_{i, T}$, then the backward orbit of this periodic orbit is not contained in $W(T)$ and then components of $B(T)$ are contained in $\bigcup \bar{W}_{i}$ where $W_{i}$ are adjacent components of $W(T)$. This situation is clarified in Example 4.8 below. The final statement holds because if $W^{\prime}$ and $W^{\prime \prime}$ have a boundary point in common, then $T^{i}\left(W \cup W^{\prime}\right)$ is a single point for some $i>0$.
4.3. The sets $\langle\boldsymbol{T}\rangle,[\boldsymbol{T}]$ and $[\boldsymbol{T}]_{b}$. As in Definition 3.2 we say that $T$ and $\tilde{T}$ are partially conjugate if there exists an orientation preserving homeomorphism which maps $Z_{i}$ to $\tilde{Z}_{i}$, which maps $B(T)$ to $B(\tilde{T})$ and which conjugates $T$ and $\tilde{T}$ outside these sets. We define $\mathcal{P H}(T)$ to be the set of $\tilde{T} \in \mathcal{S}^{b}$ which are partially conjugate to $T$. Since such maps have plateaus, we also define

$$
\begin{equation*}
\langle T\rangle=\left\{\tilde{T} \in \mathcal{S}^{b} ; W(\tilde{T})=W(T)\right\} \quad \text { and } \quad[T]=\operatorname{closure}(\langle T\rangle) \tag{6}
\end{equation*}
$$

Note that if $\tilde{T} \in\langle T\rangle$ then $\langle\tilde{T}\rangle=\langle T\rangle$.
Of course, $\langle T\rangle$ and $\mathcal{P H}(T)$ are closely related:
Lemma 4.4. If each periodic orbit of $T$ is hyperbolic, then $\langle T\rangle \subset \mathcal{P H}(T) \subset[T]$.
Example 4.8 shows that the assumption that all periodic orbits of $T$ are hyperbolic is required.

Proof. By Lemma 4.3(6)(b), if each periodic point of $T$ is hyperbolic, then each component of $B(T)$ is a component of $W(T)$. It follows that if $W(\tilde{T})=W(T)$ then the basins of $\tilde{T}$ and $T$ are the same, and therefore $\tilde{T}$ and $T$ are partially conjugate. Now assume that $\tilde{T} \in \mathcal{P} \mathcal{H}(T)$. Since $\tilde{T}$ and $T$ agree outside plateaus, and periodic points are dense outside the basins, the partial conjugacy outside $B(\tilde{T})$ and $B(T)$ has to be the identity map. It follows that if $\tilde{T}$ also has only hyperbolic periodic orbits, then $B(\tilde{T})$ is also open and $W(T)=W(\tilde{T})$. If $\tilde{T}$ has one or more nonhyperbolic periodic orbits, then by widening the corresponding plateaus one obtains a sequence of maps $\hat{T}_{n} \in\langle T\rangle \cap \mathcal{P} \mathcal{H}(T)$ with $\hat{T}_{n} \rightarrow \tilde{T}$. This implies the lemma.

Below we shall show that $\langle T\rangle$ is contained in a hyperplane $V_{T}$, and that each map $[T] \backslash\langle T\rangle$ either has touching plateaus or an orbit of one of the following special types:

Definition 4.5 (Homoclinic orbit). We say that $T$ has a homoclinic orbit (hc), if some iterate of a plateau is mapped to the boundary of a component of $W(T)$.

Definition 4.6 (Hyperbolic/Saddle-node/Period Doubling/Pitchfork). Assume that $p$ is an attracting periodic point of $T$ (of minimal period s) and therefore the orbit of $p$ enters a plateau of $T$. We say that $p$ is hyperbolic if its orbit enters the interior of a plateau of $T$. Otherwise, we say $p$ is non-hyperbolic and a
(1) saddle-node ( $s n$ ) if $p$ is in the boundary of a plateau of $T^{s}$ and $S_{0}^{s}$ is orientation preserving near $p$; if there exists a neighborhood $U$ of $p$ so that both components of $U \backslash\{p\}$ are in the immediate basin of a periodic attractor, then we say that $p$ is a saddle-node merging two basins, see the left panel of Figure 7;
(2) period-doubling ( $p d$ ) if $p$ is in a plateau of $T^{s}$ and $S_{0}^{s}$ is orientation reversing near p, see Figure 8:




Figure 7. Left: a saddle-node merging two basins. Middle: a map with two adjacent plateaus with a pitch-fork fixed point; in this case, there are nearby maps with two attracting periodic points of the same period. Right: a map with a homoclinic orbit.


Figure 8. Left: a period doubling fixed point; nearby maps have a periodic point of period two. Right: another period doubling fixed point, this time between two plateaus.
(3) pitch-fork ( $p f$ ) if $p$ is in the interior of a plateau of $T^{s}$ and $S_{0}^{s}$ is orientation preserving near $p$, see the middle panel of Figure 7 .

In the latter case, $p$ is either the common boundary of two touching plateaus of $T$, or the orbit of $p$ hits at least twice a boundary point of a plateau.

Definition 4.7. ( $\min [T]$ and $\left.[T]_{b}\right)$ Given $T \in \mathcal{S}^{b}$ we define the sets

$$
\min [T]=\left\{T^{\prime} \in[T] \text {; there is no } T^{\prime \prime} \in[T] \text { with } T^{\prime \prime}<T^{\prime}\right\} .
$$

Here $<$ is the partial ordering on the space $\mathcal{S}^{b}$ defined above Proposition 4.1. To deal with the situation that plateaus of $T$ touch, we also define

$$
[T]_{b}=\left\{T^{\prime} \in \min [T] \quad ; \quad \begin{array}{l}
\text { no plateau } Z^{\prime} \text { of } T^{\prime} \text { is eventually mapped } \\
\text { into } \cup_{i=1}^{b} \operatorname{int}\left(Z_{i}^{\prime}\right) \text { by some positive iterate } .
\end{array}\right\}
$$

Clearly

$$
\emptyset \neq[T]_{\mathrm{b}} \subset \min [T] \subset[T] .
$$

We illustrate these definitions in Examples 4.84.11, which helps in obtaining a general description of the set $\langle T\rangle$. We should emphasize that some parts of the boundary of $\langle T\rangle$ are contained in $\langle T\rangle$ whereas others are not. Figures 9,10 and 11 give explicit descriptions of $\langle T\rangle$ in a number of situations.

Example 4.8. Figure 9 illustrates the definition of $\langle T\rangle$ in the unimodal case. Let $p \in(-e, e)$ be the orientation reversing fixed point of $S_{0}$ and $q \in(0, e)$ the periodic point of $S_{0}$ of period two. For $\zeta \in[-e, p)$, we have $W\left(T_{\zeta}\right)=(-e, e)$. For $\zeta \in[p, q)$, $W\left(T_{\zeta}\right)$ consists of a countable number of adjacent intervals: $W\left(T_{\zeta}\right)=(-e, e) \backslash$ $\bigcup_{n \geq 0} T^{-n}(p)$. When $\zeta=q, W\left(T_{\zeta}\right)$ consists of a countable number of adjacent intervals: $W\left(T_{\zeta}\right)=(-e, e) \backslash \bigcup_{n \geq 0}\left(T^{-n}(p) \cup T^{-n}(q)\right)$. In particular,

$$
\left\langle T_{-e}\right\rangle=[-e, p), \mathcal{P} \mathcal{H}\left(T_{-e}\right)=\{-e\}, \min \left[T_{-e}\right]=\left[T_{-e}\right]_{b}=\{-e\},
$$

where $[-e, p)$ stands for the set of maps $T_{\zeta}$ with $\zeta \in[-e, p)$. Moreover,

$$
\left\langle T_{p}\right\rangle=[p, q), \mathcal{P H}\left(T_{p}\right)=(-e, p], \min \left[T_{p}\right]=[T]_{b}=\{p\} .
$$

This example shows that the assumption in Lemma 4.4 that $T$ has only hyperbolic periodic orbits is essential.


Figure 9. The set $\langle T\rangle$ in the unimodal case, see Example 4.8.
Example 4.9. In Figure 10 we illustrate the definition of $\langle T\rangle$ in the case when there is a component $W_{i}=(a, b)$ of the $W(T)$ containing two plateaus with the corresponding set $\langle T\rangle$ drawn in parameter space on the right. The map $T$ depicted in the figure, is the first return map to $(a, b)$ and has the property that $W(T) \cap(a, b)=$ $(a, b)$. The set $\langle T\rangle$ for this map is shown in Figure 10 on the right, and is equal to the union of the two open triangles with the open interval $\Delta$ connecting the points marked 0 and 2 (corresponding to maps $T_{0}$ and $T_{2}$ ). To see this for $T_{0}$, note that $W\left(T_{0}\right) \cap(a, b)$ is equal to $(a, b) \backslash Q$ where $Q$ is a countable set made up of backward iterates of the left boundary point of the left plateau. Therefore $\langle T\rangle$ is neither open nor closed. Note that $\min [T]=\bar{\Delta}$ and $[T]_{b}=\left\{T_{0}, T_{1}, T_{2}\right\}$. Furthermore, $\left\langle T_{2}\right\rangle=T_{2}$ since taking $\zeta_{2}>0$ (while $\zeta_{1}$ is left unchanged) results in the right endpoint of $Z_{1}$ no longer belonging to $W(T)$. The open interval in Figure 10 connecting 2 to 3 is a single $\langle\tilde{T}\rangle$, whereas the line segment connecting 1 to 3 consists of countably many different cells $\langle\tilde{T}\rangle$ consisting of half-open line segments with endpoints corresponding to maps for which one plateau is mapped into the boundary of the other plateau.

Example 4.10. Figure 11 illustrates the definition of $\langle T\rangle$ in the bimodal case when there exists a periodic component $W$ of $W(T)$ of period $s_{1}+s_{2}$ so that $W$ and the component $W^{\prime}$ of $W(T)$ containing $T^{s_{1}}(W)$ both contain a plateau. In this case $\min [T]=[T]_{b}=\left\{T_{1}\right\}$ where $T_{1}$ is the map corresponding to 1 in the figure. Note that only the left and bottom boundary is contained in $\langle T\rangle$ (not including the endpoints of these lines).


Figure 10. The set $\langle T\rangle$ in the case where two plateaus lie in one component of the basin, see Example 4.9. The • on the right indicates the parameter of the map $T$ on the left.


Figure 11. The set $\langle T\rangle$ in the bimodal case when an attractor contains two plateaus in its basin, see Example 4.10 is depicted on the left. For each of the 6 indicated parameters on the left, we draw on the right side, the corresponding graphs of $T^{s_{1}}$ and return map $T^{s_{2}} \circ T^{s_{1}}$ to $W_{1}$ (respectively $T^{s_{2}}$ and the return map $T^{s_{1}} \circ T^{s_{2}}$ to $W_{2}$ ). Here the graph with the larger slope corresponds to the first return map. Of course the diagonal has only meaning for the maps $T^{s_{2}} \circ T^{s_{1}}$ and $T^{s_{1}} \circ T^{s_{2}}$. The part of the boundary of the polygon which is contained in $\langle T\rangle$ is marked by (sn).

Example 4.11. In this example we show why we consider $\min [T]$; it is possible that (with the analogous definition) $\min \langle T\rangle=\emptyset$. Consider $T \in \mathcal{S}^{b}$ so that there exists an interval $\left[Z_{i}, Z_{i+1}\right]$ which is mapped into another plateau $Z_{k}$ with $T\left(Z_{i}\right) \subset \partial Z_{k}$, with $T\left(Z_{i+1}\right)$ contained in the interior of $Z_{k}$ and so that no iterate of $Z_{k}$ is contained in a plateau. Then $\partial Z_{k}$ is not contained $W(T)$. This implies that the $i$-th projection of $W(T)$ is a point while the $i+1$-th projection of $W(T)$ is an open interval. Therefore $\min \langle T\rangle=\emptyset$. (However, $\min [T] \neq \emptyset$; it consists of maps for which $Z_{i}$ and $Z_{i+1}$ touch.
4.4. Wandering pairs and the space $\mathcal{S}_{*}^{b}$. The space $\mathcal{S}^{b}$ contains maps that are of no use to us because they possess wandering intervals, a phenomenon that does not occur in $P^{b}$. For this reason, we define a subset $\mathcal{S}_{*}^{b}$ to be used as a more faithful parameter space of $P^{b}$ than $\mathcal{S}^{b}$.

An important property of $\mathcal{S}_{*}^{b}$, which will be used in the proof of the Main Theorem, is that isentropes $\left\{T \in \mathcal{S}_{\epsilon, *}^{b} ; h_{\text {top }}(T)=h\right\}$ are contractible in $\mathcal{S}_{*}^{b}$. The proof of this is deferred to Section 7. In this section, we will define $\mathcal{S}_{*}^{b}$ and show some of its basic properties.
Definition 4.12. A pair of plateaus $\left(Z_{i}, Z_{j}\right)$ is called wandering if there exists $n \geq 1$ such that $T^{n}(\mathcal{J})$ is a point, where $\mathcal{J}:=\left[Z_{i}, Z_{j}\right]$ is the convex hull of $Z_{i}$ and $Z_{j}$. We say that $T$ is non-degenerate if for every wandering pair $\left(Z_{i}, Z_{j}\right)$, the corresponding interval $\mathcal{J}$ belongs to the closure of a component of the basin of a periodic plateau. Let $\mathcal{S}_{*}^{b}$ denote the set of non-degenerate maps $T \in \mathcal{S}^{b}$.
Remark 4.13. If $T \in \mathcal{S}_{*}^{b}$ then $[T] \subset \mathcal{S}_{*}^{b}$.
Remark 4.14. Blocks of touching plateaus occur only at the boundary of the parameter space, and bimodal maps with touching plateaus have an attracting fixed point. For this reason, wandering pairs don't occur in Milnor ${ }^{6}$ Tresser's paper [31], and $\mathcal{S}_{*}^{b}=\mathcal{S}^{b}$ for $b \leq 2$.
Remark 4.15. For $b \geq 3, \mathcal{S}_{*}^{b}$ is clearly not closed, but it is not open either. Indeed, take $T$ with a pair of adjacent plateaus $Z_{i}, Z_{i+1}$ so that $\left[Z_{i}, Z_{i+1}\right]$ is mapped into the interior of a plateau $Z_{j}$ and $T(U) \cap Z_{j} \subset \partial Z_{j}$ for some small open neighborhood $U$ of $Z_{j}$. Since $Z_{j}$ is in the basin of an attracting fixed point, such a map $T$ exists in $\mathcal{S}_{*}^{b}$. Moreover, there exist maps $\tilde{T}$ arbitrary close to $T$ so that $\tilde{Z}_{j}$ is no longer contained in the basin of a periodic attractor (just choose $\tilde{\zeta}_{j}<\zeta_{j}$ appropriately). It follows that there exists maps $\tilde{T} \notin \mathcal{S}_{*}^{b}$ arbitrarily close to $T$.

Nevertheless, the space $\mathcal{S}_{*}^{b}$ has the following useful property:
Lemma 4.16. (a) Take $T \in \mathcal{S}^{b}$ and an interval $U$ which is not eventually mapped into a plateau and is also not contained in the basin of a periodic attractor. Then there exists $n>m$ so that $T^{n}(U) \cap T^{m}(U) \neq \emptyset$ and $\liminf _{j}\left|T^{j}(U)\right|>0$.
(b) Take $T \in \mathcal{S}_{*}^{b}$ and consider adjacent plateaus $Z_{i}, Z_{i+1}$ so that the convex hull [ $Z_{i}, Z_{i+1}$ ] is not contained in the closure of a component of the basin of a periodic attractor. Then for each map $\tilde{T} \in \mathcal{S}^{b}$ sufficiently close to $T$, the adjacent plateaus $\tilde{Z}_{i}, \tilde{Z}_{i+1}$ do not form a wandering pair either.

Proof. Take $T \in \mathcal{S}_{*}^{b}$ and interval $U$ as in statement (a). Since $T$ is expanding outside its plateaus, there exists a sequence $n_{i} \rightarrow \infty$ so that $T^{n_{i}}(U)$ intersects one of the plateaus of $T$ (but is not contained in a plateau). It follows $T^{n_{i}}(U)$ contains a neighborhood of one of the endpoints of a plateau for infinitely many $i$. This means that there exist $n>m$ so that $T^{n}(U)$ and $T^{m}(U)$ intersect, and
therefore $T^{n+(j+1) k}(U) \cap T^{n+j k}(U) \neq \emptyset$ for $k:=n-m$ and all $j \geq 0$. It follows that $V=\cup_{j \geq 0} T^{n+j k}(U)$ is an interval and $T^{k}(V) \subset V$. So either $V$ contains a repelling fixed point of $T^{k} \mid V$ or the interval $V$ is the finite union of basins of basins of plateaus, separated by one-sided fixed points of $T^{k} \mid V$. Since $U$ is not eventually mapped into the basin of a periodic attractor, $V$ contains in its interior a periodic point $p$ of period $k$ which is not attracting from both sides. Hence $U$ contains in its interior a point $x$ so that $T^{l}(x)=p$ for some $l$. The conclusion of statement (a) of the lemma follows.

Next consider $\mathcal{J}:=\left[Z_{i}, Z_{i+1}\right]$ as in statement (b). By definition of $\mathcal{S}_{*}^{b}$, the interval $U$ is never mapped into another plateau. Hence, by the proof of statement (a), some iterate of $\mathcal{J}$ contains a repelling periodic point $p$ of $T$ in its interior, where $p$ is either (i) repelling from both sides or (ii) $p$ is attracting from one side and separates the immediate basin of two adjacent fixed points of $T^{k}$. Therefore, if (i) holds and $\tilde{T}$ is sufficiently close to $T$, then $p$ is still contained in the interior of some iterate of $\left[\tilde{Z}_{i}, \tilde{Z}_{i+1}\right]$ and $p$ is still repelling for $\tilde{T}$. If (ii) holds then the two basins could merge, but will still contain the iterate of $U$. Hence in any case $\left[\tilde{Z}_{i}, \tilde{Z}_{i+1}\right]$ also does not form a wandering pair for $\tilde{T}$.
4.5. Further properties of the sets $\langle\boldsymbol{T}\rangle,[\boldsymbol{T}]$ and $[\boldsymbol{T}]_{b}$. Given $T \in \mathcal{S}^{b}$, we define the following equivalence class $\sim_{T}$ on $\{1,2, \ldots, b\}: i \sim_{T} j$ if and only if $Z_{i}$ and $Z_{j}$ are both in the immediate basin of the same periodic attractor. (We do not require that $Z_{i}$ and $Z_{j}$ are in the same component of the immediate basin.) Let $J_{1}, J_{2}, \ldots, J_{s} \subset$ $\{1, \ldots, b\}$ be the corresponding equivalence classes. For $J=\left\{i_{1}, \ldots, i_{j}\right\} \subset\{1, \ldots, b\}$ define $\pi_{J}: \mathcal{S}^{b} \rightarrow \mathbb{R}^{\# J}$ be the projection of $\zeta=\left(\zeta_{1}, \ldots, \zeta_{b}\right) \in \mathcal{S}^{b}$ to $\left(\zeta_{i_{1}}, \ldots, \zeta_{i_{j}}\right) \in$ $\mathbb{R}^{\# J}$.

Theorem 4.17. There exists an affine space $V_{T}$ so that, with respect to the coordinates $\left(\zeta_{1}, \ldots, \zeta_{b}\right) \in \mathbb{R}^{b}$, $\langle T\rangle$ is a connected polygonal region in $V_{T}$. Moreover, the following properties hold:
(1) If $\left[T_{1}\right]=\left[T_{2}\right]$ then $\left\langle T_{1}\right\rangle=\left\langle T_{2}\right\rangle$.
(2) If $T_{1}<T_{2}<T_{3}$ and $T_{1}, T_{3} \in\langle T\rangle$ then $T_{2} \in\langle T\rangle$;
(3) $\langle T\rangle$ has a product structure: if we take $J_{1}, \ldots, J_{s} \subset\{1, \ldots, b\}$ as above the statement of the theorem, then there exist connected (polygonal) sets $A_{i} \subset$ $\mathbb{R}^{\# J_{i}}, i=1, \ldots, s$ so that

$$
\langle T\rangle=\left\{\zeta ; \pi_{J_{i}}(\zeta) \in A_{i} \text { for each } i=1, \ldots, s\right\}
$$

If $\# J_{i}=1$, then the corresponding set $A_{i}$ is an interval $\left[a_{i}, b_{i}\right)$ or a point $\left\{a_{i}\right\}$ whereas if $\# J_{1}>1$, then $A_{i}$ has a polygonal shape (see Figures 10 and 11 for representative examples).
(4) For each $T \in \mathcal{S}_{*}^{b}$, there exists $T^{\prime} \in \mathcal{S}_{*}^{b}$ so that $T^{\prime} \in\left[T^{\prime}\right]_{b}$ and $T \in\left[T^{\prime}\right]$.
(5) If $\left[T_{1}\right] \neq\left[T_{2}\right]$ and $\left[T_{1}\right] \cap\left[T_{2}\right] \neq \emptyset$, then there exists $T^{\prime} \in\left[T^{\prime}\right]_{b}$ with $T^{\prime} \in$ $\min \left(\left[T_{1}\right] \cap\left[T_{2}\right]\right)$.
(6) If $\tilde{T} \in[T] \backslash\langle T\rangle$, then one of the following properties holds:

- $\tilde{T}$ has a saddle-node merging two basins;
- $\tilde{T}$ has a period doubling orbit;
- $\tilde{T}$ has a pitch-fork orbit;
- $\tilde{T}$ has a homoclinic orbit.

