

# MRST2001: partons and $\alpha_S$ from precise deep inelastic scattering and Tevatron jet data

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## Abstract

We use all the available new precise data for deep inelastic and related hard scattering processes to perform NLO global parton analyses. These new data allow an improved determination of partons and, in particular, the inclusion of the recent measurements of the structure functions at HERA and of the inclusive jets at the Tevatron help to determine the gluon distribution and  $\alpha_S$  better than ever before. We find a somewhat smaller gluon at low  $x$  than previous determinations and that  $\alpha_S(M_Z^2) = 0.119 \pm 0.002$  (expt.)  $\pm 0.003$  (theory).

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# 1 Introduction

Recently a great deal of new data have become available which help to determine the parton distributions of the proton. In particular we have new measurements of the structure functions by the H1 [1, 2, 3] and ZEUS [4] collaborations at HERA, and of the inclusive jet distribution by the D0 [5] and CDF [6] collaborations at the Tevatron. These new data are both more precise and extend the kinematic range, and thus constrain the parton distributions, and the strong coupling  $\alpha_S$ , more tightly than ever before. ZEUS have also released a new measurement of the charm contribution to the structure function [7] which, although it still has large errors, covers a wider kinematic range than previously. In addition, the CCFR collaboration [8] have re-analysed their neutrino data in a model independent way and the discrepancy with the NMC data for  $x \lesssim 0.1$  has been largely resolved. Also NuTeV data are becoming available [9], both for single and double muon production, and are improving the constraints from the neutrino sector. The E866 collaboration [10] also have increased statistics for  $pp$  and  $pn$  Drell-Yan production and improve the determination of the difference between the  $\bar{u}$  and  $\bar{d}$  distributions. Finally we note that as well as the usual data sets used in our previous fits, i.e. BCDMS [11] and SLAC [12] proton data, NMC proton and deuterium data [13], E665 proton data and deuterium data [14], CCFR data on  $F_3^{\nu(\bar{\nu})N}(x, Q^2)$  [15], E605 Drell-Yan data [16] and CDF  $W$ -asymmetry data [17] we have also included BCDMS [18] and SLAC deuterium data [12] in order to obtain as precise a determination of the separate contributions of the up and down valence quarks at high  $x$  as possible.<sup>2</sup> We no longer include prompt photon data due to theoretical problems and possible inconsistencies between data sets, and instead allow the high  $x$  gluon to be determined by the vastly improved Tevatron jet data.

We note that both H1 [3] and ZEUS [4] have recently performed NLO DGLAP fits to their respective data, supplemented in the former case by BCDMS data with  $y_\mu > 0.3$ , and in the latter case by BCDMS, NMC, E665 and CCFR  $F_3^{\nu(\bar{\nu})N}(x, Q^2)$  data. In particular, the H1 analysis determines  $\alpha_S$  and the gluon simultaneously. A value

$$\alpha_S(M_Z^2) = 0.1150 \pm 0.0017(\text{expt.}) \begin{array}{l} +0.0009 \\ -0.0005 \end{array} \text{ (model)} \quad (1)$$

is obtained, with an additional uncertainty of about  $\pm 0.005$ , mainly due to the uncertainty in the renormalization scale. A preliminary ZEUS analysis, reported at DIS2001, quoted  $\alpha_S(M_Z^2) = 0.1172 \pm 0.0008(\text{uncor.}) \pm 0.0054(\text{cor.})$  [20]. There is also an analysis including some of both the recent H1 and ZEUS data along with NMC, SLAC and BCDMS data, and which allows higher twist contributions, which obtains  $\alpha_S(M_Z^2) = 0.1171 \pm 0.0015(\text{expt.})$  [21]. We will find that the inclusion of additional data sets tend to increase these values somewhat. As an example of this we emphasize that the Tevatron jet data are an important ingredient in pinning down the value of  $\alpha_S$  from deep inelastic scattering (DIS) and related data, since they provide the dominant constraint on the gluon at large  $x$ . In fact, the inclusion of new jet data [5, 6] into

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<sup>2</sup>For all deuterium structure functions we correct for shadowing effects [19].

the global analysis considerably improves the determination of the gluon. For example, we find an uncertainty of about 15% on the gluon distribution at  $x = 0.4$  and  $Q^2 = 20 \text{ GeV}^2$ , but this is correlated with the value of  $\alpha_S(M_Z^2)$ .

## 2 The new optimum parton set

We perform a global NLO DGLAP analysis incorporating all the high-precision data mentioned in the Introduction. The evolution begins at  $Q_0^2 = 1 \text{ GeV}^2$  where there are three active quark flavours. We work in the  $\overline{\text{MS}}$  renormalization scheme and use the Thorne-Roberts [22] variable flavour number procedure to evolve through the charm and bottom thresholds. As well as deciding on data sets, we have to decide on a set of cuts in the usual variables  $Q^2$ ,  $W^2$  and  $x$ . In order to investigate this we made a study of the sensitivity of the analysis to variation of these data cuts. We discovered that there was only marginal evidence for an improvement in quality if the  $Q^2$  cut was raised from  $2 \text{ GeV}^2$  to  $3 \text{ GeV}^2$  and no marked improvement above this. There was a marked improvement in quality if  $W^2$  is raised from our previous cut of  $10 \text{ GeV}^2$  until we reach  $12.5 \text{ GeV}^2$ , which may easily be interpreted as due to the influence of higher twist and/or large  $\ln(1-x)$  terms in the perturbative expansion. Hence, for the global fit presented below DIS data with  $Q^2 > 2 \text{ GeV}^2$  and  $W^2 > 12.5 \text{ GeV}^2$  are included, in order to exclude regions where higher twist and/or higher orders in  $\alpha_S$  are expected to play an important role. We also found that if a lower cut in  $x$  was introduced there was continual improvement in the quality of the fit until  $x$  reached a value of about 0.005, suggesting that  $\ln(1/x)$  terms in the perturbative series may be important. The results and consequences of these cuts, particularly that in  $x$ , will be dealt with in a future paper [23], but for the present analysis we take the conventional approach of not using any  $x$  cut and investigate/suffer the consequences.

The optimum global NLO fit is obtained with the starting parameterizations of the partons at  $Q_0^2 = 1 \text{ GeV}^2$  given by

$$xu_V = 0.158x^{0.25}(1-x)^{3.33}(1+5.61x^{0.5}+55.49x) \quad (2)$$

$$xd_V = 0.040x^{0.27}(1-x)^{3.88}(1+52.73x^{0.5}+30.65x) \quad (3)$$

$$xS = 0.222x^{-0.26}(1-x)^{7.10}(1+3.42x^{0.5}+10.30x) \quad (4)$$

$$xg = 1.90x^{0.09}(1-x)^{3.70}(1+1.26x^{0.5}-1.43x)-0.21x^{-0.33}(1-x)^{10}. \quad (5)$$

The flavour structure of the light quark sea is taken to be

$$2\bar{u}, 2\bar{d}, 2\bar{s} = 0.4S - \Delta, \quad 0.4S + \Delta, \quad 0.2S \quad (6)$$

with  $s = \bar{s}$ , as implied by the NuTeV data [9], and where

$$x\Delta = x(\bar{d} - \bar{u}) = 1.195x^{1.24}(1-x)^{9.10}(1+14.05x-45.52x^2) \quad (7)$$

The masses of the quarks are taken to be  $m_c = 1.43 \text{ GeV}$  and  $m_b = 4.3 \text{ GeV}$ , the former giving the best fit to the charm structure function data. The optimum fit corresponds to

$\alpha_S(M_Z^2) = 0.119$ , i.e.  $\Lambda_{\overline{\text{MS}}}(n_f = 4) = 323$  MeV, in very good agreement with the world average.<sup>3</sup> We estimate the error in  $\alpha_S(M_Z^2)$  due to the errors on the data fitted in the global analysis to be about  $\pm 0.002$ , as will be explained in detail later. The resulting partons are shown in Fig. 1.

The improved HERA data greatly increase the constraints on the gluon at small  $x$ . The extra term in (5) is required to achieve an acceptable fit, and allows the starting gluon to become negative at small  $x$ . A fixed  $(1-x)^{10}$  behaviour is incorporated in this extra term so that only the small  $x$  form of the gluon is affected. Not including this additional term, which allows the input gluon to be negative at small  $x$ , would lead to the input gluon being strongly valence-like and to a global increase in  $\chi^2$  of about 100. Note that only half this increase comes from the data points at very low  $x$  (say  $x < 0.001$ ), the rest coming from the HERA and NMC points in the region  $0.001 < x < 0.1$ , as will be discussed in the next section. The gluon in the present analysis becomes positive for all  $x > 10^{-5}$  and  $Q^2 > 5$  GeV<sup>2</sup>, and for  $Q^2 > 2 - 3$  GeV<sup>2</sup> for  $x > 10^{-4}$ .

Recall that in the MRST99 analysis [26] the uncertainties in the gluon were illustrated by presenting the optimal fit  $g$  together with two fits  $g \uparrow$  and  $g \downarrow$ , with larger and smaller gluons at large  $x$ , which represented the extremes of acceptable descriptions of the data. The present analysis, with greatly improved data, significantly reduces the uncertainty in the gluon distribution and yields an optimal solution with a large  $x$  gluon nearer to  $g \uparrow$  than to  $g$ . For this reason in Fig. 2 we compare the present partons with those of the  $g \uparrow$  set of MRST99 [26]. We see that the major difference is in the gluon, or is a consequence of this changed gluon. First, we note the extended parameterization for the gluon, required by the new HERA data, leads to a far smaller gluon at the lowest  $x$  and  $Q^2$ . Since the quarks are determined by evolution driven by the gluon in this range, they are also smaller than their MRST99 counterparts. The gluon is also smaller in the range  $x \sim 0.3$  than that for MRST99( $g \uparrow$ ) — the jet data requiring less gluon in this range than the prompt photon data with no intrinsic  $k_T$  included. Both these reductions in the gluon allow for slightly more gluon in the range  $x \sim 0.1$ , giving an increased  $dF_2(x, Q^2)/d \ln Q^2$  for  $x$  a little below this. The shape of the charm (and, to a lesser extent, the strange) distribution simply follows the gluon since it is generated mainly by evolution from the gluon. Finally we note that the down quark is slightly smaller at high  $x$  than in MRST99 due to the effect of the extra deuterium data included in the present fit, and as a consequence it is slightly larger for values of  $x$  in the region of 0.01.

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<sup>3</sup>We use the matching between the  $n_f$ - and  $n_{f+1}$ -flavour couplings calculated in [24], and corrected in [25], up to NLO in  $\alpha_S$ . At this order the coupling is continuous across threshold but the derivative is discontinuous. More details may be found in section 3 of the first of [22].

### 3 The description of the DIS data

A good description of the HERA data is obtained, as can be seen from Figs. 3-6.<sup>4</sup> Compared to MRST99, the curves for  $F_2$  are flatter in  $Q^2$  for  $x < 0.001$ , but slightly steeper at higher  $x$ . In fact for  $0.001 < x < 0.01$  the data, particularly the high  $Q^2$  H1 and the NMC data, would prefer a higher  $dF_2/d\ln Q^2$ , as can be seen in Figs. 5 and 6. The systematic failure in this region is a cause for concern regarding the validity of an NLO fit. In fact we note that the aforementioned improvement to the fit in this region, which comes from allowing the negative input gluon at small  $x$ , is because this form allows there to be more gluon in the moderate  $x$  region (from the momentum sum rule), and hence a larger value of  $dF_2(x, Q^2)/d\ln Q^2$ , as preferred by the data. We also note that although at the lowest  $x$  and  $Q^2$  our gluon distribution is considerably negative,  $dF_2/d\ln Q^2$  is quite clearly positive. This highlights the fact that the frequently quoted relationship  $dF_2/d\ln Q^2 \propto \alpha_S(Q^2)xg(x, Q^2)$  is not even approximately true at small  $x$  and  $Q^2$  when one works beyond leading order in perturbative QCD. Qualitative arguments about evidence for saturation *etc.* which rely on this relationship should be treated with caution.

The fits to HERA data have been performed using the simplistic procedure of adding the systematic and statistical errors in quadrature. However, we have performed an analysis of the effects of the correlated errors. Fits where we fit first to the data with only uncorrelated errors, and then let the contributions of the correlated errors come into effect, have been investigated. We find that the absolute value of the  $\chi^2$  using this procedure increases for the ZEUS data and stays more or less constant for the H1 data, but that the incremental changes in  $\chi^2$  when comparing different theoretical results stays much the same as when using the more simplistic addition in quadrature of all errors. Hence, we have decided to use this simpler procedure in determining the partons. We note that we do however let the ZEUS normalization go to its lower limit of 98% in order to obtain our best fit.

The comparison with the charm data [7, 27] can be seen<sup>5</sup> in Fig. 7. As one can see it is of a perfectly acceptable quality, and the errors on this data are still large. There is, however, a slight tendency to undershoot the data at the lowest values of  $x$  and  $Q^2$ , and this may be a sign of the need to improve the theoretical treatment in this region. In this region of low  $x$  and  $Q^2$  the prediction for the charm structure function is a little smaller than that for the MRST99 partons, which is entirely due to the smaller gluon we now find in this region. Finally we note that the fit to the higher  $x$  EMC charm data [29] is very similar to that for the default set in Fig. 27 of [30] (the high  $x$  gluon is now a little larger, but  $m_c$  is 1.43 GeV rather than 1.35 GeV), and hence is perfectly acceptable.

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<sup>4</sup>Note that as in previous fits we have effectively fit the published cross section rather than  $F_2(x, Q^2)$  at the larger values of  $Q^2$ , i.e. we have corrected for our own values of  $F_L(x, Q^2)$  rather than use those obtained by the fits by the experiments.

<sup>5</sup>Updated charm data from the H1 collaboration have recently become available [28], but the results depend on which Monte Carlo is used to extrapolate over the full range of phase space. Since these data are similar to their previous charm data [27], we show only the latter in Fig. 7.

The prediction for  $F_L(x, Q^2)$  is shown in Fig. 8, which also shows the prediction of the MRST99 partons. We see that the increased precision of the HERA data and the increased flexibility of the gluon parameterization have led to a significant decrease in the prediction for  $F_L(x, Q^2)$  at low  $x$  and  $Q^2$ , replaced by a slight increase for  $x \sim 0.05$ . Indeed, it now seems as though  $F_L(x, Q^2)$  is taking a distinctly unphysical form for low  $x$  and  $Q^2 < 5 \text{ GeV}^2$ , and for part of this range is negative, and therefore certainly disallowed.<sup>6</sup> This may be taken as another clear sign that the standard NLO fit is not working completely properly at small  $x$ .<sup>7</sup>

At higher  $x$  the main change in our fit is due to the reanalysis of the CCFR data [8]. Their reanalysis no longer extracts  $F_2(x, Q^2)$  by modeling both  $F_L(x, Q^2)$  and  $\Delta x F_3(x, Q^2)$ , but now extracts either  $F_L(x, Q^2)$  or  $\Delta x F_3(x, Q^2)$  (there is a high degree of correlation between these) and  $F_2(x, Q^2)$  separately in a Physics Model Independent manner. This has gone a long way towards resolving the apparent discrepancy between CCFR and NMC data on  $F_2(x, Q^2)$ , where it had previously been impossible to simultaneously fit both for  $x < 0.1$ . The quality of the fit to the new CCFR  $F_2(x, Q^2)$  data is shown in Fig. 9. Overall the fit is very good. One might argue that there is still a systematic problem at the lowest  $Q^2$ , but this is far less pronounced than with previous analyses. There is also potentially a small error associated with the shadowing corrections (details of which are found in [30]) which we do not account for. The reanalysis has also established the validity of the previous  $F_3(x, Q^2)$  neutrino data. These data are essentially unaffected by the reanalysis, but we are now confident in using the data over the whole  $x$  range, rather than just for  $x \geq 0.1$ . The fit is good over the whole range of  $x$ . Note that we normalize the complete set of CCFR data up by 1% in order to obtain the best fit.

Other than this, the other new DIS data (at least new for our fit) at high  $x$  are simply the SLAC and BCDMS deuterium data, which we have introduced for the first time. The fit to these data is shown in Fig. 10. It is of a perfectly acceptable quality, and one can see that, as with the proton data, the SLAC deuterium measurements prefer a rather steeper fall with  $Q^2$  than the BCDMS data, and consequently a larger  $\alpha_S$ . The SLAC data are normalized up by 2.5% and the BCDMS down by 2%.

## 4 Tevatron Jet Data and the Gluon

One of the major differences between the MRST2001 partons and our previous parton sets is the manner in which the high  $E_T$  Tevatron jets have been included. In the past [30] we have simply checked that there is reasonable agreement with our predictions and the jet data. The difficulties in using the prompt photon data in order to determine the high  $x$  gluon combined

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<sup>6</sup>In principle, it is internally inconsistent to fit to  $F_2(x, Q^2)$  data in a region where the predicted values of  $F_L(x, Q^2)$  are negative, namely  $x < 10^{-4}$  and  $Q^2 \gtrsim 2 \text{ GeV}^2$ . Since only 6 points are affected, carrying practically no weight in the fit, we do not remove these points.

<sup>7</sup>We do not compare to the H1 extraction of  $F_L(x, Q^2)$  [31] since the different assumptions used in our fit lead to significantly different forms for the gluon and for  $\alpha_S(M_Z^2)$ , and hence different extrapolations into the high  $y$  region.

with the considerable improvement in Tevatron jet data [5, 6] has led to a change in emphasis. Besides the increase in precision, the D0 jet data are available in a range of rapidity intervals and so constrain the partons, and the gluon in particular, over a much wider  $x$  range. We, therefore, now include the D0 and CDF jet data in the global fit on an equal footing with all other data sets.<sup>8</sup> However, because in this case the correlated systematic errors are the dominant source of error, being much larger than the uncorrelated errors, it is imperative to deal with these in a correct manner. In fact we adopt the same method of fitting to the data as do the respective experimental collaborations when describing their own data. Note, however, that rather than using some NLO prescription such as JETRAD [33] or EKS [34] to generate an NLO correction for each point, we derive a smooth NLO K-factor by fitting to a set of such points. Since the NLO corrections generated from the above programs have some error and scatter, this means that our value of  $\chi^2$  will not be identical to that obtained by the experiments themselves for the same parton set (though it will be very close). In particular our values of  $\chi^2$  for the CDF1B jet data are a little higher, while for the D0 jet data they are slightly lower.

We find that for our best global fit we obtain a reasonable fit to the combined jet data with a  $\chi^2$  of 170 for 113 points. The quality of the fits is shown in Figs. 11 and 12 — the error bars account for uncorrelated errors alone. In both cases it is clear that while at the low  $E_T$  end the normalization is about correct, at higher  $E_T$  (and rapidity) the theory lies below the data. An acceptable fit is then obtained by accounting for the correlated systematic errors. For the CDF1B fit the (data–theory) obtained is allowed to move by letting data move relative to theory by application of each of the sources of correlated error at the cost of an increase in  $\chi^2$  of  $s_k^2$  for each source of correlated error  $k$ , where  $s_k$  is the fraction of one standard deviation which has been used for the given error source.<sup>9</sup> Hence in Fig. 12 we see that the data at higher  $E_T$  have effectively been brought down to match the shape of the theory prediction. For the D0 jet data the fit is performed using the full error matrix which accounts for the correlated systematic errors in a less transparent manner, but which clearly obtains a good fit in much the same way.

At the central value of  $\alpha_S(M_Z^2)$  of 0.119, this global fit (including jet data) allows a variation in  $g(x, Q^2)$  of about 5% for  $x > 0.1$  and  $Q^2 \sim 2000 \text{ GeV}^2$ , which corresponds to 10 – 15% accuracy for  $x > 0.1$  and  $Q^2 \sim 20 \text{ GeV}^2$ . This is a factor of 3 or so less than the MRST99 variation of the gluon ranging from the  $g \uparrow$  to the  $g \downarrow$  gluon, and hence we do not provide parton sets with gluon extremes. Since the body of jet cross section data is (very roughly)  $\propto \alpha_S(E_T^2/4)g(x, E_T/2)$ , then  $g(x, \mu^2)$  for  $\mu^2$  of order  $10^3 \text{ GeV}^2$  is roughly inversely proportional to  $\alpha_S(M_Z^2)$ . However, at high  $x$  the gluon distribution decreases more rapidly with increasing  $Q^2$ , the larger the value of the coupling. This increase in speed of evolution with increasing  $\alpha_S$

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<sup>8</sup>The D0 collaboration have recently produced data using the  $k_T$  algorithm rather than the more usual cone algorithm [32]. The agreement between the two methods is moderate, the major difference being at low  $E_T$ . We use the original data since we feel these have been more extensively studied, and because they also cover a much wider range in pseudorapidity.

<sup>9</sup>Since these correlated errors are expected to cut off rather more sharply than Gaussian errors we limit each  $s_k$  to 1. This does not affect any results at all significantly.

more than compensates for the decrease in the high-scale gluon required by the jets with  $\alpha_S$ , and for low  $Q^2$  ( $Q^2 \sim 10 \text{ GeV}^2$ ) the high  $x$  gluon increases as  $\alpha_S(M_Z^2)$  increases.

We note, however, that these optimum global fits are not the best possible fits to the high  $E_T$  jet data. The fit is only achieved by compensating for the smallness of the theory at high jet  $E_T$  and  $\eta$  (both of which probe the highest  $x$ ) using the correlated systematic errors. Hence, the fit can be improved by an increase in the size of the high  $x$  gluon. In principle it is possible to obtain a fit with a  $\chi^2$  of about 120 for the 113 points (see below) rather than  $\chi^2 = 170$ , the scatter of data points making this  $\chi^2$  value about the lowest that is achievable. At the central  $\alpha_S(M_Z^2)$  value of 0.119 it is possible to raise the high  $x$  gluon sufficiently to improve the quality of the jet fit to  $\chi^2 = 135$ , but only at the cost of  $\Delta\chi^2 = 60$  for the rest of the data. This is mainly at the expense of the description of the moderate  $x$  DIS data, i.e. H1, ZEUS and NMC data, since the increase of gluon at high  $x$  is countered by a decrease at intermediate  $x$ , and hence a decrease in  $dF_2(x, Q^2)/d\ln Q^2$ . At lower  $\alpha_S$  the price is even higher since the lower  $\alpha_S$  already impacts upon the behaviour of  $dF_2(x, Q^2)/d\ln Q^2$ .

However, as one goes to higher  $\alpha_S(M_Z^2)$  the situation changes. At  $\alpha_S(M_Z^2) = 0.121$  one can obtain a fit to the jet data with  $\chi^2 = 118$ , and where this improvement is only marginally overcompensated by the deterioration in the rest of the fit compared to the best global fit. For this set of partons, denoted by MRST2001J or simply J, the fit to the jet data is shown in Figs. 13 and 14. For the fit to the D0 data the shape is obviously greatly improved, both as a function of  $E_T$  and of  $\eta$ , demonstrating that the apparent excess  $\chi^2$  in the description of the Tevatron jet data is either a problem of parton distributions or of systematic errors, but is unlikely to be a sign of new physics. The normalization of the theory is a little high, but this is easily accounted for by the systematic error in normalization. The fit to the CDF1B data actually gives a slightly worse  $\chi^2$  than before. But now the fit does not rely on a large shift of data due to systematic errors, and is perhaps more satisfactory in this sense. The problem with this J set of partons is the behaviour of the gluon. The input form at  $Q_0^2 = 1 \text{ GeV}^2$  is

$$xg = 123.5x^{1.16}(1-x)^{4.69}(1-3.57x^{0.5}+3.41x) - 0.038x^{-0.5}(1-x)^{10}, \quad (8)$$

which is shown in Fig. 15, together with the behaviour at  $Q^2 = 20 \text{ GeV}^2$ . We see that the input shape has a rather worrying “kink” which results in the distinct “shoulder” at  $Q^2 = 20 \text{ GeV}^2$ .<sup>10</sup> We do not deem this to be an acceptable gluon (admittedly a subjective decision), and rule this fit out, although we do make the MRST2001J set of partons available. At  $\alpha_S(M_Z^2) = 0.120$  we obtain a similar result, i.e. the best overall fit gives about  $\chi^2 = 135$  for the jet data, but has a gluon with the same type of peculiarities (though less severe). Again we rule this fit out. Hence, in our fits we impose the condition that  $d^2(xg(x, Q_0^2))/dx^2$  does not change sign in the

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<sup>10</sup>This seems to be made possible by the interplay between a very large coefficient for the first term in Eq.(8), the large power of  $x$ , i.e.  $x^{1.16}$  in this term, and the extra effect of the second term controlling the very small  $x$  behaviour. This second term then effectively frees one parameter in the first term, which for previous parton sets represented the full parameterization of the input gluon, allowing more flexibility in the high  $x$  form of the gluon.



region of high  $x$  which rules out the possibility of both kinks and shoulders in the high  $x$  gluon distribution. Imposing this condition results in a fit to the jet data within the global fit which is roughly independent of  $\alpha_S(M_Z^2)$ .<sup>11</sup>

## 5 Description of other data

The fit to much of the rest of the data is very much along the lines discussed in detail in [30]. The NuTeV data [9] on single and double muon production do not qualitatively change the conclusions regarding the strange contribution to the sea already indicated by the CCFR dimuon data [35], i.e. that the strange distribution is acceptably obtained from half the average of the  $\bar{u}$  and  $\bar{d}$  distributions at  $Q_0^2 = 1 \text{ GeV}^2$ . Similarly the E866 collaboration [10] have provided new data on the Drell-Yan asymmetry which is more accurate and extends the kinematic range slightly, but does not really change the relative  $\bar{u}, \bar{d}$  behaviour of the partons. In particular, our simple parameterization of  $\bar{d} - \bar{u}$  still suggests that  $(\bar{d}/\bar{u}) \leq 1$  for  $x > 0.35$ , but there is no evidence whether this is really true or not. The lepton rapidity asymmetry data from CDF [17] (related to the  $W$  rapidity asymmetry) also continue to give us important information on the  $u/d$  ratio.

Finally the E605 Drell-Yan data [16] still play an important role in pinning down the form of the sea quarks at high  $x$ . However, they also play an important and unexpected role in influencing the fit to the jet data and determining  $\alpha_S(M_Z^2)$ . As we will see below, the quality of the description of these data deteriorates as  $\alpha_S(M_Z^2)$  increases. This is actually an indirect effect. As  $\alpha_S(M_Z^2)$  increases, the high  $x$  gluon at lower  $Q^2$  increases so as to give the correct gluon normalization when evolved up to the scales appropriate for the description of the jet data. This larger high  $x$  gluon (and larger  $\alpha_S(M_Z^2)$ ) drives a positive evolution of the high  $x$  sea quarks. As this effect becomes more significant it distorts the shape of the sea quark distribution in the range relevant for fitting Drell-Yan data, worsening the fit. Therefore the E605 data prefer lower value of  $\alpha_S(M_Z^2)$  and a lower high  $x$  gluon. Indeed, for the MRST2001J type partons the fit to the Drell-Yan data deteriorates quite seriously compared to the best global fit. Hence, these data have assumed a more important role in the context of the whole global fit than previously.

## 6 Quality of Fit and Determination of $\alpha_S(M_Z^2)$ .

The quality of the central fit for the major data sets is shown in Table 1 below. For each of the smaller data sets, e.g. CDF  $W$ -asymmetry [17] and E866 Drell-Yan asymmetry [10], the  $\chi^2$  per degree of freedom is about 1 per point. For all the DIS data sets the numbers are

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<sup>11</sup>We note that a very good fit to the jet data could be achieved for  $\alpha_S(M_Z^2) < 0.118$  with a gluon without peculiarities, but that this results in a fit to the rest of the data which is very poor. This problem is improved, though not completely rectified, when an  $x$ -cut is applied (see [23]).

quoted for statistical and systematic errors added in quadrature. The quality of the fits to the individual data sets is satisfactory. For the E605 data the systematic errors are quoted in a slightly ambiguous manner, and are generally subdominant, and so we fit to statistical errors alone. Hence, the quite large  $\chi^2$  in this case. The treatment of the correlated systematic errors for the Tevatron jet data has been discussed in Section 4.

Table 1: Quality of the fit for MRST2001 partons to different data sets. The first MRST column shows the  $\chi^2$  values of the optimum fit with  $\alpha_S(M_Z^2) = 0.119$ . Also shown are the values for parton sets obtained from fits with  $\alpha_S(M_Z^2) = 0.117$  and  $0.121$ , as well as those for parton set J which has structure in the high  $x$  gluon.

| Data set           | No. of data pts | MRST | MRST 0.117 | MRST 0.121 | MRST J |
|--------------------|-----------------|------|------------|------------|--------|
| H1 $ep$            | 400             | 382  | 386        | 378        | 377    |
| ZEUS $ep$          | 272             | 254  | 255        | 258        | 253    |
| BCDMS $\mu p$      | 167             | 193  | 182        | 208        | 183    |
| BCDMS $\mu d$      | 155             | 218  | 211        | 226        | 219    |
| NMC $\mu p$        | 126             | 134  | 143        | 127        | 135    |
| NMC $\mu d$        | 126             | 100  | 108        | 95         | 100    |
| SLAC $ep$          | 53              | 66   | 71         | 63         | 67     |
| SLAC $ed$          | 54              | 56   | 67         | 47         | 58     |
| E665 $\mu p$       | 53              | 51   | 50         | 52         | 51     |
| E665 $\mu d$       | 53              | 61   | 61         | 61         | 61     |
| CCFR $F_2^{\nu N}$ | 74              | 85   | 88         | 82         | 89     |
| CCFR $F_3^{\nu N}$ | 105             | 107  | 103        | 112        | 110    |
| NMC $n/p$          | 156             | 155  | 155        | 153        | 161    |
| E605 DY            | 136             | 232  | 229        | 247        | 273    |
| Tevatron Jets      | 113             | 170  | 168        | 167        | 118    |
| Total              | 2097            | 2328 | 2346       | 2345       | 2337   |

The way in which the quality of the fit to both the total and to each data set varies with  $\alpha_S(M_Z^2)$  is shown in detail in Fig. 16. It must be remembered that the quality of the fit for a single data set within the context of a global fit is not the same thing as the quality of the fit for that set alone, and many sets influence each other strongly. Nevertheless, one can pick out some interesting facts from Fig. 16.

For the DIS data sets it is clear that only the two BCDMS sets strongly prefer lower values of  $\alpha_S(M_Z^2)$ . These are more than compensated by the SLAC and NMC data sets, which for both proton and deuterium structure functions strongly prefer higher values of  $\alpha_S(M_Z^2)$ . Both CCFR data sets are relatively insensitive to the value of the coupling, at least for  $0.116 < \alpha_S(M_Z^2) < 0.122$ . This also appears to be true for the H1 and ZEUS data sets. However, this

latter apparent insensitivity is due to the fact that the combined HERA data sets carry a lot of weight in the fit, and the gluon distribution at small  $x$  is largely determined by ensuring that these data are fit well. This is therefore just a manifestation of the long-established fact that the small  $x$  gluon and the value of  $\alpha_S(M_Z^2)$  are completely correlated in fits to the HERA data, and without any additional handle on the gluon<sup>12</sup> there is no way to remove this. It is clear from Fig. 16 that if one takes only BCDMS data as well as HERA data, as in the H1 analysis [3], one will determine a low value of  $\alpha_S(M_Z^2)$ , but taking SLAC or NMC as the additional set a very different conclusion will be reached.

The combined Tevatron jet data behaves similarly to the HERA data, i.e. the gluon conspires with  $\alpha_S(M_Z^2)$  to give roughly the same  $\chi^2$  for all  $\alpha_S$ . Interestingly the jet data and HERA data manage to conspire with each other so that the sum of their  $\chi^2$  remains roughly constant. As  $\alpha_S(M_Z^2)$  increases the gluon at moderate  $Q^2$  and high  $x$  increases to maintain the fit to the jet data. From the momentum sum rule this leaves less gluon at small  $x$ , but the larger  $\alpha_S(M_Z^2)$  manages to keep the value of  $dF_2(x, Q^2)/d\ln Q^2$  acceptable. The fact that this trade-off between HERA  $F_2(x, Q^2)$  data and Tevatron jet data results in almost complete insensitivity to  $\alpha_S(M_Z^2)$  strikes us as remarkable. Note, however, that the constant total  $\chi^2$  for the jets is made up of a contribution from the CDF1B data which increases sharply with increasing  $\alpha_S(M_Z^2)$ , and a rapidly falling contribution from the D0 jet data. In detail one finds that the general normalization and shape of the theory compared to data improves with increasing  $\alpha_S(M_Z^2)$ . This leads to the improvement in the fit to D0 data. However, the precise shape of the CDF1B data seems easiest to achieve by obtaining a poor comparison between theory and data which is then compensated for by quite large movements coming from the correlated errors. When the shape and size is nearly correct to begin with this seems to leave less room for manoeuvre for the correlated errors to produce exactly the correct shape (note that the  $\chi^2$  is better in Fig. 12 than in Fig. 14). Hence, if one is uncomfortable about letting the correlated errors conspire to move the data by a large amount the high  $\alpha_S(M_Z^2)$  fits are better.

Finally, as discussed in the last section, the E605 Drell-Yan data prefer a low value of  $\alpha_S(M_Z^2)$ . However, this is mainly due to the correlation between the value of  $\alpha_S(M_Z^2)$  and the high  $x$  gluon brought about by the jet and HERA data.

Putting all the contributions together we obtain a total  $\chi^2$  which has quite a sharp minimum at  $\alpha_S(M_Z^2) = 0.119$ . We then adjudge the error in this best value of  $\alpha_S(M_Z^2)$ , within the context of an NLO-in- $\alpha_S$  fit, by letting the  $\chi^2$  increase by about 20 units. Clearly it is inappropriate to base the error on the increase of a single unit, for a variety of reasons. First, the treatment of the errors in this analysis is far from statistically rigorous, and even if it were, the errors themselves are far from having a true Gaussian distribution. Also, we have made many decisions in performing this analysis, such as data cuts, the choice of parameterizations of partons, etc. Changing any of these and refitting would lead to changes in  $\chi^2$  of about 5–10 for the remaining data, and so our increase in  $\chi^2$  should be at least this value. Making our choice of an increase

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<sup>12</sup>The charm structure function is strongly correlated to the evolution of the total structure function, and therefore does not provide an independent constraint.

of 20 we obtain  $\alpha_S(M_Z^2) = 0.119 \pm 0.002(\text{expt.})$ . We see from Fig. 16 that beyond these limits the global  $\chi^2$  increases very quickly.<sup>13</sup>

Thus, we present our determination of  $\alpha_S(M_Z^2)$  as

$$\alpha_S(M_Z^2) = 0.1190 \pm 0.002(\text{expt.}) \pm 0.003(\text{theory}). \quad (9)$$

We do not adopt the traditional, but *ad hoc*, manner of obtaining the theoretical error by varying renormalization and factorization scales up and down by factors of 2 (or of 4). This takes no account of the errors attributable to higher order logarithmic enhancements. For example, in DIS there are additional logarithms in  $(1-x)$  and  $1/x$  in the coefficient functions and splitting functions at higher orders in  $\alpha_S$  which variations in scale tell us nothing about. Similar logarithmic enhancements also exist for the other quantities fitted, such as data near threshold. Hence, we obtain our theory error by comparing with alternative theoretical treatments which do tell us something more concrete about the missing corrections, i.e. approximate NNLO fits e.g. [36], or fits which attempt a resummation of  $\ln(1/x)$  and  $\ln(1-x)$  terms [37]. These suggest that 0.003 is an appropriate theoretical error.<sup>14</sup>

As regards the errors on the partons themselves, in a separate study we will present the uncertainties in the predictions of key observables, and show how they reflect the uncertainties on the parton distributions. An example of this is seen in [38]. However, as in the case of  $\alpha_S(M_Z^2)$ , we believe the theoretical errors to be generally more important than the experimental errors, particularly in some regions of parameter space.