Proof. Let $T=T_{\zeta}$ be the map corresponding to $\zeta=\left(\zeta_{1}, \ldots, \zeta_{b}\right)$, let $\tilde{\zeta}=\left(\tilde{\zeta}_{1}, \ldots, \tilde{\zeta}_{b}\right)$, and assume that $\tilde{T}=T_{\tilde{\zeta}}$ is so that $\langle\tilde{T}\rangle=\langle T\rangle$, i.e., $W(\tilde{T})=W(T)$. Consider a plateau $Z_{k}$ of $T$ and in order to be definite assume that $S_{0}$ has a maximum in $Z_{k}$. Let $W_{k}$ be the component of $W(T)$ intersecting $Z_{k}$. We denote the $k$-th plateau of $\tilde{T}$ by $\tilde{Z}_{k}$.

Step 1: First consider the case that $T\left(Z_{k}\right)$ is not contained in the closure of a component of $W(T)$. Let us show that in this case $\tilde{\zeta}_{k}=\zeta_{k}$. Indeed, $W_{k}$ is equal to the interior of $Z_{k}$ (or equal to the interior of the union of all plateaus which touch $\left.Z_{k}\right)$. If $\tilde{\zeta}_{k}<\zeta_{k}$ then $\tilde{Z}_{k}$ strictly contains $W_{k}$ and in particular the corresponding component of $W(\tilde{T})$ strictly contains $W_{k}$, contradicting $\langle\tilde{T}\rangle=\langle T\rangle$. Hence $\tilde{\zeta}_{k} \geq \zeta_{k}$. By assumption, there is a sequence of points converging from the right to $T\left(Z_{k}\right)$ which are not in $W(T)$. Hence if $\tilde{\zeta}_{k}>\zeta_{k}$ then the component of $W\left(T_{\tilde{\zeta}}\right)$ intersecting $\tilde{Z}_{k}$ is equal to $\tilde{Z}_{k}$ and so is strictly inside $Z_{k}$, contradicting $W(\tilde{T})=W(T)$. It follows that in this case $\tilde{\zeta}_{k}=\zeta_{k}$. Note that the freedom of choice of $\zeta_{k}$ does not depend on any of the other coordinates.

Step 2. Next consider the case $W_{k}$ is not periodic and $T\left(Z_{k}\right)$ is contained in the closure of a component $W=(a, b)$ of $W(T)$.

Step 2a. Let us first assume that $T\left(Z_{k}\right) \subset[a, b)$. If $T\left(Z_{k}\right)=\{a\}$ then $W_{k}$ is equal to the interior of $Z_{k}$ and otherwise $W_{k}$ is equal to the component of $T^{-1}(a, b)$ containing $Z_{k}$. In either case, $W_{k}$ is equal to the component of $S_{0}^{-1}(a)$ containing $Z_{k}$. If $\tilde{T}\left(\tilde{Z}_{k}\right) \geq b$ then there exists $x \in Z_{k} \backslash \tilde{Z}_{k}$ so that $\tilde{T}(x)=S_{0}(x)=b \notin W$ and hence $x \notin W(T)$, contradicting that $W(\tilde{T})=W(T)$. If $\tilde{T}\left(\tilde{Z}_{k}\right)<a$, then the interior of $\tilde{Z}_{k}$ strictly contains $W_{k}$ which is not possible either. On the other hand, since $W_{k}$ is not periodic, changing $\tilde{\zeta}_{k} \in[a, b)$ does not change the component of $W(\tilde{T})$ containing $\tilde{T}\left(\tilde{Z}_{k}\right)$. It follows that in this case a necessary and sufficient condition on $\tilde{\zeta}_{k}$ for $W(\tilde{T})=W(T)$ to hold is that $\tilde{\zeta}_{k} \in[a, b)$ and $\tilde{T} \in \mathcal{S}_{*}^{b}$.

Step 2b. If $W_{k}$ is not periodic and $T\left(Z_{k}\right)=\{b\}$ where $W=(a, b)$ is a component of $W(T)$ then $W_{k}$ is the interior of $Z_{k}$ and as in Step 1 we get $\tilde{\zeta}_{k}=\zeta_{k}$. Again, the freedom of choice of $\zeta_{k}$ does not depend on any of the other coordinates.
Step 3. Now consider the case that $W_{k}$ is periodic, i.e., $T^{s}\left(W_{k}\right) \subset \overline{W_{k}}$ for some $s>0$. By Lemma 4.3 either $T^{s}\left(W_{k}\right) \subset \partial W_{k}$ or $T^{s}: W_{k} \rightarrow W_{k}$. In the latter case $T^{s}: W_{k} \rightarrow W_{k}$ has a unique fixed point $p \in W_{k}$ (and $T^{j}(p) \in \operatorname{int}\left(\cup_{i=1}^{b} Z_{k, T}\right)$ for some
$0 \leq j<s)$ and $T^{s}\left(\partial W_{k}\right) \subset \partial W_{k}$. Because of Lemma 4.3(5)-(6), the component of $W\left(T_{\tilde{\zeta}}\right)$ intersecting $\tilde{Z}_{k}$ is equal to $W_{k}$ if and only if
(i) $\tilde{T}^{s}\left(W_{k}\right) \subset \overline{W_{k}}$;
(ii) $\tilde{T}^{s}\left(W_{k}\right) \cap \partial W_{k} \neq \emptyset$ implies that $T^{s} \mid W_{k}$ is constant, and
(iii) $\tilde{T}^{2 s} \mid W_{k}$ has at most one fixed point (which is attracting).

Let $(a, b)$ be the component of $W(T)$ which contains $T\left(W_{k}\right)$ in its closure. Properties (i), (ii) and (iii) persist while decreasing $\tilde{\zeta}_{k} \in[a, b)$, by Lemma 4.3(5)-(6). While increasing $\tilde{\zeta}_{k}$ in $[a, b)$ these properties are preserved until $\partial \tilde{Z}_{k}$ first hits a fixed point $q$ of $S_{0}^{2 s}$. If $S_{0}^{j}(q) \in \operatorname{int}\left(\cup_{i=1}^{b} Z_{i, T}\right)$ for at least one $0 \leq j<s$, then one can continue to increase $\tilde{\zeta}_{k}$ until one $\partial \tilde{Z}_{k}$ hits another fixed point $q$ of $S_{0}^{2 s}$. In this way, we can keep increasing $\zeta_{k}$ (i.e., shrink the width of the plateau), until either $\partial Z_{i}$ contains a fixed point $q^{\prime}$ of $S_{0}^{2 s}$ for which $S_{0}^{j}\left(q^{\prime}\right) \notin \operatorname{int}\left(\cup_{i=1}^{b} Z_{i, T}\right)$ for all $0 \leq j<s$ or until we no longer have $\tilde{T}^{s}\left(Z_{k}\right) \subset W_{k}$.

If follows that, fixing all $\zeta_{j}$ with $j \neq k$, the set $\langle T\rangle$ is equal to a line segment of the form $a_{k} \leq \zeta_{k}<b_{k}$. Here the left boundary $a_{k}$ does not depend on the choice for the other $\zeta_{j}$ 's, because one can decrease $\tilde{\zeta}_{k} \in\left[a_{k}, b_{k}\right)$ and the latter interval does not depend on $\zeta_{j}, j \neq k$. However, the right hand boundary $b_{k}$ in general will depend on parameters $\zeta_{j}, j \neq k$ for which $Z_{j}$ is contained in (the closure of) one of the components of $W(T)$ containing $W_{k}, T\left(W_{k}\right), \ldots, T^{s-1}\left(W_{k}\right)$. The location is determined by the position of the fixed points of $S_{0}^{2 s}$.

Thus we have proved that $\langle T\rangle$ is contained in a polygonal region in some hyperplane $V_{T}$ and that properties (1), (2) and (3) in the theorem are satisfied.

Step 4. We claim that we can choose $T^{\prime} \in[T]_{b}$ (by only changing $T$ inside the basin of periodic attractors) so that $T \in\left[T^{\prime}\right]$ and $T^{\prime} \in\left[T^{\prime}\right]_{b}$. Indeed, if the number of plateaus within a component of the basin is even, then we choose $T^{\prime}$ analogously to the map $T_{1}$ in Example 4.9 and when the number of plateaus is odd then all these plateaus of $T^{\prime}$ necessarily touch and the attracting periodic point is at the boundary point of the union of the touching plateaus. This implies that, for the coordinates corresponding to attracting plateaus, $T \in\left[T^{\prime}\right]$ and $T^{\prime} \in\left[T^{\prime}\right]_{b}$. If $Z_{i}$ is contained in the basin, but not in the immediate basin of a periodic attractor, then for any map $\hat{T} \in[T]_{b}$ the corresponding plateau $\hat{Z}_{i}$ is either mapped into the boundary of this component, or $Z_{i}$ touches with one of its neighbouring plateaus. From this description it follows that $T \in\left[T^{\prime}\right]$ and $T^{\prime} \in\left[T^{\prime}\right]_{b}$. Since $T$ and $T^{\prime}$ only differ on the basin of periodic attractors, we will still have $T^{\prime} \in \mathcal{S}_{*}^{b}$.

Step 5. To prove Property (5), assume that $\left[T_{1}\right] \neq\left[T_{2}\right]$ and $\left[T_{1}\right] \cap\left[T_{2}\right] \neq \emptyset$. Property (2) gives that $\min \left(\left[T_{1}\right] \cap\left[T_{2}\right]\right) \subset\left[T_{1}\right] \cap\left[T_{2}\right]$. It is possible that $T \in \min \left(\left[T_{1}\right] \cap\left[T_{2}\right]\right)$ has two or more touching plateaus $Z_{i}, Z_{i+1}$ which are mapped by some iterate $T^{s}$ into the interior of another plateau $Z_{j}$. In this case choose a continuous deformation $T_{t}$
with $T_{0}=T$ so that for each such pair of touching plateaus, $T_{t}^{s}\left(\left[Z_{i}, Z_{i+1}\right]\right) \subset Z_{j}$ for all $t \in[0,1]$ (leaving all other plateaus unchanged). Then $T^{\prime}:=T_{1} \in \min \left(\left[T_{1}\right] \cap\left[T_{2}\right]\right)$ has the required properties.

Step 6. Property (6) holds because $T^{\prime} \in[T] \backslash\langle T\rangle$ implies that at least one periodic orbit is in the boundary of a plateau. By Property (2) this periodic orbit cannot be a saddle-node, unless it corresponds to a situation where two basins are merged.

## 5. The map $\boldsymbol{\Psi}: \boldsymbol{P}^{\boldsymbol{b}} \rightarrow \mathcal{S}^{\boldsymbol{b}}$.

Let us review some basic kneading theory, see [30], and also [26, 31]. Given an interval $I$ and piecewise monotone $b$-modal map $f: I \rightarrow I$ with turning points $c_{1}<\cdots<c_{b}$ in the interior of $I$, one can associate to each point $x \in I$ an itinerary $\underline{i}_{f}(x)$ consisting of a sequence $\left(i_{0}, i_{1}, \ldots\right)$ of symbols from the alphabet $\left\{I_{0}, c_{1}, I_{1}, c_{2}, \ldots, c_{b}, I_{b}\right\}$. Here $I_{0}, \ldots, I_{b}$ are the components of $I \backslash\left\{c_{1}, \ldots, c_{b}\right\}$ ordered from left to right. It is well-known that $x \mapsto \underline{i}_{f}(x)$ is monotone w.r.t. the signed lexicographic ordering and that therefore the $i$-th kneading sequence

$$
\nu_{i}:=\lim _{x \downarrow c_{i}} \underline{i}_{f}(x)
$$

is well-defined. Note that the sequence $\nu_{i}$ does not contain any of the symbols $c_{1}, \ldots, c_{b}$. Let $\sigma$ be the left shift on the space of symbol sequences. The kneading invariant $\nu(f)$ of $f$ is defined as

$$
\nu(f):=\left(\nu_{1}, \ldots, \nu_{b}\right) .
$$

Any kneading invariant which is realized by some piecewise monotone $b$-modal map is called admissible.

To each map $f \in P^{b}$ one can associate uniquely a stunted sawtooth map as follows. Let $\nu(f)=\left(\nu_{1}, \ldots, \nu_{b}\right)$ be the kneading invariant of $f$, and let $s_{i}$ be the unique point in the $(i+1)$-th lap $I_{i}$ of $S_{0}$ such that

$$
\begin{equation*}
\lim _{y \downarrow s_{i}} \underline{i}_{S_{0}}(y)=\nu_{i}:=\lim _{x \downarrow c_{i}} \underline{i}_{f}(x) . \tag{7}
\end{equation*}
$$

Let $Z_{i}$ be the symmetric interval around the $i$-th turning point of $S$ with right endpoint $s_{i}$. Let us define

$$
\begin{equation*}
\Psi: P^{b} \rightarrow \mathcal{S}^{b}, \quad f \mapsto \Psi(f) \tag{8}
\end{equation*}
$$

by associating to $f$ the unique stunted sawtooth map $\Psi(f)$ which agrees with $S_{0}$ outside $\cup_{i=1}^{b} Z_{i}$ and which is constant on $Z_{i}$ with value $S_{0}\left(s_{i}\right)$.
5.1. Some good properties of $\Psi$. The main reason for working with the map $\Psi$ is that it allows us to work with the Euclidean space $\mathcal{S}^{b}$ rather than with the space of kneading invariants.

Lemma 5.1. The map $\Psi: P^{b} \rightarrow \mathcal{S}^{b}$
(1) is well-defined;
(2) the kneading invariant of $f$ and $T:=\Psi(f)$ are the same in the sense that $\lim _{y \downarrow Z_{i}} i_{T}(y)=\nu_{i}$.
(3) $f$ and $\Psi(f)$ have the same topological entropy;
(4) $\Psi\left(P_{\epsilon}^{b}\right) \subset \mathcal{S}_{\epsilon, *}^{b}$.

Proof. Since $S_{0}$ allows every sequence in $\left\{I_{0}, \ldots, I_{b}\right\}^{\mathbb{N}}$ as itinerary, we can always find a stunted version $T$ with the required kneading sequences. In fact, because $S_{0}$ is expanding and so distinct points have different itineraries, the stunted version is unique, so $\Psi$ is well-defined. It also follows that the orbits of the boundary points of $Z_{i}$ under $T$ and the sawtooth map $S_{0}$ agree and therefore statement (2) holds. Entropy is fully determined by kneading sequences, so $\Psi$ preserves entropy. For the last statement, assume by contradiction that $\Psi(f) \in \mathcal{S}^{b} \backslash \mathcal{S}_{*}^{b}$. Since $f$ has no wandering intervals, $\Psi(f) \in \mathcal{S}^{b} \backslash \mathcal{S}_{*}^{b}$ implies that there exists an interval connecting two adjacent critical points, so that the $n$-th iterate of this interval is another critical point. Clearly this is impossible.
5.2. Some bad properties of $\Psi$. The next example shows that $f \mapsto \Psi(f)$ is neither continuous, nor injective nor surjective, but later on we shall see that this map is 'almost' continuous, injective and surjective. It also shows that $\Psi(f)=\Psi(\tilde{f})$ does not imply that $f, \tilde{f}$ are partially conjugate.

Example 5.2. Consider the family $f_{\lambda}(x)=\lambda x(1-x)$ and let $T_{\zeta}$ be as in Example 4.8. Then there are parameters $0<\lambda_{1}^{\prime}=2<\lambda_{1}=3<\lambda_{2}^{\prime}<\lambda_{2}<\lambda_{3}^{\prime}$ so that $\lambda_{1}, \lambda_{2}$ are the first two period doubling parameters, and $\lambda_{1}^{\prime}, \lambda_{2}^{\prime}, \lambda_{3}^{\prime}$ are the first three parameters at which the critical point of $f_{\lambda}$ is periodic. Then

$$
\Psi\left(f_{\lambda}\right)= \begin{cases}T_{-e} & \text { for } \lambda \in\left[0, \lambda_{1}^{\prime}\right] \\ T_{p} & \text { for } \lambda \in\left(\lambda_{1}^{\prime}, \lambda_{2}^{\prime}\right] \\ T_{q} & \text { for } \lambda \in\left(\lambda_{2}^{\prime}, \lambda_{3}^{\prime}\right]\end{cases}
$$

see Figure 12, As in Example 4.8, $T_{p}$ and $T_{q}$ are the maps for which the right endpoint of the plateau is a fixed point and has period two respectively. The discontinuities of $f \mapsto \Psi(f)$ occur when the critical point becomes periodic, rather than when a period doubling bifurcation occurs. In particular, for $\delta>0$ small, $\Psi\left(f_{\lambda_{1}+\delta}\right)=\Psi\left(f_{\lambda_{1}-\delta}\right)=T_{p}$ but $f_{\lambda_{1}+\delta}$ and $f_{\lambda_{1}-\delta}$ are not partially conjugate. On the other hand, for each $\lambda \in\left[0, \lambda_{1}\right], \quad f_{\lambda} \in \mathcal{P H}\left(f_{0}\right)$ and $\Psi\left(f_{\lambda}\right) \in\left\{T_{-e}, T_{p}\right\} \subset\left[T_{-e}\right]=$ $\left[\Psi\left(f_{\lambda_{1}^{\prime}}\right)\right]$; note that $T_{-e}$ and $T_{p}$ are partially conjugate.

### 5.3. The definition of $\mathcal{A}_{b}$.

Definition 5.3. The set $\mathcal{A}_{b}$ is the collection of polynomials $f \in P^{b}$ such that for each component $B$ of its basin $B(f)$ containing a critical point the following holds:
(1) if $f(\partial B)$ consists of a single point, then each points in $f(B)$ has the same itinerary.

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Figure 12. The map $\lambda \mapsto \Psi\left(f_{\lambda}\right)$ is discontinuous at parameters $\lambda_{i}^{\prime}$ where the critical point of $f_{\lambda}$ is periodic, see Example 5.2. These parameters are alternated with period doubling parameters $\lambda_{i}$, and at $\lambda=1$, the stability of 0 changes, so $\mathcal{P H}$ changes too. The equivalence classes $\mathcal{P H}$ (of partially conjugate maps) and $\mathcal{H}$ (of maps with the same kneading invariant) are also shown. The segments with an additional line correspond to maps in the set $\mathcal{A}_{b}$ defined in Section5.3.
(2) if $f(\partial B)$ consists of two points, then (i) each point in $f\left(B^{\prime}\right)$ has the same itinerary where $B^{\prime}$ be the convex hull of the critical points in $B$ and (ii) if $B$ contains an attracting periodic point $p$, then $p$ is in the interior of $B^{\prime}$.

Note that $f(\partial B)$ consists of a single point if and only if the interior of $B$ contains an odd number of critical points. The reason for introducing $\mathcal{A}_{b}$ is:
Proposition 5.4. If $f \in \mathcal{A}_{b}$ then $\Psi(\mathcal{P H}(f)) \subset[\Psi(f)]$.
Before proving this proposition we will motivate the definition of $\mathcal{A}_{b}$ by considering three examples.
Example 5.5. If $f_{\lambda}(x)=\lambda x(1-x)$ is a quadratic map with an attracting fixed point, then $\Psi\left(f_{\lambda}\right)$ is a stunted sawtooth map which is either equal to the constant map $T_{0}$ corresponding to the parameter $\zeta=-e$ or to the map $T_{p}$ which has plateau $[-p, p]$, see Example 4.8 and Figure 9. Now $T_{p} \in\left[T_{-e}\right]$ while $T_{-e} \notin\left[T_{p}\right]$. Moreover $\mathcal{P H}\left(f_{3 / 2}\right)=\left\{f_{\lambda} ; \lambda \in(1,3]\right\}$ and $\Psi\left(\mathcal{P H}\left(f_{3 / 2}\right)\right) \subset\left\{T_{-e}, T_{p}\right\} \subset\left[T_{-e}\right]$. Note that $f_{\lambda} \in \mathcal{A}_{b}$ when $\lambda \in[0,2]$ and then $\Psi(f)=T_{-e}$ and so $\Psi(\mathcal{P H}(f)) \subset[\Psi(f)]$. If $f_{\lambda} \notin \mathcal{A}_{b}$ then this inclusion does not hold, so the assumption that $f \in \mathcal{A}_{b}$ is essential for Proposition 5.4 to hold. Note that it also not true that $\Psi(\mathcal{P H}(f)) \subset \mathcal{P} \mathcal{H}(\Psi(f))$; Example 4.8 (see also Figure 12) is the simplest counter-example.

The next example, shows why in the definition of $\mathcal{A}_{b}$ we add condition 2(ii) if $f(B)$ consists of two points,

Example 5.6. Assume that $f$ is a cubic map with an attracting fixed point which attracts both critical points (say with the left critical point a maximum). Then $\Psi(f)$
is equal to one of the following five maps $T_{0}, T_{1}, T_{2}, T_{3}, T_{4}$ determined by $\left(\zeta_{1}, \zeta_{2}\right)$ equal to $(-e, e),(0,0),(e,-e),(e, 0)$ or $(0, e)$ as indicated in Example 4.9 and Figure 10 (when taking $a=-e$ and $b=e$ ). Note that $T_{0}, T_{1}, \ldots, T_{4} \in\left[T_{i}\right]$ when $i=1$ but not when $i=0,2,3,4$. Also note that $\Psi(\mathcal{P H}(f)) \subset\left\{T_{0}, T_{1}, T_{2}\right\} \subset\left[T_{1}\right]$. If $f \in \mathcal{A}_{b}$ then $\Psi(f)=T_{1}$ and so $\Psi(\mathcal{P H}(f)) \subset[\Psi(f)]$.