## 7 Conclusions

In this paper we have performed global analyses of all the most up-to-date data on deep inelastic scattering and related processes in order to best determine the parton distributions and the value of  $\alpha_S(M_Z^2)$  within the context of a conventional NLO fit. This is an improvement on our previous analyses mainly because of some very important new sets of data. In particular the new HERA data [1, 2, 3, 4] are far more precise than previously and cover an extended range in  $x$  and  $Q^2$ . Also the new D0 and CDF Tevatron jet data are again more precise, with systematic errors which are better understood, and which extend their previous kinematic ranges. These new HERA and Tevatron data sets together impose far more stringent limits on the parton distributions than ever before. We also obtain a tight constraint on the value

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<sup>13</sup>We have also investigated the fits without the Tevatron jet data included. Even though removing this constraint allows gluons to migrate to lower  $x$  and in principle fit the HERA data with lower  $\alpha_S$ , the overall impact on the global fit is not large. The minimum moves down by  $\Delta\alpha_S(M_Z^2) \sim 0.0002$ . For rather low values of  $\alpha_S(M_Z^2)$ , e.g. 0.116 or lower, the removal of the high  $x$  gluon constraint does allow an improvement in the fit to HERA and NMC data, but at this value of  $\alpha_S(M_Z^2)$  the global fit has become much worse anyway, and all one would obtain with jet data removed would be a shape for the total  $\chi^2$  like that in Fig. 16, but with the slope on the left-hand side a little more shallow.

<sup>14</sup>These investigations suggest that  $\alpha_S(M_Z^2)$  might move down slightly from 0.119.

of  $\alpha_S(M_Z^2)$ . Our best overall fit corresponds to  $\alpha_S(M_Z^2) = 0.119$ , and investigating variations about this minimum we obtain  $\alpha_S(M_Z^2) = 0.119 \pm 0.002(\text{expt.}) \pm 0.003(\text{theory})$ . The quality of the fits for the two experimental limits of 0.117 and 0.121 can be seen in Fig. 16, and are also detailed in Table 1.

The new data sets have a particularly strong impact on the gluon distribution. In order to fit the new HERA data well we have been forced into an extension of our previous input gluon parameterization, allowing it to become negative at small  $x$ . Indeed at  $Q_0^2 = 1 \text{ GeV}^2$  it behaves like  $xg(x, Q_0^2) \sim 0.2x^{-0.33}$  for  $x < 0.001$ , and this is necessary not only to obtain a good fit at low  $x$  and  $Q^2$ , but also to allow enough gluon at higher  $x$  to obtain large enough  $(dF_2/d \ln Q^2)$  for  $x \sim 0.01$ . This result of a negative gluon at low  $Q^2$  and  $x$  has been confirmed by backwards evolution in [20] (and to a lesser extent in [3]). It will be interesting to see whether a similar conclusion is obtained by other analyses [39, 40]. We anticipate that the evolution from positive definite parton distributions at very low scales [40] will be very difficult to sustain.

The Tevatron jet data constrain the high  $x$  gluon (though from the momentum sum rule and convolutions performed in evolution equations it also affects lower  $x$ ). These provide a far better constraint than any previous data, and from the best global fit we now estimate the uncertainty in the gluon distribution for  $x > 0.1$  and  $Q^2 \sim 20 \text{ GeV}^2$  to be 10 – 15%, with the error decreasing with increasing  $Q^2$ . This removes the need to produce the sets of parton distributions with extreme gluons that were presented in the MRST98 [30] and MRST99 [26] analyses. There is, however, the caveat that if we were to allow the input gluon to have a rather unusual shape and also let  $\alpha_S(M_Z^2)$  be  $\geq 0.120$  we can produce a fit which is much better for the jet data, without too great an expense in  $\chi^2$  for the rest of the data. The best possible jet fit is obtained for  $\alpha_S(M_Z^2) = 0.121$ , and the parton set is denoted by MRST2001J. The gluon is shown in Fig. 15, and the quality of the fit using this set of partons is shown in Table 1. We see that the total  $\chi^2$  is not much higher than the central fit, and is better than the nominal fit for  $\alpha_S(M_Z^2) = 0.121$ . In particular the large high  $x$  gluon helps to counter the deterioration in the fit to BCDMS data for increasing coupling, but has a poor effect on the Drell-Yan E605 data (as discussed in Section 5). However, we reject this as an acceptable set because of the structure in the form of the high  $x$  gluon at low scales, but make it available as an alternative set. It is in some senses similar to the CTEQHJ parton set [41] obtained by forcing the best fits to previous high  $E_T$  jet data, but does not seem to have quite the same features and moreover, we find that our very good fit to jet data can only be achieved without a huge cost in  $\chi^2$  to the fit to other data for  $\alpha_S(M_Z^2) > 0.120$ .

Other than the gluon, and the heavy quark distributions which are generated entirely by evolution and mainly from the gluon, there are no really dramatic changes in our parton distributions, as can be seen in Fig. 2. The inclusion of more deuterium data has caused a slight decrease in the high  $x$  down quark distribution, and a corresponding increase for  $x \sim 10^{-2}$ , but there is nothing else too significant. Indeed, our central value of  $\alpha_S(M_Z^2)$  has not changed much either. Though the changes are small, they are important for the precision predictions of the  $W$  and  $Z$  production cross sections at the LHC and the Tevatron. Table 2 shows the changes

Table 2: Predictions in  $nb$  for  $W$  and  $Z$  production at the Tevatron and LHC, compared with those of MRST99 [26].

|          |                                 | MRST99 | MRST2001 |
|----------|---------------------------------|--------|----------|
| Tevatron | $B_{\ell\nu} \cdot \sigma_W$    | 2.45   | 2.49     |
|          | $B_{\ell+\ell-} \cdot \sigma_Z$ | 0.226  | 0.228    |
| LHC      | $B_{\ell\nu} \cdot \sigma_W$    | 20.3   | 20.8     |
|          | $B_{\ell+\ell-} \cdot \sigma_Z$ | 1.87   | 1.91     |

in the predictions for these cross sections when going from the default MRST99 parton set [26] to the present set where, for ease of comparison, we have kept the electroweak parameters unchanged. We see that the predicted cross sections have increased by about 1–2%, caused in large part by the increase in the down quark distribution in the relevant  $(x, Q^2)$  range, see the second plot in Fig. 2. The uncertainty of such predictions and the influence on the parton uncertainty will be the subject of a future paper.

As a final point we note that the overall quality of the NLO-in- $\alpha_S$  fit remains fairly good. The raised cut in  $W^2$ , from 10 GeV<sup>2</sup> to 12.5 GeV<sup>2</sup> has removed some deficiencies in the high  $x$  slope which may be due to higher twist or higher orders in  $\alpha_S(Q^2)$ . Also, it is noticed that investigating cuts in  $Q^2$  implies little evidence for higher twist at general  $x$ . However, some of the previous areas of concern have been increased rather than reduced. It is a worrying point that the minimum  $\chi^2$  values for many individual data sets within the global fit (Fig. 16) lie outside the range  $\alpha_S(M_Z^2) = 0.116 - 0.122$ . Also, as in previous fits, there is also still a struggle to get a steep enough evolution of  $F_2(x, Q^2)$  in the region  $x \sim 0.01$  as is seen in Figs. 5 and 6. Moreover, it is also difficult to obtain enough high  $x$  gluon to get a very good fit to the jet data. These two points, coupled with the rather slow evolution of  $F_2(x, Q^2)$  at the lowest  $x$ , combine to produce a gluon which has gone from being valence-like to very negative at the input scale  $Q_0^2 = 1$  GeV<sup>2</sup>. While this is not necessarily a problem in itself, it has resulted in a prediction for  $F_L(x, Q^2)$  that is worryingly small at very small  $x$  (Fig. 8). Hence there are implications of problems at small  $x$ . We have not really considered the effect of a lower  $x$  cut in this paper, but will demonstrate in a future paper [23] that investigating fits with low  $x$  data cut out does have a serious effect on the partons and has strong implications on the real success of the standard NLO-in- $\alpha_S$  fit at low  $x$ . Remember that in the same way that the high  $x$  form of the gluon imposed by jet data influences small  $x$  via the sum rule and convolutions, cutting out small  $x$  data can influence the fit and partons at higher  $x$ . However, the effect of varying the  $x$  cut on the value of  $\alpha_S$  is minimal, since the major constraint comes from the evolution of the high  $x$  partons. It is, therefore, not surprising that the value remains well within our quoted experimental error,  $0.119 \pm 0.002$ . On a related point we have already noticed that extending the theory to (an approximate) NNLO-in- $\alpha_S$  does lead to a general improvement in the quality of the fit, and to some significant changes in partons and predictions, particularly

at small  $x$ . We have not considered NNLO at all here, but will produce detailed results in a forthcoming paper [36].

In summary, in this paper we have used all deep inelastic and hadron collider data available in order to obtain the most accurate and precise determination of the NLO parton distributions currently in existence, and have also determined the value of the strong coupling constant  $\alpha_S(M_Z^2)$  with tight constraints. This enables us to probe the success of the conventional NLO perturbative QCD framework in describing hadronic collider physics, and we find that overall it is still working well. This then provides us with the necessary starting point for predicting and explaining new physics coming from present and future particle colliders.

The FORTRAN code for the four NLO parton sets mentioned in Table 1 can be found at <http://durpdg.dur.ac.uk/hepdata/mrs>

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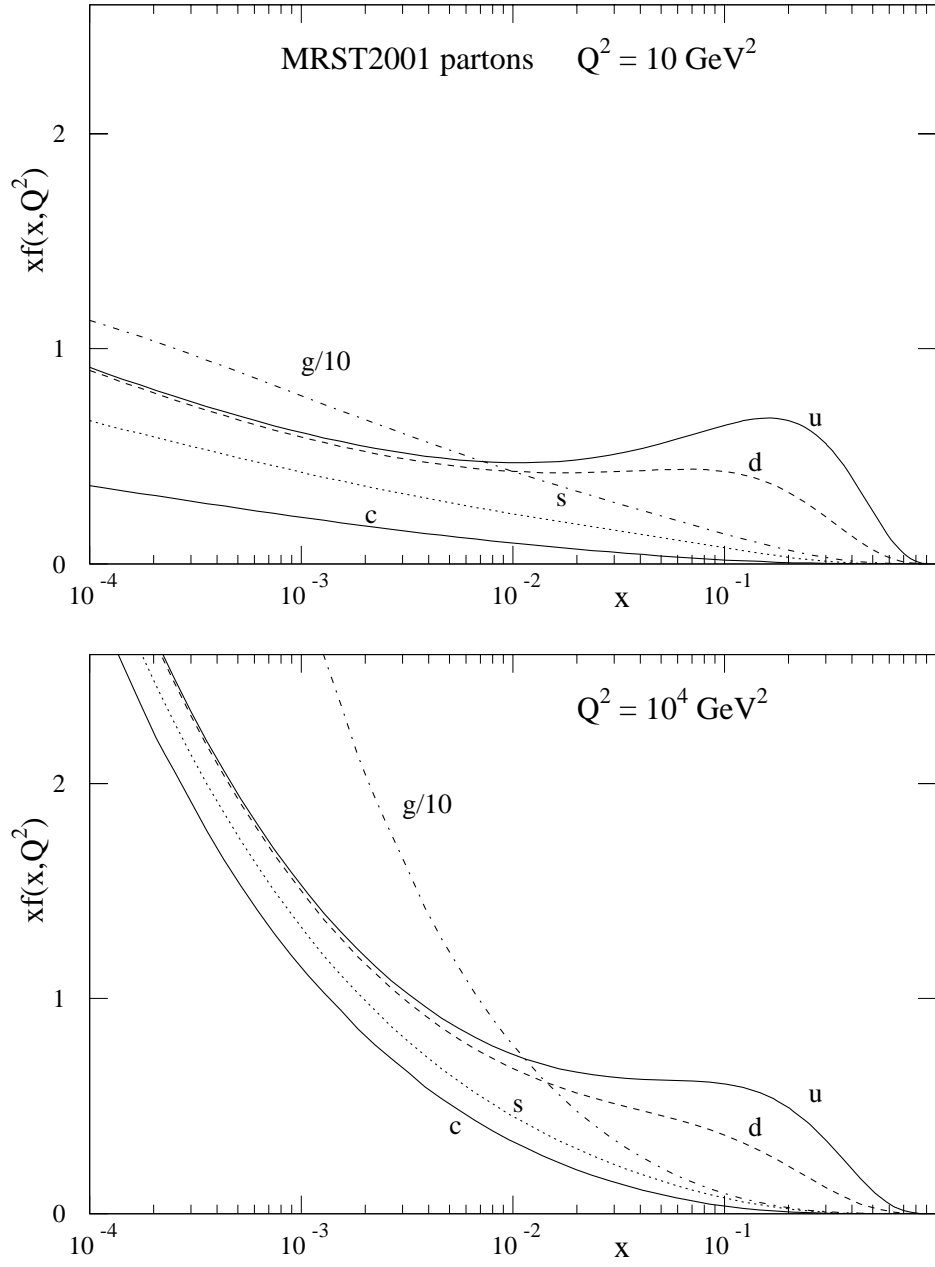


Figure 1: MRST2001 partons at  $Q^2 = 10 \text{ GeV}^2$  and  $Q^2 = 10^4 \text{ GeV}^2$ .

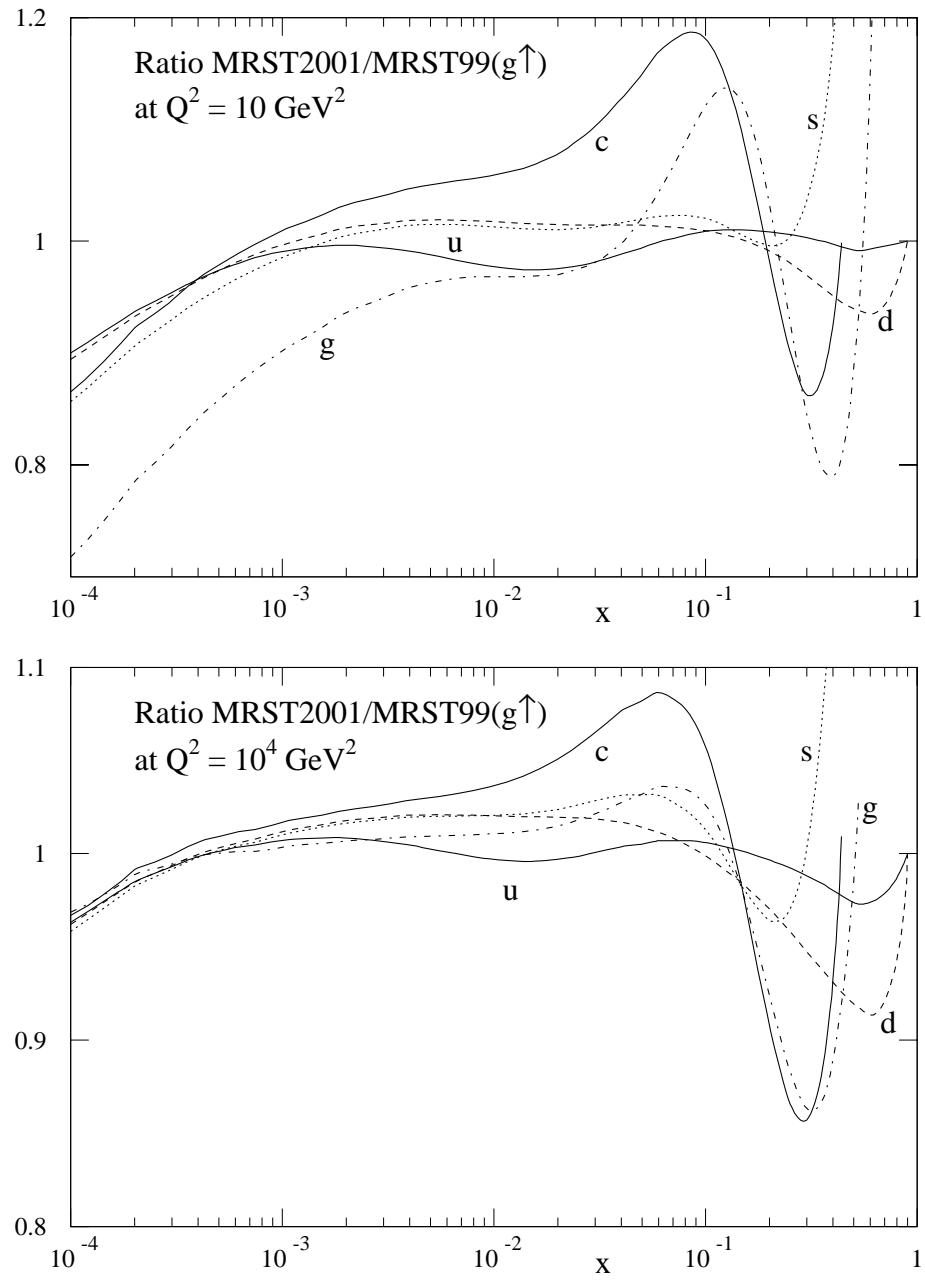


Figure 2: Comparison of the MRST2001 partons with those of MRST99( $g \uparrow$ ) [26] at  $Q^2 = 10 \text{ GeV}^2$  and  $Q^2 = 10^4 \text{ GeV}^2$ .

MRST(2001) NLO fit ,  $x = 0.00005 - 0.00032$

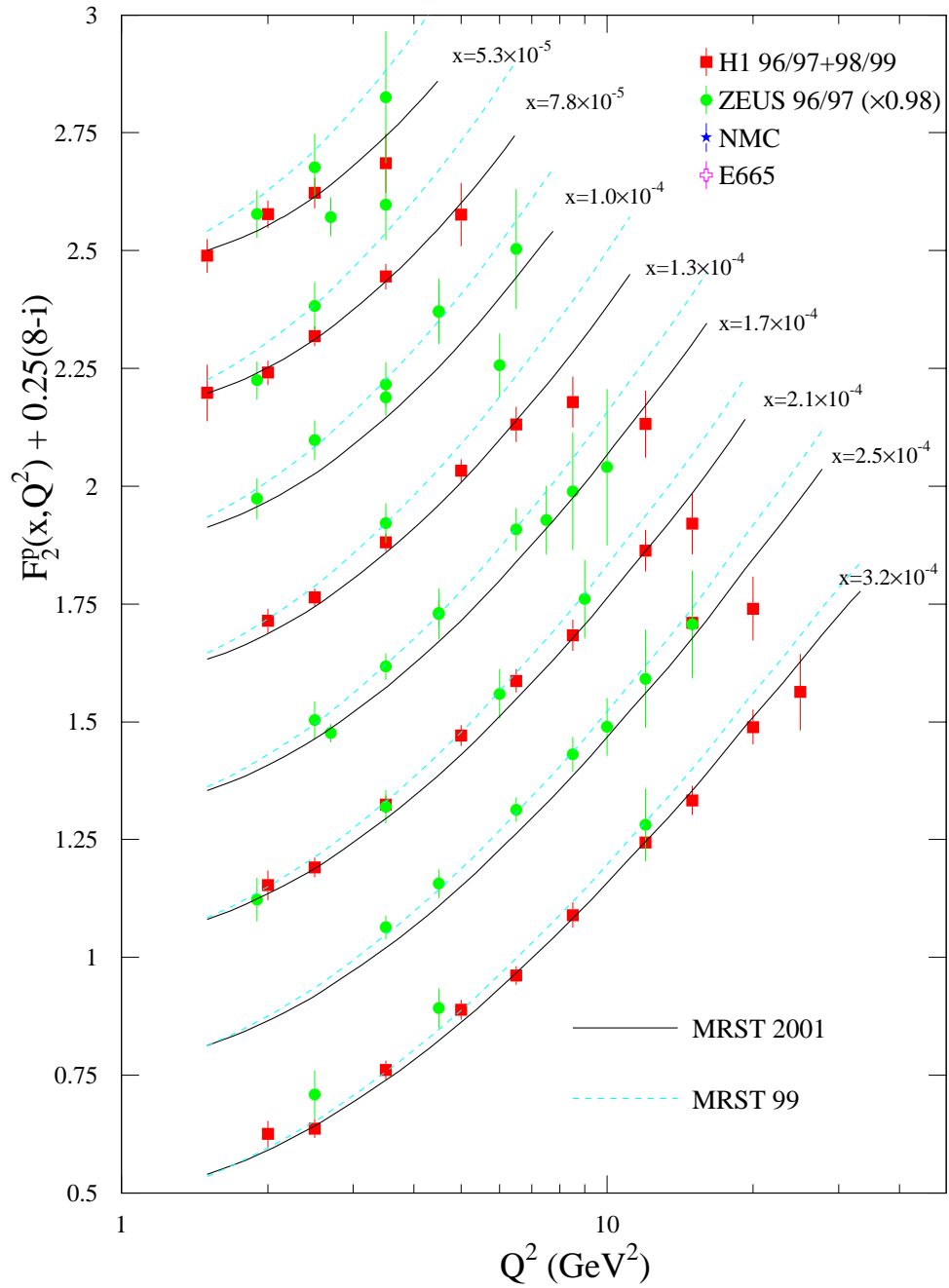


Figure 3: Comparison of the MRST2001 prediction of  $F_2(x, Q^2)$  with data and with of MRST99 for  $x = 0.00005 - 0.00032$ .

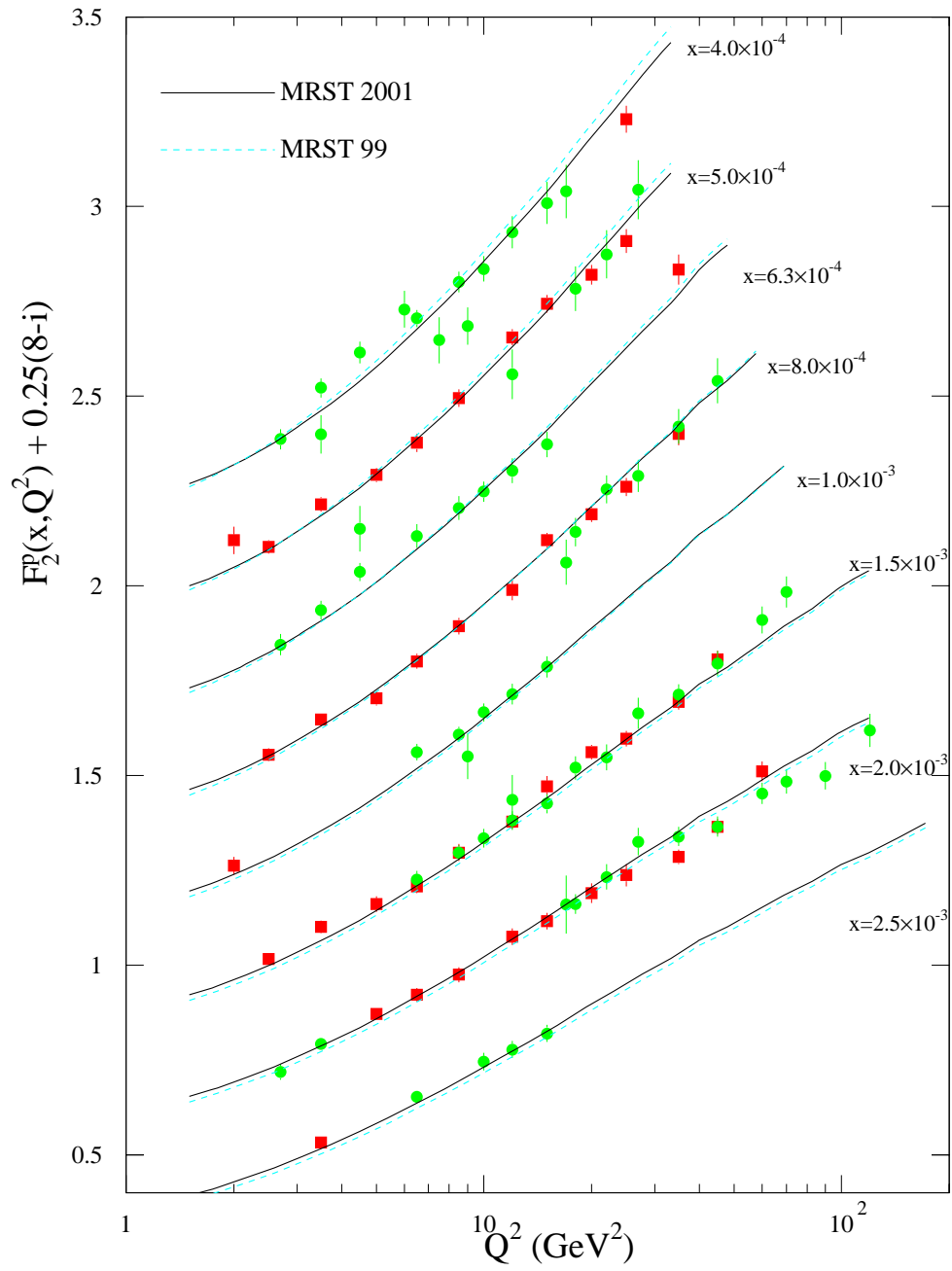


Figure 4: Comparison of the MRST2001 prediction of  $F_2(x, Q^2)$  with data and with of MRST99 for  $x = 0.0004 - 0.0025$ .

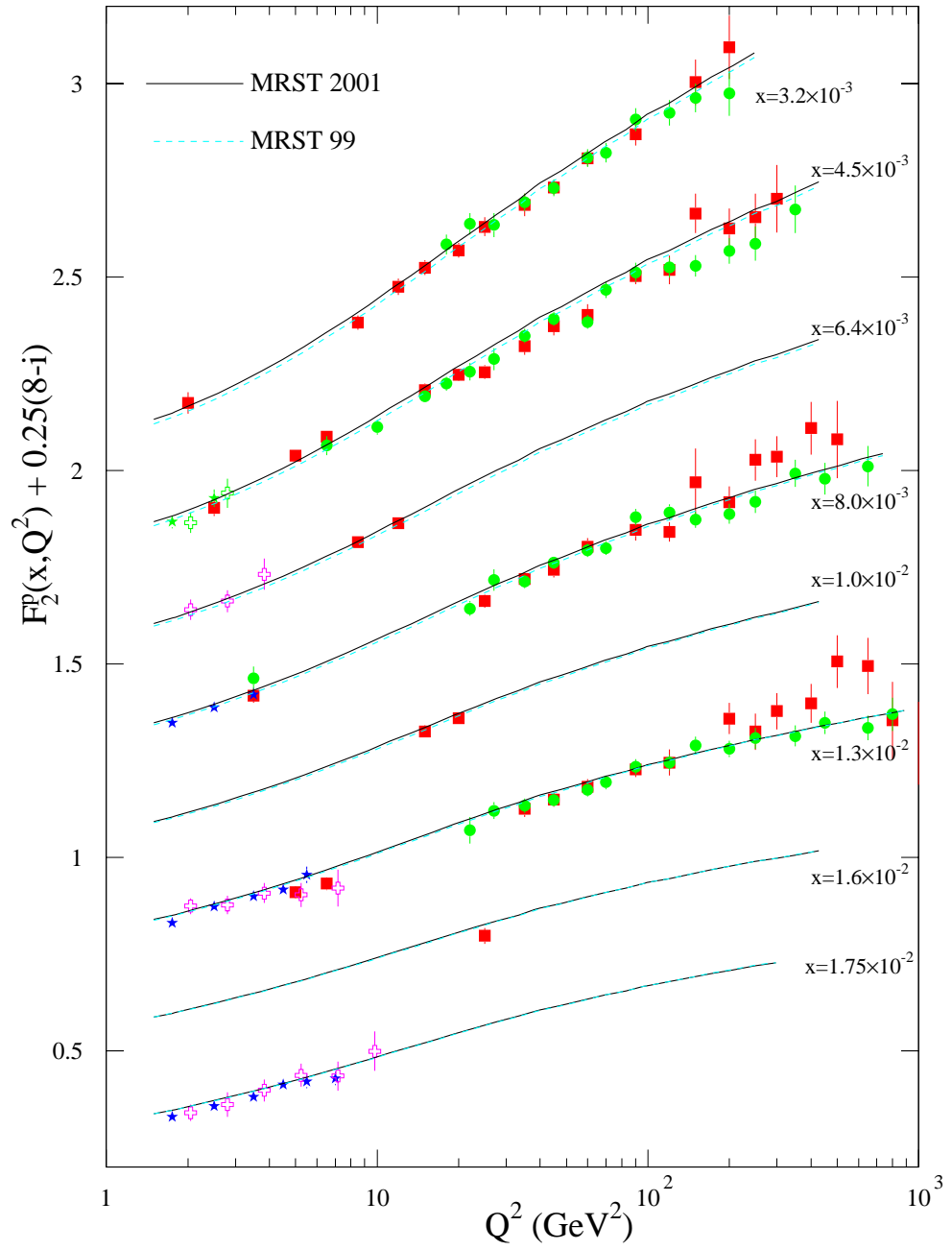


Figure 5: Comparison of the MRST2001 prediction of  $F_2(x, Q^2)$  with data and with of MRST99 for  $x = 0.0032 - 0.0175$ .

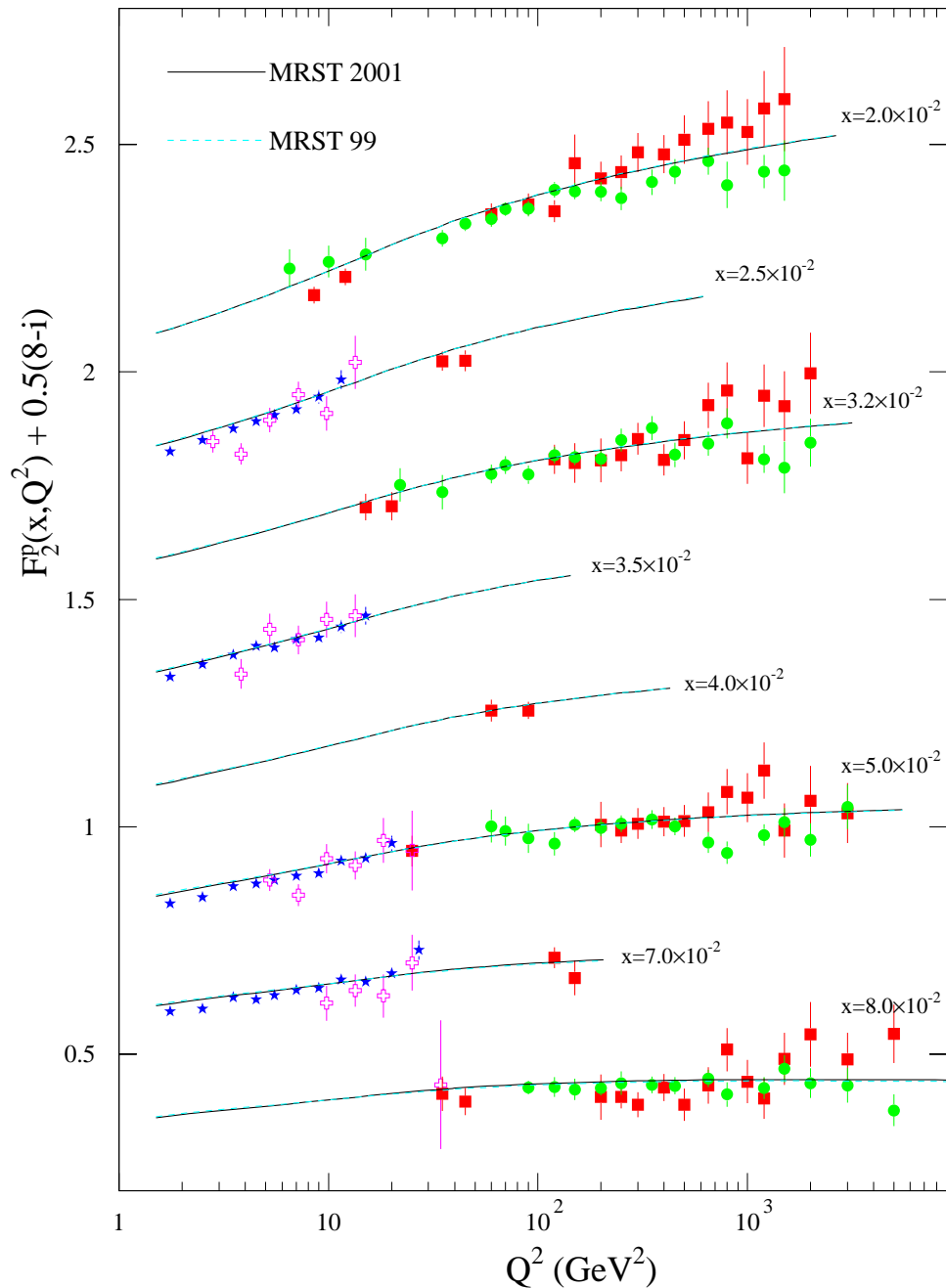


Figure 6: Comparison of the MRST2001 prediction of  $F_2(x, Q^2)$  with data and with of MRST99 for  $x = 0.02 - 0.08$ .

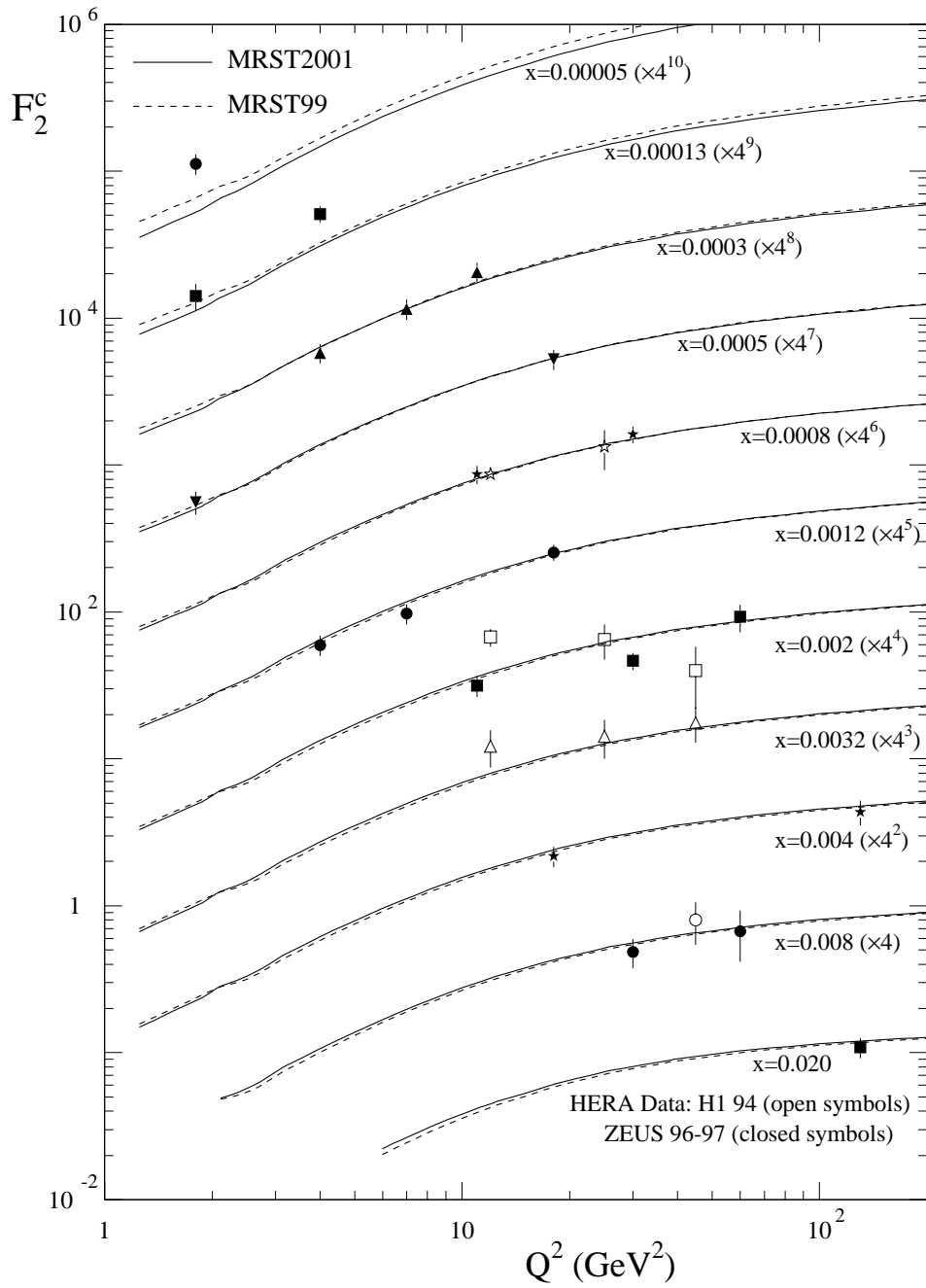


Figure 7: Comparison of the MRST2001 prediction of  $F_2^c(x, Q^2)$  with data [7, 27].