Example 5.7. Finally consider the example of a bimodal map such that there exist $s_{1}, s_{2} \in \mathbb{N}$ so that $f^{s_{1}}\left(B_{1}\right) \subset B_{2}$ and $f^{s_{2}}\left(B_{2}\right) \subset B_{1}$ where $B_{i}$ are distinct components of $B(f)$ and so that $B_{1}$ and $B_{2}$ both contain exactly one critical point. Then the situation is as in Example 4.9 and $\Psi(f)$ is contained in the polygon drawn in Figure 11. Maps in $\{\Psi(\tilde{f}) ; \tilde{f} \in \mathcal{P} \mathcal{H}(f)\}$ correspond to the maps indicated by $1,2,4,6$ in Figure 11 and their symmetric counterpart under reflexion in the diagonal of the $\left(\zeta_{1}, \zeta_{2}\right)$-space. There are six such maps, all belonging to $\partial[\Psi(f)]$. The fact that $f \in \mathcal{A}_{b}$ ensures that $\Psi(f)$ corresponds to the lower corner of the region, i.e., the map $T$ denoted by 1 in Figure 11 . This map has the property that $[T]$ is equal to this polygon (this is false for maps denoted by hc and pd).

Proof of Proposition 5.4. Let $T=\Psi(f)$ for an arbitrary $f \in P^{b}$. The definition of $\Psi$ ensures that the orbit of $T\left(Z_{i}\right)$ under $T$ is the same as an orbit of $S_{0}$, and therefore no plateau can be mapped into the interior of a plateau by $T$. On the other hand, if $T^{k}\left(Z_{i}\right) \in \partial Z_{j}$ for some minimal $k$, then a small changes of $\zeta_{i}$ can move $T^{k}\left(Z_{i}\right)$ into the interior of $Z_{j}$. Therefore $T \in \partial\langle T\rangle$, and in particular, $\Psi(f) \in[\Psi(f)]_{\text {b }}$ for $f \in \mathcal{A}_{\text {b }}$.

Take $\tilde{f} \in \mathcal{P} \mathcal{H}(f)$. For each $i$ so that $c_{i}$ is not in the basin of a periodic attractor, $c_{i}$ and $\tilde{c}_{i}$ have the same kneading invariant and so the $i$-th component of $\Psi(f)$ and $\Psi(\tilde{f})$ agree.

Now let us consider critical points in the basin of a periodic attractor. Although kneading sequences within $\mathcal{P} \mathcal{H}(f)$ are not constant, all itineraries in the basin of a periodic attractor (in the limit sense of (7)) are (pre)periodic to the same periodic sequence in $\left\{I_{0}, \ldots, I_{b}\right\}^{\mathbb{N}}$. Since every periodic attractor of $\tilde{f}$ has a critical point in its immediate basin, there is $\tilde{T} \in \mathcal{S}^{b}$ that realizes the corresponding periodic itineraries by orbits that intersect the interior of a plateau. Additionally, $\tilde{T}$ can be chosen such that all kneading sequences of $c_{i} \in B(f)$ are indeed achieved by points in the interior of plateaus. For this $\tilde{T}$ and the semiconjugacy $h$ between $\tilde{T}$ and $\tilde{f}$, we have $h^{-1}(B(\tilde{f})) \subset W(\tilde{T})$. Also $W(\tilde{T})=W(T)$, so $\tilde{T} \in\langle T\rangle$. Naturally, the periodic attractors of $\Psi(\tilde{f})$ lie on the boundary of plateaus, but that still means that $\Psi(\tilde{f}) \in[\tilde{T}]=[T]$. Since $\tilde{f} \in \mathcal{P H}(f)$ is arbitrary, $\Psi(\mathcal{P H}(f)) \subset[\Psi(f)]$.
5.4. Almost surjectivity of $\boldsymbol{\Psi}$. In Proposition 5.9 below, we shall prove that $\Psi$ is almost surjective. In order to prove this we need a result from [26] concerning full families. Let us say that a periodic attractor of a piecewise monotone interval map
$g: I \rightarrow I$ is essential if it contains a turning point in its immediate basin. We say that $g$ has no wandering intervals, if each interval $J$ for which $J, g(J), g^{2}(J), \ldots$ are all pairwise disjoint necessarily intersects the basin of some periodic attractor. If $g$ has no wandering interval, then each interval $J$ for which $g^{n} \mid J$ is a homeomorphism for all $n$ is necessarily contained in the basin of periodic attractor, see [26]. It is wellknown, see Theorem IV.A in [26], that maps in $P^{b}$ do not have wandering intervals and that all their attractors are essential.

Theorem 5.8 (Fullness of Families). Each piecewise monotone map $g$ with $b$ turning points is topologically conjugate to a polynomial in $P^{b}$, provided the following two properties are met:
(1) $g$ has no wandering intervals and no inessential attractors;
(2) each periodic turning point is an attractor (this is automatically satisfied if $g$ is $C^{1}$ ).
Moreover, assume that $g$ has an attracting periodic point, then one can find $f, \tilde{f} \in$ $P^{b}$ which are both topologically conjugate to $g$ so that the corresponding attracting periodic point is hyperbolic for $f$ and parabolic for $\tilde{f}$.

Proof. The first part of this theorem is Theorem II.4.1 in [26]. The second part requires a slightly modifying the proof in [26] on page 124-125.

The next proposition gives the required surjectivity:
Proposition 5.9 ( $\Psi$ is almost surjective). For each $T \in \mathcal{S}_{*}^{b}$ there exists a polynomial $f \in P^{b} \cap \mathcal{A}_{b}$ such that $T \in[\Psi(f)]$.

Proof of Proposition 5.9. By Theorem 4.17,4), for each $T \in \mathcal{S}_{*}^{b}$ there exists $T^{\prime} \in \mathcal{S}_{*}^{b}$ so that $T^{\prime} \in\left[T^{\prime}\right]_{b}$ and $T \in\left[T^{\prime}\right]$. Therefore, if we can prove that there exists a polynomial $f \in P^{b} \cap \mathcal{A}_{b}$ such that $T^{\prime}=\Psi(f)$, then the proposition follows as well.

Since $T$ is not piecewise monotone (because of its plateaus), we cannot apply Fullness Theorem 5.8 directly. In order to obtain a piecewise monotone map, we first replace $T$ on each component $B$ of its basin which contains plateaus, by an affinely scaled copy of a map $L_{q}$ as in Figure 13 of the appropriate type. Here $q$ is the number of plateaus in $B$. Let us call the resulting map $T^{\prime \prime}$. We can choose $L_{q}$ so that $T^{\prime}$ and $T^{\prime \prime}$ have the same kneading invariants (and hence $\left[T^{\prime}\right]_{b}=\left[T^{\prime \prime}\right]_{b}$ ), and so that if $T \in\left[T^{\prime \prime}\right]$ has an attracting periodic point in the common boundary point of two plateaus, then $T^{\prime \prime}$ has an attracting periodic point between the two corresponding turning points.

Since $T^{\prime \prime}$ may still have plateaus, we next define $x \sim y$ for $x, y \in[-e, e]$, if there exists $n \geq 0$ so that $T^{\prime \prime n}$ maps the convex hull $[x, y]$ into one of the (remaining) plateaus of $T^{\prime \prime}$. Collapse each such interval $[x, y]$ to a point and let $T^{\prime \prime \prime}$ be the corresponding map. From the definition it follows that $T^{\prime \prime \prime}$ is continuous and since $T \in \mathcal{S}_{*}^{b}$, it follows that $T^{\prime \prime \prime}$ is $b$-modal and has no wandering intervals. It also has


Figure 13. The map $L_{q}$ where $q$ is the number of touching plateaus in the component $B$ of the basin of $T \in\left[T^{\prime \prime}\right]_{b}$. We choose $L_{q}$ so that it is continuous, piecewise affine and so that the slope between its turning points is at most $1 / 4$.
no inessential attractors (since $T$ does not have these either). Hence we can apply Theorem 5.8 and there exists $f \in P^{b}$ that is topologically conjugate to $T^{\prime \prime \prime}$, and therefore have the same kneading invariants. The choice of the $L_{q}$ 's corresponds exactly with the definition of $\mathcal{A}_{b}$, so indeed $f \in \mathcal{A}_{b}$.

### 5.5. Almost injectivity of $\Psi$.

Proposition 5.10 ( $\Psi$ is almost injective). The map $\Psi: P^{b} \rightarrow \mathcal{S}_{*}^{b}$ is 'almost injective' in the sense that if $f, \tilde{f} \in \mathcal{A}_{b}$ and $[\Psi(f)] \cap[\Psi(\tilde{f})] \neq \emptyset$, then at least one of $\mathcal{P} \mathcal{H}(f) \cap$ $\overline{\mathcal{P H}(\tilde{f})}$ and $\overline{\mathcal{P H}(f)} \cap \mathcal{P H}(\tilde{f})$ is non-empty.
Remark 5.11. Here the assumption that $f, \tilde{f} \in \mathcal{A}_{b}$ is needed. Indeed, using the notation from Example 5.2, take $f=f_{\lambda_{1}^{\prime}+\epsilon}$ and $\tilde{f}=f_{\lambda_{3}^{\prime}}$. Then $\Psi(f)=T_{p}$ and $\Psi(\tilde{f})=T_{q}$, and therefore $[\Psi(f)] \cap[\Psi(\tilde{f})] \neq \emptyset$ while $\overline{\mathcal{P} \mathcal{H}(f)} \cap \overline{\mathcal{P} \mathcal{H}(\tilde{f})}=\emptyset$.

Proof. If $\Psi(f)=\Psi(\tilde{f})$ then $f, \tilde{f} \in P^{b}$ have the same kneading invariants. Hence there exists an order preserving bijection $h: \cup_{c} \cup_{n \in \mathbb{Z}} f^{n}(c) \rightarrow \cup_{\tilde{c}} \cup_{n \in \mathbb{Z}} \tilde{f}^{n}(\tilde{c})$ (where the outer union runs over the critical points $c$ of $f$ and $\tilde{c}$ of $\tilde{f}$ ), so that $h \circ f=\tilde{f} \circ h$. Here $h$ maps each critical points of $f$ to the corresponding critical point of $\tilde{f}$. It follows that if $f$ has no periodic attractors, then $\tilde{f}$ also has no periodic attractors and so $f$ and $f$ are topologically conjugate. Rigidity Theorem 3.7 then gives that $f=\tilde{f}$.

If $f$ has a periodic attractor $p$, then define $H \ni p$ to be the largest interval such that $f^{s}(H) \subset H$ and $f^{s} \mid H$ preserves orientation. That is, $s$ is either the period of $p$ if $p$ is orientation preserving and twice the period otherwise. Since $f \in P^{b}$, by taking an iterate of $p$ is necessary, we can assume that $H$ contains a critical point $c$, and $\left(f^{k s}(c)\right)_{k \geq 1}$ is a monotone sequence in $H$ converging to $p$. In this case $\tilde{f}$ also has a periodic attractor in a corresponding interval $\tilde{H}$. However, it is possible
that $f^{s} \mid H$ has a unique attracting fixed point, while $\tilde{f}^{s} \mid \tilde{H}$ has two attracting fixed points, or vice versa. It follows that either $\mathcal{P} \mathcal{H}(f)=\mathcal{P} \mathcal{H}(\tilde{f})$ or there exists a map $g \in P^{b}$ with at least one parabolic periodic point such that $g \in \mathcal{P} \mathcal{H}(f) \cap \overline{\mathcal{P H}(\tilde{f})}$ or $g \in \overline{\mathcal{P H}(f)} \cap \mathcal{P} \mathcal{H}(\tilde{f})$, see the last part of Theorem 5.8. Since $\mathcal{P} \mathcal{H}(f)$ and $\mathcal{P} \mathcal{H}(\tilde{f})$ are connected, in particular it follows that the space $\mathcal{H}(f)$ from Theorem 1 is connected.

Assume now that $[\Psi(f)] \cap[\Psi(\tilde{f})] \neq \emptyset$ and $[\Psi(f)] \neq[\Psi(\tilde{f})]$. By Theorem 4.17(6), there exists a map $T_{*} \in\left[T_{*}\right]_{b}$ so that $T_{*} \in \min ([\Psi(f)] \cap[\Psi(\tilde{f})])$. According to Theorem 4.17(6) there are four possibilities:
(1) $T_{*}$ has a saddle-node merging two basins (or splitting one into two);
(2) $T_{*}$ has a period doubling orbit;
(3) $T_{*}$ has a pitch-fork orbit;
(4) $T_{*}$ has a homoclinic orbit.

Proposition 5.9 produces a map $f_{*} \in P^{b}$ such that $\Psi\left(f_{*}\right)=T_{*}$ and with the corresponding bifurcations.

Claim: In all these cases, $f_{*} \in \mathcal{P H}(f) \cap \overline{\mathcal{P H}(\tilde{f})}$ or $f_{*} \in \overline{\mathcal{P} \mathcal{H}(f)} \cap \mathcal{P H}(\tilde{f})$.
Proof of Claim. This follows from Theorem 3.6. Indeed, let $i=1, \ldots, b$ and assume that the coordinate $\zeta_{i}(\Psi(f))<\zeta_{i}(\Psi(\tilde{f}))$. Since $f \in \mathcal{A}_{b}$, by Proposition 5.4, $\Psi(\mathcal{P H}(f)) \subset[\Psi(f)]$.

Claim: One can find a continuous path $f_{t}$ in $P^{b}(f), t \in[-1,0]$ so that $f_{t} \in \mathcal{P} \mathcal{H}^{o}(f)$ for $t \in(-1,0), f_{-1}=f, f_{0}=f_{*}$.
Proof of Claim: This follows from Theorem 3.6. Indeed, consider $T=\Psi(f)$ and for each of its non-hyperbolic basins, consider what type of bifurcations $T_{*} \in[T]$ undergoes: merging or splitting components of touching basins, see (1)-(4) above. Then choose the sign $\sigma_{i}$ for each neutral orbit of $f_{t}$ so that Theorem 3.6 ensures that $f_{t}$ undergoes the corresponding bifurcation as $t \uparrow 0$ : it splits into two components by a saddle-node, a period doubling bifurcation, a pitch-fork bifurcation or a critical point moves from the boundary into the interior of one of the components of the basin.

Similarly, since $\tilde{f} \in \mathcal{A}_{b}$ we have $\zeta_{i}(\Psi(\tilde{f}))=\zeta_{i}\left(T_{*}\right)=\zeta_{i}\left(\Psi\left(f_{*}\right)\right)$ (in fact $\zeta_{i^{\prime}}(\Psi(\tilde{f}))=$ $\zeta_{i^{\prime}}\left(T_{*}\right)=\zeta_{i^{\prime}}\left(\Psi\left(f_{*}\right)\right)$ holds for each $i^{\prime} \in\{1, \ldots, b\}$ when $c_{i^{\prime}}$ is a critical point attracted to the same periodic orbit as $c_{i}$.) Applying Theorem 3.6 again, one can find a continuous path $f_{t}$ in $P^{b}(f), t \in[0,1]$ so that $f_{t} \in \mathcal{P} \mathcal{H}^{o}(f)$ for $t \in(0,1), f_{0}=f_{*}$ and $f_{1}=\tilde{f}$. Since $\zeta_{i}(\Psi(\tilde{f}))=\zeta_{i}\left(T_{*}\right)=\pi_{i}\left(\Psi\left(f_{*}\right)\right)$, we can even make sure that $\zeta_{i}\left(\Psi\left(f_{t}\right)\right)=\zeta_{i}\left(T_{*}\right)=\zeta_{i}\left(\Psi\left(f_{*}\right)\right)$ for each $t \in[0,1]$ (so the kneading invariant of the $i$-th critical point does not change as $t$ varies in $[0,1])$. If $\zeta_{i}(\Psi(f))>\zeta_{i}(\Psi(\tilde{f}))$, these bifurcations occur in reverse. If $\zeta_{i}(\Psi(f))=\zeta_{i}(\Psi(\tilde{f}))$, then the $i$-th kneading invariant of $f$ and $\tilde{f}$ are the same, and the argument from the beginning of the proof
applies for this critical point. Since one can apply this argument for each periodic attractor simultaneously, the proposition follows.

The previous proof in particular showed:
Theorem 5.12. Given a map $f \in P^{b}$, the set $\mathcal{H}(f)$ of maps $\tilde{f} \in P^{b}$ with the same kneading invariant as forms a connected set. Moreover, if $\tilde{f} \in \mathcal{H}(f)$ then $\overline{\mathcal{P H}(f)} \cap \overline{\mathcal{P H}(\tilde{f})} \neq \emptyset$.

### 5.6. Almost continuity of $\Psi$.

Proposition 5.13 ( $\Psi$ is almost continuous). $\Psi: P^{b} \rightarrow \mathcal{S}_{*}^{b}$ is 'almost continuous' in the following sense. Assume that $f_{n} \rightarrow f$ where $f_{n}, f \in P^{b}$ and $f_{n} \in \mathcal{A}_{b}$. Then there exists $T_{*} \in \mathcal{S}_{*}^{b}$ so that any limit of $T_{n} \in\left[\Psi\left(f_{n}\right)\right]$ is contained in $\left[T_{*}\right]$ and so that $\Psi(f) \in\left[T_{*}\right]$.
Remark 5.14. It is not necessarily true that the limit of $T_{n} \in\left[\Psi\left(f_{n}\right)\right]$ is contained in $[\Psi(f)]$. Indeed, let $f$ be a cubic map so that its left critical point is a fixed point (and is a maximum). Then $\Psi(f)$ is equal to the map $T_{0}$ from Figure 10 . Since $\left[T_{0}\right]=\left\{T_{0}\right\}$ whereas for maps $f_{n} \rightarrow f$ with $f_{n} \in \mathcal{A}_{b}$ one has that $\left[\Psi\left(f_{n}\right)\right]$ is equal to the union of the two triangles in the figure. Note that $f \notin \mathcal{A}_{b}$.

Proof of Proposition 5.13. Take $T_{n}=\Psi\left(f_{n}\right)$ and $T=\Psi(f)$. By taking a subsequence, we can assume that $T_{n}$ converges to some map $\tilde{T}$. Let $Z_{i}, Z_{i, n}$ be the plateaus associated to $T$ and $T_{n}$ respectively. Note that $\langle T\rangle$ has a product structure, see Theorem 4.17 (3) and let $\pi_{i}$ and $\pi_{J}$ be the projections as defined in that theorem. If $c_{i}$ is not eventually mapped onto another critical point, then $\nu_{i}\left(f_{n}\right) \rightarrow \nu_{i}(f)$ as $n \rightarrow \infty$ (in the usual topology on sequence spaces) and so $Z_{i, n} \rightarrow Z_{i}$ as $n \rightarrow \infty$. That is, $\pi_{i}\left(\Psi\left(f_{n}\right)\right) \rightarrow \pi_{i}(\Psi(f))$. In this case define $\zeta_{i}:=\zeta_{i}(\Psi(f))$. If $c_{i}$ is mapped onto another critical point, say $f^{k}\left(c_{i}\right)=c_{j}$ but is not in the immediate basin of a periodic attractor, then $T^{k}\left(Z_{i}\right) \subset \partial Z_{j}$ and for a sufficiently small neighborhood $U_{i}$ of $Z_{i}$ one has $T^{k}\left(U_{i}\right) \cap Z_{j} \subset \partial Z_{j}$. That is, if $T$ has a maximum (minimum) at $Z_{i}$ then $T\left(Z_{i}\right)$ is the left (respectively right) endpoint of a component of $W(T)$. Since $f_{n} \rightarrow f$ and $c_{i}$ is not in the basin of a periodic attractor, one has that $\tilde{T}^{k}\left(\tilde{Z}_{i}\right) \subset \partial \tilde{Z}_{j}$ where $\tilde{T}^{k}\left(\tilde{Z}_{i}\right)$ is possibly the 'other' endpoint of $\tilde{Z}_{j}$. It follows that $\pi_{i}([\Psi(f)])$ is equal to an interval of the form $\left[a_{i}, b_{i}\right]$ and that $\pi_{i}\left(\Psi\left(f_{n}\right)\right)$ converges to an endpoint of this interval. In this case define $\zeta_{i}:=a_{i}$. If $c_{i}$ is periodic, then $f_{n}$ also has an attracting periodic point $p_{n}$ near $c_{i}$. Let $B_{n}$ be the component of the immediate basin containing $p_{n}$ and let $B$ be the component of the immediate basin containing $c_{i}$ and we define $\zeta_{i}:=\zeta_{i}(\Psi(f))$. If $f_{n}\left(\partial B_{n}\right)$ consists of one point, then $f_{n} \in \mathcal{A}_{b}$ implies that the itinerary w.r.t. $f_{n}$ of each critical point in $B_{n}$ agrees with the itinerary w.r.t. $f$ of the corresponding critical point in $B$. In this case $\pi_{i}\left(\Psi\left(f_{n}\right)\right)=\pi_{i}(\Psi(f))$. If $f_{n}\left(\partial B_{n}\right)$ consists of two points, then this no longer needs to be the case, see Example 5.6 and Remark 5.14. In this case $\pi_{i}\left(\left[\Psi\left(f_{n}\right)\right]\right)$ corresponds to a set as in Figure 10 and Figure 11 and $\pi_{i}(\Psi(f))$ is in the closure of this set. Since $f_{n} \in \mathcal{A}_{b}$ it follows that
$\Psi\left(f_{n}\right)$ is equal to the point marked 1 in these figures. Since $\Psi\left(f_{n}\right) \in\left\langle\Psi\left(f_{n}\right)\right\rangle_{b}$, we get that $\left[\Psi\left(f_{n}\right)\right]$ is equal to the closure to this set and we set $\zeta_{i}=\Psi\left(f_{n}\right)$. If $c_{i}$ is mapped to another critical point and in the immediate basin of a periodic attractor, then the same argument goes through. The map $T_{*}$ for which $\zeta_{t}\left(T_{*}\right)=\zeta_{i}$ for $i=1, \ldots, d$ and where $\zeta_{i}$ is chosen as above, has the required properties.

## 6. Proof of the Main Theorem

In this section we shall prove the Main Theorem, assuming Theorem 2.4 (which will be proved as Theorem 7.1 in the next section). First let us prove the following

Theorem 6.1 ('Fibers' of $\Psi$ are connected). Assume that $K$ is a closed and connected subset of $\mathcal{S}_{*}^{b}$ with the property that if $T \in K$ then $[T] \subset K$. Then $\Psi^{-1}(K)$ is connected.

Remark 6.2. Note that $\mathcal{S}_{*}^{b}$ is not a closed subset of $\mathcal{S}^{b}$. Therefore we merely assume that $K$ is a closed subset in the relative topology of $\mathcal{S}_{*}^{b}$ meaning that if $T_{n} \in K$ converges to $T \in \mathcal{S}_{*}^{b}$ then $T \in K$.

Proof. Take a closed connected set $K \subset \mathcal{S}_{*}^{b}$, and assume by contradiction that $\Psi^{-1}(K)$ is not connected. This means that there are disjoint open sets $U_{1}, U_{2} \subset P^{b}$ so that $U_{1} \cup U_{2} \supset \Psi^{-1}(K)$ and $C_{i}:=U_{i} \cap \Psi^{-1}(K) \neq \emptyset, i=1,2$. Write $\left[\Psi\left(C_{i}\right)\right]:=$ $\cup_{f \in C_{i}}[\Psi(f)]$.
Claim 1: $\left[\Psi\left(C_{1}\right)\right] \cup\left[\Psi\left(C_{2}\right)\right] \supset K$. Indeed, it follows from surjectivity (Proposition 5.9) that for every $T \in K$ there exists $f \in P^{b}$ such that $T \in[\Psi(f)]$. Since $[\Psi(f)] \cap K \neq \emptyset$ we have by assumption that $[\Psi(f)] \subset K$ and therefore $f \in \Psi^{-1}(K)$. Therefore $f \in C_{1} \cup C_{2}$ and since $T \in[\Psi(f)]$, the claim follows.
Claim 2: $\left[\Psi\left(C_{i}\right)\right] \cap K$ is closed (again in the relative topology of $\left.\mathcal{S}_{*}^{b}\right)$. To see this, take a sequence $T_{n} \in\left[\Psi\left(f_{n}\right)\right] \cap K$ with $f_{n} \in C_{i}$. By Theorem 4.17 (4) and Proposition 5.9 we can assume that $f_{n} \in \mathcal{A}_{b}$. By considering subsequences we may assume that $T_{n} \rightarrow T$ for some $T \in K$ and $f_{n} \rightarrow f$ for some $f \in U_{i}$. Since $\left[\Psi\left(f_{n}\right)\right] \cap K \neq \emptyset$ we have $\left[\Psi\left(f_{n}\right)\right] \subset K$. Because of this and because $f_{n} \in \mathcal{A}_{b}$, continuity (Proposition 5.13) implies that $[\Psi(f)] \subset K$. Hence $f \in \Psi^{-1}(K) \subset U_{1} \cup U_{2}$ and, since $f_{n} \in C_{i} \subset U_{i}$ converges to $f$, also $f \in U_{i} \cap \Psi^{-1}(K)=C_{i}$. This completes the proof of Claim 2 .
Claim 3: $\left[\Psi\left(C_{1}\right)\right] \cap\left[\Psi\left(C_{2}\right)\right] \neq \emptyset$. This follows from the connectedness of $K$ and Claims 1 and 2.

By Proposition 5.9, there exist therefore $f_{i} \in C_{i} \cap \mathcal{A}_{b}$ such that $\left[\Psi\left(f_{1}\right)\right] \cap\left[\Psi\left(f_{2}\right)\right] \neq \emptyset$. By injectivity (Proposition 5.10), this implies that

$$
\mathcal{P H}\left(f_{1}\right) \cap \overline{\mathcal{P H}\left(f_{2}\right)} \neq \emptyset \quad \text { or } \quad \overline{\mathcal{P} \mathcal{H}\left(f_{1}\right)} \cap \mathcal{P H}\left(f_{2}\right) \neq \emptyset .
$$

Moreover, by Proposition 5.4, $\Psi\left(\mathcal{P H}\left(f_{i}\right)\right) \subset\left[\Psi\left(f_{i}\right)\right] \subset K$. Hence $\mathcal{P H}\left(f_{i}\right) \subset$ $\Psi^{-1}\left(\left[\Psi\left(f_{i}\right)\right]\right) \subset C_{i}$. Since $\mathcal{P} \mathcal{H}\left(f_{1}\right)$ and $\mathcal{P} \mathcal{H}\left(f_{2}\right)$ are both connected, this contradicts that $C_{i} \subset U_{i}$ with $U_{1}, U_{2}$ disjoint.

Proof of Main Theorem. By Theorem 7.1(and the remark below) it follows that level sets of $h_{\text {top }}: \mathcal{S}_{*}^{b} \rightarrow \mathbb{R}$ are connected. Moreover, $h_{\text {top }}: P^{b} \rightarrow \mathbb{R}$ agrees with $h_{\text {top }} \circ \Psi$. Because the topological entropy of each map in $[T]$ is the same, Theorem 6.1 shows that the isentropes lift to connected sets in $P^{b}$. Similarly, the set $I\left(h_{0}\right.$ is connected.

## 7. Isentropes in $\boldsymbol{\mathcal { S }}_{*}^{\boldsymbol{b}}$ are contractible.

Recall from Definition 4.12 that $\mathcal{S}_{*}^{b}$ is the collection of non-degenerate stunted sawtooth maps $T \in \mathcal{S}^{b}$. That is, by definition, if $\mathcal{J}:=\left[Z_{i}, Z_{j}\right]$ is the convex hull of $Z_{i}$ and $Z_{j}$, and there is $n \geq 0$ such that $T^{n}(\mathcal{J})$ is a point (so $\left(Z_{i}, Z_{i+1}\right)$ forms a wandering pair), then $T \in \mathcal{S}_{*}^{b}$ means that $T^{n}(\mathcal{J})$ is eventually mapped into the closure of the immediate basin a periodic plateau.

Of course, if $\left(Z_{i}, Z_{j}\right)$ is a wandering pair, then all plateaus between $Z_{i}$ and $Z_{j}$ form wandering pairs as well. The subset $\mathcal{S}_{*}^{b} \subset \mathcal{S}^{b}$ is chosen because $\Psi: P^{b} \rightarrow \mathcal{S}^{b}$ fails to be surjective in a serious way (whereas $\Psi: P^{b} \rightarrow \mathcal{S}_{*}^{b}$ is almost surjective in the sense of Proposition 5.9). Indeed, if $T \in \mathcal{S}^{b} \backslash \mathcal{S}_{*}^{b}$ has a non-preperiodic wandering pair $\left(Z_{i}, Z_{j}\right)$ and $\Psi(f) \in[T]$, then $f$ has a wandering interval $\left[c_{i}, c_{j}\right]$. However, it is well-known (see e.g. [26]) that polynomials, and in fact $C^{2}$ interval maps with non-flat critical points, have no wandering intervals.

This section is devoted to proving Theorem 2.4, which we restate as
Theorem 7.1. Let $L(h)=\left\{T \in \mathcal{S}_{\epsilon}^{b} \quad ; \quad h_{\text {top }}(T)=h\right\}$ and $L_{*}(h)=L(h) \cap \mathcal{S}_{\epsilon, *}^{b}$. Then

- for every $h_{0} \in[0, \log (b+1)]$, the level set $L_{*}\left(h_{0}\right)$ is a contractible subset of $L\left(h_{0}\right)$;
- $\left.L_{*}\left(h_{0}^{+}\right):=L_{*}\left(h_{0}\right) \cap \operatorname{closure}\left(\left\{T \in \mathcal{S}_{*}^{b} ; h_{\text {top }}(T)\right)>h_{0}\right\}\right)$ is contractible.

Remark 7.2. The sets $L_{*}\left(h_{0}\right)$ has the property that if $T$ is contained in one of these sets, then $[T]$ is contained also in this set. This property holds for $L_{*}\left(h_{0}\right)$ since each map in $[T]$ has the same topological entropy. Since $L_{*}\left(h_{0}^{+}\right)$is connected and each set of the form $[T]$ is connected, the set

$$
\left[L_{*}\left(h_{0}^{+}\right)\right]:=\cup_{T \in L_{*}\left(h_{0}^{+}\right)}[T]
$$

is also connected.

Throughout this section we fix $h_{0} \in[0, \log (b+1)]$, although we separate the easier cases $h_{0}=0$ and $h_{0}=\log (b+1)$, see Sections 7.4 and 7.5 .

That $L\left(h_{0}\right)$ is contractible was already proved in Theorem 6.1 in [31]. Contractibility of $L_{*}\left(h_{0}\right)$ is much more difficult, and we have to adjust the proof of 31] in a delicate way. The proof involves the construction of a retract $R$ composed of entropy decreasing deformations (to contract $L(h)$ to a single point) and entropy increasing deformations (to keep $L(h)$ within itself). The problem is to keep $L_{*}(h)$ within itself under continuous action of the retract. To this end we are forced to compose $R$ of altogether five deformations, with some auxiliary deformations. We use the letters $\Gamma, \gamma, \hat{\Gamma}$ to indicate entropy increasing deformations, and $\delta, \hat{\delta}, \Delta$ for entropy decreasing deformations. The deformation $\beta$ will not change entropy. The letters $R$ and $r$ stand for retract.

Before we are able to give the proof of this theorem we will develop the necessary ingredients.
7.1. The piecewise affine case. An interval $K$ is a renormalization interval an interval map $f$ if $f^{n}(K) \subset K$ for some $n \geq 1$ and $f^{n}(\partial K) \subset \partial K$. If $n=1$, and $K=I$, then this is a renormalization interval only in a trivial sense, but we still want to consider it as such. The set $\operatorname{orb}(K)=K \cup f(K) \cup \cdots \cup f^{n-1}(K)$ is called a renormalization cycle.

It is well-known [30], that every interval map of entropy $h>0$ is semi-conjugate to a piecewise affine interval map with slope $\pm e^{h}$. The semi-conjugacy is a monotone map, and collapses every interval that doesn't contribute to the exponential growth rate of the lapnumber; these are wandering intervals, basins of periodic attractors and possibly renormalization intervals, as well as intervals that map into those.

The following lemma will be used at several places in the rest of the proof.
Lemma 7.3. Assume that $F:[0,1] \rightarrow[0,1]$ is a continuous, piecewise affine map with at most finitely many plateaus $Z_{i}$, and slope $\left|F^{\prime}\right|>1$ outside these plateaus. Suppose that $Z$ is a turning plateau (or point) in a minima renormalization cycle $\operatorname{orb}(K)$, such that no neighborhood of $Z$ is ever mapped into a plateau. Let $J$ be a neighborhood of $Z$ so that $F(\partial J)$ is a single point, then the function $\tilde{F}$ defined by

$$
\tilde{F}(x)= \begin{cases}F(\partial J) & \text { if } x \in \bar{J} \\ F(x) & \text { if } x \notin \bar{J}\end{cases}
$$

satisfies $h_{\text {top }}(\tilde{F} \mid \operatorname{orb}(K))<h_{\text {top }}(F \mid \operatorname{orb}(K))$.
Remark 7.4. The origins of the following proof are somewhat nebulous to us. Jozef Bobok drew our attention to the argument, ascribing it to Sasha Blokh, but we haven't been able to locate a precise source. Related results were proved by Boyland [5] and Block ${ }^{6}$ Ledis [2].

[^1]Proof. The inequality $\tilde{h}:=h_{\text {top }}(\tilde{F}) \leq h:=h_{\text {top }}(F)$ follows directly from the definition of $\tilde{F}$; in fact $|J| \mapsto h_{\text {top }}(\tilde{F})$ is a decreasing function. However, we need to prove strict inequality. Let $Y=\operatorname{orb}(K) \backslash \cup_{n} F^{-n}\left(\operatorname{int}\left(\cup_{i} Z_{i}\right)\right)$. Since $\operatorname{orb}(K)$ is a minimal cycle and the derivatives $\left|F^{\prime}\right|>1$ on $Y$, the restriction $F: Y \rightarrow Y$ is transitive and supports a unique measure of maximal entropy, see [17]. The assumption on $Z$ implies that $\partial Z \subset Y$ and in fact $\mu(J \backslash Z)>0$.

Now $\tilde{F} \mid \operatorname{orb}(K)$ is entropy-preservingly semi-conjugate (say via $\psi$ ) to a map with slope $\pm e^{\tilde{h}}$. Let $\tilde{\nu}$ be the measure of maximal entropy of this map, and $\nu=\tilde{\nu} \circ \psi$. Then $0=\nu(\tilde{F}(J)) \geq \nu(J)$, because $\nu$ is non-atomic. It follows that $\operatorname{supp}(\nu) \cap \operatorname{int}(J)=\emptyset$, and definitely $\nu \neq \mu$, whilst at the same time $\nu$ is not only $\tilde{F}$-invariant, but $F$ invariant as well. Since $\mu$ is the unique measure of maximal entropy of $F \mid \operatorname{orb}(K)$, it follows that $\tilde{h}<h$.
7.2. Increasing the entropy of maps in $\mathcal{S}: \gamma_{t}$ and $\Gamma_{\boldsymbol{t}}$. The stunted seesaw map $T:[-e, e] \rightarrow[-e, e]$ is uniquely determined by the parameters $\left(\zeta_{1}, \ldots, \zeta_{b}\right)$, and so we can define the norm

$$
\begin{equation*}
\operatorname{dist}(T, \tilde{T})=\max _{i \in\{1, \ldots, b\}}\left|\zeta_{i}-\tilde{\zeta}_{i}\right| \tag{9}
\end{equation*}
$$

Write $T \leq \tilde{T}$ if the parameters satisfy $\zeta_{i} \leq \tilde{\zeta}_{i}$ for all $i$. Similarly $T<\tilde{T}$ if $T \leq \tilde{T}$ and $\zeta_{i}<\tilde{\zeta}_{i}$ for at least one $i$. Notice that

$$
\begin{equation*}
T \leq \tilde{T} \text { implies } h_{\text {top }}(T) \leq h_{\text {top }}(\tilde{T}) \tag{10}
\end{equation*}
$$

Construction of $\gamma_{t}$ : Let $\gamma_{t}$ linearly increase all parameters: $\gamma_{t}: \zeta_{i} \mapsto \zeta_{i}+2 e t$, as long as they do not map to $\pm e$.

Let $W^{o}(T)$ be the components of the domain of the first entry map to the interior of plateaus of $T$ :

$$
\begin{aligned}
W^{o}(T):= & \{x ; \exists j \in\{1, \ldots, b\} \text { and } k \geq 0 \text { such that } \\
& \left.T^{k}(x) \in \operatorname{int}\left(Z_{j}\right) \text { and } x, T(x), \ldots, T^{k-1} \notin Z_{j}\right\} .
\end{aligned}
$$

We say that $T$ satisfies the $\beta$-property if

$$
\begin{equation*}
\text { no interval of the form } \mathcal{J}_{i}=\left[Z_{i}, Z_{i+1}\right] \text { is contained in } W^{o}(T) . \tag{11}
\end{equation*}
$$

Lemma 7.5. Assume that $T$ satisfies the $\beta$-property (11). Then $\gamma_{t}(T) \in \mathcal{S}_{*}^{b}$ for all $t>0$.

Proof. For every fixed integer $u \geq 0$, as $t$ increases, the image of the interval $\mathcal{J}_{i, t}=$ [ $Z_{i, t}, Z_{i+1, t}$ ] under the $u$-th iterate of the map $\gamma_{t}(T)$ becomes larger while the sizes of plateaus shrink. If the $\beta$-property (11) holds then it follows that $\mathcal{J}_{i, t}$ cannot be mapped into a non-periodic plateau for $t>0$ and therefore no new wandering pairs can be created by the deformation $\gamma_{t}$ as $t$ increases. Also $\gamma_{t}(T)$ has no non-trivial blocks of touching plateaus for $t>0$. This implies the lemma.

Construction of $\boldsymbol{\Gamma}_{\boldsymbol{t}}$ : The deformation $\Gamma_{t}$ uses $\gamma_{t}$ and the following observation. If $\operatorname{dist}(T, \tilde{T})<\varepsilon /(2 e)$, then $\operatorname{dist}\left(\gamma_{t}(T), \gamma_{t}(\tilde{T})\right)<\varepsilon$ for every $t>0$. It follows that

$$
\gamma_{t-\varepsilon}(T)<\gamma_{t}(\tilde{T})<\gamma_{t+\varepsilon}(T)
$$

and so by (10),

$$
h_{t o p}\left(\gamma_{t-\varepsilon}(T)\right) \leq h_{t o p}\left(\gamma_{t}(\tilde{T})\right) \leq h_{t o p}\left(\gamma_{t+\varepsilon}(T)\right) .
$$

Hence the function

$$
\begin{equation*}
t_{\max }(T):=\max \left\{t \geq 0 \quad ; \quad h_{\text {top }}\left(\gamma_{t}(T)\right)=h_{0}\right\} \tag{12}
\end{equation*}
$$

is continuous in $T$ provided $h_{t o p}(T) \leq h_{0}$. In particular, $\Gamma_{t}(T):=\gamma_{t \cdot t_{\max }(T)}(T)$ is continuous in $t$ and $T$ as well. By construction,

$$
\begin{align*}
& \left.\left\{T \in \mathcal{S}^{b} ; h_{\text {top }}(T)<h_{0}, T \text { satisfies the } \beta \text {-property } \sqrt{11}\right)\right\} \Longrightarrow  \tag{13}\\
& \Gamma_{t}(T) \in \cup_{s \leq h_{0}} L_{*}(s), \forall t \in(0,1] .
\end{align*}
$$

Moreover, if $T$ satisfies the $\beta$-property (11) then $\Gamma_{t}(T)$ also satisfies the $\beta$-property (11) for $t \in(0,1]$.
7.3. Decreasing entropy of maps in $\mathcal{S}^{\boldsymbol{b}}: \boldsymbol{\delta}_{\boldsymbol{t}}, \hat{\boldsymbol{\delta}}_{\boldsymbol{t}}$ and $\boldsymbol{r}_{\boldsymbol{t}}$. In this section we define the basic operations to decrease entropy, although later we will need refined versions of them.

Construction of $\boldsymbol{\delta}_{\boldsymbol{t}}$ : Define the 'sign'

$$
\operatorname{sgn}\left(Z_{i}\right)= \begin{cases}0 & \text { if } Z_{i} \text { touches another plateau }  \tag{14}\\ 1 & \text { otherwise }\end{cases}
$$

Now deform $T$ according to the flow defined by the system of differential equations on the parameters, where we will now indicate the $t$-dependence by $\zeta_{i, t}(T)$.

$$
\begin{aligned}
\frac{d \zeta_{i, t}}{d t}(T) & = \begin{cases}-2 \operatorname{sgn}\left(Z_{i}\right) e & \text { if } \zeta_{i, t}(T) \in(-e, e) \\
0 & \text { otherwise }\end{cases} \\
\left.\zeta_{i, t}(T)\right|_{t=0} & =
\end{aligned}
$$

Let us denote the resulting deformation by $\delta_{t}$; it continuously decreases/increases the height of a plateau if it is local maximum/minimum of $T$ until this plateau touches a neighboring plateau, or reaches the boundary of $I$. Furthermore, if $T$ has no touching plateaus then $h_{\text {top }}\left(\gamma_{s} \circ \delta_{t}(T)\right) \leq h_{\text {top }}(T)$ for $s \in(0, t)$ and $t>0$ sufficiently small. Note, however, that $\delta_{t}(T)=\delta_{s}(T)$ for $s, t \geq 1$ but this might mean that each plateau of $\delta_{t}(T)$ touches another plateau, and therefore this is not a guarantee that $h_{\text {top }}\left(\delta_{t}(T)\right)=0$. Therefore this deformation, although necessary (see Section 7.9), is not sufficient for our purposes.

Construction of $\hat{\delta}_{t}$ : A natural variant is the deformation $\hat{\delta}_{t}$ which will widen plateaus in order to decrease entropy to 0 . The difference in the deformations $\delta_{t}$ and $\hat{\delta}_{t}$ is that $\delta_{t}$ decreases the parameter $\zeta_{i}$ until the corresponding plateau $Z_{i}$ touches another plateau (or reaches $e$ ); blocks (consisting of more than one plateau) do not
move under $\delta_{t}$. By contrast, the deformation $\hat{\delta}_{t}$ will also move blocks of plateaus, provided they form a local extremum, which happens whenever the block consists of an odd number of plateaus. Blocks of an even number of plateaus are not moved by $\hat{\delta}_{t}$ (unless an extra plateau joins the block). As it turns out, this may introduce new wandering pairs.

Define the 'sign' for plateaus that are part of a block of plateaus:

$$
\widehat{\operatorname{sgn}}\left(Z_{i}\right)=\left\{\begin{array}{cc}
0 & \text { if } Z_{i} \text { is part of a block of an even number of plateaus, }  \tag{15}\\
1 & \text { if } Z_{i} \text { touches no other plateau or is an odd-numbered } \\
-1 & \text { plateau in a block of an odd number of plateaus, } \\
\text { if } Z_{i} \text { is an even-numbered plateau in a block of an odd } \\
\text { number of plateaus. }
\end{array}\right.
$$

Note that $\widehat{\operatorname{sgn}}\left(Z_{i}\right)$ depends not only on $i$ but also on $T$, and $\widehat{\operatorname{sgn}}\left(Z_{i}\right)= \pm 1$ means that $T$ has a local extremum at the block of plateaus that $Z_{i}$ is part of. We deform $T$ according to the flow defined by the system of differential equations on the parameters.

$$
\begin{aligned}
\frac{d \zeta_{i, t}}{d t}(T) & = \begin{cases}-2 \cdot \widehat{\operatorname{sgn}}\left(Z_{i}\right) \cdot e & \text { if } \zeta_{i, t}(T) \in(-e, e) \\
0 & \text { otherwise }\end{cases} \\
\left.\zeta_{i, t}(T)\right|_{t=0} & =\zeta_{i}(T)
\end{aligned}
$$

The differential equation defines a continuous deformation $\hat{\delta}_{t}$ with the property that $t \mapsto \hat{\delta}_{t}(T)$ (not necessarily strictly) decreases the topological entropy. During the deformation blocks can collide, and then the combined larger blocks are deformed according to the same rule. (As a result $\widehat{\operatorname{sgn}}\left(Z_{i}\right)$ can change during the deformation.)