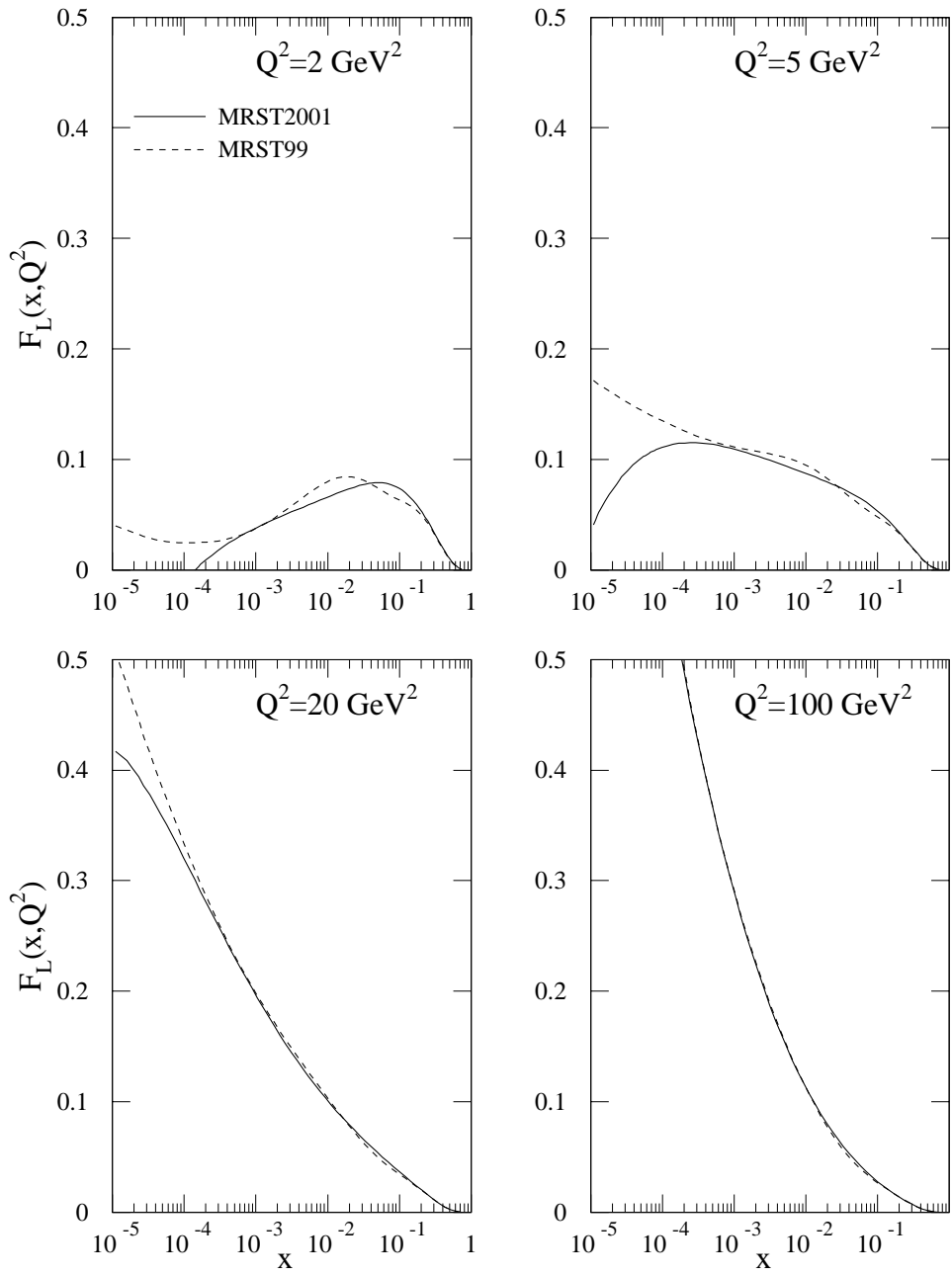


Figure 8: Comparison of the MRST2001 prediction of  $F_L(x, Q^2)$  with that of MRST99 for various values of  $Q^2$ .

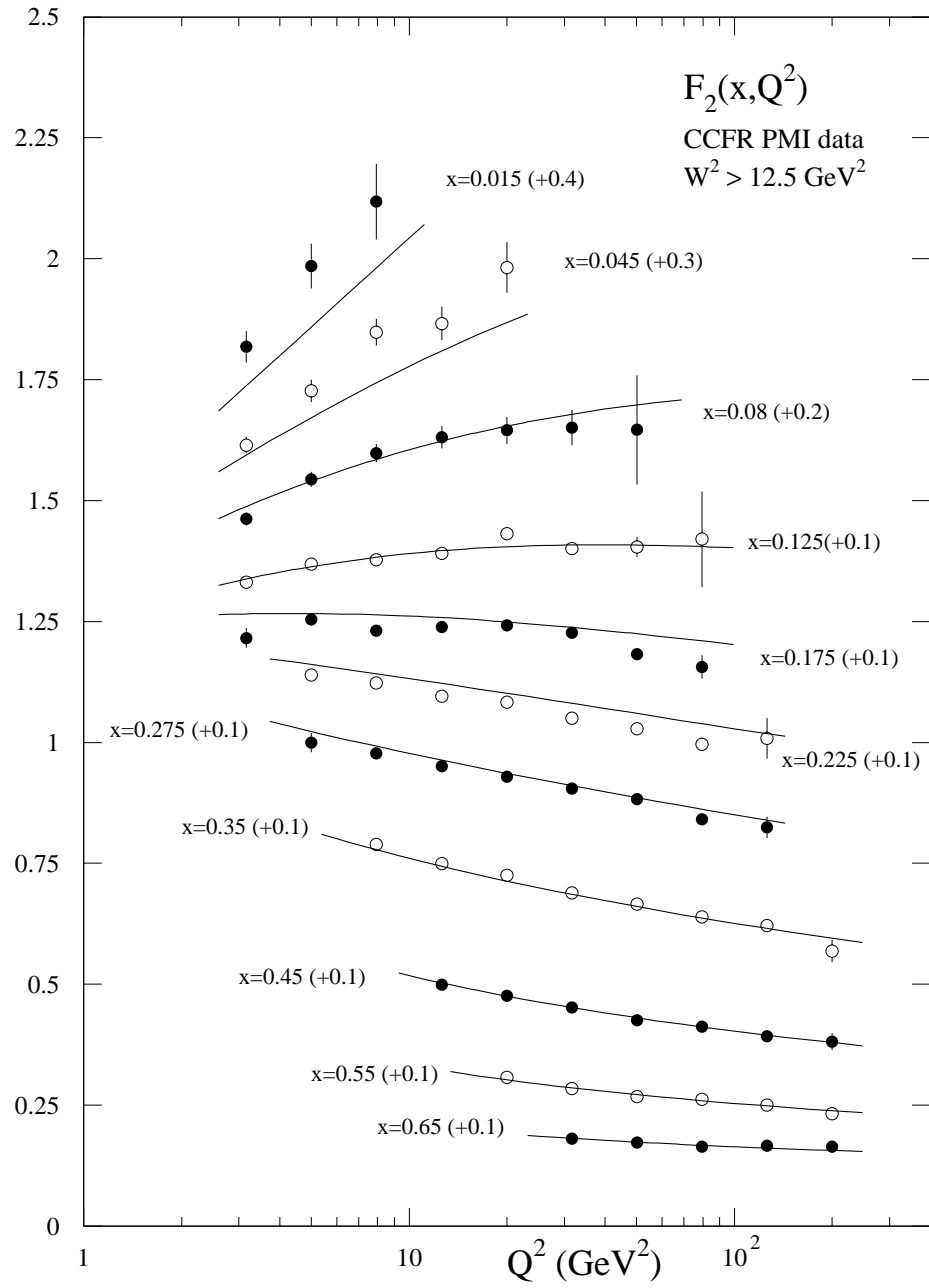


Figure 9: Quality of the MRST2001 fit to the CCFR  $F_2^{\nu(\bar{\nu})N}(x, Q^2)$  PMI data [8].

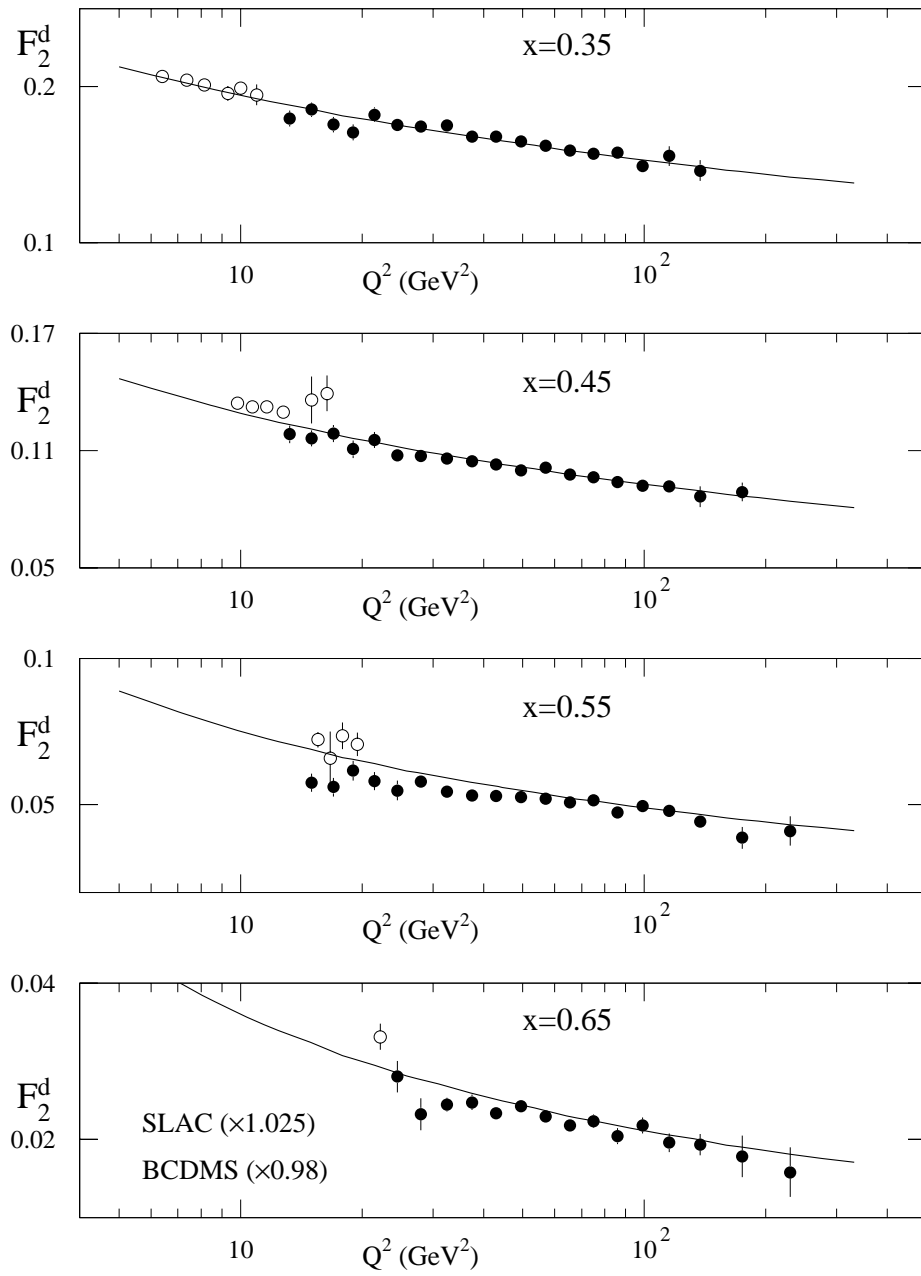


Figure 10: Quality of the MRST2001 fit to the SLAC and BCDMS deuterium structure function data [12, 18] at high  $x$ .

MRST 2001 and D0 jet data,  $\alpha_s(M_Z)=0.119$ ,  $\chi^2=106/82$  pts

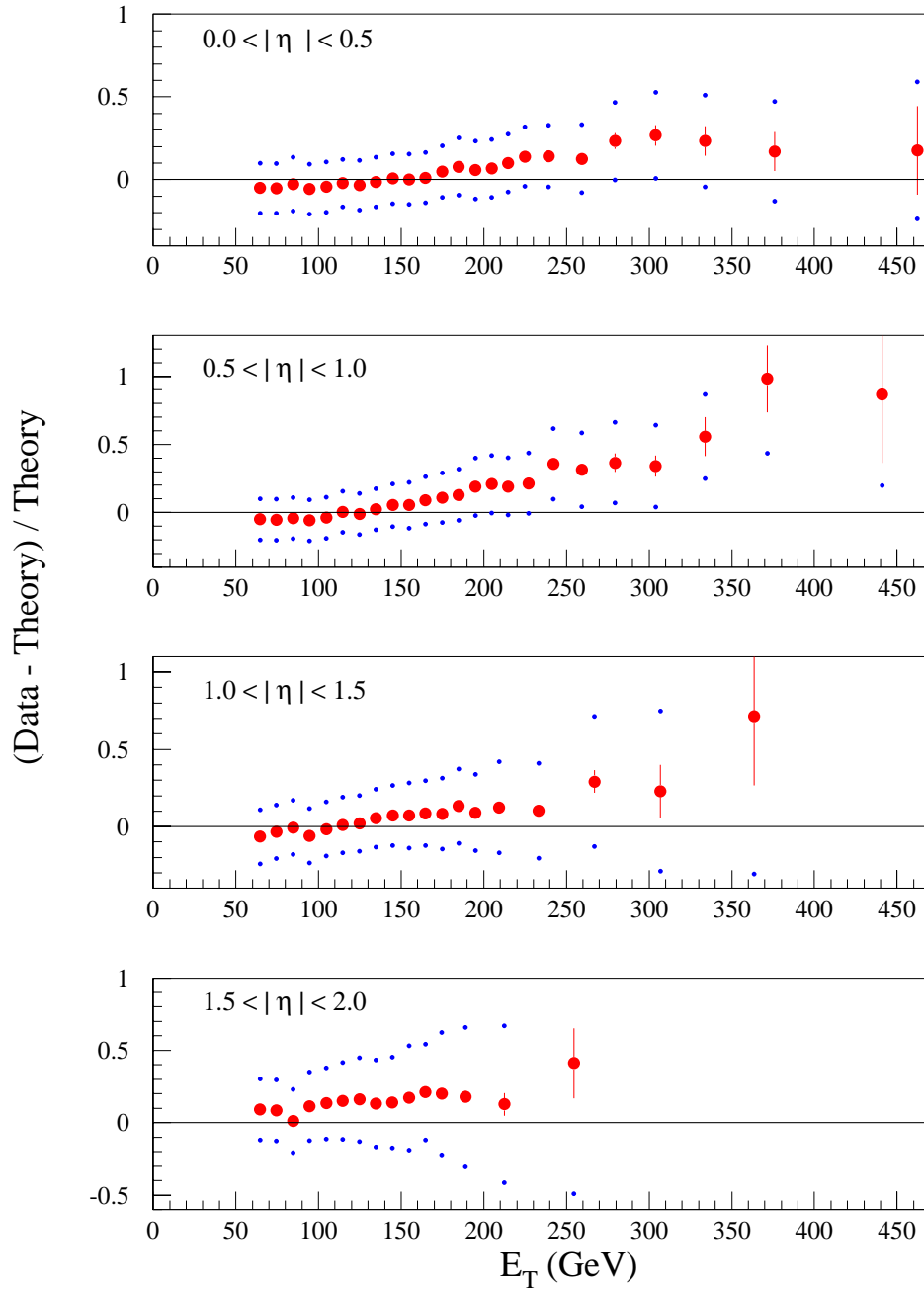


Figure 11: Quality of the MRST2001 fit to the D0 high  $E_T$  jet data in different  $\eta$  bins [5]. The band shows the size of the correlated errors.