Construction of the retract $\boldsymbol{r}_{\boldsymbol{t}}$ : If $b$ is odd and $t \geq 1$, then all plateaus of $\hat{\delta}_{t}(T)$ will touch and the map $\hat{\delta}_{t}(T)$ is constant $\pm e$. If $b$ is even and $t \geq 1$, then each map $\hat{\delta}_{t}(T)$ will be monotone (with some blocks of touching plateaus). More precisely, if $t \geq 1$ then $\hat{\delta}_{t}(T) \in \Sigma_{\epsilon}^{b}$ where

$$
\Sigma_{\epsilon}^{b}= \begin{cases}\left\{T_{0}(x) \equiv \pm e\right\} & \text { if } b \text { is odd and } \epsilon=\mp 1 \\ \left\{\text { monotone maps in } \mathcal{S}_{\epsilon}^{b}\right\} & \text { if } b \text { is even } .\end{cases}
$$

Since $\Sigma_{\epsilon}^{b}$ is a singleton in the first case and a simplex in the second case, there exists a continuous retract $r_{t}: \Sigma_{\epsilon}^{b} \rightarrow \Sigma_{\epsilon}^{b}$ with $r_{0}=i d$ and $r_{1} \equiv T_{0}$ for $T_{0}$ some map in $\Sigma_{\epsilon}^{b}$.

Construction of a retract $R_{t}$ of an isentrope of $\mathcal{S}^{b}$ : If we only had to construct a retract of an isentrope of $\mathcal{S}^{b}$ then we could finish the construction of a deformation $R_{t}$ as follows. Define

$$
R_{t}= \begin{cases}\Gamma_{3 t} & \text { for } t \in[0,1 / 3] \\ \Gamma_{1} \circ \hat{\delta}_{3 t-1} & \text { for } t \in[1 / 3,2 / 3] \\ \Gamma_{1} \circ r_{3 t-2} \circ \hat{\delta}_{1} & \text { for } t \in[2 / 3,1]\end{cases}
$$

Obviously this defines the required retract of an isentrope with the space $\mathcal{S}^{b}$ to a point. However, as we will see, this is still insufficient for our purposes because it is not a retract of an isentrope within the smaller space $\mathcal{S}_{*}^{b}$. In the remainder of this section we will show how to modify this construction to obtain a retract of an isentrope in $\mathcal{S}_{*}^{b}$.
7.4. The case $h_{0}=\log (b+1)$. There is only one stunted sawtooth map with entropy $h_{\text {top }}(T)=\log (b+1)$ (just as there is only one polynomial of give $\epsilon$ with entropy $h_{\text {top }}(f)=\log (b+1)$, namely the Chebyshev polynomial. Hence this level is trivial.
7.5. The retract $\boldsymbol{R}_{\boldsymbol{t}}$ for the case $\boldsymbol{h}_{\mathbf{0}}=\mathbf{0}$. Let us give the proof in the case that $h_{0}=0$, as the construction will be much easier in this case. Consider a map $T \in \mathcal{S}_{*}^{b}$ with zero topological entropy. Let us first review some results on the renormalization structure of such maps.

We say that $T \in \mathcal{S}^{b}$ has a 2 -renormalizable, if there exists an interval $K$ so that $K$ and $T(K)$ have disjoint interiors and $T^{2}(K) \subset K$.

Lemma 7.6. Let $T \in \mathcal{S}^{b}$ and assume that $h_{\text {top }}(T)=0$. If $T$ is not 2 -renormalizable, then the $\omega$-limit set of each point in $I$ is a fixed point of $T$.

Proof. Since $T$ has zero topological entropy, it follows that there exists no interval $J$ so that $T^{2}(\partial J) \subset \partial J$ and $T^{2}(J) \supset J$. From this, and since $T$ is not 2-renormalizable, this well-known lemma easily follows.

Remark: If $T$ is 2-renormalizable then one can apply the lemma again to $T^{2}: K \rightarrow$ $K$. If one can repeat this infinitely often, then the map is infinitely renormalizable. In this case, for each $k \geq 0$, the map $T$ has one or more periodic points of period $2^{k}$ and no other periodic points. If $T$ is not infinitely often renormalizable, then each point is eventually mapped into a periodic point of $T$ or is in the basin of a periodic orbit of period $2^{k}$.

Lemma 7.7. Let $T \in \mathcal{S}_{*}^{b}$ with $h_{0}=h_{\text {top }}(T)=0$ and assume that $t>0$. Then there exists $k_{0}<\infty$ such that all periodic attractors of $\hat{\delta}_{t}(T)$ have period $2^{k}, k \leq k_{0}$. Moreover, each point is either (pre-)periodic or in the basin of one of the periodic attractors of $\hat{\delta}_{t}(T)$.

Proof. From the previous lemma it follows that each interval map of zero entropy and finite modality has only periodic points of period $2^{k}$ for $k \in \mathbb{N}$. Take $T \in \mathcal{S}_{*}^{b}$ such that $h_{\text {top }}(T)=0$. If $Z_{i}$ is a plateau with an infinite orbit, then $T$ must be infinitely renormalizable, i.e., there exists a sequence of periodic intervals $K_{u}, u \in \mathbb{N}$, with period $2^{u}$ such that $\cap_{u} K_{u} \supset Z_{j}$ for some $j$, and $\omega\left(Z_{i}\right)=\omega\left(Z_{j}\right)$. In fact, orb $\left(K_{u}\right)$ can contain more plateaus, but since the period of $K_{u}$ tends to infinity as $u \rightarrow \infty$,
and there are only $b$ plateaus, we can assume (by an appropriate choice of $K_{u}$ ) that there exists $n_{u} \rightarrow \infty$ as $u \rightarrow \infty$ such that $T^{n}\left(K_{u}\right)$ does not intersect any plateau for $0<n<n_{u}$. Therefore $\left|T\left(K_{u}\right)\right| \rightarrow 0$ as $u \rightarrow \infty$.

Since $T \in \mathcal{S}_{*}^{b}$ there exists $t_{1} \in(0, t)$ and $\eta>0$ such that all plateaus of $T$ move at least $\eta$ when $t^{\prime}$ moves from 0 to $t_{1}$ (for $t^{\prime}$ small $\hat{\delta}_{t^{\prime}}$ agrees with $\delta_{t^{\prime}}$ ). For $u$ so large that $\left|T\left(K_{u}\right)\right|<\eta$, this means that $K_{u}$ is no longer invariant under $\hat{\delta}_{t_{1}}(T)$ and so this map is not infinitely renormalizable anymore. Instead, there is $k_{0}$ such that every plateau of $\hat{\delta}_{t_{1}}(T)$ is (eventually) periodic with period $2^{k}$ for some $k \leq k_{0}$. If we increase $t^{\prime}$ further from $t_{1}$ to $t$, each periodic attractor remains but can undergo period halving bifurcations. So all periodic orbits of $\hat{\delta}_{t}(T)$ have period $2^{k}, k \leq k_{0}$.

Now define a retract $R_{t}$ of the zero-entropy level set of $\mathcal{S}_{*}^{b}$ as follows:

$$
R_{t}= \begin{cases}\hat{\delta}_{2 t} & \text { for } t \in[0,1 / 2] \\ r_{2 t-1} \circ \hat{\delta}_{1} & \text { for } t \in[1 / 2,1]\end{cases}
$$

Lemmas 7.7 and 7.6 (and the remark above) imply that under $\hat{\delta}_{2 t}(T)$ (resp. $r_{2 t-1} \circ$ $\hat{\delta}_{1}(T)$ ), each plateau is contained in the closure of a component of the basin of one of the periodic attractors of $\hat{\delta}_{2 t}(T)$ (resp. $r_{2 t-1} \circ \hat{\delta}_{1}(T)$ ). Hence $\hat{\delta}_{t}(T), r_{t} \circ \hat{\delta}_{1}(T) \in \mathcal{S}_{*}^{b}$. Thus we obtain a retract of $\mathcal{S}_{*}^{b}$ and proved Theorem 7.1 in the zero entropy case.

The remainder of this section will deal with the case $h_{0}>0$, which is plagued with additional difficulties.
7.6. The retract $R_{t}$ when $h_{0}>0$ and the trouble with $\hat{\delta}_{t}$. As mentioned at the end of Subsection 7.3 the retract is insufficient for our purposes: we need to construct a retract of an isentrope of $\mathcal{S}_{*}^{b}$ (so the deformation is not allowed to leave the space $\left.\mathcal{S}_{*}^{b}\right)$. The hurdle we have to overcome is that if $T \in \mathcal{S}_{*}^{b}$, then $\hat{\delta}_{t}(T)$ need no longer be in $\mathcal{S}_{*}^{b}$ for $t>0$, because the deformation $t \mapsto \hat{\delta}_{t}(T)$ can create wandering pairs $\left(Z_{i}, Z_{j}\right)$. To resolve this issue, the aim is to ensure that the deformation $\Gamma_{t}$ (or a similar deformation) will be able to 'undo' these wandering pairs. In view of property (13) on page 47 we will construct a deformation $\beta_{t}$ with the property that if $T \in \mathcal{S}^{b}$ then $\beta_{1}(T) \in \mathcal{S}_{*}^{b}$. It does this by deforming $T$ in such a way that $\beta_{1}(T)$ never eventually maps an interval of the form $\left[Z_{i}, Z_{i+1}\right]$ into the interior of another plateau. However, (13) only applies to maps $T \in \overline{\mathcal{S}_{*}^{b}}$ with topological entropy $h_{0}$. For this reason we need to define a more subtle way of 'decreasing' and 'increasing' the map $T$ while the topological entropy remains equal to $h_{0}$. These analogues of $\hat{\delta}$ and $\Gamma_{1}$ may act on some of the plateaus while leaving some others alone. The challenge will be to define this as a continuous deformation. To achieve this, we will introduce three additional deformations.
$\hat{\Gamma}_{t}$ : To increase the topological entropy more carefully by increasing some (but possibly not all) $\zeta_{i}$ 's, so we essentially increase each $\zeta_{i}$ 'as far as possible'. This is the purpose of $\hat{\Gamma}_{t}$ defined in Section 7.8.
$\hat{\Delta}_{t}$ : To decrease the topological entropy more carefully in such a way that if $\hat{\Delta}_{t}\left(T^{\prime}\right)$ or $\Gamma_{t}\left(T^{\prime}\right)$ does not move certain plateaus (because otherwise the entropy would become too large), then we 'can assume' that $T^{\prime} \in \mathcal{S}_{*}^{b}$. This is the purpose of the deformations $\Delta_{t}$ and $\hat{\Delta}_{t}$ defined in Section 7.9 .
$\beta_{t}$ : Finally, we want to ensure that we only need to apply $\hat{\Gamma}_{s}$ to maps $T^{\prime} \in \mathcal{S}^{b}$ with the property that a convex hull $\left[Z_{i}, Z_{i+1}\right]$ is never eventually mapped into the interior of another plateau (i.e., only to maps with $T^{\prime} \in \overline{\mathcal{S}_{*}^{b}}$ ). This means that $\hat{\Gamma}_{s}\left(T^{\prime}\right) \in \mathcal{S}_{*}^{b}$ for every $s>0$. Unfortunately, this may not be enough because it may happen that $h_{\text {top }}\left(\hat{\Gamma}_{s}\left(T^{\prime}\right)\right)>h_{0}$ for any $s>0$. The deformation $\beta_{t}$, defined in Section 7.7, prevents this, but will not change the entropy.
7.7. The construction of $\boldsymbol{\beta}_{\boldsymbol{t}}$. Consider a map $T^{\prime}=\hat{\delta}_{t}(T)$ so that the convex hull $\mathcal{J}_{i}:=\left[Z_{i}, Z_{i+1}\right]$ of two neighboring plateaus is eventually mapped into the interior of another plateau. If this happens then for $s>0$ small, $\gamma_{s}\left(T^{\prime}\right)$ will still have this property and so in general $\gamma_{s}\left(T^{\prime}\right) \notin \mathcal{S}_{*}^{b}$. To overcome this problem we define another deformation $\beta_{t}$. This deformation does not change topological entropy, and only moves (certain) plateaus which are mapped into other plateaus.

As before, let

$$
\begin{aligned}
W^{o}(T)= & \{x ; \exists j \in\{1, \ldots, b\} \text { and } k \geq 0 \text { such that } \\
& \left.T^{k}(x) \in \operatorname{int}\left(Z_{j}\right) \text { and } x, T(x), \ldots, T^{k-1} \notin Z_{j}\right\} .
\end{aligned}
$$

Let $\mathbf{I}(T)$ be the collection of integers $i \in\{1, \ldots, b\}$ such that $T\left(Z_{i}\right) \in W^{o}(T)$ and $Z_{i}$ is not periodic. Next we say that $T$ has a local maximum (resp. minimum) at $Z_{i}$ if the sawtooth map $S_{0}$ has a local maximum (resp. minimum) at the midpoint of $Z_{i}$. Moreover, given an interval $J$, we define $\partial_{l} J$ and $\partial_{r} J$ to be the left and right endpoint of $J$ respectively. For $i \in \mathbf{I}(T)$ define $v_{i}=T\left(Z_{i}\right)$, let $W_{i}=W_{i}(T)$ be the component of $W^{o}(T)$ containing $v_{i}$ and

$$
\tau_{i}(T)= \begin{cases}\frac{\operatorname{dist}\left(v_{i}, \partial_{l} W_{i}\right)}{\left|W_{i}\right|} & \text { if } T \text { has a maximum at } Z_{i},  \tag{16}\\ \frac{\operatorname{dist}\left(v_{i}, \partial_{r} W_{i}\right)}{\left|W_{i}\right|} & \text { if } T \text { has a minimum at } Z_{i}\end{cases}
$$

We define the deformation $T_{t}=\beta_{t}(T)$ with parameters $\left(\zeta_{1, t}, \ldots, \zeta_{b, t}\right)$ of $T$ as the flow of the following differential equation:

$$
\frac{d \zeta_{i, t}}{d t}=\left\{\begin{array}{ll}
4 \cdot \tau_{i}\left(T_{t}\right) \cdot e & \text { if } i \in \mathbf{I}\left(T_{t}\right), \\
0 & \text { if } i \notin \mathbf{I}\left(T_{t}\right),
\end{array} \quad i=1, \ldots, b .\right.
$$

Note that if $Z_{i}$ contains a periodic attractor or $Z_{i}$ is never mapped into another plateau, then $Z_{i, t}$ does not vary with $t$. As $W^{o}(T)$ consists of preimages of the
interiors of such plateaus, it follows that $W^{o}\left(T_{t}\right)$ is independent of $t$ except if some plateau $Z_{i, t}$ is mapped by $T_{t}$ into the boundary of a component $W$ of $W^{o}\left(T_{t}\right)$ as in the right panel of Figure 14 . Also note that $W^{o}(T)$ depends continuously on $T$ except if some plateau $Z_{i}$ maps into the boundary of a component of $W^{o}(T)$, i.e., eventually maps to the boundary of another plateau. In fact only if it is only discontinuous if $Z_{i}$ 'arches over' this component as in the right panel of Figure 14.

Even though $W_{i}$ does not necessarily depend continuously on $T$, the deformation $(T, t) \mapsto \beta_{t}(T)$ turns out to be well-defined and continuous, see the lemma below. Since $T_{t}$ and $T$ agree outside the union of the closures of components of $W^{o}(T)$, we have $h_{\text {top }}\left(\beta_{t}(T)\right)=h_{\text {top }}(T)$.

Lemma 7.8. The deformation $(T, t) \mapsto \beta_{t}(T)$ is well-defined. Furthermore:
(a) $(T, t) \mapsto \beta_{t}(T)$ is continuous.
(b) The map $\beta_{1}(T)$ satisfies the $\beta$-property (11).

Proof. Part (a): Well-defined and continuity. If $i \in \mathbf{I}(T)$ and plateaus are mapped into the interior of components of $W_{i}(T)$, then $T \mapsto \tau_{i}(T)$ is locally Lipschitz and there is a unique solution. Otherwise, if $i \notin \mathbf{I}(T)$ the following considerations show that $t \mapsto \beta_{t}(T)$ is well-defined. Note that it suffices to consider $t \approx 0$ and maps $\tilde{T}$ near some $T \in \mathcal{S}^{b}$. Let $\tilde{\zeta}_{i, t}$ denote the $i$-th parameter of $\beta_{t}(\tilde{T})$. If $T\left(Z_{i}\right)$ is contained in a component of $W^{o}(T)$, then continuity of $(t, \tilde{T}) \mapsto \tilde{\zeta}_{i, t}$ obviously holds. Next assume that $v_{i}:=T\left(Z_{i}\right)$ is contained in the boundary of a component $W$ of $W^{o}(T)$. By definition this means that $t \mapsto \zeta_{i, t}$ is constant for $t \geq 0$. Moreover, for a nearby map $\tilde{T}$, there are two possibilities: (i) $\tilde{T}\left(\tilde{Z}_{i}\right)$ is not contained in $W^{o}(\tilde{T})$, which means that $t \mapsto \tilde{\zeta}_{i, t}$ is also constant, or (ii) $\tilde{T}\left(\tilde{Z}_{i}\right)$ is contained in a component $\tilde{W}$ of $W^{o}(\tilde{T})$. In this case, since $\tilde{T}$ is close to $T$, either (iia) $\tilde{T}\left(\tilde{Z}_{i}\right)$ is near the boundary of a component $\tilde{W}$ of $W^{o}(\tilde{T})$ or (iib) there exists $k$ so that $T^{k}\left(v_{i}\right) \in \partial Z_{j}$ and $T^{k+1}\left(v_{i}\right) \in \partial W^{\prime}$ where $W^{\prime}$ is a component of $W_{\tilde{\sim}}^{o}(T)$. In the latter case $\tilde{T}\left(\tilde{Z}_{j}\right)$ is near the boundary of a component $\tilde{W}^{\prime}$ of $W^{o}(\tilde{T})$ and $\tau_{j}(\tilde{T}) \approx 1$. These two situations are illustrated in Figures 14 and 15.

If (iia) holds, then either $\tau_{i}(\tilde{T}) \approx 0$ which means that $\tilde{\zeta}_{i, t}$ remains close to $\tilde{\zeta}_{0, i}$ for all $t \in[0,1]$ (see the left panel in Figure 14 ) or $\tau_{i}(\tilde{T}) \approx 1$ which means that $\tilde{\zeta}_{i, t}$ is moving with speed $\approx 4 e$ towards the nearest boundary point of $\tilde{W}$ (see the right panel in Figure 14). Therefore there exists $\tilde{t}>0$ close to zero, so that $\tilde{T}_{\tilde{t}}\left(\tilde{Z}_{i}\right) \in \partial \tilde{W}$ and therefore $\tilde{\zeta}_{i, t}$ remains constant for $t \geq \tilde{t}$.
If (iib) holds, see Figure 15 , then there exists $\tilde{t}>0$ small so that $\tilde{T}_{\tilde{t}}\left(\tilde{Z}_{j, t}\right) \in \partial \tilde{W}^{\prime}$ and $\tilde{W}^{*}$ no longer is a component of $W^{o}\left(\tilde{T}_{\tilde{t}}\right)$ (the interval $\tilde{W}^{*}$ is split into three components of $\left.W^{o}\left(\tilde{T}_{\tilde{t}}\right)\right)$. This means that $\tilde{T}_{\tilde{t}}\left(Z_{i}\right)$ is near a boundary of a component $\tilde{W}_{-}$of $W^{o}\left(\tilde{T}_{\tilde{t}}\right)$ and we can argue as before.


Figure 14. Case (iia) in the proof of continuity. On the left the situation where $\tilde{T}\left(\tilde{Z}_{i}\right)$ is contained in a component $W$ of $W^{o}(T)$ with $\tau_{i}(\tilde{T}) \approx 0$. In this case, $\tilde{\zeta}_{i, t}$ will increase only very slowly with $t$. On the right the situation where $\tilde{T}\left(\tilde{Z}_{i}\right)$ is close to the boundary of an endpoint of $W^{o}(T)$ when $\tau_{i}(\tilde{T}) \approx 1$. In this case $\tilde{\zeta}_{i, t}$ increases with speed $\approx 4 e$ which means that there exists $t>0$ close to 0 so that $\tilde{T}_{t}\left(\tilde{Z}_{i}\right)$ is contained in the boundary of $W$ and then stops. In this figure we set $W=\tilde{W}$, but also when these are different intervals, the same argument holds. Note that in the situation on the right, the component of $W^{o}(T)$ containing $Z_{i}$ does not depend continuously on $T$.


Figure 15. Case (iib) in the proof of continuity. In this case $W$ is a component of $T^{-k}\left(W^{*}\right)$ where $W^{*}$ is a component of $W^{\circ}(T)$ containing a plateau $Z_{j}$ and $T\left(Z_{j}\right)$ is close to the boundary of the component $W^{\prime}$ containing $T\left(Z_{j}\right)$. There exists $t>0$ close to zero, so that $T_{t}\left(Z_{j, t}\right)$ hits the boundary of $W^{\prime}$ and then $W_{t}^{*}$ splits into three components: to the left and right of $Z_{j, t}$ and the interior of $Z_{j, t}$. This means that $W_{t}$ also splits into three components, and at this moment $Z_{i, t}$ moves until $T_{t}\left(Z_{i, t}\right)$ belongs to the boundary of one of these plateaus (which in this situation means close to $v_{i}$ ).

Finally, assume that $T\left(Z_{i}\right)$ is not contained in the closure of a component of $W^{o}(T)$. If $\tilde{T}$ is near $T$ while $\tilde{T}\left(\tilde{Z}_{i}\right)$ is contained in a component $\tilde{W}$ of $W^{o}(\tilde{T})$, then this
component $\tilde{W}$ is small and so for $t>0$ small $\beta_{t}(\tilde{T})\left(\tilde{Z}_{i}\right)$ is contained in the boundary of $\tilde{W}$ (and for $t^{\prime}>t$ this plateau no longer moves).

It follows that in all cases, $(T, t) \mapsto \beta_{t}(T)$ is continuous.
Part (b). If $Z_{i}$ is non-periodic and $\mathcal{J}_{i}$ is mapped into a component $W$ of $W^{o}(T)$, then either $\tau_{i}(T) \geq 1 / 2$ or $\tau_{i+1}(T) \geq 1 / 2$, see Figure 16. By construction, $\tau_{i}\left(\beta_{t}(T)\right) \geq$ $1 / 2$ or $\tau_{i+1}\left(\beta_{t}(T)\right) \geq 1 / 2$ for all $t \geq 0$. Hence $\frac{d \zeta_{i}}{d t}\left(\beta_{t}(T)\right) \geq 2 e$ or $\frac{d \zeta_{i+1}}{d t}\left(\beta_{t}(T)\right) \geq 2 e$, for all $t \in\left[0, t_{0}\right]$ where $t_{0}$ is chosen so that $\beta_{t}(T)\left(\mathcal{J}_{i, t}\right)$ is still contained in $W$ for all $t \in\left[0, t_{0}\right]$. Since for each map $T \in \mathcal{S}^{b}$, the corresponding $\zeta_{i}$ can be at most $2 e$, there exists $t \leq 1$ so that $\beta_{t}(T)\left(\mathcal{J}_{i, t}\right) \in \partial W$. Note that $x \in \partial W$ implies that no iterate of $x$ can be mapped into the interior of a plateau. Hence the claim follows.


Figure 16. In this example, $T\left(Z_{i}\right)$ and $T\left(Z_{i+1}\right)$ are both contained in the same component $W$ of $W^{o}(T)$. We show a situation where $\tau_{i}(T) \approx 1$ whereas $\tau_{i+1}(T)$ is much smaller. As a result $\zeta_{i, t}$ changes much faster than $\zeta_{i+1, t}$ under the deformation, but as soon as $\beta_{t}(T)\left(Z_{i, t}\right)$ hits the boundary of $W$ then $\zeta_{i, t}$ stops moving. The component of $W^{o}\left(T_{t}\right)$ containing $Z_{i}$ splits in three as soon as $T_{t}\left(Z_{i, t}\right)$ hits the boundary of $W$ : one to the left of $Z_{i}$, one to the right of $Z_{i}$ and the interior of $Z_{i}$.
7.8. Increasing entropy of maps more carefully: $\hat{\boldsymbol{\Gamma}}_{\boldsymbol{t}}$. We use the entropy increasing deformation $(T, t) \mapsto \Gamma_{t}(T)=\gamma_{t \cdot t_{\max }(T)}(T)$ (with $t_{\max }$ as in (12) on page 47) until $h_{\text {top }}\left(\Gamma_{t}(T)\right)=h_{0}$. But it is possible that only part of the phase space is responsible for reaching this entropy bound, while in other parts (namely in renormalization cycles), plateaus have not been lifted 'sufficiently'yet. It is essential to perform some version of $\gamma_{s}$ for at least some time $s>0$, so as to resolve (destroy) wandering pairs that may have been created by $\delta_{t}$. Thus in the presence of renormalization intervals, we may need to lift some plateaus faster than others. This subsection explains how this is done.
Remark 7.9. Recall that $T \in \mathcal{S}^{b}$ has renormalization interval $K$ if there is $n \geq 1$ such that $T^{n}(K) \subset K$ and $T^{n}(\partial K) \subset \partial K$. The orbit $\operatorname{orb}(K)=K \cup T(K) \cup \cdots \cup$ $T^{n-1}(K)$ is called a renormalization cycle. Note that $\partial K$ consists of (pre)periodic points which do not depend on $T$, unless they disappear in a saddle node bifurcation.

The only other way by which $K$ can disappear is when $T^{n}(Z) \subset \partial K$ for some plateau $Z$ compactly contained in $K$.

Fix $h_{0} \in(0, \log (b+1))$. Let $K_{i}(T)$ be the smallest renormalization interval for $T$ which contains the omega-limit set $\omega\left(Z_{i}\right)$. If there exists no smallest renormalization interval then we take $K_{i}(T)=\emptyset$. Let

$$
\begin{cases}L_{i}\left(h_{0}^{-}\right)=\operatorname{closure}\left(\left\{T \in \mathcal{S}^{b} \quad ; \quad h_{\text {top }}\left(T \mid \operatorname{orb}\left(K_{i}(T)\right)\right) \leq h_{0}\right\}\right), \\ L_{i}\left(h_{0}^{+}\right)=\operatorname{closure}\left(\left\{T \in \mathcal{S}^{b} \quad ; h_{\text {top }}\left(T \mid \operatorname{orb}\left(K_{i}(T)\right)\right)>h_{0}\right\}\right) .\end{cases}
$$

Note that in this definition $\operatorname{orb}\left(K_{i}(T)\right)$ does indeed depend on $T$. Consequently, the common boundary of these sets contains maps $T$ for which $h_{\text {top }}\left(T \mid \operatorname{orb}\left(K_{i}(T)\right)\right) \leq h_{0}$ but which can be increased by an arbitrarily small change in parameter $\zeta_{j}$ for some $Z_{j} \subset \operatorname{orb}\left(K_{i}(T)\right)$. It does not follow, however, that $h_{\text {top }}\left(T \mid \operatorname{orb}\left(K_{i}(T)\right)=h_{0}\right.$, see Figure 17. In fact, $T \mapsto h_{\text {top }}\left(T \mid \operatorname{orb}\left(K_{i}(T)\right)\right.$ is discontinuous because if $Z_{i}$ is not in the basin of an attractor, a small change in $T$ can make $Z_{i}$ periodic itself, and then $h_{\text {top }}\left(T \mid \operatorname{orb}\left(K_{i}(T)\right)=0\right.$.


$$
\begin{aligned}
& T \in L_{2}\left(h_{0}^{-}\right) \cap L_{2}\left(h_{0}^{+}\right) \text {for } \\
& h_{0}=\log (1+\sqrt{2}), \text { but } \\
& h_{\text {top }}\left(T \mid \operatorname{orb}\left(K_{2}(T)\right)=\log 2 \neq h_{0}\right.
\end{aligned}
$$

Figure 17. For this $T$ with $h_{\text {top }}(T)=h_{0}=\log (1+\sqrt{2})$ we have $h_{\text {top }}\left(T \mid \operatorname{orb}\left(K_{2}(T)\right)\right)=\log 2<h_{0}$, but lifting $Z_{2}$ by any amount pushes $h_{\text {top }}(T)=h_{\text {top }}\left(T \mid \operatorname{orb}\left(K_{2}(T)\right)\right)$ above $h_{0}$.

If the period of $K_{i}(T)$ is $m$, then since the intervals $K_{i}(T), \ldots, T^{m-1}\left(K_{i}(T)\right)$ have disjoint interiors, the first return map of $\operatorname{orb}\left(K_{i}(T)\right)$ has at most $2^{b}$ branches. It follows that if $h_{\text {top }}(T)=h_{0}$ then

$$
\begin{equation*}
0<h_{0}=h_{t o p}(T \mid \operatorname{orb}(K)) \leq(b \log 2) / m \tag{17}
\end{equation*}
$$

which gives the upper bound $m \leq(b \log 2) / h_{0}$.
Define

$$
\Phi_{j}(T)=\operatorname{dist}\left(T, L_{j}\left(h_{0}^{-}\right) \cap L_{j}\left(h_{0}^{+}\right)\right)
$$

where dist is as in (9) on page 46. Define the following deformation of maps $T \in \mathcal{S}^{b}$ :

$$
\begin{equation*}
\hat{\Gamma}_{t}\left(\zeta_{1}, \ldots, \zeta_{b}\right)=\left(\min \left(\zeta_{1}+\Phi_{1}(T) t, e\right), \ldots, \min \left(\zeta_{b}+\Phi_{b}(T) t, e\right)\right) \tag{18}
\end{equation*}
$$

Let us say that $T \in \mathcal{S}_{*, j}^{b}$ if whenever $\mathcal{J}=\left[Z_{j-1}, Z_{j}\right]$ or $\mathcal{J}=\left[Z_{j}, Z_{j+1}\right]$ (assuming $1 \leq j-1$ and $j+1 \leq b$ respectively) is eventually mapped into a plateau, then $\mathcal{J}$ is contained in the closure of a component of the basin of $T$.

Lemma 7.10. Let $\hat{\Gamma}_{t}$ be as above and let $T \in \mathcal{S}^{b}$ be such that $h_{\text {top }}(T) \leq h_{0}$. Then the following hold:
(1) The deformation $(T, t) \mapsto \hat{\Gamma}_{t}(T)$ is continuous in $T$ and $t$, and $t \mapsto h_{\text {top }}\left(\hat{\Gamma}_{t}(T)\right)$ is non-decreasing.
(2) $h_{\text {top }}\left(\hat{\Gamma}_{t}(T)\right) \leq h_{0}$ for all $0 \leq t \leq 1$.
(3) Assume that $T$ satisfies
(a) the $\beta$-property (11) on page 46 and
(b) for any $j=1, \ldots, b$,

$$
T \in L_{j}\left(h_{0}^{-}\right) \cap L_{j}\left(h_{0}^{+}\right) \text {implies } T \in \mathcal{S}_{*, j}^{b} .
$$

Then $\hat{\Gamma}_{t}(T) \in \mathcal{S}_{*}^{b}$ for each $t>0$.

Proof. The continuity and monotonicity of statement (1) are obvious.
For statement (2), take $t \in[0,1)$ and $T \in \mathcal{S}^{b}$ with $h_{\text {top }}(T) \leq h_{0}$. Let $j_{1}$ be such that $m_{1}:=\Phi_{j_{1}}(T)$ is maximal among $\left\{\Phi_{1}(T), \ldots, \Phi_{b}(T)\right\}$, and let $M_{1}$ be the open $b$-dimensional $m_{1}$-cube centered at $T$, parallel to the coordinate hyperplanes and of side length $2 m_{1}$. Then $M_{1}$ is disjoint from $L_{j_{1}}\left(h_{0}^{-}\right) \cap L_{j_{1}}\left(h_{0}^{+}\right)$, and in particular disjoint from $L_{j_{1}}\left(h_{0}^{+}\right)$. Therefore $h_{\text {top }}\left(T^{\prime} \mid \operatorname{orb}\left(K_{j_{1}}\left(T^{\prime}\right)\right)\right) \leq h_{0}$ for any $T^{\prime} \in M_{1}$ and in particular for $\hat{\Gamma}_{t}(T)$.

Now let $j_{2}$ be such that $m_{2}:=\Phi_{j_{2}}(T)$ is second largest among $\left\{\Phi_{1}(T), \ldots, \Phi_{b}(T)\right\}$. The corresponding $m_{2}$-cuboid $M_{2}$ is the set of $T^{\prime}$ with parameters $\left\{\zeta_{1}, \ldots, \zeta_{b}\right\}$ such that $\left|\zeta_{j}-\zeta_{j}(T)\right|<m_{2}$ for all $j \neq j_{1}$ and $\left|\zeta_{j_{1}}-\zeta_{j_{1}}(T)\right|<m_{1}$. (This is the Cartesian product of a $b$-1-dimensional cube and an arc of length $2 m_{1}$ in the $\zeta_{j_{1}}$-direction.)

Claim 1: $M_{2}$ is disjoint from $L_{j_{2}}\left(h_{0}^{+}\right)$. To prove this claim, consider $T^{\prime} \in M_{2}$. Let $M_{2}^{\prime \prime}$ be the $m_{2}$-cube centered at $T$, and choose $T^{\prime \prime} \in M_{2}^{\prime \prime}$ be so that $T^{\prime}$ and $T^{\prime \prime}$ agree except at $Z_{j_{1}}$. Let $T_{t}, t \in[0,1]$ be the one-parameter family of maps connecting $T^{\prime \prime}$ to $T^{\prime}$ corresponding to maps for which the parameter $\zeta_{t, j_{1}}$ varies linearly and so that $\zeta_{t, i}=\zeta_{0, i}$ for all $t \in[0,1]$ and $i \neq j_{1}$. Denote $K_{j_{2}, t}^{*}:=\operatorname{orb}\left(K_{j_{2}}\left(T_{t}\right)\right)$ and $K_{j_{1}, t}^{*}:=\operatorname{orb}\left(K_{j_{1}}\left(T_{t}\right)\right)$. By definition of $m_{1}$ and $m_{2}$, we have $h_{\text {top }}\left(T_{t} \mid K_{j_{1}, t}^{*}\right) \leq h_{0}$ for each $t \in[0,1]$ and $h_{\text {top }}\left(T_{0} \mid K_{j_{2}, 0}^{*}\right) \leq h_{0}$.

Next let

$$
X_{1}=\left\{t \in[0,1] ; K_{j_{2}, t}^{*} \cong K_{j_{1}, t}^{*}\right\}, X_{2}=\left\{t \in[0,1] ; K_{j_{2}, t}^{*} \supsetneqq K_{j_{1}, t}^{*}\right\}
$$

and

$$
X_{3}=\left\{t \in[0,1] ; K_{j_{2}, t}^{*} \cap K_{j_{1}, t}^{*}=\emptyset\right\} .
$$

From the above properties we obtain $h_{\text {top }}\left(T_{t} \mid K_{j_{2}, t}^{*}\right) \leq h_{\text {top }}\left(T_{t} \mid K_{j_{1}, t}^{*}\right) \leq h_{0}$ for each $t \in X_{1}$. Consider a component $C$ of $X_{3}$ or of $X_{2}$. If there exists $t \in C$ so that $h_{\text {top }}\left(T_{t} \mid K_{j_{2}, t}^{*}\right) \leq h_{0}$ then $h_{\text {top }}\left(T_{s} \mid K_{j_{2}, s}^{*}\right) \leq h_{0}$ for each $s \in C$; here we use that $h_{\text {top }}\left(T_{t} \mid K_{j_{1}, t}^{*}\right) \leq h_{0}$ for each $t \in[0,1]$ if $C$ is a component of $X_{2}$.

Claim 1 follows if there exists a component of $X_{1}, X_{2}$ or $X_{3}$ containing both 0 and 1. So from now on we assume that this is not the case, and we can assume that $t^{\prime}$ in the next claim is not 0 or 1 .

Subclaim 1: for each boundary point $t^{\prime}$ of $C$ there exists $t_{n} \in X_{1}$ so that $t_{n} \rightarrow t^{\prime}$. In fact, $K_{j_{1}, t_{n}} \supset K_{j_{1}, t_{n}}=K_{j_{1}, t^{\prime}}$.

To prove this, let $\tilde{K}_{j_{1}, t^{\prime}}$ be the maximal renormalization interval containing $K_{j_{1}, t^{\prime}}$ which is either contained in $K_{j_{2}, t^{\prime}}$ (if $t^{\prime} \in X_{2}$ ) or which is disjoint form $K_{j_{2}, t^{\prime}}$ (if $t^{\prime} \in X_{3}$ ). Observe that the maximality of $\tilde{K}_{j_{1}, t^{\prime}}$ implies that there are backwards iterates of $K_{j_{2}, t^{\prime}}$ accumulating to the boundary points of $\tilde{K}_{j_{1}, t^{\prime}}$. That $t^{\prime}$ is a boundary point of $C$ means that some iterate of $Z_{j_{1}, t^{\prime}}$ under $T_{t^{\prime}}$ is mapped into the boundary of a component of $\tilde{K}_{j_{1}, t^{\prime}}$. Because of the above observation this implies that there exists a sequence of points $x_{n}$ converging to either boundary point $\tilde{K}_{j_{1}, t^{\prime}}$ such that the omega-limit set $\omega\left(x_{n}\right) \subset K_{j_{2}, t^{\prime}}^{*}$. One can choose $x_{n}$ even so that $K_{j_{2}, t^{\prime}}$ is the smallest renormalization interval whose orbit contains $\omega\left(x_{n}\right)$. It follows that there exists a sequence $t_{n} \rightarrow t^{\prime}$ with $x_{n} \in \operatorname{orb}\left(Z_{j_{1}, t_{n}}\right)$ so that $K_{j_{1}, t_{n}} \supseteqq K_{j_{2}, t_{n}}=K_{j_{2}, t^{\prime}}$. This completes the proof of Subclaim 1.

Proof of Claim 1 continued. Let us say that a renormalization $K$ of $T_{t}$ is created at parameter $t$, if $K$ is not a renormalization interval for $T_{s}$ for all $s \in(t-\delta, t)$ with $\delta>0$ small. Since $\zeta_{j_{1}}$ is the only parameter moving, this implies that

$$
\begin{equation*}
Z_{j_{1}, t}=K \text { has a one-sided attracting periodic boundary point. } \tag{20}
\end{equation*}
$$

Again, since $\zeta_{j_{1}}$ is the only parameter moving, $Z_{j_{1}, t}$ cannot be part of block of plateaus if $K$ is created.

First consider the case that $\zeta_{j_{1}, 1}>\zeta_{j_{1}, 0}$. If $t \in\left(X_{2} \cup X_{3}\right)$, then $Z_{j_{2}, t}$ is never mapped into $Z_{j_{1}, t}$. It follows that for each $s \in[t, 1]$ no iterate under $T_{s}$ of $Z_{j_{2}, s}=Z_{j_{2}, t}$ is mapped into $Z_{j_{1}, s}$. Hence $\omega\left(Z_{j_{2}, s}\right)$ remains the same for each $s \in\left[t^{\prime}, 1\right]$ where we can take $t^{*}:=\inf \left(X_{2} \cup X_{3}\right)$. If $K$ is a renormalization interval for $T_{s^{\prime}}$ which is created at parameter $s^{\prime}$ and $K$ intersects $\omega\left(Z_{j_{2}, s^{\prime}}\right)$, then by remark (20) it follows that $\omega\left(Z_{j_{2}, s^{\prime}}\right)$ would have to intersect $Z_{j_{1}, s^{\prime}}$, which is impossible when $s \geq t^{\prime}$. Hence, if the renormalization interval $K_{j_{2}, s}$ changes at some $s \in\left[t^{*}, 1\right]$, then it is replaced by some larger renormalization interval $K_{j_{2}, s^{\prime}}$ for $T_{s^{\prime}}$ with $s^{\prime}>s$ close to $s$. Thus the set $K_{j_{2}, s}$ only gets larger as $s$ increases from $t^{*}$ to 1 . Let $C$ be a component as before and let $t^{\prime}$ be an endpoint of $C$ (with $t^{\prime} \neq 0,1$ ). By Subclaim 1 there exists $t_{n} \rightarrow t^{\prime}$ with $t_{n} \in X_{1}$ and so $K_{j_{1}, t_{n}} \supset K_{j_{1}, t_{n}}=K_{j_{1}, t^{\prime}}$. This, together with the fact that $h_{\text {top }}\left(T_{t_{n}} \mid K_{j_{1}, t_{n}}^{*}\right) \leq h_{0}$, implies that $h_{\text {top }}\left(T_{t^{\prime}} \mid K_{j_{2}, t^{\prime}}^{*}\right) \leq h_{0}$. Using the first part of the proof (above Subclaim 1), Claim 1 follows.

Now consider the case that $\zeta_{j_{1}, 1}<\zeta_{j_{1}, 0}$. In this case, $K_{j_{2}, s}$ can shrink as $s$ increases, namely when $Z_{j_{1}, s}$ is contained in $K_{j_{2}, s}^{*}$ and the image of $Z_{j_{1}, s}$ is mapped to a boundary point of $K_{j_{2}, s}^{*}$. Since in this case $K_{j_{2}, s}^{*}$ can only shrink as $s$ increases, and $Z_{j_{2}, s}$ only gets wider as $s$ increases, it follows that $h_{\text {top }}\left(T_{t^{\prime}} \mid K_{j_{2}, s}^{*}\right)$ is a decreasing
function of $s$ in this case. The set $K_{j_{2}, s}^{*}$ can also change if at some parameter $s^{\prime}$ it is destroyed (i.e., created in reverse direction), but remark (20) implies that some iterate of $Z_{j_{2}, s^{\prime}}$ is contained in $Z_{j_{1}, s^{\prime}}$, so $\omega\left(Z_{j_{2}, s^{\prime}}\right)=\omega\left(Z_{j_{1}, s^{\prime}}\right)$ and we are in the parameter set $X_{1}$ and in particular $h_{\text {top }}\left(T_{s^{\prime \prime}} \mid K_{j_{2}, s^{\prime \prime}}^{*}\right)=h_{\text {top }}\left(T_{s^{\prime \prime}} \mid K_{j_{1}, s^{\prime \prime}}^{*}\right) \leq h_{0}$ for all $s^{\prime \prime} \geq s^{\prime}$. Thus Claim 1 holds also if $\zeta_{j_{1}, 1}<\zeta_{j_{1}, 0}$.