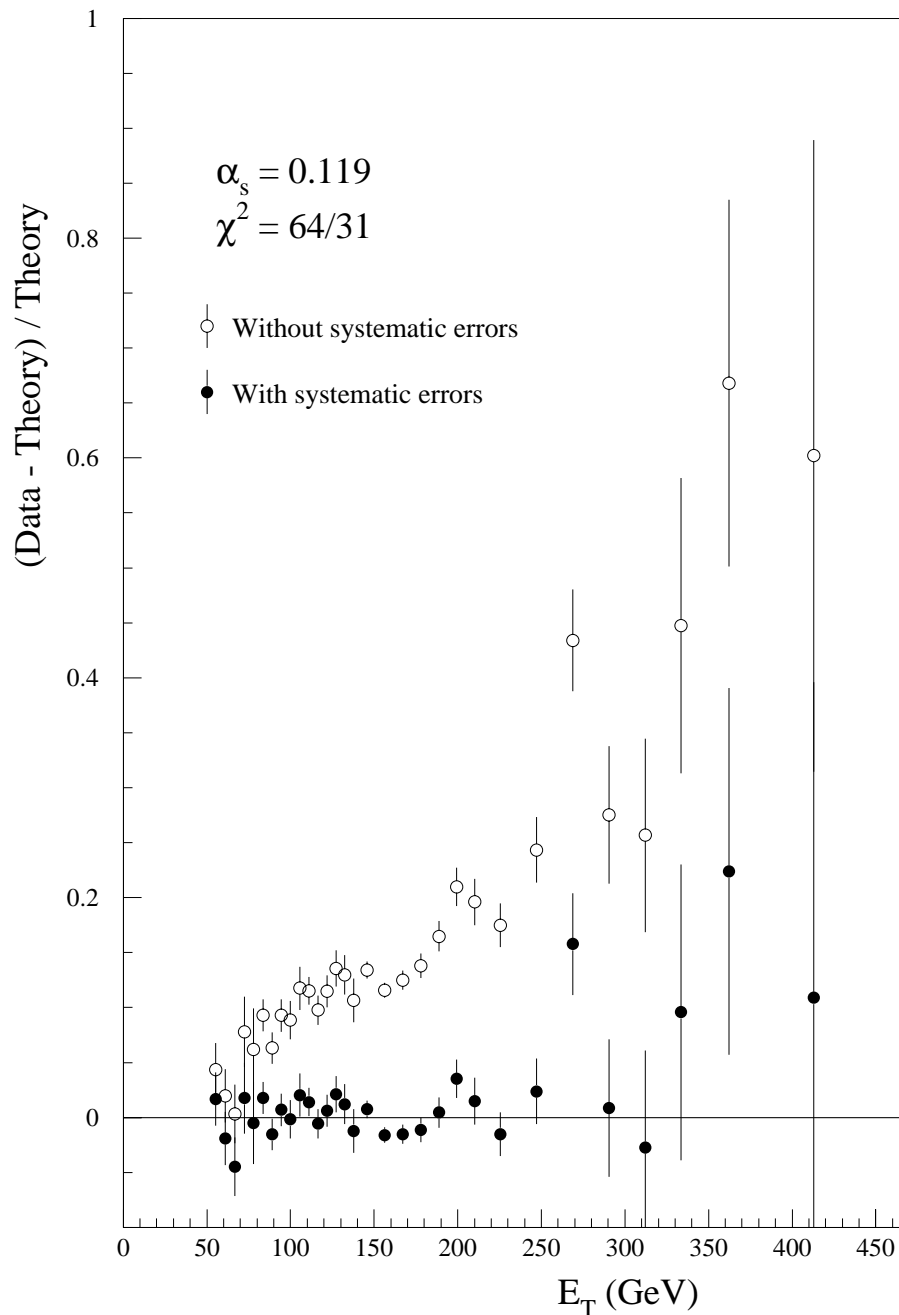


Figure 12: Quality of the MRST2001 fit to the CDF1B high  $E_T$  jet data [6]. The open points are before correlated systematic errors have been considered, while the solid points are after the correlated errors have allowed the data-theory comparison to move at some cost to the total  $\chi^2$  (shown on the plot).

MRST 2001J and D0 jet data,  $\alpha_s(M_Z)=0.121$ ,  $\chi^2=45/82$  pts

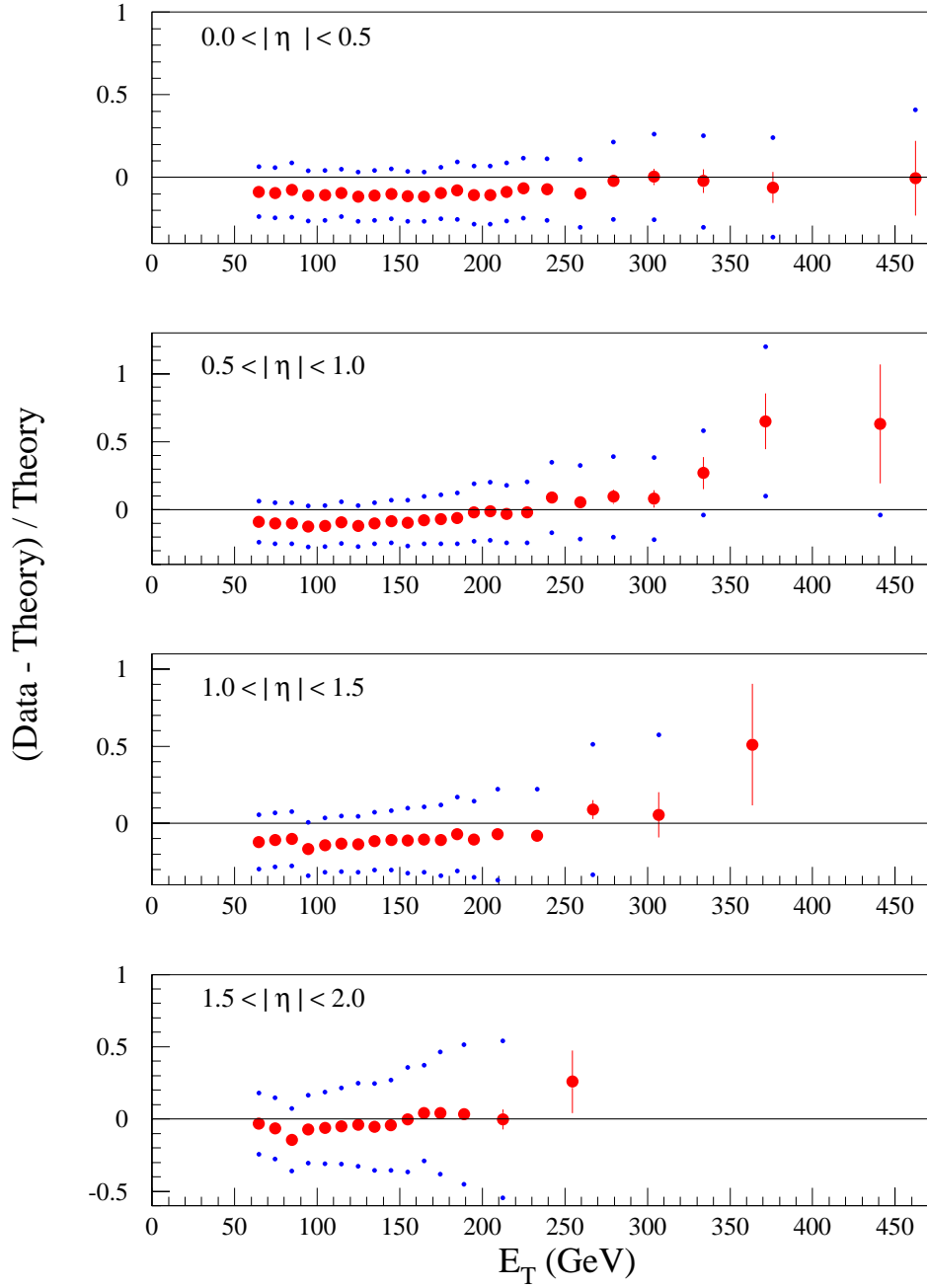


Figure 13: Quality of the MRST2001J fit to the D0 high  $E_T$  jet data [5] in different  $\eta$  bins. The band shows the size of the correlated errors.

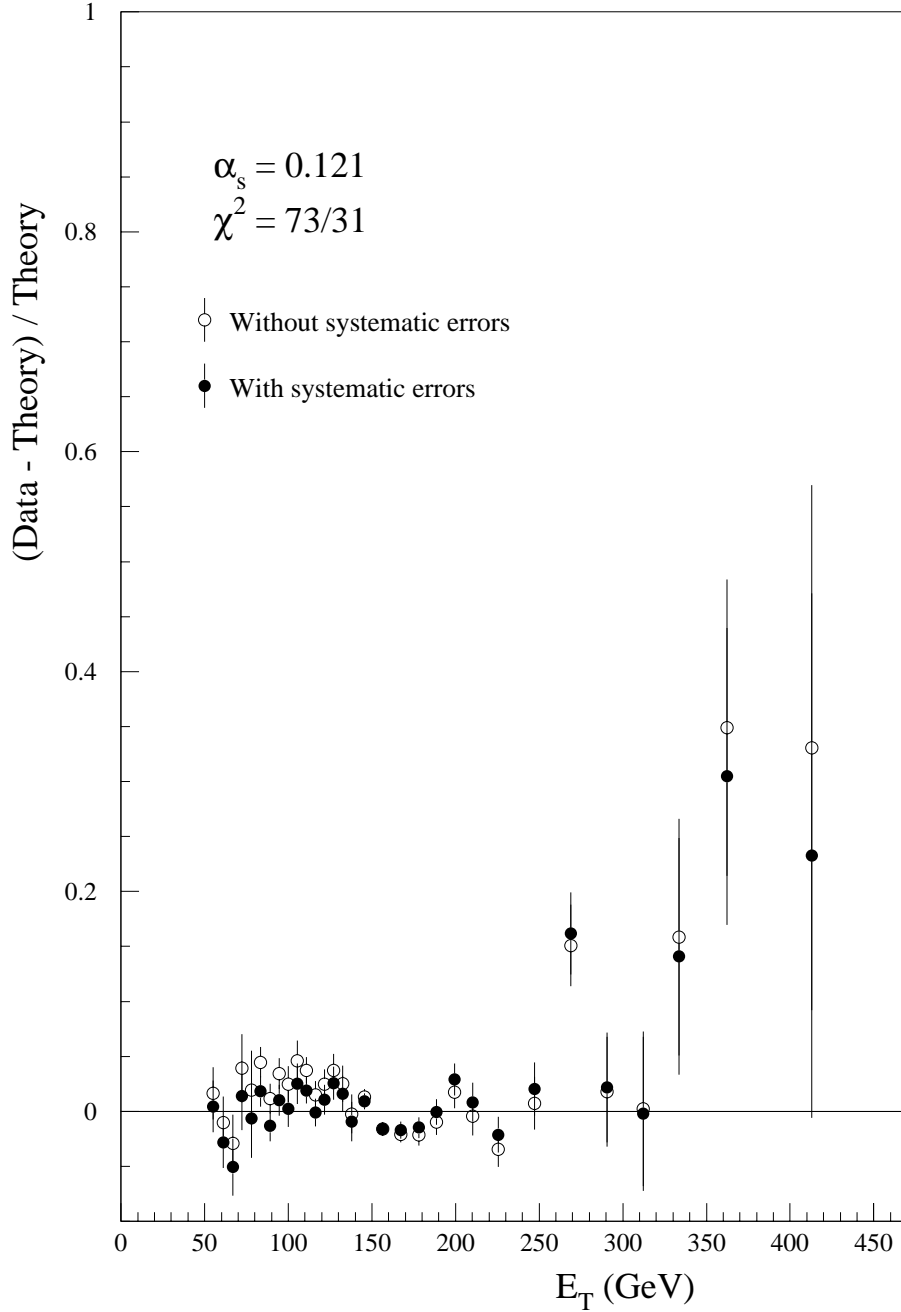


Figure 14: Quality of the MRST2001J fit to the CDF1B high  $E_T$  jet data [6]. The open points are before correlated systematic errors have been considered, while the solid points are after the correlated errors have allowed the data-theory comparison to move at some cost to total  $\chi^2$  (shown on the plot).

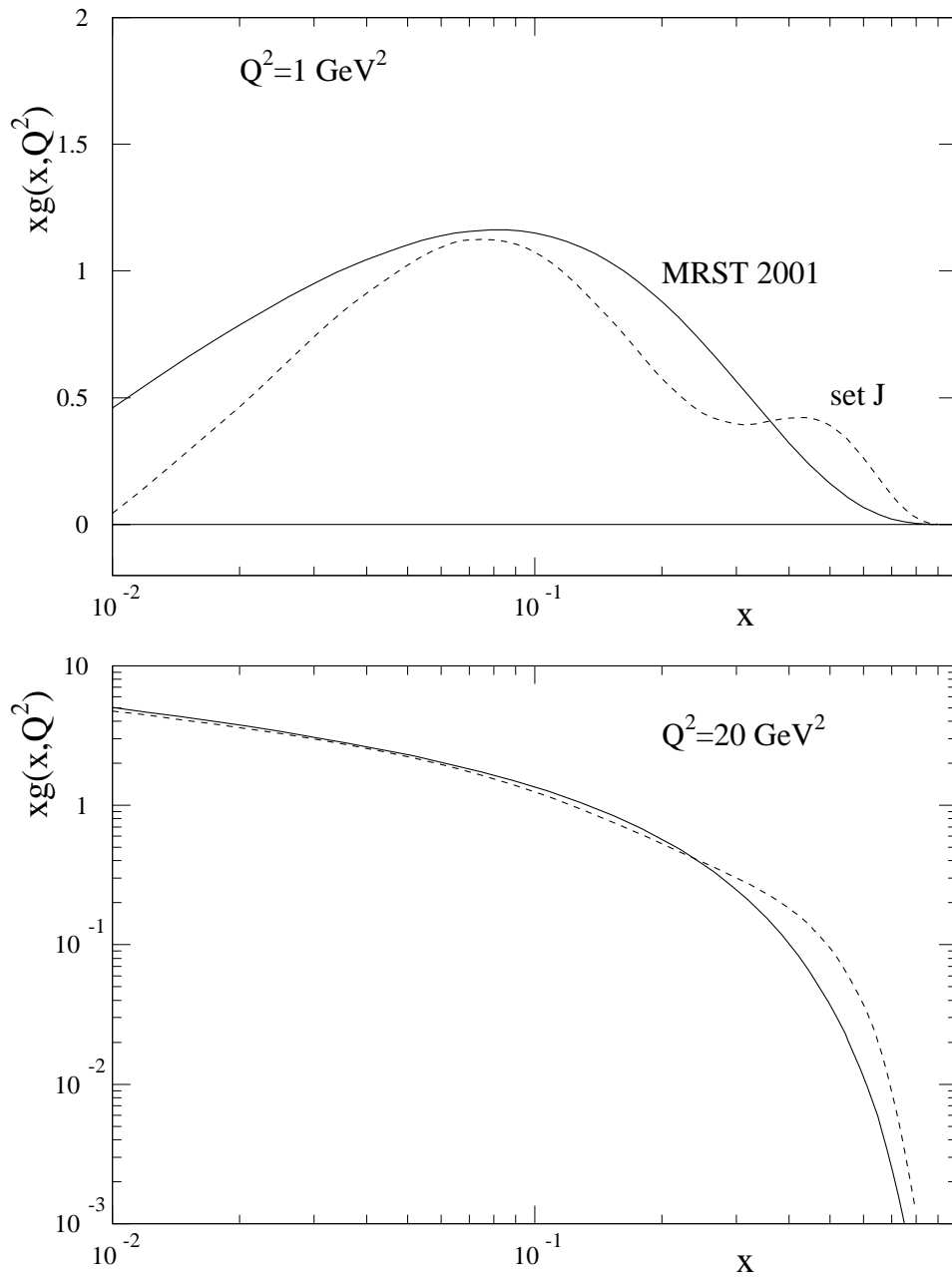


Figure 15: Comparison of the MRST2001J gluon with the MRST2001 gluon at high  $x$  and  $Q^2 = 1 \text{ GeV}^2$  and  $Q^2 = 20 \text{ GeV}^2$ .