Continuing inductively, we see that if $T^{\prime} \in \cap_{k=1}^{b} M_{k}$, then $h_{\text {top }}\left(T^{\prime} \mid \operatorname{orb}\left(K_{j}\left(T^{\prime}\right)\right)\right) \leq h_{0}$ for each $j$, and this holds in particular for $\hat{\Gamma}_{t}(T)$ for each $t \in[0,1]$. If $K_{j}\left(\hat{\Gamma}_{t}(T)\right)=$ $[-e, e]$ for some $j$ (i.e., if $\hat{\Gamma}_{t}(T)$ has no renormalization interval), then this proves statement (2).

If, however, every plateau belongs to some renormalization cycle and the entropy of $\hat{\Gamma}_{t}(T)$ is carried by the Cantor set of points that never enter these renormalization cycles, then we argue as follows. Write $T_{t}=\hat{\Gamma}_{t}(T)$ and assume by contradiction that $h_{\text {top }}\left(T_{t_{1}}\right)>h_{0}$ for some $t_{1} \in[0,1)$. Take $t_{0} \in\left[0, t_{1}\right)$ maximal such that $h_{\text {top }}\left(T_{t_{0}}\right) \leq h_{0}$. Since $T_{t} \in \cap_{k=1}^{b} M_{k}$ for each $t \in[0,1)$, the first part of the proof gives

$$
\begin{equation*}
h_{\text {top }}\left(T_{t} \mid \operatorname{orb}\left(K_{j}\left(T_{t}\right)\right) \leq h_{0} \text { for each } t \in[0,1] \text { and each } j=1, \ldots, b\right. \tag{21}
\end{equation*}
$$

Now we need the following
Claim 2: There exists $j \in\{1, \ldots, b\}$ and $t^{\prime} \in\left(t_{0}, t_{1}\right)$ so that $h_{\text {top }}\left(T_{t^{\prime}} \mid \operatorname{orb}\left(K_{j}\left(T_{t^{\prime}}\right)\right)\right)>$ $h_{0}$ for the minimal interval of renormalization $K_{j}\left(T_{t^{\prime}}\right)$ containing $\omega\left(Z_{j, T_{t^{\prime}}}\right)$.

Proof of Claim 2: For each $t \in\left(t_{0}, t_{1}\right)$ consider the semi-conjugacy of $T_{t}$ with the $\operatorname{map} F_{t}$ with slope $\pm \exp \left(h_{\text {top }}\left(T_{t}\right)\right)$ as above Lemma 7.3. Since $h_{\text {top }}\left(F_{t}\right)=h_{\text {top }}\left(T_{t}\right)$ depends continuously on $t$, and is not constant on $\left(t_{0}, t_{1}\right), h_{\text {top }}\left(F_{t}\right)$ assumes uncountably many values as $t$ moves through $\left(t_{0}, t_{1}\right)$. But there are only countably many slopes for which all turning points of $F_{t}$ are periodic. Therefore there exists $t^{\prime} \in\left(t_{0}, t_{1}\right)$ so that at least one of the turning points of $F_{t^{\prime}}$, say $c_{j}$, is non-periodic. Let $X_{j, t^{\prime}}$ be the smallest renormalization interval of $F_{t^{\prime}}$ containing this turning point. Since $F_{t^{\prime}}$ has constant slope, $h_{\text {top }}\left(F_{t^{\prime}} \mid \operatorname{orb}\left(X_{j, t^{\prime}}\right)\right)=h_{\text {top }}\left(F_{t^{\prime}}\right)>h_{0}$. Since the $j$-th critical point of $F_{t^{\prime}}$ is not periodic, the smallest renormalization $K_{j}\left(T_{t^{\prime}}\right)$ of $T_{t^{\prime}}$ containing $\omega\left(Z_{j, t^{\prime}}\right)$ is mapped by the semi-conjugacy onto $X_{j, t^{\prime}}$. It follows that $h_{\text {top }}\left(T_{t^{\prime}} \mid \operatorname{orb}\left(K_{j}\left(T_{t^{\prime}}\right)\right)\right)>h_{0}$, completing the proof of the claim.

Obviously the claim contradicts (21), so we can conclude that $h_{\text {top }}\left(T_{1}\right) \leq h_{0}$. Statement (2) follows.

Finally we turn statement to statement (3). For this we need to show that $\hat{\Gamma}_{t}(T) \in$ $\mathcal{S}_{*, j}^{b}$ for each $j$ and each $t>0$. If $T \in L_{j}\left(h_{0}^{-}\right) \cap L_{j}\left(h_{0}^{+}\right)$then this holds by assumption (19). So assume that $T \notin L_{j}\left(h_{0}^{-}\right) \cap L_{j}\left(h_{0}^{+}\right)$. Then the $\beta$-property (11) implies that whenever $T^{n}\left(\left[Z_{j}, Z_{j+1}\right]\right)$ is contained in a plateau $Z_{k}$ then $T^{n}\left(\left[Z_{j}, Z_{j+1}\right]\right) \cap \partial Z_{k} \neq \emptyset$. Since the plateau $Z_{j, t}$ shrinks at $t$ increases (as $\left.T \notin L_{j}\left(h_{0}^{-}\right) \cap L_{j}\left(h_{0}^{+}\right)\right)$, this means that for $t>0,\left(\hat{\Gamma}_{t}(T)\right)^{n}$ maps $\left[Z_{j}, Z_{j+1}\right]$ at least partially outside $Z_{k}$. So again $\hat{\Gamma}_{t}(T) \in \mathcal{S}_{*, j}^{b}$. This proves statement (3).
7.9. Decreasing the entropy more carefully: $\boldsymbol{\Delta}_{\boldsymbol{t}}$. Take $T \in \mathcal{S}_{*}^{b}$ with $h_{\text {top }}(T)=$ $h_{0}>0$. Even though $t \mapsto h_{\text {top }}\left(\hat{\delta}_{t}(T)\right)$ is non-increasing, it is possible that for fixed $t>0, h_{\text {top }}\left(\gamma_{s} \circ \hat{\delta}_{t}(T)\right)>h_{0}$ for all $s>0$. The reason is that (in the notation of Section 7.3) $\widehat{\operatorname{sgn}}\left(Z_{i}\right)$ can change from 1 to -1 (or vice versa) during the deformation. To explain what can happen, let us discuss two examples.

Example 7.11. Consider the map from Figure 18 on page 60. Although $T \in \mathcal{S}_{*}^{b}$, the map $T^{\prime}=\hat{\delta}_{t}(T)$ has a wandering pair that does not map into a periodic basin (so $T^{\prime}$ is no longer in $\mathcal{S}_{*}^{b}$ ). There is a periodic interval $\hat{K}$ (here of period 1) and $T^{\prime}$ maps the convex hull $\mathcal{J}=\left[Z_{1}^{\prime}, Z_{3}^{\prime}\right]$ into $\partial \hat{K}$. (Note that, in this example, $\hat{\delta}_{t}(T)$ first decreases $\zeta_{2}$ and then, after the plateaus $Z_{1}$ and $Z_{3}$ touch $Z_{2}$, increases it again. Even though $\gamma_{t}$ initially is 'the inverse' of the deformation $\hat{\delta}_{t}$, the map $T^{\prime}=\hat{\delta}_{t}(T)$ will have some touching plateaus.) Because the entropy within the renormalization interval $\hat{K}$ is $\leq h_{0}$, the movements of plateaus $Z_{1}, Z_{2}, Z_{3}$ under $\hat{\delta}_{t}$ have no effect on the global entropy. Therefore $h_{\text {top }}\left(T^{\prime}\right)=h_{\text {top }}(T)$, whereas $h_{\text {top }}\left(\gamma_{s}\left(T^{\prime}\right)\right)>h_{\text {top }}(T)$ for any $s>0$, because $\gamma_{s}$ decouples the plateaus $Z_{1}, Z_{2}, Z_{3}$ again and $\gamma_{s}\left(T^{\prime}\right)\left(\left[Z_{1}, Z_{3}\right]\right)$ is a closed neighbourhood of the left endpoint of (and therefore not entirely contained in) $\hat{K}$. The effect is that within $\hat{\Gamma}_{s}$, the deformation $\gamma_{s}$ will not be applied at all, and hence it will not be able to remove the wandering pair created by $\hat{\delta}_{t}$.

Example 7.12. In fact, a similar problem can occur even when we consider the family $\beta_{1} \circ \delta_{t}(T)$. It is possible that for some $t \in(0,1)$ a pair of plateaus $\left[Z_{1, t}, Z_{2, t}\right]$ of $\delta_{t}(T)$ is mapped by $\delta_{t}(T)$ into a plateau $Z_{3, t}$, which in turn is mapped into a renormalization interval $\hat{K}_{t}$. As $t$ increases, the parameters $\zeta_{t}$ associated to $\beta_{1} \circ \delta_{t}(T)$ no longer decrease with $t$ and a similar situation as in Figure 18 can arise (but with $Z_{3, t}$ the image of $\left[Z_{1, t}, Z_{2, t}\right]$ and $Z_{3, t}$ mapped into the boundary of a periodic plateau).

To overcome the issues caused by such examples, we introduce the deformation $\Delta_{t}$.
Construction of $\Delta_{t}$ : Let us define a modification $\Delta_{t}$ of the deformation $\hat{\delta}_{t}$, which allows some of the plateaus (namely those within renormalization intervals of 'low entropy') to move before others.

Fix $h_{0} \in(0, \log (b+1))$.
Definition 7.13. Given a periodic interval $K$ of period $m$, we can find intervals $K_{i} \supset T^{i}(K)$ such that $T\left(\partial K_{i}\right) \subset \partial K_{i+1 \bmod m}$ for all $0 \leq i<m$. For the minimal choice of such intervals $K_{i}$, we write cycle $(K)=\cup_{i=0}^{m-1} K_{i}$. Given a renormalization interval $K$ of $T$, we define

$$
\begin{aligned}
\Omega(K)=\{z \in \operatorname{cycle}(K) ; \quad & z, T(z), T^{2}(z), \ldots \notin \hat{K} \text { for any } \\
& \text { renormalization interval } \hat{K} \varsubsetneqq K\} .
\end{aligned}
$$

We say that $K$ has entropy $h_{0}($ for $T)$ if $h_{\text {top }}(T \mid \Omega(K))=h_{0}$.


Figure 18. The maps $T$ and $T^{\prime}:=\hat{\delta}_{t}(T)$ (in dotted lines). For $t>0$ small, $t \mapsto \hat{\delta}_{t}(T)$ increases the height of the plateau $Z_{2}$ and decreases those of $Z_{1}, Z_{3}$. Once they merge, this deformation decreases the height of all of them together. In this example, the plateaus $Z_{1}, Z_{2}$ and $Z_{3}$ are mapped into $\partial \hat{K}$, i.e., $T^{\prime}\left(Z_{1}\right)=T^{\prime}\left(Z_{2}\right)=T^{\prime}\left(Z_{3}\right) \in \partial \hat{K}$. The map is constructed so that $T^{\prime} \mid \hat{K}$ is unimodal with entropy $h_{\text {top }}\left(T^{\prime} \mid \hat{K}\right)<h_{\text {top }}(T)$. In this case, $h_{\text {top }}\left(\gamma_{s}\left(T^{\prime}\right)\right)>h_{\text {top }}(T)$ for any $s>0$, because points near $Z_{2}$ will then be mapped outside (i.e., to the left of) $\hat{K}$. Therefore $\hat{\delta}_{t}(T) \in L_{2}\left(h_{0}^{-}\right) \cap L_{2}\left(h_{0}^{+}\right)$ where $h_{0}=h_{\text {top }}(T)$ and $\hat{\delta}_{t}(T) \notin \mathcal{S}_{*}^{b}$ unless the left boundary point of $\hat{K}$ is in the boundary of a component of the basin of a periodic attractor.

It is possible that $h_{\text {top }}(T \mid \Omega(K))<h_{0}$ but that there exists a renormalization interval $\hat{K} \varsubsetneqq K$ so that $h_{\text {top }}(T \mid \Omega(\hat{K}))=h_{\text {top }}(T \mid \operatorname{cycle}(\hat{K}))=h_{0}$.

Definition 7.14. Given an interval L, we say that a plateau $Z_{j}$ arches over $\partial L$ if there exists $m$ so that $T^{m}\left(Z_{j}\right) \subset \partial L$ and if $k<m$ is maximal such that $T^{k}\left(Z_{j}\right)$ is contained in a plateau $Z$, then $T^{-(m-k)}(L)$ strictly contains the maximal block of touching plateaus containing $Z$, see Figure 19 .


Figure 19. Three basic possibilities how a block of plateaus can map onto $\partial L$. In the situations described in the left two panels, the interval $\mathcal{J}$ arches over $\partial L$.

Definition 7.15. We say that $T \in \xi_{i}$ if there exists renormalization interval $K$ of entropy $h_{0}$ and a convex hull $\mathcal{J} \subset K$ of plateaus which is non-trivial, i.e., $T(\mathcal{J})$ is not a singleton and so that
(1) there exists a maximal renormalization interval $\hat{K} \varsubsetneqq K$ so that $T^{m}(\mathcal{J}) \subset \partial \hat{K}$ for some $m \geq 2$ (in particular $T^{m}(\mathcal{J})$ is a singleton);
(2) the first return map to $\hat{K}$ is non-monotone;
(3) $T^{j}(\mathcal{J})$ intersects $Z_{i}$ for some $0 \leq j<m-1$;
(4) $Z_{i} \cap \operatorname{cycle}(\hat{K})=\emptyset$.

It will be useful to stratify the space $\mathcal{S}^{b}$, dividing the boundary of this space into subspaces according to which plateaus touch. To do this, let $\Xi$ be a partition $\{1, \ldots, b\}$. Then we define the stratum $\mathcal{S}^{b}(\Xi) \subset \mathcal{S}^{b}$ as follows: $T \in \mathcal{S}^{b}(\Xi)$ if and only if for each $i \in\{1, \ldots, b-1\}$ the plateaus $Z_{i}$ and $Z_{i+1}$ touch whenever $i, i+1$ are in the same partition element of $\Xi$. For example $\mathcal{S}^{b}(\{1\},\{2\}, \ldots,\{b\})$ consists of all maps for which none of the plateaus touch, and $\mathcal{S}^{b}(\{1,2\},\{3\}, \ldots,\{b\})$ consists of the space of maps for which the first two plateaus touch (and no other two plateaus do). Note that when $\Xi_{1}, \Xi_{2}$ are two distinct partitions, then $\mathcal{S}^{b}\left(\Xi_{1}\right)$ and $\mathcal{S}^{b}\left(\Xi_{2}\right)$ are disjoint. Clearly $\mathcal{S}^{b}$ is the disjoint union of $\mathcal{S}^{b}(\Xi)$ where the union runs over all partitions $\Xi$ of $\{1, \ldots, b\}$, and each $T \in \mathcal{S}^{b}$ is associated to a partition $\Xi(T)$ of $\{1, \ldots, b\}$.

Note that in Definition 7.15, a non-trivial convex hull $\mathcal{J}$ always contains two plateaus $Z_{i}, Z_{j}$ where $i, j$ are in distinct subsets from the partition $\Xi(T)$. Also note that property (2) Definition 7.15 in particular implies that the first return map to $\hat{K}$ is non-constant.

For each partition $\Xi$ of $\{1, \ldots, b\}$, choose a $C^{\infty}$ function $\rho_{i}^{\Xi}: \mathcal{S}^{b}(\Xi) \rightarrow[0,1]$ which is zero on $\overline{\xi_{i} \cap \mathcal{S}^{b}(\Xi)}$ and positive elsewhere. Next define $\rho_{i}: \mathcal{S}^{b} \rightarrow[0,1]$ by $\rho_{i}(T)=$ $\rho_{i}^{\Xi}(T)$ whenever $T \in \mathcal{S}^{b}(\Xi)$ for some partition $\Xi$ of $\{1, \ldots, b\}$.
Define the modification $\Delta_{t}$ of $\hat{\delta}_{t}$ as the flow of the differential equation

$$
\begin{align*}
\frac{d \zeta_{i, t}}{d t} & = \begin{cases}-\rho_{i}\left(T_{t}\right) \cdot \widehat{\operatorname{sgn}}\left(Z_{i, t}\right) & \text { when } \zeta_{i, t} \in(-e, e), \\
0 & \text { otherwise }\end{cases}  \tag{22}\\
\left.\zeta_{i, t}\right|_{t=0} & =\zeta_{i}(T)
\end{align*}
$$

Here $\widehat{\operatorname{sgn}}$ is defined as in (15) on page 48 and $T_{t}$ is the map corresponding to $\zeta_{i, t}$ and so $T_{t}=\Delta_{t}(T)$.
Proposition 7.16. For each $T \in \mathcal{S}^{b}, \Delta_{t}(T)$ exists for all $t>0$ and

$$
\mathcal{S}_{*}^{b} \times \mathbb{R}^{+} \ni(T, s) \mapsto \Delta_{s}(T)
$$

is continuous. Moreover:
(a) For each $T \in \mathcal{S}_{*}^{b}$ there exists $t_{\Delta}(T) \in(0, \infty)$ so that the map $\Delta_{t_{\Delta}(T)}(T)$ is trivial (i.e., monotone).
(b) $\mathcal{S}_{*}^{b} \ni T \mapsto t_{\Delta}(T)$ is continuous.
(c) For each $t \in[0, t(T)], \Delta_{t}(T) \in L_{j}\left(h_{0}^{-}\right) \cap L_{j}\left(h_{0}^{+}\right)$implies that if $Z_{j}$ is part of a wandering pair then $Z_{j}$ is contained in the closure of a component of the basin of a periodic attractor of $\Delta_{t}(T)$. In particular, $\Delta_{t}(T) \in L_{j}\left(h_{0}^{-}\right) \cap L_{j}\left(h_{0}^{+}\right)$ implies that $\Delta_{t}(T) \in \mathcal{S}_{*, j}^{b}$.

Proof. Note that the right hand side of differential equation (22) is smooth on each stratum $\mathcal{S}^{b}(\Xi)$. The discontinuities of the right hand side occur when two plateaus start to touch, and the nature of the equation is that once they touch they remain touching. From this and the existence and uniqueness theorem of differential equations (applied to each stratum separately), it follows that the flow of the differential equation is well-defined. However, continuity of $\mathcal{S}_{*}^{b} \times \mathbb{R}^{+} \ni(T, t) \rightarrow \Delta_{t}(T)$ still needs to be proved.

For simplicity write $T_{s}=\Delta_{s}(T)$ and let $\Xi_{s}$ be the partition associated to $T_{s}$. As $s$ increases, each plateau only widens under the flow $T_{s}$, until it touches another plateau in which case these plateaus widen jointly unless the corresponding block has an even number of touching plateaus (or touches $\pm e$ ). In this sense each coordinate of $T_{s}$ depends monotonically on $s>0$, and therefore the limit $\tilde{T}:=\lim _{s \rightarrow \infty} T_{s} \in \mathcal{S}^{b}$ exist. Also $s \mapsto h_{\text {top }}\left(T_{s}\right)$ is non-increasing in $s$.

Now take $T \in \mathcal{S}_{*}^{b}$ with $h_{\text {top }}(T)=h_{0}$, and assume by contradiction that property (a) does not hold.

Step 1. Assume that $K$ is a renormalization interval of entropy $h_{0}$ for $T_{t}$ so that there exists no renormalization interval $K^{\prime} \supsetneqq K$ of the same period. We claim that $K$ is a renormalization interval for $T_{s}$ for each $s \in[0, t]$. Indeed, otherwise there would exist $0 \leq s_{0}<t$ so that one of the plateaus $Z_{j, s} \subset K$ of $T_{s_{0}}$ arches over $\partial K$ under the first return map of $T_{s_{0}}$ to $K$. If $Z_{j, s}$ is independent of $s \in\left[0, s_{0}\right]$, then $K$ is still a renormalization interval for $T_{s}$ and there is nothing to show. If $Z_{j, s}$ does depend on $s$, then there exists $s^{\prime} \in\left[0, s_{0}\right)$ so that the first return map to $K$ under $T_{s^{\prime}}$ has another branch, see the left panel in Figure 20. Here we use that if for $T_{s_{0}}$ there exists a plateau outside $K$, but which touches an endpoint of $K$ (as is shown in the right panel of Figure 20), then there would have existed a renormalization interval $K^{\prime} \supsetneqq K$ with the same return time, contradicting the choice of $K$. Since the first return map to $K$ under $T_{s^{\prime}}$ has another branch, it follows by Lemma 7.3 that $h_{0} \geq h_{\text {top }}\left(T_{s^{\prime}} \mid \Omega(K)\right)>h_{\text {top }}\left(T_{t} \mid \Omega(K)\right)=h_{0}$, contradicting the assumption made in this step.

Step 2. Take $t>0$. We claim that there exists $\epsilon>0$ so that for each $s \in[t-\epsilon, t+\epsilon]$ there exists $i \in\{1, \ldots, b\}$ with $T_{s} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$ and $Z_{i} \subset \operatorname{cycle}(K)$.

We may assume that $K$ is a renormalization interval of entropy $h_{0}$ for $T_{t}$ because otherwise this claim holds trivially. To prove this claim, we first note that by Step


Figure 20. The situation considered in Step 1 of the proof of Proposition 7.16. The first return map to $K$ for $\Delta_{s}(T)$ and (schematically) for $\Delta_{s^{\prime}}(T)$ for $s^{\prime}<t$ is drawn in the left panel. The situation as in the right cannot occur, see the proof in Step 1.
$1, K$ is a renormalization interval for $T_{s}$ for each $s \in[0, t]$. Let us consider all the possible situations.
Case A. $T_{t}$ does not have a renormalization interval $\hat{K} \varsubsetneqq K$. We claim that in this case $h_{\text {top }}\left(T_{t} \mid \Omega(K)\right) \leq h_{\text {top }}\left(T_{s} \mid \Omega(K)\right)<h_{\text {top }}\left(T_{0} \mid \Omega(K)\right) \leq h_{0}$. By the definition of $\xi_{i}$ this implies that $T_{s} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$ for each $0 \leq s \leq t$ and each $i$ so that $Z_{i, t} \subset K$; therefore, in this case Step 2 follows from the claim. To prove this, first note that $T_{s}$ has no renormalization interval $\hat{K} \varsubsetneqq K$ for any $s \in[0, t]$. Indeed, if there exists $s \in[0, t]$ so that $T_{s}$ does have a renormalization interval $\hat{K} \varsubsetneqq K$, then by Lemma 7.3 we obtain $h_{\text {top }}\left(T_{t} \mid \operatorname{cycle}(K)\right)<h_{\text {top }}\left(T_{s} \mid \operatorname{cycle}(K)\right) \leq h_{0}$ for any $s<t$, which gives a contradiction. Next take $T_{s}$ with $s>0$ small, and let $T_{s}(n) \in \mathcal{S}^{b}\left(\Xi_{s}\right)$ be a sequence of maps so that $T_{s}(n) \rightarrow T_{s}$. Since $T_{t}$ has no renormalization intervals $\hat{K} \varsubsetneqq K$, the map $T_{t}$ and therefore $T_{0}=T$ has no periodic attractors in $K$ either (and since $T \in \mathcal{S}_{*}^{b}$, therefore no two plateaus of $T$ touch). Since $T \in \mathcal{S}_{*}^{b}$, the convex hull $\mathcal{J}_{n}$ of two adjacent plateaus $Z_{i, n}, Z_{i+1, n}$ for $T_{s}(n)$ for $n$ large and $s>0$ small, does not form a wandering pair, see Lemma 4.16 (b). It follows that $T_{s} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}}$ for each $i$ and each $s \geq 0$ small, and therefore for each $i$ so that $Z_{i} \subset K$ we have that $\rho_{i}\left(T_{s}\right) \cdot \widehat{\operatorname{sgn}}\left(Z_{i, s}\right) \neq 0$ for each $s \geq 0$ small. But this implies by Lemma 7.3 that $h_{\text {top }}\left(T_{t} \mid \Omega(K)\right) \leq h_{\text {top }}\left(T_{s} \mid \Omega(K)\right)<h_{\text {top }}\left(T_{0} \mid \Omega(K)\right)$ as claimed.

Case B. $T_{t}$ has a maximal renormalization interval $\hat{K} \varsubsetneqq K$ on which the map $T_{t}$ is non-monotone. We claim that $T_{s} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$ for each $s \in[0, t]$ and for each $i$ so that $Z_{i} \subset \hat{K}$. To see this, first note that we may assume that each maximal renormalization interval $\hat{K} \varsubsetneqq K$ is a renormalization interval for $T_{s}$ for all $s \in[0, t]$. Indeed, otherwise by Lemma 7.3, $h_{\text {top }}\left(T_{t} \mid \Omega(K)\right)<h_{\text {top }}\left(T_{0} \mid \Omega(K)\right) \leq h_{0}$, and then the claim follows immediately. To prove the claim, we consider the following cases:
(i) Each plateau $Z_{i, t}$ in $\hat{K}$ is mapped by $T_{t}$ into the interior of cycle $(\hat{K})$. Then a nearby map $\tilde{T}$ also has $\tilde{Z} \subset$ cycle $(\hat{K})$ and so part (4) of Definition 7.15 fails, and $\tilde{T} \notin \xi_{i}$. It follows $T_{s} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$ for each $i$ associated to a plateau $Z_{i, s} \subset \operatorname{cycle}(\hat{K})$ and for each $s$ near $t$.
(ii) One of the plateaus $Z_{i, t}$ in $\hat{K}$ arches by $T_{t}$ over the boundary of a component of cycle $(\hat{K})$. Then consider the following two subcases:
(iia) There exists $\epsilon>0$ so that the cardinality of the block of plateaus touching $Z_{i, s}$ does not change for $s \in[t-\epsilon, t]$. Next take a convex hull $\mathcal{J}_{s}$ connecting $Z_{i, s}$ and another plateau (outside the block of plateaus touching $Z_{i, s}$ ). (By Definition 7.15 we need to take $\mathcal{J}_{s}$ non-trivial, so $\mathcal{J}_{s}$ contains two plateaus $Z_{i, s}, Z_{j, s}$ from distinct subsets of the partition $\Xi_{s}$.) Because $Z_{i, s}$ arches, either $\mathcal{J}_{s}$ is contained in $\hat{K}$ or the block of plateaus touching $Z_{i, s}$ (and that of $Z_{i, t}$ ) contains an even number of plateaus. In the former case for $s \in[t-\epsilon, t)$ the width of this block is smaller, and in the latter case, $\widehat{\operatorname{sgn}}\left(Z_{i, s}\right)=0$ for each $s \in[t-\epsilon, t]$. It follows that in both cases $T_{s}\left(\mathcal{J}_{s}\right)$ contains a repelling (pre-)periodic boundary point of $\hat{K}$ for each $s \in[t-\epsilon, t]$. The same holds for each $\tilde{T} \in \Xi\left(T_{s}\right)$ near $T_{s}$ and therefore $\tilde{T} \notin \xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)$ for each $s \in[t-\epsilon, t]$. It follows that $T_{s} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$ for each $s \in[t-\epsilon, t]$. Provided $\epsilon>0$ is small, for each $s \in[t, t+\epsilon]$ the cardinality of the block of plateaus touching $Z_{i, s}$ does not change either, and for the same reason $T_{s} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$ for each $s \in[t, t+\epsilon]$.
(iib) There exists no such $\epsilon>0$. Then there exists a plateau $Z_{j, s} \subset \hat{K}$ (so that $Z_{j, t}$ is in the block of plateaus touching $Z_{i, t}$ ) which creates an extra branch for the return map of $T_{s}$ to $K$ (compared to the return map of $T_{t}$ ). Hence, by Lemma 7.3. $h_{\text {top }}\left(T_{t}, \Omega(K)\right)<h_{\text {top }}\left(T_{s}, \Omega(K)\right) \leq h_{0}$, which gives a contradiction.
(iii) $T_{t}$ is constant on one of the components of cycle $(\hat{K})$ and for simplicity assume that $\hat{K}$ is this component. Let $Z_{i, t}$ be contained in $\hat{K}$.
(iiia) There exists $\epsilon>0$ so that the cardinality of the block of plateaus touching $Z_{i, s}$ (in $\hat{K}$ ) does not change for $s \in[t-\epsilon, t]$. In this case the number of plateaus in this block is odd and so this block of plateaus is mapped into the interior of $\hat{K}$ for $s \in[t-\epsilon, t)$. By case (i) it follows that for $s \in[t-\epsilon, t)$ one has that $T_{s} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$. To prove this conclusion also for $s=t$, consider a sequence of maps $T_{t}(n) \rightarrow T_{t}$ where $T_{t}(n) \in \mathcal{S}^{b}\left(\Xi_{t}\right)$. By maximality of $\hat{K} \varsubsetneqq K$ it follows that each boundary point of $\hat{K}$ is a repelling periodic or pre-periodic point which is not the common boundary point of two basins. Hence there exists a sequence of repelling periodic points accumulating on $\partial \hat{K}$ (from outside $\hat{K}$ ). It follows that any non-trivial convex hull $\mathcal{J}(n)$ connecting $Z_{i, t}$ with another plateau (outside this component of $\hat{K}$ ) will contain a repelling periodic point of $T_{t}(n)$ and therefore iterates of $\mathcal{J}(n)$ are not singletons. It follows that $T_{t}(n) \notin \xi_{i}$ and therefore $T_{t} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{t}\right)}$. Moreover, for $s \in(t, t+\epsilon]$ one has $h_{\text {top }}\left(T_{s} \mid \operatorname{orb}(K)\right)<h_{\text {top }}\left(T_{t} \mid \operatorname{orb}(K)\right)=h_{0}$, where we choose $\epsilon>0$ so that $K$ remains periodic for $T_{s}$ for all $s \in[t, t+\epsilon]$. It again follows that $T_{s} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$ for $s \in(t, t+\epsilon]$ for each $i$ with $Z_{i} \subset K$.
(iiib) There exists no such $\epsilon>0$. Then there exists a plateau $Z_{j, s} \subset \hat{K}$ (so that $Z_{j, t}$ is in the block of plateaus touching $Z_{i, t}$ ) which creates an extra branch for the
return map of $T_{s}$ to $K$ (compared to the return map of $T_{t}$ ), and so Lemma 7.3 yields $h_{\text {top }}\left(T_{t}, \Omega(K)\right)<h_{\text {top }}\left(T_{s}, \Omega(K)\right) \leq h_{0}$, which again shows that $T_{s} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$ for $s \in[t, t+\epsilon]$.
Case C. $T_{t}$ is monotone on each component of $\operatorname{cycle}(\hat{K})$ for each maximal renormalization interval $\hat{K} \varsubsetneqq K$. We claim that in this case there exists $\epsilon>0$ so that $T_{s} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$ for $s \in[t, t+\epsilon]$ and for each $i$ so that $Z_{i} \subset K$. Indeed, in this setting there exists $\epsilon>0$ so that the cardinality of the block of plateaus touching $Z_{i, s}$ in $\operatorname{cycle}(\hat{K})$ is constant for $s \in[t, t+\epsilon]$. Moreover, for each sequence $T_{s}(n) \rightarrow T_{s}$ with $T_{s}(n) \in \mathcal{S}^{b}\left(\Xi_{s}\right)$, the map $T_{s}(n)$ is also monotone on each component of cycle $(\hat{K})$. It follows that for each plateau $Z_{i, s}$ which is not contained in a renormalization interval $\hat{K} \varsubsetneqq K$ one has $T_{s}(n) \notin \xi_{i}$ (because of Definition $7.15(2)$ ). Hence $T_{s} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$ for each such $i$ and $s \in[t, t+\epsilon]$. Note that for $s \in[t-\epsilon, t]$ either $T_{s}$ is also monotone on each component of $\operatorname{cycle}(\hat{K})$ (and therefore $T_{s}(n) \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$ for each $i$ as above) or one of the plateaus $Z_{i, s} \subset \hat{K}$ corresponds to a local extremum and then $T_{s}(n) \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$.

Step 3. We claim that if $h_{\text {top }}\left(T_{t} \mid \Omega(K)\right)=h_{0}$ for some $t>0$, then for each $Z_{i} \subset K$ and $s \in[0, t]$ some iterate of $Z_{i, s}$ is contained in a renormalization interval $\hat{K} \varsubsetneqq K$ (and this renormalization interval $\hat{K}$ does not depend on $s$ ). Indeed, if $Z_{i, t}$ is not eventually mapped into some renormalization interval $\hat{K} \subset K$ but this is the case for some $s \in[0, t]$ then Lemma 7.3 implies that $h_{0} \geq h_{\text {top }}\left(T_{s} \mid \Omega(K)\right)>h_{\text {top }}\left(T_{t} \mid \Omega(K)\right)$, and so we are done. Therefore it suffices to consider the case that no iterate of $Z_{i, s}$, $s \in[0, t]$, is contained in a renormalization interval $\hat{K}$. Let us show that this implies that $Z_{i, s}$ moves when $s$ small. Indeed, consider a sequence of maps $T_{s}(n) \rightarrow T_{s}$ as $n \rightarrow \infty$ and a non-trivial convex hull $\mathcal{J}_{s}(n)$ containing the plateau corresponding to $Z_{i}$. Since $T \in \mathcal{S}_{*}^{b}$, Lemma 4.16(b) implies that if $\mathcal{J}_{s}(n)$ forms a wandering pair for $T_{s}$ for $s>0$ small, then the corresponding convex hull $\mathcal{J}$ is contained in the basin of a periodic attractor for $T$. Since we assumed that $Z_{i}$ is not eventually mapped into a renormalization interval $\hat{K}$, this is impossible. It follows that $T_{s}(n) \notin \xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)$ and therefore $T_{s} \notin \overline{\xi_{i} \cap \mathcal{S}^{b}\left(\Xi_{s}\right)}$ when $s$ is small. Hence $Z_{i, s}$ moves when $s$ small, and so Lemma 7.3 implies again that $h_{0} \geq h_{\text {top }}\left(T_{s} \mid \Omega(K)\right)>h_{\text {top }}\left(T_{t} \mid \Omega(K)\right)$.

Step 4. We claim that for each $T \in \mathcal{S}_{*}^{b}$ there exists $t>0$ so that $T_{t}$ is trivial. Indeed, consider a renormalization interval $K$ for $T$. By Step 3, either $h_{\text {top }}\left(T_{t} \mid \Omega(K)\right)<h_{0}$ or each plateau $Z_{i, s}$ in $K$ is eventually mapped into a renormalization interval $\hat{K} \varsubsetneqq K$ for each $s \in[0, t]$. If $h_{\text {top }}\left(T_{t} \mid \Omega(K)\right)<h_{0}$ for each $t>0$, then each plateau in $K$ will move with positive speed and since $h_{\text {top }}\left(T_{t} \mid \Omega(K)\right)$ is decreasing, this speed will not tend to zero as $t$ increases (unless this plateau becomes part of a block consisting of an even number of plateaus). If each plateau $Z_{i, s}$ in $K$ is eventually mapped into a renormalization interval $\hat{K} \varsubsetneqq K$ for each $s \in[0, t]$, then the period of these intervals $\hat{K}$ will only depend on $T$ (and not on $t$ ). It follows by Step 2, Case B that the speed
of at least one plateau in $\hat{K}$ is bounded away from zero (until all plateaus are in blocks consisting of an even number plateaus), and so in finite time the first return map to $\hat{K}$ will be either monotone or constant. In the former case other plateaus in $K$ start moving, see Step 2, Case C. In the latter case the renormalization interval $\hat{K}$ disappears, see Step 2 Case C(iii), and $h_{\text {top }}\left(T_{t} \mid \Omega(K)\right)$ becomes $<h_{0}$.

Note that $h_{\text {top }}\left(T_{t} \mid \Omega(K)\right)=h_{0}$ implies that the period of $K$ is bounded from below, and so there are only a bounded number of intervals $K$ to be considered in the previous paragraph. Step 4 and therefore part (a) of the proposition follow.

Step 5. From Step 4 it follows that if $T \in \mathcal{S}_{*}^{b}$ then $T_{t}$ never enters the set where the speed of a plateau is actually zero, unless plateaus touch. It follows that if $T$ and $T^{\prime}$ are nearby maps in $\mathcal{S}_{*}^{b}$ so that for at time $s=t$, two plateaus of $T_{s}$ start to touch, then the same two plateaus will start touching for the map $T_{s^{\prime}}^{\prime}$ for some $s^{\prime} \approx t$. It follows that the map $(T, t) \rightarrow \Delta_{t}(T)$ is continuous and that the map $\mathcal{S}_{*}^{b} \ni T \mapsto t_{\Delta}(T)$ is continuous. This finished part (b) of the proposition.

Step 6. Let us now prove part (c) of the proposition and assume that $t>0$ and $\Delta_{t}(T) \in L_{j}\left(h_{0}^{-}\right) \cap L_{j}\left(h_{0}^{+}\right)$. Amongst other things, this implies that there exists a renormalization interval $K$ for $\Delta_{t}(T)$ on which the entropy is $h_{0}$. By Lemma 7.3, this implies that each plateau is either contained in a maximal renormalization interval $\hat{K} \varsubsetneqq K$ or is mapped into such an interval. Since $\Delta_{t}(T) \in L_{j}\left(h_{0}^{-}\right) \cap L_{j}\left(h_{0}^{+}\right)$, the plateau $Z_{j, t}$ is eventually mapped into the boundary of such a renormalization interval $\hat{K}$. If $Z_{j, t}$ maps in an arching way over $\partial \hat{K}$, then this implies that either there exists $s \in[0, t)$ so that $Z_{j, s}$ is not mapped into $\hat{K}$, contradicting that the entropy of $T_{t}$ on $K$ is $h_{0}$ (using Lemma 7.3), or $Z_{j, s}$ maps to $\partial \hat{K}$ for each $s \in[0, t]$. Since $T \in \mathcal{S}_{*}^{b}$, this implies that $Z_{j, t}$ maps to the boundary of the basin of a periodic attractor of $T_{t}$, and so we are done. Next assume that $Z_{j, t}$ maps in a non-arching way over $\partial \hat{K}$ and that $Z_{j, t}$ is part of a wandering pair. In this case, either the first return map to $\hat{K}$ is monotone or $T_{t} \in \xi_{j}$, in which case the speed of $Z_{j, s}$ at $s=t$ is zero. If the former holds, each point in $\hat{K}$ is in the boundary of a component of the basin of a periodic attractor and we are done. If the latter holds, the speed of $Z_{j, s}$ is zero for each $s \in[0, t]$ and $T_{s} \in \xi_{j}$ for each $s \in[0, t]$. But since $T \in \mathcal{S}_{*}^{b}$, this implies that a boundary point of $\hat{K}$ for $T$ (and therefore for $T_{s}$ for each $s \in[0, t]$ ) is in the boundary of a component of the basin of a periodic attractor, and again we are done.

Let us define

$$
\hat{\Delta}_{t}(T)=\Delta_{t_{\Delta}(T) \cdot t}(T)
$$

so that $\hat{\Delta}_{1}(T)$ is a trivial map for each $T \in \mathcal{S}_{*}^{b}$.
7.10. The proof of Theorem 7.1. Now that we have developed the ingredients of the proof, we can define the retract for a fixed $h_{0} \in(0, \log (b+1))$. (The cases
$h_{0}=\log (b+1)$ and $h_{0}=0$ were dealt with in Sections 7.4 and 7.5.)

$$
R_{t}= \begin{cases}\beta_{5 t} & \text { for } t \in\left[0, \frac{1}{5}\right], \\ \Gamma_{5 t-1} \circ \beta_{1} & \text { for } t \in\left[\frac{1}{5}, \frac{2}{5}\right], \\ \Gamma_{1} \circ \hat{\Gamma}_{5 t-2} \circ \beta_{1} & \text { for } t \in\left[\frac{2}{5}, \frac{3}{4}\right], \\ \Gamma_{1} \circ \hat{\Gamma}_{1} \circ \beta_{1} \circ \hat{\Delta}_{(5 t-3)} & \text { for } t \in\left[\frac{3}{5}, \frac{4}{5}\right], \\ \Gamma_{1} \circ \hat{\Gamma}_{1} \circ \beta_{1} \circ r_{5 t-4} \circ \hat{\Delta}_{1} & \text { for } t \in\left[\frac{4}{5}, 1\right] .\end{cases}
$$

Obviously, $R_{0}(T)=T$, and since for $t=1$, the retract $r_{6 t-5}$ has been carried out completely, $R_{1}(T)$ is the same map for each $T \in L_{*}\left(h_{0}\right)$ of the same shape $\epsilon$. All components of $R_{t}$ are continuous in $t$ and $T$, so the same holds for $R_{t}$.

Let us show that $R_{t}$ keeps maps within $\mathcal{S}_{*}^{b}$. First note that the only deformation which takes a map outside the space $\mathcal{S}_{*}^{b}$ is $\hat{\Delta}_{t}$. Take $T^{\prime}$ of the form $T^{\prime}=\hat{\Delta}_{t}(T)$ or $T^{\prime}=r_{t} \circ \hat{\Delta}_{1}(T)$. The deformation $\beta_{t}\left(T^{\prime}\right)$ moves plateaus $Z_{i}, Z_{i+1}$ whose convex hull is mapped into other plateaus. It does so in such a way that $\beta_{1}\left(T^{\prime}\right)$ never eventually maps $\left[Z_{i}, Z_{i+1}\right]$ into the interior of another plateau and so the $\beta$-property (11) on page 46 will hold.

If $h_{\text {top }}\left(T^{\prime}\right)<h_{0}$ then $h_{\text {top }}\left(\beta_{1}\left(T^{\prime}\right)\right)<h_{0}$ and of course $\beta_{1}(T)$ satisfies the $\beta$-property (11). Because of (13) this gives that $\hat{\Gamma}_{1} \circ \beta_{1}\left(T^{\prime}\right) \in \mathcal{S}_{*}^{b}$. In particular, we are done if $T^{\prime}=r_{t} \circ \hat{\Delta}_{1}(T)$ satisfies $h_{\text {top }}\left(T^{\prime}\right)<h_{0}$.

Let us now consider the case that $h_{\text {top }}\left(T^{\prime}\right)=h_{0}$. By the third part of Lemma 7.10, $\Gamma_{1} \circ \hat{\Gamma}_{1} \circ \beta_{1}\left(T^{\prime}\right) \in \mathcal{S}_{*}^{b}$ provided that $T^{\prime}=\beta_{1}\left(T^{\prime}\right) \in L_{j}\left(h_{0}^{-}\right) \cap L_{j}\left(h_{0}^{+}\right)$implies $T^{\prime} \in \mathcal{S}_{*, j}^{b}$. But in Proposition 7.16 it is shown that any map $T^{\prime}$ of the form $T^{\prime}=\beta_{1} \circ \hat{\Delta}_{t}(T)$ with $t>0$ and $T \in \mathcal{S}_{*}^{b}$ has indeed this property, and so again the resulting map belongs to $\mathcal{S}_{*}^{b}$. This concludes the proof of the first part of Theorem 7.1. The 2nd part follows from the construction of $\Gamma_{t}$ and $\beta_{t}$.

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Henk Bruin, Faculty of Mathematics, University of Vienna, Oskar Morgenstern Platz 1, A-1090 Vienna, Austria.
Sebastian van Strien, Department of Mathematics, Imperial College, 180 Queen's Gate, London SW7 2AZ, UK.

E-mail address: henk.bruin@univie.ac.at and s.van-strien@imperial.ac.uk
URL: http://www.mat.univie.ac.at/~bruin/ and http://www2.imperial.ac.uk/~svanstri/


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[^1]:    ${ }^{1}$ i.e., $K$ contains no strictly smaller renormalization interval.