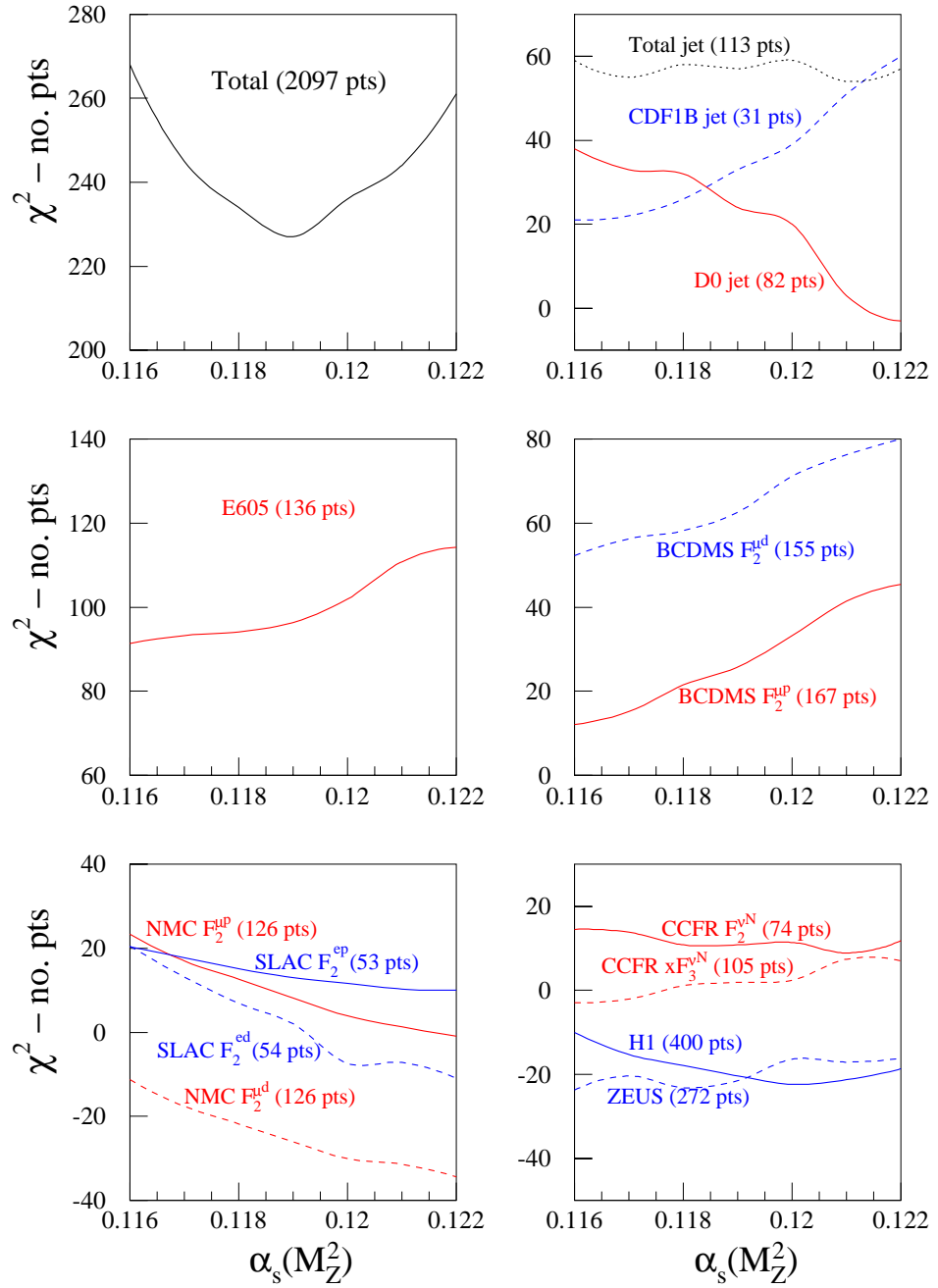
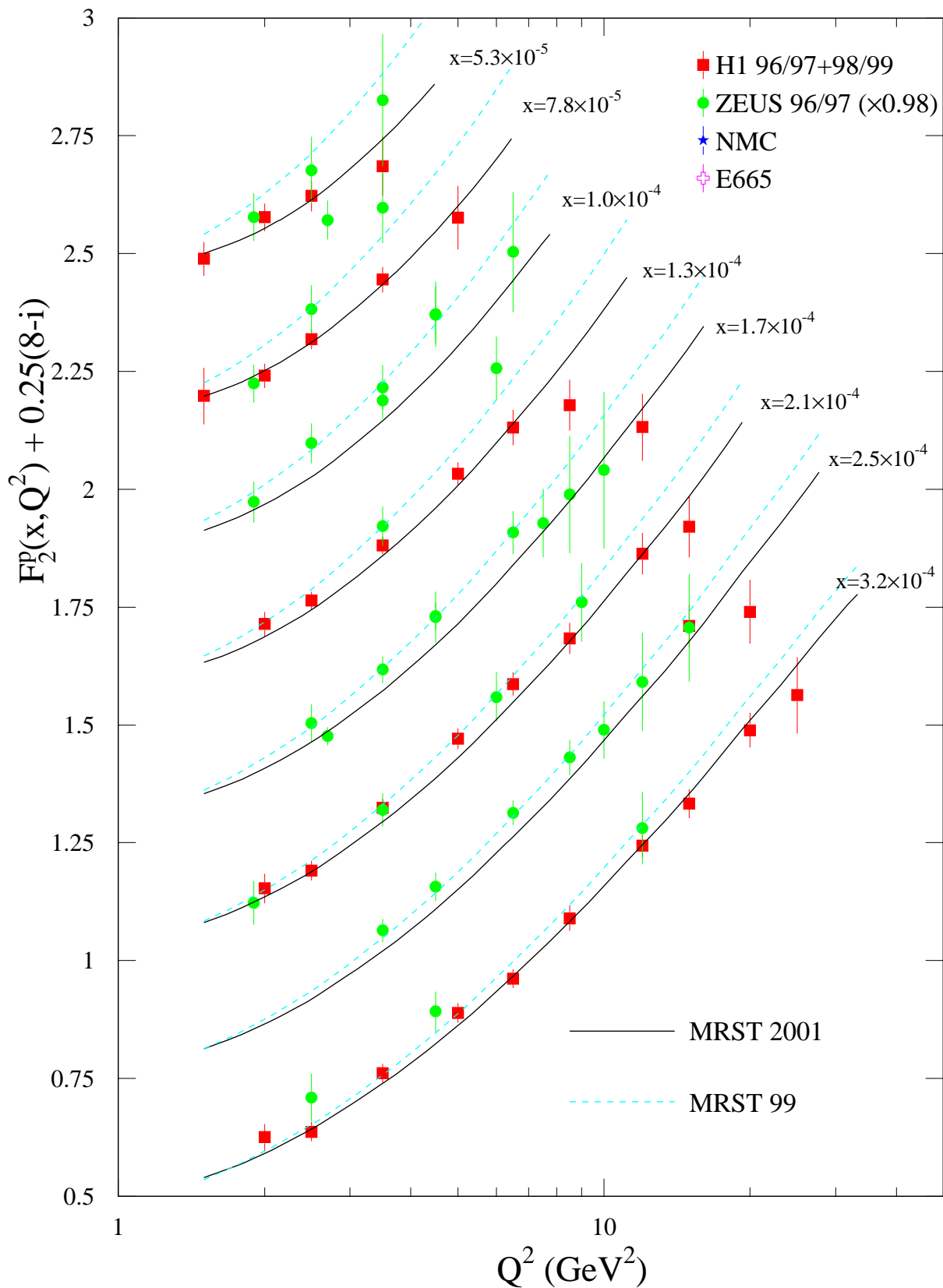
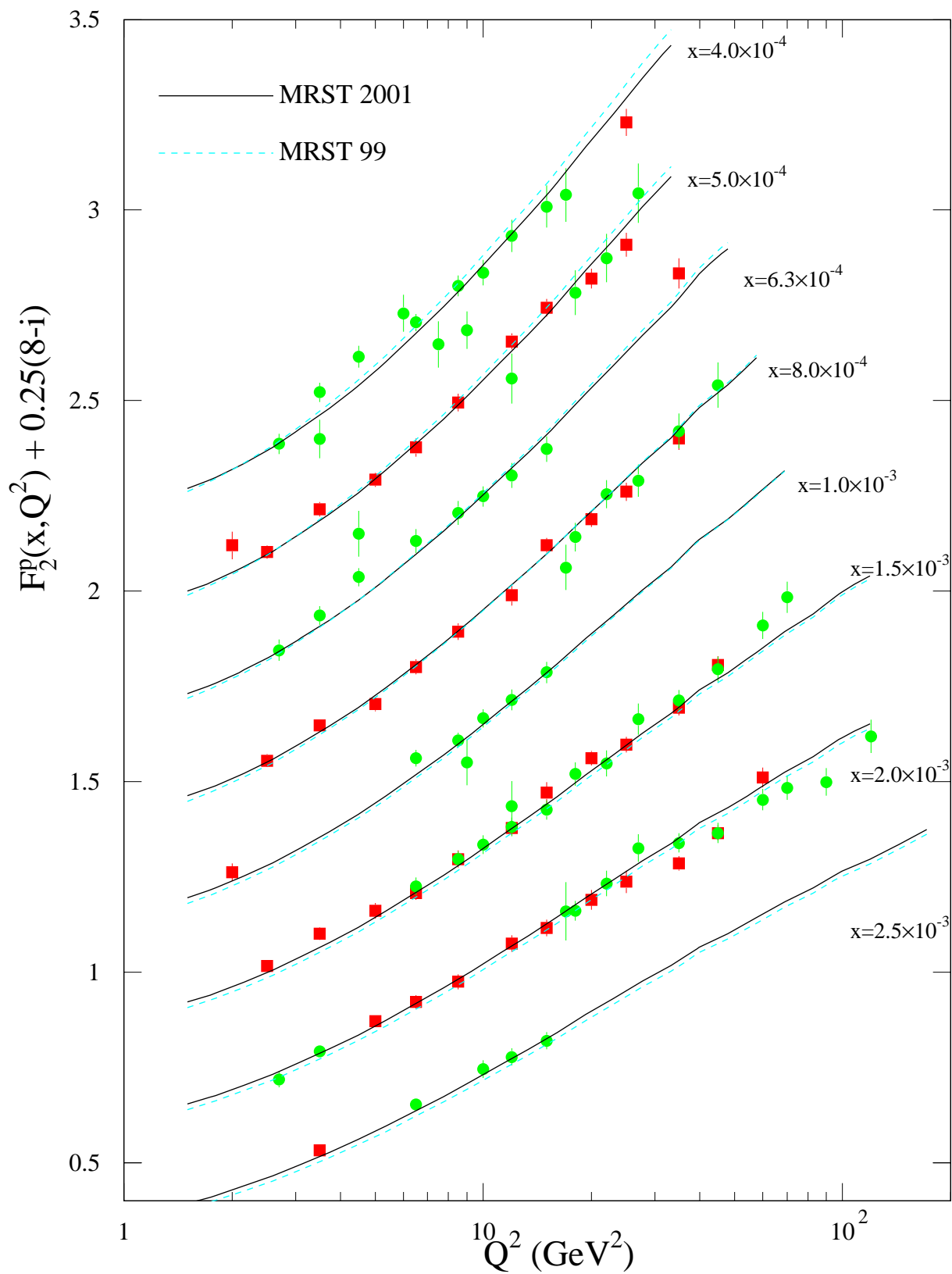


Figure 16: The quality of the fit to the individual data sets included in the global analysis, shown together with the grand total  $\chi^2$ , as a function of  $\alpha_s(M_Z^2)$ .

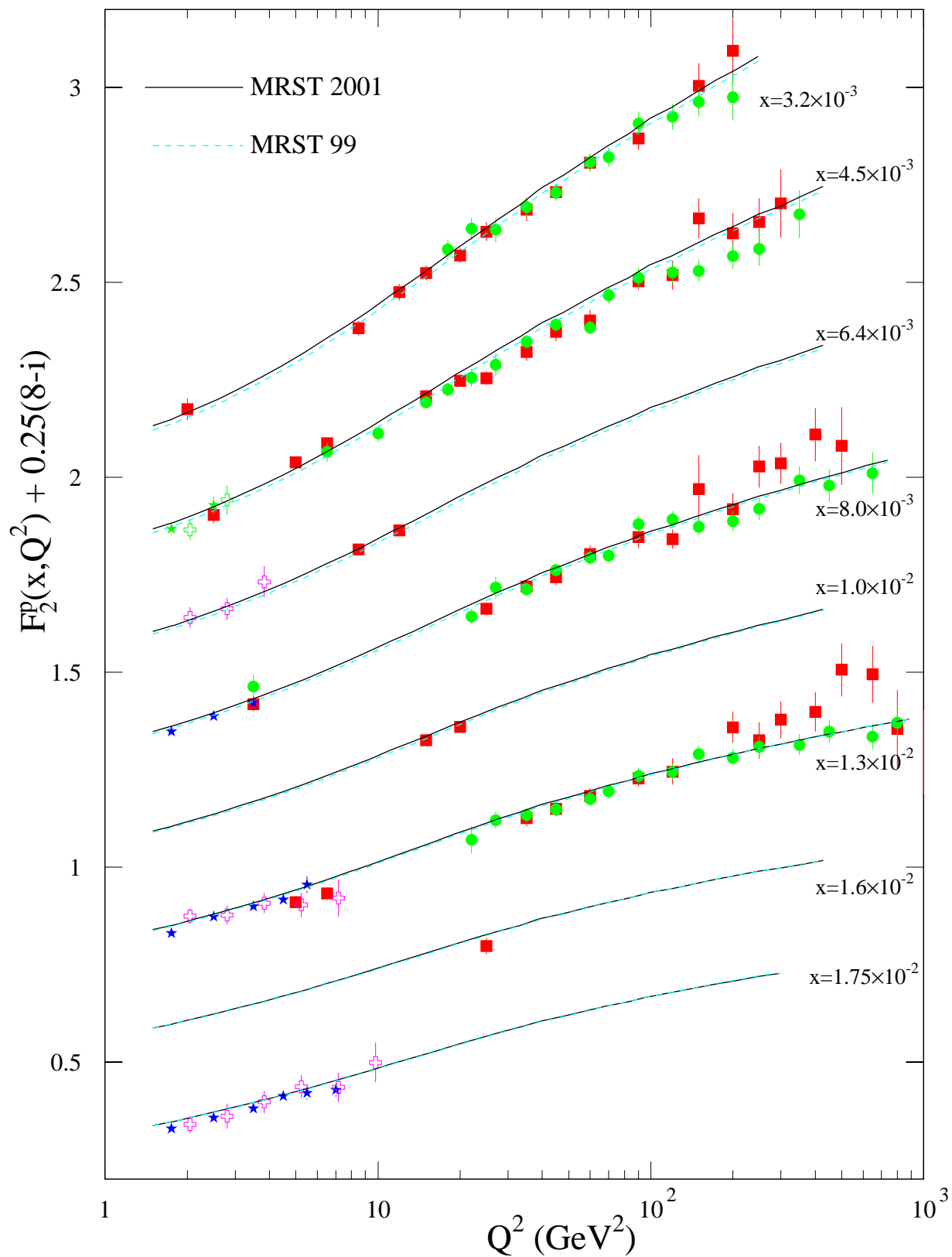
# MRST(2001) NLO fit , $x = 0.00005 - 0.00032$



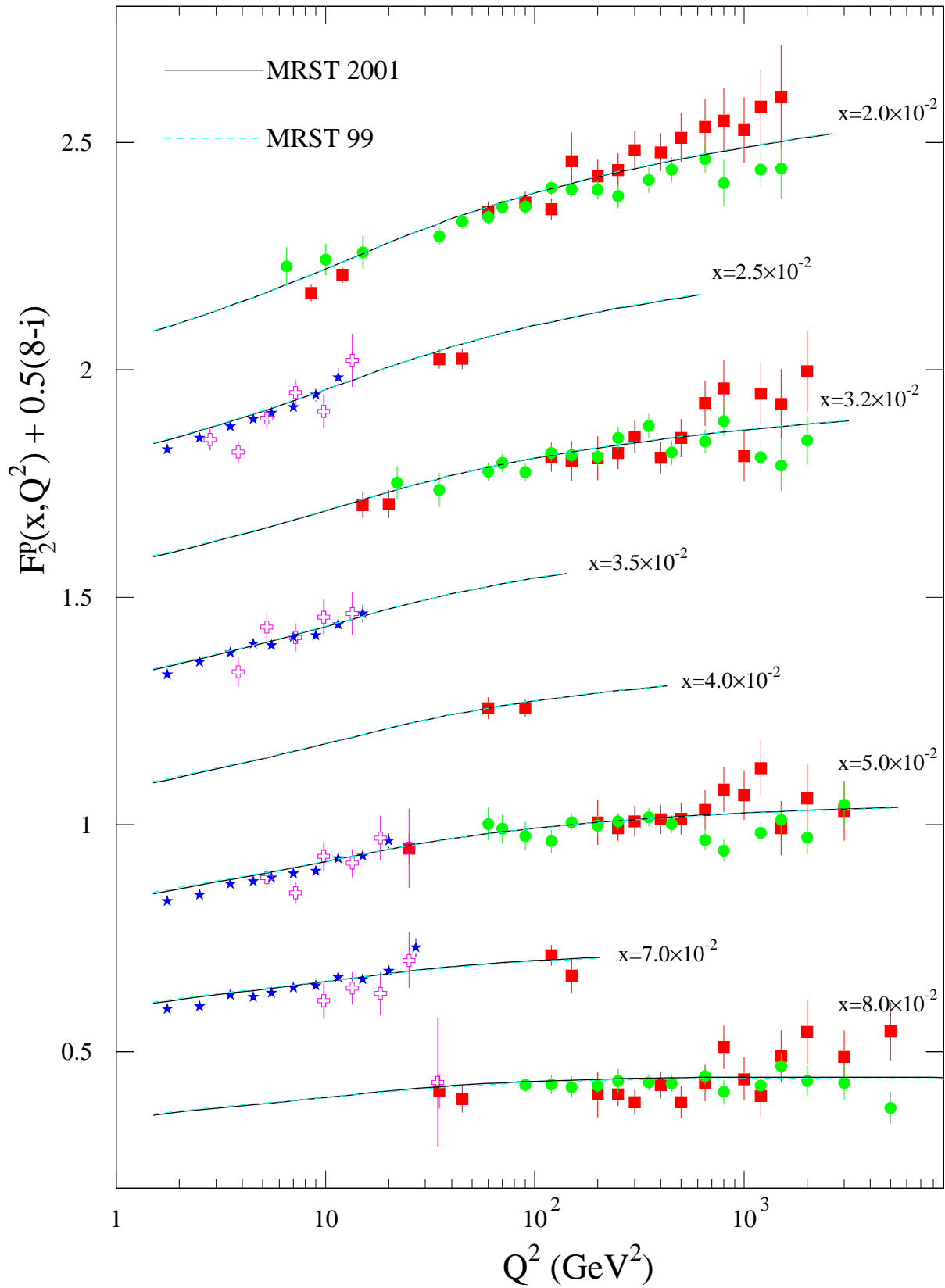
# MRST(2001) NLO fit , $x=0.0004 - 0.0025$



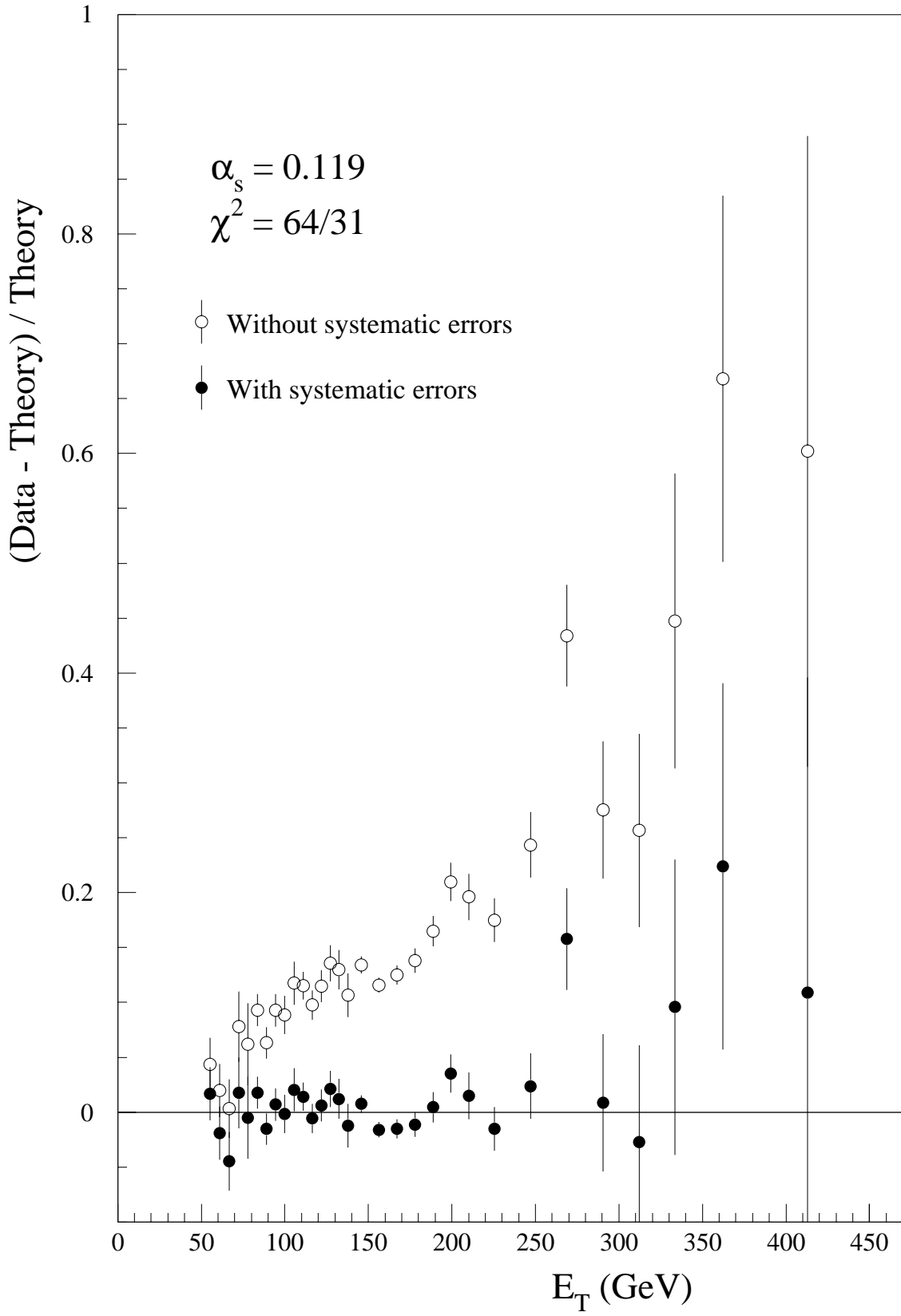
# MRST(2001) NLO fit , $x= 0.0032 - 0.0175$

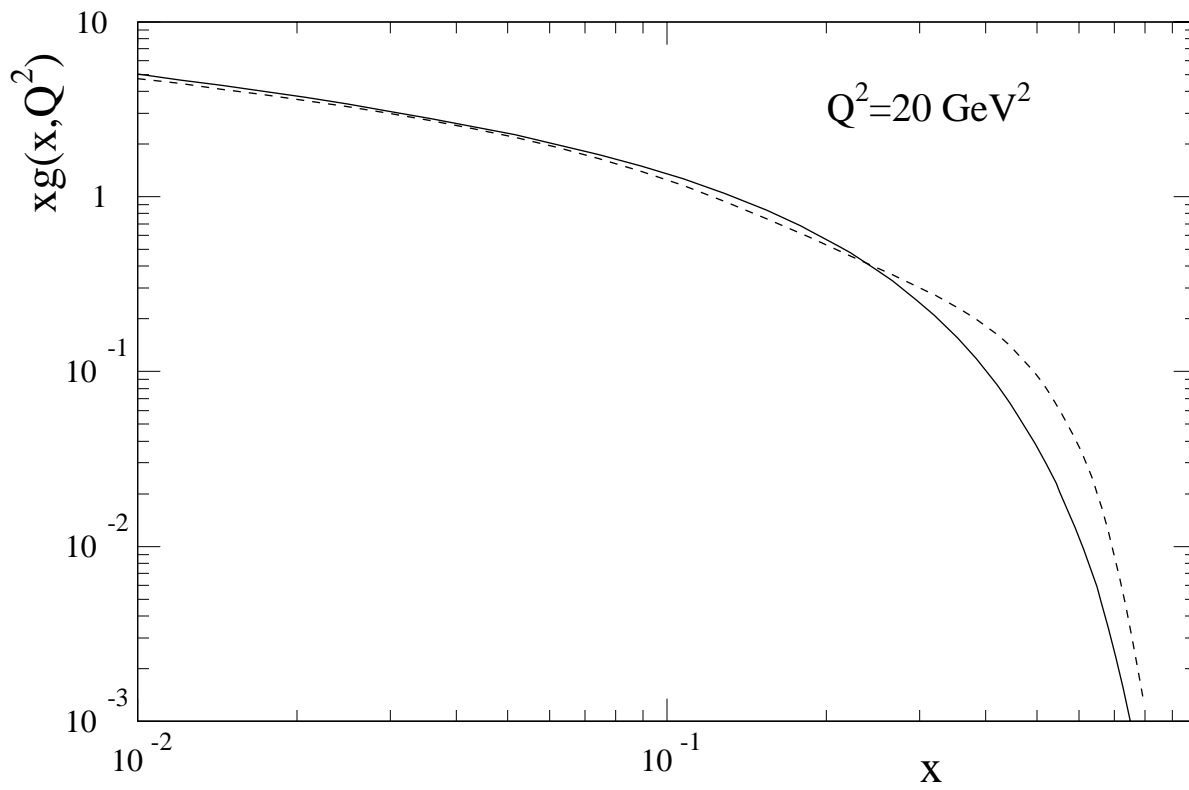
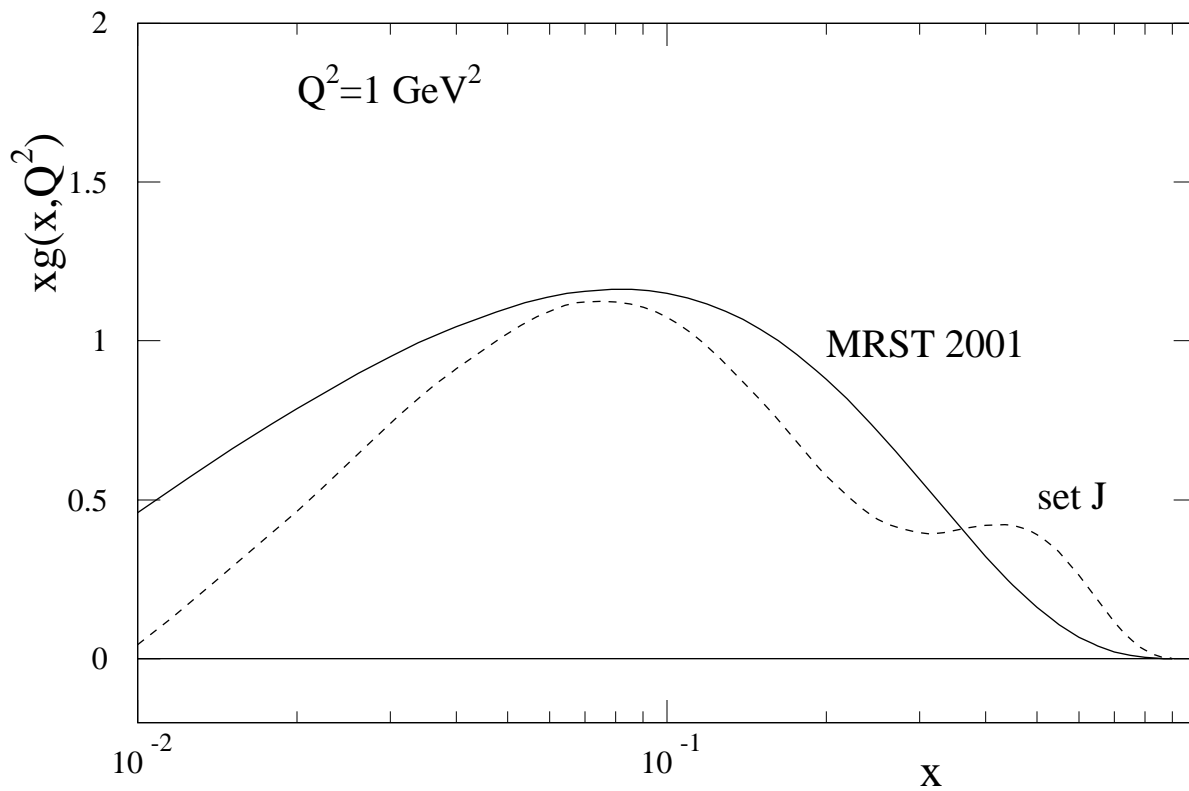


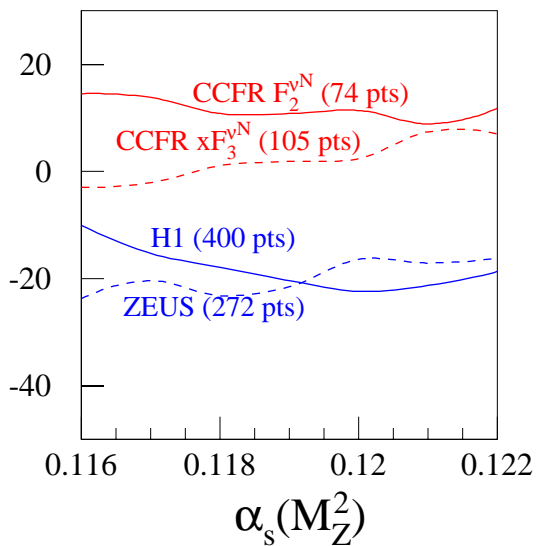
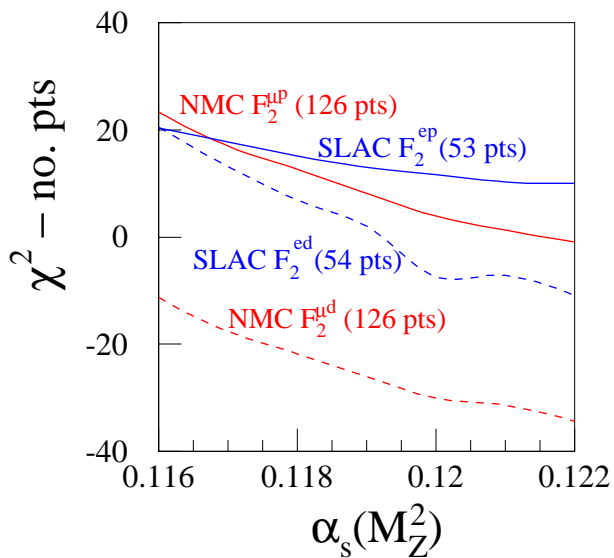
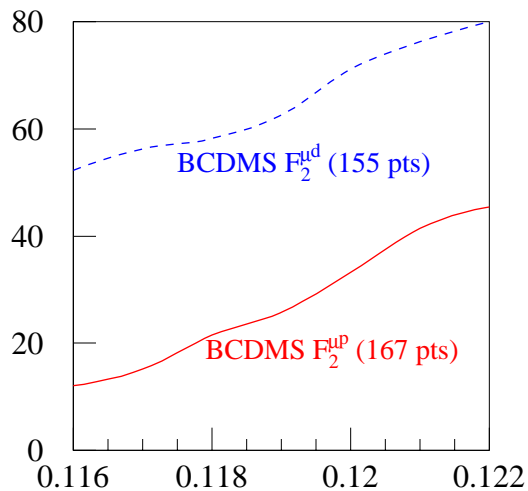
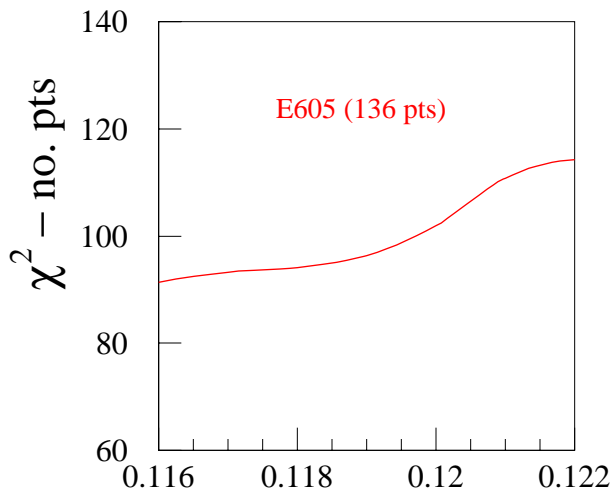
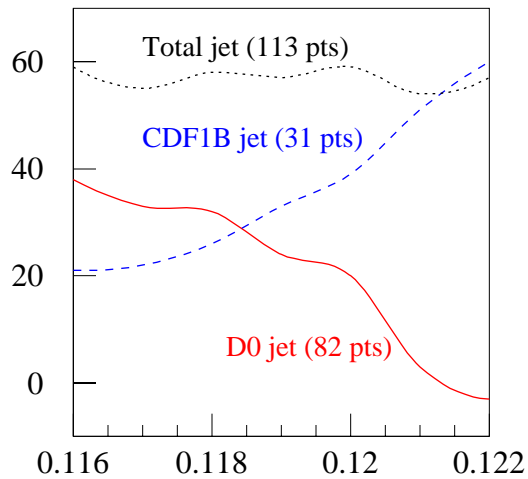
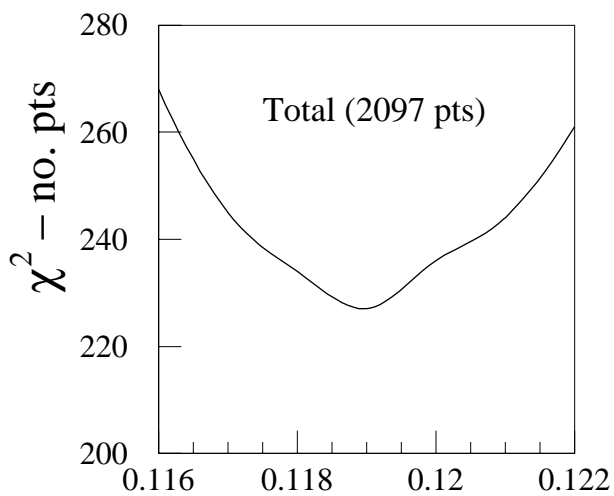
# MRST(2001) NLO fit , $x=0.02 - 0.08$



MRST 2001 and CDF1B jet data

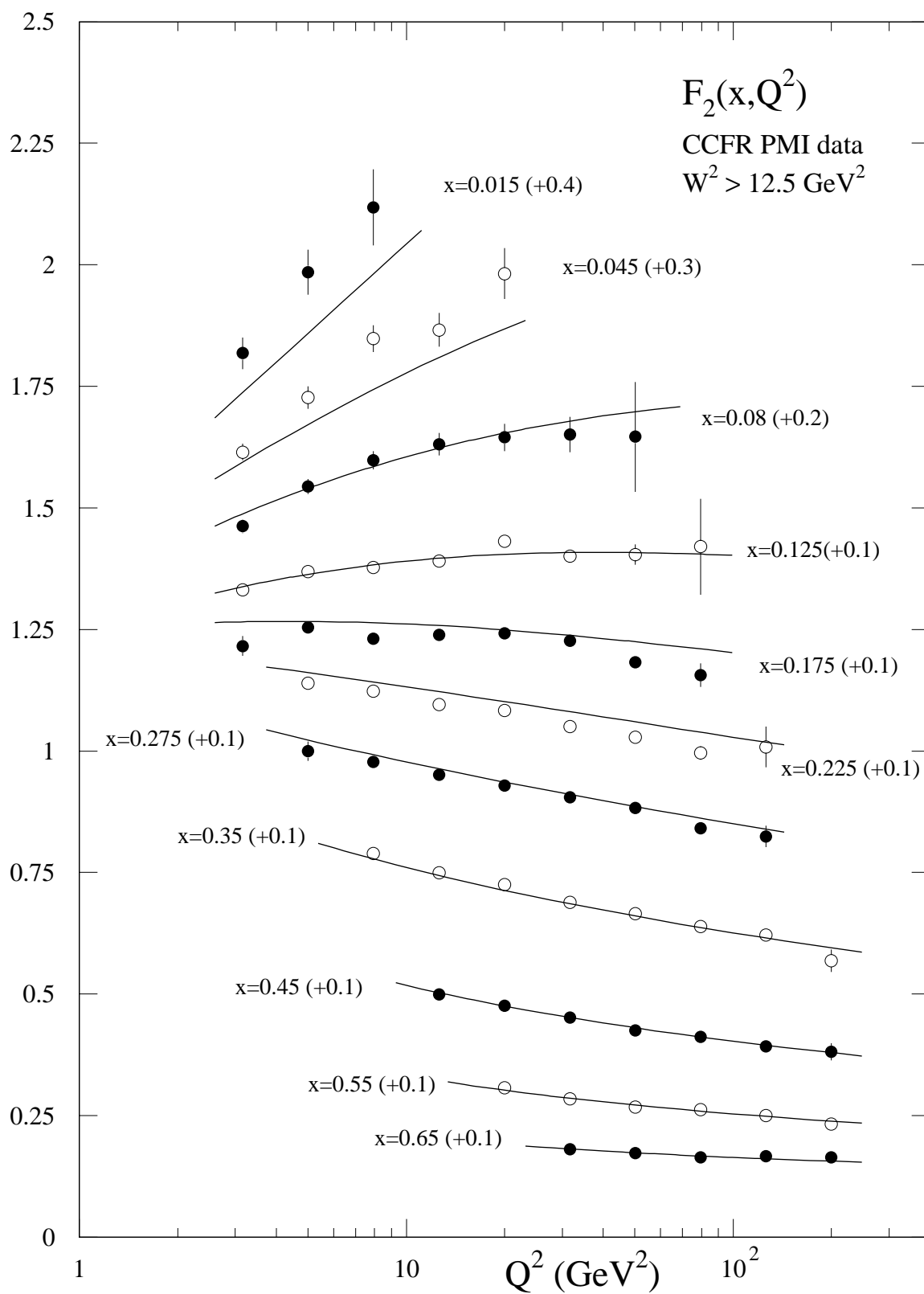








New CCFR data comparison with MRST 2001



Comparison of MRST2001 partons with MRST99 partons

